

The shelf-to-basin transport of iron from the Northern U.S West Coast to the Pacific Ocean

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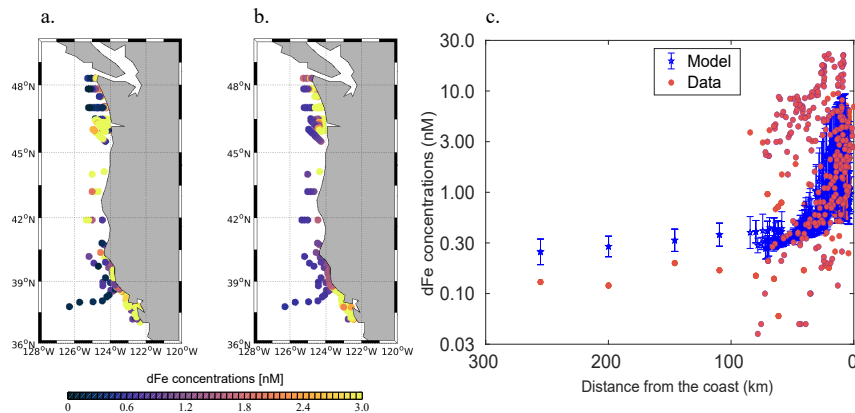
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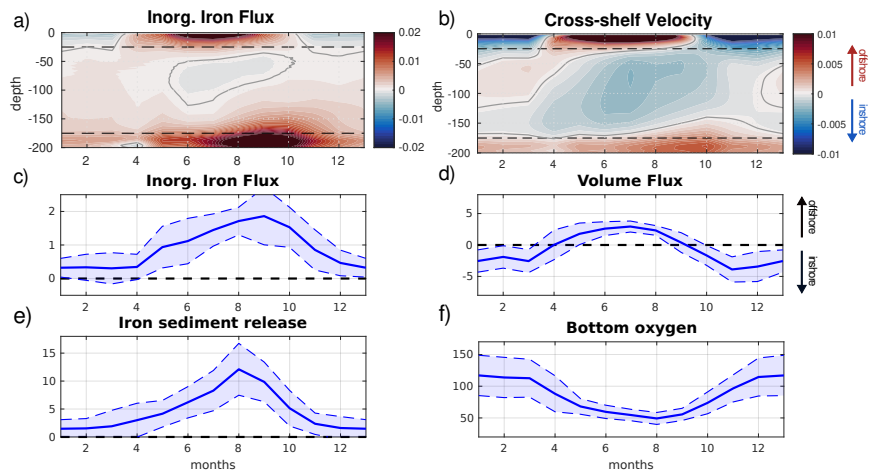
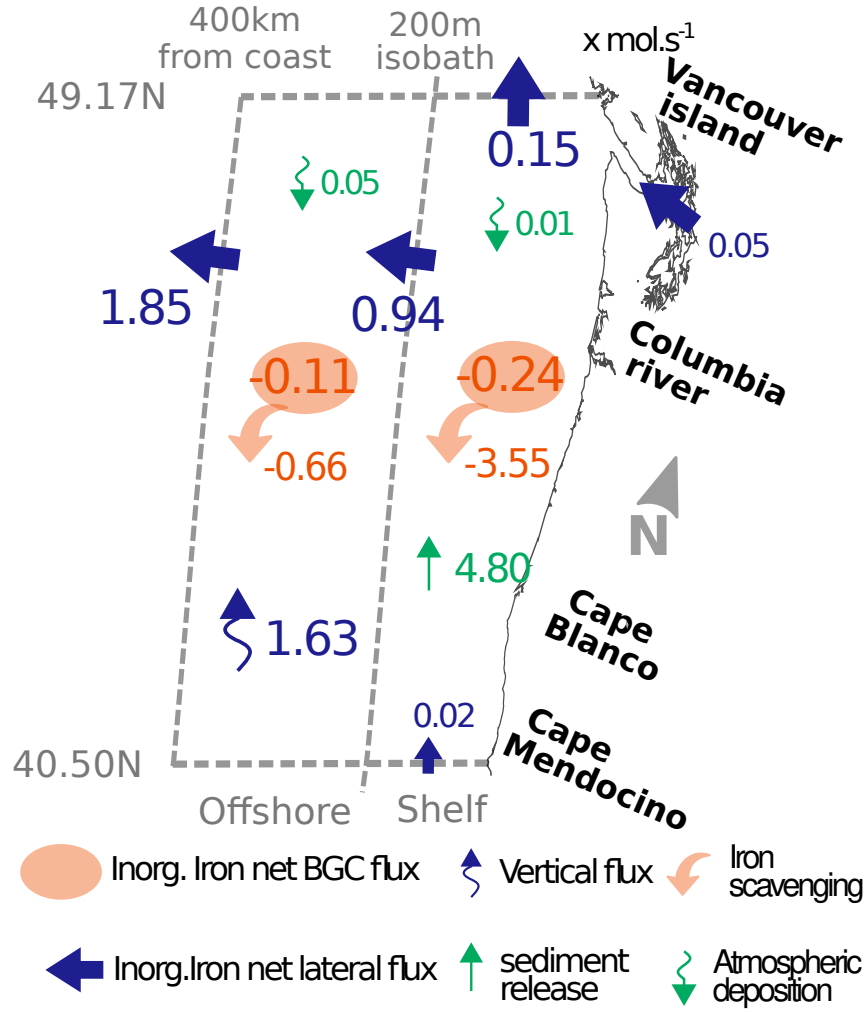
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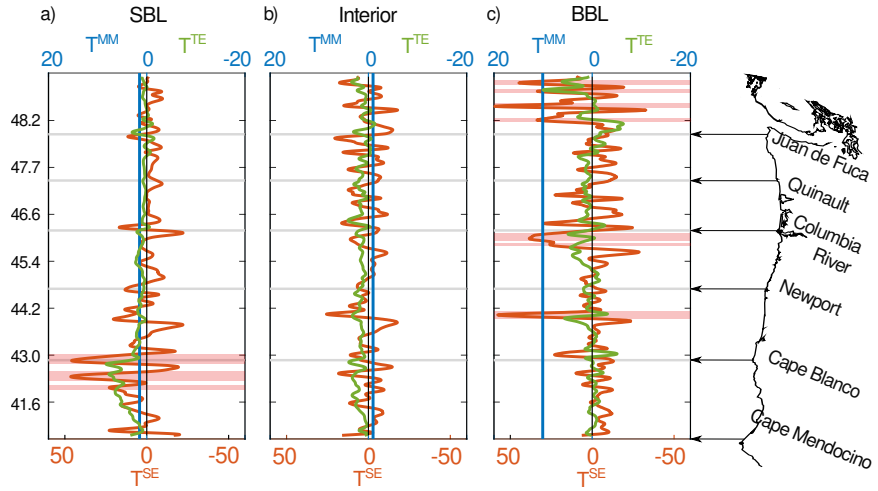
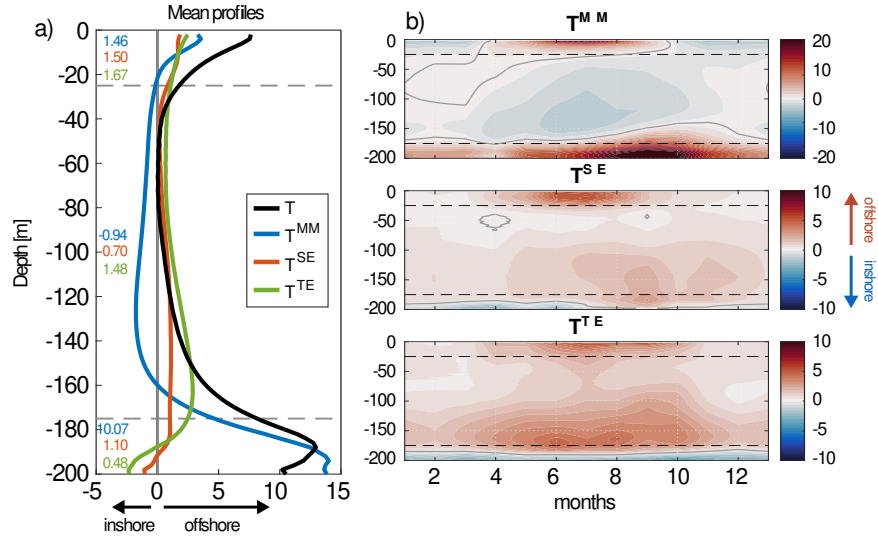
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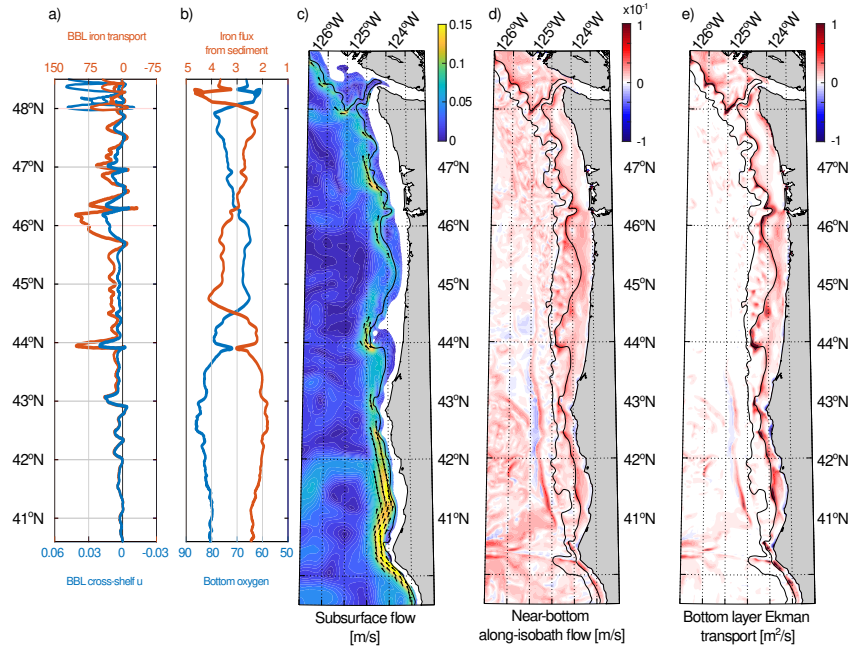
Abstract

Release of iron (Fe) from continental shelves is a major source of this limiting nutrient for phytoplankton in the open ocean, including productive Eastern Boundary Upwelling Systems. The mechanisms governing the transport and fate of Fe along continental margins remain poorly understood, reflecting interaction of physical and biogeochemical processes that are crudely represented by global ocean biogeochemical models. Here, we use a submesoscale-permitting physical-biogeochemical model to investigate processes governing the delivery of shelf-derived Fe to the open ocean along the northern U.S. West Coast. We find that a significant fraction (20%) of the Fe released by sediments on the shelf is transported offshore, fertilizing the broader Northeast Pacific Ocean. This transport is governed by two main pathways that reflect interaction between the wind-driven ocean circulation and Fe release by low-oxygen sediments: the first in the surface boundary layer during upwelling events; the second in the bottom boundary layer, associated with pervasive interactions of the poleward California Undercurrent with bottom topography. In the water column interior, transient and standing eddies strengthen offshore transport, counteracting the onshore pull of the mean upwelling circulation. Several hot-spots of intense Fe delivery to the open ocean are maintained by standing meanders in the mean current and enhanced by transient eddies and seasonal oxygen depletion. Our results highlight the importance of fine-scale dynamics for the transport of Fe and shelf-derived elements from continental margins to the open ocean, and the need to improve representation of these processes in biogeochemical models used for climate studies.









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2 **U.S West Coast to the Pacific Ocean**

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9 **Key Points:**

- 10 • Shelf-to-basin transport is a major source of iron to the North East Pacific Ocean.
11 • Half of this transport occurs year-round by downslope Ekman flow in the bottom
12 boundary layer, sustained by the poleward California Undercurrent.
13 • Regional hot-spots of iron transport are sustained by meanders in the poleward
14 undercurrent and local bottom oxygen depletion

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Abstract

Release of iron (Fe) from continental shelves is a major source of this limiting nutrient for phytoplankton in the open ocean, including productive Eastern Boundary Upwelling Systems. The mechanisms governing the transport and fate of Fe along continental margins remain poorly understood, reflecting interaction of physical and biogeochemical processes that are crudely represented by global ocean biogeochemical models. Here, we use a submesoscale-permitting physical-biogeochemical model to investigate processes governing the delivery of shelf-derived Fe to the open ocean along the northern U.S. West Coast. We find that a significant fraction ($\sim 20\%$) of the Fe released by sediments on the shelf is transported offshore, fertilizing the broader Northeast Pacific Ocean. This transport is governed by two main pathways that reflect interaction between the wind-driven ocean circulation and Fe release by low-oxygen sediments: the first in the surface boundary layer during upwelling events; the second in the bottom boundary layer, associated with pervasive interactions of the poleward California Undercurrent with bottom topography. In the water column interior, transient and standing eddies strengthen offshore transport, counteracting the onshore pull of the mean upwelling circulation. Several hot-spots of intense Fe delivery to the open ocean are maintained by standing meanders in the mean current and enhanced by transient eddies and seasonal oxygen depletion. Our results highlight the importance of fine-scale dynamics for the transport of Fe and shelf-derived elements from continental margins to the open ocean, and the need to improve representation of these processes in biogeochemical models used for climate studies.

Plain Language Summary

Iron is an essential nutrient that supports the life of marine organisms. In the ocean, large quantities of iron are released by sediments found along the continents. However, this iron is not very soluble, and it tends to precipitate back to the sediments close to where it is released. In fact, we still struggle to understand how enough of this iron makes its way to the open ocean, where it fertilizes phytoplankton and sustains fisheries. In this study, we use a sophisticated computer simulation of the ocean currents and chemistry of the Northern U.S. West Coast to study the transport of iron released along the continent to the open ocean. This computer simulation is able to reproduce the currents observed along the coast with high realism, including small swirls, eddies, and meanders that constantly mix coastal waters with open ocean waters. We found that about one-fifth of all the iron released by sediments along the coast is transported offshore, where it can sustain the life of marine organisms. This transport from the coast to the open ocean mostly takes place near the surface, reflecting upwelling caused by winds in the summer, and near the bottom, reflecting transport caused by the friction of the poleward California Undercurrent with the seafloor. We also found that episodic swirls, eddies, and meanders reinforce this delivery of iron to the open ocean, in particular along several “hot-spots” of intense transport along the coast. Our results suggest that global computer simulations used to study how marine ecosystems respond to climate change should improve how they represent small-scale currents and their effects on the cycle of iron, in particular along the ocean’s coastlines.

1 Introduction

The micronutrient iron (Fe) limits primary production in about half of the world’s oceans, regulating past and future changes in marine ecosystems and the global carbon cycle (C. M. Moore et al., 2013; Boyd & Ellwood, 2010; Tagliabue et al., 2017). The mechanisms controlling the oceanic Fe cycle have been studied extensively. However, because of complex interactions between external sources and internal cycling processes that are still poorly constrained, our ability to model the Fe cycle remains limited, hindering projections of future oceanic productivity and ecosystem change (Tagliabue et al., 2016, 2017).

65 Major sources of Fe in the ocean include atmospheric deposition (Duce & Tindale,
66 1991; Jickells et al., 2005), benthic release from continental shelves (Elrod et al., 2004;
67 Johnson et al., 1999), hydrothermal vents (Fitzsimmons et al., 2014; Resing et al., 2015;
68 Tagliabue et al., 2010), sea-ice melting (Person et al., 2021), and river runoff (Vieira et
69 al., 2020; Krachler & Krachler, 2021). Once in the ocean, Fe exists mostly as the poorly-
70 soluble ferric Fe, which quickly precipitates or is rapidly scavenged onto marine parti-
71 cles (X. Liu & Millero, 2002; Tagliabue et al., 2019; Honeyman et al., 1988). Fe can be
72 protected from removal by forming complexes with organic ligands produced by a myri-
73 ad of biological processes (Moffett & Boiteau, 2024; van den Berg, 1995; Buck et al., 2010).
74 Fe sources, sinks, and protection mechanisms are not fully understood and constrained
75 by *in situ* measurements and ocean biogeochemical models (Tagliabue et al., 2016, 2017).

76 Among major sources of Fe to the ocean, the release from continental shelves and
77 slope sediments has received particular attention (Dinniman et al., 2020; St-Laurent et
78 al., 2019, 2017; Elrod et al., 2004; Johnson et al., 1999), especially in Eastern Bound-
79 ary Upwelling System (EBUS), where Fe often limits primary production (Billier et al.,
80 2013; Messié & Chavez, 2015; Till et al., 2019; Hogle et al., 2018) despite significant ben-
81 thic fluxes (Elrod et al., 2004; Severmann et al., 2010; Dale et al., 2015). In EBUS, low
82 concentrations of dissolved oxygen (O_2) in bottom waters and high rates of organic mat-
83 ter deposition enhance Fe release from the sediments (Severmann et al., 2010; Conway
84 & John, 2014; Dale et al., 2015; Lam et al., 2020). Benthic Fe sources can exceed Fe deli-
85 very by atmospheric deposition and rivers (Deutsch et al., 2021; Robinson et al., 2022),
86 supporting the high primary production observed in EBUS (Johnson et al., 1999; Elrod
87 et al., 2004; Messié & Chavez, 2015). Benthic Fe that escapes phytoplankton uptake and
88 scavenging on the shelf can be transported to the open ocean by subsurface currents (Siedlecki
89 et al., 2012). In the presence of upwelling and vertical mixing, this Fe can fertilize phy-
90 toplankton in remote ocean regions, linking the cycles of carbon, O_2 , and Fe (Johnson
91 et al., 1999; Rapp et al., 2020; Wallmann et al., 2022).

92 Despite intense release from EBUS margins, it is unclear how much Fe reaches the
93 open ocean (Scholz et al., 2016; Lam et al., 2020). Global modeling studies suggest that
94 most benthic Fe is quickly removed by scavenging close to where it is released, restrict-
95 ing its impact to coastal waters (J. K. Moore & Braucher, 2008; J. K. Moore et al., 2004).
96 In contrast, *in situ* measurements and satellite observations indicate that Fe originat-
97 ing from continental margins can influence primary production hundreds to thousands
98 of kilometers offshore (Elrod et al., 2004; Lam et al., 2006; J. K. Moore & Abbott, 2000).

99 The export of sediment-derived Fe to the open ocean, here referred to as the shelf-
100 to-basin transport, occurs by a variety of physical processes, including upwelling and ver-
101 tical mixing, Ekman transport in the surface and bottom boundary layers (SBL and BBL
102 respectively), and eddies (Fiechter & Moore, 2012; Keith Johnson et al., 2005; Lam et
103 al., 2020; Tagliabue, Sallee, et al., 2014; Conway et al., 2018). These circulation patterns
104 are not adequately represented in global ocean Fe models (Tagliabue et al., 2016) because
105 of their coarse resolution, which limits their ability to capture fine scale currents in the
106 regions where most benthic Fe release and transport occur. These include shallow and
107 narrow shelves, where submesoscale processes dominate (McWilliams, 2016; Dauhajre
108 et al., 2017), but also the transition between continental margins and the open-ocean.
109 In this region, transport reflects mesoscale processes such as meanders in the prevailing
110 currents, eddies, and zonal jets (Conway et al., 2018; Uchida et al., 2020; Cravatte et al.,
111 2017), and finer-scale currents, such as submesoscale coherent vortices (McWilliams, 1985).
112 These subsurface coherent eddies are pervasive along EBUS (McCoy et al., 2020), includ-
113 ing the California Current System (Pelland et al., 2013), where they form by instabil-
114 ity of the poleward Undercurrent (Molemaker et al., 2015) and cause offshore transport
115 of biogeochemical tracers (Frenger et al., 2018). By trapping BBL waters, subsurface co-
116 herent eddies likely provide an efficient Fe delivery pathway from continental margins
117 to the ocean interior. Given the importance of mesoscale and submesoscale currents, the

118 role of the shelf-to-basin transport is likely underestimated in current global biogeochem-
119 ical models.

120 The seminal study by Siedlecki et al. (2012) provided an expanded picture of the
121 processes that ultimately deliver Fe to the open ocean in EBUS. By using an idealized
122 numerical model of the California Current, they showed that, on average, Fe accumu-
123 lating in the BBL during upwelling periods is transported to the ocean interior when wind
124 reverses to downwelling-favorable conditions. During wind reversals, thickening of BBL
125 waters, favored by intense vertical shear, generates a Fe-rich plume that detaches from
126 the upper continental slope and delivers Fe to the open ocean via isopycnal transport.
127 The study indicated that, depending on the frequency of wind reversals, between 10 and
128 50% of sediment-derived Fe could be delivered offshore, with the remaining part primar-
129 ily consumed on the shelf by biological activity. A stronger and shallower poleward Un-
130 dercurrent would further increase the fraction of Fe exported off the shelf. However, the
131 study was idealized in nature, and not designed to reproduce realistic Fe concentrations
132 and their long-term distribution under competing physical-biogeochemical dynamics.

133 Here, we build on this work to examine the fate of benthic Fe in a realistic, sub-
134 mesoscale permitting simulation of the U.S. West Coast circulation and biogeochemistry
135 (Kessouri et al., 2020; Damien et al., 2023). We use the model to elucidate the seasonal
136 balances behind the shelf-to-basin Fe transport and its variability along the Northeast
137 Pacific continental margin, with particular focus on interactions between sedimentary
138 Fe release and bottom O₂ and the frictional dynamics of the California Undercurrent un-
139 der realistic topography, wind, and fine-scale circulation patterns. While we focus on the
140 well-studied California Current system, our findings shed light on dynamics that are likely
141 to occur in other EBUS and, more broadly, along continental shelves and slopes, with
142 the potential to inform global models of Fe and other trace elements (Conway et al., 2018;
143 Jersild et al., 2021; St-Laurent et al., 2017).

144 The rest of the paper is organized as follows. In Section 2, we describe the model
145 configuration, Fe budget analysis, and framework to identify Fe transport processes and
146 contribution of eddies. In Section 3, we evaluate the model against observations and dis-
147 cuss the Fe balance and transport rates along the northern U.S. West Coast. In Section
148 4, we discuss the implications of these results for the global Fe cycle and conclude the
149 paper.

150 2 Methods

151 2.1 The Regional Oceanic Modeling System (ROMS)

152 We use a well-established physical-biogeochemical model of the California Current
153 system (Deutsch et al., 2021; Renault et al., 2021; Damien et al., 2023). The physical
154 component consists of the Regional Ocean Modeling System, ROMS (Shchepetkin & McWilliams,
155 2005; Shchepetkin, 2015), a primitive-equation, hydrostatic, topography-following (i.e.,
156 σ coordinate) ocean model. We focus on a Northern U.S. West Coast configuration ex-
157 tending from 36.8°N to 49.8°N and 112.5°W to 144.7°W. The grid is composed of 630
158 \times 1,340 cells with a horizontal resolution of approximately 1 km and 60 terrain-following
159 vertical levels (Damien et al., 2023). This horizontal grid spacing resolves the narrow to-
160 pography of the shelf in this region, which rarely exceeds a width of 20 km. The model
161 grid is stretched in the vertical to provide higher resolution near the surface and the seafloor,
162 and better capture boundary layer dynamics. On the shelf (depth shallower than 200
163 m), the vertical grid includes at least 13 grid points in the upper 25 m of the water col-
164 umn and 6 in the deeper 25 m. The high resolution is essential to resolve mesoscale and
165 submesoscale variability (Capet et al., 2008; Kessouri et al., 2020), the intensified wind-
166 driven overturning circulation on the continental shelf (Damien et al., 2023), and the re-
167 sulting cross-shore exchange of nutrients and Fe. The model is run for the 1997-2017 pe-

168 riod, using a time step of 150 s, and initial and boundary conditions from a “parent” con-
 169 figuration at 4 km resolution (Deutsch et al., 2021). The numerical setup does not in-
 170 clude tidal forcings. Model output is saved as daily averages for tracers and monthly av-
 171 erages for physical and biogeochemical fluxes. A detailed discussion of the model con-
 172 figuration, setup, forcings, and initialization is presented in previous publications (Deutsch
 173 et al., 2021; Renault et al., 2021; Kessouri et al., 2020; Damien et al., 2023).

174 **2.2 The Biogeochemical Elemental Cycling (BEC) ecosystem model**

175 ROMS is coupled online to the Biogeochemical Elemental Cycling model, BEC (J. K. Moore
 176 & Braucher, 2008; J. K. Moore et al., 2004). BEC represents the cycles of nutrients (in-
 177 cluding Fe), O₂, carbon, alkalinity, and organic matter as driven by three phytoplank-
 178 ton and one zooplankton functional groups (J. K. Moore & Braucher, 2008; J. K. Moore
 179 et al., 2004; Deutsch et al., 2021). Remineralization of sinking organic matter is param-
 180 eterized according to the mineral ballast model of Armstrong et al. (2001). Air-sea fluxes
 181 of O₂ and carbon dioxide are based on the gas-exchange model of Wanninkhof (1992).

182 The Fe cycle in BEC includes four separate pools: dissolved inorganic Fe (dFe), Fe
 183 scavenged onto sinking particles, including lithogenic minerals from dust deposition, and
 184 Fe associated with organic matter pools. Only the dFe and organically bound Fe pools
 185 are modeled as state variables (J. K. Moore & Braucher, 2008; J. K. Moore et al., 2004).
 186 For the dFe pool, four processes are considered: atmospheric deposition, biological up-
 187 take and remineralization, scavenging by sinking particles, and release by sediments. At-
 188 mospheric Fe deposition is based on the dust climatology of Mahowald et al. (2005). The
 189 model does not include river sources for Fe, except through the Juan De Fuca Strait. While
 190 river inputs are important in the northern California Current (Wetz et al., 2006; Chase
 191 et al., 2007), prior work suggests that benthic sources are likely dominant in this region
 192 (Deutsch et al., 2021; Severmann et al., 2010).

Release of Fe from the sediment follows the O₂-dependent parameterization by Deutsch
 et al. (2021), obtained by a fit to a compilation of benthic flux chamber measurements
 in the California margin (Severmann et al., 2010):

$$\log_{10}\Phi(Fe) = 2.5 - 0.0165 \cdot O_2, \quad (1)$$

193 where O₂ is in mmol m⁻³ and the benthic Fe flux $\Phi(Fe)$ in $\mu\text{mol m}^{-2} \text{d}^{-1}$ (Deutsch
 194 et al., 2021). We find this parameterization adequate to capture the observed range of
 195 benthic Fe fluxes and their seasonal and alongshore variability (Severmann et al., 2010).

196 Additional details on the BEC configuration, forcings, coupling with ROMS, and
 197 validation of the model solution are presented in Deutsch et al. (2021). A discussion of
 198 the wind-driven shelf circulation and biogeochemistry is presented in Damien et al. (2023).
 199 In the following, we further evaluate the solution against available measurements of dis-
 200 solved Fe concentrations along the Northern U.S. West Coast (Section 2.3)

201 **2.3 Dissolved Fe measurements along the Northern U.S. West Coast**

202 We compiled a dataset of dFe measurements from Baja California to Vancouver
 203 Island from 23 published studies between 1987 and 2022, resulting in approximately 4,068
 204 individual dFe measurements (Johnson et al., 2003; King & Barbeau, 2011; Hogle et al.,
 205 2018; Bundy et al., 2016). A list of the studies and references is presented in Supplemen-
 206 tary Table 1. These include a global compilation (Tagliabue et al., 2016), cruises from
 207 regional observational programs (Hogle et al., 2018; King & Barbeau, 2007; Johnson et
 208 al., 2003; Bundy et al., 2016), and independent studies (Biller & Bruland, 2013; Boiteau
 209 et al., 2019; Chappell et al., 2019; Chase et al., 2002, 2005; Till et al., 2019; Firme et al.,
 210 2003; John et al., 2012; Hawco et al., 2021; Bundy et al., 2014, 2015).

For this compilation, we define dFe as the sum of both the truly dissolved Fe and the dissolvable Fe, based on the definitions used in the original references. Different studies used varying filter sizes to characterize the dissolved and dissolvable Fe pools, with 0.2 μm , 0.4 μm , and 0.45 μm as the most common. Measurement methods also vary between studies. In general, measurements were taken with bottles, pump system, and/or by surface tows. In some cases, samples were acidified for short periods of time before analysis. Despite variability in sampling approaches, we find a good agreement between different sets of observations and consider the merged dataset as representative of the dFe distribution in the California Current. The final compilation includes observations from 1980 to 2021, with most of the data between 1997 and 2015 and in the upper 100 m of the water column.

2.4 Model Analysis and Diagnostics

2.4.1 Dissolved Fe Budget Analysis

To investigate the fate of the shelf-derived Fe, we analyze the balance between transport fluxes and biogeochemical sources and sinks based on the model conservation equation for dFe, i.e., the model dFe budget. We use output from a 9-year period from 2008-2016, spanning the region from Cape Mendocino to Vancouver Island. We focus on this region because it encompasses the Washington and Oregon continental shelves, where strong bottom water hypoxia occurs, leading to high benthic Fe release fluxes from low- O_2 sediments. The dFe budget includes external sources from atmospheric deposition and benthic inputs, biological cycling processes consisting of uptake by phytoplankton and remineralization and recycling, scavenging and burial in the sediment, and physical transport by currents and vertical mixing. We calculate biogeochemical rates and physical fluxes online in the model at each time step and average them in time to construct monthly climatologies. The budget of all tracers is closed to numerical precision and scaled up from individual grid cells to three-dimensional regions by numerical integration.

2.4.2 Eddy Decomposition of the Shelf-to-Basin Fe Transport

We define the shelf-to-basin transport as the advective flux of dissolved Fe across a vertical surface that intersects the 200 m depth isobath, here taken as the boundary between the shelf and the open ocean (Laruelle et al., 2013), using the approach by Damien et al. (2023), which naturally incorporates variations in the topography of the western U.S. coastline. Accordingly, Fe transport is calculated as:

$$T = u \cdot Fe, \quad (2)$$

where Fe is the dFe concentration, and u the cross-shelf horizontal current.

To elucidate the mechanisms responsible for the delivery of Fe from the shelf to the open ocean, we decompose the cross-shelf Fe transport T into contributions from the large-scale mean cross-shelf current, standing eddies and meanders, and transient eddies. The large-scale mean transport includes seasonal upwelling and Ekman transport in the SBL and BBL that occur broadly along the coast. Standing eddies and meanders represent circulation patterns that are relatively persistent in time but cause meridional variability along the coast and support localized hot-spots of shelf-to-basin exchange. In contrast, transient eddies represent the contribution of time-varying flow driven by high-frequency fluctuations in winds and instabilities in the mesoscale and submesoscale circulation.

We quantify these components by following the transport decomposition by Lee and Coward (2003). This approach involves two low-pass filters, one in time and one in space. Accordingly, an arbitrary model variable X can be decomposed into time mean and fluctuating components as:

$$X = \bar{X} + X', \quad (3)$$

257 where $\bar{X} = \frac{1}{\tau} \int_{t-\frac{\tau}{2}}^{t+\frac{\tau}{2}} X dt$, with $\tau = 1$ month. Here, \bar{X} is the monthly climatological mean
 258 of the variable and $X' = X - \bar{X}$ is the transient eddy fluctuations on time scales shorter
 259 than a month.

260 Similarly, X can be decomposed into an alongshore mean plus a spatial variation
 261 around this mean as:

$$X = [X] + X^*, \quad (4)$$

262 where $[X] = \frac{1}{L} \int_0^L X dL$, with L is the length of the shelf along the U.S. West Coast.
 263 Here, $[X]$ is the alongshore mean of the variable along the 200 m isobath and $X^* = X -$
 264 $[X]$ is the alongshore fluctuations.

265 By applying the spatial and temporal filters to the cross-shelf Fe flux, $T = u \cdot Fe$,
 266 the mean Fe transport across the shelf break can be expressed as:

$$[\overline{T}] = \overline{[u Fe]} = [\bar{u}][\overline{Fe}] + [\bar{u}^* \overline{Fe}^*] + [\overline{u'}][\overline{Fe'}] + [\overline{u'^* Fe'^*}], \quad (5)$$

267 where the full derivation is shown in Appendix A. We further set:

$$[\overline{T^{MM}}] = [\bar{u}][\overline{Fe}], \quad (6)$$

$$[\overline{T^{SE}}] = [\bar{u}^* \overline{Fe}^*], \quad (7)$$

$$[\overline{T^{TE}}] = [\overline{u'}][\overline{Fe'}] + [\overline{u'^* Fe'^*}] = [\overline{u' Fe'}]. \quad (8)$$

268 Here, $[\overline{T^{MM}}]$ is the cross-shelf Fe transport driven by the climatological mean cir-
 269 culation and Fe distribution, $[\overline{T^{SE}}]$ is the cross-shelf transport driven by meanders and
 270 standing eddies that cross the shelf break, and $[\overline{T^{TE}}]$ is the cross-shelf transport driven
 271 by transient eddies. Note that all quantities in Equations (5)–(8) represent temporal and
 272 alongshore averages. Furthermore, to facilitate the interpretation of transient terms, we
 273 combine the contribution of time-dependent, coast-wide fluctuations ($[\overline{u'}][\overline{Fe'}]$) with tran-
 274 sient fluctuations that also vary in the alongshore direction ($[\overline{u'^* Fe'^*}]$). As such, both
 275 $[\overline{T^{SE}}]$ and $[\overline{T^{TE}}]$ arise from correlations between u and Fe , solely in space for the for-
 276 mer and in time for the latter. To simplify the notation in the rest of the paper, we drop
 277 the averaging operators (i.e., the square brackets and over-bars) and refer to the aver-
 278 aged transport terms in Equations (6)–(8) as T^{MM} (mean transport), T^{SE} (standing
 279 eddies), and T^{TE} (transient eddies), respectively, unless specified otherwise. We estimate
 280 these quantities from model output, following the approach detailed in Appendix A.

281 3 Results

282 3.1 Evaluation of the model, with a focus on dissolved Fe

283 An extensive validation of the model against *in-situ* and remote sensing data, based
 284 on the “parent” 4-km U.S. West Coast simulation, is presented in Deutsch et al. (2021)

Figure 1. Observed and modeled dissolved iron (dFe, nM) in the upper 100 m depth. (a) Observations from a U.S. West Coast Fe data compilation (see Section 2.3). (b) Annual mean average dFe concentrations from the model, sampled at the same location as for observations. (c) Observed vs. modeled dissolved Fe concentrations at the ocean surface as a function of the distance from the coast line. Model output is sampled at the same locations and months as for observations. Error bars around the blue-dot points show the standard deviation of the model results.

285 and Renault et al. (2021), showing good agreement for both hydrographic and dynam-
 286 ical properties. A further assessment of the characteristics and dynamics of the Califor-
 287 nia Undercurrent is presented in Chen et al. (2021). The 1-km resolution simulations dis-
 288 cussed here are broadly similar to the “parent” simulation; in the Supplementary Infor-
 289 mation we provide an additional validation for physical and biogeochemical properties
 290 that directly influence the shelf-to-basin Fe transport: dissolved O₂ patterns (Figure S2
 291 and S3) and the California Undercurrent (Figure S4).

292 Overall, the model spatial patterns and seasonal variability for dissolved O₂ con-
 293 centrations fall within the range of the observations (Figure S2 and Figure S3). The Un-
 294 dercurrent flows faster in the southern part of the domain (in particular, south of Hec-
 295 eta Banks, Supplementary Figure S4) and flows nearly continuously along the coast, as
 296 suggested by observations (Pierce et al., 2000). We also note a significant seasonal vari-
 297 ability, with shoaling and strengthening of the Undercurrent through the Fall and Win-
 298 ter (Thomson & Krassovski, 2010) (Supplementary Figure S5). In the following, we fur-
 299 ther provide an in-depth evaluation of the model dFe against the new U.S. West Coast
 300 dFe compilation.

301 In general, the model captures the magnitude and spatial variability of observed
 302 Fe (Figure 1). Along the coast, both model and observations show high dFe concentra-
 303 tions in the upper ocean, reflecting intense benthic release from the shallow bathymetry.
 304 In the open ocean, dFe concentrations are low in both model and observations, reflect-
 305 ing limited external Fe inputs, sustained biological uptake, and scavenging. Away from
 306 the coast, model dFe concentrations are low at the surface (~ 0.3 nM) because of bio-
 307 logical uptake, and gradually increase in subsurface waters following remineralization ($>$
 308 1nM).

309 Compared to observations, the model underestimates the sharp gradient and vari-
 310 ability in dFe between coastal and open ocean waters (Figure 1c). Along the shelf, ob-
 311 served dFe values vary between 0.03 and 30 nM, while model dFe values show a some-
 312 what reduced range between 0.2 and 10 nM. In the open ocean, the model tends to over-
 313 estimate dFe concentrations, while it struggles to capture extremely low dFe in shallow
 314 layers, producing a more uniform dFe distribution than observed. This bias is likely caused
 315 by the model Fe scavenging scheme, which represents Fe protection by organic ligands
 316 with a uniform ligand concentration of around 0.6 nM, a common assumption for ocean
 317 biogeochemical models (Tagliabue et al., 2016).

318 While low compared to typical values for macro-nutrients, the correlation coeffi-
 319 cient between modeled and observed dFe ($R = 0.25$) is in the upper range of global ocean
 320 biogeochemical models (Tagliabue et al., 2016), which can show correlation coefficients
 321 as low as 0.10 or even negative, reflecting limited ability to capture the complexity of
 322 Fe cycling. Comparing to global models, our high-resolution regional model produces more
 323 realistic cross-shore gradient in dFe, despite somewhat underestimating the range of ob-
 324 servations (Figure 1c). Multiple reasons are likely behind this bias — including simpli-
 325 fications in the Fe protection and scavenging schemes. We also note that comparison of

Figure 2. Schematic of the balance between physical transport and biogeochemical processes for dissolved Fe (dFe) along the Northern U.S. West Coast. The balance reflects an average over the 9-year period from 2008-2016. Biogeochemical sources and sinks and transport rates of dissolved Fe are integrated over the shelf region (defined by the 200 m isobath) and the open-ocean region (up to a distance of 400 km from the coast, between the surface and 200 m depth). Green arrows show external sources of dFe (atmospheric deposition and benthic release); blue arrows show horizontal and vertical physical transport; and red circles show net biogeochemical transformation rates, i.e., the sum of uptake by phytoplankton, and release by remineralization. Removal by scavenging, which ultimately transports Fe to the sediment as sinking particles, is indicated by a red arrow. Units are mol s^{-1} . The integrated cross-shore transport of dFe is 0.94 mol s^{-1} .

326 climatological model values vs. instantaneous measurements is likely to underestimate
 327 the range captured by observations, which are taken at different times of the year, and
 328 are affected by eddies and other sources of variability.

3.2 Fate of sediment-derived Fe

329
 330 Analysis of the model Fe budget shows that benthic release is the dominant source
 331 of Fe along the Northern U.S. West Coast (Figure S1). On the shelf, 97% of external Fe
 332 inputs is from the sediment. Benthic release is then the dominant source of dFe (4.80
 333 mol s^{-1} vs. 0.01 mol s^{-1} from atmospheric deposition), the majority of which (3.79 mol
 334 s^{-1}) is retained on the shelf by a combination of biological uptake, scavenging, and trans-
 335 port to the sediment by sinking particles. Scavenging onto particles removes the major-
 336 ity of Fe inputs (74%), consistently with global model results (J. K. Moore & Braucher,
 337 2008) (Figure S1).

338 Our results update the picture presented in Siedlecki et al. (2012), which suggested
 339 a dominant role for biological uptake in consuming benthic-derived dFe ($\sim 80\%$), and lim-
 340 ited scavenging ($\sim 0.05\%$) on the shelf. These differences reflect the increased realism and
 341 complexity of Fe biogeochemistry — in particular scavenging — in our model (J. K. Moore
 342 et al., 2004).

343 Analysis of the dFe balance over the shelf and open ocean (up to a distance of 400
 344 km from the coast) (Figure 2) shows that the shelf-to-basin transport is one of the two
 345 major inputs of dFe to the open ocean (0.94 mol s^{-1}), second only to vertical supply (1.63
 346 mol s^{-1}), despite the narrow surface area of the shelf compared to the open-ocean. The
 347 cross-shelf flux represents 36% of the total dFe input to the surface open ocean, far ex-
 348 ceeding dust deposition (0.05 mol s^{-1}). It is also larger than the net biological uptake
 349 offshore (0.77 mol s^{-1}), suggesting that the continental shelf is a source of Fe to the broader
 350 Northeast Pacific.

3.3 The shelf-to-basin transport of Fe

3.3.1 Structure and drivers of the shelf-to-basin Fe transport

351
 352
 353 Two pathways dominate Fe transport from the shelf to the open ocean: the first
 354 in the SBL and the second in the BBL (Figure 3a). These pathways are the result of the
 355 wind-driven overturning circulation that develops on the continental shelf (Figure 3b)
 356 (Damien et al., 2023). The patterns and seasonal variability of this cross-shelf overturn-
 357 ing show a complex three-dimensional structure (Figure S5) that is influenced both by
 358 variations in surface winds and by interactions between the poleward California Under-
 359 current (Figure S4) and bottom topography.

Figure 3. Seasonal cycle of the processes governing the transport of dissolved Fe from the shelf to the open ocean. (a) Climatological cross-shore Fe transport ($10^{-6} \text{ mol m}^{-2}\text{s}^{-1}$) as a function of depth and time, averaged along the 200 m isobath. This transport flux is calculated as $[\overline{T}] = [\overline{uFe}]$, following the approach discussed in Section 2.4.2. Positive values (red colors) show transport directed offshore; negative values (blue colors) show transport directed inshore. (b) Same as (a) but for the cross-shore current velocity (m s^{-1}). (c) Climatological cross-shore Fe transport (mol s^{-1}) integrated by depth and distance along the 200 m isobath. (d) Climatological cross-shore volume transport ($10^5 \text{ m}^3\text{s}^{-1}$). (e) Climatological benthic Fe flux (mol s^{-1}) integrated over the continental shelf. (f) Bottom O_2 concentration (mmol m^{-3}) averaged over the continental shelf. Positive values in (c) and (d) show transport directed offshore.

360 In the SBL, the seasonal cross-shelf overturning consists of offshore transport dur-
 361 ing upwelling and inshore transport during downwelling and relaxation events, balanced
 362 by currents with opposite directions in the water column interior (Figure S5). In con-
 363 trast, the mean circulation is directed offshore year-round in the BBL. This can be at-
 364 tributed to the frictional dynamics of the California Undercurrent, which, despite a shoal-
 365 ing and intensification in the fall (Chen et al., 2021), continuously moves waters pole-
 366 ward through the year (Figures S4 and S5). Transport in the BBL emerges as a down-
 367 hill flow with both offshore and downward components, which peaks in magnitude over
 368 the outer shelf. Notably, the magnitude of the California Undercurrent and of the down-
 369 hill flow in the BBL are strongly correlated, especially at the location of the 200 m iso-
 370 bath, reflecting their dynamical connection.

371 Despite the good correlation between the cross-shore Fe flux (Figure 3a) and the
 372 cross-shelf currents (Figure 3b), notable differences between the water volume and Fe
 373 transport are observed, mainly in the water column interior outside the boundary lay-
 374 ers. While the SBL Fe transport pathway is active mostly during the upwelling season
 375 from spring to summer, reversing sign in winter, the BBL pathway is active throughout
 376 the year, although it is intensified in the wake of upwelling events (August to October,
 377 Figure 3a). Relative to the two boundary layers, the mean Fe transport in the water col-
 378 umn interior remains small, and mostly directed offshore, in contrast with the mean vol-
 379 ume transport (Figure 3b).

380 Our results also suggest that intense dFe export follows upwelling events, but the
 381 BBL pathway is active year-round on the continental shelf and upper slope, driven by
 382 the persistent downhill tilt of the bottom current (Chen et al., 2021; Damien et al., 2023).
 383 As discussed further in Section 3.3.3, this finding is also consistent with the results of
 384 Siedlecki et al. (2012), which suggest that when the wind relaxes and reverses after up-
 385 welling events, the Fe-laden BBL is mixed into the water column and transported off-
 386 shore along isopycnals. Integrated along the shelf break, using an average thickness for
 387 both boundary layers of 25 m, we find that 19% of the cross-shore Fe flux takes place
 388 in the SBL (0.18 mol s^{-1}) and 46% (0.43 mol s^{-1}) in the BBL.

389 Vigorous cross-shore Fe transport takes place during the entire upwelling season.
 390 However, the peak Fe transport is observed in September (Figure 3c), a two-month lag
 391 from the peak in the cross-shore volume transport (Figure 3d) and one-month lag from
 392 the peak in the benthic dFe release associated with a minimum in bottom O_2 concen-
 393 trations on the shelf (Figure 3e,f). More generally, the seasonal cycle of the shelf-to-basin
 394 Fe transport arises from the interaction between the volume transport flux (Figure 3b),
 395 benthic dFe release (Figure 3e), and bottom O_2 concentrations (Figure 3f). In partic-
 396 ular, bottom O_2 concentrations on the continental shelf dictate the magnitude of the ben-

Figure 4. Role of mean circulation and eddies for the shelf-to-basin Fe transport. (a) The total cross-shelf Fe flux (T , black) is decomposed into mean (T^{MM} , blue), standing eddy (T^{SE} , red), and transient eddy (T^{TE} , green) components. All fluxes have been averaged over time and along the 200 m isobath. Section 2.4.2, Equations 5-8 provide definitions for the different flux terms, with further details in Appendix A. Positive values indicate transport directed offshore, and negative values transport directed inshore. Units for all fluxes are 10^{-9} mol $m^{-2}s^{-1}$. The average value of the different components within the surface boundary layer (SBL, top), interior (center), and bottom boundary layer (BBL, bottom) are annotated on the figure. (b) Seasonal cycle of the different components of the cross-shore Fe flux averaged along the 200 m isobath (10^{-9} mol $m^{-2}s^{-1}$): T^{MM} (upper panel), T^{SE} (middle panel), and T^{TE} (lower panel).

Figure 5. Alongshore variability and hot-spots in the shelf-to-basin Fe transport. Vertically and time-averaged cross-shelf Fe transport (a) over the surface boundary layer (SBL, 0-25 m depth), (b) the water column interior (25-175 m depth), and (c) the bottom boundary layer (BBL, 175-200 m depth). The total cross-shelf Fe flux is decomposed into its monthly mean (T^{MM} , blue), standing eddy (T^{SE} , red), and transient eddy (T^{TE} , green) components. Note that, in this figure, to visualize the alongshore variability of the eddy components, the alongshore averaging filter is not applied. More precisely, $\overline{T^{MM}}$, $\overline{T^{SE}}$, and $\overline{T^{TE}}$ are shown. Units for all fluxes are 10^{-9} mol $m^{-2} s^{-1}$. Positive values indicate offshore transport; negative values inshore transport. The shaded red boxes highlight hot-spots of cross-shore Fe transport (here defined as total Fe flux larger than $30 \cdot 10^{-9}$ mol $m^{-2} s^{-1}$). In all panels, the thin vertical black line show zero fluxes.

397 thic dFe release, which in combination with the magnitude and timing of the volume trans-
 398 port flux, shape the seasonal cycle of the shelf-to-basin Fe transport.

399 **3.3.2 Contribution of transient and standing eddies to the variability** 400 **of the shelf-to-basin Fe transport**

401 Additional differences between the mean Fe and volume transports, in particular
 402 in the water column interior, can be attributed to the role of transient and standing ed-
 403 dies. A decomposition of the cross-shore Fe transport (Figure 3a) into different compo-
 404 nents reflecting respectively the large-scale mean circulation (T^{MM}), standing eddies and
 405 meanders (T^{SE}), and transient eddies (T^{TE}) reveals a substantial influence of fine-scale
 406 circulation on the exchange of Fe from the shelf to the open ocean (Figure 4a). As a
 407 whole, standing and transient eddies intensify the mean offshore Fe transport at the sur-
 408 face, and oppose it near the bottom. Notably, eddies also oppose the mean Fe transport
 409 in the water column interior, where their combined effect — always directed offshore —
 410 overcomes the net inshore transport by the mean circulation, resulting in a net Fe ex-
 411 port from the shelf to the open ocean.

412 In the SBL (0-25 m), the eddy terms T^{TE} and T^{SE} are positive, and, when com-
 413 bined, larger than T^{MM} . Each contributes to about one-third of the total SBL Fe trans-
 414 port (Figure 4a). In contrast, T^{MM} dominates the Fe transport in the BBL (approx-
 415 imately 175-200 m depth), where T^{SE} and T^{TE} are nearly one order of magnitude smaller.
 416 However, the direction of eddy transport terms changes between the upper part of the
 417 BBL, where they are directed offshore, and the lower part, where they are directed in-
 418 shore, and thus oppose T^{MM} . In the upper BBL, both T^{SE} and T^{TE} represent a sub-
 419 stantial contribution to the total Fe flux from the shelf to the open ocean ($\sim 50\%$). In
 420 contrast, in the lower BBL, they reduce the offshore Fe transport by $\sim 30\%$. Away from

421 the boundary layers, T^{SE} and T^{TE} oppose the T^{MM} , resulting in a weakly positive off-
 422 shore transport of Fe in the water column interior.

423 The mean and eddy components of the cross-shore Fe transport show significant
 424 temporal variability following the seasonal cycle of upwelling. This seasonal variability
 425 is more pronounced in the boundary layers, in particular for the mean component T^{MM}
 426 (Figure 4b), which shows reversals in direction similar to the volume transport (Figure 3b).
 427 In contrast, eddy components are generally weakly enhanced during upwelling (Figure 4c-
 428 d). The nearly year-round activity of the T^{MM} and T^{SE} in the BBL can be attributed
 429 to the downslope tilt of the California Undercurrent (Chen et al., 2021; Damien et al.,
 430 2023). In addition, the intensification of the offshore BBL flow from August to October
 431 coincides with the seasonal intensification of the California Undercurrent (Chen et al.,
 432 2021; Thomson & Krassovski, 2010).

433 A coast-wide analysis of the cross-shore Fe flux (Figure 5) reveals the presence of
 434 significant local hot-spots of cross-shelf Fe exchange, in particular near the surface and
 435 the bottom. These hot-spots, shown by local maxima in T^{SE} , are mainly driven by lo-
 436 cally intensified transport by standing eddies (red lines in Figure 5), reflecting the pres-
 437 ence of meanders or separations of the alongshore current that enhance exchange across
 438 the shelf break. To a lesser extent, we also observe hot-spots driven by locally intensi-
 439 fied transient eddies, although their magnitude and spatial variability are smaller than
 440 for standing eddies.

441 In the SBL, we highlight the region south of Cape Blanco as the main pathway of
 442 surface Fe export, with two large peaks near Cape Sebastian (42.35°N) and Cape Blanco
 443 (42.83°N). This standing-eddy export pathway, further intensified by transient eddies,
 444 is likely fueled by strong, localized upwelling along the Northern California coast (Damien
 445 et al., 2023). A secondary region of enhanced Fe export is located near the Columbia River
 446 estuary.

447 In the BBL, multiple hot-spots of Fe transport are linked to the coastal bathymetry,
 448 which in turns dictates the location of meanders in the mean bottom current. Three ma-
 449 jor hot-spots are located at the southern edge of the Heceta Banks (43.95°N), the Ne-
 450 halem Banks (45.95°N – 46.40°), and the Northern Juan de Fuca Canyon (48.23N – 48.50N).
 451 The cross-shore Fe flux by standing eddies there exceeds $30 \cdot 10^{-9} \text{ mol m}^{-2} \text{ s}^{-1}$. Two
 452 secondary hot-spots with smaller offshore transport ($\sim 20 \times 10^{-9} \text{ mol m}^{-2} \text{ s}^{-1}$) are lo-
 453 cated north of Cape Blanco (43°N) and Grays Canyon (47.04°N).

454 We note that hot-spots of eddy transport could be related to regions of formation
 455 of subsurface coherent eddies, which are commonly observed in this region (Pelland et
 456 al., 2013; McCoy et al., 2020). The model indeed shows a rich population of subsurface
 457 eddies just offshore of the continental slope, particularly during the fall season (Figure
 458 S5). Consistent with their formation by topographic interactions (Molemaker et al., 2015),
 459 these subsurface eddies trap Fe-rich BBL waters in their cores (Figure S6), subsequently
 460 transporting Fe offshore, thus contributing to Fe export from the continental shelf.

461 ***3.3.3 Contribution of the California Undercurrent to the shelf-to-basin*** 462 ***Fe transport***

463 The importance of standing eddies and meanders in driving hot-spots of Fe exchange
 464 across the shelf break suggests a major role for the circulation dynamics of the Califor-
 465 nia Undercurrent. Following upwelling, when both the offshore Fe transport and the un-
 466 dercurrent are seasonally the strongest (Figure 3a), we observe a high correlation ($R=0.93$)
 467 between the alongshore Fe and volume transports along the shelf (Figure 6a). This in-
 468 dicates that alongshore variability of the cross-shore Fe transport, on spatial scale of the
 469 order of 10 to 100 km, is primarily driven by mean current patterns.

Figure 6. Role of the California Undercurrent for cross-shelf Fe transport in the bottom boundary layer (BBL). (a) Time-averaged Fe transport (10^{-9} mol $m^{-2}s^{-1}$, solid red line) and cross-shelf velocity ($m s^{-1}$, solid blue line) in the BBL between August and October. (b) Time averaged Fe release from the sediment (10^{-10} mol $m^{-2}s^{-1}$, solid red line) and bottom O_2 concentrations ($mmol m^{-3}$, solid blue line) along the 200 m isobath between August and October. (c) Mean horizontal currents averaged in the 100-300 m layer between August and October. Colors show the current velocity; arrows the current direction, for speed above $0.08 m s^{-1}$. The black solid line shows the 200 m isobath. (d) Along-isobath component of the mean flow ($m s^{-1}$) near the bottom. Positive values (red colors) indicate a flow with the slope to its right. (e) Ekman transport ($m^2 s^{-1}$) in the BBL. Positive values (red colors) indicate downhill Ekman transport.. In (d) and (e), the black solid lines show the 200 m and 1000 m isobaths.

470 However, alongshore variability in dFe concentrations on the shelf is also impor-
 471 tant. The progressive increase of the cross-shore Fe flux relative to the volume flux north
 472 of Cape Blanco (Figure 6a) reflects a concomitant increase in dFe concentrations on the
 473 shelf. These in turn can be explained by larger benthic inputs (Figure 6b) to low- O_2 bot-
 474 tom waters. North of Cape Blanco, the ratio between the Fe and volume flux is much
 475 larger for the offshore transport than the inshore transport, confirming the Fe enrich-
 476 ment of shelf waters transported to the open ocean. Thus, recurring low- O_2 conditions
 477 on the shelf support high benthic fluxes that increase dFe concentration on the shelf and
 478 enhance offshore dFe transport. The influence of high benthic dFe release can be observed
 479 on relatively large spatial scales, on the order of 100 km, in particular between $43.5^\circ N$
 480 and $47.5^\circ N$. In contrast, downhill flow in the BBL shapes the shelf-to-basin dFe trans-
 481 port on spatial scales of tens of km.

482 These considerations apply to the three major hot-spots of Fe export located near
 483 the southern Heceta and Nehalem Banks, and the Northern Juan de Fuca Canyon. Here,
 484 vigorous dFe export is mainly driven by enhanced volume transport associated with stand-
 485 ing eddies and meanders, but with a significant contribution from locally intense release
 486 of Fe from low- O_2 sediment. In contrast, at two secondary hot-spots north of Cape Blanco
 487 and north of Grays Canyon, benthic Fe inputs are not particularly enhanced relative to
 488 the surrounding waters.

489 On the outer shelf, enhanced volume fluxes correspond to local intensification of
 490 the California Undercurrent (Figure 6c). Episodes of detachment of the California Un-
 491 dercurrent from the continental shelf facilitate bottom Fe export in the southern part
 492 of the Heceta Banks and the Northern Juan de Fuca Canyon. In contrast, the hot-spot
 493 of Fe transport at the Nehalem Banks does not correspond to a significant detachment
 494 of the California Undercurrent from the upper slope, suggesting a possible important role
 495 for BBL dynamics.

496 A more direct assessment of the importance of Undercurrent-topographic interac-
 497 tions is provided by analysis of the along-isobath component of the mean flow near the
 498 bottom, which we diagnose as $\hat{z} \cdot (\vec{U}_b \times \vec{\nabla} H) / |\vec{\nabla} H|$, where \hat{z} is the vertical unit vector,
 499 \vec{U}_b the velocity at the bottom, and H the bottom topography. This quantity is positive
 500 nearly everywhere, i.e., the bottom current flows with the slope to its right (Figure 6 d).
 501 Larger bottom velocities occur around or just shallower of the 200 m isobath, where the
 502 core of the Undercurrent is found closer to the seafloor. These local maxima are closely
 503 correlated with intensification of the Undercurrent (Figure 6 c). Frictional interactions
 504 and mixing generate an Ekman transport in the BBL. We diagnose the cross-isobath com-
 505 ponent of this Ekman transport as $-\frac{1}{\rho f} \hat{z} \cdot \vec{\tau}_b \times \vec{\nabla} H / |\vec{\nabla} H|$ (Figure 6e), where ρ is the
 506 density, f the Coriolis term, and $\vec{\tau}_b$ the mean current stress at the bottom. Along the

shelf, the mean Ekman transport in the BBL is directed downhill, and peaks around or just above the 200 m isobath, reflecting a more intense drag from the poleward Undercurrent. Averaged between isobaths 150 m and 250 m, the alongshore bottom poleward current and downhill Ekman transport are significantly correlated ($R=0.91$). Furthermore, local maxima in the downhill Ekman transport are generally co-located with Fe transport hot-spots (Figure 6 a). In addition, considering a BBL thickness of 25 m, the average mean velocity within the BBL at the hot-spot locations ranges from 0.02 to 0.04 m s^{-1} . These relatively large velocities suggest that the downslope Ekman transport caused by the Undercurrent plays a major role in transport of Fe away from the continental shelf.

4 Discussion and Conclusions

Our high-resolution simulations reveal that the continental shelf of the U.S. West Coast is a major source of Fe for the North Pacific Ocean. The majority of this Fe is locally released by the sediment under low- O_2 conditions. While most of this Fe (77%) is scavenged by sinking particles, or biologically taken up on the shelf, a substantial fraction (17%) escapes removal and is transported offshore by a combination of mean and eddy currents, ultimately enhancing primary production in the strongly Fe-limited open-ocean environment.

We highlight two main pathways for the shelf-to-basin Fe transport along the northern U.S. West Coast. The first, in the SBL, is responsible for 19% of the total transport, and is mostly active during upwelling. The second, in the BBL, is responsible for 46% of the total transport and is active year-round. These findings are consistent with results by Siedlecki et al. (2012), but emphasize the dynamics of the California Undercurrent and its interactions with the seafloor. Amplification of BBL Fe transport in the wake of upwelling reflects a combination of more vigorous bottom currents, and more intense benthic Fe release under seasonal hypoxia (Severmann et al., 2010).

Several hot-spots of cross-shelf Fe exchange can be observed along the coast, driven by meanders of the California Undercurrent, and enhanced by local maxima in benthic Fe release in low- O_2 bottom waters. Standing and transient eddies contribute each to about one-third of the total Fe transport in the SBL, and increase transport in the BBL by nearly 50%. In the water column interior, eddies oppose the mean circulation, sustaining net offshore Fe transport despite the mean upwelling circulation directed inshore.

While model biases remain, a realistic representation of the cross-shore dissolved Fe gradients and fine-scale circulation along the shelf (Damien et al., 2023) suggests that our findings are robust to model assumptions and can be qualitatively extended to other productive margins with similar circulation and low- O_2 conditions. Release from the sediment is a dominant source of Fe not only in the California Current, but also in other EBUS and low- O_2 continental margins (Elrod et al., 2004; Johnson et al., 1999; Severmann et al., 2010; Lam et al., 2006; T. Liu et al., 2022). Because of the widespread distribution of coastal low- O_2 environments (Breitburg et al., 2018; Fennel & Testa, 2019) benthic Fe fluxes from continental margins likely impacts primary production in both nearshore and offshore waters over vast parts of the global ocean (Elrod et al., 2004; Johnson et al., 1999; Severmann et al., 2010; Lam et al., 2006; T. Liu et al., 2022). Extrapolating our findings for the U.S. West Coast, we argue that a better characterization of the fine-scale mechanisms controlling the shelf-to-basin Fe transport is crucial for understanding variability in oceanic primary production and to improve ocean biogeochemical models from regional to global scales (Tagliabue, Williams, et al., 2014; Boyd & Tagliabue, 2015; Tagliabue et al., 2016).

Recent studies based on observations (Wong et al., 2022) and models (Misumi et al., 2021) suggest the presence of preferential export pathways for shelf-derived Fe in the North Pacific, facilitated by protection by organic ligands and interaction with slowly

sinking particles. A more sophisticated representation of Fe sources and ligand dynamics in our model would likely improve its ability to capture the observed range of Fe variability. This would lead to higher Fe concentrations and stronger gradients, and potentially greater transport than currently found, in particular by eddies. Thus, the quantitative results presented here should be interpreted cautiously and are likely to represent a lower bound for the magnitude of the shelf-to-basin Fe transport of the real ocean.

More importantly, our results expand previous findings on the importance of fine-scale circulation from other basins, such as subarctic Pacific, subtropical Atlantic, and the Southern Ocean, which highlighted the role of eddies in the delivery of dFe and macronutrients from continental margins to the open ocean (Fiechter & Moore, 2012; Conway et al., 2018; Jersild et al., 2021; Lovecchio et al., 2022; Frenger et al., 2018; St-Laurent et al., 2019). We emphasize the importance of mesoscale and submesoscale currents to enhance or counterbalance the mean circulation, not just in the open ocean, but specifically on the continental shelf where Fe inputs are greatest. These fine-scale circulation patterns not only include transient eddies and cross-shelf meanders in the prevailing currents but also subsurface coherent eddies (Pelland et al., 2013) that trap Fe-laden BBL waters (Molemaker et al., 2015) and transport them offshore, often hundreds to thousands km away from the margins where they originate from (McCoy et al., 2020). A clear separation of the transport by subsurface coherent eddies from that of other transient eddies would require detecting and tracking them in model output (Frenger et al., 2018). We also note that our model configuration does not include tidal motions, which would cause a tidal rectification of advective Fe transport — i.e., a net Fe flux averaged over a tidal cycle — and increased vertical mixing near the bottom, potentially enhancing the shelf-to-basin transport. We leave a detailed analysis of these dynamics to future work.

The physical and biogeochemical processes that govern the shelf-to-basin Fe transport, such as benthic release and biological uptake, are likely to undergo significant changes in a future climate, following widespread oceanic warming, changes in circulation, and O₂ loss (T. Liu et al., 2022; Wallmann et al., 2022; Kwiatkowski et al., 2020). In addition, the shelf-to-basin pathways described here are expected to affect the cross-shore transport of non-dissolved Fe pools (in particular suspended Fe minerals), as well as other trace elements (Weber et al., 2018; Richon et al., 2020), which can also modulate primary production in far-field regions. We argue that future studies should combine an improved representation of Fe cycle processes — in particular ligand and particulate dynamics — with a more accurate understanding of the small-scale physical circulation along continental margins, as informed by high-resolution models and observational campaigns. In turn, a correct representation of the shelf-to-basin transport of Fe, trace metals, and other limiting nutrients is essential to improve current models used for ocean biogeochemistry and climate change studies (Tagliabue et al., 2016, 2020).

5 Open Research

The ROMS-BEC model code used to generate the simulations can be found manuscript submitted to *Global Biogeochemical Cycles* on the Zenodo repository, under the following DOI: (Kessouri et al., 2022) <https://doi.org/10.5281/zenodo.3988618>. Because of the size, model outputs can not be deposited on a public repository. The simulations are reproducible using the setup and forcing described . The compilation of Fe observations along the US West Coast can be found at the Biological and Chemical Oceanography Data Management Office (BCO-DMO) under the following DOI: (link will be provided before publication).

604 **Appendix A Eddy decomposition of the Fe transport and online com-**
 605 **putation**

606 Following the approach by Lee and Coward (2003), we diagnose the spatial and tem-
 607 poral variabilities of the shelf-to-basin Fe transport flux by depositing it into a mean flux
 608 plus two eddy-induced fluxes. The approach involves two low-pass filters, one in time and
 609 one in space, as described in the main text. First, the cross-shelf Fe flux, $T = u Fe$ is
 610 decomposed into time mean and fluctuating components as:

$$u = \bar{u} + u' \quad \text{and} \quad Fe = \overline{Fe} + Fe', \quad (\text{A1})$$

We then apply a spatial decomposition into alongshore mean and alongshore variations
 to the temporal mean component (A2) or to both components (A3) as:

$$u = [\bar{u}] + \bar{u}^* + u' \quad \text{and} \quad Fe = [\overline{Fe}] + \overline{Fe}^* + Fe' \quad (\text{A2})$$

or

$$u = [\bar{u}] + \bar{u}^* + [u'] + u'^* \quad \text{and} \quad Fe = [\overline{Fe}] + \overline{Fe}^* + [Fe'] + Fe'^* \quad (\text{A3})$$

Here, by construction, $\bar{u}' = 0$, $\overline{Fe}' = 0$, $[u^*] = 0$, and $[Fe^*] = 0$, so that the transport
 $T = u Fe$ has the following property:

$$\overline{u Fe} = \bar{u} \overline{Fe} + \overline{u' Fe'} \quad (\text{A4})$$

$$[u Fe] = [u] [Fe] + [u^* Fe^*] \quad (\text{A5})$$

$$(\text{A6})$$

Therefore, the alongshore and temporal mean Fe transport, $[\overline{T}]$, is simply :

$$[\overline{T}] = [\overline{u Fe}] = [\bar{u}] [\overline{Fe}] + [\bar{u}^* \overline{Fe}^*] + [\overline{u' Fe'}] \quad \text{from A2} \quad (\text{A7})$$

or

$$[\overline{T}] = [\overline{u Fe}] = [\bar{u}] [\overline{Fe}] + [\bar{u}^* \overline{Fe}^*] + [\overline{u'}] [\overline{Fe'}] + [\overline{u'^* Fe'^*}] \quad \text{from A3} \quad (\text{A8})$$

We further define the mean, standing eddy, and transient eddy transport terms as:

$$[\overline{T^{MM}}] = [\bar{u}] [\overline{Fe}] \quad (\text{A9})$$

$$[\overline{T^{SE}}] = [\bar{u}^* \overline{Fe}^*] \quad (\text{A10})$$

$$[\overline{T^{TE}}] = [\overline{u' Fe'}] = [\overline{u'}] [\overline{Fe'}] + [\overline{u'^* Fe'^*}] \quad (\text{A11})$$

611 As discussed in the main text, $[\overline{T^{MM}}]$ is the cross-shelf Fe transport caused by the
 612 climatological mean circulation and Fe distribution, $[\overline{T^{SE}}]$ is the cross-shelf transport
 613 caused by meanders and standing eddies that intersect the shelf break, and $[\overline{T^{TE}}]$ is the
 614 cross-shelf transport caused by transient eddies.

A caveat of this approach is the need to save u and Fe at a frequency high enough
 to capture the variability allowed by the temporal and spatial resolution of the model,
 (i.e., submesoscale). Considering the long duration of the simulations needed to achieve
 statistically robust results, on the order of several years or longer, this implies high out-
 put storage requirements that are currently unfeasible. To overcome this limitation, we
 compute and average T , u , and Fe online, as the model runs, and save them at daily fre-
 quency. This allows us to capture signals at frequencies as high as captured by the model.

Then, we compute the terms in the eddy decomposition offline by difference, applying the flux formulation on filtered u and Fe as:

$$\begin{aligned} T &= T^{MM} + T^{SE} + T^{TE} \\ T^{MM} &= [\bar{u}] [\overline{Fe}] \\ T^{SE} &= \bar{u} \overline{Fe} - [\bar{u}] [\overline{Fe}] \\ T^{TE} &= T - \bar{u} \overline{Fe} \end{aligned}$$

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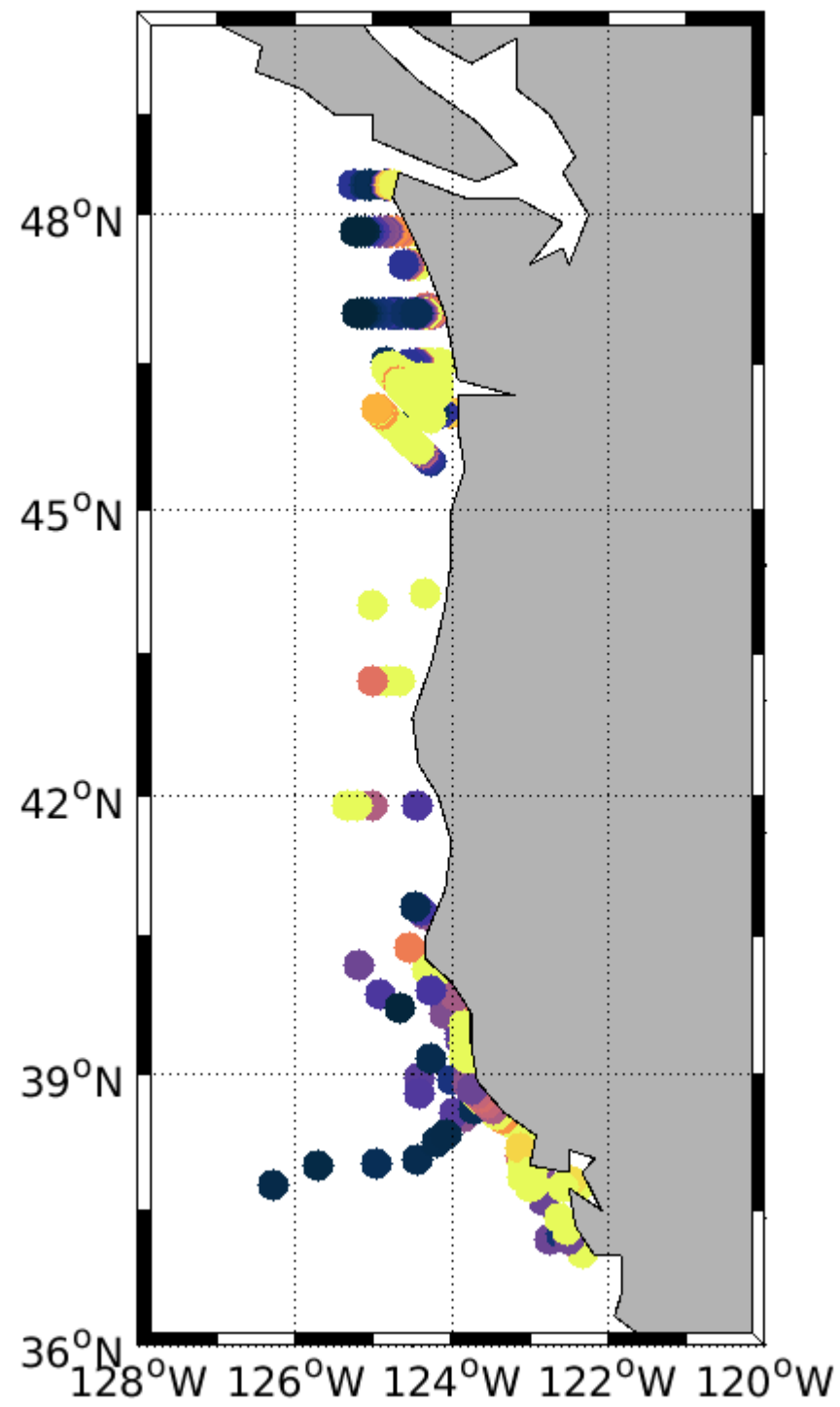
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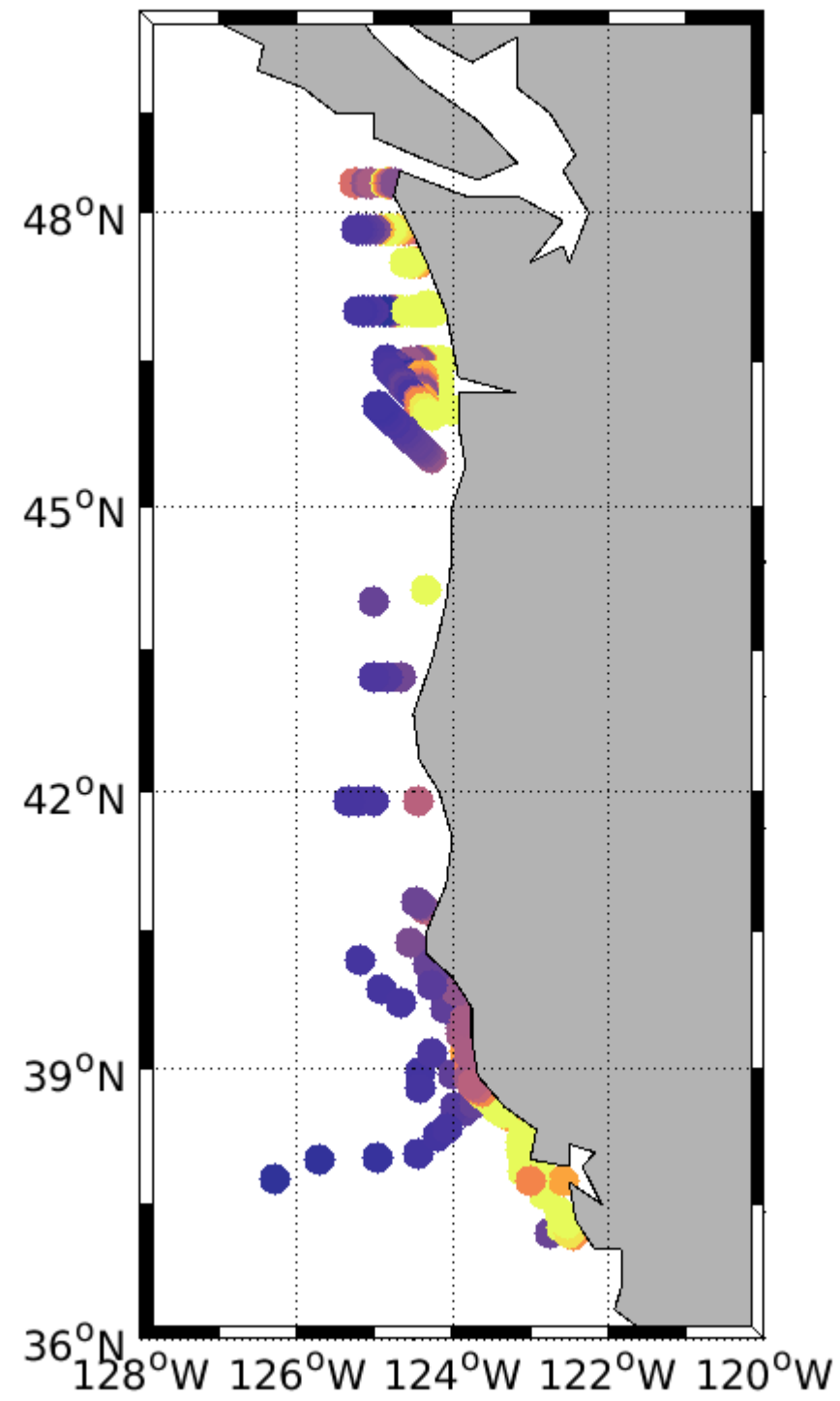
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- 1087 Wong, K. H., Nishioka, J., Kim, T., & Obata, H. (2022). Long-range lateral trans-
 1088 port of dissolved manganese and iron in the subarctic pacific. *Journal of Geo-*
 1089 *physical Research: Oceans*, 127(2), e2021JC017652. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1029/2021JC017652)
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Figure 1.

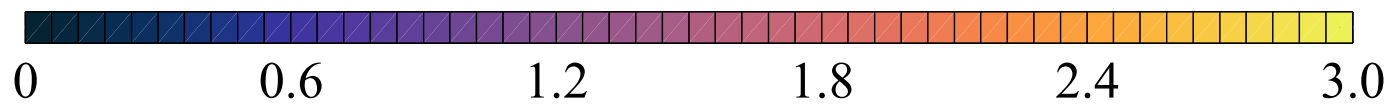
a.



b.



dFe concentrations [nM]



c.

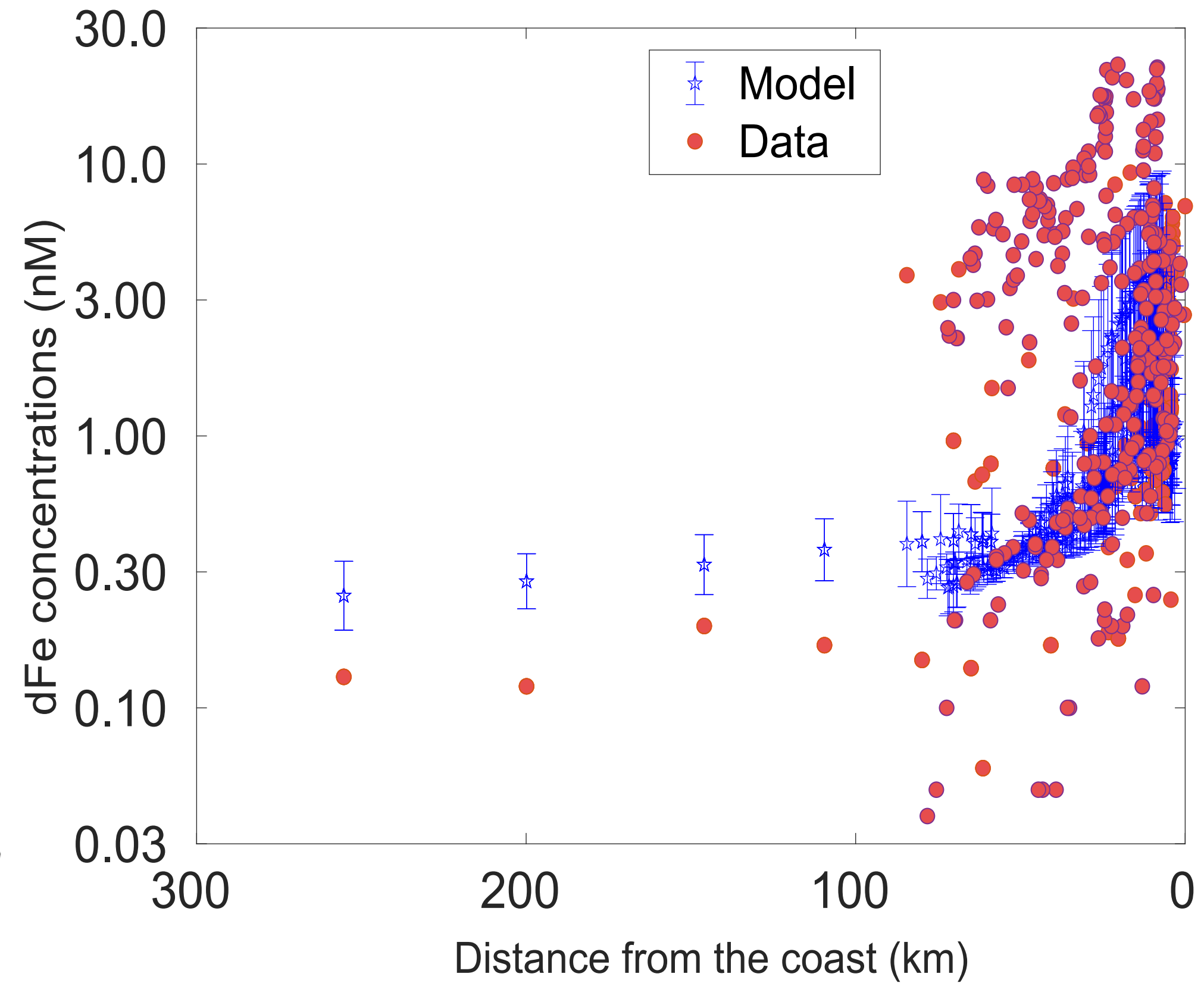


Figure 2.

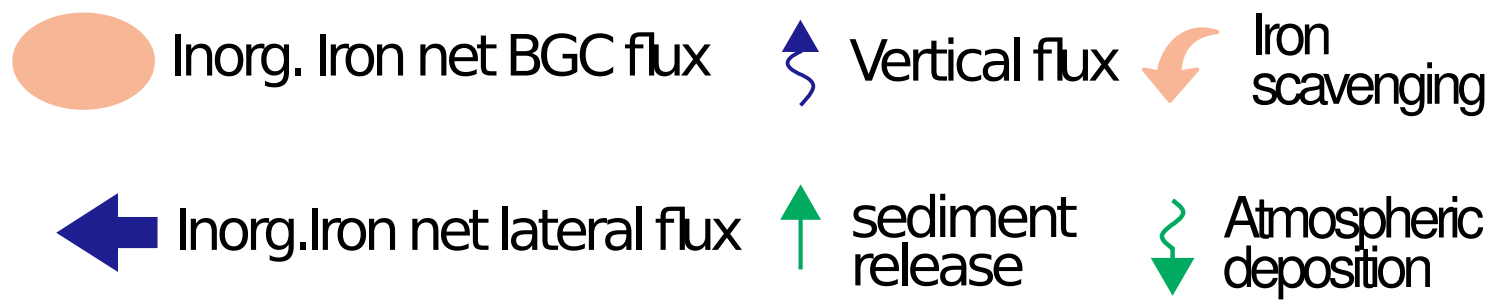
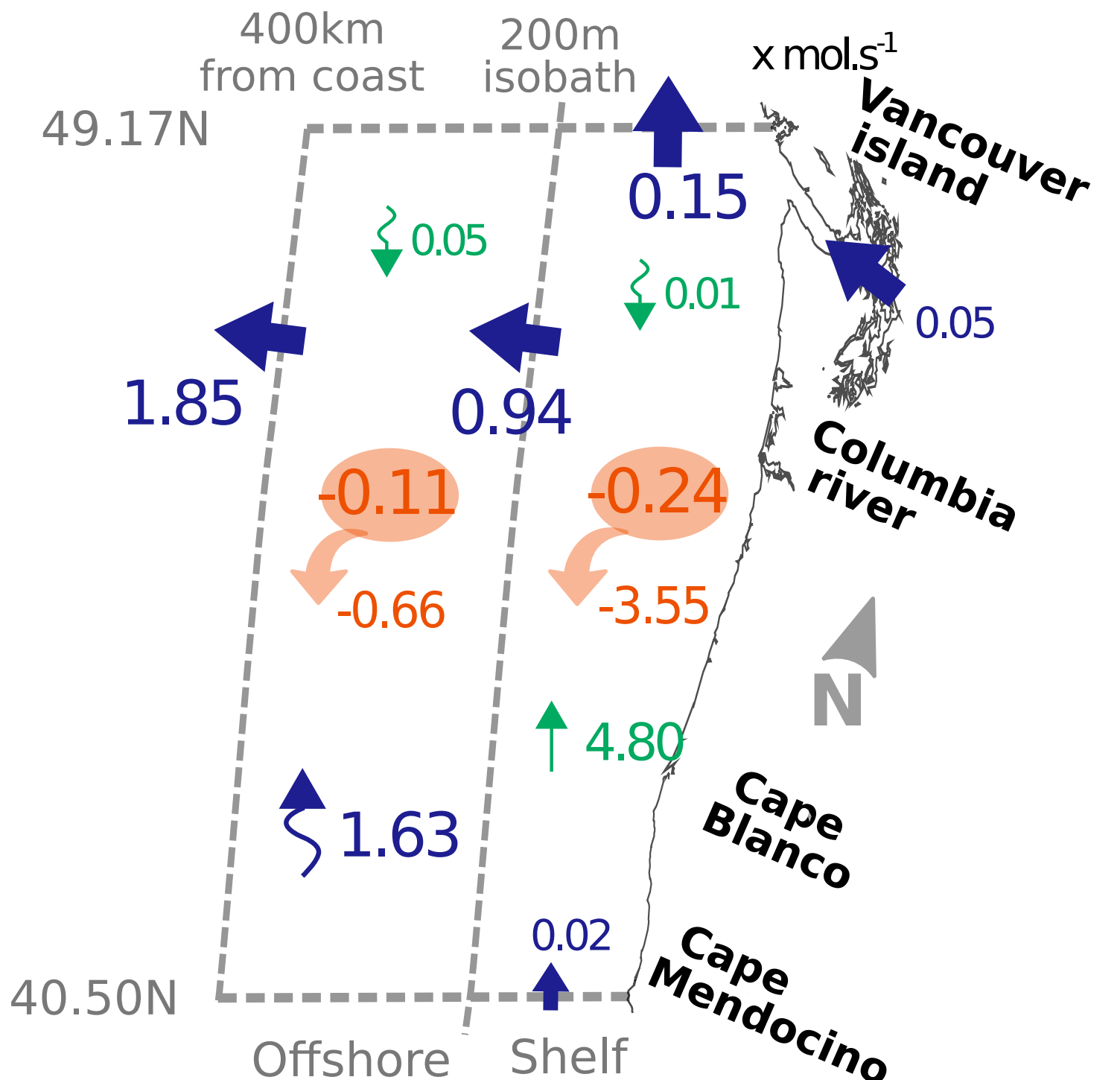


Figure 3.

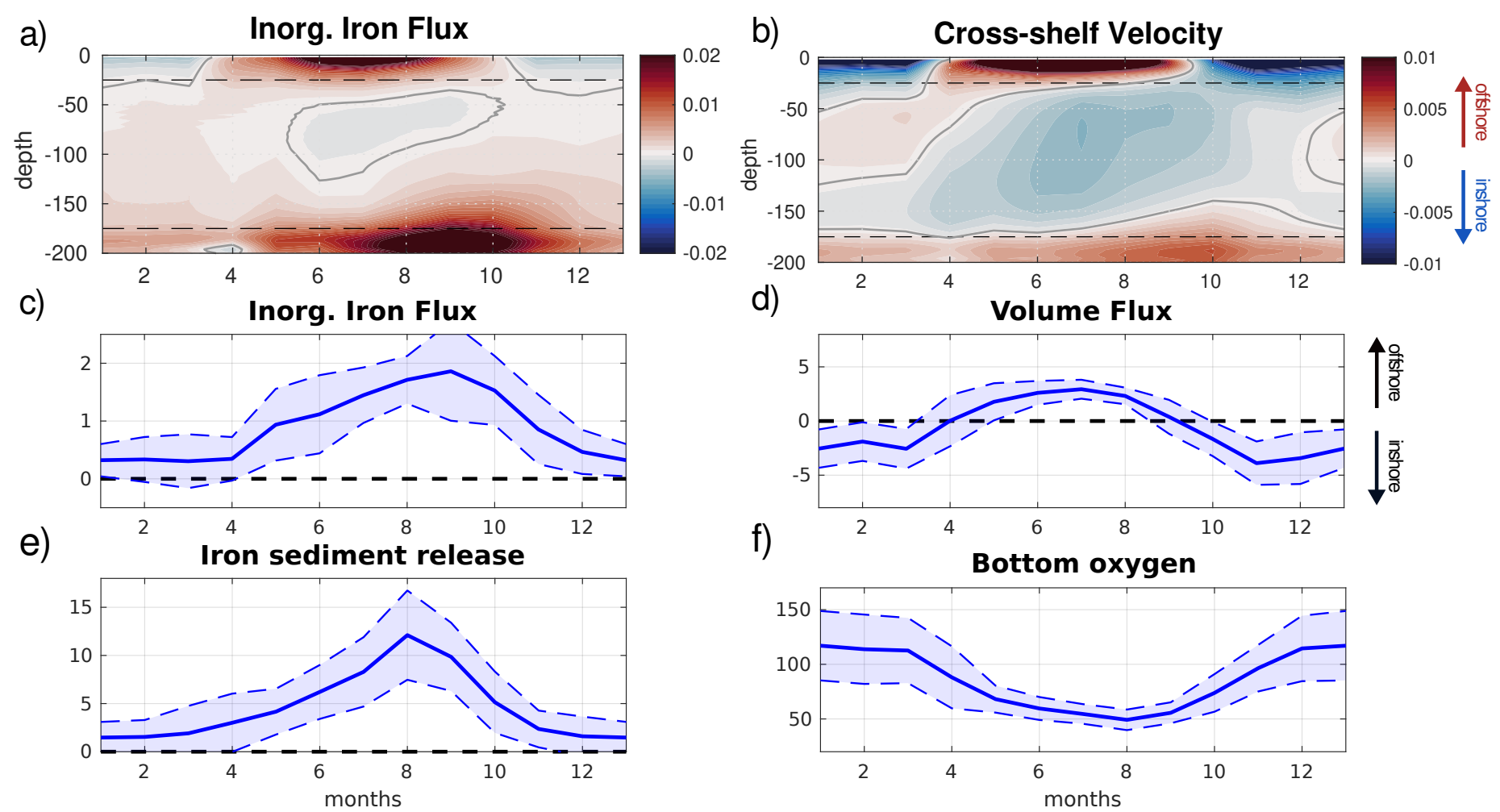


Figure 4.

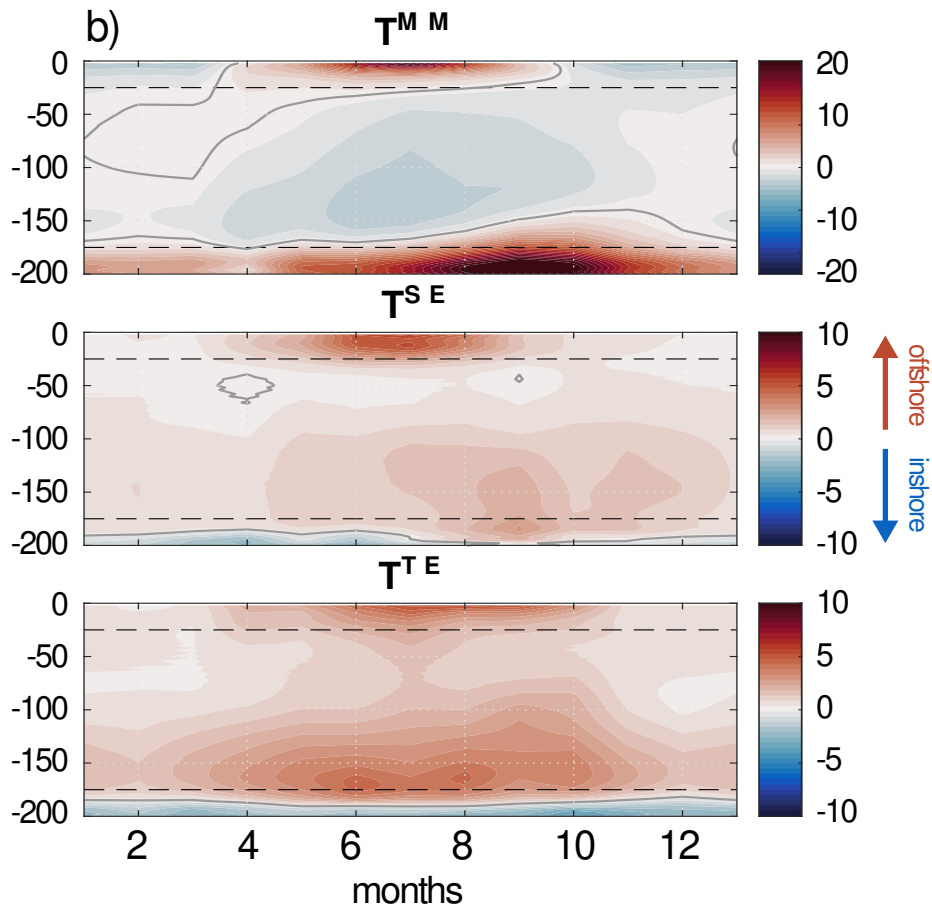
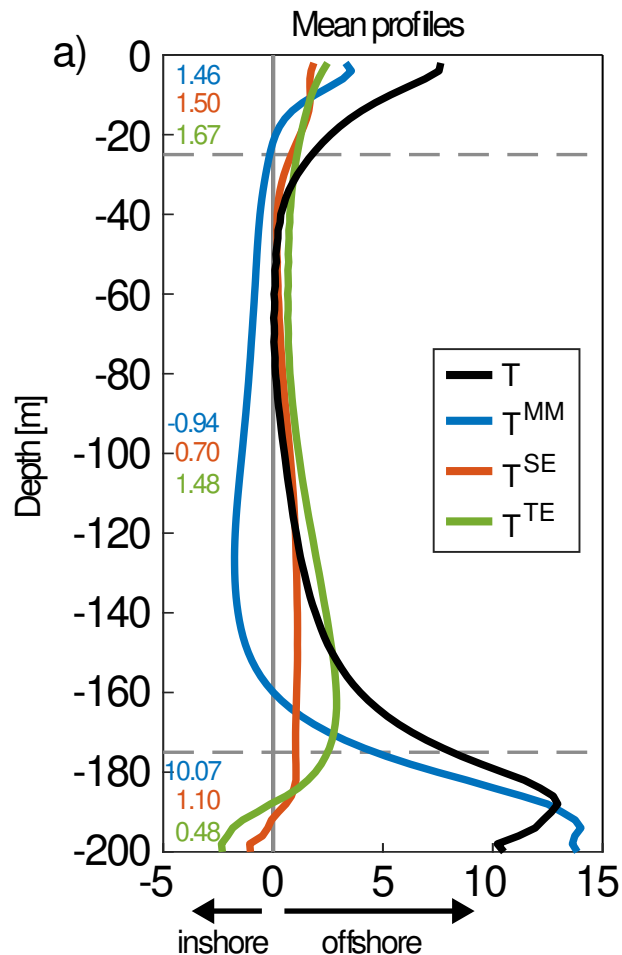


Figure 5.

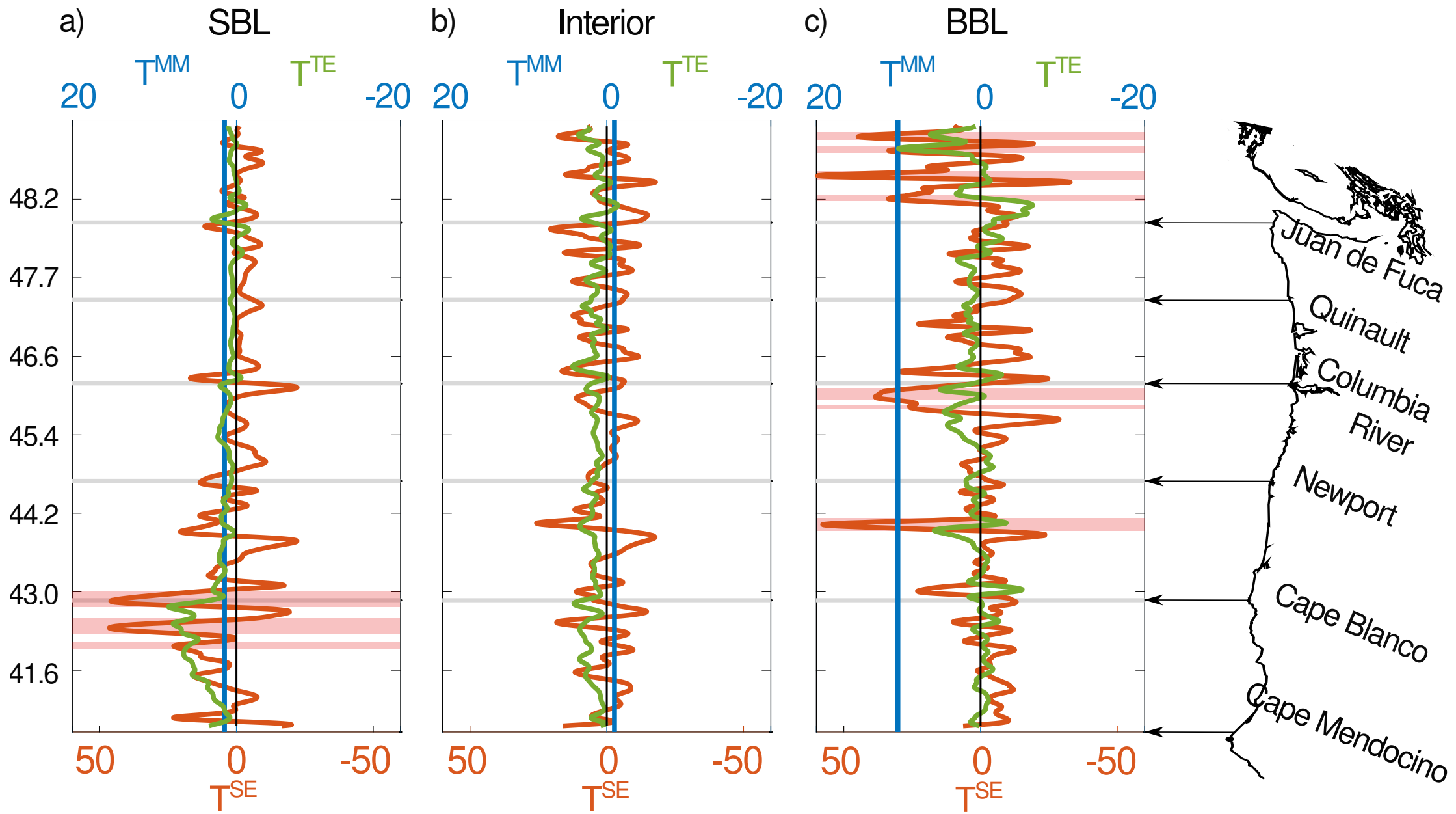
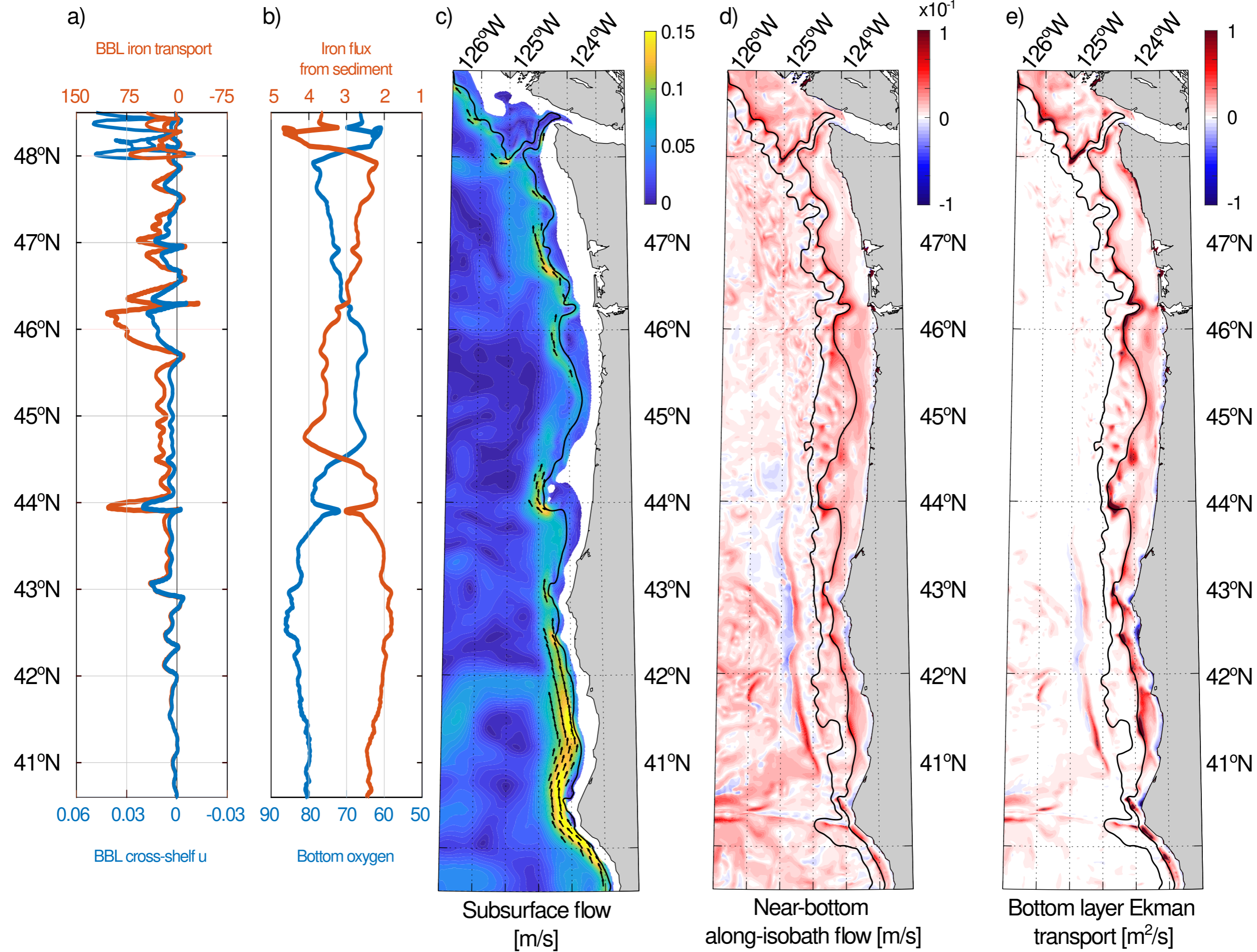


Figure 6.



Supporting Information for ”” The shelf-to-basin transport of iron from the Northern U.S West Coast to the Pacific Ocean”

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Contents of this file

1. Model validation

(i) Model validation: Iron

- Figure S1
- Table S1

(ii) Model validation: Oxygen

- Figure S2
- Figure S3

(iii) Model validation: California Undercurrent

* A.L.P. and P.D. contributed equally to
this work

- Figure S4

2. Seasonal variability of the California Undercurrent

- Figure S5

3. Spatial variability of iron concentration offshore

- Figure S6

1. Model validation

1.1. Model validation: Iron

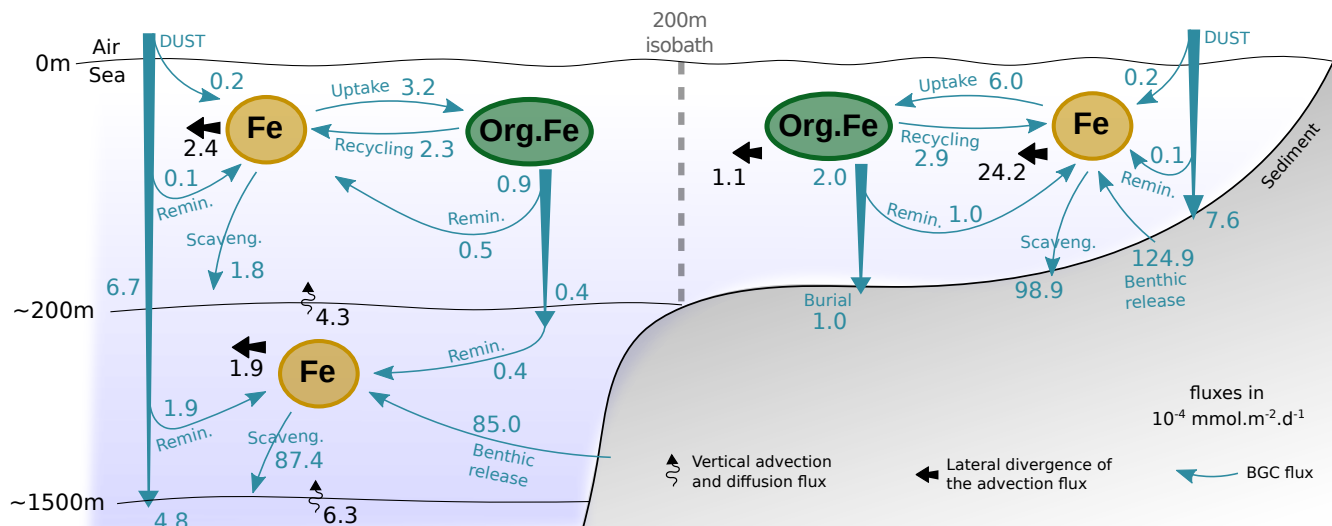


Figure S1. Budget analysis for the ocean Fe cycling along the Northern US West Coast (from Cape Mendocino to Vancouver Island) over a 9-year period from 2008-2016. Dark cyan arrows are biogeochemical fluxes; black curly arrow is vertical advection and diffusion; and black arrow is divergence of lateral advection. The boundary between coastal and open ocean areas is defined as the 200m isobath. In this figure, physical and biogeochemical terms are vertically integrated over a surface layer (0-200m) and a deep layer (200-1500m). All terms are expressed per unit area ($10^{-4} \text{ mmol/m}^2/\text{day}$) to facilitate the comparison of Fe cycling and transport terms between different regions of the domain. The transport terms are calculated as the divergence of horizontal and vertical fluxes, which, for vertical transport, include turbulent diffusion processes. Note that sedimentary dFe release in the offshore region occurs on the continental slope, where the average value is $727 \cdot 10^{-4} \text{ mmol/m}^2/\text{d}$. However, in order to represent a “closed” dFe budget (where sources and sinks balance), this figure reports the average terms normalized by the oceanic area extending up to 400 km offshore, thus much larger than the area of the continental slope, resulting in a lower average value.

Table S1. Sources of our dFe data compilation along with information on the type of Fe being measured

Data Sources	Iron Measurement Type
Landing and Bruland (1987)	dFe (0.3 μ m filter) and particulate Fe
Martin and Michael Gordon (1988)	dFe (0.4 μ m filter) and particulate Fe
Johnson et al. (2001)	dFe (0.45 μ m filter)
Johnson et al. (2003)	dFe (0.45 μ m filter)
Chase, van Geen, Kosro, Marra, and Wheeler (2002)	dFe (unfiltered, non acidified sample stream)
Fitzwater et al. (2003)	Dissolvable (Fe(III) detected after seawater is held at pH \sim 3 for 1 min.) and particulate Fe
Firme, Rue, Weeks, Bruland, and Hutchins (2003)	dFe (0.2 μ m filter)
Chase et al. (2005)	Dissolvable Fe: Fe species passed through a 20 μ m filter and were acidified inline to pH 3.4 for 1 minute prior to analysis
Lohan and Bruland (2008)	dFe (0.4 μ m filter)
Elrod, Johnson, Fitzwater, and Plant (2008)	dFe (0.5 μ m filter). Dissolvable Fe is defined as Fe leached from particles at pH 3.2
Severmann, McManus, Berelson, and Hammond (2010)	dFe (0.45 μ m filter)
King and Barbeau (2011)	dFe (0.4 μ m filter)
John, Mendez, Moffett, and Adkins (2012)	dFe (0.45 μ m filter)
Biller, Coale, Till, Smith, and Bruland (2013)	dFe (0.2 μ m filter)
Bundy, Biller, Buck, Bruland, and Barbeau (2014)	dFe (0.2 μ m filter)
Bundy et al. (2015)	dFe (0.45 μ m filter)
Bundy, Barbeau, Carter, and Jiang (2016)	dFe (0.2 μ m filter)
Hogle et al. (2018)	dFe (0.4 μ m filter)
Boiteau et al. (2019)	dFe (0.2 μ m filter)
Till et al. (2019)	dFe (0.2 μ m filter)
Kelly et al. (2021)	dFe (0.2 μ m filter)
Wong, Nishioka, Kim, and Obata (2022)	dFe (0.2 μ m filter)
Abdala et al. (2022)	dFe (0.2 μ m filter)

1.2. Model validation: Oxygen

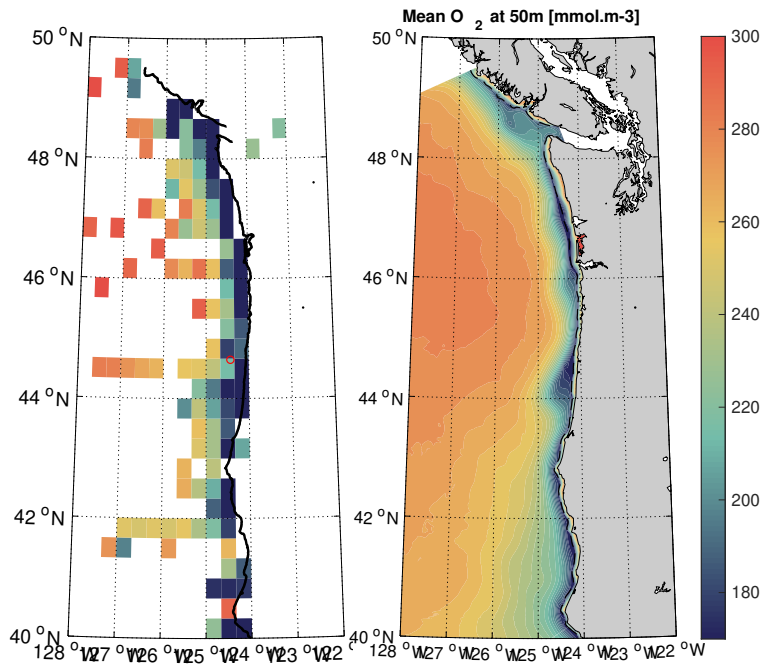


Figure S2. Maps of mean oxygen at 50m retrieved from (left) in-situ observations and (right) ROMS solutions from December 1999 to November 2017. Observations were collected from various sources and gridded at 1/3 degree resolution. For the right panel, the (dashed line) 50 m isobath is superimposed and, for waters shallower than 50 m, the mean oxygen field at bottom is shown. Taking into account that the lower mean oxygen concentration observed on the shelf is partially due to the larger amount of oxygen measurement during summer, the mean distribution of the modeled oxygen in ROMS is overall in good agreement with observed concentration.

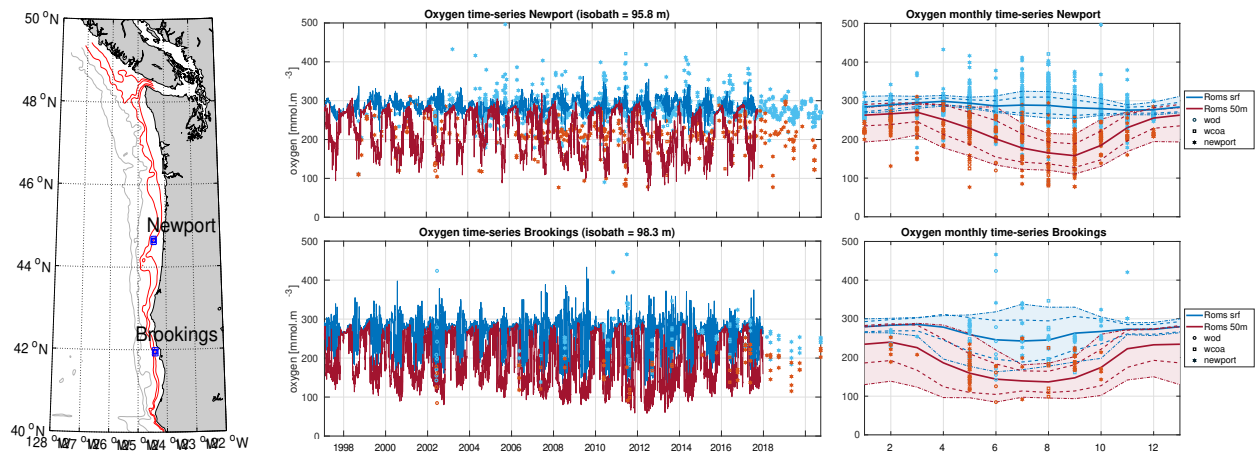


Figure S3. (Left panel) Location of the 2 geographical points selected for time-series comparisons in central and left panels. Red contours represents the 100 and 200m isobath that define the shelf and the grey contours are the 1000 and 2000 m isobath. (Central panels) Simulated (lines) and observed (dots) oxygen concentration time series at (upper) Newport and (lower) Brookings. The surface time series is displayed in blue and the 50m time series in red. For observations, each marker shape corresponds to a dataset source. We collected, compiled and merged data from the West Coast Ocean Acidification Cruises (WCOA) (Feely et al., 2016), the World Ocean Database (WOD) (Garcia et al., 2009), and the Newport hydrographic line (Risien et al., 2023). (right panels) Corresponding seasonal cycles with the (full line) mean, the (dashed line) daily rms and the (dots line) 5th and 95th percentiles of the monthly distribution. Despite an underestimation of the temporal variability evidenced by the wider spread of observations measurements, the seasonal oxygen variability is overall well reproduced.

1.3. Model validation: California Undercurrent

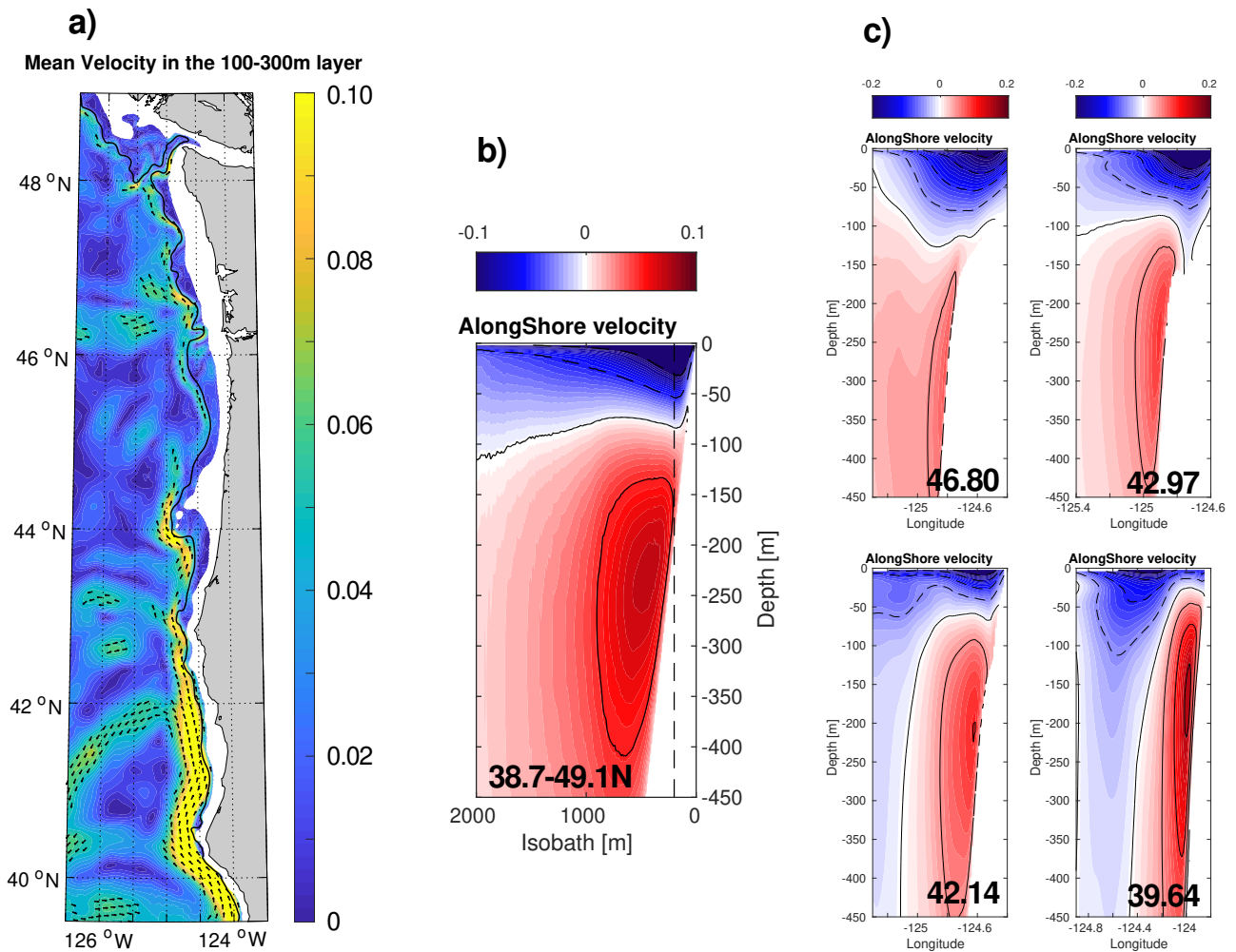


Figure S4. (a) Depth-averaged subsurface flow between 100 and 300 m in m/s . (b) Spatial mean section of alongshore flow between 38.7 and 49.1 °N, shown in an isobath/depth coordinate system. Positive velocity are in the poleward direction. Black lines evidence velocity contours in steps of 0.05 m/s , the dashed line stand for negative alongshore velocity (or in the equatorward direction). (c) Selected sections of the alongshore flow. This figure compares the modeled CUC with the ADCP observations reported in Pierce et al. (2000). Panel a compares with figure 1, panel b with figure 3 and panel c with figure 2. Despite some differences likely related to the interannual variability of the CUC, the model velocity and position are overall within the range of the observations. Thus, the model produces a realistic CUC.

2. Seasonal variability of the California Undercurrent

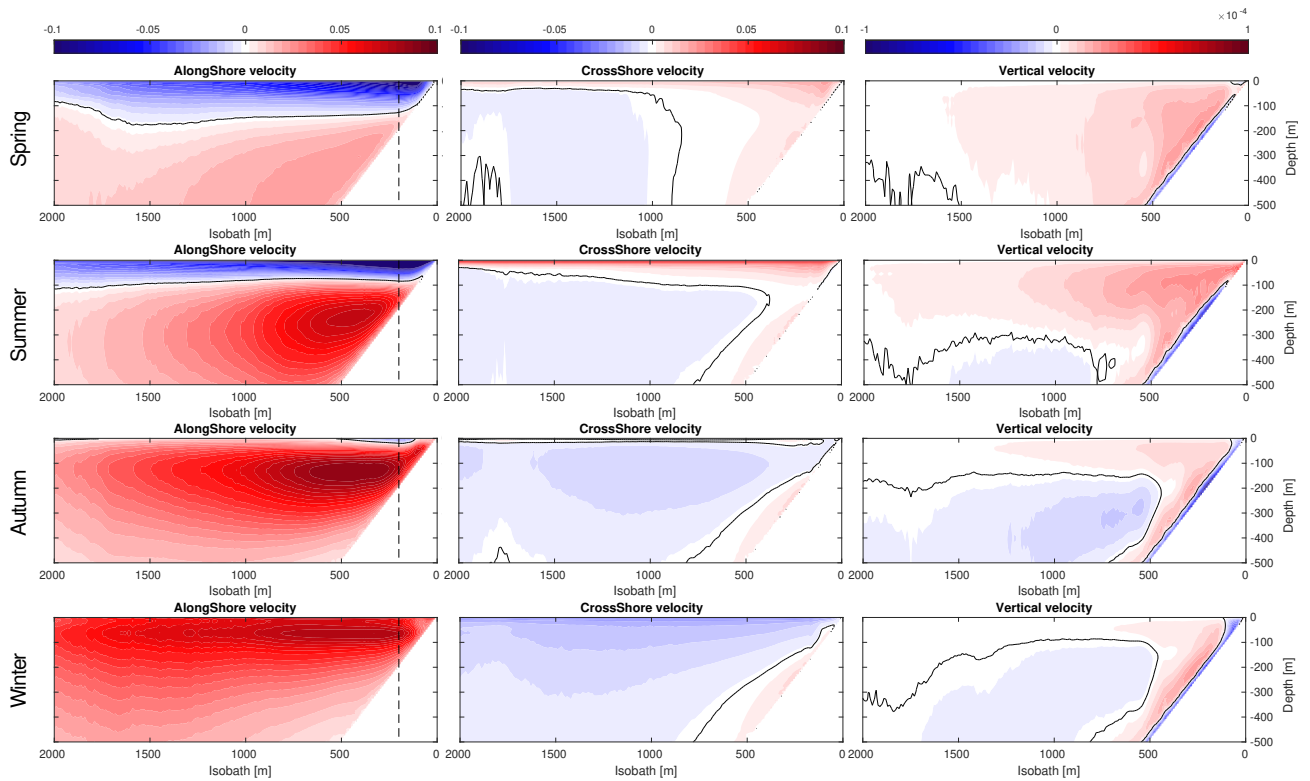


Figure S5. Seasonal variability of the nearshore flow shown in an isobath/depth coordinate system. Velocity were projected into the alongshore (positive directed poleward), crossshore (positive directed offshore), and vertical (positive directed upward) directions. Black line evidence the zero seasonal mean motion. The CUC is evidenced by the subsurface alongshore velocity maximum (right panel). It is present throughout the whole year but intensified in Autumn. The CUC-topographic interaction produces an Ekman flow in the bottom mixed layer directed offshore (central panels) and downwad (right panel). Similar to the CUC, this bottom-confined downhill flow is active during the whole year but intensified in Autumn.

3. Spatial variability of iron concentration offshore

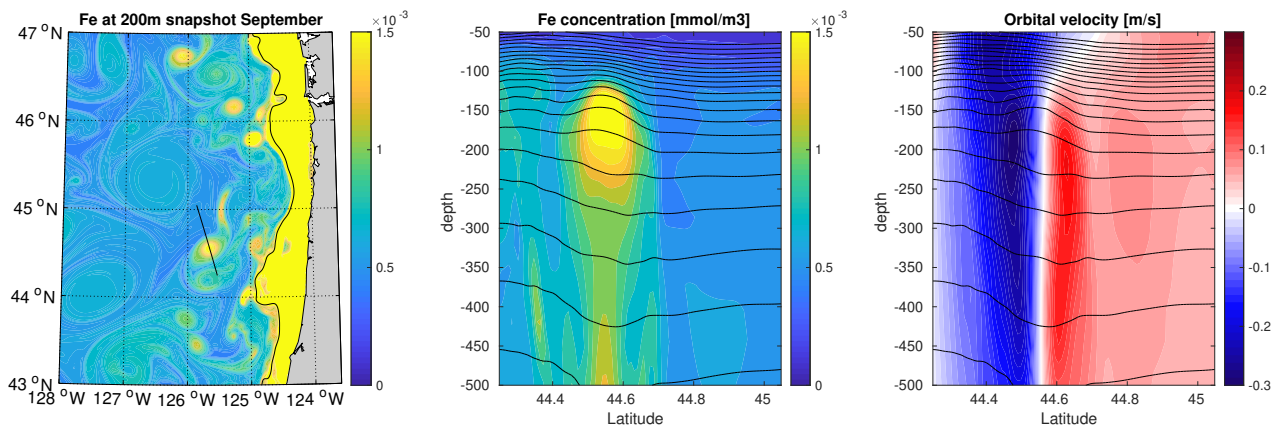


Figure S6. (a) Snapshot of the iron concentration at 200m depth in September. The black line highlights the 200m isobath and the bottom concentration is shown for depth shallower than 200m. Offshore, a significant amount of the iron released from the shelf is found in lenslike eddies characterized by a large positive iron anomaly. A crosssection (black transect on figure 1) of (b) iron and (c) velocity across one eddy evidences a subthermocline, low-stratified, and anticyclonic structure characteristic of the "Cuddies" described in Pelland et al. (2013). The large iron concentration in cuddies evidence the shelf origin of the water trapped in their cores where it undergoes mixing and stirring rates much lower than the background flow. This suggest Cuddies as a significant and effective mechanism for the iron shuttle from the shelf to remote offshore regions. This requires a closer inspection in future studies.

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