Framework for an ocean-connected supermodel of the Earth System

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Abstract

A supermodel connects different models interactively so that their systematic errors compensate and achieve a model with superior performance. It differs from the standard non-interactive multi-model ensembles (NI), which combines model outputs a-posteriori. We formulate the first supermodel framework for Earth System Models (ESMs) and use data assimilation to synchronise models. The ocean of three ESMs is synchronised every month by assimilating pseudo sea surface temperature (SST) observations generated from them. Discrepancies in grid and resolution are handled by constructing the synthetic pseudo-observations on a common grid. We compare the performance of two supermodel approaches to that of the NI for 1980—2006. In the first (EW), the models are connected to the equal-weight multi-model mean, while in the second (SINGLE), they are connected to a single model. Both versions achieve synchronisation in locations where the ocean drives the climate variability. The time variability of the supermodel multi-model mean SST is reduced compared to the observed variability; most where synchronisation is not achieved and is bounded by NI. The damping is larger in EW than in SINGLE because EW yields additional damping of the variability in the individual models. Hence, under partial synchronisation, the part of variability that is not synchronised gets damped in the multi-model average pseudo-observations, causing a deflation during the assimilation. The SST bias in individual models of EW is reduced compared to that of NI, and so is its multi-model mean in the synchronised regions. The performance of a trained supermodel remains to be tested.

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Key Points:

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10	•	Data assimilation can be used to synchronise different Earth System models.
11	•	Partial synchronisation is achieved and bias is reduced in key regions by exchanging monthly
12		sea surface temperature.
13	•	Synchronising the models towards a weighted-mean causes a deflation of variability.

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14 Abstract

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³³ Plain Language Summary

Supermodelling is a novel approach in which different models are run interactively and are 34 connected to achieve a model with superior performance. The method exploits model diversity 35 and reduces model error by training the connection terms using observations. Structural differ-36 ences between the models and the amount of data exchange among models caused technical chal-37 lenges and have limited the applicability of supermodelling to Earth system models. Here, we 38 show that data assimilation can handle these discrepancies. A demonstration is done with three 39 Earth system models having different resolution and structural differences and using a limited 40 amount of data exchange (monthly ocean surface temperature). Synchronisation is achieved in 41 several key regions. There, the ocean surface temperature error is smaller than by taking the a-42 posteriori average of the non-interactive multi-model ensemble, an approach commonly used for 43 handling Coupled Model Intercomparison Project simulations. Connecting the model via their 44 weighted mean (one of the two approaches tested) causes a spurious deflation of variability. This 45 study opens the application and training of supermodelling to Earth system models. 46

47 **1 Introduction**

Climate models have been key tools for understanding fundamental questions about our 48 climate systems. However, large uncertainties exist with many key processes being parameterised 49 and biases in several of the Earth system components (e.g., ocean, atmosphere, sea ice and land) 50 being larger than the projected changes in climate and larger than the variability being predicted 51 (Palmer & Stevens, 2019). While models have improved through the successive generations of 52 coordinate model intercomparison project (CMIP) — version 6 being the latest (Eyring et al., 2016) 53 - many of the large biases persist (e.g., the double-Intertropical Convergence Zone (ITCZ) prob-54 lem (Tian & Dong, 2020); the tropical Atlantic bias (Richter et al., 2014) and the signal-to-noise 55 paradox (Scaife & Smith, 2018). Although some simulations are being tested at a breakthrough 56 resolution (Zhongming et al., 2020), we are still decades away from being able to perform op-57 erational climate simulations with models that can explicitly resolve the most important phys-58 ical processes. In the meantime, one can test alternate methods to understand and mitigate these 59 biases using the current generation of models. 60

The classical approach to mitigate model error is to take the multi-model average of independent model simulations (a-posteriori) so that errors in the different models cancel. This approach is standard for climate projection but can also improve predictions (Branicki & Majda, 2015). One can refine the post-processing by taking a weighted mean of the different runs — an
 approach referred to as a super ensemble (Krishnamurti et al., 2016). However, this approach has
 limitations as most models share the same deficiencies (e.g., double ITCZ, warm bias in the trop ical Atlantic) and because linear post-processing does not necessarily correct non-linear responses,
 such as climate sensitivity.

Supermodel builds on the interactive ensemble (Kirtman & Shukla, 2002) where multiple 69 realizations of the same atmospheric general circulation model are simultaneously coupled to a 70 single ocean general circulation model through averaging their air-sea fluxes. Supermodels cou-71 72 ple simultaneously different models and take advantage of model diversity to compensate for their errors (Duane et al., 2018). Models are connected as they run via their state variables or their ten-73 dencies. Models can either be connected to each other (e.g. Mirchev et al., 2012; Smith, 2001) 74 or towards their weighted mean (Wiegerinck et al., 2013; Schevenhoven & Carrassi, 2021). Dur-75 ing a training phase, the connection terms are optimised to formulate a new synchronised dynam-76 ical system that achieves enhanced performance. Supermodels rely on two important properties: 77 first, it is possible to synchronise different models through a few variables — an approach referred 78 to as chaos synchronisation of non-linear dynamical systems (Pecora et al., 1997; Duane & Trib-79 bia, 2001) — and second, model diversity can encompass the true behaviour of the dynamical 80 system. 81

Supermodels have been tested with various models and experimental designs of different 82 complexity. Idealised framework experiments (or observing system simulation experiments, Halem 83 & Dlouhy, 1984) are convenient because they allow controlling challenges faced with the real framework and because the truth (constructed from a model) is known. One can introduce model 85 error (e.g., by perturbing parameter values or using a different model) and disclose only part of 86 the true model state as observations (perfect or not). Supermodelling was successfully demon-87 strated for parametric model error forming a convex envelope around the truth and that with low 88 dimensional dynamical systems (e.g., Lorenz 63, Lorentz 96, Rossler systems; see Mirchev et 89 al., 2012; Van den Berge et al., 2011; Du & Smith, 2017), quasi-geostrophic atmospheric mod-90 els (Schevenhoven & Selten, 2017; Wiegerinck & Selten, 2017), and the global atmosphere-ocean-land 91 model of intermediate complexity SPEEDO (Schevenhoven & Selten, 2017; Selten et al., 2017). 92 However, when the model error does not cancel out (parameters do not form a convex envelop 93 around the truth), one can use negative weights (Schevenhoven et al., 2019; Schevenhoven & Car-94 rassi, 2021), which raises new challenges. Furthermore, the supermodel can degrade performance 95 at a different time scale than it was trained for when the imperfect models do not fully resolve 96 the processes of the truth (e.g., in Wiegerinck & Selten, 2017, using models at a coarser resolu-97 tion). Nevertheless, the first demonstration of supermodelling with real data successfully mit-98 igated the double ITCZ bias and improved the representation of the dynamics in the equatorial Pacific (Shen et al., 2016, 2017). These results were achieved with two versions of the ECHAM5 100 Atmospheric General Circulation Models (AGCM) — each of them using a different convection 101 scheme — providing the weighted average fluxes to a single Oceanic General Circulation Model 102 (OGCM), MPIOM (Shen et al., 2016). One may expect to enhance the performance of super-103 modelling by broadening model diversity that expands the convex envelop so that it may enclose 104 the true Earth system – to the extent possible, recognising that the real climate system is far more 105 complicated than any numerical model. 106

Building a supermodel with fundamentally different ESMs raises new challenges. Mod-107 els do not share the same model-state space, nor do they have comparable resolution and repre-108 sentativity (Janjić et al., 2018). For this purpose, Du and Smith (2017) suggested formulating pseudo-109 observations of the different models and assimilating them back into the different models. The 110 approach was demonstrated with low-dimensional systems and in an idealised framework. An-111 other challenge is practical limitations with feasible data volume exchanged among models and 112 frequency of the synchronisation steps. Here, we aim to assess whether a supermodel can achieve 113 synchronisation (a requirement of supermodelling) in a configuration where data exchange is sparse, 114 synchronisation steps infrequent, and models have a radically different structure. We will apply 115 part of the formalism proposed by Du and Smith (2017)— connecting models via assimilation 116



Figure 1. Schematics of the supermodel. The green arrows denote dynamical coupling in the individual models, and the yellow band denotes synchronisation done in the different ocean models.

of synthetic pseudo-observations generated from the multi-model ensemble. We will examine the impact of generating the pseudo-observations from a single model or constructing it from an equal-weighted mean. In future work, we will implement training and analyse the resulting supermodel's performance.

This paper is organised as follows. Section 2 presents the practical implementation of the supermodelling framework: the description of the individual ESMs, the synchronisation methodology and the data assimilation method. Section 3 introduces the validation data sets and metrics, and Section 4 presents the result of two prototype supermodelling approaches, first globally with a focus on the damping of internal variability and impact on the SST bias and second with a focus on the ENSO region.

2 Supermodelling framework

This section describes the practical implementation of the supermodel framework for Earth system models using ocean connection. It combines the Norwegian Earth System Model (NorESM), the Community Earth System Model (CESM) and the Max Planck Institute Earth System Model (MPI-ESM) (see Figure 1). The three ESMs are connected via their SST every month. We implement an individual assimilation system based on the Ensemble Optimal Interpolation. The data assimilation method updates the whole water column based on the synthetic SST pseudo-observations constructed from the multi-model ensemble mean or from a single model.

2.1 CESM

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The Community Earth System Model is a global, fully coupled model for climate simu-136 lations. We used the CESM Large Ensemble Project (LENS) version based on CESM1.1.2 (Kay 137 et al., 2015), with all components at approximately 1° horizontal resolution. External forcing com-138 plies with CMIP5's historical experiment. The atmospheric component is the Community At-139 mosphere Model version 5, (CAM5, Hurrell et al., 2013) with 30 vertical levels and a finite-volume 140 grid (f09, i.e., approximately 1°). The Community Land Model, version 4 (CLM4), is on the same 141 grid as the atmosphere. The Parallel Ocean Program, version 2 (POP), is run with 60 vertical lev-142 els. The horizontal resolution of the ocean is approximately 1°, but it is enhanced in the merid-143 ional direction around the equator and both in zonal and meridional directions at high latitudes 144 (g16 grid). The sea ice [Los Alamos Sea Ice Model (CICE)version 4] component model is on 145 the same grid as the ocean model. We use historical forcing from 1920 to 2005 (Lamarque et al., 146 2010) and retrieve the initial conditions from NCAR repository (b.e11.B20TRC5CNBDRD.f09_g16.001) 147 in 1950. As POP2 only allows modification for one time level of the leapfrog scheme, we fol-148 low the approach of DART (Anderson et al., 2009). It uses a forward Euler scheme for the first 149

time step but reverts to the leapfrog scheme afterwards. The barotropic velocities and surface pres sure gradients are adjusted to preserve the ocean volume. The flag POPDART can activate this
 option in POP2.

2.2 NorESM

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We use the medium-resolution NorESM1-ME (Tjiputra et al., 2013) that contributed to the 154 Coupled Model Intercomparison Project Phase 5 (CMIP5). It is based on the Community Earth 155 System Model version 1.0.3 (CESM1, Vertenstein et al., 2012). However, the atmospheric com-156 ponent (CAM4-OSLO) features an advanced aerosol chemistry scheme (Kirkevåg et al., 2013), 157 and the ocean is an updated version of the isopycnal coordinates ocean model MICOM (Bentsen 158 et al., 2012). The atmosphere and land components are configured on a finite-volume grid with 159 a latitude and longitude resolution of $1.9 \times 2.5^{\circ}$. The atmosphere component uses 26 hybrid sigma-160 pressure levels with a model top at approximately three hPa. The horizontal resolution of the ocean 161 and sea-ice model is approximately 1°. The ocean uses 51 isopycnal layers and two layers rep-162 resenting the bulk mixed layer with time-evolving thicknesses and densities. The initial condi-163 tion is taken from a historical simulation in 1980 that started from a stable pre-industrial simu-164 lation in 1850. 165

166 **2.3 MPI-ESM**

We use the MPI-ESM1-LR (Block & Mauritsen, 2013; Dunstone et al., 2018; Giorgetta 167 et al., 2013) model that contributed to the CMIP5. The AGCM of MPI-ESM1-LR is the 6^{th} gen-168 eration European Centre Hamburg general circulation model (ECHAM6 Stevens et al., 2013), 169 and the OGCM is the Max Planck Institute Ocean Model (MPIOM) (Marsland et al., 2003; Jung-170 claus et al., 2013). The land model (JSBACH, Reick et al., 2013; Schneck et al., 2013), which 171 includes vegetation, and the marine biogeochemistry model (HAMOCC5, Ilyina et al., 2013) are 172 considered as subsystems of ECHAM6 and MPIOM, respectively. ECHAM6 employs T63 spec-173 tral resolution (approximately 1.9° horizontal resolution) and 47 vertical levels, and MPIOM em-174 ploys a rotated curvilinear grid with an approximate 1.5° horizontal resolution and 40 z-levels. 175 The poles of MPIOM are moved to Greenland and the Weddell Sea. 176

Table 1. We summarise the key characteristics of the different ESMs used. The first column details the model versions. The second column reports the name of the ocean models, their resolution and coordinate system in the vertical — i.e., isopycnal (σ -coordinate) or geopotential depth (*z*-coordinate). The third column is the name of the atmospheric models and their discretisation scheme. The last column provides a reference to the model version.

Model version	Ocean	atmosphere	reference
NorESM1-ME	MICOM(σ; 1°)	CAM4 (finite-volume, 2°)	Tjiputra et al. (2013)
CESM1.1.2	POP2 (z, 1°)	CAM5 (finite-volume, 1°)	Kay et al. (2015)
MPI-ESM1-LR	MPIOM (z,1.5°)	ECHAM6 (spectral, 2°)	Block and Mauritsen (2013)

177 2.4 Synchronisation methodology

The three ESMs use different models, grids, coordinates, and resolutions (Table 1). NorESM and CESM may be more similar, but they use a fundamentally different ocean model (in geopotential depth for CESM and isopycnal coordinates for NorESM), and the atmospheric model in CESM is a more advanced version (CAM5 vs CAM4) and has a higher resolution (1° vs 2°). Consequently, one cannot simply use state replacement or nudging, as they will generate imbalances and may lead the models to crash. Data assimilation (DA) can estimate the best possible (most likely) state based on observations, a dynamical model, and their uncertainties. It is designed to
 preserve the dynamical consistency of the individual models and optimally handle observation
 uncertainty (Carrassi et al., 2018).

A limitation of DA methods for this application is that they are (with few exceptions, e.g., S. Zhang et al., 2007; Nerger et al., 2020) working offline - meaning that the model is stopped, the state written on disk, data assimilation applied on the files and the model restarted. With such large systems as the ESMs, the time required for initialising the model and writing the input/output is burdensome (see, e.g., Karspeck et al., 2018), limiting the feasible frequency of the synchronisation. Similarly, one needs to limit the number of variables that will be synchronised to keep the cost of the DA-step low.

As a first attempt before developing a more advanced connected supermodel of ESMs, we try to synchronise the three models through their SST at a monthly frequency. SST is sufficient to constrain the variability in many regions of the earth system, particularly in the tropics (Shukla, 1978; Zhu et al., 2017; Wang et al., 2019). It is observed over a long period with a good level of accuracy, enabling the possibility to effectively train our supermodel and validate it for an independent period. With monthly synchronisation, the additional computational cost of the monthly assimilation remains small.

We test two supermodelling approaches that differ in their formulation of the pseudo-observations. 201 The first scheme belongs to the category of state-constrained weighted supermodel (Schevenhoven 202 & Carrassi, 2021). The three models are integrated forward for one month, and their SSTs are 203 interpolated to a common 1° grid. The pseudo-observations are the equal-weighted mean of these 204 outputs (referred to as EW in the following). Weights should be trained using observations to op-205 timise model performance, but this optimisation is out of the scope of the paper. The pseudo-observations 206 are then assimilated into the individual models using the EnOI method (see Section 2.5), and the 207 models are then restarted for the next cycle. The first synchronisation step in the two supermodel 208 framework started on the first of February 1980 for practical reasons. We do not synchronise mod-209 els under a the union of all three models sea ice mask. 210

In the second scheme, the workflow is similar, but the pseudo-observations are formulated 211 from a single model (hereafter referred to as SINGLE). This approach is, for instance, used in 212 the cross-pollination in time method (Smith, 2001; Du & Smith, 2017). Another objective of this 213 experiment is that it should not be affected by variability damping and can serve as a benchmark 214 for the EW. It can serve to assess the potential synchronisation that can be achieved with SST and 215 monthly synchronisation steps. We have selected arguably CESM that has higher resolution and 216 provides overall the best performance. We interpolate the CESM pseudo-observations onto the 217 common grid to have a comparable interpolation error between the two schemes. Formulating 218 the pseudo-observations in the native CESM grid enhances the performance slightly but does not 219 change any conclusion unless reported. 220

The performances of the two supermodel approaches are compared to a non-interactive multimodel ensemble (hereafter referred to as **NI**) which start from the same initial condition than the two supermodels, but do not stop for the synchronisation steps. All models start from a historical simulation in 1980 using CMIP5 historical forcing, and RCP8.5 is used for 2006 (Taylor et al., 2012).

The system runs on 11 nodes (1408 CPU) on the Norwegian high performance computer 226 Betzy (a BullSequana XH2000) and can achieve about 10 model-year per day. CESM runs on 227 eight nodes and runs one month in approximately seven minutes, NorESM runs on two nodes 228 and performs the 1-month simulation in approximately seven minutes, MPI-ESM uses one node 229 and performs a 1-month simulation in approximately four and half minutes. The model integra-230 tion accounts for approximately seven minutes, while the DA step accounts for approximately 231 two minutes. The assimilation is currently performed sequentially for each model, and this step 232 could have been reduced to 40 seconds if parallelised. 233



Figure 2. Pseudo-observation error standard deviation (in °C) used for assimilation.

234 **2.5 Ensemble Optimal Interpolation**

The Ensemble Optimal Interpolation (EnOI, Oke et al., 2002; Evensen, 2003) is a computationally cheap sequential data assimilation method derived from the Ensemble Kalman Filter (EnKF, Evensen, 2003).

The EnOI provides multivariate updates based on the model's historical covariance. Formulating the covariance from the same model ensures the preservation of linear quantities (such as geostrophic balance) and limits initialisation shock (Counillon & Bertino, 2009). The covariance matrix is constructed based on an ensemble of N model snapshots (each of dimension n, where n is the state dimension):

$$\mathbf{X}^{s} = [\mathbf{x}_{1}, \dots, \mathbf{x}_{N}] \in \mathcal{R}^{n \times N}.$$
(1)

The static ensemble anomaly \mathbf{A}^s is calculated so that $\mathbf{A}^s = \mathbf{X}^s - \overline{\mathbf{X}^s} \mathbb{1}^T$, with $\overline{\mathbf{X}^s}$ being the static ensemble mean and $\mathbb{1}_m = [1, 1, ..., 1] \in \mathcal{R}^{1 \times N}$.

To correct a forecast \mathbf{x}^{f} , using the observation vector \mathbf{d} , one can estimate a new analysis state \mathbf{x}^{a} as follows:

$$\mathbf{x}^a = \mathbf{x}^f + \mathbf{K}(\mathbf{y} - \mathbf{H}\mathbf{x}^f). \tag{2}$$

where **H** is the observation operator that relates the prognostic model state variables to the measurements. The Kalman Gain **K**, is computed as follows:

$$\mathbf{K} = \mathbf{A}^{s} \mathbf{A}^{sT} \mathbf{H}^{T} \left(\mathbf{H} \mathbf{A}^{s} \mathbf{A}^{sT} \mathbf{H}^{T} + \mathbf{R} \right)^{-1}.$$
 (3)

For each model, the static ensemble is composed of the monthly snapshot outputs from a 240 stable pre-industrial control run. Sampling the model states from a pre-industrial condition was 241 preferred over sampling it from a historical run because composing the static ensemble from a 242 model with transient forcing can introduce spurious correlation. The monthly static ensemble com-243 prises 72 members for MPI-ESM, 80 for NorESM, and 80 for CESM. We use a local framework 244 analysis (Evensen, 2003) with a radius of 235 km without tapering — this radius ensures at least 245 one observation per grid cell. We limit the local observations to one (retaining only the nearest). 246 We update all prognostic state variables in the vertical. The update is done in the model's native 247 coordinates — i.e., in geopotential depth for MPI-ESM and CESM and isopycnal coordinates 248 for NorESM. We use the k-factor formulation (Sakov et al., 2012) which artificially inflates the 249 observation error if the assimilation pushes the update beyond twice the ensemble spread. We 250 set the pseudo-observation error for both supermodel approaches as the pointwise de-seasoned 251 (with the mean seasonal cycle removed) time standard deviation of the three models divided by 252 30 (see Figure 2). We have played with the scaling factor (from 3 to 100), but the results were 253 not very sensitive to this choice (not shown) due to the k-factor formulation. 254

3 Validation data sets and metrics

For validating the simulations, we focus on SST and use the NOAA OI-SST V2 (Reynolds et al., 2002) analysis data set. We use the monthly averaged product available on a 1° grid, which extends back to 1982. We assess performance by comparing the climatological difference between the models and the observation and calculating grid cell area-weighted root mean square error (RMSE).

Internal variability can be suppressed when combining models, as for example through the averaging of multi-instance fluxes in the interactive ensemble framework (Kirtman & Shukla, 2002; Kirtman et al., 2004; W. Zhang & Kirtman, 2019). Hence, we introduce two metrics to investigate how internal variability is affected by the connection. The parameter δ is the ratio between the de-seasoned time standard deviation of the multi-model mean and the time average of the inter-model standard deviation. This metric is computed for every grid cell (see Equation 4). It can be demonstrated that this quantity should be equal to $\sqrt{\frac{1}{N_s-1}}$. With three models, this is about 0.7. If δ is larger than 0.7, some synchronisation is achieved. For example, a value of three indicates that the time standard deviation of the multi-model mean is three times larger than the inter-model standard deviation. Values lower than 0.7 can occur when the bias of the individual models is larger than the time variability of the multi-model mean (i.e., models are strongly attracted to their bias and have little variability).

$$\delta = \frac{\sqrt{\frac{1}{N_t} \sum_{i=1}^{N_t} (\overline{\mathbf{x}_i}^s - \overline{\overline{\mathbf{x}}^s}^t)^2}}{\frac{1}{N_t} \sum_{i=1}^{N_t} \sqrt{\frac{1}{N_s} \sum_{j=1}^{N_s} (\mathbf{x}_i^j - \overline{\mathbf{x}_i}^s)^2}}$$
(4)

With **x** being the model SST. The superscript *j* refers to the model indices ($N_s=3$; between 1–3),

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the superscript *i* is the time indices (N_t =312; i.e., 26*12). The overbar with superscript "s" denotes the multi-model mean, and the overbar with superscript "t" denotes the time average.

A second metric, called λ , quantifies the ratio between the multi-model mean time variability (denominator) and the mean of the variability in the unconnected individual models (numerator; Eq. 5). In a perfect synchronisation regime, there is no damping (Duane & Tribbia, 2001). Assuming that the variability among the individual models is comparable, an equal-weight multimodel mean would have comparable variability to the mean variability of the individual models (λ should be one). If only partial synchronisation is achieved, the value gets larger than one. For example, λ equal to 2 means that the standard deviation of the multi-model mean is half that of the unconnected model.

$$\lambda = \frac{\sqrt{\frac{1}{N_t} \sum_{i=1}^{N_t} (\mathbf{x}_i^j - \overline{\mathbf{x}_i}^t)^2}}{\sqrt{\frac{1}{N_t} \sum_{i=1}^{N_t} (\overline{\mathbf{x}_i}^s - \overline{\overline{\mathbf{x}}^s}^t)^2}}$$
(5)

²⁶⁴ **4 Supermodel results**

We compare the performance of EW, SINGLE and NI (see Section 2.4) over the period 1982– 2006. In Figure 3, we show the de-seasoned time standard deviation of the multi-model mean of NI, EW and SINGLE and the OISST observations. The time variability in the NI multi-model mean is substantially lower than in the observations. In EW, it is also reduced—albeit less than in the NI—, while SINGLE has nearly comparable amplitude to the observations.

In Figure 4, we are analysing the properties of each system in achieving synchronisation and causing damping. We present individual 2-dimensional maps of δ and λ (metrics introduced in Section 3) and the 2-dimensional probability density function (PDF) of δ versus λ . An ideal supermodel will have the highest probability centred on the brown line and on the right-hand side of the red line (i.e., $\delta > .7$).

We can see that δ for NI is lower or equal to 0.7 in most places, and values of λ are above one and mostly close to two. There is barely any synchronisation, and taking the multi-model



Figure 3. De-seasoned time standard deviation of SST in the NOAA OISST2 observations (upper left panel), the multi-model average of NI (upper right), EW (lower left) and SINGLE (lower right).



Figure 4. The first line is the synchronisation metric δ with NI mean, EW and SINGLE. The second line is the same for the synchronisation metric λ , and the last line is the 2-dimensional PDF of δ vs λ in the different systems. The red and brown lines highlight the threshold of 1.

mean causes damping nearly identical to a random process (i.e., $\sqrt{3}$ with three models). Hence, 277 in NI, models are only connected via their historical CMIP5 forcing, while the internal variabil-278 ity overpowers the de-seasoned time variability compared to the climate change signal from 1980-

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2005.280 In the EW supermodel, there is some synchronisation ($\delta > 0.7$) and the damping λ is mostly 281 between [1.1 1.5]. The synchronisation and the damping are improved compared to NI, and the 282 maximum likelihood moves towards the optimal. Synchronisation is maximum in the equato-283 rial Pacific, reaching a value of 3.5 and is above 1 in the tropical Pacific, part of the North Pa-284 cific and part of the North Atlantic (entrance of the Nordic Seas). In those regions, λ lowers towards 1 (albeit remaining above). However, δ remains below or close to 0.7 in several regions: 286 e.g., in the equatorial Atlantic and the eastern boundary upwelling system and most of the South-287 ern oceans. No synchronisation is achieved there, and internal variability of the multi-model mean 288 is damped and of comparable amplitude to NI. The shape of δ in the equatorial Pacific shows max-289 ima on either side of the equator but achieves only moderate synchronisation at the equator. Hence, 290 the fast Kelvin waves are driven by wind bursts in the equatorial Pacific, while atmospheric vari-291

ability is poorly constrained in our ocean-constrained system. Slower Rossby waves control the 292 thermocline variability off the equator, and the effect of winds is weaker and slower. Ocean data 293 assimilation with monthly synchronisation steps can control such a process better. 294

The pattern of the value of δ with SINGLE resembles that of EW, but the values are con-295 siderably larger. Synchronisation δ is greater than 1.5, and the damping λ is reduced (approach-296 ing one and rarely exceeding 1.5). Actually, the CESM model (to which all models are connected) has a lower time standard deviation than MPI-ESM and NorESM, while the denominator in λ 298 is the average of the time standard deviation of the three models. It partly explains the slight damp-299 ing in λ . Synchronisation is now achieved in most parts of the southern oceans, the Atlantic Ocean, 300 and the Indian Ocean. In the equatorial Atlantic and the eastern boundary upwelling system, syn-301 chronisation is still not achieved, and the damping is substantial. This region is notorious for be-302 ing very challenging for models (Richter et al., 2014) and a considerable fraction of the bias re-303 lates to atmospheric origin (e.g., wind bias Koseki et al., 2018).

The above analysis focused on the time standard deviation of the multi-model mean. We 305 analyse now the de-seasoned time standard deviation of the individual models. In Figure 5 we 306 show the quantile-quantile plot of CESM-NI versus CESM-EW and CESM-SINGLE. Focusing 307 on CESM allows us discarding discrepancies of variability between models, e.g., CESM has weaker 308 temporal variability than the two other models. In EW, the time variability of CESM is reduced. The other models show comparable damping (not shown). Contrarily, SINGLE does not cause 310 deflation, and the regression line yields a slight overestimation. It may relate to the fact that as-311 similating pseudo-observations in a common grid adds energy to the system (because of the im-312 balance). A perfect fit is obtained if pseudo-observations are formulated in the CESM native grid 313 (not shown). 314

In Figure 6, we analyse the SST bias of the individual models compared to OISSTV2 ob-315 servations. In NI, CESM has a lower bias than MPI-ESM and NorESM. NorESM has a pronounced 316 cold bias. In the EW supermodel, the bias of all models is reduced compared to that of NI (al-317 beit identical in MPI-ESM). Achieving such a result is highly promising, considering that weights 318 have not been trained. In SINGLE, the bias structure in all models resembles that of CESM in 319 NI. It yields a bias reduction for MPI-ESM and NorESM but increases the bias in CESM com-320 pared to NI. The interpolation to the common grid causes the degradation (not shown). 321

The bias of the multi-model mean of EW and NI (see Figure 7) are comparable in pattern 322 and amplitude (NI has a slightly lower area-weighted RMSE). Still, some discrepancies exist. 323 Figure 8 shows the absolute difference between the two. The bias of EW reduces in the tropics, 324 except in the equatorial Atlantic and the eastern boundary upwelling system. Improvement tend 325 to coincides with regions where synchronisations is achieved (i.e., $\delta > 1$ and $\lambda < 1.5$, see Fig-326 ure 4). In the equatorial Pacific ($5^{\circ}S-5^{\circ}N$ and $150^{\circ}W-90^{\circ}W$), the error reduces from 0.91 to 0.67°. 327 A bias reduction compared to the a-posteriori average results from the non-linear response of the 328



Figure 5. Quantile-quantile plot of the pointwise de-seasoned time standard deviation of CESM in NI versus that of EW (red) and SINGLE (blue). The dashed colour line show the regression line and the solid black line shows the perfect regression line.

climate system to model bias. Contrarily, performance degrades where synchronisation is poor 329 (Antarctic Circumpolar current, Equatorial Atlantic and eastern boundary upwelling system). There, 330 ocean synchronisation is quickly lost, and noise adds energy to the system. In the north Pacific, 331 synchronisation is achieved, but the bias is degraded. The reason is unclear, but we propose sev-332 eral hypotheses. First, variability there during winter is driven by storm systems mixing the ocean. 333 The EW experiment without atmospheric synchronisation effectively damps the impact of fluxes 334 in the ocean, which can cause a bias. Second, the Pacific Decadal Oscillation (PDO) over 1982-335 2006 is almost solely positive, while the interannual variability of our models is not synchronised 336 with the observations. The bias may relate to a modulation of the PDO. Third, we did not apply 337 synchronisation under sea ice, which may cause some artefacts near the ice edge. 338

The bias in the SINGLE multi-model mean has a comparable spatial pattern than in the CESM model of NI, but is larger because of the interpolation error in the pseudo-observation grid. It does not outperform the multi-model mean of NI.

We further analyse the results in the ENSO region, which stands out as one of the regions 342 where our ocean-connected supermodel reaches good synchronisation and reduced bias. The time 343 series of the Niño 3.4 index (SST anomaly in the region, 5°S-5°N and 120°W-170°W) is pre-344 sented in Figure 9. In the unconnected model (Figure 9), the individual models produce large ENSO 345 variability that is not synchronised. The ensemble mean shows one prominent peak when a strong 346 El Niño event occurs (by chance) in phase in 2003 in all models. The probability of such a co-347 incidental occurrence scale with the power of the number of models. In our example (with three 348 models and El Niño occurring typically every four years), such an event will happen every 64 349 years. Consequently, the NI methods will underestimate extreme events, particularly with more 350 models (as in CMIP). Contrarily, in EW and SINGLE, all models are evolving in phases. 351

To analyse the representation of extreme events of ENSO, we compare the PDF of the Niño 3.4 index for 1982-2006 in Figure 10. The tropical Pacific variability is notoriously asymmetric, with the magnitude of SST anomalies over the eastern equatorial Pacific being more prominent during the warm phase than during the cold phase (e.g., T. Zhang et al., 2017). By taking the mean of unconnected and independent models, one expects the PDF to show less variabil-



Figure 6. SST climatological bias computed over 1982-2006 for MPI-ESM, CESM and NorESM with the NI ensemble (first row) in the EW (second row) and the SINGLE (last row). The quantity in red reports the global spatial RMSE normalised by grid cell area.



Figure 7. Climatological SST bias computed over 1982–2006 for the multi-model mean of NI ensemble (left), EW (middle) and SINGLE (right). The quantity in red reports the global spatial RMSE normalised by grid cell area.



Figure 8. Difference of the absolute (EW-NI) of the climatological SST bias error of the multi-model mean over 1982–2006. Negative values indicate that the error in EW is smaller than in NI.



Figure 9. Niño 3.4 time series of the three models in the NI ensemble (left), the EW (middle) and the SINGLE (right). The black line shows the multi-model ensemble mean.



Figure 10. Histogram of Niño 3.4 over 1982–2006. The blue bars are for OISSTV2; the green bars are for the multi-model mean of NI (left), the red for EW (middle) and the yellow for SINGLE. The standard deviation and skewness of the distribution are reported in parenthesis for observations and the multi-model mean.

ity (being steeper around 0) and becoming more Gaussian (getting less skewed). The latter is be-357 cause the average of skewed distributions will converge towards a Gaussian distribution (central 358 limit theorem). As expected, there is damping of variability in the NI model mean, while in EW 359 and SINGLE, the standard deviation reasonably matches the observed PDF. All multi-model means 360 show a more skewed PDF than in the observations. However, the period 1982-2006 was quite 361 anomalous in the observations, as no significant ENSO events occurred beyond 1998. In com-362 parison, the standard deviation and skewness computed for 1950-2005 were respectively 0.87, 363 and 0.89 (T. Zhang et al., 2017). We do not see the reduction of skewness expected in NI, but we 364 think it is artificially high here because of the coincidental El Niño event referred to above. Ob-365 taining more statistically robust results requires much longer runs and thus we limit our analy-366 sis to these basic metrics (Wittenberg, 2009). 367

5 Conclusions and future perspectives

This paper investigates critical characteristics for developing the first supermodel of ESMs with ocean synchronisation. Synchronisation with models having different grids and structures is handled by assimilating synthetic pseudo-observations of SST constructed from the multi-model every month. We show that such a framework can achieve partial synchronisation in distinct regions where the ocean drives the climate variability — where the impact of atmospheric fluxes is limited in relation to the oceanic timescales.

We compared two methodologies for constructing the pseudo-observations, either from a 375 single model or an equal-weight mean. The latter tends to reduce model bias in the synchronised 376 region compared to the unconnected version. The variability of the multi-model mean in the two 377 connected models is smaller than in the observations. This damping is reduced in synchronised 378 regions and converges to the non-interactive multi-model mean in the unsynchronised regions. 379 The damping is more pronounced in the equal-weight version than in the version connected to 380 a single model, as in the former, the variability of the individual models is also reduced. The multi-381 model averaging reduces the unsynchronised variability (e.g., driven by chaotic atmospheric vari-382

ability) in the assimilated pseudo-observations causing a deflation when updating the model snap shots.

We will investigate in a future study how model error can reduce with the training of the weights. A first supermodel with monthly ocean connection via SST and with trained weights (varying spatially and monthly) will be presented in Schevenhoven et al. (in prep). The system constrains the SST bias and the double ITCZ problem in the equatorial Pacific. However, the system faces similar challenges regarding variability damping to those presented in this paper.

The system presented only uses a minimal amount of data exchange (monthly SST). Increasing the frequency of the synchronisation steps and assimilating more pseudo-observations will enhance synchronisation and may help to reduce the damping in the weighted mean supermodel framework. Ongoing works include a complete ocean pseudo-observations network (sea surface elevation, 3D hydrography), increasing the frequency of the ocean connection (to weekly synchronisation step) and complementing the system with synchronisation of other components of the Earth System Model. A supermodel with complementary synchronisation of the atmospheric component is also in development.

It is indeed unclear to speculate which of the two framework presented here will, at term, 398 achieve the best performance. Weighted-mean supermodels have drawbacks in a partial synchro-300 nisation regime (variability damping) not faced with a connection to a single model. However, 400 a weighted supermodel can provide locally a more accurate fit to the observation than a version 401 where models are connected to a single model because — in the latter, skill is bounded by the 402 best model. However, the single model to which we connect can vary spatially and for different 403 variables (Smith, 2001; Du & Smith, 2017; Schevenhoven & Carrassi, 2021). Furthermore, if one 404 can afford an ensemble of supermodels (with several members for each model), models could 405 be synchronised from a randomly drawn single member/model every time, so that the frequency 406 of the optimal weight is satisfied. Such a scheme resembles the cross-pollination in time, envisioned in Du and Smith (2017). 408

We will explore other alternatives to reduce the spurious damping in the weighted mean supermodel version by 1) adding back the reduced atmospheric driven variability term in the pseudoobservations or 2) isolating the part of variability that can be synchronised from the model snapshots (observation operator). In particular, machine learning techniques have emerged as powerful tools to carry process-identification (e.g., Sonnewald et al., 2019).

6 Open Research

⁴¹⁵ Data presented in this article has been organised and made available. It contains a sepa-⁴¹⁶ rate for each of the experiments presented (SINGLE, EW and NI). Each folder contains the multi-⁴¹⁷ model mean and the individual SST outputs from the individual models provided on a common ⁴¹⁸ grid and in netcdf format. The full simulations will be made available on https://archive.sigma2.no ⁴¹⁹ with a specific doi upon acceptance of the manuscript. To retrieve the simulations, one can use ⁴²⁰ the following link (457 MB):

wget -c http://ns9039k.web.sigma2.no/Synchronisation_supermodel.tar.gz

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