

# Toward a Better Understanding of Wildfire Behavior in the Wildland-Urban Interface: A Case Study of the 2021 Marshall Fire

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## Abstract

On 30 December 2021, the Marshall Fire devastated the Boulder, Colorado region. The fire initiated in fine fuels in open space just southeast of Boulder and spread rapidly due to the strong, downslope winds that penetrated into the Boulder Foothills. Despite the increasing occurrence of wildland-urban interface (WUI) disasters, many questions remain about how fires progress through vegetation and the built environment. To help answer these questions for the Marshall Fire, we use a coupled fire-atmosphere model and Doppler on Wheels (DOW) observations to study the fire's progression as well as examine the physical drivers of its spread. Evaluation of the model using the DOW suggests that the model is able to capture general characteristics of the flow field; however, it does not produce as robust of a hydraulic jump as the one observed. Our results highlight limitations of the model that should be addressed for successful WUI simulations.

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2 **the Wildland-Urban Interface: A Case Study of the**  
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10 **Key Points:**

- 11 • Complex meso- and micro-scale meteorology, along with fire ember spotting, were  
12 responsible for rapid spread of the Marshall Fire  
13 • Radar observations from “Doppler on Wheels” elucidates three-dimensional flow  
14 structures that impact fire and plume evolution  
15 • Initial fire propagation in dry, fine fuels is well-represented by the coupled WRF-  
16 Fire model, but urban spread remains a challenge

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## Abstract

On 30 December 2021, the Marshall Fire devastated the Boulder, Colorado region. The fire initiated in fine fuels in open space just southeast of Boulder and spread rapidly due to the strong, downslope winds that penetrated into the Boulder Foothills. Despite the increasing occurrence of wildland-urban interface (WUI) disasters, many questions remain about how fires progress through vegetation and the built environment. To help answer these questions for the Marshall Fire, we use a coupled fire-atmosphere model and Doppler on Wheels (DOW) observations to study the fire’s progression as well as examine the physical drivers of its spread. Evaluation of the model using the DOW suggests that the model is able to capture general characteristics of the flow field; however, it does not produce as robust of a hydraulic jump as the one observed. Our results highlight limitations of the model that should be addressed for successful WUI simulations.

## Plain Language Summary

Wildland-urban interface, or WUI, fires are increasing in the United States and around the world as the built environment continues to expand into the wildland. To better inform real-time management of active wildfires, it is critical that the scientific community can better predict WUI fire spread. In this study, we rely on multiple observational platforms, including the “Doppler on Wheels” radar, to investigate the performance of a state-of-the-art, coupled fire-atmosphere model during the Marshall Fire, which was a recent WUI fire that occurred in Colorado. While the modeling system performs well during the fire’s initial propagation in fine fuels, it is unable to accurately predict spread in the built environment. While mesoscale to microscale simulations can accurately represent atmospheric flow features, more reliable predictability of wildfire behavior in the WUI will require consideration of urban fuels and fire ember spotting.

## 1 Introduction

Wildfire activity in the United States (U.S.) and across the globe has increased markedly over the last several decades (e.g., Westerling et al., 2006; Dennison et al., 2014; Balch et al., 2017; Iglesias et al., 2022). In the U.S., many of the recent wildfire seasons have involved long and intense burning periods, leading to the loss of life and property, as well as poor air quality (e.g., Buchholz et al., 2022) and reductions in solar energy production as a result of smoke generation (e.g., T. W. Juliano et al., 2022). Global climate models suggest that the recent trend of more large-scale fire events will continue and even increase in the future (e.g., Yue et al., 2013; Yoon et al., 2015; Abatzoglou & Williams, 2016). As part of the complexity, the so-called wildland-urban interface (WUI) is rapidly growing (e.g., Radeloff et al., 2018; Burke et al., 2021) and, therefore, amplifying the direct threat of wildfires on daily human activities.

Several notable WUI fires have occurred over the past decade in the U.S., primarily in California, including the Tubbs Fire (2017; e.g., Coen et al., 2018; Mass & Ovens, 2019), Camp Fire (2018; e.g., Brewer & Clements, 2019; Mass & Ovens, 2021), Woolsey Fire (2018; e.g., Keeley & Syphard, 2019), and Thomas Fire (2017; e.g., Fovell & Gallagher, 2018). One commonality between these wildfires is that they were considered wind-driven fires, allowing them to expand rapidly and wreak havoc on communities in their paths. In California, so-called “Santa Ana” (e.g., Randles et al., 2003, and references therein) or “Diablo” (e.g., Liu et al., 2021) wind events are often associated with destructive wind-driven wildfires (e.g., Nauslar et al., 2018; Smith et al., 2018). Wind-driven WUI fires are also a concern in regions outside of the U.S., including Australia (e.g., Cruz et al., 2012), France (e.g., Ganteaume, 2020), and Greece (e.g., Efthimiou et al., 2020).

The Marshall Fire is an example of a catastrophic wind-driven, WUI fire event that occurred just outside of Boulder, Colorado, U.S. on 30 December 2021 (Fovell et al., 2022),

causing two deaths and destroying more than 1,000 buildings, leading to over \$500 M in damages. The fire ignited near the Marshall Mesa during strong wind conditions (wind gusts over  $40 \text{ m s}^{-1}$ ), and it began spreading rapidly in dry, fine fuels driven by intense, westerly winds. Approximately one hour after ignition, the fire transitioned into an urban conflagration, including “hopping” a six-lane interstate, Highway-36, via ember spotting. The large-scale meteorological setup favored a downslope windstorm along the Front Range (Fovell et al., 2022), which is a relatively common occurrence in this geographical region during the cold season (e.g., Whiteman & Whiteman, 1974; Durran, 1990).

In this study, we use observations and numerical simulations to examine the impact of the meso- and micro-scale meteorology on the Marshall Fire behavior. Specifically, we use a state-of-the-art numerical framework, the Weather Research and Forecasting (WRF) model with a fire behavior model (WRF-Fire), as well as measurements from the Doppler on Wheels (DOW) radar system, to address the following fundamental questions related to the topic of wildfire-weather: (1) *What were the observed and modeled atmospheric flow characteristics during the Marshall Fire?* and (2) *How well does the WRF-Fire model reproduce the Marshall Fire spread in the WUI?*

## 2 Data and Methods

### 2.1 WRF-Fire Model

The WRF model is a numerical weather prediction system used widely by research and operational forecasting communities alike (Skamarock et al., 2019). WRF has proven to be a powerful tool for simulating the full range of atmospheric scales, including meso- and micro-scales (e.g., Mazzaro et al., 2017; Muñoz-Esparza et al., 2017; Rai et al., 2019). Here, we utilize WRF in a one-way nested, mesoscale to microscale configuration (e.g., Haupt et al., 2019) whereby the inner domain is turbulence-resolving. The WRF domains are positioned to capture the westerly inflow that plays an important role on the wildfire propagation (Fig. S1). Our model setup and physics options closely follow those used in a recent study of the East Troublesome Fire by our team (DeCastro et al., 2022), and additional details may be found in Text S1.

To examine the Marshall Fire evolution, we conduct WRF simulations with a fire behavior model based on the Coupled Atmosphere-Wildland Fire Environment (Clark et al., 2004; Coen, 2013). This coupled fire-weather model is called WRF-Fire (Coen et al., 2013; Shamsaei et al., 2022). In addition to the meteorological grid that is defined in a standard WRF simulation, a fire grid is also required for a WRF-Fire simulation. The fire grid is refined by a factor of four relative to the meteorological grid ( $\Delta x = \Delta y = 27.78 \text{ m}$ ) to track the evolution of the fire perimeter via an improved level-set method (Muñoz-Esparza et al., 2018) and compute small-scale changes in the fuel properties. The fire mesh is assigned fuel properties according to the Anderson 13 class fuel model (Anderson, 1981, Fig. S2) and topography according to 3-arc second Shuttle Radar Topography Mission terrain data. Using the parameterization developed by Rothermel (1972), the near-surface atmospheric winds, fuel characteristics, and terrain slope dictate the fire rate of spread. After a fire is ignited in the model, the fuel burn rate is calculated according to Albini and Reinhardt (1995). The amount of heat and moisture released back into the atmosphere is computed as a function of fuel properties, allowing for full coupling between the fire and atmosphere. Text S2 provides discussion on the model ignition times and locations, including the importance of ember spotting.

### 2.2 DOW Measurements

The DOW platform was deployed during the Marshall Fire to capture the three-dimensional smoke/ash plume and flow structures. Operating at 3-cm wavelength, the DOW is a mobile/quick deployable Doppler radar with high spatial resolution ( $50 \times 160 \times 160 \text{ m}$

115 at 10 km range), allowing it to measure microscale structures (Wurman et al., 1997, 2021).  
 116 During this deployment, the DOW operated mostly in a rastered Plan Position Indica-  
 117 tor (PPI) scanning mode, with elevation scans ranging from  $\sim 0.5 - 23$  degrees (adjusted  
 118 throughout the deployment) above the horizon. Parameters derived from the DOW ob-  
 119 servations and relevant to this study include reflectivity, radial velocity, and spectrum  
 120 width. Reflectivity scans from the DOW provide information about the number and size  
 121 of scatterers [fire-generated debris, or pyrometeors (McCarthy et al., 2019)] present in  
 122 a retrieval volume, and, therefore, they may be used as a proxy for understanding the  
 123 intensity of combustion and relative concentration of pyrometeors. The radial velocity  
 124 product reveals flow moving away and toward the radar location (along the radar beam’s  
 125 path), while the spectrum width field indicates the variability of scatterer velocities in  
 126 the retrieval volume and, therefore, it may be used as a proxy for turbulence levels.

### 127 **3 Fire Spread in the WUI**

128 The Marshall Fire had two reported initial ignition points, occurring at 18:08 and  
 129 19:00 UTC and approximately several 100s of meters apart (see Text S2 for more infor-  
 130 mation). Thus, in our WRF-Fire simulations, we first ignite two separate fires. Both ig-  
 131 nition points were in dry, short grass fuels. During the early stages of the fire, the com-  
 132 bination of fuels and strong ( $\sim 25 \text{ m s}^{-1}$ ), westerly winds supported rapid fire growth in  
 133 the Marshall Mesa area (magenta star in Fig. 1). At 19:00 UTC, the initial burn region  
 134 in the model takes on a finger-like structure, with spotting on the east and southern flanks  
 135 as it approaches Highway 36 (Fig. 1).

136 According to Visible Infrared Imaging Radiometer Suite (VIIRS) fire detections  
 137 at 19:25 UTC, the fire had spotted across Highway-36 to cause secondary ignitions (Fig. S3),  
 138 but the simulated fire does not cross the highway via spotting until 19:45 UTC (cf. Fig. 1),  
 139 at which time another burning lobe to the south originating from the second ignition has  
 140 nearly reached the interstate. Two snapshots from VIIRS shortly after, at 20:15 UTC  
 141 and 21:00 UTC (Fig. S3), show that the modeled leading edge is too slow and the north-  
 142 south expansion is too narrow. We note that during this  $\sim 1.5$  h period, the intense west-  
 143 erly winds continue across much the region; however, the model shows weaker wester-  
 144 lies and even “reversed” flow intruding (further discussed in Section 4).

145 By mid-afternoon (22:05 UTC), the rapid fire spread is halted in the model as it  
 146 reaches the urban region (Fig. 1), where non-burnable fuels are present in the model fuel  
 147 layer. The westerly, low-level flow is confined to only the western portion of the area, with  
 148 winds near the fire front opposing the original fire spread direction. This flow transition  
 149 will be discussed further in Section 5. Nonetheless, around this time, the DOW reflec-  
 150 tivity isosurfaces show active plume cores in two of the main fingers further north and  
 151 east (Fig. 2a), confirming that the model is not able to capture the rapid and consequen-  
 152 tial propagation across the Highway-36.

153 Between 22:05 and 23:00 UTC, the model shows generally slower fire spread com-  
 154 pared to previous hours, as it expands the burned region mostly to the north and south  
 155 due to the relatively weak, variable winds (Fig. 1). During this time, and over the next  
 156 couple of hours, the radar reflectivity isosurfaces indicate that the fire becomes increas-  
 157 ingly active in the middle finger (Fig. 2b,c) before dissipating, while a new southern fin-  
 158 ger becomes more active (Fig. 2d). Only by the evening (02:30 UTC) does the simulated  
 159 burn area finally spread into Louisville on the north side of Highway-36 and toward the  
 160 southernmost observed finger (Fig. 1). In Section 5, we will discuss potential sources of  
 161 error in the WRF-Fire simulations.

## 4 Horizontally Heterogeneous Wind Field

The synoptic-scale and mesoscale meteorology during the Marshall Fire event fostered intense downslope winds along the Front Range (Fovell et al., 2022). A north-south band of strong, westerly flow (gusts  $>30 \text{ m s}^{-1}$ ) was positioned along the Boulder Foothills, where the plunging, downslope flow remained attached at the surface, including in the vicinity of the Marshall Fire ignitions. In contrast, many locations not too far to the east generally experienced weaker winds (gusts  $<20 \text{ m s}^{-1}$ ) where the flow detached from the surface, as shown in Fig. 1 and in agreement with Fovell et al. (2022, their Fig. 1). To evaluate the WRF-Fire model’s ability to represent the spatially variable, low-level flow during the event, in Fig. S4 we compare observed and modeled time series of the surface stations shown in Fig. S1 and described in Text S3. By and large, the model performs better at the western stations, where strong, westerly winds persisted before the flow transitioned. Even still, WRF tends to underestimate the strongest wind gusts of  $40\text{-}50 \text{ m s}^{-1}$  that were observed at CO1 and BLD. This result supports the aforementioned underestimation in model fire spread. Compared to the western area, both observations and WRF show much more variable wind speeds and directions toward the east of the fire.

The horizontal structure and variability in the wind field is captured by the DOW radial velocity observations. Figure 3a shows the time-mean radial velocity for scans below 5 degrees, revealing (i) the strong west-southwest winds across the fire, (ii) a region of reversed flow, especially over the southern portions of the fire area, and (iii) a subsequent return to west winds aloft and to the east. In Fig. 3b, we also show the fraction of the time the radial wind is positive. These data show that within the time-mean reversed flow regions, many locations experience positive winds  $\sim 50\%$  of the time, suggesting that the winds were substantially variable. As we will discuss in the next section, the flow variability is related to the presence of a hydraulic jump. Shown in both plots are also station observations (colored circles) that indicate the radial wind component and the vector wind during the averaging period (Fig. 3a), as well as the fraction of time with positive winds at each site (Fig. 3b). Overall, the agreement between the radar and near-surface observations is good; however, some differences are expected because the height of the radar retrieval volume increases as the radial distance increases according to the DOW scan angle (not shown).

## 5 Vertical Structure and Flow Evolution

Based on quasi-idealized simulations, Fovell et al. (2022) suggest that a “hydraulic jump-like feature” was present downwind of the strongest winds in the Boulder Foothills. To further explore this aspect of the atmospheric flow, we use model output and DOW observations. In Fig. 4, we present east-west vertical cross-sections of the zonal wind component. Each panel represents a different snapshot in time, with the times corresponding to those shown in Fig. 1. Throughout the event, the low-level, downslope winds upstream of the fire (west of  $\sim 105.3^\circ \text{W}$ ) are consistently strong and capped by a strong inversion where winds diminish quickly with height. This band of intense winds continues eastward, bringing strong westerlies into the Boulder Foothills during the early stages of the Marshall Fire, and rising with height toward the east (Fig. 4). As a result, the atmospheric flow supports the fire’s rapid advancement around 1900 UTC (cf. Fig. 1). Over the ensuing hours, the wind maxima retreats westward, and eventually the well-defined vertical structure breaks down into a more chaotic structure further east (Fig. 4). The strong inversion erodes where the intense winds diminish, as the wavy isentropes (solid green lines) suggest strong vertical mixing within the lower-levels. In the transition zone, a hydraulic jump is evident with a sharp decrease, and even complete reversal, in the zonal winds and vertical displacement of the isentropes.

212 The transition from strong flow in a shallow boundary layer to weaker winds as the  
 213 boundary layer deepens further downwind, with turbulence production in between, are  
 214 classical characteristics of a hydraulic jump (e.g., Ball, 1956; T. W. Juliano et al., 2017).  
 215 To probe the dynamical support for the presence of a hydraulic jump, we conduct a Froude  
 216 Number ( $Fr$ ) analysis along the vertical cross-sections shown in Fig. 4. Results presented  
 217 in Fig. S5 indicate a transition from supercritical ( $Fr > 1$ ) to subcritical flow ( $Fr < 1$ )  
 218 – a well-known requirement for the presence of a hydraulic jump. Upstream  $Fr$  values  
 219 between 2 and 4 (Fig. S5 and Text S4) suggest a hydraulic jump with a roller (e.g., Chan-  
 220 son, 2009), whereby much of the mean kinetic energy is converted into turbulent kinetic  
 221 energy (TKE). In this particular case, the WRF model simulates an extraordinary transi-  
 222 tion, with maximum TKE values exceeding  $200 \text{ m}^2 \text{ s}^{-2}$  due to the strong decay in in-  
 223 tense westerly winds (Fig. S6).

224 The hydraulic jump and subsequent gravity wave structures in Fig. 4 are readily  
 225 apparent in a cross section of the radar reflectivity spanning 2202-2226 UTC (Fig. 5).  
 226 Specifically, the DOW data show a leading wave linked to the fire’s updraft that is em-  
 227 bedded in the hydraulic jump region followed by a subsidence region (i.e., diminishing  
 228 plume heights) and a second wave crest (Fig. 5a). Spectrum width measurements (Fig. S7)  
 229 show maximum values in the primary plume with a secondary maximum in the down-  
 230 stream wave (qualitatively similar to the TKE field from WRF; cf. Fig. S6). The con-  
 231 temporaneous isentropes extracted from WRF suggest that the simulation underestimates  
 232 the injection height of smoke and ash in the leading wave. This discrepancy may be due  
 233 to a lack of urban fuels in the model: the combustion of urban fuels, which have high  
 234 fuel loads and long residence times, may have produced more intense heat release in re-  
 235 ality compared to what was simulated. Nonetheless, the structure of the second wave  
 236 agrees fairly well between observations and simulations. Also shown is the downwind vari-  
 237 ation of the column maximum radar reflectivity (Fig. 5b), which is a measure of plume  
 238 dilution and debris fall out. The maximum reflectivity (uncorrected) is  $\sim 30 \text{ dBZ}$ , with  
 239 a logarithmic decay to the east. The sharpest reduction in reflectivity is close to the main  
 240 updraft, suggesting the potential for ember fall out in this region.

## 241 6 Discussion and Conclusions

242 In this article, we present observations and numerical model simulations of the Mar-  
 243 shall Fire in December 2021, which spread rapidly in the WUI due to strong, westerly  
 244 winds along with dry, fine fuels and ember spotting. Observations from surface stations  
 245 near the fire show that the WRF-Fire model generally underestimates the strongest recorded  
 246 wind speeds, leading to a slightly slower propagation through the wildland fuels west of  
 247 Highway-36. Satellite measurements at the beginning of the event confirm the model’s  
 248 slower spread prior to the fire reaching the towns of Superior and Louisville.

249 We also rely on data from the DOW radar – deployed several hours after the ini-  
 250 tial ignition – to highlight the three-dimensional atmospheric flow structure during the  
 251 wildfire. The radar retrievals illustrate substantial horizontal variability in the low-level  
 252 wind field, in addition to vertical plume structure embedded in a robust hydraulic jump.  
 253 Turbulence-resolving output from WRF-Fire suggests that the highly variable, low-level  
 254 winds are related to the hydraulic jump. In this jump region, the flow transitions from  
 255 intense westerlies to much weaker westerlies or even a shift to easterlies, ultimately af-  
 256 fecting the Marshall Fire spread rate and direction.

257 Even though the model produces generally encouraging results, there are two main  
 258 shortcomings related to the fire module in WRF that should be discussed. First, while  
 259 the most up-to-date version of WRF-Fire as of this writing (version 4.4) contains a fire-  
 260 brand parameterization, it does not ignite spot fires, but rather provides only a likeli-  
 261 hood of spot fire ignition. Rapid wildfire spread is often caused by embers generating  
 262 new ignitions ahead of the main fire front (e.g., N. P. Lareau et al., 2022). The Marshall

263 Fire was able to cross Highway-36, which is a six-lane interstate. Such advancement is  
 264 possible only through ember spotting, and, therefore, a WRF-Fire simulation – without  
 265 additional manual ignitions such as in this study – is not able to produce further fire spread.

266 Second, the WRF-Fire model must be improved to account for WUI fuels and re-  
 267 lated fire propagation in the built environment. During post-fire investigations of the Mar-  
 268 shall Fire, the Institute for Business and Homes Safety found evidence that wooden fences  
 269 falling between homes in Superior and Louisville were a primary cause of fire spread (Reppenhagen,  
 270 2022). At present, the WRF-Fire model contains fuel categories (based on Anderson 13)  
 271 strictly intended for fires in the wildland and a rate of spread parameterization (based  
 272 on Rothermel) developed without considering fire-atmosphere coupling. However, given  
 273 the increasing trend in WUI fire occurrence, fuel maps including WUI materials, as well  
 274 as improved representation of fire spread, should be developed for coupled fire-atmosphere  
 275 models. The WUI challenge highlights the urgent need to better understand the com-  
 276 plex interactions between humans and the built environment, weather, and wildfire, and  
 277 ultimately develop more effective solutions to predict wildfire behavior and risk.

## 278 7 Open Research

### 279 7.1 Data Availability Statement

280 Surface weather station data and model outputs (<https://doi.org/10.7910/DVN/M6WCBT>; T. Juliano et al., 2022), as well as VIIRS fire detections (<https://doi.org/10.7910/DVN/PR6XDM>; T. Juliano & Shamsaei, 2022), used in this study are stored on  
 281 Harvard Dataverse and will be finalized prior to publication. DOW measurements are  
 282 available at (<https://doi.org/10.48514/JKJO-TE44>; Wurman & Kosiba, 2022). Users  
 283 should contact Josh Wurman ([jwurman@illinois.edu](mailto:jwurman@illinois.edu)) or Karen Kosiba ([kakosiba@illinois.edu](mailto:kakosiba@illinois.edu))  
 284 to gain access to the data.  
 285  
 286

### 287 7.2 Software Availability Statement

288 The WRF v4.4 code used for the simulations is publicly available on Github (<https://github.com/wrf-model/WRF/tree/release-v4.4>). Codes for the model (<https://doi.org/10.7910/DVN/M6WCBT>; T. Juliano et al., 2022) and DOW (<https://doi.org/10.7910/DVN/KYSLUE>; N. Lareau, 2022) analyses are available on Harvard Dataverse and  
 289 will be finalized prior to publication.  
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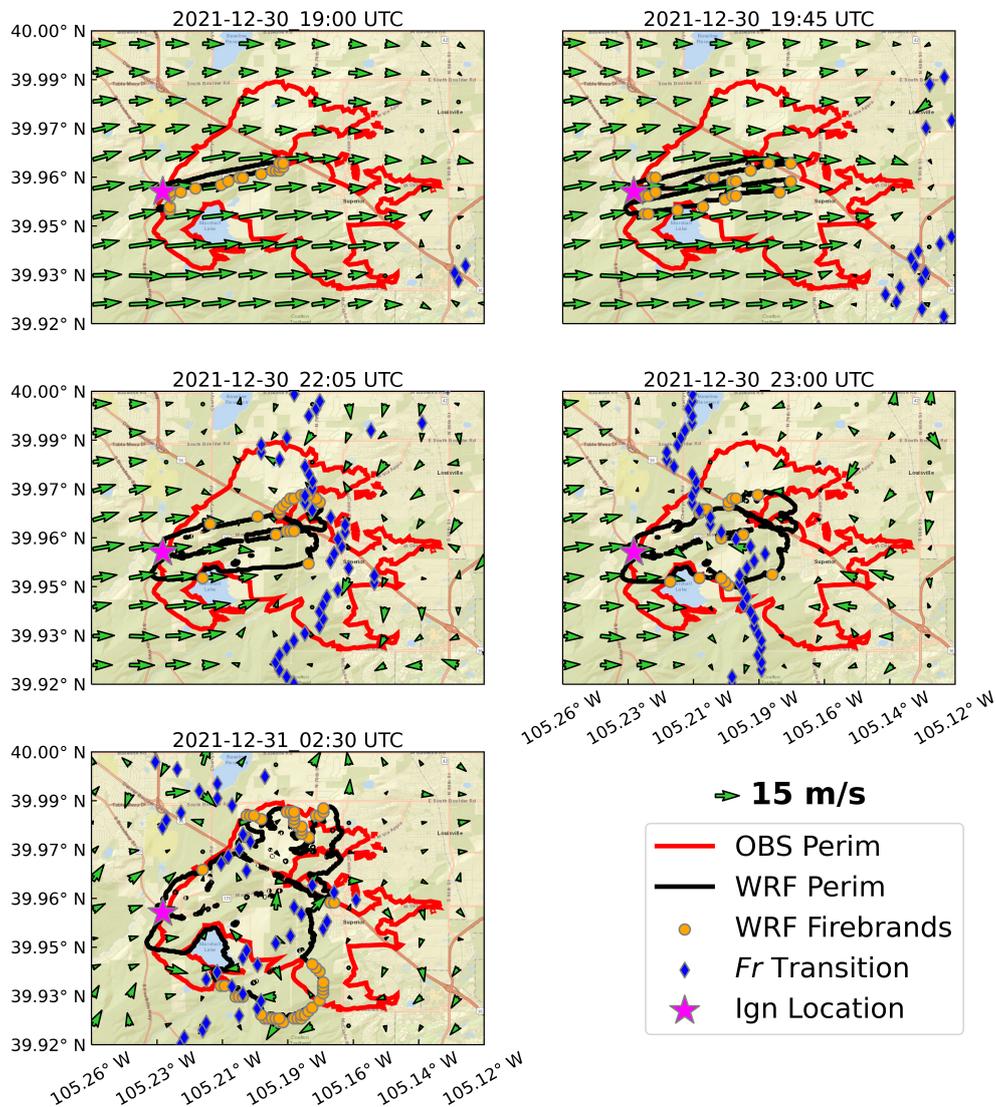
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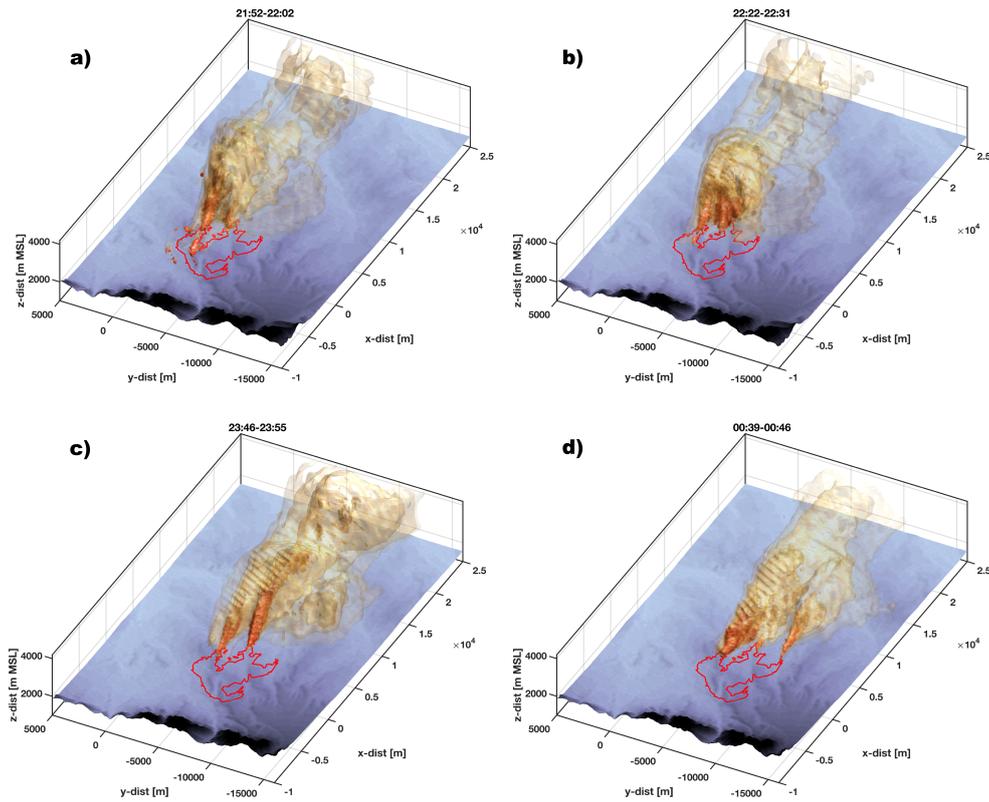
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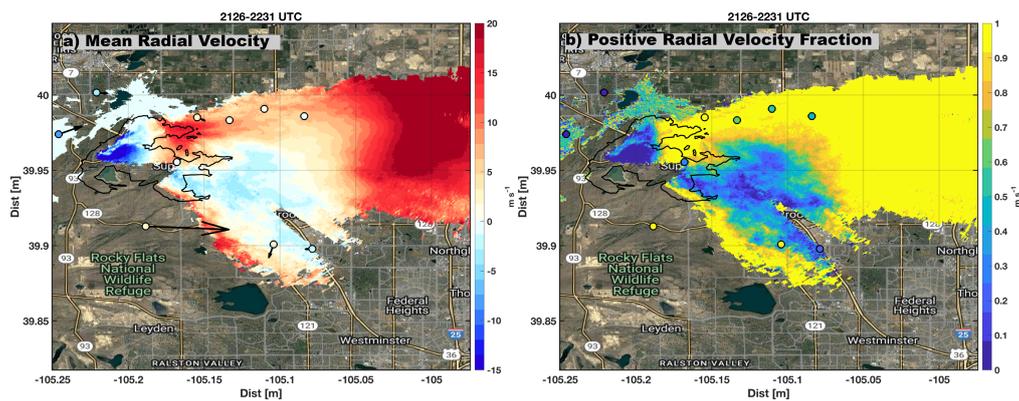
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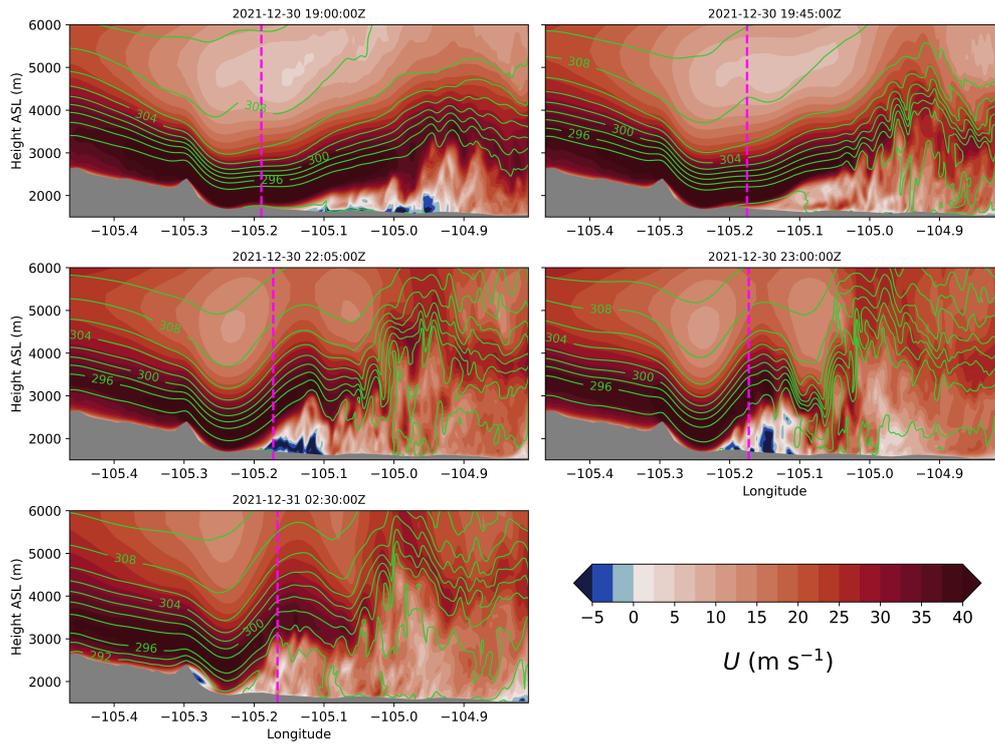
**Figure 1.** Temporal progression of the Marshall Fire spread. The magenta star represents the approximate location of the initial fire ignitions, while the red line represents the final observed perimeter and the black line represents the WRF-Fire perimeter at the indicated time. The orange circles represent firebrand landing locations according to WRF-Fire. Flow transitions from supercritical to subcritical are shown by the blue diamonds. Also shown are 10 m wind arrows according to the key.



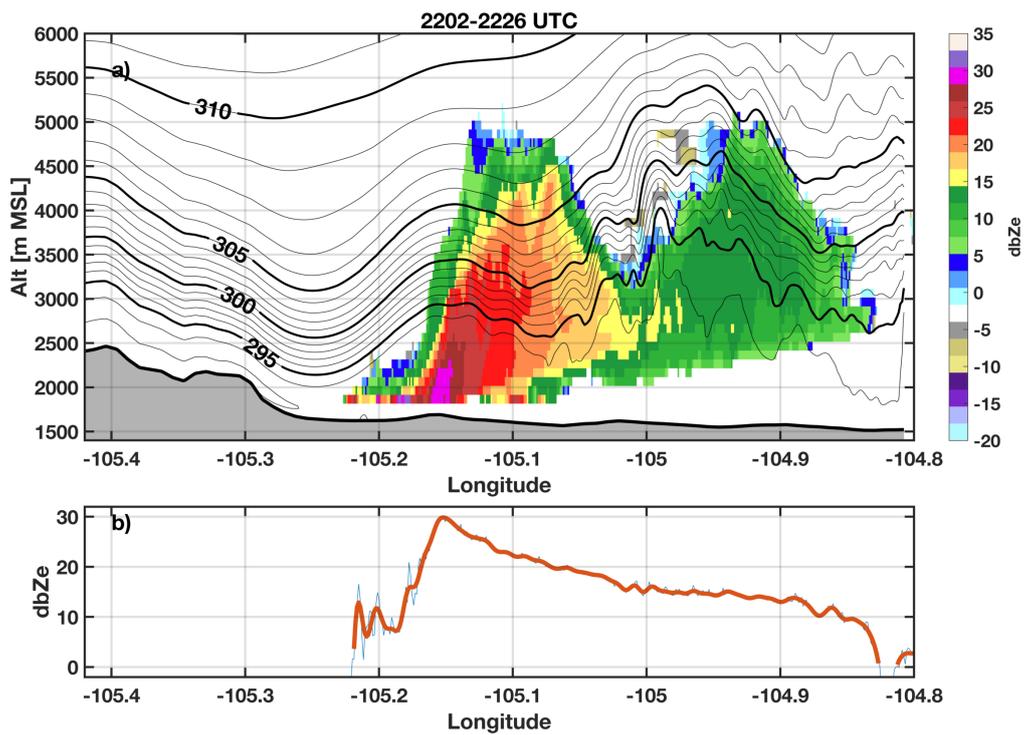
**Figure 2.** Radar reflectivity isosurfaces showing plume evolution. Transparent Isosurfaces are rendered at -10, 10, 15, 20, 23, 26, and 27 dbZ with colors becoming increasingly red for higher values. The data window (UTC), is shown at the top of each panel. Also shown are the IR fire perimeter (red contour) and terrain elevation (gray shaded relief).



**Figure 3.** (a) Time-mean radial velocity data with station observations showing mean wind vectors and mean radial velocity (color shaded). (b) Fraction of the time with a positive radial wind component. Both figures also show the final observed perimeter (solid black contour).



**Figure 4.** East-west vertical transects showing the  $U$ -component of the wind speed according to the colorbar, along with isentropes (potential temperature contours) every 2 K in green. The vertical magenta line represents the furthest eastward progression of the fire front in the whole domain. Gray shading represents the terrain profile. Times shown are the same as in Fig. 1.



**Figure 5.** (a) Time and meridional maximum radar reflectivity cross section for the 2202-2226 UTC interval. Reflectivity values are shaded, with potential temperature contours from WRF (contours every 1 K, bold and labeled every 5 K). Gray shading represents the terrain profile. (b) Column maximum radar reflectivity as a function of longitude.

# Supporting Information for ”Toward a Better Understanding of Wildfire Behavior in the Wildland-Urban Interface: A Case Study of the 2021 Marshall Fire”

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## Contents of this file

1. Text S1-S4
2. Figures S1 to S7

**S1: WRF model setup** In our two domain configuration, the outer and inner domains are resolved using a horizontal grid cell spacing of  $\Delta x = \Delta y = 1000$  m and 111.11 m, respectively. We use a total of 45 model grid cells in the vertical column and set the grid cell spacing,  $\Delta z \approx 25$  m adjacent to the surface before stretching with increasing height. In the outer domain, we activate the Yonsei University (YSU) planetary boundary layer

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(PBL) parameterization (Hong et al., 2006) to handle vertical turbulent mixing, and 2D horizontal diffusion is computed by Smagorinsky (1963). In the inner domain, we turn off the PBL scheme and use WRF’s large-eddy simulation capability by activating the three-dimensional turbulent kinetic energy (TKE)-based sub-grid scale (SGS) model of Deardorff (1980). Details about the resolved TKE calculation is described in Text S4.

**S2: Marshall Fire ignition approach** The first ignition, initially reported by a 911 call at 18:08 UTC, was slightly northeast of the intersection of Marshall Road and Highway-93 (approximate location:  $39.956127^{\circ}\text{N}$ ,  $105.230487^{\circ}\text{W}$ ). The 911 call indicated that a structure was burning uncontrollably in the intense winds. At 19:00 UTC, a Boulder park ranger on the scene noticed a second ignition location to the southwest, near the Marshall Mesa Trailhead parking lot (approximate location:  $39.951209^{\circ}\text{N}$ ,  $105.231441^{\circ}\text{W}$ ). This ignition occurred in dry, fine fuels and began spreading rapidly toward the northeast. In the WRF-Fire simulations, we ignite these two sources according to the above information.

Other ignitions also occurred later during the Marshall Fire likely due to ember spotting. Without any conclusive information, we hypothesize that ember spotting is the only plausible way that the fire was able to “hop” across Highway-36. Therefore, we use the firebrand spotting parameterization (Frediani et al., 2021) in WRF-Fire to produce spotting likelihood maps. We first run a simulation with only the two aforementioned primary ignitions and allow the fire to approach Highway-36. At this stage, the fire is not able to propagate further because any roadways in the fuel model are considered non-burnable fuel. However, we use the spotting parameterization to estimate where and when new ignitions are most likely to occur. While this approach is somewhat subjective, we believe

that it is reasonable, especially because the modeled fire spread appears realistic in the fine fuels. Technical details about the WRF-Fire spotting parameterization may be found in Appendix A of the WRF User Guide starting in v4.4 ([https://www2.mmm.ucar.edu/wrf/users/docs/user\\_guide\\_v4/v4.4/users\\_guide\\_chap-fire.html#firebrand](https://www2.mmm.ucar.edu/wrf/users/docs/user_guide_v4/v4.4/users_guide_chap-fire.html#firebrand)).

Two additional fire ignition locations and times, based on the spotting parameterization, are estimated: (39.9638°N, 105.180°W) at 19:40 UTC and (39.966248°N, 105.185050°W) at 20:30 UTC. We then conduct a second WRF-Fire simulation with these additional spotting ignitions and present the results in the main manuscript.

**S3:  $Fr$  calculation** The dimensionless  $Fr$ , which is often used to examine atmospheric flow adjustment, may be interpreted as the ratio of the mean planetary boundary layer (PBL) wind speed to the fastest possible gravity wave traveling along the fluid interface between the PBL and free troposphere. For the  $Fr$  analysis, we follow a similar approach as in (Juliano et al., 2017). Here the PBL height is determined based on the sharpest vertical gradient in the potential temperature field. Reduced gravity is computed as  $g' = g \frac{\theta_t - \overline{\theta_{PBL}}}{\overline{\theta_{PBL}}}$  where  $\theta_t$  is the free troposphere potential temperature,  $\overline{\theta_{PBL}}$  is the mean PBL potential temperature (and so  $\theta_t - \overline{\theta_{PBL}}$  represents the change in  $\theta$  across the PBL inversion). For  $\theta_t$ , we use the  $\theta$  value from two grid cells above the defined PBL top. The  $Fr$  is then calculated as  $Fr = \frac{V}{\sqrt{g'H}}$  where  $H$  is the PBL depth and  $V$  is the mean PBL wind speed.

**S4: TKE calculation** In this study, we use the TKE field to better understand the presence of a hydraulic jump (i.e., supercritical to subcritical flow transition). To compute the resolved TKE from the WRF output, we follow steps outlined in previous studies (e.g.,

Schmidli, 2013; Wagner et al., 2014). First, we decompose a fully turbulent variable,  $\phi$ , into its model grid cell average (represented by the instantaneous model output),  $\bar{\phi}$ , and unresolved,  $\phi'$ , components:

$$\phi(x, y, z, t) = \bar{\phi}(x, y, z, t) + \phi'(x, y, z, t)$$

where  $(x, y, z, t)$  represents the space and time dimensions. The unresolved component is computed by the large-eddy simulation SGS model and output at each  $(x, y, z)$  grid cell and time stamp (1 min sampling interval). The resolved turbulent component,  $\phi''$ , is then calculated as

$$\phi''(x, y, z, t) = \bar{\phi}(x, y, z, t) - \langle \bar{\phi}(x, y, z, t) \rangle$$

where  $\langle \rangle$  represents a temporal average defined as

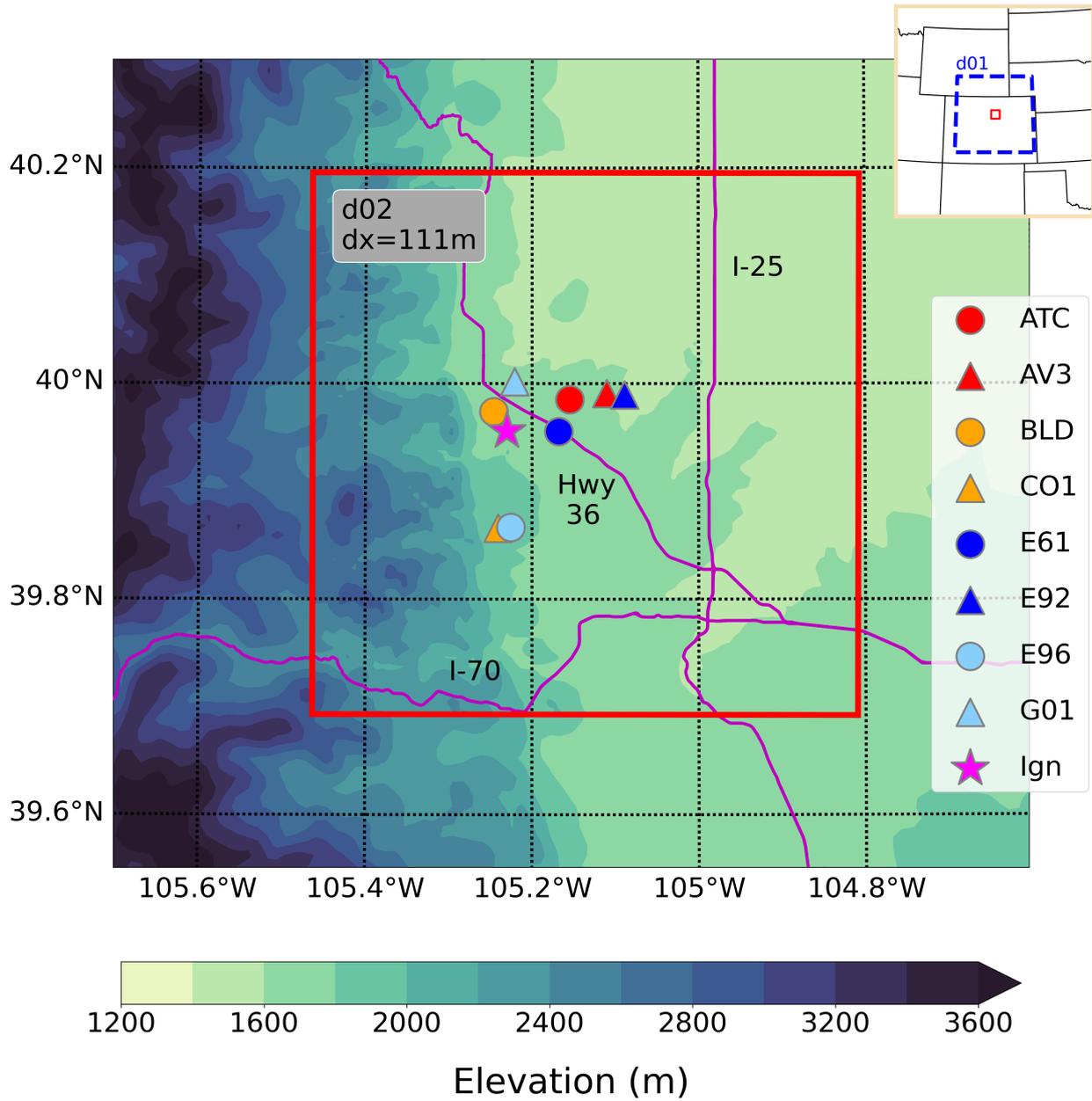
$$\langle \rangle = \frac{1}{T} \int_{t-T/2}^{t+T/2} \bar{\phi}(x, y, z, t) dt$$

with  $T$  representing the time-averaging interval. We choose  $T = 40$  min with a 1 min sampling interval, which is similar to other studies (e.g., Juliano et al., 2017, 2022). Results are relatively insensitive when using  $T = 20$  min with a 1 min sampling interval (not shown).

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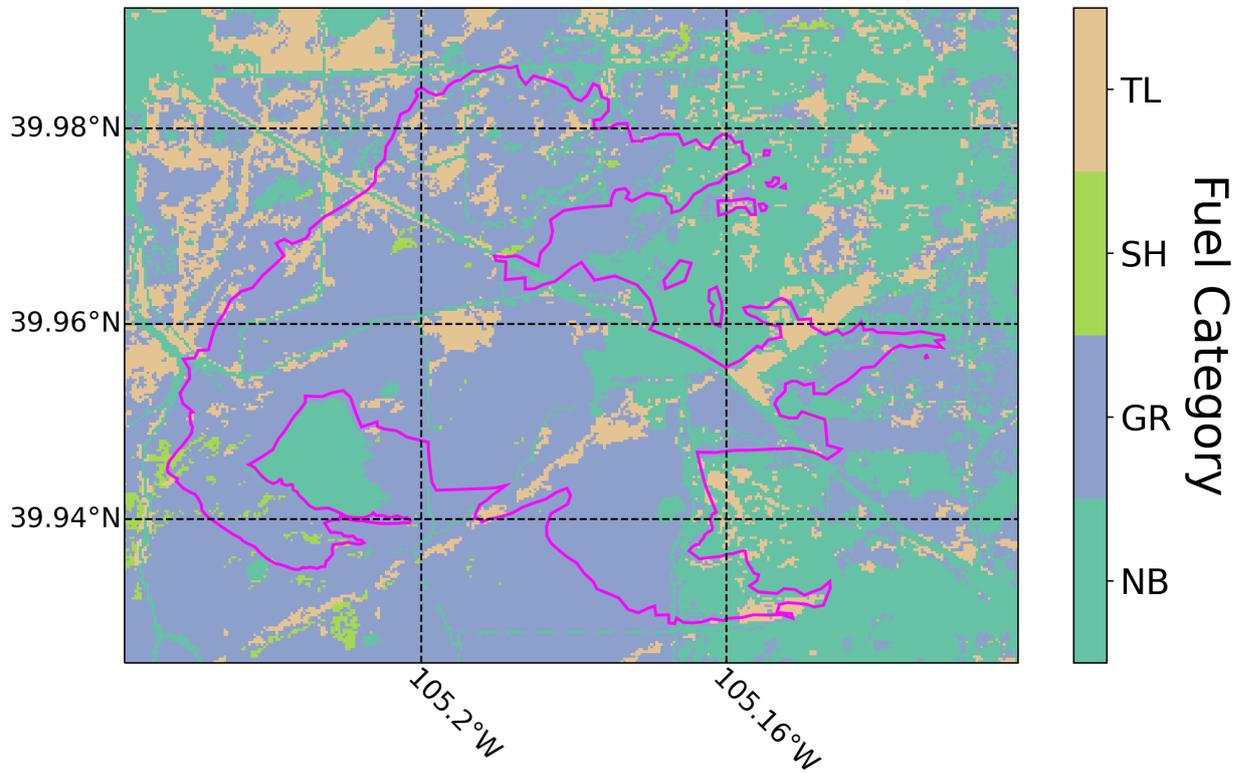
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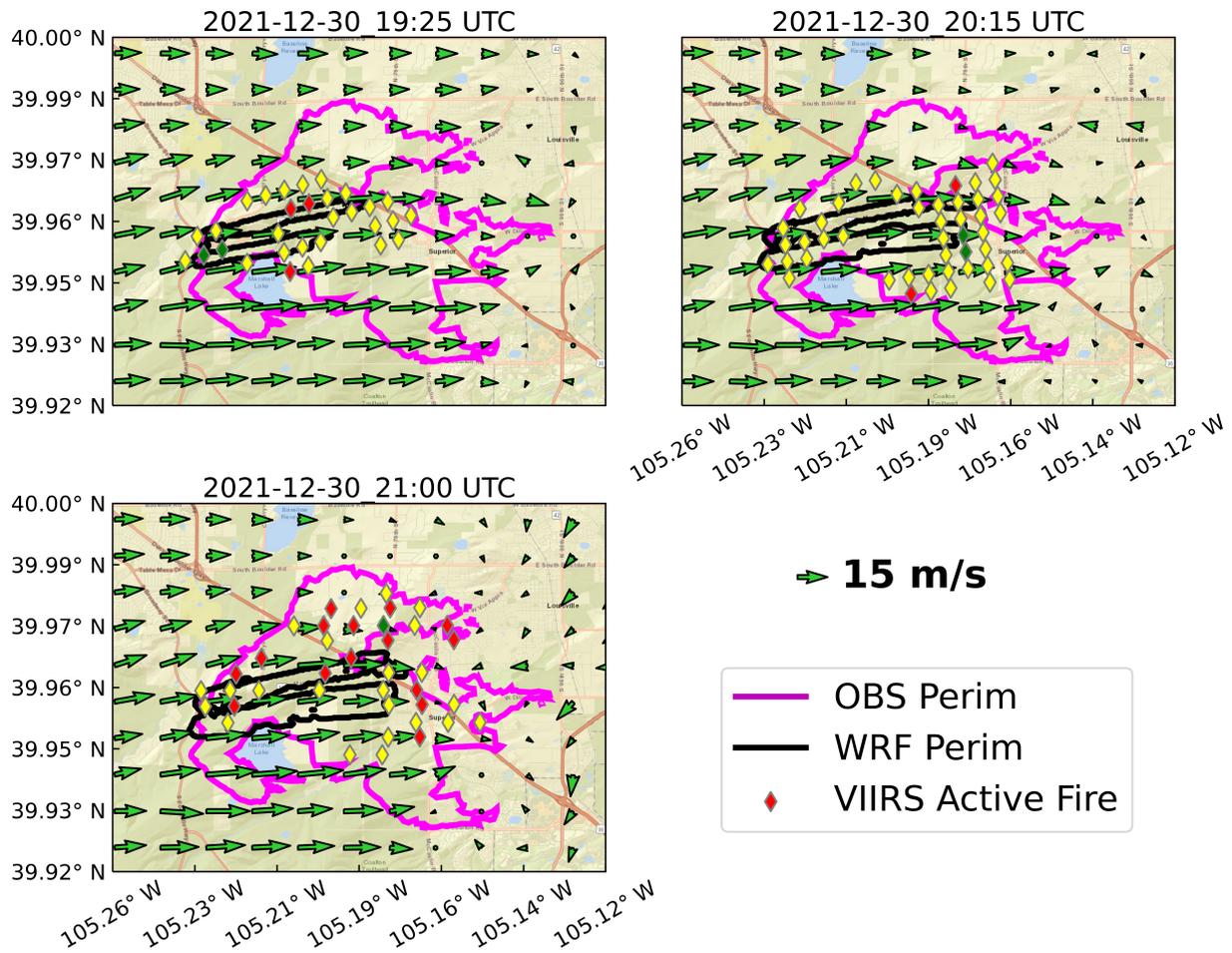


**Figure S1.** WRF-Fire domain configuration. The outer and inner domains are abbreviated as d01 and d02, respectively. The colored red, orange, and blue symbols represent the various surface weather stations. The magenta star represents the approximate locations of the two initial fire ignitions, which are ~550 m apart.

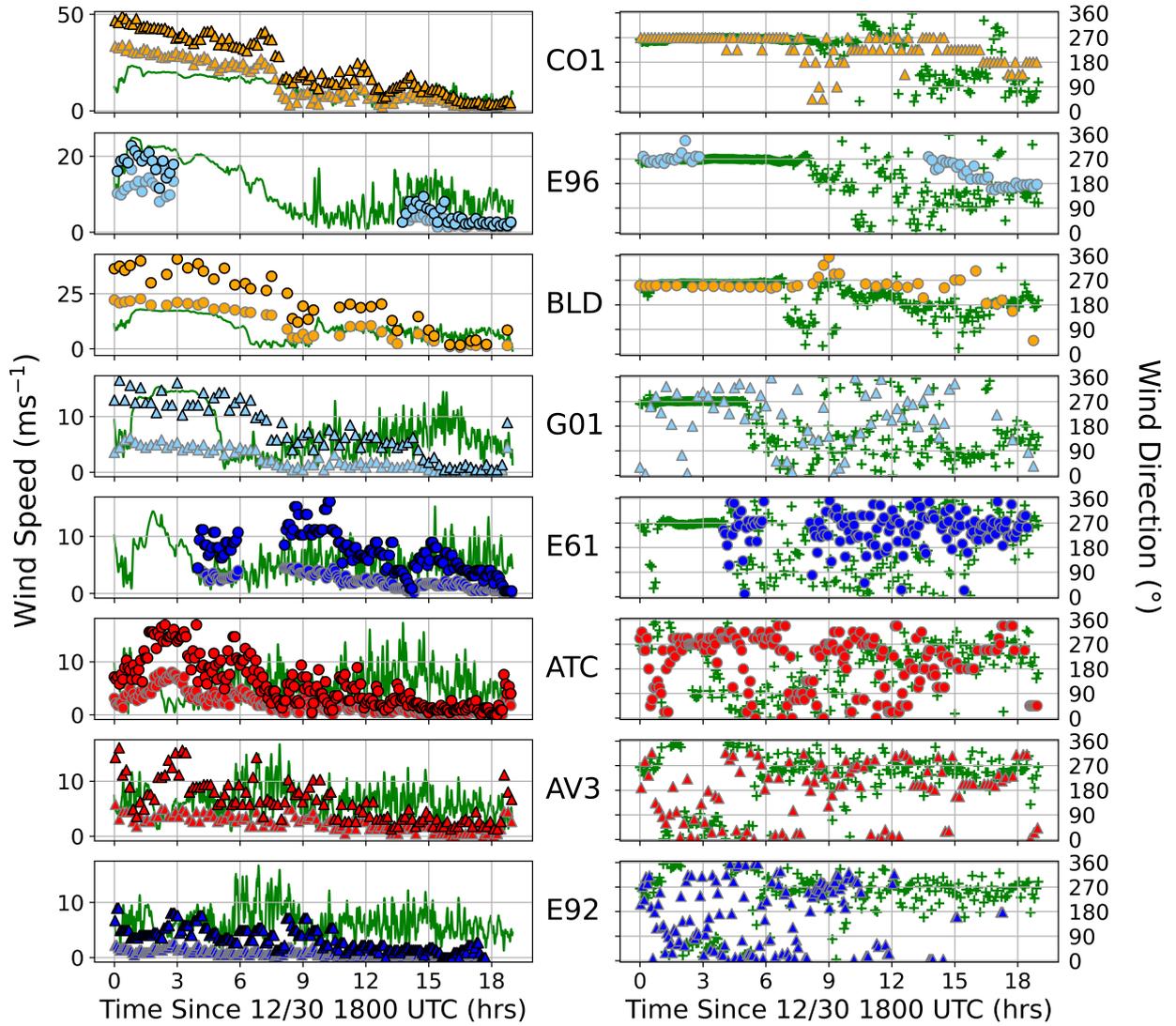
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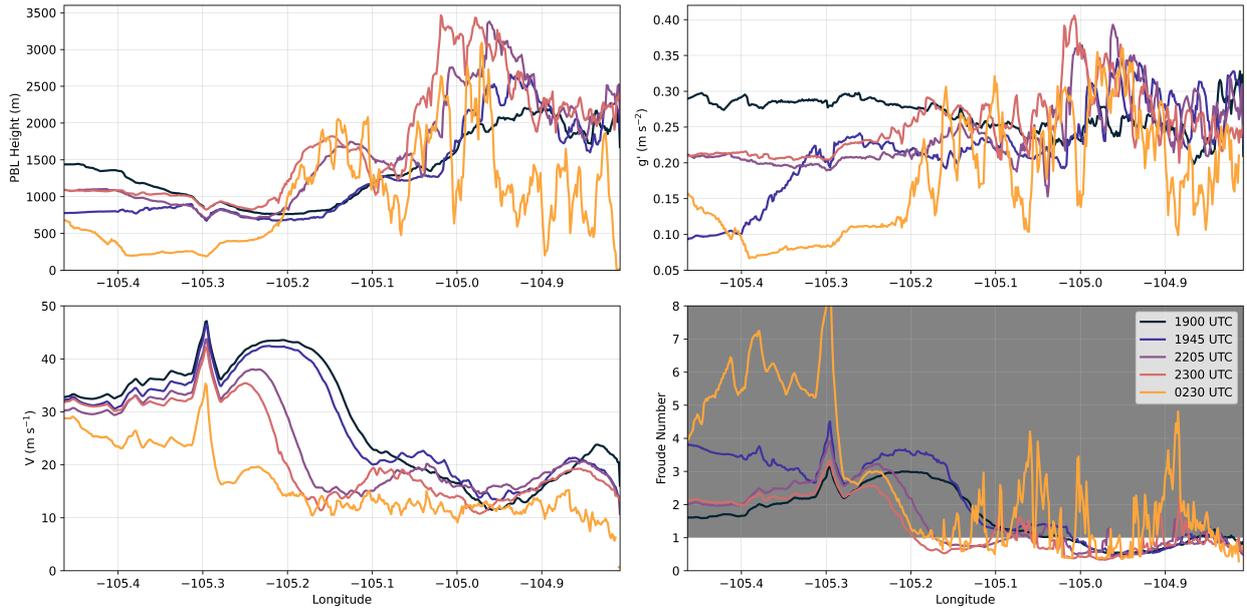
**Figure S2.** The Anderson 13 fuel model layers used in the WRF-Fire simulation. The final observed fire perimeter is shown in magenta. Fuel categories represent the fuel type of the fuel model; timber litter (TL), shrub (SH), grass (GR), and nonburnable (NB). Note that the slash-blowdown category is not shown due to the absence of this fuel type in our domain.



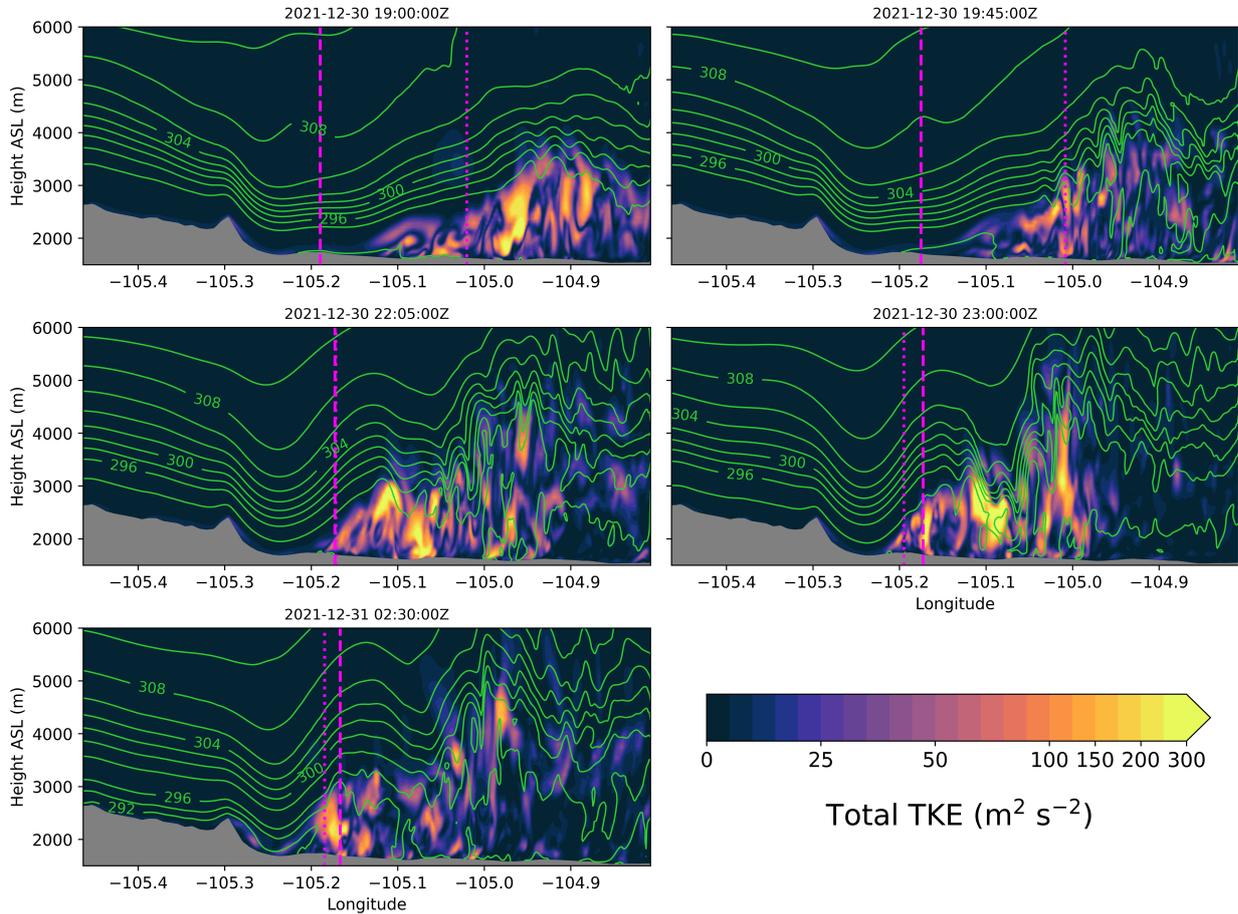
**Figure S3.** Temporal progression of the Marshall Fire spread. The magenta line represents the final observed perimeter and the black line represents the WRF-Fire perimeter at the indicated time. The colored diamonds represent VIIRS fire detections, with red, yellow, and green representing low, nominal, and high confidence intervals, respectively. Also shown are 10 m wind arrows with wind speed according to the key.



**Figure S4.** Time series of (left panels) wind speed and (right panels) wind direction for the various surface weather station locations. Green lines represent high-frequency ( $\frac{2}{3}$  s) output from WRF, while the symbols represent observations. For the left panels, the light and dark colored symbols show the observed wind speeds and gusts, respectively.



**Figure S5.**  $Fr$  analysis along the east-west transects in Fig. 4 showing: (top left) PBL height, (top right) reduced gravity,  $g'$ , (bottom left) mean PBL wind speed, and (bottom right)  $Fr$ . The shaded region in the bottom right panel represents the supercritical flow region ( $Fr > 1$ ). Details about the  $Fr$  calculation are outlined in Text S3.



**Figure S6.** As in Fig. 4, except showing total (resolved plus SGS) TKE according to the colorbar. Also plotted is the location of the flow transition ( $Fr = 1$ ; dotted magenta line) tr Details about the TKE calculation are outlined in Text S4.

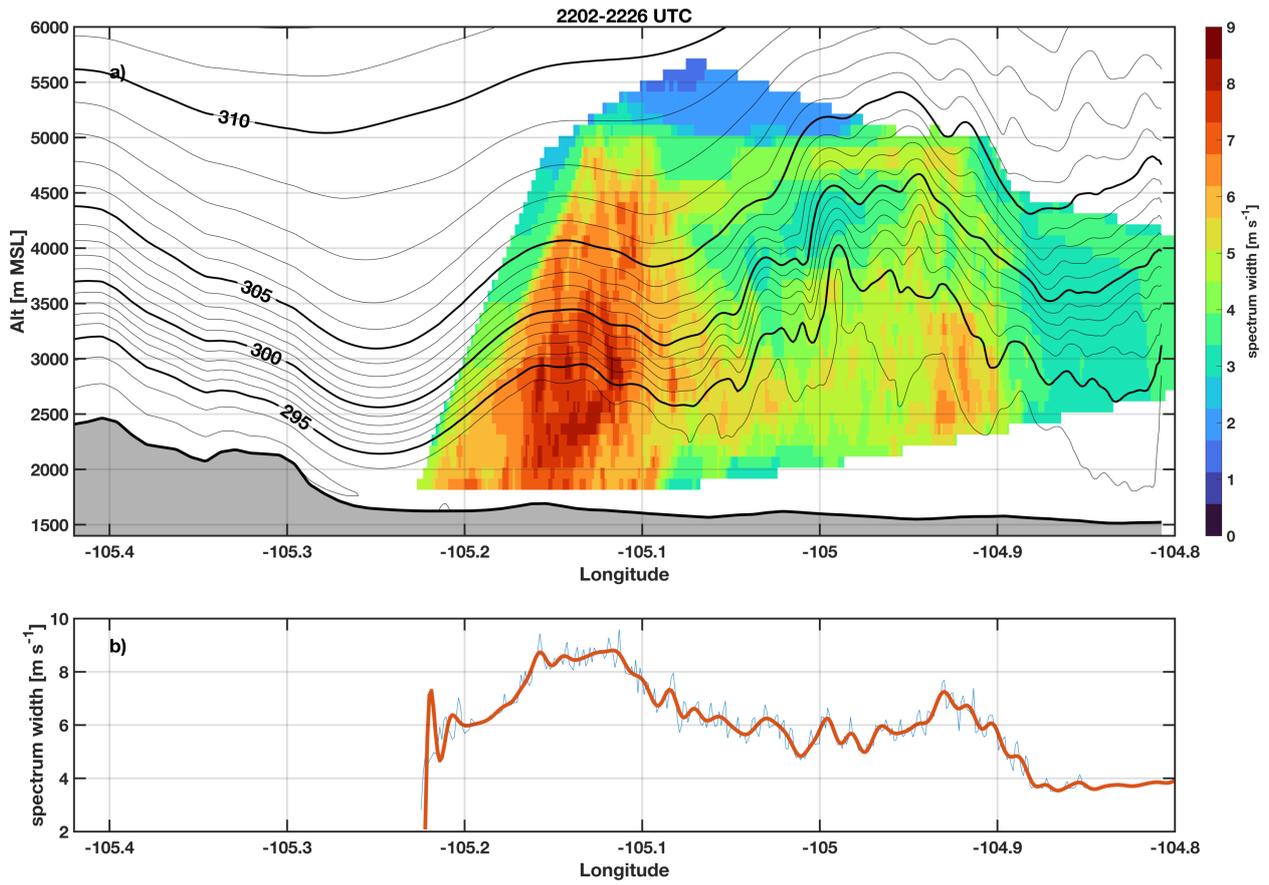


Figure S7. As in Fig. 5, except showing spectrum width.