

Application of the Pseudo-Global Warming Approach in a Kilometer-Resolution Climate Simulation of the Tropics

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Abstract

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2 **in a Kilometer-Resolution Climate Simulation of the**
3 **Tropics**

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6 **Key Points:**

- 7 • We perform a kilometer-resolution climate simulation over the tropical Atlantic
8 for current and future climate conditions using the PGW approach
- 9 • We find an accurate representation of the annual cycle of shallow cumulus clouds
10 and a realistic structure of the ITCZ, without double ITCZ
- 11 • The ITCZ intensifies in a warming climate while the narrowing typically seen in
12 GCMs is not visible

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Abstract

Clouds over tropical oceans are an important factor in Earth’s response to increased greenhouse gas concentrations, but their representation in climate models is challenging due to the small-scale nature of the involved convective processes.

We perform two 4-year-long simulations at kilometer-resolution (3.3 km horizontal grid spacing) with the limited-area model COSMO over the tropical Atlantic on a $9000 \times 7000 \text{ km}^2$ domain: A control simulation under current climate conditions driven by the ERA5 reanalysis, and a climate change scenario simulation using the Pseudo-Global Warming (PGW) approach. We compare these results to the changes projected in the CMIP6 scenario ensemble.

We find a good representation of the annual cycle of albedo, in particular for trade-wind clouds, even compared to the ERA5 reanalysis. Also, the vertical structure and annual cycle of the marine intertropical convergence zone (ITCZ) is accurately simulated, and the simulation does not suffer from the double ITCZ problem commonly present in global climate models (GCMs). The ITCZ responds to warming through a vertical extension and intensification primarily at high levels, as well as a slight southward extension of the annual mean ITCZ, while the narrowing typically seen in GCMs is not visible.

1 Introduction

Clouds over tropical oceans are among the most uncertain factors controlling Earth’s temperature response to anthropogenic greenhouse gas emissions (Forster et al., 2021). They form along the branches of the Hadley circulation (HC, e.g. Held & Hou, 1980), for instance, in the form of deep convection at the intertropical convergence zone (ITCZ) (Waliser & Gautier, 1993) and shallow convection in the marine boundary layer (MBL) in the Trades (e.g. Stevens, 2007; Wood, 2012; Vial et al., 2017). Tropical clouds have the potential for a strong radiative feedback in a warming climate (Bony & Dufresne, 2005; Zelinka et al., 2016). Yet, their evolution with climate change is uncertain (e.g. Bretherton, 2015), making them a prime focus of current climate change research.

Model intercomparison projects of global climate models (GCMs) such as the fifth or sixth phase of the Coupled Model Intercomparison Project (CMIP5, CMIP6, Taylor et al., 2012; Eyring et al., 2016) allow for an assessment of the magnitude and inter-model variability of cloud changes in a large ensemble of state-of-the-art GCMs. With respect

44 to tropical deep convection at the ITCZ, many GCMs project that the upper part of the
45 clouds (i.e. the anvils) will rise in a warming atmosphere and remain at approximately
46 the same temperature, according to the fixed anvil temperature (FAT) hypothesis (Hartmann
47 & Larson, 2002). As the anvils rise, they find themselves in a more stable environment
48 which reduces the anvil cloud fraction according to the stability iris hypothesis (Bony
49 et al., 2016). There is observational evidence supporting these hypotheses (Saint-Lu et
50 al., 2020). High-resolution simulations in aqua-planet and slab-ocean configuration mostly
51 reproduce this result of GCMs (Wing et al., 2020), even though there are exceptions (Satoh
52 et al., 2012; Singh & O’Gorman, 2015; Ohno & Satoh, 2018).

53 GCMs also project a narrowing of the annual mean ITCZ with stronger convec-
54 tive ascent near the equator (Huang et al., 2013; Byrne & Schneider, 2016; Byrne et al.,
55 2018), and a drying and widening of the subtropics, which together have been illustra-
56 tively termed the ”deep-tropics squeeze” (Lau & Kim, 2015). These projected changes
57 of tropical deep-convection are statistically robust among GCMs (Lau & Kim, 2015), even
58 though a non-negligible amount of models projects ITCZ changes of opposite sign (Byrne
59 et al., 2018). Yet GCMs do not agree on the representation of the ITCZ under current
60 climate conditions, for example, many models exert a double ITCZ structure (Mechoso
61 et al., 1995; Zhang et al., 2019). While observations show one single annual mean ma-
62 rine ITCZ rain band north of the equator, many GCMs simulate an additional rain band
63 south of the equator at certain locations and seasons. This so-called ”double ITCZ prob-
64 lem” has existed for more than two decades (e.g. Fiedler et al., 2020) and is thought to
65 be linked, among other factors, to air-sea interaction (e.g. Lin, 2007; Li & Xie, 2014) and
66 aspects of convective parameterizations (e.g. Lin, 2007; Bellucci et al., 2010; Song & Zhang,
67 2018). The narrowing and intensification of the convective regions in the deep tropics
68 in a warming climate found in GCMs is supported by observations (Wodzicki & Rapp,
69 2016; Byrne et al., 2018) and thermodynamic arguments (Jenney et al., 2020; Lau et al.,
70 2020). However, it has been argued that the observed narrowing of the ITCZ refers to
71 the width of the seasonal ITCZ band, while the deep-tropics squeeze is evident in the
72 width of the annual-mean zonal-mean tropical ascent region (Zhou et al., 2020). No clear
73 signal of a reduced mid-cloud fraction with warming was found in high-resolution sim-
74 ulations in aqua-planet configurations (Wing et al., 2020). Yet, aqua-planet configura-
75 tions show a large degree of idealization compared to the real world. Comparably lit-

76 tle is known about changes in the structure of the ITCZ from high-resolution climate sim-
77 ulations in real-world application (e.g. Satoh et al., 2012; Tsushima et al., 2014).

78 With respect to tropical low cloud changes, GCMs overall project a reduction of
79 the low-cloud albedo, but the inter-model spread is much larger than in projections of
80 deep convection (e.g. Zelinka et al., 2017; Vial et al., 2017). Also, there is a notorious
81 negative cloud bias in subtropical low-cloud regions in GCMs (e.g. Noda & Satoh, 2014;
82 Kawai & Shige, 2020). Large-eddy simulations (LES) show a more consistent climate change
83 response of low clouds (e.g. Blossey et al., 2013), but given their small domain sizes and
84 idealized setups, generalization of LES results to the entire planet introduces new un-
85 certainties.

86 The fundamental problem behind the representation of convective clouds in GCMs
87 is that a high horizontal and vertical resolution is required to resolve the small-scale con-
88 vective circulations that drive clouds. Convective circulations represent the primary mode
89 of vertical transport in the tropical atmosphere. If unresolved, these circulations, the clouds,
90 as well as the vertical transport of heat and moisture associated with them have to be
91 represented by convective parameterization schemes (e.g. Kawai & Shige, 2020). These
92 schemes introduce substantial uncertainty in the simulation of deep-convective clouds
93 (Suhas & Zhang, 2015), low-level clouds (Vial et al., 2016), and in how these clouds re-
94 spond to climate change (Sherwood et al., 2014; Vial et al., 2017). With higher model
95 resolution, convective parameterizations become less important and can eventually be
96 switched off, which reduces the degree of parameterization and allows for a model for-
97 mulation closer to physical first principles. For deep convective clouds, this threshold is
98 reached at kilometer-resolution (Prein et al., 2015) which is why kilometer-resolution cli-
99 mate simulations are increasingly considered a major milestone towards more confident
100 climate projections (e.g. Schneider et al., 2017; Satoh et al., 2019; Stevens et al., 2020;
101 Schär et al., 2020). Precipitation statistics in the deep tropics have been found to be largely
102 improved at kilometer-resolution compared to coarser models (Klocke et al., 2017; Stevens
103 et al., 2020; Hohenegger et al., 2020).

104 Global kilometer-resolution multi-year climate simulations are not yet feasible due
105 to computational cost (Schär et al., 2020), although rapid progress is evident (e.g. Satoh
106 et al., 2012, 2019; Stevens et al., 2019). Instead, multi-year kilometer-resolution simu-
107 lations are typically run on limited-area domains using boundary conditions from reanal-

108 ysis data sets for evaluation runs (e.g. Ban et al., 2021), and from GCMs for climate change
109 scenario simulations (e.g. Pichelli et al., 2021). Usually, a historical control simulation
110 and a future scenario simulation are compared to extract the climate change signal. An
111 alternative to this dynamical downscaling approach is the pseudo-global warming (PGW)
112 approach (Adachi & Tomita, 2020; Brogli et al., 2022) in which reanalysis boundary con-
113 ditions are used for both the control and the scenario simulation. The climate change
114 signal is obtained by imposing large-scale changes in the climate system on the reanal-
115 ysis boundary fields of the scenario simulation. Doing so has the advantage that the bi-
116 ases from the historical GCM run do not enter the limited-area simulation, and that rel-
117 atively short simulation periods can be used (Brogli et al., 2022). The PGW approach
118 has extensively been applied in the mid-latitudes (Schär et al., 1996; Wu & Lynch, 2000;
119 Sato et al., 2007; Rasmussen et al., 2011; Kröner et al., 2017; Brogli et al., 2019). We
120 are aware of applications in the subtropics (Chen et al., 2020; Nakamura & Mäll, 2021),
121 but to our knowledge, this study represents the first application of a PGW simulation
122 at kilometer-resolution covering the full extent of the HC including the deep tropics.

123 We run a 4-year-long limited-area atmospheric simulation at 3.3 km resolution over
124 the tropical Atlantic with the goal to (i) evaluate how well the tropical climate and the
125 associated distribution of clouds are represented in a kilometer-resolution atmospheric
126 model, and (ii) compare the climate change response of the HC in terms of its structure,
127 dynamics and clouds to the projections from the CMIP6 models. In a subsequent pa-
128 per, a systematic analysis of the ensuing radiative feedbacks in this simulation will be
129 presented. The following Section describes the modelling framework. Section 3 presents
130 the results which are discussed in Section 4 and concluded in Section 5.

131 **2 Materials and Methods**

132 **2.1 Experimental Setup**

133 The limited-area model COSMO (see Section 2.3) is used in two 4-year-long sim-
134 ulations. The first one (CTRL) serves as a control simulation and represents current cli-
135 mate conditions. It is initialized and driven at the boundaries by the European Center
136 for Medium Range Weather Forecast (ECMWF) ERA5 Re-Analysis (Hersbach et al.,
137 2020). CTRL is used to evaluate the COSMO model against observations, and serves
138 as a baseline for comparison with the second simulation. The second simulation is a cli-

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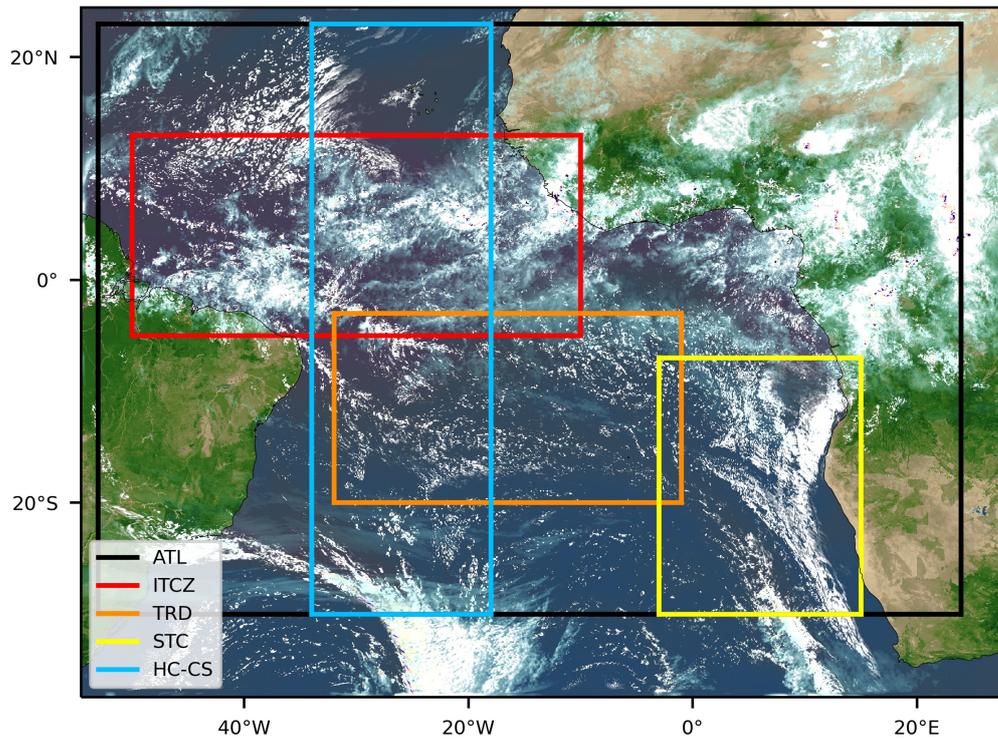


Figure 1. Simulation visualization and analysis domains. The COSMO simulations are run on the entire domain shown and the rectangles indicate different analysis domains. ATL covers the South Atlantic deep tropics and subtropics (30°S-23°N). The three subdomains ITCZ, TRD and STC comprise the regions of the three major tropical marine cloud regimes (deep convection at the ITCZ, trade-wind cumulus and stratocumulus). The HC-CS is used to compute altitude-latitude cross-sections to visualize the structure of the Hadley cell.

139 mate change scenario simulation (PGW) obtained with the pseudo-global warming ap-
140 proach (see Section 2.2).

141 Both, CTRL and PGW simulations are initialized on August 1, 2006 (for details
142 on the initialization see below) and the analysis is done for the years 2007-2010 and fo-
143 cused on five geographic regions (Fig. 1). The analysis period is too short to fully av-
144 erage out inter-annual variability. The effect of this is quantified in Section 3. The sim-
145 ulation domain covers 37.5° S - 24.5° N and 54.5° W - 28.0° E and consists of 2750 x 2065 x 60
146 grid points at 0.03° / 3.3 km resolution, integrated with a time step of 25 seconds. The
147 vertical grid stretches to an altitude of 30 km with a resolution of about 20 m near the
148 surface, 500 m at 5 km altitude, and 1.5 km at the model top. The domain covers the deep-
149 tropical and parts of the subtropical Atlantic (Fig. 1) encompassing the full southern hemi-
150 spheric branch of the HC. Although the focus lies on the Atlantic, the simulation domain
151 includes parts of Africa and South America to enable interaction between marine and
152 continental areas for instance through Monsoon circulations or the African easterly waves.

153 2.2 Pseudo-Global Warming Approach

154 The initial and boundary conditions of the PGW simulation are obtained follow-
155 ing Brogli et al. (2022) by adding the mean climate change signal (so-called climate deltas)
156 for temperature, relative humidity, horizontal wind, and surface temperature to the ERA5
157 boundary conditions of the CTRL simulation period. The climate deltas are a function
158 of latitude, longitude, pressure and month, and represent the mean annual cycle of the
159 spatial change pattern between two climate states, i.e. here between a historical and a
160 future scenario climatology. Note that, apart from the model initialization, the climate
161 deltas are only applied at the lateral boundary conditions of the limited-area model, and
162 at the surface for SST. The change signal PGW-CTRL in the interior of the domain
163 is thus a model-internal response to the forcing applied at the boundaries.

164 The climate deltas are computed from the CMIP6 output of the MPI-ESM1-2-HR
165 model (von Storch et al., 2017) as the difference between the Intergovernmental Panel
166 on Climate Change (IPCC) SSP5-8.5 scenario (Kriegler et al., 2017) simulation during
167 2070-2099 and the CMIP6 historical simulation during 1985-2014. The output of the
168 MPI-ESM is obtained as daily mean values from the CMIP6 output group CFday and
169 aggregated into monthly means. Since this output group was intended for the Cloud Feed-

170 back Model Intercomparison Project (Webb et al., 2017) it is provided on the native ver-
171 tical grid of the MPI-ESM model, and hence at fine vertical resolution. Fine resolution
172 is desirable to accurately represent the difference in warming across the trade-wind in-
173 version (see Brogli et al., 2022, Fig. 4 and corresponding discussion). The obtained changes
174 are displayed in the supplemental information (Figs. S1-4).

175 The monthly mean climate deltas are then linearly interpolated to the grid and time
176 of the ERA5 boundary files of the CTRL simulation where the deltas are added to ob-
177 tain the boundary files of the PGW simulation. After modifying temperature and rel-
178 ative humidity, the pressure field is adjusted to restore the hydrostatic balance. The cor-
179 responding changes in cloud and precipitation quickly adjust to the new thermodynamic
180 environment within the model domain, and are thus not otherwise accounted for in the
181 PGW methodology. The change of the soil temperature is computed based on the sur-
182 face temperature climate delta assuming an exponential decay of the annual cycle sig-
183 nal with depth. Initial soil moisture is not modified and taken from the CTRL simula-
184 tion (5 months before the analysis period begins). Greenhouse gas concentrations are
185 held fixed during the simulation and set to 530 ppm CO₂-eq during CTRL and 1100 ppm
186 CO₂-eq during PGW consistent with the SSP5-8.5 scenario. Aerosols are identical in CTRL
187 and PGW following the Tegen et al. (1997) climatology. Even though biomass burning
188 over Africa is a significant source of aerosol over the Atlantic (Zuidema et al., 2016), the
189 change of aerosol loading between CTRL and PGW is neglected here for simplicity. The
190 same is the case for ozone.

191 **2.3 COSMO Model**

192 The COSMO model is a fully compressible non-hydrostatic atmospheric model origi-
193 nally developed as a numerical weather prediction model (Baldauf et al., 2011) and later
194 evolved into a regional climate model (Rockel et al., 2008). Here a COSMO version ca-
195 pable of exploiting Graphics Processing Units is employed (Fuhrer et al., 2014; Leutwyler
196 et al., 2016). This version of COSMO has been extensively validated in kilometre-scale
197 configurations including a 10-year-long reanalysis-driven simulation over Europe (Leutwyler
198 et al., 2017), validation of clouds (Hentgen et al., 2019), and surface winds (Belušić et
199 al., 2018). The model discretizes the horizontal and vertical dimensions on a rotated latitude-
200 longitude grid and a generalised Gal-Chen coordinate, respectively. The model equations
201 are integrated in time with a split-explicit third-order Runge-Kutta scheme (Klemp &

202 Wilhelmson, 1978; Wicker & Skamarock, 2002; Baldauf et al., 2011). Horizontal advec-
 203 tion is treated with a fifth-order advection scheme except for moist quantities which are
 204 integrated using a positive-definite second-order scheme (Bott, 1989). The upper bound-
 205 ary is treated following (Klemp & Durran, 1983) and no relaxation of the model top to-
 206 wards the boundary files is performed.

207 Radiative transfer is computed following the δ -two-stream approach after Ritter
 208 and Geleyn (1992). The subgrid-scale vertical turbulent fluxes are parameterized with
 209 a TKE-based model (Raschendorfer, 2001). Cloud microphysics is parameterized using
 210 the single-moment bulk scheme after Reinhardt and Seifert (2006). The parameteriza-
 211 tions for deep and shallow convection are switched off as this was previously found to
 212 give a reasonable representation of clouds in the COSMO model at kilometer-resolution
 213 (Vergara-Temprado et al., 2020; Heim et al., 2021). At the surface, the second-generation
 214 land-surface model TERRA_ML (Heise et al., 2003) with the groundwater-runoff scheme
 215 after Schlemmer et al. (2018) is used on land grid points.

216 Soil moisture profiles are initialised based on a 12-year-long soil spin up COSMO
 217 simulation at 24 km grid spacing. The resulting soil moisture conditions serve as initial
 218 condition for a 5-month-long spin up at full (3.3 km) resolution, initialized on August
 219 1, 2006 (for CTRL and PGW). Over ocean grid points, sea-surface temperature is read
 220 in from the surface boundary fields. Lateral and surface boundary fields are updated ev-
 221 ery three hours. A number of empirical model parameters are adjusted to improve the
 222 representation of low clouds in comparison to previous simulations over the extratrop-
 223 ics: The vertical turbulent length scale is set to 200 m. The minimum threshold for eddy-
 224 diffusivity for heat and momentum under stable conditions are set to $0.25 \text{ m}^2 \text{ s}^{-1}$ (see
 225 Possner et al. (2014) for more details about these parameters).

226 **2.4 Data Sources**

227 ***2.4.1 CMIP6 Models***

228 The change signal between the future and the historical climate in COSMO is ob-
 229 tained by taking the difference between the PGW and the CTRL simulation. To put this

230 into perspective, the change signal of the ensemble mean of 26 CMIP6 models¹, here-
 231 after referred to as CMIP6-EM, is computed as the difference between the SSP5-8.5 ex-
 232 periment during 2070-2099 (SCEN), and the historical experiment during 1985-2014 (HIST).
 233 SCEN–HIST is thus consistent with (and in the case of the MPI-ESM model equiva-
 234 lent to) the climate delta of the PGW simulation. The CMIP6-EM is computed using
 235 output of the Amon group, thus with monthly frequency and on 11 pressure levels be-
 236 low 100 hPa. One exception is the cloud fraction which is provided on the native verti-
 237 cal grid. All analyses are performed in geometric altitude space, and the CMIP6 data
 238 is vertically interpolated on a z-coordinate.

239 *2.4.2 Observational Data Sets*

240 The following observational data sets are used to evaluate the simulations:

- 241 • The Clouds and the Earth’s Radiant Energy System (CERES) Energy Balanced
 242 and Filled (EBAF) Top-of-Atmosphere (TOA) Edition-4.0 Data Product (Loeb
 243 et al., 2018) provides monthly values of TOA radiation at 1° horizontal resolution.
- 244 • The Satellite Application Facility on Climate Monitoring (CM SAF) TOA radi-
 245 ation (Clerbaux et al., 2013), based on the Geostationary Earth Radiation Bud-
 246 get (GERB) instrument, provides monthly values of TOA radiation at 45 km hor-
 247 izontal resolution.
- 248 • The global Precipitation Measurement (GPM) Integrated Multi-satellitE Retrievals
 249 for GPM (IMERG) data set (Huffman et al., 2019): provides precipitation obser-
 250 vations at daily frequency and 0.1° horizontal resolution.
- 251 • The ERA5 reanalysis (Hersbach et al., 2020) is a gridded reanalysis data set. It
 252 is obtained from the CDS data store (Copernicus Climate Change Service (C3S),
 253 2017) and used at 3-hourly frequency and 0.25° horizontal resolution.

¹

¹ The analysed models include: ACCESS-CM2, ACCESS-ESM1-5, CAMS-CSM1-0, CanESM5, CESM2, CESM2-WACCM, CMCC-CM2-SR5, CMCC-ESM2, CNRM-CM6-1, CNRM-ESM2-1, E3SM-1-1, FGOALS-f3-L, FGOALS-g3, GFDL-CM4, GFDL-ESM4, GISS-E2-1-G, HadGEM3-GC31-LL, MIROC6, MIROC-ES2L, MPI-ESM1-2-HR, MPI-ESM1-2-LR, MRI-ESM2-0, NorESM2-LM, NorESM2-MM, TaiESM1, UKESM1-0-LL

254 **3 Results**

255 We start by looking at a cloud visualization to provide an overview of the cloud
 256 phenomena occurring within the domain. Figure 2 shows snapshots during boreal sum-
 257 mer (top) and winter (bottom). The variety of shapes and scales of tropical cloud phe-
 258 nomena and their representation in the model is remarkable. We list some of the key phe-
 259 nomena in order of decreasing size and show close-up views of them in the small pan-
 260 els (i)-(vi): (i) A mid-latitude frontal system moving eastward across the southern sub-
 261 tropical Atlantic. Such extra-tropical disturbances can reach far into the southern At-
 262 lantic subtropics during boreal summer and alter the properties of the atmosphere and
 263 subsequent formation of MBL clouds (e.g. Schulz et al., 2021). (ii) to the North of the
 264 domain, a tropical cyclone with multiple rain bands has formed and makes its way to-
 265 wards north-west. (iii) Large mesoscale convective systems travelling westward are pro-
 266 ducing heavy rainfall over the African tropical belt. (iv) Deep convection at the marine
 267 ITCZ. (v) The vast region of the Namibian stratocumulus decks (visible in both pan-
 268 els, but with larger extent during boreal winter). Finally, the stratocumulus topped MBL
 269 transitioning into (vi) the trade-wind-cumulus topped MBL on its way towards the deep
 270 tropics. Hereby, different modes of mesoscale cloud aggregation are producing regional
 271 differences in cloud cover.

272 The horizontal and vertical circulations underlying the clouds shown in Fig. 2 – from
 273 the large-scale tropical overturning HC down to small-scale convective MBL circulations
 274 – are all represented explicitly on the model grid, even though many of the circulation
 275 features are resolved only at the coarse end of the spectrum. In the following section,
 276 we are going to evaluate the simulation and compare it to the CMIP6 historical runs.

277 **3.1 Evaluation of the CTRL Simulation**

278 We start the evaluation at the large-scale with the analysis of the meridional struc-
 279 ture of the HC. Afterwards, we look at the spatial structure and annual cycle of indi-
 280 vidual cloud regimes.

281 **3.1.1 The Hadley Cell**

282 Figure 3 shows the meridional distribution of clouds and surface precipitation along
 283 the HC-CS domain for the annual mean as well as for the 3-month-periods with south-

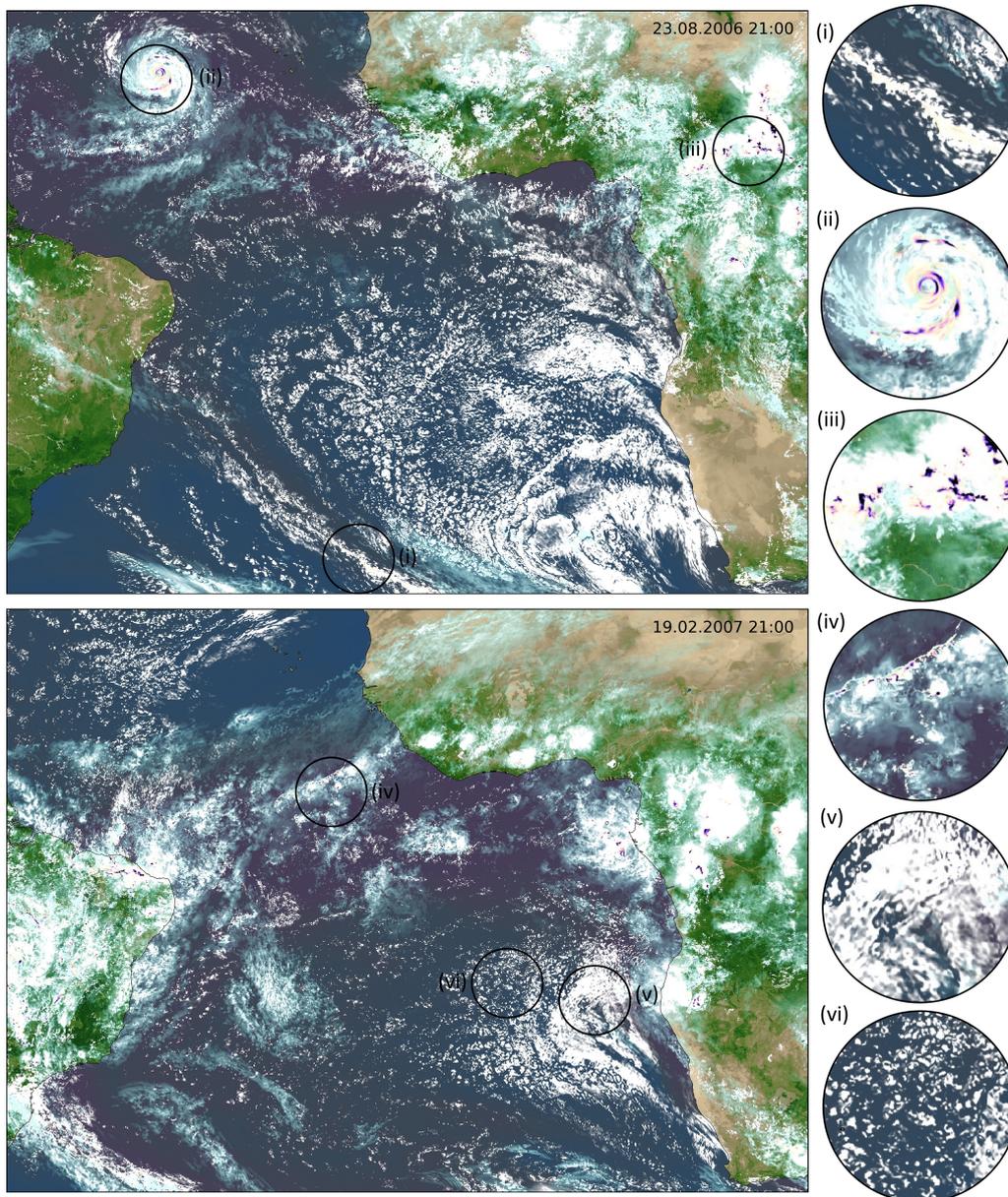


Figure 2. Snapshots of the CTRL simulation obtained during boreal summer on 23.08.2006 21:00 UTC (top) and winter on 19.02.2007 21:00 UTC (bottom). The visualization shows atmospheric cloud liquid and ice water content in white and light-blue-to-white colors, respectively, as well as surface precipitation in yellow-to-blue colors. Areas of high atmospheric water vapor content over oceans are visualized using purple shading. The land surface is rendered based on the model surface albedo and vegetation types, with a desert-to-green color gradient that is modulated by the soil moisture content to imitate the seasonal cycle of vegetation density. The panels on the right-hand side show close-up views of (i) a mid-latitude frontal system, (ii) a tropical cyclone, (iii) a mesoscale convective system, (iv) deep convection at the marine ITCZ, (v) marine stratocumulus clouds, (vi) marine shallow cumulus (or trade-wind cumulus) clouds. An animation of this visualization can be obtained via <https://doi.org/10.3929/ethz-b-000568941>.

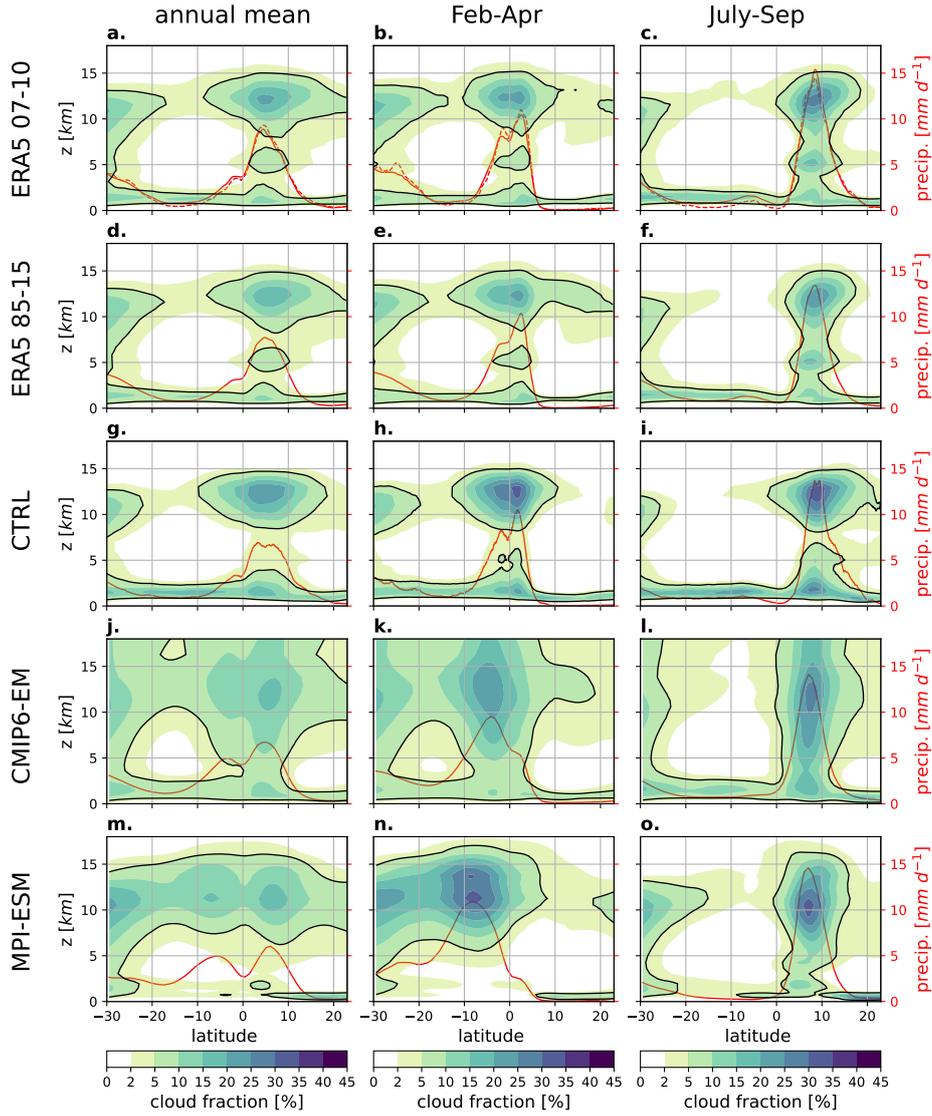


Figure 3. Altitude-latitude cross-sections of the cloud fraction [%] averaged along the longitudes of the HC-CS domain. The black contour lines denote the 5% cloud fraction level. The solid red lines show surface precipitation [mm d^{-1}]. The corresponding y-axis is located on the right-hand side of the panels. The panels show (a-c) ERA5 with the same period as the CTRL simulation (2007-2010), (d-f) an extended period of ERA5 corresponding to HIST (1985-2014), (g-i) CTRL (2007-2010), (j-l) CMIP6-EM HIST (1985-2014), and (m-o) MPI-ESM HIST (1985-2014). The three columns represent multi-year averages (left panels) during the entire year, as well as (center panels) during February–April and (right panels) July–September when the marine ITCZ reaches its southernmost and northernmost extent, respectively. The cloud fraction is obtained from the respective model output, except for COSMO where grid points with a specific cloud liquid plus ice water content $\geq 0.01 \text{ g m}^{-3}$ are considered cloudy while the remaining grid points are considered cloud-free. As a complementary reference observation, GPM IMERG precipitation is shown as a red dashed line in panels (a-c).

ernmost (February-April) and northernmost (July-September) extent of the marine ITCZ. The ERA5 record (Fig. 3a-c) indicates that the annual mean cloud fraction and precipitation have their peak at 4°N. The peak shifts to 8°N during boreal summer, and splits into a primary peak persisting at 3°N, and a secondary peak at about 2°S during boreal winter. Comparing the 4-year-long (Fig. 3a-c) and the 30-year-long (Fig. 3d-f) ERA5 cross-sections shows that the climatological distributions of cloud and precipitation are well represented by the 4-year-long simulation period used in CTRL (see Section 2.1). The comparison of surface precipitation between ERA5 and GPM IMERG indicates a close agreement between these two reference data sets (Fig. 3a-c).

The zonal mean precipitation is well reproduced in CTRL with respect to ERA5 with the exception of a slightly underestimated annual mean peak (Fig. 3g-i). The CMIP6-EM captures the northward shift of the ITCZ during boreal summer, but largely overestimates the boreal winter secondary peak in the southern hemisphere (Fig. 3j-l). The latter is a manifestation of the double ITCZ problem (Fig. 3j). Besides an overestimation of precipitation and clouds in the southern hemispheric deep tropics, the double ITCZ also results in too frequent subtropical high clouds. We further show the cross-sections for the MPI-ESM model individually (Fig. 3m-o), as it is used to compute the climate delta for the PGW simulation. The double ITCZ problem is more pronounced in the MPI-ESM model than in the CMIP6-EM and results in a bimodal annual mean precipitation distribution that is almost symmetric about the equator.

In ERA5, the cloud field at the ITCZ consists of (i) a concentration of low-level clouds, (ii) a secondary liquid cloud maximum at around 5 km (which appears to be related to an elevated inversion layer), and (iii) the deep-convective anvil clouds between 8-15 km altitude. In the subtropics, the free troposphere contains virtually no clouds below 10 km as a result of the stable and dry conditions in the downward branch of the HC. At the surface, low clouds are topping the MBL. The MBL is less shallow south of the equator than north of the equator. In the former case, the MBL is located further off the coastal upwelling regions of Africa, and thus experiences warmer sea surface temperature (SST) favoring the development of a deep MBL (e.g. Bretherton & Wyant, 1997). Beyond 25°S, clouds of extra-tropical origin penetrate into the subtropical atmosphere, in particular at high altitudes, where they contribute to the subtropical high-cloud fraction.

316 The annual mean and seasonal structure of ITCZ clouds in CTRL corresponds well
 317 with ERA5 (Fig. 3a-c,g-i). The main difference is that in the annual mean and during
 318 Feb-Apr, the extra-tropical clouds reach less far into the subtropics in CTRL. Overall,
 319 the differences are not fundamental and we conclude that CTRL simulates the cloud field
 320 along the HC consistent with ERA5. In contrast, the CMIP6-EM does not reproduce
 321 the vertical cloud structure at the ITCZ as seen in ERA5 and CTRL (Fig. 3d-f,j-l). In-
 322 stead, many of the analysed CMIP6 members simulate a too coherent cloud field through-
 323 out the tropical tropospheric column which also penetrates too high up into the trop-
 324 ical tropopause layer in some models. Although the focus of this study lies on marine
 325 clouds, the structure of the HC over land is shown in supplementary Fig. S6. While the
 326 vertical cloud structure in CTRL is comparable to ERA5, the high-cloud fraction and
 327 surface precipitation is significantly larger than in ERA5 and GPM IMERG. These quan-
 328 tities are both related to deep convection and thus indicate that deep convection at the
 329 continental ITCZ may be overestimated in CTRL, as will be shown later.

330 The first two columns of Fig. 4 show the large-scale overturning motion of the HC
 331 in terms of the meridional and vertical mass flux along the HC-CS domain. Air rises at
 332 the surface of the ITCZ and diverges above an altitude of 10 km towards the poles. The
 333 poleward (i.e., the elevated) branch of the HC converges with the northward branch of
 334 the Ferrel cell at 15°S in ERA5 (Fig. 4a) which sets the latitude of strongest subtrop-
 335 ical subsidence (Fig. 4b). The HC is closed by the trade winds at the surface that are
 336 largely confined to the MBL and converge at 5°N (Fig. 4a).

337 Compared to ERA5, the poleward branch of the HC in the southern hemisphere
 338 reaches further south in CTRL and the CMIP6-EM (Fig. 4a,e,i). Consequently, the sub-
 339 tropical subsidence extends further south in CTRL and the CMIP-EM than in ERA5
 340 (Fig. 4b,f,j). This dynamical difference between CTRL and ERA5 is consistent with the
 341 differences in the high-cloud fraction of extra-tropical origin at around 20°S between CTRL
 342 and ERA5 (Fig. 3a,g). A further difference between CTRL and ERA5 is related to the
 343 ITCZ outflow in the lower troposphere. ERA5 shows a pronounced shallow circulation
 344 between $2 \text{ km} < z < 6 \text{ km}$ (Fig. 4a) which is much weaker in CTRL (Fig. 4e). Thus, while
 345 the HC in ERA5 has a pronounced dual circulation structure (i.e. a deep circulation and
 346 a shallow circulation), the shallow circulation is almost absent in CTRL and the deep
 347 circulation is stronger (Fig. 4m). In line with that, the subtropical subsidence profile in
 348 ERA5 shows a pronounced maximum at 3 km (Fig. 4b), while it is more constant through-

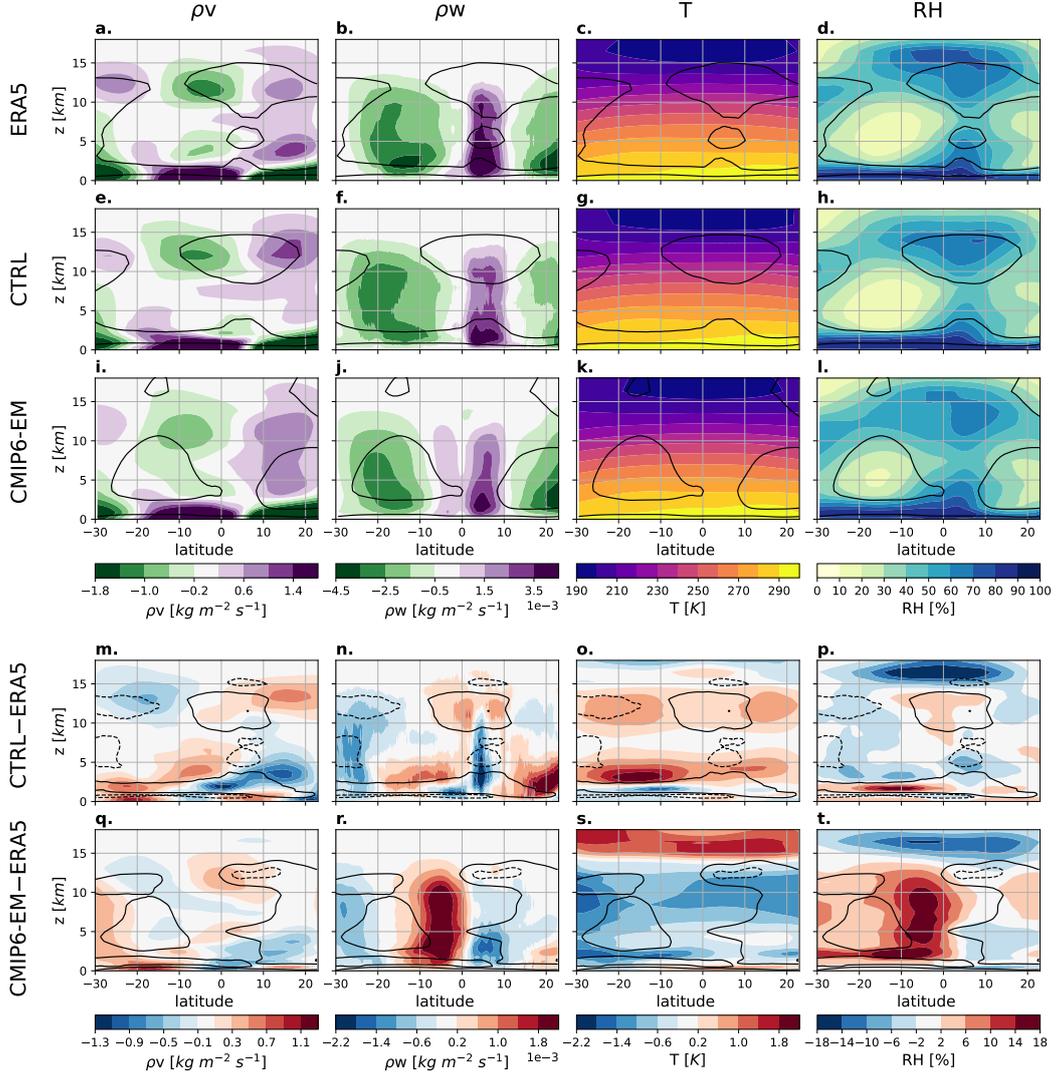


Figure 4. Annual mean altitude-latitude cross-sections of (first column) meridional mass flux [$\text{kg m}^{-2} \text{s}^{-1}$], (second column) vertical mass flux [$\text{kg m}^{-2} \text{s}^{-1}$], (third column) temperature [K], and (fourth column) relative humidity [%] averaged along the longitudes of the HC-CS domain. The panels show (a-d) ERA5 (2007-2010), (e-h) CTRL (2007-2010), (i-l) CMIP6-EM HIST (1985-2014), and the difference between (m-p) CTRL and ERA5 (2007-2010), and (q-t) CMIP6-EM HIST and ERA5 (1985-2014). The black contour lines indicate (a-l) the 5% cloud fraction level and (m-t) the 2% level of difference in the cloud fraction level where solid (dashed) lines represent a positive (negative) difference.

349 out the troposphere in CTRL (Fig. 4f). Differences in subsidence can result from dif-
 350 ferences in the radiative cooling rate or temperature stratification. The weaker subtrop-
 351 ical subsidence at low levels in CTRL appears to be due to weaker radiative cooling rate
 352 compared to ERA5 (see supplementary Fig. S7). The meridional outflow of the ITCZ
 353 in the CMIP6-EM (Fig. 4i) is more evenly distributed over the free-tropospheric column
 354 compared to CTRL and ERA5, in line with the evenly distributed cloud fraction (Fig. 3j-
 355 1). Further, the imprint of the double ITCZ is well visible in the bias of the vertical wind
 356 field, showing an anomalous upward motion south of the equator compared to ERA5 (Fig. 4r).

357 The third and fourth columns of Fig. 4 show the thermodynamic structure of the
 358 HC along the HC-CS domain. Temperature in CTRL does not deviate from ERA5 by
 359 more than 1 K except in the subtropical lower troposphere (Fig. 4o). The differences in
 360 relative humidity between CTRL and ERA5 (Fig. 4p) are also small except for altitudes
 361 above 15 km where temperatures are very low and small differences in the amount of deep-
 362 convective outflow have a large effect on the relative humidity. Overall, the subtropical
 363 troposphere is slightly drier in CTRL than in ERA5, and (as for temperature) the dif-
 364 ferences are largest in the lower troposphere. The trade-wind inversion in CTRL is more
 365 elevated than in ERA5 which explains the lower temperature and enhanced humidity
 366 in CTRL in between (i.e. between 1-2 km). Above the inversion, the differences may be
 367 related to lower-tropospheric mixing which alters the moisture content of the free tro-
 368 posphere and thus modulates the clear-sky radiative cooling rate. The drier free tropo-
 369 sphere in CTRL (Fig. 4p) may thus explain the weaker radiative cooling rate at these
 370 levels (Fig. S7), and consequently the warmer temperature (Fig. 4o) and weaker subsi-
 371 dence (Fig. 4n) in the subtropical lower troposphere in CTRL. The biases in tempera-
 372 ture and relative humidity in the CMIP6-EM (Fig. 4s,t) are larger than in CTRL. This
 373 is expected since CTRL is driven by ERA5 at its boundaries while the CMIP6 simula-
 374 tions are global. Tropospheric temperature is lower in the CMIP6-EM than in ERA5 while
 375 stratospheric temperature is higher (Fig. 4s). Further, the deep tropics in the southern
 376 hemisphere are much moister than in ERA5, due to the double ITCZ (Fig. 4t).

377 **3.1.2 Tropical Cloud Regimes**

378 We continue with a more detailed evaluation of clouds and precipitation. Figure 5
 379 shows the annual mean spatial pattern of the TOA albedo, surface precipitation, and TOA
 380 OLR. The low-cloud albedo is substantially overestimated in CTRL (Fig. 5c). Surface

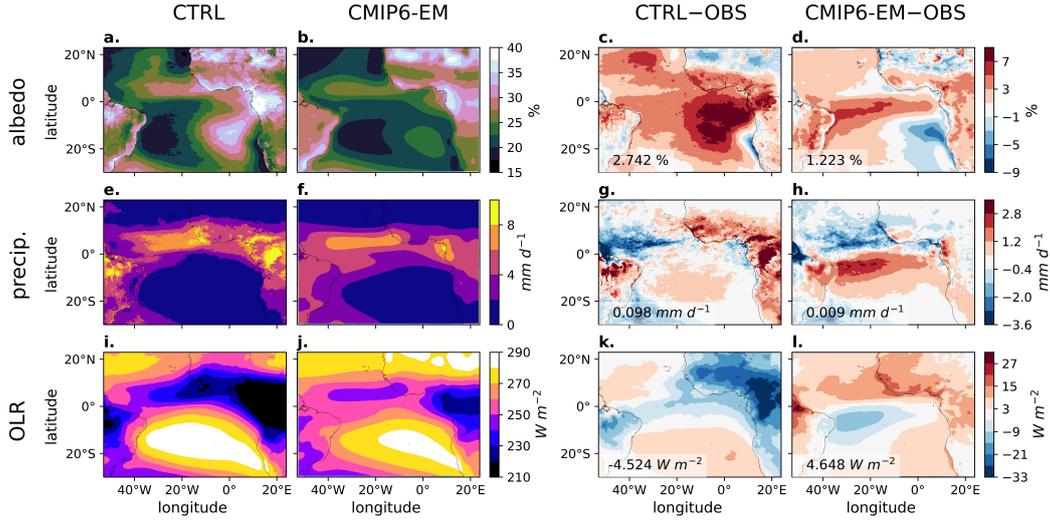


Figure 5. Evaluation of (a-d) TOA albedo [%], (e-h) surface precipitation [mm d^{-1}], and (i-l) TOA outgoing longwave radiation (OLR) [W m^{-2}]. The panels show (first column) CTRL (2007-2010), (second column) CMIP6-EM HIST (1985-2014), (third column) CTRL - OBS (2007-2010), and (fourth column) CMIP6-EM HIST (1985-2014) - OBS, where OBS is (c,d,k,l) the CM SAF record during (c,k) 2007-2010 and (d,l) 2004-2010, and (g,h) the GPM IMERG record during (g) 2007-2010 and (h) 2001-2014. The comparison between the CMIP6-EM and the OBS is thus based on the longest available observational period overlapping with HIST. The labels in the lower-left corners show domain average biases.

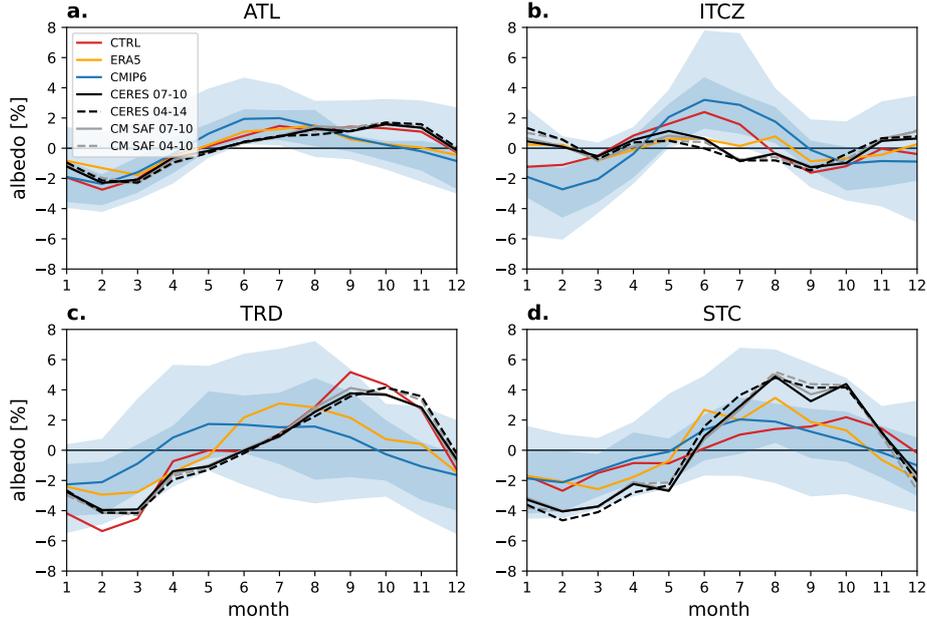


Figure 6. Evaluation of the mean annual cycle of TOA albedo shown for the four marine analysis domains (a) ATL, (b) ITCZ, (c) TRD and (d) STC. The annual cycle, expressed as the deviation from the annual mean, is shown for CTRL (2007-2010, red), ERA5 (2007-2010, yellow), CMIP6-EM HIST (1985-2014, blue), CERES EBAF (2007-2010, solid black), CERES EBAF (2004-2014, dashed black), CM SAF (2007-2010, solid gray), and CM SAF (2004-2010, dashed gray). The shading shows the CMIP6 ensemble spread between the 10th and 90th percentile (light blue) and the interquartile range (dark blue). Only ocean grid points are used in this figure.

381 precipitation in CTRL is overestimated over land in comparison to the GPM IMERG
 382 data set (Fig. 5g). The precipitation of the marine ITCZ is well represented to the East,
 383 but its westward extent is underestimated. OLR is far too low over land (Fig. 5k), in line
 384 with the precipitation bias. Over sea, the underestimation of OLR in the deep tropics
 385 is smaller, but subtropical OLR is overestimated. The CMIP6-EM bias patterns of these
 386 three variables (Fig. 5d,h,l) reveal the two well-known deficits of GCMs over low-latitude
 387 oceans: The double ITCZ problem and the underestimation of stratocumulus clouds. Also
 388 note that over land, the CMIP6-EM shows much smaller biases than CTRL.

389 We continue with the evaluation of the annual cycle of clouds on the four marine
 390 analysis domains ATL, ITCZ, TRD and STC. Figure 6 shows the mean annual cycle of
 391 TOA albedo in CERES EBAF, CM SAF, CTRL, ERA5, and the CMIP6 mean and en-

Table 1. Domain and time average values of TOA albedo [%], precipitation [mm d^{-1}] and OLR [W m^{-2}] on the four marine analysis domains ATL, ITCZ, TRD and STC shown for CTRL (2007-2010), ERA5 (2007-2010), CMIP6-EM HIST (1985-2014) and the satellite observations (CERES EBAF and CM SAF for albedo and OLR, and GPM IMERG for precipitation). As in Figs. 6-8, the satellite observations are listed during the CTRL period (2007-2010) and during the longest period overlapping with the HIST period (2004-2014 for CERES EBAF, 2004-2010 for CM SAF, and 2001-2014 for GPM IMERG).

albedo. [%]	ATL	ITCZ	TRD	STC
CTRL	24.5	24.5	23.8	27.7
ERA5	21.5	22.0	20.1	21.6
CMIP6-EM	22.2	23.1	20.6	22.2
CERES EBAF 07-10	21.1	21.4	18.8	24.4
CERES EBAF 04-14	21.2	21.3	18.8	24.7
CM SAF 07-10	20.9	20.9	18.8	24.5
CM SAF 04-10	21.0	21.0	18.9	24.9

precip. [mm d^{-1}]	ATL	ITCZ	TRD	STC
CTRL	2.10	4.15	0.95	0.37
ERA5	2.45	5.12	0.91	0.30
CMIP6-EM	2.46	4.58	1.72	0.52
GPM IMERG 07-10	2.44	5.01	0.57	0.20
GPM IMERG 01-14	2.33	4.78	0.49	0.19

OLR [W m^{-2}]	ATL	ITCZ	TRD	STC
CTRL	262.4	248.0	280.4	274.4
ERA5	269.5	259.4	285.0	279.3
CMIP6-EM	264.2	255.4	273.9	275.8
CERES EBAF 07-10	267.4	255.6	283.7	277.3
CERES EBAF 04-14	267.7	256.4	284.2	277.5
CM SAF 07-10	261.3	250.5	277.2	269.7
CM SAF 04-10	261.5	251.0	277.8	270.0

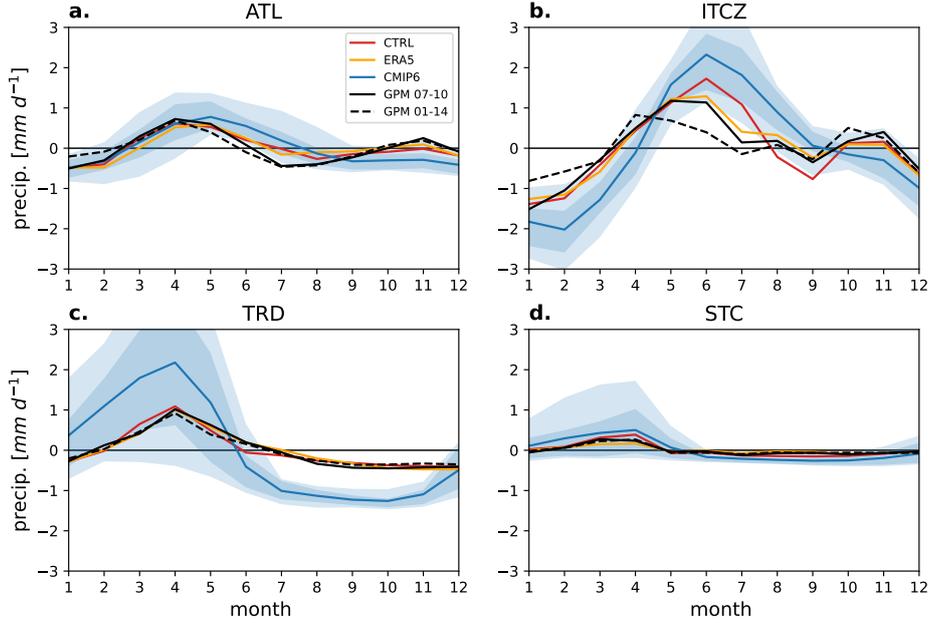


Figure 7. As Fig. 6 but showing surface precipitation. The black lines show GPM IMERG (2007-2010, solid) and GPM IMERG (2001-2014, dashed).

392 ensemble spread. The annual cycle is expressed as deviations from the annual mean. Ta-
 393 ble 1 lists the annual mean values for all data sets and analysis domains. Both obser-
 394 vational records (CERES EBAF in black and CM SAF in grey) show very similar re-
 395 sults for albedo, and are shown during the CTRL period (2007-2010, solid lines), but also
 396 during the longest available period overlapping with HIST (i.e. 2004-2014 and 2004-2010,
 397 dashed lines) to assess the effect of inter-annual variability. The annual cycle of the 4-
 398 year-period is very similar as in the extended period, indicating that the former is rep-
 399 resentative of the long-term conditions. Table 1 shows that the marine albedo in CTRL
 400 is overestimated by approximately 3.5%. However, the timings of annual maximum and
 401 minimum cloud cover as well as the amplitude of the annual cycle are much improved
 402 in CTRL compared to the CMIP6-EM on the ATL (Fig. 6a) and the TRD (Fig. 6c) do-
 403 mains, where CTRL even outperforms the ERA5 record. On the ITCZ (Fig. 6b) and the
 404 STC (Fig. 6d) domains, on the other hand, similar (though mitigated) deficiencies as in
 405 the CMIP6-EM are visible, i.e., an overestimation and underestimation of the ITCZ albedo
 406 during boreal summer and winter, respectively, as well as an underestimation of the an-
 407 nual cycle on the STC domain.

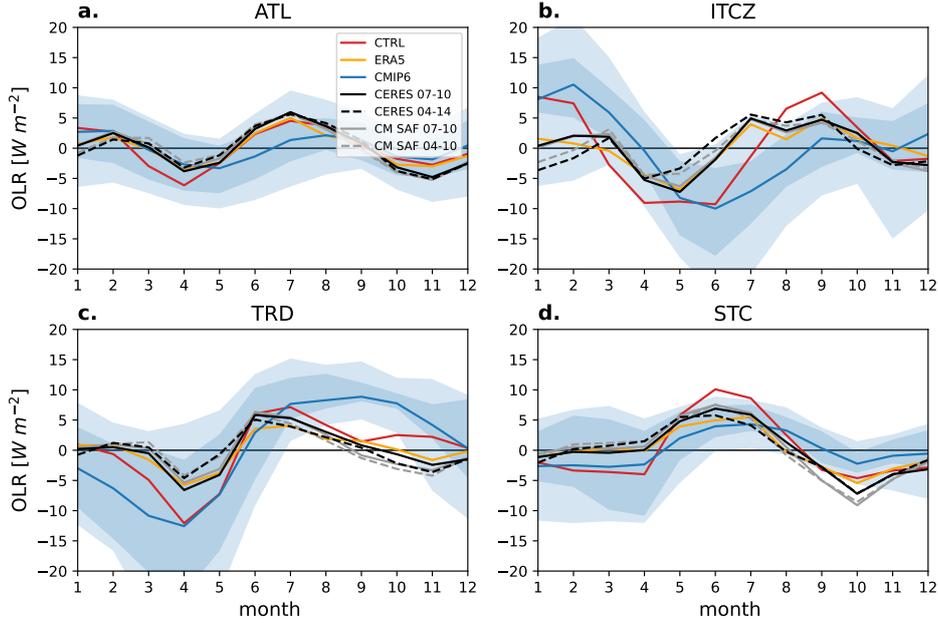


Figure 8. As Fig. 6 but showing TOA outgoing longwave radiation (OLR).

408 Figure 7 shows the annual cycle of surface precipitation over the marine analysis
 409 domains. The annual cycle is overall well represented in CTRL with the largest devia-
 410 tions found over the ITCZ domain (Fig. 7b), where the simulated timing of maximum
 411 precipitation lags one month behind the observed due to an overestimation of the bo-
 412 real summer precipitation (similar as for the albedo in Fig. 6). The relative difference
 413 in precipitation amount between the different domains is well simulated in CTRL, but
 414 with slightly more precipitation on the TRD and the STC domains compared to GPM IMERG
 415 (Table 1). In the CMIP6-EM, the precipitation amount over the TRD is overestimated
 416 more strongly due to the double ITCZ problem.

417 The annual cycle of OLR is shown in Fig. 8. Unlike for the albedo, there is a sur-
 418 prisingly large difference between CERES EBAF and CM SAF of about 6 W m^{-2} (Ta-
 419 ble 1). ERA5 is closer to CERES EBAF. The amplitude of the annual cycle of OLR in
 420 CTRL is overestimated on the three small analysis domains (ITCZ, TRD and STC; Fig. 8b-
 421 d). This appears to be mainly due to an overestimated high-cloud fraction originating
 422 at the ITCZ during the first half of the year and resulting in too much downwelling long-
 423 wave radiation. We further see a signal of too high tropospheric water vapor content (not
 424 shown) originating from the African ITCZ which contributes to the opacity of the at-

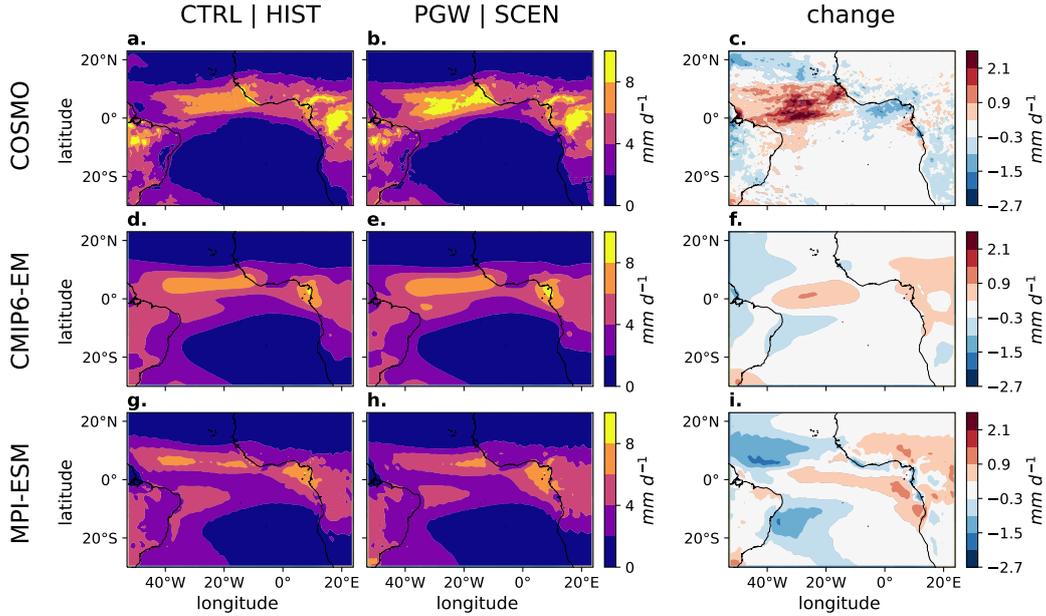


Figure 9. Annual mean climate-change signal of mean surface precipitation [mm d^{-1}]. The panels show (a-c) CTRL (2007-2010) and PGW (2007-2010), (d-f) the CMIP6-EM for HIST (1985-2014) and SCEN (2070-2099), and (g-i) the MPI-ESM model for HIST and SCEN. The columns show the simulated precipitation during (first column) CTRL and HIST, (second column) PGW and SCEN, and (third column) the change between CTRL and PGW, and between HIST and SCEN, respectively. CTRL and PGW are remapped to a 50 km grid.

425 atmosphere over the Atlantic. Similar as for precipitation, the error of CTRL is largest over
 426 the ITCZ domain (Fig. 8b).

427 3.2 Application of PGW

428 We continue with the analysis of the climate change signal obtained from the PGW
 429 simulation (see Section 2.2). Figure 9 shows the annual mean spatial distribution of sur-
 430 face precipitation change between CTRL and PGW, and between HIST and SCEN. Ma-
 431 rine ITCZ precipitation strongly increases in PGW, in some locations by up to 50% (Fig. 9c).
 432 Averaged over the ITCZ and ATL domains, precipitation increases by $7\% \text{K}^{-1}$ and $2\% \text{K}^{-1}$
 433 respectively. The temperature change for this computation was evaluated at 1 km alti-
 434 tude which roughly corresponds to the cloud base (Fig. 3g). Consistent with the CMIP6-
 435 EM (Fig. 9f), the precipitation change in COSMO (Fig. 9c) is most pronounced in the
 436 center of the Atlantic, rather than along the West-African coastline (as, e.g., in the MPI-

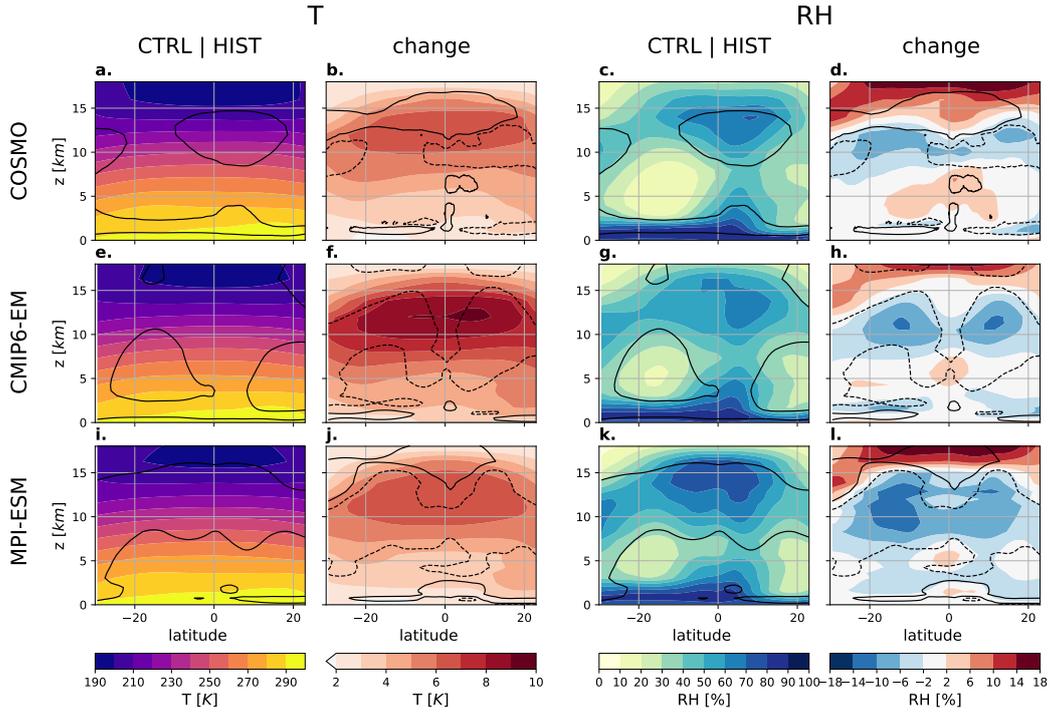


Figure 10. Annual mean altitude-latitude cross-sections of the climate-change signal of (left panels) temperature [K] and (right panels) relative humidity [%] averaged along the longitudes of the HC-CS domain. The panels show (a-d) CTRL and PGW (2007-2010), (e-h) CMIP6-EM HIST (1985-2014) and SCEN (2070-2099), and (i-l) MPI-ESM HIST and SCEN. The first and third columns show CTRL and HIST, while the second and fourth columns show the respective changes PGW–CTRL and SCEN–HIST. Panels (j,l) correspond to the climate delta from the MPI-ESM model used to derive the PGW simulation. The black contour lines indicate (first and third columns) the 5% cloud fraction level and (second and fourth columns) the level of 1% cloud fraction change where solid (dashed) lines represent a positive (negative) change.

437 ES; Fig. 9i). Also consistent is the southward propagation of the precipitation max-
 438 imum, i.e., the most pronounced change is located to the South of the precipitation max-
 439 imum in CTRL/HIST (see also Fig. 11). However, unlike in the CMIP6-EM, there is no
 440 substantial precipitation reduction in the West Atlantic trades, and precipitation over
 441 land is reduced, instead of increased. Finally, while the precipitation changes in the CMIP6-
 442 EM and the MPI-ESM associated with the ITCZ are relatively symmetric about the equa-
 443 tor as a result of the double ITCZ, this is not the case in COSMO which does not show
 444 a double ITCZ.

445 Figure 10 shows the change in the thermodynamic structure of the HC. The tem-
 446 perature change PGW–CTRL (Fig. 10b) is similar to SCEN–HIST of the MPI-ESM
 447 (Fig. 10j) but slightly smaller overall. The similarity is expected since the latter is the
 448 climate delta used to derive the PGW boundary conditions. Tropospheric relative hu-
 449 midity decreases in the CMIP6 models (Fig. 10h,l) which is a reflection of the overall dry-
 450 ing of the tropics, with the exception of a moistening deep-tropical lower troposphere
 451 (e.g. Lau & Kim, 2015). COSMO projects a qualitatively similar humidity change pat-
 452 tern (Fig. 10d) as the CMIP6 models, but with a weaker drying of the upper troposphere,
 453 a stronger moistening of the lower troposphere in the deep tropics, and – unlike in the
 454 CMIP6 models – this signal of increased humidity reaches the subtropics. Note that the
 455 relative humidity increase in the tropopause layer in all models appears to be associated
 456 with a comparably weak temperature increase due to enhanced longwave radiative cool-
 457 ing (Shine et al., 2003) and enhanced vertical moisture transport (Lau & Kim, 2015).

458 Figure 11 shows the simulated changes in the cloud field along the HC-CS domain.
 459 The signal PGW–CTRL (Fig. 11c) shows a rise of the anvil clouds at the ITCZ accom-
 460 panied by a strong increase in the high-cloud fraction. In the CMIP6-EM (Fig. 11f), the
 461 rise of the high clouds is barely visible in the cloud field change, but will be visible in
 462 the meridional wind change (see Fig. 12f). In contrast to COSMO, both the CMIP6-EM
 463 (Fig. 11f) and the MPI-ESM (Fig. 11i) exhibit a deep-tropics squeeze, i.e. a reduction
 464 of the cloud fraction at the poleward margins of the annual mean ITCZ. Note that this
 465 reduction is visible at both instances of the double ITCZ (the real one north of the equa-
 466 tor and the spurious one south of the equator). As a result of the deep-tropics squeeze,
 467 the ITCZ deep convection and precipitation in SCEN (Fig. 11e,h) is slightly more con-
 468 centrated around the equator than in HIST (Fig. 11d,g). Finally, we note that the change
 469 PGW–CTRL (Fig. 11c) in trade wind clouds exhibits an opposite sign in the North and
 470 South Atlantic, unlike in SCEN–HIST (Fig. 11f,i) where shallow cloud cover decreases
 471 in both hemispheres.

472 The circulation changes along the HC-CS domain are shown in Fig. 12 in terms of
 473 the meridional and vertical mass fluxes. COSMO simulates an upward shift (maxima rise
 474 from approximately 12 km to 14 km) and a shallower upper-level meridional outflow (lower
 475 boundary rises more than upper boundary) of the ITCZ in PGW compared to CTRL
 476 (Fig. 12a-c). This change pattern qualitatively agrees with the CMIP6-EM (Fig. 12d-
 477 f) and the MPI-ESM (Fig. 12g-i), but the change in magnitude is slightly stronger com-

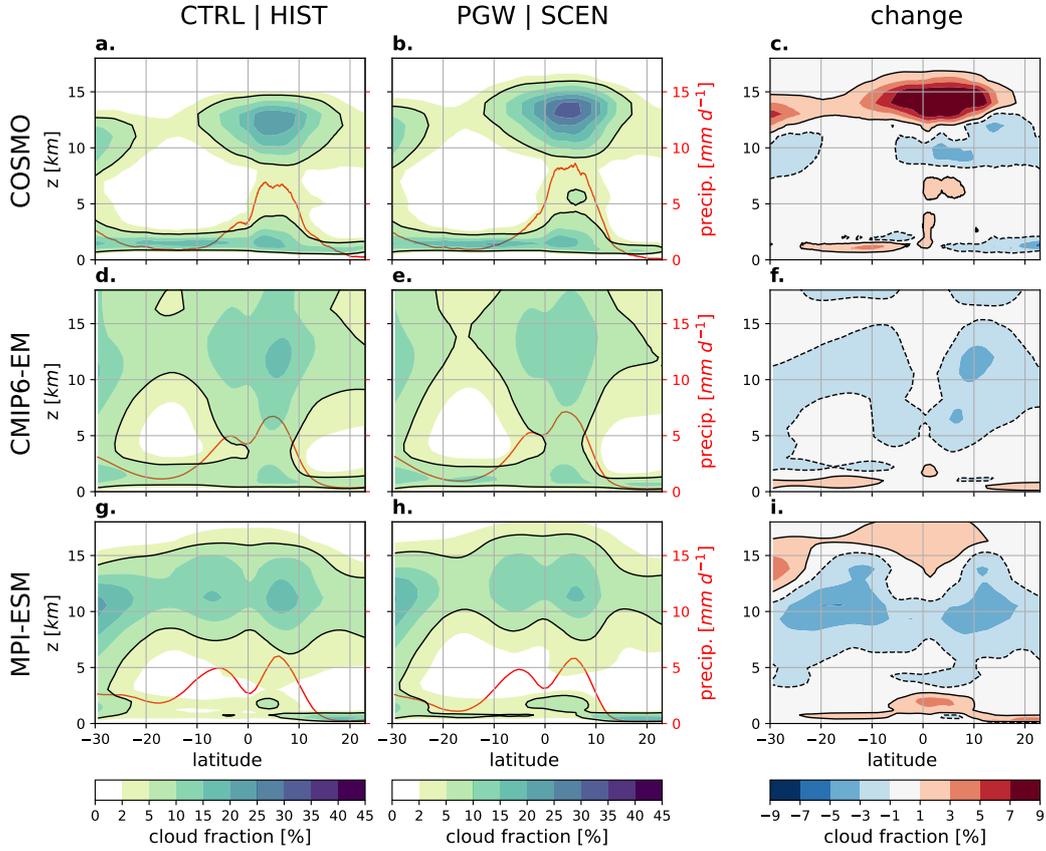


Figure 11. Annual mean altitude-latitude cross-sections of the cloud fraction [%] averaged along the longitudes of the HC-CS domain shown for (a-b) CTRL and PGW (2007-2010), (c) PGW–CTRL, (d) CMIP6-EM HIST (1985-2014), (e) CMIP6-EM SCEN (2070-2099), (f) CMIP6-EM SCEN–HIST, (g) MPI-ESM HIST, (h) MPI-ESM SCEN, and (i) MPI-ESM SCEN–HIST (i.e. corresponding to the climate delta used to derive the PGW simulation). (left and middle panels) The black contour lines high-light the 5% cloud fraction level. The red lines show surface precipitation [mm d^{-1}] represented on the scale of the right y-axis. (right panels) The black contour lines show the 1% level of cloud fraction change where solid (dashed) lines represent a positive (negative) change.

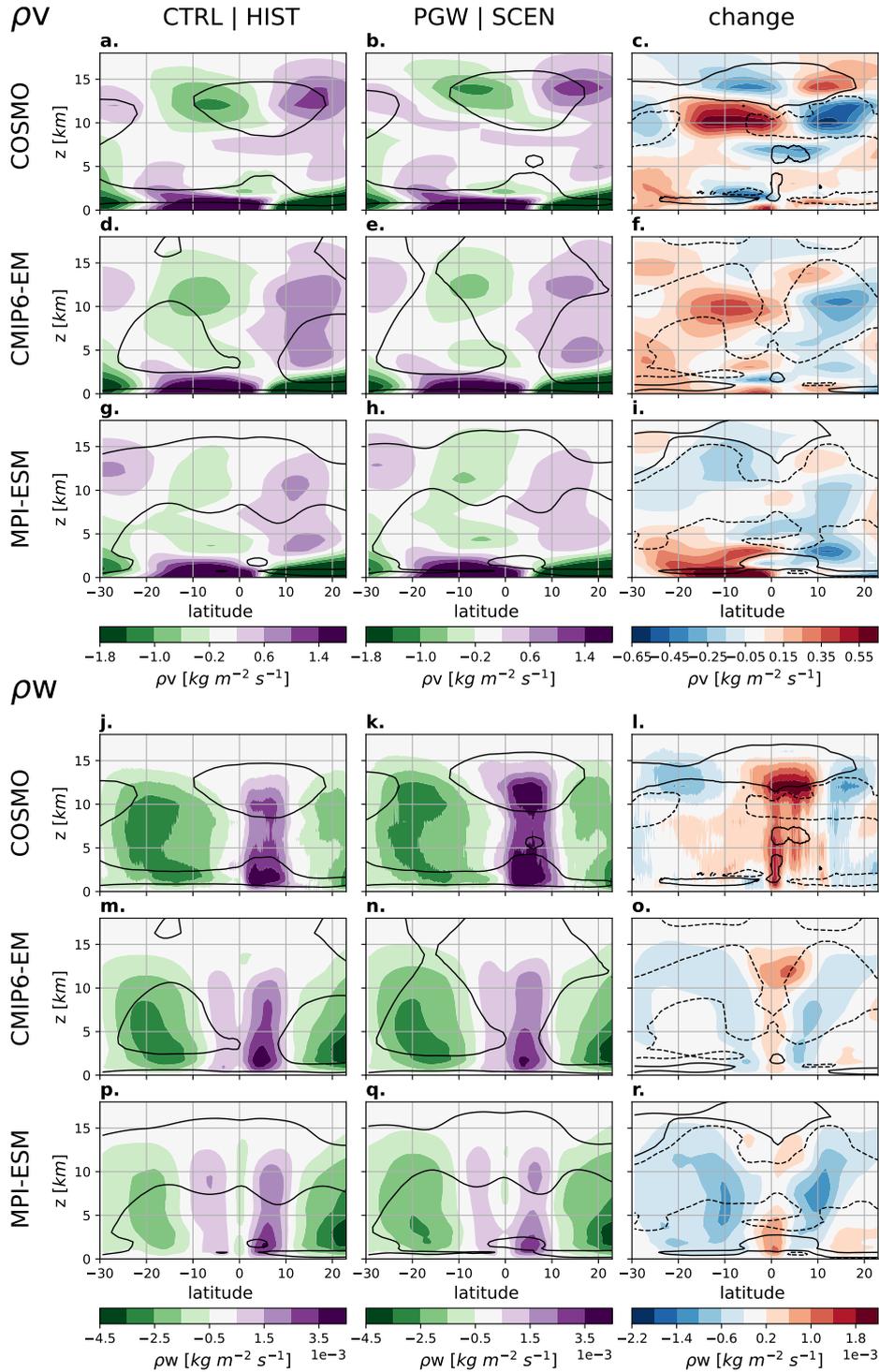


Figure 12. Annual mean altitude-latitude cross-sections of the climate change signal of (a-i) meridional mass flux ρ_v and (j-r) vertical mass flux ρ_w [$\text{kg m}^{-2} \text{s}^{-1}$] averaged along the longitudes of the HC-CS domain. The panels show (a-c,j-l) CTRL and PGW (2007-2010), (d-f,m-o) CMIP6-EM HIST (1985-2014) and SCEN (2070-2099), and (g-i,p-r) MPI-ESM HIST and SCEN. The first column shows CTRL and HIST, the second PGW and SCEN, and the third column shows the respective changes PGW–CTRL and SCEN–HIST. Panels (i,r) correspond to the climate delta from the MPI-ESM model used to derive the PGW simulation. The black contour lines indicate (first and second columns) the 5% cloud fraction level and (third column) the level of 1% cloud fraction change where solid (dashed) lines represent a positive (negative) change.

478 pared to the CMIP6-EM, and substantially stronger compared to the MPI-ESM. Along
 479 with the change in the meridional wind, the upward motion at the ITCZ in COSMO reaches
 480 higher levels and intensifies (Fig. 12j-l). The intensification occurs over the entire tro-
 481 pospheric column, but most pronounced above 10 km altitude. This response of the ITCZ
 482 to warming in COSMO shows remarkable differences to the CMIP6 models (Fig. 12m-
 483 r): First, the intensification of the ITCZ above 10 km is significantly stronger in COSMO
 484 than in the CMIP6 models (compare Figs. 12l and 12o,r), and – with a vertical exten-
 485 sion of about 2 km, i.e. from around 13 km to 15 km altitude (Fig. 12j,k) – the deepening
 486 of the ITCZ is larger compared to the CMIP6-EM (about 1 km, from around 12 km
 487 to 13 km, Fig. 12m,n). Second, the change of the ITCZ below 10 km represents a response
 488 that differs from the deep-tropics squeeze. While the CMIP6 models simulate a weak-
 489 ening of the upward motion at the margins of the deep-tropics and only a weak inten-
 490 sification at the equator (i.e. the deep-tropics squeeze; Fig. 12o,r), COSMO simulates
 491 an extension of the ITCZ towards south and an intensification over the entire meridional
 492 extent of the CTRL ITCZ (Fig. 12j,l).

493 With respect to subtropical subsidence, the response of COSMO also differs from
 494 the CMIP6 models. First, the strengthening of subsidence is mostly confined to the edge
 495 of the cloud anvils above 10 km and extends less prominently through the tropospheric
 496 column than in the CMIP6-EM and the MPI-ESM. In the northern hemisphere, the sub-
 497 sidence intensification below 10 km is still comparable to the CMIP6 models (even though
 498 confined to the subtropics), but in the southern hemisphere, there is an overall weaken-
 499 ing of annual mean subsidence in COSMO, as opposed to the strengthening in the CMIP6
 500 models. Finally, the intensification of subtropical subsidence above 10 km is substantially
 501 larger in COSMO than in the CMIP6-EM, consistent with the more pronounced deep-
 502 ening and upper-level intensification of the ITCZ deep convection.

503 **4 Discussion**

504 **4.1 Evaluation of CTRL**

505 In Section 3, we discussed the realism of the ERA5-driven CTRL simulation in com-
 506 parison to the CMIP6-EM and found significant differences. As the two underlying sim-
 507 ulation strategies differ strongly, it is not feasible to disentangle effects due to compu-
 508 tational resolution (3 km versus 50-200 km) and simulation setup (ERA5 driven atmo-

509 spheric simulations versus free-running coupled simulations). The main purpose of the
510 following discussion is thus to summarize the differences between the CTRL simulation
511 and the CMIP6 ensemble, and to determine whether the ERA5-driven simulations are
512 credible enough to serve as the basis of climate-change simulations using the PGW ap-
513 proach.

514 The improved representation of the annual cycle of the albedo, in particular on the
515 TRD analysis domain (representative of shallow cumulus clouds), as well as the accu-
516 rate vertical structure and meridional position of the ITCZ (i.e. no double ITCZ) are per-
517 haps the most promising improvements compared to the CMIP6-EM. The prescribed SST
518 obtained from ERA5 likely has a beneficial impact on the properties of the MBL and
519 the position of the ITCZ in CTRL. For instance, the double ITCZ problem of the CMIP6-
520 EM is thought to be related to air-sea interaction, among other factors (Lin, 2007; Li
521 & Xie, 2014). It would therefore be interesting to test if for instance a coupled model
522 setup at kilometer-resolution or a GCM-driven kilometer-resolution simulation were to
523 suffer from the double ITCZ problem. Under the assumption that the improved repre-
524 sentation of the ITCZ in our limited-area CTRL simulation is due to the forcing from
525 ERA5, our application demonstrates one benefit of the PGW approach compared to con-
526 ventional downscaling, i.e. that GCM circulation biases are not propagated to the limited-
527 area simulation. We argue that this realistic representation of the ITCZ location is a good
528 starting point to study its climate change signal.

529 Concerning the simulation of low clouds, the representation of the annual cycle of
530 the albedo in CTRL is better on the TRD domain than on the STC domain. This dis-
531 crepancy may relate to the type of clouds most prevalent on the two domains. The TRD
532 domain is predominantly covered by trade-wind cumulus clouds while stratocumulus clouds
533 are more frequent on the STC domain (Warren et al., 1988). The difficulty to represent
534 the annual cycle of stratocumulus clouds in a kilometer-resolution model with 60 ver-
535 tical levels is not unexpected since a firm representation of the stratocumulus-topped MBL
536 with its very shallow inversion cloud layer is challenging even in LES (e.g. Stevens et al.,
537 2005). Nevertheless, the fact that the COSMO simulations yield stratocumulus decks
538 already at kilometer-resolution, notably without any shallow convection scheme, is very
539 promising. In the trade-wind cumulus regime clouds often aggregate into clusters that
540 frequently exceed the kilometer-scale (e.g. Bony et al., 2020). The CTRL simulation in-
541 deed produces such clusters (see Fig. 2) suggesting that some of the dominant mesoscale

542 patterns of MBL circulations and clouds in the Trades are at least partially resolved. Sim-
 543 ilar results have been found in previous studies using kilometer-resolution models (Klocke
 544 et al., 2017; Heim et al., 2021; Caldwell et al., 2021). It is interesting to note that the
 545 annual cycle of albedo in CTRL on the TRD domain is actually better simulated than
 546 in ERA5. This result suggests that the improved representation of these clouds is not
 547 primarily a result of the prescribed SST, but portrays the added value of explicit con-
 548 vection and fine model resolution.

549 On the other hand, we find a mean bias in the low-cloud albedo in the CTRL sim-
 550 ulation compared to satellite observations (Fig. 5). This bias was found to be caused by
 551 an overestimation of cloud water (i.e., cloud opacity) rather than cloud fraction (Heim
 552 et al., 2021). As shown by Liu et al. (2022), this bias of the COSMO model at kilometer-
 553 resolution can be reduced through systematic model calibration. The model version used
 554 here is still based on a set of empirical parameters that were calibrated for applications
 555 over continental regions of the mid-latitudes (Bellprat et al., 2016). We also find a bias
 556 in quantities related to deep convection at the continental ITCZ over Africa (Fig. 5). Com-
 557 pared to the well calibrated COSMO simulations in the mid-latitudes (e.g. Leutwyler
 558 et al., 2017; Vergara-Temprado et al., 2020; Ban et al., 2021; Zeman et al., 2021), the
 559 bias in precipitation and OLR is still quite substantial. The set of empirical parameters
 560 used in this study differs from other COSMO setups that have been used over Africa (Bucchignani
 561 et al., 2016; Sørland et al., 2021). A calibration effort similar as it was done for the trop-
 562 ical Atlantic in Liu et al. (2022), but for continental Africa would likely result in a sim-
 563 ulation setup with less biased deep convection overall. Note, it is possible that the poor
 564 representation of the continental ITCZ could affect the representation of the marine ITCZ
 565 via the lower-tropospheric mean easterly flow or via gravity waves (e.g. Leutwyler & Ho-
 566 henegger, 2021).

567 **4.2 Climate Change Signal PGW–CTRL**

568 The changes SCEN–HIST in wind and humidity at the Atlantic HC compare qual-
 569 itatively well to the the global CMIP5 models (Lau & Kim, 2015). This agreement in-
 570 dicates that, despite the local computational domain employed, the obtained results may
 571 be indicative of the global patterns. Concerning the change signal in COSMO (PGW–CTRL),
 572 the tropospheric warming profile closely follows the climate delta (SCEN–HIST) of the
 573 MPI-ESM simulation (Fig. 10). This similarity is expected since the temperature change

574 is a large-scale signal that enters the model at the lateral boundaries (see Sec. 2.2). The
575 change signal PGW–CTRL for humidity shows a qualitatively similar change pattern
576 as the CMIP6-EM and the MPI-ESM, however, with an overall weaker drying of the trop-
577 ical atmosphere (Fig. 10). The distribution of humidity is tied to the representation of
578 deep convection and how it changes between CTRL and PGW (or HIST and SCEN).
579 Since domain-average convection at the ITCZ intensifies in COSMO but weakens in the
580 MPI-ESM model, some differences in the humidity change are expected.

581 The circulation changes, on the other hand, differ quite substantially between COSMO
582 and the CMIP6 models. The intensification of deep convection at the ITCZ is remark-
583 ably strong and accompanied by a widening of the ITCZ in the presented kilometer-resolution
584 simulation. This result is novel, since GCM projections show an anti-correlation between
585 strengthening and widening of the ITCZ between models (Byrne et al., 2018). Also, the
586 rise of the anvil clouds is more pronounced than in the CMIP6 models, and the increase
587 in the anvil cloud fraction is even contrary to the expectation of the stability iris hypoth-
588 esis (Bony et al., 2016). In this respect, our simulation qualitatively differs from high-
589 resolution simulations of radiative-convective equilibrium in aqua-planet configurations,
590 whereof a majority shows a reduction in the high-cloud fraction with warming (Wing
591 et al., 2020). Yet, the increase in tropical high clouds shows similarities to the response
592 Satoh et al. (2012) found in their global kilometer-resolution short-term climate simu-
593 lation. For this simulation, Tsushima et al. (2014) determined that the change in high
594 ice clouds is sensitive to the formulation of subgrid turbulent mixing. The work of Tsushima
595 et al. (2014); Ohno and Satoh (2018); Ohno et al. (2019, 2021) demonstrates that even
596 at kilometer-resolution the response of tropical deep convection to warming may be sub-
597 ject to extensive inter-model variability, and that the here presented results require cor-
598 roborations from kilometer-resolution climate simulations employing other model codes,
599 microphysics schemes, and downscaling approaches.

600 An often discussed hypothesis on the change in the dynamics of the HC is the promi-
601 nent deep-tropics squeeze, i.e. the narrowing of the annual mean ITCZ, detectable in GCMs
602 (e.g. Lau & Kim, 2015; Byrne & Schneider, 2016). In our CMIP6 ensemble, the squeeze
603 is clearly evident in the form of a strengthening and narrowing of the deep-tropical con-
604 vection and a corresponding reduction of cloud fraction at the edges of the ITCZ (Fig. 11
605 and Fig. 12). However, this narrowing of the annual mean ITCZ seems to be enhanced
606 by the fact that the CMIP6-EM projects a similar but mirrored change signal at both

607 branches of the ITCZ (i.e. the one north of the equator, and the spurious one south of
608 the equator – the double ITCZ). This perception is supported by the fact that the nar-
609 rowing of the ITCZ in GCMs is associated mainly with a northward shift of the south-
610 ern edge (Byrne & Schneider, 2016). The deep-tropics squeeze can not be visually de-
611 tected in the kilometer-resolution simulation which does not produce a double ITCZ in
612 CTRL. So, the question arises whether the narrowing of the deep tropics in the CMIP6-
613 EM would be equally pronounced if it did not exhibit the double ITCZ in HIST. The
614 circulation changes projected by COSMO differ more prominently from CMIP6-EM at
615 the southern edge of the ITCZ, suggesting that the double ITCZ may indeed contribute
616 to the differences in the projected change. The double ITCZ was found to relate to the
617 strength of the low-cloud feedback in GCMs (Tian, 2015) which was argued to be driven
618 by differences in the lower-tropospheric stability depending on the strength of the dou-
619 ble ITCZ (Webb & Lock, 2020). Whether and how the double ITCZ responds to warm-
620 ing and how this relates to radiative feedbacks is thus of high relevance for climate pro-
621 jections and requires further research.

622 There are some limitations of the model setup presented in this study. The COSMO
623 model was originally designed as a weather prediction model, and aerosols and ozone are
624 represented in a simplified manner compared to comprehensive climate models. Further,
625 the one-moment microphysics scheme assumes a constant cloud-droplet number concen-
626 tration. Changes in aerosol concentrations therefore do not directly alter the properties
627 of the simulated clouds. Keeping ozone and aerosol concentrations constant between CTRL
628 and PGW is thus a pragmatic choice for the given model configuration. Still, account-
629 ing for such effects might alter the simulated response to warming. For instance, the MPI-
630 ESM shows an increase and slight upward shift of the ozone maximum between HIST
631 and SCEN. Another simplification of the modelling setup in this study is the use of a
632 limited-area model and the PGW approach. Given that the same weather enters the model
633 domain at the boundaries in CTRL and PGW, large-scale circulation changes from the
634 GCM may be restrained by the persistence of the weather phenomena at the lateral bound-
635 aries. Specifically, at the boundary between the subtropics and the mid-latitudes, it is
636 unclear how the extension of the HC towards South with warming (e.g. Lau & Kim, 2015)
637 is restrained by the fact that the mid-latitude frontal systems enter the PGW simula-
638 tion at the same latitudes as in CTRL.

639 An interesting extension of this study would be to repeat the analysis using PGW
640 simulations derived with climate deltas of different GCMs to test the sensitivity of the
641 change signal PGW–CTRL to the climate delta. The role of SST warming patterns ap-
642 pears to be of particular interest here. Given the importance of the SST pattern on changes
643 of the ITCZ (Huang et al., 2013), it would not be surprising to find differences in the
644 change PGW–CTRL in terms of structure and location of the ITCZ for different climate
645 deltas.

646 5 Conclusion

647 In this study, we conducted what is, to our best knowledge, the first application
648 of the pseudo-global warming (PGW) approach on a marine tropical domain that con-
649 tains the entire Hadley circulation. We performed two 4-year-long simulations at 3.3 km
650 horizontal resolution with the limited-area model COSMO over the tropical Atlantic. The
651 analysis includes an evaluation of the structure of the Hadley circulation and tropical
652 clouds under current climate conditions (CTRL), and a comparison of the obtained cli-
653 mate change signal (PGW–CTRL) to that of a CMIP6 model ensemble (SCEN–HIST).
654 The radiative feedback between CTRL and PGW will be analysed in a follow-up study.
655 The main analysis findings include:

- 656 1. An improved representation of the vertical structure and seasonal cycle (in terms
657 of the meridional location) of the Atlantic ITCZ compared to the CMIP6 ensem-
658 ble. In particular, our limited area simulation with explicit convection does not
659 suffer from the double ITCZ problem.
- 660 2. An improved representation of the annual cycle of the TOA albedo compared to
661 the CMIP6 ensemble, in particular in the trade-wind cumulus region where CTRL
662 even outperforms the ERA5 reanalysis. This suggests that kilometer-resolution
663 simulations are a suitable tool to study cloud feedbacks in the trade-wind region.
664 Despite disabling the models shallow convection scheme, stratocumulus clouds are
665 evident, albeit somewhat too frequent, and with an underestimated amplitude of
666 the annual cycle.
- 667 3. The dynamics of the ITCZ respond to warming in a different way in our kilometer-
668 resolution simulation compared to the analysed GCMs. While the CMIP6 ensem-
669 ble shows a narrowing and central intensification of the ITCZ, i.e. a prominent

670 deep-tropics squeeze, the kilometer-resolution simulation shows an overall inten-
671 sification of the ITCZ, most pronounced at high altitudes, and a slight extension
672 towards south.

673 Overall, our results demonstrate the merit of high-resolution climate simulations in a real-
674 world configuration to compare against GCM projections. kilometer-resolution models
675 enable an unprecedented view on tropical clouds and circulations from the large-scale
676 tropical overturning circulation down to small-scale convective MBL circulations and clouds.
677 Even though global kilometer-resolution climate simulations are not yet feasible, our study
678 demonstrates that downscaling strategies like the PGW approach allow to gain insights
679 from these models already today. We presented one such simulation that, compared to
680 GCMs, produces a remarkably different climate-change response for the HC and in par-
681 ticular for the ITCZ. The realism of this response is difficult to assess as long as such
682 simulations remain a rarity. We will analyse in more detail the cause of the response in
683 upcoming work.

684 **6 Data Availability**

685 The CERES EBAF TOA radiation data are available at [https://ceres-tool.larc.nasa](https://ceres-tool.larc.nasa.gov/ord-tool/jsp/EBAF_TOA41Selection.jsp)
686 [.gov/ord-tool/jsp/EBAF_TOA41Selection.jsp](https://ceres-tool.larc.nasa.gov/ord-tool/jsp/EBAF_TOA41Selection.jsp) via DOI:10.5067/TERRA-AQUA/CERES/
687 EBAF-TOA_L3B004.1.

688 The CM SAF TOA radiation data are available at [https://wui.cmsaf.eu/safira/action/](https://wui.cmsaf.eu/safira/action/viewProduktList?dId=3)
689 [viewProduktList?dId=3](https://wui.cmsaf.eu/safira/action/viewProduktList?dId=3) via DOI:10.5676/EUM_SAF_CM/TOA_GERB/V002.

690 The GPM IMERG precipitation data are available at <https://disc.gsfc.nasa.gov>
691 via DOI:10.5067/GPM/IMERGDF/DAY/06.

692 The ERA5 reanalysis data are available at the Copernicus Climate Change Service (C3S)
693 Climate Data Store via DOI:10.24381/cds.bd0915c6.

694 The CMIP6 data are available at the <https://esgf-node.llnl.gov/projects/cmip6/>.
695 The software to prepare PGW simulations can be obtained from [https://github.com/](https://github.com/Potopoles/pgw-python)
696 [Potopoles/pgw-python](https://github.com/Potopoles/pgw-python) via DOI:10.5281/zenodo.6759029.

697 The weather and climate model COSMO is free of charge for research applications (for
698 more details see: <http://www.cosmo-model.org>).

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Supporting Information for ”Application of the Pseudo-Global Warming Approach in a Kilometer-Resolution Climate Simulation of the Tropics”

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Contents of this file

1. Figures S1 to S7

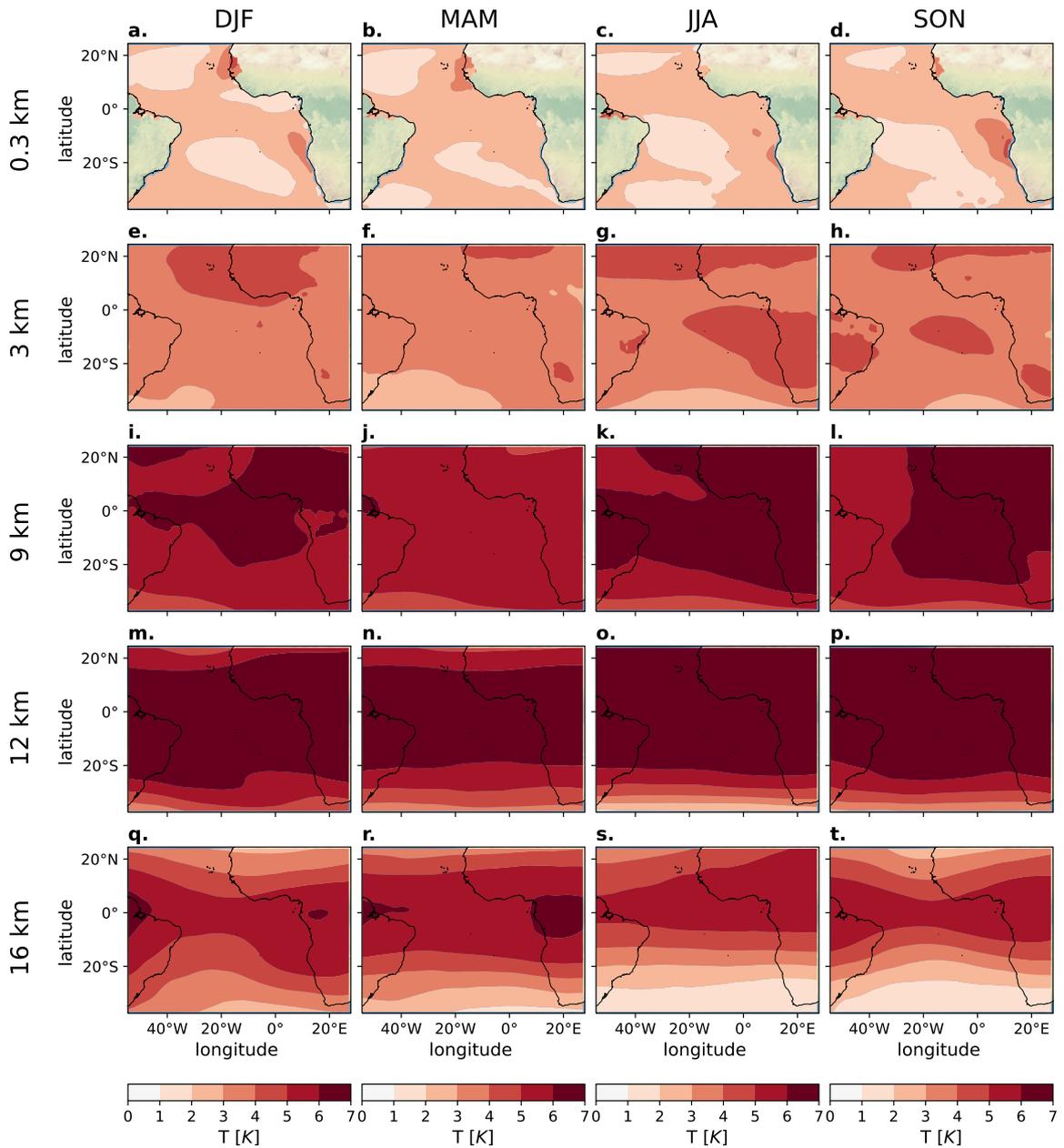


Figure S1. Climate delta for temperature shown at (first row) 0.3 km, (second row) 3 km, (third row) 9 km, (fourth row) 12 km, and (fifth row) 16 km altitude for the seasons (first column) December-February, (second column) March-May, (third column) June-August, and (fourth column) September-November. The delta is computed as the difference between SCEN (SSP5-8.5, 2070-2099) and HIST (historical, 1985-2014) for the MPI-ESM1-2-HR model.

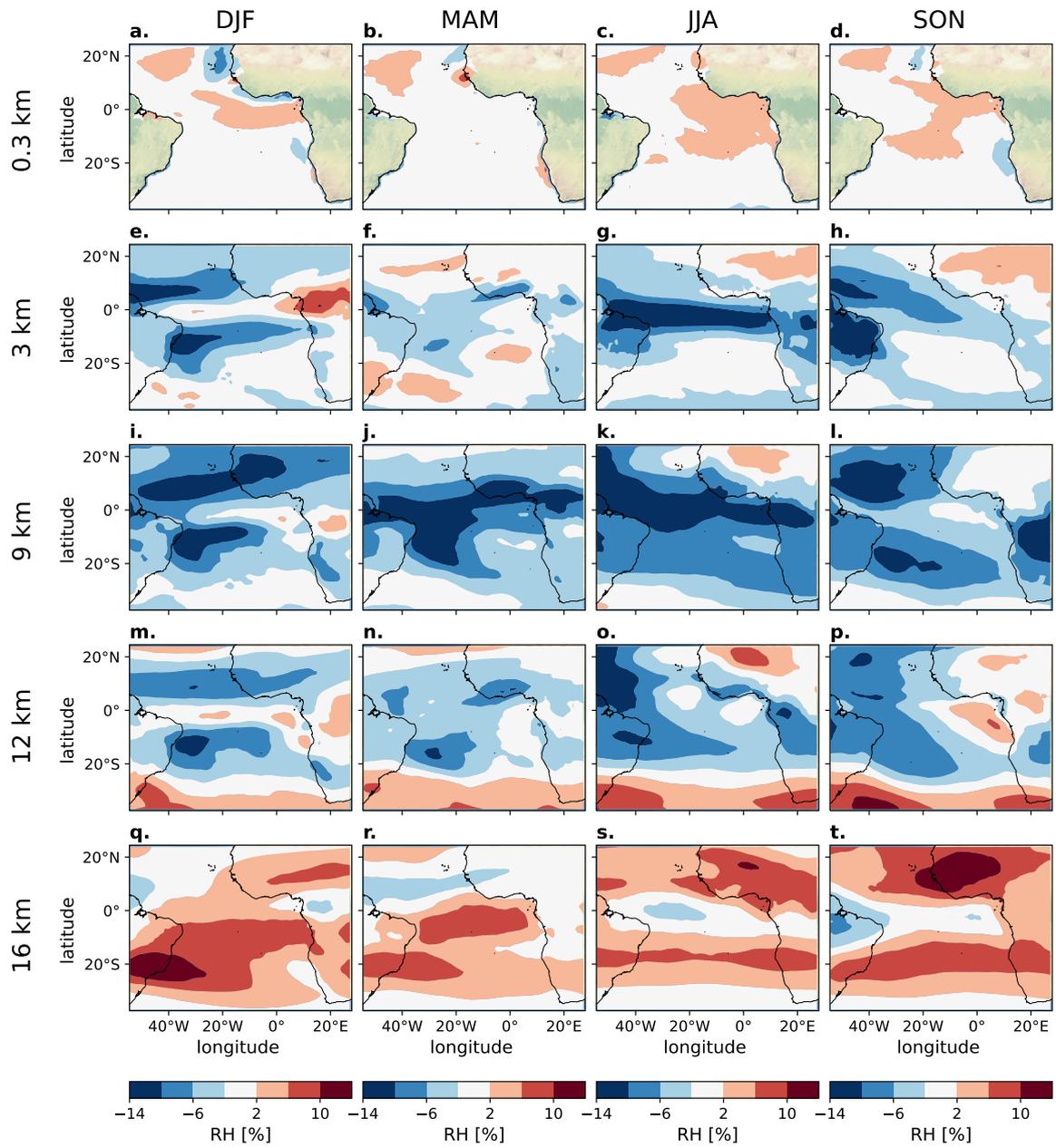


Figure S2. As Fig. S1 but shown for the relative humidity.

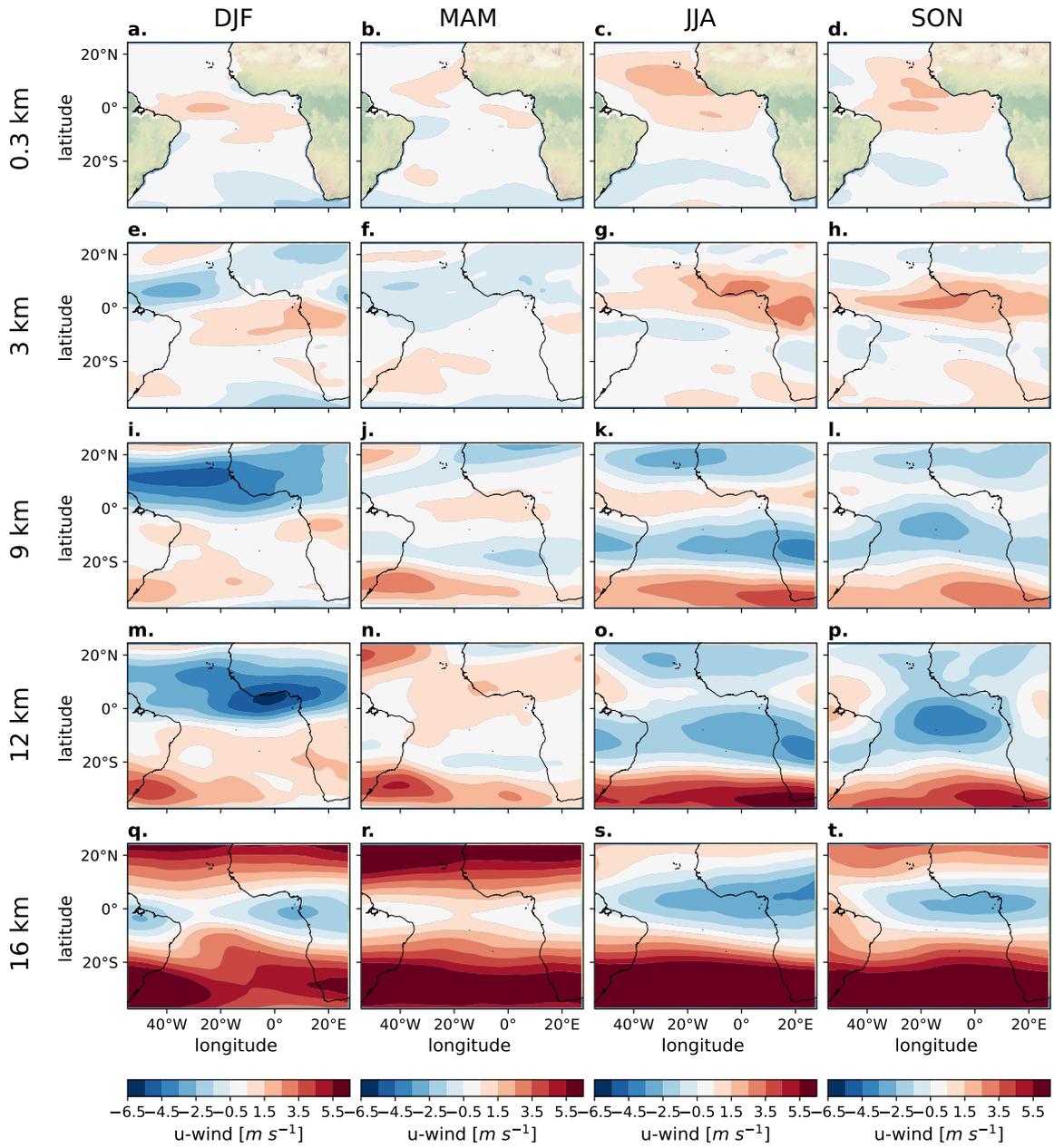


Figure S3. As Fig. S1 but shown for the zonal wind.

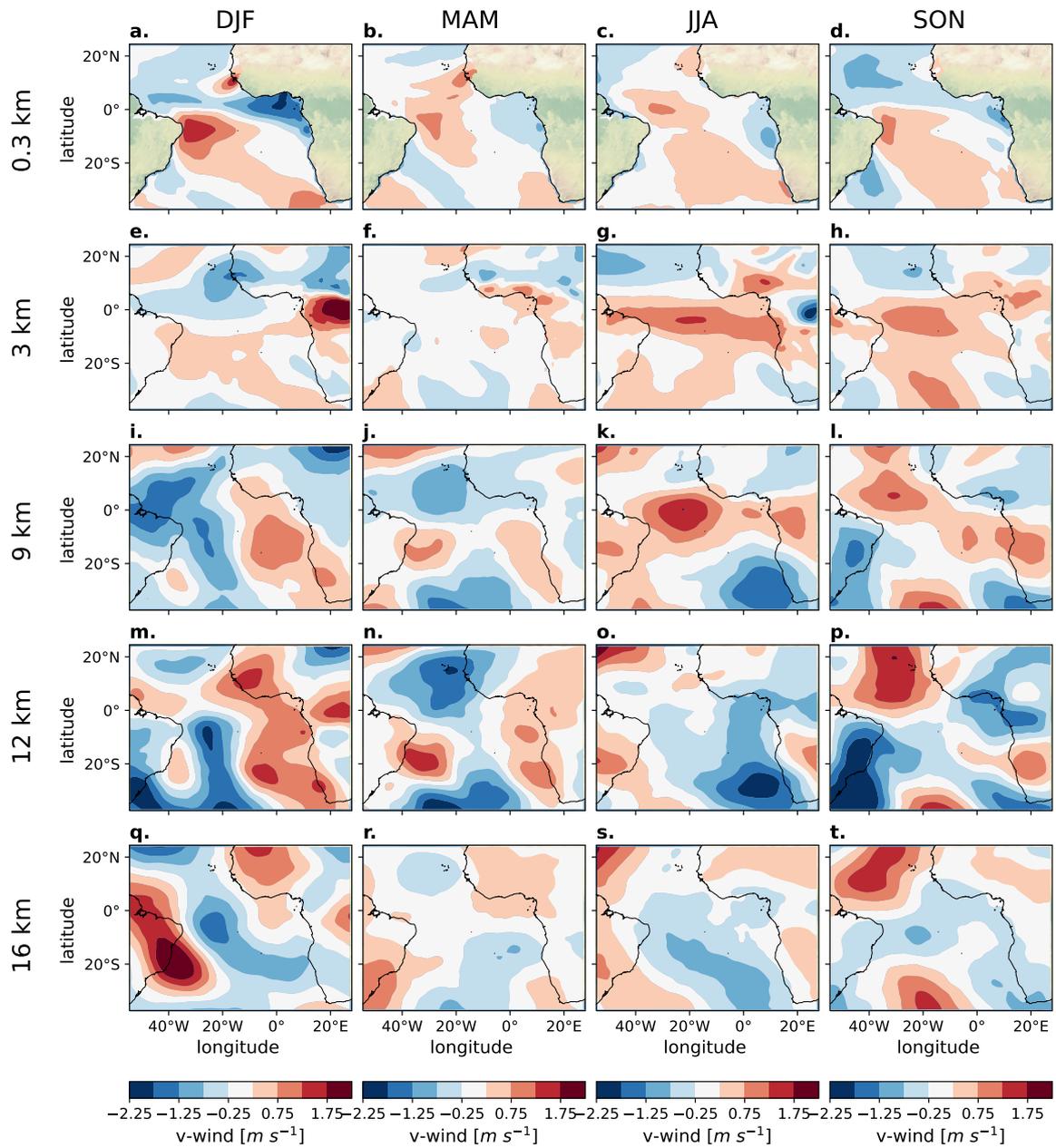


Figure S4. As Fig. S1 but shown for the meridional wind.

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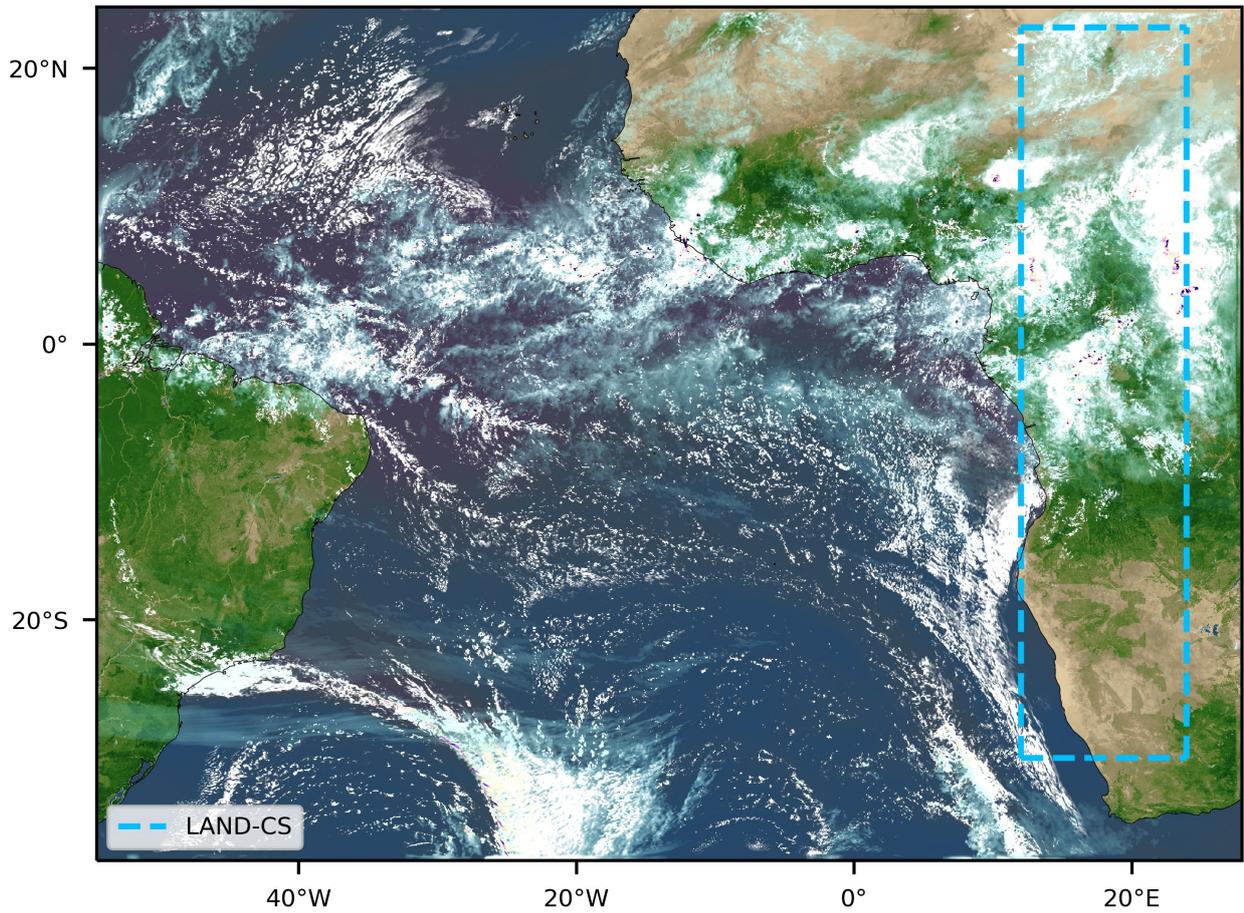


Figure S5. Simulation visualization and LAND-CS analysis domain used in supplementary Fig. S6.

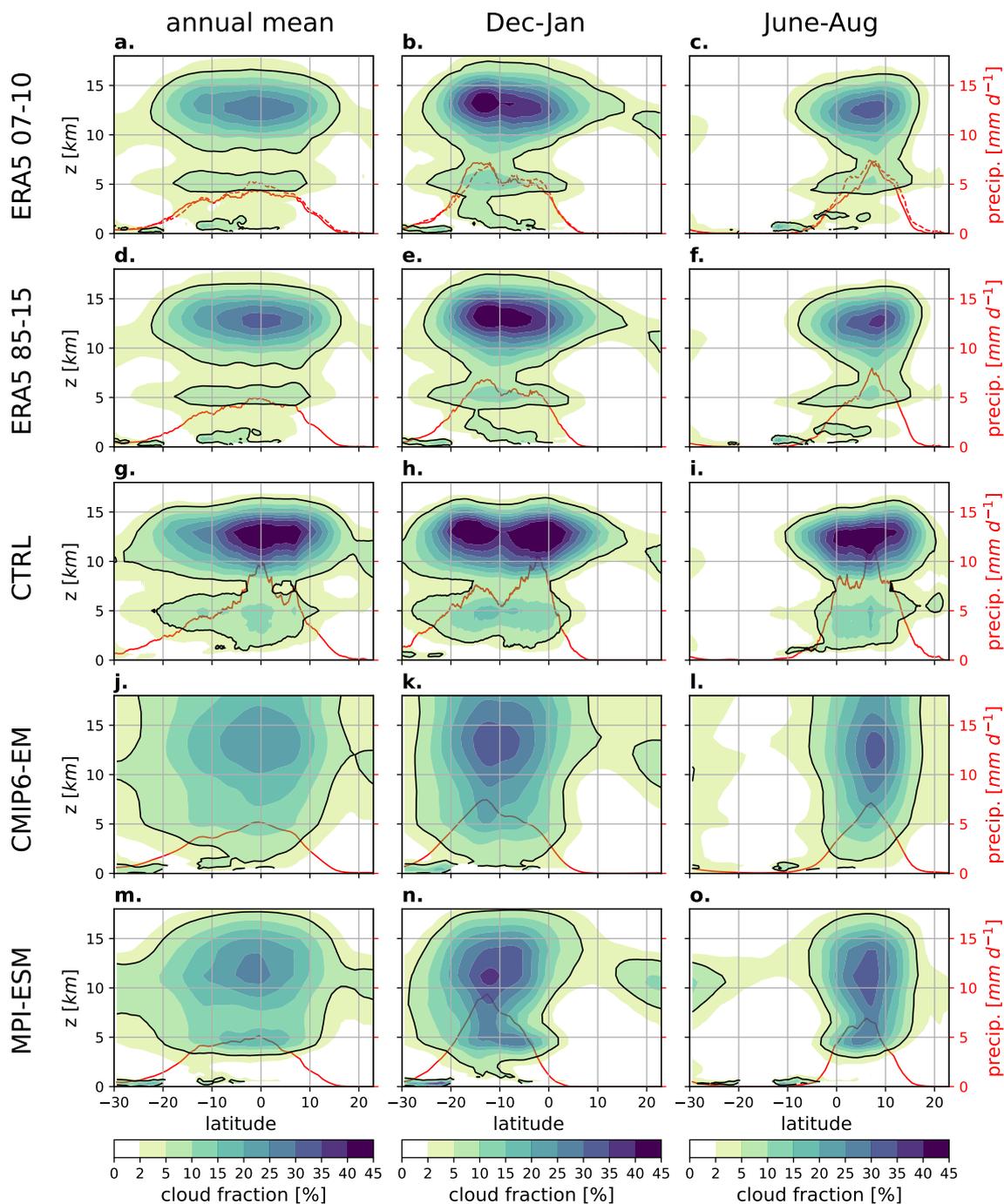


Figure S6. Altitude-latitude cross-sections of the cloud fraction [%] (green-to-blue contours) averaged along the longitudes of the LAND-CS domain. The black contour lines locate the 5% cloud fraction level. The solid red lines show surface precipitation [mm d^{-1}] represented on the right y-axis. The panels show (a-c) ERA5 (2007-2010), (d-f) ERA5 (1985-2014), (g-i) CTRL (2007-2010), (j-l) CMIP6-EM HIST (1985-2014), and (m-o) MPI-ESM HIST (1985-2014). The values represent multi-year averages (left panels) during the entire year, as well as (middle panels) during December, January, February and (right panels) June, July, August when the continental ITCZ reaches its southernmost and northernmost extent, respectively. As a complementary reference observation, GPM IMERG precipitation is shown as a red dashed line in panels (a-c).

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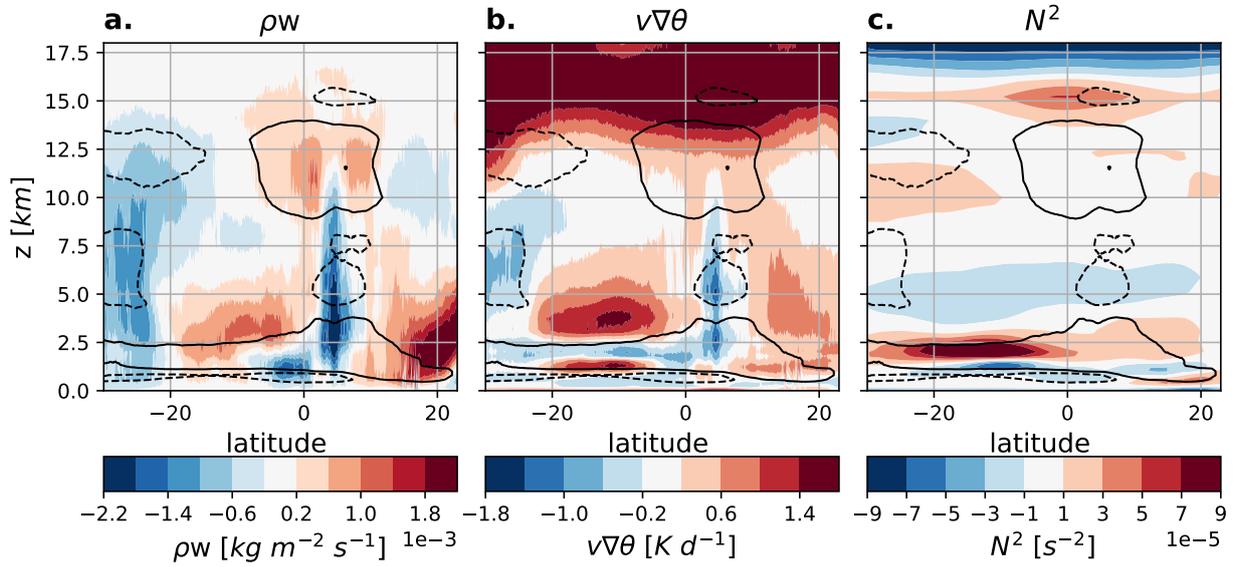


Figure S7. Altitude-latitude cross-sections of the difference CTRL-ERA5 (2007-2010) in (a) the vertical mass flux [$\text{kg m}^{-2} \text{s}^{-1}$], (b) the diabatic heating rate $\dot{\theta}$ [K d^{-1}] approximated by $\dot{\theta} \approx \mathbf{v} \cdot \nabla \theta$, and (c) the Brunt-Väisälä frequency [s^{-1}] along the longitudes of the HC-CS domain. The black contour lines indicate the 2% level of difference in the cloud fraction level where solid (dashed) lines represent a positive (negative) difference. The comparison of all panels reveals that the weaker subsidence in the lower subtropical free troposphere in CTRL compared to ERA5 results from a weaker diabatic cooling rather than increased stability. The stability of the lower free troposphere is lower in CTRL than in ERA5.