

Annually-resolved propagation of CFCs and SF6 in the global ocean over eight decades

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Abstract

Oceanic transient tracers, such as chlorofluorocarbons (CFCs) and sulfur-hexafluoride (SF6), trace the propagation of intermediate-to-abyssal water masses in the ocean interior. Their temporal and spatial sparsity, however, has limited their utility in quantifying the global ocean circulation and its decadal variability. The Time-Correction Method presented here is a new approach to leverage the available CFCs and SF6 observations to solve for the Green's functions describing the steady-state transport from the surface to the ocean interior. From the Green's functions, we reconstruct global tracer concentrations (and associated uncertainties) in the ocean interior at annual resolution (1940 to 2021). The spatial resolution includes 50 neutral density levels that span the water column along WOCE/GO-SHIP lines. The reconstructed tracer concentrations return a global view of CFCs and SF6 spreading into new regions of the interior ocean, such as the deep north-western Pacific. For example, they capture the southward spreading and equatorial recirculation of distinct NADW components, and the spreading of CFC-rich AABW out of the Southern Ocean and into the North Pacific, East Indian, and West Atlantic. The reconstructed tracer concentrations fit the data in most locations (~75%), indicating that a steady-state circulation holds for the most part. Discrepancies between the reconstructed and observed concentrations offer insight into ventilation rate changes on decadal timescales. As an example, we infer decadal changes in Subantarctic Mode Water (SAMW) and find an increase in SAMW ventilation from 1992 to 2014, highlighting the skill of the time-correction method in leveraging the sparse tracer observations.

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2 **global ocean over eight decades**

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7 **Key Points:**

- 8 • New “time-correction” method permits an annually-resolved global view of CFCs
9 and SF₆ over eight decades
- 10 • Steady-circulation solution is simultaneously consistent with atmospheric histo-
11 ries and 75% of the $\sim 10^6$ total CFC and SF₆ observations
- 12 • CFCs and SF₆, now detected in most of the global ocean, allow for a synopsis of
13 the interior ocean age, ventilation time, and variability

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Abstract

Oceanic transient tracers, such as chlorofluorocarbons (CFCs) and sulfur-hexafluoride (SF_6), trace the propagation of intermediate-to-abyssal water masses in the ocean interior. Their temporal and spatial sparsity, however, has limited their utility in quantifying the global ocean circulation and its decadal variability. The *Time-Correction Method* presented here is a new approach to leverage the available CFCs and SF_6 observations to solve for the Green's functions describing the steady-state transport from the surface to the ocean interior. From the Green's functions, we reconstruct global tracer concentrations (and associated uncertainties) in the ocean interior at annual resolution (1940 to 2021). The spatial resolution includes 50 neutral density levels that span the water column along WOCE/GO-SHIP lines. The reconstructed tracer concentrations return a global view of CFCs and SF_6 spreading into new regions of the interior ocean, such as the deep north-western Pacific. For example, they capture the southward spreading and equatorial recirculation of distinct NADW components, and the spreading of CFC-rich AABW out of the Southern Ocean and into the North Pacific, East Indian, and West Atlantic. The reconstructed tracer concentrations fit the data in most locations ($\sim 75\%$), indicating that a steady-state circulation holds for the most part. Discrepancies between the reconstructed and observed concentrations offer insight into ventilation rate changes on decadal timescales. As an example, we infer decadal changes in Subantarctic Mode Water (SAMW) and find an increase in SAMW ventilation from 1992 to 2014, highlighting the skill of the time-correction method in leveraging the sparse tracer observations.

Plain Language Summary

The penetration of chlorofluorocarbons (CFCs) and sulfur-hexafluoride (SF_6) into the oceans represents an opportunity to estimate the ventilation rate of the global ocean more directly than other seawater properties. Properties like temperature and salinity are nearly in balance with the ocean circulation and thus are challenging for finding rates of motion, but CFCs and SF_6 are transiently evolving due to a great increase in atmospheric concentrations since 1940. However, the analysis of CFCs and SF_6 also poses challenges, as they are observed infrequently and they have not had time to permeate the entire global ocean. Here, we analyze over $\sim 10^6$ CFC and SF_6 observations that have been taken over four decades, permitting the largest fraction to date of the global ocean to be analyzed using these tracers. A time-correction method is developed to address the temporal sparsity of the observations and an eight-decade annually-resolved picture of CFCs and SF_6 evolution is produced. Roughly 75% of the observations can be explained by the large-scale, statistically-steady ocean circulation acting on the anthropogenically-driven time varying atmospheric concentration.

1 Introduction

The formation and sinking of intermediate-to-abyssal waters, their pathways in the ocean interior and the amount of time they remain sequestered below the mixed layer are key aspects of the ocean circulation that regulate the exchange of heat, anthropogenic carbon and other tracers between the atmosphere and the deep ocean. Since the 1930s, the atmospheric concentrations of chlorofluorocarbons (CFCs), including CFC-11 and CFC-12, and sulfur-hexafluoride (SF_6) have varied significantly following their usage as industrial compounds (Figure 1). Their time-dependent atmospheric input histories are well known and they are conservative in the ocean interior, two characteristics that make them an excellent tool to trace water formation and pathways and to infer ocean ventilation rates (J. Bullister & Weiss, 1983; Smethie Jr, 1993; Rhein, 1994; Orsi et al., 1999; Rhein et al., 2004; Smethie Jr & Jacobs, 2005; LeBel et al., 2008; Tanhua et al., 2009; Rhein et al., 2015; Rivaro et al., 2015; Purkey et al., 2018). Furthermore, transient tracers have been used to quantify the sequestration and spreading of other climatically im-

portant tracers, such as anthropogenic carbon (Sarmiento & Sundquist, 1992; Matear & McNeil, 2003; Khatiwala et al., 2009; Fine, 2011; Ríos et al., 2012; Murata et al., 2019; Mahieu et al., 2020).

CFCs and SF₆ have been measured in the ocean since the 1980s along the major hydrographic sections (Gouretski & Koltermann, 2004; Hood et al., 2010) and used to infer deep-water formation rates and pathways in the Southern Ocean (Orsi et al., 1999, 2001, 2002; Smethie Jr & Jacobs, 2005; Rivaro et al., 2015) and North Atlantic (Weiss et al., 1985; Smethie et al., 2000; Smethie Jr & Fine, 2001; Doney & Bullister, 1992; Azetsu-Scott et al., 2003; LeBel et al., 2008; Tanhua et al., 2009; Rhein et al., 2015). In addition, a number of previous studies have used CFCs to constrain how ocean properties at the surface are connected with those in the ocean interior. This connection is sometimes represented by a Transit Time Distribution, i.e., the probability density function of the transit times since a water parcel was last in the surface mixed layer. The Transit Time Distribution defines the water age at a specific location by representing the advection, mixing and turbulent diffusion of the water parcels that start somewhere at the surface and get to a given interior location. It is often referred to as age spectrum, age distribution, transit-time probability density function, boundary propagator or the Green's function for boundary conditions (Hall et al., 2002). Hereafter, we refer to the Green's function (GF) as the distribution connecting a density class surface points with the ocean interior.

The GF is the solution of an underdetermined problem, and thus requires information beyond that from transient tracer observations. One solution method is found by requiring the shape of the GF to be an Inverse Gaussian function described by just two parameters (mean and width), and the resulting solution is consistent with advection and diffusion in one or two dimensions (Waugh et al., 2003, 2004; Hall et al., 2002, 2007). The assumption of an Inverse Gaussian, however, does not strictly hold in three dimensions when disparate water masses mix (Trossman et al., 2014). The Maximum Entropy method addresses this issue by adding an additional entropy constraint that renders the solution unique while providing as little non-observational information to the problem as possible (Khatiwala et al., 2009; Holzer et al., 2010; Khatiwala et al., 2012), but at the cost of a large computational burden. The limitations listed above call for a method that minimizes the effect of using an Inverse Gaussian as initial guess while analyzing a few hundreds thousand observations simultaneously and while preserving computational efficiency.

This study presents a novel method that solves for the GF and reconstructs the time evolution of CFCs and SF₆ concentrations along hydrographic sections in the global ocean. With now four decades of data, for any given year in which CFCs or SF₆ observations are available, they provide useful information that can be used to constrain the ocean circulation. The result is a steady-state GF that minimizes the misfit between reconstructed and observed tracer concentrations. This solution incorporates not only the large-scale geostrophic flows that can be inferred from inverse solutions (e.g. Reid (1994, 1997, 2003); Talley et al. (2003); Lumpkin and Speer (2007)) but also the net effect of any small-scale advective and diffusive transports. Moreover, the TCM allows us to quantify the reconstruction uncertainty, which have often not been accounted for in inverse solutions (e.g Reid (1994, 1997, 2003)). Finally, through the assumption of a steady-state GF, we can reconstruct tracer concentrations from the first time they entered the ocean in 1940s even if observations were not available. We will refer to this method as the *Time-Correction Method* (TCM), in the vein of Orsi et al. (1999) and as proposed by Purkey et al. (2018). Orsi et al. (1999) produced a map of CFC-11 concentrations by “correcting” CFC-11 observations over the 1984-1996 period, i.e. normalizing them to a mid-term year to minimize the temporal biases.

The annually-resolved reconstruction of tracer concentrations addresses the temporal sparsity of anthropogenic tracer ocean data, which poses large uncertainty in the

117 assessment of the timescales of physical processes and the ocean circulation variability.
 118 For example, formation rates inferred from tracer inventories are inevitably inaccurate
 119 because they need to assume observations are synoptic (e.g. Orsi et al. (1999); Willey
 120 et al. (2004)). Furthermore, only a few studies (Huhn et al., 2013; Waugh et al., 2013)
 121 to date have used CFCs to infer variability in ocean circulation and ventilation, partly
 122 owing to the temporal sparsity of observations, with most sections having decadal oc-
 123 cupations at best. With the TCM, we investigate how well a steady-state circulation ex-
 124 plains the observed tracer concentrations and gain insights into the decadal variability
 125 of water mass properties in the ocean interior.

126 The TCM theory and its skill (and error estimates) in reconstructing CFCs and
 127 SF₆ concentrations across the global ocean are presented. A description of the data used
 128 and the TCM theory are described in Sections 2 and 3, respectively. The results (Sec-
 129 tion 4) (i) test the TCM with synthetic data, (ii) present the time corrected CFCs ob-
 130 servations in the global ocean including a discussion of deep water pathways, and (iii)
 131 discuss the error and potential for this method to constrain decadal variability in ven-
 132 tilation rates. We conclude by discussing future applications of the TCM.

133 2 CFC and SF₆ observations

134 2.1 Atmospheric histories

135 Atmospheric concentrations of anthropogenic tracers CFC-11, CFC-12, and SF₆
 136 have been well monitored since the late 1970s, with reliable reconstructions extending
 137 the time series back to the 1930s (Figure 1a). CFC-11 and CFC-12 atmospheric concen-
 138 trations increased almost linearly between the late 1930s and the 1990s, from when they
 139 slowly started decreasing following the regulation by the Montreal Protocol in 1987 (Fig-
 140 ure 1a). Emissions of SF₆ started in 1950s and have been increasing steadily since, with
 141 a linear increase since mid-1980s (Figure 1a). The atmospheric concentrations of CFC-
 142 11 and CFC-12 have been directly measured since 1979, and they were reconstructed prior
 143 to that by knowing the industrial production data and the atmospheric lifetimes of their
 144 compounds, with uncertainties being a few percent (Prinn et al., 2000; Walker et al., 2000;
 145 Fine, 2011). Similarly, atmospheric concentrations of SF₆ have been measured since 1953,
 146 and reconstructed by Maiss and Brenninkmeijer (1998) prior to that (Fine, 2011). Here,
 147 we use the annual global mean reconstructed atmospheric histories by J. L. Bullister (2015)
 148 up to 1977 (1995 for SF₆) and the measured atmospheric values from the NOAA Halo-
 149 carbons and other Atmospheric Trace Species (HATS) monitoring laboratory (<https://gml.noaa.gov/hats/>)
 150 through 1978-2021.

151 2.2 Oceanic observations

152 A global data set of hundreds of thousands of high quality oceanic dissolved tracer
 153 measurements has been taken since the late 1970s. Many of these observations were taken
 154 along repeated transoceanic full depth hydrographic sections, occupied originally as part
 155 of the World Ocean Circulation Experiment (WOCE, Gouretski and Koltermann (2004)),
 156 then repeated roughly once every 5-10 years thereafter under the Global Ocean Ship-
 157 Based Hydrographic Investigations Program (GO-SHIP, Hood et al. (2010)). The loca-
 158 tions of the hydrographic sections discussed or mentioned throughout the manuscript
 159 are shown in Figure 1b. Data used here is sourced primarily from the second update of
 160 the Global Ocean Data Analysis Project (GLODAPv2.2021) data compilation, which com-
 161 piles in a single product most of the available shipboard measurements (Lauvset et al.,
 162 2021). The GLODAPv2.2021 data is quality controlled but no adjustments have been
 163 made to the transient tracer observations. We add data collected during the AnSlope
 164 program along the Ross Shelf break (Gordon et al., 2009). Transient tracer observations
 165 are accurate within a 5% error (Lauvset et al., 2021) with a CFCs detection limit of 0.01-
 166 0.001 pmol/kg and SF₆ detection limit of 0.1-0.01 fmol/kg (Stöven et al., 2015).

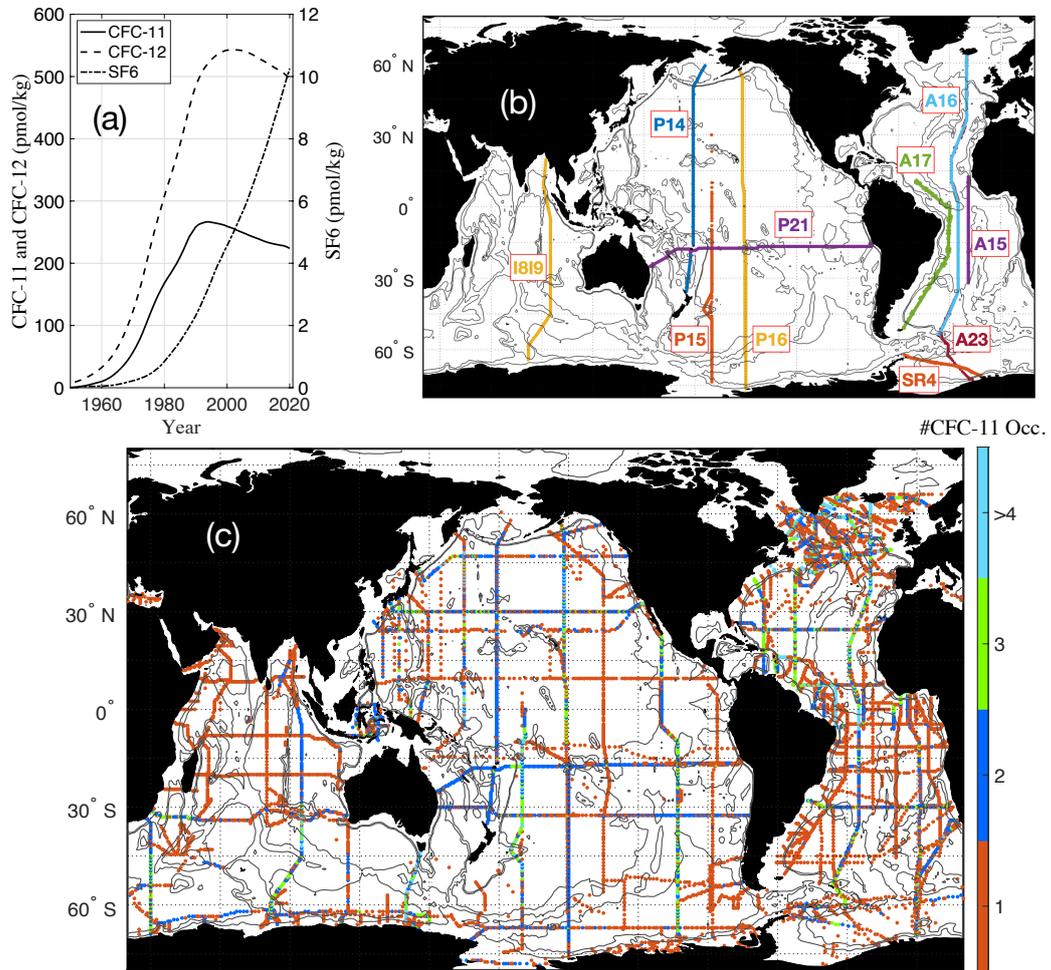


Figure 1. (a) Atmospheric history of chlorofluorocarbon 11 (CFC-11; solid), CFC-12 (dashed) and sulfur-hexafluoride (SF₆; dashed-dotted). (b) Location of key WOCE/GO-SHIP cruises referenced within this manuscript. (c) Number of years with at least one CFC-11 measurement within each 1/2 degree grid cell in waters denser than $\gamma^n = 27$. Thin gray lines indicate the 3000 and 4000 m isobaths.

167 CFCs and SF₆ observations are distributed globally (Figure 1c). Many sections have
 168 been sampled only once, yet a valuable number of observations are available in all ocean
 169 basins. In particular, our dataset includes 21,793 and 23,391 casts with CFC-11 and CFC-
 170 12 observations, respectively, and 5,724 casts with SF₆ observations. Most stations have
 171 24-36 vertical bottles with higher sampling in the upper ocean. Bottle samples near the
 172 seafloor are typically 100 m apart but vary by cruise. Although the method presented
 173 here can be applied to any water mass, Figure 1c shows the locations and number of CFC-
 174 11 occupations of the observations below neutral density $\gamma^n = 27$ (Jackett & McDougall,
 175 1997), which lies at ~ 500 m (global average) and approximately separates the shallower
 176 thermocline waters from the deep ocean.

177 3 Theory

178 At each location, the tracer surface boundary conditions (C_s) and the tracer con-
 179 centration in the ocean interior (C) are given by

$$180 C_s(\mathbf{r}_s, t) = \int_0^\infty C_{atm}(t - \tau_e) \mathcal{E}(\mathbf{r}_s, \tau_e) d\tau_e, \quad (1)$$

181 and

$$182 C(\mathbf{r}, t) = \int_0^\infty C_s(\gamma^n, t - \tau) \mathcal{G}(\mathbf{r}, \tau) d\tau, \quad (2)$$

183 where all terms are detailed below. Equation (1) returns the tracer surface boundary con-
 184 ditions at location \mathbf{r}_s , where C_{atm} is the tracer atmospheric concentration and τ_e is the
 185 air-sea equilibration time. The Equilibration Time Distribution (ETD), $\mathcal{E}(\mathbf{r}_s, \tau_e)$, describes
 186 the partitioning of the surface waters according to when they were last in equilibrium
 187 with the atmospheric tracer concentration, taking into account the lag-time in surface
 188 equilibration due to mixed layer dynamics (see Section 3.1).

189 Equation (2) describes the propagation of the surface tracer concentration in the
 190 ocean interior location \mathbf{r} , where τ is the surface-to-interior transit time and $(t - \tau)$ is
 191 when a fluid element was last in the surface mixed layer. $\mathcal{G}(\mathbf{r}, \tau)$ is the GF, which par-
 192 titions the concentration in the ocean interior according to the time when its fluid el-
 193 ements were last in the surface mixed layer. The surface boundary conditions are cal-
 194 culated per each density surface, $C_s(\gamma^n, t - \tau)$, as the average surface condition at the
 195 density level outcrop points (see Section 3.1).

196 Wherever CFCs or SF₆ observations are available, Equations (1) and (2) allow for
 197 the reconstruction of the time-dependent interior tracer concentration $C(\mathbf{r}, t)$. The next
 198 two Sections describe how to solve for the probability density functions \mathcal{E} and \mathcal{G} .

199 3.1 Solving for the Equilibration Time Distribution

200 Tracer saturation of the surface mixed layer is temporally variable both in the short
 201 and long term; changes in temperature (and thus in solubility), sea-ice formation, and
 202 mixed layer processes control the seasonal variations of saturation values, while the long-
 203 term growing atmospheric concentration leads to a lag in the equilibration of the sur-
 204 face mixed layer (Rodehacke et al., 2010; Shao et al., 2013). For example, Rodehacke et
 205 al. (2010) showed that ignoring the temporal trends in CFC saturation values can lead
 206 to up to 10% error in the inventory-based ventilation time-scales.

207 The ETD partitions the surface tracer concentration according to when it was last
 208 in equilibrium with the atmosphere. We parameterize the ETD with a gamma distribu-
 209 tion, following DeVries and Primeau (2010), that defines \mathcal{E} as

$$210 \mathcal{E}(\mathbf{r}_s, \tau_e) = \tau_e^{\theta_e/\beta_e - 1} \beta_e^{-\theta_e/\beta_e} \left[\Gamma\left(\frac{\theta_e}{\beta_e}\right) \right]^{-1} \exp\left(-\frac{\tau_e}{\beta_e}\right), \quad (3)$$

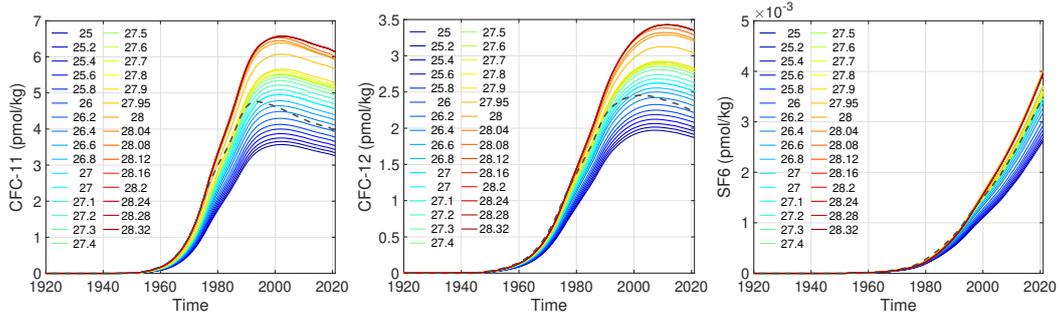


Figure 2. Surface boundary conditions for several density classes for (left) CFC-11, (middle) CFC-12 and (right) SF₆ calculated with the ETD solution (Equation 3). The gray dashed line represent the surface boundary condition obtained with a fixed saturation level (92% and 80% for CFCs and SF₆, respectively) and a fixed solubility calculated with $S = 35$ and $T = 5^\circ\text{C}$. Note the different scale of the panels.

211 where Γ is the Gamma function, θ_e is the mean value of the ETD, and β_e is the ratio
 212 of the variance of ETD over the mean.

213 We solve for θ_e and β_e values that minimize the discrepancy between the estimated
 214 surface boundary conditions and the tracer observations in the mixed layer. In partic-
 215 ular, first the maximum tracer concentration is calculated within the surface mixed layer,
 216 defined by using the maximum mixed layer depth (MLD) in the global Argo mixed layer
 217 climatology (Holte et al., 2017). This concentration reflects the surface equilibration time
 218 under the assumption that winter waters can hold the greatest concentration of the trac-
 219 ers and are preferentially subducted. Such a step is necessary because the seasonal cy-
 220 cle is not resolved by the sparse data. These observations are then used to invert for θ_e
 221 and β_e at each location, hence to calculate ETD (Equation 3) and the surface bound-
 222 ary conditions C_s (Equation 1) at each surface location \mathbf{r}_s (wherever tracer data are avail-
 223 able).

224 For each neutral density surface, the average surface boundary condition is calcu-
 225 lated by averaging C_s in the outcrop points. For a density level γ_*^n , the outcrop points
 226 are defined as all the locations where γ_*^n is lighter than the density at the maximum MLD
 227 (see supplementary Figure S1). Note that for waters denser than $\gamma^n = 28.20$, the num-
 228 ber of outcrop points with available tracer data rapidly drops (from $\sim 10^2$ points to $<$
 229 10), and the same surface condition as the closest density layer available is used.

230 The ETD-based surface boundary conditions used here span a wide range of val-
 231 ues (Figure 2). This range is explained by the strongly temperature-dependent solubil-
 232 ity used to calculate the surface boundary condition at each location. Solubility is cal-
 233 culated using the average observed temperature and salinity between the surface and the
 234 maximum MLD, and using the solubility coefficients and methods reported in Warner
 235 and Weiss (1985) for CFCs and J. L. Bullister et al. (2002) for SF₆. Moreover, the roll-
 236 off of CFC-11 and CFC-12 surface boundary conditions, following their decrease in at-
 237 mospheric histories, is delayed in the ETD-based solution when compared with a solu-
 238 tion based on a fixed saturation and solubility (dashed gray line in Figure 2). This de-
 239 lay confirms that the ETD approach takes into account the adjustment time of the mixed
 240 layer.

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3.2 Solving for the Green's function

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The TCM is based on solving a non-negative, weighted least-squares problem to find the Green's function, which forces the solution to be strictly positive. Note that we solve for the discrete (annually-resolved) Green's function, \mathbf{G} , which is constrained with CFCs and SF₆ up to the age that can be inferred from transient tracers (~ 80 years ago). We use a first-guess of the GF, \mathbf{G}_0 , to constrain our solution for waters older than 80 years, up to 3000 years (see next paragraph).

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At each location, the least-squares problem is solved for the optimal \mathbf{G} that fits the existing GLODAPv2.2021 data when convolved with the atmospheric history C_s (see Supplementary Information for more details about the least-squares method used here). We solve for the solution that minimizes the cost function

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$$J = \|\mathbf{C} - C_s \mathbf{G}\|_2^2 + \|\mathbf{G} - \mathbf{G}_0\|_2^2 = \quad (4)$$

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$$= (\mathbf{C} - C_s \mathbf{G})^T W^{-1} (\mathbf{C} - C_s \mathbf{G}) + (\mathbf{G} - \mathbf{G}_0)^T S^{-1} (\mathbf{G} - \mathbf{G}_0). \quad (5)$$

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At each location, \mathbf{C} is the vector of available tracer observations, C_s is the matrix of the corresponding surface boundary conditions, and \mathbf{G}_0 is the first-guess Green's Function. Equation (4) indicates that we look for the solution that (i) minimizes the discrepancy between the GLODAPv2.2021 tracer observations \mathbf{C} and the predicted concentration, i.e. the convolution of C_s and \mathbf{G} , and (ii) minimizes the discrepancy between the first guess \mathbf{G}_0 and the solved function \mathbf{G} . W^{-1} and S^{-1} are the matrices of the weights associated with each one of the terms, respectively (see Supplementary Information for more details).

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In particular,

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1. W is a diagonal matrix, with $1/\max(5\% \mathbf{C}, C_{lim})^2$ on the diagonal. C_{lim} is the observation detection limit (Section 2), for which we choose 0.01 pmol/kg for CFCs and 0.1 fmol/kg for SF₆, a conservative choice of the limits presented in Stöven et al. (2015). 5% of \mathbf{C} represents a 5% error, which is the measurement error associated with transient tracer observations (Lauvset et al., 2021).
2. Our first guess of the GF, \mathbf{G}_0 , is an Inverse Gaussian whose characteristics are based on the solution of the Total Matrix Intercomparison (TMI) approach described in Gebbie and Huybers (2010, 2012). In the TMI method, several tracers (temperature, salinity, nitrate, phosphate, silica and oxygen isotope) are used to invert for the ocean circulation and radiocarbon is used to infer water ages. For clarity, radiocarbon is used only in the first guess \mathbf{G}_0 and not to deconvolve our final GF. Given that our final GF uses 'young' tracers, our solution tends to relax to the first guess GF for ages older than the eight decades constrained by CFCs and SF₆.
To construct the first guess GF at each location, we use the water age estimated by the TMI method (Γ_{TMI}) to produce an ensemble of Inverse Gaussians: the ensemble is generated by varying Γ_{TMI} with a normal distribution, assuming a factor of 2 of confidence of Γ_{TMI} and a ratio $\Delta/\Gamma = 1$, as consistent with tracer observations (Waugh et al., 2003, 2004; Hall et al., 2004). Errors in Γ_{TMI} are probably no larger than 100-200 yrs for the oldest waters and less for younger waters (Gebbie & Huybers, 2012), but here we conservatively assume that Γ_{TMI} is uncertain to a factor of 2 to avoid over-constraining our solution of \mathbf{G} . The first guess \mathbf{G}_0 is an Inverse Gaussian calculated as the ensemble average. This approach is summarized in Supplementary Figure S2.
3. The level of confidence of the first guess \mathbf{G}_0 is given by the covariance of the Inverse Gaussians ensemble. It follows that $S^{-1} = (\mathbf{G}' \mathbf{G}'^T)^{-1}$, where $\mathbf{G}' = \mathbf{G}_i - \langle \mathbf{G}_i \rangle = \mathbf{G}_i - \mathbf{G}_0$, \mathbf{G}_i is the i -th Inverse Gaussian of the ensemble, and $\langle \mathbf{G}_i \rangle = \mathbf{G}_0$ is the ensemble average. In solving for \mathbf{G} , we do not want our first guess to constrain the solution to a prescribed shape. In other words, our priority is to

293 fit the available observations rather than fitting \mathbf{G}_0 (Equation 4). Therefore, if
 294 our solution does not fit the observations within a 5% error, the weight of the first
 295 guess S is reduced until the mismatch error is reduced within the 5% error limit
 296 or when the data fitting does not improve anymore.

297 In summary, for each location where tracer concentrations are available, we make
 298 use of knowing (i) the tracer observations \mathbf{C} and (ii) the surface boundary conditions C_s ,
 299 to solve for the boundary GF \mathbf{G} at annual resolution.

300 3.3 Applying the TCM to oceanic tracer data

301 At each station, the GLODAP oceanic tracer data (Section 2) is first vertically in-
 302 terpolated onto a uniform neutral density grid with spacing of $\Delta\gamma^n = 0.01$. The inter-
 303 polation uses a piecewise cubic Hermite interpolating polynomial and depth-separation
 304 limits to avoid interpolating between points too far away for the interpolated value to
 305 be deemed acceptable, following Key et al. (2010) (see their Table 4; depth-separation
 306 limits used here are the more conservative values reported for the Arctic Ocean). All sta-
 307 tions taken within one year and within a single $1/2^\circ$ horizontal resolution WOCE grid
 308 are binned by taking the mean at each density level.

309 The TCM is applied to every $1/2^\circ \times 1/2^\circ$ grid box and density level where at least
 310 one transient tracer observations exists (Figure 1c). Note this analysis does not return
 311 a three-dimensional gridded product yet, but the spatial interpolation in locations where
 312 tracer data is missing is part of ongoing work. All available tracer observations at each
 313 grid box is used to constrain our solution for GF (see details in the Supplementary Text
 314 S1). Then the GF is used to reconstruct the time-dependent tracer concentration in the
 315 ocean interior following Equation (2) by convolving the GF with the surface boundary
 316 conditions resulting in an estimate of the CFC-11, CFC-12 and SF₆ concentration for
 317 each year between 1940 and 2021. Error estimates of the time-corrected tracer concen-
 318 trations are also calculated as 95% confidence limits.

319 4 Results

320 4.1 Testing TCM skill with synthetic data

321 Synthetic data is used to test the TCM performance in (i) deconvolving the GF
 322 and (ii) predicting the tracer concentration and its associated error, and to analyze its
 323 sensitivity to the number of observations and tracers available. We start by assuming
 324 a known true GF, for which we use an Inverse Gaussian (Figure 3, black curve) and a
 325 bimodal function generated by combining two different Inverse Gaussian functions (not
 326 shown but results are qualitatively similar to Figure 3). In these tests, we assume a fixed
 327 surface boundary condition, as represented by the gray dashed lines in Figure 2 and use
 328 the known GF to predict the *true* tracer concentrations in the ocean interior (thick col-
 329 ored lines in the right sub-panels). To test the TCM, we sub-sample the true tracer con-
 330 centrations, add random error (up to 5%) that represents contamination of the signal
 331 by noise, and deconvolve the GF. We use an initial-guess GF as described in Section 3.2
 332 (dashed line). The resulting estimated \mathbf{G} (dash-dotted line) and predicted tracer con-
 333 centrations (thin colored curves) for various sampling scenarios are shown in Figure 3.

334 The TCM reproduces the known tracer concentrations within error for all tested
 335 scenarios (Figure 3). We start by showing the limiting case of only one tracer (CFC-11)
 336 observation at one time (panels a and b). The TCM prediction fits the available (con-
 337 taminated) observation perfectly, i.e. tends to overfit the solution given the underdeter-
 338 mined nature of the problem. The TCM-solved GF is smooth and very close to fitting
 339 the true GF, indicating that the first guess is helping constrain the solution. The tracer
 340 prediction is also consistent with the true tracer concentration for all the tracers, includ-

341 ing the tracers not used to solve the GF (CFC-12 and SF₆). The prediction is closest
 342 to the truth near the observation, with a slight underestimate toward the end of the time
 343 period. However, the true concentrations are within the errorbar limits of the TC solu-
 344 tion at all times. Note that even the error estimates (95% confidence limits) of CFC-12,
 345 which is not used in this scenario, are nicely constrained by the CFC-11 error estimates
 346 due to the similar atmospheric histories of these transient tracers.

347 The second scenario, again tests the method with only one observation, but now
 348 considers when the observation is only available either early (blue) or late (red) in the
 349 time frame (Figure 3b). The TCM's GFs are only slightly different and the tracer pre-
 350 dictions are within the error bars of the true concentration in both cases. As in the first
 351 scenario, the error bars become larger the further away from the time of the observation.
 352 The bias between the reconstructed and 'true' tracer concentrations, i.e. whether the re-
 353 construction is larger or smaller than the truth, is determined by the sign of the contam-
 354 ination of the 'true' observations and by the fixed surface boundary conditions. In the
 355 examples of Figure 3b, when the observation is available only late in the predicted time
 356 frame, the method slightly overestimates early tracer values; and when the observation
 357 is only available early, the method slightly underestimates true concentrations toward
 358 the end of the time frame.

359 In scenario three, again we allow only observations at one instance, but provide the
 360 TCM with three tracers (CFC-11, CFC-12, and SF₆) (panel c). Having all three tracers
 361 at the same time improves the predictive power of the method and reduces the error es-
 362 timates, in particular for the newly added tracer (SF₆ in this example, compare the er-
 363 rorbars in panels a and c).

364 Finally, in scenario 4, we show that the method's skill improves as we increase the
 365 number of available observations (panel d). When two (or more observations, not shown)
 366 samples are taken 1 decade or more apart, error bars are reduced by 30-50% (panel d
 367 VS a), showing increased level of confidence in our reconstruction.

368 4.2 Reconstructed transient tracer distributions

369 The reconstructed global data set of annual concentrations of CFC-11, CFC-12, and
 370 SF₆ paints a more detailed picture of ocean ventilation and deep water circulation than
 371 previously possible with observations alone. High CFC-11, CFC-12, and SF₆ values are
 372 found in the upper ocean due to the effective ventilation by the wind-driven circulation,
 373 as expected. Regions of the deep and abyssal waters also have high values, especially those
 374 that have been in contact with the atmosphere within the last decades. Analyses of the
 375 time varying concentrations following these CFC-rich water-masses meridionally provide
 376 direct observations of time scales of ocean transport. In addition, the solved GFs allow
 377 for direct assessment of the mean age and age distribution of interior water throughout
 378 the global ocean. We focus on a global description of the data set, highlighting its key
 379 scientific applications. Given the assumptions in the boundary conditions, this product
 380 is expected to best reproduce large-scale patterns, although the same method could be
 381 applied in the future to regional problems.

382 The reconstructed tracer concentrations along three WOCE meridional sections that
 383 span the Atlantic (A16+A23), Pacific (P16) and Indian (I08S+I09N) are used to demon-
 384 strate the time evolution of the tracers meridionally and to allow direct comparison be-
 385 tween basins by adjusting GO-SHIP data to the same year. Along each of the sections,
 386 grid boxes nominally following the cruise tracks were selected (see Figure 1b for their
 387 location). For each year, the reconstructed values of CFC-11, CFC-12 and SF₆ were
 388 objectively mapped in latitude-depth space using decay scales of 40 km and 1 km in the
 389 horizontal and vertical directions, following Roemmich (1983). The reconstruction of CFC-
 390 11 concentration through 2021 is shown in Figure 4, while complementary reconstruc-

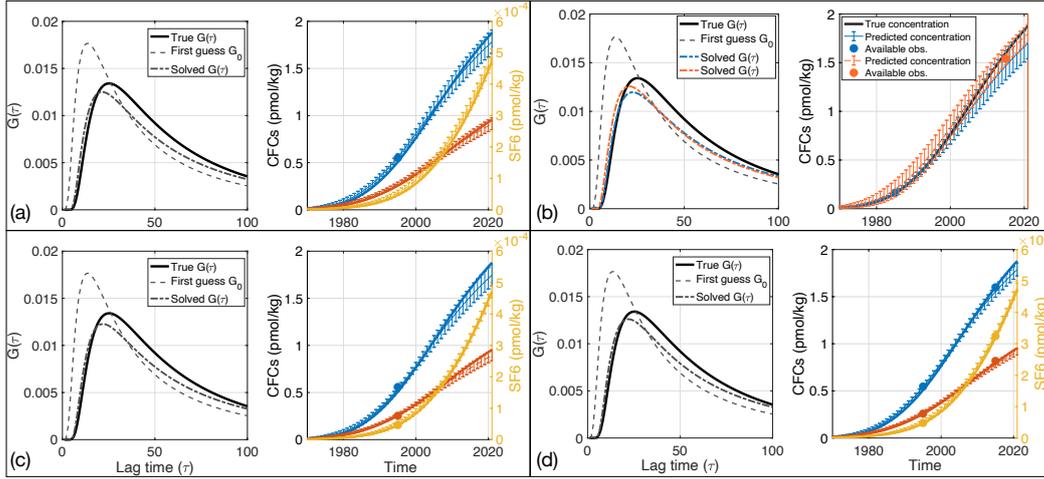


Figure 3. Results from four TCM performance tests with synthetic data. Per each sub-panel: right panels show the GFs, including the True GF (black), first guess \mathbf{G}_0 (dashed gray line) and predicted GF solved using the TCM (dash-dotted black line); left panels show the tracer concentrations, CFC-11 (blue), CFC-12 (red), SF₆ (yellow). Thick tracer lines are the true tracer concentrations and the thin lines are the predicted values using the TCM with associated error estimates (95% confidence limits). Colored dots indicate the sub-sampled “observations” used for the TCM prediction.

391 tions of CFC-12 and SF₆ concentrations are shown in Supplementary Figures S3 and S4,
 392 respectively.

393 **4.2.1 Antarctic Abyssal Water**

394 Antarctic Bottom Water (AABW) is primarily produced along the Antarctic shelf
 395 where dense shelf waters entrain ambient waters as it flows down the continental slopes
 396 (e.g. Jacobs et al. (1985); Gordon (2019)), and is returned to shallower depth via cross-
 397 density mixing, which leads to a net dense-to-light water mass transformation. High tracer
 398 values are observed at the bottom near and directly downstream of all four AABW pro-
 399 duction regions (Figure 5). The Weddell Sea, one of the primary sites of export of dense
 400 shelf waters, is highlighted by the large CFC-11 values at 70-55°S in the Atlantic Ocean
 401 along the northern edge of the gyre where recently formed water on the shelf are advected
 402 clockwise around the gyre (Figure 4a and Figure 5; (Orsi et al., 1999)). Somewhat large
 403 values are also found in the southern edge of the gyre located east of the primary AABW
 404 formation region, indicating newly formed deep waters either coming from further east
 405 near Cape Darnley or the return of the Weddell Sea Deep Waters having made a full cir-
 406 cle around the gyre, while minimum CFC-11 concentrations are found in the middle of
 407 the gyre. The reconstructed CFC-11 concentrations along SR03 to the east(not shown,
 408 see Figure 1b for location) include values up to 0.5 pmol/kg everywhere below 2000 m
 409 across the gyre, and even larger values along the northern boundary, confirming the sig-
 410 nificance of the gyre circulation in transporting AABW, as shown by the deep pathways
 411 suggested in Figure 4a.

412 As AABW flows northward in all ocean basins, it is strongly constrained by topog-
 413 raphy and often channeled through passages or fracture zones. In the South Atlantic,
 414 CFC-11-rich AABW has propagated northward all the way to the equator. The signal
 415 in Figure 4a is clear only up to 10°N, but A16 does not well capture the core of the north-
 416 ward flow along the west boundary. The northward propagation of CFC-11 within AABW

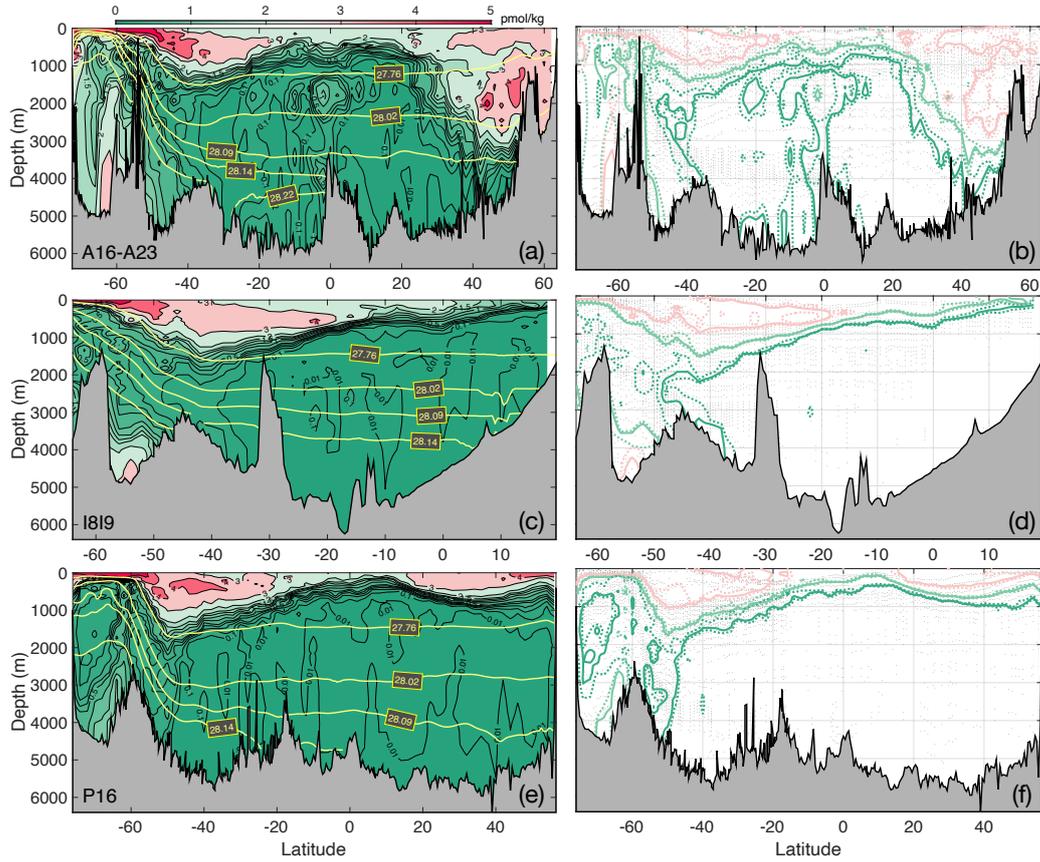


Figure 4. Reconstructed CFC-11 concentrations along A16+A23 (panels a,b), I819 (c,d), and P16 (e,f) in 2021. Panels in left column show the reconstructed concentration and neutral density contours (yellow) along each section. Panels in right column show the observation locations throughout all decades of available data (gray dots) and the uncertainty of the time corrected CFC-11 concentrations for three contours (0.2, 1 and 3 pmol/kg). The solid line shows our best estimate, while the dotted lines show the 95% confidence limits.

417 density layers is better seen along A17 (not shown), although with some more noise due
 418 to the low number of observations available along this section, in particular at depth.
 419 In previous estimates from 1990s data (Orsi et al., 2002), the detection limit contour of
 420 0.01 pmol/kg did not extend farther than 40-35°S, where now we reconstruct values up
 421 to 0.15-0.2 pmol/kg. The only locations where CFC-11 reconstruction is predicted to be
 422 below detection limit in 2021 within the AABW layer in the Atlantic Ocean is east of
 423 the Mid-Atlantic ridge in the very north and southern parts of the basin (Figures 4a and
 424 5). Here, the first few hundred meters above the seafloor are mainly occupied by AABW
 425 (G. C. Johnson, 2008), and hence our result is consistent with a slow abyssal circulation
 426 in the eastern Atlantic after crossing the Romanche Fracture Zone. This result is high-
 427 lighted by the CFC-11 bottom arrival time i.e., the year in which CFC-11 concentrations
 428 first exceeded the detection limit at the ocean bottom (Figure 5).

429 In the Indian Ocean, high CFC-11 concentrations are found in the Princess Eliz-
 430 abeth Trough and Antarctic-Australian Basin, which I819 crosses between 50-80°E. Here,
 431 our reconstruction captures high CFC-11 both on the southwest and northeast flanks of
 432 the basin, consistent with Ross-formed and Adelie-formed AABW flowing in from the
 433 southeast (i.e. directed northwest), looping within the basin and flowing back southeast
 434 (Figure 5). We also find values well above detection limit up to 30°S, indicating that CFC-
 435 11-rich AABW waters have crossed the Antarctic-Australian Discordance at 50°S and
 436 are moving northward along the Southeast Indian ridge. Further north, our reconstruc-
 437 tion displays minimum values between 2000-4000 m, and values above detection limit
 438 both above and below, suggesting AABW waters carrying CFC-11 could have already
 439 travelled all the way to the equator. Note however we only rely on a limited number of
 440 observations north of 30°S at depth (Figure 4d), so the conclusion on how far north AABW
 441 has spread in the Indian Ocean still needs caution.

442 The deep Pacific Ocean is mostly filled with Ross Sea-formed and Adelie Coast-
 443 formed AABW (Solodoch et al., 2022), consistent with the high CFC-11 values through
 444 the Bellinghausen Basin (between 60-65°S in Figure 4e) and its spreading through the
 445 Southwest Pacific Basin, with values up to 0.1 pmol/kg up to 50°S. Within the Pacific,
 446 the highest deep CFC-11 values are found along the western boundary of the basin sam-
 447 pled by the P15 section, tracing the northward path of AABW that flows as a Deep West-
 448 ern Boundary Current hugging the Campbell Plateau (see Figure 1b for location and Fig-
 449 ure 5 for the earlier arrival time of AABW along P15). The core of northward AABW
 450 flow in the western Pacific around Campbell Plateau continues northward offshore of the
 451 Kermadec and Tonga trenches and flows in the north Pacific channel through the Samoan
 452 Passage (Figure 5). Our reconstruction returns a clear signal of AABW crossing the Samoan
 453 Passage (P21, Figure 6a) in the late 90s-early 2000s and moving northward (P14, Fig-
 454 ure 6b) in the early-to-late 2000s.

455 **4.2.2 North Atlantic Deep Water**

456 The Atlantic Ocean interior is mostly filled with North Atlantic Deep Water (NADW),
 457 formed in the Labrador Sea and Nordic Seas of the North Atlantic. NADW flows south-
 458 ward above the denser AABW throughout the Atlantic Ocean (spanning depths of ~1000-
 459 4000 m) and re-surfaces in the Southern Ocean mostly through wind-induced Ekman pump-
 460 ing (Talley et al., 2003; Lumpkin & Speer, 2007; Talley, 2008; Smeed et al., 2014; Cessi,
 461 2019; H. L. Johnson et al., 2019). The North Atlantic, through the formation and south-
 462 ward spreading of NADW, constitutes the largest CFC-11 inventory in the global ocean
 463 (Supplementary Figure S5; see also Willey et al. (2004)). A clear CFC-rich NADW core
 464 is observed down to 40°N along A16 and 30°N within the western NADW boundary cur-
 465 rent to the west (Figure 7). These high CFC values (up to 3 pmol/kg) were not found
 466 South of 50°N in previous reconstructions of Orsi et al. (2002), who used all the avail-
 467 able observations up to the late 1990s.

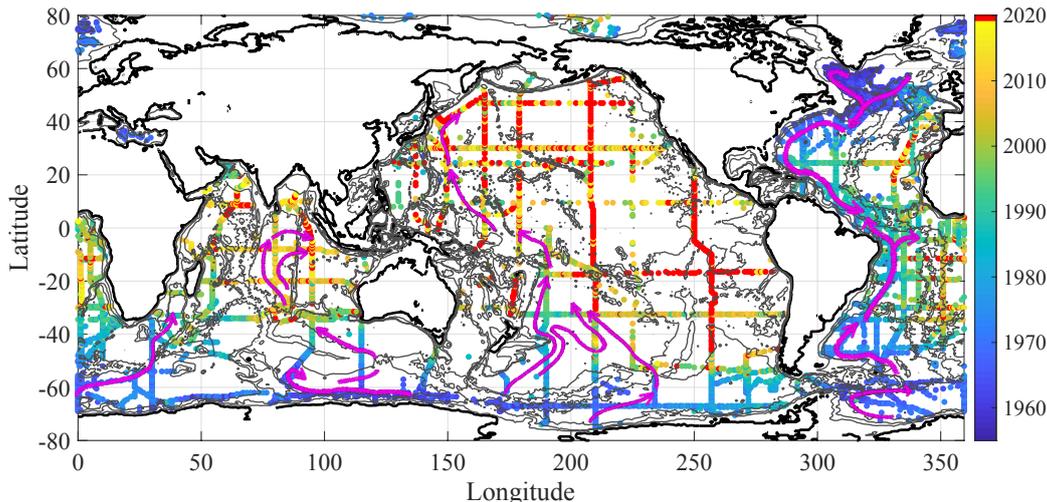


Figure 5. (a) Reconstructed CFC-11 arrival time at the ocean floor. The arrival time is defined as the year when CFC-11 first exceeded the detection limit of 0.01 pmol/kg. Values in red indicate areas where CFC-11 values are still below detection limit. Magenta arrows indicate some of the major AABW and NADW pathways taken from Reid (1994, 1997, 2003).

468 The Deep Western Boundary Current (DWBC) pathway of the southward flow-
 469 ing NADW is the fastest pathway for NADW to reach the southern hemisphere. To fur-
 470 ther quantify the time scales of advection, here we show the 10% ventilation time, t_{10} ,
 471 defined as the time it took to ventilate 10% of the water at each location, namely $t_{10}(\mathbf{r}, t) =$
 472 $\int_0^{t_{10}} G(\mathbf{r}, \tau) d\tau = 0.1$. Figure 7 shows t_{10} on $\gamma^n = 27.90$, which lies at about 2000 m
 473 in the middle of the Atlantic Basin and crosses the core of the Labrador Sea Water com-
 474 ponent of NADW. The 10% ventilation time is better suited to describe younger waters
 475 than the mean age, $t_m(\mathbf{r}, t) = \int_0^\infty \mathcal{G}(\mathbf{r}, \tau) d\tau$, and so is a better suited expression of the
 476 water age on the decadal timescales described by CFCs and SF₆. The broad region where
 477 a fast arrival time is found in the subpolar North Atlantic reflects the multiple varieties
 478 of dense waters being formed here. While their pathways are not necessarily restricted
 479 to the western boundary (Lozier, 1997; Lozier et al., 2022), the 10% ventilation time is
 480 shorter (10-30 years) in the western North Atlantic, compared to the eastern Atlantic,
 481 where CFCs and SF₆ have barely arrived as of 2021 ($t_{10} \simeq 80$ years).

482 CFCs and SF₆ are effective in constraining the GFs and minimizing the effect of
 483 using an Inverse Gaussian as the first guess. Following NADW as it moves southward,
 484 the GFs become increasingly wide, indicating that waters have moved away from the source
 485 regions and mixed along the way (Figure 7b). The solved GFs can differ quite substan-
 486 tially from the first guess, and even return bimodal solutions that indicate the contri-
 487 bution of different water masses (solid vs dashed lines in Figure 7b). The 10% ventila-
 488 tion time progressively increases from just 2 years at 56°N to 32 years at the equator.
 489 The reconstructed CFC-11 concentration reveals that the high latitudes in the North At-
 490 lantic are already experiencing the roll-off in CFC-11 atmospheric history, and even val-
 491 ues at 35°N have started plateauing, while CFC-11 are still steadily increasing further
 492 south (Figure 7c).

493 Further downstream, subsurface tracer maxima are found in the equatorial deep
 494 Atlantic, where part of the southward flowing NADW turns eastward along the equa-
 495 tor (Weiss et al., 1985; Schott et al., 2003; Rosell-Fieschi et al., 2015; Herrford et al., 2017).
 496 We observe a CFC-11 minimum sandwiched between a shallow maximum at 1500-2500

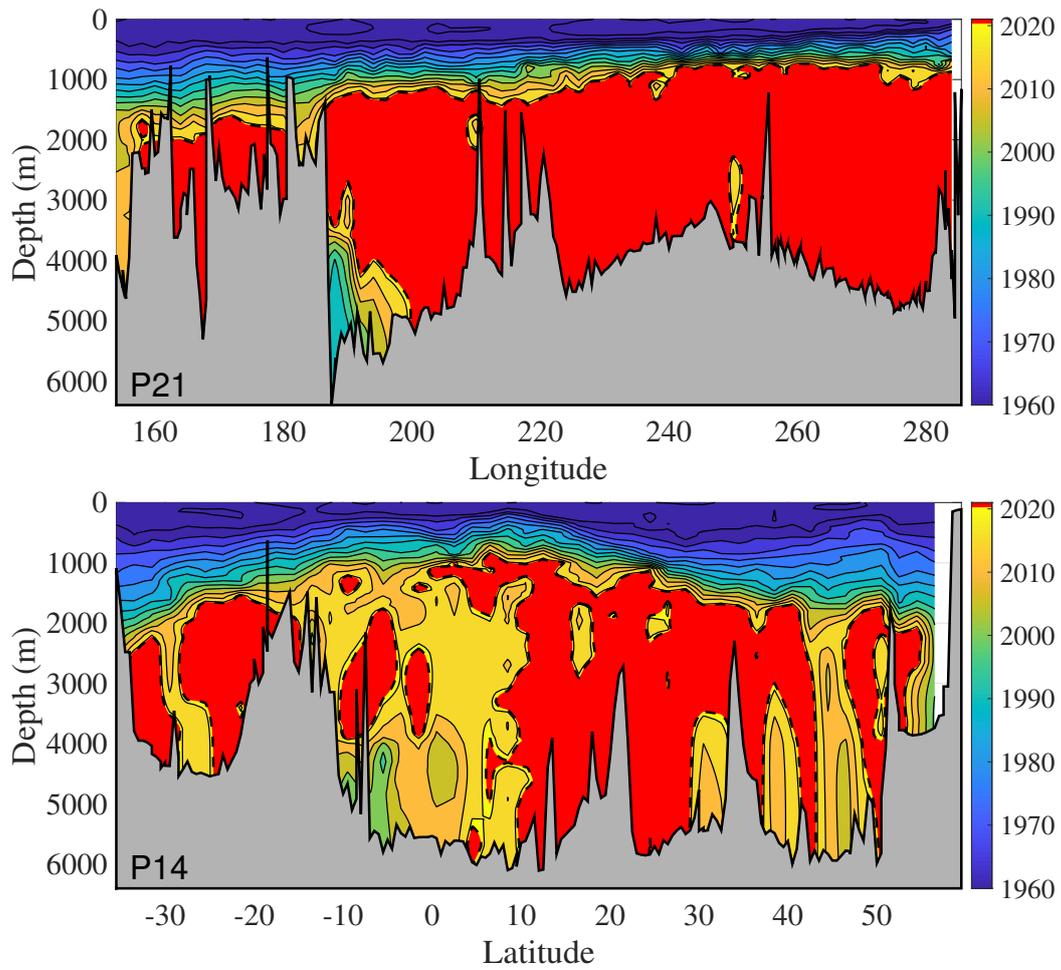


Figure 6. Reconstructed CFC-11 arrival time along P21 and P14 in the Pacific Ocean (see Figure 1b for the location of the sections).

497 m and a deep maximum at 3200-4500 m. This CFC-11 pattern is consistent with oxy-
 498 gen patterns and transport patterns found in previous studies (Schott et al., 2003; Rosell-
 499 Fieschi et al., 2015; Herrford et al., 2017). The maxima correspond to the lighter and
 500 denser variations of NADW, known as upper NADW (uNADW) and lower NADW (lNADW),
 501 respectively (the neutral density levels indicating the interface between the different com-
 502 ponents of NADW are displayed in Figure 4a). In between, the CFC-11 minimum is likely
 503 caused by an older water mass coming from the southeast, possibly modified or recir-
 504 culated NADW water (Friedrichs et al., 1994; Rhein et al., 1995; Herrford et al., 2017).
 505 The time-corrected CFC-11 concentrations along A17 and A15 sections, which cross the
 506 equator at 30°W and 19°W (while A16 crosses it at 25°W), respectively show that the
 507 two maxima are more evident further eastward (A15) and less separated further west-
 508 ward (A17), in agreement with tracer dilution being caused by older NADW recirculat-
 509 ing from the southeast.

510 The shallow equatorial CFC-11 maximum (1500-2500 m) was first observed in 1983
 511 (Weiss et al., 1985) and consistently makes its first appearance in the early 1980s in our
 512 reconstruction (not shown). The smaller deep maximum (3200-4500 m) was first observed
 513 in 1988 (Doney & Bullister, 1992), and again consistently first appears in the mid-1980s
 514 in the tracer reconstruction. The core of the deep maximum initially sits within lNADW
 515 and spreads between lNADW and AABW within just a few years, potentially because
 516 of the large mixing between these two water masses across this region (Herrford et al.,
 517 2017). As time progresses, the deep maximum crosses denser density classes, suggest-
 518 ing a delayed contribution of AABW-carried CFC-11. The high CFC-11 equatorial con-
 519 centrations were only partially evident in Orsi et al. (2002) and in the WOCE atlas (Koltermann
 520 et al., 2011), with values still very close to the detection limit used here (0.01 pmol/kg)
 521 in the 1990s. The deep CFC-11 maximum, in particular, was nearly undetectable and
 522 confined to 4200-4500 m in both estimates, suggesting it was mainly associated with lNADW
 523 flow rather than AABW.

524 4.3 Variability

525 The TCM finds the best Green’s function fit for the whole period of available at-
 526 mospheric and oceanic observations, implicitly characterizing the circulation as steady.
 527 A steady-state circulation, i.e. one for which observations are within the uncertainty er-
 528 rorbars of our reconstruction, is globally capable of explaining 73% of the points within
 529 the 95% confidence errorbars, varying some by density class (Figure 8b). Individually
 530 considered, the different tracers used here return different levels of misfit between the
 531 reconstructed and observed concentrations. About 80% of CFC-11, 70% of CFC-12 and
 532 55% of SF₆ observations are within the 95% confidence errorbars. SF₆ observations are
 533 rarely available independently and it could prove more challenging for the method to find
 534 a solution that satisfies tracers with different atmospheric histories at the same time, pos-
 535 sibly explaining the lower percentage of SF₆ data comprised by the errorbars.

536 If the TCM solution was always consistent with the available observations, 90% of
 537 the reconstructions should be within errorbars, given that errorbars are calculated as the
 538 5% - 95% confidence limits. The lower fit found could be caused by temporal variabil-
 539 ity in the circulation or by sensitivity to the initial conditions used in the method. An
 540 analysis of the spatial (horizontal and vertical) and time distributions of the misfit be-
 541 tween the reconstructed and observed concentrations offers insight into where the cir-
 542 culation is not in steady state and can suggest either a slowdown or speed up in venti-
 543 lation of particular water masses. However, given the spatial sparsity of available obser-
 544 vations, a regional approach would be better suited to investigate changes in different
 545 water masses where local dynamics can be examined, thus we leave this analysis mostly
 546 to future studies, with only one example following on previous studies discussed below.
 547 Here, we simply show that the misfit is normally-distributed around zero, i.e. the TCM

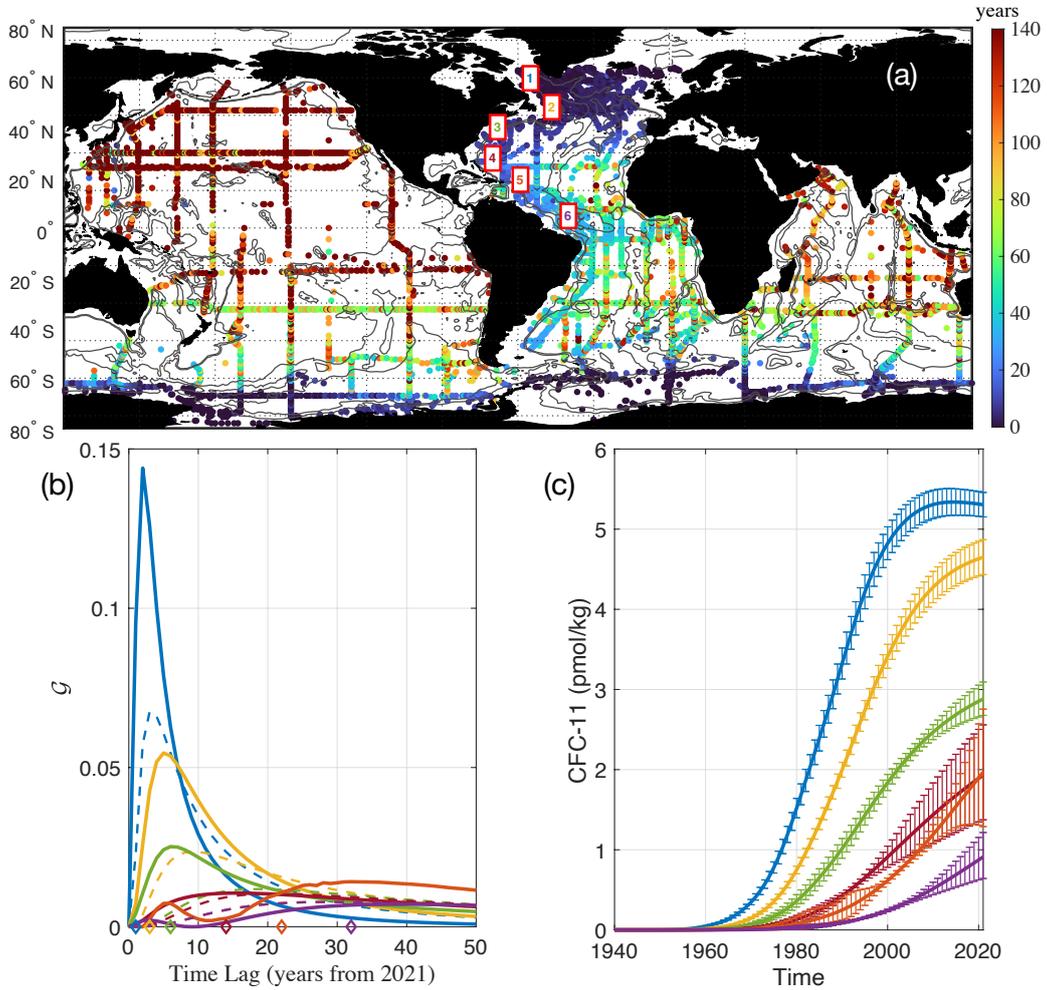


Figure 7. (a) 10% ventilation time on neutral density surface $\gamma^n = 27.90$, the density class that crosses the core of the Labrador Sea Water component of NADW. The Green's functions (b) and the reconstructed tracer concentrations (c) at 6 points (color-coded labels in panel a) across the Atlantic Ocean. The color-coded diamonds at the bottom on panel b indicate the 10% ventilation time (t_{10}) for each of the displayed Green's functions. Dashed lines in panel (b) show the first guess Green's function used per each location (color-coded).

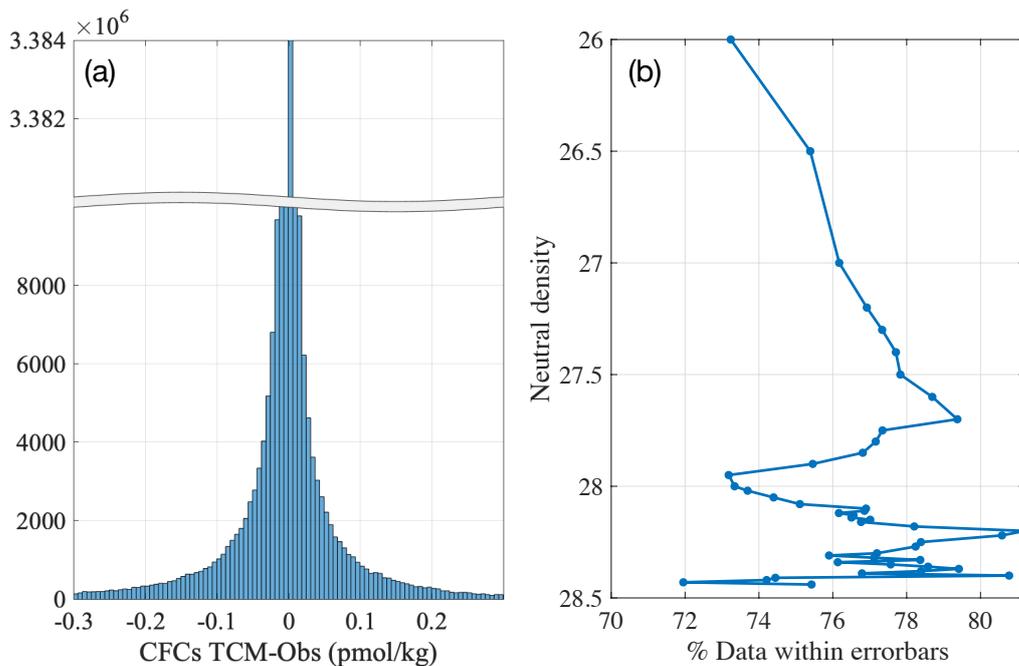


Figure 8. (a) Distribution of the misfit between the reconstructed and observed tracer concentrations, for all tracers and all ocean basins. (b) Globally-averaged percentage of observations within the errorbars of our reconstruction for each density level used in this study.

548 tracer reconstruction fits the data in most locations, confirming that the steady-state holds
 549 for the most part and that there are no biases in the TCM reconstruction (Figure 8a).

550 *4.3.1 Variability in Subantarctic Mode Water*

551 The TCM not only informs us about the steady-state ocean circulation, but offers
 552 a quantitative method for assessing decadal ocean variability. The TCM reconstruction
 553 fits most of the data (Figure 8a), and it tends to overfit the data when we have only one
 554 occupation due to the large underdetermined nature of the problem (Figure 3). There-
 555 fore, when we have multiple occupations and the reconstruction does not agree with ob-
 556 servations within the error estimates, the null hypothesis of steady-state circulation is
 557 rejected and the temporal variability can be assessed.

558 As an example, decadal changes in the ventilation of Subantarctic Mode Water (SAMW)
 559 are discussed. Transient tracer vertical inventories (Figure S5) reveal that a primary area
 560 of transient tracer accumulations is the Southern Ocean north of the Sub-Antarctic Front
 561 (around 50°S), co-located with the maximum wind-stress curl where SAMW and Antarc-
 562 tic Intermediate Waters (AAIW) are formed. SAMW and AAIW flow northward through
 563 the sub-tropical thermocline in all ocean basins, and contribute to the closure of the At-
 564 lantic Meridional Overturning Circulation by returning water to the North Atlantic. High
 565 CFC-11 concentrations downstream of SAMW and AAIW formation sites are evident
 566 across all ocean basins, with values up to 3 pmol/kg spreading up to 20°S in 2021 (Fig-
 567 ure 4), a northward propagation of almost 20 degrees latitude from the 1990s (Orsi et
 568 al., 2002).

569 Several studies have used repeat measurements of transient tracers to estimate changes
 570 in ventilation rates in the Southern Ocean thermocline waters, and have generally shown
 571 a decrease in SAMW age, implying an increase in the ventilation rates (Tanhua et al.,

572 2013; Waugh et al., 2013; Fine et al., 2017; Ting & Holzer, 2017; Morrison et al., 2022).
 573 Here, we compare the reconstructed and observed concentrations along section P16, in
 574 analogy with the past studies cited above.

575 The comparison of the reconstructed and observed tracer concentrations suggest
 576 an increase in SAMW ventilation rate from 1992 to 2014. The TCM analysis solves for
 577 the GF that best fits all occupations (1992, 2005 and 2014). If the circulation is chang-
 578 ing linearly in time, the solution would be roughly centered in the mid-term of the con-
 579 sidered period, in this example ~ 2003 , explaining why the differences between reconstructed
 580 and observed concentrations are smallest in 2005 (Figure 9b). The reconstructed con-
 581 centrations in SAMW ¹ in 1992 are slightly higher than observations (Figure 9a), while
 582 they are lower than observations in 2014 (Figure 9c). On the other hand, the reconstructed
 583 CFC-11 concentrations in the Circumpolar Deep Waters (CDW ²) reveal the opposite
 584 pattern, namely they are lower than observations in 1992, and higher than observations
 585 in 2014 (Figures 9a,c).

586 To further confirm the differences described above, we repeat the TCM analysis
 587 using only the 1992 data to predict the 2005 and 2014 (Figure 9d-f), as it has been done
 588 in Waugh et al. (2013). While the GFs used to predict the tracer concentrations are dif-
 589 ferent, the idea is the same: the prediction assumes steady-state circulation and diver-
 590 gences between reconstructed and observed concentrations suggest ventilation rate changes.
 591 In this scenario, the anomalies between the reconstructed and observed CFC-11 in 1992
 592 are naturally small (close to zero), confirming that the TCM finds a solution that fits
 593 the available data (Figure 9d). The tracer differences in 2005 confirm the patterns shown
 594 in Waugh et al. (2013), i.e. lower (larger) reconstructed concentrations in SAMW (CDW)
 595 than in observations (Figure 9e). The tracer differences in 2014 shows a further decrease
 596 (increase) in SAMW (CDW) ventilation, showing a continuation of the trend discussed
 597 in Waugh et al. (2013). In 2014, the differences between reconstructed and observed con-
 598 centrations are qualitatively similar whether we employ all available occupations simul-
 599 taneously or 1992 observations only (Figure 9c,f).

600 The increase in SAMW ventilation rate has been linked to the intensification and
 601 poleward shift of the westerly winds around Antarctica (Waugh et al., 2013). Stronger
 602 winds would lead to an increase of the upper overturning circulation, i.e. a stronger ex-
 603 port of SAMW and more upwelling of the older, tracer-poor CDW. While this mecha-
 604 nisms could explain the observed changes, other processes could have caused similar changes
 605 in ventilation (e.g. changes in the subtropical gyres, see Morrison et al. (2022)), chal-
 606 lenging the attribution to a single mechanism. Despite these attribution difficulties, the
 607 changes have been determined with a number of different methods in different studies,
 608 suggesting that they are not the result of uncertainties in estimates of the tracer con-
 609 centration (or age changes), but rather the result of variability.

610 5 Summary and Discussion

611 Transient tracers have played a critical role in advancing our understanding of ocean
 612 circulation and ventilation, but direct quantitative interpretation of their presence can
 613 be challenging owing to (i) the non-linearity of their time varying atmospheric history,
 614 (ii) unknown surface saturation, and (iii) the complicated competition between advec-
 615 tion and mixing in the ocean interior (Purkey et al., 2018). Here, we present a new math-
 616 ematical framework that leverages the global data set of CFC-11, CFC-12 and SF₆ data
 617 to solve for the steady-state age distribution and the annual concentration of each tracer
 618 between 1940 and 2021 at every location where tracer data exists in the interior ocean.

¹ SAMW: $\gamma^n = 26.6 - 27.2$ and north of the Sub-Antarctic Front at $\sim 50^\circ\text{S}$

² CDW: $\gamma^n = 27.2 - 27.6$ south of the Sub-Antarctic Front

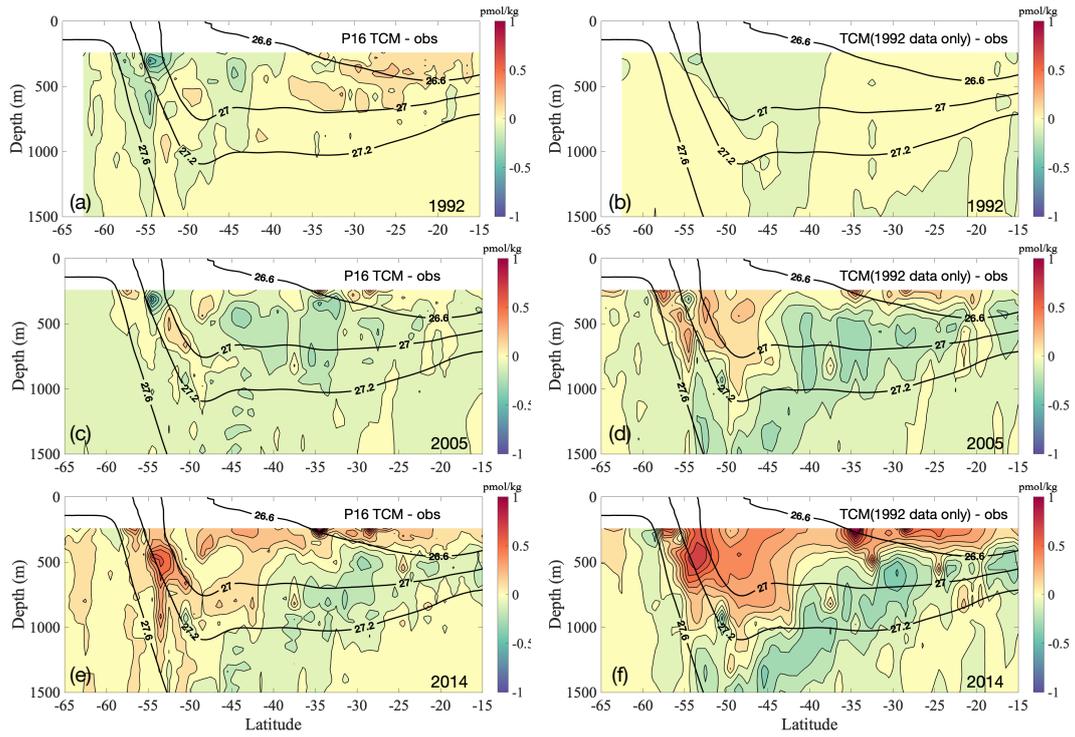


Figure 9. Depth-latitude cross sections of the difference between reconstructed (TCM) and observed CFC-12 concentrations for repeat cruises along P16 (shading) and potential density (σ_0) contours at the section location (black contours). Left panels: reconstructed concentrations are estimated using all available data (three occupations, in 1992, 2005 and 2014). Right panels: reconstructed concentrations are estimated using 1992 data only.

619 This global, annual, tracer dataset allows for evaluation of the mean circulation of the
 620 deep ocean and time scales of ventilation. In addition, this new dataset can be used to
 621 validate ocean models (Solodoch et al., 2022). Furthermore, comparison of the steady-
 622 state reconstructed tracer concentration with the observed concentrations allows for direct
 623 assessment of where ocean ventilation has varied over the 1940-2021 period, with
 624 potential implications for the rate of ocean heat and carbon uptake.

625 The Time Correction Method presented here leverages the longer record of CFCs
 626 and SF₆ that is now available. The further expansion of these tracers into deep ocean
 627 permits a more direct inference of ventilation time for a greater fraction of the global ocean.
 628 In addition, the atmospheric histories of CFC-11, CFC-12, and SF₆ have diverged over
 629 the last few decades, and thus these tracers now provide more independent constraints
 630 regarding ocean circulation. The time correction method doesn't restrict the age distribution
 631 of the interior ocean to follow a prescribed Inverse Gaussian distribution that describes
 632 the evolution of a single water mass. The combination of the time-correction method
 633 with a longer timeseries of data leads to more complicated descriptions of the circulation
 634 to be inferred, such as bimodal age distributions that correspond to the isopycnal
 635 mixing of two water masses with different histories.

636 The findings presented here come with some known caveats, including some due
 637 to the use of surface boundary conditions in discrete density classes constrained by the
 638 available observations within the geographical region where that density outcrops. Especially
 639 for some of the deep density layers formed seasonally under ice where there is extremely
 640 limited data, the saturation value is uncertain, however, an inference based on the closest
 641 observations is likely more accurate than assuming that the surface is perfectly saturated
 642 or that the entire sea surface has a uniform saturation rate. Increasing the number of
 643 tracer observations in high latitudes in both winter and summer could help constrain the
 644 boundary conditions in the future. We do not account for the small lag in atmospheric
 645 CFC concentration between the northern and southern hemispheres which introduces a
 646 ~6 month offset in waters of southern origin. In addition, a number of SF₆ tracer
 647 release experiments have introduced an artificial (i.e., non-atmospheric) source into
 648 the interior ocean over the last three decades. We assume this is a relatively small source
 649 of SF₆ in the ocean and note the ratio of CFCs to SF₆ would not be the same as the
 650 atmospheric history.

651 The oceanic histories of CFCs and SF₆ are less than a century old and thus their
 652 distributions likely still reflect fast pathways that result from advective transport. As
 653 advective pathways induce a tracer transport that is primarily isopycnal, it is a good
 654 assumption that the boundary condition, $C_{atm}(t)$ in Equation 1, is given by following the
 655 isopycnal surface of any interior point back to the surface. In regions such as the deep
 656 Pacific, diapycnal transport can no longer be assumed to be small, and it is no longer
 657 clear which value for $C_{atm}(t)$ is best. For the deep Pacific, in particular, ventilation
 658 occurs by horizontal advection of abyssal waters followed by slow upward diapycnal
 659 transport, and thus the most appropriate boundary condition may come from a denser
 660 density surface. Fortunately for this analysis of CFCs and SF₆, there is little penetration
 661 into the deep Pacific at this time.

662 At some point in the future, analyses of CFCs and SF₆ will have to take into
 663 account the effect of diapycnal transport for the inference of proper boundary conditions.
 664 Other tracers that have a longer history of variability, such as anthropogenic carbon, may
 665 already require such accounting. Methods that simultaneously invert tracer data for water-
 666 mass mixing and aging (DeVries & Primeau, 2011; Gebbie & Huybers, 2012) are a natural
 667 way to solve this problem, at the cost of requiring a global inversion with commensurate
 668 computational costs.

669 Here, we have demonstrated the ability to use tracers to evaluate variability in ocean
 670 ventilation and assess advection timescales in the deep ocean. In addition, our Green's

671 Functions can be convolved with any surface boundary conditions to assess ocean up-
 672 take of other quantities, including anthropogenic carbon and heat. Continuing to mea-
 673 sure transit tracers through international programs such as GO-SHIP will be a key tool
 674 for assessing the uptake of anthropogenic heat and carbon by the ocean and monitor-
 675 ing variability in ocean ventilation and circulation.

676 **6 Open Research**

677 Data used here is publicly available from GLODAP (<https://www.glodap.info/>).
 678 AnSlope tracer data will be made available before publication via the CCHDO database.

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Supporting Information for ” Annually-resolved propagation of CFCs and SF₆ over eight decades”

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Contents of this file

1. Text S1
2. Figures S1 to S5

Introduction

The supplementary information in Text S1 reports the details of how to construct and solve the modified, non-negative least squares problem to deconvolve the discrete Green’s function \mathbf{G} (together with the information reported in Section 3.2 of the main manuscript).

Supplementary Figures S1 and S2 also provide additional details about the Time-Correction Method, in particular about the decisions made in this study to calculate the surface boundary conditions (Figure S1) and the first guess Green’s function (Figure S2).

Figures S3 and S4 report the reconstructed CFC-12 and SF₆ concentrations in 2021 for the hydrographic sections A16-A23, P16 and I8I9. The reconstructed CFC-11 concentrations are shown in the main manuscript (Figure 4).

Figure S5 shows the vertically-integrated CFC-11 concentration reconstructed with the TCM in 1994 and 2020. The 1994 year was chosen for a direct comparison to the CFC-11 inventory by Willey et al. (2004). Note that our product is not gridded on a uniform horizontal and vertical grid, such as the CFC-11 inventory by Willey et al. (2004); as such, in the vertical integration of the reconstructed CFC-11 we examine only the locations where observations describe at least 66% of the water column. This choice explains the lack of data in the southwest Indian Ocean, in the Atlantic sector of the Southern Ocean and partially in the east Pacific Ocean.

Text S1. To solve the least squares problem to find \mathbf{G} , at each location we define:

$$\hat{\mathbf{C}} = \begin{bmatrix} \mathbf{C}_1 \\ \mathbf{C}_2 \\ \vdots \\ \mathbf{C}_{N_t} \\ \mathbf{G}_0 \end{bmatrix}, \quad \hat{\mathbf{C}}_s = \begin{bmatrix} C_{s1} \\ C_{s2} \\ \vdots \\ C_{sN_t} \\ \mathbf{I} \end{bmatrix}, \quad \hat{W}^{-1} = \begin{bmatrix} W_1^{-1} & 0 & 0 & 0 & 0 \\ 0 & W_2^{-1} & 0 & 0 & 0 \\ 0 & 0 & \ddots & 0 & \\ 0 & 0 & 0 & W_{N_t}^{-1} & 0 \\ 0 & 0 & 0 & 0 & S^{-1} \end{bmatrix}.$$

N_t is the number of tracers used (depending on the available observations), and \mathbf{C}_i represents the array of observations available at one location in the ocean interior for each tracer. The C_{si} matrix returns the surface boundary conditions per each tracer, while W_i^{-1} and S^{-1} indicate the weight matrices for the tracer observations and the Green's function first guess \mathbf{G}_0 , respectively. We re-arrange the terms reported above to use the non-negative least-squares method of Lawson and Hanson (1995), defining

$$\begin{aligned} \hat{\mathbf{C}} &= \hat{W}^{-1/2} \hat{\mathbf{C}} \\ \hat{\mathbf{C}}_s &= \hat{W}^{-1/2} \hat{\mathbf{C}}_s \end{aligned} \tag{1}$$

which we can use to rewrite the cost function as

$$J = (\hat{\mathbf{C}} - \hat{C}_s \mathbf{G})^T (\hat{\mathbf{C}} - \hat{C}_s \mathbf{G}), \quad (2)$$

which is equivalent to solving the equation $\hat{C}_s \mathbf{G} = \hat{\mathbf{C}}$, going back to the canonical form of a least-squares problem.

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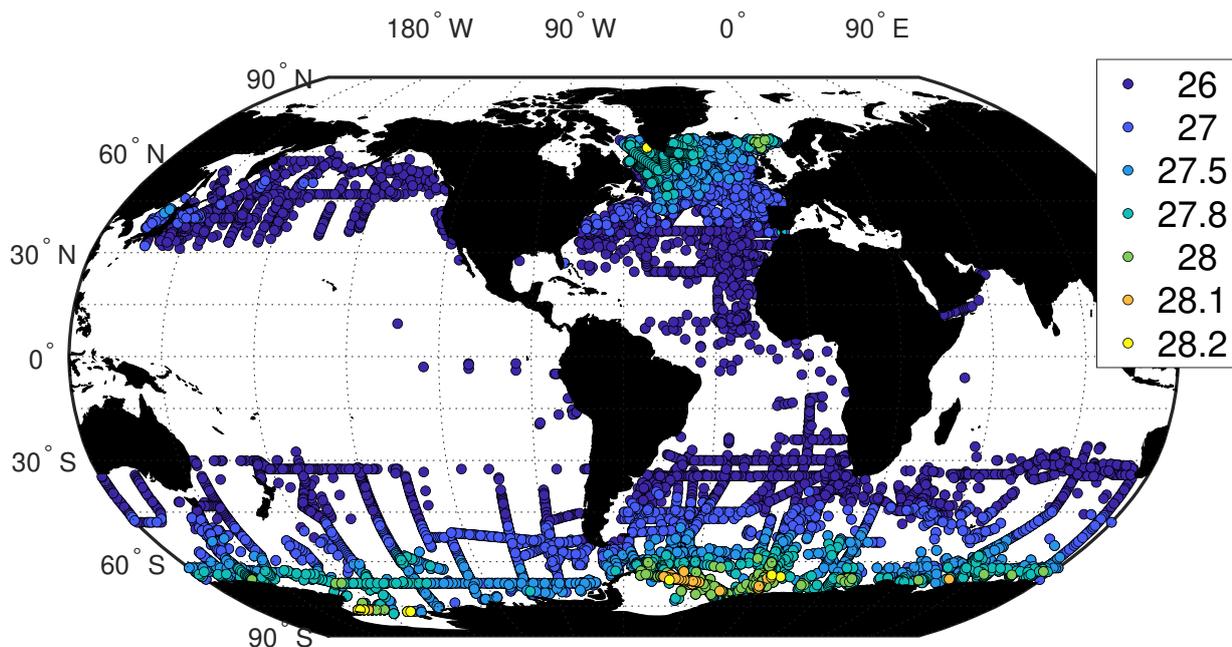


Figure S1. Outcrop points for the neutral density levels indicated in the legend. Outcrop points are defined as locations where the maximum mixed layer depth (from Holte et al. (2017)) is found at a density larger than a specific density level.

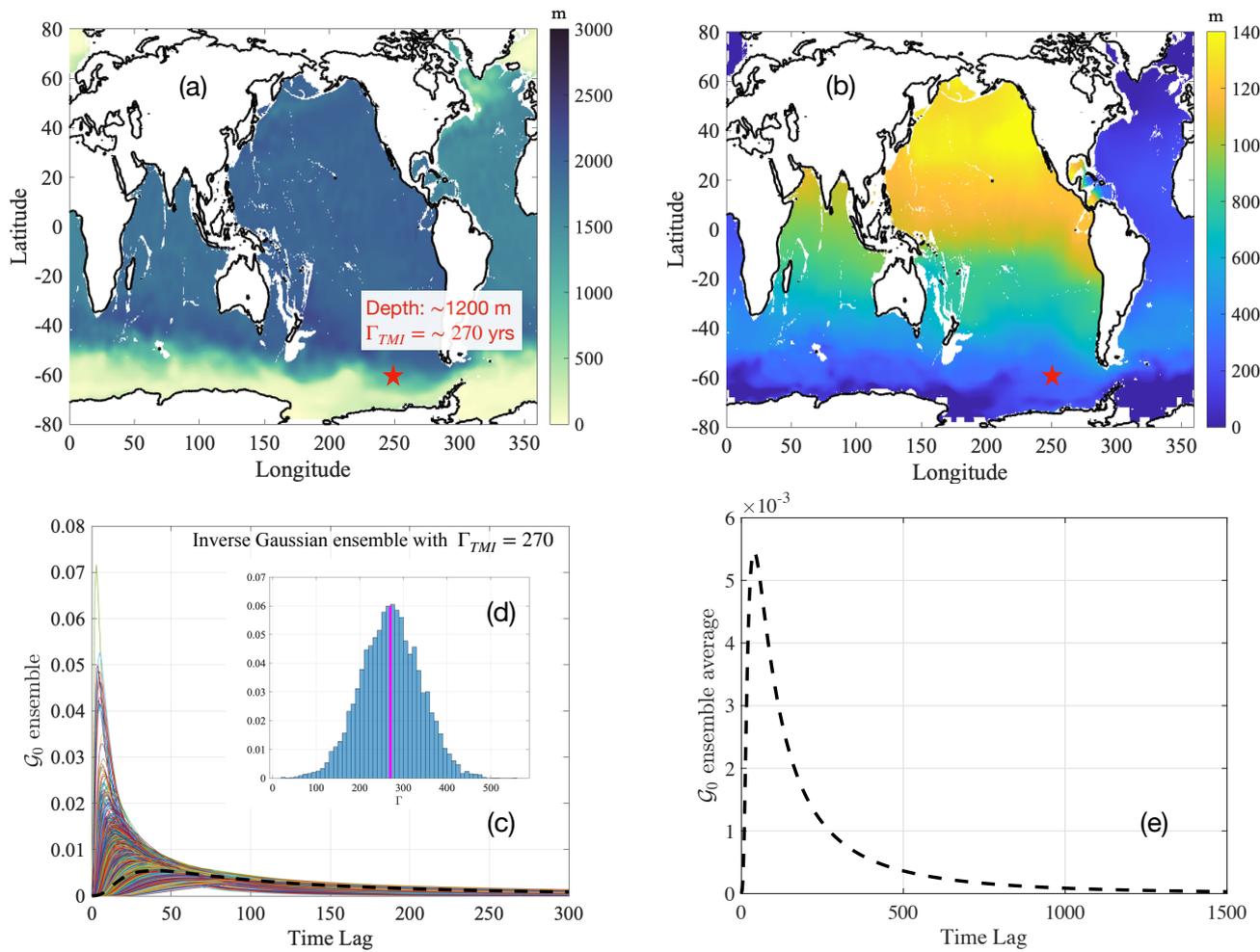


Figure S2. Steps to define the first guess \mathbf{G}_0 : example at 150°W, 60°S on neutral density surface $\gamma^n = 27.90$ (red star in panels a, b). For each location, we find the closest TMI water age estimate. (a)-(b) Depth and Γ_{TMI} of the density surface. At the starred location, $\gamma^n = 27.90$ is about 1200 m deep and the TMI water age estimate is $\Gamma_{TMI} = 270$ years. (c) Ensemble of Inverse Gaussians with age within the range $\Gamma_{TMI}/2 : \Gamma_{TMI} \times 2$ (see subset histogram in panel (d)). The ensemble average returns the kernel first guess used here, \mathbf{G}_0 (indicated in panel d, and shown in panel e).

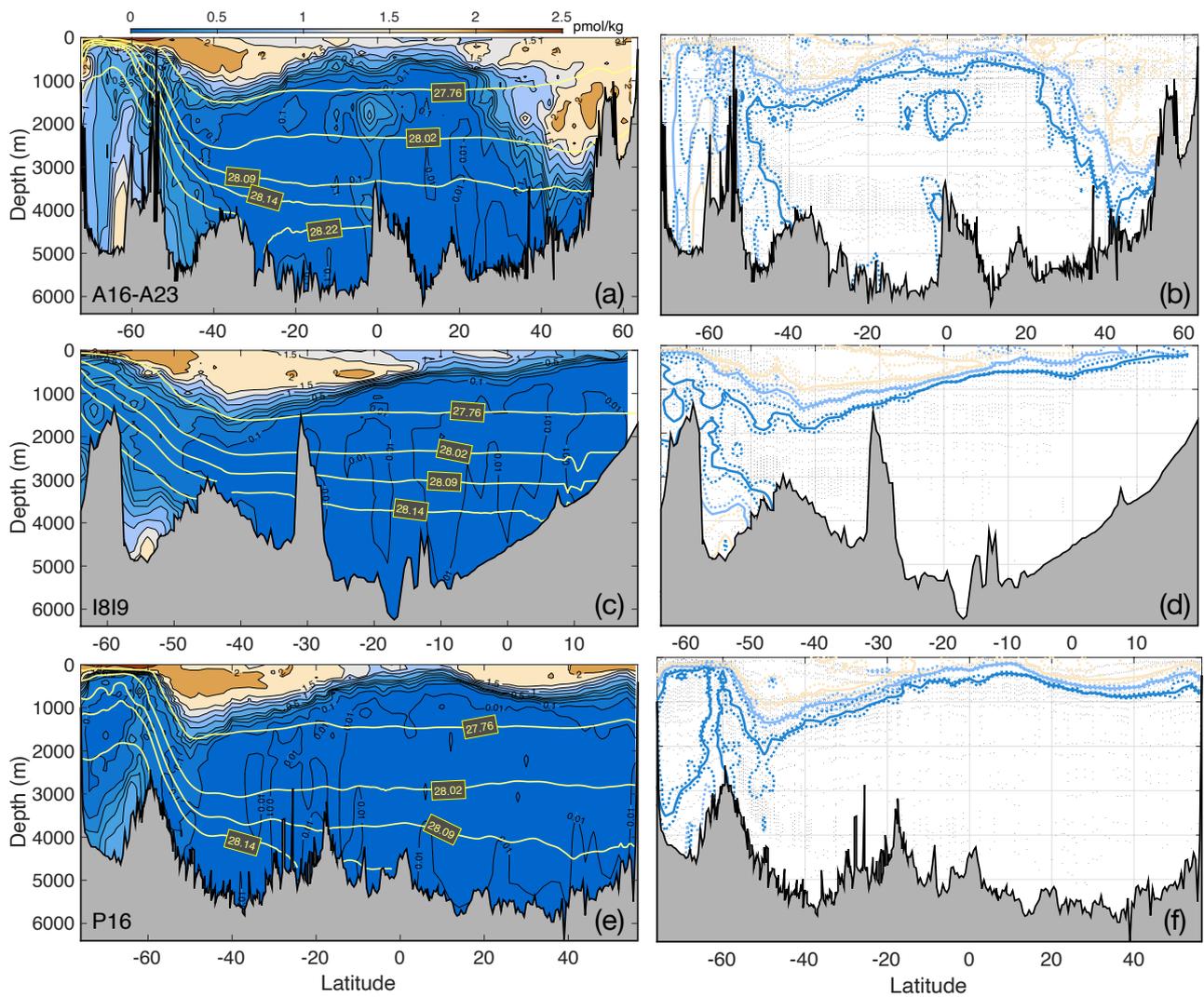


Figure S3. Reconstructed CFC-12 concentrations along A16+A23 (panels a,b), I8I9 (c,d), and P16 (e,f) in 2021. Panels in left column show the reconstructed concentration and neutral density contours (yellow) along each section. Panels in right column show the observation locations throughout all decades of available data (gray dots) and the uncertainty of the time corrected CFC-12 concentrations for three contours (0.2, 0.75 and 1.5 pmol/kg). The solid line shows our best estimate, while the dotted lines show the 95% confidence limits.

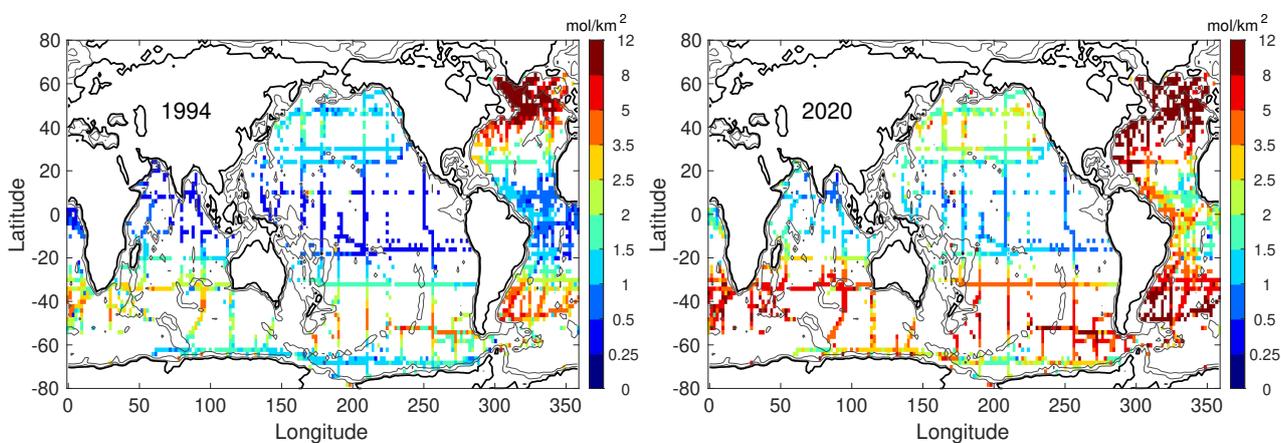


Figure S5. Vertically-integrated CFC-11 concentrations in 1994 (which can be compared with the inventory presented in Willey et al. (2004)) and 2020.