On the role of ULF waves in the spatial and temporal periodicity of energetic electron precipitation

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Abstract

Energetic electron precipitation to the Earth's atmosphere is a key process controlling radiation belt dynamics and magnetosphereionosphere coupling. One of the main drivers of precipitation is electron resonant scattering by whistler-mode waves. Lowaltitude observations of such precipitation often reveal quasi-periodicity in the ultra-low-frequency (ULF) range associated with whistler-mode waves, causally linked to ULF-modulated equatorial electron flux and its anisotropy. Conjunctions between ground-based instruments and equatorial spacecraft show that low-altitude precipitation concurrent with equatorial whistlermode waves also exhibits a spatial periodicity as a function of latitude over a large spatial region. Whether this spatial periodicity might also be due to magnetospheric ULF waves spatially modulating electron fluxes and whistler-mode chorus has not been previously addressed due to a lack of conjunctions between equatorial spacecraft, LEO spacecraft, and ground-based instruments. To examine this question, we combine ground-based and equatorial observations magnetically conjugate to observations of precipitation at the low-altitude, polar-orbiting CubeSats ELFIN-A and -B. As they sequentially cross the outer radiation belt with a temporal separation of minutes to tens of minutes, they can easily reveal the spatial quasi-periodicity of electron precipitation. Our combined datasets confirm that ULF waves may modulate whistler-mode wave generation within a large MLT and \$L\$-shell domain in the equatorial magnetosphere, and thus lead to significant aggregate energetic electron precipitation exhibiting both temporal and spatial periodicity. Our results suggest that the coupling between ULF and whistler-mode waves is important for outer radiation belt dynamics.

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8 Abstract

Energetic electron precipitation to the Earth's atmosphere is a key process controlling q radiation belt dynamics and magnetosphere-ionosphere coupling. One of the main drivers 10 of precipitation is electron resonant scattering by whistler-mode waves. Low-altitude ob-11 servations of such precipitation often reveal quasi-periodicity in the ultra-low-frequency 12 (ULF) range associated with whistler-mode waves, causally linked to ULF-modulated 13 equatorial electron flux and its anisotropy. Conjunctions between ground-based instru-14 ments and equatorial spacecraft show that low-altitude precipitation concurrent with equa-15 torial whistler-mode waves also exhibits a spatial periodicity as a function of latitude over 16 a large spatial region. Whether this spatial periodicity might also be due to magneto-17 spheric ULF waves spatially modulating electron fluxes and whistler-mode chorus has 18 not been previously addressed due to a lack of conjunctions between equatorial space-19 craft, LEO spacecraft, and ground-based instruments. To examine this question, we com-20 bine ground-based and equatorial observations magnetically conjugate to observations 21 of precipitation at the low-altitude, polar-orbiting CubeSats ELFIN-A and -B. As they 22 sequentially cross the outer radiation belt with a temporal separation of minutes to tens 23 of minutes, they can easily reveal the spatial quasi-periodicity of electron precipitation. 24 Our combined datasets confirm that ULF waves may modulate whistler-mode wave gen-25 eration within a large MLT and L-shell domain in the equatorial magnetosphere, and 26 27 thus lead to significant aggregate energetic electron precipitation exhibiting both temporal and spatial periodicity. Our results suggest that the coupling between ULF and 28 whistler-mode waves is important for outer radiation belt dynamics. 20

Key Points:

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31	•	We report quasi-periodic energetic electron precipitation observed by low-altitude
32		ELFIN CubeSats
33	•	Quasi-periodicity of the precipitation is due to periodicity of whistler-mode wave
34		generation at the equator
35	•	Whistler-mode wave generation is modulated by ultra-low-frequency waves seen

within a large MLT, L-shell domain

37 1 Introduction

Resonant scattering of energetic electrons from Earth's radiation belts by electro-38 magnetic whistler-mode waves is the main driver of energetic electron precipitation from 39 the outer radiation belt to the atmosphere (e.g., Millan & Thorne, 2007; W. Li & Hud-40 son, 2019; Thorne et al., 2021). Statistical analyses of observed wave characteristics (e.g., 41 Meredith et al., 2012; W. Li, Thorne, Bortnik, Shprits, et al., 2011; O. V. Agapitov et 42 al., 2013), coupled with global numerical simulations of resonant scattering, can provide 43 estimates of long-term electron losses due to precipitation (e.g., Mourenas et al., 2014; 44 Orlova et al., 2016; Ma et al., 2018, 2020; Hsieh et al., 2021, and references therein). Con-45 versely, case studies using equatorial spacecraft observations of wave and plasma char-46 acteristics specific for each event can describe well the dynamics of resonant electron fluxes, 47 which are often localized (e.g., O. V. Agapitov et al., 2015; Foster et al., 2014; Gan et 48 al., 2020; Capannolo et al., 2019). Neither statistical nor equatorial spacecraft case stud-49 ies can separate the temporal and spatial (mesoscale, $\sim R_E$ at the equator) variations 50 of precipitation. However, multi-spacecraft measurements at low altitudes or correlative 51 studies using low-altitude and equatorial measurements can be used to infer such scales. 52 These two approaches were employed to study the most intense but highly transient and 53 localized precipitation events, microbursts (O'Brien et al., 2004; Douma et al., 2017), by 54

(Shumko et al., 2018) and (e.g., Breneman et al., 2015; Mozer et al., 2018; Shumko et al., 2021), respectively. Mesoscale precipitation events, with temporal and spatial scales comparable to those of equatorial chorus waves (O. V. Agapitov et al., 2017; O. Agapitov et al., 2018), have yet to be investigated with similar methods.

A classical example of dynamic precipitation is the pulsating aurora (Belon et al., 59 1969; Coroniti & Kennel, 1970; Johnstone, 1978; McEwen et al., 1981), which is asso-60 ciated with $\sim 10 \text{ keV}$ electron precipitation as a result of quasi-periodic whistler-mode 61 (chorus) wave scattering (Nishimura et al., 2010; Kasahara et al., 2018). The quasi-periodicity 62 63 of these wave emissions may be caused by ultra-low-frequency (ULF) waves modulating plasma and energetic electron fluxes around the equator (Coroniti & Kennel, 1970; 64 Bryant et al., 1971; W. Li, Thorne, Bortnik, Nishimura, & Angelopoulos, 2011; W. Li, 65 Bortnik, et al., 2011; Motoba et al., 2013; Jaynes, Lessard, et al., 2015). Recent optical 66 and low-altitude measurements showed that ≤ 10 keV precipitation forming the pul-67 sating aurora may be accompanied by precipitation of relativistic electrons (Miyoshi et 68 al., 2020; Shumko et al., 2021), which makes such a quasi-periodic precipitation pattern 69 particularly important in the context of energetic electron losses and altering of atmo-70 sphere properties (Miyoshi et al., 2021). Incoherent scatter radar and ionospheric total 71 electron content also show that ULF-modulated precipitation can significantly alter iono-72 spheric conductance, potentially affecting the dynamics of a range of magnetosphere-ionosphere 73 current systems and ULF waves (e.g., Buchert et al., 1999; Pilipenko, Belakhovsky, Ko-74 zlovsky, et al., 2014; Pilipenko, Belakhovsky, Murr, et al., 2014; Wang et al., 2020). Us-75 ing ground-based ULF and equatorial whistler wave observations to characterize precip-76 itation can at most confirm the temporal periodicity of its equatorial source, but can-77 not resolve the spatial periodicity of the precipitation nor the spatial periodicity of its 78 equatorial sources. The most promising way to reveal both the temporal and spatial scales 79 of electron precipitation patterns is to combine low-altitude, near-equatorial, and ground-80 based measurements. 81

In this study we analyze three events of quasi-periodic electron precipitation driven 82 by near-equatorial electron scattering due to whistler-mode waves modulated by com-83 pressional ULF waves. We combine ground-based magnetometer measurements of ULF 84 waves, low-altitude ELFIN-A and -B (Angelopoulos et al., 2020) measurements of > 50keV 85 electron precipitation, and near-equatorial Time History of Events and Macroscale In-86 teractions during Substorms (THEMIS; (Angelopoulos, 2008)) measurements of whistler-87 mode and ULF waves. Ground-based measurements localize the L-shell and MLT sec-88 tor of ULF waves during the entire interval. Multi-spacecraft THEMIS measurements 89 provide us with estimates of ULF wavelength (spatial scale), which serve as a good proxy 90 for the spatial periodicity scale of whistler-mode wave modulation. Low-altitude ELFIN 91 measurements of quasi-periodic electron precipitation show a spatial periodicity simi-92 lar to the estimated ULF wavelength. Finally, the combination of near-equatorial mea-93 surements of whistler-mode wave characteristics and background plasma properties pro-94 vides typical resonant energies of precipitating electrons, which agree well with the pre-95 cipitating energy spectra in ELFIN measurements. In summary, these three selected events 96 confirm that compressional ULF waves can modulate whistler-mode waves and result 97 in spatially quasi-periodic precipitation of energetic electrons over a large L-shell and 98 MLT domain. 99

The paper is organized as follows: in Sect. 2 we describe available datasets and methods of data analysis; in Sect. 3 we describe three events with combined ground-based, THEMIS, and ELFIN measurements; and in Sect. 4 we discuss our results and their possible application in radiation belt modeling.

¹⁰⁴ 2 Spacecraft Instruments and Dataset

Investigation of electron precipitation requires pitch-angle and energy resolved elec-105 tron distributions within the loss cone, which is almost impossible near the equator due 106 to the small loss cone size there (see, e.g., Kasahara et al., 2018). However, the much larger 107 loss cone size at low altitudes allows polar-orbiting ionospheric spacecraft to measure trapped 108 and precipitating (those within the loss cone) electron fluxes. In this study, we employ 109 energetic (> 50 keV) electron precipitation measurements by the low-altitude (~ 450 110 km) twin ELFIN CubeSats (ELFIN-A and ELFIN-B), which provide electron pitch-angle 111 distributions between 50 keV to 6 MeV, with energy resolution < 40%, angular reso-112 lution of $\sim 22.5^{\circ}$ and temporal resolution of 2.8s (spin period) (Angelopoulos et al., 2020). 113 We use the ratio of j_{loss} (pitch-angle-averaged flux within the loss cone) to j_{trap} (pitch-114 angle-averaged flux outside the loss cone) to study enhancements of precipitation driven 115 by near-equatorial scattering of energetic electrons by whistler waves (see detailed anal-116 ysis of ELFIN measurements of whistler-driven precipitation events in Artemyev et al., 117 2021; Zhang, Artemyev, et al., 2022; Mourenas et al., 2022; Tsai et al., 2022). 118

ELFIN measurements of j_{loss}/j_{trap} are supplemented by ULF wave observations 119 from the THEMIS ground-based magnetometer (GMAG) network ((Russell et al., 2008); 120 FYKN, BRW), USGS magnetometer network (CMO, SHU), AUTUMNX magnetome-121 ter network (PUVR), and University of Iceland magnetometer (LRV). The main advan-122 tage of ground-based observations is the absence of spatio-temporal ambiguity inherent 123 in spacecraft measurements. ELFIN flux ratio j_{loss}/j_{trap} measurements and ground-based 124 ULF observations will be compared with equatorial THEMIS (Angelopoulos et al., 2008) 125 measurements of ULF and VLF (whistler-mode frequency range) fields. The THEMIS 126 fluxgate magnetometer (FGM; (Auster et al., 2008)) provides magnetic field with a 1/128s, 127 1/4s or 3s (spin period) sample rate (depending on availability and choice of data prod-128 uct), whereas the THEMIS search-coil magnetometer (SCM; (Le Contel et al., 2008)) 129 provides < 8 kHz waveform measurements that are further converted to on-board pro-130 cessed, Fast-mode Fast Fourier series, "FFF", spectra data over 32 or 64 frequency bands 131 (Cully et al., 2008). 132

We select three quasi-periodic precipitation events at ELFIN, which are in good conjunction with THEMIS, providing equatorial measurements of ULF and whistler waves, and with ground-based magnetometers provding measurements of ULF waves. We show detailed analysis of the first event, and then reinforce the conclusions using the two other events.

¹³⁹ **3.1 Event #1**

The first event happened on June 17, 2021. Figure 1 shows the projections of THEMIS (three spacecraft, TH-A, TH-D, and TH-E) orbits from 03:30 UT to 06:00 UT, and the projection of ELFIN-A orbits from 05:32:30 UT to 05:33:30 UT. We also include two ground stations (PURV and LRV) located near the ELFIN-A (EL-A) and THEMIS footpoints.

Figures 2(a-c) show the wavelet analysis of the ULF magnetic field component ob-144 served by THEMIS-E and two ground stations: the parallel (compressional) component 145 is shown for THEMIS observations, and the East-West component (corresponding to poloidal 146 oscillations in space) for the ground stations. The ground stations and THEMIS-E ob-147 served compressional ULF waves with similar frequency ranges, indicating that the same 148 ULF waves existed over a large MLT (22-03) and L-shell $\in [6, 9]$ domain (L shells here 149 are calculated using the T89 model (Tsyganenko, 1989)). During this event, THEMIS-150 E was in fast survey mode (when FFF data is available) and SCM measured quasi-periodic 151 whistler wave bursts from 03:30 UT to 04:30 UT at $L\sim6-9$ (see Panel (d)). Within the 152



Figure 1. Projection of ELFIN-A orbits and THEMIS orbits to the ground, and the location of two ground stations on June 17, 2021, from 03:30 UT to 06:00 UT. The red trace along ELFIN-A orbit shows the sub-interval (05:32:30 UT to 05:33:30 UT) analyzed in Figure 3. The dots mark the start time of each orbit.

same L-shell range, THEMIS-A and -D also observed similar whistler wave bursts (not 153 shown here), consistent with expectation from the ground stations that the periodic whistler 154 emission extended over a wide MLT range in space. Panel (e) shows the THEMIS-E whistler 155 wave spectrum (in color) and the line-plot (black trace) of the field-aligned magnetic field 156 component variation ($\delta B z = \delta B_{\parallel}$) of ULF waves between ~9-12mHz. ULF waves in this 157 frequency range are visible in THEMIS-E, consistent with the wave-power in the ground-158 based station measurements (see Panels (a-c)). There is a reasonably good correlation 159 of whistler wave bursts and local δB_{\parallel} minima (see expanded view of Panel (e) in Panel 160 (g)). These observations are consistent with a scenario of a ULF-wave modulated ther-161 mal electron anisotropy, resulting in the observed quasi-periodic generation of whistler 162 waves (as previously reported by W. Li, Bortnik, et al., 2011; Xia et al., 2020; Zhang et 163 al., 2019; L. Li et al., 2022). 164

To determine the typical energies of electrons precipitated by the observed quasiperiodic whistler waves, we calculate the mean wave frequency $\langle f \rangle$, and the frequency width Δf , using the following equations:

$$\begin{split} \langle f \rangle &= \frac{\int_{f_{lh}}^{f_{ce}/2} B_w^2(f) f df}{\int_{f_{lh}}^{f_{ce}/2} B_w^2(f) df} \\ \Delta f &= \left(\frac{\int_{f_{lh}}^{f_{ce}/2} B_w^2(f) \left(f^2 - \langle f \rangle^2\right) df}{\int_{f_{lh}}^{f_{ce}/2} B_w^2(f) df}\right)^{1/2} \end{split}$$

where B_w^2 is the wave intensity at a specific frequency, f_{lh} is the lower hybrid frequency, and f_{ce} is the electron cyclotron frequency. Figure 2(d) shows that the mean wave frequency $\langle f \rangle$ (depicted by the red crosses), is between 200-800 Hz (0.2-0.4 f_{ce}) and decreases with increasing *L*-shell. We calculate the cyclotron resonance energy of electrons for fieldaligned whistler waves at the estimated mean frequency, $\langle f \rangle$, and at the minimum frequency, $f_{min} = \langle f \rangle - \Delta f$, using the equation (Kennel & Petschek, 1966):

$$E_{res} = \frac{B^2}{2\mu_0 n_e} \frac{f_{ce}}{f} \left(1 - \frac{f}{f_{ce}}\right)^3$$

where B is the magnetic field strength, estimated using the dipole model scaled at the equatorial field intensity from THEMIS-E measurements, and n_e is the electron density



Figure 2. Observations of ULF and whistler waves. Wavelet power spectra of the East-West component magnetic field at two ground stations (Panels a, b) and of the parallel magnetic field component at TH-E (Panels c,d, covering the ULF and VLF range respectively). Panels (a-c) also demarcate, in white rectangles, the ULF band and time range of enhanced magnetic field fluctuations of interest. Panel (d) denotes peak whistler wave power in red crosses at 1min cadence. Over-plotted in solid and dashed lines are f_{ce} and $0.5f_{ce}$. Panels (e, f) show the same overview of the VLF waves at TH-E, containing the same information as Panel (d), except only showing intensities greater than 10^{-7} nT²/Hz. Overploted on them are the band-pass filtered waveforms of δB_{\parallel} of ~10 mHz ULF waves, and f_{pe}/f_{ce} , respectively. Panel (g) is a zoomed-in view, in time, of Panel (e).

from THEMIS-E equatorial measurements assumed to be constant along magnetic field 167 lines. As the resonance energy increases with latitude, it is important to constrain waves 168 to a reasonable latitudinal extent: $|\lambda| \sim 30^{\circ}$ is used here as derived from the empiri-169 cal whistler wave model (see Meredith et al., 2001, 2003; O. V. Agapitov et al., 2018) 170 at the MLT sector of our observations. The resultant resonance energies are in the range 171 of \sim 50-300 keV, showing that the observed whistlers can provide scattering and lead to 172 precipitation of electrons in this energy range. It is worth noting that for the observed 173 $f_{pe}/f_{ce} \sim 6-12$, typical electromagnetic ion cyclotron (EMIC) waves can only lead to the 174 loss of >1 MeV electrons (e.g., Kersten et al., 2014), and therefore precipitation of 100s 175 of keV electrons should be mostly attributed to scattering by whistler waves. We antic-176 ipate that the quasi-periodic whistler waves, modulated by ULF waves, will lead to quasi-177 periodic energetic electron precipitation. 178

To test this hypothesis, we examine ELFIN-A observations of precipitating elec-179 tron fluxes j_{loss} , trapped electron fluxes j_{trap} , and their ratio j_{loss}/j_{trap} (Figure 3). At 180 the same L-shell and MLT sector where THEMIS observed ULF-modulated whistler waves, 181 ELFIN-A indeed captured quasi-periodic precipitation of $\sim 50-200$ keV electrons with 182 average peaks of j_{loss}/j_{trap} greater than 0.3, indicative of fast scattering caused by whistler 183 waves. Note that moving along a low-altitude orbit, ELFIN crosses the entire L-shell range 184 of precipitation within a couple of minutes, and thus the periodicity at ELFIN obser-185 vations represents a spatial periodicity of the scattering process at and near the equa-186 tor. 187

It is evident that these periodic j_{loss}/j_{trap} peaks contribute to most of the precipitating flux observed by ELFIN at L = 6 - 9. This indicates that such periodic precipitation can play a major role in energetic electron losses during intervals with ULFmodulated whistler bursts. The highest energy channel showing the periodic j_{loss}/j_{trap} peaks is around 200 keV (Figure 3), located in Figures 3b, 3c) between the upper (black crosses) and mean (white crosses) resonance energies corresponding to $\langle f \rangle$ and f_{min} , respectively. Thus, ELFIN measurements confirm that quasi-periodic whistler wave bursts observed by THEMIS are indeed responsible for quasi-periodic electron precipitation.

However, the periodicity seen by THEMIS is temporal, whereas the periodicity seen 196 by ELFIN is spatial. ULF waves observed by THEMIS extend over a large L-shell and 197 MLT sector (as revealed by ground-based observations), and thus ELFIN likely measures 198 precipitation from spatially periodic whistler bursts with the periodicity comparable to ULF wavelengths. To confirm this, we compare the spatial scale of periodic precipita-200 tion peaks (δL) and our estimate of ULF wavelengths $\delta \lambda$. ELFIN moves along its or-201 bit at a nearly constant geomagnetic longitude. Therefore, when mapped to the equa-202 tor, ELFIN's trajectory, and the measured quasi-periodic precipitation, correspond ap-203 proximately to a radial sampling of the equatorial magnetosphere. The spatial period-204 icity (δL) of the precipitation can be estimated by calculating the spatial separations be-205 tween the peaks of j_{loss}/j_{trap} . For three observed j_{loss}/j_{trap} peaks, the average spatial 206 scale is $0.85 \pm 0.28 R_E$ (Figure 3e, where ELFIN time-series of precipitation were converted 207 to L-shell profiles of precipitation using the T89 model). If the precipitation is driven 208 by ULF-modulated whistler waves, then this δL should be comparable to the ULF wave-209 length $\delta\lambda$ in the radial direction (or, more precisely, the ULF field δB_{\parallel} periodicity in the 210 radial direction). The wavelength in the radial direction can be calculated by using cross-211 correlation analysis on the ULF fields measured from two THEMIS spacecraft at sim-212 ilar MLT (mostly separated in the radial direction) to obtain the phase difference for the 213 ULF wave. If the two spacecraft are radially separated by a distance r, and δB_{\parallel} (at a 214 particular frequency f) at the two spacecraft has a phase difference $\delta\phi$, then the wave-215 length can be estimated as $\lambda/2\pi = r/\delta\phi$. The phase difference between two spacecraft 216 can be calculated using $\delta \phi = 2\pi \delta t/T$, where δt is the lag time as inferred from the peak 217 cross-correlation between the ULF wave field measured at the two spacecraft and T is 218 the period of the wave. We use THEMIS-A and THEMIS-E measurements separated mostly 219

along the radial direction over distances comparable with the expected ULF wavelength. 220 We examined the phase difference of 9-12 mHz ULF waves between THEMIS-A and 221 THEMIS-E; Figure 3f shows that the cross-correlation peaks at a lag time of $\delta t \sim 26.5$ s 222 (the comparison of observed ULF waveforms is shown in the Supplementary Informa-223 tion). The corresponding wavelength is $\delta\lambda \sim 0.78R_E$, which is very close to the $\delta L \approx 0.85 \pm 0.28R_E$ 224 derived above from ELFIN measurements. These spatial dimensions are consistent with 225 past studies of ULF-induced precipitation (e.g., Baddeley et al., 2017), and provide fur-226 ther support for the idea that the periodic precipitation is driven by ULF-modulated pe-227 riodic whistler waves. 228

229 **3.2 Event #2**

The second conjunction event occurred on May 13, 2021. Figure 4 shows the or-230 bit projections of THEMIS-A, -D, and -E in the northern hemisphere from 17:00 UT to 231 19:00 UT, and the ELFIN-B orbit projection from 17:12 UT to 17:15 UT. Two ground 232 stations (BRW and FYKN) were located near the ELFIN-B and THEMIS footpoints at 233 the time. The ground stations and three THEMIS spacecraft observed ULF waves over 234 a similar frequency range (Figures 5a-c). We are mostly interested in ULF waves from 235 3 mHz to 5 mHz (depicted by the rectangle in Figure 5c), because THEMIS-E observed 236 quasi-periodic whistler wave bursts (from 17:00 UT to 19:00UT, see Fig. 5(d)) modu-237 lated at a frequency $\sim 4 \text{mHz}$; see Fig. 5(e) for the correlation of the peaks of whistler 238 wave intensity and δB_{\parallel} of compressional ULF wave component band-pass filtered at 3– 230 5mHz. The whistler waves' mean frequency increases from 150 Hz to 700 Hz (0.1-0.2 f_{ce}) 240 as THEMIS-E moves to lower L-shells. Such whistler waves can resonate with \sim 50-1000 241 keV electrons for the observed $f_{pe}/f_{ce} \sim 6 - 20$ (see Fig. 5(f)). Note that for such a 242 low f_{pe}/f_{ce} , we can ignore the contribution of EMIC waves to the scattering of $\lesssim 1 \text{MeV}$ 243 electrons (Summers et al., 2007). 244

Around 17:13 UT, ELFIN-B traversed the same L-shell and MLT region where THEMIS 245 captured quasi-periodic whistler waves. ELFIN-B observed strong bursts of electron pre-246 cipitation with the precipitating to trapped flux ratio reaching (or even slightly exceed-247 ing) one for 10s-100s of keV electrons (see Fig. 5h-i). The periodic peaks of j_{loss}/j_{trap} 248 are observed up to $\sim 500-600$ keV. As in the previous event studied, this energy range 249 is consistent with the resonance energies of electrons for the observed whistler waves (with 250 an assumption of wave propagation up to $\sim 30^{\circ}$ of magnetic latitude; the typical lat-251 itudinal spread of whistler waves at this MLT sector, see Meredith et al. (2001, 2003); 252 O. V. Agapitov et al. (2018)). Within the quasi-periodic precipitation, ELFIN also ob-253 served spin-resolution bursts with $j_{loss}/j_{trap} \ge 1$ up to 600 keV, around the mean res-254 onance energies estimated for THEMIS measurements of whistler waves and background 255 plasma. These are likely very short, microburst-like precipitation lasting less than one 256 spin of ELFIN (see detailed analysis of such microbursts observed by ELFIN in Zhang, 257 Angelopoulos, et al., 2022). Miyoshi et al. (2020); Shumko et al. (2021) previously re-258 ported that relativistic microbursts may be embedded into quasi-periodic precipitation 259 of low-energy (< 50 keV) electrons. Our observations of strong precipitation at the min-260 imum detectable energy ($\approx 50 \text{keV}$), are suggestive that precipitation extends to < 50261 keV and support the idea that < 50 keV and > 500 keV losses can occur simultaneously. 262 Furthermore they indicate that such broad energy precipitation can occur during ULF-263 driven, quasi-periodic precipitation.

Figure 5(k) marks the local precipitation maxima (peaks) in the plot of j_{loss}/j_{trap} as a function of *L*. Only peaks with $j_{loss}/j_{trap} > 0.3$ have been considered. The spatial periodicity of these peaks is $\delta L = 0.5 \pm 0.17R_E$. The phase difference of 3-5 mHz ULF waves between THEMIS-A and THEMIS-E using the method described for Event #1 (see Figure S1 (b) in Supplementary information) results in an inferred equatorial ULF wavelength of $0.62R_E$, approximately in the radial direction, which is close to the δL derived from ELFIN measurements.



Figure 3. ELFIN-observed trapped electron flux j_{trap} (a), precipitating flux j_{loss} (b), and the ratio j_{loss}/j_{trap} (c and d, in spectral- and line-plot format respectively). Black and white crosses in Panels (b, c) show the upper and mean resonance energy of the observed whistler waves (see text for details). (e) The average j_{loss}/j_{trap} of the first four energy channels. (f) Cross-correlation between the ULF wave field from THEMIS-A and THEMIS-E as a function of lag time.



Figure 4. Projection of ELFIN-B and THEMIS orbits to the ground, and the location of two ground stations on May 13, 2021, from 17:00 UT to 19:00 UT. The red trace along ELFIN-B's orbit shows the sub-interval (17:12:30 UT to 17:14:40 UT) analyzed in Figure 5. The dots mark the start time of each orbit.

272 **3.3 Event #3**

The third event occurred on Oct 22, 2021. Figure 6 shows the orbit projections of 273 THEMIS from 10:00 UT to 12:00 UT and the ELFIN-A orbit projection from 11:48 UT 274 to 11:51 UT. THEMIS-E observed quasi-periodic whistler waves modulated by 10-20 mHz 275 compressional ULF waves during 10:20 UT to 10:50 UT (Figure 7(c)). For this event, 276 THEMIS and ELFIN footpoints in the north hemisphere are located near MLT = 23. 277 No ground-based stations are available in the same MLT, however, ground-based sta-278 tions (SHU and CMO) at MLT = 0-1 detected the same, 10-20 mHz frequency ULF 279 wave (Figures 7a, 7b)), therefore the ULF waves covered a large MLT, L-shell domain. 280 ELFIN-A observed quasi-periodic electron precipitation within $L \in [5.5, 7.5]$ (correspond-281 ing to whistler wave bursts observed by THEMIS) with $\delta L = 0.33 \pm 0.045$ (δL is the 282 average spatial scale between the precipitation peaks shown in Figure 7(k). The pre-283 cipitating to trapped flux ratio of 10s-100s of keV electrons exhibits peaks at around 0.1-284 0.3. The highest energy channel showing the periodic peaks, ~ 300 keV, is between the 285 upper and mean resonance energies estimated for THEMIS measurements of whistler waves 286 using the local plasma density (assumed to be the same as at the equator). The wave-287 length of the 10-20 mHz ULF waves (estimated by comparing THEMIS-A and THEMIS-288 E ULF measurements) is $\delta\lambda \sim 0.27R_E$ (the ULF fields are shown in Figure S1 (c) in 289 the Supplementary information), consistent with the spatial periodicity of electron pre-290 cipitation. Both Events #2 and #3 support the hypothesis that the ULF-modulated pe-291 riodic whistler waves lead to the observed periodic precipitation. 292



Figure 5. Observations of ULF and whistler waves. East-West component magnetic field at two ground stations (Panels a, b) and of the parallel magnetic field component at TH-E (Panels c,d, covering the ULF and VLF range respectively). Panels (a-c) also demarcate, in white rectangles, the ULF band and time range of enhanced magnetic field fluctuations of interest. Panel (d) denotes peak whistler wave power in red crosses at 1min cadence. Over-plotted in solid, dashed, and dashed-dotted lines are f_{ce} , $0.5f_{ce}$, and f_{lh} , respectively. Panels (e, f) show the same overview of the VLF waves at TH-E, containing the same information as Panel (d), except only showing intensities greater than $10^{-7}nT^2/Hz$. Overploted on them are the band-pass filtered waveforms of δB_{\parallel} of ~4 mHz ULF waves, and f_{pe}/f_{ce} , respectively. The vertical red line indicates the time of ELFIN-B pass. Panels (g, h and i) show ELFIN observations of trapped electron flux j_{trap} , precipitating flux j_{loss} , and the ratio j_{loss}/j_{trap} , respectively. Black and red crosses show the upper and mean resonance energy of whistler waves. Panel (j) is the same information as Panel (i) except only for the 4 lowest energies, with the ratios depicted as line-plots. Panel (k) is the average j_{loss}/j_{trap} of the first four energy channels.



Figure 6. Projection of ELFIN-A orbits and THEMIS orbits to the ground, and the location of two ground stations on Oct 22, 2021, from 10:20 UT to 10:50 UT. The red trace along the ELFIN-A orbit projection shows the sub-interval (11:48:50 UT to 11:49:40 UT) analyzed in Figure 7. The symbols mark the start time of each orbit.



Figure 7. ULF wavelet spectra at two ground stations and THEMIS-E, spectra of whistler waves at THEMIS-E, and electron precipitation observed at ELFIN-A, in the same format as Figure 5.

²⁹³ 4 Discussion and Conclusions

Compressional ULF waves are known to modulate equatorial electron populations 294 causing whistler-mode wave generation (W. Li, Thorne, Bortnik, Nishimura, & Angelopou-295 los, 2011; Xia et al., 2016, 2020; Zhang et al., 2019), resulting quasi-periodic electron pre-296 cipitation (Nishimura et al., 2010; Kasahara et al., 2018). Such precipitation is not only 297 modulated by ULF waves but actually generated by such waves: the ULF waves can drive 298 marginally unstable plasmas of the inner magnetosphere to whistler-mode instability which 200 can cause subsequent electron precipitation (Xia et al., 2016; Zhang et al., 2019). The 300 subject ULF waves are typically generated at the magnetopause by solar wind transients 301 and propagate towards the inner magnetosphere (e.g., O. V. Agapitov et al., 2009; Hartinger 302 et al., 2013, 2014; Hwang & Sibeck, 2016). Both the ULF waves (Klimushkin et al., 2019; 303 Wright & Elsden, 2020; Elsden et al., 2022; Di Matteo et al., 2022) and the correlated 304 whistler-mode waves (Zhang et al., 2020) have been observed to extend over a large L-305 shell and MLT range. However, the efficiency of electron precipitation by such ULF-driven 306 whistler-mode waves (e.g., the precipitating electron energy range, flux magnitude and 307 net contribution to the ionospheric energy input) cannot be reliably determined from equa-308 torial spacecraft measurements alone. In this study we use low-altitude observations of 309 such precipitation in conjunction with ground based and equatorial measurements of the 310 ULF and whistler waves to show that during three events: 311

- Electron precipitation exhibits spatial periodicity (in *L*-shell) with a scale com-312 parable to that of equatorial ULF wavelengths. ULF waves and the associated whistler-313 mode wave driven precipitation extend over L-shells from the plasmasphere to the 314 magnetopause (in agreement with Wright & Elsden, 2020; Zhang et al., 2020; Sandhu 315 et al., 2021). The spatially and temporally periodic ULF waves and whistler-mode 316 waves, and lead to similar periodic scattering of energetic electrons across the en-317 tire L-shell region and, by inference, the wide MLT region where ULF waves are 318 present and can generate whistlers. 319
- The energy of the precipitating electrons can range from below the low energy limit 320 of the ELFIN instrument (50keV) to upwards of 100 keV, to as high as ~ 1 MeV. 321 The upward limit is consistent with the estimated maximum resonance energy of 322 the observed whistler-mode waves, for the observed equatorial plasma conditions, 323 and empirical constraints on the wave distribution along magnetic field lines (see 324 O. V. Agapitov et al., 2018). This energy range of precipitation covers the entire 325 radiation belt "seed" electron population, and extends upwards into the low-energy 326 portion of radiation belt electrons (Javnes, Baker, et al., 2015; Boyd et al., 2018; 327 Turner et al., 2021). 328
- The observed quasi-periodic precipitation of $\sim 500 \text{keV-1}$ MeV electrons which is 329 squarely attributed to quasi-periodic whistler-mode waves, can only be the result 330 of resonance at middle to high magnetic latitudes (where the local f_{pe}/f_{ce} can be 331 sufficiently low). For the waves to propagate to such high latitudes (along the field 332 line) without becoming oblique and Landau-damped by the thermal electrons, it 333 must be that they are ducted (see discussion in Artemyev et al., 2021) by plasma 334 density gradients (e.g., Hanzelka & Santolík, 2019; Streltsov & Bengtson, 2020; 335 Chen et al., 2021). The ducts themselves must be also be set up by the compres-336 sional character of the ULF waves, since the observed relativistic electron precip-337 itation exhibits the spatial structure of the ULF waves. 338
- These results confirm the important role of ULF waves in precipitating not only auroral (< 10 keV) electrons (Nishimura et al., 2010; Kasahara et al., 2018), but also energetic (~ 100 keV) and even relativistic (> 500 keV) electrons. Therefore, whistlermode waves generated by ULF-modulated thermal electron anisotropy can contribute significantly to radiation belt dynamics. The generation of whistler-mode waves by ULFwave modulation of the plasma sheet electrons is very different from the classical mechanism of whistler-mode wave generation by electron injections during substorms (e.g.,

Thorne et al., 2010; Tao et al., 2011; Fu et al., 2014). Thus, the electron precipitation 346 discussed herein does not necessary coincide with geomagnetic activity associated with 347 substorms (identified using the AE and AL indices). Such a mechanism of energetic elec-348 tron precipitation during instantaneously low AE intervals can still significantly deplete 349 the remnant radiation belts that may have been built up by prior activity. Yet, this mech-350 anism may be underestimated or completely absent in models of inner magnetosphere 351 dynamics and magnetosphere-ionosphere coupling. While ULF waves can also be gen-352 erated by substorm-time injections (e.g., Runov et al., 2014; Liu et al., 2017), magne-353 topause buffeting by solar wind pressure variations and by ion foreshock transients (e.g., 354 Hartinger et al., 2013, 2014; Hwang & Sibeck, 2016) is thought to result in the largest 355 amplitude ULF waves (e.g., Di Matteo et al., 2022). Examination of a wider variety of 356 events using the methods presented herein can determine how the spatial scale of ULF-357 modulated precipitation varies according to the ULF wave driver. Future parameteri-358 zation of ULF-driven whistler-mode waves and inclusion of this wave population in ra-359 diation belt models can improve our understanding of solar wind transients as a driver 360 of electron precipitation, through ULF wave generation at the dayside magnetopause and 361 subsequent whistler wave generation as ULF waves propagate inward as well as towards 362 the flanks. 363

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³⁷⁵ Open Research

ELFIN data is available at https://data.elfin.ucla.edu/. Data analysis was done using SPEDAS V4.1 Angelopoulos et al. (2019) available at https://spedas.org/.

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Supporting Information for "On the role of ULF waves in the spacial and temporal periodicity of energetic electron precipitation"

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1. Figures S1

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Figure S1. Comparisons of ULF waves magnitude δB_{\parallel} measure by THEMIS-A and THEMIS-E. Blue lines show the ULF waves measured by THEMIS-E. Red lines show the THEMIS-A observations after shifted by the lag time with maximum correlation.

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