# De-tuning a coupled Climate Ice Sheet Model to simulate the North American Ice Sheet at the Last Glacial Maximum

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# Abstract

The maximum extent of the last North American ice sheet is well constrained empirically, but it has proven to be challenging to simulate with coupled Climate-Ice Sheet models. Coupled Climate-Ice Sheet models are often too computationally expensive to sufficiently explore uncertainty in input parameters, and it is unlikely values calibrated to reproduce modern ice sheets will reproduce the known extent of the ice at the Last Glacial Maximum. To address this, we run a series of ensembles with a coupled Climate-Ice Sheet model (FAMOUS-ice), simulating the final stages of growth of the last North American Ice Sheets' maximum extent. Using this large ensemble approach, we explore the influence of uncertain ice sheet, albedo, atmospheric, and oceanic parameters on the ice sheet extent. We find that albedo parameters determine the majority of uncertainty when simulating the Last Glacial Maximum North American Ice Sheets. Importantly, different albedo parameters are needed to produce a good match to the Last Glacial Maximum North American Ice Sheets than have previously been used to model the contemporary Greenland Ice Sheet, due to differences in cloud cover over ablation zones. Thus calibrating coupled climate-ice sheet models solely for present day strongly biases simulations of past and future climates different from today.

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# De-tuning a coupled Climate Ice Sheet Model to simulate the North American Ice Sheet at the Last Glacial Maximum

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# Key Points:

14	•	Simulating the Last Glacial Maximum Laurentide Ice Sheet with the coupled Climate-
15		Ice Sheet model FAMOUS-ice requires re-tuning
16	•	We efficiently rule out unsuitable parameter values using a Gaussian Process em-
17		ulator to design waves of short and long simulations
18	•	The different cloud cover between North America and Greenland makes the Last
19		Glacial Maximum a good test for climate-ice-sheet models

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# 20 Abstract

The maximum extent of the last North American ice sheet is well constrained em-21 pirically, but it has proven to be challenging to simulate with coupled Climate-Ice Sheet 22 models. Coupled Climate-Ice Sheet models are often too computationally expensive to 23 sufficiently explore uncertainty in input parameters, and it is unlikely values calibrated 24 to reproduce modern ice sheets will reproduce the known extent of the ice at the Last 25 Glacial Maximum. To address this, we run a series of ensembles with a coupled Climate-26 Ice Sheet model (FAMOUS-ice), simulating the final stages of growth of the last North 27 28 American Ice Sheets' maximum extent. Using this large ensemble approach, we explore the influence of uncertain ice sheet, albedo, atmospheric, and oceanic parameters on the 29 ice sheet extent. We find that albedo parameters determine the majority of uncertainty 30 when simulating the Last Glacial Maximum North American Ice Sheets. Importantly, 31 different albedo parameters are needed to produce a good match to the Last Glacial Max-32 imum North American Ice Sheets than have previously been used to model the contem-33 porary Greenland Ice Sheet, due to differences in cloud cover over ablation zones. Thus 34 calibrating coupled climate-ice sheet models solely for present day strongly biases sim-35 ulations of past and future climates different from today. 36

# <sup>37</sup> Plain Language Summary

At the peak of the last ice age, an ice sheet covered much of North America. The extent of this ice sheet is well-understood after decades of intensive data collection, but producing a computer simulation of the ice sheet which matches our observations has been a challenge. This is partly because of uncertainty about the "correct" model setup to create the best simulation, and partly because the computer models used in the simulations require large computing resources.

In this paper, we present a series of simulations of the North American ice sheet 44 at the peak of the last ice age using a fast-running computer model in which the atmo-45 sphere and ice sheets interact. We run hundreds of simulations to tackle the uncertainty 46 about the optimum values for unknown input parameters. We find that the model's rep-47 resentation of how reflective the ice sheet surface is has the most impact on the size and 48 shape of the simulated ice sheet. Importantly, the parameter values that produce the best 49 simulations of modern-day Greenland produce poor simulations of the North American 50 ice sheets during the last ice age, calling into question whether the parameters chosen 51 for modern Greenland will produce reasonable simulations of future ice sheet change and 52 sea level rise. 53

#### 54 1 Introduction

Accurately estimating future changes in ice sheets is crucial for producing mean-55 ingful projections of future sea level rise (IPCC, 2021). Ice sheets interact with the at-56 mosphere and ocean, and are vulnerable to instabilities in their growth and retreat (e.g. 57 Shepherd et al., 2012; J. Gregory et al., 2012). These instabilities, along with the un-58 certainties in processes of climate and ice sheet evolution, make future projections us-59 ing numerical models difficult, and the accuracy of any future simulations is a challenge 60 to assess. For the Greenland ice sheet, one of the main sources of uncertainty is the fu-61 ture changes in surface mass balance (the balance of accumulation and melt of snow and 62 ice at the surface) (Fettweis et al., 2011). This surface mass balance is highly dependent 63 on both the climate and the ice sheet topography, as well as the strong interactions be-64 tween the two. Thus projections of future Greenland evolution need to account for climate-65 ice sheet interactions (Goelzer et al., 2017). However, there are major challenges in rep-66 resenting surface mass balance and climate ice sheet interactions in models: (i) climate 67 models often have large biases in ice sheet regions (Davy & Outten, 2020), (ii) these re-68

gions are difficult environments to work in, limiting the observations we have of the cli mate and surface mass balance (Vernon et al., 2013) and (iii) surface melt occurs in nar-

- <sup>71</sup> row steep regions at the edge of the ice sheets that are difficult processes to capture or
- <sup>72</sup> represent in global climate models.

Major progress has been made to tackle these challenges and there are now sev-73 eral earth system models that include interactive ice sheets in Greenland and/or Antarc-74 tica (e.g. Danabasoglu et al., 2020; R. S. Smith et al., 2021). However, there are numer-75 ous challenges in simulating climate-ice sheet interactions. Perhaps most acutely, there 76 77 is a mismatch of spatial and temporal scales between typical ice sheet and global climate models. Spatially, kilometer (or sub-kilometer) processes are important for accurate sim-78 ulation of ice sheet processes, such as grounding line migration, and margin precipita-79 tion gradients (Franco et al., 2012; Cornford et al., 2013), so recent ice sheet models have 80 been developed to simulate ice sheets at this scale, e.g. through adaptive mesh refine-81 ment, and they require climate (or surface mass balance inputs) at that scale. On the 82 other hand, Atmosphere Ocean General Circulation Models (AOGCMs) have a grid-box 83 size  $10-1000 \times$  larger than the scale of modelled ice sheet processes and inputs (Sellar et al., 2019; Danabasoglu et al., 2020). This conundrum of mistmatching scales is flipped 85 in the time domain. Temporally, AOGCMs require sub-daily timesteps to accurately sim-86 ulate the climate system, while ice sheet change is usually a multi-centennial process. 87 The spatial-temporal mismatch of scales creates a problem for computational efficiency, 88 since high resolution is needed for both and isn't easily compromised for one in favour 89 of the other. Consequently, most AOGCMs cannot practically simulate interactive ice 90 sheet change, instead prescribing the ice sheet extent as a boundary condition (e.g. Kageyama 91 et al., 2017; Ivanovic et al., 2015; Menviel et al., 2019) and updating the ice sheet pe-92 riodically for palaeo runs where significant ice sheet change occurs. Similarly, ice sheet 93 simulations often rely on prescribed climate or surface mass balance fields (e.g. L. J. Gre-94 goire et al., 2016; Patton et al., 2013; Gandy et al., 2021). 95

Nonetheless, recent technical advances have allowed coupled simulations of the cli-96 mate and ice sheets, through a combination of model development and increasing com-97 pute power. The spatial mismatch between ice sheet and climate model grids can be ad-98 dressed by calculating ice-sheet relevant processes, such as albedo calculations, with sub-99 gridscale parameterisations (Ganopolski et al., 2010; Vizcaíno et al., 2013; Ziemen et al., 100 2014; R. S. Smith et al., 2020, 2021). One way to solve the problem of time scale (sub-101 daily timesteps for multimillennial integrations) is to couple the ice sheet to climate mod-102 els which are typically computationally efficient, in part due to having relatively low spa-103 tial resolution, meaning the the simulations can be run for the length of time needed to 104 spin-up and simulate the co-evolution of ice sheets and climate. 105

A remaining challenge is how coupled climate-ice sheet models should be calibrated 106 and tested. For many models, uncertain parameters are hand-tuned to produce stable 107 modern ice sheets of the right shape and size compared to observations. However, this 108 only represents one point in time. The recent past, for which we have direct observations 109 of ice sheets has seen relatively small changes compared to those expected in the next 110 centuries. These observations thus provide poor constraints on the strength of climate-111 ice sheet feedbacks and there is a danger of over-fitting the model to modern conditions 112 and compensating for biases in the simulated climate. To have confidence in the abil-113 ity of coupled climate-ice sheet models we need to test them under conditions different 114 from today where we have sufficient observational constraints on the climate and ice sheets. 115

We propose to use the Last Glacial Maximum as a benchmark for coupled climateice sheet models. This period, which occurred around 21,000 years ago, has been a focus of the Palaeoclimate Model Intercomparison Project since the 1990s (Kageyama, Abe-Ouchi, et al., 2021), because it was a period of relatively stable and well documented climate, with CO2 concentrations much lower than today (180 ppm). The North American ice sheet is thought to have reached a relatively stable maximum extent, that is very well reconstructed (e.g. Dyke et al., 2002; Peltier et al., 2015; Gowan et al., 2021). It is
thus possible to run equilibrium simulations under LGM conditions with an interactive
North American ice sheet until a stable maximum ice extent is reached, which can be
meaningfully compared to reconstructions. Ice at the LGM reached much lower latitudes
than today, providing a way to test the ability of models to represent SMB and climateice interactions under energy balance conditions different than modern Greenland.

We use FAMOUS-Ice (R. S. Smith et al., 2020), a coarse resolution, fast running 128 AOGCM, which has been used in long-palaeo simulations (R. S. Smith, 2012; J. Gregory 129 130 et al., 2012; Roberts et al., 2014; Dentith et al., 2019; L. J. Gregoire et al., 2012) and uncertainty quantification (L. Gregoire et al., 2010), coupled to the Glimmer Ice Sheet 131 model by downscaling SMB calculations (R. S. Smith et al., 2020), rather than previ-132 ous work using a PDD SMB scheme (J. Gregory et al., 2012). This coupled model has 133 been used to simulate present and future Greenland Ice Sheet evolution (J. M. Gregory 134 et al., 2020). We start the manuscript by presenting the first attempt at simulating the 135 LGM with the FAMOUS-ice model with interactive ice sheets in Greenland and North 136 America. Here, we use an atmosphere-only version of FAMOUS, prescribing Sea Sur-137 face Temperatures (SSTs) and Sea Ice Concentrations (SICs) in order to minimise bi-138 ases in surface climate. We found that with the standard tuning that reproduced the mod-139 ern Greenland shape and size (R. S. Smith et al., 2020) the model produces a collapse 140 of the North American ice sheet under LGM climatic conditions. We then present the 141 efficient method we developed to generate large ensembles of coupled climate ice sheet 142 simulations simultaneously varying uncertain parameters controlling the climate, ice sheet 143 and surface mass balance. Finally we show that albedo parameters are the primary con-144 trol on uncertainty in our ensemble and we investigate the link between, cloudiness, albedo 145 and surface mass balance under modern and glacial conditions. 146

# <sup>147</sup> 2 Model description and setup

FAMOUS-ice is a fast climate model coupled to the Glimmer ice sheet model (R. S. Smith 148 et al., 2020). The atmospheric component is that of FAMOUS, a fast low resolution gen-149 eral circulation model designed for running simulations of the climate on multi-millennial 150 timescales (e.g. Dentith et al., 2019; L. J. Gregoire et al., 2012) and large ensembles for 151 calibration or uncertainty quantification purposes (L. Gregoire et al., 2010). The Glim-152 mer ice sheet model is a fast simplified 3D dynamical ice sheet model based on the shal-153 low ice approximation that is used to simulate continental ice sheets over glacial inter-154 glacial cycles (L. J. Gregoire et al., 2015; Rutt et al., 2009). A multilayer surface snow 155 scheme is used in the atmosphere model to calculate surface mass balance (SMB) at ten 156 different elevation levels within each grid cell that contains part of an ice sheet. The SMB 157 calculated by the atmosphere model is regridded from the coarse FAMOUS-ice grid  $(7.5^{\circ})$  longitude 158 by 5° latitude) onto the surface of the Glimmer ice sheet model (in this case  $40 km \times 40 km$ ) 159 each model year. The use of ten elevation levels on which SMB is calculated provides 160 a mean to effectively "downscale" the SMB from the coarse atmospheric grid to the finer 161 ice sheet grid. Coupled climate-ice sheet simulations would usually be too computation-162 ally expensive to run as part of large multi-millennial ensembles, but the speed of both 163 FAMOUS and Glimmer make these experiments possible. 164

During the development of FAMOUS-ice, albedo parameters were manually tuned 165 to simulate a stable Greenland ice sheet at present day (R. S. Smith et al., 2020). The 166 model has been applied to evaluate the long-term future decline of the Greenland Ice Sheet 167 (J. M. Gregory et al., 2020). In both cases, sea surface temperatures and sea ice concen-168 trations in FAMOUS were prescribed from the output of higher resolution and complex-169 ity climate models to allow some control of the climate evolution and reduce the impact 170 of biases resulting from atmosphere-ocean and sea ice interactions. FAMOUS can also 171 be used with a dynamical ocean (e.g. Dentith et al., 2019) and the Glimmer ice sheet 172

<sup>173</sup> component can be replaced with the more complex and computationally expensive BISI-<sup>174</sup> CLES ice sheet model (Cornford et al., 2013; Matero et al., 2020; Gandy et al., 2018).

We setup FAMOUS-ice to simulate the climate of the Last Glacial Maximum, with 175 interactive North American and Greenland ice sheets. We follow the PMIP4 LGM pro-176 tocol (Kageyama et al., 2017) to setup most of the climate boundary conditions, includ-177 ing the  $CO_2$ ,  $CH_4$ , and  $N_2O$  concentrations. Global orography and the land-sea mask 178 was taken from the 21 ka BP Glac-1D reconstruction (Tarasov et al., 2012b), and prein-179 dustrial (PI) vegetation distribution was implemented and kept constant throughout the 180 run. Thus our simulation neglect the effects of climate-ice-vegetation feedbacks that can 181 affect ice sheet evoluation (Stone & Lunt, 2013). The orbital configuration is set to 23 182 ka BP, rather than 21 ka BP as in the PMIP4 protocol. This is 2,000 years prior to the 183 ice sheet maximum extent to represent an orbit closer to point of maximum volume for 184 the North American Ice Sheet (Peltier et al., 2015; Tarasov et al., 2012b). This slightly 185 inhibits ice sheet growth as shown by a sensitivity experiment included in the supple-186 mentary information. 187

SSTs and sea ice concentrations are taken from the statistical reconstruction of Astfalck et al. (2021), combining information from the PMIP LGM multi-model ensemble (Kageyama, Harrison, et al., 2021) and compilations of proxy data (Kucera et al., 2005), and their associated uncertainties. The method is able to generate ensembles of plausible SST and SIC pairs that can be used to drive atmosphere models. The simulated SSTs are in good agreement with the proxy-based reconstruction of Paul et al. (2021) with significantly warmer tropical SSTs than the data assimilation product of Tierney et al. (2020).

The interactive ice sheet model domain is set to cover North America, Greenland, 195 and Iceland. All other ice sheets are fixed to match the Glac-1D reconstruction. The Glim-196 mer initial condition is taken from a previous ensemble of North American Ice Sheet deglacia-197 tions (L. J. Gregoire et al., 2016); specifically, ensemble member Cano3-022 at 18.2 ka 198 BP. This was chosen to represent an intermediate sized ice sheet resembling the likely 199 extent during Marine Isotope Stage 3 (Gowan et al., 2021), from which to grow the North 200 American ice sheet to an equilibrium ice sheet volume. In FAMOUS-ice, ice is able to 201 grow by flowing onto a gridcell not previously covered in ice, but ice is not able to form 202 from the accumulation of snow into an unglaciated gridcell. We thus chose an initial con-203 dition with a Cordilleran ice sheet. We also chose to start with ice already covering the 204 Hudson Bay as the Glimmer ice sheet model does not represent the complex processes 205 of grounding line migration. Simulations are run with  $10 \times ice$  sheet acceleration; i.e., 206 for every climate year simulated, Glimmer integrates for 10 ice sheet years with the same 207 surface mass balance inputs (from the climate model). After this, the new Glimmer ice 208 sheet surface elevation is passed back to the climate model, regridded and processed to 209 update its orography and ice fraction fields. 210

We first run an initial standard experiment, using the atmosphere model parameters from simulations that produce a stable contemporary Greenland ice sheet (J. M. Gregory et al., 2020) (apart from boundary conditions altered according to the PMIP4 LGM protocol), and ice sheet model parameters from previous simulations of the North American ice sheet with glimmer (L. J. Gregoire et al., 2016). Parameters that may have an impact on the ice sheet evolution are listed in Table 1.

# <sup>217</sup> 3 Collapse of the LGM ice sheet with a standard setup

Unexpectedly, with the the standard parameter values from (R. S. Smith et al., 2020), instead of growing from a mid-glacial ice sheet size, the North American ice sheet rapidly deglaciates in our standard experiment, losing half its volume in 2,500 years. This eventually results in a simulation with LGM climate conditions, but no North American Ice Sheet (Figure 1). The deglaciation is driven by ablation across the North American Ice

Parameter	Standard Value	Ensemble Range	Units	Notes
Lapse rate	6	2-10	$K \ km^{-1}$	Prescribed lapse rate for air tem- perature used to downscale FA- MOUS near-surface ice sheet climate onto surface elevation tiles. Downwelling longwave radiation is also adjusted for consistency
Daice	-0.35	-0.4-0	$K^{-1}$	Sensitivity of bare-ice albedo to surface air temperatures once the surface is in a melt regime
AVGR	0.007	0.001-0.01	$\mu m^{-1}$	Sensitivity of the snow albedo to variation in surface grain size
Fsnow	600	350-800	$kg \ m^{-3}$	The threshold in surface snow den- sity at which the FAMOUS albedo scheme switches from a scattering paradigm appropriate for a con- glomeration of snow grains to one more appropriate for a solid surface
Flow factor	3	1-10		The softness of ice. Increasing the factor makes the ice softer and more deformable
Mantle relax- ation time	3000	500-9000	yr	The relaxation time of the mantle, a lower value essentially making the mantle less viscous, thus allow- ing a quicker topographic rebound.
Basal sliding	10	0.5-20	$mm \ yr^{-1}$ $Pa^{-1}$	The basal sliding rate. A higher value allows increased ice velocity.
RHCRIT	0.85	0.6-0.9		The threshold of relative humidity for cloud formation (R. Smith, 1990).
VF1	1.882	1-2	$ms^{-1}$	The precipitating ice fall-out speed (Heymsfield, 1977).
СТ	0.000302	$5 \times 10^{-5} - 4 \times 10^{-4}$	$s^{-1}$	The conversion rate of cloud liquid water droplets to precipitation (R. Smith, 1990).
CW	0.001688	0.0001-0.002	$kg \ m^{-3}$	The threshold values of cloud liquid water for formation of pre- cipitation (R. Smith, 1990). Only the value for the land is varied.
Entrainment rate	3	1.5-6		Entrainment rate coefficient Con- vection Scales rate of mixing be- tween environmental air and con- vective plume.
Alpham	0.2152	0.2-0.65		The sea ice low albedo (Crossley & Roberts, 1995).

# Table 1. Key climate and ice sheet parameters in the simulation



Figure 1. LGM North American Ice Sheet evolution in the standard setup. a) The ice sheet volume, normalised by the initial volume (green line) compared to the PI Greenland simulations (red line). b-d) Ice Thickness at 0, 2000 and 4000 years into the run. e-g) The SMB at 0, 2000 and 4000 years into the run.

Sheet, including the ice sheet interior (Figure 1e), which is present from the start of the simulation and causes a rapid retreat of the southern margin northward through Hudson Bay (Figure 1b-d). The Greenland Ice Sheet on the other hand maintains its initial extent, which corresponds to full glacial conditions, in good agreement with observations (Simpson et al., 2009).

We know from geologic constraints (Dyke et al., 2002) that the ice sheet should be considerably larger than simulated here (i.e. it should cover the whole of Canada). We therefore conclude that parameters previously tuned to simulate the present day Greenland ice sheet well (as in (R. S. Smith et al., 2020)) are not suitable for the LGM North American Ice Sheet. To find a reasonable simulation, we thus need to fully explore the uncertainty in model input parameters controlling the surface mass balance, ice sheet dynamics and climatic conditions over the ice sheets in FAMOUS-ice as described below, in essence to "de-tune" the model and find parameter combinations that produce
 good representations of the LGM North American Ice Sheet.

<sup>237</sup> 4 The Ensemble Approach

238

# 4.1 An Initial Pass: Wave 1

We started by running a large 280-member ensemble of simulations varying param-239 eter values of Table 1. These values were sampled using k-extended Latin Hypercube sam-240 pling (Williamson, 2015) within ranges derived from previous uncertainty quantification 241 work with FAMOUS (L. Gregoire et al., 2010) and Glimmer (L. J. Gregoire et al., 2016), 242 with the addition of the Entrainment rate coefficient. We purposefully chose wide but 243 plausible ranges (Table 1) with the aim to identify a region of the parameter space that 244 would produce a reasonable North American ice extent at the LGM. We expected that 245 similar to the standard simulation (Section 3), many runs would fully deglaciate. Thus, 246 to optimise our use of computing resources, simulations were stopped at 2,500, 5,000, 247 and 10,000 ice sheet years if they lost more than 25% of the initial ice sheet area. 248

In the majority of the 280 ensemble members, the North American ice volume re-249 duces dramatically and thus all but 18 simulations were terminated early (figure 2a). How-250 ever, five simulations remain relatively stable at the initial volume and area, and ice vol-251 ume grows in twelve simulations. This deglaciation could be caused by errors in the sim-252 ulated climate. We thus compared the ensemble results to prior simulations from PMIP3 253 and PMIP4 (Figure 3). We find that global mean temperature is broadly in line with 254 previous PMIP3 and PMIP4 simulations - on the warmer end around 284K. The range 255 of global mean temperatures in the ensemble is limited by the prescribed SSTs, which 256 have a variability in the temperatures spatial distribution, but the same global mean SST 257 (within 0.5K). 258

To assess our ice sheet results, we compare them to the ice extent reconstructed 259 by Dyke et al. (2002). We calculate ice extent error as in Gregoire et al. (2016); by sum-260 ming up the number of gridcells where the ice does not match the reconstruction. The 261 maximum allowed error extent was chosen to align with the maximum LGM extent er-262 ror from the NROY ensemble of L. J. Gregoire et al. (2016). L. J. Gregoire et al. (2016) 263 applied a cumulative extent error over the whole deglaciation to identify their NROY 264 set of simulations, we translated this into a maximum bound for our LGM extent error 265 metric by applying our metric to their final NROY set and identifying the maximum value 266 obtained. Constraints on North American ice volume are not as well known as ice ex-267 tent, but can provide a useful metric for ruling out simulations. We chose to set a min-268 imum threshold of  $2.1 \times 10^7 km^3$  for ice volume as in L. J. Gregoire et al. (2016), based 269 on a variety of individual reconstructions (Clark & Tarasov, 2014; Lambeck et al., 2014; 270 Peltier et al., 2015; Tarasov et al., 2012a). Only a small subset of the simulations (18 271 out of 280) terminate the simulation within this accepted criteria for volume and extent, 272 highlighted by the green box in Figure 2b. We refer to the parameters of these simula-273 tions as Not Ruled Out Yet (NROY). 274

The average spatial extent of the ensemble is a poor fit to reconstructed ice extent 275 (Figure 2c). Very few simulations approach the southern margin. In fact, despite stop-276 ping deglaciating simulations early only around 40% of simulations cover Hudson Bay, 277 which should be in the ice sheet interior. Of the 18 simulations that do have a larger ex-278 tent, there is a variety of ice configurations, but some consistent model-data mismatch. 279 All the ice sheets that grow from the initial extent include extensive ice in Alaska, which 280 was mostly ice free at the LGM (Dyke et al., 2002). After this first wave of 280 runs, more 281 simulations were required to determine if these errors are systematic in the model set-282 up, or an artefact of the small number of simulations with a larger ice sheet volume. 283



Figure 2. Wave 1 of the ensemble. a) North American Ice volume evolution for each ensemble member. b) The final ice volume and extent error (compared to the Dyke et al. (2002) margin) for each ensemble member. c) The % of simulations with ice cover compared to the Dyke et al. (2002) margin. The red line and point on panels a and b shows the control run, as shown in Figure 1.



Figure 3. The surface temperature climatology of the Wave 1 ensemble compared to previous simulations as part of PMIP3 and PMIP4.

# 4.2 Refining the Ensemble: Wave 2

The fact that only 18 of 280 simulations in Wave 1 sustain a large enough North 285 American ice sheet indicates that the majority of the parameter space, defined in Ta-286 ble 1, produce inappropriate conditions for the maintenance of a North American Ice Sheet. 287 Searching through the parameter space to find realistic simulations thus required the use 288 of statistical emulation and an intelligent and efficient iterative ensemble design to iden-289 tify and target the space of 'reasonable' parameters. In order to provide optimal (pa-290 rameter) space coverage of the design, whilst respecting the simulations that are already 291 292 run, we use a stratified k-extended Latin Hypercube design (details follow).

Reaching an equilibrium in ice volume requires long simulations, particularly for 293 the good simulations that take 10,000 years to reach a reasonable ice volume (Figure 294 4). In order to reduce computational cost, we investigated if it was possible to guess whether 295 a simulation would be good or bad based on information from the start of the simula-296 tion. For the ice sheet to grow to a glacial maximum volume, it needs to accumulate more 297 snow than the snow and ice lost in the ablation zone. In other words, the surface mass 298 balance (SMB) at least needs to be positive. This is a conservative minimum require-299 ment; in reality an ice sheet with a positive SMB may not reach the target reconstructed 300 LGM extent, but this requirement excludes simulations that certainly cannot reach the 301 target extent. 302

To efficiently select inputs for Wave 2 we generate a candidate set of runs by first 303 identifying, from Wave 1, initial SMB values that result in plausible equilibrium ice-sheet 304 areas; and second, emulating these initial SMB values as a function of the input param-305 eters. We identified a strong relationship between the equilibrium ice-sheet area and the 306 average SMB value from the first 20 years. Denote by b the 20 year averaged SMB value, 307 and by A the "equilibrium" ice-sheet area after 10,000 ice sheet years. A predictive model 308 of equilibrium ice-sheet area takes the form  $A = f(b) + \epsilon$  where  $f(\cdot)$  may be any func-309 tion. We considered f to be either linear or a Gaussian process and found that the lin-310 ear model gave more conservative in the uncertainty estimates, which was desired as we 311 want the Wave 2 runs to bound the NROY space. Define a predictive interval P(b) =312  $[f(b) + 3\sqrt{\operatorname{var}(\epsilon)}, f(b) - 3\sqrt{\operatorname{var}(\epsilon)}]$  from our predictive model. We target equilibrium 313 ice-sheet areas in the interval  $T = [1.5 \times 10^7 km^2, 2 \times 10^7 km^2]$  and consider the b (and 314 the corresponding input parameters) such that the intersection  $P(b) \cap T$  is non-zero as 315 plausible for design in Wave 2. 316

To find combinations of input parameters that produce plausible values of b (and 317 hence equilibrium ice-sheet areas), we ran three sub-waves of 20 year-long simulations 318 to obtain the average SMB in the first 20 years. Note that simulations in the first sub-319 wave were run for 50 years to examine relationships other than 20-year SMB average that, 320 in the end, were not used. Running these shorter simulations is still somewhat costly, 321 and so we utilised a Gaussian Process (GP) emulator to design for parameterisations that 322 was likely to produce desirable values for b. Define by x the multivariate vector of pa-323 rameters that we build the emulator over: here x comprised of the 4 most influential pa-324 rameters FSNOW, AVGR, DAICE, and FLOW FACTOR (see Table 1). We model b via a stochas-325 tic GP,  $b \sim \mathcal{GP}(x)$ , where the effects of the parameters not explicitly represented in x 326 is handled by the stochasticity of the process. The specific set-up of this GP is provided 327 in code as supplementary material. 328

From our three sub-waves of 20-50 year-long simulations, we are able to extract a candidate set of simulations for the Wave 2 ensemble. The first sub-wave (Wave 1.1) samples 200 ensemble members, which are predicted from the emulator to have non-negligible probability of positive SMB. This results in around 50% of simulations in this subwave having a positive SMB, an increase from 15% in the original wave (Figure 4b, Wave 1.1). We attempt to refine the predictive bounds on the GP model twice more (Figures 4cd, Wave 1.2 and 1.3), with no improvement. This is likely due to the inherent stochas-



**Figure 4.** Ice Volumes simulated in the successive ensemble subwaves of simulations sampled to have a positive initial surface mass balance using the Gaussian Process emulator.

ticity of the climate model and cumulative effects of the parameters that we absorb into 336 the predictive error term. At the end of this process of iterative short waves, our can-337 didate set contains over 1000 20-year long simulations that have a positive SMB over the 338 North American ice sheet. From this candidate set, we select an optimal (with respect 339 to space-filling and accounting for the previous Wave 1 runs) design of 200 ensemble mem-340 bers to continue for a full 10,000 years to an equilibrium North American Ice Sheet. These 341 200 simulations make up our Wave 2. For context, this process of GP model subwaves 342 saved around 230,000 core hours (or about 2 months of real time) compared to running 343 a full second ensemble wave. 344

## 4.3 Wave 2 Results

345

By sampling a second wave of long simulation from the shorter subwaves, we de-346 sign an ensemble of simulations that produce more NROY ice sheet extents, with 120/200347 simulations maintaining or growing volume beyond the initial extent (Figure 5a). 176 348 more simulations out of the 200 wave 2 simulation are considered to be NROY based on 349 the volume and extent error thresholds previously described (Figure 5b), a factor of ten 350 increase from the 18 out of 280 simulations in Wave 1. This demonstrates the success 351 of our approach in efficiently identifying good combinations of parameters values. Some 352 simulations (56/200) shrank in volume slightly, but expended in area to meet both the 353 volume and area constraints to be considered NROY. The definition of an NROY en-354 semble member is lenient, but could be further constrained (for example, by adding a 355

target for the top of the atmosphere energy balance) depending on future research ques-tions or aims.

The mean ice sheet extent of the second wave is close to the southern ice sheet mar-358 gin (Figure 5c). However, simulations that meet this margin still show some consistent 359 model-data mismatch. All simulations with a low Laurentide Ice Sheet extent error have 360 ice that is too extensive in Alaska. This is also common in ice sheet-only simulations with 361 Glimmer (L. J. Gregoire et al., 2016; Ji et al., 2021) driven by climate forcing from FA-362 MOUS and from the higher resolution CCSM3 model, so is likely a systematic bias re-363 sulting from the climate model. Alaskan ice extent was limited by the wider ice sheet disrupting atmospheric circulations (Löfverström & Liakka, 2016; Tulenko et al., 2020). 365 The likelihood of matching observations is not helped by the climate coupling; a low res-366 olution of FAMOUS struggles to simulate the temperature and precipitation gradients 367 caused by steep topography such as the Aleutian Range (Abe-Ouchi et al., 2007). Pre-368 scribing the SMB forcing in an uncoupled model nudges the simulated ice sheets towards 369 the ice sheet prescribed (usually from reconstructions) in the climate-only model, encour-370 aging a better match between modelled and reconstructed ice geometry, but not neces-371 sarily for predictive reasons. Instead, in this case, a coupled climate-ice sheet model es-372 sentially introduces additional freedom into the simulations to produce a greater vari-373 ety of ice sheets. 374

At this point, we could repeat the subwave emulation step from the first wave of simulations in an attempt to further narrow the parameter range and produce a third wave of simulations with a greater proportion of NROY simulations. However, this second wave has already produced 176 more NROY simulations, and running large waves of simulations is computationally costly. We deem that there would be diminishing benefits to running subsequent waves of simulations beyond this point.

## <sup>381</sup> 5 Importance of Albedo Values

While 13 parameters are varied in the ensemble, it is clear that only three to four of these parameters explain the majority of the variation in the model outputs. This is demonstrated in Figure 6, where the parameter ranges for Wave 1 and Wave 2 are compared. In Wave 2, the statistical emulator has removed regions of the parameter space that do not produce reasonable ice sheet extents. There are large changes to the ranges of FSNOW, DAICE, and AVGR, in the NROY ensemble, but limited changes for the other ten parameters.

The parameters that are most influential on the simulated ice sheet volume all con-389 trol the ice sheet surface albedo, these are the snow-ice density threshold (FSNOW), the 390 maximum albedo of bare ice (DAICE), and the snow grain size (AVGR) (Table 1, see 391 Smith et al 2020 for full details of how these terms are used in their respective param-392 eterisations). As expected, the largest ice volumes result from having more reflective snow 393 and bare ice (from AVGR and DAICE respectively), and a high density threshold (FS-394 NOW) to start considering snow to be ice (keeping a surface classed as more reflective 395 snow, rather than ice, for longer). The average surface albedo of the NROY simulations 396 is 0.1 higher than from the ruled-out simulations (Figure 5). 397

398

# 5.1 Why does the ice sheet deglaciate at low albedos?

We have shown that albedo parameters that have been manually tuned to produce a reasonable contemporary Greenland ice sheet (J. M. Gregory et al., 2020; R. S. Smith et al., 2020) produce a collapsed North American Ice Sheet at the LGM. Here, we explore the reasons why the present day Greenland ice sheet may not be a sufficient target for calibrating a coupled climate-ice sheet model. The contrasting behaviour between present day Greenland and North American LGM ice sheets is associated with differ-



Figure 5. Wave 2 of the ensemble. a) Ice volume evolution for each ensemble member. b) The final ice volume and extent error (compared to the Dyke et al. (2002) margin) for each ensemble member. The % of simulations with ice cover shown as shades of blue compared to the Dyke et al. (2002) margin plotted as the red contour for the whole Wave 2 ensemble (c), and just for the NROY members of the Wave 2 ensemble (d).



Figure 6. The resulting initial 20-year SMB and equillibrium Ice Volumes as a function of normalised albedo and flow factor parameter combinations in the ensembles. Wave 2 ensemble members are circled in black, and are clustered in the parameter space as a result of the GP emulator aiming for a positive SMB. For reference, the control simulation is marked by an orange star.



Figure 7. Average Ice Sheet surface albedos for NROY and Ruled Out ensemble members across both Wave 1 and Wave 2. a) The distribution of ice surface albedos for NROY and Ruled Out ensemble members. b) The relationship between final ice sheet volume and the ice sheet surface albedo. c) The relationship between the 30-year ice sheet SMB and the ice sheet surface albedo. All ice sheet surface albedos are taken as a 30-year average of the July albedo at the start of a simulations, so that ice sheet extents are comparable between ensemble members.







Figure 8. July cloud cover and downward shortwave radiation for the PI simulation, Wave 1, and Wave 2.

ences in the magnitude of downward shortwave radiation at surface and cloud cover between the two ice sheets. In the modern Greenland simulations, the ablation zone at the ice sheet margin is often covered by thick clouds (Figure 8).

Simply, the relationship between cloud and ice sheet surface albedo can be split into 408 three groups (Figure 9). At high surface elevations there is lower temperatures and cloud 409 cover; a "sunny cold" regime that can sustain an ice sheet. At lower elevations there can 410 be sufficient cloud cover to reduce downward shortwave radiation, and therefore limit 411 ice sheet melt. In this "cloudy warm" regime the cloud cover means the ice sheet is in-412 sensitive to its surface albedo. Finally, at lower elevations and limited cloud cover we can 413 see a "sunny warm" regime, where the ice sheet is highly sensitive to surface albedo pa-414 rameters. Figure 9 shows downward shortwave radiation and ice surface albedo values 415 for each ice sheet surface gridbox in the control runs for PI Greenland and LGM North 416 America. While both runs occupy the "sunny cold" and "cloudy warm" regime, the "sunny 417 warm" regime is dominated by the LGM run. In this manner, the North American ice 418 sheet becomes very sensitive to low ice albedo parameters. 419



**Figure 9.** Comparison of energy balance between LGM North American and modern Greenland ice sheets: a) Downward Shortwave radiation vs ice surface albedo in each ice sheet gridcell for Greenland in the "standard" pre-industrial simulation (orange) and for the North American ice sheet in the "standard" LGM simulation (blue). Shaded ovals show the proposed surface albedo and radiation behaviour groups for the Greenland and Laurentide Ice Sheets. b) A schematic of the different radiative effect of clouds on the PI Greenland and LGM Laurentide Ice Sheets.

We further examine outputs from PMIP3 models (Brady et al., 2013; Voldoire et 420 al., 2013; Ullman et al., 2014; Sueyoshi et al., 2013; Adloff et al., 2018) to verify whether 421 the very strong downward shortwave radiation over the North American ice sheet ob-422 served in FAMOUS-ice is a common feature among other models. The area-averaged val-423 ues over the North American ice sheet are summarised in Table 2. The results show that 424 other PMIP models simulate stronger downward shortwave radiation at the southern mar-425 gin of the ice sheet, and that FAMOUS-ice shows the smallest value for this together with 426 CCSM4 (Table 2). In other words, other models may impose an even stronger melt on 427 the ice sheet than in our simulations. This is associated with the strongest cloud radia-428 tive effect in shortwave radiation in FAMOUS-ice and CCSM4. These results show that 429 the North American ice sheet would also be sensitive to low albedo values in other PMIP 430 models; if they allow a very small minimum bare ice albedo, a very large amount of so-431 lar energy would be absorbed. Therefore, the larger sensitivity of the LGM North Amer-432 ican ice sheet to low albedo parameters is likely not a unique feature of FAMOUS-ice, 433 but seems to be a common feature among other climate models. 434

There is a large uncertainty and variety in minimum ice sheet surface albedo among 435 PMIP models ranging from 0.2 in MRI and FAMOUS-ice to 0.7 in MPI (Alder & Hostetler, 436 2019). These differences in albedo values are induced by the combined effects of discrep-437 ancies in the physics of the albedo scheme and biases in AGCMs. For the latter, biases 438 in cloud radiative effects (R. S. Smith et al., 2020) and horizontal resolution (Kapsch et 439 al., 2021) can affect the choice of albedo values. Importantly, the albedo values selected 440 are often strongly tuned to reproduce the modern SMB. This is sensible when the fo-441 cus is on future change of Greenland ice sheet in the next few decades, since changes in 442 SMB are the dominant driver on this time scale. However, ice sheets evolve (e.g. in re-443 sponse to climate) over much longer timeframes, and on a longer time-scale, when the 444

PMIP3 Model	Surface albedo	Sfc down s/w radiation	Absorbed s/w at sfc
FAMOUS-ice	$0.68 \ (0.56 - 0.75)$	296 (265-324)	96 (71-126)
CCSM4	0.70	295.18	87.40
CNRM	0.34	307.67	204.47
GISS	0.66	341.52	116.34
IPSL	0.75	343.80	86.58
MIROC	0.75	351.76	86.51
MPI	0.83	348.92	58.07
MRI	0.57	318.19	138.25

**Table 2.** Albedo and radiative characteristics of PMIP3 models and FAMOUS-ice. FAMOUS-ice values are taken from the Wave 2 ensemble, showing the mean value and ensemble range (in brackets).

ice sheet is subject to larger instabilities and more pronounced climate interactions, tuning the albedo parameters may cause an unrealistic relationship between surface albedo
and cloud cover. Moreover, changes in the clouds over the next century could have pronounced effects on SMB with unrealistic surface albedo.

In the case of FAMOUS-ice, the surface albedo parameters used for the contem-449 porary Greenland ice sheet were originally tuned to a low value to compensate for an 450 excessive reflection of shortwave radiation by clouds (R. S. Smith et al., 2020). This was 451 performed to better simulate a stable and realistic Greenland ice sheet geometry under 452 modern day climate. However, the resulting ice albedo parameter sets are too low to pro-453 duce a realistic North American ice sheet at the LGM due to the different cloud cover 454 and downward shortwave radiation over the ablation zone. Therefore we are in an un-455 desirable situation where a model with good SMB in modern climate is unable to sim-456 ulate the LGM North American Ice Sheets, and could also be unable to simulate other 457 ice sheets and time periods. This result suggests that overtuning the albedo to compen-458 sate for biases in other components under modern climate may cause degraded simula-459 tion results under different climate states, when the cloud properties and downward short-460 wave radiation over ablation zones are different from modern. Most concerning, this sug-461 gests that simulations projecting the future of the Greenland ice sheet are particularly 462 vulnerable to uncertain future cloudiness over the ice sheet. 463

# **6** Conclusions

We have applied a new coupled Climate-Ice Sheet model (FAMOUS-ice) to sim-465 ulate the maximum extent of the last North American Ice Sheets. The standard model 466 setup manually tuned for modern day Greenland resulted in a collapsed ice sheet at the 467 LGM. We underwent a process of "detuning" the model, running hundreds of simula-468 tions to produced a range of reasonable equilibrium ice sheets, enabling us to explore the 469 influence and importance of key uncertain parameters. Large parts of the parameter space 470 produced collapsed ice sheets at the LGM. Efficiently scanning the parameter space for 471 good input parameter combinations thus required an iteration of short simulations in-472 formed by emulation of equilibrium ice volume as a function of initial surface mass bal-473 ance. The results show that FAMOUS-ice is able to simulate the maximum extent of the 474 LGM North American Ice Sheet, with particularly good match to the southern Lauren-475 tice limits, though some systematic ice overgrowth remains in Alaska. 476

From our results, we are able to identify that the parameters controlling ice sheet 477 surface albedo dominate the simulated variability in ice sheet geometry. Importantly, com-478 binations of albedo parameter values that produced a reasonable contemporary Green-479 land ice sheet do not necessarily produce a reasonable LGM North American Ice Sheet. 480 This is because albedo parameter can be overtuned to compensate for biases in modern 481 clouds over Greenland. The different cloud distribution over the Southern Laurentide 482 ice sheet at the Last Glacial Maximum provide a useful "stress test" for coupled climate-483 ice sheet models. This highlights the potential problems of relying solely on contempo-484 rary observations for model tuning. Efforts to find a region of the parameter space that 485 produce reasonable simulations of contemporary and glacial ice sheets are important, and 486 could lead to improved confidence in future ice sheet projections. 187

# 488 Author Contributions

Niall Gandy (principle researcher - climate ice sheet modelling): Methodology, Soft-489 ware, Investigation, Formal Analysis, Data Curation, Writing - Original Draft, Visual-490 ization. Lachlan Astfalck (co-researcher - statistics): Methodology, Software, Investiga-491 tion, Formal Analysis, Writing - Original Draft. Lauren Gregoire (principle investiga-492 tor): Conceptualization, Methodology, Writing - Review and Editing, Supervision, Project 493 Administration, Funding Acquisition. Ruza Ivanovic (co-investigator - climate modelling): 494 Conceptualization, Writing - Review and Editing, Supervision. Violet Patterson (contributing researcher - 21 ka orbital simulation): Investigation. Sam Sherriff-Tadano (con-496 tributing researcher - albedo and cloud analysis): Investigation, Formal analysis, Writ-497 ing - Original Draft. Robin Smith (collaborator - climate-ice sheet modelling): Method-498 ology, Software, Writing - Review and Editing. Daniel Williamson (co-investigator - statis-499 tics): Methodology, Writing - Review and Editing, Supervision. Richard Rigby (software 500 scientist): Software. 501

502 Data Availability

Output from a Wave 2 simulation is archived online (10.17632/8kswwpnjyz.1). All the other simulations can be accessed by contacting the author.

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