An Automatic Differentiation Method for Surface Carbon Flux Inversion

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Abstract

We attempt, for the first time, to estimate the surface carbon flux by the method of automatic differentiation. The atmospheric transport model is developed using a deep learning framework and is validated against standard approaches. Depends on the built-in automatic differentiation feature of the deep learning framework, the system derivatives/gradients are readily available without any extra effort. We then formulate the surface carbon flux estimation as an inverse problem using the variational approach, driven by back-propagated objective gradients. The feasibility of the automatic differentiation method is demonstrated in identical-twin observing system simulation experiments (OSSEs). The proposed framework shows favorable accuracy and great efficiency in both fully and partially observable scenarios. The present study establishes a link between machine learning frameworks and general data assimilation or inverse modeling problems, and the promising results encourage more investigations in incorporating machine learning techniques in inverse carbon modeling.

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5 Key Points:

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6	• A fully-differentiable atmosphere transport model is developed, which encoded the
7	PDE into a deep learning computational graph.
8	• Using the framework's automatic differentiation capabilities, the carbon inversion
9	problem can be solved.
10	• The proposed framework has been tested in identical-twin OSSEs and has achieved
11	favorable performance.

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12 Abstract

We attempt, for the first time, to estimate the surface CO_2 flux by the method of auto-13 matic differentiation. The atmospheric transport model is developed using a deep learning 14 framework and is validated against standard approaches. Depends on the built-in automatic 15 differentiation feature of the deep learning framework, the system derivatives/gradients are 16 readily available without any extra effort. We then formulate the surface carbon flux es-17 timation as an inverse problem using the variational approach, driven by back-propagated 18 objective gradients. The feasibility of the automatic differentiation method is demonstrated 19 in identical-twin observing system simulation experiments (OSSEs). The proposed frame-20 work shows favorable accuracy and great efficiency in both fully and partially observable 21 scenarios. The present study establishes a link between machine learning frameworks and 22 general data assimilation or inverse modeling problems, and the promising results encourage 23 more investigations in incorporating machine learning techniques in inverse carbon model-24 ing. 25

²⁶ Plain Language Summary

In this paper, we use the deep learning framework for the first time to solve the carbon 27 inversion problem in atmospheric transport, and verify the feasibility of the framework on 28 the inversion problem. In the scheme, The atmospheric transport model pyATM is imple-29 mented using a deep learning framework, and the performance of the model is quantitatively 30 compared with the classical transport model GEOS-Chem. Relying on the built-in auto-31 matic differentiation function of the deep learning framework, the gradient of the error to 32 the input variable can be calculated, that is, in the inversion problem, the gradient of the 33 observed error to the surface carbon flux is observed, and the carbon flux value is updated. 34 The proposed framework is demonstrated in an identical twin-observing system simulation 35 experiment (OSSE), showing good accuracy and high efficiency in both fully and partially 36 observable situations. 37

³⁸ 1 Introduction

Carbon dioxide (CO_2) is the most important long-lived greenhouse gas, which has 39 caused climate issues such as global warming and ocean acidification (Solomon et al., 2009; 40 Collins et al., 2013). Since the beginning of the Industrial Revolution (1750s), the at-41 mospheric CO_2 concentration has increased dramatically due to anthropogenic emissions, 42 mainly from fossil fuel burning. Approximately half of the emitted CO_2 was absorbed by 43 the ocean and land through the natural carbon cycle, while the rest remains in the at-44 mosphere and keeps pushing the atmospheric CO_2 level to new highs (Quéré et al., 2009; 45 Friedlingstein et al., 2019; Blunden & Arndt, 2020). Moreover, the absorbing capacity of 46 the oceanic and terrestrial ecosystems is not guaranteed to follow the fast growth of anthro-47 pogenic emissions or even remain at its present level as global warming tends to slow down 48 the carbon uptake of the land and ocean (Fung et al., 2005; Quéré et al., 2007). On the 49 other hand, the effectiveness of the measures that have been taken to reduce CO_2 emissions 50 is hard to evaluate without a good understanding of the variations of the natural carbon 51 sources/sinks. For more accurate projections of future CO₂ levels and the associated climate 52 impacts, it is therefore essential to accurately quantify the temporal and spatial pattern of 53 surface CO_2 fluxes and to better understand the natural carbon cycle. (Baker et al., 2006, 54 2010)55

The surface carbon fluxes are often estimated by either the bottom-up or the topdown approach. The former uses a combination of human activity data and inventoryspecific emission factors, process-based models, or remote-sensing approaches to estimate the fluxes (Poulter et al., 2022). The latter formulates it as an inverse problem, in which statistical methods or optimization approaches are used to invert the CO₂ transport process to find the optimal surface fluxes that best match the observed atmospheric CO₂ concentration (Thompson et al., 2022). The link between carbon dioxide concentration and surface
carbon fluxes is established through atmospheric transport models (Community, 2021; Grell
et al., 2005; Emmons et al., 2010).

The current study focuses on top-down techniques, with ensemble Kalman filters (EnKFs) (Peters 65 et al., 2005, 2007; Feng et al., 2009; Kang et al., 2011) and variational methods (Baker et 66 al., 2006, 2010; Chevallier, Engelen, et al., 2009; Chevallier, Maksyutov, et al., 2009) being 67 the most common ones. Ensemble Kalman filters use a Bayesian framework to estimate the 68 statistically optimal state given the prior or background predictions and the observations. EnKFs only rely on the model outputs and the observations, and thus are more flexible 70 and developer-friendly. However, a reasonable ensemble size is typically required to cap-71 ture accurate statistics, which requires considerable computational resources. On the other 72 hand, the 4-dimensional variational (4D-Var) approach minimizes the differences between 73 the model forecast trajectory and observations distributed over a period of time. These two 74 approaches can be made in an equivalent Bayesian form if the background errors and the 75 observational errors are both Gaussian, though the 4D-Var tends to give better results if 76 the assimilation time is short while the EnKF needs some spinup time. However, an adjoint model of the forward transport process is required in 4D-Var to back-propagate the gradi-78 ents, which is often a considerable development endeavor (Henze et al., 2007). This paper 79 aims to reduce the development effort by developing the atmospheric transport model within 80 deep learning frameworks, in which the adjoint model can be easily obtained by automatic 81 differentiation. 82

In terms of observations, both in situ observations and satellite retrieved data have been 83 used to study the carbon cycle and estimate the surface carbon fluxes. The in situ data such 84 as Total Column Carbon Observing Network (TCCON) (Toon et al., 2009) and Observation 85 Package (ObsPack) (Masarie et al., 2014) often provide pretty accurate observations and are 86 more sensitive to surface emissions. However, these observations are extremely sparse spa-87 tially, which poses difficulties to the global flux estimates. On the other hand, satellite-based 88 measurements, such as those of Greenhouse Gases Observing Satellite (GOSAT) (Yokota et al., 2009a) and the Orbiting Carbon Observatory-2 (OCO-2) (Crisp et al., 2004b), often have better spatial coverage over the globe, while they are less accurate and are sparse in 91 both space and time. In the present study, we focus on algorithmic development and evalu-92 ating the new methodology. Therefore, we use the identical twin OSSEs, in which we know 93 the "truth", to test the proposed framework, eliminating possible uncertainties associated 94 with the model errors and the observational errors. However, we do mimic the temporal 95 and spatial pattern of the satellite data to ensure its applicability in real-world systems. 96

Despite its extraordinary success achieved in computer vision and natural language pro-97 cessing, deep learning has not received much attention in computational physical modeling until recently (Zhu & Zabaras, 2018; Ravuri et al., 2021; Espeholt et al., 2021; Karniadakis 99 et al., 2021; Chen & Fidkowski, 2021; Hu et al., 2022; Pathak et al., 2022). A few attempts 100 have been devoted to data assimilation or inverse modeling. Laloyaux et al. (2022) intro-101 duced a neural network model into the 4d-Var to correct for the bias that accumulates along 102 the model trajectory. Frerix et al. (2021) learned an inverse observation operator to map 103 the observational data to physical states and then reformulated the objective function in the 104 better-behaved physics space instead of the observation space. Fablet et al. (2021) adopted 105 an end-to-end learning approach, which embedded the joint training of the forward dynam-106 ical process and the inverse assimilation in a physics-informed approach. Encoder-decoder 107 type architectures have also been used in data assimilation to reduce the problem dimen-108 sion and computational cost (Mack et al., 2020; Peyron et al., 2021; Amendola et al., 2021). 109 Nevertheless, these studies treat the machine learning models mostly as black boxes and 110 the generalization is largely limited. In this paper, we directly embedded the atmospheric 111 transport model within a deep learning framework, *i.e.*, the physics is numerically "exact" 112 instead of data-driven, which generalizes well to different initial and boundary conditions. 113 The main contribution of this work can be summarized as follows: 114

- We developed a fully-differentiable atmosphere transport model called pyATM, which encoded the partial differential equations (PDEs) solver into a computational graph based on a deep learning framework.
- Using the framework's automatic differentiation capabilities, the derivatives of the loss function with respect to the learnable parameters are obtained. When the learnable parameter is the carbon flux and the loss the function is the observation error, the carbon flux can be corrected by observations.
- The proposed framework has been tested in identical-twin OSSEs and has achieved
 favorable performance in both fully-observed and partially-observed systems with
 simulated satellite observations.

The remainder of this paper proceeds as follows. We describe in Section 2 the detailed implementations of the proposed flux inversion framework, including the forward atmospheric transport model and its backward mode, as well as the variational inversion approach. Section 3 presents the observing system simulation experiment (OSSE) and provides a comprehensive summary of the problem setup. The primary results are shown in Section 4, and Section 5 concludes the present work and discusses potential future work.

131 2 Method

Deep learning has made extraordinary achievements in the field of artificial intelligence. However, its application in numerical computation has not been studied until recently. Here, we use the deep learning framework to encode the numerical solving process of partial differential equations as computational graphs. First, we use the generalized hyperbolic parabolic differential equations to formulate the inverse problem, including the forward dynamic modeling and the corresponding derivative calculations. Then, the pyATM atmospheric transport model built with this framework is introduced, followed by its structural design and the adopted inversion algorithms.

¹⁴⁰ 2.1 Implementation of PDE

Consider a PDE on $[0, T] \times D$ with the following form:

$$\partial_t u = \mu(u, \partial_1 u, \cdots, \partial_d u) + f, \tag{1}$$

$$u(0,x) = u_0(x),$$
 (2)

where $t \in [0, T]$, $x \in D \subset \mathbb{R}^d$. u(x, t) is the state variable of the PDE, $\mu(u, \partial_1 u, \dots, \partial_d u)$ denotes the state-dependent spatial derivative term, f is the external forcing term, and u_0 represents the initial condition. The external force term f can be independent of state variables. For this set of equations, we can obtain the transition relation of the corresponding state variables between two adjacent moments t_i and t_{i+1} by the integration in time,

$$u_{i+1} = u_i + \int_{t_i}^{t_{i+1}} [\mu(u_i, \partial_1 u_i, \cdots, \partial_d u_i) + f_i] dt.$$
(3)

The dynamical process of the partial differential equation is implemented using a deep learning framework within a computational graph. The proposed architecture consists of three parts: the spatial derivative calculations, the construction of the right-hand side (RHS), and the time integration. All orders of partial derivatives are calculated first, which are then added up with the source term f to complete the RHS. In time integration, the RHS is integrated forward in time to update the system state. The spatial discretization of partial differential equations adopts the finite difference method, mainly due to its simple matrix operation form that can be favorably implemented by tensor operations in the deep learning framework. In the present study, we adopted second-order differencing for the transport problem, in which either upwind schemes or central schemes with filtering are used to stabilize the system. The spatial derivative degenerates to first-order forms at the boundaries corresponding to the north and south poles. With regards to time integration, explicit Runge-Kuta schemes are used.

¹⁵⁹ 2.2 Inverse Calculation

In traditional numerical computing, the forward PDE model is usually programmed by 160 languages such as Fortran or C for speed and efficiency. However, the reverse adjoint model 161 often requires dedicated development efforts. On the other hand, if the forward PDE is en-162 coded within a deep learning framework such as torch (Paszke et al., 2019), tensorflow (Paszke 163 et al., 2019), and JAX (Bradbury et al., 2018), it automatically stores operational dependen-164 cies and intermediate results, which then enables flexible and convenient gradient compu-165 tation between variables. Moreover, those deep learning frameworks provide user-friendly 166 interfaces for interpreted languages while they rely on backends that are written in com-167 piled languages for intensive computations to compensate for efficiency loss. Therefore, deep learning frameworks have great potential in physical modeling and are especially suitable 169 for inverse problems. 170

For inverse problems, the loss function, as shown in Eqn. 4, is usually defined as the 171 difference between the projected state variables in the observation space $\mathcal{H}(u_i)$ and observa-172 tion values o_i . Therefore, the solution process of the inverse problem is to iteratively update 173 the learnable variables such that the loss function decreases along the gradient descending 174 direction. If the learnable variable is the external force term f in Eqn. 1, then the update of 175 f depends on the gradients of the loss with respect to the external force terms as in Eqn. 5. 176 In the expression, the external force f_i applied at time i will have an impact on the states 177 of all subsequent moments. As a result, the gradient calculation also includes the gradient 178 of all the subsequent times. 179

$$\mathcal{L} = \sum_{i=1}^{T} \|\mathcal{H}(u_i) - o_i\|$$
(4)

$$\frac{\partial \mathcal{L}}{\partial f_i} = \sum_{j=i+1}^T \frac{\partial \mathcal{L}_j}{\partial f_i} = \sum_{j=i+1}^T \frac{\partial \|\mathcal{H}(u_j) - o_j\|}{\partial f_i}$$
(5)

In the forward time integration, the state change at a certain moment will affect all the states along the entire time trajectory. The perturbation of the state at each moment will gradually expand its influence range as time progresses. Conversely, the error function at some future moment will contribute to the state derivatives at all previous moments. According to the chain rule, the derivative of the error (loss) function at moment j with respect to the flux at moment i can be written as:

$$\frac{\partial \mathcal{L}_j}{\partial f_i} = \frac{\partial \|\mathcal{H}(u_j) - o_j\|}{\partial f_i} = \frac{\partial \|\mathcal{H}(u_j) - o_j\|}{\partial u_j} \frac{\partial u_j}{\partial u_{j-1}} \cdots \frac{\partial u_{i+2}}{\partial u_{i+1}} \frac{\partial u_{i+1}}{\partial f_i}$$
(6)

The above chain expression contains the product of several terms, including the derivative of the loss function with respect to the last state $\partial \|\mathcal{H}(u_j) - o_j\|/\partial u_j$, a series of derivatives of the next state with respect to the previous state $(\partial u_j/\partial u_{j-1}) \cdots (\partial u_{i+2}/\partial u_{i+1})$, and finally the derivative of the state with respect to the flux $\partial u_{i+1}/\partial f_i$. The derivative of the state variable at the next moment with respect to the state variable at the previous moment $\partial u_{i+1}/\partial u_i$, can be obtained from Eqn. 3 by taking the derivatives on both sides,

$$\frac{\partial u_{i+1}}{\partial u_i} = 1 + \frac{\partial}{\partial u_i} \int_{t_i}^{t_{i+1}} \mu(u) dt.$$
(7)

Similarly, the derivative of the current states with respect to the external force term at the previous moment can also be derived from Eqn. 3,

$$\frac{\partial u_{i+1}}{\partial f_i} = \int_{t_i}^{t_{i+1}} 1dt = t_{i+1} - t_i.$$
(8)

Eqn. 7 and Eqn. 8 assume that the external force term is independent of the state variables, which is true for the carbon flux inversion problems that are of particular interest in this paper.

The gradient of the loss function in Eqn. 5 can be easily computed through automatic differentiation in our proposed method. By using a deep learning framework, chain rules can be automatically applied by backpropagation of the computation graph, without the need to explicitly compute all derivatives individually.

¹⁹³ 2.3 Flux Inversion Framework

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The atmospheric transport model used in the inversion problem takes a special form of the hyperbolic parabolic differential equation in Eqn. 1, where only the first order derivatives (transport) are present in the system,

$$\frac{\partial \rho c}{\partial t} = -\nabla (\rho c \cdot V) + s. \tag{9}$$

The equation mainly describes the evolution of the carbon concentration field c under the transport effect of meteorological velocity field V, and the source/sink effects of the surface flux s. ρ is the density of the air which is constant by the incompressible flow assumption. In this problem, the meteorological velocity field V is considered a known quantity, which can be obtained through historical reanalysis data. In Eqn. 9, we do not consider the atmospheric chemistry-related terms. This set of transport equations applies to relatively long-lived gases in the atmospheric circulation such as CO2, since their chemical changes are small and thus negligible relative to transport effects.

Here, we use the pytorch framework to implement the atmospheric transport model, 203 denoted as pyATM. The numerical implementation of pyATM mainly refers to GEOS-Chem (Community, 2021), which is a numerical model commonly adopted by researchers 205 in a wide range of atmospheric chemistry problems. In our ATM implementation, the 206 influence of the atmospheric boundary layer is taken into account. In the atmospheric 207 boundary layer, the gas concentration field is usually mixed in a relatively short period 208 of time. The TURBDAY mixing scheme is the default atmospheric boundary layer (ABL) 209 mixing scheme in GEOS-Chem. In the TURBDAY scheme, it assumes instantaneous vertical 210 mixing from the surface through the mixing depth (Bey et al., 2001). In our pyATM, the 211 same implementation is also adopted, which assumes that the gas concentrations within 212 the boundary layer are instantaneously mixed. The boundary layer thickness data are 213 obtained from meteorological reanalysis data. We also adopt some numerical tricks employed 214 in GEOS-Chem to help stabilize the system within the polar regions. Specifically, the 215 concentration fields are regionally averaged in the two layers that are closest to the polar 216 boundaries. 217

Using the atmospheric transport model implemented within pyATM, the computational graph has the following two characteristics. First, the depth of the computational graph is proportional to the length of the time integral. The whole process consists of many



Figure 1. Schematic diagram of solving the carbon inversion problem using the automatic gradient inversion framework pyATM.

computational blocks, with each block representing the numerical evolution within a specific
time step. Second, the structure of the computation graph is exactly the same among all
blocks, only the inputs are different. In each computation block, the meteorological data
and flux data corresponding to the time step is fed into the graph, while the former is fixed
and the latter is learnable.

226 2.3.2 Flux Inversion

The atmospheric transport equation can be denoted as: $c_p = \mathcal{M}(c_0, f_p)$, where the 227 model \mathcal{M} represents the full dynamics of the concentration evolution. Conditioned on the 228 initial concentration field c_0 , as well as the surface flux f_p emitted within these periods, 229 the predicted concentration vector c_p can then be obtained. The concentration in the state 230 space can be mapped to the observation space by the observation operator $\mathcal{H}(c_p)$. In the 231 flux inversion problem, the decision variables are the gas emissions f_p specified to the model. 232 Estimation of the surface flux can be obtained by minimizing the following object function 233 $\mathcal{L}(f_p)$ with respect to f_p : 234

$$\mathcal{L}(f_p) = \mathbb{E}[\mathcal{H}(c_p) - c_o)^T R^{-1} (\mathcal{H}(c_p) - c_o)] + \mathbb{E}[(f_p - f_b)^T B^{-1} (f_p - f_b)]$$
(10)

The error function contains two parts. The first part is the observation error term, and the second part is the background error term. The observation error term describes the error between the observation c_o and the simulated concentration values in the observation space $\mathcal{H}(c_p)$, weighted by the measurement uncertainty matrix R. The second part is called the background error, which describes the difference between the predicted flux f_p and background flux f_b , weighted by the background error covariance matrix B.

Based on the above description of the atmospheric transport model pyATM, the derivative calculation for the flux update, and the definition of the error function, we summarize the inversion process in the pyATM framework. Figure 1 depicts the schematic diagram of solving the carbon inversion problem using automatic gradient in the pyATM framework. At each time step, the state variable and flux input are fed into the single-step ATM block, and the state at the next moment can be obtained. These time steps are concatenated to complete the entire numerical simulation. The states obtained by pyATM are then mapped
to the observation space in which the errors compared to the actual observations are computed. Finally, all errors are aggregated and contribute to the derivative of the total error
with respect to fluxes at each time step. After a complete forward state prediction and backward gradient propagation, the fluxes are updated by gradient descent methods to minimize
the loss function. The problem of particular interest in this study is carbon flux inversion,
while the proposed framework is formulated without loss of generality such that it is also
applicable to general PDE systems.

The proposed framework has several advantages compared to other inversion methods. 255 Using the automatic gradient scheme, the flux updating is similar to the variational method. 256 Both methods update the flux driven by the derivative of the error with respect to the flux 257 input. However, compared with traditional variational methods that rely on the adjoint 258 model, the proposed framework avoids complicated adjoint development, which is often a 259 dedicated endeavor. Moreover, the flexibility of the variational method is usually limited. 260 Once the code implementation of the forward atmospheric transport has changed, the cor-261 responding adjoint model also needs to be updated accordingly. On the other hand, deep learning frameworks provide automatic differentiation of all tensor operations. Once the 263 forward computational graph is built, it can automatically obtain the gradient relationship 264 between any two connected tensors, and the dependencies get automatically updated if the 265 forward computational graph changes. Using automatic differentiation, the computational 266 cost of the inversion algorithm is independent of the input parameter dimensions, making 267 it efficient in high-dimensional inversion problems. 268

²⁶⁹ 3 Experimental Setup

In the experimental part, this paper mainly discusses two aspects: one is to verify the effectiveness of the pyATM atmospheric transport model, and the other is to demonstrate the automatic differentiation method for solving the inversion problem through twin experiments.

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3.1 Numerical Setting

The implementation of the pyATM forward process mainly refers to the numerical 275 schemes in GEOS-Chem, such that numerical settings and the input meteorological conditions are the same as those in GEOS-Chem. The atmospheric transport models operate on a global scale, which is discretized evenly in both the longitude and latitude directions with 278 rectangular meshes. The spatial resolution is 5° (longitude) $\times 4^{\circ}$ (latitude), corresponding 279 to a horizontal mesh grid of 72×46 cells. In the vertical direction, 47 pressure layers are used 280 as those in GEOS-Chem. The time step Δt for the numerical simulation is 30 minutes. The 281 MERRA2(ref) reanalysis data are used for the meteorological fields, providing the velocity 282 vector data as well as the boundary layer thickness information. 283

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3.2 Identical-twin Observing System Simulation Experiments

In the experiments, we adopt the identical-twin observing system simulation exper-285 iments to examine the effectiveness of the inversion algorithm. The concept of twin ex-286 periments is very useful for testing the effectiveness of inversion algorithms as the "true" 287 solution is known. The schematic diagram of the identical-twin inversion problem using 288 pyATM is shown in Figure 2. In the identical-twin experimental setup, "true" observa-289 tions are obtained by running a forward numerical process with known initial conditions and "true" fluxes. In the inversion stage, the initial concentration field is fixed, and only 291 the surface flux parameters are zero initialized and then iteratively updated. Through such 292 an experiment, the influence of other interference factors can be excluded, and the back-203 ward relationship from the observation error to the flux can be reflected more clearly. The 294



Figure 2. Schematic diagram of the identical-twin observing system simulation experiments using pyATM. In the twin experimental setup, CO2 observations are obtained by running a forward numerical procedure with known initial conditions and "true" fluxes.

identical-twin problem has a "true" flux that can provide a quantitative evaluation of the
correctness and accuracy of the inversion algorithm. At the same time, the identical-twin
experiment makes it more convenient to perform some operations on the observation data,
such as adding perturbations, using sparse sampling observations, etc.

The observable variable is XCO_2 , the column-weighted average CO_2 concentration, 299 which is typically available in the satellite retrieved data. Two sets of validation under 300 different observation coverage are used, namely the full observation inversion and the partial 301 observation one. With regards to the partial observation case, the time stamps and spatial 302 coordinates of the observation points are generated based on the sampling pattern of carbon 303 satellites. In the partial observation scenarios, the inversion problem is usually ill-posed due 304 to the sparsity of valid observations, as well as the dissipative nature of the atmospheric transport. The problem of inversion underdetermination caused by sparse observations can be alleviated by making more observations or by reducing the degrees of freedom of the flux 307 parameters. 308

To verify the validity of the proposed inversion framework, we do not consider the 309 background flux error term that is commonly adopted, and thus only the first term on the right-hand side of Eqn. 10 is included in the loss function. By ignoring the prior information 311 of the background flux, the inverse problem is made more challenging as less information is 312 used. However, excluding prior information also helps to examine more clearly the carbon 313 flux updates that directly correspond to observational errors, while eliminating interferences 314 from background terms. Therefore, carbon fluxes are updated from zero initialization, and 315 no background term is added to the loss function. Here we assume that all observations are 316 uncorrelated since all sensors observe independently. In such a case, a unit diagonal matrix 317 is used as the error covariance matrix R. During the iterative process, we use the Adam optimizer with a learning rate of 1e-5. The inversion process is stopped after 300 iterative 319 steps. 320

321 4 Results

322 4.1 ATM validation

The atmospheric transport equation encoded in pyATM is implemented with reference to GEOS-Chem. To verify its effectiveness, we quantitatively verify the pyATM results against the GEOS-Chem simulations. In the experiment, the time spans over 7 days, from July 1^{st} 2019 T 00:00 UTC to July 8^{st} 2019 T 00:00 UTC. We use the same initial concen-

Table 1.MAE error under different pyATM numerical schemes, compared with the correspondingGEOS-Chem numerical simulation in the range of 7 days.

version	Spatial scheme	Vertical Transport	PBL Mixiing	XCO2 MAE(ppm)	bottom CO2 MAE(ppm)
V01	central	off	off	0.304	1.044
V02	upwind	off	off	0.229	0.949
V03	upwind	on	off	0.181	0.949
V04	upwind	on	on	0.111	0.192



Figure 3. MAE error under different pyATM numerical schemes, compared with the corresponding GEOS-Chem numerical simulation in the range of 7 days. (a) bottom CO2 MAE(ppm). (b) XCO2 MAE(ppm).

tration field and meteorological data to quantitatively compare the results of the forward 327 integrations between pyATM and GEOS-Chem. In order to examine the implementation of 328 the atmospheric transport model, the carbon flux is not included in the simulation process. 329 To obtain numerical results that are closer to GEOS-Chem, we conducted a series of explo-330 rations of numerical schemes. On the spatial discretization, we tried both the second-order 331 central scheme and the upwind scheme. The effect of adding vertical convection, as well as 332 the effect of boundary layer mixing, are also within the scope of the numerical investigation. 333 Table 1 lists different versions of the pyATM and the corresponding mean absolute error 334 (MAE) within a 7-day numerical simulation. The V01 version adopts the central difference 335 scheme with filtering while ignoring both vertical convection and boundary layer mixing. The V02 version uses the second-order upwind scheme to increase the numerical stability. 337 Based on version V02, V03 adds the convection effect in the vertical direction. Furthermore, 338 version V04 adds the implementation of boundary layer mixing compared with version V03. 339 The daily error curves for the four versions within the 7-day range are illustrated in Fig-340 ure 3. In the quantitative results, we present the bottom CO_2 and XCO_2 errors, which 341 are most relevant to the actual observations. For observation stations located close to the 342 ground, only CO_2 concentration at the bottom of the atmosphere can be obtained. On the 343 other hand, only the column weighted concentration XCO₂ can be retrieved from satellite 344 observations. 345

Considering the spatial discretization scheme, the upwind scheme outperforms the central one, reducing the XCO₂ error from 0.304 parts per million (ppm) to 0.229 ppm. The upwind scheme achieves better numerical stability compared to the central one since it respects the convective nature of the system. The V03 version takes into account convection effects in the vertical direction, making the momentum transfer three-dimensional. Since the bottom layer is not affected by the 3D effect, the error of the bottom layer CO₂ will not change. However, for a full vertical level, the error of XCO₂ is significantly reduced to 0.181 ppm. If the influence of the boundary layer mixing is further considered, the error is greatly reduced as shown in Table 1 for the version V04. After considering the boundary layer mixing, the bottom CO₂ error can be reduced to 0.192 ppm, and the error of XCO2 is further reduced to 0.111 ppm, achieving similar results compared to GEOS-Chem.

The errors during the 7-day range are plotted in Figure 3 for different numerical settings. 357 V01, V02, and V03 versions all have a monotonic increasing trend, in which the error grows 358 fast at the beginning and then gradually flattens. Among them, the V01 version exhibits 350 the highest XCO_2 error around 0.4 ppm after 7 days while the V03 achieves the lowest error level around 0.2 ppm. For the V04 version, the error peaks after about 3 days, and finally 361 tends to decrease slowly. The error of the predicted XCO_2 after 7 days is about 0.12 ppm. 362 According to the carbon satellite observations, the overall XCO₂ retrieval error is about 363 0.3ppm. Therefore, the model error of pyATM (assumes the GEOS-Chem as the "true" 364 model) can be made much smaller than the observation error, indicating its capability of 365 serving as a forward model in carbon flux inversion. For the error curves of the bottom 366 CO_2 concentration, it can be seen that error of V04 flattens right after the CO_2 has been well mixed in the boundary layer while the error accumulates but gradually flattens in the other schemes. 369

Figure 4 presents the predicted results of the pyATM model on days 1, 4, and 7, compared with the results from GEOS-Chem. The XCO₂ distribution is smoothed out under the atmospheric transport and boundary layer mixing, in which both the maximum and minimum values are attenuated. In terms of spatial distribution, the characteristics of high concentration in the northern hemisphere and low concentration in the southern hemisphere are retained. Visually, the results of pyATM are very close to those of GEOS-Chem, indicating a reliable capability of pyATM to simulate carbon atmospheric transport.

4.2 Inversion Result

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For the inversion tests, we designed experiments for both the full observation and partial observation scenarios respectively, to examine the influence of observation sparsity 379 on the inversion results. Full observation means that the concentration information on any 380 latitude and longitude coordinates can be obtained at any time. However, in the vertical 381 direction, only the column weighted concentration can be obtained (XCO_2) . The generation 382 of sparse observation data mainly refers to the samples from satellite observation data. 383 This verifies the effectiveness of solving the inversion problem at the practical spatial and 384 temporal sampling frequency of on-orbit satellites. Here, we use the combined dataset of 385 the two most commonly used carbon satellites, the Japanese Greenhouse gases Observing Satellite (GOSAT) (Yokota et al., 2009b) and the NASA Orbiting Carbon Observatory 387 (OCO-2)(Crisp et al., 2004a). Figure 6 shows the spatial sampling points for the first two 388 days of July 2019. The sampling trajectory of the satellite is different every day, and only 380 a small area can be scanned in a short period of time, and the overall sampling frequency 390 is about 1%. In addition, these two satellites cannot achieve full coverage over the globe. 391 For the month of July, we are concerned about the lack of observation data coverage in the 392 high-latitude areas in the southern hemisphere, which will be further discussed later. 393

We mainly investigate the long-period inversion capability of the proposed framework. 394 Therefore, a temporally fixed while spatially varied carbon flux field is adopted at all times-395 tamps as the time-averaged flux field. Utilizing the automatic differentiation capability of 396 deep learning frameworks to compute the derivatives, the update of carbon fluxes is then 397 performed iteratively. Figure 5 shows the statistics during the iterative process, in which (a) is the MAE at observation points, and (b) presents the MAE of the inversed flux field. For 399 the MAE error of the concentration at the observation points, it decreases rapidly with the 400 progress of the iterative step, and it presents a log-linear law in the later stage. For the full 401 observation case, the XCO₂ error decreases to around 1.6×10^{-3} ppm after 300 iterations. 402 As for the partial observation one, the observation error is slightly higher, which is about 403



Figure 4. Simulation results of XCO2 for GEOS-Chem and pyATM at days 1, 4, and 7 respectively.

 2.2×10^{-3} ppm after 300 iterations. For the inversed flux error as shown in Figure 5 (b), the 404 error decreases gradually as the iteration proceeds for both the full and the partial observa-405 tion cases. However, the error is significantly smaller under the full observations, indicating high sensitivities of the XCO_2 with respect to the fluxes. After 300 iterations, the flux errors 407 of the full observation and the partial observation cases are around $2.59 \times 10^{-8} mol \cdot m^{-2} s^{-1}$ 408 and $6.10 \times 10^{-8} mol \cdot m^{-2} s^{-1}$, respectively. In the case of sparse observations, due to the 409 data sparseness in both space and time, only the flux within a small influence range of the 410 observation point can be effectively updated. As discussed earlier, the observation in the 411 high-latitude areas of the southern hemisphere is obviously insufficient, so the nearby flux 412 cannot be effectively updated. 413

The update history of the inversed flux during the optimization is shown in Figure 7 414 and Figure 8. These can further help to investigate the capability as well as the effect of the 415 inversion algorithm. Figure 7 shows the inversed fluxes after 1, 5, 25, 100, and 300 iterations, 416 along with the ground truth under the full observation scenario. After one iteration step, 417 the flux source and sink patterns can be seen. In the northern part of Eurasia and North 418 America, the flux sinks due to vegetation absorption are manifested. While in East Asia, South Asia, Central America, North America, and South America, there exhibit carbon 420 sources. In the fifth iteration step, the intensities of carbon sources and sinks are further 421 enhanced, showing a large carbon sink in Siberia, as well carbon sources corresponding to 422 several urban areas in South Asia. After 25 iteration steps, the overall flux distribution 423 is well captured, and all major carbon sources and sinks in the ground truth are reflected 424 in the inversion results, although small-scale differences still exist. Iterating from the 25^{th} 425 to the 100^{th} iteration, the major carbon sources and sinks are further refined. The peak 426 flux sources corresponding to East and South Asian urban regions are more clearly and 427 accurately depicted. From the 100^{th} iteration to the 300^{th} iteration, the update slowly 428 saturates, and the flux field at the 300^{th} iteration is very close to the ground truth. 429

For partially observed scenarios, the flux update is consistent with the full observation 430 scenario, the flux gradually converges as the optimization proceeds. However, the flux convergence in the partial observation scenario is less effective compared to that of the 432 full observation one. After 300 iterations, there are still noticeable discrepancies between 433 the inversed flux and the ground truth, which is particularly prominent in the southern 434 hemisphere with little observation coverage. The inversion results failed to recover oceanic 435 carbon sinks in the southern hemisphere. Meanwhile, the inversed flux shows abnormal 436 sources in the Antarctic region. These deficiencies are sort of expected due to inadequate 437 observations in the high-latitude areas of the southern hemisphere. As for the northern 438 hemisphere, the overall flux distribution is well captured, although the carbon sources in East and South Asia are underestimated. Moreover, more spikes in the flux distributions 440 can also be observed due to sparse observations. 441

In summary, the spatial-temporal distribution pattern and the magnitudes are better retrieved in fully-observed scenarios. However, considering the difficulties of practical data acquisition, the system is only partially observable most of the time. Despite some flaws in data-poor regions, the inversion results employing the synthesized satellite data show good agreement with the ground truth, indicating the applicability of the proposed pyATM framework in real-world applications. Furthermore, we anticipate that prior knowledge of the flux distributions can also help regularize the flux inversion in practice.

449 5 Conclusions

In this paper, we developed a fully-differentiable atmospheric transport model called pyATM, which encodes the PDE solver into computational graphs based on a deep learning framework. Using the framework's automatic differentiation capabilities, the derivatives of the loss function with respect to the learnable parameters are obtained. When the learnable parameter is the carbon flux and the loss function is the observation error, the flux can be



Figure 5. Error convergence history in the inversion tests. (a) XCO2 MAE in the observation points. (b) Inversed flux MAE.



Figure 6. Satellite observation sampling points of 2 typical days. (a) July 1^{st} 2019 (b) July 2^{nd} 2019.

455 corrected by observation. The proposed framework has been demonstrated in identical-twin
 456 OSSEs and has achieved favorable performance in both fully-observed and partially-observed
 457 systems.

Based on the proposed framework, the connection between the machine learning frame-458 work and inverse problem solving is established. Since the PDE solver is embedded in the 459 machine learning framework, popular machine learning models can be easily added to the 460 PDE solving procedure, making it possible to seamlessly integrate machine learning models 461 within PDE-constrained inverse problems in future research. Furthermore, the application 62 of automatic differentiation is not limited to carbon flux inversion problems but also ap-463 plies to inverse/optimization problems constrained on more general PDEs, such as data 464 assimilation in numerical weather predictions and shape optimization in engineering design. 465

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Figure 7. The flux update history in the full observation case. The 6 images present the inversed flux fields at the 1^{st} , 5^{th} , 25^{th} , 100^{th} , 300^{th} iteration steps, and the ground truth.



Figure 8. The flux update history in the partial observation case. The 6 images present the inversed flux fields at the 1^{st} , 5^{th} , 25^{th} , 100^{th} , 300^{th} iteration steps, and the ground truth.

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