# Restratification Structure and Processes in the Irminger Sea

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#### Abstract

The Irminger Sea is one of the few regions in the ocean where deep (> 1000 m) convection occurs. Convection is followed by restratification during summer, when the stratification of the water column is reestablished and the convectively formed water is exported to the lower limb of the Atlantic Meridional Overturning Circulation (AMOC). We investigate the interannual variability and physical drivers of restratification in the upper 600 m of the central Irminger Sea using reanalysis data for the years 1993–2019. We find that there are distinctly different restratification processes in the upper 100 m of the water column (the upper layer) and the water below it (the lower layer). In the upper layer, the stratification is dominated by a strong seasonal cycle that matches the cycle of the surface heat flux. In 2010 and 2019, there were peaks in upper layer stratification, which could be related to strong atmospheric heat and freshwater fluxes. By contrast, in the lower layer the seasonal cycle is weaker and there is strong interannual variability. Restratification can continue for up to 5 months after the surface heat flux has become negative, indicating a role for lateral advection. The strength of the restratification is strongly correlated with the eddy kinetic energy in the eastern Irminger Sea. This suggests the lateral advection is driven by warm, saline eddies from the Irminger Current. In the future, surface warming and freshening of the Irminger Sea due to anthropogenic climate change are expected to increase stratification.

# Restratification Structure and Processes in the Irminger Sea

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## Key Points:

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| 8  | • | Restratification in the Irminger Sea is distinctly different in the upper 100 m of     |
|----|---|--|
| 9  |   | the water column than below  |
| 10 | • | In the upper layer, restratification is driven by atmospheric fluxes and freshwa-      |
| 11 |   | ter export from Greenland  |
| 12 | • | In the lower layer, restratification is driven by lateral advection of warm and saline |
| 13 |   | waters through Irminger Current eddies   |

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#### 14 Abstract

The Irminger Sea is one of the few regions in the ocean where deep (> 1000 m) convec-15 tion occurs. Convection is followed by restratification during summer, when the strat-16 ification of the water column is reestablished and the convectively formed water is ex-17 ported to the lower limb of the Atlantic Meridional Overturning Circulation (AMOC). 18 We investigate the interannual variability and physical drivers of restratification in the 19 upper 600 m of the central Irminger Sea using reanalysis data for the years 1993–2019. 20 We find that there are distinctly different restratification processes in the upper 100 m 21 of the water column (the upper layer) and the water below it (the lower layer). In the 22 upper layer, the stratification is dominated by a strong seasonal cycle that matches the 23 cycle of the surface heat flux. In 2010 and 2019, there were peaks in upper layer strat-24 ification, which could be related to strong atmospheric heat and freshwater fluxes. By 25 contrast, in the lower layer the seasonal cycle is weaker and there is strong interannual 26 variability. Restratification can continue for up to 5 months after the surface heat flux 27 has become negative, indicating a role for lateral advection. The strength of the restrat-28 ification is strongly correlated with the eddy kinetic energy in the eastern Irminger Sea. 29 This suggests the lateral advection is driven by warm, saline eddies from the Irminger 30 Current. In the future, surface warming and freshening of the Irminger Sea due to an-31 thropogenic climate change are expected to increase stratification. 32

#### <sup>33</sup> Plain Language Summary

In the Irminger Sea, cold winters can cause the surface ocean to cool enough that 34 they start mixing downward. The mixing can continue for months and eventually form 35 a dense water mass with uniform properties of more than 1 km deep. This process is called 36 deep convection. During summer, the dense water gets exported to large depths, and the 37 water column will go back to its original structure with lighter waters over denser wa-38 ters. This is called restratification. The formation and export of dense waters in the Irminger 39 Sea plays an important role in the global ocean circulation. Here we study what drives 40 the restratification in the central Irminger Sea, and how it varies in time. We find that 41 in the surface layer (0-100 m) the restratification is driven by heating from the atmo-42 sphere, which warms the surface waters and makes them lighter. Additionally, rainfall 43 or ice melt which make the surface waters fresher contribute to restratification. Below 44 the surface layer, restratification happens through transport of warm and saline waters 45 from the eastern Irminger Sea. With rising ocean temperatures due to climate change, 46 the Irminger Sea might see more restratification and less deep convection in the future. 47

#### 48 1 Introduction

The Atlantic Meridional Overturning Circulation (AMOC) plays an important role 49 in the global climate system (e.g. Pérez et al., 2013; McCarthy et al., 2017; Fox-Kemper 50 et al., 2021). In the Atlantic upper ocean, warm and saline waters are transported north-51 ward in the AMOC's upper limb from the equator to high latitudes, where wintertime 52 deep convection transforms them into colder and denser waters. These dense waters can 53 spread to the AMOC's lower limb in the deep ocean, where they remain isolated from 54 the atmosphere for hundreds to thousands of years. Thus the AMOC has an important 55 effect on the ocean's uptake and storage of heat and anthropogenic  $CO_2$  from the atmo-56 sphere (Fröb et al., 2016; Brown et al., 2021). Deep convection can theoretically be de-57 scribed as a three-phase process (Jones & Marshall, 1997; Marshall & Schott, 1999; Gelder-58 loos et al., 2011; de Jong et al., 2018). It starts with the preconditioning phase, which 59 consists of a large-scale (order of 100 km) cyclonic gyre circulation with 'doming' isopy-60 cnals, bringing weakly stratified waters of the interior close to the surface. The second 61 phase is the deep convection phase, during which strong winter buoyancy loss gradually 62 deepens the mixed layer, forming a large and deep 'mixed patch' or 'convective patch' 63

with a diameter in the order of 100 km and a depth in the order of 1 km. The deep con-64 vection is followed by the restratification phase which occurs after winter, in spring and 65 summer, on a time scale of weeks to months. During this phase the stratification of the 66 convective patch is reestablished by surface buoyancy gain and by lateral exchange between the convective patch and its stratified environment. The restratification determines 68 the stratification at the start of the next deep convection phase and thus affects precon-69 ditioning. There are only a few areas in the world's oceans where deep convection (more 70 than 1 km) occurs. Two important deep convection regions are the Labrador Sea (e.g. 71 Lilly et al., 1999; Lazier et al., 2002; Straneo, 2006a; Yashayaev & Loder, 2017) and the 72 Irminger Sea (e.g. Pickart et al., 2003; De Jong et al., 2012; de Jong & de Steur, 2016; 73 Piron et al., 2016; de Jong et al., 2018), located west and east of Greenland, respectively. 74 The convectively formed dense waters in these basins can spread to the seas' perimeter 75 and enter the lower limb of the AMOC (Straneo, 2006b; Katsman et al., 2018; Georgiou 76 et al., 2019; Le Bras et al., 2020). 77

Multiple studies have addressed the restratification phase in the Labrador Sea. It 78 was found that there is interannual variability in the deep convection as well as in the 79 restratification in the Labrador Sea (e.g. Lazier et al., 2002; Straneo, 2006a; de Jong et 80 al., 2016). Straneo (2006a) found that there are distinct temporal behaviours in the sur-81 face layer between 0-200 m and the lower layer between 200-1300 m in the Labrador Sea 82 interior. The surface layer properties follow a strong seasonal cycle, whereas the lower 83 layer properties show strong interannual variability. Moreover, the surface layer fresh-84 ens during summer and becomes more saline during winter; these signs in the seasonal 85 cycle are reversed for the lower layer. Many studies have focused on the lateral exchange 86 between the interior, where deep convection occurs, and the buoyant cyclonic boundary 87 current in the Labrador Sea. Katsman et al. (2004) and Gelderloos et al. (2011) found 88 that Irminger Rings (IRs) shed by the West Greenland Current (WGC) near the west 89 coast of Greenland are an important contributor to restratification. Observations from 90 Lilly et al. (1999), Hátún et al. (2007) and de Jong et al. (2014) show that IRs consist 91 of a warm and saline core topped by a colder and fresher surface layer. This can be re-92 lated to the structure of the WGC, which consists of warm and saline Irminger Sea Wa-93 ter (ISW) below colder and fresher polar waters (Cuny et al., 2002). The IR vertical struc-94 ture contributes to the distinct behaviours in the upper and lower layer in the interior 95 Labrador Sea found by Straneo (2006a). 96

Deep convection also occurs in the cyclonic Irminger Gyre (Våge et al., 2011; De Jong 97 et al., 2012; de Jong & de Steur, 2016; de Jong et al., 2018). Recent studies by Lozier 98 et al. (2019), Petit et al. (2020) and Li et al. (2021) found that deep convection in the Irminger Sea is a major contributor to overturning in the North Atlantic subpolar gyre. 100 Moreover, the Irminger Sea is found to be an important region for anthropogenic car-101 bon uptake (Fröb et al., 2016). While the relation between surface buoyancy forcing and 102 Irminger Sea deep convection has been described (De Jong et al., 2012; de Jong & de 103 Steur, 2016; de Jong et al., 2018), the relation to changes in stratification is less explored. 104 Since deep convection is predicted to weaken under climate change as a result of increas-105 ing stratification, it is important to study processes that contribute to stratification in 106 the Irminger Sea. Therefore, this study aims to provide a description of the vertical struc-107 ture and temporal variability of restratification in the Irminger Sea, as well as of the phys-108 ical processes that contribute to it. 109

This paper is organised as follows: in Section 2, the data sets and methods used for the study are introduced. We employ an ocean reanalysis, validated against local mooring data, as well as surface flux data from atmospheric reanalysis. In Section 3, we describe the general hydrographic properties in the Irminger Sea and the interannual variability in stratification; moreover, we investigate the physical processes, such as surface forcing and eddy processes, that contribute to restratification. Finally, a summary and discussion of the results are presented in Section 4.

#### <sup>117</sup> 2 Data and Methods

#### 2.1 Data

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We used the GLORYS12V1 global ocean physics reanalysis from the Copernicus 119 Marine Environment Monitoring Service (CMEMS; Fernandez & Lellouche, 2018) to in-120 vestigate hydrographic variability in the Irminger Sea. GLORYS12V1 uses the NEMO 121 model component (Madec, 2016) and is forced at the surface by ECMWF ERA-Interim 122 reanalysis from 1993–2017, which was replaced by ERA5 reanalysis for recent years. For 123 this study, we used the monthly mean data for the period from January 1993 until De-124 cember 2019. The data set has a spatial resolution of  $1/12^{\circ}$  and has 50 vertical levels, 125 with vertical intervals of 1 m at the surface increasing to 450 m at the bottom. The vari-126 ables obtained from the data set are the potential temperature (thetao), practical salin-127 ity (so), eastward and northward ocean current velocity (uo and vo), and mixed layer 128 depth (mlotst). The mixed layer depth (henceforth abbreviated as MLD) in CMEMS 129 is defined as the depth where the density difference compared to a reference level is 0.01 130  $kg/m^3$  (Madec, 2016, p. 246). The TEOS-10 software was used to recompute practical 131 salinity to absolute salinity  $S_A$ , and potential temperature to conservative temperature 132  $\Theta$  (McDougall & Barker, 2011). From these properties, the potential density anomaly 133 referenced to the surface,  $\sigma_0$ , could then be computed. From now on, unless explicitly 134 stated otherwise, every mention of salinity, temperature, and density from the CMEMS 135 data will refer to  $S_A$ ,  $\Theta$ , and  $\sigma_0$ , respectively. The current velocities were used to com-136 pute eddy kinetic energy (EKE), defined as 137

EKE = 
$$\frac{1}{2} \left[ (u - \bar{u})^2 + (v - \bar{v})^2 \right],$$
 (1)

where  $\bar{u}$  and  $\bar{v}$  are the time-mean zonal and meridional velocity components over the whole time series.

To validate the CMEMS data, it was compared to data from the Long-term Ocean 140 Circulation Observations (LOCO) mooring, which was located at 39.5°W, 59.2°N (Fig-141 ure 1) from September 2003 until June 2018. The processed data from September 2003 142 until July 2015 as published by De Jong et al. (2012), de Jong and de Steur (2016) and 143 de Jong et al. (2018) were used. The LOCO data is not assimilated in the CMEMS re-144 analysis and therefore provides an independent comparison of the reanalysis data with 145 observations. For the comparison, the original daily resolution of the LOCO data was 146 reduced to monthly means. This time series is compared to the CMEMS data interpo-147 lated at the LOCO mooring location. The comparison between the two data sets is de-148 scribed in the Supplementary Material. 149

To study atmospheric fluxes at the ocean surface, we used the ERA5 reanalysis from 150 the Copernicus Climate Change Service (C3S) Climate Data Store (CDS) (Hersbach et 151 al., 2020). The monthly averaged data at the surface was selected for January 1993 un-152 til December 2019. The data set has a horizontal resolution of  $0.25^{\circ}$ . For freshwater fluxes 153 the evaporation (e) and total precipitation (tp) were studied; the sum of these two yields 154 the net atmospheric freshwater flux  $\mathcal{F}$ . For heat fluxes the surface net solar radiation 155 (ssr), surface net thermal radiation (str), surface latent heat flux (slhf), and surface 156 sensible heat flux (sshf) were used to obtain the net atmospheric heat flux Q. 157

#### 158 **2.2** Methods

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## 2.2.1 Defining the Deep Convection Area

Because the aim of this study is to describe restratification after deep convection, we needed to identify the area within the Irminger Sea where deep convection occurs. We focused on the southwestern part of the Irminger Sea, where the centre of the cyclonic Irminger Gyre is located. Out of the 27 years within the CMEMS data set, there are 28 winter months in which MLDs of more than 1 km deep occurred in this region. The MLD field was averaged over these months and the result is shown in Figure 1. The 650 m contour of this MLD field smoothed with a Gaussian filter ( $\sigma = 5$  grid cells = 5/12°) was chosen as the boundary of the deep convection area, which we will abbreviate as DCA. The data from the CMEMS and ERA5 reanalyses were averaged over the DCA to get time series of surface freshwater and heat fluxes, MLD, and salinity, temperature and density profiles in the DCA.

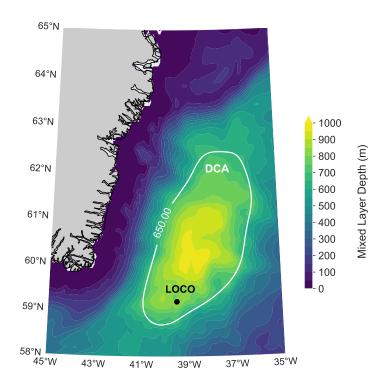


Figure 1. The mixed layer depth (MLD) field averaged over the months from the monthly 1993–2019 CMEMS data set in which the maximum MLD within the region  $[45^{\circ}W, 35^{\circ}W] \times [58^{\circ}N, 65^{\circ}N]$  exceeds 1000 m. In white, the 650 m contour of this MLD field smoothed with a Gaussian filter ( $\sigma = 5$  grid cells) is shown, representing the boundary of the Deep Convection Area (DCA). The black dot shows the location of the LOCO mooring used for data validation.

#### 2.2.2 Quantifying Stratification and Restratification

As a quantitative measure of the stratification of the water column between  $z = -h_1$  and  $z = -h_2$  (where  $0 \le h_1 < h_2$ ), we used the stratification index (SI), defined as the integral of the potential density gradient times depth:

$$SI = \int_{-h_2}^{-h_1} \frac{\partial \sigma}{\partial z} z \, \mathrm{d}z,\tag{2}$$

with units of kg/m<sup>2</sup>. Partial integration of Equation (2) yields an expression of the SI in terms of the density profile  $\sigma(z)$  (rather than the density gradient profile):

$$SI = -h_1 \sigma (-h_1) + h_2 \sigma (-h_2) - \int_{-h_2}^{-h_1} \sigma(z) \, dz.$$
(3)

The SI is the amount of buoyancy that must be removed in order for the water column to mix down from depth  $-h_1$  to depth  $-h_2$  with subsequent uniform density. A strat-

<sup>179</sup> ified layer has a positive SI, a perfectly mixed layer has an SI of zero, and an unstable

layer about to overturn has a negative SI. The SI has been used as a measure for stratification in various Mediterranean studies, where it is usually defined with a factor  $g/\rho_0$ so that it is the integral of the Brunt-Väisälä frequency (e.g. Herrmann et al., 2010; L'Hévéder et al., 2013; Somot et al., 2018; Margirier et al., 2020). The partial integration Equation (3) multiplied by -1 has also been used in North Atlantic Studies; Bailey et al. (2005) referred to it as the integrated buoyancy anomaly, and Frajka-Williams et al. (2014) called it the Convection Resistance.

<sup>187</sup> To study the influence of temperature and salinity on the stratification, two ad-<sup>188</sup> ditional values of the SI were computed:  $SI_T$  and  $SI_S$ , which only consider the contri-<sup>189</sup> bution of temperature and salinity to stratification, respectively. This was done by us-<sup>190</sup> ing the linear approximation for the density of seawater (Dijkstra, 2008):

$$\rho = \rho_0 \left[ 1 - \alpha \left( \Theta - \Theta_0 \right) + \beta \left( S_A - S_{A,0} \right) \right],\tag{4}$$

where  $\Theta_0$  and  $S_{A,0}$  are a reference temperature and salinity and  $\rho_0 = \rho(\Theta_0, S_{A,0})$  is 191 a reference in-situ density. Furthermore,  $\alpha = -\partial \rho / \partial \Theta$  is the thermal expansion coef-192 ficient with units 1/K, and  $\beta = \partial \rho / \partial S_A$  is the saline contraction coefficient in kg/g. 193 Based on Equation (4), we can view  $-\alpha\Theta$  as a first approximation of the temperature 194 contribution to density, and similarly  $\beta S_A$  as the salinity contribution. Hence, replac-195 ing  $\sigma$  in Equation (3) by  $-\alpha\Theta$  and  $\beta S_A$ , respectively, yields an approximation of the tem-196 perature and salinity contribution to stratification, which we denote as  $SI_T$  and  $SI_S$ . Note 197 that  $-\alpha\Theta$  and  $\beta S_A$  are unitless, so that the stratification contributions SI<sub>T</sub> and SI<sub>S</sub> are 198 given in m instead of  $kg/m^2$ . From the time series of salinity, temperature and density 199 profiles averaged over the DCA, monthly values of the SI,  $SI_T$  and  $SI_S$  were computed. 200

Using the SI, a quantitative measure for restratification could be created as well. The restratification in a certain year was defined to be the increase in SI from the minimum in winter to the maximum in summer. Thus the monthly time series of the SI could be used to compute annual time series of restratification.

#### 205 3 Results

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#### 3.1 The Deep Convection Area

Figure 2 shows maps of the time-mean hydrographic properties in the Irminger Sea. 207 All of these properties are depth-averaged over the upper 600 m, where the variability 208 in density related to convection and restratification is strongest (Figure 3). From Fig-209 ure 2 we can identify the main currents and their characteristics. On the eastern side 210 of the Irminger Sea we find the northward-flowing Irminger Current (IC) transporting 211 warm and saline water originating from the North Atlantic Current (de Jong et al., 2020; 212 Fried & Jong, 2022). The IC follows the bathymetry, so around 65°N, it veers west and 213 continues southward along the East Greenland Shelf. Alongside this part of the IC we 214 find two currents of Arctic origin, which are thus fresher and colder: the East Greenland 215 Current (EGC) along the shelfbreak and the East Greenland Coastal Current (EGCC) 216 on the inner shelf (e.g. Le Bras et al., 2018; Duyck & De Jong, 2021). The EKE in the 217 Irminger Sea is highest in the IC, in the EGCC and to the northeast of the DCA. These 218 regions are in agreement with the strong EKE regions that were derived from surface 219 drifter data by Fratantoni (2001) and from satellite altimetry data by Volkov (2005) and 220 Fan et al. (2013). Eddies observed in the central Irminger Sea travel westward and have 221 a warm and saline core, which suggests that they are shed by the IC (Fan et al., 2013). 222

To study the interannual variability and vertical structure of hydrographic prop-223 erties in the DCA, we consider the full 1993–2019 time series of spatially averaged pro-224 files of DCA salinity, temperature and density (Figure 3). The mean and maximum MLD 225 over the DCA are also shown. The temperature and density have monotonously strat-226 ified structures, with warm, buoyant layers above cold, dense ones. By contrast, salin-227 ity has a mid-depth minimum in the layer between 500 and 1500 m, which is the signa-228 ture of convectively formed water. From the evolution of the MLD we see that there is 229 strong interannual variability in the convection strength. There is also interannual vari-230 ability in stratification: some years have much weaker vertical density gradients than other 231 years. In the early 1990s, cold and stormy winters forced the deepest mixed layers of the 232 time series (Våge et al., 2011), which led to weak vertical density gradients and outcrop-233 ping of isopycnals. In the 2000s, mixed layers were shallower, and the DCA was more 234 stratified. From 2015 until 2018, the MLD increased again, resulting in a thickening of 235 deep cold and dense layers and thus a decrease in stratification. The largest temporal 236 fluctuations in hydrographic properties are found in the upper 600 m of the water col-237 umn (right column of Figure 3), on which we will focus. Below this depth the tempo-238 ral variations in density become very small; at 600 m the standard deviation in density 239 is only 7.3% of that at the surface. 240

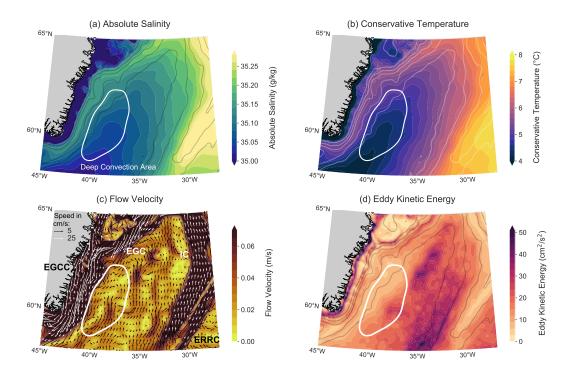


Figure 2. The 1993–2019 mean from the CMEMS reanalysis of the (a) absolute salinity, (b) conservative temperature, (c) flow velocity and (d) eddy kinetic energy in the Irminger Sea, averaged over the upper 600 m of the ocean. The white contour indicates the boundary of the Deep Convection Area. In panels (a), (b) and (d), smoothed contours of the bathymetry between 500 m and 2500 m are shown at 500 m intervals. In panel (c), velocity direction and magnitude are indicated by vectors. Velocities smaller than 0.07 m/s are indicated by black arrows, larger velocities by white arrows. The black and white vectors have different scalings, as shown by the reference vectors in the top right of the panel. In this panel the different currents are indicated: EGCC = East Greenland Coastal Current; EGC = East Greenland Current; IC = Irminger Current; ERRC = East Reykjanes Ridge Current.

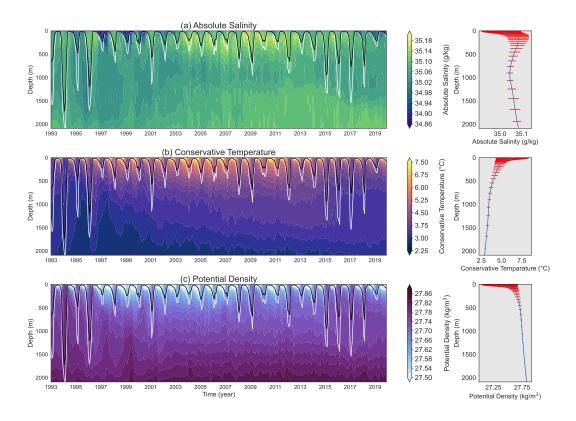
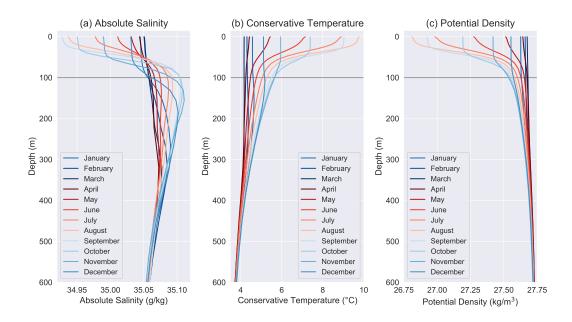


Figure 3. The temporal evolution of profiles of (a) absolute salinity, (b) conservative temperature and (c) potential density from the CMEMS reanalysis averaged over the Deep Convection Area (DCA). The mixed layer depth averaged over the DCA is shown in black; the maximum mixed layer depth reached within the DCA is shown in white. The panels on the right show the 1993–2019 time-mean (blue curve) and standard deviations (red bars) of each variable.

#### 3.2 Structure and Variability of Restratification

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To see how properties within the DCA generally vary throughout the year, we com-242 posed climatologies of the salinity, temperature and density profiles in the upper 600 m 243 of the DCA (Figure 4). A seasonal cycle consisting of deepening convection in autumn 244 and winter, followed by restratification in spring and summer, is clearly visible in all three 245 variables. The deepest mixed layers are typically reached in February through April (al-246 though a few years with shallower convection imprint as a slight non-zero vertical gra-247 dient in the climatology). Restratification sets in in April and May, with the halocline, 248 249 thermocline and pycnocline steepening through to August (for temperature and density) and September (for salinity). The uppermost 30–50 m of the water column is typically 250 characterised by a wind-mixed layer (Sampe & Xie, 2007). The salinity climatology (Fig-251 ure 4a) shows different behaviour in the upper and lower parts of the water column. In 252 summer, the surface layer freshens, likely through precipitation and freshwater from the 253 Greenland shelf (Duyck & De Jong, 2021; Duyck et al., 2022). The deeper part of the 254 water column is dominated by lateral exchange with the perimeter of the Irminger Sea, 255 where salinity is generally higher than in the DCA (Figure 2a). In winter, convection 256 homogenises the salinity over the upper and lower layer, making the former saltier and 257 the latter fresher. The temperature and density climatologies (Figure 4b,c) exhibit the 258 same sign in change over the upper 600 m, but a difference in timing and magnitude in 259 restratification in the upper and lower layer can still be observed. In the upper layer, the 260 increase in temperature is much stronger and maximum temperatures are reached in Au-261 gust. This is likely the effect of heat gain through surface fluxes (Figure 5). In the lower 262 layer, the temperature gradient remains significantly weaker and warming continues un-263 til November, as is expected from warming by lateral exchange with somewhat warmer 264 waters (Figure 2b). The net effect of temperature and salinity is a strong increase in sta-265 ble density gradient in the upper layer during summer and a weaker, delayed increase 266 below (Figure 4c). Given the different behaviour in restratification in the upper layer 267 down to 100 m and the lower layer below (down to 600 m), we will study the interan-268 nual variability in these layers separately. 269



**Figure 4.** Climatologies of the (a) absolute salinity, (b) conservative temperature and (c) potential density profiles averaged over the Deep Convection Area, from the monthly 1993–2019 CMEMS reanalysis.

The stratification index of the two layers is shown in Figure 5, together with the 270 climatology of the MLD, the net surface heat flux and the net surface freshwater flux. 271 The SI climatology shows the difference in timing described above. Additionally, the lower 272 layer SI has a larger standard deviation in the lower layer, indicating stronger interan-273 nual variability. Increasing upper layer SI matches with periods of surface heating and 274 decreasing upper layer SI with surface cooling. Timing in lower layer SI matches with 275 the MLD as expected. The net surface freshwater flux is highly variable, but generally 276 is higher in summer than in winter, which can add to the summer freshening of the up-277 per layer. This is discussed further in relation to the interannual variability in Section 278 3.3. 279

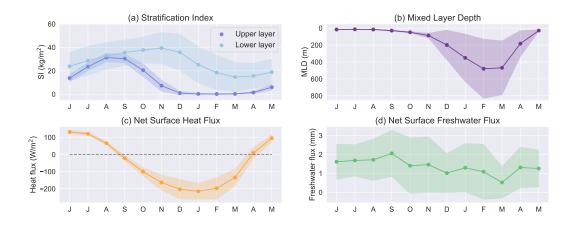
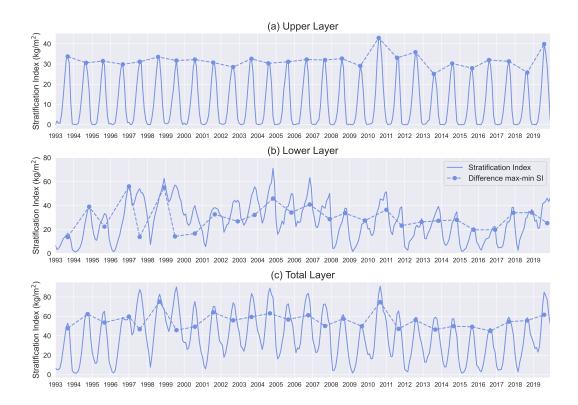


Figure 5. Climatologies of the (a) stratification index in the upper and lower layer, (b) mixed layer depth, (c) net surface heat flux and (d) net surface freshwater flux averaged over the Deep Convection Area. The data of (a) and (b) are from the monthly 1993–2019 CMEMS reanalysis, the data of (c) and (d) are from the monthly 1993–2019 ERA5 reanalysis. The shaded areas indicate the standard deviations.

The time series of SI and restratification from 1993 to 2019 (Figure 6) show the 280 difference in interannual variability in the upper, lower and total (upper+lower) layer. 281 The upper layer stratification is dominated by the seasonal cycle, with very little inter-282 annual variability in either SI or restratification. Two years (2010 and 2019) stand out 283 as having higher summer maxima and we will come back to these in Section 3.3.2. In 284 the lower layer interannual variability dominates over seasonality. Unlike in the upper 285 layer, where the SI is set to zero each winter, the lower layer has a memory of stratifi-286 cation and mixing over previous years. The stratification over the total layer (from 0– 287 600 m) is the superposition of the strong seasonal cycle in the upper layer and the in-288 terannual variability in the lower layer. 289



**Figure 6.** The stratification index (SI) time series in the Deep Convection Area for (a) the upper layer from 0–100 m, (b) the lower lower from 100–600 m, and (c) the total (upper+lower) layer from 0–600 m. The dashed line shows the annual values of the restratification, defined as the increase in SI from the minimum in winter to the maximum in summer.

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#### 3.3 Physical Processes Contributing to Restratification

#### 3.3.1 Role of Salinity and Temperature in Restratification

The climatology in Figure 4 suggested there are different roles for temperature and 292 salinity in restratification. To study these further, we investigate time series of the  $SI_T$ 293 and  $SI_S$  as introduced in Section 2.2.2 (Figure 7). In the upper layer, the  $SI_T$  and  $SI_S$ 294 both follow a strong seasonal cycle. Values are slightly negative in winter indicating un-295 stable stratification during deep convection and positive for the remainder of the year 296 due to higher temperatures and lower salinities. In the lower layer the  $SI_S$  is always neg-297 ative, indicating the destabilising effect of salinity (saline water over fresher water) seen 298 also in Figure 4a. Despite the stable, fresh upper layer, the effect of salinity over the to-299 tal 0–600 m layer is destabilising. Since the temperature exhibits the same sign gradi-300 ent over both layers, the effect of temperature is always stabilising except during some 301 winter months with deep mixed layers (cf. Figure 3). In the lower layer, temperature and 302 salinity co-vary as the restratifying waters from the perimeter are warm and saline. This 303 is not necessarily so for the upper layer, where heat and freshwater fluxes can vary in-304 dependently. Even so, the correlation between the variability in  $SI_S$  and  $SI_T$  is quite strong. 305 From the yearly restratification contributions in Figure 7 it can be seen that the restrat-306 ification contribution of temperature (light dashed line) is always larger than that of salin-307 ity (dark dashed line) in absolute value. 308

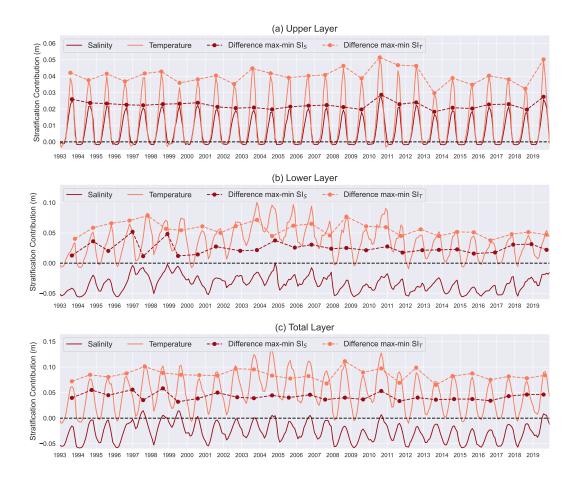
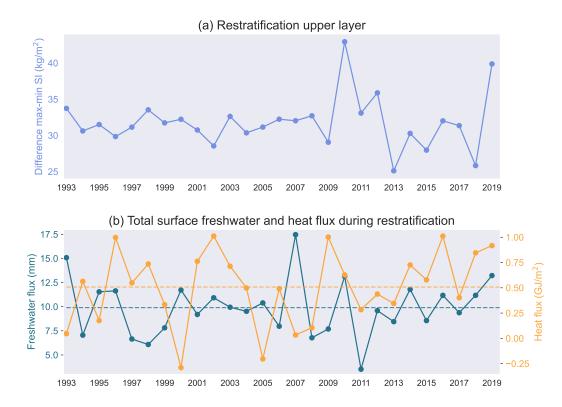


Figure 7. The stratification contribution of salinity and temperature in the Deep Convection Area for (a) the upper layer from 0–100 m, (b) the lower lower from 100–600 m, and (c) the total water layer from 0–600 m. The dashed lines show the annual values of the restratification, defined as the increase in stratification contribution from the minimum in winter to the maximum in summer.

#### 3.3.2 Role of Surface Fluxes and Lateral Fluxes in Restratification

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The climatology from Figure 5 suggested a relation between upper layer stratifi-310 cation and surface fluxes. Therefore we investigate interannual variability in surface heat 311 and freshwater fluxes over the period of restratification, from the winter minimum SI in 312 the upper layer to the next summer maximum of each year. Figure 8 shows the spring/summer 313 accumulated atmospheric freshwater and heat fluxes together with the annual upper layer 314 restratification. The restratification is not significantly correlated with either the sur-315 face freshwater flux nor with the heat flux separately. However, in years when both are 316 stronger than average, they may reinforce each other. This appears to be the cases for 317 the years with the highest restratification values, 2010 and 2019. Both years were char-318 acterised by above-average freshwater and heat gain during the restratification period, 319 with the high freshwater gain mainly due to low evaporation (not shown here). The pres-320 ence of a fresh surface layer, which already increased stratification, restricted warming 321 by the high heat flux to a layer, resulting in higher temperatures and very high strat-322 ification by the end of summer. The reverse does not appear to be true; years with below-323 average heat and freshwater fluxes do not exhibit much weaker restratification. 324



**Figure 8.** Annual time series of (a) restratification and (b) total accumulated surface freshwater and heat flux during the restratification period in the upper layer over the Deep Convection Area. The restratification period starts in the month after the occurrence of the minimum winter stratification index (SI) and ends in the month of the next maximum SI in the upper layer. The dashed horizontal lines show the time-mean values of the depicted time series.

In the lower layer, as opposed to the upper layer, there is strong interannual vari-325 ability in both the summer maximum SI and the winter minimum SI (Figure 6b). In win-326 ter, the stratification in the lower layer is directly affected by the MLD (significant cor-327 relation of -0.68 at a 95% significance level). Figure 9a shows the annual restratifica-328 tion in the lower layer together with the maximum MLD in the preceding winter. Years 329 with high maximum stratification, such as 2004 and 2006, were preceded by winters with 330 shallow mixed layers (MLD < 600 m) so that some stratification remained at the end 331 of winter. By contrast, years in which the maximum stratification was low, e.g. 2009 and 332 2015 (Figure 6b), were preceded by mixed layers of more than 750 m deep. However, there 333 is no significant correlation between the MLD and the restratification in the lower layer. 334 On the longer time scales (order of 10 years) they are correlated, but this may also be 335 caused by both processes responding to changes in North Atlantic Oscillation phases. 336

During spring and summer the lower layer is isolated from the upper layer, so the 337 lower layer restratification is less affected by surface fluxes and more by lateral advec-338 tion of buoyant water from the perimeter of the Irminger Sea. Since this advection must 339 cross the streamlines of the Irminger Gyre to enter the DCA, it is likely facilitated by 340 eddies. This appears to be supported by the relation between the time series of lower 341 layer restratification and mean surface EKE during restratification (Figure 9b; signif-342 icantly correlated at r = 0.66). The EKE time series is averaged over the area of the 343 WOCE AR7E line (Koltermann et al., 2011) that lies east of the eastern DCA bound-344 ary between 315 and 700 km (Figures 9c and d), where the highest EKE values are found. 345

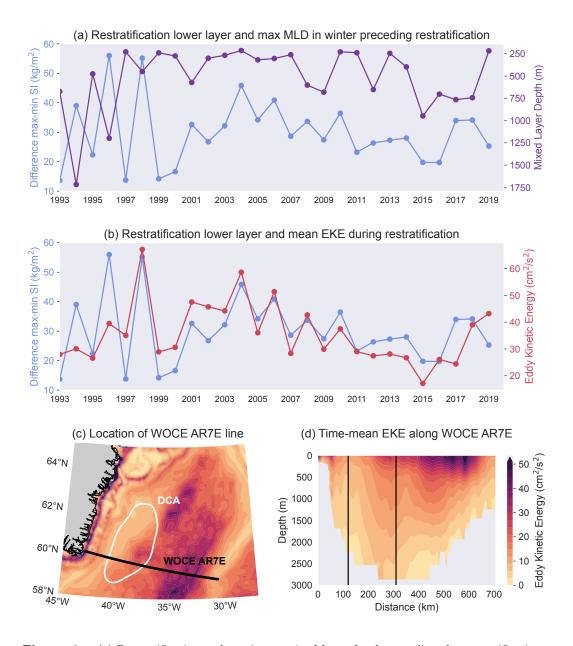


Figure 9. (a) Restratification and maximum mixed layer depth preceding the restratification period in the lower layer of the Deep Convection Area (DCA). (b) Restratification in the lower layer of the DCA and surface eddy kinetic energy (EKE) along the WOCE AR7E section east-ward of the DCA, averaged over the restratification phase. The restratification period starts in the month after the occurrence of the minimum winter stratification and ends in the month of the next maximum stratification in the upper layer. (c) The WOCE AR7E line in the Irminger Sea (Koltermann et al., 2011) and the boundary of the DCA, plotted on top of the 1993–2019 mean EKE field averaged over the upper 600 m. (d) The 1993–2019 mean from the CMEMS reanalysis of the EKE along the WOCE AR7E line. The vertical black lines mark the points where the line intersects the DCA boundary.

<sup>346</sup> The eddies from this region bring warm and saline Irminger Current water into the DCA

(Fan et al., 2013); this further supports the notion that eddies drive lower layer restrat ification.

#### <sup>349</sup> 4 Discussion and Conclusions

To summarise our results, most of the variability in stratification in the central Irminger 350 Sea is located in the upper 600 m of the water column. We found distinctly different vari-351 ability in stratification in the upper 100 m and the 100–600 m layer. In the upper layer, 352 there is a strong seasonal cycle with freshening and warming in spring and summer and 353 very little interannual variability. Exceptionally high upper layer stratification may be 354 reached in years when both heat and freshwater fluxes are above average. In the lower 355 layer, there is substantial interannual variability in stratification. The winter minimum 356 stratification is determined by the strength of convection: deeper mixed layers lead to 357 weaker stratification. The restratification after winter is strongly linked to the eddy ki-358 netic energy in the eastern Irminger Sea, suggesting that warm, saline eddies shed by 359 the Irminger Current play an important role in restratifying the DCA. 360

This two layer structure in restratification in the Irminger Sea is very similar to 361 what is observed in the Labrador Sea (Straneo, 2006a), which also displays surface fresh-362 ening during restratification while the layer below gets more saline. However, there is 363 an important difference. At least part of the two layer structure of restratification in the 364 Labrador Sea is derived from the two layer eddies shed from the boundary current. These 365 eddies have a cold and fresh upper layer extending until 100 to 400 m depth and a warm 366 and saline layer below (Lilly et al., 1999; Hátún et al., 2007; de Jong et al., 2014). How-367 ever, in the Irminger Sea no eddies with a fresh top are observed, due to their origin in 368 the Irminger Current which is warm and saline throughout (Fan et al., 2013). Instead, 369 the upper layer freshening in the Irminger Sea is due to surface freshwater gain and thin 370 layers of freshwater driven off the East Greenland shelf during strong wind events (Duyck 371 et al., 2022). 372

Amplification of upper layer stratification through co-occurring above normal sur-373 face heat and freshwater gain was seen in 2010 and 2019. The 2010 event was described 374 by Oltmanns et al. (2018). They speculate that the fresh surface layer might have had 375 an additional continental origin. They found that surface temperatures were high in a 376 broad area near the southeast Greenland shelf, where glacier melt would be accelerated. 377 Freshening due to ice melt is a plausible cause for freshwater anomalies near the East 378 Greenland shelf, which we know can be exported to the DCA (Foukal et al., 2020; Duyck 379 & De Jong, 2021; Duyck et al., 2022) and thus increase the stratification there. In 2019 380 runoff from Greenland was high was well (Slater et al., 2021), but there was also an ad-381 ditional cause for the high stratification: a strong freshwater anomaly that formed be-382 tween 2012 and 2016 in the eastern subpolar North Atlantic, as described by Holliday 383 et al. (2020). This anomaly went on to propagate northwards into the Irminger Sea, the 384 Labrador Sea and the Nordic Seas. Its effects can be seen in the freshening of the up-385 per and intermediate layers in the Irminger Sea from 2017 onwards (Figure 3a) as well 386 as in the high stratification in 2019 (Figure 6a). In the long term, anthropogenic climate 387 change is expected to further enhance Greenland ice melt. Upper layer stratification in-388 creases as seen in 2010 and 2019 may occur more frequently and stronger. 389

The lower layer displays much stronger interannual variability, partly through longer 390 ocean memory (set by convection in previous winters) and partly through differences in 391 restratification. We found a strong relation between the strength of the annual restrat-392 ification and the EKE in the eastern Irminger Sea. Anticyclonic eddies with warm and 393 saline cores shed by the Irminger Current move westward into the DCA, thus restrat-394 ifying it. In the Labrador Sea, eddies play a crucial role both in determining the loca-395 tion of the deep convection region (Chanut et al., 2008) as well as in the downwelling 396 and export of convective waters (Georgiou et al., 2019). As the Irminger Sea is an im-397

<sup>398</sup> portant contributor to the overturning in the subpolar gyre (Lozier et al., 2019; Petit

et al., 2020; Li et al., 2021) it is imperative that eddy processes and their contribution

 $_{400}$  to restratification in the Irminger Sea are better understood.

## 401 Data Availability Statement

The CMEMS monthly mean data is available at https://resources.marine.copernicus
.eu/product-detail/GLOBAL\_MULTIYEAR\_PHY\_001\_030/DATA-ACCESS. The files for 19932012 were last accessed on 29-10-2020; the files for 2013-2019 were last accessed on 2105-2021.

The ERA5 monthly averaged data on single levels is available at https://cds.climate .copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels-monthly-means ?tab=overview and was last accessed on 16-09-2021.

Data was processed using Python 3.6.13. All code used for processing data and making figures for this study is available at https://github.com/MiriamSterl/RestratificationIrmingerSea.

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### 415 **References**

423

- Bailey, D. A., Rhines, P. B., & Häkkinen, S. (2005). Formation and pathways of North Atlantic Deep Water in a coupled ice-ocean model of the
  Arctic-North Atlantic Oceans. *Climate Dynamics*, 25, 497-516. doi:
  10.1007/s00382-005-0050-3
- Brown, P. J., Mcdonagh, E. L., Sanders, R., Watson, A. J., King, B. A., Smeed,
  D. A., ... Meinen, C. S. (2021). Anthropogenic carbon transports and uptake.
  - doi: 10.1038/s41561-021-00774-5
- Chanut, J., Barńier, B., Large, W., Debreu, L., Penduff, T., Molines, J. M., & Mathiot, P. (2008). Mesoscale eddies in the Labrador Sea and their contribution
  to convection and restratification. *Journal of Physical Oceanography*, 38, 1617-1643. doi: 10.1175/2008JPO3485.1
- Cuny, J., Rhines, P. B., Niiler, P. P., & Bacon, S. (2002). Labrador Sea boundary
   currents and the fate of the Irminger Sea Water. Journal of Physical Oceanog raphy, 32, 627-647. doi: 10.1175/1520-0485(2002)032(0627:LSBCAT)2.0.CO;2
- <sup>431</sup> De Jong, M. F., Aken, H. M. V., Våge, K., & Pickart, R. S. (2012). Convective mix <sup>432</sup> ing in the central Irminger Sea: 2002-2010. Deep-Sea Research Part I: Oceano <sup>433</sup> graphic Research Papers, 63, 36-51. doi: 10.1016/j.dsr.2012.01.003
- de Jong, M. F., Bower, A. S., & Furey, H. H. (2014). Two years of observations of
  warm-core anticyclones in the Labrador Sea and their seasonal cycle in heat
  and salt stratification. Journal of Physical Oceanography, 44, 427-444. doi:
  10.1175/JPO-D-13-070.1
- de Jong, M. F., Bower, A. S., & Furey, H. H. (2016). Seasonal and interannual variations of Irminger ring formation and boundary-interior heat exchange in FLAME. Journal of Physical Oceanography, 46, 1717-1734. doi:
  10.1175/JPO-D-15-0124.1
- de Jong, M. F., & de Steur, L. (2016). Strong winter cooling over the Irminger
  Sea in winter 2014–2015, exceptional deep convection, and the emergence of
  anomalously low SST. *Geophysical Research Letters*, 43, 7106-7113. doi:
  10.1002/2016GL069596

| 446 | de Jong, M. F., de Steur, L., Fried, N., Bol, R., & Kritsotalakis, S. (2020). Year-   |
|-----|---|
| 447 | Round Measurements of the Irminger Current: Variability of a Two-Core                 |
| 448 | Current System Observed in 2014–2016. Journal of Geophysical Research:                |
| 449 | Oceans, 125. doi: 10.1029/2020JC016193  |
| 450 | de Jong, M. F., Oltmanns, M., Karstensen, J., & Society, T. O. (2018). Deep con-      |
| 451 | vection in the Irminger Sea observed with a dense mooring array. Oceanogra-           |
| 452 | phy, 31, 50-59.   |
| 453 | Dijkstra, H. A. (2008). Dynamical oceanography. Springer US. doi: 10.1007/978-3       |
| 454 | -540-76376-5  |
| 455 | Duyck, E., & De Jong, M. F. (2021). Circulation Over the South-East Greenland         |
| 456 | Shelf and Potential for Liquid Freshwater Export: A Drifter Study. Geophysi-          |
| 457 | cal Research Letters, 48, 1-9. doi: 10.1029/2020GL091948                              |
| 458 | Duyck, E., Gelderloos, R., & Jong, M. F. (2022, 5). Wind-Driven Freshwater Ex-        |
| 459 | port at Cape Farewell. Journal of Geophysical Research: Oceans, 127. doi: 10          |
| 460 | .1029/2021jc018309  |
| 461 | Fan, X., Send, U., Testor, P., Karstensen, J., & Lherminier, P. (2013). Observa-      |
| 462 | tions of Irminger Sea anticyclonic eddies. Journal of Physical Oceanography,          |
| 463 | 43, 805-823. doi: 10.1175/JPO-D-11-0155.1   |
| 464 | Fernandez, E., & Lellouche, J. M. (2018). Product User Manual For the Global          |
| 465 | Ocean Physical Reanalysis product GLOBAL_REANALYSIS_PHY_001_030.                      |
| 466 | Foukal, N. P., Gelderloos, R., & Pickart, R. S. (2020). A continuous pathway for      |
| 467 | fresh water along the East Greenland shelf. Science Advances, 6. doi: 10.1126/        |
| 468 | sciadv.abc4254  |
| 469 | Fox-Kemper, B., Hewitt, H. T., Xiao, C., Adalgeirsdóttir, G., Drijfhout, S. S., Ed-   |
| 470 | wards, T. L., Yu, Y. (2021). Ocean, Cryosphere and Sea Level Change.                  |
| 471 | Cambridge University Press. In Press.   |
| 472 | Frajka-Williams, E., Rhines, P. B., & Eriksen, C. C. (2014). Horizontal stratifica-   |
| 473 | tion during deep convection in the Labrador Sea. Journal of Physical Oceanog-         |
| 474 | raphy, 44, 220-228. doi: 10.1175/JPO-D-13-069.1                                       |
| 475 | Fratantoni, D. M. (2001). North Atlantic surface circulation during the 1990's ob-    |
| 476 | served with satellite-tracked drifters. Journal of Geophysical Research, 106,         |
| 477 | 22,067-22,093.  |
| 478 | Fried, N., & Jong, M. F. (2022, 3). The Role of the Irminger Current in the           |
| 479 | Irminger Sea Northward Transport Variability. Journal of Geophysical Re-              |
| 480 | search: Oceans, 127. Retrieved from https://onlinelibrary.wiley.com/                  |
| 481 | doi/10.1029/2021JC018188 doi: 10.1029/2021JC018188                                    |
| 482 | Fröb, F., Olsen, A., Våge, K., Moore, G. W., Yashayaev, I., Jeansson, E., & Ra-       |
| 483 | jasakaren, B. (2016). Irminger Sea deep convection injects oxygen and an-             |
| 484 | thropogenic carbon to the ocean interior. <i>Nature Communications</i> , 7. doi:      |
| 485 | 10.1038/ncomms13244   |
| 486 | Gelderloos, R., Katsman, C. A., & Drijfhout, S. S. (2011). Assessing the roles of     |
| 487 | three eddy types in restratifying the Labrador Sea after deep convection. Jour-       |
| 488 | nal of Physical Oceanography, 41, 2102-2119. doi: 10.1175/JPO-D-11-054.1              |
| 489 | Georgiou, S., van der Boog, C. G., Brüggemann, N., Ypma, S. L., Pietrzak, J. D., &    |
| 490 | Katsman, C. A. (2019). On the interplay between downwelling, deep convec-             |
| 491 | tion and mesoscale eddies in the Labrador Sea. Ocean Modelling, 135, 56-70.           |
| 492 | Retrieved from https://doi.org/10.1016/j.ocemod.2019.02.004 doi:                      |
| 493 | 10.1016/j.ocemod.2019.02.004  |
| 494 | Herrmann, M., Sevault, F., Beuvier, J., & Somot, S. (2010). What induced the          |
| 495 | exceptional 2005 convection event in the northwestern Mediterranean basin?            |
| 496 | Answers from a modeling study. Journal of Geophysical Research: Oceans,               |
| 497 | 115, 1-19. doi: 10.1029/2010JC006162  |
| 498 | Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J., |
| 499 | Thépaut, J. N. (2020). The ERA5 global reanalysis. Quarterly Journal of               |
| 500 | the Royal Meteorological Society, $146$ , 1999-2049. doi: $10.1002/qj.3803$           |

| 501<br>502<br>503 | <ul> <li>Holliday, N. P., Bersch, M., Berx, B., Chafik, L., Cunningham, S., Florindo-López,</li> <li>C., Yashayaev, I. (2020). Ocean circulation causes the largest freshening<br/>event for 120 years in eastern subpolar North Atlantic. Nature Communica-</li> </ul> |
|-------------------|---|
| 504               | tions, 11. doi: $10.1038/s41467-020-14474-y$  |
| 505               | Hátún, H., Eriksen, C. C., & Rhines, P. B. (2007). Buoyant eddies entering  |
| 506               | the Labrador Sea observed with gliders and altimetry. Journal of Physical   |
| 507               | Oceanography, 37, 2838-2854. doi: 10.1175/2007JPO3567.1   |
| 508               | Jones, H., & Marshall, J. (1997). Restratification after deep convection. Journal of  |
| 509               | Physical Oceanography, 27, 2276-2287. doi: 10.1175/1520-0485(1997)027(2276:   |
| 510               | $RADC\rangle 2.0.CO;2$  |
| 511               | Katsman, C. A., Drijfhout, S. S., Dijkstra, H. A., & Spall, M. A. (2018). Sink-   |
| 512               | ing of dense North Atlantic waters in a global ocean model: Location and  |
| 513               | controls. Journal of Geophysical Research: Oceans, 123, 3563-3576. doi:   |
| 514               | $10.1029/2017 \mathrm{JC}013329$  |
| 515               | Katsman, C. A., Spall, M. A., & Pickart, R. S. (2004). Boundary current eddies  |
| 516               | and their role in the restratification of the Labrador Sea. Journal of Phys-  |
| 517               | ical Oceanography, $34$ , 1967-1983. doi: $10.1175/1520-0485(2004)034(1967:$  |
| 518               | $BCEATR \rangle 2.0.CO;2$   |
| 519               | Koltermann, K., Gouretski, V., & Jancke, K. (2011). Hydrographic Atlas of the   |
| 520               | World Ocean Circulation Experiment (WOCE), Volume 3: Atlantic Ocean   |
| 521               | (J. G. M. Sparrow P. Chapman, Ed.).   |
| 522               | Lazier, J., Hendry, R., Clarke, A., Yashayaev, I., & Rhines, P. (2002). Convection  |
| 523               | and restratification in the Labrador Sea, 1990-2000. Deep-Sea Research Part I:  |
| 524               | Oceanographic Research Papers, 49, 1819-1835. doi: 10.1016/S0967-0637(02)   |
| 525               | 00064-X   |
| 526               | Le Bras, I. A., Straneo, F., Holte, J., de Jong, M. F., & Holliday, N. P. (2020).   |
| 527               | Rapid Export of Waters Formed by Convection Near the Irminger Sea's West-   |
| 528               | ern Boundary. Geophysical Research Letters, 47. doi: 10.1029/2019GL085989   |
| 529               | Le Bras, I. A., Straneo, F., Holte, J., & Holliday, N. P. (2018). Seasonality of Fresh-   |
| 530               | water in the East Greenland Current System From 2014 to 2016. Journal of  |
| 531               | Geophysical Research: Oceans, 123, 8828-8848. doi: 10.1029/2018JC014511   |
| 532               | L'Hévéder, B., Li, L., Sevault, F., & Somot, S. (2013). Interannual variability of  |
| 533               | deep convection in the Northwestern Mediterranean simulated with a coupled  |
| 534               | AORCM. Climate Dynamics, 41, 937-960. doi: 10.1007/s00382-012-1527-5  |
| 535               | Li, F., Lozier, M. S., Holliday, N. P., Johns, W. E., Le Bras, I. A., Moat, B. I.,  |
| 536               | de Jong, M. F. (2021). Observation-based estimates of heat and freshwa-   |
| 537               | ter exchanges from the subtropical North Atlantic to the Arctic. <i>Progress</i>  |
| 538               | in Oceanography, 197, 102640. Retrieved from https://doi.org/10.1016/   |
| 539               | j.pocean.2021.102640 doi: 10.1016/j.pocean.2021.102640  |
| 540               | Lilly, J. M., Rhines, P. B., Visbeck, M., Davis, R., Lazier, J. R., Schott, F., &   |
| 541               | Farmer, D. (1999). Observing deep convection in the Labrador Sea dur-   |
| 542               | ing winter 1994/95. Journal of Physical Oceanography, 29, 2065-2098. doi:   |
| 543               | 10.1175/1520-0485(1999)029(2065:odcitl)2.0.co;2   |
| 544               | Lozier, M. S., Li, F., Bacon, S., Bahr, F., Bower, A. S., Cunningham, S. A.,  |
| 545               | Zhao, J. (2019). A sea change in our view of overturning in the subpolar North  |
| 546               | Atlantic. Science, 363, 516-521. doi: 10.1126/science.aau6592   |
| 547               | Madec, G. (2016). NEMO Ocean Engine. Retrieved from http://www.nemo-ocean   |
| 548               | .eu/About-NEMO/Reference-manuals%5Cnpapers2://publication/uuid/   |
| 549               | 73E7FF17-99BE-4B10-A823-0037C823EF6E doi: 10.5281/zenodo.3248739  |
| 550               | Margirier, F., Testor, P., Heslop, E., Mallil, K., Bosse, A., Houpert, L., Tail-  |
| 551               | landier, V. (2020). Abrupt warming and salinification of intermediate waters  |
| 552               | interplays with decline of deep convection in the Northwestern Mediterranean  |
| 553               | Sea. Scientific Reports, 10, 1-11. Retrieved from https://doi.org/10.1038/  |
| 554               | s41598-020-77859-5 doi: 10.1038/s41598-020-77859-5  |

- Marshall, J., & Schott, F. (1999). Open-Ocean Convection: Observations, Theory, and Models. *Reviews of Geophysics*, 37, 1-64.
- McCarthy, G. D., Smeed, D. A., Cunningham, S. A., & Roberts, C. D. (2017). At lantic Meridonal Overturning Circulation. Marine Climate Change Impacts
   *Partnership: Science Review*, 15-21. Retrieved from www.ncdc.noaa.gov.
   doi: 10.14465/2017.arc10.002-atl
- McDougall, T. J., & Barker, P. (2011). Getting started with TEOS-10 and the Gibbs Seawater (GSW) Oceanographic Toolbox. SCOR/IAPSO WG127. Retrieved from http://www.teos-10.org/pubs/gsw/v3\_04/pdf/Getting\_Started.pdf
- Oltmanns, M., Karstensen, J., & Fischer, J. (2018). Increased risk of a shutdown
   of ocean convection posed by warm North Atlantic summers. Nature Cli mate Change, 8, 300-304. Retrieved from http://dx.doi.org/10.1038/
   s41558-018-0105-1 doi: 10.1038/s41558-018-0105-1
- Petit, T., Lozier, M. S., Josey, S. A., & Cunningham, S. A. (2020). Atlantic Deep
   Water Formation Occurs Primarily in the Iceland Basin and Irminger Sea
   by Local Buoyancy Forcing. *Geophysical Research Letters*, 47, 1-9. doi: 10.1029/2020GL091028
- Pickart, R. S., Straneo, F., & Moore, G. W. (2003). Is Labrador Sea Water formed
   in the Irminger basin? Deep-Sea Research Part I: Oceanographic Research Papers, 50, 23-52. doi: 10.1016/S0967-0637(02)00134-6
- Piron, A., Thierry, V., Mercier, H., & Caniaux, G. (2016). Argo float observations
   of basin-scale deep convection in the Irminger sea during winter 2011–2012.
   *Deep-Sea Research Part I: Oceanographic Research Papers*, 109, 76-90.

578

579

584

585

586

587

588

595

596

603

- Retrieved from http://dx.doi.org/10.1016/j.dsr.2015.12.012 doi: 10.1016/j.dsr.2015.12.012
- Pérez, F. F., Mercier, H., Vázquez-Rodríguez, M., Lherminier, P., Velo, A., Pardo,
   P. C., ... Ríos, A. F. (2013). Atlantic Ocean CO2 uptake reduced by weaken ing of the meridional overturning circulation. *Nature Geoscience*, 6, 146-152.
   doi: 10.1038/ngeo1680
  - Sampe, T., & Xie, S. P. (2007). Mapping high sea winds from space: A global climatology. Bulletin of the American Meteorological Society, 88, 1965-1978. doi: 10 .1175/BAMS-88-12-1965
  - Slater, T., Shepherd, A., McMillan, M., Leeson, A., Gilbert, L., Muir, A., ... Briggs, K. (2021, 12). Increased variability in Greenland Ice Sheet runoff from satellite observations. *Nature Communications*, 12. doi: 10.1038/s41467-021-26229-4
- Somot, S., Houpert, L., Sevault, F., Testor, P., Bosse, A., Taupier-Letage, I.,
  ... Herrmann, M. (2018). Characterizing, modelling and understanding the climate variability of the deep water formation in the NorthWestern Mediterranean Sea. *Climate Dynamics*, 51, 1179-1210. doi: 10.1007/s00382-016-3295-0
  - Straneo, F. (2006a). Heat and freshwater transport through the central Labrador Sea. Journal of Physical Oceanography, 36, 606-628. doi: 10.1175/JPO2875.1
- Straneo, F. (2006b). On the connection between dense water formation, overturn ing, and poleward heat transport in a convective basin. Journal of Physical
   Oceanography, 36, 1822-1840. doi: 10.1175/JPO2932.1
- Volkov, D. L. (2005). Interannual variability of the altimetry-derived eddy field
   and surface circulation in the extratropical North Atlantic Ocean in 1993-2001.
   *Journal of Physical Oceanography*, 35, 405-426. doi: 10.1175/JPO2683.1
  - Våge, K., Pickart, R. S., Sarafanov, A., Øyvind Knutsen, Mercier, H., Lherminier,
- P.,... Bacon, S. (2011). The Irminger Gyre: Circulation, convection, and
   interannual variability. Deep-Sea Research Part I: Oceanographic Research
   Papers, 58, 590-614. doi: 10.1016/j.dsr.2011.03.001
- Yashayaev, I., & Loder, J. W. (2017). Further intensification of deep convection in the Labrador Sea in 2016. *Geophysical Research Letters*, 44, 1429-1438. doi: 10 .1002/2016GL071668

# Supporting Information for "Restratification Structure and Processes in the Irminger Sea"

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1. Figure S1

# Introduction

To validate the time series from the CMEMS reanalysis, we compare them with time series from the LOCO mooring in Figure S1. Values of all properties are very similar in the CMEMS and LOCO time series. In both time series, the salinity has a mid-depth minimum, whereas the temperature and density have monotonously stratified structures. The MLD follows the same temporal trends in both data sets, with local maxima in 2008, 2009 and 2015 observable in both time series. Also, features of individual years such as the salinity and density maximum in the 2010–2011 winter are visible in both the CMEMS and the LOCO data. Thus, we conclude that the CMEMS reanalysis accurately captures the interannual variability of hydrographic properties in the Irminger Sea.

July 15, 2022, 8:45am



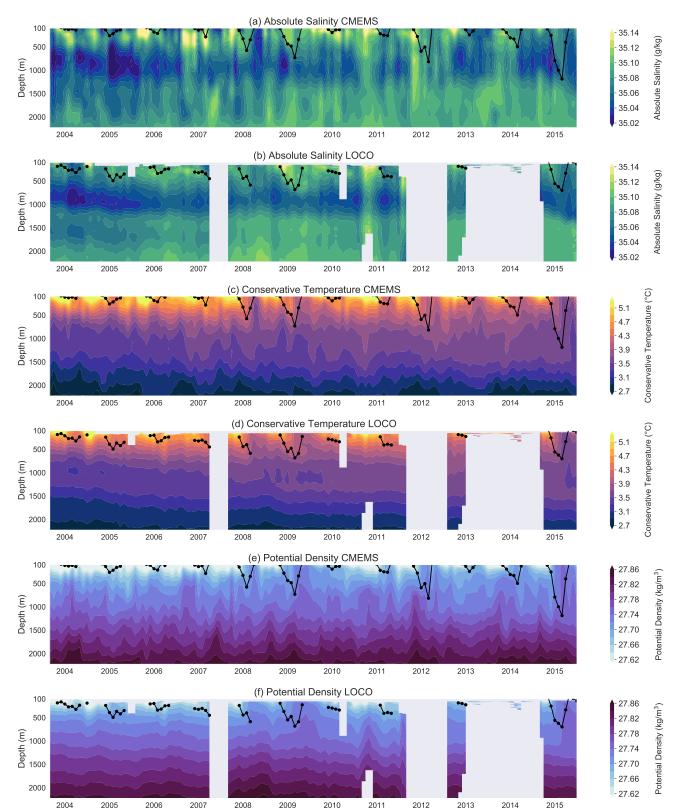


Figure S1. The temporal evolution of profiles of (a)-(b) absolute salinity, (c)-(d) conservative temperature and (e)-(f) potential density at the LOCO mooring. The mixed layer depth is shown in black. Panels (a), (c) and (e) show the data from the CMEMS reanalysis interpolated to the mooring location. Panels (b), (d) and (f) show the data from the LOCO mooring.

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