

SANS: Publicly Available Daily Multi-Scale Seismic Ambient Noise Source Maps

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Abstract

Seismic ambient noise sources have received increased attention recently, creating new possibilities to study the Earth's subsurface and the atmosphere-ocean-solid Earth coupling. In efforts to locate such noise sources using nonlinear finite-frequency inversions, methodological developments such as pre-computed wavefields and spatially variable grids were necessary. These make inversions feasible for the secondary microseismic sources in a frequency range up to 0.2 Hz on a daily basis. By obtaining a starting model for the inversion using Matched Field Processing (MFP) we are able to further steer the inversion towards acceptable global noise source models and improve the final result. Analysis of one year of daily inversions shows the seasonal variations of the secondary microseisms and their dependence on the atmosphere-ocean-solid Earth coupling due to storm-induced ocean waves. We present a web framework, SANS (Seismic Ambient Noise Sources, sans.ethz.ch), where daily regional- to global-scale seismic ambient noise source maps are made available to the public. This eases the implementation of time-variable noise source distributions into full-waveform ambient noise tomography and time-dependent subsurface monitoring methods. Additionally, it encourages other studies to verify if changes in the seismic data are due to changes in the subsurface velocity or spatio-temporal variations of noise sources.

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Key Points:

- We combine Matched Field Processing and nonlinear finite-frequency inversion to locate ambient noise sources
- We present a web framework for publicly available daily regional to global scale seismic ambient noise source maps (sans.ethz.ch)
- Analysis of a full year of daily inversions shows the seasonal variations of the secondary microseisms

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30 Plain Language Summary

31 The Earth is constantly vibrating due to various man-made and natural sources.
 32 One of the main sources of natural background noise is the ocean, specifically when ocean
 33 waves come into contact with the solid Earth. The strength of these sources strongly de-
 34 pends on the wave height, which changes constantly due to atmospheric phenomena such
 35 as storms. We study these seismic waves, so-called seismic ambient noise, to analyse the
 36 spatial and temporal variations which allow us to study the interaction between the at-
 37 mosphere, ocean, and solid Earth, as well as imaging the subsurface. In this study, we
 38 combine two different methods, namely Matched Field Processing (MFP) and nonlin-
 39 ear finite-frequency inversions, to create regional to global scale seismic ambient noise
 40 source maps on a daily basis. By looking at a full year of daily noise source maps we can
 41 observe the seasonal variations of noise sources. These daily noise source maps are pre-
 42 sented on a website (sans.ethz.ch) where anyone can download the results and imple-
 43 ment them in their own research. We hope that this will aid others by simplifying the
 44 implementation of noise source information which should make tomography and moni-
 45 toring methods more accurate.

46 1 Introduction

47 Seismic ambient noise sources have been studied thoroughly over the last few decades.
 48 Specifically, since studies showed that they could be used to study the Earth’s interior
 49 (Aki, 1957; Shapiro & Campillo, 2004; Shapiro et al., 2005; Sabra et al., 2005) further
 50 research was performed to understand the generation of ambient vibrations (e.g. Ard-
 51 huin, Stutzmann, et al., 2011; Arduin & Herbers, 2013; Arduin et al., 2015; Gualtieri
 52 et al., 2014, 2015), and new methods were developed to help locate these sources (e.g.
 53 Gerstoft & Tanimoto, 2007; Retailleau et al., 2017; Retailleau & Gualtieri, 2019; Gal et
 54 al., 2018; Sager, Ermert, et al., 2018; Ermert et al., 2020; Igel et al., 2021). More data-
 55 driven methods like correlation-based beamforming have been used to obtain the direc-
 56 tionality (e.g. Bucker, 1979; Hinich, 1979; Ruigrok et al., 2017) and physical location (e.g.
 57 Ishii et al., 2005; Meng et al., 2012; Retailleau et al., 2017; Retailleau & Gualtieri, 2019)
 58 of noise sources.

59 In theory, cross-correlation functions approach the Green’s functions for homoge-
 60 neously distributed, uncorrelated, random noise sources and an equipartioned wavefield.

61 Many ambient noise interferometry studies assume that the noise sources are sufficiently
62 homogeneous for the cross-correlations to converge to Green’s functions (e.g. Nakata et
63 al., 2019; Snieder & Wapenaar, 2010; Weaver et al., 2009; Wapenaar, 2004; Wapenaar
64 & Fokkema, 2006). However, several studies have shown that the omni-present ambient
65 noise wavefield changes on a daily basis (e.g. Bertelli, 1872; Longuet-Higgins, 1950; Ard-
66 huin et al., 2015) and the cross-correlation and Green’s function diverge if more realis-
67 tic constraints - such as global or local energy and directionality constraints - are im-
68 plemented into the modelling (Tsai & Sager, 2022). The heterogeneity of noise source
69 distributions can have a significant effect on travel times, particularly for monitoring ap-
70 plications (Zhan et al., 2013; Delaney et al., 2017). Prior knowledge of the noise source
71 locations can help distinguish if changes in the cross-correlations are due to changes in
72 the noise source distribution or subsurface velocities.

73 Inspired by work in helioseismology (Woodard, 1997), recent works introduced the
74 direct numerical modelling of noise cross-correlations for any heterogeneous noise source
75 distribution on Earth (Tromp et al., 2010; Hanasoge, 2013b; Fichtner, 2014; Ermert et
76 al., 2017; Sager, Ermert, et al., 2018; Datta et al., 2019). This has resulted in several stud-
77 ies using adjoint techniques (e.g. Fichtner et al., 2006) and sensitivity kernels (e.g. Tromp
78 et al., 2010; Fichtner, 2014) to invert for the seismic ambient noise source distribution
79 for different frequency ranges on various scales (Ermert et al., 2017; Xu et al., 2019; Igel
80 et al., 2021). Expanding on these adjoint and sensitivity kernel techniques, Bowden et
81 al. (2021) showed that certain beamforming algorithms are mathematically similar to
82 the first iteration of nonlinear finite-frequency inversions.

83 The direct forward-modelling of ambient noise cross-correlations allows us to cir-
84 cumvent common assumptions in ambient noise studies - e.g. wavefield equipartition-
85 ing and a quasi-random noise source distribution - that are necessary for Green’s func-
86 tion retrieval (e.g. Shapiro & Campillo, 2004; Shapiro et al., 2005; Wapenaar & Fokkema,
87 2006; Sánchez-Sesma & Campillo, 2006). Additionally, full-waveform ambient noise to-
88 mography methods are capable of directly implementing information about the noise source
89 distribution (Sager, Ermert, et al., 2018). Recent developments have made the compu-
90 tation of cross-correlations for ambient noise source inversions more efficient by using spa-
91 tially variable grids and pre-computed Green’s function databases (Ermert et al., 2020;
92 Igel et al., 2021); particularly for the frequency range of secondary microseismic sources
93 (between 0.1 and 0.2 Hz). This allows us to rapidly invert for the noise source distribu-
94 tion on a regional to global scale with reasonable computational cost by taking advan-
95 tage of high performance computing (HPC) resources.

96 Building on these various developments, we present a web framework to make daily
97 Seismic Ambient Noise Source (SANS) maps available to the public (sans.ethz.ch). In-
98 depth knowledge of the ambient noise source distribution should help to improve am-
99 bient noise tomography and imaging methods; particularly to ensure that changes in the
100 subsurface are not confused with the spatio-temporal variations of the microseismic noise
101 source distribution. Providing these maps should ease the implementation in full wave-
102 form ambient noise tomography methods and encourage future studies to take the noise
103 source distribution into account.

104 In this paper, we focus on the combination of Matched Field Processing (MFP) with
105 nonlinear finite-frequency inversions to improve the inversion results, and present the web
106 framework SANS, where the daily ambient noise maps are made publicly available. Since
107 both methods have previously been described individually in detail, we refer the inter-
108 ested reader to earlier publications for more in-depth derivations and explanations (Bowden
109 et al., 2021; Ermert et al., 2020; Igel et al., 2021).

2 Methodology

In the following section, we will explain the main steps of the two methods: nonlinear finite-frequency inversions and Matched Field Processing (MFP). Despite the differences in the approaches taken, Bowden et al. (2021) show that these methods are well connected. Both have their advantages and disadvantages: MFP is an efficient, data-driven approach that works on any cross-correlation data. An inversion is computationally more expensive but - in contrast to MFP - models the wave propagation more accurately and allows us to account for the nonlinearity by using an iterative approach.

More importantly, an inversion allows for prior knowledge to be implemented. Hence we use the more efficient, data-driven MFP algorithm to compute a starting model for the nonlinear finite-frequency inversion, to avoid local minima and accelerate the convergence towards an acceptable model. Both methods rely on the fact that vertical-component seismic ambient noise data in the frequency range of 0.1 to 0.2 Hz are dominated by surface waves.

2.1 Nonlinear Finite-Frequency Inversion

The inversion method is based on a concept from helioseismology (Woodard, 1997) which enables the direct modelling of cross-correlations for any noise source power-spectral density (PSD). The work has been adapted for applications to Earth by several authors (e.g. Tromp et al., 2010; Hanasoge, 2013a; Fichtner, 2014; Ermert et al., 2017; Sager, Ermert, et al., 2018) with some additional implementations of pre-computed wavefields (Ermert et al., 2020) and spatially variable grids (Igel et al., 2021) to improve efficiency and make inversions feasible for higher frequencies. In the following section we will provide a short overview of the gradient-based iterative inversion method. For more details, the reader is referred to the aforementioned publications.

2.1.1 Cross-correlation modelling

The following equation allows us to forward model the cross-correlation wavefield \mathcal{C}_{ij} , for two stations at locations \mathbf{x}_1 and \mathbf{x}_2 , for an arbitrary noise source PSD S_{nm} , at points $\boldsymbol{\xi}$ on the Earth's surface $\partial\oplus$, using the Green's functions \mathbf{G} , in the frequency domain (Fichtner, 2014; Ermert et al., 2017; Sager, Boehm, et al., 2018; Igel et al., 2021):

$$\mathcal{C}_{ij}(\mathbf{x}_1, \mathbf{x}_2) = \int_{\partial\oplus} G_{in}^*(\mathbf{x}_1, \boldsymbol{\xi}) G_{jm}(\mathbf{x}_2, \boldsymbol{\xi}) S_{nm}(\boldsymbol{\xi}) d\boldsymbol{\xi}. \quad (1)$$

We imply the Einstein summation convention for repeated indices, and * indicates the complex conjugate. To reduce the computational cost we pre-compute the Green's functions \mathbf{G} using the time-domain spectral-element codes for spherically symmetric Earth models AxiSEM to model wave propagation (Nissen-Meyer et al., 2014), and Instaseis to extract seismograms (van Driel et al., 2015) with 1-D isotropic PREM (Dziewonski & Anderson, 1981) as underlying velocity model. The Green's function database is then re-used for subsequent iterations and inversions. Additionally, we implement spatially variable grids with regional dense areas and no grid points on land, to reduce the number of possible noise sources, and thus the modelling parameters for regional applications (Igel et al., 2021). The combination of pre-computed wavefields and spatially variable grids allows us to efficiently invert for the noise source distribution of the secondary microseisms in a frequency range of 0.1 to 0.2 Hz (Ermert et al., 2020; Igel et al., 2021).

2.1.2 Inversion

Once we have modelled the synthetic cross-correlations, we measure the difference to the observed cross-correlations using the logarithmic energy ratio (e.g. Ermert et al.,

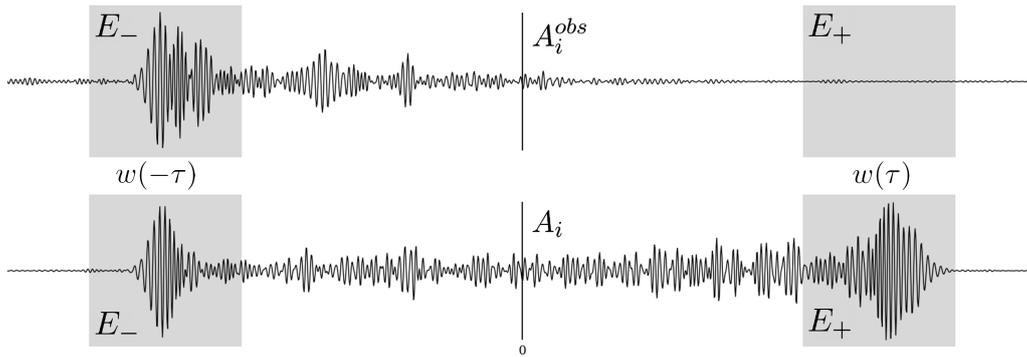


Figure 1. Our chosen measurement is the logarithmic energy ratio as previously used by Ermert et al. (2020) and Igel et al. (2021) (see Equation 2). To obtain a misfit we take a measurement of synthetic A_i (bottom) and observed A_i^{obs} (top) cross-correlations and compute the L_2 -norm (see Equation 3).

154 2017). This measurement quantifies the asymmetry of the cross-correlation which arises
 155 from a heterogeneous noise source distribution as illustrated in Figure 1. The logarithmic
 156 energy ratio computes the ratio of the energies E_+ and E_- of the expected surface
 157 wave arrival window $w(\tau)$ in the causal and acausal parts of the cross-correlation $C(\tau)$:

$$A = \ln \left(\frac{\int [w(\tau)C(\tau)]^2 d\tau}{\int [w(-\tau)C(\tau)]^2 d\tau} \right) = \ln \left(\frac{E_+}{E_-} \right) \quad (2)$$

In contrast to full-waveform misfits, the logarithmic energy ratio aims to match the energy in the causal and acausal expected surface wave arrival windows. Although this measurement contains less information, it is much more robust, and relatively insensitive to unknown 3-D Earth structure (Sager, Boehm, et al., 2018) due to only comparing energies in a certain time window. Consequently, this allows us to use a simple 1-D PREM (Dziewonski & Anderson, 1981) velocity model to compute synthetic cross-correlations. During the inversion we aim to minimise the squared L_2 -norm, i.e. the misfit χ , of the measurements A_i and A_i^{obs} on the synthetic and observed cross-correlations, respectively:

$$\chi = \frac{1}{2} \sum_{i=1}^N [A_i - A_i^{obs}]^2 \quad (3)$$

158 where N is the number of measurements.

159 Adjoint techniques (e.g. Fichtner et al., 2006) allow us to compute source sensi-
 160 tivity kernels (e.g. Tromp et al., 2010; Hanasoge, 2013b; Fichtner, 2014) which provide
 161 a spatial reference of where an increase or decrease in noise source strength should de-
 162 crease the misfit. By compiling the gradient, i.e. the sum of all sensitivity kernels, we
 163 can update the noise source distribution and continue with the next iteration by re-computing
 164 the cross-correlations, misfits, and sensitivity kernels. To minimise the misfit we adopt
 165 a gradient-based iterative scheme using the steepest descent method, including regular-
 166 isation and step-length tests. Several synthetic and real-data tests have shown that there
 167 are usually no significant improvements in the noise source distribution after roughly 5
 168 iterations. Hence, we run 8 iterations of the inversion to ensure that we have converged
 169 to a model that explains the data based on our measurement.

170 In previous research (Igel et al., 2021), we used a homogeneous distribution in the
 171 ocean as the initial noise source distribution. For a gradient-based iterative inversion method

172 like ours, a good initial model can be helpful in steering the inversion towards an accept-
 173 able global noise source model and avoid local minima.

174 By introducing a different method to locate noise sources - namely Matched Field
 175 Processing - we are able to efficiently create a more realistic initial model from the same
 176 observed cross-correlations. This is similar to full-waveform inversions, where starting
 177 models are often constructed with more efficient methods such as ray-based travel time
 178 tomography or dispersion curve analysis (e.g. Virieux & Operto, 2009; Teodor et al., 2021).

179 **2.2 Matched Field Processing**

180 Matched Field Processing, in this context, may be considered similar to beamform-
 181 ing and backprojection methods, where time-shifts are applied to the data and rays are
 182 backprojected to obtain a source location. However, whereas beamforming generally as-
 183 sumes plane waves arrive at an array of sensors, MFP directly considers sources anywhere
 184 within a computational domain and estimates travel times accordingly. This makes it
 185 very suitable for global ambient noise source studies, where stations from all over the globe
 186 may be used. Additionally, it is able to map noise sources on any grid, which allows us
 187 to use the same source grid for MFP and the inversion.

188 MFP algorithms of varying complexity have been developed, for example: to lo-
 189 cate hydrothermal acoustic sources (Cros et al., 2011); microseismic sources in exploration
 190 geophysics (Corciulo et al., 2012); glacial tremors (Umlauf et al., 2021); or applied to
 191 three-component seismic array data for microseisms (Gal et al., 2018). The algorithm
 192 could also be adapted to be nearly identical to full-waveform methods by including syn-
 193 thetic Green’s functions (Bowden et al., 2021; Schippkus & Hadziioannou, 2022). Although
 194 there may be some value to more complex MFP implementations, we prefer the com-
 195 putationally efficient version described below, as the subsequent inversion iterations will
 196 add further complexity.

197 **2.2.1 Constant Velocity MFP**

Our MFP algorithm is based on the assumption that a point source at a proposed
 noise source location for a set surface wave group velocity will lead to signal in the cross-
 correlation at a certain lag. The first step is to compute all cross-correlations and cre-
 ate a grid of possible noise sources. Subsequently, we iterate over all possible noise sources
 and compute the travel times to the stations based on a constant surface wave speed v
 of $2,900 \frac{m}{s}$. This surface wave speed is roughly the average Rayleigh wave group veloc-
 ity in the 0.1 to 0.2 Hz frequency range in PREM (Dziewonski & Anderson, 1981) and
 has provided good results for synthetic and real-data applications. For such a narrow
 frequency band we can consider the group velocity to be roughly constant. The travel
 time difference Δt_{ij} between arrivals t_i and t_j at the two receiver locations \mathbf{x}_i and \mathbf{x}_j
 determines the lag at which the current noise source location \mathbf{x} would result in a signal
 in the cross-correlation:

$$\Delta t_{ij} = t_i - t_j = \frac{\|\mathbf{x} - \mathbf{x}_i\|}{v} - \frac{\|\mathbf{x} - \mathbf{x}_j\|}{v}. \quad (4)$$

198 Note that $\|\cdot\|$ denotes the vector norm for a 2-D example as illustrated in Figure 2. For
 199 our applications we extend the vector norm from 2-D to a sphere where it is adapted to
 200 be the great circle distance between the noise source locations and stations.

201 Finally, the corresponding value of the cross-correlation - or in our case the value
 202 of the square envelope of the cross-correlation as explained in section 2.2.2 - is added to
 203 the ‘power’ of that grid point, and we repeat the process for the next possible noise source
 204 location. This is equivalent to applying phase shifts to the raw signals and then mea-
 205 suring coherencies, as MFP or beamforming is often described. An illustration of this

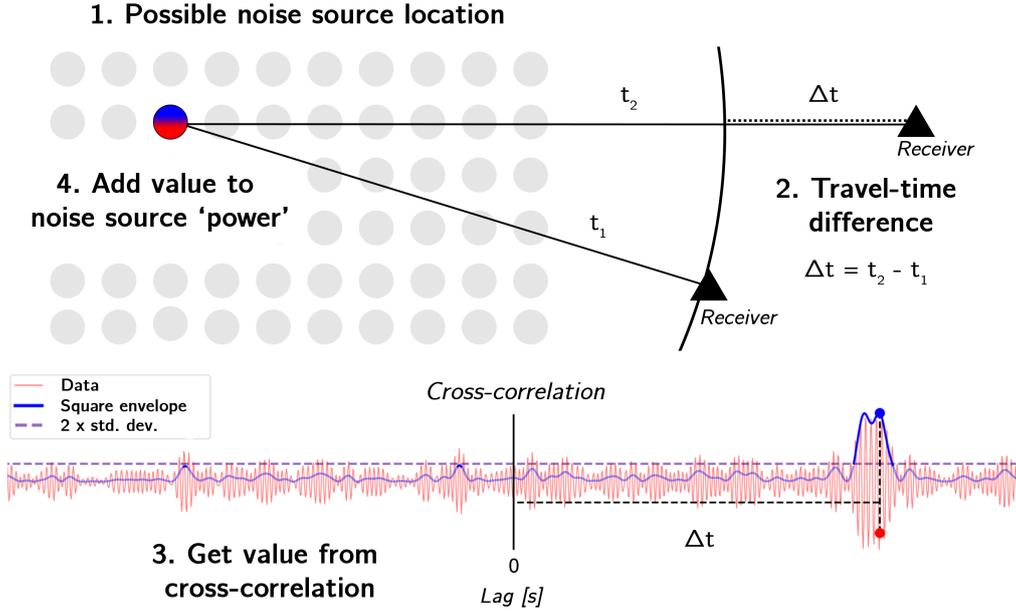


Figure 2. Illustration of the Matched Field Processing algorithm. We iterate over all possible noise sources, calculate the travel time difference Δt and finally obtain the value from the cross-correlation, e.g. red dot for the actual waveform value or blue dot for the value of the square envelope. Additionally, we calculate the standard deviation of the envelope and set everything below twice the standard deviation to 0. This increases the signal-to-noise ratio and improves the final MFP power map. This process is repeated for every possible noise source location and every station pair.

206 algorithm can be seen in Figure 2. By using a larger array of stations, we are able to spa-
 207 tially restrict the locations of noise sources and obtain a map of noise source 'power'.

208 *2.2.2 Square Envelope Measurement*

209 The simplest MFP method uses the value of the cross-correlation waveform to ob-
 210 tain the 'power' for each noise source location. However, since this often results in strong
 211 fluctuations of the noise source power due to the oscillatory nature of the waveforms and
 212 struggles with low signal-to-noise ratios we instead take the value of the square envelope
 213 $S(C(\tau))$ of the cross-correlation $C(\tau)$.

$$S(C(\tau)) = C^2(\tau) + \mathcal{H}(C(\tau))^2, \quad (5)$$

214 where \mathcal{H} indicates the Hilbert transform. Additionally, we compute the standard devi-
 215 ation $\sigma(S(C(\tau)))$ of the square envelope and set all values below twice the standard de-
 216 viation to 0, i.e. we do not add any 'power' for those noise source locations.

217 Besides increasing the signal-to-noise ratio when a signal is present, this also smooths
 218 the resulting noise source and avoids the fluctuations of noise source power. This does
 219 mean that for cross-correlations with no clear signal, i.e. where the square envelope is

220 nearly constant, nothing is removed and the signal-to-noise ratio can not be increased.
 221 This has little effect on the final MFP power distribution as it would add a near con-
 222 stant value.

$$P(C(\tau_i)) = \begin{cases} 0 & \text{if } S(C(\tau_i)) < 2 * \sigma(S(C(\tau))) \\ S(C(\tau_i)) & \text{else,} \end{cases} \quad (6)$$

223 where $P(C(\tau_i))$ is the MFP power for the time lag τ at source location i and cross-correlation
 224 $C(\tau_i)$. Synthetic and real data tests show that using the square envelope with a cut-off
 225 threshold greatly increases the contrast of the final MFP maps and ensures that we mainly
 226 use the signal from the cross-correlations.

227 To account for geometric spreading we multiply each value of the square envelope
 228 of the cross-correlations with an amplitude decay factor D_i as introduced by previous
 229 studies (e.g. Corciulo et al., 2012; Bowden et al., 2021), which depends on the surface
 230 wave group velocity v , the average frequency of our bandpass filter f , and the average
 231 distance of the station pair to the proposed source location r_i :

$$D_i = \sqrt{\frac{2v}{\pi f r_i}} \quad (7)$$

232 This process is repeated for all possible noise source locations and cross-correlations,
 233 and the values of the square envelope of the cross-correlations are added up as illustrated
 234 in Figure 2. Of course, more sophisticated methods to model either the travel times or
 235 amplitude decays and attenuation can be implemented in MFP (Bowden et al., 2021;
 236 Schippkus & Hadziioannou, 2022). Such modelling is precisely the point of subsequent
 237 iterations of the full-waveform approach, whereas the MFP is only intended to give a com-
 238 putationally efficient initial model.

239 In contrast to other array-based beamforming methods, MFP works best when the
 240 stations surround the noise source location. We illustrate this in Figure 3 by running a
 241 synthetic example, where we forward model cross-correlations using the pre-computed
 242 Green’s function database and cross-correlation model code from the inversion with a
 243 dominant noise source blob within the domain, and a frequency content of 0.1 to 0.2 Hz.
 244 We apply the MFP algorithm to two sets of stations: 6 stations in closer proximity and
 245 35 stations spread out in the whole domain. If the dominant noise source is outside the
 246 array we see strong smearing and MFP is only able to give us a direction of the dom-
 247 inant noise source. On the other hand, if the dominant noise source is surrounded by sta-
 248 tions, MFP is able to constrain the spatial extent of the dominant noise sources more
 249 accurately.

250 2.3 MFP starting model

251 Thanks to MFP being most capable when the dominant noise sources are within
 252 the array, it is a useful method to locate noise sources on a regional to global scale. MFP
 253 and the finite-frequency inversion use slightly different information from the cross-correlation
 254 to obtain a noise source distribution. The logarithmic energy ratio is largely insensitive
 255 to unknown Earth structure, as it only takes the energy in a given window but ignores
 256 the actual waveform. On the other hand, MFP with the square envelope measurement
 257 uses more information from the waveform itself but does not properly account for wave
 258 propagation. Additionally, the resulting MFP maps are harder to interpret in terms of
 259 physical units as they are not an actual model of a physical quantity but rather an im-
 260 age of the noise source distribution.

261 To combine the two methods we normalise a smoothed MFP noise source map and
 262 set it as the initial power-spectral density model for the finite-frequency inversion. In con-

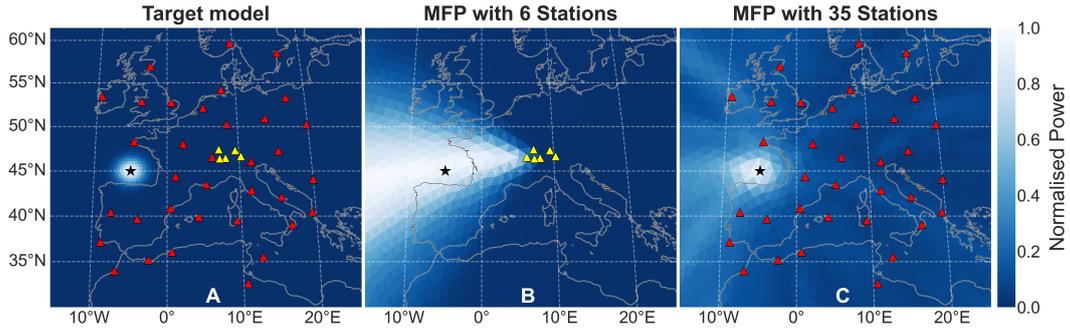


Figure 3. Example of the MFP algorithm for synthetic cross-correlations modelled with a dominant noise source blob (\star) within the domain (A). If the dominant noise source is outside the array (\blacktriangle), MFP is mainly able to provide a direction (B). If the stations surround the dominant noise source (\blacktriangle), MFP is able to constrain the spatial extent of the dominant noise source much more accurately (C).

263 contrast to the previously used homogeneous starting model, this greatly reduces the presence of inversion artefacts. Synthetic and real-data tests have shown that this improves the final noise source maps without significantly increasing the computational cost. Figure 4 shows a regional synthetic comparison of two inversions with a homogeneous and an MFP starting model. The synthetic cross-correlations were modelled using the noise source distribution on the left with added random noise to make them more realistic. The random noise is introduced by normalising a random time series, multiplying it by the maximum amplitude of the cross-correlation and a scaling factor of 1.5, and finally adding it to the cross-correlation. Comparisons show that this resembles our real ambient noise cross-correlations more closely.

273 The inversion with the homogeneous starting model does contain the most dominant noise sources but shows a strong tendency to move noise sources closer to the coast, especially for the large dominant noise source area off the European coast. In contrast, the inversion with an MFP starting model does not lead to strong coastal sources, and better represents the spatial distribution of the dominant noise sources in the actual model. This is particularly useful for global inversions where MFP can help to avoid inversion artefacts due to lack of data by increasing the probability of noise sources in certain areas before the first iteration.

281 MFP introduces new information to the inversion, as it actually uses the cross-correlation waveforms, as opposed to the finite-frequency inversion where we measure the energy in the expected surface wave arrival time windows. Hence, we expect this to reduce the null space of the inversion and produce a more accurate noise source map. Despite the clear differences in the resulting inversion models, the misfits of the final iterations shown in Figure 4 are very similar. However, it is clear that the inversion with the MFP starting model is visually more similar to the target model than the inversion with a homogeneous starting model. This indicates that including the additional waveform information via the MFP starting model does reduce the null space and helps to steer the inversion in a direction that is more closely aligned with the actual noise source distribution.

291 3 SANS: Daily Seismic Ambient Noise Sources

292 In light of the recent developments that have significantly decreased the computational cost of ambient noise source inversions for the secondary microseisms on a global scale (Ermert et al., 2020; Igel et al., 2021), we introduce a new web framework, SANS,

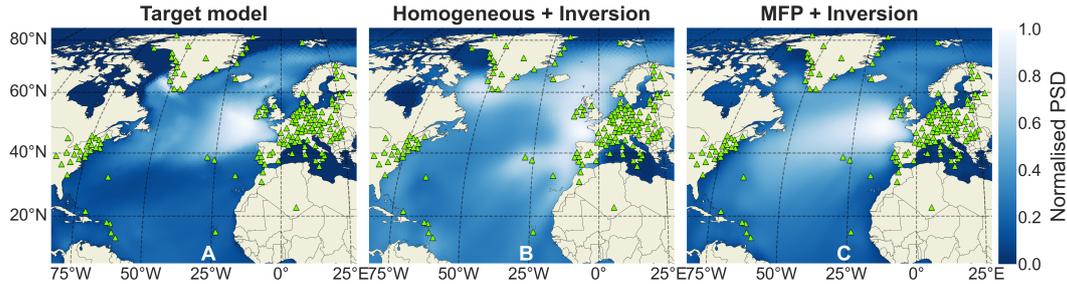


Figure 4. Synthetic inversions using cross-correlations with added random noise that were forward-modelled with the noise source distribution on the left (A) for stations surrounding the North Atlantic (\blacktriangle). A homogeneous starting model introduces stronger noise sources along the coast during the inversion (B). In contrast, an inversion with an MFP starting model results in a noise source distribution that is closer to the target model (C).

295 where daily Seismic Ambient Noise Sources are made available to the public (`sans.ethz`
 296 `.ch`). Currently, we run two inversions every day: one for a global station list and one
 297 more regional with stations surrounding the North Atlantic. A regional inversion allows
 298 for a higher spatial resolution of the noise source distribution in that area.

299 Users can obtain the inversion results by either directly looking at a plot of the noise
 300 source distribution maps online or downloading the inversion output and analysing it them-
 301 selves, e.g. for implementation in other studies.

302 3.1 Data selection and processing

303 We download and process the seismic ambient noise data automatically every morn-
 304 ing at 4 am (CET) using ObsPy (Krischer et al., 2015). All stations within a chosen sta-
 305 tion list are checked for available data. The station lists are based on globally available
 306 broadband sensors but limit the minimum distance between stations to roughly 1° (=
 307 111 km) to avoid smaller arrays. Dense station arrays would lead to high local sensitiv-
 308 ities that would distort the final noise source distribution. This results in 414 stations
 309 for the global and 153 stations for the North Atlantic station list. The global distribu-
 310 tion of stations is illustrated in the resolution analysis in Figure 5 and both station lists
 311 can be downloaded from the website. The data availability changes on a daily basis, lead-
 312 ing to roughly 70% of these stations having data available on average.

313 After downloading all available data, we remove the instrument response, down-
 314 sample to 1 Hz, segment the data into 2 h windows, and remove any windows contain-
 315 ing earthquakes that are in the GCMT catalogue (Ekström et al., 2012) with a minimum
 316 magnitude of 5.6. Occasionally this can lead to all windows being removed if there was
 317 one strong or several smaller earthquakes in a day. Subsequently, we compute the daily
 318 cross-correlations of the windowed seismic ambient noise data by stacking the individ-
 319 ual cross-correlations of the 12 windows. This helps to increase the signal-to-noise ra-
 320 tio of the final daily ambient noise cross-correlations.

321 Similar to Igel et al. (2021), we ignore cross-correlations with a signal-to-noise ratio
 322 below 3.5. The signal-to-noise ratio is determined by dividing the maximum ampli-
 323 tude within the expected surface wave arrival window by the standard deviation of the
 324 whole time series. Hence, we define a signal as a clear surface wave arrival within the
 325 expected window. This is usually the case if the dominant noise source is in-line with
 326 the station pair. Due to our chosen measurement being the logarithmic energy ratio, cross-

327 correlations with little signal - i.e. asymmetry - would not contribute much to the final
 328 gradient and update of the noise source distribution. Besides improving the final result
 329 of the inversion, ignoring cross-correlations below a signal-to-noise threshold also decreases
 330 the computational cost, as fewer cross-correlations have to be modelled during the in-
 331 version. During the inversion we apply a band-pass filter between 0.1 and 0.2 Hz as we
 332 focus on the secondary microseismic noise sources.

333 3.2 Web framework

334 After the data have been downloaded, processed, and correlated we run 8 iterations
 335 of the inversion on Piz Daint, a supercomputer at the Swiss National Supercomputing
 336 Centre (CSCS). The computational cost of the inversions varies with the number of avail-
 337 able cross-correlations for each day. However, we greatly reduce the computational cost
 338 since we have already pre-computed the wavefield and extracted the Green’s function
 339 database which is re-used every day. We run both inversions, one global and one regional
 340 surrounding the North Atlantic, on 600 cores, with the usual computational times be-
 341 ing 60 minutes (50 node hours) and 30 minutes (25 node hours), respectively. We use
 342 two different spatially variable grids, with a more homogeneous distribution of about 29,000
 343 grid points for the global inversion and a locally dense grid in the North Atlantic with
 344 roughly 21,000 grid points for the regional inversion. Once the inversions are done, we
 345 plot the output and copy all relevant files to the ETH web server where the website is
 346 hosted. These files are then made available to the public on sans.ethz.ch.

347 The web framework allows users to look through the iterations of all available in-
 348 version results and compare them to significant wave height maps (Tolman & Chalikov,
 349 1996; *WAVEWATCH III*, 2005) of that day. Note that the generation mechanism of the
 350 secondary microseisms requires ocean waves travelling in opposite directions to overlap
 351 (Nakata et al., 2019); therefore the wave height maps are merely a reference as to where
 352 the areas of dominant noise sources may be and should not be directly compared. Users
 353 can download the full inversion output folder including the parameter file, station list,
 354 source grid, further plots such as the gradients, misfit reduction and other relevant files.
 355 We provide code that helps a user to plot and analyse these results themselves. In the-
 356 ory, the inversions are reproducible as the inversion code is made available on github.
 357 However, this does require the additional computation or download of an AxiSEM wave-
 358 field and access to HPC facilities. The global inversion requires roughly 50 node hours
 359 which includes the data download, processing, and 8 iterations of the inversion but ex-
 360 cludes the extraction of a Green’s function database.

361 3.3 Resolution analysis

362 Recent efforts have estimated the resolution and covariance of noise source full wave-
 363 form inversions by treating it as a linear problem and using singular value decomposi-
 364 tion (Xu & Mikesell, 2022). However, due to our nonlinear measurement of the logarith-
 365 mic energy ratio this is not applicable to our inversion method. To show the effect of the
 366 changing data and station availability on the resolution of the inversions, we forward model
 367 cross-correlations with added noise for 414 stations around the globe and perform inver-
 368 sions with different station lists in Figure 5. The noise source distribution that we use
 369 to forward-model the data is an adapted significant wave height map from the WaveWatch
 370 III model (Tolman & Chalikov, 1996; *WAVEWATCH III*, 2005). The inversions are run
 371 with the same parameters as the daily SANS inversions. We choose different station lists
 372 from daily inversions to give a realistic station distribution that would be used for real-
 373 data inversions.

374 The inversion with 414 stations shows the model that we are able to recover us-
 375 ing all potential stations. Due to the much higher station density in the Northern Hemi-
 376 sphere the resolution is higher and we are able to recover the dominant noise sources more

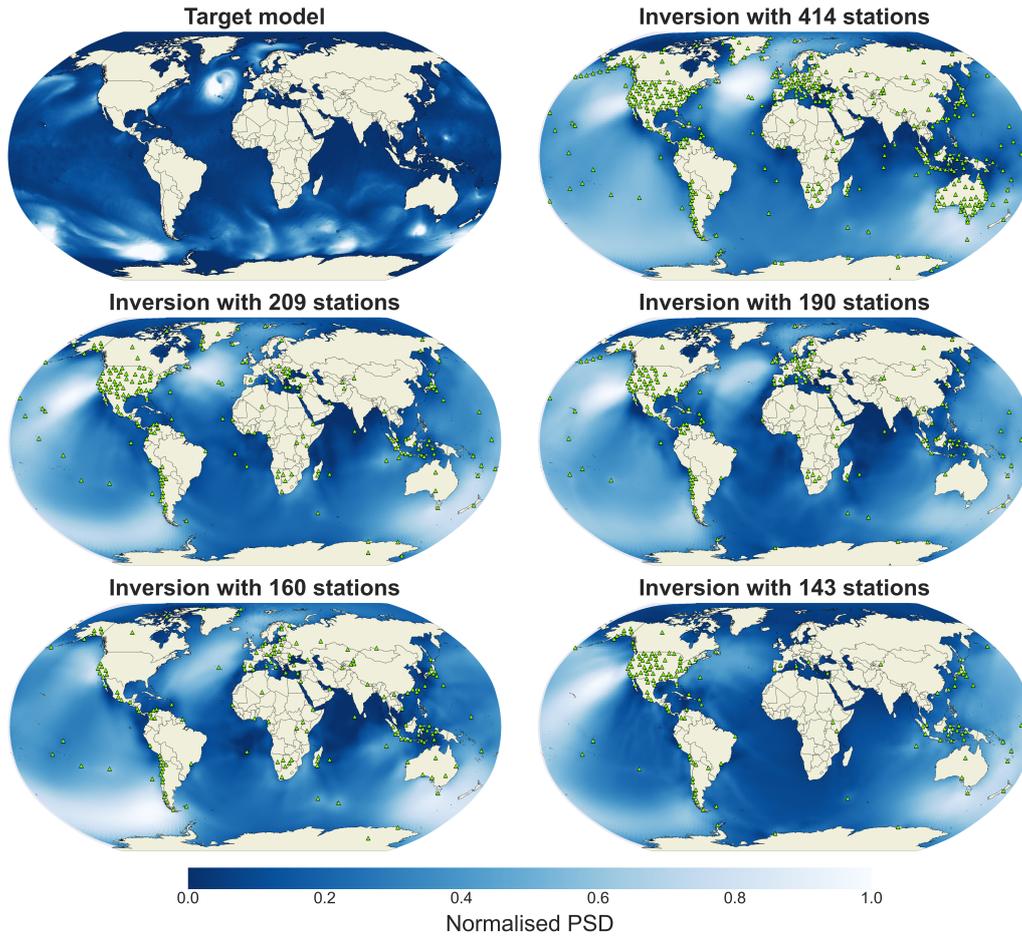


Figure 5. Synthetic inversions for the target model to show the effect of different stations (\blacktriangle) on the final inversion models. Inversions with fewer stations still capture the most dominant noise sources in the final model.

377 accurately than in the Southern Hemisphere. As we decrease the number of stations we
 378 can see how the recovered model changes, especially when the stations are predominantly
 379 in Europe (160 stations) or North America (143 stations). However, even the inversions
 380 with fewer stations still include the most dominant noise sources from the target model.
 381 In that sense, the daily global inversions should not necessarily be seen as the global noise
 382 source distribution for each day, but rather the noise source distribution that the given
 383 station list is able to observe. Generally, the resolution in the Southern Hemisphere is
 384 lower due to the lack of station coverage and the North Atlantic usually has the high-
 385 est resolution since it is surrounded by stations in Europe and North America.

386 3.4 Example applications

387 In the following section we present two example applications of the daily SANS maps.
 388 Firstly, we take the average of the daily inversions for Northern Hemisphere summer and
 389 winter to study the seasonal variations of the secondary microseisms. Secondly, we model
 390 cross-correlations for different noise source distribution models to illustrate the effect of
 391 a changing noise source distribution on the cross-correlation waveforms.

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3.4.1 Seasonal analysis

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Secondary microseismic sources are generated when two oceans travelling in opposite directions overlap, which in turn creates a vertical pressure wave. This induces seismic waves at the ocean bottom. The strength of these sources is directly related to the wave height of the overlapping waves (Longuet-Higgins, 1950; Hasselmann, 1963; Nakata et al., 2019; Ardhuin, Hanafin, et al., 2011).

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Due to the seasonal variations in significant wave height we would expect a similar pattern for the noise source distribution of the secondary microseisms, which has already been observed over a century ago (Klotz, 1910; Burbank, 1912; Banerji, 1925). This relationship has recently been studied for various different frequency bands of microseisms ranging from the hum to secondary microseismic sources (Nishida & Fukao, 2007; Rhie & Romanowicz, 2006; Ermert et al., 2017; Stutzmann et al., 2012; Landés et al., 2010; Gualtieri et al., 2021). Thus, we would expect similar patterns to emerge if we average the daily inversions generated by the SANS workflow.

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We include 335 daily inversions from the 17th May 2021 to the 2nd May 2022 in the analysis and choose to define Northern Hemisphere summer (21st April to 21st October) and winter (21st October to 21st April) based on the Icelandic first day of summer in 2022. The final iterations of all inversions within those two time ranges are averaged, resulting in 164 inversions for the summer and 171 for the winter months. Before averaging, we smooth the noise source model with a 4° Gaussian smoothing filter to avoid any artefacts from small changes in the inversion parameters during that time period.

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Similarly, we average the significant wave height maps from the WaveWatch III model (Tolman & Chalikov, 1996) as a comparison. This should not be taken as a direct comparison but more as a reference of where the probability of more dominant noise sources is higher. The actual mechanism of generation of secondary microseismic sources is more complicated and requires more complex modelling (Ardhuin, Stutzmann, et al., 2011; Ardhuin & Herbers, 2013; Nakata et al., 2019). Figure 6 shows the comparison of the normalised average significant wave height with the normalised PSD of the average SANS inversions for Northern Hemisphere summer and winter.

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The average inversions show clear seasonal variations that are in-line with our expectations. During the Northern Hemisphere summer the more dominant noise sources are in the Southern Hemisphere, specifically the South Pacific. As shown by previous studies on seasonal noise source variations (Landés et al., 2010; Gualtieri et al., 2021; Stutzmann et al., 2012), the Northern Hemisphere winter is dominated by noise sources in the North Atlantic. This supports the result of previous studies and shows that the SANS inversions are able to observe the spatio-temporal variations of the secondary microseismic sources on various timescales.

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3.4.2 Cross-correlation modelling

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A changing noise source distribution has a significant effect on the cross-correlations, particularly on a global scale. A common assumption is a homogeneous noise source distribution which, in theory, results in a symmetric cross-correlation. However, the noise source distribution is often strongly heterogeneous and changes constantly. We forward model cross-correlations using our modelling code from the inversion for three different noise source distributions to illustrate the changes: (i) homogeneous distribution everywhere, (ii) homogeneous distribution in the ocean, and (iii) the final SANS inversion model for the 9th March 2022.

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In Figure 7 we plot the cross-correlations for 6 station pairs and the three different models. The cross-correlations are filtered between 0.1 and 0.2 Hz. As the noise source

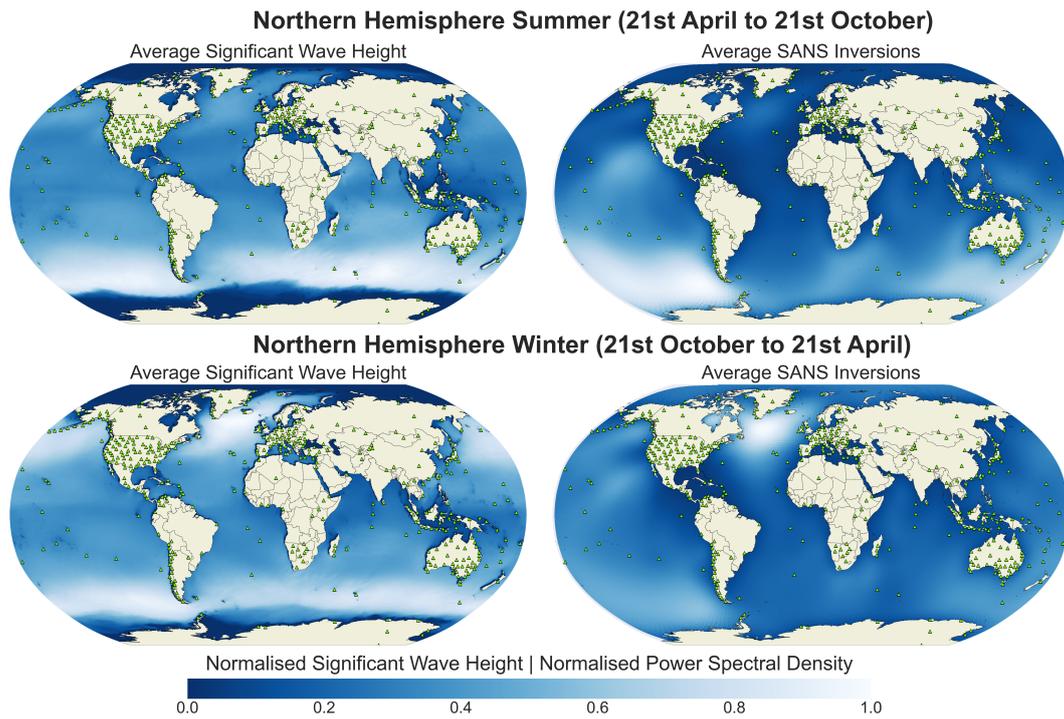


Figure 6. Comparison of the normalised average significant wave height and normalised PSD of the SANS daily inversion results for Northern Hemisphere summer (21st April to 21st Oct) and winter (21st Oct to 21st April) using a global station distribution (\blacktriangle). Northern Hemisphere summer is dominated by sources in the Southern Hemisphere and Northern Hemisphere winter is dominated by sources in the North Atlantic.

441 distribution becomes more realistic, the changes to the cross-correlations become more
 442 and more significant. Especially for full waveform ambient noise studies, the influence
 443 of a changing noise source distribution should not be ignored. In many cases the main
 444 arrival also shifts significantly which makes travel time picking more difficult. We en-
 445 courage future ambient noise studies to consider including information about the noise
 446 source distribution.

447 4 Discussion

448 With the daily computation of seismic ambient noise source maps we are able to
 449 study the interaction between the atmosphere, the ocean, and the solid Earth in near
 450 real-time. The daily maps show the clear heterogeneous nature of secondary microseis-
 451 mic noise sources and their strong spatio-temporal variations. Since the generation mech-
 452 anism of secondary microseismic sources is quite well understood (Longuet-Higgins, 1950;
 453 Ardhuin, Stutzmann, et al., 2011; Ardhuin & Herbers, 2013; Ardhuin et al., 2015), these
 454 variations can also be studied by computing ocean surface pressure maps (Ardhuin, Stutz-
 455 mann, et al., 2011) using significant wave height and bathymetry data. A comparison
 456 in our previous research (Igel et al., 2021) shows that these ocean surface pressure maps
 457 and finite-frequency inversions coincide quite well with similar areas of dominant noise
 458 sources present in both. Similarly, we observe correlation between the location of dom-
 459 inant noise sources in our daily SANS maps with respect to the higher amplitude sig-
 460 nificant wave heights. Due to the generation mechanism requiring two overlapping waves
 461 travelling in opposite directions, this comparison should only be considered as a rough
 462 reference of where there is a higher probability of dominant noise sources.

463 Particularly for ambient noise tomography and monitoring, knowledge of the noise
 464 source distribution is vital to circumvent common assumptions like the quasi-randomness
 465 of the noise field and equipartitioning of the wavefield. For these methods, daily maps
 466 can help reduce the misinterpretation of noise distribution changes as subsurface veloc-
 467 ity changes. (Sager, Ermert, et al., 2018) inverted for both the noise source distribution
 468 and subsurface structure at the same time. However, this comes at an increased com-
 469 putational cost. By already having knowledge of the noise source distribution beforehand,
 470 we can reduce the complexity and computational cost of such full-waveform ambient noise
 471 tomography methods. As we illustrate in Figure 7, a heterogeneous noise source distri-
 472 bution can have a significant effect on the cross-correlations which should not be neglected,
 473 especially in full waveform ambient noise studies.

474 To make our inversion process as efficient as possible we use a simple 1-D PREM
 475 Earth model to simulate the Green’s functions and cross-correlations. Despite (Sager,
 476 Boehm, et al., 2018) showing that our measurement of the logarithmic energy ratio is
 477 largely insensitive to unknown 3-D Earth structure, this simplification could have an
 478 effect on the inversion. However, seismic studies within our frequency range of 0.1 to 0.2
 479 Hz are generally considered less sensitive to small heterogeneities in the crust. Future
 480 studies might incorporate more complex Earth models (e.g. Fichtner et al., 2018) by pre-
 481 computing the Green’s function database using a wavefield solver like Salvus (Afanasiev
 482 et al., 2019). This would also allow the implementation of a fluid ocean layer and 3-D
 483 structure, albeit at the cost of increased computation time.

484 Furthermore, since the availability of ambient noise data changes daily, the num-
 485 ber of stations included in the daily inversions can fluctuate greatly. This has an effect
 486 on the spatial sensitivity of the inversion, as dominant noise sources cannot be resolved
 487 without data from surrounding stations. In combination with the lack of grid points on
 488 land due to our parameterisation, this can lead to inversion artefacts in areas where we
 489 would not necessarily expect dominant noise sources; for example in marginal seas like
 490 Hudson Bay or the Mediterranean Sea. This also happens when there is a lack of coher-
 491 ent signals in the cross-correlations which are then ignored due to our data selection based

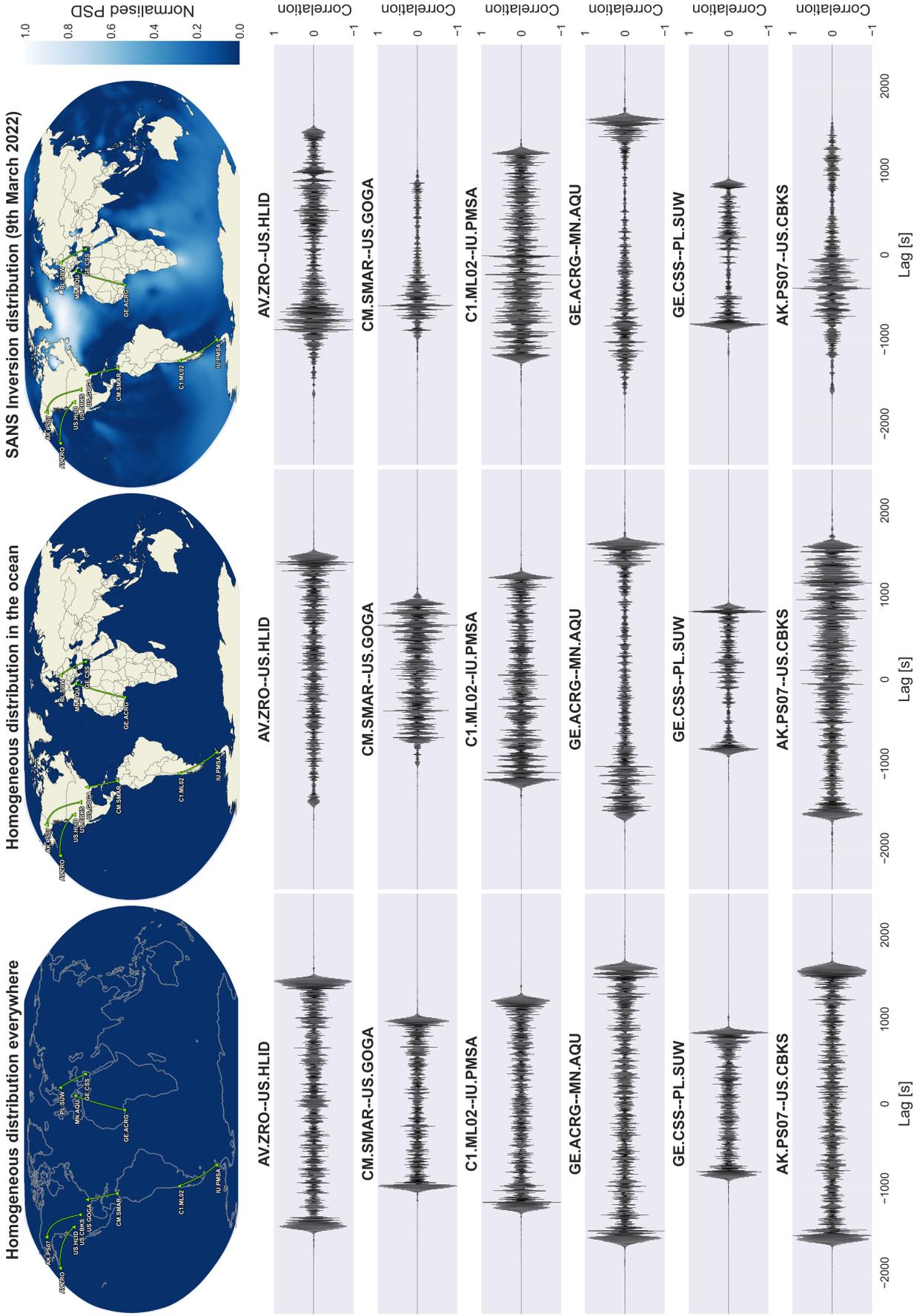


Figure 7. Comparison of modelled cross-correlations for 6 station pairs (\blacktriangle) and three different noise source distributions: (i) homogeneous everywhere, (ii) homogeneous in the ocean, and (iii) final SANS inversion model for the 9th March 2022. As the noise source distribution becomes more realistic, the changes in the cross-correlations become more and more significant. This effect should not be ignored in ambient noise studies. Note that the noise source distribution for the homogeneous forward models are set to 1 everywhere, not 0 as indicated by the colourbar.

492 on the signal-to-noise ratio. To analyse the effect of the changing data and station avail-
 493 ability on the resolution, we perform a synthetic test where we forward model cross-correlations
 494 with added noise and invert for them using different station lists. Generally, the inver-
 495 sions with fewer stations still include the most dominant noise sources, albeit with a lower
 496 spatial resolution.

497 5 Conclusions and Outlook

498 We present a new web framework SANS (sans.ethz.ch) where daily seismic am-
 499 bient noise source maps for the secondary microseisms on a regional to global scale are
 500 made available to the public. Two methods are combined to improve the final noise source
 501 distribution: Matched Field Processing (MFP) and a gradient-based iterative finite-frequency
 502 inversion. The efficient data-driven MFP approach provides a starting model to steer the
 503 inversion in the right direction. Pre-computed wavefields and spatially variable grids have
 504 decreased the computational cost of the inversions, allowing us to run the inversions ev-
 505 ery morning for the previous days' data and presenting the results shortly after. Users
 506 are able to download the inversion results and we provide code to ease the implemen-
 507 tation of the noise source distribution maps into other workflows. Comparisons to sig-
 508 nificant wave height maps do show that areas with high waves and strong dominant noise
 509 sources often coincide, which is in-line with the generation mechanism of secondary mi-
 510 croseisms. Furthermore, we compute the averages of the noise sources maps for North-
 511 ern Hemisphere summer and winter and compare them to the averages of the significant
 512 wave height maps. These show very similar areas of stronger activity which are in-line
 513 with other studies: Northern Hemisphere summer has more dominant sources in the South-
 514 ern Hemisphere and Northern Hemisphere winter is dominated by noise sources in the
 515 North Atlantic.

516 We hope that making the noise source distribution data readily available to the pub-
 517 lic encourages new tomographic studies and methods exploiting seismic ambient noise
 518 vibrations. The accuracy of tomographic models could be improved by implementing knowl-
 519 edge of the noise sources. Specifically studies that make assumptions about a homoge-
 520 neous or quasi-random noise source distribution would benefit and this may lead to more
 521 accurate velocity models. Studies that focus on time-dependent velocity changes in the
 522 subsurface often try to observe changes on the order of 1% or less (e.g. Zhan et al., 2013;
 523 Delaney et al., 2017). Particularly for such monitoring purposes, it is important to ver-
 524 ify that these changes are not a result of a changing noise source distribution. The near
 525 real-time seismic ambient noise source maps we present here are a crucial tool to pro-
 526 vide this verification. Future applications could also make this approach feasible for more
 527 local studies like the near real-time monitoring of avalanches and landslides.

528 Acronyms

529 **CSCS** Centro Svizzero di Calcolo Scientifico
 530 **ETH** Eidgenössische Technische Hochschule Zürich
 531 **GCMT** Global Centroid Moment Tensor
 532 **HPC** High-Performance Computing
 533 **MFP** Matched Field Processing
 534 **PREM** Preliminary Reference Earth Model
 535 **PSD** Power-Spectral Density
 536 **SANS** Seismic Ambient Noise Sources

Data availability

The website introduced here can be found on <https://sans.ethz.ch/>. The inversion code is available on github: https://github.com/jigel/noisi_inv and is based on previous work (Ermert et al., 2020; Igel et al., 2021). Within the repository is a Jupyter Notebook Tutorial on how to run an inversion, including downloading, processing, and cross-correlating the data. Pre-computed AxiSEM wavefields that can be downloaded and implemented are available online at <http://ds.iris.edu/ds/products/syngine>.

The seismic data was collected from multiple data centers using ObsPy (Krischer et al., 2015) and the authors thank everyone involved in setting up and maintaining these: IRIS (<http://service.iris.edu>), GEOFON (<http://geofon.gfz-potsdam.de>), ORFEUS (<http://www.orfeus-eu.org>), NIEP (<http://eida-sc3.infp.ro>), RESIF (<http://ws.resif.fr>), INGV (<http://webservices.ingv.it>), SCEDC (<http://service.scedc.caltech.edu>), BGR (<http://eida.bgr.de>), ETH (<http://eida.ethz.ch>), KOERI (<http://eida.koeri.boun.edu.tr>), LMU (<http://erde.geophysik.uni-muenchen.de>), NCEDC (<http://service.ncedc.org>).

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