

Mesoscale and Large-Eddy Simulation of the Boundary Layer Process of Cumulus Development over Naqu, Tibetan Plateau Part A: Comparison Between Simulation and Observation

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Abstract

Cumulus clouds are of great interest in numerical weather prediction. However, the scarcity of observed data on the Tibetan Plateau (TP) has not allowed correct interpretation of their development. The Third TP Atmospheric Science Experiment provided experimental data to address this challenge. This study used a combined weather research and forecasting large-eddy simulation (WRF-LES) model and final reanalysis data from the Global Forecast System to simulate cumulus clouds over southern TP on July 19, 2014. We applied observation nudging and one-way nesting strategies to influence the optimality of WRF-LES runs. The study performed simulations with six different scenarios in comparison with observation data. Results showed that cumulus clouds locally initiated and grew upscale but were however influenced by large-scale forcing. Compared to the observations, simulations with observation nudging provided more accurate and reliable results than the simulations without nudging. LES with mesoscale forcing yielded a relatively good atmospheric boundary-layer (ABL) water vapor profile and a similar microphysical evolution to the observations but misled the observed surface variables. Without mesoscale forcing, LES provided the best ABL water vapor and sensible heat flux, however, failed to provide a good microphysics field. In this aspect large-scale forcing played an important role in cumulus development on July 19, 2014. The study recommended observation nudging and one-way nesting strategies in separate iterations to be improved by focusing on the model's response to the terrain and boundary conditions.

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1 **Mesoscale and Large-Eddy Simulation of the Boundary Layer Process of Cumulus**
2 **Development over Naqu, Tibetan Plateau**
3 **Part A: Comparison Between Simulation and Observation**

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20 **Key Points:**

- 21 • The ability of the WRF-LES model to reproduce the atmospheric boundary-layer
22 processes of cumulus development over complex orography;
- 23 • Simulated ABL, reflectivity, and Liquid water content compared with observed Lidar
24 measurements, radar reflectivity, and satellite image;
- 25 • WRF-LES simulations with observation nudging and one-way nesting strategies better
26 replicated the observed clouds pattern.

27 **Abstract**

28 Cumulus clouds are of great interest in numerical weather prediction. However, the scarcity
29 of observed data on the Tibetan Plateau (TP) has not allowed correct interpretation of their
30 development. The Third TP Atmospheric Science Experiment provided experimental data to
31 address this challenge. This study used a combined weather research and forecasting large-eddy
32 simulation (WRF-LES) model and final reanalysis data from the Global Forecast System to
33 simulate cumulus clouds over southern TP on July 19, 2014. We applied observation nudging
34 and one-way nesting strategies to influence the optimality of WRF-LES runs. The study
35 performed simulations with six different scenarios in comparison with observations data. Results
36 showed that cumulus clouds locally initiated and grew upscale but were however influenced by
37 large-scale forcing. Compared to the observations, simulations with observation nudging
38 provided more accurate and reliable results than the simulations without nudging. LES with
39 mesoscale forcing yielded a relatively good atmospheric boundary-layer (ABL) water vapor
40 profile and a similar microphysical evolution to the observations but misled the observed surface
41 variables. Without mesoscale forcing, LES provided the best ABL water vapor and sensible heat
42 flux, however, failed to provide a good microphysics field. In this aspect large-scale forcing
43 played an important role in cumulus development on July 19, 2014. The study recommended
44 observation nudging and one-way nesting strategies in separate iterations to be improved by
45 focusing on the model's response to the terrain and boundary conditions.

46 1. Introduction

47 The atmospheric boundary layer (ABL) processes of Cumulus Cloud (CC) development are
48 small-scale unresolved motions and have significant effects on larger-scale resolved signals.
49 These processes have particular consideration in numerical weather prediction (NWP) (Emanuel,
50 1997; Ravichandran & Narasimha, 2020). Meso- and micro-scales modeling (MMM) is being
51 actively used to solve the physics of the ABL processes of CC development over complex
52 terrains (Kane & Klein, 2005; Krueger, 1988; Mechem & Oberthaler, 2013), with the purpose to
53 enrich the phenomenological basis for MMM. Unfortunately, despite an increase in computing
54 power, the latter modeling strategy comes with a host of other challenges, such as the scarcity of
55 observation data over complex orography to simulate the development mechanism of CC.
56 Ultimately, the availability of appropriate boundary conditions to replicate the development of
57 the ABL processes (P. Ray, 2015) is a limit to MMM.

58 As such, Tibetan Plateau (TP) region is one of the most complex terrains in the world.
59 Human life and the ecosystem of South East Asia (SEA) depend on water from the major rivers
60 such as the Brahmaputra, Ganges, Irrawaddy, Mekong, Salween, Yangtze, and Yellow River,
61 which headwaters are located on TP. The water supply in these rivers is strongly related to the
62 ABL processes of the cumulus convection that further produces precipitation fall (Zhao et al.,
63 2018, 2019). In addition, the ABL processes over TP are well known to hydrating the global
64 atmosphere (R. Fu et al., 2006; IPCC, 2014). However, the predictability of the latter process on
65 TP relied on sparse and scarce observations that did not provide the required accuracy, spatial
66 density, and temporal frequency (Y. Liu et al., 2020; Zhao et al., 2019), posing difficult
67 challenges when applying numerical modeling to understand the CC development in that region.
68 The latest campaign work, the Third TP Atmospheric Science Experiment (TIPEX-III),
69 conducted from July 1 to August 31, 2014 (L. Liu et al., 2015; Zhao et al., 2018, 2019), provided
70 comprehensive experimental data, which motivated our interest in studying ABL process of CC
71 development over Naqu in south TP (STP).

72 Typically, CCs are detached and dense, with sharp outlines. They develop vertically in the
73 form of rising towers. CCs often begin to form in sunny and fair weather as soon as the
74 ascending air cools to the point where vapor becomes supersaturated. Subsequently, the water
75 vapor condenses into liquid water droplets or solid ice crystals. Congestus CCs may turn into
76 cumulonimbus and produce thunderstorms when influenced by mesoscale instability, humidity,
77 and temperature gradient. Cumulonimbus grows vertically and may penetrate at greater heights
78 from 300 to 12,000 m AGL (Cotton & Anthes, 1992). The large CCs and thunderstorms are
79 classified into meso- γ atmospheric processes (Anthes, 1986; Cotton & Anthes, 1992; Emanuel,
80 1993; Fujita, 1986). CCs are precursors of other cloud types, while meteorologists may refer to
81 CCs underway to determine the kind of weather that will occur. However, the ABL processes of
82 CC developments have limited predictability due to uncertainty associated with the initial
83 conditions. According to the Fifth Assessment Report of the Intergovernmental Panel on Climate
84 Change (IPCC, 2014), the numerical representation of the ABL clouds is an open problem in
85 cloud modeling. This points out a need for the numerical study of ABL clouds development.

86 The broad classes of mesoscale phenomena influencing the ABL processes of CC
87 development are the internally and externally forced-mesoscale processes. The former class
88 derives structure and circulation within the atmosphere. The latter class results from the earth's

89 surface–atmosphere interaction, including either thermal forcing due to the differential heating
90 of the coupled earth’s surface–atmosphere, or mechanical forcing due to the atmospheric
91 response to the irregular topography (Arya, 2001; T. Wang et al., 2002; Wu et al., 2007). Each
92 area of TP, including STP, has particular ABL processes of CC development, interacting with
93 large–scale forcing. The diurnal surface heating over TP in summer reflects an external forcing
94 element, which interacts with a large–scale forcing, and mesoscale moisture transport then
95 influences the ABL process of cumulus development. There is fifty years back, (Flohn, 1968)
96 highlighted that during the summer the TP act as a heat engine with an enormous convective
97 chimney in the southeastern sector where giant cumulonimbus cells play a major role in
98 continuously carrying heat upward into the high troposphere. Li et al., (2010) and Ding et al.,
99 (2018) revealed that warm and wet events have notably increased over the region and altered
100 hydrological processes (X. Liu et al., 2006; J. Wang et al., 2020). As a result, various weather
101 conditions influence billions of people living in the SEA region, particularly in the downstream
102 sector (Lei et al., 2021; T. Wang et al., 2002; K. Yang et al., 2004).

103 The effect of the diurnal temperature variation on TP has been observed in the diurnal
104 evolution of the TP’s ABL structure, the most characteristic in the world. In dry and warm land
105 surface–atmosphere conditions, the top of the ABL over TP could reach approximately 5 km of
106 mean height AGL, higher than any reported ABL worldwide while a shallow ABL top is
107 observed only about 2 km mean of height AGL in the moist condition of the coupled system (Ao
108 et al., 2017; Chen et al., 2013, 2016; Y. Li & Gao, 2007; Sato, 2009; Yanai & Li, 1994; Zhao et
109 al., 2018). According to previous studies, ABL on TP exerts a profound thermal and dynamic
110 influence on the cumulus development (Chen et al., 2016; Y. Li & Gao, 2007; Slättberg & Chen,
111 2020; K. Yang et al., 2004). Many studies showed that the radiative budget of the ABL has a
112 direct effect on the variation in CC cover (Ao et al., 2017; Betts & Ball, 1994; Margulis &
113 Entekhabi, 2004). However, the upper–level potential vorticity structures and the meridional
114 position of the subtropical jet (STJ) also influence the feature of the ABL on the TP (Chen et al.,
115 2016). In midsummer, STJ and the South Asian High (SAH), coexist and intensify the
116 upper–level subsidence in some STP’s zones, characterized by less cumulus development (Chen
117 et al., 2016; Sato, 2009). The mean position of the STJ was 40° north between 1979 and 2003
118 (Sato, 2009).

119 From the 1970s to 2014, three scientific experiments, including TIPEX-III, promoted the
120 understanding of the cloud processes on TP (Y. Fu et al., 2020; L. Liu et al., 2015; Zhao et al.,
121 2018, 2019). However, the pursuit of in-depth knowledge of water and energy cycles on TP still
122 raises questions related to multi-scale CC interactions. Three decades ago, Tingyang & Reiter,
123 (1990) found that the condensation rate of water vapor in clouds, the clouds’ liquid water
124 content (LWC), and the precipitation efficiency in clouds are lower than those in surrounding
125 regions through observation analysis and model simulations. Based on CloudSat and Cloud-
126 Aerosol Lidar and Infrared Pathfinder Satellite Observations (CALIPSO), (Luo et al., 2011)
127 further showed that CC convections were shallower over TP than those over the other two
128 subregions of the TP–South Asian monsoon region. By contrast, the TIPEX-III campaign work
129 revealed that most congestus CCs developed on TP, preferably in the afternoon, some of which
130 penetrated vertically as high as 16.5 km AGL (L. Liu et al., 2015; Zhao et al., 2018).

131 Most studies mentioned above focused on model simulations of the large-scale and
132 mesoscale cloud processes on TP (Couvreur et al., 2009; Gao et al., 2016; Larson et al., 1999;

133 Tingyang & Reiter, 1990; K. Yang et al., 2004). In line with this, Fu et al., (2020) reviewed the
134 progress of land-atmosphere interactions on TP. The authors suggested conducting an
135 observational analysis coupled with numerical simulations to further understand the effect of
136 ABL processes. Couvreur et al., (2009) pointed out that ABL water vapor variability is the main
137 trigger of mesoscale convective clouds. Accordingly, this study analyzed the ABL features that
138 may influence CC development.

139 Because of the extensive limitations in understanding the cloud processes on TP, Gao et al.,
140 (2016) showed that clouds and precipitation over TP have not been studied sufficiently due to the
141 lack of observations over the harsh mountainous areas and the poor representation, in NWP
142 models, of the CC processes over TP. Many studies emphasized that many modeling studies of
143 CC processes on TP were probably inadequate to capture the complex interactions of the
144 physical processes (Gao et al., 2016; L. Liu et al., 2015; Zheng et al., 2015, 2016). Sato, (2009)
145 investigated the resolution dependency of the diurnal cycle of convective clouds on TP in a
146 mesoscale model. They pointed out that convective clouds over TP during the day tend to have a
147 small horizontal scale rather than a resolvable mesoscale resolution, and suggested further study
148 with finer resolutions of less than 7 km. However, they did not specify how fine is enough to
149 resolve the inaccuracy associated with the orography feature. With advances in the application of
150 the coupled WRF–large eddy simulation (WRF-LES), it is now easy to understand the link
151 between multi–scale CC development and determine the dominant trigger of CCs (Chow et al.,
152 2005; Takemi & Rotunno, 2003; Zhu et al., 2010).

153 This study answered the question of how the model can reproduce the CC process by
154 performing six scenarios with WRF-LES in a real-case mode. To this end, this study organized
155 the other steps as follows: Section 2 describes the experimental data and modeling framework.
156 Section 3 gives a comparison between the simulation and observation. The summary and
157 conclusions are presented in Section 4.

158 **2. Data, Modeling Framework**

159 **2.1. Experimental Data**

160 The study used the TIPEX-III dataset from 8 meteorological stations, where the surface
161 measurements and L-band radiosonde soundings were used as input for the observation nudging
162 to improve the initial and lateral boundary conditions (LBCs) and gradually adjust the model
163 state toward observations during its integration. The purpose was to link the model’s prognostic
164 variables to the observed quantities to capture the atmospheric flow around the Naqu area,
165 including Naqu1 and Naqu2. We can approximate the distance between Naqu1 and Naqu2
166 stations as 5.4 km. The stations consist of active remote sensing sites, automatic meteorological
167 towers, ABL towers, and 2 surface weather observation stations, Naqu Plateau Cold Climate and
168 Environment 1 and 2 (NPCE1 and NPCE2). These stations follow the recommendations of
169 WMO, and all data from TIPEX-III are quality controlled before publication. The measurement
170 devices used during TIPEX-III included operational radiosondes, cloud radars, laser ceilometer
171 31 (CL31), and other in situ observation devices. The Cold and Arid Regions Environmental and
172 Engineering (CARE) research institute and the Chinese Academy of Sciences (CAS) support
173 these stations in providing unique observation data. The data are distributed by the China

174 Meteorological Administration (CMA). Liu et al., (2015) describes most devices used during the
 175 TIPEX-III.

176 **Table 1.** Experimental data stations; asterisks indicate sounding sites; double asterisks indicate
 177 the source of radar reflectivity; AGL, above ground level.

Station Name	Longitude	Latitude	Station Height (m) AGL	Station ID
LHASA*	91.133	29.667	3649	55591
LINZHI*	94.340	29.667	2992	56312
NAQU1*	92.067	31.483	4508	55299
NAQU2**	92.01	31.48	4507	55299
NPCE1	91.900	31.300	4534	-
NPCE2	90.900	31.300	4508	-
RIKAZE*	87.080	28.633	4302	55664
YITUOHE*	92.433	34.210	4534	56004

178

179 2.1.1. Surface and Tower-Based Measurements

180 Surface measurements from NPCE1 and NPCE2 and tower-based measurements from
 181 Naqu1 include air temperature (T), relative humidity (RH), air pressure (P), vapor pressure (e),
 182 and wind speed (U_z). Tower-based measurements were available at 0.75, 1.5, 3, 6, 12, and 22 m
 183 AGL. Data from each site were time series recorded from July 17 to 20, 2014. Each item in the
 184 series represents the mean of the data recorded over a 30-min period. The tower-based
 185 measurements were used to compute the mean wind speed U_z , the potential temperature
 186 difference $\Delta\theta = \theta(T_z) - \theta(T_s)$, and specific humidity difference $\Delta q = q_s - q_z$, where subscript
 187 s represents 0.75 m AGL and z represents 3 m AGL for temperature and specific humidity, and
 188 12 m AGL for wind speed. This study estimated HFX and the latent heat flux (LH) as given in
 189 Equations (1) and (2) (Gavilán & Berengena, 2006; Verma et al., 1986):

$$\text{HFX} = C_H \cdot \rho_a \cdot C_p \cdot U_z \cdot \Delta\theta \quad (1)$$

$$\text{LH} = C_E \cdot \rho_a \cdot L_v \cdot U_z \cdot \Delta q \quad (2)$$

190 where $L_v = 2.501 - (2.361 \cdot 10^{-3}) \cdot T$ is the latent heat of vaporization, C_p is the specific heat of
 191 the air, ρ_a is the mean air density, and C_H and C_E are the bulk transfer coefficients for heat and
 192 moisture, respectively. The estimation was done assuming the bulk transfer coefficients equal to
 193 the drag coefficient (C_D). Li et al., (1996) and Zhao et al., (2018) found $C_D = 4.4 \cdot 10^{-3}$ for the
 194 Naqu area, and there was no need to perform this calculation in this study.

195 2.1.2. L-Band Radiosonde Sounding Measurements

196 The L-band radiosonde sounding stations are indicated with an asterisk in Table 1. The
 197 variables used in this study include air temperature, wind speed, wind azimuth, air pressure,
 198 relative humidity, and geographical height. The data were recorded each minute for about 1.5
 199 hours and were available at the predetermined synoptic times, except on July 18, 2014, where 3
 200 soundings were launched, at 0615UTC, 1115UTC, and 2315UTC, over Naqu1. Wind speed and

201 azimuth were used to decouple the wind field into meridional and zonal wind components. The
202 latter was used as input for observation nudging.

203 **2.1.3. Automatic Laser Ceilometer Measurements**

204 The ABL height data were determined over Naqu2 with a Vaisala CL31 automatic laser
205 ceilometer. The CL31 has a time resolution of 16 s and a vertical resolution of 5 m. The CL31
206 used in this study is a mini-Lidar made in Finland and maintained by China Ocean University.
207 The CL31 is used for active remote sensing measurements to characterize the ABL height with
208 the backscatter signal and has good accuracy (Kotthaus & Grimmond, 2018; Sokół et al., 2014).
209 This study used the measurements from the CL31 as a reference for the simulated ABL. The
210 purpose of using the CL31 measurement was to avoid a false ABL depth estimation.

211 **2.1.4. Ka-Band Millimeter-Wave Radar and Fengyun 2D Satellite Images**

212 The cloud row data were determined over Naqu2 using a Ka-band millimeter–wave cloud
213 radar. The device has a time resolution of 0.85 s and a vertical resolution of 30 m. The data from
214 July 19, 2014, were used to retrieve the radar reflectivity. The radar reflectivity mentioned here
215 is a measure of the fraction of the precipitation intensity reflected from the cloud surface. The
216 millimeter-wave cloud radar data browsing software, HMB-Disp, provided by Naqu2 was used
217 to extract the radar reflectivity. Visible light 2D images from the Fengyun (FY-2D) geostationary
218 meteorological satellite at 16:45 local standard time (LST) on July 19, 2014, were also used as a
219 reference to simulated reflectivity. The original data has a horizontal resolution of 5600×4800
220 pixels.

221 **2.2. WRF Modeling Framework**

222 **2.2.1. WRF Model Description**

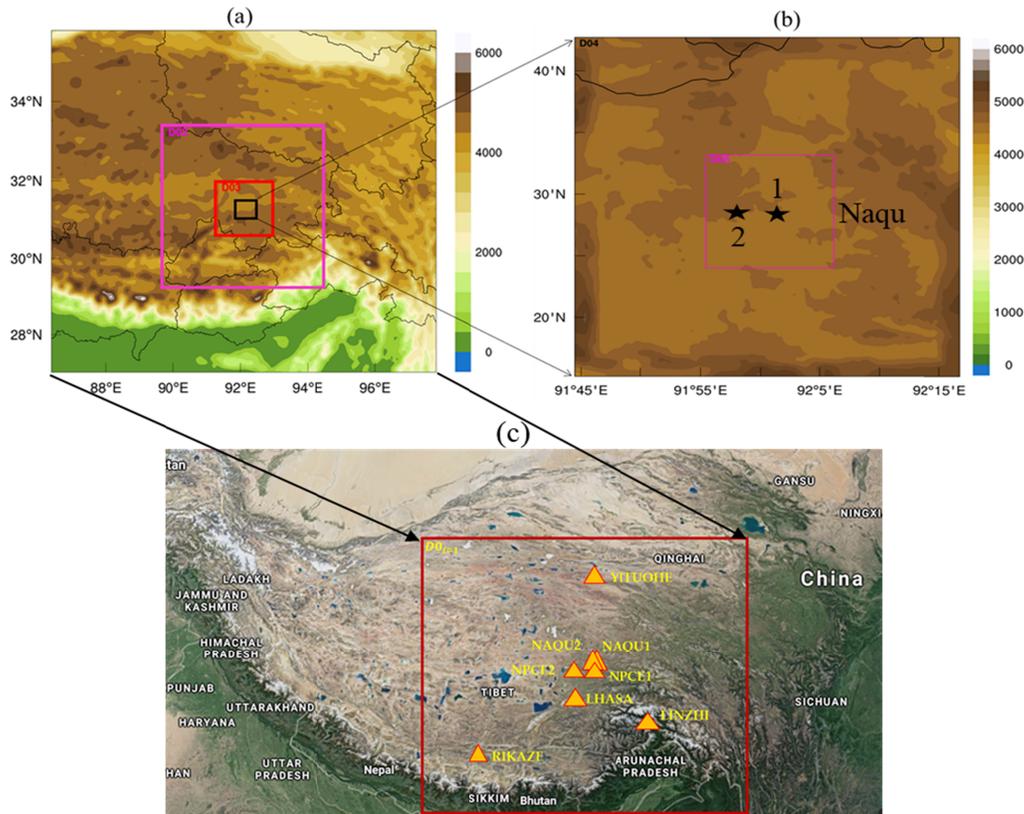
223 The WRF model is a non–hydrostatic, compressible atmospheric model, which is the most
224 widely used in NWP for research and operational needs (Powers et al., 2017). The WRF model
225 has grid nesting capability, which allows simultaneous multiscale simulation (Powers et al.,
226 2017; Skamarock et al., 2008). The model has several initialization programs for idealized and
227 real–data simulation cases and provides several parameterization options, such as land surface,
228 surface layer, planetary boundary layer (PBL), microphysics, CC parameterization (Gentry &
229 Lackmann, 2010), and adaptive subgrid–scale (SGS) mixing parameterization (Canuto &
230 Cheng, 1997; Y. Liu et al., 2011; Takemi & Rotunno, 2003). In this model, the turbulent kinetic
231 energy and non–local closure schemes enable the PBL to develop with entrainment (Dai et al.,
232 2014; Dudhia, 2014; Shin & Dudhia, 2016; Skamarock et al., 2008).

233 As mentioned above, the WRF model has more computing power and allows the coupled
234 WRF–LES to perform real case simulations at meso– and micro–scales resolutions (Y. Liu et
235 al., 2020; Talbot et al., 2012). The purpose of applying the LES is to implicitly calculate the
236 small-scale turbulence from the flow field and explicitly resolve the largest scales of energy
237 production, while the basis of the LES strategy in the WRF model is low–pass filtering (Chow et
238 al., 2005; Chow & Street, 2009; Mirocha et al., 2010; Shin & Dudhia, 2016). The application of

239 coupled WRF-LES in real case mode uses a surface layer scheme to connect the lower boundary
 240 and the atmosphere.

241 **Table 2.** Five nesting domains, $\Delta x_{i=1-5} = \Delta y_{i=1-5}$, represent grid spacing (m) where i ranges
 242 from 1 to 5 to represent the domains $D0_{i=1}, D0_{i=2}, D0_{i=3}, D0_{i=4},$ and $D0_{i=5}$; $\partial t(s)$, the time
 243 step in seconds; \mathcal{A} is the area of the domain (km^2).

Mesoscale Simulation Runs				
Simulation Type	$\Delta x_{i=1-5} = \Delta y_{i=1-5}$ (m)	Horizontal Grids Points	\mathcal{A} (km^2)	$\partial t(s)$
$D0_{i=1}$ Mesoscale- γ	9,234	150×150	1385.1×1385.1	15
$D0_{i=2}$ Mesoscale- γ	3,078	150×150	461.7×461.7	5
$D0_{i=3}$ Microscale- α	1,026	150×150	153.9×153.9	5/3
Large-Eddy Simulation Runs				
$D0_{i=4}$ Microscale- α	342	150×150	51.3×51.3	3/20
$D0_{i=5}$ Microscale- β	114	150×150	17.1×17.1	1/20



244
 245 **Figure 1:** WRF domain configuration: (a) mesoscale domains, (b) microscale domains; shaded area represents terrain height (m)
 246 above sea level (ASL). The black stars 1 and 2 are respectively the Naqu1 and Naqu 2 stations (c) google map with yellow
 247 triangles indicating the location of each station.

248 The coupled WRF-LES is constructed by applying a one-way nesting strategy in two
 249 separate iterations. This strategy is an option in the WRF model, defined as a finer-grid-
 250 resolution run. This strategy is used as a subsequent run after the coarser-grid-resolution run,
 251 where the program (ndown) is run in between the two simulations. The coarse-grid run provides

252 the initial and lateral boundary conditions (LBCs) to the finer-grid run, with data from higher-
 253 resolution land and masked surface fields. The latter strategy only feeds suitable information
 254 from the outer domain to the inner one. The one-way nesting strategy has been used in many
 255 studies with good results (Chow & Street, 2009; Y. Liu et al., 2011; Moeng & Wyngaard, 1989;
 256 Soriano et al., 2004; Zhu et al., 2010).

257 Soriano et al., (2004) suggested that for mesoscale simulation runs, the one-way nesting
 258 technique should be used because two-way nesting gave the worst results in their case study. Zhu
 259 et al., (2010) suggested that using one-way nesting runs allows simulation of complicated and
 260 heterogeneous forcings, and recommended that this strategy be used for cloud cases that require
 261 extremely high resolution.

262 The WRF model system also includes the observation nudging strategy used in this study.
 263 This strategy uses Newtonian relaxation to improve either the dynamics initialization or analysis,
 264 respectively (Cheng et al., 2009; Reen et al., 2016; Yesubabu et al., 2014). However,
 265 assimilation of moisture fluxes may pose particular difficulties when applying observation
 266 nudging due to the high spatial variability of the variables such as the water vapor mixing ratio
 267 and specific humidity relative to their absolute values (Reen et al., 2016). For simplicity, this
 268 study applied observation nudging to wind and temperature to avoid excessive drying that can be
 269 caused by negative water vapor values occurring within the model domain.

270 2.2.2. WRF Model Domain Configuration

271 This study used version 3.8 of the WRF-ARW dynamical solver installed on the Tianhe
 272 high-performance computing (HPC) system and assisted by the Sugon HPC system from the
 273 State Key Laboratory of Disaster Prevention and Reduction for Power Grid Transmission and
 274 Distribution Equipment (SKL) at Changsha, China. To get a suitable design for the multiscale
 275 atmospheric reanalysis, we designed 5 WRF model domains, $DO_{i=1-5}$, where $DO_{i=1} = D01$,
 276 $DO_{i=2} = D02$, ..., $DO_{i=5} = D05$ by assigning $DO_{i=1-3}$ to mesoscale simulations, and 2
 277 microscale domains $DO_{i=4,5}$ all in $DO_{i=3}$, assigned to LES runs. The model domain
 278 configuration is given in Table 2 and depicted in Figure 1. As shown in Figure 1a, Naqu1 is the
 279 center of all domains. The parent domain $DO_{i=1}$ includes China's 6 administrative prefectures,
 280 Naqu, Linzhi, Lhasa, Rikaze, Shannan, Chamdo, and Yituohe, and neighboring southern
 281 countries. As shown in Figure 1b, $DO_{i=4}$ was configured to cover Naqu1 and Naqu2
 282 observation stations, while Naqu1 is the center of $DO_{i=5}$.

283 As shown in Table 2, each domain has 150 horizontal grids. In this configuration, the
 284 innermost domain has 114m horizontal grid resolution with an area of $114 \times 114 \text{ km}^2$ including
 285 Naqu1 and Naqu2. In this configuration, some vertical layers was set to 50 full sigma levels up
 286 to 50 hPa, except in one scenario namely scenario B1 wherethe vertical levels have been
 287 increased and reorganized manually to be many within the ABL. In the domain configuration for
 288 scenario B1, the vertical layers are composed of 56 full sigma levels up to 50 hPa, while the
 289 vertical increment $\Delta z_{i=0-55}$ between two consecutive levels increased linearly following the
 290 hyperbolic cosine function:

$$291 \left\{ \Delta z_i = Ch(\lambda_i) = \frac{e^{\lambda_i} + e^{-\lambda_i}}{2}; \lambda_{i=0-55} \in [2.797324366; 8.217039] \right\}$$

292 In the latter configuration, the first model level's height is 8 m AGL, while the first 40
293 model levels are below 4000 m height AGL (~ 300 hPa). The domain $DO_{i=3}$ provides the initial
294 conditions for $DO_{i=4}$ to optimize LES results. The horizontal and vertical interpolations are
295 performed with an overlapping quadratic approach and linear log pressure, respectively. These
296 interpolations allow the boundary forcing to vary in time and space. Mesoscale simulations use
297 1D PBL parameterization to fully parameterize the anisotropic turbulent motion using the
298 Reynolds-averaged Navier–Stokes (RANS) technique (Powers et al., 2017), whereas LES runs
299 treat unresolved isotropic turbulence using SGS schemes (Chow & Street, 2009; Deardorff,
300 1980; Moeng & Wyngaard, 1989; Shin & Dudhia, 2016; Takemi & Rotunno, 2003). In fact,
301 different ranges of scales from microscale to mesoscale must be solved to obtain a complete
302 representation of the superimposed physical mechanisms involved in the ABL process of CC
303 development. However, in the coupling model domain configuration, there is a range of grid
304 resolutions where certain processes are neither sub-grid nor resolved, the so-called “grey zone”
305 or “terra incognita”. These domains may be considered either mesoscale using 1D PBL or
306 microscale requiring LES strategy. The test experiment of (Talbot et al., 2012) successfully used
307 $DO_{i=4} = 450\text{m}$ for LES runs of the ABL. This study increased the resolution for $DO_{i=4}$, be closer
308 to the LES standard.

309 Our study focused on ABL processes of CC development on July 19, 2014. However, all
310 simulations were performed over 72 hours from July 17 to 20, 2014. This is because TP is at a
311 very high altitude and has complex orography, which causes the model to be statically unstable.
312 Therefore, after several text experiments, this study assumed that 24 hours are enough to assess
313 the effect of the model’s spin-up during model domain initialization. The history interval was set
314 to 15 minutes to create high time resolution data. We recorded 48 hours series (from day 2 to day
315 3) for model evaluation.

316 The traditional United States Geological Survey dataset (USGS) collected from 1992 to
317 1993 was applied to account for the influence of terrain and its related water bodies on WRF-
318 LES surface heat fluxes and meteorological variables. This study selected USGS 30 arc-second
319 (~ 900 m) Details on the data are given at the USGS website (<http://www.usgs.gov/>). Liu et al.,
320 (2020) and Sertel et al., (2010) demonstrate that the use of the default WRF dataset may cause
321 misrepresentation of the study region, while, there are spectral mixing problems between classes
322 in land use. Therefore, this study tested the high-resolution Shuttle Radar Topography Mission
323 (SRTM) dataset, 3 arc-second ($\sim 90\text{m}$) for comparison with USGS dataset (Not shown). The
324 data can be downloaded at (<http://www2.jpl.nasa.gov/srtm/>). However, there needed to be no
325 improvement with SRTM dataset. Some simulation results can be seen in Figure 1 of the
326 supplementary file.

327 Initial and lateral boundary conditions were generated from the National Centers for
328 Environmental Prediction (NCEP) Final Analysis (FNL) derived from the Global Forecast
329 System (GFS) and were accessed on September 05, 2019. FNL data have a resolution of $1^\circ \times 1^\circ$
330 grids at six-hourly time steps with 26 vertical levels from 10 to 1000 hPa of isobaric surface data
331 and can be downloaded at <https://doi.org/10.5065/D6M043C6>.

332 2.2.3. Flow Parameterization Options

333 This section gives details of the parameterization schemes used in this study and describes
 334 the experimental design. It also presents the one-way nesting strategy that coupled mesoscale
 335 simulations with microscale one. Table 3 summarizes the parameterization options. As shown in
 336 Table 3, the coupled land-atmosphere fluxes were computed using the Unified Noah Land
 337 Surface Model (Niu et al., 2011; Z.-L. Yang et al., 2011). Surface boundary conditions used
 338 Monin–Obukhov logarithmic similarity theory to prescribe fluxes of heat, moisture, and
 339 momentum (Chen et al., 1997; van den Hurk & Holtslag, 1997; Jiménez et al., 2012). The
 340 Mellor–Yamada scheme (Janjic, 2002) was selected for mesoscale simulations to account for
 341 turbulent kinetic energy (TKE) in local vertical mixing. The Thompson scheme was selected to
 342 parameterize the microphysical process. Short– and longwave radiation were integrated from the
 343 radiative transfer scheme (RRTMG) (Iacono et al., 2008; Mlawer et al., 1997). Kain, (2004)
 344 (hereafter referred to as KF) proposed a deep and shallow CC convection scheme, which was
 345 selected to resolve CC processes only in domain $D0_{i=1}$. Jeworrek et al., (2019) suggested
 346 combining the KF and Thompson microphysics schemes to improve the high-resolution
 347 numerical simulation results. However, the KF scheme was useless for $D0_{i=2-5}$, respectively,
 348 because these domains fell into the grey zone. On the other hand, a test simulation with the KF
 349 scheme in LES runs provided too much precipitation. Some pictures are shown in Figure 3 of the
 350 supplementary file.

351 Table 3: Summary of flow parameterization. SGS, subgrid-scale; $D0_{i=1-5}$.

Parameterization options	Mesoscale domains			Microscale domains	
	$D0_{i=1}$	$D0_{i=2}$	$D0_{i=3}$	$D0_{i=4}$	$D0_{i=5}$
Cumulus	Kain–Fritsch	Useless			
Planetary boundary layer	Mellor–Yamada			Real case LES mode	
Surface layer	Monin–Obukhov				
Land-surface	Unified Noah LSM				
Cloud microphysics	Thompson				
SGS stress model	Useless			TKE1.5	

352 For the real–case LES in the high mountain range, the recommended SGS turbulence
 353 model is a 1.5–order of TKE energy (TKE1.5) closure model (Skamarock et al., 2008). The
 354 TKE1.5 model accounts for diffusive transport of the TKE and allows more uniform diffusivity
 355 and entrainment through the convective ABL (Canuto & Cheng, 1997; Chow & Street, 2009; Y.
 356 Liu et al., 2020; Zheng et al., 2015). In real–case LES mode, the use of SGS TKE1.5 needs a
 357 filter (C_k) in the inertial subrange to filter the SGS noise (Deardorff, 1980; Takemi & Rotunno,
 358 2003). The filter was set to a default value ($C_k = 0.15$) in this study. The advection options of
 359 order $O(h) = 5$ and $O(h) = 3$ were used to compensate for the coarser horizontal resolution
 360 (Skamarock et al., 2008).

361 **2.2.4. Experimental Design**

362 We designed six scenarios (case experiments A, B, B1, C, D, and E), based on Talbot et al.,
 363 (2012) and Heinze et al., (2017), to account for the reliability of simulation results compared to
 364 observations, the efficiency of physical parameterization, and the skill of the strategy used in
 365 each scenario. These allowed us to prescribe dynamics associated with ABL process of CC
 366 development over TP. We first assessed the mesoscale simulations, then performed the LES runs
 367 through a one-way nesting strategy, as presented in Table 4.

368 **Table 4**, : Controlled experiment for mesoscale simulations and LES; observation data were used
 369 through Observation Nudging strategy; FNL, final reanalysis; TKE1.5, one and a half order
 370 turbulent kinetic energy closure model, used in anisotropic turbulence for mesoscale simulations
 371 and isotropic turbulence for LES; $R_{i=1-3}$ represent the radii of influence values from $DO_{i=1}$ to
 372 $DO_{i=3}$, and $R_{i=4,5}$ are the values from $DO_{i=4}$ to $DO_{i=5}$, Eta levels were generated automatically
 373 for cases A, B, and C but customized for case B1. SGS, subgrid-scale.

Mesoscale Simulations	Case	Input Data	Horizontal Turbulence	Eta levels	Horizontal Grid Spacing $\Delta x_{i=1-3} = \Delta y_{i=1-3}$ and $R_{i=1-3}$		
					$DO_{i=1}$ $\Delta x1 = \Delta y1$ $= 9,234$ m	$DO_{i=2}$ $\Delta x2 = \Delta y2$ $= 3,078$ m	$DO_{i=3}$ $\Delta x3 = \Delta y3$ $= 1,026$ m
					R_1 is useless	R_2 is useless	R_3 is useless
A	FNL	TKE1.5	50, 1rst model level: 65 m with 16 layers < 4000m	$R_1 = 980,000$ m	$R_2 = 330,000$ m	$R_3 = 120,000$ m	
B	FNL, Obs.		56, 1rst model level: 8 m with 40 layers < 4000m				
B1							
Microscale Simulations	Case	Input Data	Parent Mesoscale Turbulence Model	SGS Turbulence Model	Horizontal Grid Spacing $\Delta x_{i=4,5} = \Delta y_{i=4,5}$ and $R_{i=4,5}$		
					$DO_{i=4}$ $\Delta x4 = \Delta y4$ $= 342$ m	$DO_{i=5}$ $\Delta x5 = \Delta y5$ $= 114$ m	
	C	Input from case B, FNL, and observation	TKE1.5	TKE1.5	$R_4 = 40,000$ m	$R_5 = 15,000$ m	
	D	Observation and FNL	The parent mesoscale turbulence model (TKE1.5) is useless				
E	FNL						

374

375 As presented in Table 4, we designed a mesoscale simulation, scenarios A as a benchmark
 376 experiment to help determine the optimality of observation nudging in scenario B. We applied
 377 observation nudging in scenarios B, B1, C, and D. The nudging strength in these simulations was
 378 set to $120 \times e^{-4} s^{-1}$. As shown in Table 4, the radii of influence were set to $R_{i=1-5}$, respectively
 379 for $DO_{i=1-5}$. Each radius of influence on a given domain was slightly greater than a half-
 380 diagonal of that domain. Therefore, each point of the WRF domain at any distance within the
 381 domain will be influenced at least by one observation station. Observation stations are less dense
 382 and sparse (Figure 1c). We wanted every observation at any station to influence the whole
 383 domain at maximum. Therefore, scenario B determines the impact of observation nudging on the
 384 simulation results. Scenario B1 determines the optimality of the external forcing on the ABL
 385 such as surface drag and heating caused by the infrared radiation divergence imposed to the

386 atmosphere. Note that surface characteristics directly affect ABL. Therefore, increasing the
 387 layers in this part accentuates the effect of surface characteristics on the simulation results.
 388 Scenario C is an LES that determines the effect of large-scale forcing on microscale circulation.
 389 LES is useful for understanding the specific processes underlying the ABL, clouds, etc. In this
 390 scenario, a one-way nesting strategy was performed between $D0_{i=3}$ and $D0_{i=4}$ where the former
 391 provides the boundary condition for the latter with data from higher resolution land and masked
 392 surface fields. As mentioned earlier, this strategy only feeds suitable information from $D0_{i=3}$ to
 393 $D0_{i=4}$ and there is no feedback between these two domains. Scenario C differs from scenario D
 394 only by the use of the one-way nesting strategy. Indeed, scenario D is a control experiment for
 395 scenario C, which helps to determine the optimality of applying the one-way nesting strategy.
 396 We designed the scenario E as a benchmark experiment for LES run to help determine the
 397 optimality of observation nudging in scenario D. The vertical resolution has not been tested for
 398 scenarios C, D, and E due to their very high computational cost. Briefly, in this approach, we
 399 focused on scenario B to achieve our objective of using observation nudging. Scenario C utilized
 400 the input from scenario B as the boundary condition, and we were expecting the scenario C to
 401 reproduce the results from B but is much improved. A study by Heinze et al., (2017) successfully
 402 used the same approach to evaluate the mean ABL quantities and turbulence statistics.

403 **2.2.5. Evaluation metrics**

404 The model evaluation is based on the mean bias (MB) and the root mean squared error
 405 (RMSE) statistics (Stanski et al., 1989; Willmott et al., 1985) as given in Equations (12) and
 406 (13), respectively:

$$\mathbf{MB} = \frac{\sum_{j=1}^n (f_j - o_j)}{n} \quad (12)$$

407 where $f_{j=1,\dots,n}$ are the simulated values from the model, $o_{j=1,\dots,n}$ are the observation values, and n
 408 is the number of data points used in the calculation; and

$$\mathbf{RMSE} = \sqrt{\frac{\sum_{j=1}^n (f_j - o_j)^2}{n}} \quad (13)$$

409 Mean bias represents a gross measure of reliability, while RMSE represents a measure of the
 410 spread of differences between the forecast and observed values with the same units of
 411 measurement.

412 **3. Comparisons Between Simulations and Observations**

413 **3.1. Model Evaluation Results**

414 The evaluation scores are shown in Table 5, and Figures 2a,d illustrates the time series used
 415 for model evaluation. We considered the period after 24 hours of spin-up time for evaluation,
 416 that is from day 2 to 3. Table 5 shows that WRF-LES had a general tendency to overestimate the
 417 observation field except for the temperature, where the model underestimated the observation in
 418 all scenarios.

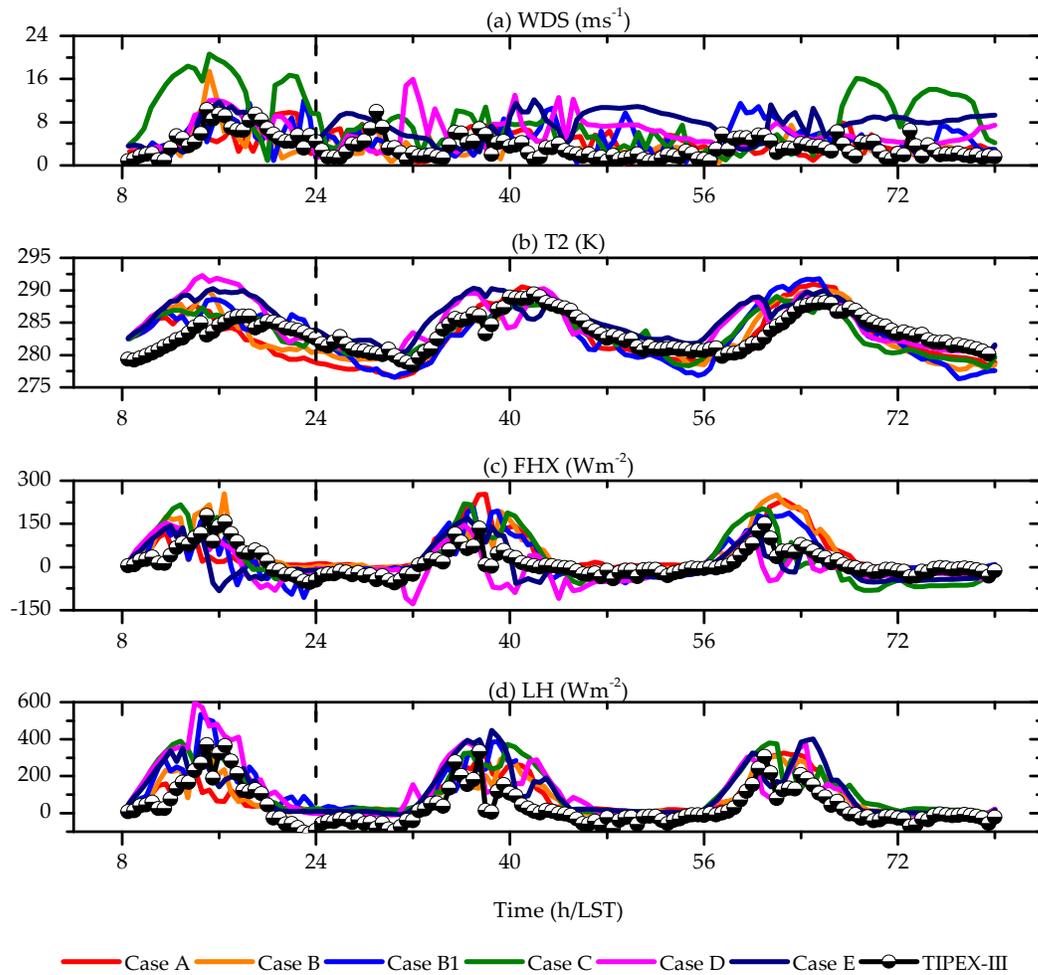
419 Table 5 shows that scenario A has a more accurate result in temperature than those in other
 420 simulations. The lowest score value is 1.8 K for RMSE. The largest values found in scenario A
 421 are LH scores of 61.64 w/m^3 for MB and 72.21 w/m^3 for RMSE. Large LH indicates large
 422 precipitation. As shown in Table 5, scenario B overestimated the observation except for the
 423 temperature. However, it tended to record the best lowest score values for wind speed (0.16 m/s
 424 for MB and 1.05 m/s for RMSE) and LH (69.57 w/m^2 for RMSE). The release of LH plays a role
 425 in heating the air, which rises. The air cools while the water vapor condenses, gradually forming
 426 clouds. Therefore, scenario B should provide the best distribution of cloud patterns compared to
 427 observation.

428 The scenario C, as shown in Table 5 overpredicted the observation, with a tendency to
 429 record the largest score values, especially for WSD (4.89 m/s for MB, 5.94 m/s for RMSE) and
 430 LH (61.98 w/m^2 for MB and 111.44 w/m^2 for RMSE). As mentioned earlier, WRF-LES
 431 overestimated observation in LH and predicted excessive precipitation. From the visual
 432 inspection, unexpected high wind speed values in scenario C (Figure 2a) contributed to high LH.
 433 Thus, the overprediction of wind speed in scenario C may be related to lateral boundary noise in
 434 the WRF-LES model runs, specifically when using a one-way nesting strategy in separate
 435 iteration steps. In fact, the one-way nesting strategy in separate iterations for the wind speed
 436 simulation over unresolved topography may increase errors in simulation results.

437 As shown in Table 5, scenario D provided a more accurate result in HFX compared to other
 438 scenarios, with the lowest score value of 40 w/m^3 . Also, scenario D had reliable results in
 439 temperature and LH, with the lowest score values of -0.16 K and 44.98 w/m^3 , respectively.
 440 Referring to a study by Zhao et al., (2018), WRF-LES runs in scenario D can predict the best
 441 ABL because it has better results in HFX.

442 Table 5: The mean bias (MB) and the root mean square error (RMSE) for two-meter temperature (T2), 10 m wind
 443 speed (WSD), the sensible heat flux (HFX), and the latent heat flux (LH); four scenarios are denoted as cases A, B,
 444 C, and D. $D0_{i=3}$ is the third domain and $D0_{i=5}$ is the fifth domain.

Score	Boundary-Layer Variables	Case A	Case B	Case B1	Case C	Case D	Case E
MB	Temp (K)	-0.41	-0.42	-0.15	-1.01	-0.16	2.02
	WSD (m/s)	0.83	0.07	1.32	4.75	4.56	3.80
	HFX (w/m^2)	42.24	37.88	19.26	-4.57	-23.87	-5.50
	LH (w/m^2)	61.64	56.26	97.41	-21.36	44.98	78.59
RMSE	Temp (K)	1.8	1.86	2.83	3.31	3.11	3.12
	WSD (m/s)	1.56	1.20	3.12	6.27	6.14	5.20
	HFX (w/m^2)	56.89	53.62	56.54	44.07	40.00	54.78
	LH (w/m^2)	72.21	69.57	132.71	103.20	89.45	116.00



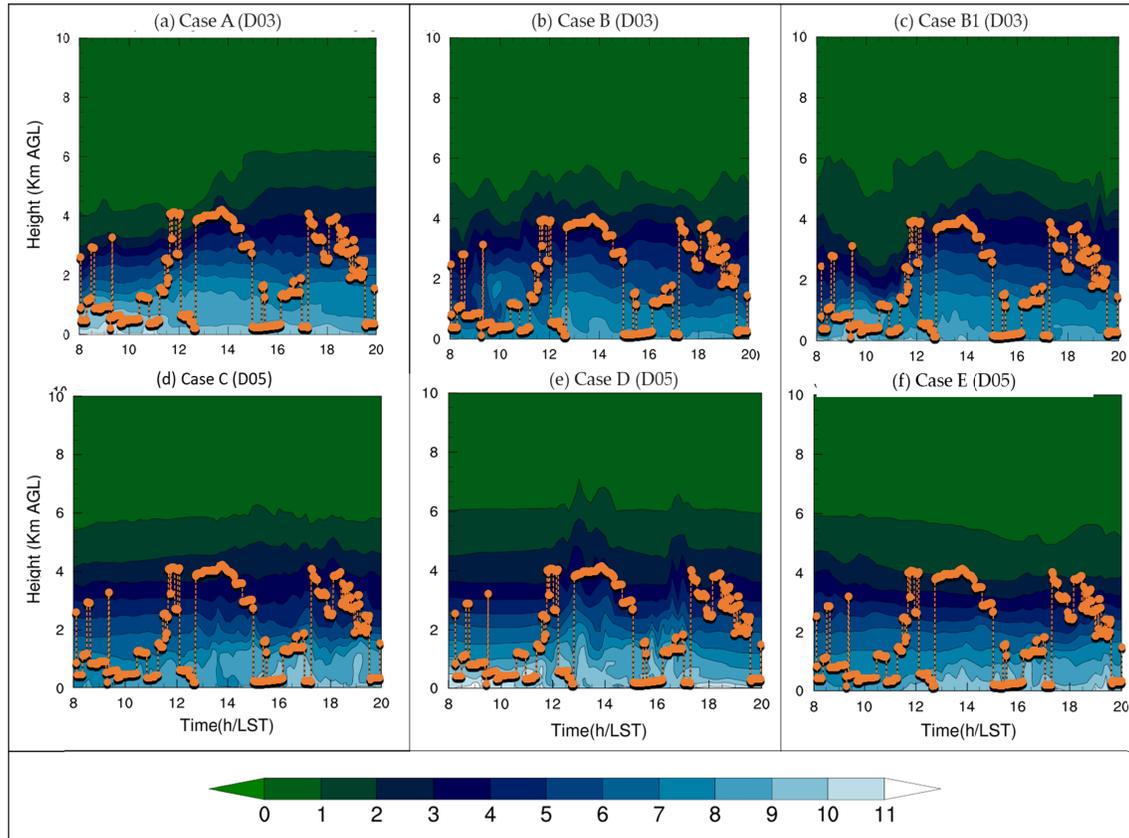
445

446 **Figure 2:** Time series of surface variables from $D0_{i=3}$ of the scenarios A, B and B1, and from
 447 $D0_{i=5}$ of the scenarios C, D and E (Table 4), compared to time series from Naqul (92.067,
 448 31.483): (a) wind speed (WSD); (b) Two-meter temperature (T2); (c) the sensible heat flux
 449 (HFX) at the surface; (d) the latent heat flux (LH) at the surface. Dashed lines indicate the upper
 450 limit of spin-up time.

451 Overall, for surface meteorological variables and heat fluxes, considering the set of results
 452 in each scenario, scenarios B and D had more accurate and reliable results than scenarios A, B1,
 453 and C. However, compared to the scenario A we found that the observation nudging strategy
 454 improved WSD, FHX, and LH results in the scenario B and HFX results in the scenario C and D.
 455 In fact, some results from the scenario C such as MB relative to HFX and LH were least scores
 456 and have a relatively good RMSE. However, we believe that noise occurred while transitioning
 457 through the ‘ndown’ program during the performance of the one-way nesting strategy. This is
 458 because

459 **3.2. The Simulated and Observed Boundary Layer Height**

460 The ABL height is a fundamental parameter characterizing the depth of atmospheric
 461 mixing near the earth's surface. It is critical for understanding the cloud processes related to ABL
 462 features and their feedback on the weather and climate system. The detection of ABL height
 463 from Lidar is as follows: at the top of the ABL, water vapor decreases abruptly and affects the
 464 Lidar signal to change rapidly around the ABL top.



465

466 **Figure 3:** Time-high cross-section of water vapor mixing ratio (g/kg) (shaded) over Naqu1
 467 predicted on July 19, 2014, at local standard time and compared to laser ceilometer
 468 measurements of the boundary layer height (ABL) from the Third Tibetan Plateau Atmospheric
 469 Science Experiment (TIPEX-III). Cases A, B and B_1 from $D0_{i=3} = D03$, C, D and E from
 470 $D0_{i=5} = D05$ represent scenarios in Table 4.

471 In this study, we compared the simulated ABL to that from the CL31 measurements based
 472 on the time–height cross-section of the water vapor mixing ratio. We found three peaks of the
 473 ABL from CL31 measurements, at 12:00, 14:00, and 18:00 LST. The peak found at 12:00 LST
 474 may be related to intermittent turbulence. The diurnal peak was reached around 14:00 LST. The
 475 peak found at 18:00 LST may have arisen from the high HFX released after intense convection
 476 developed in the early evening. We can also look for the last two peaks in the time series from
 477 temperature and heat fluxes shown in Figure 2 between the 56th and 72nd hours of the simulation
 478 period. The simulated ABL shows undulations and plumes of the water vapor mixing ratio with
 479 the mixing decreasing with height. We depicted the results of the comparison in Figure 3. Upon

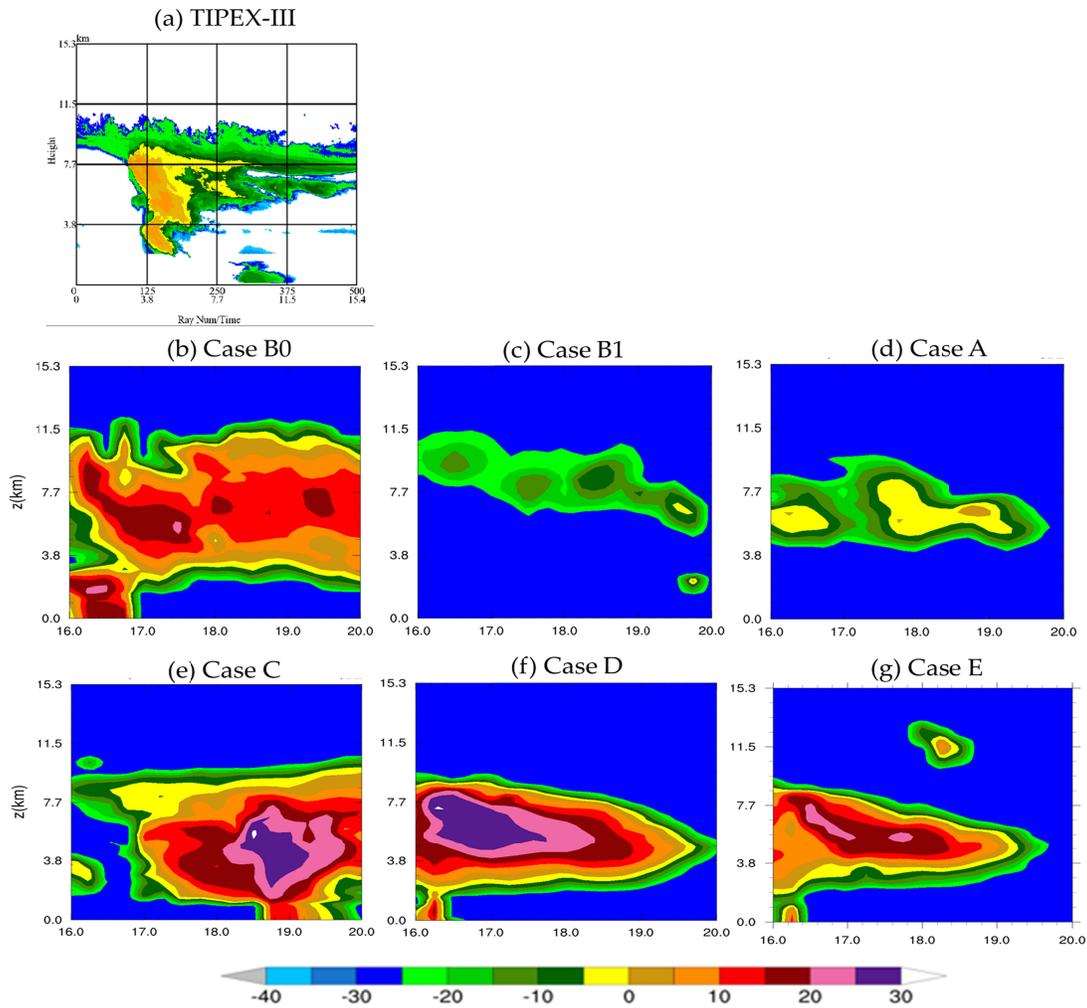
480 visual inspection, we can estimate the simulated ABL height as the height at which the water
481 vapor mixing ratio is 4 g/kg. This is because the CL31 measurements fit better with that height in
482 scenarios C and D (Figures 3d,e).

483 In the overall mesoscale simulation results, scenarios A, B, and B1 (Figures 3a,b,c)
484 presented a convective ABL between 10:00 and 20:00 LST, slightly higher than the observed.
485 We note that in dry and warm land surface-atmosphere conditions, the top of the ABL over TP is
486 deep, while a shallow ABL characterizes the moist condition of the coupled system. Under
487 current conditions, it is perceived that all mesoscale simulations will reproduce little clear skies
488 with few cumulus clouds (examples: case B and B1 of Figures 5c,d). We note also that each of
489 the scenarios A, B, B1 capture the observation field of the diurnal peak of the ABL. However,
490 the model failed to reproduce clearly the ABL peak around 18:00 LST. This is probably related
491 to the PBL scheme, which cannot treat the surface heterogeneity well. In scenarios C and D
492 (Figures 3d,e), the ABL peaks collapsed slightly compared to the CL31 measurements, but the
493 simulated ABL has a similar trend to CL31. The input from the mesoscale simulation forced LES
494 from scenario C, which affected the simulated ABL structure. We can therefore consider that
495 applying the one-way nesting strategy could increase the noise in the ABL structure due to
496 complex topography that is not resolved in the scenario C. Particularly, in the scenario D (Figure
497 3d), the boundary between the free and moist atmospheric layers followed the trend of the CL31
498 measurement. Scenario E presented a relatively shallow ABL. According to our hypothesis, it is
499 likely that scenario E will produce a cloudier sky. The ABL from each scenario, depending on its
500 features, developed CC convection more or less consistent with the observed. To set this idea
501 down, we discuss the convection that occurred between 16:00 and 20:00 LST in the next
502 subsections.

503 3.3. Pattern Reflectivity and FY-2D Satellite Image

504 We compared the simulated reflectivity (Figures 4b,e) to the radar reflectivity from TIPEX-
505 III performed over Naqu2 (Figure 4a). The maximum intensity of reflectivity in scenarios A and
506 B1 was similar to the intensity (10 dBZ) of the observed reflectivity, whereas in scenarios B, C,
507 D and E, the model overpredicted the intensity of the observed reflectivity by 30 dBZ. We
508 verified from the observation that very weak or almost no precipitation occurred over Naqu2
509 between 15:04 and 18:46 LST. By contrast, the model presented maximum reflectivity before
510 18:00 LST in all scenarios except scenario C. Comparing the numerical pattern reflectivity in
511 each scenario to the observed, the model anticipated the observed reflectivity over Naqu2 by at
512 least three hours, which can be seen in the time series of surface variables shown earlier in
513 Figure 3. The reflectivity in scenario C reflects the observed better than in other scenarios. The
514 reflectivity in scenarios D and E (Figures 4f,g) reached a height above that observed in scenario
515 C (Figure 4d) with similar intensity. But the maximum precipitation (the orange shaded
516 reflectivity) observed from TIPE-III (Figure 4a) around 19:00 LST at the Naqu site was rather
517 predicted just after 16:00 LST. The little improvement relative to the prediction time of the
518 precipitation intensity in scenario C may be related to the combination of the one-way nesting
519 and observation nudging strategies. Note that scenarios D and E have not applied a one-way
520 nesting strategy. On the other hand, the earlier evaluation result showed that scenarios C, D and
521 E overpredicted the surface wind speed and hence the latent heat flux and intense precipitation.
522 Wind speed depends on the large-scale pressure gradient force and the local geography. In
523 addition, as the wind speed is modulated by the large-scale forcing, it is obvious that scenarios D

524 and E presented different results from scenario C because they performed LES without a one-
 525 way nesting strategy. Furthermore, the difference between scenario D and E is because scenario
 526 E has not used observation nudging strategy.



527

528 **Figure 4:** Vertical cross-section of pattern reflectivity on day 3 of simulation period: (a) the
 529 observed radar reflectivity from TIPEX-III between 18:45 and 20:00 LST on July 19, 2014; (b–
 530 e) simulated reflectivity. Cases A, B and B1 from $D0_{i=3} = D03$, C, D and E from $D0_{i=5} = D05$
 531 represent scenarios in Table 4.

532 3.4. The Meso- and Micro-scales Liquid Water Content and The Vertical Wind Shear

533 We investigated the types of CCs and the effect of the vertical wind shear on their
 534 development in each scenario, based on the amount of the liquid water content (LWC). Vertical
 535 wind shear is a factor that determines the origin of clouds forming over a given region, while the
 536 amount of LWC in the atmosphere determines what types of clouds form. We illustrated the
 537 LWC in Figure 5 at 200, 300 and 400 hPa levels, respectively, each overlapped with the vertical
 538 wind shear. Note that Fengyun's satellite images (Figures 5a,b) are plane projections including
 539 clouds from the low level to the tropopause. The white parts of Fengyun's satellite images

540 represent the CCs. The dark parts of the images represent no clouds. We depict LWC and wind
541 shear at 16:48 LST. This analysis emphasizes CC and cumulonimbus clouds development. A
542 study by Hess et al., (1998) classified the types of clouds associated with the amount of LWC
543 over land and recommended $0.25 \text{ g/m}^3 \leq \text{LWC} \leq 0.3 \text{ g/m}^3$ for CC and stratocumulus clouds, and
544 $1 \text{ g/m}^3 \leq \text{LWC} \leq 3 \text{ g/m}^3$ for cumulonimbus. Scenarios A and E were useless in this analysis due
545 to their uncertainty with the observed clouds. Moreover, these scenarios were not the target
546 scenarios so far.

547 In scenario, B, Figures 5c,d shows a spatial distribution of LWC similar to that of the cloud
548 distribution in Fengyun's satellite image (Figure 5a). However, compared to the satellite image,
549 scenario B shows scattered small CCs and some Cumulonimbus with $\text{LWC} \geq 1 \text{ g/m}^3$ from the
550 mid-level (Figure 5d) to the upper level (Figure 5e) at the northwest and southeast flanks,
551 respectively. Scenario B also presents a strong vertical wind shear. As we can see in Figure 5c,
552 there is a southwesterly low-level jet (LLJ) at 400hPa. The LLJ is a monsoonal wind that
553 provides moist air favorable for cloud formation. The southeasterly wind dominates the middle-
554 level (300 hPa). The upper-level (200 hPa) is dominated by the easterly wind with strong
555 horizontal wind shear. We can also see at 200 hPa level that the southeast wind is abruptly
556 deflected west by the northeast wind.

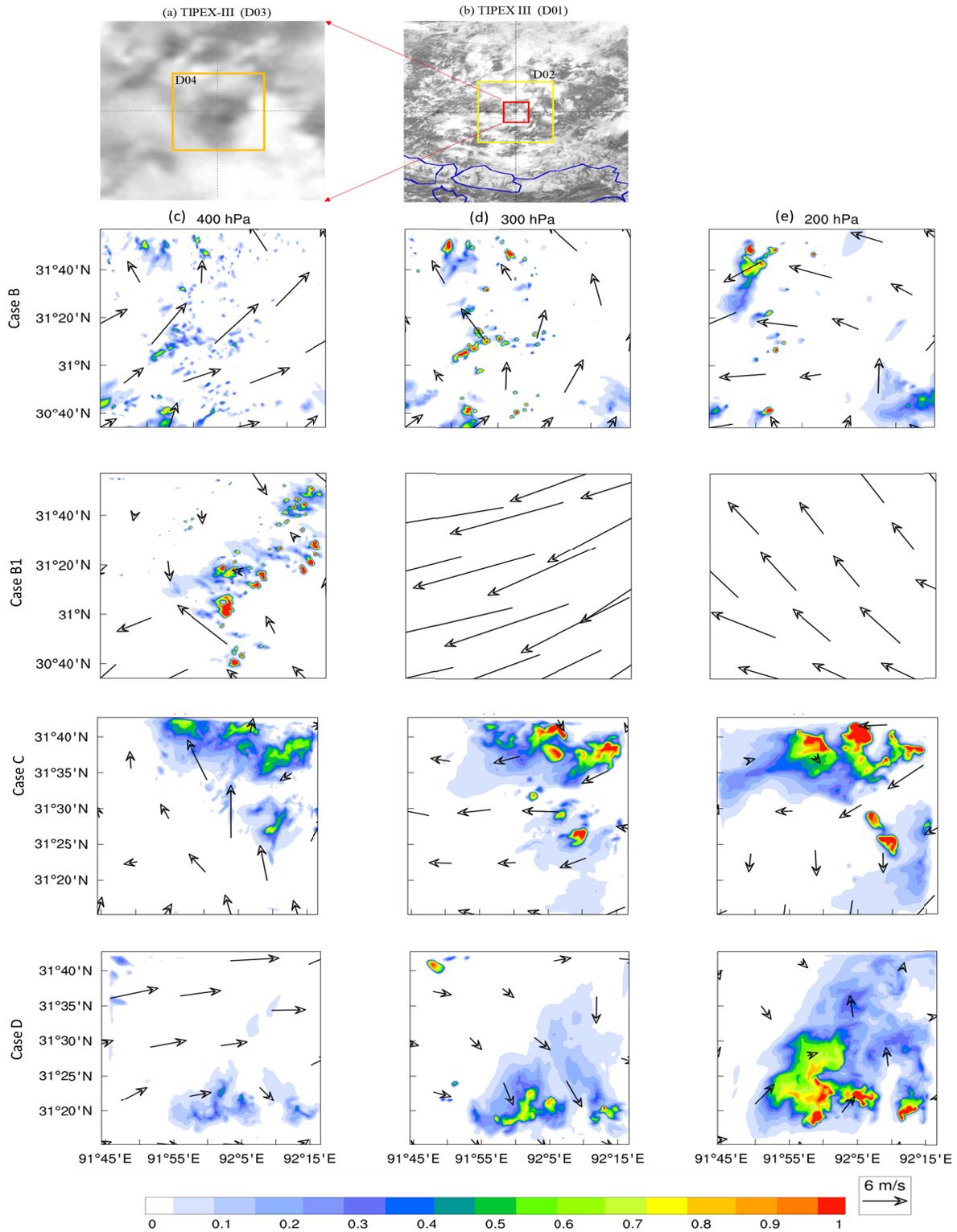
557 Scenario B1 presents an intense mesoscale convective system at the 400 hPa level (Figure
558 5c), but subsidence dominates at 300 hPa levels. In the mesoscale convective system, Figure 5c
559 of scenario B1 shows favorable conditions ($\text{LWC} > 1 \text{ g/m}^3$) for cumulonimbus clouds. However,
560 the subsidence limits the cumulonimbus to expand vertically due to a strong northwesterly
561 upper-level jet (ULJ). In this sense, the large-scale forcing played an important role in cumulus
562 development on July 19, 2014. Indeed, in scenario B1, we have emphasized external forcing by
563 increasing the number of model levels in the ABL, but the results are not consistent with the
564 observation.

565 The scenario C presents the results of LES with one-way nesting from $DO_{i=4}$. This scenario
566 has developed cumulonimbus with $\text{LWC} \geq 1 \text{ g/m}^3$ and has a spatial distribution of the LWC
567 (Figure 5e) similar to the observed clouds (Figure 5a). Noted that this scenario presented the
568 ABL water vapor relatively consistent with the observation. Scenario C predicted low and upper-
569 level wind differently from scenario B. There is southeasterly wind at 400 hPa level (Figure 5c),
570 easterly wind at 300 hPa level (Figure 5d), and northeasterly wind that predominates at the 200
571 hPa level. This study assumes that the wind field in scenario C from 200 to 400 hPa levels is
572 more consistent with the observed than in scenario B because scenario C presented better clouds
573 pattern than scenario B. Thus, the biased surface wind from scenario C discussed in section 3.1)
574 did not influence the spatial distribution of the LWC. This is true when referring to the results
575 from scenario B1, meaning that the internal forcing dominated the CC development on July 19,
576 2014 at the Naqu site.

577

578

579



580

581 **Figure 5:** Liquid water content (LWC) and vertical wind shear at 16:48 local standard time
 582 (LST); vector field stands for wind direction and shaded represents LWC (g/m^3). Classification
 583 cloud types associated with the amount of LWC can be found in (Hess et al., 1998). Cases B
 584 and B1 from $D0_{i=3} = D03$, and C and D from $D0_{i=4} = D04$ represent the scenarios in Table 4.

585 The scenario D (Figures 5d,e), LES without a one-way nesting strategy presents
586 cumulonimbus clouds from $D0_{i=4}$ with $LWC \geq 1 \text{ g/m}^3$. The atmosphere was quite turbulent due
587 to pronounced wind shear. For example, the west wind at the 400 hPa level (Figure 5d) changed
588 to a north wind at the 300 hPa level. By contrast, the wind at the 200 hPa level (Figure 5e) has no
589 specific direction. We note that scenario D has the best distribution of the water vapor mixing
590 ratio in the ABL. When compared to scenarios B, B1, and C, scenario D was able to handle a
591 complex orography. However, this scenario failed to reproduce the observed clouds.

592 4. Summary and Conclusions

593 We explored the ability of the combined weather research and forecasting large-eddy
594 simulation (WRF-LES) in the simulation of the atmospheric boundary-layer (ABL) processes of
595 cumulus clouds (CCs) development over the south Tibetan plateau (STP). The goal of this study
596 was to increase our understanding of the WRF-LES's ability to predict the ABL processes of CCs
597 development over the complex orography, which is an unresolved problem in numerical
598 modeling. We carried out simulations of six scenarios with different setups and compared the
599 results to observations from the third Tibetan Plateau Atmospheric Science Experiment (TIPEX-
600 III), including Lidar ceilometer measurements, Fengyun 2D satellite images, radar reflectivity,
601 sounding, and surface-base measurements such as temperature, wind, and heat fluxes. The new
602 findings were obtained by applying observation nudging strategies and one-way nesting in two
603 separate iterations.

604 Six scenarios, with mesoscale simulation, runs in scenarios A, B and B1, and large-eddy
605 simulation (LES) runs in scenarios C, D and E, are discussed. Scenarios A and E were
606 benchmark experiments to help determine the optimality of observation nudging used in
607 scenarios B and D, respectively. We applied observation nudging in scenarios the B, B1, C, and
608 D. scenario B determined the impact of observation nudging on the simulation results. Scenario
609 C determined the effect of large-scale forcing on microscale circulation. Scenario C differed
610 from scenario D only in the use of the one-way nesting strategy, while scenario D was a control
611 experiment to help determine the optimality of the one-way nesting strategy used in scenario C.

612 Among the six scenarios carried out, the model setting with observations nudging yielded
613 better results than those without nudging. Compared to the observations, the control scenarios A,
614 B1 and E do not well reflect the field's observations. Scenario B1, which emphasized an
615 external forcing, did not yield consistent results with the observations. As a result, large-scale
616 forcing played an important role in CCs development over TP. The evaluation results on one
617 hand also showed that scenario B perfectly reproduced the surface variables nevertheless with
618 low microphysical particles while on the other hand scenario D reproduced an ABL, which is
619 consistent with the observations however this scenario could not correctly reproduce the
620 microphysics patterns. The scenario C had a relatively good representation of ABL and better
621 reproduced the microphysical pattern, which also confirms the role of the large-scale forcing in
622 cumulus development over STP.

623 For CC simulations, we recommend scenarios B and C, which generally captured the
624 overall joint effects of ABL processes and cumulus cloud development, observed on July 19,
625 2014. However, the WRF-LES presented some biases since the reliable data are still insufficient
626 such as sounding that take place twice a day, which is a very low temporal frequency. Therefore,

627 we suggest that for a scientific purpose and clouds simulation over STP, scenarios B and C can
628 still be improved, by focusing on the model's response to the terrain and meteorological initial
629 and boundary conditions. There is an interest in statistically stabilizing the model on the spin-up
630 because the coupled WRF-LES would require the mesoscale simulation data to be reliable. In
631 this case, this study performed a series of simulations by varying the spin-up time as preliminary
632 work to determine the time interval that best matched this simulation. Based on these results, part
633 B of this paper will discuss the development mechanism of the deep cumulus convection over
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