# Lithosphere Removal in the Sierra Nevada de Santa Marta, Colombia

David Ernesto Quiroga<sup>1</sup>, Claire A. Currie<sup>1</sup>, and Jillian Pearse<sup>2</sup>

<sup>1</sup>University of Alberta <sup>2</sup>Universidad de Los Andes

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#### Abstract

The Sierra Nevada de Santa Marta (SNSM) in northwestern Colombia is one of the world's highest coastal mountains, with an elevation above 5.7 km. Gravity measurements show that the SNSM has a high Bouguer anomaly (>+130 mGal), indicating that the mountain lacks a crustal root. In this work, we test the hypothesis that these observations can be explained by gravitational removal of the dense lower lithosphere. We use 2D numerical models to examine the dynamics of lithosphere removal and its effect on surface elevation, gravity and heat flow. The models consist of continental lithosphere that includes a pre-thickened crustal region, representing the SNSM. In our preferred model, the dense mantle lithosphere instability and crustal root are gravitationally unstable and undergo removal as local drips within 10 Ma from the onset of foundering. This creates an area of thinned crust (38 km) underlain by a buoyant sublithospheric mantle where melting and low seismic velocities are predicted. Subsequent non-isostatic forces maintain a topography of 3.3 km with a Bouguer gravity anomaly of +103 mGal. Parameter tests show that a strong lower-crustal rheology provides greater support for the high topography and that a weak mantle lithosphere rheology produces faster removal. The models demonstrate that local lithosphere dynamics can explain the first-order observations in the SNSM. We propose that lithosphere removal could have occurred at 40-50 Ma, possibly inducing anomalous short-lived Eocene magmatism, or more recently (2 Ma), explaining the localized low seismic velocity zone below the SNSM.

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## D. E. Quiroga<sup>1</sup>, C. A. Currie<sup>1</sup> and J. Pearse<sup>2</sup>

<sup>1</sup>Department of Physics, University of Alberta, Edmonton, AB T6G 2E1, Canada <sup>2</sup>Department of Geociences, Universidad de los Andes, Bogotá, Colombia

## Key Points:

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7	•	Mantle lithosphere and crustal root are gravitationally removed as local drips, gen-
8		erating a thinned crust within 25 Ma
9	•	Removal event can explain the observed high topography and positive Bouguer
10		gravity anomaly for the Sierra Nevada de Santa Marta, Colombia
11	•	Non-isostatic forces have a significant role in the support of high topography dur-
12		ing removal

 $Corresponding \ author: \ David \ Quiroga, \ \texttt{dquirogaQualberta.ca}$ 

#### 13 Abstract

The Sierra Nevada de Santa Marta (SNSM) in northwestern Colombia is one of the 14 world's highest coastal mountains, with an elevation above 5.7 km. Gravity measurements 15 show that the SNSM has a high Bouguer anomaly (>+130 mGal), indicating that the 16 mountain lacks a crustal root. In this work, we test the hypothesis that these observa-17 tions can be explained by gravitational removal of the dense lower lithosphere. We use 18 2D numerical models to examine the dynamics of lithosphere removal and its effect on 19 surface elevation, gravity and heat flow. The models consist of continental lithosphere 20 21 that includes a pre-thickened crustal region, representing the SNSM. In our preferred model, the dense mantle lithosphere instability and crustal root are gravitationally unstable and 22 undergo removal as local drips within  $\sim 10$  Ma from the onset of foundering. This cre-23 ates an area of thinned crust ( $\sim$ 38 km) underlain by a buoyant sublithospheric mantle 24 where melting and low seismic velocities are predicted. Subsequent non-isostatic forces 25 maintain a topography of 3.3 km with a Bouguer gravity anomaly of +103 mGal. Pa-26 rameter tests show that a strong lower-crustal rheology provides greater support for the 27 high topography and that a weak mantle lithosphere rheology produces faster removal. 28 The models demonstrate that local lithosphere dynamics can explain the first-order ob-29 servations in the SNSM. We propose that lithosphere removal could have occurred at 40-30 50 Ma, possibly inducing anomalous short-lived Eocene magmatism, or more recently 31  $(\sim 2 \text{ Ma})$ , explaining the localized low seismic velocity zone below the SNSM. 32

#### <sup>33</sup> Plain Language Summary

Most high mountain regions are formed during continental deformation near tec-34 tonic plate boundaries, creating an area of thick low-density crust. The Sierra Nevada 35 de Santa Marta (SNSM) is an unusual mountain in northwestern Colombia. It is isolated 36 from the main Andes Mountain belt, but it has an elevation above 5.7 km and geophys-37 ical measurements indicate that the crust is unexpectedly thin. Here, we test the idea 38 that the SNSM once had a thick crust, but the lower part of this layer was anomalously 39 dense and sank into the deep mantle. We use 2D computer simulations to investigate 40 this phenomenon and its effect on the Earth's surface. The models start with a  $\sim 125$ 41 km wide area with a thicker crust and a deflection in the underlying rigid mantle litho-42 sphere. Because the mantle lithosphere is cooler and denser than the mantle below, it 43 sinks as a drip. This process induces an increase in the density of the lower crust such 44 that it also sinks as a drip. The models show that the removal leaves a high elevation 45 region with a locally thin crust that is underlain by a low seismic wave velocity zone and 46 melting, consistent with SNSM observations. 47

## 48 1 Introduction

The Sierra Nevada de Santa Marta (SNSM) in northern Colombia has a maximum 49 height of  $\sim 5.7$  km (Figure 1a). This triangular shaped mountain region is isolated from 50 the Andes Cordillera, and yet it has the highest peak in Colombia. Early geophysical 51 observations show that this region has a positive Bouguer gravity anomaly of more than 52 +130 mGal (Case & Macdonald, 1973) (Figure 1). This is unexpected because high moun-53 tains are typically supported by thick crustal roots, producing a negative gravity anomaly. 54 Consequently, it is argued that the high elevations in the SNSM massif are not supported 55 by an Airy-type crustal root (e.g., Ceron-Abril, 2008; Montes et al., 2005; Villagómez 56 et al., 2011) and instead, that there is a thin crust to explain the positive gravity anomaly 57 (Case & Macdonald, 1973). Indeed, crustal thickness estimations based on wide-angle 58 seismic refraction, seismic receiver functions, and surface geology, suggest a locally thin 59 crust (i.e., 30-35 km) for the SNSM (Sanchez-Rojas & Palma, 2014). This raises the ques-60 tions of how the massif is mechanically supported and why there is no crustal root. 61



Figure 1. (a) Topographic map of the study region in northern Colombia. BGB: Baja Guajira Basin, OF: Oca Fault, SNSM: Sierra Nevada de Santa Marta, LMB: Lower Magdalena Basin, CRB: Cesar-Rancheria Basin, PR: Perijá Range, SMBF: Santa Marta-Bucaramanga Fault. The barbed line is the plate margin where the Caribbean plate (CAR) converges with South America. (b) Bouguer gravity anomaly map from the EIGEN-GL04C Global gravity field model (Förste et al., 2008). Profiles of topography and gravity anomaly from A to A' are used to compare the models with observations.

One possibility is that the underthrusting of the buoyant Caribbean plateau and 62 oceanic seamounts could provide support for the SNSM topography (e.g., Case & Mac-63 donald, 1973; Montes et al., 2005). In addition, Montes et al. (2005) propose that the 64 Caribbean slab could have caused an upward bulge of the continental Moho, creating 65 a thin SNSM crust and positive gravity anomaly. However, if the SNSM is directly un-66 derlain by the Caribbean slab, a locally low surface heat flow would be expected because 67 of conductive cooling by the slab. However, the surface heat flow for the SNSM is 60 to 68  $80 \text{ mW} \cdot \text{m}^{-2}$  (Quintero et al., 2019), indicating a hot lithosphere. Studies by Quintero 69 et al. (2019) and Vargas et al. (2009) using Curie point depth estimations and direct bore-70 hole temperature measurements, respectively, also suggest that the lithosphere is rela-71 tively hot. 72

Another possibility is that the SNSM topography is laterally supported by a rigid and strong continental lithosphere, based on elastic thickness values of more than 26 km (Arnaiz-Rodríguez & Audemard, 2014). Tassara et al. (2007) finds elastic thickness values of 30 to 40 km for northern Colombia. While the presence of a rigid lithosphere could provide mechanical support for the high topography, it does not explain the high magnitude positive Bouguer gravity anomaly, meaning that there must be an additional factor that creates an excess of mass.

Another consideration is the timing of development of the high SNSM elevations). Villagómez et al. (2011) used fission track thermochronology to determine the timing of rock exhumation. They find that between the Late Eocene and Early Oligocene (40-30 Ma) there was exhumation at a rate  $<0.5 \text{ mm}\cdot\text{yr}^{-1}$ . The last period of exhumation occurred between the Late Oligocene and Miocene (30-16 Ma), where high exhumation rates (>0.5 mm\cdotyr^{-1}) are attributed to underthrusting of the Caribbean plate below northern South America produced by increased convergence rates between the North and South American plates. Additionally, Villagómez et al. (2011) find no evidence of significant exhumation in the last 16 Ma, which is inconsistent with the high topography and high annual rainfall (1000-2000 mm·yr<sup>-1</sup>; Peraza (2014)) in this region, and therefore they argue that this reflects uplift in the last 1-2 Ma. From their work, the exhumation between 40-30 Ma and the recent uplift phase ( $\sim$ 2 Ma) are not related to regional tectonics and the driving forces remain unclear.

Here, we explore the idea that the SNSM observations can be explained by local 93 gravitational lithosphere thinning. Previous work has shown that gravitational founder-94 ing of the deep lithosphere as a local drip can simultaneously produce lithosphere and/or 95 crustal thinning, surface uplift, and symmetric high topography (e.g., Göğüş & Pyskly-96 wec, 2008; H. Wang & Currie, 2017). This mechanism has been proposed to explain un-97 usual surface expressions in various mountain ranges, such as the Sierra Nevada, Cal-98 ifornia, the Puna Plateau of South America, the Wallowa Mountains in northeast Ore-99 gon, and the central Anatolian plateau (e.g., Saleeby et al., 2012; DeCelles et al., 2015; 100 Hales et al., 2005; Gögüş et al., 2017). Removal of the lower lithosphere can be produced 101 by the growth of perturbations in the cold, dense mantle lithosphere (Houseman et al., 102 1981) or by gravitational instability of localized magmatic/metamorphic eclogite (Jull 103 & Kelemen, 2001; H. Wang et al., 2015). As discussed below, the SNSM region has ex-104 perienced significant shortening, and therefore crustal thickening could have fostered meta-105 morphic eclogitization in the deep crust, simultaneously producing a mantle lithosphere 106 instability. Alternatively, previous stages of arc magmatism in the SNSM region (e.g., 107 Cardona et al., 2010; Duque-Trujillo et al., 2019; Ramírez et al., 2020) could have resulted 108 in the formation of dense residues (also called arclogites) after partial melting or frac-109 tionation in the lower crust (Ducea, Chapman, Bowman, & Triantafyllou, 2021). The 110 removal of the dense lower lithosphere would explain the absence of a thick crustal root 111 in the SNSM. 112

In this study we use 2D numerical models to investigate if the observations of high 113 surface topography, positive Bouguer gravity anomaly and high surface heat flow can be 114 explained by lithosphere removal. Our models explore the effect of different lithosphere 115 rheologies on removal dynamics and the surface expressions. We also examine the effect 116 of denser upper crustal rocks (e.g., Sanchez-Rojas & Palma, 2014) and their influence 117 in the gravity anomaly. We calculate the non-isostatic topography in the models and show 118 that non-isostatic forces may provide significant support for the load of the massif. Ad-119 ditionally, we use the models to predict melting patterns and seismic wave velocities in 120 the upper mantle, as other observations that can be used to identify lithosphere removal. 121

#### <sup>122</sup> 2 Tectonic/Geologic Background

The SNSM is part of a larger crustal fragment known as the Maracaibo block and 123 is located at its northwestern corner, approximately 100 km southeast of the tectonic mar-124 gin between the Caribbean and South American plates (Figure 1a). At present, the Caribbean 125 plate is obliquely converging with South America at a rate of 10-20 mm  $\cdot$  vr<sup>-1</sup> toward the 126 southeast (e.g., Freymueller et al., 1993). Regional seismological studies show that the 127 Caribbean plate subducts (or underthrusts) the South American plate at a shallow an-128 gle  $(<30^{\circ})$  and the slab is at a depth of 80 to 120 km below the SNSM (e.g., Cornthwaite 129 et al., 2021; Londoño et al., 2020; Taboada et al., 2000; Van Der Hilst & Mann, 1994; 130 Vargas, 2020). 131

The triangular SNSM massif is bounded by major faults and basins on its three sides (Figure 1*a*). At the north, the SNSM is limited by the right lateral Oca fault and the Baja Guajira basin. At the southwest, it is bounded by the left-lateral Santa Marta-Bucaramanga fault and the Lower Magdalena basin. At the southeast, it is bounded by the Cesar-Rancheria basin and the Perijá range. The Santa Marta-Bucaramanga and Oca faults exceed 200 km in length and have experienced significant lateral displacements since the Paleocene of 100 km and 65 km, respectively (e.g., Case & Macdonald, 1973; Tschanz et al., 1974).

The SNSM is commonly divided into three geo-tectonic provinces (Tschanz et al., 140 1974): the Sierra Nevada, Sevilla, and Santa Marta provinces. According to Ramírez et 141 al. (2020) the oldest rocks in the SNSM are in the Sevilla province (The Buriticá and 142 Los Muchachitos gneisses) and Sierra Nevada province (Los Mangos granulite and Dibulla 143 gneiss). These form the Precambrian metamorphic basement (e.g., Sanchez & Mann, 2015). 144 In the Sierra Nevada province, the basement is overlain by Devonian-Carboniferous and 145 late Triassic-Early Jurassic sedimentary rocks (e.g., Cuchilla de Carbonal and Los In-146 dios formation, respectively). These rocks are mostly covered by outcropping Jurassic 147 igneous bodies (Tschanz et al., 1974), which have intermediate to felsic compositions and 148 mafic intrusions (Ramírez et al., 2020). (Quandt et al., 2018) show that these are con-149 tinental arc plutons associated with a Jurassic magmatic arc. In the Sevilla Province the 150 basement is mostly overlain by Paleozoic metamorphic rocks (e.g., El Encanto orthogneiss, 151 the Sevilla Metamorphic complex). The Santa Marta province is constituted by Creta-152 ceous metamorphic rocks (Santa Marta metamorphic belt) which are intruded by Pa-153 leogene arc-type igneous rocks (e.g., Santa Marta Batholith). These rocks show the evo-154 lution of northwest South America after the Late Cretaceous. This region was influenced 155 by the collision of the Caribbean large igneous province, starting 75 Ma ago (Spikings 156 et al., 2015). This diachronous collision progressed towards the north during the west-157 ward displacement of the South American plate relative to the Caribbean plate (e.g., Bay-158 ona et al., 2011; Pindell et al., 2005). The collision was followed by the onset of the shal-159 low subduction of the Caribbean plate below South America in the early Paleocene (Villagómez 160 et al., 2011), which produced magmatism in the NW corner of the SNSM, forming the 161 Santa Marta batholith and other arc-type plutons (Cardona et al., 2010). 162

#### 163 **3 Methods**

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## 3.1 Governing Equations and Modeling Approach

We use 2D numerical models to study the dynamics of lithosphere removal for the 165 SNSM. The models use the ASPECT code version 2.2.0 (Bangerth, Dannberg, Gassmoeller, 166 Heister, & Others, 2020; Heister et al., 2017). The code solves the governing equations 167 of incompressible mass conservation, the momentum conservation, and the heat conser-168 vation, using the Boussinesq approximation formulation and assuming plane strain (Kronbichler 169 et al., 2012). We model the thermal-mechanical evolution of the lithosphere and upper 170 mantle, where the thermal and mechanical fields are coupled using temperature-dependent 171 material properties. 172

We compare the model results to three key surface observables: surface topography, surface heat flow and Bouguer gravity anomaly. The topography is determined by using a stress-free upper boundary that allows dynamic surface deformation throughout the model evolution (Rose et al., 2017). The surface heat flow is calculated by multiplying the vertical temperature gradient at the surface by the thermal conductivity.

The Bouguer gravity anomaly calculation at each timestep, for each horizontal position along the profile, is based on the density differences (anomalies) with respect to the reference (non-perturbed) density structure in Figure 2b. The contribution to the total gravity anomaly of the density difference in each element in the domain is approximated as the anomaly of a finite slab (Telford et al., 1990) (futher details in Quiroga, 2022).

#### 3.2 Initial Setup and Material Properties

The 2D model domain (Figure 2) has a width of 1320 km and height of 660 km rep-185 resenting the continental lithosphere and the upper mantle in the SNSM region. Our model 186 does not include the subducting Caribbean plate because the goal is to model a litho-187 spheric drip and its isolated contribution to surface observables. The mesh has square 188 elements with lengths and widths of 2.58 km in the upper 200 km and 10 km below. The 189 upper region has a finer mesh to properly resolve lithospheric deformation. The side bound-190 ary conditions are free slip, the bottom boundary condition is no slip, assuming that the 191 upper mantle is coupled to the high-viscosity lower mantle, and the top boundary has 192 a free surface. 193



**Figure 2.** (a) Initial model setup. Only the enclosed region is shown in the following figures. (b) Non-perturbed density structure. (c) Initial temperature conditions. The continental geotherm is stretched in the thickened region keeping the Moho temperature constant. (d) Reference viscosity structure at a constant strain rate of  $10^{-15} s^{-1}$  (black). Black dotted lines show the reference viscous rheologies above plastic yielding including: Wet Quartzite (WQ), Dry Maryland Diabase (DMD), and Dry Olivine (DO). Coloured lines show other viscous rheologies tested including: Eclogite (Ec) with f = 0.1, Mafic Granulite (MG) and Wet Olivine (WO). See Table 1 for references.

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The continental crust is 40 km thick, with a 20 km upper crust and 20 km lower 194 crust. This is underlain by a 60 km mantle lithosphere. An area of initially thick lithosphere is placed in the middle of the model (Figure 2a) to represent the SNSM, assuming the mountain region experienced an earlier episode of shortening. Ramírez et al. (2020) 197 report a continental crust of about 64  $\sim$ km during the Jurassic arc, and subsequent col-198 lision of the Caribbean plateau in the Late Cretaceous-Erly Paleocene is inferred to have 199 produced further thickening (Villagómez et al., 2011). In the model, the thickened region consists of an initial surface topography and crustal root, where the deflections have 201 gaussian shapes with half-widths of 62.5 km, resembling the width of the SNSM. At its 202 maximum, the total crustal thickness in this region is about 35 km thicker than the undis-203 turbed crust. The amplitude of the initial topography is set to be in isostatic balance 204 under the assumption that equilibrium occurred right after or during shortening. The 205 crustal thickening results in a perturbation in the lithosphere-asthenosphere boundary 206 (LAB). In this work, we do not model shortening but rather we look at the dynamic re-207 sponse after crustal thickening. The zero topography corresponds to the non-perturbed 208 surface (upper boundary) at the sides of the thickened region in the modeling setup (Fig-209 ure 2a). 210

All materials have temperature-dependent densities and viscous-plastic rheologies 211 using the material properties given in Table 1. The reference model (model A) is based 212

Parameter	Upper Crust Lower Crust		Eclogite	Mantle Lithosphere	Sublithospheric Mantle		
$Density^{a} \\ \rho_{0} (kg m^{-3}) \\ \alpha (K^{-1})$	$\begin{array}{c} 2800\\ 3\times10^5 \end{array}$	$\begin{array}{c} 2900\\ 3\times10^5\end{array}$		$\begin{array}{c} 3550\\ 3\times 10^5 \end{array}$	$\frac{3400}{3\times 10^5}$	$\begin{array}{c} 3400\\ 3\times10^5\end{array}$	
$\begin{array}{l} Plastic \ rheology^b \\ \phi \\ C \ (MPa) \end{array}$	$\frac{15^\circ - 2^\circ}{20 - 2}$	$\begin{array}{c} 15^{\circ}-2^{\circ}\\ 20-2 \end{array}$		$\begin{array}{c} 15^{\circ}-2^{\circ}\\ 20-2 \end{array}$	$\frac{15^{\circ}-2^{\circ}}{20-2}$	$\begin{array}{c} 15^{\circ}-2^{\circ}\\ 20-2 \end{array}$	
$ \begin{array}{l} Viscous\ rheology^c \\ f \\ A_{ps}\left(Pa^{-n}s^{-1}\right)\ ^d \\ n \\ E\left(kJ\ mol^{-1}\right) \\ V\left(m^3\ mol^{-1}\right) \end{array} $	$\begin{array}{c} Wet \ Quartzite \\ 1 \\ 8.57 \times 10^{-28} \\ 4.0 \\ 223 \\ 0 \end{array}$	$Dry \ Diabase \\ 1 \\ 5.78 \times 10^{-27} \\ 4.7 \\ 485 \\ 0 \\ 0$	$\begin{array}{c} {\it Mafic \ Granulite} \\ 1 \\ 3.17 \times 10^{-21} \\ 3.2 \\ 244 \\ 0 \end{array}$	$\begin{array}{c} Dry \ Eclogite \\ 0.1 \\ 1.18 \times 10^{-17} \\ 3.5 \\ 403 \\ 2.72 \times 10^{-5} \end{array}$	$\begin{array}{c} Dry \ Olivine \\ 1 \\ 1.43 \times 10^{-16} \\ 4.0 \\ 540 \\ 1.5 \times 10^{-5} \end{array}$	$\begin{array}{c} Wet \ Olivine \\ 1 \\ 1.76 \times 10^{-14} \\ 3.0 \\ 430 \\ 1.0 \times 10^{-5} \end{array}$	
$\begin{array}{l} Thermal \; Parameters^{e} \\ k \; (Wm^{-1}  K^{-1}) \\ C_{p} \; (Jkg^{-1}  K^{-1}) \\ Q_{p} \; (W \; m^{-3}) \end{array}$	2.25 1250 $1.0 \times 10^{-6}$	2 11 0.4 >	2525 $250 \times 10^{-6}$	2.25 1250 $0.4 \times 10^{-6}$	$\begin{array}{c} 2.25\\ 1250\\ 0\end{array}$	$\begin{array}{c} 2.25\\ 1250\\ 0\end{array}$	

 Table 1.
 Model parameters for the reference model (model A). The reference model does not include eclogitization. The paremeters for eclogite are used in all the other models in this work.

<sup>*a*</sup> The density ( $\rho$ ) is given by:  $\rho = \rho_0 [1 - \alpha (T - T_0)]$ , where  $\alpha$  is the thermal expansion coefficient, *T* is the temperature (in °*C*) and *T*<sub>0</sub> is the surface temperature ( 0 °*C*).

<sup>b</sup> The frictional-plastic yield stress  $(\sigma_y)$  is based on a Drucker Prager yield mechanism:  $\sigma_y = Psin(\phi) + Ccos(\phi)$ , where  $\phi$  is the angle of internal friction, and C is the cohesion.

 $^{c}$  Where stresses are below the plastic yield stress viscous deformation follows a dislocation creep power law, with an effective viscosity ( $\eta_{eff}$ )

given by:  $\eta_{eff} = \frac{1}{2} f A_{ps}^{-\frac{1}{n}} \varepsilon_{II}^{\frac{1-n}{n}} e^{\frac{E+PV}{RT_k}}$ , where  $A_{ps}$  is the pre-exponential factor scaled to plane strain, n is the stress exponent,  $\varepsilon_{II}'$  is the second invariant of the deviatorie strain rate tensor, E is the activation energy, V is the activation volume, P is the pressure  $T_k$  is the absolute temperature, and R is the gas constant (8.3145 J mol<sup>-1</sup>  $K^{-1}$ ).

<sup>d</sup> The experimental uniaxial strain viscosity pre-factor  $A_{uni}$  is scaled to plane strain  $A_{ps}$  using a scaling factor of  $3^{\frac{n+2}{2}}2^{-1}$ .

<sup>e</sup> k is the thermal conductivity,  $C_p$  is the specific heat capacity, and  $Q_p$  is the internal radiogenic heat production.

on the setup in Figure 2a. The upper crust, lower crust, mantle lithosphere and sub-lithospheric 213 mantle have viscous rheologies of wet quartzite (Gleason & Tullis, 1995), dry Maryland 214 diabase (Mackwell et al., 1998), dry olivine and wet olivine (Karato & Wu, 1993), respec-215 tively. Frictional-plastic deformation includes linear strain weakening, where all mate-216 rials have a reduction in the friction angle and cohesion of  $15^{\circ} - 2^{\circ}$  and 20 - 2 MPa, re-217 spectively, for cumulative plastic strain between 0.5 and 1.5. Other models (sets B, C, 218 and D) examine variations in crustal density and/or lithosphere rheology, and also in-219 clude the effect of lower crustal eclogitization. All the modifications in the models pre-220 sented with respect to the reference model A, are listed in (Table 2). 221

The initial thermal structure is based on a conductive geotherm in the lithosphere, calculated using a surface heat flow of 54 mW·m<sup>-2</sup>, thermal conductivity of 2.25 W·m<sup>-1</sup>·K<sup>-1</sup>, and a radiogenic heat production in the upper and lower crust of 1 and 0.4 W $\mu$ ·m<sup>-3</sup>, respectively; there is no heat production in the mantle (Table 1). This geotherm has an initial Moho temperature of 650 °C and it intersects the mantle adiabat at the LAB (100 km depth) for the undisturbed lithosphere.

Below the LAB, the mantle is adiabatic, with a potential temperature of 1292  $^{\circ}$ C and adiabatic gradient of 0.4  $^{\circ}$ C·km<sup>-1</sup>. The initial geotherm is shown in Figure 2c. In the central region, the continental geotherm is linearly stretched following the gaussianshaped thickening, such that the initial Moho temperature is constant everywhere. During the model run, the model boundaries have fixed temperature conditions using the initial temperature.

#### <sup>234</sup> 4 Eclogitization

Eclogitization is considered to affect the lower crust in areas of thickened crust and can result in a density increase that destabilizes the crustal root, leading to its detach-

Table 2. List of models with parameter modifications with respect to the reference values given in Table 1. The mantle lithosphere rheology is varied locally (in the central region) but remains as dry olivine at the sides for all models.

Model	Eclogitization Lower Crust Local Mantle Lithosphe Rheology Rheology		Local Mantle Lithosphere Rheology	${\rho_{0}^{*}}^{a} {}^{a} (kg  m^{-3})$	Figure No.
Reference					
A	×	$\mathrm{DMD}^{b}$	$\mathrm{DO}^{c}$	2800	4
Set B					
B1	1	DMD	DO	2800	5a
B2	1	$LMG^{e}$	DO	2800	5b
B3	1	$\mathrm{MG}^d$	DO	2800	5c
Set $C$					
C1	1	DMD	$WO^{f}$	2800	6a
C2	1	LMG	WO	2800	6b
C3	1	MG	WO	2800	6c
Set D					
D1	1	DMD	WO	2900	8a
D2	1	LMG	WO	2900	86
D3	1	MG	WO	2900	8c

 $^a~\rho^*_{\bf 0}$  is the reference density in the upper crust in the perturbed region.  $^b$  DMD=Dry Maryland Diabase (Mackwell et al., 1998)

<sup>c</sup> DO=Dry Olivine (Karato & Wu, 1993)

<sup>d</sup> MG=Mafic Granulite (Y. F. Wang et al., 2012)

<sup>e</sup> LMG=Local Mafic Granulite (MG in the central region and DMD at the sides)

 $^f$  WO=Wet Olivine (Karato & Wu, 1993)

ment (e.g., Leech, 2001). Metamorphic eclogitization can occur at temperatures and pres-237 sures greater than  $600-800^{\circ}$ C and 1.0-2.0 GPa, respectively, and the phase change is pro-238 moted by the presence of water (e.g., Ahrens & Schubert, 1975; Austrheim et al., 1997; 239 Leech, 2001). In addition, in continental arcs, eclogite facies rocks can form as dense residues 240 after magmatic differentiation in the crustal root; these are sometimes called arclogites 241 (e.g., Ducea, Chapman, Bowman, & Triantafyllou, 2021). Therefore, it is possible that 242 either metamorphic or magmatic eclogitization affected the SNSM, as it is an area of ear-243 lier shortening and it originally formed as a continental magmatic arc (e.g., Cardona et 244 al., 2010; Ramírez et al., 2020). 245

Model A does not include eclogitization. In model sets B, C, and D, eclogitization in the lower crust occurs when its temperature is above 680°C and its pressure is above 1.2 GPa. The eclogite has a density of 3550 kg·m<sup>-3</sup>, representing full eclogitization of a basaltic lower crust (Austrheim et al., 1997). The rheology of eclogitized lower crust is poorly constrained. Laboratory measurements using a dry synthetic eclogite suggest relatively high strengths (e.g., Zhang & Green, 2007).

On the other hand, eclogitization may be accompanied by rheological weakening, 252 suggesting that eclogites could be weaker than their protoliths and that the overall strength 253 of eclogites is conditioned by the extent and rate of the reaction, which is strongly de-254 pendent on the amount of water available to trigger the phase change or the mechanisms 255 by which water enters the rock (e.g., Austrheim et al., 1997; Leech, 2001). In our mod-256 els, we use the viscous rheology of eclogite of Zhang and Green (2007) with a scaling fac-257 tor (f) of 0.1 which reduces its viscosity by a factor of 10, to represent rheological weak-258 ening relative to the initial diabase rheology due to the phase change (Figure 2d). Ad-259 ditional model tests show that variations in eclogite rheology produce changes in the tim-260 ing of crustal root detachment/removal; weaker eclogite rheologies produce a faster root 261 removal (and vice-versa). However, these variations do not affect the magnitudes and 262 general behavior of surface expressions (Figures S1 and S2 in supplementary material). 263

#### 264 5 Results

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## 5.1 Reference Model (Model A)

The results of model A are shown in Figures 3 and 4 and in Animation 1 (supple-266 mentary material). In this model, the imposed perturbation of the lithosphere-asthenosphere 267 boundary (LAB) causes the lower part of the mantle lithosphere to undergo dripping (Fig-268 ure 3). As the instability at the center founders, the mantle lithosphere is entrained from 269 the sides and pulled towards the center. This results in thinning of the mantle lithosphere 270 by 18 km at the sides. The drip detaches at  $\sim 40$  Ma, removing the lower 30 km of the 271 mantle lithosphere in the central region. The shallower lithosphere is cool and therefore 272 is too viscous to be removed (e.g., Conrad & Molnar, 1999). The crustal root is too buoy-273 ant to be removed but the drip causes a downward pull of the Moho by <3 km as it founders. 274

The evolution of topography on top of the perturbation is characterized by a to-275 tal subsidence of 0.12 km produced by the pull of the dripping mantle lithosphere (Fig-276 ure 4a). The drip detachment at  $\sim$ 40 Ma has a minimal effect on topography, and it is 277 followed by minor topographic changes (<0.1 km) related to continued mantle convec-278 tion. Throughout the model evolution there is a gradual decrease in the gravity anomaly 279 of  $\sim 20$  mGal (Figure 4b). This is mostly attributed to progressive mantle lithosphere thin-280 ning at the sides of the drip as the cold, dense mantle lithosphere is replaced with hot, 281 buoyant sub-lithospheric mantle. There may also be a minor contribution from the down-282 ward deflection of the Moho during the drip. As the drip detaches at 40 Ma, there is an 283 abrupt gravity anomaly increase of 28 mGal produced by the fast foundering of the cold, 284 dense mantle lithosphere. Its effect on the surface gravity anomaly decreases rapidly as 285 the detached mantle lithosphere sinks into the deeper mantle. The surface heat flow in-286



Figure 3. Evolution of the reference model (model A) at different times, showing the effect of the mantle lithosphere drip. (a) Surface Topography profiles (H). (b) Bouguer gravity anomaly profiles  $(g_B)$ . (c) Surface heat flow profiles  $(q_0)$ . (d) Snapshots of the density and thermal structure. Zero topography corresponds to the surface of the modeling domain.

creases from 50 to 58 mW·m<sup>-2</sup> during the model evolution, due to internal radiogenic heating and conductive heating of the crust following lithosphere thinning (Figure 4c).

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#### 5.2 Effect of Lower Crustal Eclogitization (Model B1)

The results of model B1 are shown in Figures 4 and 5a and in Animation 2 (sup-290 plementary material). This model considers eclogitization of the deep crust; all other pa-291 rameters are the same as in model A. The mantle lithosphere perturbation grows, drips, 292 and detaches in a similar manner to model A (Figure 5a). However, the phase of insta-293 bility growth prior to detachment is accompanied by progressive eclogitization of the low-294 ermost crust. The deep crust is not initially in the eclogite stability field. Rather the phase 295 change conditions are induced by the pull of the foundering mantle lithosphere and the 296 temperature increase during lithosphere removal. The negative buoyancy of both the eclog-297 ite and mantle lithosphere perturbation results in detachment of the mantle lithosphere 298 drip at 38 Ma, which is 2 Ma earlier than in model A. Eclogitization continues after drip 299 detachment, and the crustal root starts to founder at  $\sim$ 55 Ma because of its higher den-300 sity. From 55 to 71 Ma, both the crustal root and underlying mantle lithosphere drip 301 into the deeper mantle. The eclogitized crustal root detaches in two pulses (Animation 302



Figure 4. Time variation of surface observables averaged over a 50 km width centered over the initial perturbation, for the reference model (black) and model sets B (blue) and C (red). (a) Surface topography (H). (b) Bouguer gravity anomaly ( $g_B$ ). Green arrows indicate the gravity response to the mantle lithosphere drip and black brackets indicate the response produced by crustal root removal. (c) Surface heat flow ( $q_0$ ). Surface heat flow increases after the removal episode (indicated by yellow arrows) as the heat transfers through the crust by conduction.

2). At 59 Ma there is detachment of the lowermost root, which sinks together with the underlying mantle lithosphere. This leaves a  $\sim 4$  km thick eclogite layer that fully detaches from the non-eclogitized crust at 71 Ma and sinks to the deep mantle. Hereafter we refer to the time of root removal as the time when the eclogitized crustal root detaches and is no longer in contact with the non-eclogitized crust. Root removal leaves a thin crust (38 km) in the perturbed region, underlain by the upwelling sublithospheric mantle.

The evolution of topography in the perturbed region is influenced by the combined 310 effects of the mantle lithosphere drip, eclogitization, and removal of the crustal root. From 311 0 to 71 Ma, the negative buoyancy of the growing mantle lithosphere instability, and the 312 progressive eclogitization and foundering of the lower crust produce  $\sim 2.4$  km of surface 313 subsidence (Figure 4a). Subsidence is not steady, but includes a period of higher sub-314 sidence rate  $(0.2 \text{ mm} \cdot \text{yr}^{-1})$  at ~36 Ma preceding detachment of the mantle lithosphere 315 drip. This is followed by relaxation of the downward pull that results in a subsidence rate 316 of  $0.04 \text{ mm} \cdot \text{yr}^{-1}$  by 38 Ma. The pulses of crustal root detachment at 59 and 71 Ma are 317 also associated with temporary decreases in subsidence lasting 3 and 1 Ma, respectively. 318 This is explained by the upwelling of hot and buoyant sublithospheric mantle as the crustal 319 root founders, which allows for temporary support of elevated topographies of 2.5 and 320 2.2 km at 59 and 71 Ma, respectively. After root removal, subsidence continues as the 321 model comes into isostatic equilibrium with the new lithosphere structure, reaching an 322 elevation of  $\sim 1 \text{ km}$  at 100 Ma. 323

Initially, the central region has a Bouguer gravity anomaly of -169 mGal due to the 324 non-eclogitized crustal root (Figure 4b). As the crustal root undergoes eclogitization, the 325 gravity anomaly increases and has a positive value between 47 to 72 Ma. As in model 326 A, the mantle lithosphere drip produces a rapid peak of 12 mGal at 38 Ma. At the on-327 set of crustal root dripping, there is an increase in the Bouger gravity anomaly, reach-328 ing a maximum of +30 mGal at  $\sim 55$  Ma. Detachment of the crustal root between 59 329 and 71 Ma causes small fluctuations in the gravity anomaly between 4 and 16 mGal. Af-330 ter root detachment, the gravity anomaly decreases and becomes negative as the eclog-331 itized root sinks and is replaced by the buoyant sublithospheric mantle. The absence of 332

high-density mantle lithosphere in the central region and the decrease in crustal density
 due to localized heating result in a final Bouguer gravity anomaly of -56 mGal.

Between 0 and 70 Ma, the surface heat flow in the perturbed region is similar to 335 that in model A (Figure 4c). The only difference is that the topographic changes and 336 shallow deformation produced by crustal root removal generate a local decrease of <5337  $mW \cdot m^{-2}$  (Figure 4c). From 70 to 100 Ma, model B1 exhibits a progressive increase in 338 surface heat flow (up to 84.3  $\text{mW}\cdot\text{m}^{-2}$ ). This is caused by upwelling of the hot sublitho-339 spheric mantle and conductive heating of the overlying crust following root removal at 340 71 Ma. The largest increase in heat flow occurs after  $\sim$ 79 Ma because the crust must 341 be heated conductively from below. 342

This model shows that a positive gravity anomaly and a topographic high are simultaneously produced during crustal root removal between 46 and 72 Ma (Figure 4*a*). This trend is observed in the models discussed below, but the timescales and magnitudes vary with different rheologies and/or density parameters.

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## 5.3 Lower Crustal Rheology (Set B)

Lithospheric rheology strongly affects drip dynamics (e.g., Göğüş & Pysklywec, 2008; 348 H. Wang & Currie, 2017). In model set B, we test different lower crustal rheologies (Fig-349 ure 2d; Table 2). The lower crust in model B1 has the rheology of dry Maryland diabase, 350 as in reference model A. Model B2 uses the lower crustal rheology of mafic granulite (Y. F. Wang 351 et al., 2012) in the thickened region and has dry Maryland diabase at the sides. We re-352 fer to this as local mafic granulite, assuming it represents a locally weaker composition 353 associated with the previous magmatic arc. In model B3, the rheology of the whole lower 354 crust is mafic granulite, producing a weak lower crust along the full profile. 355

The general dynamics of the three models are similar, with an initial mantle lithosphere drip, followed by removal of the eclogitied lower crust (Figure 5). Compared to model B1, the weaker lower crust allows for a slight increase in the removal rate, where detachment occurs at 64 Ma and 60 Ma for models B2 and B3, respectively, which are 7 Ma and 11 Ma earlier than in model B1. In models B2 and B3, elevations are consistently lower than in model B1, suggesting less lateral support for the topography owing to the weaker crust (Figure 4*a*).

There are also differences between the models after crustal root detachment. In model 363 B2 after 80 Ma, there is crustal thinning of up to 6 km induced by the upwelling of the 364 sublithospheric mantle because of the locally weak lower crustal rheology. The crustal 365 thinning results in a higher Bouguer gravity anomaly, resulting in a final value of -25 mGal 366 at 100 Ma (about 25 mGal higher than in model B1; Figure 4b). In addition, the final 367 surface heat flow in model B2 is 15 mW m<sup>-2</sup> higher than in model B1 (Figure 4c). This 368 is attributed to the combined effects of the thinner crust and the longer heating time due 369 to earlier root detachment in model B2. 370

In model B3, crustal root removal and localized ascent of sublithospheric mantle 371 in the perturbed region triggers lithosphere delamination in a similar fashion to the mod-372 els of Krystopowicz and Currie (2013) (Figure S3). This occurs because the weak mafic 373 granulite rheology allows for decoupling along the interface between the lower crust and 374 mantle lithosphere. As the mantle lithosphere detaches laterally and sinks, a wide area 375 of thinned lithosphere is produced, where the crust is thinned by  $\sim 6$  km and there is sub-376 lithospheric mantle upwelling. Consequently, Model B3 exhibits an uplift pulse of  $\sim 0.3$ 377 km after lithosphere removal in the perturbed region, fluctuations in the gravity anomaly 378 of up to 50 mGal and surface heat flow that is  $\sim$ 30 mW·m<sup>-2</sup> higher than in model B2. 379

Figure 5. Evolution of (a) model B1, (b) model B2 and (c) model B3, including pro les of surface topography (H), Bouguer gravity anomaly ( $g_B$ ) and surface heat ow ( $q_0$ ), and snapshots of the density and thermal structure, at times before, during and after the main lithosphere removal episode. Model set B includes eclogitization and explores the e ect of di erent lower crustal rheologies.

380 5.4 Mantle Lithosphere Rheology (set C)

Model set C tests the e ect of mantle lithosphere rheology on drip dynamics. Model A and model set B use a dry olivine rheology for the mantle lithosphere. In set C, the mantle lithosphere in the perturbed region has the rheology of wet olivine. This is a simple approach to test a weaker mantle lithosphere (e.g., due to arc-related processes such as hydration and presence of melt). Models C1, C2 and C3 include lower crustal eclogitization and use the same modi cations in the lower crust as models B1, B2 and B3, respectively (Table 2).

The results of model set C are shown in Figures 4 and 6. Model C1 is shown in Animation 3 as a representative model for the set. As in model set B, lithosphere removal is induced by both the mantle lithosphere perturbation and formation of dense eclogite. In model set C, the presence of the weaker mantle lithosphere speeds up the removal process. Root detachment occurs at 25, 26 and 28 Ma for C1, C2, and C3 respectively, 35-42 Ma earlier than in the comparable set B models. In set B, the mantle lithosphere



Figure 10. Comparison between topography and Bouguer gravity anomaly profiles of the preferred model (model D1) during root removal (~2 Ma before full detachment at t = 23 Ma) with observed data. Initial topography (H) and gravity anomaly ( $g_B$ ) profiles (t = 0 Ma) are shown for comparison (blue). (a) Modeled topography (red). Data from the ASTER global digital elevation model (black). Non-isostatic topography (orange). This shows that non-isostatic effects are localized in the elevated region. (b) Modeled Bouguer gravity anomaly (red). Data from the EIGEN-GL04C global gravity field model (black). Observed profiles are extracted along profile line A - A' in Figure 1a.

at least 15 Ma to match the observations of heat flow. However, heat flow measurements in the study region are sparse and have large uncertainties, and additional data should be collected.

#### 6.2 Isostatic vs. Non-Isostatic Topography

The models exhibit an overall subsidence in the central region as the result of the 524 combined effects of the evolving density structure and mantle dynamics. To separate these effects, we calculate the isostatic topography, which is the topography that would be ob-526 served if there was isostatic equilibrium based on the density structure at each time (Fig-527 ure 9d). Figure 9e shows the non-isostatic topography based on the difference between 528 the total topography and isostatic topography. Non-isostatic effects involve the contri-529 bution to topography produced by both lateral lithosphere strength and vertical stresses 530 from the dynamics of the deep lithosphere and underlying mantle. For models D1 and 531 D2, the isostatic topography is significantly lower than the total topography through-532 out the removal episode, whereas model D3 is closer to isostatic balance throughout its 533 evolution. This shows that the non-isostatic forces are more significant in models D1 and 534 D2, compared to model D3. This is possibly due to the support provided by a rheolog-535 ically stronger lower crust in model D1 (e.g., dry diabase), compared to the weaker crust 536 in model D3 (e.g., mafic granulite). In model D2, the perturbed region has a mafic gran-537 ulite lower crust, with dry diabase at the sides, and the role of non-isostatic forces is in-538 termediate between models D1 and D3. 539

In model D1, the non-isostatic topography reaches a maximum of 3.5 km at 23 Ma (~2 Ma before full detachment). Figure 10a shows that at 23 Ma the topography is predominantly non-isostatic. This suggests that the coexistence of a positive Bouguer gravity anomaly and high topography at this time is possible due to the delay in subsidence produced by non-isostatic effects.

#### 6.3 Melting

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Our study primarily considers the effect of lithosphere removal on surface topog-546 raphy, gravity and heat flow. The models also make additional predictions. As shown 547 above, a lithosphere drip results in thinning of the lithosphere and causes upwelling of 548 the sublithospheric mantle to fill the gap created by the drip. Decompression of upwelling 549 mantle and simultaneous lower crustal heating may result in melting of both the man-550 tle and crust, and therefore magmatism may accompany lithosphere removal (e.g., Kay 551 & Mahlburg Kay, 1993; Mahlburg Kay et al., 1994). To assess this, we use the pressure-552 temperature conditions in model D1 to assess whether melting is predicted. Our calcu-553 lations use the dry granite solidus from Elkins-Tanton (2005), and the solidus of partially 554 hydrated peridotite with 0.15 wt% H2O (Katz et al., 2003), to consider lower crustal and 555 mantle melting, respectively. The models do not include the transport of melt to the sur-556 face, nor the effects of melt on rheology and density. 557



Figure 11. (a) Density structure of the preferred model (model D1) at t = 33 Ma (8 Ma after removal). Contours enclose regions where the *P*-*T* conditions are above the wet peridotite and dry granite solidi, for the mantle (purple infill) and lower crust (white infill), respectively. (b) Temperature profile at x=230 km and t = 33 Ma, with the dry granite (Elkins-Tanton, 2005) and the wet peridotite solidus (Katz et al., 2003). (c) Calculations of melt volume per unit length along strike for the lower crust. Local mantle melt volume is the melt volume in a 60 km width centered in the elevated region.

Figure 11a - b shows model D1 at 10 Ma after root detachment (t = 33 Ma) to allow enough time for significant melt production. In the central region, melting of the sublithospheric mantle and deep crust are predicted. There is also a zone between 80 to 100 km depth outside of the main drip region where lower mantle lithosphere has been partially removed and conditions favor mantle melting.

Figure 11c shows the temporal evolution of mantle melt volume (per unit along strike) 563 for both the full model domain and for a local 60 km wide region centered in the elevated 564 zone. This is calculated by monitoring the regions of the model that are above the solidus, 565 assuming a melt fraction per Kelvin above the solidus of 0.3% for the mantle (Katz et 566 al., 2003) and 0.75% for the crust (Annen & Sparks, 2002). The onset of melting in the 567 central region occurs at about 23 Ma, corresponding to mantle upwelling after root de-568 tachment and crustal heating. These calculations suggest that mantle and crustal melt 569 volumes in the central region could be significant, with approximately equal amounts of 570 each. The model also predicts that melts may be generated in the side regions, due to 571 lithosphere thinning induced by lithosphere foundering in the central region. 572

#### 573 6.4 Seismic Velocity Structure

In addition to melts, our models can be used to assess how lithosphere removal may affect the mantle seismic structure. For this, we use the temperature, pressure, and composition of model D1 to calculate the expected seismic P- and S-wave velocity structure. This uses Perple\_X (Connolly, 2005) to obtain the bulk and shear moduli, assuming a pyrolite composition, with a correction for attenuation (van Wijk et al., 2008; Quiroga, 2022). The crust is excluded from this analysis, owing to uncertainties in crustal composition.

The calculated P- and S-wave velocities and their corresponding anomalies are shown 581 in Figure 12 at a time of 23 Ma (time of the maximum Bouguer gravity anomaly dur-582 ing root removal). After removal of the dense crustal root, the gap created in the man-583 tle lithosphere has a clear low velocity anomaly of up to -9.7 % in P-waves and -19.1 % 584 in S-waves, with absolute values of 6.9 and  $3.5 \text{ km} \cdot \text{s}^{-1}$ , respectively. The anomaly cor-585 responds to the percentage difference compared to the average velocity at each depth. 586 Note that the calculations do not include the effects of melt. Figure 12a shows where 587 melt is predicted and therefore the seismic wave velocities could be even lower. These 588 calculations indicate that lithosphere removal may be associated with a significant low-589 velocity anomaly. This could persist for >10 Ma following the removal event, as the litho-590 sphere gap remains hot (Figure 12c). 591



Figure 12. (a) Calculated seismic structure from model D1 at t = 23 Ma, showing P-wave velocities (left), S-wave velocities (center), and vertical profiles at the side (x = 34 km) and center (x = 230 km). Black contours enclose regions where the pressure and temperature conditions are above the peridotite solidus, which could have lower velocities. (b) P and S wave velocity anomalies and profiles at the side and center.

## <sup>592</sup> 7 Discussion

593

## 7.1 Origin of the High Gravity Anomaly in the SNSM

The models presented above demonstrate that gravitational removal of an eclog-594 itized crustal root can produce both a positive Bouguer gravity anomaly and an elevated 595 topography. In model D1, this results in a gravity anomaly and topography of +103 mGal 596 and 3.3 km, suggesting that a lithosphere drip is a viable mechanism to explain the SNSM 597 observations. Based on the model results and geological observations, we suggest that 598 a drip could have occurred at two possible times: in the Eocene (50-40 Ma) or within 599 the last 2 Ma. These possibilities are discussed separately below. However, it is also pos-600 sible that there was lithosphere removal at both times. 601

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## 7.1.1 Lithosphere removal in the Eocene (50-40 Ma)

The first possibility is that lithosphere removal occurred in the Eocene. At this time, the SNSM was located in a volcanic arc region with subduction of the Caribbean plate below South America. In principle, there are not enough constraints to determine a specific time for this lithosphere removal episode. However, a removal event at 50-40 Ma would predate the slab flattening that has been proposed to explain the cessation of Paleocene-Eocene magmatic activity (e.g., Taboada et al., 2000), and therefore this timing has the advantage that a low-angle Caribbean slab would not block lithospheric foundering.

We suggest that removal at this time could have been driven by metamorphic eclogitization of thickened SNSM crust. Alternatively, root densification could have been produced by magmatic processes. Ducea, Chapman, Bowman, and Balica (2021) argue that melt extraction at an arc creates a lower layer of high-density "arclogite" rocks through magmatic differentiation. This process has been shown to promote gravitational removal even in garnet-free arcs where eclogite is not produced (e.g., Jull & Kelemen, 2001), which is possibly the case in the SNSM (e.g., Duque-Trujillo et al., 2019).

There were two main Paleocene-Eocene magmatic events that were separated by 617 a magmatic gap. The first event produced the Paleocene leucogranites (age  $\sim 65$  Ma). 618 e.g., Duque-Trujillo et al. (2019) argues that the protolith of this rocks had a mafic com-619 position, under the stability field of amphibole, garnet, and rutile, which are all common 620 minerals within arclogite assemblages (Ducea, Chapman, Bowman, & Triantafyllou, 2021). 621 The second event created the Eocene Santa Marta Batholith (SMB) at 56 Ma. Duque-622 Trujillo et al. (2019) report that the SMB has substantial ultramafic garnet-free cumu-623 lates. Some of these have been classified as pegmatitic pyroxenites, perhaps showing ev-624 idence of a densified lower crust outside the garnet stability field, which could be grav-625 itationally unstable. Duque-Trujillo et al. (2019) suggest that the SMB magmatism does 626 not correspond to typical Andean-type subduction magmatism because it was localized, 627 low-volume and short term ( $\sim 10$  Ma). Instead, they propose that the melting of thick-628 ened lower crust could produce the required parental magma. A removal event at this 629 time may explain the trends in magmatism. In our models crustal root removal takes 630 place on timescales of  $\sim 10$  Ma following onset of root densification, and that lower-crustal 631 and shallow sublithospheric mantle melts can be generated by decompression and heat-632 ing after removal (Figure 11). This provides a mechanism to explain the anomalous Paleocene-633 Eocene magmatism. Continued heating can also explain the observations of high sur-634 face heat flow today (Quintero et al., 2019). 635

However, one problem with an Eocene removal event is that the surface expressions
must persist for >40 Ma in order to explain the present-day high elevation and positive
Bouguer gravity anomaly for the SNSM. Because the models in this study show only a
short-lived period (~5 Ma) of simultaneous high elevation and positive gravity anomaly,
this may require additional factors such as elastic stresses, mantle upwelling, and support from the present-day Caribbean slab.

## 7.1.2 Recent lithosphere removal (<2 Ma)

The other possibility is that lithosphere removal occurred recently. Results from 643 model D1 show that lithosphere removal produces a simultaneous positive Bouguer grav-644 ity anomaly and high topography similar to those observed in the SNSM. According to 645 the models, the peak in the Bouguer gravity anomaly after lithosphere removal is shortlived (< 5 Ma) suggesting that the current observations are compatible with recent re-647 moval. Additional evidence that favors a recent lithospheric drip comes from seismic ob-648 servations. Seismic tomography studies consistently show low shear wave velocities (<3.8649  $km \cdot s^{-1}$ ) in the SNSM region in the deep crust and shallow mantle (Syracuse et al., 2016; 650 Poveda et al., 2018). These values are consistent with the shear wave velocities predicted 651 for model D1 (Figure 12a). 652

The problem with a recent removal episode is that according to the models, the high 653 Bouguer gravity anomaly and surface topography must be sustained for at least 15 Ma 654 to allow enough heating time to match the observed surface heat flow values of  $60-80 \text{ mW} \cdot \text{m}^{-2}$ 655 (Quintero et al., 2019). Lithosphere removal within the last 2 Ma is predicted to have 656 a lower surface heat flow (50-60  $\mathrm{mW \cdot m^{-2}}$ ), although we note that there are uncertain-657 ties in both the observations and in the model thermal parameters. On the other hand, 658 the elastic thickness in the SNSM is inferred to be 26-40 km (e.g., Arnaiz-Rodríguez & 659 Audemard, 2014; Tassara et al., 2007), which suggests a cool lithosphere and relatively 660 lower surface heat flow ( $<60 \text{ mW} \cdot \text{m}^{-2}$ ) (e.g., Hyndman et al., 2009). Another issue is 661 that our calculations predict that lithosphere removal would induce lower crustal and 662 upper mantle melting, whereas the last pulse of magmatism in the SNSM was at about 54-49 Ma (Duque-Trujillo et al., 2019). However, Figure 11c shows that the melt vol-664 ume produced by lithosphere removal in the perturbed region is only significant >10 Ma 665 after the removal episode. This means that if there was recent lithosphere removal there 666 is possibly a low amount of melting in the mantle and deep lower crust that has not yet 667 reached the surface. 668

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#### 7.2 Origin of High Topography in the SNSM

The models in this study show that a lithosphere removal episode produces sub-670 sidence rather than uplift. Upward forces from the upwelling sublithospheric mantle and 671 lateral strength are only enough to delay subsidence providing some support. This is be-672 cause our models start with a thick crust with initially high topography in the perturbed 673 region and any non-isostatic forces have to overcome the load of the initial topography. 674 However, the topographic variations in the models depend on the initial conditions, which 675 are not well-known for the SNSM. For instance, lithosphere removal can induce uplift in models with an initially flat topography because the upward non-isostatic forces do 677 not have to overcome the load of the initial topography (e.g., Göğüş & Pysklywec, 2008; 678 H. Wang & Currie, 2017). Figure 9e shows that lithosphere removal in our models can 679 create >3 km of non-isostatic topography meaning that non-isostatic effects can poten-680 tially produce uplift depending on the initial conditions. This may explain the observa-681 tions of unexplained exhumation periods determined by Villagómez et al. (2011). 682

683

## 7.3 Additional Considerations

The models presented here investigate the dynamics of lithosphere removal and the associated surface expressions. One limitation is the 2D nature of the models. Previous studies show that 3D lithospheric drips occur faster and may produce greater topographic deflections, (e.g., Houseman & Gemmer, 2007; Pysklywec & Cruden, 2004). Also, Arnaiz-Rodríguez and Audemard (2014) suggest that elastic strength in the crust is important for supporting the SNSM massif. Kaus and Becker (2007) show that the inclusion of crustal elasticity in lithosphere drip models result in larger surface deflections and larger downwelling rates. Future 3D models including elasticity are needed to assess the implications for the general dynamics and surface expressions of lithosphere removal.

Furthermore, the models address only the dynamics of the continental lithosphere 693 and do not include the subducting Caribbean slab or regional tectonics. Seismic stud-694 ies show that the SNSM region is underlain by the Caribbean slab, but there are uncertainties in the slab location and depth. It is unclear how the presence of the Caribbean 696 plate may affect lithosphere removal, especially if the removal event occurred in the last 697 2 Ma. On one hand, the Caribbean plate may hinder the foundering of a crustal root; 698 on the other hand, subduction may induce a lateral drag that promotes removal (e.g., 699 Currie et al., 2015). Future models should include Caribbean plate subduction to assess 700 its effect on the dynamics and structure of the overlying continent, including the SNSM. 701

#### 702 8 Conclusions

This work uses numerical models of lithosphere removal to study the surface ob-703 servations in the SNSM region of northwest Colombia. The models examine localized 704 lithosphere removal below an area of pre-thickened crust and the associated variations 705 in surface topography, Bouguer gravity, and surface heat flow. We find that foundering 706 of a mantle lithosphere perturbation can induce eclogitization in the overlying crustal 707 root, which is then gravitationally removed. This leaves a locally thinner crust ( $\sim 38$  km) 708 underlain by the hot sublithospheric mantle, and results in a local excess of mass, a pos-709 itive Bouguer gravity anomaly, and a progressive increase in surface heat flow as the crust 710 heats through conduction. During removal, the formation of eclogite and the downward 711 pull of the drip produce surface subsidence. During and after removal, the overall topog-712 raphy also depends on non-isostatic forces, such as lateral lithosphere strength and up-713 welling of the sublithospheric mantle. The timescales of removal and topographic changes 714 depend on lithosphere rheology. Firstly, a local strong mantle lithosphere rheology (e.g., 715 dry olivine) results in removal within 70 Ma, while a weak local rheology (e.g., wet olivine) 716 results in removal within 25 Ma and allows less subsidence time than the stronger rhe-717 ology. Also, a strong lower crustal rheology (e.g., dry diabase) provides more lateral sup-718 port and sustains the initial topography for longer compared to a weak rheology (e.g., 719 mafic granulite). We find that crustal eclogitization significantly modifies the drip dy-720 namics. If it is not included (e.g., model A), the instability only involves the lowermost 721 mantle lithosphere, there is no significant thinning of the lithosphere or changes in surface topography, and the Bouguer gravity anomaly remains negative (-150 to -200 mGal) 723 due to the thick low-density crust. The gravity anomaly results confirm earlier studies 724 that argue that the SNSM is not supported by an Airy-type crustal root (e.g., Case & 725 Macdonald, 1973). 726

We propose that the high gravity anomaly and high topography in the SNSM re-727 gion was produced by a lithosphere removal episode involving the eclogitization and re-728 moval of the crustal root. Our models show that the SNSM observations are compat-729 ible with a strong lower crust and a locally weak mantle lithosphere rheology (model D1), 730 where crustal root removal produces a simultaneous positive Bouguer gravity anomaly 731 (+103 mGal) and high topography (3.3 km), with a maximum non-isostatic topography 732 of 3.5 km. In this case, the observed anomaly in the SNSM (>+100 mGal) requires both 733 the high-density shallow crustal rocks (e.g., Sanchez & Mann, 2015) and gravitational 734 lithosphere/crustal thinning. We suggest that lithosphere removal could have occurred 735 either during the Paleocene-Eocene (40-50 Ma), possibly affecting regional magmatism, 736 or very recently (in the last 2 Ma), creating an area of low seismic velocity in the shal-737 low mantle. We prefer the interpretation of a more recent event because this can explain 738 the observed low mantle seismic velocity, high topography and positive gravity anomaly. 739 In either case, the models presented demonstrate that a gravitational removal event is 740 a viable explanation for many of the uncommon observations in the SNSM. 741

## 742 9 Open Research

Our models use the Aspect code version 2.2.0 (Bangerth, Dannberg, Gassmoeller,
 & Heister, 2020). This software is open source and is availeable at: https://github.com/
 geodynamics/aspect/tree/aspect-2.2.

The parameter and input files developed for this study are also availeable at https:// github.com/d-quiroga/Lithosphere-removal-in-the-SNSM. These correspond to models A and B1 as representative examples. The other models use the same inputs with the modifications expained in the text.

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