Decadal and annual variations in meteoric flux from Ulysses, Wind, and SOFIE observations

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November 26, 2022

Abstract

Our solar system is filled with meteoric particles, or cosmic dust, which is either interplanetary or interstellar in origin. Interstellar dust (ISD) enters the heliosphere due to the relative motion of the sun and the interstellar flow. Interplanetary dust (IPD) comes primarily from asteroid collisions or comet sublimation, and comprises the bulk of material entering Earth's atmosphere. This study examines variations in ISD and the IPD flux at Earth using observations from three different satellite techniques. First are size-resolved in situ meteoroid detections by the Ulysses spacecraft, and second are in situ indirect dust observations by Wind. Third are measurements of meteoric smoke in the mesosphere by the Solar Occultation For Ice Experiment (SOFIE). Wind observations are sorted into the interstellar and interplanetary components. Wind ISD show the anticipated correlation to the 22-yr. solar magnetic cycle, and are consistent with model predictions of ISD. Because Wind does not discriminate particle size, the IPD measurements were interpreted using meteoric mass distributions from Ulysses observations and from different models. Wind observations during 2007-2020 indicate a total meteoric influx at Earth of 22 metric tons per day (t d⁻¹), in reasonable agreement with long-term averages from SOFIE (25 t d⁻¹) and Ulysses (32 t d⁻¹). The SOFIE and Wind influx time series both show an unexpected correlation to the 22-yr. solar cycle. This relationship could be an artifact, or may indicate that IPD responds to changes in the solar magnetic field.

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15	Main Points:	
16	1) SOFIE, Wind, and Ulysses give consistent estimates of the meteoric influx at Earth.	
17	2) Annual and decadal variations in Wind interstellar dust observations agree with simulations.	
18	3) Both interstellar and interplanetary dust are correlated to the 22-yr. solar magnetic cycle.	
19		
20	Keywords: Meteoric influx, SOFIE, Wind, Ulysses, interstellar dust, meteoric smoke	

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21 Abstract. Our solar system is filled with meteoric particles, or cosmic dust, which is either 22 interplanetary or interstellar in origin. Interstellar dust (ISD) enters the heliosphere due to the 23 relative motion of the sun and the interstellar flow. Interplanetary dust (IPD) comes primarily from 24 asteroid collisions or comet sublimation, and comprises the bulk of material entering Earth's 25 atmosphere. This study examines variations in ISD and the IPD flux at Earth using observations 26 from three different satellite techniques. First are size-resolved in situ meteoroid detections by the 27 Ulysses spacecraft, and second are in situ indirect dust observations by Wind. Third are 28 measurements of meteoric smoke in the mesosphere by the Solar Occultation For Ice Experiment 29 (SOFIE). Wind observations are sorted into the interstellar and interplanetary components. Wind ISD show the anticipated correlation to the 22-yr. solar magnetic cycle, and are consistent with 30 31 model predictions of ISD. Because Wind does not discriminate particle size, the IPD 32 measurements were interpreted using meteoric mass distributions from Ulysses observations and 33 from different models. Wind observations during 2007-2020 indicate a total meteoric influx at Earth of 22 metric tons per day (t d⁻¹), in reasonable agreement with long-term averages from 34 SOFIE (25 t d⁻¹) and Ulysses (32 t d⁻¹). The SOFIE and Wind influx time series both show an 35 unexpected correlation to the 22-yr. solar cycle. This relationship could be an artifact, or may 36 37 indicate that IPD responds to changes in the solar magnetic field.

38 **1. Introduction**

39 The solar system is filled with meteoric particles, or cosmic dust, which is interplanetary 40 or interstellar in origin. Interplanetary dust (IPD) comes from asteroid collisions or comet 41 sublimation, and is typically bound to solar orbits on the ecliptic plane. Interstellar dust (ISD) 42 enters the heliosphere due to the relative motion of the sun within the local interstellar cloud [Krueger et al., 2019; Sterken et al., 2019]. Meteoroids are constantly entering Earth's atmosphere, 43 44 with larger IPD particles dominating the mass influx. Frictional heating during entry vaporizes a 45 fraction of the particles at altitudes from ~80 to 100 km, and ablation products combine to form 46 nanometer sized meteoric smoke particles that reside in the mesosphere and stratosphere [Plane et 47 al., 2012; Hervig et al., 2017]. Smoke in the mesosphere is enhanced during polar winter due to 48 transport by the mesospheric meridional circulation, as indicated by models [Megner et al., 2008] and satellite observations [Hervig et al., 2009]. Recent estimates of the meteoric influx into Earth's 49 50 atmosphere range from ~15 - 60 metric tons per day (t d⁻¹) [e.g., *Carrillo-Sánchez et al.*, 2020], an 51 improvement over the bewildering range of previous decades (1 - 300 t d⁻¹) [e.g., *Plane et al.*, 52 2012]. The meteoric influx at Earth has implications for atmospheric chemistry, aerosol processes, 53 and ocean productivity [e.g., *Rudraswami et al.*, 2021], motivating further improvements to the 54 understanding of IPD influx and its variability.

The present study examines meteoric flux in the near-Earth environment using observations from Ulysses, Wind, and the Solar Occultation For Ice Experiment (SOFIE). The Wind and Ulysses spacecraft offer long-term records of in situ dust measurements, which are related here to meteoric smoke measurements from the SOFIE satellite instrument. Interpreting the Ulysses, Wind, and SOFIE measurements requires an understanding of the meteoric mass distribution, and dust enhancement due to Earth's gravity and size. The Ulysses and Wind observations contain both IPD and ISD, and separating these is important for estimating the
meteoric influx at Earth. The Wind ISD results are validated through comparisons with ISD
simulations from the Interplanetary Meteoroid environment for EXploration (IMEX) model
[*Sterken et al.*, 2015; *Strub et al.*, 2019].

65 **2. SOFIE observations**

66 SOFIE has conducted solar occultation measurements from the Aeronomy of Ice in the 67 Mesosphere (AIM) satellite since 2007 [Russell et al., 2009]. The measurements are used to 68 retrieve vertical profiles of temperature, five gases (O₃, H₂O, CO₂, CH₄, and NO), polar 69 mesospheric cloud (PMC) extinction at 11 wavelengths, and meteoric smoke extinction at three 70 wavelengths (330 - 1037 nm). SOFIE observes primarily polar latitudes, with the exception of 71 2017 - 2019 when orbital progression caused an excursion through the tropics and a change from 72 sunsets in the Southern Hemisphere (SH) to the Northern Hemisphere (NH) (vice versa for 73 sunrises, see *Hervig et al.*, 2021 for details). The current SOFIE data is version 1.3 which is 74 available online (sofie.gats-inc.com).

75 SOFIE smoke measurements have been used to characterize the variation of smoke in 76 height and time, and revealed the chemical composition of smoke [Hervig et al., 2009; 2017]. The 77 smoke extinctions used here are monthly zonal means, avoiding summer measurements when 78 PMCs contaminate the smoke signal [Hervig et al., 2012]. Extinction is converted to volume 79 density for a smoke composition of olivine (Mg_{2x}Fe_{2-2x}SiO₄, x = 0.4), which is optically detected by SOFIE. Volume density is then used to derive the ablated meteoric influx through comparisons 80 81 with smoke simulations from the Whole Atmosphere Community Climate Model (WACCM) 82 [Bardeen et al., 2008; Hervig et al., 2017, 2021]. SOFIE results during 2007-2020 indicate a global 83 mean ablated influx into Earth's atmosphere of 7 ± 2 metric tons per day (t d⁻¹). Since only ~30%

of incoming meteoroids are ablated [*Carrillo-Sánchez et al.*, 2020], the corresponding total influx (ablated plus surviving material) is 25 ± 7 t d⁻¹. The influx from SOFIE observations are shown in Figure 1, where the results indicate year-to-year variations and greater influx in the NH than in the SH. The hemispheric difference is still not understood, but could indicate an asymmetry in meteoric influx that is not represented in WACCM. The results below consider the meteoric influx from SOFIE as the average of the NH and SH values.



measurements in the Northern and Southern Hemispheres during winter months (Nov-Feb in the NH and May-Aug in the SH). The annual mean for both hemispheres is also shown.

90 **3. Wind observations**

The Wind spacecraft was launched in 1994 to quantify the dynamics of the Sun-Earth system [*Wilson et al.*, 2021]. Wind operates within 1° of the ecliptic plane and has orbited the first Lagrange point (L1) since July 2004 (~1.5×10⁶ km sunward from Earth). Prior to 2004, Wind conducted a variety of orbital maneuvers, including petal orbits through the magnetosphere, lunar flybys, and an excursion to the second Lagrange point [*Malaspina and Wilson*, 2016; *Wilson et* *al.*, 2021]. Many of these periods are not useful for dust measurements, and were screened from
the analyses presented here.

98 Meteoroids are detected when they collide with Wind and a fraction of the spacecraft body 99 is vaporized and ionized [Mann, 2019]. The resulting plasma perturbs the electric potential of 100 spacecraft surfaces [Shen et al., 2021], which is observed by the WAVES electric field antennas 101 and recorded by the Time Domain Sampler (TDS) [Bougeret et al., 1995]. A similar approach has 102 been used for dust measurements by other spacecraft including Voyager [Gurnett et al., 1983] and 103 the Mars Atmosphere and Volatile Evolution Mission (MAVEN) [Andersson et al., 2015]. The 104 Wind dust detector area is the cross-sectional area of the cylindrical spacecraft body (1.8 m height \times 2.4 m diameter), or 4.3 m². Wind is estimated to be sensitive to meteoroids with radii (r) of 0.1 105 - 11 µm. or 10^{-14} to 10^{-8} g in mass (*m*) for a dust density of $\rho = 2.65$ g cm⁻³, but cannot resolve the 106 107 size of individual impactors. Malaspina et al. [2014] noted that the lower and upper mass bounds 108 are uncertain by a factor of 10 or more, due to observational uncertainties and assumptions in the 109 measurement interpretation. Wind reports the number of dust detections per day, which represents particles with m from 10^{-14} to 10^{-8} g. The Wind dust observations are discussed in detail by 110 111 *Malaspina et al.* [2014], *Meyer-Vernet et al.* [2014], *Kellogg et al.* [2016], and *Wood et al.* [2015]. Malaspina and Wilson [2016] describe the archived data which are available online 112 113 (cdaweb.gsfc.nasa.gov). The Wind dust record is summarized in Figure 2, where the monthly 114 average meteoric flux is shown. The results indicate pronounced decadal and annual variations 115 (e.g., more dust in March than September), which are investigated below.



116 4. Ulysses Observations

117 Ulysses operated during 1990 - 2007, and was the first spacecraft to conduct a polar orbit 118 around the Sun. An in situ dust detector used impact ionization to measure the mass of individual particles with m from ~10⁻¹⁶ - 10⁻⁶ g [Grün et al., 1992; Krüger et al., 2006; 2019]. The detector 119 120 sensitive area is quoted as a maximum of 0.02 m² by *Krüger et al.* [2015], which is the value used 121 here. The Ulysses dust observations are illustrated in Figure 3a, where the mass and radius of each 122 particle are shown. The reported particle mass uncertainties are typically a factor of 5 to 10, due 123 to measurement errors combined with uncertainties in the interpretation. Due to the high-124 inclination polar orbit of the Sun, the Ulysses data set consists mostly of ISD. Strub et al. [2015] 125 describe the criteria for identifying ISD in Ulysses, and the present study considered these (in 126 reverse) to find IPD in the Ulysses record. Particles detected at high ecliptic latitudes (b, Figure 127 3b) are most likely interstellar in origin [Krüger et al., 2006], since interplanetary dust is 128 concentrated near the ecliptic plane [e.g., Soja et al., 2019]. Exceptions can occur, however, as 129 dust from Halley type comets and Oort cloud comets has been detected far above the ecliptic plane. 130 IPD in the near-Earth environment was identified by considering 1) spacecraft - sun distances (D_S) 131 of less than 1.5 AU, and 2) relatively low ecliptic latitudes ($|b| < 30^{\circ}$). These criteria limit the 132 Ulysses observations to those near perihelion and also exclude the Jupiter flybys (see Figure 3).

Note that filtering by rotation angle will not separate IPD from ISD because the interstellar flow direction was aligned with the heliocentric prograde motion (i.e., the dominant IPD direction) near perihelion [*Strub et al.*, 2015]. The selection criteria here yields dust observations at an average D_S of 1.36 AU. These observations occurred during four periods (see Figure 3) comprising a total of 0.76 years, which is the observing interval used for the Ulysses IPD flux quantities below.



Figure 3. a) Time series of individual Ulysses meteoroid detections in terms of the particle mass or radius (for $\rho = 2.65$ g cm⁻³). b) The Ulysses ecliptic latitude and longitude. c) Distance from the spacecraft to the Earth, Sun, and Jupiter. Observations used to identify as IPD in the near-Earth environment are indicated by dots ($D_S < 1.5$ AU and $|b| < 30^\circ$, see text for details).

138 **4. Meteoric mass distributions**

139 Interpreting the various measurements and relating them to each other requires an 140 understanding of the meteoric mass distribution. The mass of meteoroids spans many orders of magnitude, with ISD ranging from 10^{-16} - 10^{-10} g ($r \approx 0.1$ - 2 µm, for $\rho = 2.65$ g cm⁻³) and IPDs 141 spanning roughly 10^{-15} to 10 g ($r \approx 0.1 \text{ }\mu\text{m}$ to 1 cm) [e.g., Krüger et al., 2019; Sterken et al., 2015]. 142 Visual meteors are roughly 10^{-2} to 10^3 g ($r \approx 0.1 - 5$ cm) but contribute little to the total meteoric 143 144 mass influx at Earth, and larger bodies (m > 1 kg) appear only on geologic time scales. Grün et al. 145 [1985] (G85) described a meteoric mass distribution based on spacecraft in situ observations, lunar crater analysis, and photometric measurements of the Zodiacal light. The G85 expression yields 146 the cumulative dust flux in free space, $n_{c}(m)$ (g m⁻² s⁻¹), for a given m (g) (i.e., the number of 147 148 particles with mass > m),

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$$n_{C}(m) = (2.2 \times 10^{3} \ m^{0.306} + 15)^{-4.38} + 1.3 \times 10^{-9} \ (m + 10^{11} \ m^{2} + 10^{27} \ m^{4})^{-0.36} +$$

 $1.3 \times 10^{-16} (m + 106 m^2)^{-0.85}$

(1)

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151 The number of particles with a given *m*, n(m), is found from $n(m_i) = n_C(m_i) - n_C(m_{i+1})$. The number 152 distribution is easily converted to a mass distribution, f(m) = m n(m) [*Grün et al.*, 1985].

153 The G85 expression and observations from Wind and Ulysses describe the meteoric flux 154 in free space. Relating this to the flux into Earth's atmosphere requires consideration of Earth's 155 gravity and surface area. The focusing effect of planetary gravity on meteoric dust was described 156 by *Drolshagen et al.* [2017] as the enhancement factor

157 $H_F = v^2 / (v^2 - v_{esc}^2)$ (2)

where *v* is the dust velocity (far from Earth) and v_{esc} is the escape velocity (11.1 km s⁻¹ for Earth) [see also *Jones and Poole*, 2007]. The present study assumed a mean dust velocity of 17 km s⁻¹ [e.g., *Borin et al.*, 2017] giving $H_F = 1.74$, which is the value used below. The statistical uncertainty 161 in H_F is ~20%, for considering v in the range of 14 - 28 km s⁻¹. Earth's surface area is calculated 162 at the typical meteor ablation altitude of 100 km ($S_E = 5.26 \times 10^{14} \text{ m}^2$). With these considerations 163 the total meteoric influx at Earth, F_E , is given by

164

$$F_E = F H_F S_E \tag{3}$$

where *F* is the total flux in free space, $F = \sum f(m)$. Note that this expression is adaptable to the numeric or mass flux distributions as well (e.g., $f_E(m) = f(m) H_F S_E$, assuming constant *v*).

167 Ulysses observations in the near-Earth environment (see Section 4) were used to 168 characterize the IPD mass distribution. The numeric flux distribution from Ulysses is compared to 169 G85 in Figure 4a, along with results from Carrillo-Sánchez et al. [2020] (CS20) which describe 170 f(m) at Earth based on observations and models. The CS20 distribution (their Figure 1a) was related 171 to free space here using equation (3). The distributions from Ulysses and CS20 are in good agreement in the range where they overlap, and a log-linear fit to CS20 for $m < 10^{-7}$ g is very close 172 to the Ulvsses values for smaller particles. The meteoric mass distributions from G85 and Ulysses 173 174 were related to Earth using equation (3), and are compared to the CS20 curve in Figure 4b. The 175 distributions in terms of mass reveal some differences that are not apparent in the number 176 distributions. In particular, both Ulysses and CS20 indicate lower influx than G85 for the smallest particles ($m < 10^{-10}$ g), while none of the distributions agree for the largest meteoroids ($m > 10^{-7}$ 177 g). In the case of Ulysses, there are no observations of $m > 10^{-6}$ g, although the meteoric flux in 178 the interval 10^{-7} to 10^{-6} g is ~4 times greater than for CS20 or G85. Note that the meteoric mass 179 influx at Earth is dominated by particles in the range of 10⁻⁸ to 10⁻³ g. The sum of the distributions 180 in Figure 4b gives the total meteoric influx at Earth, with $F_E = 37$ t d⁻¹ for G85 and 28 t d⁻¹ from 181 182 CS20. Because Ulysses resolves the meteoroid mass, the individual observations can be added to give $F_E = 32$ t d⁻¹. This is slightly more accurate than summing the histogram values, which assign 183

the average mass to each interval. The reported Ulysses particle mass uncertainties are typically a factor of 5 to 10, and this error dominates the uncertainty in F_E . Even with the advantage of counting statistics (error reduction by N^{-2}) the Ulysses F_E uncertainty is large at 76 t d⁻¹. The above estimates are in good agreement with SOFIE observations which give $F_E = 25 \pm 7$ t d⁻¹ on average [*Hervig et al.*, 2021].



Figure 4. a) Meteoric flux in free space versus particle mass (or radii for $\rho = 2.65 \text{ g cm}^{-3}$), as the daily number of meteoroids, per sample area, per mass interval (decade). Results are from Ulysses, Grün et al. [1985], and Carrillo-Sánchez et al. [2020]. The Ulysses distribution is for observations at $D_S < 1.5$ AU and $|b| < 30^\circ$, and a fit to values with $m > 10^{-14}$ g is shown. The CS20 distribution was related to free space using equation (3), and a fit to values for $m < 10^{-7}$ g is shown. The mass range detected by Wind is indicated. b) As in panel (a) except the distributions are in terms of the mass influx at Earth. The fit to Ulysses, $f_E(m) = 10^{8.1} m^{0.95}$ (t d⁻¹, m in g), is similar to the fit to CS20 for $m < 10^{-7}$ g, $f_E(m) = 10^{7.7} m^{0.93}$ (t d⁻¹).

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The meteoric mass and number distributions (Figure 4) were used to interpret the Wind observations, which do not resolve the mass of an impactor, but rather indicate the total flux of particles within the Wind mass detection range, $N_W = \sum_{m_1}^{m_2} n(m)$, with $m_1 = 10^{-14}$ g and $m_2 = 10^{-8}$ g. N_W can be computed from the meteoric size distributions (e.g., Figure 4a), and used to provide and an estimate of the total influx at Earth from Wind,

195
$$F_E = N_W \sum f_E(m) / \sum_{m_1}^{m_2} n(m) = N_W C$$
(4)

In this expression N_W is the Wind observation, and $C = \sum f_E(m) / \sum_{m_1}^{m_2} n(m)$ can be computed from 196 197 a given modeled or measured mass distribution. Note that the summation of $f_E(m)$ is for the relevant range of IPD (roughly 10^{-9} to 10^{-2} g). For the mass distributions in Figure 4, C = 2.1 (t m²) for G85 198 199 and C = 11.3 (t m²) for CS20. One error component in C arise from the factor of 10 uncertainties 200 in the Wind detection limits $(m_1 \text{ and } m_2)$ [Malaspina et al., 2014]. This error component was 201 estimated to be $\sim 29\%$, using the Ulysses mass distributions and perturbing m₁ and m₂ alternately 202 by factors of 10 and 0.1. The total uncertainty in C from the G85 distribution is difficult to estimate, 203 but it is at least 35% due to the combined errors in H_F (20%) and in m₁ and m₂ (29%). The 204 uncertainty in C from CS20 is 70% due to errors in F_E , H_F , and m_1 and m_2 . The value for Ulysses was determined directly by summing the individual observations, giving C = 9.9 (t m²). The 205 uncertainty in C from Ulysses is close to a factor of ~3, due to the uncertainty in F_E (discussed 206 207 above) combined with the other terms. The value of C based on G85 is lower than from Ulysses 208 or CS20 because the G85 curve indicates more dust at the smallest sizes. For comparison, the G85 distribution gives $N_W = 21.7 \text{ d}^{-1} \text{ m}^{-2}$, compared to 4.1 d⁻¹ m⁻² from Ulysses and 2.4 d⁻¹ m⁻² from 209 210 CS20. While the small particles contribute little to the total mass, they contribute greatly to the total number of meteoroids (Figure 4). Equation (4) is used below with Wind observations toexamine the IPD influx at Earth.

213 **5. Decadal and annual variations in meteoric flux**

214 Wind shows an annual cycle in meteoric flux, that varies in amplitude on a decadal time 215 scale (e.g., Figure 2). The flux annual cycle is examined in greater detail in Figure 5, where the 216 amplitude in 2009 is much larger than in 2020. This annual variation has been discussed by others [Kellog et al., 2016; Malaspina et al., 2014; Wood et al., 2015; Zaslavsky et al., 2012], who 217 218 concluded that higher flux near the Vernal equinox is associated with ISD. The reason is that 219 Earth's ram direction is into the interstellar flow in March, and away from it in September [see 220 also *Malaspina et al.*, 2014]. Because the Earth's orbital velocity (~30 km s⁻¹) is similar in magnitude to that of ISD (~26 km s⁻¹), the ISD flux detected by Wind approaches zero during 221 222 September, when the velocities are nearly parallel (Figure 6). Given this pattern, Wind 223 observations near the autumn equinox should represent mostly IPD, and an estimate of the IPD 224 flux was determined from Wind as the mean of observations surrounding the autumn equinox (September 22 ± 45 days). It should be noted that the assumption of a constant ISD velocity here 225 226 is a simplification [e.g., Sterken et al., 2012], and that faster and slower ISD can exist. It is thus 227 possible that Wind measurements in September contain a few ISD, although a more detailed 228 treatment of ISD velocities is beyond the scope of this paper. The ISD flux from Wind was 229 subsequently obtained by subtracting the IPD flux from the monthly mean values during the rest 230 of the year (Figure 5). The approach taken here for isolating IPD and ISD in Wind observations is 231 similar to that used by Wood et al. [2015]. The Wind ISD flux shows a strong annual variation, 232 that is consistent with ISD flux simulations from the IMEX model [Sterken et al., 2015; Krüger et 233 al., 2019]. The IMEX results shown here are for the near Earth environment (1 AU), and ISD with





Figure 5. Meteoric flux as monthly means during a) 2009 and b) 2020. Results are for all observations from Wind, and Wind ISD values which are the monthly mean minus the Wind IPD estimate (the average for September 22 ±45 days). Results from the IMEX model [*Sterken et al.*, 2015] are near 1 AU for ISD with radii from 0.1 - 0.73 μ m.



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Figure 6. Diagram of Earth's position and velocity during the equinoxes and solstices (blue). The interstellar flow direction (259° ecliptic longitude) and velocity are indicated (red).

244 Monthly mean ISD flux during 1995 - 2021 from Wind compare favorably to IMEX 245 simulations (Figure 7a), with both indicating a stronger annual variation and greater overall ISD 246 flux during 2004 - 2016. The trajectories of ISD in the solar system are controlled by solar gravity, 247 solar radiation pressure, and magnetic field strength. Theory predicts that ISD flux should correlate 248 to changes in the solar magnetic field (SMF), due to the Lorentz forces experienced by charged 249 particles in motion [e.g., Landgraf et al., 2000; Sterken et al., 2015]. The SMF varies with a 22-yr 250 periodicity (the Hale cycle, see Figure 7c), where the net effect on ISD is a focusing of particles 251 in the inner heliosphere during solar south pole positive phases. Variations in the ISD flux are 252 indeed coincident with changes in SMF, and linear regression to annual means gives correlation 253 coefficients (p) of 0.74 for Wind and 0.67 for IMEX (Figure 7b). The heliospheric current sheet 254 (HCS) tilt was also examined, which is a representation of the interplanetary magnetic field that 255 exhibits an 11-yr cycle (in phase with Lyman- α flux). The HCS is considered a factor in ISD 256 trajectories [Sterken et al., 2015], although the Wind and IMEX ISD are only weakly correlated 257 with the HCS tilt (p = 0.1 for Wind and 0.34 for IMEX). The good agreement between Wind and 258 IMEX confirms that Wind observations are dominated by ISD during spring. This agreement furthermore suggests that using Wind measurements near the autumn equinox provides a good approximation of IPD, which is important for characterizing the influx at Earth. Note that the Wind and IMEX results indicate an ISD influx at Earth (using equation (3)) of less than 0.01 t d⁻¹, which is insignificant compared to the IPD influx of ~25 t d⁻¹.



Figure 7. a) Time series of interstellar dust flux as monthly means from Wind observations compared to simulations from the IMEX model (divided by 5). b) Solar magnetic field strength at both poles and the heliospheric current sheet (HCS) tilt angle, from the Wilcox Solar Observatory.

263 Meteoric influx at Earth was estimated from the Wind IPD flux determined as above. The 264 Wind F_E estimates use a constant (equation (4)) that is based on integrating the mass distributions 265 from either Ulysses, G85, or CS20 (see Section 4). Wind F_E determined from the yearly IPD flux 266 values are compared to SOFIE, Ulysses, and CS20 in Figure 8a. Wind F_E using C determined from 267 G85 are much lower than when C from Ulysses or CS20, with the later providing Wind influx 268 estimates that are in reasonable agreement with SOFIE, Ulysses, and CS20. Note that both Wind 269 and SOFIE show similar year-to-year influx variations. The influx time series are compared to the 270 solar magnetic field and HCS tilt angle in Figure 8b. Linear regression indicates a strong F_E - SMF 271 correlation for both Wind (p = 0.70) and SOFIE (p = 0.62). The correlation with HCS tilt angle is 272 weaker, however, with p = 0.01 for Wind and 0.42 for SOFIE. The Wind and SOFIE F_E time series 273 are shown in Figure 8c with a fit to the results based on linear regression to the SMF. These results 274 suggest that changes in the SMF can alter the flow of IPD in the inner heliosphere. This result is 275 not anticipated, however, because IPD have small charge-to-mass ratios (Q/m), unlike ISD. While 276 the IPD response to changing SMF would not be instantaneous, it is possible that a cumulative 277 effect is realized after many orbits of the sun. This idea is supported somewhat by the apparent 278 time delay in changing F_E with respect to SMF. Indeed, the highest F_E - SMF correlation is found 279 for a time lag of 14 months, with p = 0.76 for Wind and 0.70 for SOFIE. The agreement between 280 SOFIE and Wind concerning year-to-year and decadal changes in meteoric influx is encouraging, 281 and furthermore suggests that variability in IPD appears in Earth's mesosphere.



Figure 8. a) Time series of total meteoric influx (F_E) at Earth. Wind results are shown for conversion factors based on the dust distribution from either Ulysses, *Carrillo-Sánchez et al.* [2020] (CS20), or *Grün et al.* [1995] (G85). SOFIE results are the average for winter months near 67°N (Nov-Feb) and 67°S (Jun-Sep) [*Hervig et al.*, 2021]. The average influx based on Ulysses IPD observations during 1990 - 2007 is indicated ($|b| < 30^\circ$ and $D_S < 1.5$ AU). The CS20 F_E based on observations of iron and sodium in the mesosphere and cosmic spherules collected at South Pole is shown. b) The North-South average solar polar magnetic field strength and HCS tilt angle, as annual means. c) SOFIE and Wind F_E (*C* from Ulysses), with regressions to the average SMF time series and regression coefficients as listed.

282	The median F_E from Wind during 2007 - 2020 for C determined from the Ulysses, CS20,
283	and G85 results are 22, 33, and 6 t d ⁻¹ , respectively, each with a standard deviation of \sim 14% during
284	the time period. The median F_E based on the three Wind time series during 2007 - 2020 (i.e., Figure
285	8a) is 22 t d ⁻¹ , with a standard deviation of 11 t d ⁻¹ (50%) which is due mostly to the spread in C .
286	The total uncertainty in Wind F_E includes contributions from the error in H_F and uncertainties in
287	the IPD flux measurements, and is closer to 13 t d ⁻¹ (60%). Table 1 summarizes the meteoric influx
288	estimates from this work and some recent publications.

Table 1. Total meteoric influx (ablated + surviving material) at Earth from different sources.					
Source	Method	Meteoric influx (t d ⁻¹)			
Wind	Satellite in situ dust detection (2007-2020), this work	22 ± 13			
Ulysses	Satellite in situ dust detection (1990-2007), this work	32 ± 76			
SOFIE	Satellite remote measurements of meteoric smoke in the mesosphere (2007-2021) [<i>Hervig et al.</i> , 2021]	25 ± 7			
Carrillo-Sánchez et al. [2020]	Model scaled to measurements of Fe & Na in the mesosphere and cosmic spherules from the South Pole	28 ± 16			
<i>Borin et al.</i> [2017]	Long Duration Exposure Facility (LDEF) satellite in situ dust measurements (1984-1990)	15 ± 3			
Gardner et al. [2014]	Lidar measurements of Na in the mesosphere	60 ± 16			

289 **5. Summary**

290 This work examined meteoric influx using in situ dust detection by the Wind and Ulysses 291 spacecraft, and observations of meteoric smoke in the mesosphere by the SOFIE satellite 292 instrument. Wind does not resolve the mass of the detected meteoroids, but rather reports the total number of particles with mass from 10⁻¹⁴ - 10⁻⁸ g. The Wind measurements were separated into 293 294 the interstellar and interplanetary components. The Wind ISD estimates are in good agreement 295 with simulations from the IMEX model, in terms of the both the annual and decadal flux variations. 296 The decadal ISD variation is correlated to the 22-year solar magnetic cycle, as anticipated by 297 theory and predicted by IMEX. The Wind IPD observations were related to the total meteoric

influx at Earth using IPD mass distributions from Ulysses and previous publications. The resulting Wind influx estimates are in good agreement with SOFIE and Ulysses. The SOFIE and Ulysses influx time series show similar year-to-year and decadal changes in meteoric influx. The decadal influx variation exhibits an unanticipated correlation to the 22-year solar magnetic cycle. This relationship may be an artifact, or could indicate that changes in the solar magnetic field can alter the trajectories of interplanetary dust.

Acknowledgements. This work was funded in part by the AIM mission through NASA contract NAS5-03132. V. J. Sterken received funding from the European Union's Horizon 2020 research and innovation program under grant agreement N 851544. L. B. Wilson was partially supported by Wind MO&DA funds. SOFIE data are available online (sofie.gats-inc.com). WIND meteoric dust data are available online (cdaweb.gsfc.nasa.gov). Wilcox Solar Observatory data are available online (http://wso.stanford.edu/#MeanField). Ulysses dust observations are available online (cdaweb.gsfc.nasa.gov).

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