## Drivers and reversibility of abrupt ocean state transitions in the Amundsen Sea, Antarctica

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November 22, 2022

#### Abstract

Ocean warming around Antarctica has the potential to trigger marine ice-sheet instabilities. It has been suggested that abrupt and irreversible cold-to-warm ocean tipping points may exist, with possible domino effect from ocean to ice-sheet tipping points. A  $1/4^{\circ}$  ocean model configuration of the Amundsen Sea sector is used to investigate the existence of ocean tipping points, their drivers, and their potential impact on ice-shelf basal melting. We apply idealized atmospheric perturbations of either heat, freshwater or momentum fluxes, and we characterize the key physical processes at play in warm-to-cold and cold-to-warm climate transitions. Relatively weak perturbations of any of these fluxes are able to switch the Amundsen Sea to an intermittent or permanent cold state, i.e., with ocean temperatures close to the surface freezing point and very low ice-shelf melt rate. The transitions are reversible, i.e., cancelling the atmospheric perturbation brings the ocean system back to its unperturbed state within a few decades. All the transitions are primarily driven by changes in surface buoyancy fluxes over the continental shelf, as a direct consequence of the freshwater flux perturbation, or through changes in net sea-ice production resulting from either heat flux perturbations or from changes in sea-ice advection for the momentum flux perturbation. These changes affect the vertical ocean stratification and thereby ice-shelf basal melting. For warmer climate conditions than presently, the surface buoyancy forcing becomes less important as there is a decoupling between the surface and subsurface layers, and ice-shelf melt rates appear less sensitive to climate conditions.

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## Key Points:

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8	•	The currently warm ice-shelf cavities of the Amundsen sector could become or have
9		been cold for slightly colder climatic conditions.
10	•	The transitions are reversible:cancelling the atmospheric perturbation brings the
11		ocean back to its unperturbed state within a few decades.
12	•	All the transitions are primarily driven, at multi-decadal scale, by changes in sur-
13		face buoyancy fluxes over the continental shelf.

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#### 14 Abstract

Ocean warming around Antarctica has the potential to trigger marine ice-sheet insta-15 bilities. It has been suggested that abrupt and irreversible cold-to-warm ocean tipping 16 points may exist, with possible domino effect from ocean to ice-sheet tipping points. A 17  $1/4^{\circ}$  ocean model configuration of the Amundsen Sea sector is used to investigate the 18 existence of ocean tipping points, their drivers, and their potential impact on ice-shelf 19 basal melting. We apply idealized atmospheric perturbations of either heat, freshwater 20 or momentum fluxes, and we characterize the key physical processes at play in warm-21 to-cold and cold-to-warm climate transitions. Relatively weak perturbations of any of 22 these fluxes are able to switch the Amundsen Sea to an intermittent or permanent cold 23 state, i.e., with ocean temperatures close to the surface freezing point and very low ice-24 shelf melt rate. The transitions are reversible, i.e., cancelling the atmospheric pertur-25 bation brings the ocean system back to its unperturbed state within a few decades. All 26 the transitions are primarily driven by changes in surface buoyancy fluxes over the con-27 tinental shelf, as a direct consequence of the freshwater flux perturbation, or through changes 28 in net sea-ice production resulting from either heat flux perturbations or from changes 29 in sea-ice advection for the momentum flux perturbation. These changes affect the ver-30 tical ocean stratification and thereby ice-shelf basal melting. For warmer climate con-31 ditions than presently, the surface buoyancy forcing becomes less important as there is 32 a decoupling between the surface and subsurface layers, and ice-shelf melt rates appear 33 less sensitive to climate conditions. 34

## 35 Plain Language Summary

The West Antarctic Ice Sheet is under the threat of a partial collapse, which would 36 induce rapid global sea level rise. This threat is partly related to the thinning of float-37 ing ice shelves, and the consequent retreat of the grounding line, which is a self-sustained 38 ice dynamics process. It is triggered by increased basal melting of the ice shelves, which 39 results from enhanced flow of relatively warm waters onto the continental shelf. It has 40 been suggested that self-sustained ocean processes may lead to abrupt changes in the 41 flow of warm water into ice-shelf cavities, which could facilitate the tipping to a marine 42 ice-sheet instability. Here, we analyze whether such abrupt ocean changes can occur un-43 der cold-to-warm or warm-to-cold transitions in the Amundsen Sea, West Antarctica. 44 We use a regional ocean model with a set of idealized local atmospheric perturbations 45 to characterize the thresholds and reversibility of ocean abrupt changes. We find that 46 the currently warm Amundsen Sea could switch intermittently or permanently to a cold 47 state for relatively weak atmospheric perturbations and could be slightly warmer in the 48 future. All transitions are reversible. The main mechanism involved on decadal scale is 49 related to a change in the surface buoyancy fluxes. 50

#### 51 **1** Introduction

The West Antarctic Ice Sheet has lost mass over the last few decades and has thus 52 contributed significantly to global sea level rise. Warming of the oceanic sub-surface seems 53 to have caused an increase in melting under floating ice shelves, particularly in the Amund-54 sen Sea (Jenkins et al., 2018). Depending on the bedrock slope direction (Schoof, 2007; 55 Pattyn et al., 2012) and ice-shelf lateral buttressing (Gudmundsson, 2013), a sufficiently 56 strong and persistent increase in basal melting can lead to a Marine Ice-Sheet Instabil-57 ity (MISI), resulting in a self-sustained retreat of the glacier's grounding line and to the 58 acceleration of its flow (Favier et al., 2014; Joughin et al., 2014). 59

Instabilities are triggered above a certain oceanic warming (critical threshold or tipping point), with the possible existence of multiple thresholds. Thus, Rosier et al. (2021) estimated that Pine Island Glacier would undergo a MISI and major mass loss for an oceanic warming of +1.2°C relative to the present. Garbe et al. (2020) estimated that a tipping point of +2°C global warming relative to pre-industrial could cause a MISI of
the entire West Antarctic Ice Sheet. Tipping points are characterized by a hysteresis,
i.e., restoring the forcing to before the occurrence of the tipping point is not sufficient
to restore the system to its original state. Identifying these tipping points precisely and
linking them to climate projections would allow the effects of future rapid sea level rise
to be anticipated and possibly mitigated (Hinkel et al., 2019).

The abrupt nature of these ice tipping points in West Antarctica could be enhanced if ocean warming itself is subject to a tipping point. This would be a cascading tipping point, or domino effect (Dekker et al., 2018; Brovkin et al., 2021; Wunderling et al., 2021). It has been suggested, that beyond a certain threshold of melting, the Greenland Ice Sheet could induce a sudden weakening of the Atlantic Meridional Overturning Circulation, which, in turn, would lead to ocean warming around Antarctica (Turney et al., 2020; Wunderling et al., 2021).

Another type of oceanic tipping point has been highlighted in the Weddell Sea (Hellmer 77 et al., 2012, 2017). Reduced sea-ice formation under continued global warming, a fresh-78 ening of the continental shelf, and increased ocean surface stress could cause the slope 79 current to diverge in the southeast Weddell Sea. The reorientation would facilitate the 80 entry of Warm Deep Water, a cooler variant of Circumpolar Deep Water (CDW), onto 81 the continental shelf and significantly increase basal melting, which would lead to a self-82 reinforcing process due to the injection of meltwater. The process is irreversible with the 83 twentieth-century atmospheric forcing: only an imposed decrease in basal melt rate can 84 hinder the self-sustaining process. 85

The Amundsen Sea environment is very different as relatively warm cavities already 86 exist (Jacobs et al., 1996, 2012). Paleoclimatic indicators suggest that the entire Amund-87 sen continental shelf was covered by an ice sheet (either resting or floating) at the Last Glacial Maximum (Larter et al., 2014). A particularly large retreat of the ice-sheet front 89 and grounding line occurred between 20,000 and 10,000 years BP (Larter et al., 2014), 90 with further smaller retreats occurring thereafter, notably around 1945 and then 1970 91 (Smith et al., 2017). Ocean temperatures and warming rates during these transitions are 92 not known, but it is possible that oceanic tipping points similar to those reported by Hellmer 93 et al. (2012, 2017) for the Weddell Sea occurred in the Amundsen Sea area as well. 94

In this paper, we analyze under which atmospheric forcing conditions warm-to-cold, cold-to-warm and warm-to-warmer ocean transitions in the Amundsen Sea have occurred or could occur, and we test the reversibility of these transitions, i.e., the presence of hysteresis. We use a regional ocean modelling approach with a set of idealized atmospheric perturbations.

## <sup>100</sup> 2 Materials and Methods

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#### 2.1 Model and configuration

The Nucleus for European Modelling of the Ocean (NEMO) model, version 3.6, in-102 cluding the OPA ocean model (Madec & the NEMO Team, 2016) and the Louvain-la-103 Neuve sea-ice model LIM-3.6 (Rousset et al., 2015), is used in a regional configuration 104 of the Amundsen Sea (Fig. 1). Our model parameters are similar to Jourdain et al. (2019), 105 with a representation of ice-ocean exchange beneath static ice shelves, with melt rate 106 depending on ocean velocity, temperature and salinity (Mathiot et al., 2017; Jourdain 107 et al., 2017), and barotropic tides prescribed as lateral boundary conditions from seven 108 constituents of the FES2012 tidal model (Carrère et al., 2012; Lyard et al., 2006). 109

Compared to Jourdain et al. (2019), the domain is slightly extended, now covering from 142°W to 85°W and from 76.3°S to 59.8°S, and the resolution is reduced to 1/4° in longitude, i.e., a quasi-isotropic resolution ranging from 14 km at the northern bound-



Figure 1. Used regional configuration of the Amundsen Sea. Bathymetry and ice-shelf draft are from the second version of the BedMachine Antarctica dataset (Morlighem et al., 2020). Grounded ice is shaded in white, ice shelves are colored in blue and main tabular icebergs in light cyan. The general view is drawn from the geospatial data package Quantarctica (Matsuoka et al., 2021).

ary to 6.5 km in the southernmost part of the domain. Bathymetry, as well as surface
and lateral boundary conditions also differ from Jourdain et al. (2019) and cover the period 1958-2018 in this study. The period 1958-1968 is left for spin-up and discarded in
our analyses.

The bathymetry and ice-shelf draft interpolated on the model grid are from the sec-117 ond version of the BedMachine Antarctica dataset (Morlighem et al., 2020). This recent 118 dataset represents Thwaites Ice Shelf after its partial collapse. The B22A iceberg as well 119 as other very large tabular icebergs, absent from BedMachine Antarctica, are represented 120 as static flat ice shelves in the middle of the ocean (assumed to be grounded by the subgrid-121 scale bathymmetry when no grounded area is explicitly represented). Their shape and 122 location are derived from a MODIS-visible image (provided by the US National Snow 123 and Ice Data Center) taken on 5th September 2003. The huge B22A iceberg calved from 124 the Thwaites ice tongue in 2002 and has drifted very slowly since then (Antarctic Ice-125 berg Tracking Database, Budge & Long, 2018). A similar calving event occurred in the 126 late 1960s (Lindsey, 1995). The resulting iceberg was eventually designated B10 in 1992 127 when it started a 15-year drift across the Amundsen Sea before breaking up and drift-128 ing further away (Budge & Long, 2018). Numerous smaller icebergs regularly drift west-129 ward in the Amundsen Sea and ground on the eastern flank of bathymetric ridges shal-130 lower than approximately 400 m (Mazur et al., 2017). We therefore artificially place a 131 wall along the 380 m isobath on the eastern flank of Bear Ridge (in a similar way as Bett 132 et al., 2020), north of Siple Island, and on the main ridge in between. These permanent 133 lines of grounded icebergs were shown to favor the formation of polynyas with impact 134 on ice-shelf melting (Nakayama et al., 2014; Bett et al., 2020). 135

The conditions at the lateral ocean and sea-ice boundaries are derived from the 5day mean outputs of a global simulation very similar to the one described in Merino et al. (2018) except that it is spun up from 1958 and that the imposed ice-shelf melt flux increases linearly from 1990 to 2005 and is constant before and after that, with values

corresponding to the FRESH+ and FRESH- reconstructions of Merino et al. (2018). 140 Here, the temperature and salinity boundary conditions are corrected by the difference 141 between the seasonal climatology of the World Ocean Atlas 2018 (WOA18) database (Garcia 142 et al., 2019) and the seasonal climatology of the global simulation. The global simula-143 tion used for boundary conditions represents melting of Lagrangian icebergs (Merino et 144 al., 2016), and the corresponding 5-day mean melt fluxes are applied as a freshwater flux 145 at the surface of our regional configuration. The atmospheric forcing data are taken from 146 the JRA55-do reanalysis (Tsujino et al., 2018) between 1958 and 2018. The fluxes be-147 tween ocean (or sea ice) and atmosphere are calculated using the CORE bulk formulae 148 described in Griffies et al. (2009); Large and Yeager (2004). 149

Some model parameters are varied to reduce biases in the reference configuration (see Supporting Information), while atmospheric forcing fields are perturbed (Section 2.2) to investigate ocean tipping points.

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## 2.2 Atmospheric forcing perturbations

In the following, we investigate three pathways to induce ocean tipping points in the Amundsen Sea through surface flux modifications of either heat, freshwater, or momentum. We decided to consider idealized atmospheric perturbations in order to identify and isolate the processes at play. Thus, each surface flux is perturbed independently.

The heat flux is perturbed through air temperature, to which the flux is particu-158 larly sensitive. To limit the impact of this perturbation on evaporation, and thus on the 159 freshwater flux, specific humidity is also modified consistently with the air temperature 160 perturbation, according to the Clausius Clapeyron law. The choice of air temperature 161 is convenient for the definition of the perturbation range, which is bounded by typical 162 conditions of the Last Glacial Maximum, i.e., approximately -10°C relative to the cur-163 rent temperature (Masson-Delmotte et al., 2010) and by typical projections at 2300 un-164 der the SSP5-8.5 scenario, i.e., about 10°C warmer than the current situation (Lee et al., 165 2021).166

For the freshwater flux, we decide to modify precipitation while maintaining the 167 ratio between solid and liquid precipitation for the sake of simplicity (the heat flux as-168 sociated with snow melting in the ocean is relatively low). Precipitation near Antarc-169 tica has been shown to evolve following the Clausius-Clapeyron law (Ligtenberg et al., 170 2013; Donat-Magnin et al., 2021). The range of variation is therefore indexed to the tem-171 perature range considered for the heat flux: precipitation is multiplied by factors between 172 0.48 and 1.99, corresponding to coldest (-10°C) and warmest (+10°C) climatic conditions, 173 respectively. 174

The momentum flux is perturbed through meridional shifting of winds. To maintain flux independence, only the wind involved in the momentum flux calculation (i.e., ocean and sea ice surface friction) is modified, while we keep the wind seen by latent and sensible heat fluxes unchanged in the bulk formulae. The applied wind shift ranges between a 4.7° northward shift for coldest (-10°C) climatic conditions (Gray et al., 2021) and a 4.7° southward shift for warmest (+10°C) conditions (extrapolated from the 2100 CMIP5-RCP8.5 sensitivity described in Spence et al., 2014).

For the three types of perturbations, we conduct simulations with intermediate perturbations between the coldest and warmest climate perturbations in order to better characterize potential tipping points. Perturbations are local, only applied on continental shelf and slope. We do not perturb lateral boundary conditions, i.e., we maintain the presence of CDW in front of the continental shelf in all our simulations. It seems clear that cold conditions would prevail if CDW stopped to exist, and it is more interesting to identify how warm-to-cold abrupt transitions could occur in the presence of CDW.



**Figure 2.** Location of the perturbed area (a) for heat and freshwater fluxes and (b) for momentum flux. The perturbed area is highlighted in light blue, the transition area in blue and the unperturbed area in dark blue. The transition area avoids the artificial formation of strong density gradient or wind curl.

Basal melting highly depends on whether the perturbation is applied only on the 189 continental shelf or on the continental shelf and slope (not shown). We decided to in-190 clude the continental slope in the perturbed area as this area is relevant for CDW in-191 truding onto the shelf. The ice shelf melt rates are not sensitive to further northward 192 extension of the perturbation area, which indicates some robustness of our methodol-193 ogy. For the heat and freshwater flux perturbations, a transition area of  $3^{\circ}$  in latitude 194 (about 340 km) and 4° in longitude limits the temperature and precipitation gradient, 195 and thus the formation of strong density gradients, between the perturbed and unper-196 turbed areas (Fig. 2a). For the momentum flux perturbation, we additionally put a coastal 197 transition area of 150 km width between the perturbed wind and the katabatic winds 198 near the ice-sheet edges to avoid creating a substantial artificial wind curl perturbation 199 (Fig. 2b). 200

#### 201 2.3 Simulations

In order to assess the model response to atmospheric perturbations, we run a 61-202 vear simulation over the period 1958-2018. The simulation length is a compromise be-203 tween computational cost and the description of the natural decadal variability of the 204 ocean system, which can potentially impact the system stability and the occurrence of 205 tipping points. The reference experiment corresponds to the configuration retained af-206 ter calibration (see Supporting Information) with natural atmospheric and oceanic forc-207 ing over the modelled period. The model calibration improves the fidelity of the refer-208 ence simulation although the interannual variability is smaller than expected (consequences 209 will be discussed in section 4). The model spin-up is achieved after 10 years, thus, only 210 the period 1968-2018 is analyzed. The perturbed runs are identical to the reference run 211 except for the atmospheric forcing. We study three possible types of transition: cold-212 to-warm transitions as reported by Hellmer et al. (2012, 2017) for the Weddell Sea, warm-213 to-cold and warm-to-warmer transitions related to ancient or distant future climate tran-214 sitions. When an abrupt ocean transition occurs, reversibility is studied, i.e., for cold-215 to-warm (C2 in Fig. 3) and warm-to-cold transitions (C3 in Fig. 3). 216



Figure 3. Simulation set-up. The annotations C1, C2 and C3 correspond to the  $1^{st}$ ,  $2^{nd}$  and  $3^{rd}$  simulation cycle, respectively. The warm-to-warmer, warm-to-cold (C1) and cold-to-warm simulations enable us to identify the possible existence of transitions and the atmospheric conditions under which they occur. The cold-to-warm (C2) and warm-to-cold (C3) simulations are used to study their reversibility.

#### 217 **3 Results**

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### 3.1 Description of the transitions and their reversibility

For the sake of clarity, this section focuses on the mean melt rate of Pine Island and Thwaites on the one hand, and of Crosson and Dotson ice shelves on the other hand. Cosgrove and Getz ice shelves undergo similar melt transitions, albeit with much lower and higher mean melt values, respectively (not shown).

For Pine Island–Thwaites, the heat flux perturbations lead to a permanent collapse 223 of ice-shelf melting for air cooled by 2.5°C or more, with average melt rates below 0.3 224 m.w.e.yr<sup>-1</sup> (meters of water equivalent per year, i.e., 1 m.w.e.yr<sup>-1</sup> = 1000 kg.m<sup>2</sup>.yr<sup>-1</sup>), 225 comparable to those experienced by the Ronne or Eastern Ross ice shelves (Rignot et 226 al., 2013) (Fig. 4a). The -1°C perturbation leads to a collapse of melt rates after year 227 2000, while the -0.5°C perturbation keeps relatively high melt rates. Cycling our sim-228 ulations by repeating the period 1958-2018 indicates that the cooler state over 2000-2018 229 is related to the forcing data and not to a slow drift of our regional system as melt rates 230 are again high before 2000 in the repeated simulations (not shown). The heat flux per-231 turbations associated with higher air temperatures lead to a limited increase in basal melt 232 rates, with no more than a 34% increase for the  $+10^{\circ}$ C perturbation. The effect of in-233 creasing air temperatures seems to saturate, with little differences between  $+2^{\circ}$ C and 234  $+10^{\circ}$ C warming. Basal melt rates beneath Crosson–Dotson show a similar behavior as 235 Pine Island–Thwaites, with intermittent periods of very low melt rates for perturbations 236 as small as -0.5°C, permanent collapse of melt rates below -2.5°C (melt rate is slightly 237 higher than that of Pine Island–Thwaites with typical values of  $0.8-1.0 \text{ m.w.e.yr}^{-1}$ , and 238 a 28% increase in melt rates for  $+10^{\circ}$ C (Fig. 5a). It can also be noted that the ampli-239 tude of the seasonal melt cycle increases in response to warm perturbations for Pine Island-240 Thwaites but not for Crosson–Dotson. 241

The freshwater flux perturbations associated with lower precipitation lead to intermittent reductions in melt rates for precipitation reduced by 30% or more for Pine Island–Thwaites (Fig. 4b). Particularly low melt rates are found in the mid 1970s, early 2000s and late 2010s, but never reach the extremely low values resulting from the heat flux perturbations. This contrasts with Crosson–Dotson for which extended periods of very low melt rates (below 1.1 m.w.e.yr<sup>-1</sup>) are found when precipitation is reduced by 20% or more (Fig. 5b). Increased precipitation does not have a strong effect on melt rates, with only 17% and 9% increase in response to doubled precipitation for Pine Island–Thwaites and Crosson–Dotson, respectively.

Finally, the momentum flux perturbations associated with northward-shifted wind 251 at Pine Island–Thwaites results in intermittent decreases in melt rates, which is notice-252 able for a 2° northward wind shift (Fig. 4c). An extended collapse of basal melting is found 253 over the period 2000-2018 for a northward wind shift of  $4.7^{\circ}$ . The extended period of low 254 melt rates matches relatively well with those found for reduced precipitation. Crosson-255 Dotson is again more sensitive, with extended periods of very low melt rates for north-256 257 ward wind shift by 1° or more (Fig. 5c). The poleward-shifted winds lead to minor changes in basal melting: less than 5% and 15% increase for Pine Island–Thwaites and Crosson– 258 Dotson, respectively. 259



Figure 4. Evolution of monthly basal melting for Pine Island and Thwaites ice shelves over the period 1958-2018 for perturbations of (a) heat flux, (b) freshwater flux, and (c) momentum flux. The black curve (reference curve) corresponds to the simulation with the JRA55 reanalysis without modification. The red and blue curves correspond to simulations with atmospheric perturbations that aim to increase and decrease basal melting, respectively. The vertical green line indicates the end of the 10-year spin up.

We have just shown that abrupt transitions from a permanently high to a perma-260 nently low melt state can exist, and we now address the reversibility of these warm-to-261 cold transitions. We focus on transitions resulting from the strongest perturbations, i.e., 262 air cooled by  $10^{\circ}$ C, precipitation decreased by 52%, and winds shifted northward by  $4.7^{\circ}$ , 263 and we revert the atmospheric forcing to zero perturbation to re-run the period 1958-264 2018 starting from the 2018 perturbed state (Fig. 3). After 14 to 21 years, all perturbed 265 melt time series go back to the unperturbed state and remain within  $\pm 5\%$  of the orig-266 inal time series (Fig. 6). We conclude that all our warm-to-cold transitions in the Amund-267 sen Sea are reversible. This also means that our description of the warm-to-cold tran-268 sitions can be reverted to describe the cold-to-warm transitions. 269



Figure 5. Evolution of monthly basal melting for Crosson and Dotson ice shelves over the period 1958-2018 for perturbations of (a) heat flux, (b) freshwater flux, and (c) momentum flux.. The black curve (reference curve) corresponds to the simulation with the JRA55 reanalysis without modification. The red and blue curves correspond to simulations with atmospheric perturbations that aim to increase and decrease basal melting, respectively. The vertical green line indicates the end of the 10-year spin-up.

We also evaluate the reversibility of cold-to-warm transitions, bearing in mind that 270 such transitions may have occurred in the past. To do this, we take the final state of the 271 2nd cycle of 1958-2018 (unperturbed warm-climate (natural) forcing following a first cold-272 climate perturbed cycle), and we run a 3rd cycle of 1958-2018 again driven by the cold-273 climate perturbed forcing (Fig. 3). The two perturbed simulations (1st and 3rd cycle) 274 converge (within  $\pm 5\%$ ) after 5-6 years for the perturbed heat flux, after 13-20 years for 275 the perturbed freshwater flux and after 24-34 years for the perturbed momentum flux 276 (Fig. 7). We conclude that the cold-to-warm transitions are also reversible in the Amund-277 sen Sea. 278



Figure 6. Evolution of the reversed warm-to-cold transition for (a) Pine Island and Thwaites ice shelves and (b) Crosson and Dotson ice shelves. Only the cold-climate perturbations of maximum amplitude are drawn. The black curve corresponds to the simulation driven by the JRA55 reanalysis without modification (natural state), with shading indicating  $\pm 5\%$ .

#### 3.2 Physical processes

From a general perspective, the main external drivers of ocean variations are (i) 280 wind stress changes and (ii) surface heat and freshwater fluxes that modify the sea sur-281 face buoyancy (e.g., Marshall & Plumb, 2008; Talley et al., 2011). Winds induce a tan-282 gential stress at the ocean surface (directly or via sea-ice advection) and, thus, induce 283 surface water transport towards the side of the wind. This transport results in areas of 284 divergence and convergence that lead, respectively, to upwelling (Ekman suction) and 285 downwelling (Ekman pumping). Surface heat and freshwater fluxes modify the sea sur-286 face buoyancy, which can affect convection and the horizontal circulation via density gra-287 dients. At high latitudes, the net sea-ice production plays a key role in these processes. 288

Here, we analyze the Ekman vertical velocity  $(w_{Ek})$  and buoyancy flux at the ocean surface  $(B_s)$  to assess the impact of atmospheric forcing perturbations on the ocean properties. They are defined as:

$$w_{Ek} = \frac{1}{\rho_0} \vec{\nabla}_z \wedge \left(\frac{\vec{\tau}}{f}\right) \tag{1}$$

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$$B_s = \frac{g\alpha}{c_p}Q + g\beta S_s F \tag{2}$$

where  $w_{Ek}$  is the upward Ekman vertical velocity,  $\rho_0$  the reference seawater density,  $\vec{\tau}$ the wind/sea-ice stress at the ocean surface and f the Coriolis parameter.  $B_s$  is the buoyancy flux at the ocean surface,  $c_p$  the specific heat, g the gravitational acceleration,  $S_s$ 



Figure 7. Evolution of the reversed cold-to-warm transition for (a) Pine Island and Thwaites ice shelves and (b) Crosson and Dotson ice shelves. Only the cold -climate perturbations of maximum amplitude are drawn. The solid line represents the melt rate after applying the cold-climate perturbations for the first time over 1958-2018 (1st cycle). The dotted line represents the melt rate under the cold-climate perturbations (3rd cycle) following a 1958-2018 cycle of unperturbed conditions (2nd cycle). The shading corresponds to the melt value of the perturbed state of the 1st cycle  $\pm 5\%$ .

the sea surface salinity,  $\alpha$  the surface thermal expansion coefficient of seawater and  $\beta$  the corresponding coefficient for salinity, F and Q are the heat and freshwater fluxes received by the ocean surface (positive downward).

A striking feature of our ensemble of experiments is that all types of perturbation 299 approximately have the same ice-shelf melt evolution as a function of the surface buoy-300 ancy flux over the continental shelf (Fig. 8). The evolution curve consists of a highly sen-301 sitive regime bounded by a low plateau with no melt variations and a high plateau with 302 lower melt sensitivity. The similarity between the three curves in Fig. 8 suggests that 303 all perturbations mostly modify melt rates through changes of the surface buoyancy fluxes. 304 Hereafter, we describe the processes that affect the surface buoyancy for the various types 305 of perturbations. 306

The freshwater flux perturbations ("PREC" in Fig. 8) are the easiest to understand 307 as precipitation directly affects the surface buoyancy. Lowering precipitation reduces the 308 vertical density gradient and thereby favors convective mixing (Fig. 9c), which extracts 309 the heat of the deep spreading CDW. A much colder water below the thermocline (Fig. 9a) 310 explains the lower melt rates in the experiments with reduced precipitation. The oppo-311 site mechanism explains higher melt rates in the presence of enhanced precipitation. A 312 small part of the freshwater flux modification is also related to minor changes in sea-ice 313 production (Fig. 10a), due to the insulating properties of snow on sea ice (not shown). 314

The heat flux perturbations ("TEMP" in Fig. 8) have a less direct effect on sur-315 face buoyancy than just thermal expansion. Modified heat fluxes indeed explain less than 316 25% of the changes in surface buoyancy fluxes, while changes in freshwater fluxes related 317 to net sea-ice production (i.e., growth minus melt) have a preponderant effect on the sur-318 face buoyancy fluxes. In the presence of colder air, the net sea-ice production increases 319 considerably over the continental shelf (Fig. 10a), mostly due to a drastic decrease in sum-320 mer melting (not shown). The case is very similar to decreased precipitation, albeit with 321 a larger amplitude: increased convective mixing and related cooling below the thermo-322



Figure 8. Mean ice-shelf melt rate in the Amundsen Sea as a function of the mean surface buoyancy flux over the Amundsen Sea continental shelf over the period 1988-2018. The green, blue and purple curves correspond to perturbations of heat, freshwater and momentum fluxes, respectively. The black star represents the reference case. The red star represents a more realistic case. It corresponds to the current climate -  $0.5^{\circ}$ C by combining the perturbation of all the fluxes (TEMP -0.5°C, PREC x0.96 and WIND +0.24°N).

cline (Fig. 9d,e,f) leads to reduced ice-shelf melting (Fig. 8). Its minimum is reached when 323 the entire water column is close to the surface freezing temperature and the ice-shelf cav-324 ities are cold, i.e., melt rates are low and only controlled by the pressure dependency of 325 the freezing point. For the warm perturbations, the opposite effect exists until there is 326 too little net sea-ice production (Fig. 10a) to induce convective mixing. Beyond that, 327 the CDW layer remains mostly unchanged and ice-shelf melt rates keep increasing only 328 because warmer surface water gets in contact with the ice-shelf base (Fig. 9d,e,f). This 329 is consistent with the aforementioned increased seasonality of the Pine Island and Thwaites 330 melt rates (Fig. 4). 331

The results of the momentum flux perturbations are probably the most surprising 332 as they affect the surface buoyancy fluxes (see "WIND" in Fig. 8), although we have been 333 cautious not to modify the wind field in the calculation of the turbulent heat and evap-334 oration fluxes. The impact of winds on sea-ice drift actually explains the variation in buoy-335 ancy flux. In the experiments with a northward wind shift, the net production increases 336 (Fig. 10a) as winter sea-ice growth increases and summer melting decreases (not shown), 337 but the sea-ice volume decreases (frozen area and thickness decrease in Fig. 10b,c). This 338 is explained by enhanced advection of thinner sea ice towards the deep ocean (Fig. 11), 339 which leaves space for more air-sea exchange on the continental shelf, i.e., more sea-ice 340 production. Therefore, it is a similar perturbation of the vertical ocean stratification as 341 in the case of the freshwater and heat perturbations. In the case of a southward wind 342 shift, the annual sea-ice characteristics are little changed (Fig. 10) and so is the mean 343 ice-shelf melt rate (Fig. 8). 344

Although surface buoyancy flux on the continental shelf appears as the major driver of ice-shelf basal melt changes, the set of curves in Fig. 8 do not exactly overlap, especially the curve associated with the momentum flux perturbation. For a given buoyancy flux, the cold-climate momentum perturbation induces a slightly higher melt rate than



Figure 9. Shelf-averaged vertical profiles of conservative temperature (left), absolute salinity (middle) and potential density (right) over the period 1988-2018 for the various atmospheric perturbations : freshwater flux (top), heat flux (middle), and momentum flux (bottom) perturbations.



Figure 10. Sea-ice characteristics on the Amundsen Sea continental shelf related to surface buoyancy flux: (a) net production (i.e growth minus melt), (b) effective thickness (mean thickness over the whole continental shelf including ice-free areas), and (c) frozen area. The star represents the reference case. The perturbed flux configurations are colored in green for heat, blue for freshwater, and purple for momentum.



**Figure 11.** Sea-ice velocity anomaly relative to the reference case for a northward wind shift of 4.7°N. The background map is identical to the one shown in Fig. 1.

those related to freshwater and heat perturbations. The difference could be explained by Ekman dynamics if northward-shifted winds were associated with stronger Ekman upwelling, which would expose the ice shelves to a thicker layer of warm water (and, thus, partially inhibit the effects of decreased surface buoyancy fluxes). However, a stronger Ekman downwelling is found when considering the average velocity over the continental shelf (Fig. 13a).

A more detailed analysis at the ice-shelf scale shows that these differences are only 355 noticeable for the eastern ice shelves (Fig. 12), i.e., for Cosgrove, Pine Island and Thwaites, 356 357 suggesting regional differences in the acting mechanisms. We, therefore, analyzed the Ekman velocity at the entrance of the Pine Island–Thwaites Troughs as in Holland et al. 358 (2019), but the most extreme point  $(+4.7^{\circ}N)$  does not match either with the expected 359 upwelling anomaly (Fig. 13b). Ekman pumping in our simulations is spatially very noisy 360 (like Fig. 2a of Dotto et al., 2019), and we acknowledge a strong sensitivity to the ex-361 act location of the box used for the spatial average. Further investigation of Ekman ve-362 locities near individual ice-shelf fronts were similarly highly dependent on the location 363 of box boundaries and therefore not conclusive. In summary, Ekman dynamics might 364 explain the small difference in the melt response to the momentum perturbation and the 365 other two perturbations, but such effect remains elusive. Other possible explanations may 366 involve changes in ocean dynamics near the shelf break influencing the water mass prop-367 erties advected onto the continental shelf. 368



Figure 12. Mean basal melt rate of individual ice shelves in the Amundsen Sea as a function of the mean surface buoyancy flux over the Amundsen Sea continental shelf for the period 1988-2018. The green, blue and purple curves correspond to perturbations of heat, freshwater and momentum fluxes, respectively. The black star represents the reference case.



Figure 13. Mean basal melt rate as a function of upward Ekman velocity averaged over (a) the Amundsen Sea continental shelf and (b) over the entrance of Pine Island–Thwaites Troughs (box similar to the one defined in Fig. 1a of Holland et al. (2019)). The green, blue and purple curves correspond to perturbations of heat, freshwater and momentum fluxes, respectively. The black star represents the reference case.

## <sup>369</sup> 4 Discussion and Conclusion

#### 370

## Robustness of the thresholds with respect to our model biases

Our regional model captures well the seasonal variability (see Section S3 in Sup-371 porting Information). Despite a calibration, however, the reference simulation still has 372 a melt rate bias. The rate is outside the range of uncertainties of oceanic or satellite-based 373 estimates, although of a similar order of magnitude to those found in other regional model 374 studies (e.g., Nakayama et al., 2014; Kimura et al., 2017; Naughten et al., 2022), and has 375 an overly low interannual variability (Figs. 4-5) compared to observational estimates (Dutrieux 376 et al., 2014; Jenkins et al., 2018). Nevertheless, a more realistic interannual variability 377 is observed for relatively small atmospheric perturbations. It should be kept in mind that 378 a small perturbation of 0.5°C of the air temperature is of the order of magnitude of the 379 reanalysis biases estimated by Jones et al. (2016) for the Amundsen Sea region. Biases 380 are also large for precipitation, which is not constrained by data assimilation (Bromwich 381 et al., 2011; Palerme et al., 2017). This means that the 'real' Amundsen Sea might cor-382 respond to a slightly cooler and drier climate than our reference state. However, the melt 383 rate vs. buoyancy-flux curve is realistic and only the position of the reference state (black 384 star) on this curve could be biased. Thus, the exact thresholds for air temperature, pre-385 cipitation and wind shift for which a transition to the cold state occurs should still be 386 considered as uncertain. 387

#### 388

#### Robustness of the reversibility of abrupt transitions

Our results show that abrupt and reversible warm-to-cold as well as cold-to-warm 389 transitions could occur in the Amundsen Sea for relatively weak regional atmospheric 390 perturbations. The reversibility found in our experiments contrasts with the irreversibil-391 ity of similar cold-to-warm transitions found in simulations of the Weddell Sea and Filchner-392 Ronne Ice-Shelf cavity (Hellmer et al., 2017; Hazel & Stewart, 2020; Comeau et al., 2022). 393 The reason why the cold-to-warm transition is reversible in the Amundsen Sea but not 394 in the Weddell Sea remains unclear. The melt-induced circulation was presented as the 395 cause of the irreversibility in Hellmer et al. (2017), but the strong melt-induced circu-396

lation in the Amundsen Sea after a cold-to-warm transition (Jourdain et al., 2017; Donat-397 Magnin et al., 2017) does not seem able to maintain the onshore flow of Circumpolar Deep 398 Water when the forcing is reverted to cold climate conditions. For the coldest pertur-300 bations, the deep Amundsen Sea is approximately at a conservative temperature of -1.9°C 400 and an absolute salinity of 34.7 g kg<sup>-1</sup> (Fig. 9d,e), i.e., typical of the High Salinity Shelf 401 Water (HSSW) produced in the Weddell Sea. Hazel and Stewart (2020) explains the Wed-402 dell Sea tipping point by a feedback of ice-shelf meltwater to the salinity of newly formed 403 HSSW. The tipping conditions and associated hysteresis may, therefore, be sensitive to 404 the ratio between the HSSW formation rate and total ice-shelf basal mass loss, which 405 could explain different regimes in the Weddell and Amundsen Seas. 406

These transitions and their reversibility may be complicated or facilitated by effects not taken into account in our simulations, such as the feedbacks with the large-scale atmospheric and oceanic circulations or the ice-sheet dynamics.

First of all, we do not change the ocean lateral boundary conditions in our sensi-410 tivity experiments, while the Amundsen Sea is also sensitive to changes of water prop-411 erties advected from remote locations (Nakayama et al., 2018). It is known that large 412 atmospheric changes over multiple decades will have global effects, and we, therefore, ac-413 knowledge that our regional point of view is somewhat limited. Furthermore, strong mod-414 ifications of ice-shelf melting in the Amundsen Sea are expected to have significant con-415 sequences at circum-Antarctic (Nakayama et al., 2020) and global scales, with some pos-416 itive feedback in which more meltwater enhances the stratification and further exposes 417 ice shelves to CDW (Merino et al., 2018; Bronselaer et al., 2018; Golledge et al., 2019). 418 Such feedback is not considered in our study and it is difficult to estimate how they would 419 affect the thresholds and reversibility of our transitions. 420

Another limitation of our study is the missing evolution of ice-sheet dynamics in 421 response to changes in ice-shelf melting. In the presence of higher melt rates, ice shelves 422 are expected to thin and their grounding line to retreat. Ice-shelf thinning may slow down 423 melting, if the ice draft raises above the thermocline (De Rydt et al., 2014). Conversely, 424 strong grounding line retreat may enhance melting by exposing a larger basal area to 425 warm water and thus favoring a stronger melt-induced sub-ice shelf circulation (Donat-426 Magnin et al., 2017). For some geometrical configurations, the retreat of the calving front 427 may also favor melting by facilitating the circulation into ice-shelf cavities (Bradley et 428 al., 2022). If ice-shelf basal melt rates increase sufficiently, the ice dynamics is likely to 429 cross tipping points (Rosier et al., 2021), which would irreversibly put the Amundsen 430 Sea in a different state due to the aforementioned feedback. Nonetheless, it is difficult 431 to quantify the exact thresholds for which irreversibility would be found without using 432 a fully coupled ocean-ice-sheet model. 433

434

## Buoyancy vs wind-stress forcing

There is a consensus that the intrusion of warm CDW on the continental shelf plays a major role in the variability of ice-shelf basal melting, but several different processes have been suggested to explain their transport. As several studies independently investigate the eastern (Pine-Island–Thwaites) and western (Dotson–Getz) parts of the shelf (Wåhlin et al., 2012; Nakayama et al., 2013; Dotto et al., 2019) and as our study identifies distinct regimes between these two parts, it seems suitable to separate the analysis of processes according to these two regions.

Our study shows that the surface buoyancy flux on the shelf is the main driver of
the multi-decadal changes in basal melting for the western Amundsen Sea (Dotson-Getz)
regardless of perturbation. They also indicate no changes in local Ekman pumping in
response to idealized wind perturbations, in contrast with the observational study by Kim
et al. (2021), in which Ekman pumping along the Dotson-Getz trough explains 43% of
the summer thermocline interannual variability. Dotto et al. (2020) have suggested that

local winds at the shelf break may affect the eastward undercurrent and thereby the heattransport onto the continental shelf.

In the eastern Amundsen Sea (Cosgrove, Pine Island, Thwaites), our results again 450 indicate that changes in the surface buoyancy fluxes are the main drivers of ice-shelf melt 451 rate variations at multi-decadal time scales. A small deviation of the wind-perturbation 452 experiments (Fig. 8 - purple line) nonetheless suggests that other wind-related processes 453 might play a role, although the exact mechanism remains elusive. Previous studies have 454 largely attributed interannual variability of the eastern Amundsen Sea to Ekman pump-455 ing at the shelf break (e.g., Holland et al., 2019; Dotto et al., 2019; Webber et al., 2019; 456 Naughten et al., 2022), although sea-ice formation can also play a role in some specific 457 years (St-Laurent et al., 2015; Webber et al., 2017). 458

In summary, changes in the surface buoyancy forcing appear to be the dominant 459 driver of the variations in ice-shelf melting in all our experiments, whereas most previous studies have emphasized the direct role of wind stress, in particular through Ekman 461 pumping. Part of the apparent discrepancy may be related to the multi-decadal time scale 462 of our perturbations, which can slowly induce a change in the baroclinic balances that 463 could overwhelm the relatively fast Ekman dynamics. We acknowledge, however, that 464 our wind perturbations are highly idealized and may not capture the full complexity of 465 wind changes at the continental shelf break, although we do have increasing Ekman ve-466 locities at the shelf break for the transition from cold to warm climate. 467

### Implications for past and future climates

Our results indicate cold Amundsen Sea cavities (close to surface freezing point) 469 for conditions of the Last Glacial Maximum (Fig. 8). This is consistent with grounding 470 lines of paleo-ice streams near the continental shelf break during the last glacial period 471 (Larter et al., 2014). Combined heat, freshwater and momentum perturbations main-472 tained cold-cavities for climate conditions typical of -0.5 °C compared to present day even 473 in the presence of CDW at the continental shelf break (red star in Fig. 8). This suggests that pre-industrial conditions (approximately 1°C colder than present day (IPCC, 2021)) 475 were associated with cold cavities in the Amundsen Sea. The transition to warm cav-476 ities may have occurred or be occurring as multi-year oscillations between cold and warm 477 periods (Figs. 4-5). For conditions warmer than today, the decadal variability is relatively 478 weak and cavities remain permanently warm. Our idealized experiments suggest a grad-479 ual but limited increase in ice-shelf basal melting in response to global warming beyond 480 present levels. 481

#### 482 Data and softwares

468

The model version and set of parameters used to run our experiments are provided
 in https://github.com/Astrolabe-JC/Simulations\_NEMO. THE GITHUB REPOS ITORIES WILL BE ARCHIVED ON http://zenodo.org AFTER ACCEPTANCE.

#### 486 Acknowledgments

The authors would like to acknowledge Clara Burgard for her valuable advice on figures and presentations as well as Hélène Seroussi for her constructive feedback as a member of the PhD committee.

This study was funded by the European Union's Horizon 2020 research and innovation programme under grant agreements No 820575 (TiPACCs) and by the French National Research Agency under grant No ANR-19-CE01-0015 (EIS). This work was granted access to the HPC resources of CINES under the allocation A0100106035 attributed by GENCI.

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## Supporting Information for Drivers and reversibility of abrupt ocean state transitions in the Amundsen Sea, Antarctica

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Introduction This supplementary document describes the calibration and evaluation of the reference simulation. It includes a description of both the observational data sets used to validate the configuration (see Section S1) and sensitivity experiments carried out in this study (see Section S2 and Table S1). The evaluation of these sensitivity experiments are made through the analysis of sea-ice extent (see Section S3 and Figure S1), temperature profiles (see Section S4 and Figure 2) and basal melt rates (see Section S5 and Figure S3).

## S1. Observational Data Sets

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The model is evaluated over the period 2012-2018 through a comparison with available observational datasets relative to sea-ice concentration, temperature profiles, and iceshelves melt rates.

We used 1607 Conductivity-Temperature-Depth (CTD) profiles collected during summer campaigns from 1994 to 2018 (Dutrieux et al., 2014; Heywood et al., 2016) to evaluate the simulated temperature field. Individual CTD profiles are interpolated vertically onto the 75 vertical levels of our model configuration. The comparison to CTD data is performed by sampling model outputs in space (nearest profile) and time (linear interpolation between monthly outputs) following the actual CTD station distribution.

To evaluate the simulated sea-ice cover, the model outputs are compared to the NOAA/NSIDC Climate Data Record (CDR) of the Passive Microwave Sea-Ice Concentration dataset, version 4 (Meier et al., 2022), which provides daily and monthly estimates of sea-ice concentration in the polar regions for the period 1987-2020 on a  $25 \times 25$  km stereographic grid. The observational data are interpolated onto the model grid and only concentration values above 0.15 are considered in both the model and observations due to observational uncertainties for lowest concentrations.

Finally, the simulated ice-shelf basal melt rates over the period 2012-2018 are compared to estimates based on satellite observations (Adusumilli et al., 2020).

## S2. Sensitivity experiments

Preliminary to our main study, several sensitivity experiments were carried out in order to improve the fidelity of the model, in particular its ability to correctly represent iceshelf basal melting. Their main characteristics are detailed in Tab. S1. In view of the shortcomings of the initial simulation (INITIAL), we explored three distinct ways

to improve the representativeness of the model: (i) vertical mixing parameterization, (ii) forcing, and (iii) parameterization of melting under ice shelves.

The diffusivity and vertical turbulent viscosity coefficients are derived from a turbulent closure model that does not allow the thermocline depth to be represented correctly, particularly in the Southern Ocean (Rodgers et al., 2014). The thermocline depth is often too high in summer and when it is windy, as observed in INITIAL. To compensate for the lack of representation of some processes, an ad hoc parameterization exists in NEMO (see "TKE scheme" in Madec & the NEMO Team, 2016) to inject an amount D of additional turbulent kinetic energy below the mixed layer, with D defined as:

$$D = (1 - f_i) f_r e_s e^{-\frac{z}{h_\tau}} \tag{1}$$

where  $f_i$  is the sea-ice concentration,  $f_r$  the fraction of surface turbulent kinetic energy penetrating below the mixed layer,  $e_s$  the surface boundary condition of the turbulent kinetic energy (diagnosed by the TKE scheme), z the depth, and  $h_{\tau}$  the vertical mixing length scale. To lower the depth of the thermocline, we modify the  $f_r$  parameter in EFR by taking the maximum value suggested by Madec (2008); Heuzé et al. (2015). In HTAU, we modify the vertical mixing length to  $h_{\tau} = 60$  m, which is appropriate for high latitudes (Rodgers et al., 2014).

One of the main sources of divergence between our simulations and the observations is due to uncertainties in the forcing. We tested the model sensitivity to oceanic and atmospheric forcing. The BDY experiment includes the correction of oceanic lateral boundaries to match the World Ocean Atlas 2018 seasonal climatology (Garcia et al., 2019), and BDY + HTAU + FORCING takes into account another atmospheric reanalysis product, DFS5.2

(Dussin et al., 2016). The BDY + HTAU is a combination of the changes made both in BDY and HTAU.

Finally, in order to improve the basal melt rates, especially under Getz, Thwaites and Pine Island ice shelves, several corrections are explored. BDY + HTAU + GAMMA corrects basal melt rates through a modification of the heat and salt exchange coefficients ( $\Gamma_T$  and  $\Gamma_S$ , respectively). According to Jourdain et al. (2017), melt rates are proportional to  $\Gamma^{1.25}$  where  $\Gamma$  is the heat/salinity exchange coefficient. Since melting under Getz was too high by about 1/3, the coefficients were multiplied by  $\frac{2}{3}^{\frac{1}{1.25}}$  to reduce the melt rate by one third. The partial collapse of Thwaites Ice Shelf in 2015 altered the topography significantly. To assess the influence of topography on the melt rates under Thwaites Ice Shelf, the Rtopo2 topography (Schaffer et al., 2016) prior to the collapse of the ice shelf is used in BDY + HTAU + GAMMA + TOPO.

## S3. Evaluation of sensitivity experiments: Sea-ice extent

The sea-ice cover analysis gives an insight into the seasonal variability (CTD profiles are collected only in austral summer and basal melt rates are often estimated over several years). In all our simulations, the seasonal variability is well represented (see Fig. S1) although the simulated summer and winter extrema differ from those in the NOAA/NSIDC climatology. Sea-ice extent is very sensitive to lateral boundary conditions. The correction of lateral boundaries to the World Ocean Atlas 2018 (BDY) results in an overestimation of the sea-ice extent (+35% in autral summer and +17% in winter), which hides an overestimation in the deep ocean and an underestimation in front of the ice shelves. The sea-ice extent is also sensitive to the stratification of the water column (HTAU). The change in vertical mixing leads to an underestimated sea-ice extent (-2% in austral summer -16%

in winter). The choice of atmospheric forcing impacts mostly the summer sea-ice cover, while changes in the vertical mixing parameter  $f_r$ , topography, and heat and salt exchange coefficients do not affect the sea-ice extent.

## S4. Evaluation of sensitivity experiments: Temperature profiles

The CTD profiles measured in front of the Pine Island–Thwaites and Dotson–Getz ice shelves suggest a different flow pattern between the western and eastern continental shelf, consistent with Wåhlin et al. (2012); Dotto et al. (2019). The waters are generally 0.5°C warmer in the Pine Island–Thwaites area where CDW onshore flow is facilitated.

In front of the Pine Island–Thwaites ice shelves, INITIAL overestimates the subsurface temperature (0-400 m) by 1°C and the bottom temperature (800-1200 m) by  $0.5^{\circ}$ C, and the thermocline is too shallow by 200 m (Fig. S2). The ocean boundary correction with the World Ocean Atlas 2018 results in a significant change in the water column of the deep ocean. As the near-bottom temperature is essentially determined by the oceanic conditions at the shelf break, this correction lowers the temperature by about  $0.5^{\circ}$ C and thus approaches the observations in the lower part of the water column (deviation from observations less than  $0.2^{\circ}$ C below 600 m), although the thermocline is still too shallow. The increased vertical mixing length scale lowers the overestimated thermocline depth. Above 500 m, the temperature is lowered by an average of  $0.5^{\circ}$ C, but it is still higher than the CTD profiles ( $0.5^{\circ}$ C on average). The other sensitivity tests have little impact on the temperature of the water column ( $0.2^{\circ}$ C down to 300 m).

Concerning Dotson–Getz ice shelves, INITIAL overestimates the bottom temperature (800-1200 m) by 0.5°C, and the thermocline is again too shallow by 200 m (Fig. S2). The ocean boundary correction with the World Ocean Atlas 2018 lowers the temperature by

about 0.5°C, resulting in a deviation from observations of less than 0.2°C. The increased vertical mixing length scale lowers the overestimated thermocline depth by 100 m, but has also an impact on near bottom temperature as does the change in the values of the heat and salt exchange coefficients.

## S5. Evaluation of sensitivity experiments: Basal melt rates

We evaluate our basal melt rates by comparing them to satellite data estimates from Adusumilli et al. (2020) despite large uncertainties. Simulated cavity melt rates are close to published estimates for Venable, Abbot, and Cosgrove, underestimated for Pine Island and Thwaites (-30% and -77%, respectively for BDY + HTAU), although the vertical temperature distribution is overestimated by 0.5°C in front of these ice shelves, and overestimated for Crosson, Dotson and Getz (337%, 98% and 119%, respectively for BDY + HTAU) (see Fig. S3). In general, the interannual variation in melting is not as large as observed. Due to proportionality, a unique modification of the melt rates (BDY + HTAU + GAMMA) does not improve the overall results. Thus, we prefer to keep the original heat and salt exchange coefficients ( $\Gamma_T$  and  $\Gamma_S$ ) in order to simulate correct melt rates for Pine Island rather than Getz ice shelf, which was already problematic when setting up Nemo's cavity module (Mathiot et al., 2017). The gain from the change in ice topography at Thwaites (BDY + HTAU + GAMMA + TOPO) is negligible.

## S6. Conclusion on the sensitivity experiments

We selected the simulation (BDY + HTAU) as the reference simulation from which simulations with perturbations were built. The reference simulation captures well the seasonal variability with an average sea-ice extent error in austral summer and winter, of, respectively, 31% and 6%. The temperature between 700 and 1200 m depth is well represented on the shelf (deviation from observations less than 0.2°C). The thermocline is still too shallow, which corresponds to an overestimation of the temperature between 300 m and 600 m depth of about 0.5°C. Simulated basal melt rates are close to published estimates for Venable, Abbot and Cosgrove, underestimated for Pine Island and Thwaites (-30% and -77%, respectively), and overestimated for Crosson, Dotson and Getz (337%, 98% and 119%, respectively).

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 Table S1.
 Characteristics of sensitivity experiments used for model calibration (different terms are defined in section 1)

Simulation Name	$f_r$	$h_{ au}$	Boundaries	$\Gamma_T$	$\Gamma_S$	Торо	Forcing
INITIAL	0.05	$h_{ au}\left(\phi ight)$	-	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	JRA55
EFR	0.10	$h_{ au}\left(\phi ight)$	-	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	JRA55
BDY	0.05	$h_{ au}\left(\phi ight)$	WOA18 Correction	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	JRA55
HTAU	0.05	60m	-	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	JRA55
BDY + HTAU	0.05	60m	WOA18 Correction	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	JRA55
BDY + HTAU + GAMMA	0.05	60m	WOA18 Correction	$1.6010^{-2}$	$4.4810^{-4}$	BedMachine	JRA55
BDY + HTAU + GAMMA + TOPO	0.05	60m	WOA18 Correction	$1.6010^{-2}$	$4.4810^{-4}$	RTopo2	JRA55
BDY + HTAU + FORCING	0.05	60m	WOA18 Correction	$2.2110^{-2}$	$6.1910^{-4}$	BedMachine	DFS5.2



Figure S1. Variability of the sea-ice extent over the period 2012-2018 depending on the sensitivity experiments. As a reminder, only areas of sea-ice concentration greater than 0.15 were considered for the sea-ice extent estimation. The observations correspond to the NOAA/NSIDC Climate Data Record of Passive Microwave Sea-Ice Concentration dataset.

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**Figure S2.** Mean conservative temperature profiles in front of Pine Island and Thwaites (left), Dotson(middle), and Getz (right) ice shelves depending on the sensitivity experiments. The observations are CTD profiles collected during austral summer campaigns from 2012-2018.



Figure S3. 2012-2018 mean basal mass loss under various Amundsen Sea ice shelves depending on the sensitivity experiments. Simulated basal mass losses are compared to the 2010-2018 estimates from Adusumilli et al. (2020)