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Abstract

Auroral particle precipitation potentially plays a main role in ionospheric plasma structuring. The impact of auroral particle precipitation on plasma structuring is investigated using multi-point measurements from scintillation receivers and all sky cameras from Longyearbyen, Ny-Ålesund and Hornsund on Svalbard. This provides us with the unique possibility of studying the spatial and temporal dynamics of the aurora. Here we consider three case studies to investigate how plasma structuring is related to different auroral forms. We demonstrate that plasma structuring impacting the GNSS signals is largest at the edges of auroral forms. Here we studied two stable arcs, two dynamic auroral bands and a spiral. Specifically for arcs we find elevated phase scintillation indices at the pole-ward edge of the aurora. This is observed for auroral oxygen emissions (557.7 nm) at 150⁻km in the ionospheric E-region. This altitude is also used as the ionospheric piercing point for the GNSS signals as the observations remain the same regardless of different satellite elevations and azimuths. Further, there may be a time delay between the temporal evolution of aurora (f.e. commencement and fading of auroral activity) and observations of elevated phase scintillation indices. The time delay could be explained by the intense influx of particles, which increases the plasma density and causes recombination to carry on longer, which may lead to a persistence of structures - a 'memory effect'. High values of phase scintillation indices can be observed even shortly after strong visible aurora and can then remain significant at low intensities of the aurora.

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Ionospheric Plasma Structuring in Relation to Auroral Particle Precipitation

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ABSTRACT

Auroral particle precipitation potentially plays a main role in ionospheric plasma structuring. The 8 impact of auroral particle precipitation on plasma structuring is investigated using multi-point 9 measurements from scintillation receivers and all sky cameras from Longyearbyen, Ny-Ålesund 10 and Hornsund on Svalbard. This provides us with the unique possibility of studying the spatial and 11 temporal dynamics of the aurora. Here we consider three case studies to investigate how plasma 12 structuring is related to different auroral forms. We demonstrate that plasma structuring impacting 13 the GNSS signals is largest at the edges of auroral forms. Here we studied two stable arcs, two 14 dynamic auroral bands and a spiral. Specifically for arcs we find elevated phase scintillation indices 15 at the pole-ward edge of the aurora. This is observed for auroral oxygen emissions (557.7 nm) at 16 150 km in the ionospheric E-region. This altitude is also used as the ionospheric piercing point for 17 the GNSS signals as the observations remain the same regardless of different satellite elevations 18 and azimuths. Further, there may be a time delay between the temporal evolution of aurora (f.e. 19 commencement and fading of auroral activity) and observations of elevated phase scintillation 20 indices. The time delay could be explained by the intense influx of particles, which increases the 21 plasma density and causes recombination to carry on longer, which may lead to a persistence of 22 structures - a 'memory effect'. High values of phase scintillation indices can be observed even 23 shortly after strong visible aurora and can then remain significant at low intensities of the aurora. 24

Key words. particle precipitation, phase scintillation index, plasma structuring, ionospheric E-region

1. Introduction

The aurora can be seen as the signature of direct coupling of the ionosphere and magnetosphere. During high geomagnetic activity, energetic particle precipitation leads to higher intensity of the aurora resulting in different auroral forms. Dynamical processes in the E- and F- regions of the ionosphere are often associated with instabilities and turbulence which result in plasma structuring and irregularities at various scales. Such irregularities in ionospheric plasma density have impact on the propagation of radio waves (e.g., Keskinen and Ossakow, 1983; Huba et al.,

1985; Kintner and Seyler, 1985; Moen et al., 2013; Deshpande et al., 2014). Trans-ionospheric 32 radio waves propagating through regions of density irregularities undergo diffraction and refraction, 33 and they result in rapid fluctuations in phase and amplitude of the received signal, referred to as 34 scintillation (e.g., Hey et al., 1946; Kintner et al., 2007). Scintillation of the received signal affects 35 man-made systems, such as the radio communication and/or satellite based positioning systems. 36 At the same time, scintillation of the received signal can be used as an indication for ionospheric 37 plasma structuring. This will be also the approach in this work, where we will focus on the role 38 of the auroral particle precipitation during geomagnetic substorms and investigate how different 39 discrete auroral forms, i.e., stable arcs and fast moving forms (such as spirals), relate spatially and 40 temporally to structuring in the ionospheric plasma density. 41

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The Earth's ionosphere and magnetosphere are directly coupled in the polar regions via the 43 Birkeland currents, which can be seen as drivers for the aurora. There is a variety of resulting 44 auroral forms, which have been categorized over many years with regards to their shapes, process 45 and lifetime, but to this day there is no clear and well accepted definition for all of the forms. Some 46 auroral forms and their evolution can be linked to certain substorm phases (e.g., Akasofu, 1966; 47 Elphinstone et al., 1996; Partamies et al., 2015). The auroral arcs are well-studied phenomena in 48 quiet (for a review see Karlsson et al. (2020)) and active geomagnetic periods. Davis (1978) studied 49 auroral arcs and their distortions into complex forms (spirals and curls) and defined an auroral 50 arc as an recognizable luminosity resulting from the impingement of a field-aligned sheet beam 51 of charged particles upon the atmosphere. The most simple form is an east-west elongated quiet 52 discrete auroral arc. The width of the arc can be very thin (0.5-1.5 km) (Partamies et al., 2010), but 53 most mesoscale size arcs have an observed width of around 10 km - 50 km (Knudsen et al., 2001). 54 Arcs often appear in multiple- arc structures, and are found both during quiet and active periods. 55 Multiple arc structures are arrays of arcs which are near-parallel to each other and in close 56 proximity, and they are also the most common form of discrete auroral observations (Davis, 1978; 57 Gillies et al., 2014). The multiple arc structures are referred to as arc packets when they are formed 58 by splitting of the trailing arc into two, possibly due to the Alvén wave dispersion (Semeter et al., 59 2009). Stable and quiet auroral arcs are expected in the evening and at night-time. Here, we will 60 study auroral emissions in the night time sector from 21h to 03h MLT (magnetic local time). 61 Other noteable auroral distortions are curl, spiral, bands, ray forms, westward travelling surge 62 or omega bands (e.g., Elphinstone et al., 1996; Hallinan and Davis, 1970; Ivchenko et al., 2005; 63 Partamies et al., 2017). In this paper we will focus on spiral structures. Spirals are larger scale 64 auroral vertices (20-1300km), which form as the aurora twists counter-clock wise. It is suggested 65 that the Kelvin-Helmholtz instability could play a role in formation of the spiral (Davis and 66

⁶⁷ Hallinan, 1976; Hallinan and Davis, 1970; Hallinan, 1976; Partamies et al., 2001). While Davis and

⁶⁸ Hallinan (1976); Keiling et al. (2009) relate spirals to magnetically disturbed periods, Partamies

et al. (2001) relate them primarily to magnetically quiet conditions. However, fast-moving spirals,

⁷⁰ which signatures are only found in a single all-sky camera (ASC) image, seem to be related to an

⁷¹ increasing geomagnetic activity (Partamies et al., 2001).

⁷² A variety of highly dynamic phenomena (e.g., polar cap patches, field aligned currents, high

⁷³ density F-region plasma, or particle precipitation) cause plasma irregularities in the E and F- region

⁷⁴ ionosphere (e.g., Moen et al., 2013; van der Meeren et al., 2015; Spogli et al., 2016; Jin et al., 2016,

⁷⁵ 2017; Fæhn Follestad et al., 2020). Ionospheric plasma structuring can be indirectly observed by

⁷⁶ scintillation receivers. Recorded signals allow for calculating the phase and amplitude scintillation

⁷⁷ indices. At high latitudes, effect on irregularities on the trans-ionospheric wave can also be reflected

⁷⁸ in the degradation of the receiver tracking performance, Total Electron Content (TEC) jumps and

⁷⁹ cycle slips (Skone et al., 2001; Alfonsi et al., 2008; Moen et al., 2013; Chernyshov et al., 2020, &

⁸⁰ references therein).

⁸¹ Auroral particle precipitation can lead to significant plasma irregularities (Kelley et al., 1982;

⁸² Keskinen and Ossakow, 1983; Weber et al., 1985; Prikryl et al., 2011), but it is still an open

question to what extent it contributes to plasma structuring and whether it is dominant in the E or

⁸⁴ F-region ionosphere. It is observed that soft particle precipitation is unlikely the main source of

the large-scale F-region plasma structures on the nightside (Jin et al., 2016). While on the dayside

the main ionization source of the E-region is the EUV solar radiation, on the nightside Joule

⁸⁷ heating and auroral particles with energies of 10-30 keV dominate the energy input into the system

⁸⁸ (Millward et al., 1999; Wilson et al., 2006; Nikolaeva et al., 2021).

⁸⁹ There have been previous studies that tried to answer this question. Kinrade et al. (2013) studied

⁹⁰ ionospheric irregularities caused by the auroral particle precipitation using scintillation receivers

and auroral imagers located at the South Pole. They found spatially and temporally well-correlated scintillations with atomic oxygen emissions at wavelengths of λ =557.7 nm and 630.0 nm. This

⁹³ correlation is better for emissions at $\lambda = 557.7$ nm, which could be due to the tracking method's

⁹⁴ effectiveness at two emission wavelengths of different characteristic intensities. They found optical

⁹⁵ images of aurorae to be a useful spatial and temporal indicator of the GPS phase scintillations

⁹⁶ during intense and persistent auroral activity, especially for discrete arcs and at the auroral

⁹⁷ boundary. Another study, by van der Meeren et al. (2015), sheds light on the Global Navigation

Satellite System's (GNSS's) signal scintillations during intense substorm aurora. The data is recorded at GNSS receivers around Svalbard. They observed severe phase scintillation, which are following the intense pole-ward edge of the auroral oval as it expands pole-ward and show that

received signals may experience strong scintillation when they intersect oxygen (λ =557.7 nm)

emissions. The satellite systems of GPS, GLONASS and Galileo were affected similarly through the scintillations in relation to the intense line of sight auroral emissions in a highly localized

region of the sky. Discrete aurora and GPS signal corruption have been studied by Semeter et al. (2017), using a network of receivers and imagers in Alaska. The ASC sampled images of oxygen

(2017), using a network of receivers and imagers in Alaska. The ASC sampled images of oxygen emissions at $\lambda = 558$ nm. The auroral form of interest was a westward travelling surge and the loss of lock events consistently appeared at the edges of auroral forms. The scintillation regions were confined to a narrow strip (<20 km) tangential to the trailing edge of the observed aurora. As the appearance of the loss of lock events consistently followed the auroral boundaries irrespective of

the satellite elevation and azimuth, the authors suggested the source to be near the oxygen 558 nm emission line. This is within the E-region of the ionosphere.

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The questions i) in which layer of the ionosphere auroral particle precipitation plays a dominant role in plasma structuring and ii) what the spatial and temporal characteristics of the elevated scintillation indices with respect to different auroral forms are still remain open (e.g., Kinrade et al., 2013; van der Meeren et al., 2015; Jin et al., 2016; Semeter et al., 2017). In this paper, we will study how stable auroral arcs and fast moving forms (bands and spirals) relate spatially and temporally to ionospheric plasma structuring. The observations are made by using several scintillation receivers and ASCs located in Longyearbyen, Ny-Ålesund, and Hornsund on Svalbard. Case studies which

¹²⁰ consider different auroral forms during the substorm events are presented. We observe elevated val-

¹²¹ ues of the phase scintillation index pole-ward of the arcs, and on the boundaries of fast moving

¹²² forms. These characteristics have been found for the 557.7 nm oxygen emissions, indicating that

the auroral particle precipitation in the E-region can contribute to relevant plasma structuring. We

also observe a short time-delay between the onset of the aurora and observed plasma structuring.

2. Instruments and Approach

To study whether plasma structuring is driven by particle precipitation, we investigate the relative 126 location between elevated scintillation indices and the aurorae. For this purpose we use data from 127 three ASC on Svalbard. Two of the cameras are Keo Sentry 4ix Monochromatic Imagers from KEO 128 Scientific, with filters, operated by the University of Oslo (UiO). They are situated in Longyearbyen 129 (LYR, geographic coordinates: 78.15° N, 16.04° E) and in Ny-Ålesund (NYA, geographic coordi-130 nates 78.92° N, 11.93° E). The imagers record emission intensities every 30 s, with a field of view 131 (FOV) of 180°. Both imagers are equipped with the narrow band filters to monitor 557.7 nm (green) 132 and 630.0 nm (red) auroral emissions. For the analysis, the ASC images are projected to 150 km 133 and 250 km (green and red auroral emission altitudes) respectively. The imager in Hornsund (HOR, 134 geographic coordinates 77° N, 15.55° E) is a Sony A7 SII color camera with a FOV of 180°, also 135 operated by UiO. The availability of recorded images for this camera on the days of interest is every 136 15 minutes. The ASCs are used to determine the form of the aurora and its relative location to the 137 observed elevated scintillation indices of Global Navigation Satellite Systems (GNSS) signals. 138

To analyse the impact of auroral particle precipitation on the GNSS signals, five GNSS ionospheric scintillation receivers on Svalbard are used in this study. The delay Δt of a signal propagating through a plasma is dependent on the electron density n_e and the signal frequency f:

$$_{142} \quad \Delta t = \frac{40.3}{cf^2} \int n_e \, d\rho, \tag{1}$$

with c being the speed of light, ρ the ray path (e.g., Kintner et al., 2005, 2007).

The phase ϕ is connected to time delay Δt and therefore to electron density variations along the signal path (equation 1) through

$$_{146} \quad \phi = f\Delta t \tag{2}$$

147 (e.g., Yeh and Liu, 1982; Kintner et al., 2007). As the phase is affected by the time delay and

electron content variations it indicates plasma structuring processes. Phase scintillations are caused

¹⁴⁹ by large scale irregularities in a range of meters (above the Fresnel radius) to a few kilometers (e.g.,

Basu et al., 1998; Kintner et al., 2007). The phase scintillation index σ_{ϕ} ,

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$$\sigma_{\phi} = \langle \phi^2 \rangle - \langle \phi \rangle^2, \tag{3}$$

can be understood as the standard deviation of a measured phase and is dominated by large-scale fluctuations (Yeh and Liu, 1982; Kintner et al., 2007). Even though the index is not issue-free, e.g. due to the dominance of the low-frequency component of the phase power spectrum, the phase or related electron density variations nevertheless can indicate physical structuring in the ionosphere (Beach, 2006). Refractive and diffractive variations of σ_{ϕ} are not differentiated in this study, as we

use σ_{ϕ} as a measure of plasma structuring and to study the relative location to auroral forms, but are

not studying the effects on the carrier phase. Amplitude scintillations are described by the amplitude scintillation index S_4

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$$S_4 = \sqrt{\frac{\langle I^2 \rangle - \langle I \rangle^2}{\langle I \rangle^2}}.$$
 (4)

(e.g., Briggs and Parkin, 1963; Yeh and Liu, 1982). The S_4 index describes and is effected by irregularities in a range of hundreds of meters to meters (at and below the Fresnel radius) (e.g., Basu et al., 1998; Kintner et al., 2007).

In this study, we use the calibrated 60-seconds reduced data (Oksavik, 2020a) of the phase (σ_{ϕ}) and amplitude (S_4) scintillation index. A cut-off angle of 15° is used on the scintillation data to minimize the multipath effects.

The data is recorded by NovAtel GPStation-6 receivers. The receiver in Hornsund is situated about 500 meters from the ASC and is run by the University of Oslo. The receivers situated in Ny-Ålesund (co-located with the NYA ASC), Longyearbyen (co-located with the LYR ASC), Bjørnøya (BJN, geographic coordinates 74.50° N, 19° E) and Hopen (HOP, geographic coordinates 76.51° N, 25.01° E) are operated by the University of Bergen (Oksavik, 2020b). The receivers track GPS,

¹⁷² Galileo and GLONASS satellites. We use all received signals from these satellite systems as they

have previously been shown to be similarly affected by ionospheric irregularities (van der Meeren
et al., 2015).

The Interplanetary Magnetic Field (IMF) data and the solar wind conditions are downloaded 175 from NASA/GSFC's OMNI data set through OMNIWeb (King and Papitashvili, 2005). We assess 176 the IMF Bz magnetic field component, the solar wind flow speed and sym-H (1-min resolution 177 GSM (geocentric solar magnetospheric coordinates) data). A negative IMF Bz component can in-178 dicate dayside reconnection. The sym-H index is a measure of the magnetospheric ring current 179 intensity and used to quantify geomagnetic storms. The index is calculated using data from differ-180 ent magnetometer stations near the equator and describes the symmetric portion of the magnetic 181 field horizontal component (Wanliss and Showalter, 2006). These indices help to describe and filter 182 for background conditions and evaluate whether the geomagnetic conditions indicate geomagnetic 183 storms. To investigate substorm conditions the horizontal component of the local magnetic field 184 was used. It is recorded by a magnetometer network around Svalbard operated by the Tromsø 185 Geophysical Observatory (Tanskanen, 2009). The decrease in the B_x component of the magnetic 186 field at high latitudes is a signature of the enhancement of the westward electrojet and the substorm 187 current wedge in superposition with eastward electrojet enhancements (Akasofu, 1965; D'Onofrio 188 et al., 2014). Further we use the Kp index (from GFZ Potsdam and the National Geophysical Data 189 Center), which is a proxy for the energy input from the solar wind to the magnetosphere. The Kp 190 index is calculated based on 13 selected subauroral ground-based magnetic observatories and is the 191 mean value of the disturbance levels in the horizontal magnetic field components. Higher Kp values 192 correspond to stronger disturbances (Matzka et al., 2021). 193

3. Case Selection and Conditions

¹⁹⁵ Data from the ASC network were used to select events for the case study. The first season ¹⁹⁶ (2019/2020) when all three cameras (NYA, LYR, HOR) were in operation and recorded data, is

considered for this study. The season is spanning from October-March. On 73 days all three ASCs 197 have recorded data in an overlapping timespan (two hours). Days with a KP index lower than 2.5 198 were excluded. This reduced the set to 24 days. The solar wind conditions for each of the selected 199 events are shown in Figure 1. Most of these days are part of six longer lasting multi-day (5+days) 200 moderate geomagnetic storm events with high solar wind speeds (above 400 km/s; see panels b,e,h), 201 negative sym-H component (-20 nT to -40 nT; see panels c,f,i), re-occuring negative Bz (see panels 202 a,d,g) and negative drops of the local B_x component (shown in Figure 2 panels a,e,i). For compa-203 rability of the cases events with available data between 18-24 UT (21h-03h MLT nightside) were 204 selected, this means the data set has been reduced to find times at which all three ASC recorded 205 continuous data in the majority of the 18h-24h UT section. The set is now down-filtered to 12 days 206 as remaining candidates for the case study. After the filtering process, the imagery was visually 207 assessed and out of the days with intense auroral emissions the days with the least cloud cover (for 208 all three ASC) that were part of different geomagnetic storm events were selected. The selected 209 dates are the 28th October 2019, 29th January 2020 and the 22nd February 2020. The similarity 210 in the events lies in the background conditions as all selected dates are substorm events with KP 211 indices higher than 2.5. The events are in the night sector to meet similar geomagnetic conditions 212 and types of the aurora. Particle precipitation is strongly visible on all of them. The randomness 213 in the case study events is achieved through limited data availability and cloud cover on the ASC 214 images, providing us a representative set of case studies with similar conditions, but spread over the 215 whole season. 216

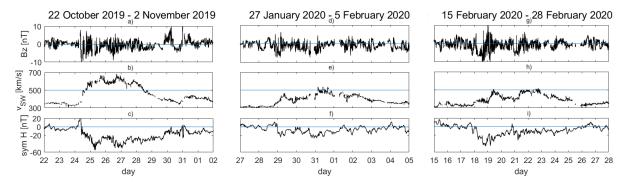


Fig. 1. Solar wind data recorded during the selected case study dates: 28th October 2019 (panels a-c), 29th January 2020 (d-f) and 22nd February 2020 (g-i). Panels a,d,g show the solar wind/ interplanetary magnetic field Bz component, indicating reconnection with Earth's magnetic field when negative. Panels b,e,h shows the solar wind flow speed which increases to over 400 km/s during geomagnetic storms, while the sym-H component (a measure of the ring current), shown in panels c,f,i, is abruptly negative during storms with a change of sym H > 20 nT.

4. Observations

Figure 2 shows data from the selected case study dates: 28th October 2019, 29th January 2020 and 22nd February 2020 from 18-24 UT. In panels (a,e,i) the horizontal magnetic field component

 B_x is shown. The decreases in B_x indicate substorm events at the respective stations (NYA-red,

LYR-green, HOP-black, BJN-orange, HOR-blue). Times of intense particle precipitation can be in-221 vestigated by presenting the center pixel column of the ASC images as a time series- this time versus 222 latitude plot is named a keogram. The keograms reveal times of auroral activity above zenith of the 223 recording ASC. The keograms presented in panels (b,f,j) show images that originate from the LYR 224 ASC using the filter for 557.7 nm (green) emissions. Here we associate the auroral activity, visible 225 bright auroral emissions, with times of particle precipitation. Panels (c,g,k) and (d,h,l) show σ_{ϕ} 226 and S4 scintillation index data recorded by the scintillation receivers in the respective stations. The 227 indices are used to quantify plasma structuring in location, intensity and scale. S4 indicates plasma 228 structuring on scales down a few meters, below the Fresnel radius, while σ_{ϕ} indicates structuring 229 above the Fresnel radius (hundreds of meters to km). 230

231 4.1. Altitude of Plasma Structuring

Different regions in the ionosphere are influenced by different phenomena characteristic for the 232 specific altitudes ranges. Determining the altitude in which the GNSS signals are disrupted is 233 crucial to relate the disturbances observed through the phase scintillation index to physical phe-234 nomena. The images of green and red auroral emissions are projected to their estimated emission 235 altitudes (150 km and 250 km respectively). The scintillation indices are projected to different al-236 titudes (piercing points) to find out whether the observed disturbances happen in the same altitude 237 range as the auroral emissions and whether they correspond better to the green or red emissions. 238 For this, different piercing points between 100 km to 350 km were tested. For the green auroral 239 emissions (150 km), it is observed that for lower (100 km) or higher (200 km) piercing point al-240 titudes, the elevated σ_{ϕ} values appear randomly placed with regards to the auroral forms. When 241 observing small patches of auroral emissions further east/west of zenith, one observes the elevated 242 σ_{ϕ} values further east/west of the patch when the ionospheric piercing point is chosen low/high. 243 However, when choosing the same altitude as piercing point as the green auroral emission altitude, 244 we find that the elevated σ_{ϕ} values at the boundaries and aligning well with the evolution of auroral 245 forms. Using imagery from all three ASC and all five receivers, elevated σ_{ϕ} values are consistently 246 found at the edges of the auroral forms for this piercing point altitude. This behaviour is observed 247 for chosen piercing point for satellites regardless of satellite azimuth or elevation, just as in Semeter

for chosen piercing point for satellites regardless of satellite azimuth or elevation, just as in Semeter et al. (2017). The elevated σ_{ϕ} values are increased at the boundaries of the auroral emissions, invariant with different satellite elevations. Another evidence for this is given when projecting data of only the co-located imager and scintillation receiver on the same plot, a case where the mapping altitude is irrelevant. We again observe elevated σ_{ϕ} values at the pole-ward edges of the aurora, see the video using data from only the NYA ASC and NYA scintillation receiver in supplementary material. This indicates that the signal is impacted right at the altitude of green auroral emission. This behavior has not been detected in the same extent with regards to red auroral emissions.

The emission altitude of green auroral emissions is used as the piercing point altitude for the study along with ASC images of the green auroral emissions.

258 4.2. Case Studies: Spatial and Temporal Evolution of Plasma Structuring

ASC imagery of the three case study events is shown in this section. The ASC images (green aurora projected to 150 km) are plotted onto the geographical coordinates and a map of Svalbard shown in

²⁶¹ contrast. On top of that the observed phase scintillation data (piercing point 150 km) is displayed.

The phase scintillation data is referred to as slightly elevated σ_{ϕ} above 0.2 rad, moderately elevated 262 σ_{ϕ} above 0.3 rad, strongly elevated σ_{ϕ} above 0.5 rad and very strongly elevated σ_{ϕ} above 0.7 rad. 263 Not all of the measured data during auroral activity of that day is shown: very faint aurora, forms 264 that are not classified as part of a process showing arcs/spirals/bands, forms that are mostly cut 265 off by the FOV or repetitive images are excluded. The auroral activity is fluctuating in intensity, 266 and can decrease to low intensity values or vanish shortly from the observation location before it 267 onsets again. Whether low intensity aurora and no aurora is observed is difficult to distinguish, as 268 it depends on the chosen intensity scale what will be visible. In the following, intensities measured 269 under 10 kR-NYA/20 kR-LYR are referred as no visible aurora. This threshold has been chosen as 270 the elevated phase scintillation index is not observed in combination with intensities under 10 kR 271 for NYA and 20 kR for LYR before the auroral onset. It can however be observed after the aurora 272 vanished, this may have other reasons than auroral emissions below 10 kR/20 kR as discussed later. 273 Only the representative images are shown in the results, but the remaining images were combined 274 to videos attached in supplementary material. In the following especially the high intensity values 275 of the aurora may be influenced by the way the camera and brightness is calibrated and calculated. 276

They are however a good measure for comparison between and relation to the σ_{ϕ} indices.

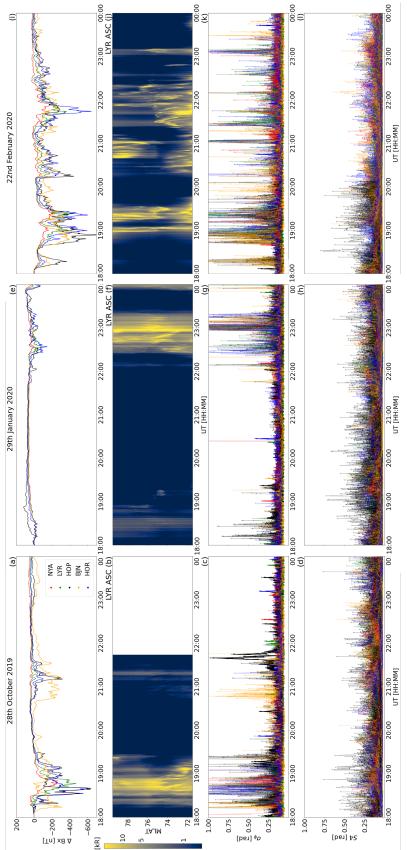


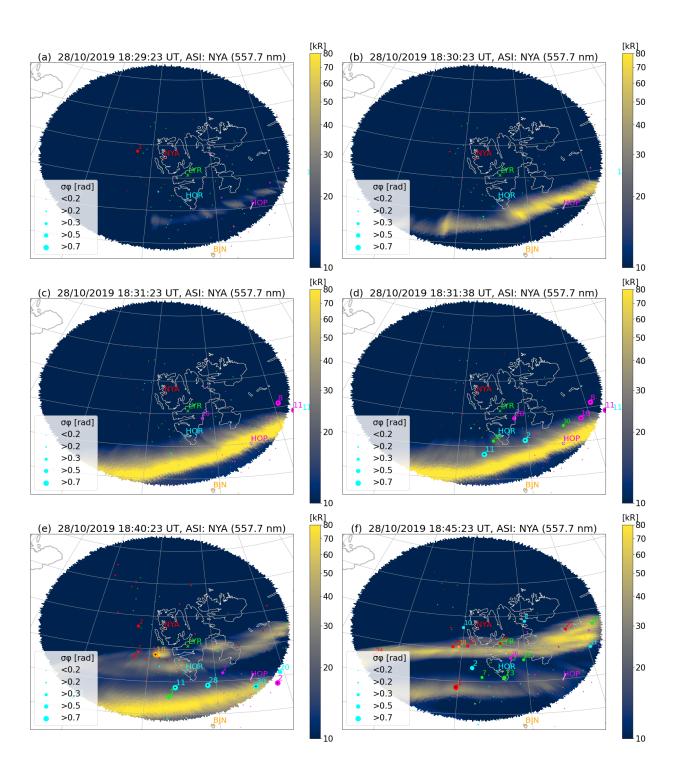
Fig. 2. Data from the selected events: 28th October 2019, 29th January 2020 and 22nd February 2020 from 18-24 UT. In panels (a,e,i) the horizontal magnetic field component B_x is shown. The dips indicate substorm events at the respective stations (NYA-red, LYRgreen, HOP-black, BJN-orange, HOR-blue). Panels (b,f,j) shows the 557.7 nm emissions intensity observed by the LYR ASC in form of a keogram. The bright auroral emissions correspond to times of particle precipitation. Panels (c,g,k) and (d,h,l) show σ_{δ} and S4 scintillation index data recorded by the scintillation receivers in the respective stations. They quantify the location, scale and intensity of plasma structuring.

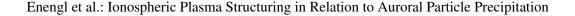
4.2.1. 28th October 2019

The first selected event of this season is the 28th October 2019, see Figure 2(a-d). The local mag-279 netic field B_x component shows a decrease between 18:10-19:30 UT measured at NYA (shown in 280 red), LYR (green), HOP (black), BJN (orange), HOR (blue) and indicating a substorm. At HOP 281 the magnitude of the decrease is especially large, up to a change of -600 nT in the B_x component, 282 followed by HOR and LYR. A less severe decrease in the B_x component is observed around 20:40-283 21:40 UT (strongest in BJN and HOP, lightest in LYR and NYA). At the same time auroral activity 284 is observed with the LYR ASC. Intense precipitation lasts from 18:15 to 19:30 UT and light pre-285 cipitation from 20:45-21:20 UT. The latter shows only faint aurora, which is also reflected in the B_x 286 component measurements, where we observed only light variations for LYR. Elevated σ_{ϕ} coincides 287 with particle precipitation between from 18:15 to 19:30 UT, this is observed for all stations. The 288 later time interval or particle precipitation is likewise accompanied by elevated σ_{ϕ} , but disturbances 289 are not in all stations recorded. BJN (panel c, shown in orange) records elevated σ_{ϕ} about 25 min-290 utes before the strongest emissions are observed at LYR (panel b), but σ_{ϕ} increases right as the B_x 29 component at BJN drops. The response of S4 to the particle precipitation is not as clear as for σ_{ϕ} , 292 but a slight increase of S4 especially in the measurements at HOP (panel a, shown in black) may 293 be observed around 19 UT. Elevated S4 values can be an indication for diffractive effects (Yeh and 294 Liu, 1982), but are not discussed here. 295

The 28th October 2019 observations indicate a substorm event, auroral emissions and distur-296 bances in the phase scintillation index between around 18:10-19:30 UT and 20:40-21:40 UT (Figure 297 2). This long and intense auroral activity can be best viewed from the NYA ASC as shown in Figure 298 3. The Figure shows a time evolution (a) to (h) from when the auroral activity commences to when 299 it fades out. This arc is also observed from the HOR ASC and LYR ASC. Videos of the arc from the 300 NYA and LYR station are shown in the supplementary material. In Figure 3a at 18:29:23 UT light 301 auroral activity (<30 kR) starts in the south-east quarter of the ASC's FOV, no elevation of σ_{ϕ} is 302 observed at this stage. A minute later, at 18:30:23 UT, the auroral activity is at double its intensity 303 (>60 kR) forming a faint arc. No elevation of σ_{ϕ} values is yet measured. At 18:31:23 UT the auroral 304 intensity is reaching its intensity maximum (>80kR) and now very strongly elevated σ_{ϕ} values are 305 observed in the east. The delay between auroral emissions until elevated σ_{ϕ} values are observed 306 is in this case 2 minutes. The underlying processes causing elevated σ_{ϕ} values/ plasma structuring 307 may come with a time delay related to the recombination rates of the precipitating electrons arriv-308 ing in the E-region ionosphere. The fast flows of the injected electrons could drive a two-stream 309 instability. The difference in drift velocity between the electrons and ions, which collide with neu-310 trals, could in specific drive the modified two-stream instability, the Farley-Buneman instability 311 (Farley Jr., 1963). Shortly after, at 18:31:38 UT, very strongly elevated σ_{ϕ} values measured pole-312 ward of the established auroral arc (with an intensity still over >80 kR) are observed. Instabilities 313 working specifically at the boundary of the particle stream may be dominant at this point, leading 314 to very localized elevated σ_{ϕ} values along the pole-ward boundary. Note that the plasma density is 315 lower at the pole-ward boundary than the equator-ward boundary. The intensity fluctuates (down to 316 <60 kR), but the σ_{ϕ} values stay elevated over the next 9 minutes. At 18:40:23 UT a second arc (with 317 an intensity <40 kR) forms in the center of the ASC image combined with very strongly elevated 318 σ_{ϕ} . The southern arc is still very intense (>80 kR) with pole-ward very strongly elevated σ_{ϕ} values. 319 At 18:45:23 UT the arc in the center intensifies (>60 kR) and moderately elevated σ_{ϕ} values are 320

observed pole-ward of the central arc and stronger elevated σ_{ϕ} values between the arcs. The south-321 ern arc has faded out (with an intensity < 45 kR). At 15:52:38 UT the arc structure is dissolved 322 and the intensity is decreased to similar values (<60 kR) as at 18:30:23 UT (before we observed 323 elevated σ_{ϕ}). Occasional very strong σ_{ϕ} values are still observed alongside strong and moderate σ_{ϕ} 324 values. This could be due to a 'memory effect', the precipitation has moved away, but the E-region 325 ionosphere has not yet restored its original state and electrons may still recombine and structures 326 persist until the amount of excess particles has declined. At 18:54:23 UT, the intensity (<30 kR) 327 has decreased to its starting values (panel a). However occasional moderately to strongly elevated 328 σ_{ϕ} values are measured. Eastwards, a low intensity auroral patch is co-located with a very strongly 329 elevated σ_{ϕ} value. The E-region ionosphere is still unstable as the structuring process continues 330 even for weak auroral emissions. 331





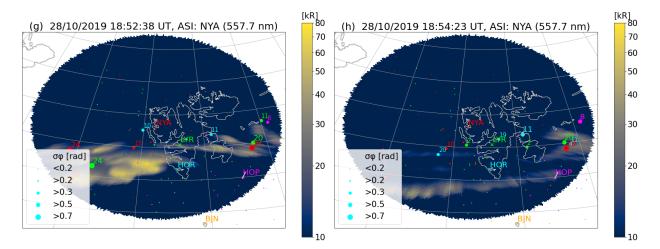


Fig. 3. The projection of the 557.7 nm emissions shown together with the phase scintillation index σ_{ϕ} over a map of Svalbard and geographical longitude/latitude on 28th October 2019 as seen from the NYA ASC for different time instances. Brighter yellow auroral emissions mean stronger intensity and bigger markers mean stronger σ_{ϕ} . The markers represent σ_{ϕ} measurements from NYA in red, LYR in green, HOR in cyan, HOP in magenta and BJN in orange. Panels a and b show the commencement of an auroral arc. Panels c and d show the arc at its maximum intensity (>80kR). In panel d we can point out the very strongly elevated σ_{ϕ} pole-ward of the arc. Panels e and f show two auroral arcs and corresponding σ_{ϕ} values. Panels g and h show fading of the aurora and decrease in σ_{ϕ} . Even though panels a and h show similar auroral intensities (<30kR) they show different levels of σ_{ϕ} : panel a - no elevated σ_{ϕ} , panel h - different levels of elevated σ_{ϕ} . The aurora shown in this figure is classified as arcs.

332 4.3. 29th January 2020

On the 29th January 2020 decreases in the local B_x component and strong auroral particle pre-333 cipitation combined with increases in σ_{ϕ} values are observed, see Figure 2(e-h). There are minor 334 variations in the B_x component (about 50 nT) and light precipitation between 18 and 19 UT (panel 335 f). After 21:50 UT a fluctuations in the local B_x components is observed. The dip in B_x reaches its lo-336 cal minimum at a change of around -200 nT at 22:20, a change much weaker than in the first studied 337 case. Nevertheless, the auroral emissions are intense, especially around 23 UT. The σ_{ϕ} values are 338 elevated strongly in the same time frame as precipitation is observed. The S4 index does not show 339 an as strong correlation to particle precipitation as σ_{ϕ} . In Figure 4 the event is shown as observed 340 by the LYR ASC. The first 30 minutes the auroral activity is fluctuating and moderately elevated 341 σ_{ϕ} values are found for faint aurora (<45 kR). One example is shown at 22:53:08 UT where mod-342 erately elevated σ_{ϕ} values are measured, but the auroral emissions are beneath (<30 kR). Whether 343 there is a time delay in this selected case between the auroral onset and the elevated σ_{ϕ} values or 344 not is difficult to determine as the auroral activity does not increase monotonically here as in the 345 case on 28th October 2019. At 22:55:53 UT the intensity increases (>80 kR) and moderately to 346 strongly elevated σ_{ϕ} values located within and surrounding the auroral form are measured. A high 347 intensity spiral-shaped aurora (>80 kR) is observed at 22:56:38 UT with very strongly elevated σ_{ϕ} . 348 A spiral-shaped aurora has previously been linked to the Kelvin-Helmholtz instability (Hallinan, 349 1976). The strongly elevated σ_{ϕ} values follow the boundary of the auroral spiral neatly, just as in 350 3d where the elevated σ_{ϕ} values follow the pole-ward boundary of the auroral arc. At 22:57:38 UT, 351 the auroral spiral and high intensity area has grown. We continue to observe very strongly elevated 352 σ_{ϕ} values at the boundary of the form, but now also within the form. At the boundaries of the spiral 353 form, we observe primarily elevated σ_{ϕ} values pole-ward of the form and only rarely equator-ward. 354 A minute later at 22:58:38 UT the form has shrank and is still surrounded by very strongly ele-355 vated σ_{ϕ} values. Elevated σ_{ϕ} values continue to be observed for over 15 minutes after the auroral 356 intensity decreased below 40 kR. At 23:14:23 UT the auroral intensity is beneath 30 kR and very 357 strongly elevated σ_{ϕ} values are still observed. This is another indication that there may be longer 358 lasting structuring processes, after the precipitation has declined. The observations are confirmed 359 by the NYA and HOR cameras. The images recorded by the LYR and NYA camera for this event 360 are attached as videos in the supplementary material. 361

362 4.4. 22nd February 2020

On the 22nd of February multiple local negative changes in B_x , auroral emissions and elevated 363 σ_{ϕ} values are measured, see Figure 2(i-1). In this event also S4 is elevated, mostly in the section 364 before 20:30 UT (panel 1). The forms before 20:30 UT are very turbulent and fast moving (video 365 of data from all three imagers shown in the supplementary material). The data before 20:30 UT 366 are due to its fast and dynamic forms not conclusive on the spatial relation between auroral forms 367 and elevated σ_{ϕ} values. Here we discuss the auroral forms after 20:30 UT as viewed from the 368 LYR ASC. In Figure 5, at 21:00:38 UT, a band has formed in the east. The band appears fast and 369 intense (>60kR), with immediate strongly to very strongly elevated σ_{ϕ} values on the north-west-370 ward boundary of the band. No time delay between the abrupt and intense auroral onset and elevated 371 σ_{ϕ} values is observed. At 21:02:08 UT the intensity has reached over 80 kR and the elevated σ_{ϕ} 372 values are found in the center of the band. Here, the elevated σ_{ϕ} values are first observed at the 373

boundary, but the form expands and intensifies so that the elevated σ_{ϕ} values move into the band-374 shape. Note that more GNSS satellites are crossing the band / are close to the band equator-ward of 375 HOR, but none of these are experiencing elevated σ_{ϕ} values. Barely any satellites are in the band's 376 vicinity above NYA latitudes. At 21:02:28 UT the auroral band starts fading out (<70 kR) leaving 377 very strongly elevated σ_{ϕ} values pole-ward of the band. This is another example of how elevated σ_{ϕ} 378 values are still measured for faint auroral activity. A single-point very strongly elevated σ_{ϕ} index 379 can still be observed until three minutes after intense aurora. Only an auroral patch (<50 kR) is 380 still visible in the east, here shown at 21:04:53 UT. After some minutes without strong activity, a 381 faint auroral band (<50 kR) moves into the FOV around 21:21:53 UT. At 21:23:23 the form has 382 increased its intensity to over 80 kR but only a slightly elevated σ_{ϕ} value is observed. Over one 383 minute later at 21:24:38 UT the high intensity area of the aurora grows bigger and very strongly 384 elevated σ_{ϕ} values are found at the boundaries, especially pole-ward. Here we see again, how the 385 auroral activity grows and expands before elevated σ_{ϕ} values are observed. At 21:25:52 UT the 386 form becomes more complex, still reaching over 80 kR, and strongly to very strongly elevated σ_{ϕ} 387 values are measured at the boundaries and within the form. At 21:32:38 UT (7 minutes later) the 388 auroral intensity has decreased (<50 kR) and only moderately to slightly elevated σ_{ϕ} values are 389 observed pole-ward of the band shaped aurora. In both here presented cases, elevated σ_{ϕ} values are 390 measured for faint aurora, suggesting plasma structuring to continue after the auroral activity has 39 declined. On the 22nd February 2020 we observe fast-moving forms, and the elevated σ_{ϕ} values are 392 not solely found pole-ward but also equator-ward. As they are changing more rapidly than the other 393 forms, the elevated σ_{ϕ} values may be following the auroral activity with a delay. 394

5. Discussion

We have studied elevated phase scintillation indices in relation to regions of auroral emission with 396 data from three different substorm events in detail. Three distinct auroral forms were considered: 397 arcs, spirals and bands. While the observed arcs were stable, the other forms were more dynamic. 398 The first selected substorm event (28th October 2019) shows arcs. For the first observed arc, we 399 found a delay in the onset of elevated σ_{ϕ} values of 2 minutes, from when auroral activity (<30 kR) 400 was observed, see again Figure 3a-c. The very strongly elevated σ_{ϕ} values are then observed at 401 the pole-ward edge of the arc (Figure 3d-f) and between multiple arcs (3e,f). Even for faint aurora, 402 very strongly elevated σ_{ϕ} values are measured for at least 2 minutes after the auroral intensity 403 has decreased to under 60 kR, Figure 3g-h, moderately elevated σ_{ϕ} values are observed for over 404 4 minutes after. In the substorm event on 29th January 2020 an auroral spiral was observed. For the 405 spiral, the intensity is not monotonically increasing as it was for the arc. Moderately elevated σ_{ϕ} 406 values are observed even for faint Aurora (<30 kR), see again Figure 4a. The elevated σ_{ϕ} values 407 are observed on the edges of the auroral spiral, see Figure 4c-e. Over 10 minutes after the auroral 408 intensity decreases under 60 kR very strongly elevated σ_{ϕ} values are still measured, see Figure 4f. 409 The last of the studied substorm events (22nd February 2020) shows auroral bands. The abrupt and 410 fluctuating intensity during the onset of the band caused auroral emissions with elevated σ_{ϕ} to be 411 observed right away without a delay, see Figure 5a. The second band first shows emissions above 412 60 kR and it takes over a minute to observe elevated σ_{ϕ} values, see Figure 5e. For both bands we 413 observe elevated σ_{ϕ} values at the pole-ward boundary of the auroral form, see Figure 5(b,d&f,g). 414 For the first band elevated σ_{ϕ} values are observed also at the west-ward boundary and within the 415

form, see Figure 5b. When the auroral bands fade out, elevated σ_{ϕ} values are still observed over 1 to 7 minutes (first and second example respectively), see Figure 5(c,d & g,h).

In summary, we observed: (1) Elevated phase scintillation indices correspond consistently well to the spatial and temporal evolution of auroral forms in the green emissions (oxygen, 557.7 nm) altitudes, which means particle precipitation into the ionospheric E-region is a driver for plasma structuring. (2) There may be a time delay between the temporal evolution of aurora (f.e. commencement and fading of auroral activity) and elevated phase scintillation index measurements. (3) The elevated phase scintillation indices are observed at the boundary of the auroral emissions (poleward for discrete and stable arcs and on all boundaries for bands and spirals).

When the ionospheric piercing point for the navigation satellites is chosen at the same altitude as 425 the projections of the green (557.7 nm) auroral emissions, the elevated σ_{ϕ} values are consistently 426 observed on the auroral edges, invariant of their elevation and azimuth. The same behaviour is 427 found when using only the co-located ASC and scintillation receiver (see the video with data from 428 the NYA ASC and scintillation receiver in supplementary material), a case where the mapping 429 altitude becomes irrelevant. Thus we suggest the σ_{ϕ} values to be effected at the altitude of the green 430 auroral emissions. This means that E-region irregularities and instabilities cause observed plasma 431 structuring. The findings are in agreement with Kinrade et al. (2013); Semeter et al. (2017), who 432 also find better correlation of observed elevated phase scintillation indices with the 557.7 nm auroral 433 emissions. 434

Previous studies, from e.g. van der Meeren et al. (2014); Jin et al. (2016), linked phase scintil-435 lation mainly to the polar cap F-region ionosphere (especially to polar cap patches and tongues of 436 ionization). The effects of intense local particle precipitation in the auroral oval on elevated σ_{ϕ} has 437 been studied and referred to as a blobs type II by Jin et al. (2016). They also state that soft F-region 438 particle precipitation does not contribute much to plasma structuring processes in the nightside au-439 roral region. Our findings are in agreement with this, but we however show that more energetic 440 particle precipitation penetrating down to the E-region may be a main source and is found co-441 located with intense elevated σ_{ϕ} . The link between elevated σ_{ϕ} or phase scintillations and E-region 442 auroral particle precipitation has also been observed by several authors (e.g., Kinrade et al., 2013; 443 Forte et al., 2017; Loucks et al., 2017; Sreenivash et al., 2020; Makarevich et al., 2021). 444

The time delay strongly depends on the auroral dynamics with regards to auroral intensity. For a 445 case where the aurora appears out of quiet conditions and the auroral intensity increases monoton-446 ically, the phase scintillation indices are elevated around 1 (2) minutes after the intensity reached 447 over 60 kR (30 kR), see again the example of the auroral arc or the second auroral band. This is at a 448 level on which simultaneous and co-located aurora and phase scintillation index enhancements are 449 commonly found, even without further increases to higher intensity levels. The duration time delay 450 depends strongly on which intensity is used as a threshold. When the auroral intensity is increasing 451 (decreasing) non-monotonically, it the time delay may be even longer (shorter). In case of the spi-452 ral, or, the first auroral band considered, the elevated σ_{ϕ} is observed right after a prompt increase in 453 auroral intensity. 454

As the auroral intensity decreases monotonically, elevated σ_{ϕ} is observed down to levels of under 20 kR. Especially when the auroral intensity is fluctuating around these levels, we observe occasionally high elevated scintillation indices. We have observed elevated σ_{ϕ} between 1 to over 7 minutes after auroral intensities dropped back under 60 kR. The effect of a time-delay after decreasing auroral emissions is stronger than the time-delay of the elevated σ_{ϕ} value onset. Both may only be

reached for certain conditions. For stable forms, the relation may be more clear than for dynamic forms.

Particle precipitation is a signature observed as a part of field-aligned currents (Carter et al., 462 2016; Xiong et al., 2020). They couple the magnetosphere to the ionosphere and can cause severe 463 phase scintillations through direct driving of the ionospheric plasma by structured precipitation or 464 electric fields resulting in elevated scintillation indices (Boström, 1964; Fæhn Follestad et al., 2020). 465 Particle precipitation is usually observed in times of high convection and convection patterns were 466 investigated with SuperDARN. We see a two cell pattern (during southward Bz) for the studied 467 dates. We do not observe strong flows on the nightside, and experience mostly growth in the cells, 468 linking to dayside reconnection (Dungey, 1961; Juusola et al., 2014). The negative By component 469 is prominent for the 29 January 2020, peaking at 22:00 UT the time we observe the spiral form, 470 and is driving an asymmetric ionospheric convection pattern (e.g., Weimer, 1995). During times of 471 particle precipitation, the twin cell convection pattern covers Svalbard and with this precipitation 472 area. 473

We propose that the time delay may occur through a 'memory effect'. Particle precipitation into the ionospheric E-region enhances conductance and causes a widespread irregularity dissipation and redistribution of energy (e.g., Ivarsen et al., 2021). The influx of particles increases the plasma density leading to ionization and prolonging recombination. When the particle precipitation declines or moves, it will still take some time for plasma structures to diffuse. The structuring process initiated during particle precipitation will carry on. Instabilities in the plasma lead to further structuring processes, irregularity dissipation and redistribution of energy.

The location of the elevated phase scintillation indices is also dependent on the spatial and tempo-481 ral auroral dynamics. When investigating a stable discrete auroral form, such as the arc, we observe 482 strongly elevated σ_{ϕ} (> 0.5 rad) solely at the pole-ward boundary. If the form is faster moving, such 483 as for spirals or bands, elevated σ_{ϕ} may be seen even westward or equator-ward, but nevertheless on 484 the boundaries. Fast-moving forms may move away quicker from the measurement location while 485 elevated σ_{ϕ} value are still observed at this location. For spirals and bands, elevated σ_{ϕ} are observed 486 also on the east and westward boundaries as well as on the boundaries within the shape. Semeter 487 et al. (2017) studied loss of lock (LL) events and their correlation with discrete auroral emissions. 488 They studied mostly fast-moving auroral forms, and found LL events on equator-ward (trailing) 489 edge of the auroral form, while the form is moving pole-ward. This agrees with our results, as we 490 propose that there is a time delay in the elevated scintillation index measurements. Therefore, for 491 fast-moving pole-ward forms, the LL events would be observed equator-ward as the form moved 492 past the measurement location further pole-ward. They note that although the pole-ward side of the 493 arc has a similar or even greater density gradient, no LL events were observed there. This explana-494 tion may be valid in our case, as we propose that the structuring process in fact happens pole-ward, 495 but the a fast-moving form has simply moved from its original place as the structuring persists. This 496 however would not explain the cases where we observe pole-ward elevated σ_{ϕ} values at multiple 497 stable arcs, as the equator-ward edge of the most pole-ward arc has a higher density gradient than 498 the pole-ward boundary of an equator-ward arc. Unless the equator-ward arcs are sufficiently more 499 intense than the pole-ward arcs, then the density gradient would be higher on the pole-ward edge of 500 equator-ward arcs (see Figure 3e and f with two arcs). 501

The energy source for driving instabilities in the E-region ionosphere can be manifold, such from as flow shears, from gradients or directly by kinetic energy. Instabilities associated with particle

precipitation are e.g. kinetic instabilities, two-stream instabilities. The flow of particles in fieldaligned currents can also produce current-driven instabilities (e.g., Kropotkin, 2016).

One instability that can be directly produced by a velocity shear (by particle precipitation/ electron beam) along the direction of or perpendicular to an externally imposed magnetic field is the Kelvin-Helmholtz instability (D'Angelo and Goeler, 1966; Hallinan and Davis, 1970; Pécseli, 2013). It can drive curls, or spirals as observed in the aurora (Hallinan and Davis, 1970).

Another relevant instability that may explain the observed plasma structuring is the Farley-Buneman instability (Farley Jr., 1963; Buneman, 1963; Treumann, 1997). The instability arises from the difference of the electron and ion velocity, caused by collisions of the ions with neutrals (Farley Jr., 1963). These conditions are given in the equatorial and polar E-region ionosphere, where this instability is typically found (Rogister and D'Angelo, 1970; Pécseli, 2013).

Various instabilities can arise in the ionosphere driven by e.g. currents, energetic particle streams or density gradients. However, the Kelvin-Helmholtz and Farley-Buneman instability do not only satisfy the encountered background conditions (E-region ionosphere, particle stream, collisions, availability of neutrals), but could also explain the observed behaviour of elevated σ_{ϕ} observations at the edges of different auroral forms (spirals, arcs, bands), the boundaries between the injected particles and the ionospheric E-region plasma.

The Kelvin-Helmholtz instability is extracting energy out of a shear flow along a boundary 521 (Treumann, 1997) and is associated with auroral spirals (Hallinan and Davis, 1970). This may ex-522 plain why we measure elevated σ_{ϕ} values on the outer boundaries and on boundaries whithin auroral 523 spirals. The Farley-Buneman instability provides a more general explanation for plasma structuring 524 observed at boundaries of different auroral forms, not only spirals. The auroral form boundaries are 525 where the electrons stream past ions, which collide with neutrals, which is what feeds the Farley-526 Buneman instability (Farley Jr., 1963; Buneman, 1963). While the Kelvin-Helmholtz instability is 527 working on bigger scales, the Farley-Buneman instability is operating on smaller scales. Multiple 528 instabilities may be responsible for the plasma structuring processes we observe. Whether it is the 529 difference in plasma density gradient that drives plasma structuring predominately at pole-ward 530 edges of auroral arcs remains remains unknown. Further case studies with even higher spatial and 531 temporal resolution are needed to understand the structuring process and to confirm which instabil-532 ity can cause delayed structuring processes in the E-region at the boundaries of auroral forms and 533 at pole-ward boundaries for auroral arcs. 534

535 6. Conclusion

In this study, the relation between auroral particle precipitation and plasma structuring was investigated. In summary, the temporal and spatial evolution of auroral forms and phase scintillation indices were studied. For this, three event days with similar background conditions (substorm events, particle precipitation and nightside events) showing clear strong auroral emissions were selected. Data was available from three ASI imagers (NYA, LYR and HOR) and five scintillation receivers (NYA, LYR and HOR, HOP, BJN). This provides us with 18 h of data in which we observed auroral forms such as arcs, spirals and bands. Our results show that:

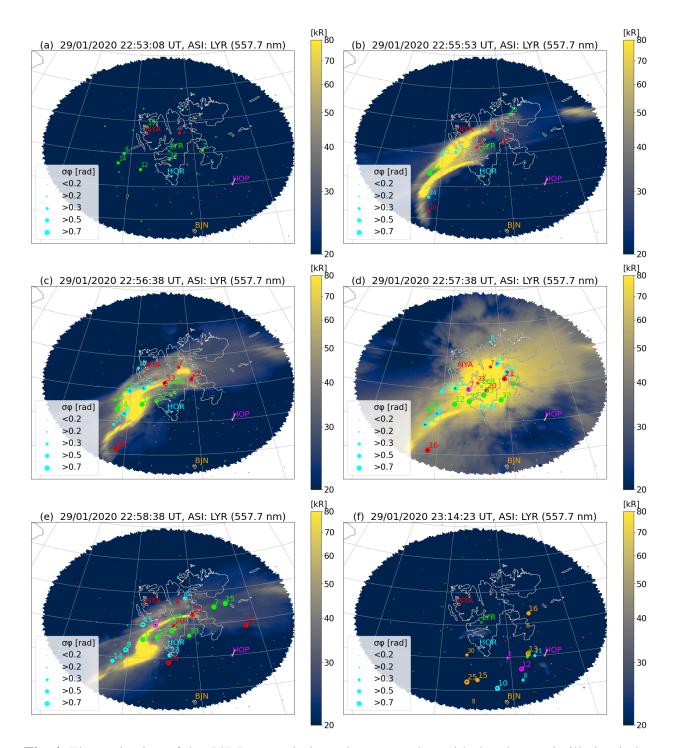
- (1) Elevated phase scintillation indices correspond consistently (invariant of satellite elevation and azimuth) to the spatial and temporal evolution of auroral forms when an ionospheric piercing

⁵⁴⁵ point for navigation satellites is chosen the same as the estimated green emissions (557.7 nm) al-⁵⁴⁶ titude (150 km). This suggests that plasma structuring in the ionospheric E-region is an important ⁵⁴⁷ driver for phase scintillations.

 - (2) There may be a time delay between the temporal evolution of aurora (e.g. commencement and fading of auroral activity) and elevated phase scintillation index measurements. Particle precipitation enhances the plasma density. When the precipitation declines or moves, it will still take some time for the plasma structures to diffuse. Until then, instabilities will further cause redistribution of energy and irregularity dissipation.

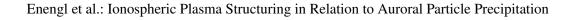
(3) The elevated phase scintillation indices are observed at the boundary of the auroral emissions. 553 For discrete and stable arcs elevated phase scintillation indices are predominately observed pole-554 ward and for faster moving shapes, including spirals and bands, on the boundaries of the form. 555 The irregularities and instabilities causing the elevated phase scintillation indices especially in 556 the E-region may be due to instabilities which are driven by energy at the boundary of auroral 557 forms, such as the Kelvin-Helmholtz instability (directly produced by a velocity shear such as 558 from particle precipitation) or Farley-Buneman instability (through fast flows at the boundaries). 559 Plasma structuring may predominately be observed on pole-ward boundary as the gradient in 560 plasma density is larger than it is on the equator-ward boundary. 561

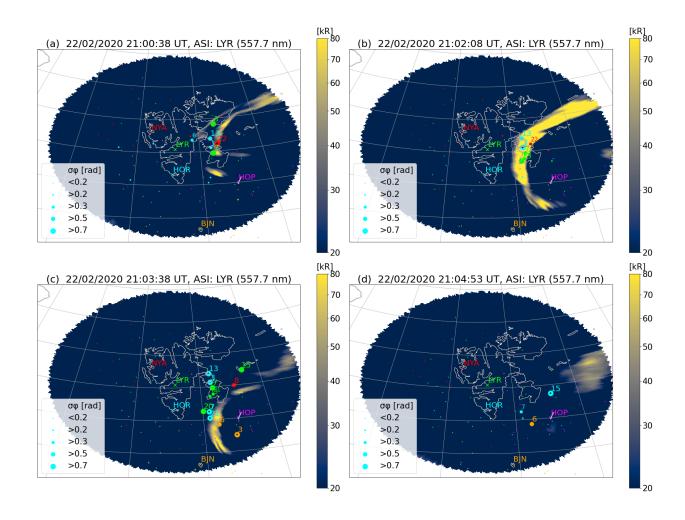
The question on why plasma structuring processes in the E-region are observed specifically at the edges of auroral forms, such as spirals and bands, and at polew-ard boundaries for auroral arcs remains open. Further case studies with even higher spatial and temporal resolution and bigger statistical studies investigating time-delay statistics are needed to understand the structuring process. In the future we also need to investigate further, which instabilities are caused during the plasma structuring processes and how they effect trans-ionospheric radio waves.



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Fig. 4. The projection of the 557.7 nm emissions shown together with the phase scintillation index σ_{ϕ} over a map of Svalbard and geographical longitude/latitude on 29th January 2020 as seen from the LYR ASC for different time instances. Brighter yellow auroral emissions mean stronger intensity and bigger markers mean stronger σ_{ϕ} . The markers represent σ_{ϕ} measurements from NYA in red, LYR in green, HOR in cyan, HOP in magenta and BJN in orange. Panel a shows elevated σ_{ϕ} values yet without strong Aurora. Panel b shows strong auroral emissions and paired with elevated σ_{ϕ} values. In panels c, d and e an auroral spiral has formed and the elevated σ_{ϕ} values are primarily observed at the boundary of the auroral spiral. In panel f elevated σ_{ϕ} values on top of the faded out form. The aurora shown in this figure is classified as a spiral. 20





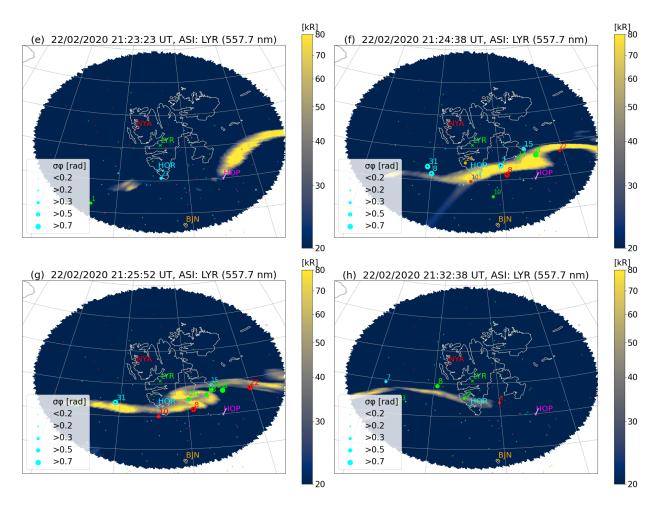


Fig. 5. The projection of the 557.7 nm emissions shown together with the phase scintillation index σ_{ϕ} over a map of Svalbard and geographical longitude/latitude on 22nd February 2020 as seen from the LYR ASC for different time instances. Brighter yellow auroral emissions mean stronger intensity and bigger markers mean stronger σ_{ϕ} . The markers represent σ_{ϕ} measurements from NYA in red, LYR in green, HOR in cyan, HOP in magenta and BJN in orange. Panel a shows a band-shaped Aurora and elevated σ_{ϕ} values west-ward. Panel b shows the intensified band, no with even stronger σ_{ϕ} values. Panels c and d show the fading of the auroral band, however σ_{ϕ} is still elevated. Panel e shows the commencement of the next auroral band, yet without strong elevated σ_{ϕ} values. Panels f and g show stronger auroral emissions paired with elevated σ_{ϕ} values at the boundary of the auroral band. In panel h elevated σ_{ϕ} is observed pole-ward of the fading auroral band. The aurora shown in this figure is classified as bands.

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