

High-resolution, basin-scale simulations reveal the impact of intermediate zonal jets on the Atlantic oxygen minimum zones

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Abstract

Intermediate, eastward zonal jets connect the oxygen-rich western boundary of the Atlantic Ocean with the oxygen minimum zones on the eastern boundary. These jets are not well represented in climate models because the low horizontal resolution of these models yields excessive viscosity. We use two physical-biogeochemical model configurations of the Tropical Atlantic Ocean to show that the increase in resolution results in more robust intermediate zonal jets and a better representation of the OMZs. The OMZ structure is distorted in the low-resolution run as westward jets advect low-oxygen waters from the eastern boundary further west than in the climatology. The emergence of more robust eastward jets in the high-resolution run alleviate this problem and provide a more realistic structure of the OMZs. The asymmetry between the effect of westward and eastward jets occurs because the former are associated with homogenous potential vorticity regions in the eastern boundary while the latter are associated with potential vorticity gradients. Intermediate, eastward jets constrain the westward expansion of the OMZs by supplying oxygen to their western edge. Within the isotropic OMZs, high resolution better represents the boundary current system and eddying processes at depth which are important in the redistribution of low oxygen values from the eastern boundary. Our results show that basin-scale, high-resolution simulations reproduce more accurately the transfer of energy across scales that results in robust zonal jets as well as their impact on the biogeochemistry. Accurate model predictions provide a pathway to disentangle natural and anthropogenic causes of ocean deoxygenation.

Abstract

Intermediate, eastward zonal jets are important conduits of oxygen across the Atlantic Ocean as they connect the oxygen-rich western boundary of the basin with the oxygen minimum zones (OMZs) on the eastern boundary. These jets are not well represented in climate models because the relatively low horizontal resolution of these models usually yields excessive viscosity. We use two physical-biogeochemical model configurations of the Tropical Atlantic Ocean to show that the increase in resolution results, on average, in more robust intermediate, eastward zonal jets and a better representation of the OMZs. The OMZ structure is distorted in the low-resolution run as surface, westward jets advect low-oxygen waters from the eastern boundary much further west than in the climatology. The emergence of more robust eastward jets in the high-resolution run alleviate this problem and provide a more realistic structure of the Atlantic OMZs. The asymmetry between the effect of westward and eastward jets occurs because the former are associated with homogenous potential vorticity regions originating in the eastern boundary while the latter are associated with potential vorticity gradients. Intermediate, eastward jets constrain the westward expansion of the OMZs by supplying oxygen to the western edge of the OMZs. Within the more isotropic OMZs, high resolution allows a better representation of the boundary current system and eddying processes at depth which are important in the redistribution of low oxygen values from the productive eastern boundary. Our results show that basin-scale, high-resolution simulations reproduce more accurately the transfer of energy across scales that results in robust zonal jets as well as their impact on the ocean biogeochemistry. Accurate model predictions provide a pathway to disentangle natural and anthropogenic causes of ocean deoxygenation.

Plain Language Summary

Long-term averages of ocean velocities reveal the existence of east-west, alternating jets along multiple latitudes. These are difficult to observe and model because of their small speeds at great depths. Despite their low intensity, in the long-term they can transport tracers across the ocean basins with oxygen being a very important one as it provides conditions for aerobic respiration in so-called oxygen minimum zones (OMZs) on the eastern side of the basin. Long-term measurements show that oxygen concentrations are decreasing in various regions of the ocean and that OMZs are expanding, which can be a problem as these regions may become inhospitable for aerobic life. That is why we

46 need to understand the processes that supply oxygen to OMZs and are important for their
47 evolution with time. Models can be used as tools for testing hypotheses regarding the
48 expansion or contraction of OMZs in the future. However, models must be shown to cor-
49 rectly simulate the dynamics and biogeochemistry of the region as a whole. Our results
50 show that zonal jets are important in structuring the OMZs and that higher-resolution,
51 basin-scale simulations are necessary to correctly simulate their impact on oxygen con-
52 centrations in the ocean.

53 1 Introduction

54 Over the last decade, several studies have shown that regions with low oxygen con-
55 centrations are expanding over the world’s oceans, a phenomenon which has been termed
56 ocean deoxygenation (Breitburg et al., 2018; Levin, 2018; Stramma et al., 2008). These
57 changes are driven by a combination of anthropogenic climate change and the natural
58 variability of the ocean. As climate change warms the upper ocean it reduces oxygen sol-
59 ubility, increases upper ocean stratification and thus reduces oxygen mixing as well as
60 induces changes in respiration rates (Levin, 2018). The continuous reduction of oxygen
61 may affect metabolic pathways by favoring denitrification or anamox as the preferential
62 conversion processes of fixed nitrogen to N_2 (Ryabenko et al., 2012). Natural ocean vari-
63 ability may also alter the oxygen supply as major climate fluctuations such as El Niño,
64 the North Atlantic Oscillation and the Atlantic Meridional Mode change atmospheric
65 and oceanic conditions over relatively long periods which makes it challenging to sep-
66 arate natural oscillations from long-term trends (Bryden et al., 2003). Disentangling the
67 natural and anthropogenically-induced oxygen variability requires the use of models as
68 prognostic or diagnostic tools, as they can be forced with different conditions which may
69 or may not include the effects of climate change and allow a detailed examination of spe-
70 cific processes. A pre-requisite for the use of models in this fashion is their ability to pro-
71 vide a realistic distribution of tracers and the velocity structure as compared to avail-
72 able observations.

73 The oxygen budget is a combination of physical and biological processes on coastal,
74 shelf and open seas which vary on spatial scales of a few to thousands of kilometers and
75 temporal scales of a few days to decades. Therefore, a thorough understanding of the
76 system requires a multi-scale approach in which most of the important processes are ad-
77 equately resolved. This is both an observational and modeling challenge. On one hand,

78 observations are limited by repeated, relatively sparse hydrographic stations or point-
79 wise long-term measurements. While extremely valuable in detecting and monitoring long-
80 term changes, they do not allow a complete understanding of specific processes that mod-
81 ulate oxygen concentrations in oxygen minimum zones (OMZs). On the other hand, hy-
82 drodynamic models are usually limited by trade-offs between domain size, horizontal res-
83 olution and subgrid-scale parameterizations which hinder their ability to appropriately
84 resolve all the important processes in the time-space spectrum. The coupling with bio-
85 geochemical models compounds the problem as more variables are simulated and a num-
86 ber of biogeochemical parameters need to be selected, adding more uncertainty to the
87 prognostics. Long-term predictions of oxygen concentrations in the ocean are based on
88 ensembles of climate models with sometimes divergent prognostics (Kwiatkowski et al.,
89 2020; Cabré et al., 2015) which may not allow a mechanistic understanding of specific
90 processes because of intrinsic differences between the models. Large biases in the mod-
91 eled oxygen distribution in the oceans have been attributed to the inability of models to
92 correctly simulate the equatorial and off-equatorial zonal, subsurface currents (Cabré et
93 al., 2015; Duteil et al., 2014; Oschlies et al., 2018).

94 Previous studies show the existence of alternating surface and deep zonal jets in
95 all ocean basins (Maximenko et al., 2008; Eden, 2006; Richards et al., 2006; Maximenko
96 et al., 2005; Nakano & Hasumi, 2005). The connection between the zonal jets and oxy-
97 gen concentrations has been mostly explored in the Pacific and Atlantic oceans, where
98 the jets are shown to modulate tracer fields by generating large-scale frontal regions in
99 the deep ocean (Delpech et al., 2020) and enhancing meridional oxygen gradients in the
100 OMZs so that mesoscale eddy mixing becomes an important supply mechanism (Lévy
101 et al., 2022; Brandt et al., 2012). Busecke et al. (2019) highlight the importance of cor-
102 rectly simulating the equatorial current system in the Pacific Ocean and how the OMZ
103 is sensitive to the magnitude and structure of the Equatorial Undercurrent. In the At-
104 lantic OMZs, eastward jets are known to be important conduits of oxygen as they con-
105 nect the more oxygenated western part of the basins with the OMZs in the eastern bound-
106 aries (Burmeister et al., 2019; Brandt et al., 2012, 2010, 2008; Stramma et al., 2008). Con-
107 sequently, the long-term variability of the jets may in turn affect oxygen concentrations
108 in the OMZs (Czeschel et al., 2015, 2011). For a recent overview on the phenomenology
109 and dynamics of zonal jets in the ocean the reader is referred to Cornillon et al. (2019)
110 and Kamenkovich et al. (2019) and references therein.

1.1 OMZs in the Atlantic Ocean

Two large-scale OMZs exist in the Atlantic Ocean. They arise as a combination of the subtropical circulation, which limits the ventilation of the eastern, tropical portion of the gyres, and the mostly wind-driven Eastern Boundary Upwelling Systems (EBUS) (Karstensen et al., 2008), characterized by high biological productivity and, consequently, high levels of respiration and remineralization in the water column, thus further decreasing oxygen concentrations.

In the Eastern Tropical North Atlantic (ETNA), a number of observations and modeling studies suggest that the equatorial current system has a large impact on the supply of oxygen-rich waters from the western part of the basin into the eastern part, thereby modulating the oxygen content of the OMZs (Duteil et al., 2014; Brandt et al., 2010). In fact, the sensitivity of the equatorial current system to natural oscillations may be a strong modulator of oxygen concentrations in the North Atlantic OMZ (Brandt et al., 2021). Therefore, a correct representation of the OMZ and its evolution relies on accurate simulations of the equatorial and off-equatorial current system. The ETNA is also characterized by a shallow OMZ, located near the coast and with a distinct vertical structure that is not captured by model simulations with resolutions up to 0.1 degrees (Duteil et al., 2014; Frenger et al., 2018; Thomsen et al., 2019).

The eastern tropical South Atlantic (ETSA) OMZ is located between the Equator and Angola-Benguela Frontal Zone (ABFZ) (Tchikalanga et al., 2018). The surface flow is characterized by the intermittent, southward Angola Current (Kopte et al., 2017), which transports water from the equatorial region along the coast until it meets the Benguela current, where the flow turns offshore forming the ABFZ (Monteiro et al., 2008). Oxygen concentrations in the OMZ core are lower than in the ETNA OMZ. This could be due to the geometry of the basin, with the OMZ core more isolated from the influence of the South Atlantic zonal jets, thus not receiving as much oxygen. Local biogeochemical processes in the water column associated with longer residence times may also lead to more oxygen consumption and result in lower oxygen concentrations.

In this work, we compare two coupled physical-biogeochemical simulations of the Tropical Atlantic ocean at different horizontal resolutions to show that the increase in resolution (i) allows the emergence of more robust intermediate zonal jets which have a major impact on the overall structure of the North and South Atlantic OMZs by lim-

143 iting their westward extent and supplying oxygen to the western edge of the OMZs and
144 (ii) provides a better representation of the coastal and shelf current system as well as
145 a more robust eddying field at depth which effectively redistribute low oxygen values that
146 originate from the productive eastern boundary upwelling regions.

147 2 Methods

148 We use version 1.1 of the Coastal and Regional Ocean Community Model (CROCO)
149 <https://www.croco-ocean.org/> coupled with the biogeochemical model PISCES in or-
150 der to investigate the sensitivity of the North and South Atlantic OMZs to the large-scale
151 ocean circulation in a climatological sense. The aim is to obtain a realistic representa-
152 tion of the oxygen structure in the OMZs of the Atlantic Ocean and its sensitivity to the
153 ocean circulation, in particular to the equatorial current system, so that the model can
154 be used for more accurate predictions of the temporal evolution of the oxygen content
155 in the basin.

156 We use initial conditions for the physical variables (Sea surface height, tempera-
157 ture, salinity and horizontal velocity) from January 1st. 2003 from the version 3.3.1 of
158 the SODA reanalysis product (Carton et al., 2018). Boundary conditions for the same
159 variables are obtained each 5 days for the same year and were downloaded from the Asia-
160 Pacific Data Research Center at [http://apdrc.soest.hawaii.edu:80/dods/public](http://apdrc.soest.hawaii.edu:80/dods/public_data/SODA/soda_3.3.1/)
161 [_data/SODA/soda_3.3.1/](http://apdrc.soest.hawaii.edu:80/dods/public_data/SODA/soda_3.3.1/). The year 2003 was chosen because it is considered a neutral
162 year both in terms of the North Atlantic Oscillation (NAO) (Chassignet & Xu, 2017)
163 and the Atlantic Meridional Mode (AMM). With neutral oceanic conditions, we used cli-
164 matological wind stress and heat fluxes from COADS as surface forcing. Initial and bound-
165 ary conditions for the biogeochemical variables were obtained from a climatological PISCES
166 run (Aumont & Bopp, 2006).

167 The full model domain encompasses the whole Tropical Atlantic Ocean, from 19°S
168 to 31°N and from 99°W to 15°W. We used two horizontal resolutions, namely 1/10° and
169 1/30°, which roughly translates into a nominal horizontal grid size of 10 km (TATL10)
170 and 3 km (TATL3), respectively.

171 Momentum advection is done with the 3rd-order upstream biased scheme and tracer
172 advection is done with the split and rotated 3rd-order upstream biased advection scheme
173 (RSUP3), where diffusion is split from advection and is represented by flow-dependent

174 hyperdiffusion and isopycnal rotation (Marchesiello et al., 2009). This scheme is shown
 175 to be more effective in reducing spurious diapycnal tracer mixing in sigma-coordinate
 176 models. In this case, the tracer mixing coefficient is a function of the absolute value of
 177 the local velocity and the grid size as $B = \frac{1}{12}|U|(\Delta x)^3$, where U is the scale of the ve-
 178 locity field and Δx is the grid size. Therefore, the smaller grid size of TATL3 yields a
 179 smaller diffusivity coefficient which may favor the emergence of zonal jets at depth.

Starting from the initial conditions from the SODA reanalysis, the physical and bio-
 geochemical model were spun up for 9 years with climatological surface forcing. Results
 shown in this study are from monthly and annual averages of the last 7 years (year 10
 to year 16) of model run. Daily averages are used in the calculation of the anisotropy
 coefficient in Section 4. The degree of anisotropy, α , is defined as (Huang et al., 2007),

$$\alpha = \frac{\langle u'^2 \rangle - \langle v'^2 \rangle}{\langle u'^2 \rangle + \langle v'^2 \rangle}, \quad (1)$$

180 where u' and v' are the zonal and meridional velocity anomalies from the averaged val-
 181 ues for the specific averaging period, and the operator $\langle \rangle$ indicates temporal averag-
 182 ing, which is done at 6, 12, 24, 36, 48 and 60 months. We calculate α for 4 subdomains
 183 which are representative of the western part of the basin adjacent to the OMZs, namely
 184 ETNA-West (36°W–23°W, 5°N–16°N), ETSA-West (30°W–10°W, 16°S–5°S), ETNA-OMZ
 185 (23°W–23°W, 5°N–16°N) and ETSA-OMZ (10°W–10°E, 16°S–5°S).

186 The isopycnic PV is calculated from the model output averaged between isopyc-
 187 nal surfaces as

$$q = \frac{1}{\rho_0}(\zeta + f)\frac{\Delta\rho}{H}, \quad (2)$$

188 where ζ is the relative vorticity, ρ is the potential density, both averages of the values
 189 on each layer, H is the layer thickness and $\Delta\rho$ the density difference between the lay-
 190 ers. Layer 1 is defined as the one between the 25.8 and 26.6 isopycnal surfaces and layer
 191 2 is the one between 26.6 and 27.4 isopycnal surfaces.

192 2.1 Oxygen Budget

193 The oxygen budget in the model is a combination of physical and biological pro-
 194 cesses as follows

$$\frac{\partial O_2}{\partial t} = \left(\frac{\partial O_2}{\partial t} \right)_{Physics} + \left(\frac{\partial O_2}{\partial t} \right)_{Biology}, \quad (3)$$

195 where

$$\left(\frac{\partial O_2}{\partial t} \right)_{Physics} = - \underbrace{u_h \cdot \nabla O_2}_{Horizontal\ Advection} - \underbrace{w \frac{\partial O_2}{\partial z}}_{Vertical\ Advection} + \underbrace{\frac{\partial}{\partial z} \left(k_z \frac{\partial O_2}{\partial z} \right)}_{Vertical\ Mixing}, \quad (4)$$

196 where K_z is the mixing coefficient, u_h are the horizontal components of the velocity vec-
197 tor and w is the vertical velocity.

198 The effect of biological processes on the oxygen concentration in PISCES is com-
199 puted as follows,

$$\begin{aligned} \left(\frac{\partial O_2}{\partial t} \right)_{Biology} = & \underbrace{(R_{o:c}^1 + R_{o:c}^2)(\mu_{NO_3}^P P + \mu_{NO_3}^D D)}_{New\ Production} + \underbrace{R_{o:c}^1(\mu_{NH_4}^P P + \mu_{NH_4}^D D)}_{Regenerated\ Production} \\ & - \underbrace{\lambda_{DOC}^* f(O_2) DOC}_{Remineralization} - \underbrace{G^Z Z - G^M M}_{Respiration} - \underbrace{R_{o:c}^2 Nitrif.}_{Nitrification}. \end{aligned} \quad (5)$$

200 Oxygen is produced during photosynthesis via new and regenerated production by diatoms
201 (D) and nanophytoplankton (P) and consumed by dissolved organic matter (DOC) rem-
202 ineralization, respiration by small (Z) and large (M) zooplankton and nitrification. λ_{DOC}^*
203 is the remineralization rate and $f(O_2)$ is a function that varies between 0 and 1 to rep-
204 resent oxic or anoxic remineralization, depending on local oxygen concentrations. Sto-
205 ichiometric oxygen to carbon ratios during ammonium conversion into organic matter
206 ($R_{o:c}^1$) and during nitrification ($R_{o:c}^2$) are set to 140:122 and 32:122, respectively. G^Z and
207 G^M are functions for zooplankton and mesozooplankton grazing whose details may be
208 found in Resplandy et al. (2011). More details about the model structure and param-
209 eterizations may be found in Aumont et al. (2003).

210 3 Results

211 3.1 Horizontal Structure of Oxygen and Velocity Fields

212 Overall, both simulations capture the climatological structure of the surface cir-
213 culation of the Tropical Atlantic. The 7-year averages of modeled surface velocities com-
214 pare equally well with climatological averages obtained from surface drifters (Laurindo

215 et al., 2017) (Fig. A1) as the main large-scale features, namely the eastward-flowing North
216 Equatorial Counter Current (NECC), the northern and southern branches of the westward-
217 flowing South Equatorial Current (SEC) and the North Brazil Current and its retroflec-
218 tion in the western part of the basin are correctly simulated.

219 A detailed model evaluation in terms of the horizontal and vertical tracer struc-
220 ture as compared to the CARS climatology is provided in the Appendix. The statistics
221 required for the Taylor and target diagrams (Figs. A2 and A3) are calculated from the
222 regridded 7-year model averages to match the resolution of the climatology. Given the
223 relatively low resolution of the climatology, these metrics serve to identify the ability of
224 the models to simulate the large-scale, climatological structure of the tracers. Both sim-
225 ulations provide a relatively good skill in the North and South Atlantic OMZs, with some
226 improvement in terms of correlation and bias in TATL3, particularly for nitrate and oxy-
227 gen, with correlation coefficients larger than 0.8, which shows that the model is able to
228 capture the correct range of variability and spatial patterns as compared to the CARS
229 climatology (Figs. A2 and A3). The improvement in TATL3 is also seen in the vertical
230 transects shown in the next section and in the Appendix. We must keep in mind, how-
231 ever, that our “truth”, namely the CARS climatology, is a product created from the in-
232 terpolation and extrapolation of sparse data. As such, it is also prone to errors, partic-
233 ularly in relatively sub-sampled regions such as the South Atlantic ocean. This can be
234 one of the reasons why the correlation between some of the model and climatology vari-
235 ables is relatively low at 650 m depth in the South Atlantic (Fig. A2). A useful discus-
236 sion on biogeochemical model evaluation is provided in (Doney et al., 2009; Jolliff et al.,
237 2009).

238 While both runs could be considered acceptable in terms of their climatological model
239 skill, a detailed visual inspection of the oxygen distribution at specific depths reveals large
240 differences both in terms of oxygen concentration as well as the geometry of the OMZs.
241 Fig. 1 shows a comparison between the 7-year average of TATL10 and TATL3 and the
242 CARS climatology at 250 m, 450 m and 650 m depths. TATL3 is more similar to the
243 the CARS climatology than TATL10 both in terms of oxygen concentrations and the struc-
244 ture of the OMZs, as seen by the $80 \mu\text{mol L}^{-1}$ isoline which is assumed to be the limit
245 of the OMZ. In TATL10, lower oxygen values at 250 m depth are advected much fur-
246 ther west than in TATL3. The westward excursions of the $80 \mu\text{mol L}^{-1}$ isoline suggest
247 a strong influence of alternating zonal jets in the equatorial/tropical region on oxygen

248 concentrations, with westward jets advecting low oxygen waters and eastward jets ad-
249 vecting high oxygen waters given the large-scale zonal oxygen gradient. TATL10 also strongly
250 underestimates oxygen concentrations in the ETSA region at 250 m with values as low
251 as $10 \mu\text{molL}^{-1}$ as opposed to TATL3 which is more similar to CARS at this depth. At
252 450 m depth, the impact of eastward equatorial currents is seen in both simulations as
253 they effectively transport high oxygen waters and separate the two OMZs. At 650 m depth,
254 the OMZ practically vanishes in TATL10 while TATL3 shows a structure more similar
255 to the climatology.

256 At 450-m depth and below, the high oxygen concentrations in the western part of
257 the basin, between 8°N and 18°N , persist in TATL10 as opposed to TATL3, which is more
258 in line with the climatology. This suggests that eddying processes at depth, responsi-
259 ble for tracer redistribution, are not properly resolved by the 10-km resolution run which
260 can have serious consequences for climate simulations unless these processes are adequately
261 parameterized.

262 The impact of the zonal jets on the OMZ structure at these depths is shown in Fig.
263 2 with the 7-year averaged zonal velocities and the OMZ boundary at 250 m, 450 m and
264 650 m depths. At 250 m, westward jets extending all the way across the basin in TATL10
265 are responsible for the unrealistic westward extension of the OMZ. In TATL3 these jets
266 become weaker on the eastern part of the basin with a reduced ability to advect low oxy-
267 gen waters westward. In addition, the magnitude of the zonal jets in TATL3 is larger
268 than in TATL10, particularly eastward jets. They not only supply oxygen-rich waters
269 from the western part of the basin but also delimit the westward extent of the OMZ bound-
270 ary as they weaken past that point both in the ETNA and the ETSA. Interestingly, the
271 number of zonal jets increases in the North Atlantic in TATL3, particularly in the west-
272 ern part of the basin. The zonal jets located at around 9°N , 12°N and 15°N seem to
273 be particularly relevant for the ETNA OMZ. At 650 m the pattern is similar, with the
274 zonal jets still robust in TATL3.

275 3.2 Vertical Structure

276 The meridional distribution of 7-year averaged oxygen concentrations at 23°W (Fig.
277 3a-c) shows two minima which correspond to the deep OMZs in each hemisphere. The
278 northern minimum, as shown in the CARS climatology, is located at 11.5°N at 450 m

279 depth with oxygen concentrations of $54.1 \mu\text{mol L}^{-1}$. The southern one is located at ~~6.5~~7.0°S
 280 at 400 m depth with oxygen concentrations of $89.7 \mu\text{mol L}^{-1}$. Because of the geometry
 281 of the basin, 23°W is further away from the ETSA than from the ETNA. The climato-
 282 logical structure is better reproduced by TATL3, which shows the two minima at approx-
 283 imately the same latitudes and depths as in the CARS climatology, with slightly larger
 284 minimum oxygen concentrations, $66.1 \mu\text{mol L}^{-1}$ in the ETNA extension and $98.2 \mu\text{mol L}^{-1}$
 285 in the ETSA extension. The OMZ is located within the limits of Central ($\sigma_\theta = 25.8-$
 286 -27.1 kg m^{-3}) and Antarctic Intermediate Waters ($\sigma_\theta = 27.1--32.15 \text{ kg m}^{-3}$), con-
 287 sistent with observations (Karstensen et al., 2008). Oxygen maxima are associated with
 288 the latitudinal positions of the intermediate eastward jets which seem to be responsible
 289 for the large-scale, meridional distribution of oxygen with a peak at around 2°S. Tay-
 290 lor and target diagrams (Figs. A4) show that oxygen concentrations in TATL3 have the
 291 highest correlation (0.94) with the CARS climatology and the smaller bias and root mean
 292 square deviation.

293 The vertical structure of the zonal currents in the upper 300 m (Fig. 3 d,e) is well
 294 reproduced by both simulations when compared to previous observational work at the
 295 same longitude (Brandt et al., 2010, 2015; Burmeister et al., 2019), with the Equatorial
 296 Undercurrent (EUC) core at the Equator, the North Equatorial Countercurrent (NECC)
 297 from roughly 4°N to 10 °N mainly in the upper 100 m with a narrow subsurface expres-
 298 sion at 5°N, which corresponds to the North Equatorial Undercurrent (NEUC) whose
 299 stronger counterpart, the South Equatorial Undercurrent (SEUC), is seen at 4°S. The
 300 main difference between the two runs in the upper 300 m is a stronger branch of the NECC
 301 at 6°N in TATL3 than in TATL10. This branch, as well as a deeper eastward current
 302 at this latitude, is consistent with the observations of Burmeister et al. (2019) from Febru-
 303 ary to March 2018, although it is not as pronounced in the zonal average obtained by
 304 Brandt et al. (2010, 2015).

305 Below 300 m, more intense intermediate zonal jets in TATL3 than in TATL10 have
 306 important consequences for the supply of oxygen to the core of the OMZs, roughly be-
 307 tween 350 and 800 m depth. Both runs show alternating jets whose wavelength is approx-
 308 imately 2.5° in latitude, consistent with observational and modeling estimates that show
 309 alternating zonal jets at mid-depth (Maximenko et al., 2005; Eden, 2006; Richards et al.,
 310 2006; Nakano & Hasumi, 2005). They reproduce the Northern and Southern Interme-
 311 diate Countercurrents (NICC and SICC) at approximately 2°N and 2°S, respectively. The

312 more robust jets in TATL3 suggest that eddy-driven processes, not fully resolved at 10
 313 km horizontal resolution, are important in their formation and maintenance, as suggested
 314 by previous works (Lévy et al., 2010; Eden, 2006; Nakano & Hasumi, 2005).

315 Both TATL10 and TATL3 show two oxygen minima equidistant from the Equa-
 316 tor at approximately 2.5°S and 2.5°N at around 250 m depth. These are shallower than
 317 the deep OMZs in each hemisphere and while they are not very evident in the CARS cli-
 318 matology, these features are consistent with the observations of Brandt et al. (2010, 2012)
 319 which show that these minima are due to the westward-flowing northern and equatorial
 320 branches of the South Equatorial Current. These are stronger in TATL10, which explains
 321 the lower oxygen concentrations (around 60 $\mu\text{mol L}^{-1}$, as opposed to 83 $\mu\text{mol L}^{-1}$ in TATL3)
 322 at approximately the same locations, as these westward currents advect low-oxygen wa-
 323 ters from the OMZs, as discussed previously (see also Figs. 2).

324 A zonal transect of the zonal velocity at 6.4°N shows larger values from 450 m down
 325 to 1000 m in TATL3 than in TATL10 (Fig. A5), with a corresponding eastward retrac-
 326 tion of the 80 $\mu\text{mol L}^{-1}$ oxygen concentration isoline. A similar pattern is seen at other
 327 latitudes (not shown).

328 The ETNA is also characterized by a shallow OMZ just below the mixed layer, be-
 329 tween 60 m and 200 m deep, which arises as a combination of remote, physical drivers,
 330 namely, transport of SACW by the boundary current system and local organic matter
 331 remineralization due to biological productivity in the area (Thomsen et al., 2019; Karstensen
 332 et al., 2008). Fig. 3 shows that the shallow ETNA OMZ is captured in both runs. Oxy-
 333 gen concentrations in the ETNA OMZ at 23°W in TATL10 are around 58.3 $\mu\text{mol L}^{-1}$
 334 reaching from 130 m depth at 13.3°N down to 510 m at 8.8°N. In TATL3 the shallow ETNA
 335 OMZ at the same longitude is located roughly at 150 m depth and the deeper, more pro-
 336 nounced minimum at around 10.5°N, at 468 m depth with oxygen concentrtaions of 62.8
 337 $\mu\text{mol L}^{-1}$. The zonal oxygen distribution at 11°N (Fig. A6), the latitude at which the
 338 ETNA OMZ is more pronounced as seen in Fig. 3, shows that the hypoxic layer reaches
 339 36°W with minimum concentrations of approximately 50 $\mu\text{mol L}^{-1}$ along the σ_θ isoline
 340 of 27.1 kg m^{-3} , consistent with observational estimates (Karstensen et al., 2008). The
 341 limitations of TATL10 in reproducing a realistic OMZ structure are clear in the zonal
 342 transect as the structure of the deeper OMZ is shallower and more confined to the coast
 343 than in the climatology. This results in a lower correlation (0.82 vs. 0.93 in TATL3) and

344 a larger normalized bias for oxygen concentrations, as seen in the Taylor and target di-
345 agrams in Fig. A4.

346 In the South Atlantic, a meridional transect at 5°E, within the Angola Gyre (Gordon
347 & Bosley, 1991) is shown in Fig. 4. This transect is characterized by a wider OMZ south
348 of 12°S, from 100 m to 400 m, and a deeper OMZ north of 12°S. These features are seen
349 in observations from a cruise in 1995, as shown by Karstensen et al. (2008) (their Fig.
350 3). Only TATL3 is able to reproduce the recirculation of the Angola Gyre, which is re-
351 sponsible for the OMZ structure south of 12°S. The correlation between the modeled oxy-
352 gen and the CARS climatology for this section is 0.96 in TATL3 against 0.85 in TATL10
353 (Fig. A4).

354 A zonal transect at the core of the ETSA OMZ at 9 °S from CARS shows the deep
355 OMZ at around 400 m depth with oxygen concentrations as low as 32.5 $\mu\text{mol L}^{-1}$ be-
356 tween 5°E and 10°E, just above the 27.05 kg m^{-3} isopycnal (Fig. 5). The climatology
357 does not show a clear shallow OMZ near the coast as in the ETNA, which is likely due
358 to the lack of observations. The deep OMZ is located between the 26.6 and 27.05 kg m^{-3}
359 isopycnals. TATL3 shows the deep OMZ at the correct depth, with a broader core as
360 compared to the climatology, and lower oxygen concentrations (19.1 $\mu\text{mol L}^{-1}$ at around
361 350 m depth) than in the climatology. The zonal extent of the OMZ is also more sim-
362 ilar to the climatology in TATL3 than in TATL10, as confirmed by the Taylor and tar-
363 get diagrams in Fig. A4.

364 Fig. 5 suggests that the source of low oxygen waters is the shelf region near the EBUS
365 which is then diffused by coastal and shelf currents/undercurrents as well as the deeper
366 circulation. Because these are not adequately resolved in TATL10, the low oxygen wa-
367 ters tend to be more confined near the source region. TATL3, on the other hand, seems
368 to be more effective in simulating the mechanisms responsible for the diffusion of low oxy-
369 gen waters, therefore providing a more realistic representation of the OMZ structure.

370 The relatively low horizontal resolution of TATL10 may not be sufficient to cor-
371 rectly simulate coastal and shelf processes that influence oxygen distribution, such as the
372 dynamics of the poleward undercurrent and the formation of coastal eddies. Also, the
373 low resolution of the wind stress forcing, which was deliberately used both in TATL10
374 and TATL3, may not generate the wind stress curl near the coast which is an important
375 local vorticity source. Nevertheless, TATL3 is overall consistently better than TATL10

376 thus highlighting the importance of high horizontal resolution on a basin-scale simula-
377 tion for a correct representation of the OMZ both in ETSA and ETNA.

378 **3.3 Impact of the jets on the OMZ Structure**

379 From Figs 1 and 2, we see that westward velocities in the equatorial region tend
380 to advect low-oxygen waters from the eastern boundary. These currents are similarly re-
381 produced in both runs (Fig. 3), but seem to have a stronger influence on the structure
382 of the OMZs in TATL10. This is clearly seen when we compare the average hypoxic layer
383 of each model run, defined as the thickness of the water column that comprises oxygen
384 concentrations smaller than $80 \mu\text{mol L}^{-1}$, to the CARS climatology (Fig. 6). The un-
385 realistic westward extend of the hypoxic layer in TATL10 is primarily driven by the west-
386 ward currents that flank the eastward-flowing EUC. In TATL3, the effect of these cur-
387 rents is counteracted by increased eastward velocities. TATL3 shows a much better agree-
388 ment with the CARS climatology than TATL10 in terms of the structure and thickness
389 of the Atlantic hypoxic layer. Despite the clear differences in the structure of the OMZs
390 in the different runs, their averaged volume is surprisingly similar to the one from the
391 climatology, as seen in Table 1. These values are consistent with the estimates of Karstensen
392 et al. (2008), who used a threshold of $90 \mu\text{mol L}^{-1}$ for their calculation. The thinning
393 of the OMZs in TATL10 is compensated by a larger areal extent induced by the west-
394 ward zonal currents. Therefore, metrics such as the volume of the OMZs should not be
395 presented without an examination of the horizontal and vertical structure as seemingly
396 correct numbers may be misleading.

397 The asymmetry between the effect of westward and eastward jets in the different
398 runs occurs because the former are associated with more homogenous potential vortic-
399 ity (PV) tongues while the latter are associated with more intense PV gradients (Delpech
400 et al., 2020; Assene et al., 2020).

401 Fig. 7 shows that the westward jets, associated with homogenized PV, are more
402 pronounced in layer 1, which is equally reproduced in both runs, as shown by the zonal
403 transport stream function in layer 1, meridionally integrated at 15°W (Fig. 8a). Homoge-
404 nous PV waters, associated with the shallow OMZ in the eastern part of the basin, are
405 advected westwards. Eastward jets, associated with higher oxygen concentrations and
406 more intense PV gradients, are stronger in TATL3 in both layers, thus effectively coun-

407 teracting the westward advection of homogenized PV from the east. The inability of TATL10
408 in simulating the intermediate jets causes the horizontal stretching of the shallow OMZ
409 and distorts its overall structure as shown in Fig. 6.

410 The larger westward volume transport in layer 2, which contains the core of the
411 OMZs, in TATL3 is confirmed by the calculation of the meridionally integrated zonal
412 transport stream function (Fig. 8b). The meridional integration started at around 9°S,
413 where both velocities were near zero. Eastward zonal jet transport peaks are seen at around
414 4°N, 7°N and 9°N, which is approximately the latitude of the intermediate jets in the layer-
415 averaged zonal velocity, zonal transport and also in Fig. 3.

416 3.4 Oxygen Budget

417 Monthly averages of the oxygen budget terms were saved for model years 13 to 16
418 (see Section 2.1). The 4-year averages of each term at 26 °W, western limb of the ETNA,
419 shows that physical and biological processes largely balance each other out in the up-
420 per 200 m (Fig. 9). This is true both for TATL3 and TATL10, with specific differences
421 that are related to the impact of horizontal resolution on primary production and up-
422 per ocean dynamics which are out of the scope of this study. Between 200 m and 800
423 m, the depth interval relevant to the deep OMZ, the tendency is slightly more positive
424 or neutral in TATL3 while for TATL10 it is mostly negative. Physical processes dom-
425 inate the oxygen budget at this depth interval, with the advection terms being the most
426 important. This is a consequence of the stronger intermediate eastward jets in TATL3.
427 Vertical mixing is important for oxygen supply below the mixed layer, which makes it
428 relevant for the oxygen balance in the shallow OMZ on short time-scales and for the deep
429 OMZ in longer time-scales as it diffuses below the thermocline.

430 At 10°W, roughly the westward limb of the ETSA OMZ, the oxygen budget be-
431 low 200 m is still dominated by advection as seen in Fig. 10. It is largely positive in TATL3,
432 particularly between 5°S and 12 °S, and largely negative in TATL10. The increased oxy-
433 gen supply at depth on the western edge of the ETSA OMZ in TATL3 results in a more
434 realistic OMZ structure as seen in Figs. 1 and 6.

435 A box average of the budget terms within the deep ETNA OMZ (9°N – 14°N ; 23°W
436 – 18°W) illustrates the importance of advection in supplying oxygen to the OMZ core
437 region. As seen in Fig. 11a,b, below 380 m the tendency in TATL3 is positive and is mostly

438 driven by advection. In TATL10 the tendency is very close to zero as physical and bi-
 439 ological processes balance each other out from 200 m down to 1000 m and the physical
 440 supply is smaller than in TATL3 below 400 m. Specifically, between 350 m and 570 m,
 441 the OMZ core, the average of the oxygen tendency term in TATL10 is $0.07 \mu\text{mol L}^{-1}\text{Year}^{-1}$
 442 while in TATL3 it is $0.82 \mu\text{mol L}^{-1}\text{Year}^{-1}$ which is due to the increased advective sup-
 443 ply in TATL3, whose average value is $3.91 \mu\text{mol L}^{-1}\text{Year}^{-1}$ against $3.0 \mu\text{mol L}^{-1}\text{Year}^{-1}$
 444 in TATL10. The averaged biological consumption is relatively similar in both runs, -3.09
 445 $\mu\text{mol L}^{-1}\text{Year}^{-1}$ in TATL10 and $-3.28 \mu\text{mol L}^{-1}\text{Year}^{-1}$ in TATL3. These values are com-
 446 parable to the estimate of Karstensen et al. (2008) for a similar region, namely $-4.1 \mu\text{mol kg}^{-1}\text{Year}^{-1}$.
 447 The magnitude of oxygen budget terms in TATL3 is much larger than the domain-averaged
 448 model drift over the simulation, namely $-0.01 \mu\text{mol L}^{-1}\text{Year}^{-1}$. In TATL10, while indi-
 449 vidual terms are much larger than the model drift ($-0.11 \mu\text{mol L}^{-1}\text{Year}^{-1}$), the trend is
 450 of the same order of magnitude, but positive (see Table B1 and Fig. A8). We note that
 451 regional trends at different depth intervals may differ from the globally-averaged trends.
 452 In addition, the vertical profile is consistent with the observational estimates of Brandt
 453 et al. (2015) for a similar region. Admittedly, 4-year averages are a relatively short time
 454 to speculate over long-term trends. Our results, however, highlight the different impact
 455 of the advection terms in the oxygen budget in the two climatological runs.

456 The box-average at the core of the ETSA OMZ ($10.5^{\circ}\text{S} - 7.5^{\circ}\text{S} ; 5^{\circ}\text{W} - 12^{\circ}\text{E}$) shows
 457 that biological consumption causes a deficit in the oxygen budget from 150 m down 400
 458 m in TATL3. As it decreases with depth, advection and mixing are able to balance it
 459 so that the trend is slightly positive between 400 m and 700 m depth, the core of the OMZ
 460 in TATL3. In TATL10 the trend is positive below 580 m, which may be explained by
 461 the fact that the OMZ in the region is actually shallower than in TATL3 and the CARS
 462 climatology, as seen in Fig. 5. Vertical mixing is a small, but positive contribution to
 463 the oxygen supply to the region from roughly 200 m down to 500 m in TATL3, and from
 464 200 m down to 450 m in TATL10, but with a smaller magnitude. While the magnitude
 465 of vertical mixing is relatively small, the positive contributions show that it may be an
 466 important component of the oxygen budget over long time-scales (Lévy et al., 2022).

467 4 Eddy-Driven Oxygen Supply

468 As shown by Hahn et al. (2014), the meridional eddy oxygen flux is also important
 469 in the OMZ. We calculated the eddy-driven meridional oxygen flux at 23°W (Fig. 12)

470 using the 7-year mean and deviations from that mean based on monthly values of the
 471 oxygen concentration and meridional velocities. The calculation was performed along isopyc-
 472 nals and the values were re-gridded onto depth coordinates. Values from both runs are
 473 consistent with those estimated by Hahn et al. (2014) using the diffusive flux parame-
 474 terization (Fig. 8 in their study) but TATL3 provides a more realistic representation of
 475 these fluxes as compared to the observational estimates. Specifically, we note the pos-
 476 itive (northward) oxygen flux between 400 m and 700 m near the ETNA boundary, be-
 477 tween 4°N and 10°N, which drops to zero and becomes negative at the core of the ETNA
 478 OMZ.

479 To have a better idea of the importance of zonal versus meridional eddy velocities
 480 in and out of the OMZs, we calculate the degree of anisotropy α (Huang et al., 2007) (See
 481 Section 2). In the ETNA, there is a noticeable decrease in α from the western part into
 482 the OMZ, which suggests that meridional anomalies become more relevant in the OMZ.
 483 In the ETSA, the west-east decrease in α is not as pronounced as in the ETNA. It de-
 484 creases in TATL3 when compared to TATL10, showing a more isotropic eddy field. Over-
 485 all, the eddy flow at depth tends to be more isotropic than at 15 m, but even after re-
 486 moving the mean currents, where the zonal dominance is overwhelming, there is a slight
 487 anisotropy in the upper ocean favoring zonal currents.

488 We note that our α values are relatively lower when compared to the ones from Huang
 489 et al. (2007), who calculated α for the Pacific ocean. This is probably because they used
 490 monthly means while we use daily means, which contain more variability in both com-
 491 ponents.

492 The increased eddying velocities in TATL3 yield larger eddy diffusivities than in
 493 TATL10 with large values usually associated with the location of the eastward zonal jets
 494 (Fig. 14). Eddy diffusivities are estimated using the formulation of (Eden & Greatbatch,
 495 2008)

$$\kappa = \sqrt{\frac{U_e^3}{2\beta}}, \quad (6)$$

496 where U_e is the r.m.s. eddy velocity and β is the meridional gradient of the Coriolis pa-
 497 rameter. U_e is estimated with the averaged velocities between the 26.6 and 27.4 kgm^{-3}
 498 isopycnals as it is representative of the OMZs core regions. The magnitude of κ values
 499 is consistent with in-situ estimates based on tracer release experiments in the North At-

500 lantic OMZ (Banyte et al., 2013). The increased values within the jets and the OMZs
501 show that the high resolution run is more capable of redistributing tracers horizontally
502 in intermediate layers than the low resolution run. This seems to be an essential for the
503 correct simulation of the vertical structure of the OMZs.

504 **5 How does a global model reproduce the Atlantic zonal jets and OMZs?**

505 In order to contextualize our results with existing global simulations, we use the
506 output of the Global Ocean Physics Reanalysis (<https://doi.org/10.48670/moi-00021>)
507 and the Global Ocean Biogeochemistry Hindcast (<https://doi.org/10.48670/moi-00019>).
508 Both products are available at the Copernicus Marine Center Service (CMEMS) at [https://](https://marine.copernicus.eu)
509 marine.copernicus.eu. The physical model has a nominal horizontal resolution of $1/12^\circ$ and
510 is forced with daily atmospheric fields from ERA-interim. In addition, observations from
511 sea level, *in-situ* vertical profiles of temperature and salinity, sea surface temperature and
512 sea ice are assimilated into the model. The biogeochemical model output is interpolated
513 into a $1/4^\circ$ global grid from an original $1/12^\circ$ horizontal resolution run. No data assim-
514 ilation is performed in the biogeochemical model. We compare our 7-year climatolog-
515 ical simulations with the average of the whole period of the CMEMS product (1993-2019
516 in this case). This is not a coupled physical-biogeochemical run, these two products were
517 chosen because of the availability of the data on the same time range. The physical model
518 used in the biogeochemical hindcast (the non-assimilative FREEGLORYS2V4 run) is
519 not available on the CMEMS website. More details about these simulations may be found
520 on their website.

521 Variables from the CMEMS model are included in the Taylor and target diagrams
522 shown in Figs. A2, A3 and A4. The physical model from CMEMS has a better skill in
523 simulating the climatological temperature and salinity structure than TATL3 and TATL10
524 as compared to CARS, particularly at depth. This is due to more realistic forcing and
525 the data assimilation of physical variables. Nevertheless, some fundamental issues still
526 remain as the relatively low horizontal resolution precludes the formation of robust, in-
527 termediate zonal jets away from the equatorial region, as seen in the average of the zonal
528 velocity at 23°W . The structure of the zonal velocity at 450 m depth shows weaker and
529 broader zonal jets when compared to TATL3, for example (Fig. A7). The biogeochem-
530 ical model overestimates the size of the Atlantic OMZs, in particular the OMZ thickness
531 in the coastal region. In addition, the impact of the westward zonal velocities in advect-

532 ing the OMZ westward persists even in the 27-year average, as seen by the alternating
533 bands of low oxygen. We note that the zonal velocities shown in Fig. A7 are from the
534 assimilative physical model. The biogeochemical model is coupled to a non-assimilative
535 physical model, as previously mentioned. The lower oxygen concentrations in the OMZs
536 than in the CARS climatology may be a consequence of the lack of oxygen supply from
537 the western part of the basin by the zonal jets, inadequate representation of the east-
538 ern boundary current system, insufficient eddy variability at depth to redistribute trac-
539 ers horizontally near the productive coastal regions and/or excessive biological produc-
540 tion and respiration, as seen from the low values near the EBUS.

541 In addition, larger oxygen concentrations on the western part of the basin as com-
542 pared to the CARS climatology or TATL3 suggests that tracers are not redistributed prop-
543 erly at depth, which seems to be an inherent limitation of the low resolution of the model,
544 as seen in TATL10.

545 In summary, the relatively low resolution precludes a realistic representation of al-
546 ternating zonal jets and deeper mesoscale structures which affect tracer redistribution
547 in the coastal and open ocean as well as the supply of oxygen to the OMZs. While data
548 assimilation is extremely useful for a dynamically- consistent monitoring of the state of
549 the ocean and for medium-term forecast, intrinsic limitations due to low model resolu-
550 tion eventually arise, thus preventing a mechanistic understanding of the coupled physical-
551 biogeochemical processes that modulate oxygen concentrations in the ocean.

552 6 Discussion and Conclusions

553 In this study, we used two physical-biogeochemical model configurations of the Trop-
554 ical Atlantic, one at $1/10^\circ$ resolution, roughly the current highest horizontal resolution
555 used in climate models and one at $1/30^\circ$ resolution. The increase in horizontal resolution
556 produces more robust intermediate zonal jets which have a substantial impact on the At-
557 lantic OMZs. The jets are more intensified on the western part of the basin, which is also
558 where higher oxygen concentrations are observed. The eastward extent of the interme-
559 diate zonal jets delimits the western border of the OMZs. Within the OMZs, increased
560 eddy variability at depth and a better representation of coastal and shelf current sys-
561 tems in the high-resolution run makes it more efficient in the horizontal redistribution
562 of low-oxygen waters from the productive eastern boundaries.

563 Increased horizontal resolution and a relatively large domain allow the interplay
564 of eddying and wave processes at multiple scales which lead to the emergence of robust,
565 intermediate zonal jets (Lévy et al., 2010; Kamenkovich et al., 2009). In addition, nu-
566 merical diffusivity is grid-size dependent, the higher the resolution the lower the numer-
567 ical diffusivity. This may facilitate the formation of more robust zonal jets in high-resolution
568 runs. Currently, the complexity of these motions and their impact on tracer transport
569 is not fully captured in subgrid scale parameterizations (Kamenkovich et al., 2019, 2021).
570 The anisotropy, spatial inhomogeneity and non-stationary of eddying motions, partic-
571 ularly at depth, seem to be relevant in this case.

572 The high resolution run TATL3 produced a more robust field of alternating zonal
573 jets down to 1000 m. The stronger eastward jets impact the vertical structure of the OMZs
574 by reducing the magnitude of the surface westward jets which in the low-resolutin run
575 carry low-oxygen waters from the OMZs and distort their vertical structure by making
576 them wider and thinner. Because the mostly wind-forced circulation in the upper 300
577 m is relatively similar in both runs, the westward jets have a strong impact on the OMZ
578 as they are not counteracted by the deeper, eastward jets. The horizontal extent of the
579 OMZs is determined by the longitude at which the eastward zonal jets weaken.

580 The oxygen budget below 200 m is dominated by advection, which is more impor-
581 tant to the core of the ETNA OMZ than the ETSA OMZ. This could be due to the fact
582 that the ETSA OMZ core is further away from the region of influence of the South At-
583 lantic zonal jets because of the geometry of the basin. The implication for this would
584 be that the core of the ETNA OMZ is more immediately sensitive to variations on the
585 transport of the zonal jets. The ETSA OMZ is also known to be influenced by the vari-
586 ability of the Angola-Benguela Frontal Zone, which we do not fully simulate because of
587 computational limitations in terms of domain size.

588 Vertical mixing is an important source of oxygen for the shallow OMZs at approx-
589 imately 100 m depth. Deeper in the water column, the dominant balance is between bi-
590 ological consumption and advection although mixing is shown to have a small, positive
591 contribution in the Atlantic OMZs. Eddy diffusivities are larger in TATL3 with larger
592 values associated with the locations of the zonal jets. While the OMZs have lower eddy
593 diffusivities, values also increase in TATL3 with similar contributions of both zonal and
594 meridional eddy velocities.

595 While the comparison between TATL3 and TATL10 is mostly based on the differ-
596 ent horizontal resolution, it is important to note that in sigma-coordinate (i.e. bathymetry-
597 following) models such as CROCO, an increase in the horizontal resolution of the model
598 grid also increases the bathymetry resolution which can impact vertical mixing rates par-
599 ticularly on the continental shelf/shelfbreak and regions with rough topography with con-
600 sequences for tracer distribution.

601 The repeating, climatological wind and heat fluxes used to force the model runs
602 in this study may cause the formation of the zonal jets at approximately the same lat-
603 itudes so that they are more easily detected on the 7-year average. More realistic forc-
604 ing with interannual variability (e.g. meridional shifts of the ITCZ) could generate merid-
605 ional variations in the zonal jets so that such defined bands may not be so clear.

606 The biogeochemical coupling with the hydrodynamical model at different horizon-
607 tal resolutions also raises questions about how resolution-dependent are biogeochemi-
608 cal parameters as well as about nonlinear feedbacks between oxygen consumption and
609 production by the biology with different rates of supply and mixing by the physics.

610 This study shows that intermediate zonal jets have a significant impact on the sup-
611 ply of oxygen and the overall structure of the Atlantic OMZs and provides a pathway
612 for an accurate representation of the OMZs and their intrinsic properties which are es-
613 sential for the investigation of coastal and open ocean deoxygenation and its sensitiv-
614 ity to natural and anthropogenic forcing as well as for reliable long-term forecasts.

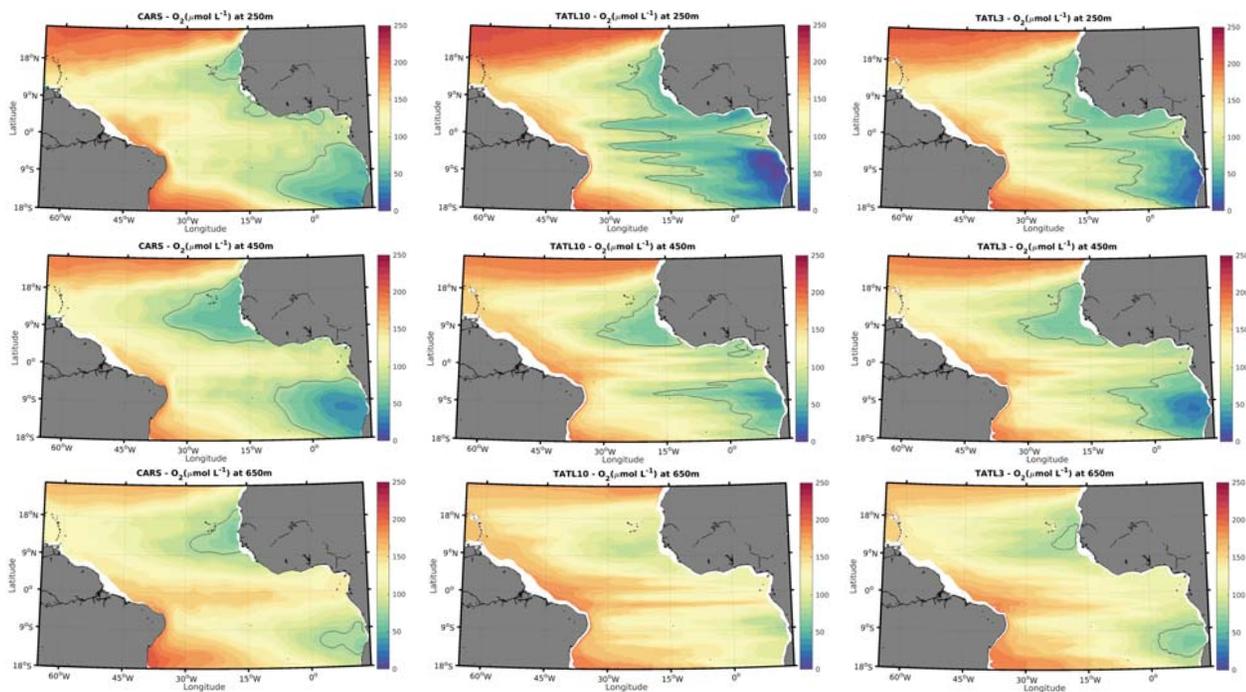


Figure 1. Horizontal maps of oxygen concentrations at 250 m , 450 m and 650 m depths from the CARS climatology (left) and 7-year averages of TATL10 (middle) and TATL3 (right). Black lines show the $80 \mu\text{mol L}^{-1}$ oxygen concentration isoline.

Table 1. Volume of hypoxic waters (m^3)

Source	ETNA	ETSA
CARS	7.82×10^{14}	1.25×10^{15}
TATL10	8.12×10^{14}	1.26×10^{15}
TATL3	8.18×10^{14}	1.25×10^{15}

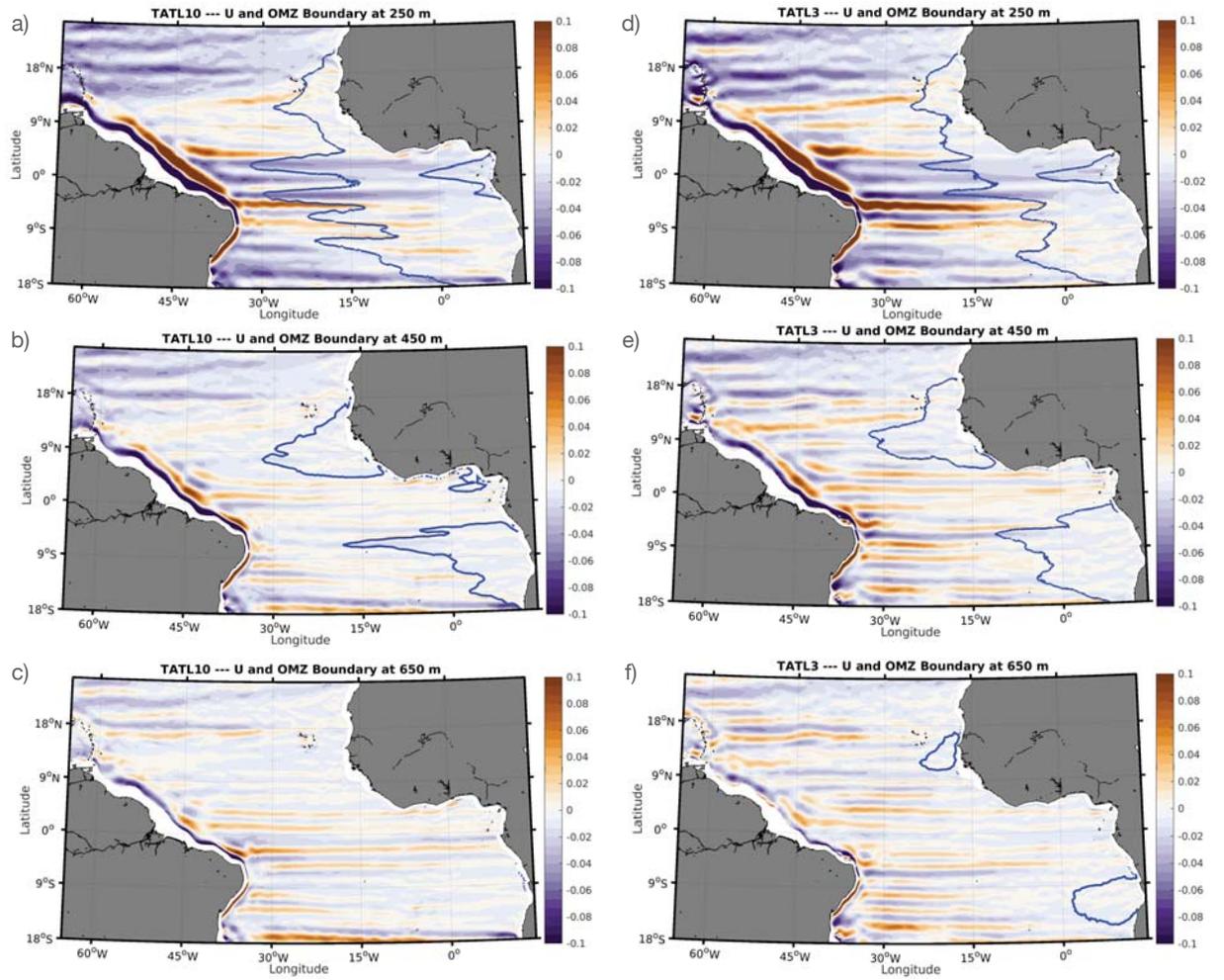


Figure 2. 7-year averages of the zonal velocities from TATL10 (left) and TATL3 (right) at 250 m, 450 m and 650 m depths. Blue lines show the $80 \mu\text{mol L}^{-1}$ oxygen concentration isoline.

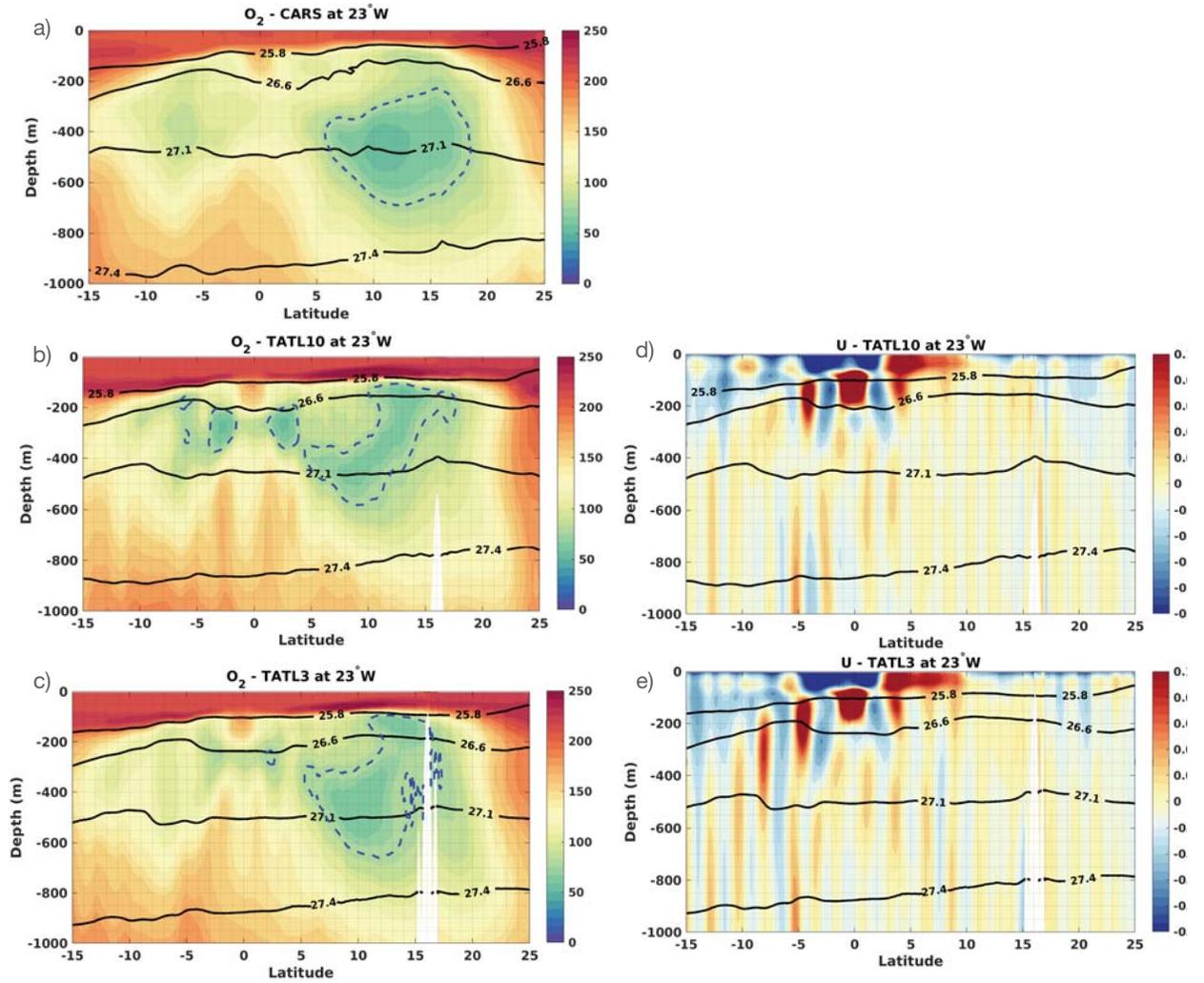


Figure 3. Meridional transects of climatological oxygen concentrations from (a) CARS and 7-year averages from (b) TATL10 and (c) TATL3 at 23°W. Right panel shows the 7-year averaged zonal velocities for (d) TATL10 and (e) TATL3 at the same longitude. Isopycnal surfaces 25.8, 26.6 and 27.1 kg m^{-3} are shown as black lines.

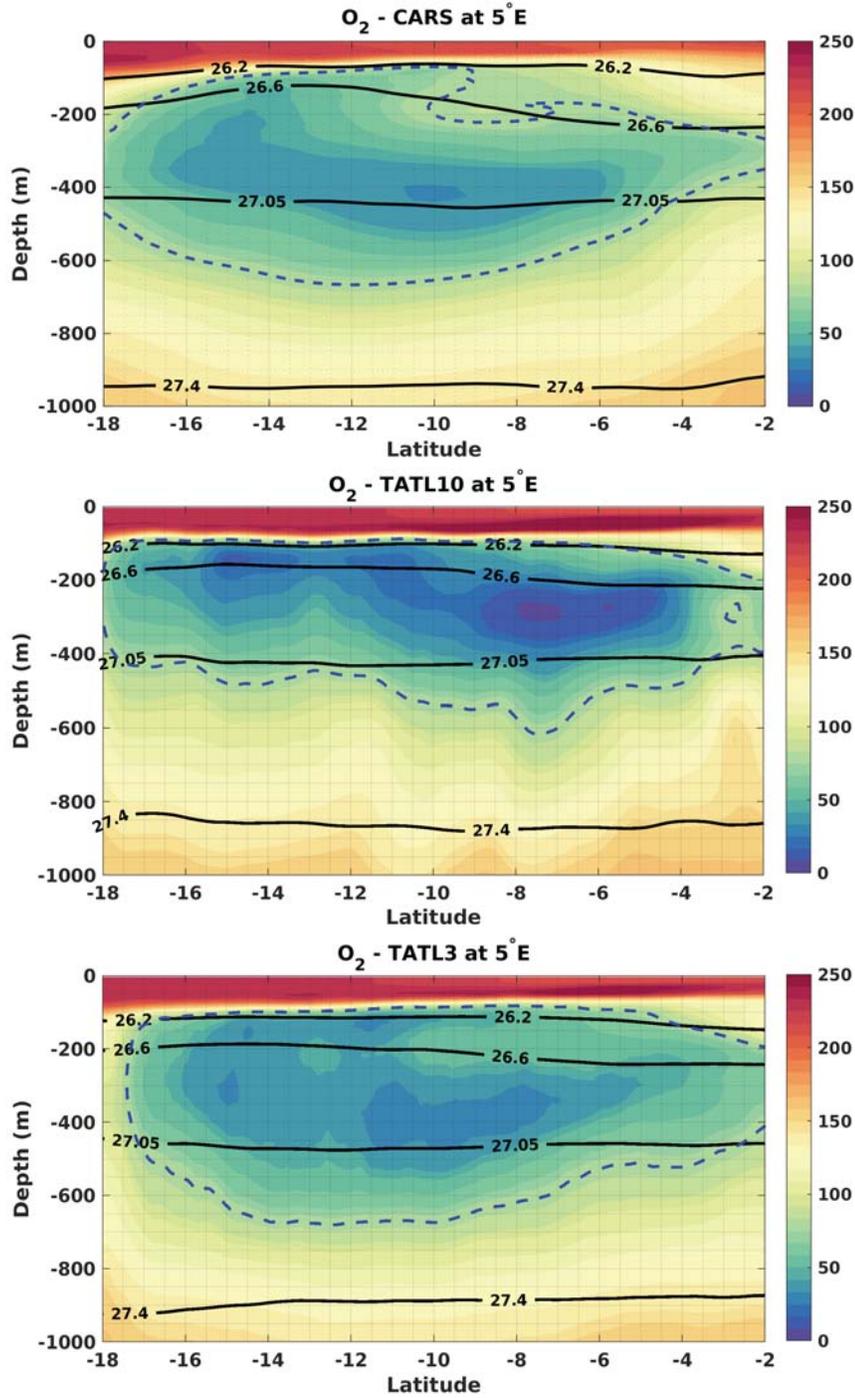


Figure 4. Meridional transect of climatological oxygen concentrations from CARS (top) and 7-year averages from TATL10 (middle) and TATL3 (bottom) at 5°E. Isopycnal surfaces 26.2, 26.6 and 27.05 kg m^{-3} are shown as black lines.

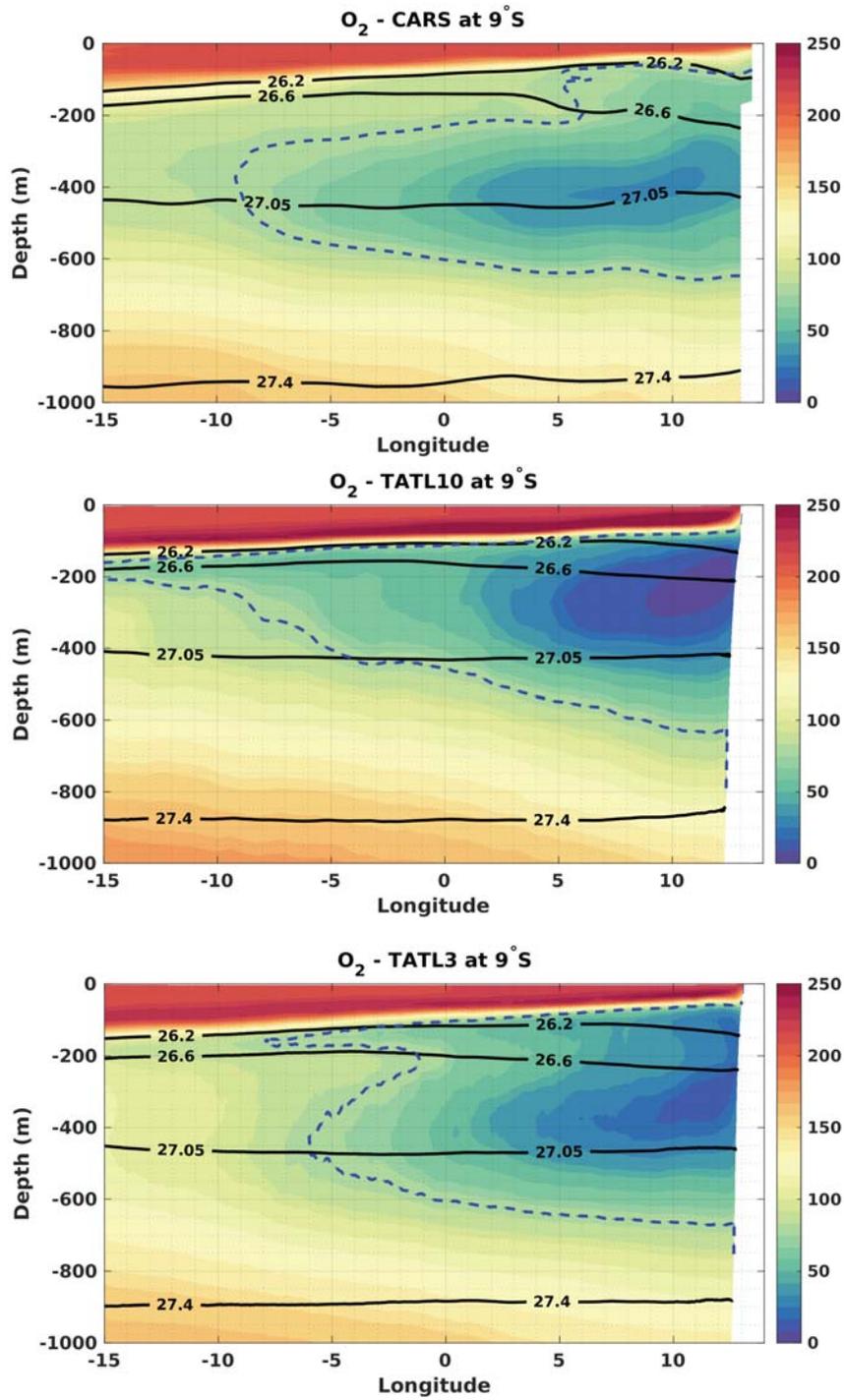


Figure 5. Zonal transect of climatological oxygen concentrations from CARS (top) and 7-year averages from TATL10 (middle) and TATL3 (bottom) at $9^\circ S$. Isopycnal surfaces 26.2 , 26.6 and 27.05 kg m^{-3} are shown as black lines.

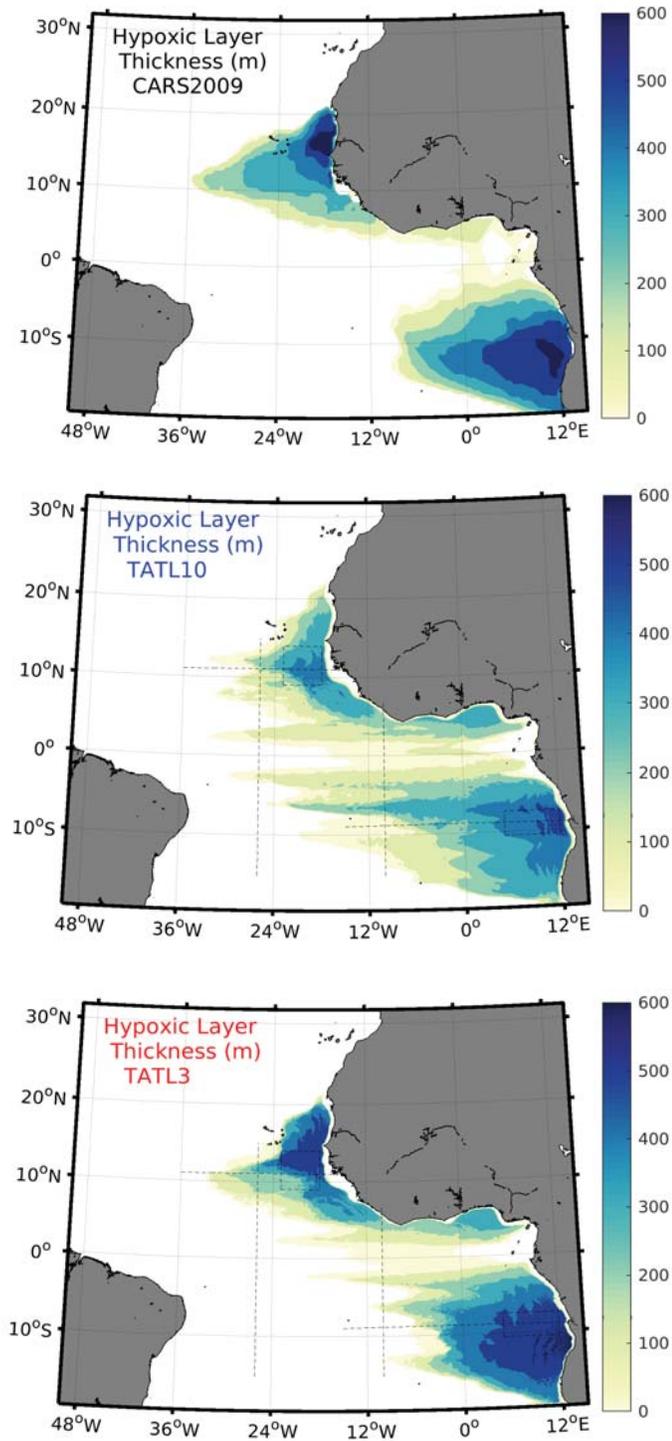


Figure 6. Thickness of the hypoxic layer, here defined as the layer which contains oxygen concentrations below $80 \mu\text{mol L}^{-1}$.

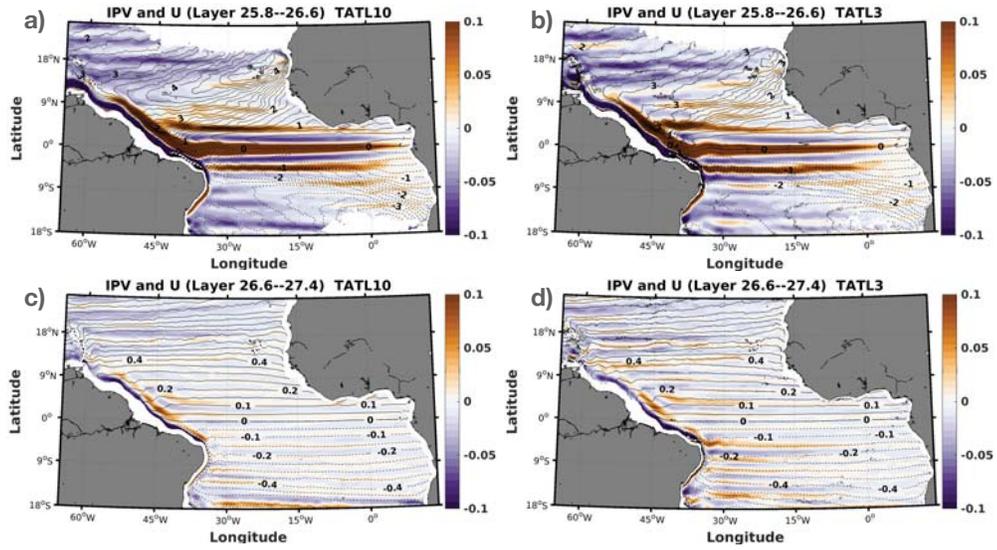


Figure 7. Contours of isopycnic potential vorticity ($\times 10^{-10} \text{ s}^{-1} \text{ m}^{-1}$) calculated from 7-year averages of the model runs for layer 1 (density classes 25.8 and 26.6 kg m^{-3}) and layer 2 (density classes 26.6 and 27.4 kg m^{-3}). Averaged zonal velocities between isopycnal surfaces are shown in color. For layer 1, the PV contour interval is $2 \times 10^{-11} \text{ s}^{-1} \text{ m}^{-1}$, for layer 2 the PV contour interval is $5 \times 10^{-12} \text{ s}^{-1} \text{ m}^{-1}$.

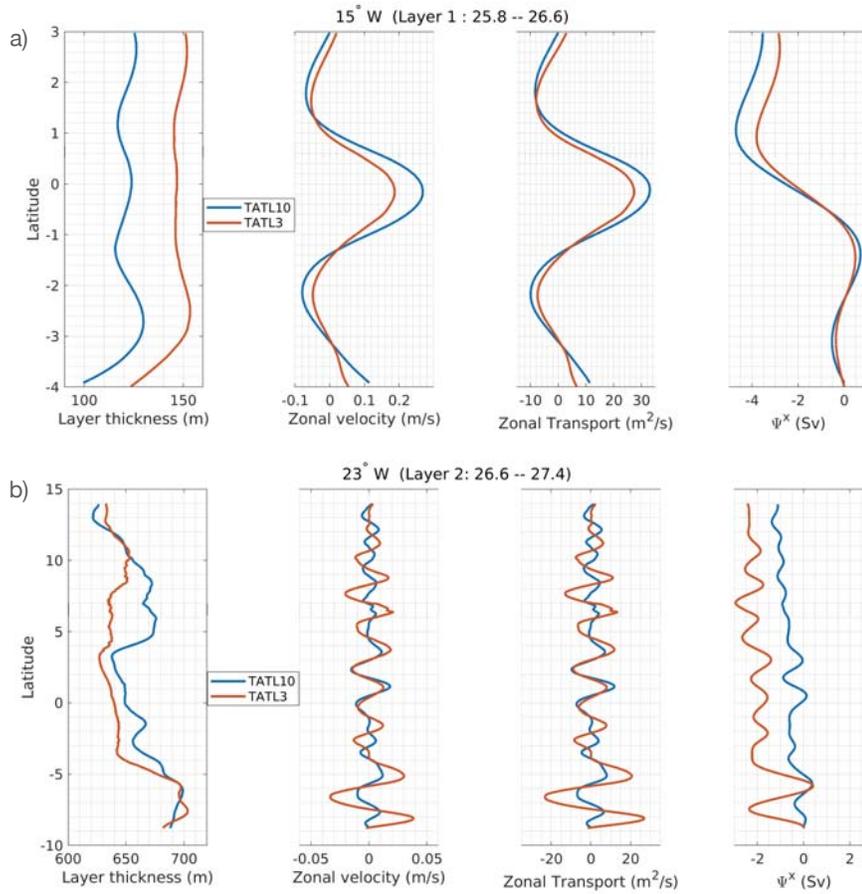


Figure 8. Layer thickness, zonal velocity, zonal transport and the meridionally-integrated zonal transport stream function for (a) layer 1 at 15°W and (b) for layer 2 at 23°W.

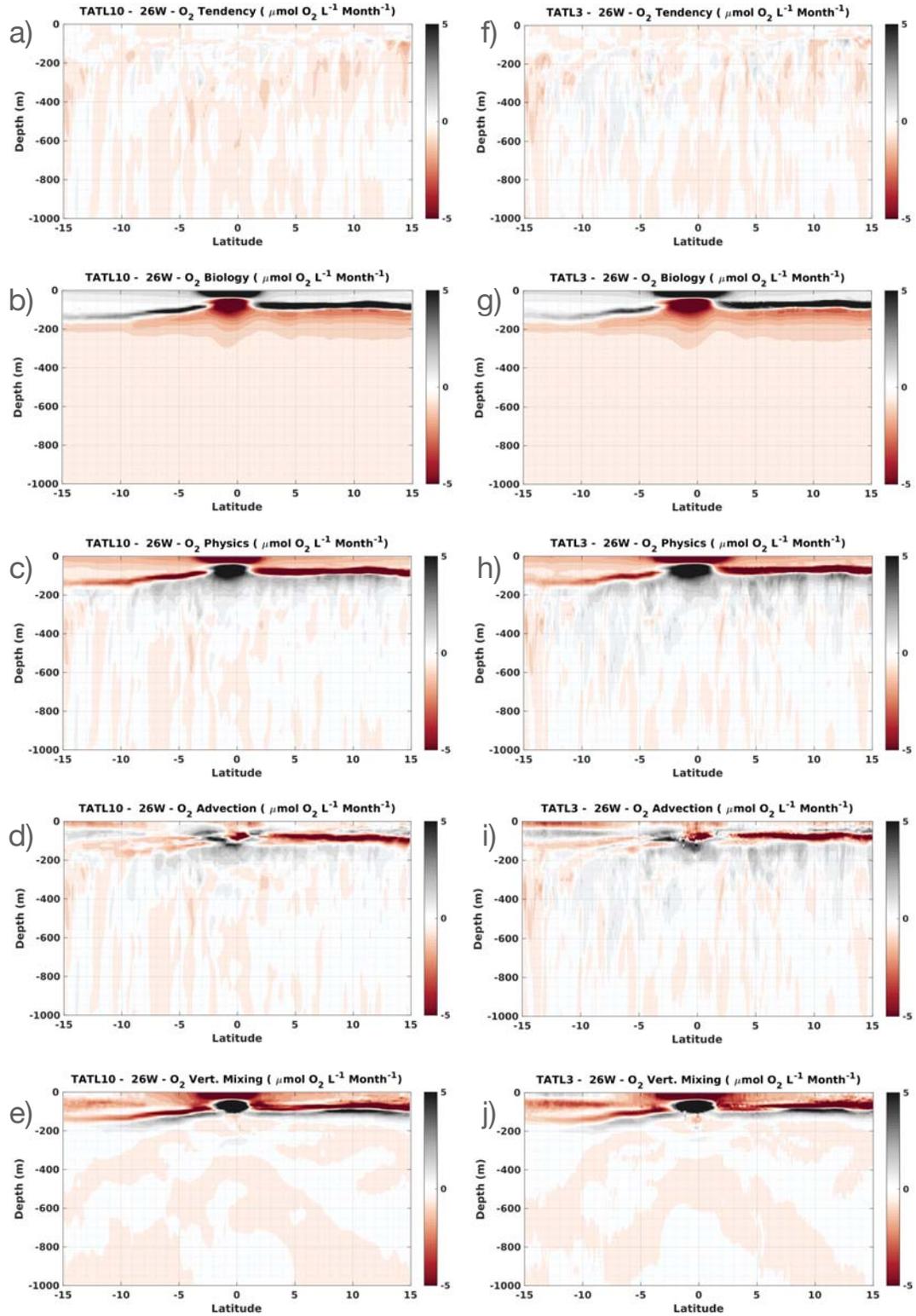


Figure 9. 4-year (Model years 13 to 16) averaged oxygen budget terms at 26°W for TATL10 (left) and TATL3 (right).

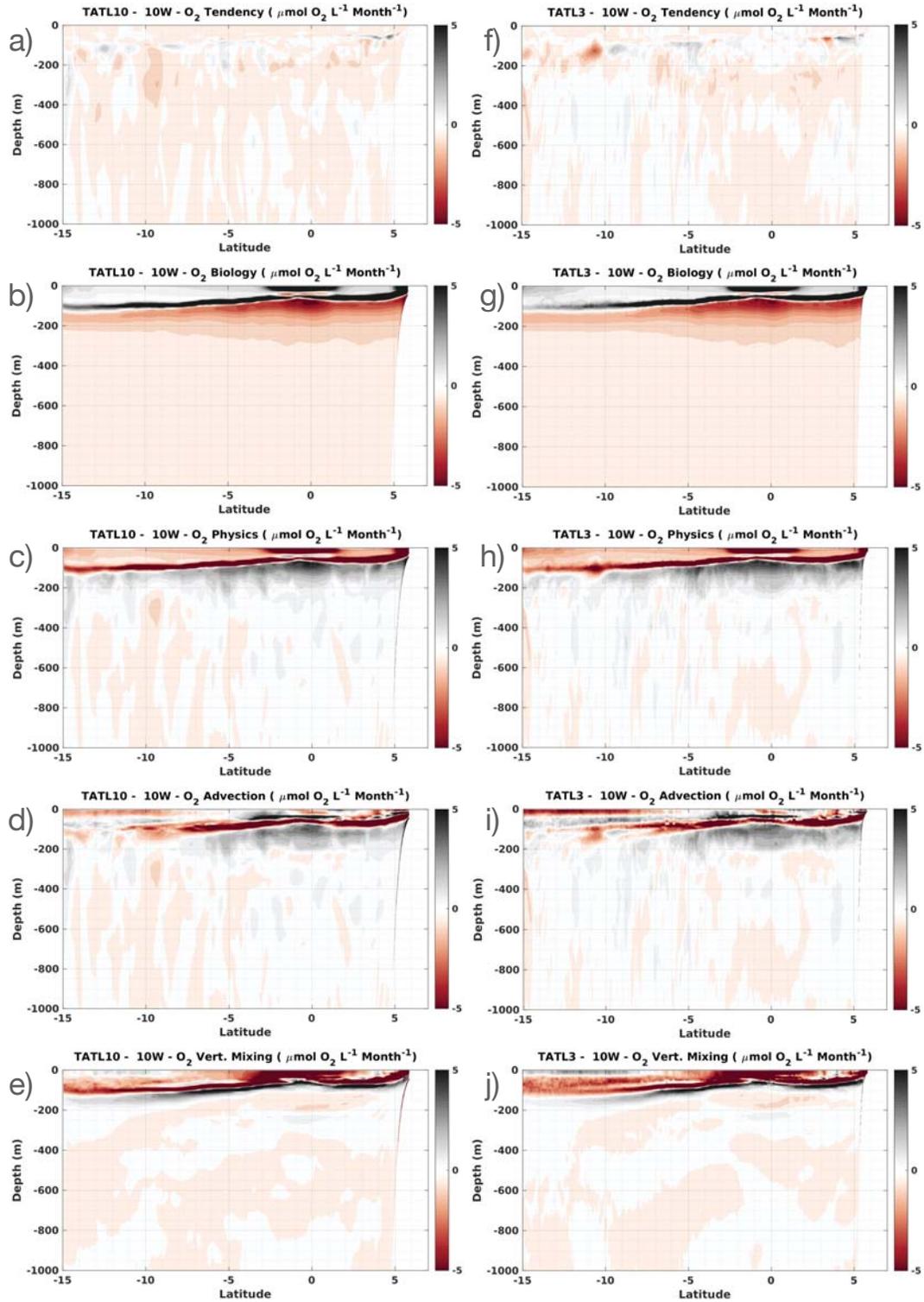


Figure 10. 4-year (Model years 13 to 16) averaged oxygen budget terms at 10 °W for TATL10 (left) and TATL3 (right).

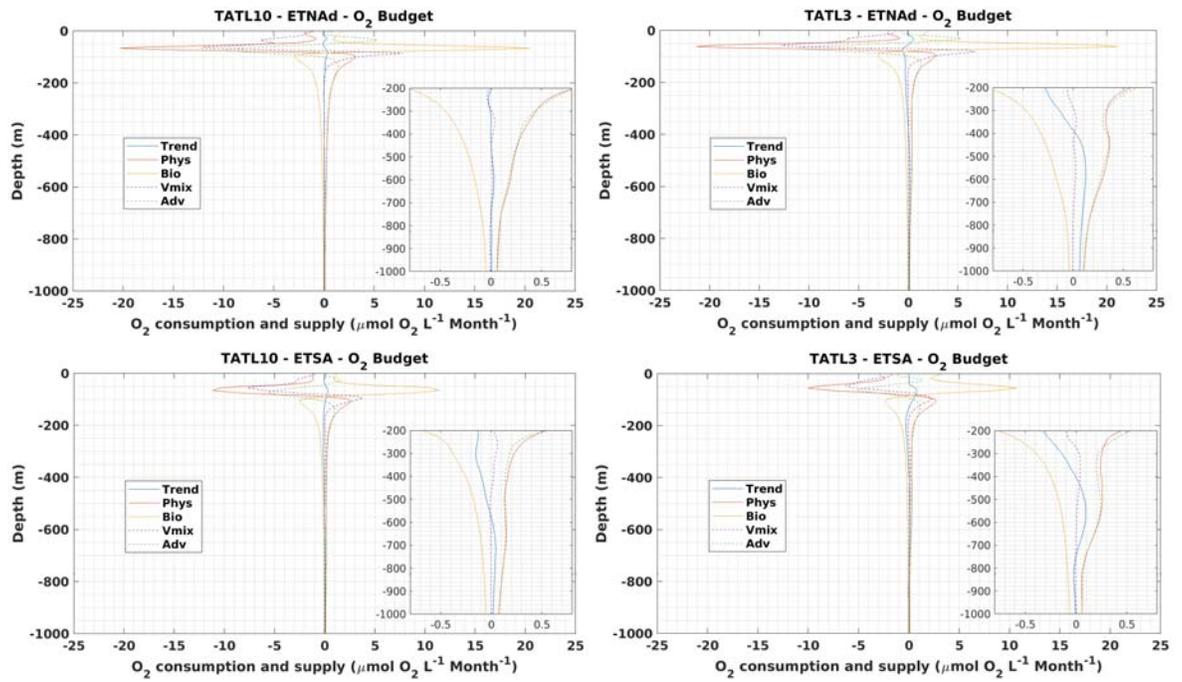


Figure 11. (a,b) Box-averaged oxygen budget terms in the deep ETNA OMZ region (9°N to 14°N and from 23°W to 18°W). (c,d) Box-averaged oxygen budget terms in the ETSA OMZ region ($10.5^{\circ}\text{S} - 7.5^{\circ}\text{S}$; $5^{\circ}\text{W} - 12^{\circ}\text{E}$).

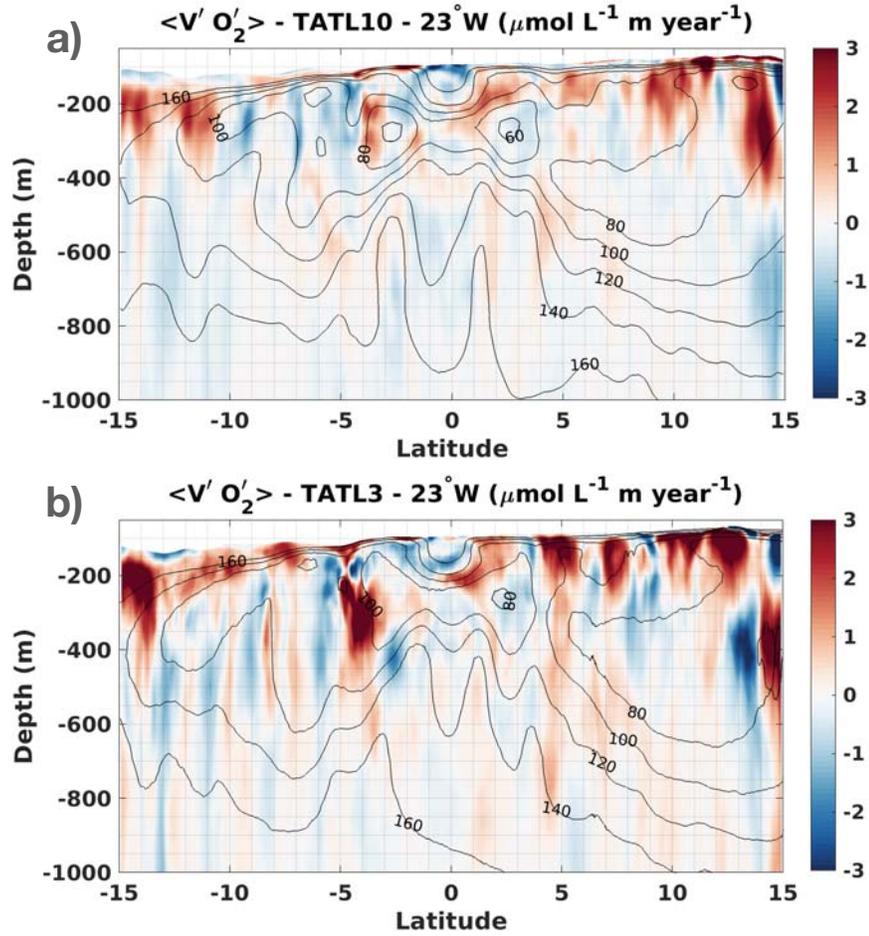


Figure 12. 7-year average of the meridional eddy oxygen flux calculated from monthly averages of meridional velocities and oxygen concentrations. The calculation was performed on isopycnal surfaces and re-mapped onto depth coordinates.

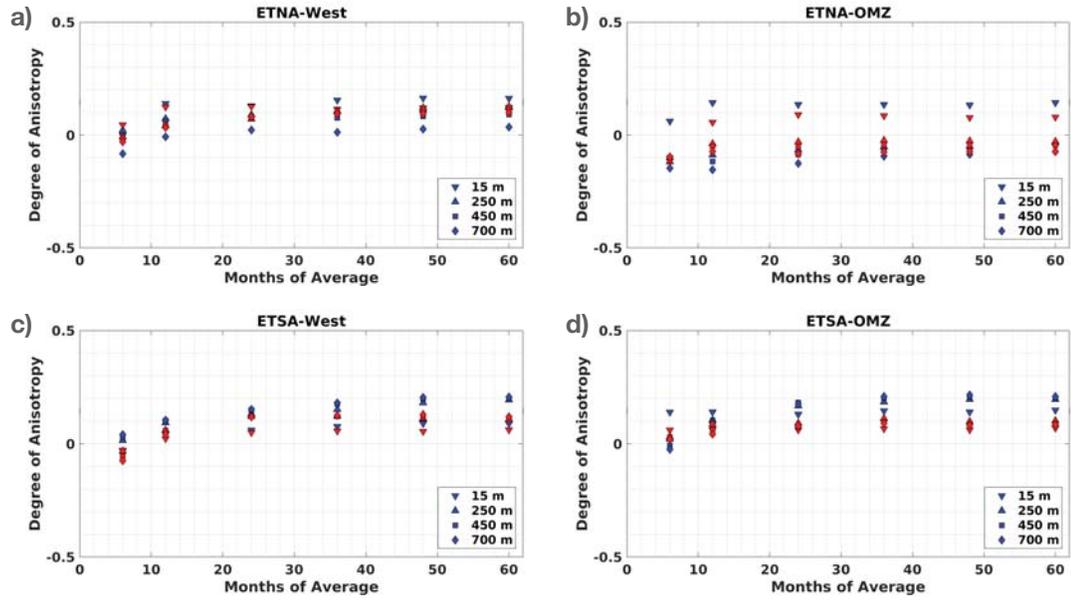


Figure 13. Degree of anisotropy (α) calculated from daily averages of 5 years of model run from TATL10 (blue) and TATL3 (red). α was calculated at different depths and for different time-averaging periods, namely 6 months, 12 months, 24 months, 36 months, 48 months and 60 months.

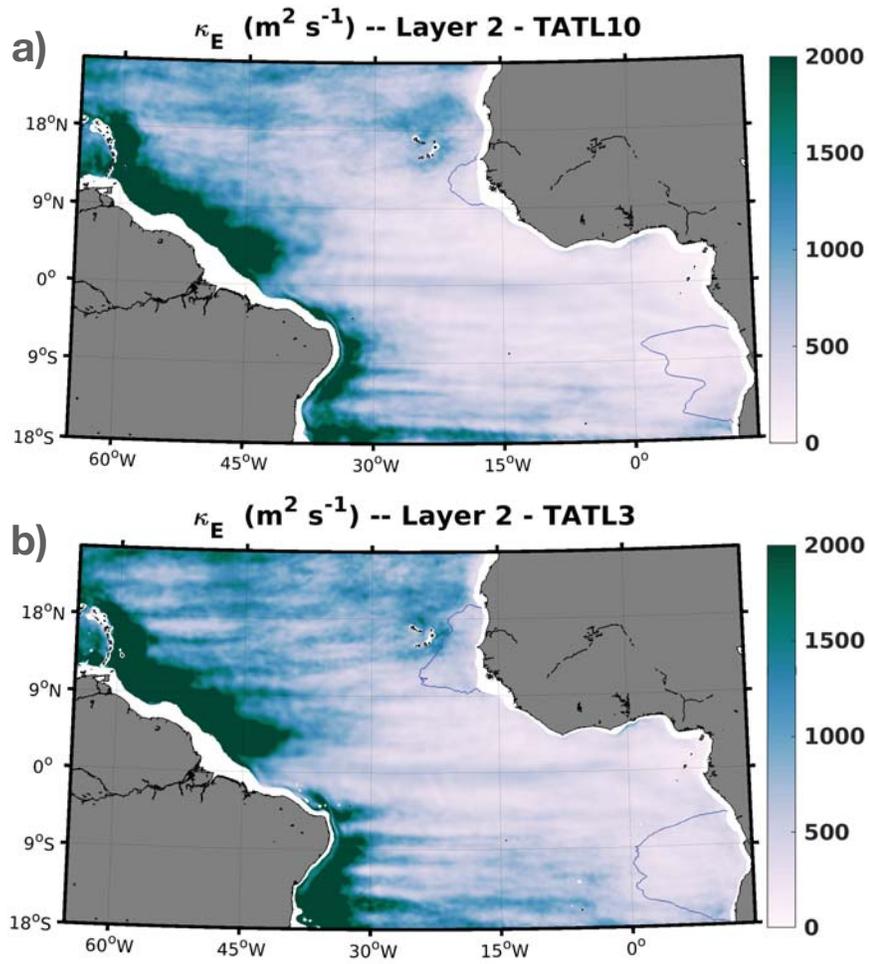


Figure 14. Eddy diffusivity calculated from averaged values in layer 2 for (a) TATL10 and (b) TATL3. The 80 $\mu\text{mol L}^{-1}$ oxygen concentration isoline interpolated at this isopycnal surface is shown as a blue line.

Table B1. Domain-Averaged Model Drift

Variable	TATL10	TATL3
Temperature ($^{\circ}\text{C}/\text{year}$)	5.0×10^{-3}	4.0×10^{-3}
Salinity (PSU/year)	4.4×10^{-4}	4.8×10^{-4}
Oxygen ($\mu\text{mol L}^{-1}/\text{year}$)	-0.11	-0.01
Nitrate ($\mu\text{mol L}^{-1}$)	-5×10^{-3}	2×10^{-2}

615 Appendix A Model Evaluation

616 Normalized Taylor and target diagrams for salinity, temperature, oxygen and ni-
617 trate averaged at 250 m, 450 m and 650 m depths are shown in Figs. A2 and A3. These
618 depths roughly correspond to the upper, core and lower parts of the OMZs, respectively.
619 The regions used for the Taylor and target diagram calculation encompass large portions
620 of the North and South Atlantic oceans, including the OMZs (the squares in Fig. A1a).
621 Overall, there is a general improvement in TATL3 for all variables, most notably for oxy-
622 gen and nitrate. The skill of the simulations is better in the North Atlantic, which may
623 be a consequence of a larger number of observations available in the climatology. In the
624 South Atlantic, these variables may not be as well constrained as in its northern coun-
625 terpart, particularly below 250 m depth. These metrics show that both simulations pro-
626 vide realistic representations of the averaged hydrographic conditions of both physical
627 and biogeochemical variables.

628 Appendix B Model Drift

629 Here we show the evolution of domain-averaged properties in the 7 years of model
630 simulation (Fig. A8). Model drift for each variable are considerably smaller than regional
631 trends as shown in Section 3.4. The domain-averaged model drift for each variable dur-
632 ing the 7-years of model run are shown in Table B1.

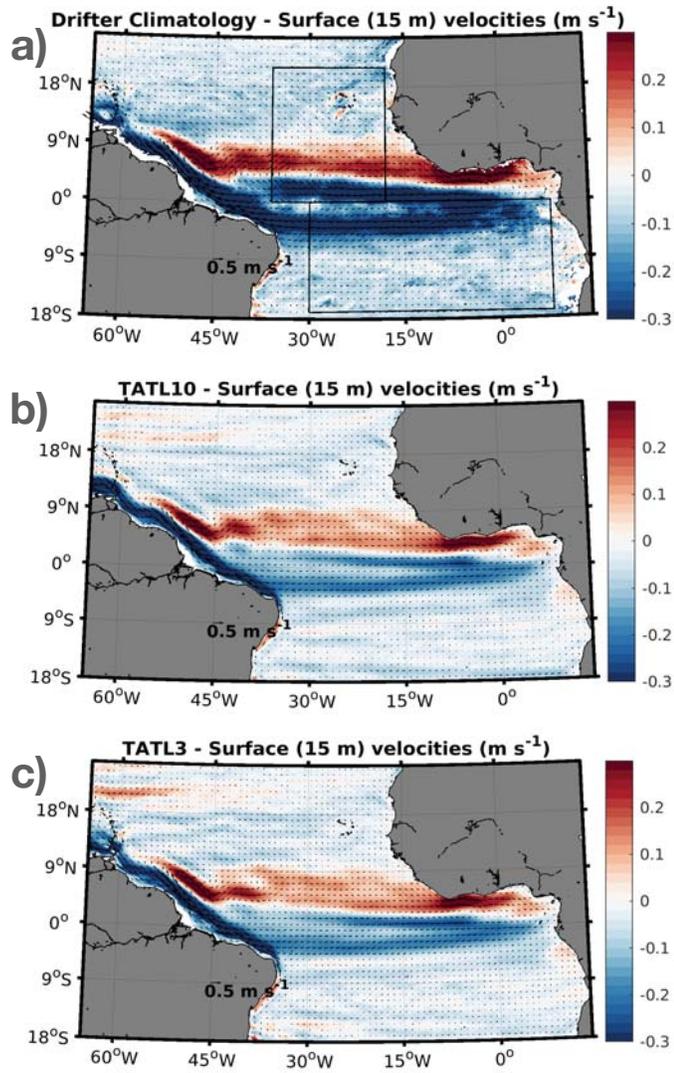


Figure A1. Comparison of surface velocities from TATL10 and TATL3 with a climatology from surface drifters obtained by Laurindo et al. (2017). Squares shown in (a) are the regions in the North and South Atlantic selected for computing Taylor and target diagrams shown in Figs. A2 and A3.

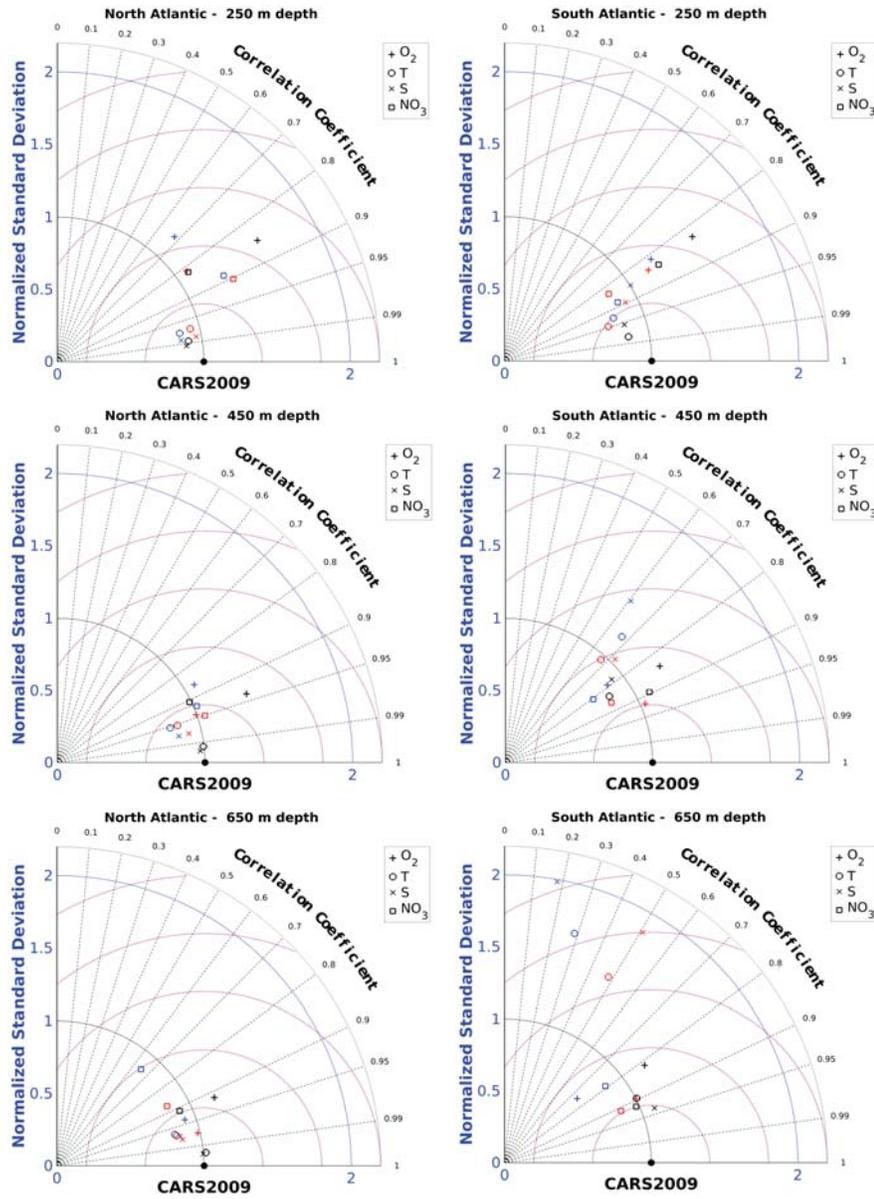


Figure A2. Normalized Taylor diagrams calculated at different depths in the North (left) and South (right) Atlantic subdomains for oxygen, nitrate, temperature and salinity in TATL10 (blue), TATL3 (red) and CMEMS (black). CMEMS models results are discussed in Section 5.

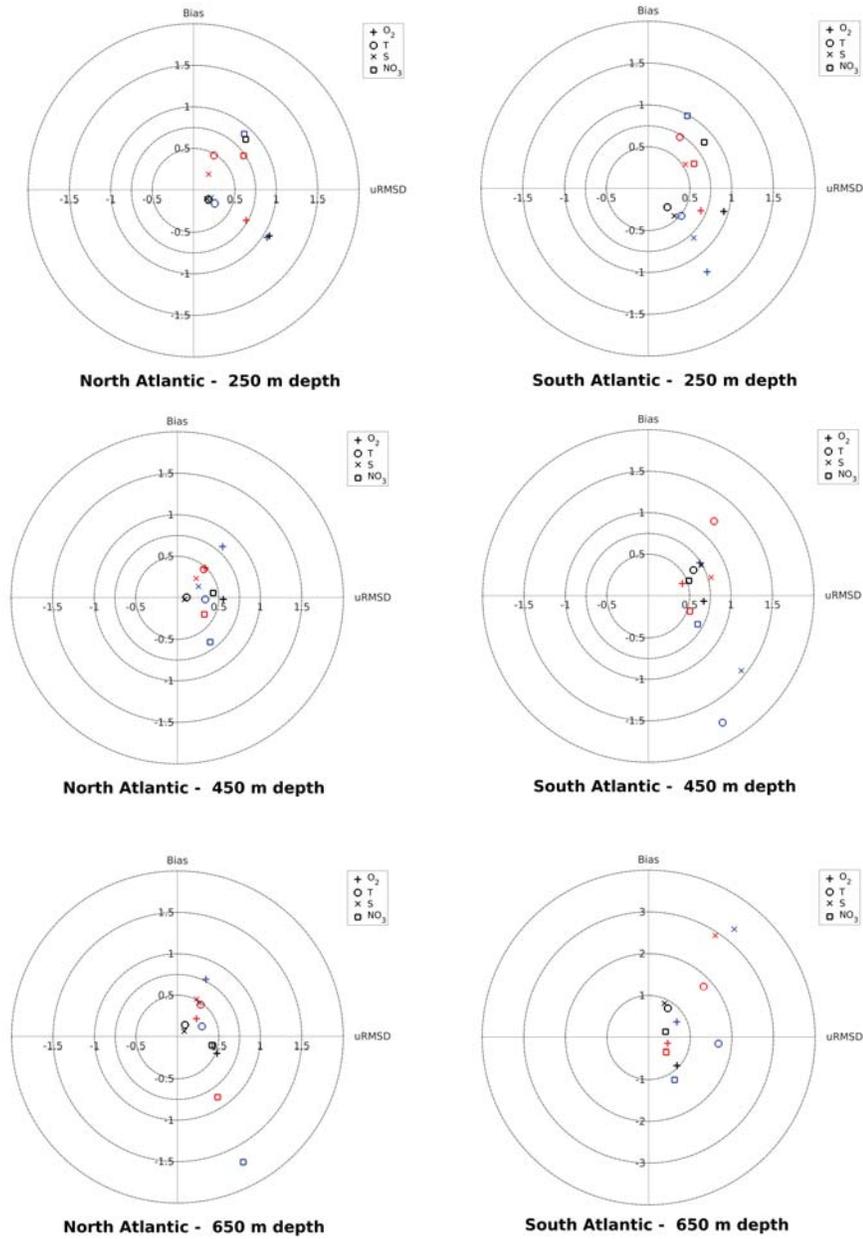


Figure A3. Target diagrams for averages of vertical sections for oxygen, nitrate, temperature and salinity in TATL10 (blue), TATL3 (red) and CMEMS (black). CMEMS models results are discussed in Section 5.

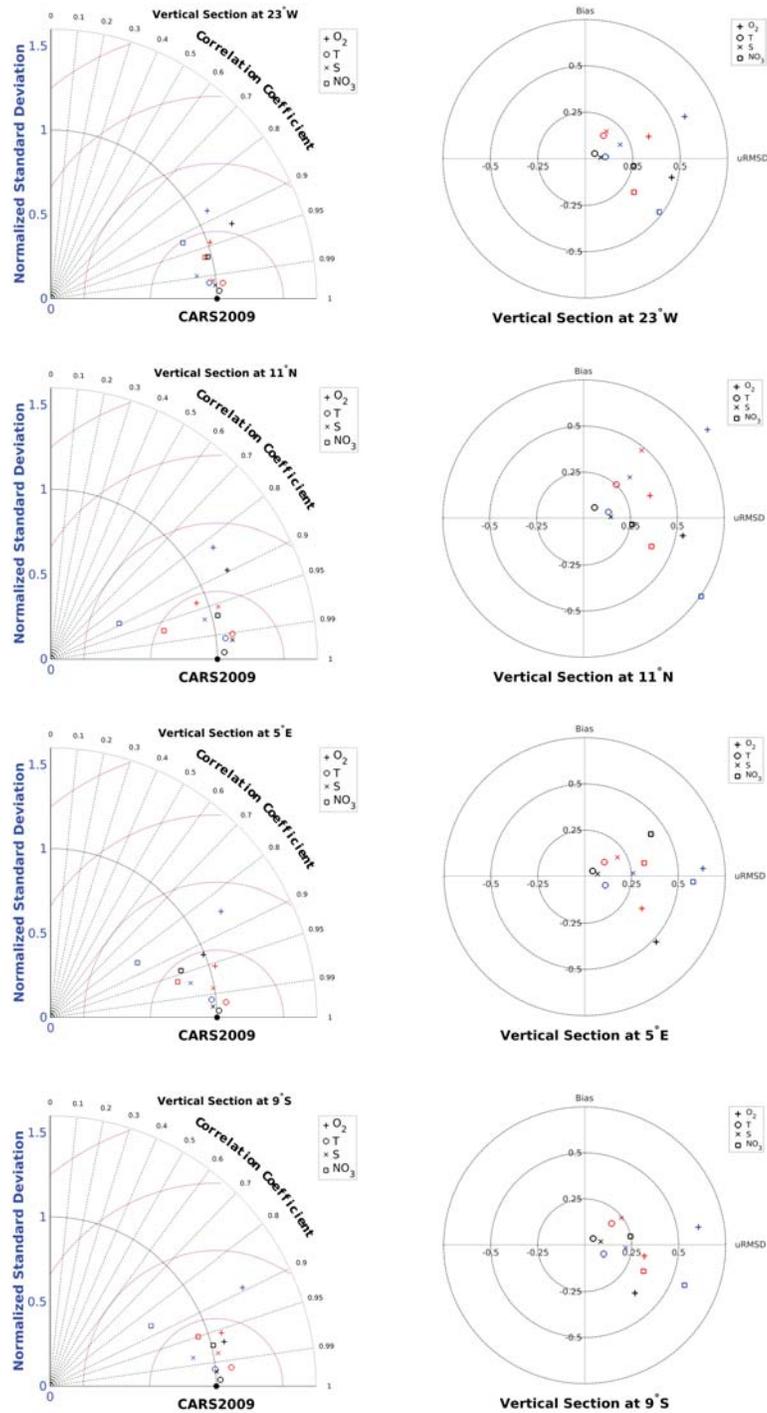


Figure A4. Target and Taylor diagrams for various vertical sections shown in the text. TATL10 (blue), TATL3 (red) and CMEMS (black). CMEMS models results are discussed in Section 5.

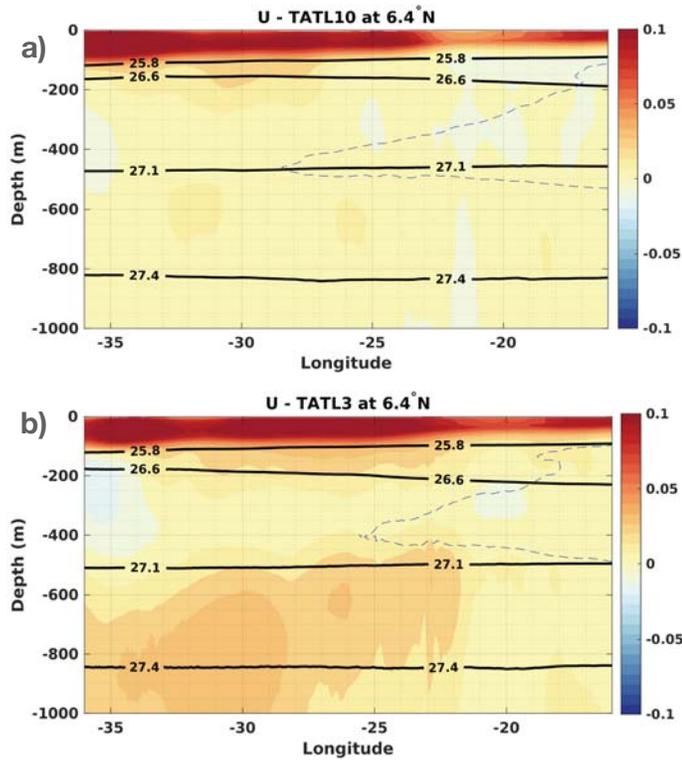


Figure A5. Zonal transect of 7-year averaged zonal velocities from (a) TATL10 and (b) TATL3 at 6.4°N. Isopycnal surfaces 25.8, 26.6, 27.1 and 27.4 kg m⁻³ are shown as black lines.

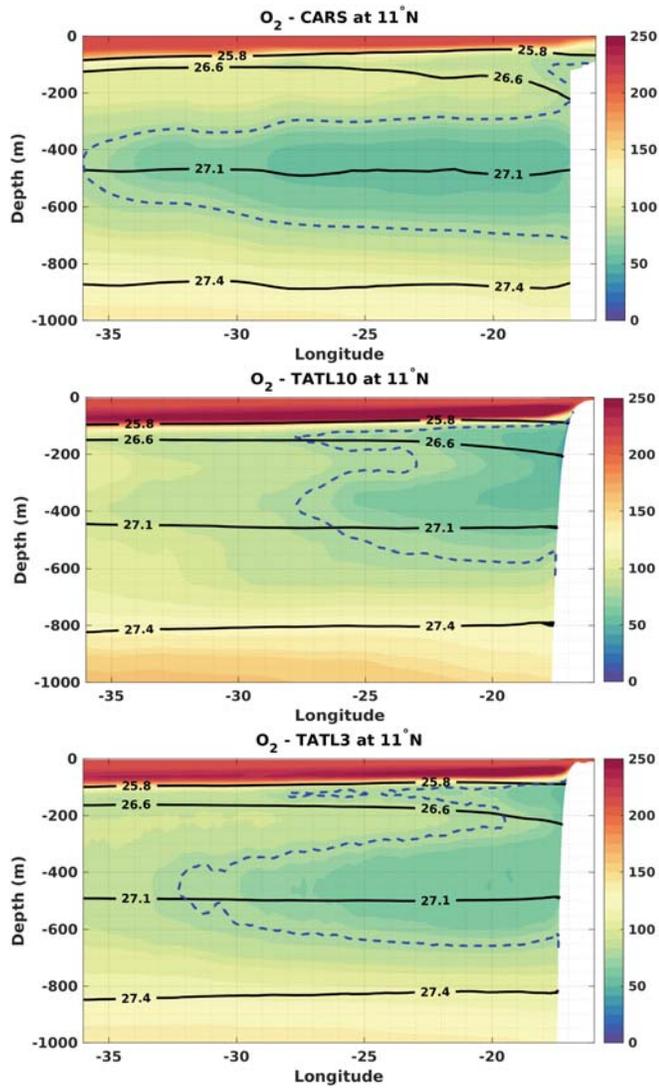


Figure A6. Zonal transect of climatological oxygen concentrations from (a) CARS and 7-year averages from (b) TATL10 and (c) TATL3 at 11°N. Isopycnal surfaces 25.8, 26.6, 27.1 and 27.4 kg m⁻³ are shown as black lines.

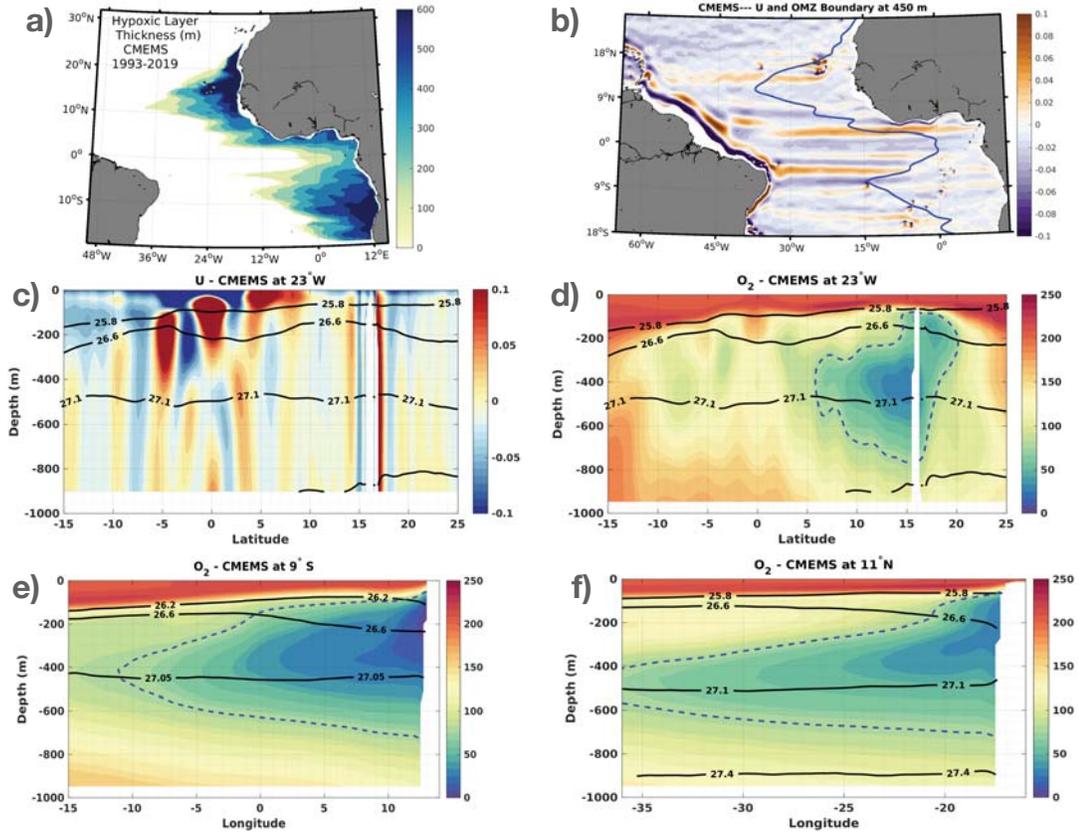


Figure A7. Averages from 1993 to 2019 obtained from two CMEMS model products: (a) hypoxic layer thickness, (b) zonal velocities at 450 m with the $80 \mu\text{mol L}^{-1}$ oxygen concentration isoline shown as a blue line, (c) zonal velocities at 23°W with isopycnal surfaces shown as black lines, (d) oxygen concentrations at 23°W with isopycnal surfaces shown as black lines, (e) averaged oxygen concentrations at 9°S and (f) 11°N.

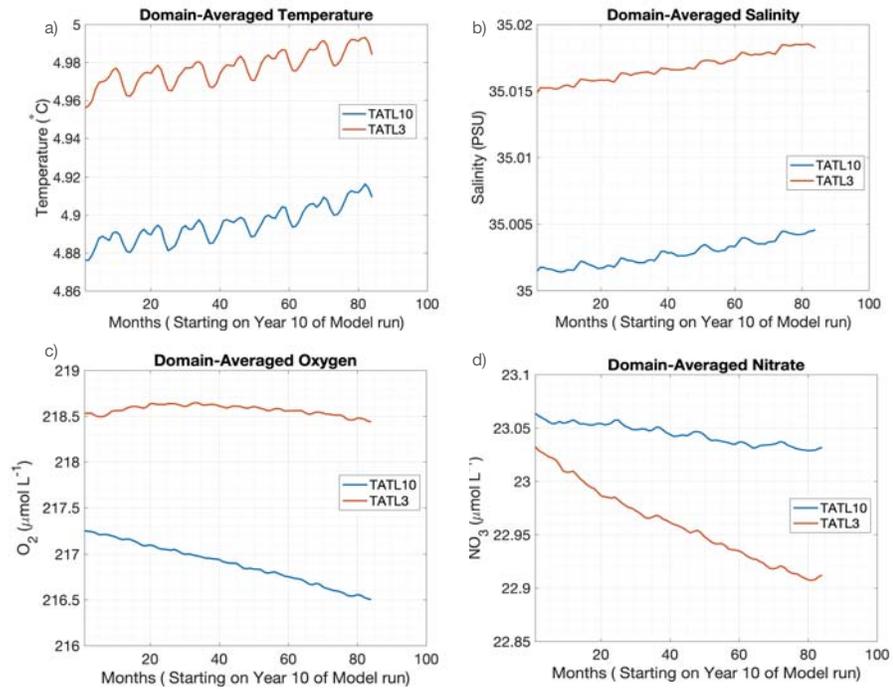


Figure A8. Domain-averaged model drift during the simulation period analyzed in this study.

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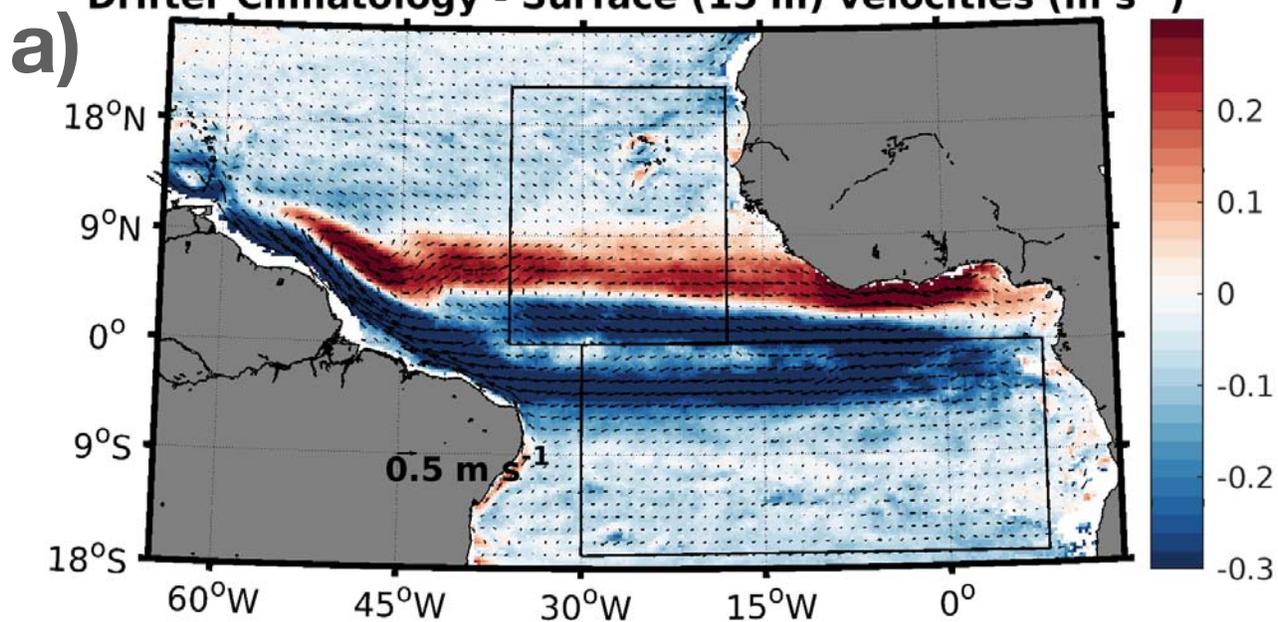
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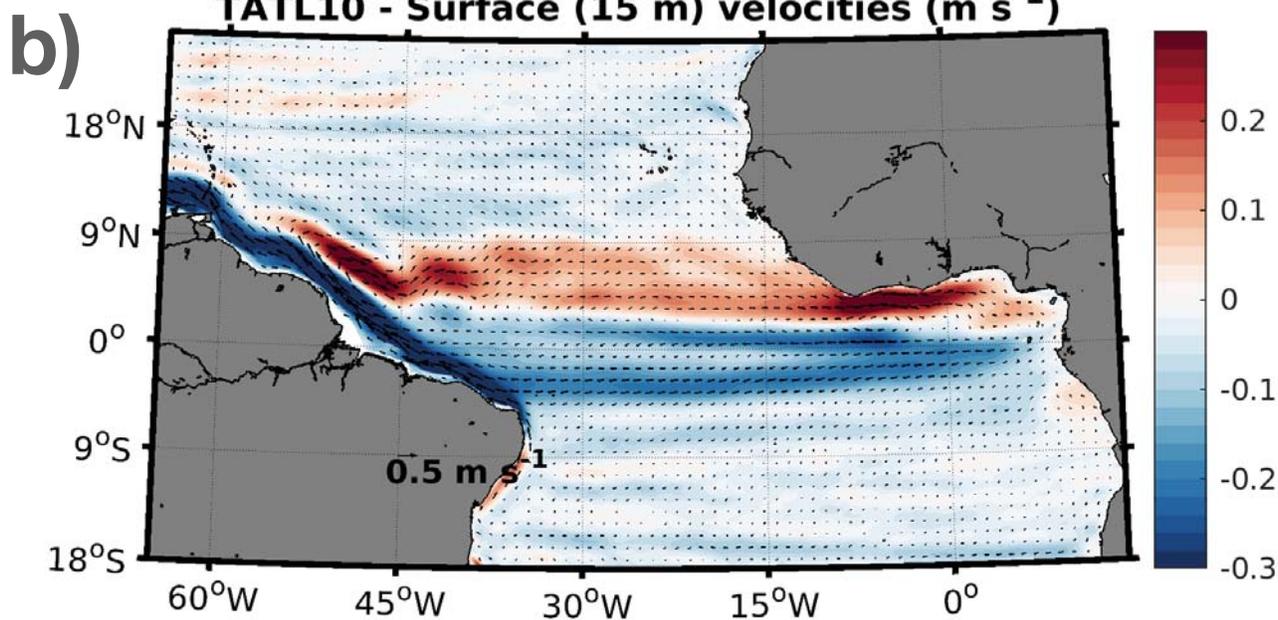
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Figure A1.

Drifter Climatology - Surface (15 m) velocities (m s^{-1})



TATL10 - Surface (15 m) velocities (m s^{-1})



TATL3 - Surface (15 m) velocities (m s^{-1})

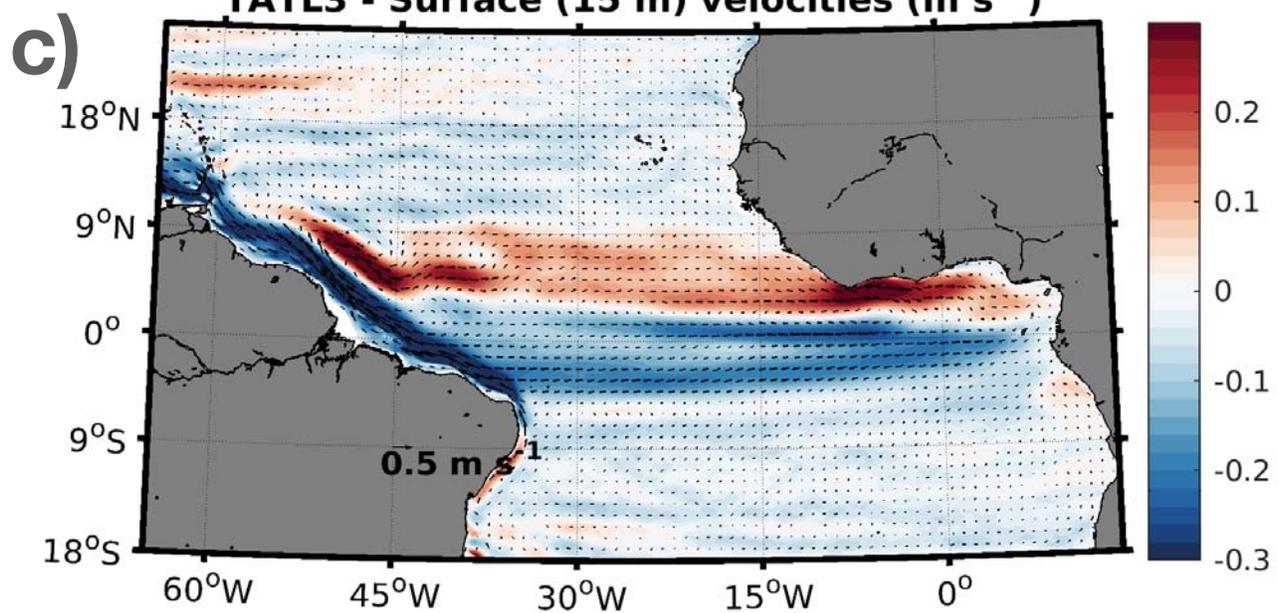


Figure A2.

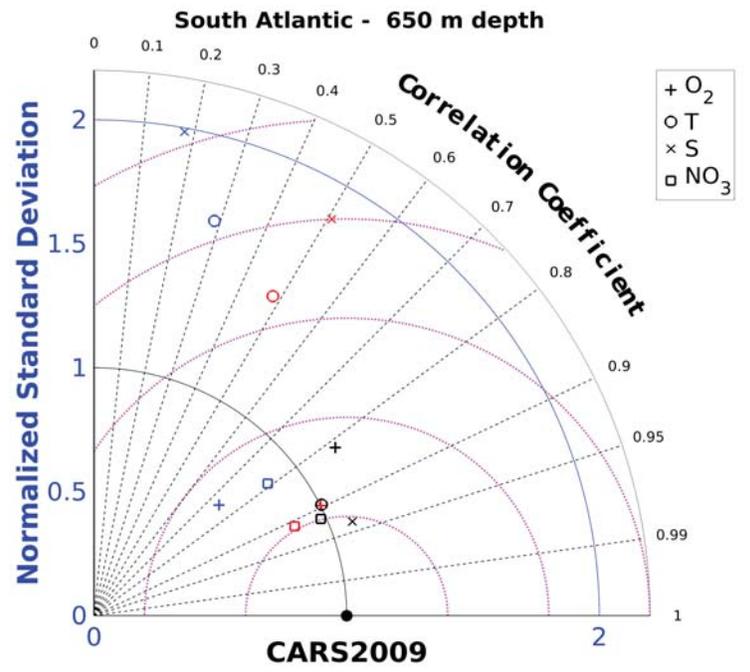
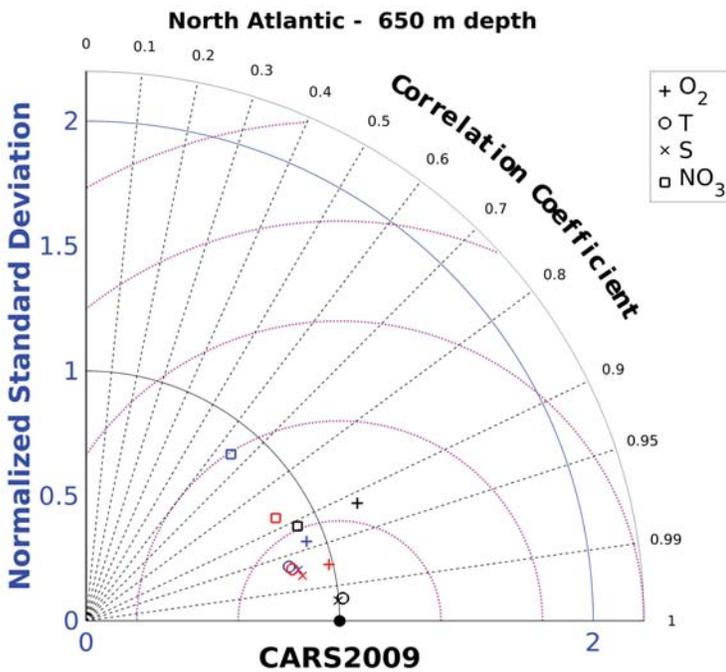
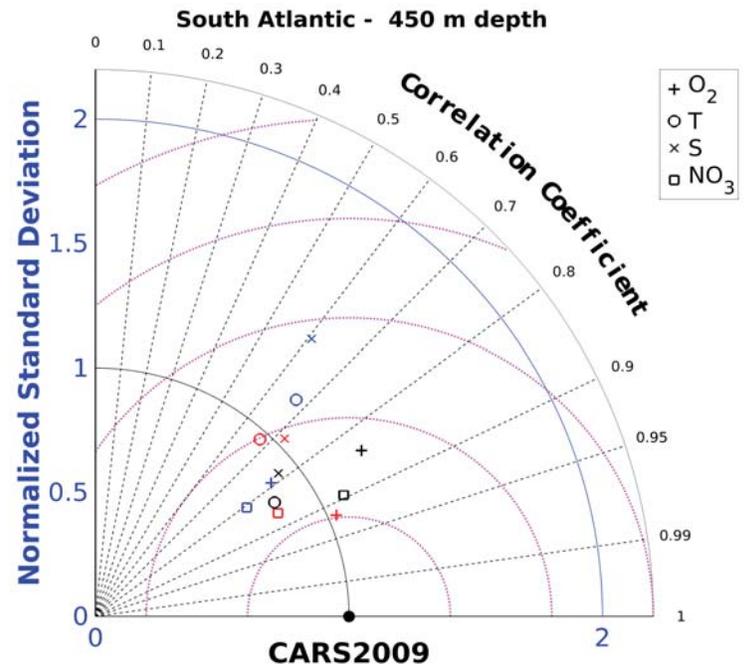
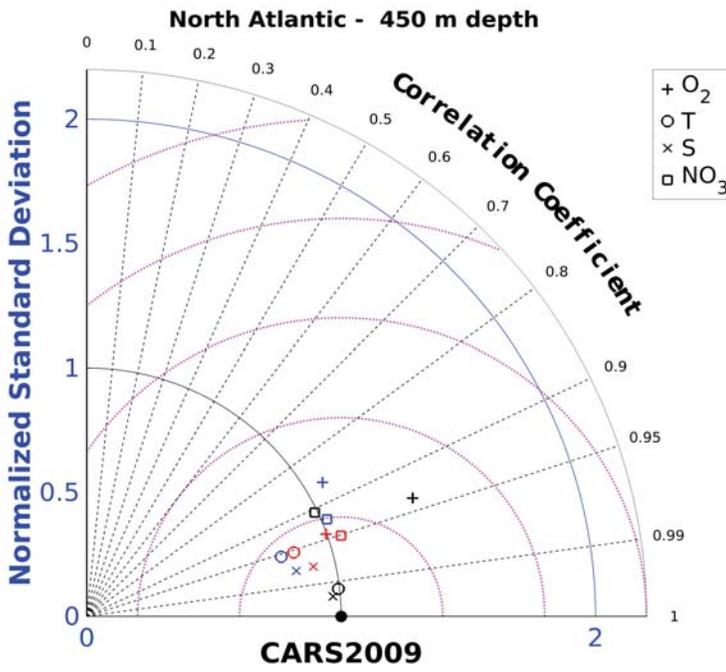
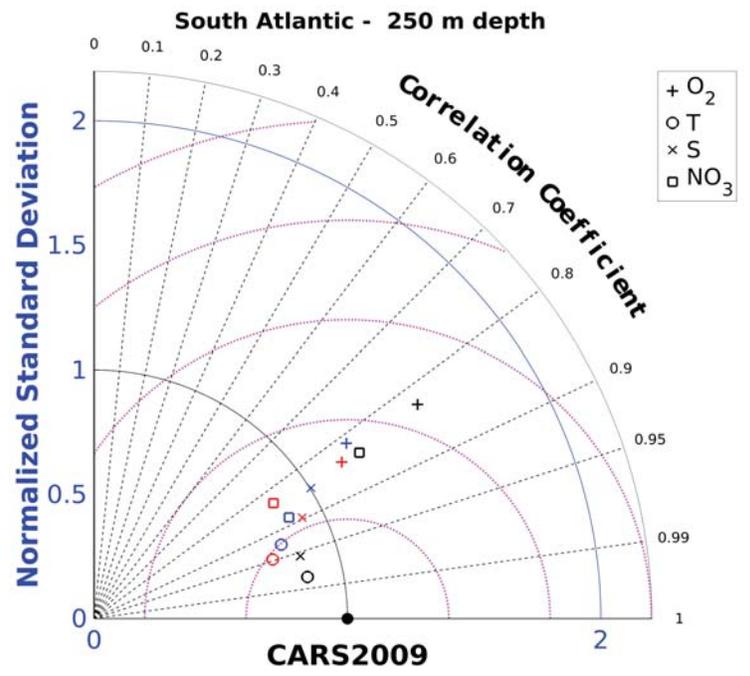
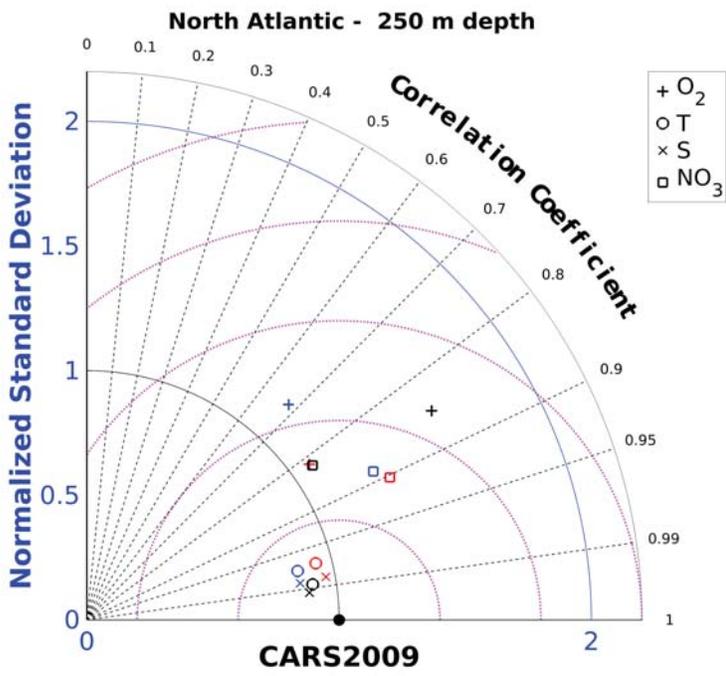
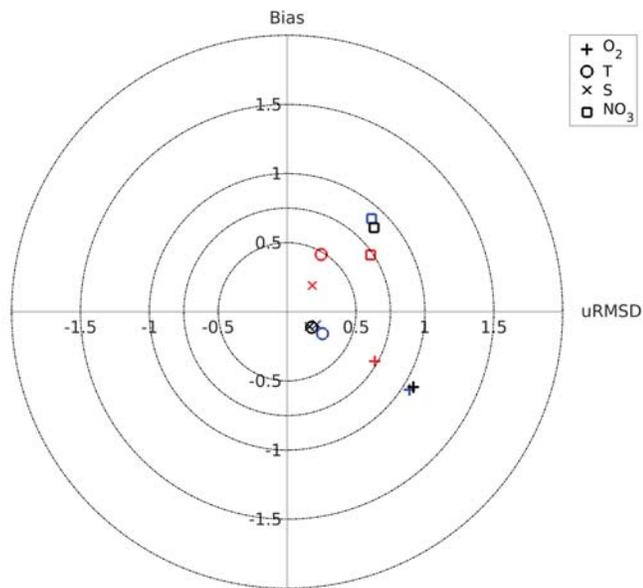
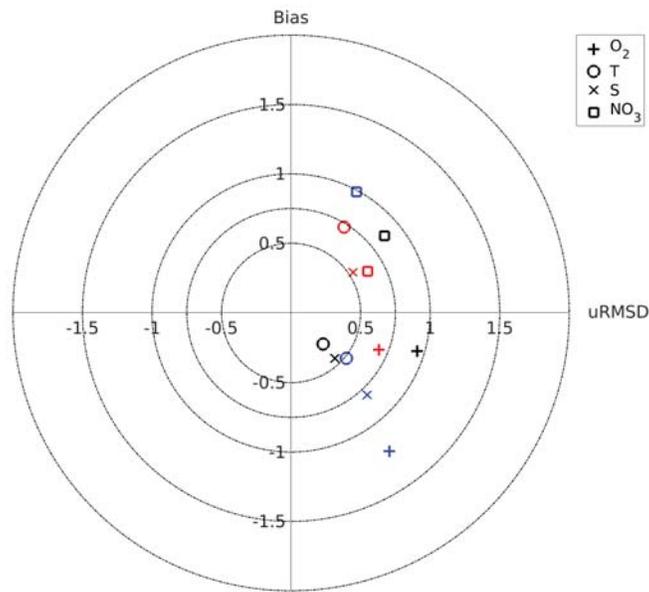


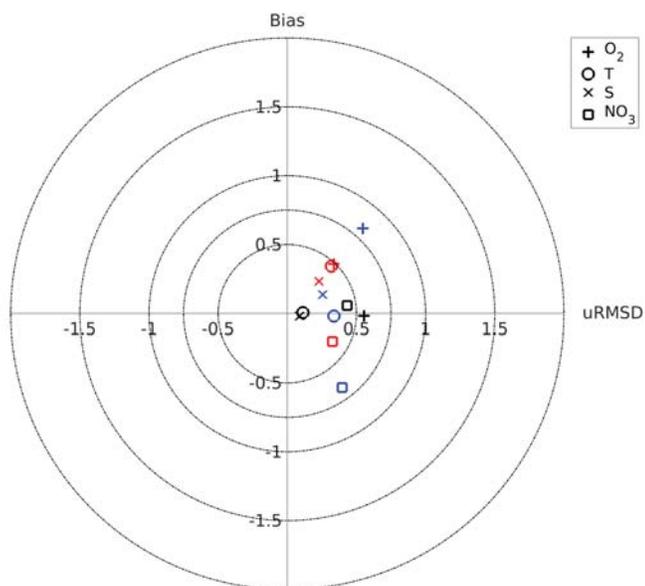
Figure A3.



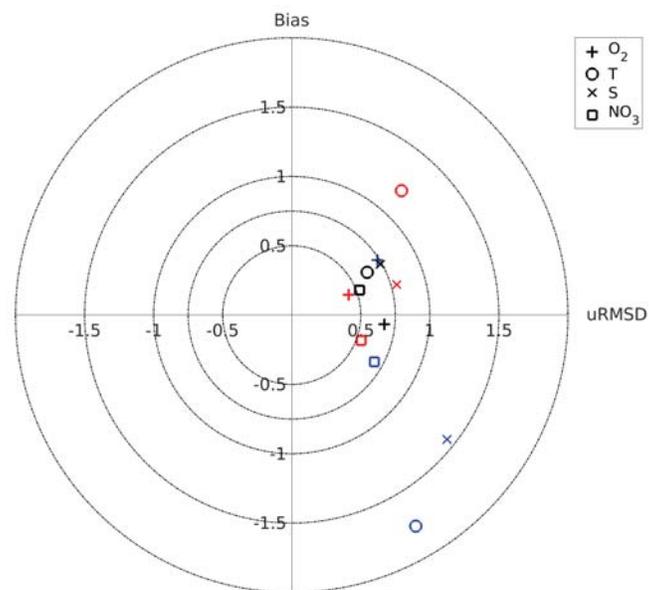
North Atlantic - 250 m depth



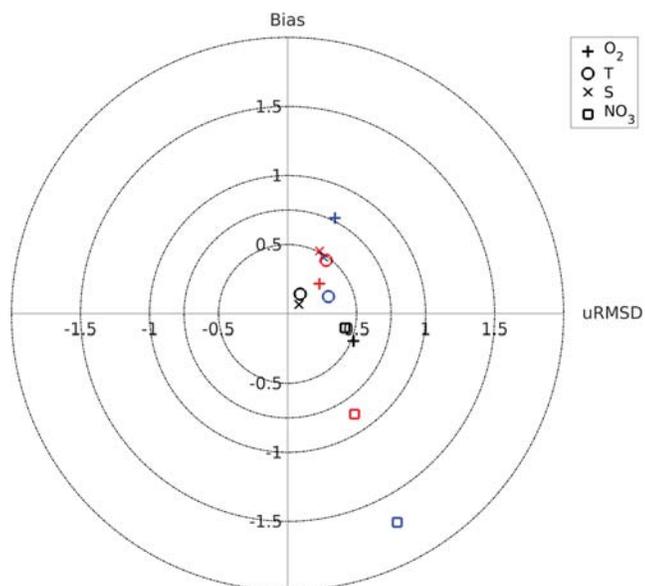
South Atlantic - 250 m depth



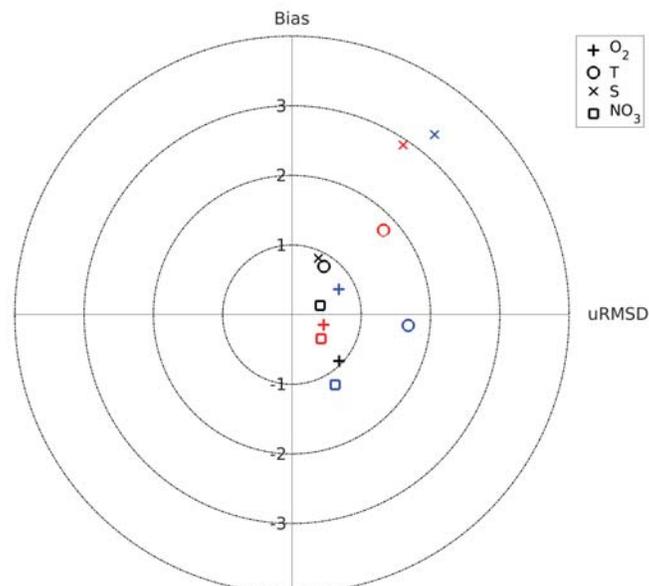
North Atlantic - 450 m depth



South Atlantic - 450 m depth



North Atlantic - 650 m depth



South Atlantic - 650 m depth

Figure A4.

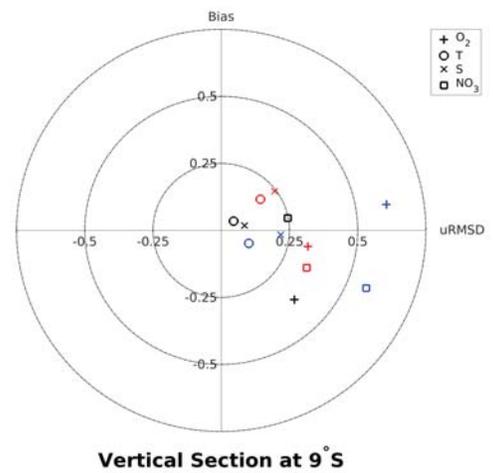
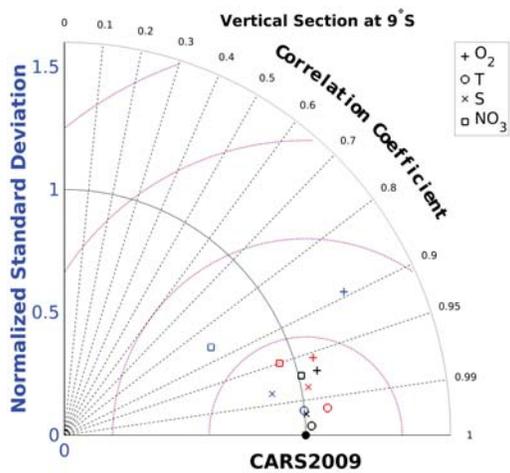
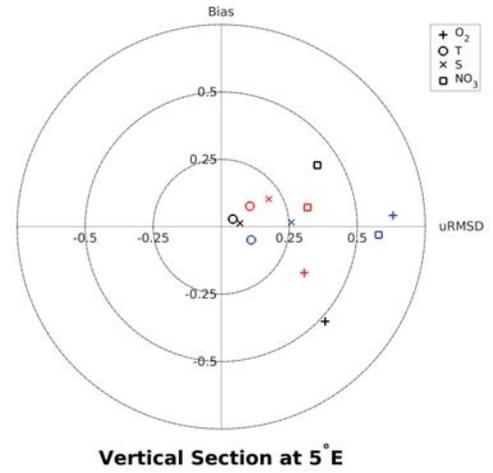
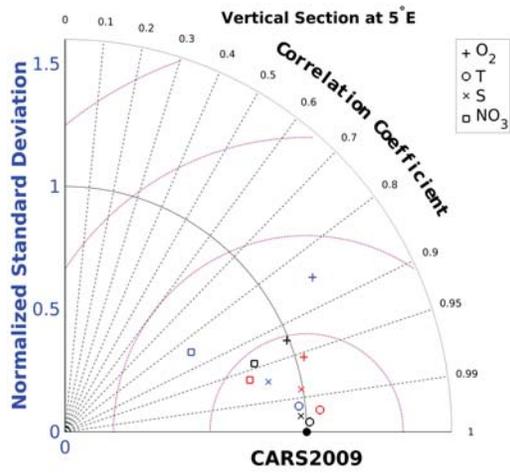
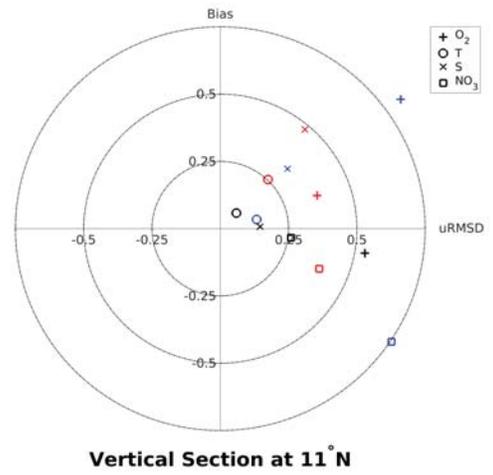
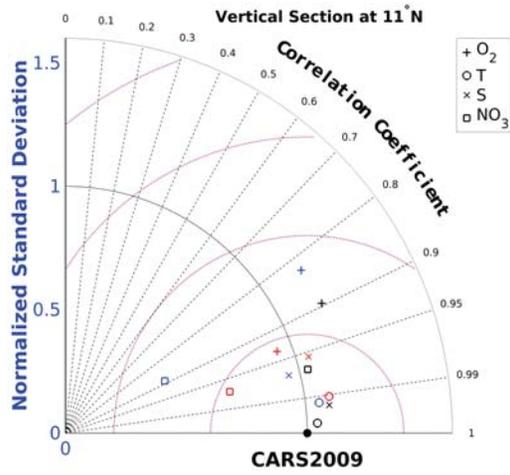
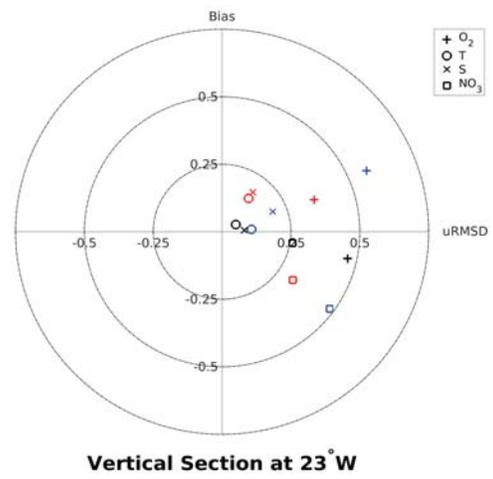
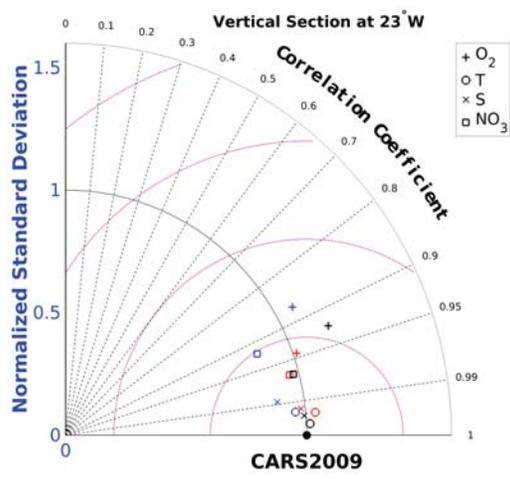


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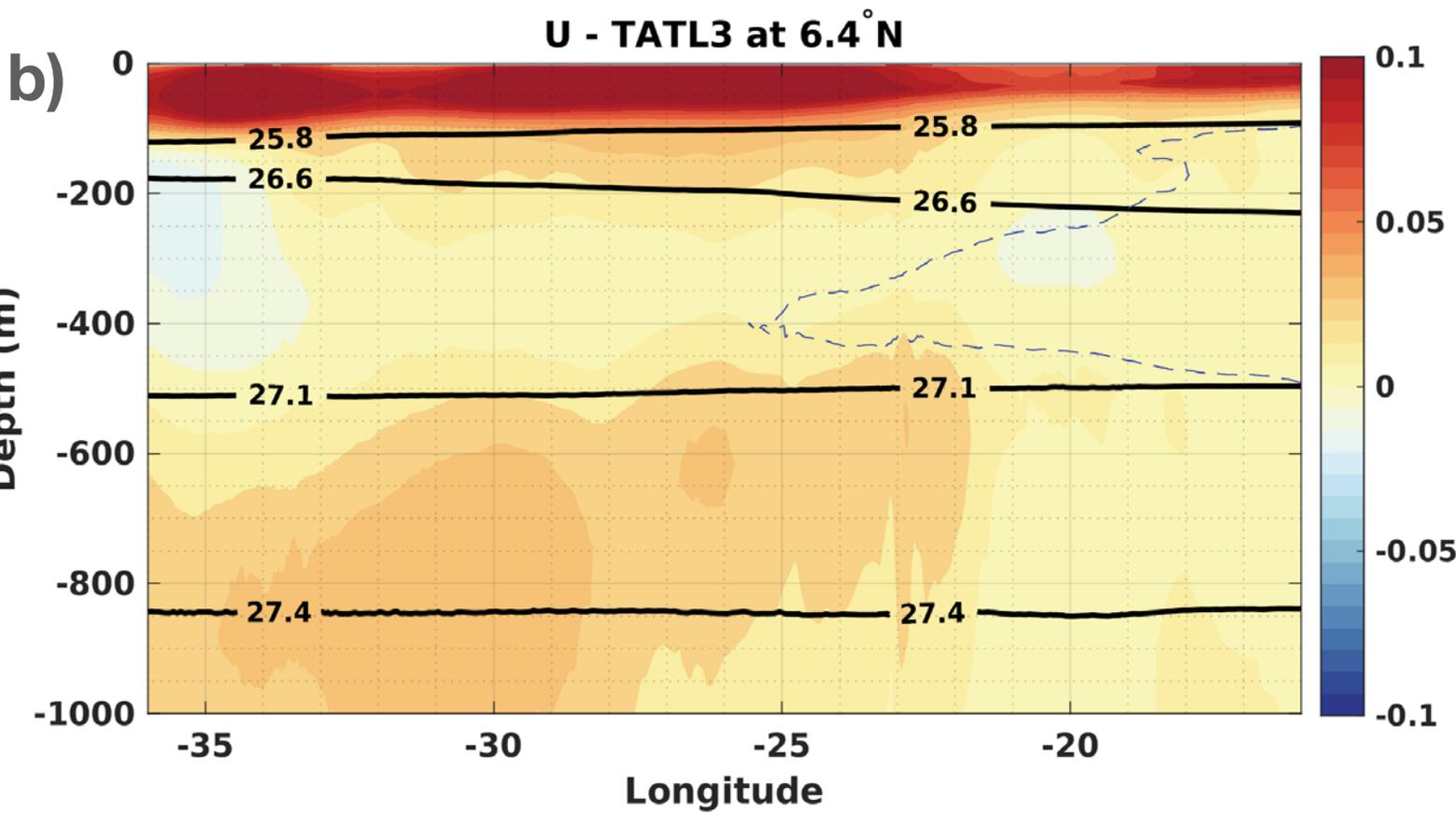
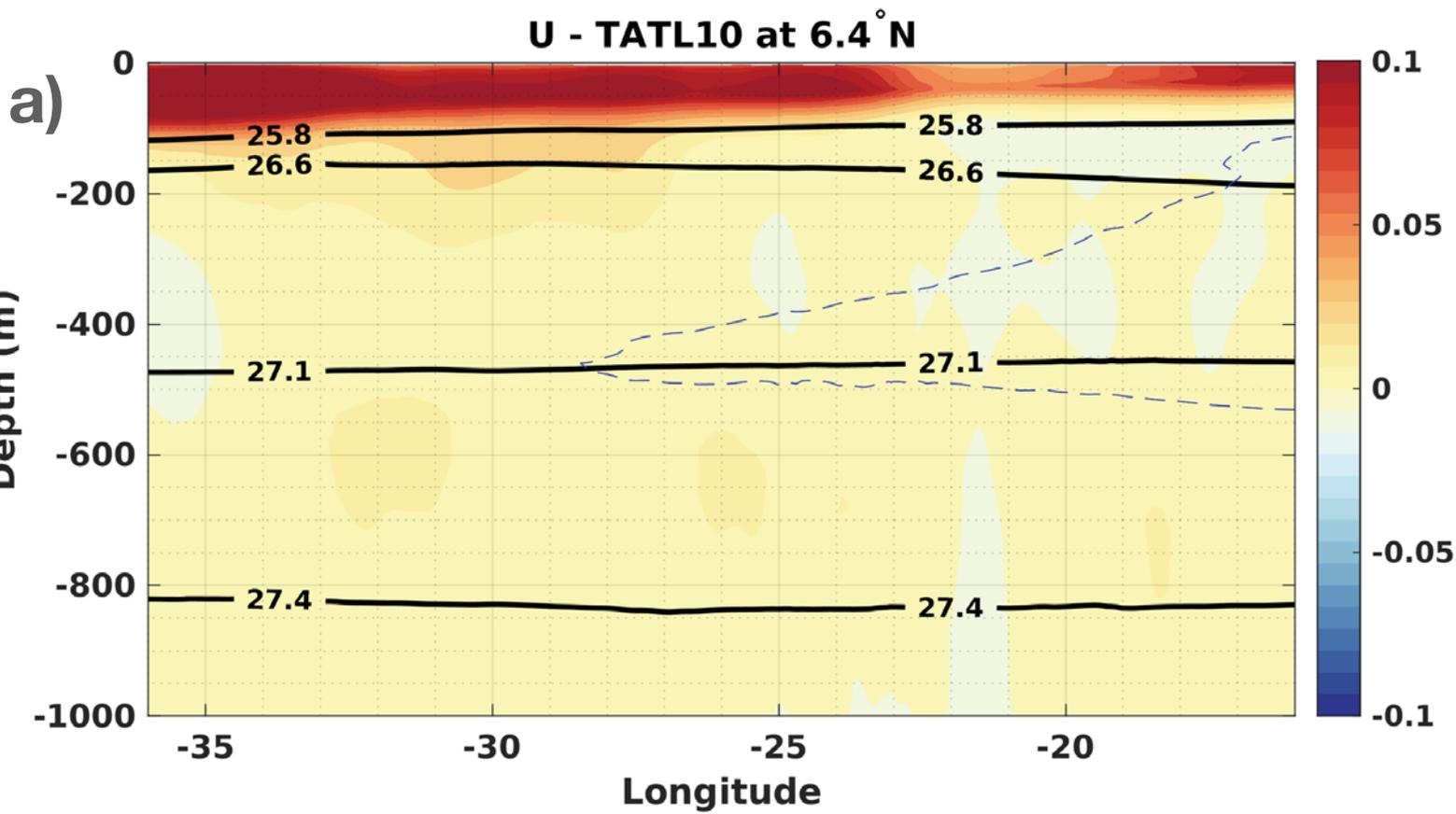


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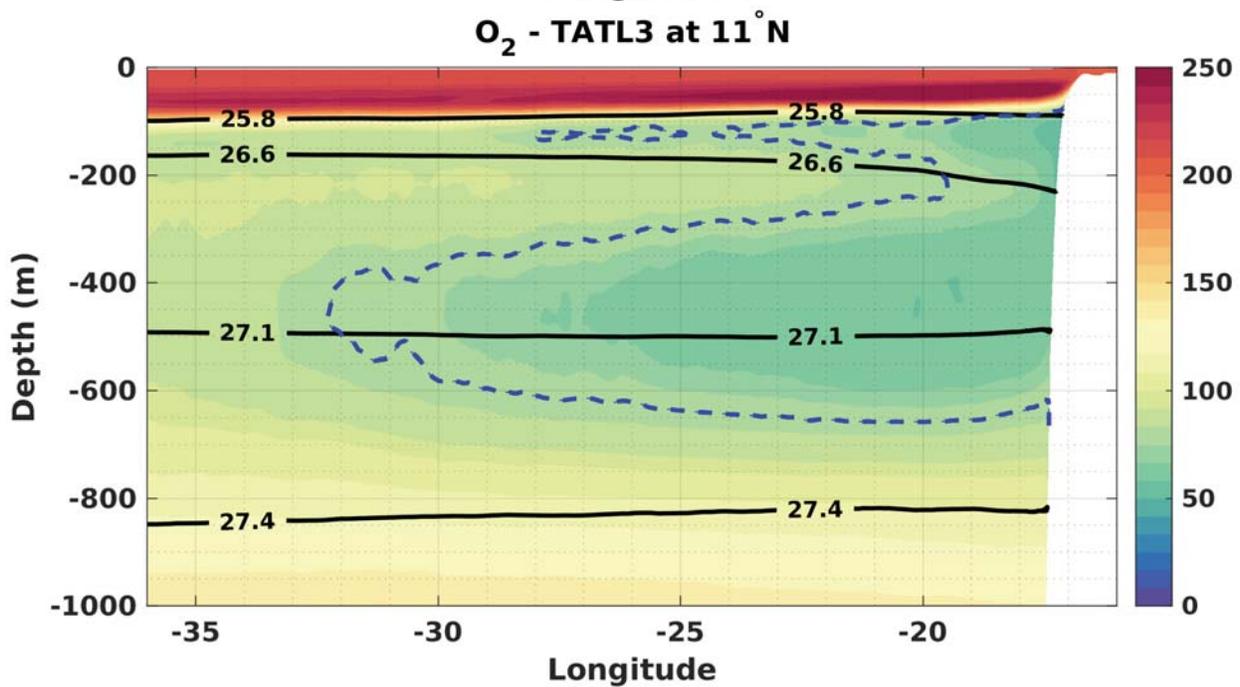
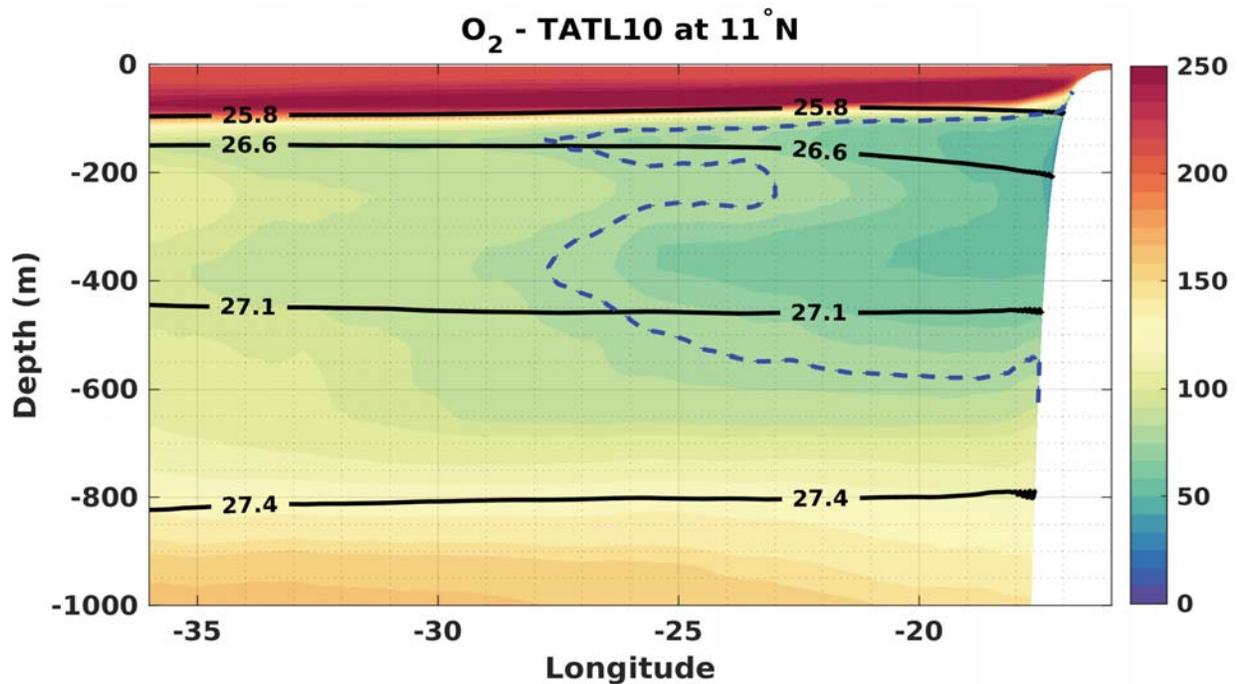
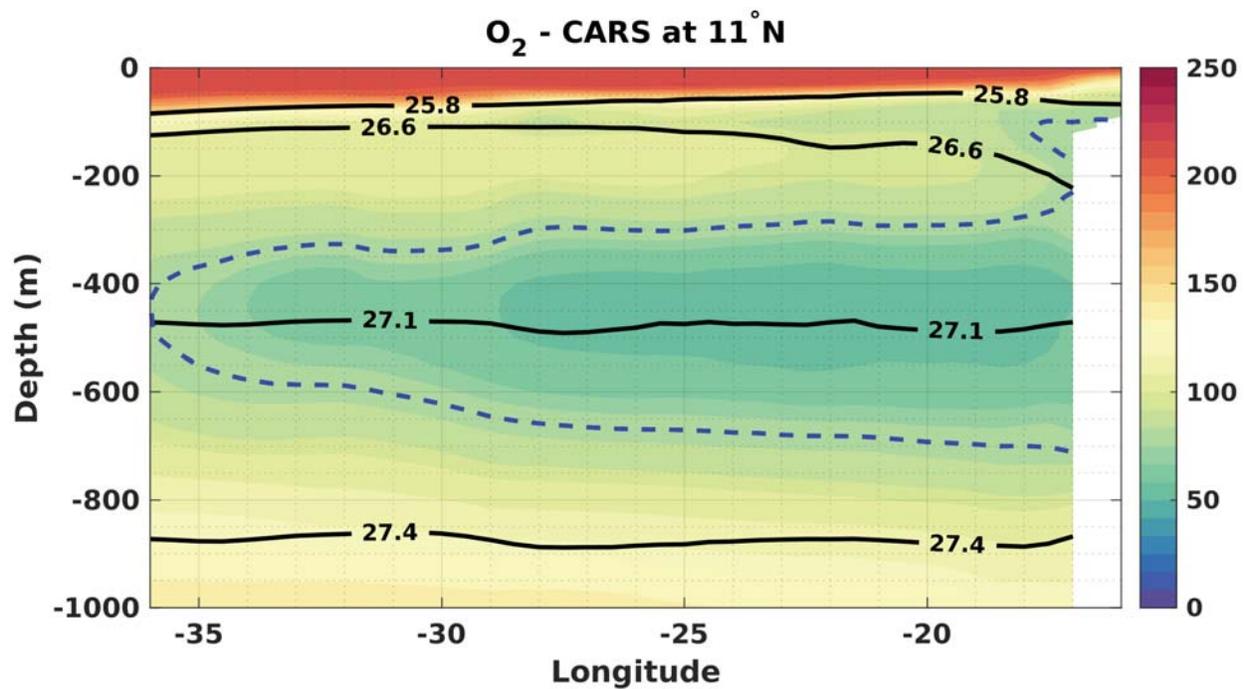


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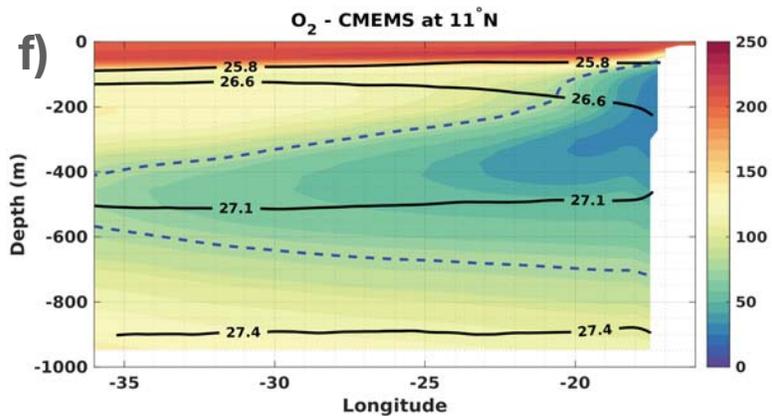
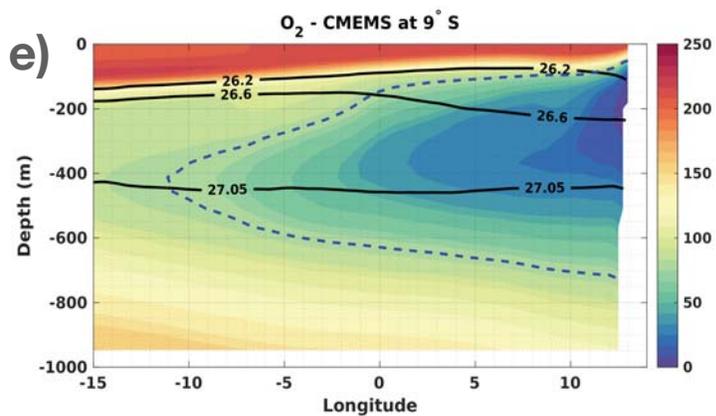
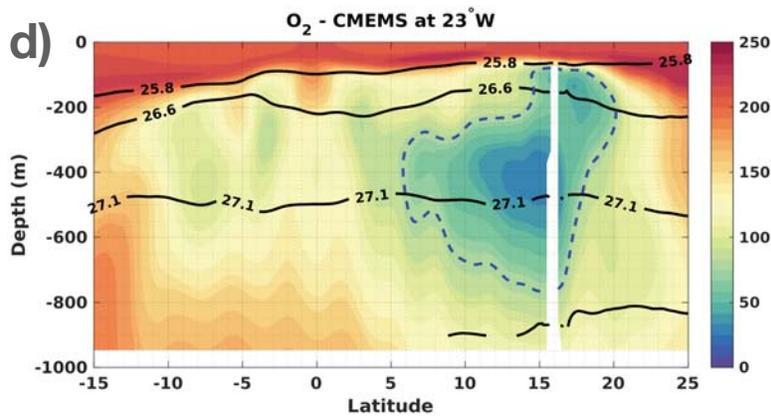
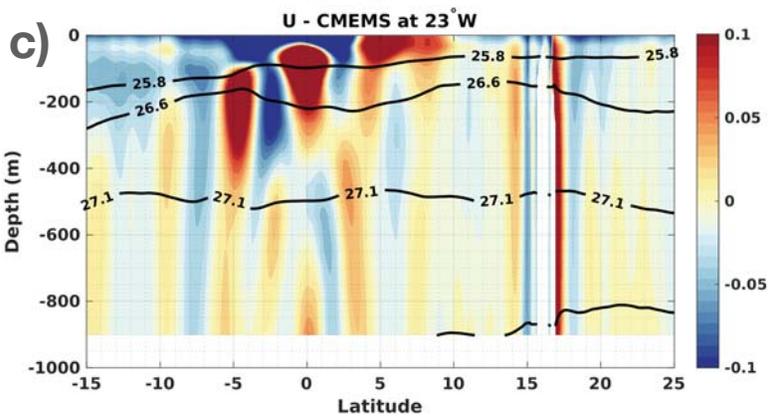
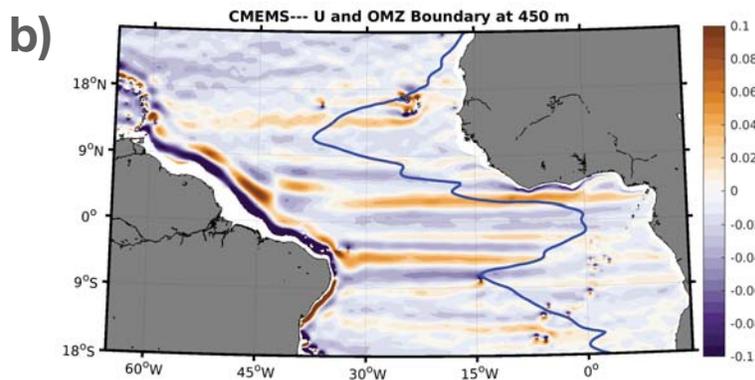
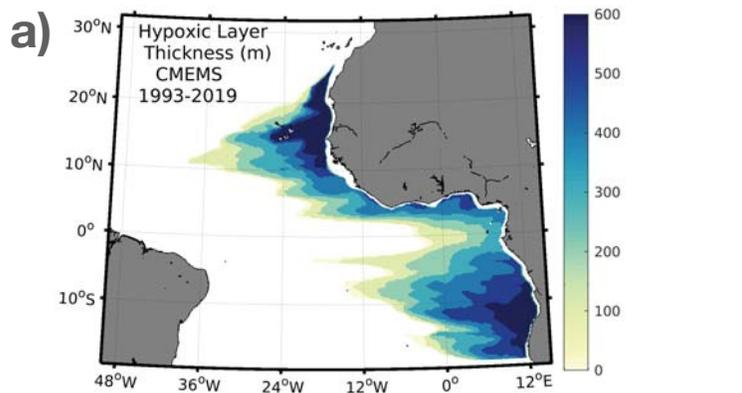
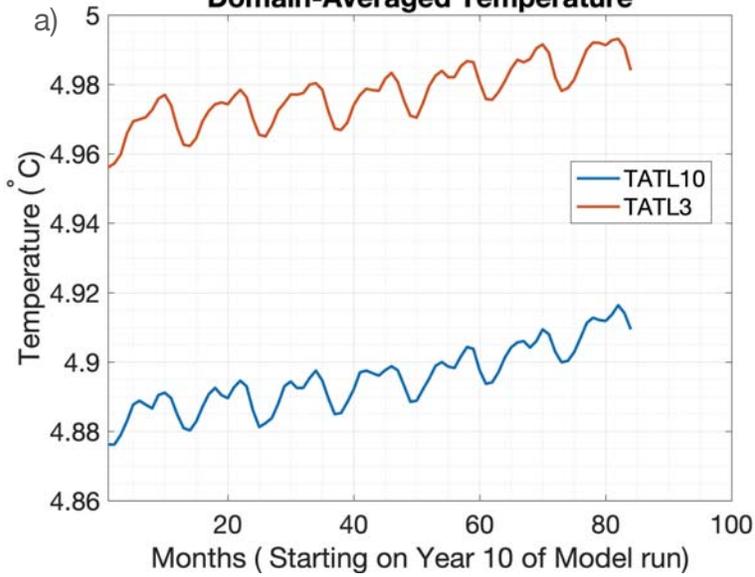
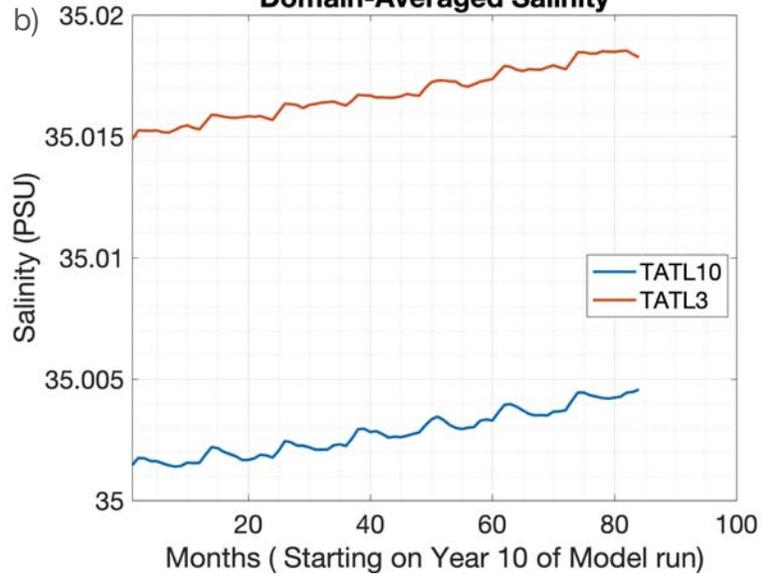


Figure A8.

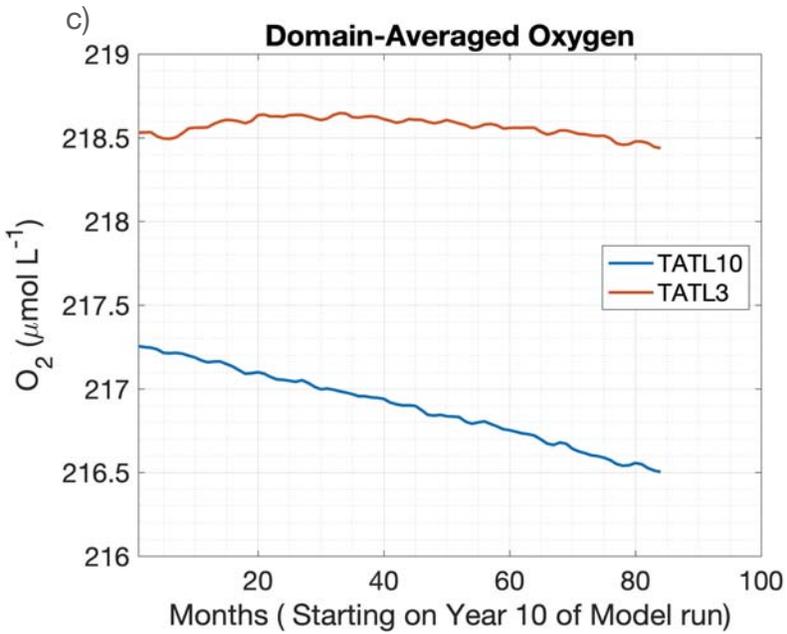
Domain-Averaged Temperature



Domain-Averaged Salinity



Domain-Averaged Oxygen



Domain-Averaged Nitrate

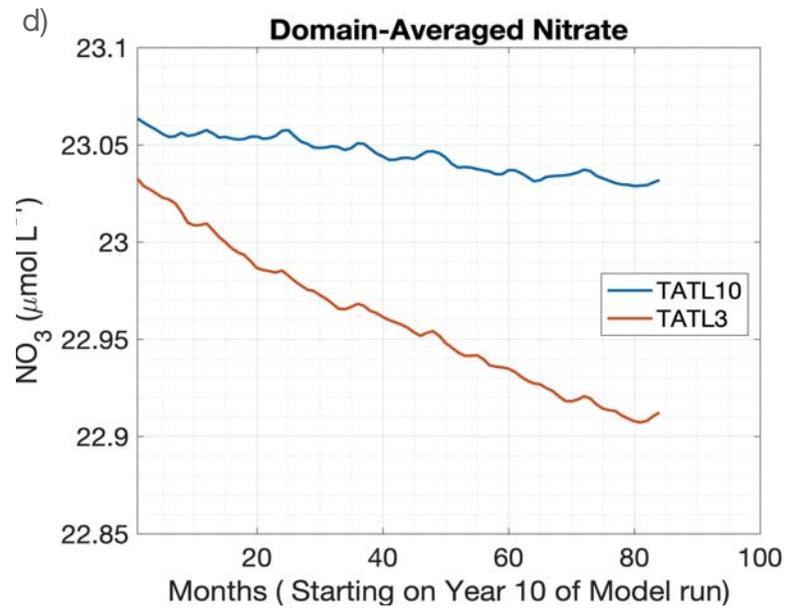


Figure 1.

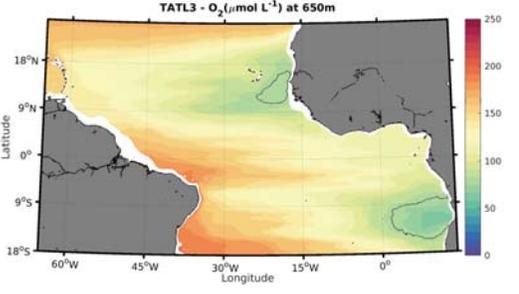
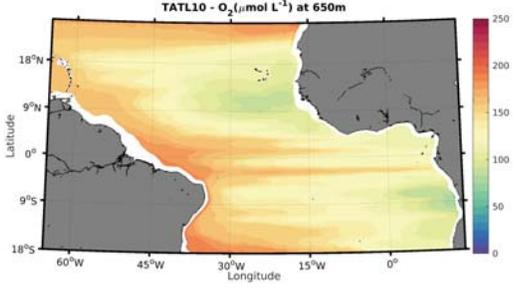
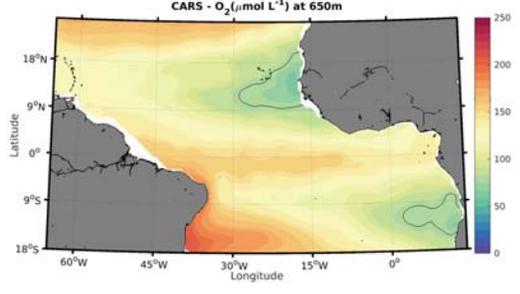
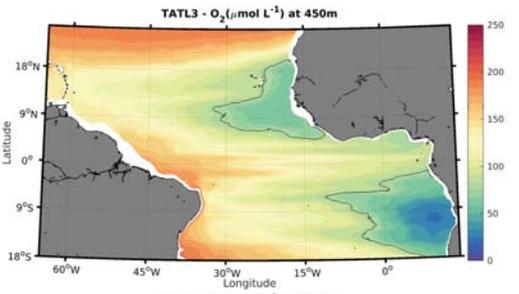
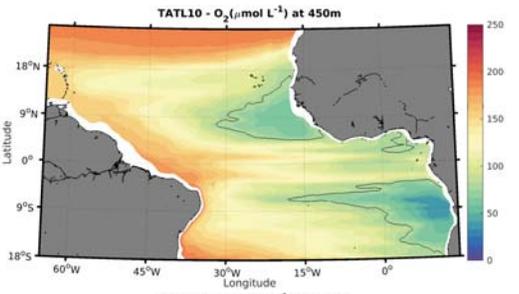
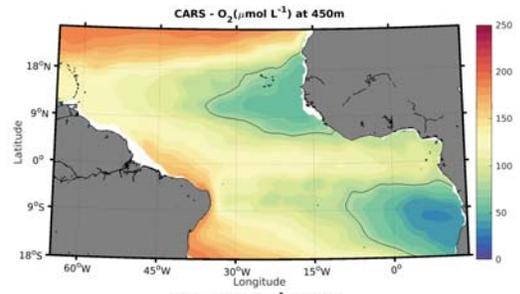
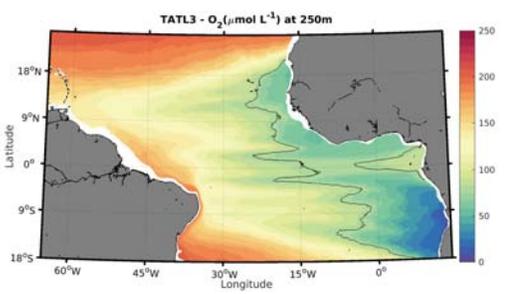
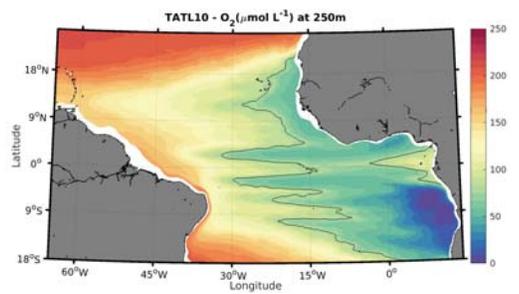
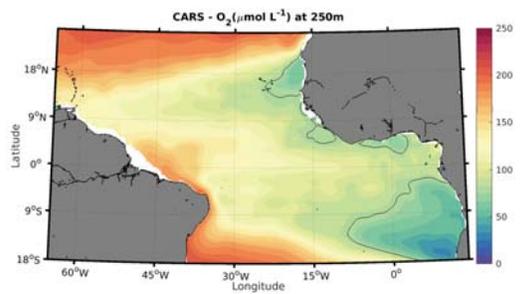


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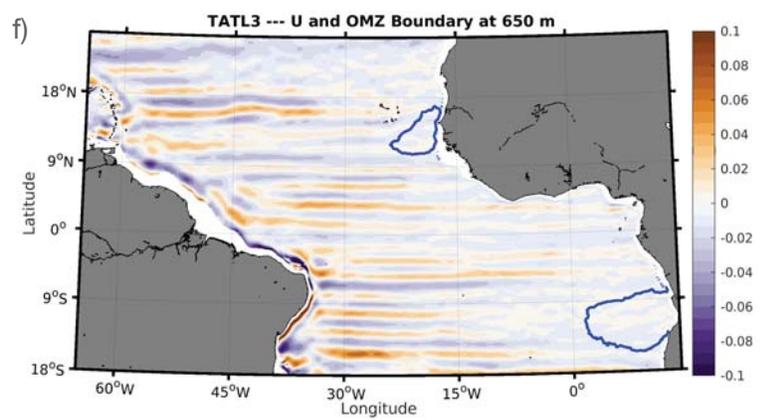
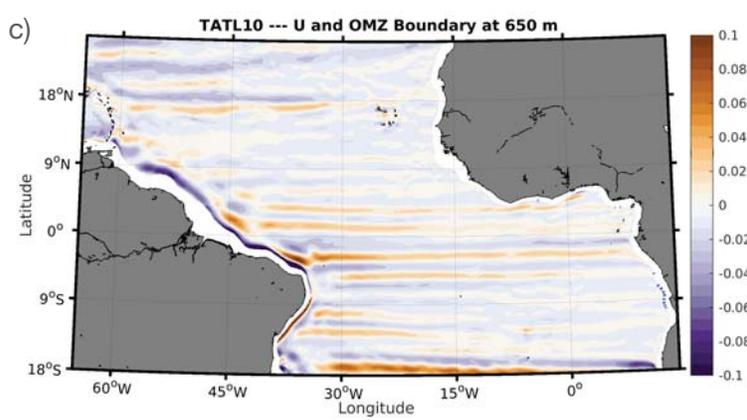
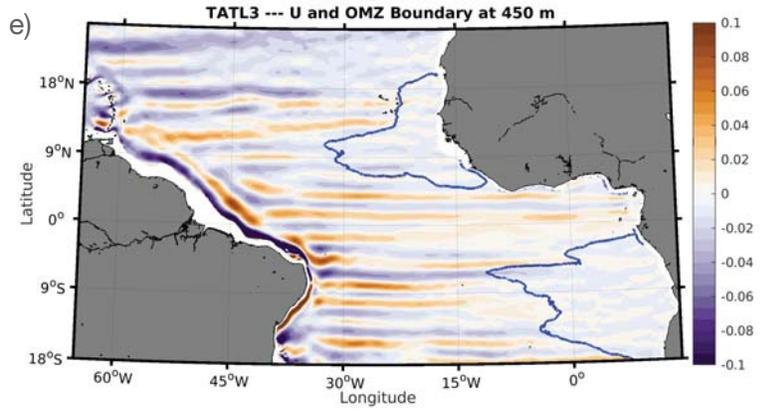
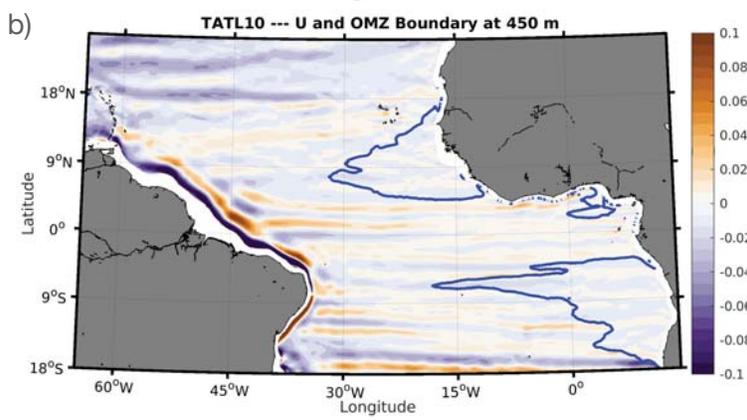
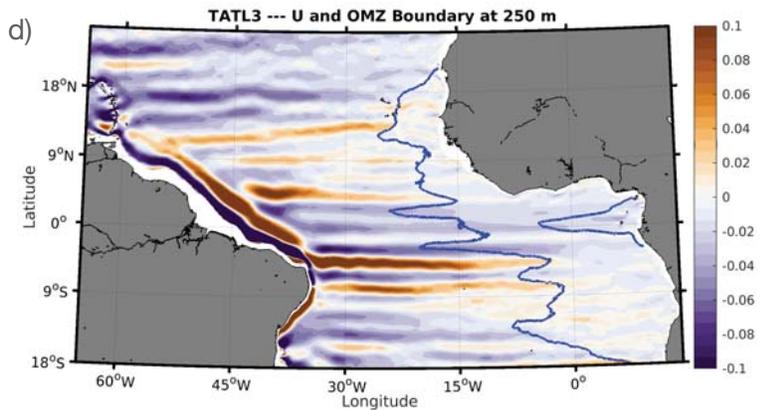
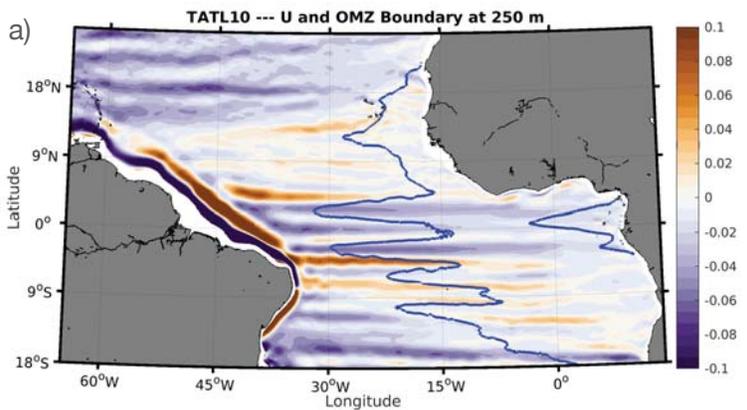


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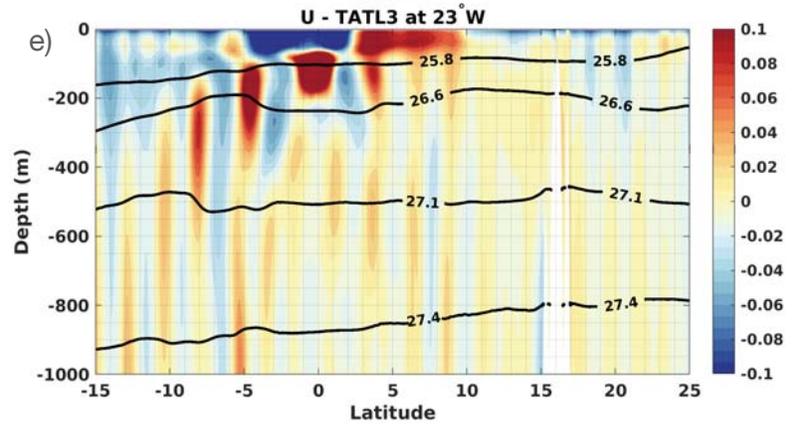
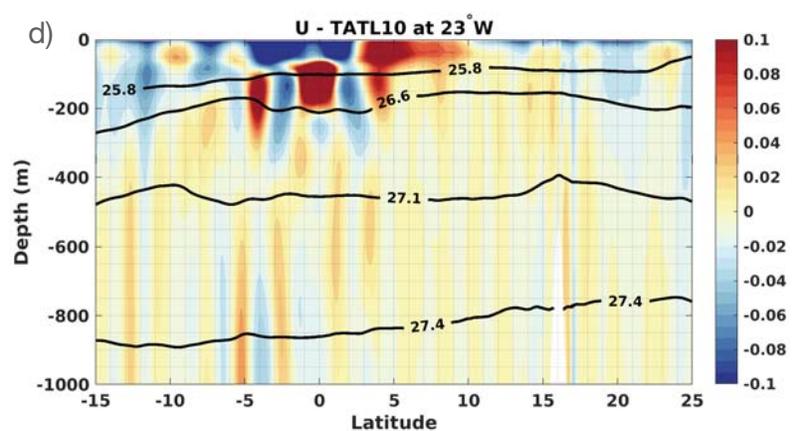
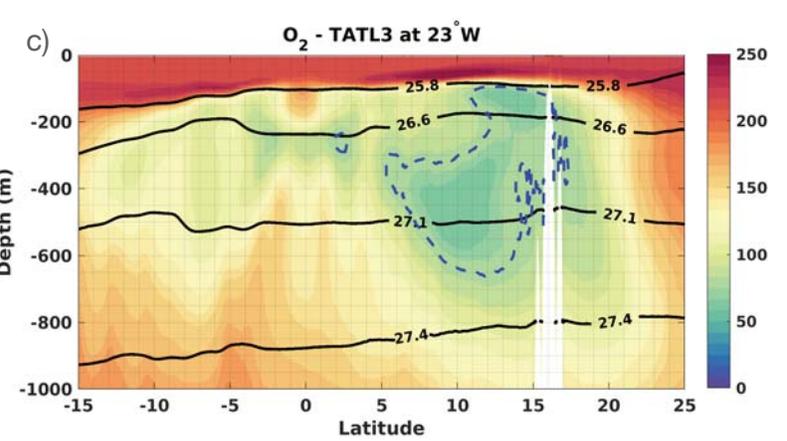
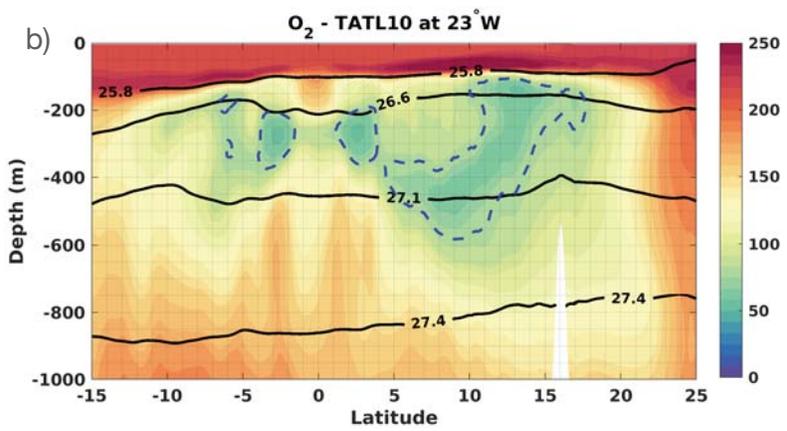
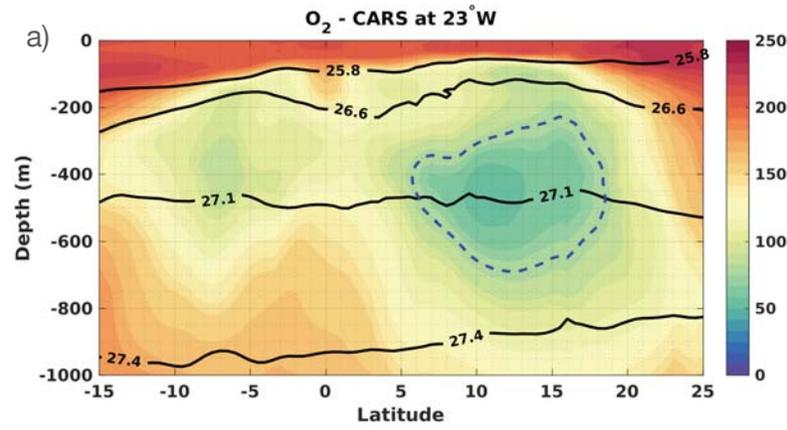


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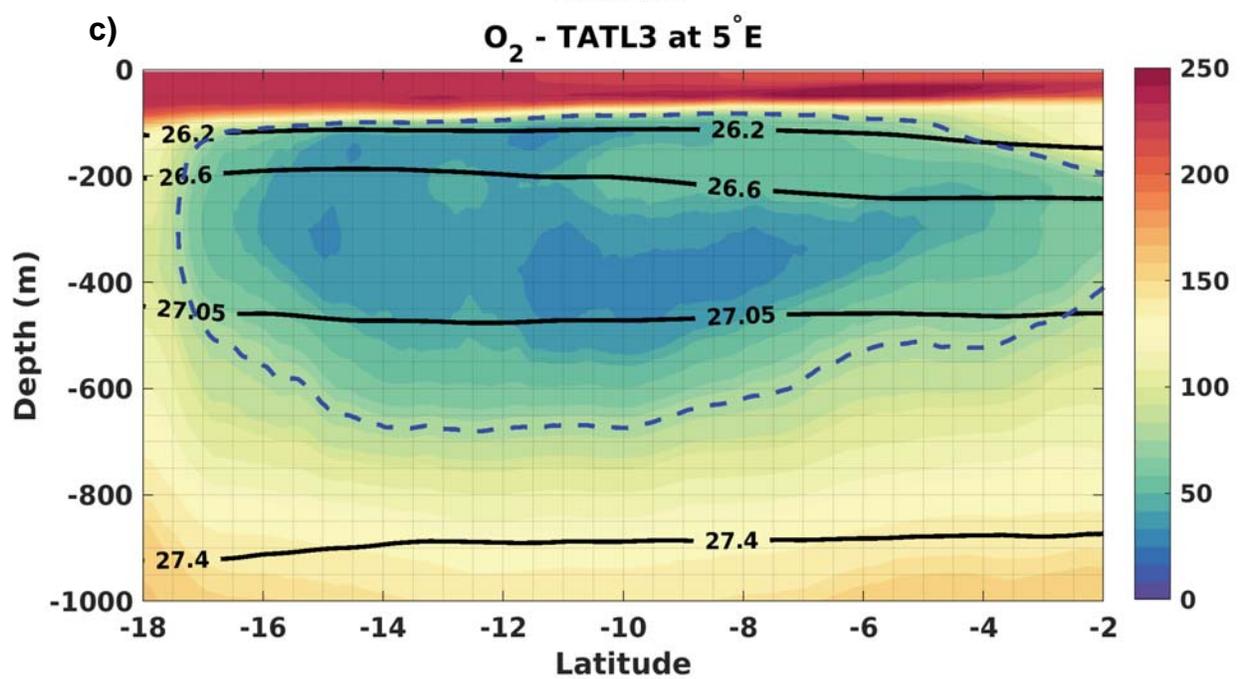
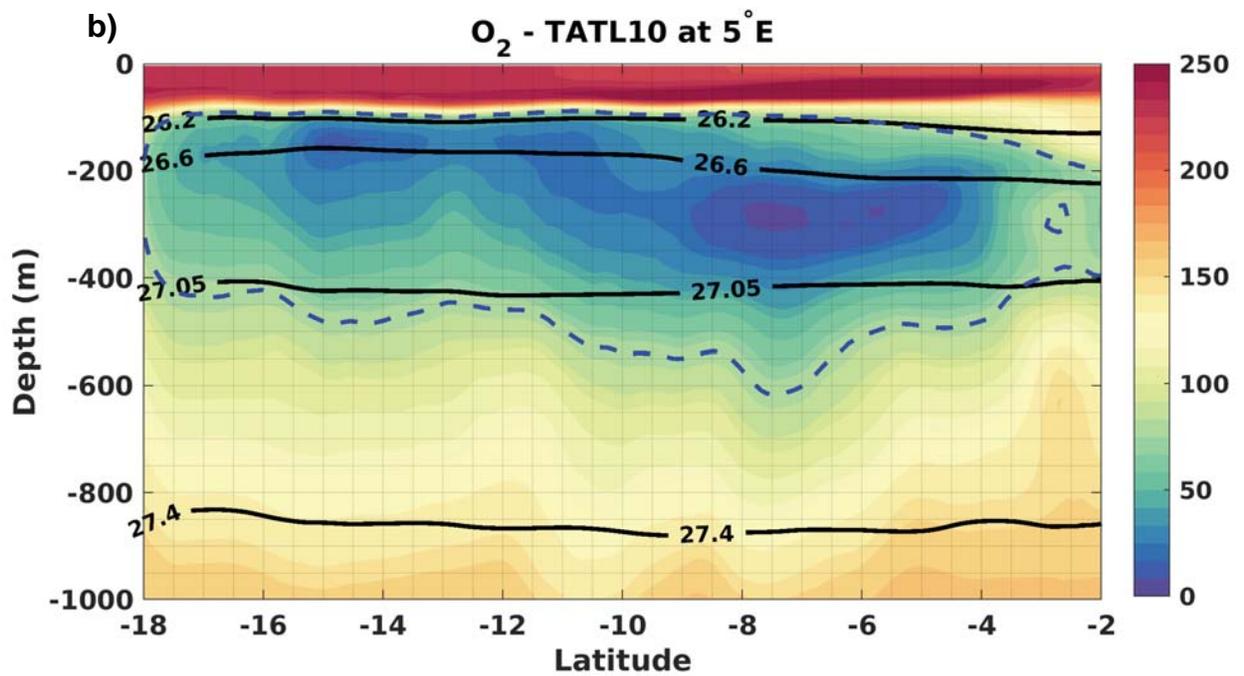
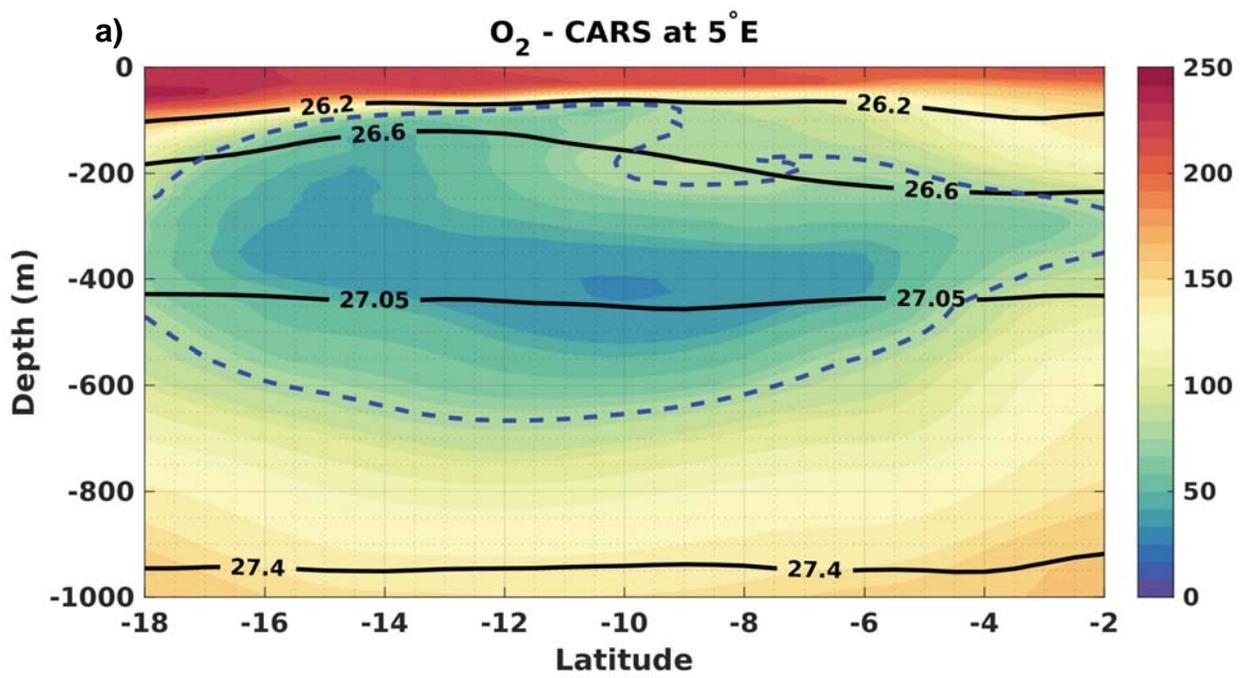


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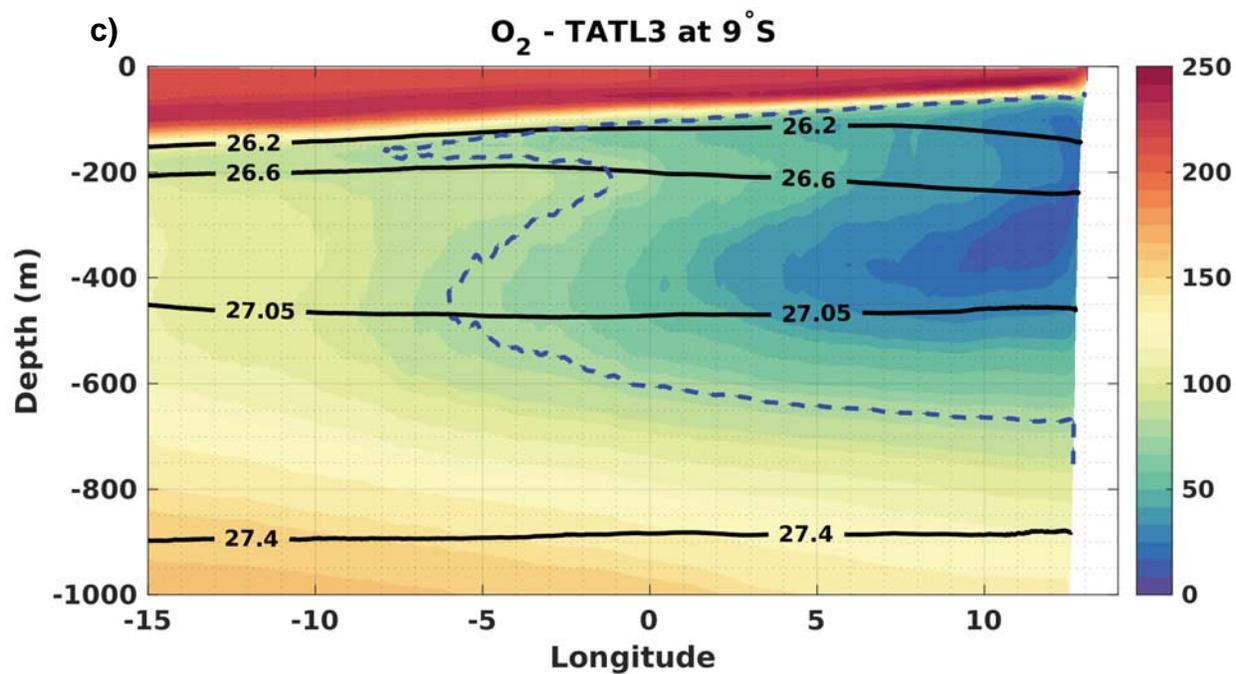
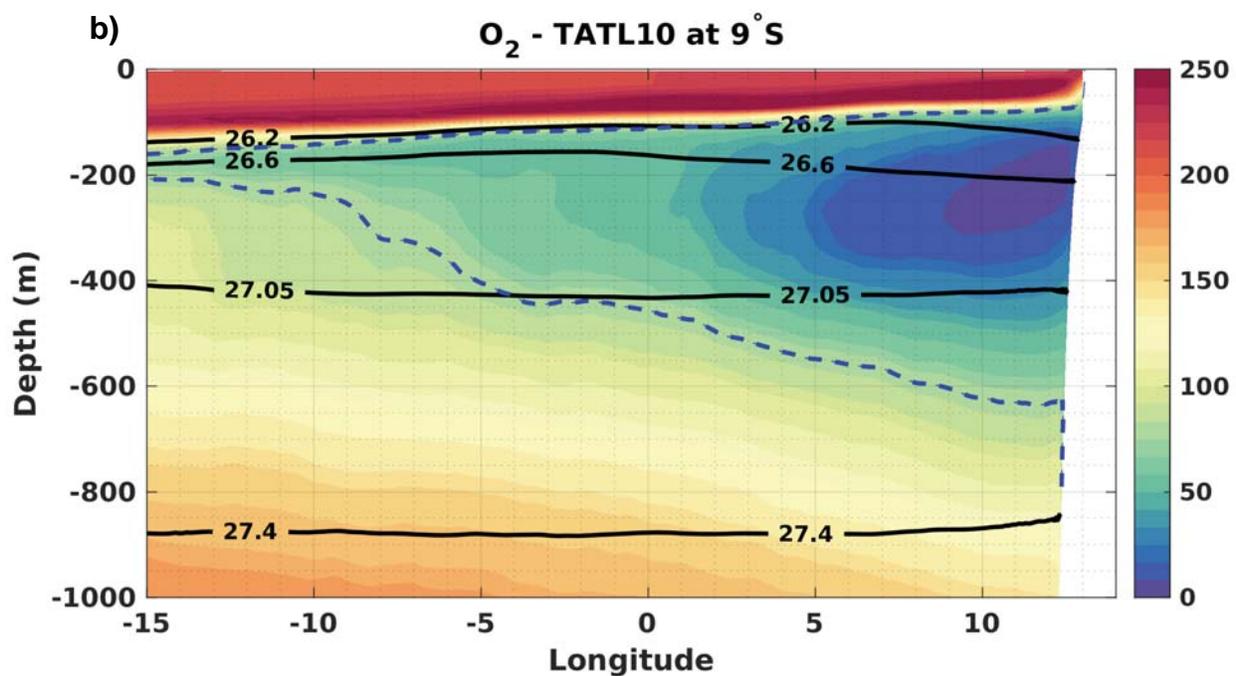
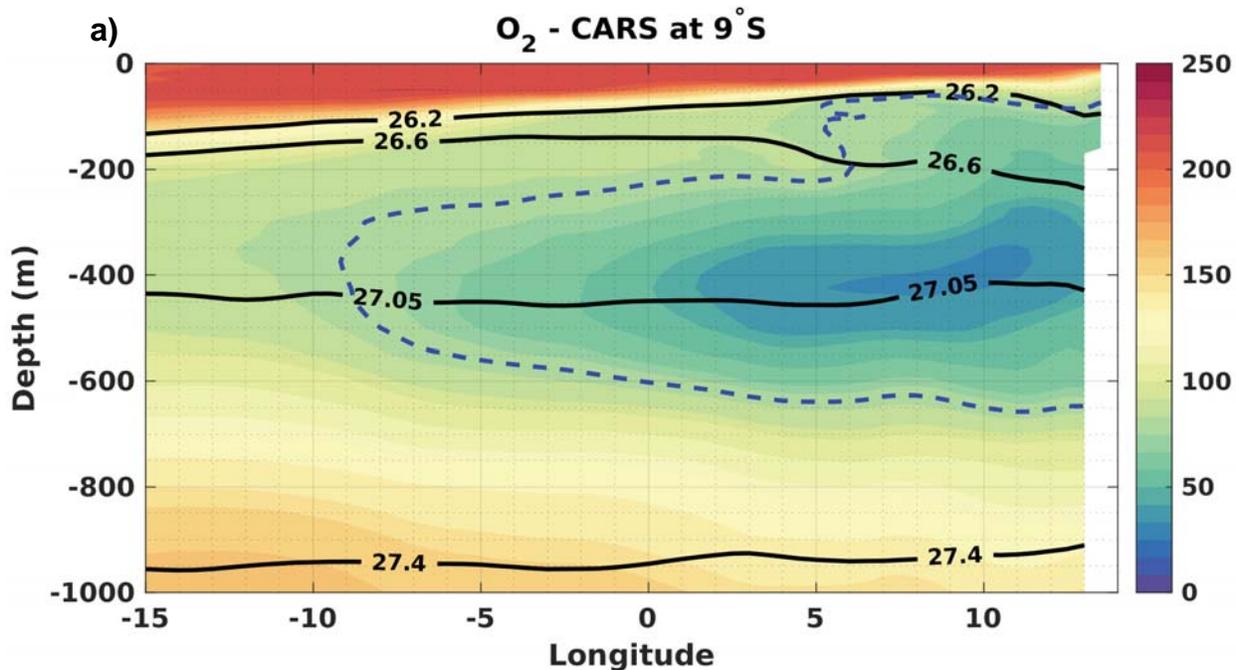


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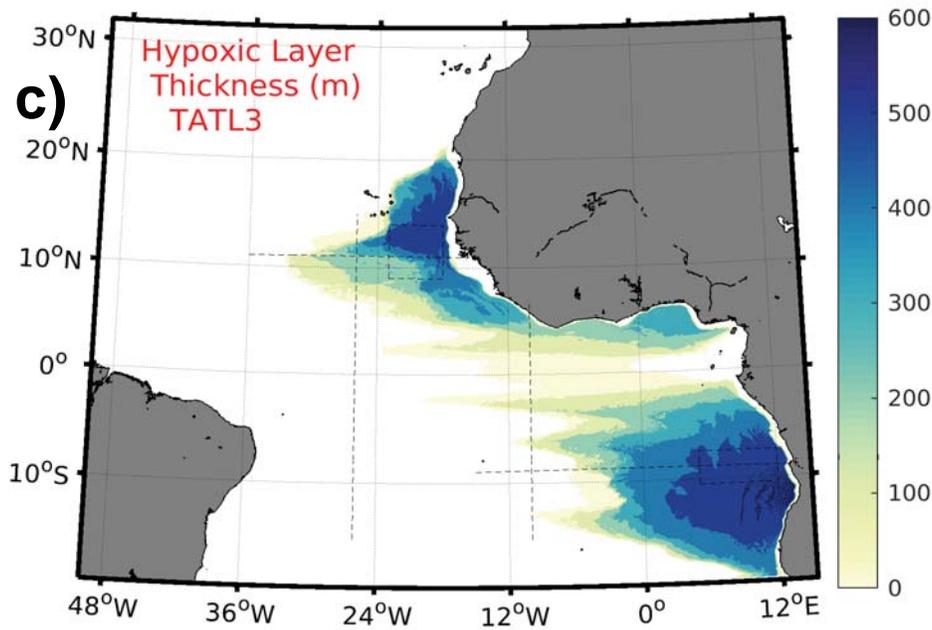
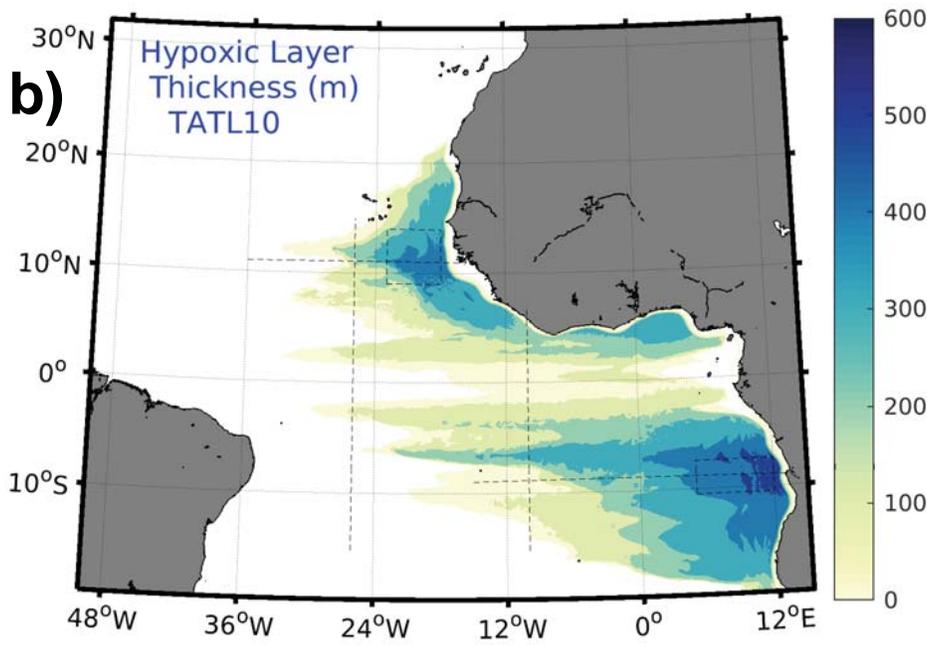
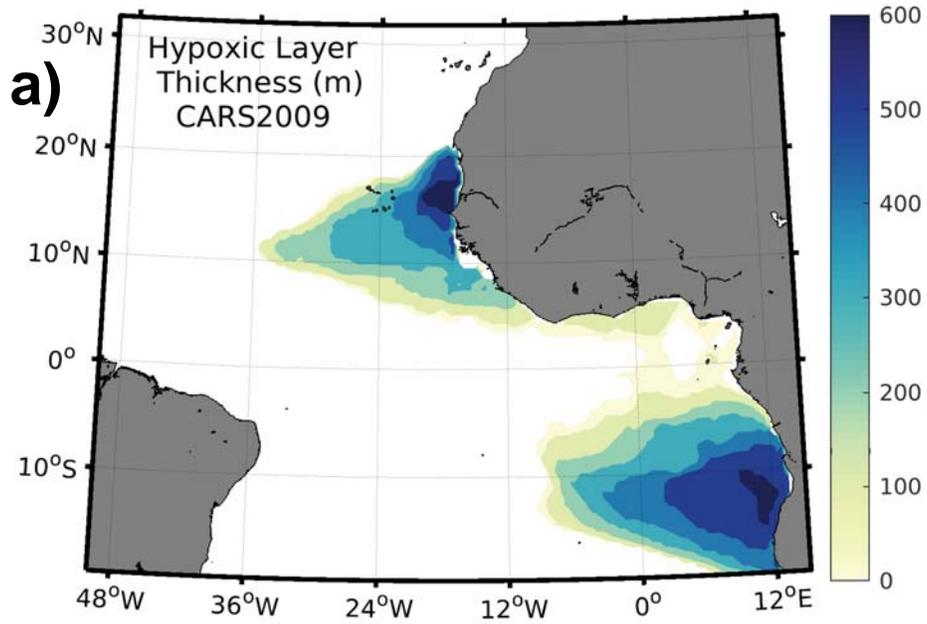


Figure 7.

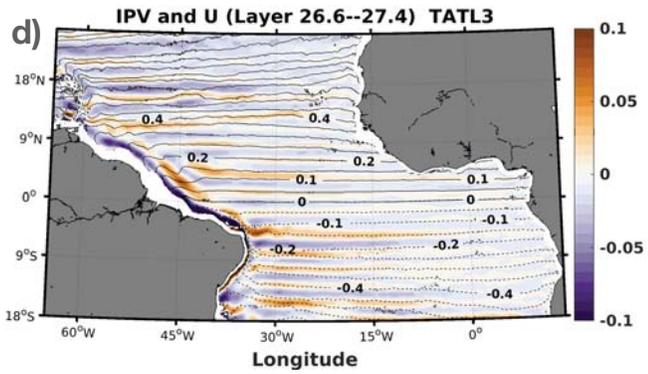
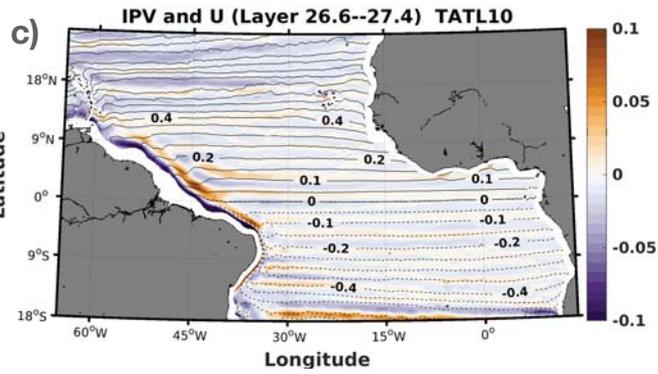
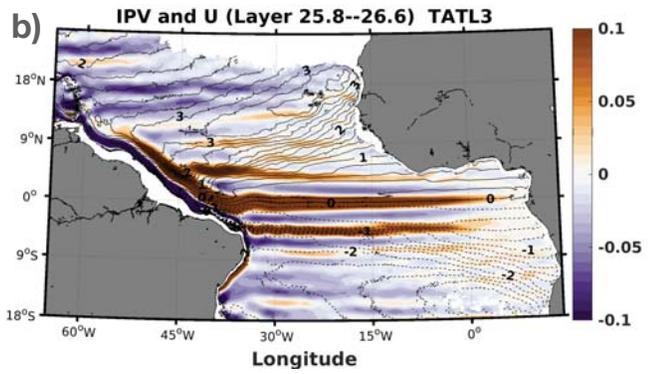
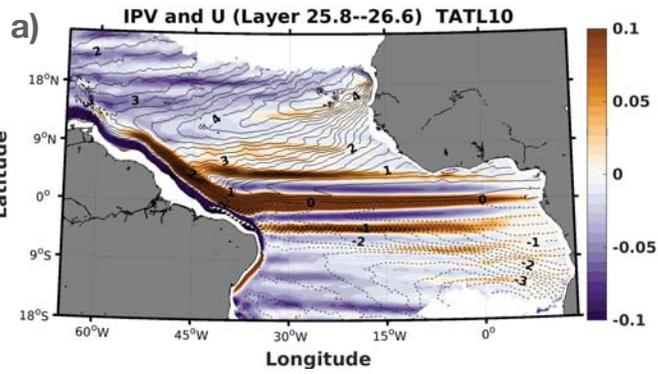


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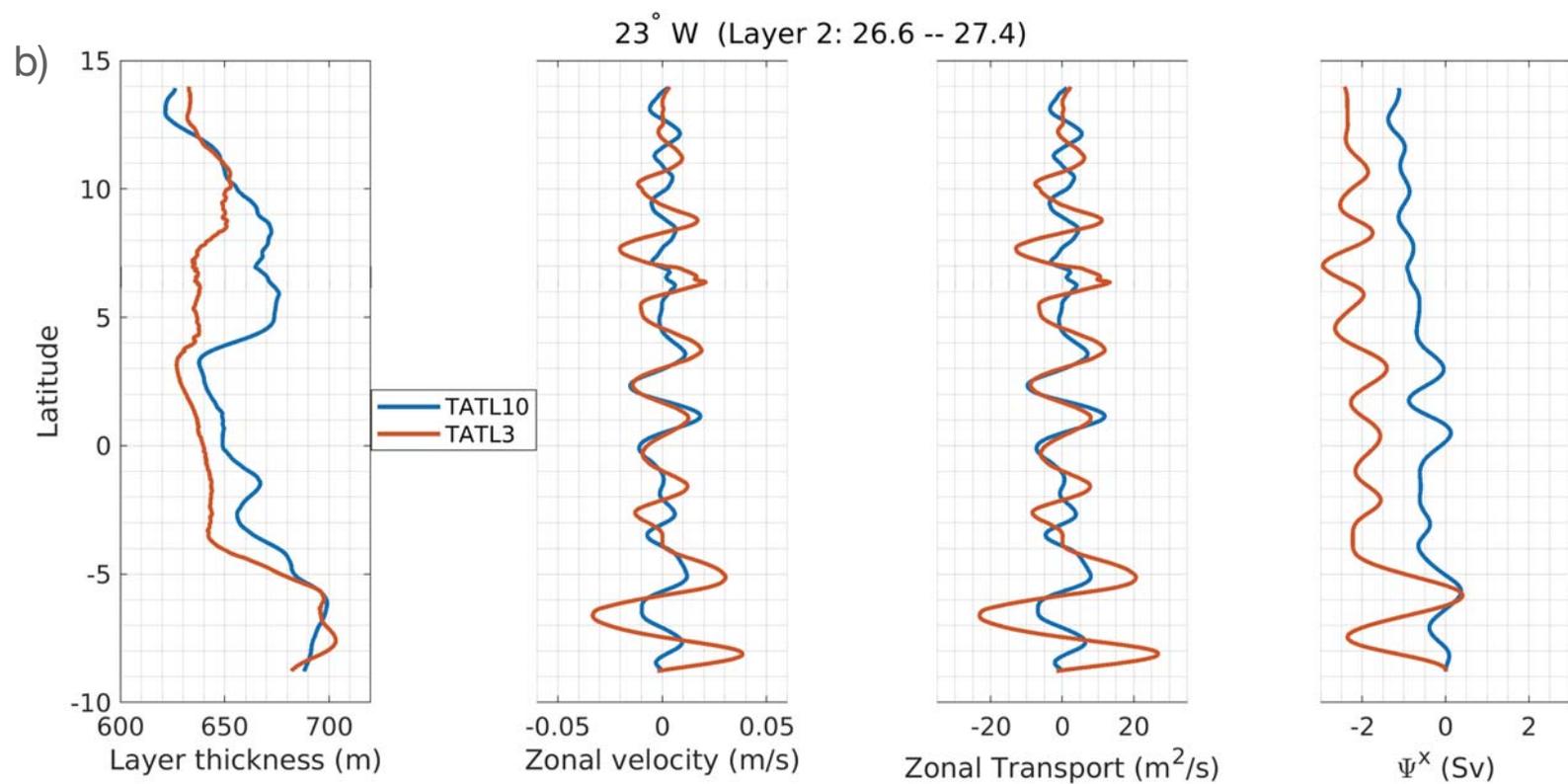
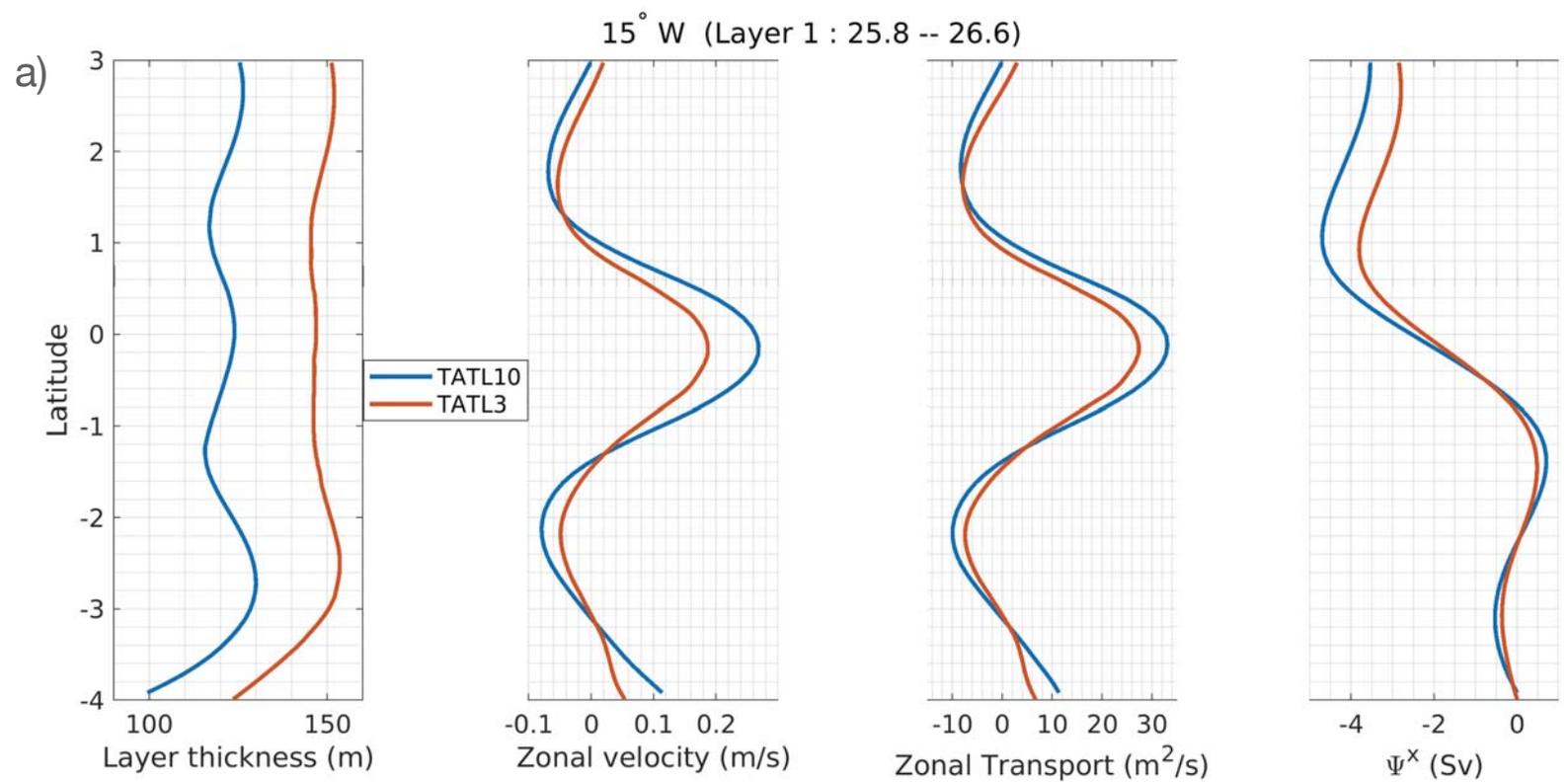


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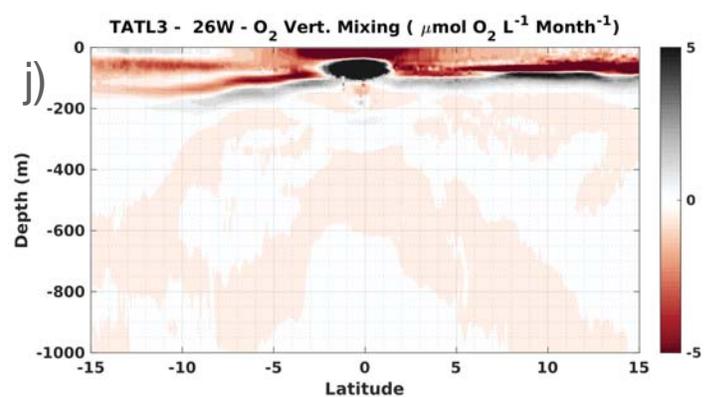
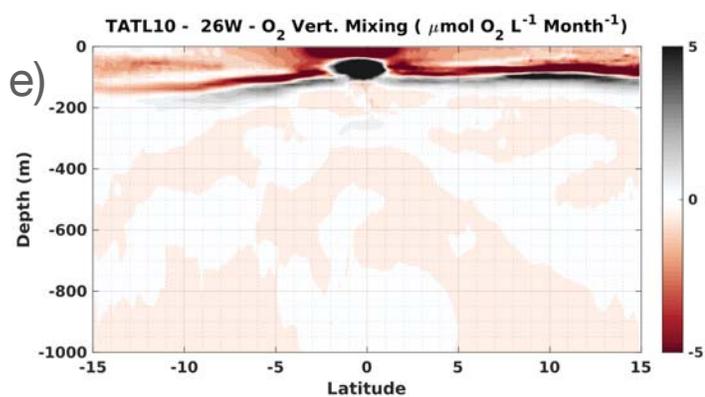
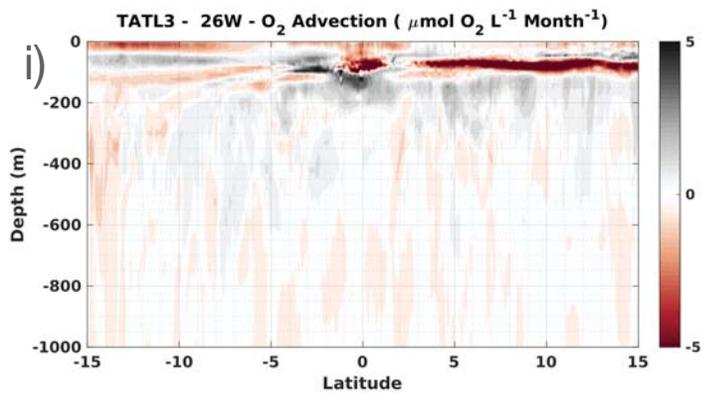
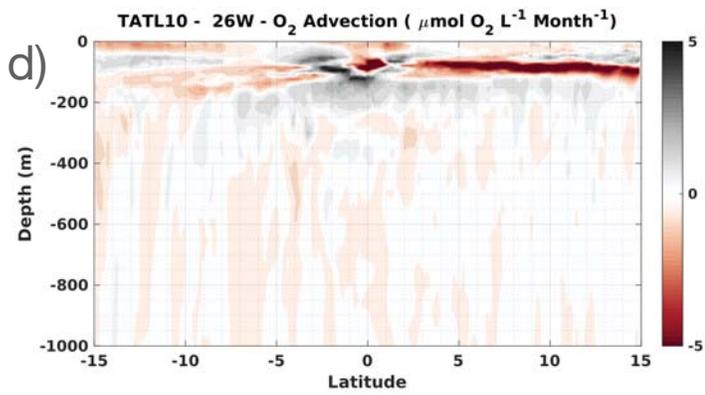
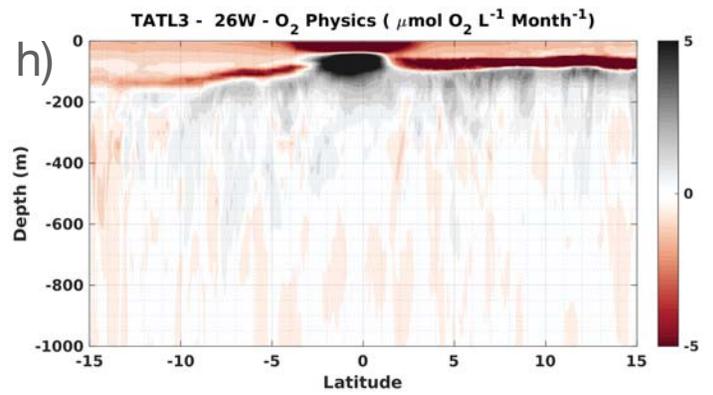
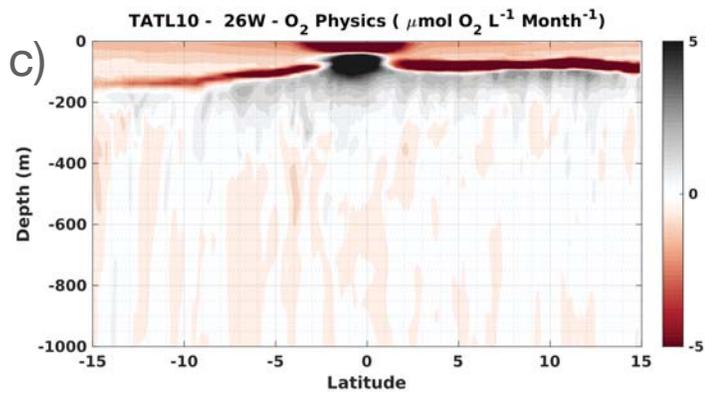
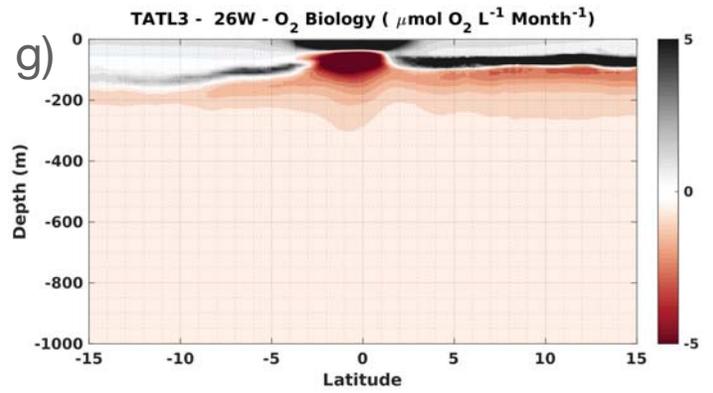
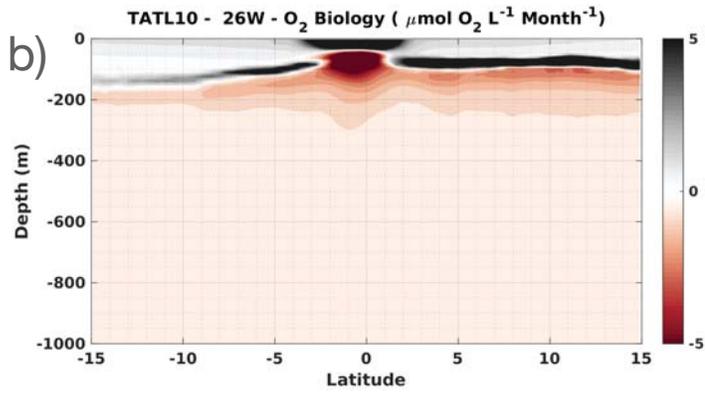
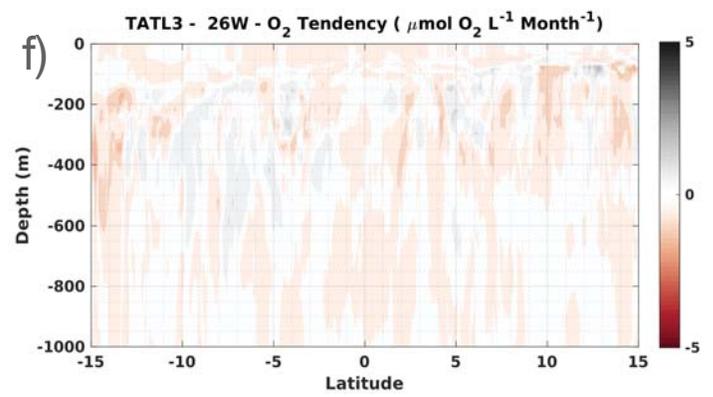
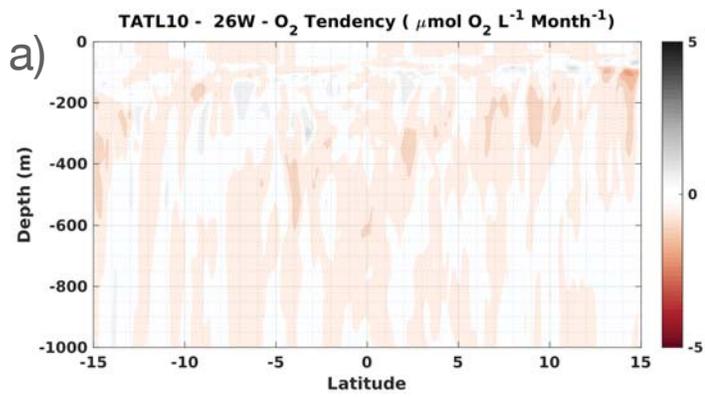


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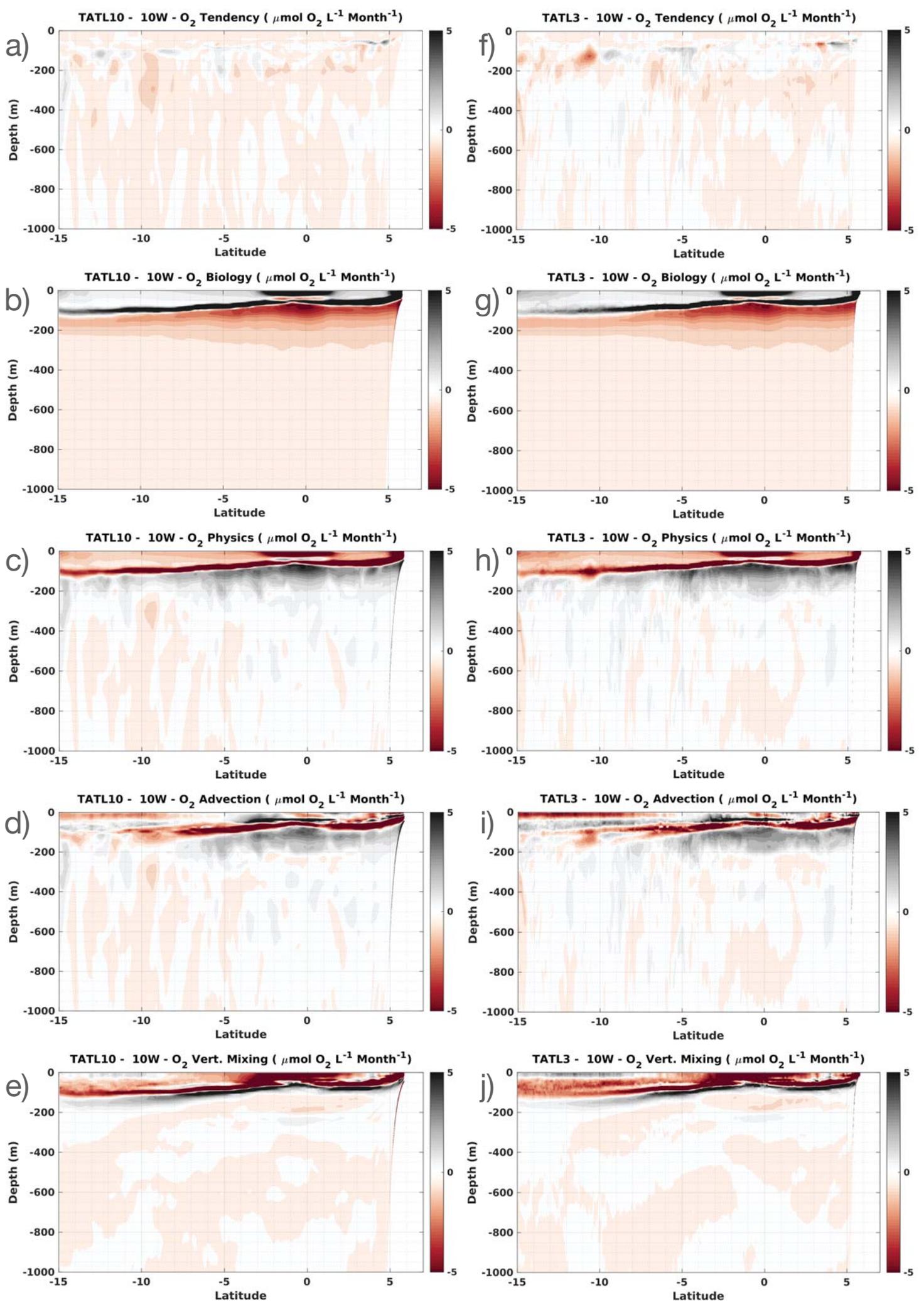
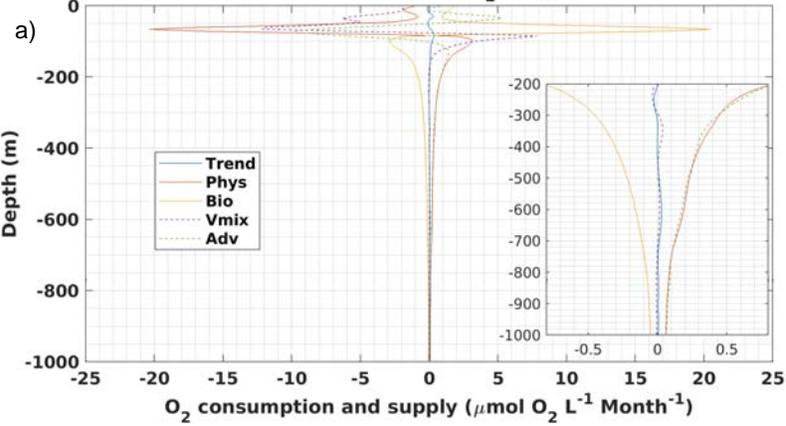
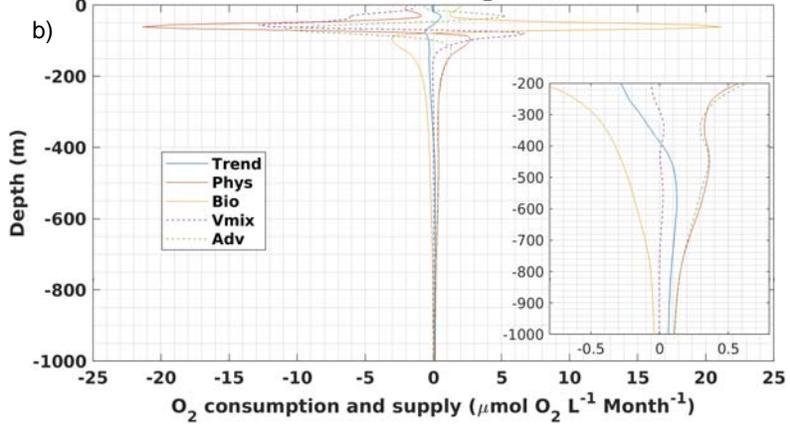


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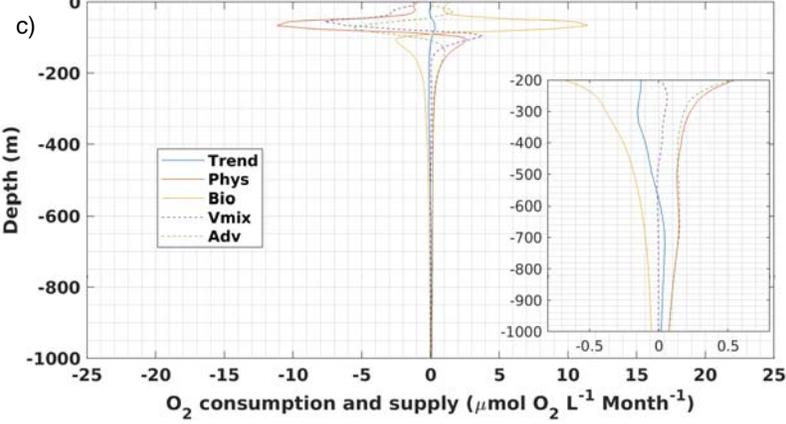
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TATL3 - ETNAd - O₂ Budget



TATL10 - ETSA - O₂ Budget



TATL3 - ETSA - O₂ Budget

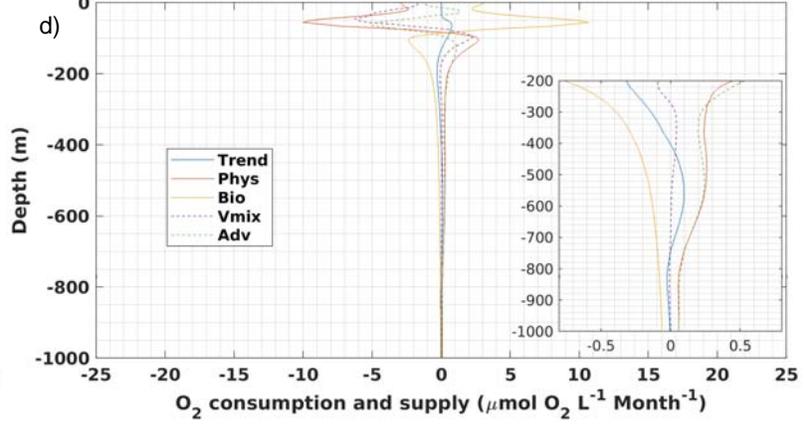
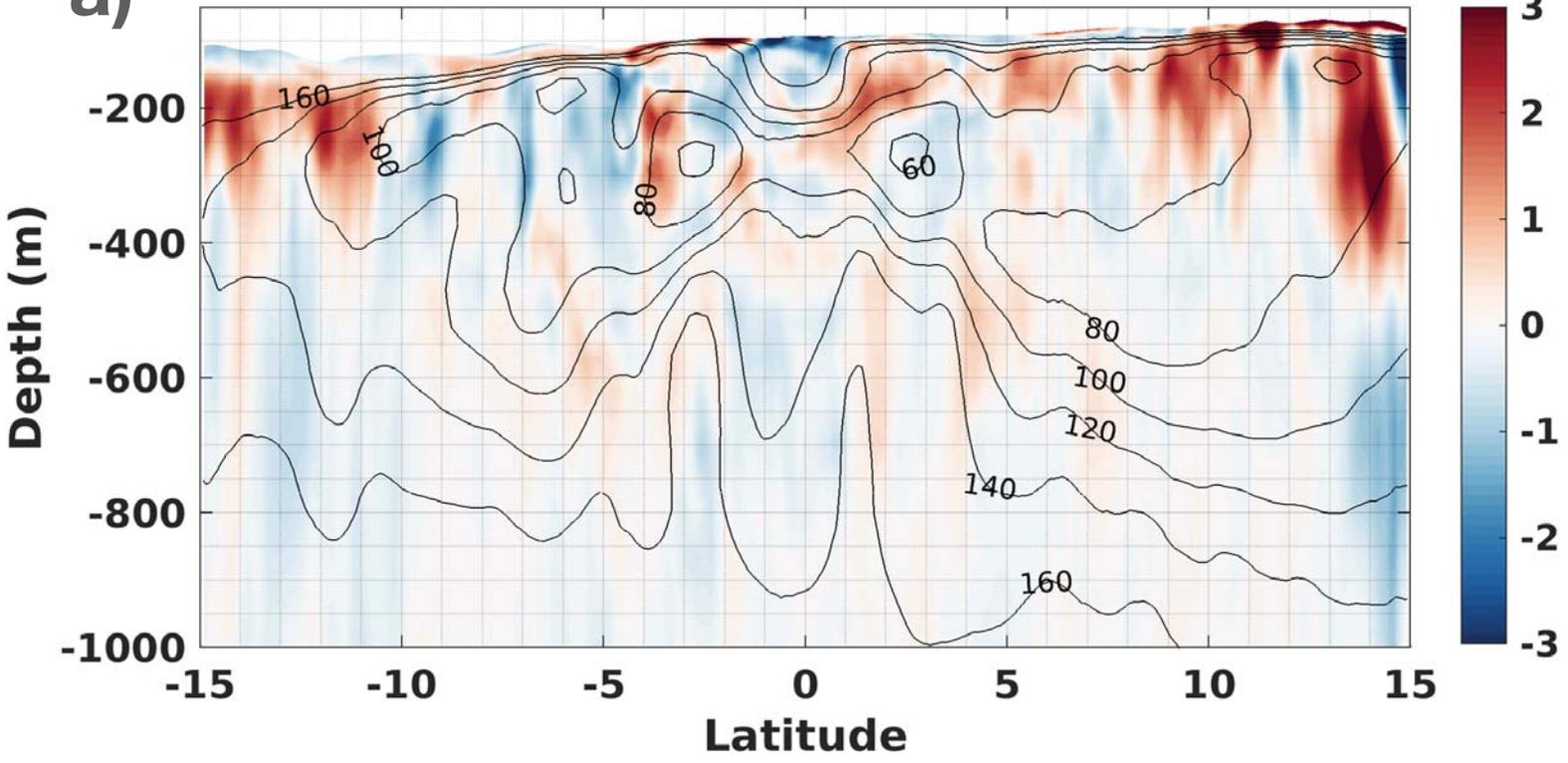


Figure 12.

$\langle V' O'_2 \rangle$ - TATL10 - 23°W ($\mu\text{mol L}^{-1} \text{m year}^{-1}$)

a)



$\langle V' O'_2 \rangle$ - TATL3 - 23°W ($\mu\text{mol L}^{-1} \text{m year}^{-1}$)

b)

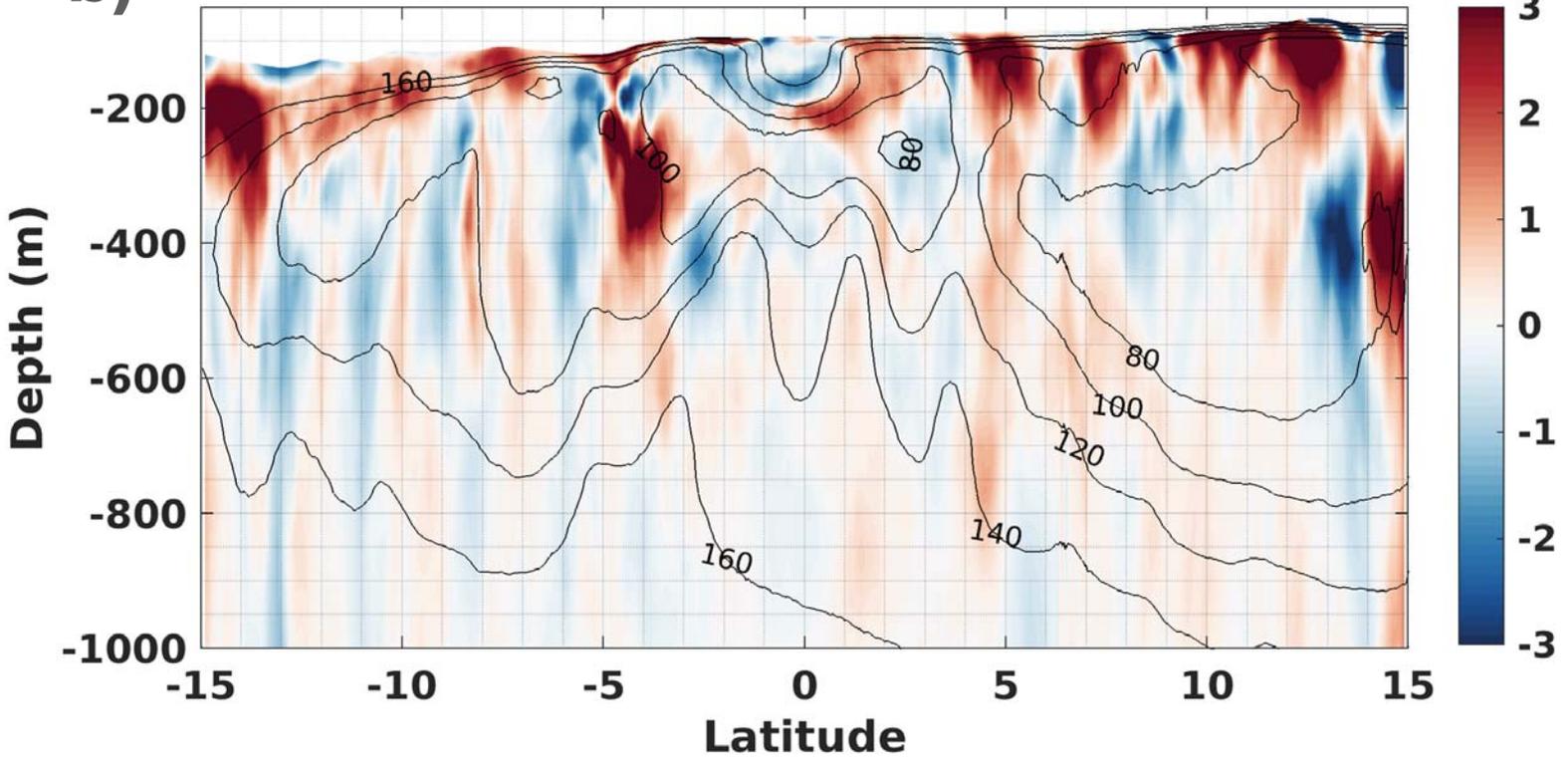


Figure 13.

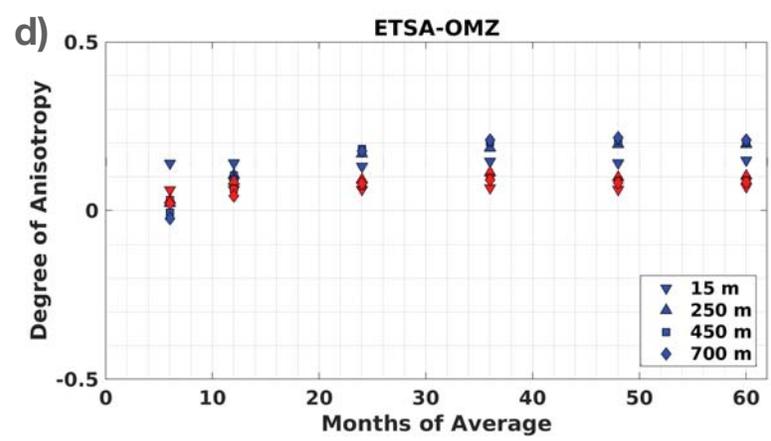
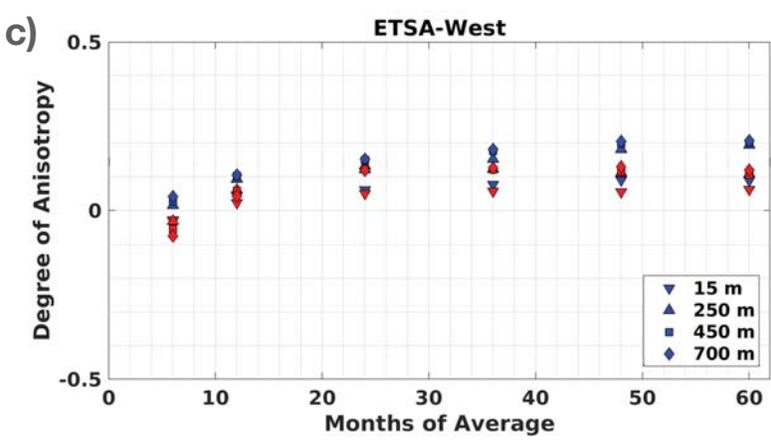
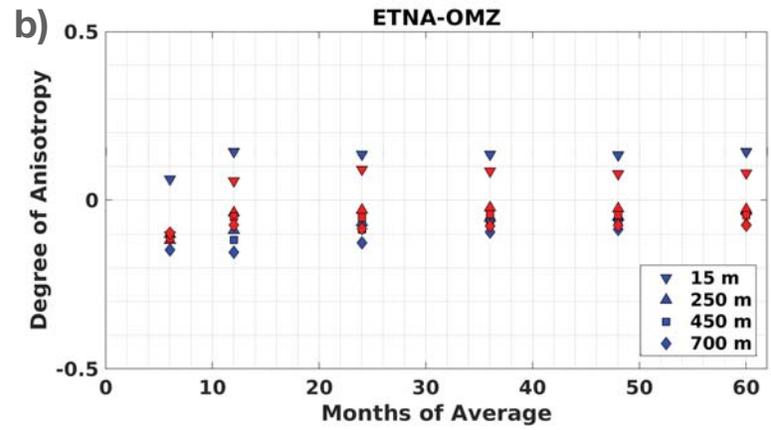
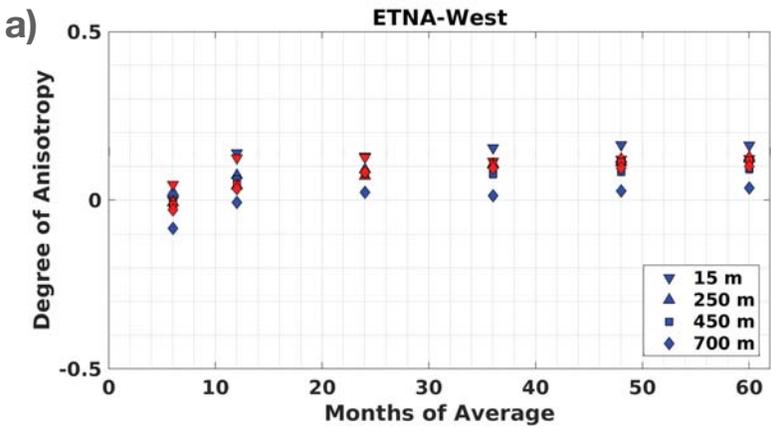


Figure 14.

