Gardening of the Martian Regolith by Diurnal CO2 Frost and the Formation of Slope Streaks

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Abstract

Before dawn on the dustiest regions of Mars, surfaces measured at or below 148 K are common. Thermodynamics principles indicate that these terrains must be associated with the presence of CO2 frost, yet visible wavelength imagery does not display any ice signature. We interpret this systematic absence as an indication of CO2 crystal growth within the surficial regolith, not on top of it, forming hard-to-distinguish intimate mixtures of frost and dust, i.e., dirty frost. This particular ice/regolith relationship unique to the low thermal inertia regions is enabled by the large difference in size between individual dust grains and the peak thermal emission wavelength of any material nearing 148 K (1-2 μ m vs. 18 μ m), allowing radiative loss (and therefore ice formation) to occur deep within the pores of the ground, below several layers of grains. After sunrise, sublimation-driven winds promoted by direct insolation and conduction create an upward drag within the surficial regolith that can be comparable in strength to gravity and friction forces combined. This drag displaces individual grains, possibly preventing their agglomeration, induration, and compaction, and can potentially initiate or sustain downslope mass movement such as slope streaks. If confirmed, this hypothesis introduces a new form of CO2-driven geomorphological activity occurring near the equator on Mars and explains how large units of mobile dust are currently maintained at the surface in an otherwise soil-encrusting world.

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12	Key	Points:
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13	•	Near dawn, diurnal frost is not apparent on cold, dusty, low thermal inertia ter-
14		rains;
15	•	These observations are consistent with a model of dirty diurnal CO ₂ frost, fluff-
16		ing up the surface layer when it sublimes;
17	•	This mechanism could trigger dynamic phenomena on the Martian surface and
18		lead to the formation of slope streaks.

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19 Abstract

20 Before dawn on the dustiest regions of Mars, surfaces measured at or below $\sim 148 \text{ K}$ are common. Thermodynamics principles indicate that these terrains must be associ-21 ated with the presence of CO_2 frost, yet visible wavelength imagery does not display any 22 ice signature. We interpret this systematic absence as an indication of CO_2 crystal growth 23 24 within the surficial regolith, not on top of it, forming hard-to-distinguish intimate mixtures of frost and dust, i.e., dirty frost. This particular ice/regolith relationship unique 25 26 to the low thermal inertia regions is enabled by the large difference in size between in-27 dividual dust grains and the peak thermal emission wavelength of any material nearing 28 148 K (1-2 μ m vs. 18 μ m), allowing radiative loss (and therefore ice formation) to oc-29 cur deep within the pores of the ground, below several layers of grains. After sunrise, sublimation-driven winds promoted by direct insolation and conduction create an up-30 ward drag within the surficial regolith that can be comparable in strength to gravity and 31 32 friction forces combined. This drag displaces individual grains, possibly preventing their 33 agglomeration, inducation, and compaction, and can potentially initiate or sustain downslope mass movement such as slope streaks. If confirmed, this hypothesis introduces a new 34 form of CO₂-driven geomorphological activity occurring near the equator on Mars and 35 36 explains how large units of mobile dust are currently maintained at the surface in an oth-37 erwise soil-encrusting world.

38 Plain Language Summary

39 Surface CO_2 ice forms at all latitudes on Mars, with a strong seasonal control. In this study, we show that diurnal CO_2 ice is not observed in visible wavelength imagery 40 in dusty terrains, where diurnal frost preferentially forms. We interpret this situation 41 as the indication of the presence of hard-to-distinguish dirty frost, where ice crystals grow 42 43 within the surficial regolith, not on top of it, resulting in apparent soil-like dark ice. At 44 sunrise, sublimation-driven winds within the regolith are occasionally strong enough to displace individual dust grains, initiating and sustaining dust avalanches on steep slopes, 45 forming ground features known as slope streaks. This model suggests that the CO_2 frost 46 47 cycle is an active geomorphological agent at all latitudes and not just at high or polar 48 latitudes, and possibly a key factor maintaining mobile dust reservoirs at the surface.

49 1 Introduction

The seasonal transfer of CO_2 between the atmosphere and the surface at high lat-50 itudes is associated with a wide range of exotic processes shaping the surface morphol-51 ogy (Diniega et al., 2021, 2013; Dundas et al., 2012; Hansen et al., 2013; Pilorget & For-52 53 get, 2015; Piqueux et al., 2003; Piqueux & Christensen, 2008), the composition of the atmosphere (Sprague et al., 2004, 2007), impacting the global climate and weather (Haberle, 54 55 1979; Haberle et al., 1979; Hourdin et al., 1993), and possibly degrading the near-surface climate record stored in the polar layered deposits (Kieffer et al., 2006; Kieffer, 2007). 56 In contrast, the existence and ubiquity of a diurnal CO_2 cycle at mid and low-latitudes 57 has only recently been exposed (Piqueux et al., 2016; Khuller et al., 2021a). Its impact 58 59 on the physical state of the regolith, if any, is speculative although Mischna and Piqueux (2020) and Piqueux et al. (2016) proposed that it could take different forms, including 60 61 initiating slope streaks (Schorghofer et al., 2002, 2007), preventing widespread dust induration/duricrust formation, and influencing the global environment over long periods 62 63 of time through the nurturing of a global reservoir of mobile dust able to be lifted in the 64 atmosphere.

65 Piqueux et al. (2016) suggest that the recurring diurnal growth and sublimation 66 of CO_2 crystals could indeed cryoturb the very surficial regolith because frost should be



Figure 1. Schematic model of the CO₂ frost (blue squares) and regolith dust grains (orange circles) relationship in dusty low thermal inertia regions, and its evolution throughout a night and sunrise. Approximate conceptual penetration depths computed with Eq. (1) (without scattering) are indicated on the right and figurated as a shade of grey: $D_{18\mu m}$ peaks at 5 μm and corresponds to the typical distance traveled by photons at a wavelength $\lambda = 18 \ \mu m$ (i.e., peak radiative loss for material approaching T_{CO_2}); D_{12.57µm} peaks near 10 µm and corresponds to the typical distance traveled by photons detected by THEMIS (band 9); $D_{32\mu m}$ peaks near 5 μm and corresponds to the typical distance traveled by photons detected by MCS (channel B1). Frost crystals and regolith grains are approximately 1 μ m in size to conform with various observations (Lemmon et al., 2019; Piqueux et al., 2016) and laboratory work (Presley & Christensen, 1997). In this model, the surficial regolith is free of diurnal frost at the beginning of the night ("Early Night"); when the surface approaches ~ 148 K, the peak radiative loss occurs at $\lambda = 18$ μ m (Wien's law) with photons emerging from ~ 5 μ m depth (see Eq. (1) and associated text), where ice crystals are predicted to form first resulting in dirty frost ("Mid-Night"); radiative loss continues during the night, forming more CO_2 ice, preferentially at 5 μ m depths ("Late Night"), and eventually in shallower and deeper layers ("Pre-Dawn") until the end of the night. At sunrise, CO_2 ice sublimates, creating an upward sublimation-driven wind that could lead to grain displacement.

present in the pores of the regolith, not exclusively on top of it. This model of diurnalice forming within the pores of the regolith is supported by the following two arguments:

69	• Low-latitude CO_2 frost-like surface temperatures have been identified at night in
70	low thermal inertia terrains using Mars Climate Sounder (MCS, (McCleese et al.,
71	2007)) data acquired in the thermal infrared at $\sim 32 \ \mu m$ (Piqueux et al., 2016),
72	and Thermal Emission Imaging System (THEMIS, Christensen et al. (2004)) data
73	acquired at 12.57 μ m (Khuller et al., 2021a). At these wavelengths, the penetra-
74	tion depths D_{λ} [m] (i.e., the distance traveled by photons in the regolith at a given
75	wavelength $\lambda[\mathbf{m}]$ are up to one order of magnitude longer than the typical regolith
76	grain size where frost is observed (1-2 μ m for atmospherically sedimented mate-
77	rial found in these low thermal inertia regions (Lemmon et al., 2019; Presley &
78	Christensen, 1997)). These depths ignore scattering and are estimated using the
79	inverse of k (i.e., the imaginary part of the refractive index of dust given by Wolff
80	et al. (2006)), following Eq. (1) (Hansen, 1997):

$$D_{\lambda} = \frac{\lambda}{4\pi k} \tag{1}$$

At THEMIS and MCS wavelengths, $D_{12.57\mu m}$ and $D_{32\mu m}$ computed with Eq. (1) 81 82 peak at 5 and 10 μ m using $k \approx 0.1$ and $k \approx 0.45$, respectively. These penetra-83 tion depths are up to one order of magnitude longer than the typical regolith grain size where frost is observed (1-2 μ m (Lemmon et al., 2019; Presley & Christensen, 84 85 (1997)). As such, the photons captured by the MCS and THEMIS detectors indi-86 cate that CO_2 frost-like surface temperatures are present at least within the top 87 few microns to tens of microns of the regolith, that is, within a regolith layer that is characterized by several dust grains in thickness. For this reason, these MCS 88 89 and THEMIS observations show that temperatures conducive to diurnal CO_2 ice formation exist within a surface layer that is many regolith grains thick, and not 90 91 exclusively on top of it. Again, this reasoning only holds for very low thermal inertia regions associated with very small and loose regolith grains, and those hap-92 93 pen to be the ones displaying overnight CO_2 ice-like temperatures (Piqueux et al., 94 2016; Khuller et al., 2021a); 95 • Wien's displacement law indicates that the peak infrared emission for $\sim 148 \text{ K}$ surfaces, i.e., the average CO₂ ice temperature on Mars (Leighton & Murray, 1966) 96 97 is ~ 18 μ m. This is the typical wavelength of photons emitted by a surface about to -or already in the process of- forming, or losing, CO_2 ice. From Eq. (1) and us-98 ing $k \approx 0.35$ (Wolff et al., 2006), $D_{18\mu m}$ is in the order of 5 microns when neglect-99 100 ing scattering, that is, many regolith grains in thickness. As such, the surface layer 101 is preferentially cooling/radiating energy from "within" multiple layers of regolith grains, not at the very top as would be the case with larger grains, and the low-102 103 est regolith temperatures may not always be encountered at the very surface, but a few grains deep. This configuration is somewhat analogous to models predict-104 105 ing peak daytime temperatures away from the topmost surface (Henderson & Jakosky, 1994), although the process involved here is different. As mentioned above, these 106 107 photons associated with a wavelength $\lambda \sim 18 \ \mu m$ can freely travel through such thin dust layers, confirming that energy is lost many grains below the surface with-108 out reabsorption on the way out. A word of caution about the length of the op-109 tical paths for thermal infrared photons calculated with Eq. (1) and reported in 110 111this paper: they could correspond to an upper limit as scattering should reduce these distances. A formal calculation of the effect of scattering in the model pro-112 posed in Fig. 1 is difficult because of the complex geometrical relationship between 113 ice crystals and dust grains, requiring advanced radiative modeling not performed 114 115 as part of this work. But because of the significant amount of dust inherent to our 116 dirty frost model, the optical properties of the dust grains are dominant in the resulting ice-dust mixture (Kieffer 2007, Langevin et al. 2006, Fig. 6 from Pilorget 117 et al., 2011), and laboratory experiments have demonstrated that at thermal-infrared 118 wavelengths, photons can penetrate thick (i.e., tens of micrometers) analog/terrestrial 119 dust lavers (Christensen & Harrison, 1993; Christensen et al., 2004; Ramsey & Chris-120 tensen, 1992; Johnson et al., 2002), even larger than those reported in Fig. 1. 121 We conclude that in very low thermal inertia terrains only, CO_2 frost is anticipated 122 123 to form in the pores of the regolith below a few dust grains (or deeper), resulting in mixtures of dust and frost, referred to in this paper as icy regolith or dirty 124 frost. Of course, when the regolith thermal inertia is associated with grains much 125 larger than a few microns in size (thermal inertia > 100 J m⁻² K⁻¹ s^{-1/2}, Presley 126 and Christensen (1997)), individual regolith grains are large enough to contribute 127 128 individually to the surface thermal emission, and ice shall form on the surface of these uppermost grains. 129

Fig. 1 provides a schematic view of this proposed dirty frost model. After the sun sets ("Early Night"), the surface temperature is above T_{CO_2} and no ice is present. Conduction and radiative cooling compete near the surface, but the very low thermal con-

ductivity in the low thermal inertia regions observed near the Equator on Mars results 133 in rapid cooling of the surface. The peak radiative loss occurs at $\lambda \sim 18 \ \mu m$ and is as-134 sociated with photons typically emitted between from 5 μ m in depth (see discussion above). 135 When the temperature reaches T_{CO_2} , the first crystals form at these depths in the pores 136 of the regolith ("Mid-Night"). Radiative loss continues, but the minimum kinetic tem-137 perature does not drop below T_{CO_2} . Instead, more CO₂ ice forms in the pores, still pref-138 erentially near 5 μ m in depth ("Late Night"). Thermal infrared spectroscopy indicates 139 that the CO₂ crystals are $\sim 1 \ \mu m$ in size (Piqueux et al., 2016). Cooling through con-140 141 duction and radiation bleeds to other depths, leading to the formation of CO_2 ice in the pores at other depths ("Pre-Dawn", with dawn loosely defined in this model as the late 142 143 night time period, before the sun rises, when some light illuminates the surface and allows visible wavelength imagery acquisition with THEMIS showing surface features). En-144 145 ergy and mass balance modeling shows that CO_2 ice typically condenses equivalent thicknesses of a few tens of microns in the form of micrometer size ice crystals (Piqueux et 146 al., 2016), and the very porous substrate associated with low inertia terrains (Presley 147 & Christensen, 1997) implies that the ice to dust ratio should be measured in %, if not 148 149 10's of %, but certainly not at the contamination level. One consequence is that the vis-150 ible wavelength albedo of this thin dusty ice unit should be that of the bare dust (Singh 151 & Flanner, 2016; Warren et al., 1990).

152 This model predicts that surface temperatures associated at surface with T_{CO_2} in 153 low thermal inertia terrains should not display any signature of ice in imagery of the ground, in contrast with other terrains at T_{CO_2} but associated with higher thermal inertia val-154 155 ues. Consequently, we hypothesize that diurnal CO₂ frost associated with dusty low thermal inertia surfaces should be exceedingly difficult to identify in visible wavelength im-156 157 agery, in contrast with diurnal or seasonal CO_2 frost present elsewhere on Mars in ter-158 rains where the surface material consists on larger regolith grains. Note that this model is not predicted to apply to the polar regions where low thermal inertia deposits are not 159 observed. In this paper, we present an analysis of THEMIS data acquired near sunrise 160 at both visible and thermal infrared wavelengths in order to constrain the nature of the 161 frost/regolith relationship on dusty terrains on Mars and test the dirty frost model (Fig. 162 163 1).

In addition, one potential implication for this recurring diurnal growth and sub-164 165 limition of CO_2 frost crystals within low thermal inertia terrains may include regular overnight surficial mechanical disruption of the soil, and possible fluffing by vertical sublimation-166 driven winds. This process could maintain high regolith porosity, and prevent compaction 167 168 or inter-grain induration/cementation that seems to be ubiquitous elsewhere on Mars (Jakosky & Christensen, 1986; Mellon et al., 2000) except in low inertia terrains (Piqueux 169 170 & Christensen, 2009a, 2009b; Putzig et al., 2005). Such recurring mechanical alteration 171 of the regolith has been proposed as a process maintaining mobile dust available for lifting (Piqueux et al., 2016) contributing to impact the global climate over long periods 172 173 of time (Mischna & Piqueux, 2020). Similar conclusions can been drawn when translucent ice forms on (or within) the Martian regolith: at sunrise, the solar energy is deposited 174 at the base of the transparent frost layer, at the interface with the regolith grains, not 175 176 at the very atmosphere/ground interface (see the abundant literature on the topic for 177 the polar regions (Diniega et al., 2013, 2021; Hansen et al., 2010, 2013; Pilorget et al., 2013; Pilorget & Forget, 2015; Pilorget et al., 2011; Piqueux et al., 2003; Piqueux & Chris-178 tensen, 2008; Piqueux et al., 2016; Pommerol et al., 2011; Pommerol et al., 2013; Portyank-179 ina et al., 2010, 2012; Thomas et al., 2010)). Basal sublimation yields winds internal to 180 the very surficial regolith, and has the potential to disturb the upper regolith. 181

182 Another potential implication of this regolith gardening model that sets individ183 ual grains in motion at sunrise could be the initiation of dynamic surface mechanisms
184 leading to the formation of slope streaks. Slope streaks are dark wedge-shaped surface
185 features on sloped terrains, associated with downslope mass movement (Chuang et al.,

2007; Ferris et al., 2002; Kreslavsky & Head, 2009; Miyamoto et al., 2004; Phillips et al., 186 2007; Schorghofer et al., 2002, 2007; Sullivan et al., 2001) that form exclusively in dusty 187 188 low thermal inertia terrains. The initiating and sustaining mechanisms of slope streaks has been attributed to wet (Kreslavsky & Head, 2009; Mushkin et al., 2010; Schorghofer 189 et al., 2002) or dry processes (Baratoux et al., 2006; Burleigh et al., 2012; Chuang et al., 190 2007, 2010; Dundas, 2020; Phillips et al., 2007; Schorghofer et al., 2002; Sullivan et al., 191 192 2001), without conclusive evidence one way or the other (Bhardwaj et al., 2019; Dun-193 das, 2020). The longer-than-expected character of slope streaks on shallow Martian slopes 194 requires some form of lubricating agent that has led Piqueux et al. (2016) to propose a connection with diurnal CO₂ ice sublimation grain motions, and slope destabilization. 195 196 We hypothesize the existence of a spatial and temporal relationship between the pres-197 ence of diurnal frost and the formation of slope streaks. The temporal relationship has 198 been explored by Heyer et al. (2019) and is inconclusive, mainly because of the difficulty to determine the time of formation of slope streaks on images acquired infrequently from 199 200 orbit. In this paper, we present an additional analysis of the spatial relationship between 201 slope streaks distribution and diurnal frost presence.

The existence of a diurnal mechanical cycle within the surficial regolith associated with the growth and sublimation of CO₂ ice crystals would have significant geomorphological implications, including for equatorial terrains. The work presented in this paper aims at characterizing this relationship through visible and infrared wavelength imagery analysis, surface features mapping, and numerical modeling.

207 2 Methods

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2.1 CO₂ Frost Optical Properties in Dusty Terrains

209The optical properties of CO_2 ice (albedo and emissivity) have been extensively210studied (Hansen, 1997; Schmitt et al., 1998) because of their control on local and global211processes, from the condensation of the atmosphere (Forget et al., 1995; Forget & Pol-212lack, 1996; Forget et al., 1998; Colaprete & Toon, 2002; Hayne et al., 2012, 2014) to the213basal sublimation and venting of the caps (Piqueux et al., 2003; Kieffer et al., 2006; Pi-214lorget et al., 2011).

215 Piqueux et al. (2016) demonstrated that thin layers (i.e., a few microns to tens of 216 microns in thickness) of micrometer size ice crystals (as opposed to translucent slab-like ice) best match the high emissivities derived from MCS data in the low thermal iner-217 218 tia dusty units. Their conclusion is further supported by energy and mass balance mod-219 eling results suggesting that condensation involves up to a few hundreds of microns of 220 CO_2 ice in thickness, much too small to form large translucent and high emissivity crys-221 tals (i.e., centimeters in sizes, (Kieffer, 2001, 2007; Kieffer et al., 2000; Langevin et al., 222 2006).

223 Here, we add that the high emissivities derived from the MCS observations at 32 μ m 224 are also consistent with frost forming within the uppermost pores of the regolith, and not just on top of it as these authors proposed. The surface emissivity is highly sensi-225 226 tive to the presence of frost at $\lambda \sim 12.57 \ \mu m$ (THEMIS band 9, (Christensen et al., 2004)), 227 even when it forms thin layers: ice thicknesses in the order of tens and hundreds of mi-228 crons should result in a noticeable decrease of the surface emissivity, from ~ 0.96 to ~ 0.9 -229 0.6 (Piqueux et al., 2016). While this reduction of the emissivity is magnified by the high 230 emission angles considered for the MCS observations in the work present by Piqueux et al. (2016), (i.e., 70° vs. 0° for THEMIS), we can still reasonably assume that frost form-231 232 ing on the top of the regolith should measurably lower the emissivity compared to an 233 ice-free surface. In contrast, frost forming within the pores, in other words highly dust-234 contaminated ice, should lead to surface emissivity approximately equal to the emissiv-235 ity of ice-free dusty terrain (~ 0.96 , (Ruff & Christensen, 2002; Piqueux et al., 2016)).

In this paper we present an analysis of the surface emissivity to further constrain the frost/regolithrelationship in the low thermal inertia regions of Mars.

Micrometer size CO_2 ice crystals are associated with high albedo values at visible 238 wavelengths (Hansen, 1997; Kieffer et al., 2000; Titus et al., 2008; Appéré et al., 2011; 239 Singh & Flanner, 2016), but dust contamination can drastically reduce it (Kieffer et al., 240 241 2000; Murchie et al., 2019; Singh & Flanner, 2016; Warren et al., 1990). For this reason, 242 the albedo (or color) of frosted terrains provides an excellent diagnostic tool of the re-243 lationship between the ice and the regolith grains, as intimate mixtures of ice and dust 244 will be associated with dust-like albedo values, whereas clean ice layers on top of the sur-245 face dust will be characterized by high albedo values. As a word of caution, the refer-246 ences mentioned hereabove generally focus on dust contamination with large ice crystals, typically hundreds of microns in size or more, relevant to the polar regions (Kieffer 247 et al., 2000). This large difference with the ice crystal sizes we consider in this work may 248 not be inconsequential as Kieffer et al. (2000) report that the impact of dust contam-249 250 ination on albedo for ice crystals smaller than 10 μ m can be more limited. We can expect that even with significant dust mixed in, very small and bright ice crystals placed 251 252 on the top of the dust grains still display high albedo values (0.7-0.9) compared to ice-253 free ground (i.e., 0.27, Putzig et al. (2005)). Furthermore, Singh and Flanner (2016) show 254 that the smaller the CO_2 ice crystals, the smaller the thickness of frost needed to con-255 trast with a frost-free surface. Even for the relatively thin diurnal frost layers anticipated 256 in the Martian mid latitudes of just a few tens of microns, an extrapolation of Singh and Flanner (2016)'s results (their Fig. 6) suggest high albedo values for the very small cry-257 258 tal sizes expected here (i.e., $\sim 1 \ \mu m$ based on Piqueux et al. (2016)).

For all these reasons, we conclude that CO_2 frost layers 10's to 100's μ m in thick-259 260 ness composed of clean $\sim 1 \ \mu m$ crystals forming on the top of dusty surfaces are expected 261 to yield a signature at visible wavelengths with a spectral slope from blue to red (Murchie 262 et al., 2019), that is, blue(ish) pixels, similar to configurations observed on the seasonal caps (Calvin et al., 2015, 2017), and low emissivities at 12.57 μ m. In contrast, CO₂ frost 263 264 layers 10's to 100's μ m in thickness mixed within the pores of the dusty regolith in the 265 low thermal inertia terrains are expected to remain free of albedo contrast at visible wave-266 lengths compared to unfrosted terrains and associated with high dust-like emissivity at $12.57 \ \mu m$. In this paper, we leverage this difference of behavior to characterize the re-267 268 lationship between diurnal CO_2 frost and surface dust and test whether this frost forms on the regolith or within the pores. 269

270 2.2 Dataset and Frost Identification

271 To identify the presence of CO_2 ice and characterize its optical properties in dusty 272 terrains, we use coincident visible and temperature observations acquired by THEMIS, a multispectral visible and thermal infrared wavelength imager observing the surface and 273 274 atmosphere of Mars at various local times depending on the phase of the mission. THEMIS visible wavelength images form a smaller footprint than the thermal infrared images (~ 18 275 vs. ~ 32 km wide swath on the ground), so the infrared data is cropped to the extent 276 277 of the overlapping visible wavelength data where the analysis is carried out. Henceforth, 278 we do not consider or discuss THEMIS thermal infrared data where simultaneous vis-279 ible wavelength observations are not available.

We only analyze data acquired near sunrise, when adequate lighting allows the acquisition of multiband visible wavelength imagery of the ground, and when diurnal frost is not only expected to be present, but near peak thickness, up to several hundreds of microns (Piqueux et al., 2016). We select data acquired at high incidence angles *i*, e.g., $70^{\circ} < i < 110^{\circ}$ (with $i < 90^{\circ}$ indicating the Sun above the horizon, $i = 90^{\circ}$ the Sun at the horizon, and $i > 90^{\circ}$ the Sun below the horizon). Such data were thus acquired typically between 6 - 8 A.M. and 5 - 6 P.M. (Local True Solar Time -LTST-). Obser-



Figure 2. Spatial and seasonal distribution of THEMIS visible/thermal infrared pairs acquired when the sun incidence angle was between 70° and 110° . The colored background indicates the seasonality of direct solar illumination (top of the atmosphere) as calculated by a Keplerian orbital model used by KRC (Kieffer, 2013). An evident data distribution bias caused by operational, local weather, and illumination constraints exist across the dataset. As a result, surface coverage in the northern high/mid latitudes is more sparse than at other seasons/locations.

vations that are acquired during the afternoon are not used because not associated with
diurnal CO₂ ice. This selection criterion yields a hemispheric bias (Fig. 2), as far fewer
visible wavelength images at dawn meeting these criteria were acquired in the Northern
hemisphere compared to the Southern hemisphere. Nonetheless, global sampling at all
seasons is available and adequate for this work.

Surface temperatures are derived from THEMIS band 9 centered at 12.57 μm as 292 293 they offer the best signal on cold surfaces (noise equivalent delta temperature of 1 K at 294 180 K (Christensen et al., 2004), estimated to 3 K at 150 K (Pilorget et al., 2013)), and 295 because they have frequently been used for polar studies (e.g., (Kieffer et al., 2006; Pilorget et al., 2013; Piqueux & Christensen, 2008; Piqueux et al., 2008)). We eliminate 296 297 image pairs clearly impacted by calibration issues (e.g., (Edwards et al., 2011)), poor ob-298 servation conditions (mainly the identifiable presence of clouds in visible wavelength imagery), or defined in the infrared by the coldest scene on the image smaller than 0.9 times 299 300 the local CO_2 frost point T_{CO_2} calculated as follow:

$$T_{CO_2} = \frac{\beta}{\gamma - \ln(P)} \tag{2}$$

with $\gamma = 23.3494$ [1], $\beta = 3182.48$ [K] (James et al., 1992), and P, the local CO₂ partial pressure taken as 0.96 (Mahaffy et al., 2013) of the total surface pressure derived from the local topography and parametrized surface pressure observations (Withers, 2012). This selection criterion stems from the fact that on Mars, atmospheric CO₂ is in quasi-

305 unlimited supply for overnight condensation at the surface and therefore any surface at

 T_{CO_2} must be associated with the presence of carbon dioxide ice. Such selection is based 306 307 on the assumption that no surface can be colder than the kinetic temperature of CO_2 , 308 and CO_2 ice of large grains has a high emissivity near 12.57 μ m (Hayne et al., 2012). However, deposits composed of ice crystals with size of nearly 1 μ m should lead to low emis-309 310 sivity surfaces, even at 12.57 μ m (Piqueux et al., 2016), and could thus induce $T_{surf} < 0.9 \times T_{CO_2}$. Nevertheless, observations in the polar regions have demonstrated that such occurrences 311 312 are rare (Forget et al., 1995; Forget & Pollack, 1996; Forget et al., 1998; Colaprete & Toon, 2002; Hayne et al., 2012, 2014) and generally associated with the presence of snow falls 313 314 that are not expected near the equator.

315 To account for the instrument noise and possible atmospheric contributions that 316 may approach a few K when looking at cold surfaces, we assign a 5 K tolerance for the identification of surface CO₂ ice. In other words, CO₂ is considered present on any sur-317 face where the local temperature is within 5 K of the predicted frost point based on the 318 319 local atmospheric pressure and Eq. (2). This is a conservative approach based on work by Pilorget et al. (2013), who found that a 7 K margin is reasonable to identify CO_2 ice. 320 321 Out of the 32,236 THEMIS images pairs inspected, 3,258 (i.e., $\sim 10\%$) are associated with temperatures requiring the presence of CO_2 ice on the ground. 322

Once an image pair is flagged for CO_2 ice based on the surface temperature, we 323 inspect the associated visible wavelength image to identify surface frost patches. This 324 325 characterization is performed using the THEMIS public viewer at viewer.mars.asu.edu/ 326 viewer/themis, with colorized THEMIS visible wavelength images, as surface frost seems 327 to be indiscernible in greyscale images alone at high solar incidence angles, without fur-328 ther processing. The THEMIS visible camera has a resolution of 18 m/pixel and has five filters with band centers located at 425 (band 1), 540 (band 2), 654 (band 3), 749 (band 329 330 4), and 860 nm (band 5) (Christensen et al., 2004). When available, we used "R2B" im-331 ages, a colorized product where of band 4 (red) and band 1 (blue) are combined using 332 $0.65 \times$ band $1 + 0.35 \times$ band 4 to generate a simulated green band used for the RGB composite (Bennett et al., 2018; Murray et al., 2016). When not available, we used con-333 334 ventional RGB composite resulting from band 4 (or, if not available, band 3), band 2, 335 and band 1 in the blue channel (Bennett et al., 2018; Murray et al., 2016). Surface frosts 336 are identified based on their blue/white hues, in stark contrast with the orange/brown/grey surrounding terrains, in spite of the challenging illumination conditions (Fig. 3). The 337 338 vast majority of these blue/white units are confidently attributed to the surface (as opposed to the atmosphere) based on their sharp boundaries following morphometric or 339 340 color units, topography, or preferential slope orientation. Some image pairs with blue/white 341 patches are disqualified because they show evident calibration issues most likely due to 342 the challenging illuminations conditions encountered at the terminator, where THEMIS 343 was not originally designed to operate. Such images show unrealistic large bright color 344 patches often associated with ghosts or individual wavelets, and/or no surface pattern. 345 Another subset of images displays blue/white signatures with blurry boundaries, pos-346 sibly indicating that hazes or clouds might be present and are not further considered for this work. Among the 3,258 image pairs associated with temperatures consistent with 347 CO_2 ice, 2,761 images (i.e., ~ 85%) show no calibration issue and contain potential sur-348 349 face frost signatures in visible wavelength. A confidence level for surface frost identifi-350 cation upon visual inspection is defined by assigning each image to one of five classes de-351 fined as follows:

<sup>Class 1 (Fig. 3a, 115 images that account for 6% of visible wavelength detections):
Excellent contrast between blue/white surface patches and nearby units charac</sup>terized by very different colors (generally orange, brown, grey), adequate illumination conditions, geomorphological surface features (i.e., craters, scarps, plains, hills, etc.) perfectly recognizable. In this class, images display no perceivable calibration artifacts, and the spatial distribution of these blue/white surface patches

358	generally shows some spatial coherence, favoring pole facing slopes and flat sur-
359	faces. Class 1 image pairs are mainly located between 50° S and 30° S (Fig. S1);
360	• Class 2 (Fig. 3b, 534 images that represent 28% of the visible wavelength detec-
361	tions): Good color contrast between the blue/white surface units and the rest of
362	the scene, and illumination conditions adequate for geological mapping but not
363	ideal, resulting in harder-to-detect blue/white units, but still perfectly distinguish-
364	able. Slope patches are generally still fairly easily identified, but blue/white units
365	on flat surfaces are more challenging. Some images display ghosts or under-saturation
366	that sometimes results in an artificial enhancement of purple hues, decorrelated
367	from surface morphology. Class 2 image pairs are preferentially found at mid and
368	tropical latitudes (Fig. S1);
369	• Class 3 (Fig. 3c, 708 images that represent 37% of the visible wavelength detec-
370	tions): Fair contrast between blue or white surface units and the surrounding back-
371	ground, patches difficult to detect. However, high coherence with specific slopes
372	and slope azimuth provides higher frost identification confidence. Image pairs in
373	this class are either 1) of high intrinsic quality (contrast, saturation), but the blue/white
374	units do not particularly stand out relative to the surrounding terrains as in Class
375	1 and 2, or 2) difficult to identify because of low contrast, obvious color artifacts,
376	ghosts, and also because of the potential presence of near-surface hazes. Class 3
377	images have been detected at mid and high latitudes (Fig. S1);
378	• Class 4 (Fig. 3d, 556 images that account for 29% of the visible wavelength de-
379	tections): Poor contrast between discreet blue/white units and the nearby terrains,

- resulting in speculative identifications, sometimes because of overall low image quality due to the extremely challenging illumination conditions encountered at the terminator, sometimes because of the subtle color contrast with the regional terrains. In Class 4, the size, shape and sometimes distribution of white/blue patches within the images allow us to confidently exclude hazes, or color processing artifacts. Class 4 images are more generally found at mid latitudes, between 25° and 30°S or polar latitudes (Fig. S1);
- Class 5 (Fig. 3e, 848 images): No identification of blue or white surface sugges tive of surface frost despite surface temperatures consistent with the presence of
 CO₂ ice.

390 This classification is designed with the underlying assumption that the blue/white 391 patches observed invisible wavelength images are indeed due to surface ice given the pres-392 ence of CO_2 ice-like surface temperatures. This assumption is reasonable given our ex-393 perience with mid-afternoon THEMIS visible wavelength imagery: high afternoon sur-394 faces temperatures inconsistent with the presence of frost do not display patches such 395 as those presented in Fig. 3.

396 3 Results

397 The spatial distribution of surface units at the CO_2 frost point where visible wave-398 length imagery is also available and without indications of calibration issues near dawn 399 is shown in Fig. 4. THEMIS surface temperatures acquired at sunrise are consistent with the presence of CO_2 frost at virtually all latitudes (black and white dots in Fig. 4), con-400 401 firming results presented by others with MCS (Piqueux et al., 2016) and THEMIS (Khuller 402 et al., 2021a) data. THEMIS images presenting surface temperature close to T_{CO_2} clus-403 ter at the seasonal polar caps, at high latitudes in the North (up to 60° N) and middle 404 latitudes in the South (down to 20°S), and the dusty terrains with low thermal inertia 405 of Tharsis, Elysium, and Arabia Terra. No thermal images show temperatures close to 406 T_{CO_2} at low latitudes in areas with high thermal inertia because these units have sur-407 faces that are too warm at night to allow CO_2 condensation. Fewer occurrences of CO_2 408 frost are reported in Fig. 4 compared to the mapping results in Khuller et al. (2021a) 409 because these authors did not use a visible wavelength image presence requirement, which



Figure 3. Examples of THEMIS visible wavelength (a-e) and corresponding thermal infrared (f-j) images acquired simultaneously near dawn. The blue/white surface patches (a-d) and low surface temperatures within 5 K of T_{CO_2} in infrared images (f-j) are interpreted as frost (Classes 1-4, see text for definition). Class 5 (no frost signature) despite surface temperatures consistent with CO₂ frost is illustrated with (e) and (j). Coordinates, solar longitude (Ls), and local true solar time (LTST) are given in the different panels. Red arrows emphasize hard-todistinguish blue/white patches. White arrow point to the position of the sun in the sky. a: Class 1, V71796004; b: Class 2, V63705007; c: Class 3, V78900003; d: Class 4, V63305007; e: Class 5, V76958011; f: I71796003 associated with a; g: I63705006 associated with b; h: I78900002 associated with c; i: I63305006 associated with d. j: I76958010 associated with e (Christensen et al., 2002). f-j underlain with a THEMIS daytime IR mosaic to enhance topography (Edwards et al., 2011). Some terrains appear black in the thermal infrared images because of the background mosaic (not because of an absence of measurement).

410 drastically limits the number of usable infrared observations, and they applied a looser 411 CO_2 ice temperature identification criterion (e.g., 7 K vs. 5 K tolerance on T_{CO_2}).

In the Southern hemisphere, blue/white patches at the CO₂ frost point are observed
on visible wavelength imagery at latitudes as low as ~ 20°S, noticeably closer to the equator than a previous analysis of bright patches in THEMIS visible wavelength imagery,
e.g., 33°S (Schorghofer & Edgett, 2006), and ~ 35°S based on Observatoire pour la Minéralogie, l'Eau, les Glaces et l'Activité (OMEGA, Bibring et al. (2004)) spectral identifications (Vincendon et al., 2010). This difference confirms the high sensitivity of our sur-

face frost identification approach. A strong dichotomy appears between the North and 418 419 the South, resulting from several factors. First, the southern winter lasts longer, favor-420 ing the accumulation of more ice at lower latitudes. Second, far fewer images have been acquired in the North than the South during the winter, for operational reasons. Third, 421 422 the southern hemisphere is much older and therefore rougher than the northern hemisphere (Aharonson et al., 2001; Kreslavsky & Head, 1999; Kreslavsky & Head, 2000), with 423 significantly more occurrences of pole facing terrains there where seasonal ice can sur-424 vive longer, thanks to reduced direct insolation. Finally, the clarity of the atmosphere 425 426 tends to be higher through the Southern polar vortex compared to the Northern one. In the North, the seasonal coverage after Ls $\sim 270^{\circ}$ is too sparse to precisely determine 427 428 the spatiotemporal envelope for frost presence. Nonetheless, frost occurrences are iden-429 tified over a wide range of latitudes.

Despite more restrictive ice detection criteria than used in previous studies (Piqueux 430 et al., 2016; Khuller et al., 2021a), our approach shows a similar or higher sensitivity at 431 432 visible wavelengths given the lower latitudes of frost identifications. Yet, we still find 848 image pairs (out of 2,761, i.e., $\sim 30\%$) that do not show a signature of frost in visible 433 wavelength imagery, while the ground is at CO_2 ice temperature. Noticeably, no iden-434 435 tification in visible wavelength imagery is located in the 45°N-15°S latitude band in the low thermal inertia terrains (i.e., purple in Fig. 4, defined as thermal inertia $<100~J~m^{-2}~K^{-1}~s^{-1/2}$), 436 where widespread diurnal CO_2 frost forms during a significant fraction of the year (Piqueux 437 438 et al., 2016). In contrast, at other latitudes, 87.5% of the image pairs at the CO₂ frost point temperature are associated with bright patches on the ground (whether contam-439 440 inated by water ice or not). This difference in morning frost identifications at visible wavelength suggests a distinct optical behavior of diurnal CO_2 frost when it is associated with 441 442 dusty low thermal inertia terrains compared to CO_2 frost elsewhere on the planet, whether 443 it is diurnal or seasonal, contaminated by water ice, sedimented dust, or either. This ab-444 sence of identifications in visible wavelength imagery on dusty CO_2 cold surfaces could potentially be biased by observational factors, including: 445

- Poor contrast between the bright Martian dust and surface frost, resulting in seem-446 ingly no frost signature on dusty regions. The Martian dust is associated with the 447 448 highest visible and near-infrared lambertian albedo units (excluding polar ice-exposing terrains), i.e., 0.27 (Putzig et al., 2005). This configuration is unlikely because the 449 450 identification criterion for frost in visible wavelength images is based on the pres-451 ence of blue/white hues, and as such, frosts seem perfectly distinguishable from 452 regular darker colored warm surfaces (Fig. 3). In addition, we have identified blue/white surface units in high thermal inertia/albedo/dust cover index areas (Ruff & Chris-453 tensen, 2002), near 180-270°E and 20-30°S (Fig. 4), further confirming that our 454 455 mapping approach has a demonstrated capability to identify surface frost using 456 visible wavelength imagery on high albedo, bright dust-like terrains. For these reasons, an absence of contrast or color between frost and the Martian dust does not 457 458 seem to explain the systematic absence of frosted surfaces in visible wavelength imagery on low thermal inertia terrains. In brief, we acknowledge that a regolith/ice 459 relationship diagnosis solely based on terrain color can be complicated by the low-460 461 ering of the ice's albedo by dust (e.g., Khuller et al. (2021b) with water ice, Singh 462 and Flanner (2016)). To mitigate this potential limit, we have derived the emissivity of these frosted surfaces at 12.57 μ m as emissivity is diagnostic of ice pres-463 ence (see section 2.1). A future study could contribute to this regolith and ice char-464 acterization by performing a thorough analysis of the spectral properties of these 465 466 terrains using more THEMIS bands.
- 467 Clean frost layers growing on the top of the regolith are optically thin at visible
 468 wavelengths, thus indistinguishable in visible wavelength imagery because too thin,
 469 similar to a model proposed by Svitek and Murray (1990) for water ice frost at
 470 the Viking 2 landing site. While it is certainly possible that an unknown fraction
 471 of the image pairs we flagged may be concerned with optically thin layers of frost,

we eliminate a systematic bias because 1) the equivalent thickness of diurnal CO₂
frost frequently approaches several tens to hundreds of microns at dawn in the low
thermal inertia regions (Piqueux et al., 2016), and 2) we have discussed in section
2.1 of this paper how such frost layers should be associated with pronounced albedo
signatures.

477 For these reasons, we conclude that diurnal CO_2 ice should generally be distinguishable at visible wavelengths, and therefore the absence of identification implies very low 478 contrast compared to unfrosted terrains. Quantitatively, the frequency of positive cor-479 relation between T_{CO_2} (Eq. (2)) and blue/white surfaces is insignificant for low thermal 480 inertia terrains where diurnal frost is present (i.e., 5.5% of the images where the ther-481 mal inertia $< 100 \text{ Jm}^{-2} \text{ K}^{-1} \text{ s}^{-1/2}$ are associated with blue/white bright surfaces), but 482 high with medium/high thermal inertia terrains (78.5% of the images where the ther-483 mal inertia > 100 J m⁻² K⁻¹ s^{-1/2} are associated with blue/white bright surfaces, Fig. 484 5a). The 5.5% blue/white surface images pairs on low thermal inertia terrains are grouped 485 486 near 15°S and 200-225°E, and are associated with unique intrinsically blue/white-colored 487 surfaces (see for instance images V71284003/I71284002; V67886003/I67886002 (Christensen 488 et al., 2002)). These occurrences seem thus to have no relationship to surface frost but 489 we have reported them in Fig. 4 for completeness (Class 2). If ignored, our survey shows that no image associated with blue/white bright surfaces are found on low thermal in-490 ertia < 100 J m⁻² K⁻¹ s^{-1/2} terrains. 491



Figure 4. Distribution of dawn THEMIS visible and infrared wavelength image pairs within 5 K of the local CO₂ frost point (see text for the list of image pair selection criteria) corresponding to classes 1-4 (see Fig. 3). Black dots indicate no signature of frost in visible wavelength images (848 images, Class 5). White dots indicate the presence of image pairs where frost is identified at visible wavelengths (1,931 cases). Colorized background is a thermal inertia map (scale bar in J m⁻² K⁻¹ s^{-1/2}), from Christensen et al. (2001) overlaid on with a MOLA shaded relief (Zuber et al., 1992), only shown outside the maximum extent of the continuous seasonal caps (Piqueux et al., 2015). During the winter, the Northern high/mid latitudes are subject to much fewer observations than the Southern high/mid latitudes, partially explaining the hemispheric asymmetry (see Fig. 2 and associated text).

In other words, when the thermal inertia $< 100 \text{ Jm}^{-2} \text{ K}^{-1} \text{ s}^{-1/2}$, i.e., when the 492 surface grain size $R < 20 \ \mu m$ in diameter (Presley & Christensen, 1997), visible wave-493 494 length images do not display a diurnal frost signature. In contrast, the ratio of images presenting a signature at visible wavelengths over the total number of images in our dataset 495 approaches 70-80% for $R>20~\mu{\rm m}$ (i.e., most of the typical Martian regolith (Kieffer, 496 2013)). Most of the missing frost signatures in visible wavelength imagery when the sur-497 498 face temperatures at T_{CO_2} at high latitudes are generally linked to poor image quality, with challenging illumination conditions that prevent us from clearly assessing the color 499 500 of the surface. This difference suggests that on low thermal inertia terrains, the visible wavelength optical properties of diurnal frost are uniquely dominated by those of the sur-501 502 face dust.

For completeness, we mention here that our mapping approach presents two important limits: 1) the impact of H_2O ice on frost signatures at visible wavelengths is not considered, and 2) there is no formal distinction between seasonal versus diurnal ices.

506 • 1: Some water ice may be cold-trapped on CO₂-cold surfaces, but we are unable 507 to characterize its effect in terms of visible wavelength imagery signature. Indeed, 508 previous work has shown that seasonal water ice deposits may be important con-509 tributors to surface blue/white hues in visible wavelength imagery (see for instance 510 Bapst et al. (2015)), if not too contaminated by dust (< 1% dust contamination, 511 Khuller et al., 2021b). Similarly, Kieffer (1968) showed that small amounts of water frost can mask CO_2 frosts. However diurnal deposits may not have adequate 512 time to form thick-enough deposits to be visible on the ground (Martínez et al., 513 2016). Therefore, it is questionable whether an absence of diurnal frost signature 514 515 in visible wavelength imagery on low thermal inertia terrains is biased by the ef-516 fective absence of water ice cold-trapped with CO_2 ice. A differential signature between seasonal and diurnal CO₂ frost due to water ice veneers is not fully sup-517 ported by OMEGA observations of low latitude seasonal CO_2 ice not obscured by 518 519 water ice veneers (Vincendon et al., 2010);

520 • 2: The lack of overlapping visible wavelength observations at different local times prevents us from determining whether blue/white units are sometimes, if ever, as-521 522 sociated with diurnal CO_2 ice. In other words, the blue/white occurrences might 523 all be associated with seasonal CO_2 ice, and diurnal CO_2 ice does not form blue/white patches, regardless of regolith grain size and properties. Nevertheless, CO_2 ice de-524 tected in thermal infrared data at low latitudes in the low thermal inertia regions 525 (purple units in Fig. 4) must be diurnal because these terrains are much too warm 526 for CO_2 ice survival during the day (Piqueux et al., 2016), at any conceivable spa-527 tial scale and season. Therefore, our approach is indeed able to single out a large 528 529 body of diurnal CO_2 frost occurrences between 45°N and 15°S.

The emissivity of frosted terrains in the low thermal inertia regions can also constrain the nature of the CO₂ ice/regolith relationship. As mentioned in section 2.1, ice deposits with micrometric grain size are expected to decrease the emissivity of the surface, especially in THEMIS band 9 data used here. As a result, brightness temperatures are expected to be lower than T_{CO_2} (Piqueux et al., 2016). For each THEMIS image (even those with $T < 0.9 \times T_{CO_2}$) acquired over the low thermal inertia terrain, we extract the averaged surface temperature T_{surf} to derive the surface emissivity ϵ (Fig. 5b) with:

$$\epsilon = \frac{T_{surf}^4}{T_{CO_2}^4} \tag{3}$$

Figure 5b shows that most of THEMIS images in dusty grounds display a high surface emissivity, close to the emissivity of ice-free dusty grounds (0.96, (Ruff & Christensen, 2002; Piqueux et al., 2016)). After removing problematic (i.e. calibration issue) images,



Figure 5. Thermal inertia values of the terrains associated with images pairs at T_{CO_2} (Christensen et al., 2001). The approximate relationship between thermal inertia and grain size is indicated in red (Presley & Christensen, 1997; Kieffer, 2013). Images pairs at T_{CO_2} on low thermal inertia terrains (i.e., $< 100 \text{ Jm}^{-2} \text{ K}^{-1} \text{ s}^{-1/2}$, $R < 20 \mu \text{m}$) are rarely associated with a frost signature at visible wavelengths, and most are located at high latitudes, see Fig. 4). b: CO₂ emissivity derived at $\lambda = 12.57 \mu \text{m}$. Dust emissivity at 12.57 μm (red line) from Ruff and Christensen (2002). Modeled emissivity of the dusty surface overlaid by 10 μm (solid line) and 100 μm (dashed line) CO₂ ice layer indicated in blue (Piqueux et al., 2016).

540 we find that more than 70% of the images show a CO_2 frosted surface emissivity between

541 0.96 ± 0.04 (uncertainty defined as the 3- σ spread from our 5 K tolerance, see section

542 2.2). Such high emissivities at 12.57 μ m indicate the absence of CO₂ frost over dusty sur-

faces, or over micrometer size thickness (Fig. 11 in Piqueux et al. (2016)). But in the
latter case the ice thicknesses would be inconsistent with mass/energy balance results
(Piqueux et al., 2016). Thus, the high surface emissivity values are most consistent with
an absence of frost on the surface.

The visible and thermal infrared observations could conceivably match other re-547 548 golith/ice configurations and formation models than the one presented here (Fig. 1); however, this regolith/ice model is based on reasonable theoretical reasonings, and our work 549 550 shows that the observations we present are consistent with this model. Pathological re-551 golith/ice configurations or alternate ice growth models could match the observations, 552 but they would require still-to-be-proposed formation models compared to what is pre-553 sented here. The observations are consistent with dirty diurnal CO_2 ice in these terrains, following the hypothesis formulated by Piqueux et al. (2016), and illustrated in Fig. 1. 554

555 4 Discussions

556

4.1 Destabilization of a dust grain trough CO₂ sublimation

The sublimation of seasonal carbon dioxide frost has been associated with numerous processes shaping the surface of Mars (Dundas et al., 2012; Dundas, 2020; Diniega et al., 2013, 2021; Hansen et al., 2013; Pilorget & Forget, 2015; Piqueux et al., 2003, 2008).
Similarly, we propose to quantitatively describe the interaction between the sublimating diurnal CO₂ frost and the surface layer.

Based on theoretical models, laboratory experiments, and field observations on Earth 562 563 (with water frost), we hypothesize that the recurring growth and removal of ice may cre-564 ate a stress cycle in the regolith leading to internal grain displacement and microscopic 565 weathering (Everett, 1961; Ruedrich & Siegesmund, 2006; Woronko & Pisarska-Jamroży, 566 2015). On Mars, at sunrise, the rapid sublimation of the diurnal ice creates a short-lived wind field within the regolith that exerts a drag on individual grains, opposite in direc-567 tion to gravity and cohesion forces, somewhat similar to models proposed for the sub-568 569 limation of translucent seasonal CO_2 ice sand grains at high latitudes (Kieffer et al., 2006; 570 Kieffer, 2007; Pilorget & Forget, 2015). However, we reiterate that this model is unique to the low latitudes of Mars where low thermal inertia terrains conducive to the forma-571 tion of diurnal ice exist. This restriction stems from the need of regolith grains smaller 572 573 in size than the typical path length of photons emitted at the peak emission wavelength 574 (Fig. 1). These conditions are met in the lowest thermal inertia terrains, not found in 575 the polar regions of Mars.

576 We note that the upward movement of dust grains in relation to CO_2 ice sublimation is a distinct process from saltation, a process well known to occur on Mars with sand-577 578 size grains (Greeley, 2002; Greeley et al., 1976; Greeley et al., 1980, 1992; Merrison et 579 al., 2007). A lift force is always orthogonal to the flow that creates it, by definition, and is the root cause of the saltation process. Saltation models consider grains subjected to 580 581 forces induced by horizontal winds. With saltation, gas flow around grains is not symmetrical, inducing a circulation that causes the lift following the Kutta-Joukowski the-582 orem (Sears, 1981). Motion starts with a rotational component. 583

In contrast, the grain levitation model we propose differs from saltation models (Fig. 584 585 7a) in that the movement caused by the drag is parallel to flow lines, not orthogonal. In contrast to the saltation model, the flow of CO_2 gas around the grain is symmetrical, 586 587 without circulation and therefore without the lifting described in saltation models by the 588 Kutta-Joukowski theorem. The drag caused by the gas itself on the grain sets the dust 589 grains in motion. When the normal component of the drag is larger than the combina-590 tion of the normal component of the gravity and cohesion forces, grain movement occurs 591 without any rotation. To avoid confusion, in this paper, we describe the movement of grains as "upward movement" as opposed to "lift" because lift implies orthogonality be-592

tween gas flow and grain movement (saltation). This model considers that CO₂ sublimation only acts one single dust grain at a time, and ignores diffusion trough the granular medium.

596 The drag F_d on individual grains created by the moving CO₂ gas is proportional 597 to the gas flow velocity w(t) [m s⁻¹]:

$$F_d = \frac{1}{2}\rho w^2 S C_d \tag{4}$$

with $S = \pi R^2 \, [m^2]$ the section of a spherical grain, C_d a dimensionless drag coefficient, 598 ρ the density of CO₂ gas (0.02 kg m⁻³ (Owen et al., 1977)), R [m] the radius of dust par-599 ticle in dusty low thermal inertia regions (~ 1 μ m). The gas flow velocity w(t) is pri-600 601 marily controlled by the solar energy input and CO_2 ice mass balance at the surface. To evaluate w(t), we first calculate the frost thickness, following the approach presented by 602 Piqueux et al. (2016). In short, when the surface temperature reaches T_{CO_2} , heat lost 603 604 at the surface through radiation to the atmosphere (minus the downwelling radiance) 605 is converted into a mass using the latent heat of condensation of CO_2 . When the sun 606 rises, heat is added to the system at the surface and leads to sublimation, creating a la-607 tent heat flux LE(t). Quantitatively, the CO₂ ice thickness δ [m], and thus the frost mass, 608 can be expressed as a function of the energy balance at the surface:

$$\rho_{i,c}L_e \frac{\mathrm{d}\delta^3}{\mathrm{d}t} = \epsilon \sigma T_{CO_2}^4 - Q_{IR}(t) - (1-A) \times S_M(t) \times \cos(i(t))$$
(5)

with $\rho_{i,c}$ the CO₂ ice density set to 1.6×10^3 kg m⁻³ (Mangan et al., 2017; Putzig et 609 al., 2005), and L_e the latent heat of sublimation of CO₂ ice (5.9 10⁵ J kg⁻¹ (Pilorget & 610 Forget, 2015)), ϵ the ice emissivity taken as 0.99 to be consistent with previous work from 611 Piqueux et al. (2016), σ the Stefan Boltzmann constant, Q_{IR} [W m⁻²] the atmospheric 612 downwelling flux, A the lambertian vis/ near-infrared albedo of dusty surface set to 0.27613 (Putzig et al., 2005), S_M [W m⁻²] the insolation of the surface after atmospheric cor-614 rection, and i [rad] the solar incidence angle. Q_{IR} and S_M are calculated using the nu-615 616 merical thermal model KRC (Kieffer, 2013). In our model (Fig. 1), the CO_2 ice crystals form at depth, within the first few microns to hundreds of microns under the sur-617 618 face, that is, at a depth three orders of magnitude smaller than the diurnal thermal skin 619 depth in these terrains (i.e., 1-3 cm, (Grott et al., 2007)). Therefore, at dawn, the heat wave induced by direct solar insolation has to travel through a negligible distance com-620 pared to a diurnal skin depth to reaches the ice, in a matter of just a few tens or sec-621 onds to a minute. This time constant is ignored in our model. 622

623 The latent heat
$$LE$$
 [W m⁻²] released at the surface can thus be expressed as

$$LE(t) = L_e \mid \frac{\mathrm{d}m_{ice}}{\mathrm{d}t} \mid = \mid \rho_{i,c}L_e \frac{\mathrm{d}\delta^3}{\mathrm{d}t} \mid = \mid \epsilon \sigma T_{CO_2}^4 - Q_{IR}(t) - (1-A) \times S_M(t) \times \cos(i(t)) \mid (6)$$

624 Latent heat is only released when $\delta(t) > 0$, i.e., when frost is still present at the 625 surface. Therefore, Eq. (5) can be numerically integrated to derive $\delta(t)$. LE(t) can then 626 be computed from Eq. (6) when $\delta(t) > 0$.

627 The CO₂ gas created by this sublimation is generated uniformly in the pore result-628 ing in a mass flow rate q [kg m⁻² s⁻¹] perpendicular to the surface, which gives a mean 629 gas velocity using Eq. (7) (Kieffer et al., 2000; Diniega et al., 2013):

$$w(t) = \frac{q}{\rho} = \frac{LE(t)}{\rho L_e} \tag{7}$$

To bound w(t) in low thermal inertia regions, we use the season-dependent CO₂ 630 frost thickness at dawn provided by Piqueux et al. (2016). Using Eq. (7), we compute 631 the wind-driven velocity at a spatial resolution of one point per degree, 36 times per year 632 to conform with the frost thickness maps of Piqueux et al. (2016). Frost thicknesses cal-633 634 culated by Piqueux et al. (2016) represent lower bounds because they assumed 0 frost thickness at 3 AM, when MCS observed the surface. They acknowledge that CO_2 frost 635 may have formed earlier in the night, resulting in non-0 thicknesses at 3 AM. This as-636 sumption results in conservative wind speeds here, as more frost might be available for 637 638 sublimation. Figure 6 shows the frequency distribution of peak wind velocity for each location on Mars where diurnal frost is expected to form. Values range from 2.3 mm s^{-1} 639 to 3.2 cm s^{-1} , with 1.7 cm s^{-1} on average and a standard deviation of 0.8 cm s^{-1} , with 640 the highest velocities associated with the thickest frost deposits and the lowest ground 641 thermal inertia values. This range of gas velocities is generally consistent (although higher) 642 than values reported by others for syn-regolith CO_2 winds, based on laboratory measure-643 644 ments (Sylvest et al., 2018) and numerical models both for seasonal ice sublimation (Kieffer, 2007) and Knudsen pumping (de Beule et al., 2013; Schmidt et al., 2017). The model 645 646 presented here assumes uniform laminar flow, but turbulent eddies could conceivably be present and necessitate an additional diffusivity coefficient term in Eq. (7) (Brutsaert, 647 1982), resulting in lower wind values, and therefore fewer occurrences of regolith grains 648 destabilization. Energy losses related to turbulent transport or the absorption of solar 649 650 energy by the medium, reducing the energy brought to the ice crystals, are also neglected, which can lead to an overestimation of w. A full model of the drag on grains would re-651 quire a significantly more complex mathematical treatment (see for instance Hu (2019)), 652 beyond the scope of this paper. The approach presented here solely aims at evaluating 653 whether diurnal frost sublimation can conceivably destabilize sloped terrains. 654

655 The resulting drag can then be calculated assuming an incompressible laminar flow.656 The Reynolds number is given by:

$$Re = \frac{\rho w 2R}{\mu} \tag{8}$$

657 where $\mu = 7.42 \times 10^{-6}$ [Pa s] is the dynamic viscosity of CO₂ gas at 148 K given by 658 Sutherland's law (Sutherland, 1893). With $w \leq 0.03$ m s⁻¹ (Fig. 6), Eq. (4) can be 659 simplified following a formulation by Yang et al. (2015):

$$C_d = \frac{24}{Re} \tag{9}$$

660 Introducing Eq. (8) and Eq. (9) in Eq. (4), and deriving S with R yield to the final form 661 of :

$$F_d = 6\pi R\mu w \tag{10}$$

662 For grain movement to occur, F_d must be larger than the combination of gravity 663 F_g and cohesion forces F_c . Hence, vertical movement is possible when Eq. (11) is ver-664 ified:

$$F_d \ge 2F_c \cos(\psi) + F_g \tag{11}$$

665 with ψ an angle describing the packing of grains (see Fig. 7a). The cohesive force be-666 tween two spherical dust particles of radius R [m] can be expressed following an expres-667 sion given by (Hartzell et al., 2013; Perko et al., 2001; Scheeres et al., 2010)



Figure 6. Distribution of the sublimation-driven wind velocities in the Martian regolith, and comparison with other published values (red) and those reported for Knudsen pump (blue). Each bin has been normalized by the total number of samples (2552884) to get a normalized distribution. The green arrow indicates the fluidization threshold beyond which avalanching becomes possible (Cedillo-Flores et al., 2011)

$$F_c = \alpha R \tag{12}$$

where α [N m⁻¹] is a cohesive parameter that depends on the physical and chemical prop-668 erties of individual dust grains. Wind tunnel experiments with frost-free dust aggregates 669 suggest that $\alpha \sim 10^{-5}$ N m⁻¹ (Merrison et al., 2007). However, the presence of the frost 670 (carbon dioxide and possibly water frosts) within the pores certainly increases the co-671 672 hesion between grains (Greenberg et al., 1995; Perko et al., 2002). Consequently, it certainly impacts α . Furthermore, ice crystals might interact with electrostatic (Merrison 673 674 et al., 2007; Sullivan et al., 2008) and/or magnetic (Kinch et al., 2006; Goetz et al., 2008) forces that hold dust grains together (none of which modeled here), possibly creating ag-675 glomerate dust forms rather than spherical grains (Kinch et al., 2015). A contrario, un-676 677 der terrestrial conditions, experimental work has demonstrated that recurring ice crys-678 tal formation and growth in porous media can generate sufficient stress leading to mi-679 croscale physical weathering on individual grain, internal grain displacement, and even

the rupture of cemented materials (Everett, 1961; Ruedrich & Siegesmund, 2006; Woronko
& Pisarska-Jamroży, 2015), confirming that the presence of ice can also be a factor of
reduced cohesion.

We can now compare the drag, cohesion and gravity forces with each other. The
normal component of the buoyancy-corrected gravity to a spherical dust grain (Phillips,
1980) is given by:

$$F_g = \frac{4}{3}\pi R^3 (\rho_g - \rho)g\cos(\theta) \tag{13}$$

686 with ρ_g the density of the grains set to 2500 kg m⁻³, g the gravity set to 3.71 m s⁻², θ 687 [degrees] the slope of the surface. Introducing Eq. (10), (12) and (13) into (11) leads to

688 the final condition for the levitation of the grain:

$$6\pi R\mu w \ge 2\alpha R\cos(\psi) + \frac{4}{3}\pi R^3(\rho_g - \rho)g\cos(\theta)$$
(14)

689 A regolith dust grain moves upward if the velocity of the sublimation-driven CO₂ gas 690 flow w is larger than a velocity threshold w_* [m s⁻¹]:

$$w \ge w_* = \frac{2\alpha R \cos(\Psi) + \frac{4\pi R^3}{3}(\rho_g - \rho)g\cos(\theta)}{6\pi R\mu}$$
(15)

691 w_* is computed for slope angles θ ranging between 0 and 30°, which corresponds to slopes 692 commonly observed on Mars (Kreslavsky & Head, 1999; Kreslavsky & Head, 2000). The 693 grain packing angle ψ theoretically ranges from 0 to 60°. As the low thermal inertia units 694 are associated with very porous media (Presley & Christensen, 1997), ψ should be high 695 and is set to 60° in the following.

The wind velocity threshold w_* necessary for grain movement is shown in Fig. 7b 696 as a function of grain size, local slope and cohesive parameter α . w_* is typically smaller 697 than 3.5 cm s⁻¹ for the 1-2 μ m grains expected in 100 J m⁻² K⁻¹ s^{-1/2} thermal iner-tia terrains with typical cohesive parameter $\alpha \sim 5 \ 10^{-5} \ N \ m^{-1}$. Under favorable con-698 699 ditions (see Fig. 7b), $w_* < 1 \text{ cm s}^{-1}$. Sublimation-driven wind velocities w can reach 700 701 this range of values (Fig. 6 and 7b), indicating that vertical drags are comparable or larger in strength to gravity and cohesive forces combined, and might be able to move individ-702 703 ual grains upward at the time of peak CO₂ sublimation. Further, Fig. 7b confirms that 704 the local slope and grain cohesive parameter also play an important role: flat or shal-705 low sloped terrains are associated by larger w_* compared to steep slopes and grains are 706 easier to move on surfaces that are already close to the angle of repose $(30-35^{\circ})$ (Kleinhans 707 et al., 2011; Atwood-Stone & McEwen, 2013)). In addition, the grain packing angle seems to display an even larger control over w_* : larger ψ , i.e., looser packing, strongly decreases 708 w_* compared to more compact grain arrangements (e.g., w_* for a micrometer dust grain 709 on a 30° sloped surfaces and $\alpha = 5 \ 10^{-6} \ \mathrm{N \ m^{-1}}$ can change from 3.1 cm s^{-1} for $\psi =$ 710 60° to 5.5 cm s^{-1} for $\psi = 30^{\circ}$). The low thermal inertia values associated with the dusty 711 712 terrains on Mars require very low-density regolith, high porosity, and therefore high ψ , 713 thus favoring grain movement. Generally, grains $\sim 20 \ \mu m$ or larger cannot be displaced by this drag mechanism (Fig. 7b), but 1-2 μ m uncohesive grains located on sloped ter-714 715 rains can be associated with a regime where frost sublimation can initiate grain movement. 716

717 Similar conclusions regarding the ability of subliming CO2 ice to displace regolith 718 grains were drawn by Kieffer et al. (2000); Kieffer (2007) and Kieffer et al. (2006) for sea-719 sonal CO₂ sublimation near the poles fracturing a pressurized CO₂ ice slab that could 720 promote the suspension of 1-2 μ m dust grains. Our modeling results are also somewhat



Figure 7. Regolith grain motion domains. a: schematic of the forces acting on a stationary regolith grain subjected to a drag caused by CO₂ frost sublimation. F_g is the gravitational force ; F_c an interparticle force; F_d is the drag. b: Velocity threshold w_* (colored lines) as a function of grain size, local slope θ , and cohesive parameter α for a packing angle $\psi = 60^{\circ}$. Each curve delineates two domains: a domain "above" where grain movement is possible ($w > w_*$), and a domain "below", where grain movement is not possible ($w < w_*$). The horizontal dashed line marks the maximum modeled wind velocity (e.g., $w = 3.2 \text{ cm s}^{-1}$, Fig. 6). Locally, dust grains can be moved upward if the wind velocity w is larger than the velocity threshold w_* . For the average wind velocity $w = 2.3 \text{ mm s}^{-1}$, grains ~ 1-5 microns in size on steep slopes are in a domain where their movement is possible if the grains are not strongly bonded to each other ($\alpha \sim 10^{-6} \text{N} m^{-1}$, $w_* \sim 1 \text{mm s}^{-1}$).

comparable to laboratory observations of subliming CO₂ frost on sloped surfaces able to trigger downslope material movement (Sylvest et al., 2018), although significant experimental and modeling differences exist (i.e., the processes we describe here exclusively involves dust, whereas those experiments involved sand-size material; we assume nearsunrise illumination conditions whereas these experiments assumed mid-day insolation relevant to seasonal processes; and we assume extremely small amounts of ice compared 727 to these experiments with CO₂ ice slabs). In addition, we occasionally derive large gas velocities $w > 0.0165 \text{ m s}^{-1}$ (Fig. 6) suggesting that fluidization of avalanching mate-728 729 rial becomes possible (Cedillo-Flores et al., 2011), indicating that under favorable conditions, dust grains put in movement along sloped terrains may therefore display a fluid-730 731 like behavior similar to that modeled by others (Miyamoto et al., 2004). We conclude that short-lived CO_2 sublimation-driven winds in the surficial regolith at sunrise can set 732 733 individual surface dust grains in motion, preferentially on sloped, porous terrains composed of poorly cohesive dust grains, and can even lead to flow-like patterns. 734

735 Although this model indicates that the conditions are met to set individual dust 736 grains in motion at the very surface of the regolith through morning CO_2 frost sublimation, vertical drag is expected to decrease shortly after the grains are displaced vertically. 737 738 In other word, we do not suggest that this mechanism sends individual dust grains on 739 ballistic trajectories above the ground, as proposed for the formation of fans and other 740 features in the polar regions; but rather that grains rearrangement can occur, internal 741 to the regolith in the top few 10's to 100's μ m. During sublimation, the equivalent an-742 gle of repose is artificially and temporarily decreased, but material is not envisioned to 743 depart from the surface.

744 4.2 The formation of Slope Streaks

Here, we hypothesize that this sublimation-driven slope destabilization can initiate and facilitate dust avalanching and the formation of slope streaks. Slope streaks (Fig. 8) are dark wedge-shaped surface features on sloped terrains, initiating at a point and
broadening downslope (Chuang et al., 2007; Ferris et al., 2002; Kreslavsky & Head, 2009;
Miyamoto et al., 2004; Phillips et al., 2007; Schorghofer et al., 2002; Sullivan et al., 2001).
Work by (Chuang et al., 2007) also shows that slope streaks are associated with mass
movement downslope, and are not limited to a simple darkening of the surface.

A survey by Bhardwaj et al. (2019) using the High Resolution Imaging Science Ex-752 753 periment (HiRISE, McEwen et al. (2007)) data shows that the length of slope streaks 754 varies from 14 to 1600 m, and their widths range from 1 to 148 m. They exhibit a wide 755 variety of morphologies, including linear, curved, and fan shaped (Bhardwaj et al., 2019). 756 The vast majority of slope streaks is found on dusty low thermal inertia surfaces (Baratoux 757 et al., 2006; Kreslavsky & Head, 2009), and work by Heyer et al. (2019) seems to highlight a loose relationship between formation rate and seasons, in contrast with a previ-758 759 ous study that did not show any seasonality in the development of new streaks (Schorghofer 760 & King, 2011).

761 The formation process for slope streaks is still debated, and generally falls into two 762 categories: wet vs. dry mechanisms, although both types can be associated with dust 763 avalanches. Wet processes involve episodic seeping of transient liquids generating lowalbedo oxides that precipitate as slope streaks (Mushkin et al., 2010), subsurface water 764 765 ice melting that triggers dust avalanches (Schorghofer et al., 2002), and regolith dark-766 ening through seasonal iron oxide precipitation associated with chloride brines resurgence 767 originating from shallow ice reservoirs (Kreslavsky & Head, 2009). In this model, brines 768 act as a lubricant facilitating the movement of dust particles. Dry mechanisms involve 769 dust avalanches created by wind-regolith interaction (Baratoux et al., 2006; Chuang et 770 al., 2007, 2010; Dundas, 2020; Sullivan et al., 2001), and surficial gravity-driven erosional 771 processes followed by transport in the downslope direction (Phillips et al., 2007). In ad-772 dition, some slope streaks seem to have been linked to specific triggering agents, including dust devils (Schorghofer et al., 2007), impact craters (Burleigh et al., 2012; Chuang 773 774 et al., 2007), and rock falls (Chuang et al., 2007). However, the initiating and sustain-775 ing mechanism for the vast majority of observed slope streaks is not identified (Bhardwaj et al., 2019; Dundas, 2020), and Earth analogs have not been recognized. We propose 776 777 a dry mechanism for the initiation and development of slope streaks involving the sub-



Figure 8. Example of a slope streak (HiRISE) ESP_053518_1955) on the Olympus Mons aureole (Lat = 15.23° N, Lon = 214.9° E, at Ls = 106.3°). The black arrow points the sun orientation.

778 limation of diurnal CO_2 ice on dusty terrains. In this model, diurnal CO_2 frost forms dirty ice within the top few tens to hundreds of microns of the dusty low thermal iner-779 780 tia regolith, where both diurnal frost and slope streaks are observed (Fig 9), but where no frost signature at visible wavelengths is recognized at dawn (Fig. 4). At sunrise, this 781 782 diurnal frost sublimates, initiating a vertical drag on individual grains able to set those 783 uncohesive in motion (Fig. 7). In this model, steep dusty terrains already close to the angle of repose (i.e., 30-35° (Kleinhans et al., 2011; Atwood-Stone & McEwen, 2013)) 784 785 are destabilized and initiate an avalanche of dust. This model is consistent with the fact 786 that typical slope associated with slope streaks is high, i.e., 25° (Brusnikin et al., 2016). 787 Warmer subsurface dust gets entrained downslope and is placed in contact with the still-788 frosted pristine regolith surface. Sublimation is further promoted, fluidizing the material moving downslope. 789

790 This model requires slope streaks to initiate and develop in a relatively short pe-791 riod of time near dawn, when CO_2 frost is present and actively subliming. This predic-



Figure 9. Distribution of slope streaks (black dots) extracted from (Schorghofer et al., 2007) and diurnal CO_2 frost distribution and frequency (adapted from (Piqueux et al., 2016)). No color is reported where no diurnal frost is modeled to be present at dawn. 0 means no diurnal frost ever, 100 means diurnal frost is present every dawn of the year. Background is a TES albedo map, Mollweide projection, 200°E prime meridian.

tion is consistent with observations by Aharonson et al. (2003) who noticed that slope
streaks completely develop within a few sols at most, in contrast with Recurring Slope
Lineae that develop over seasonal periods of time (McEwen et al., 2011). A triggering
and sustaining agent associated with a relatively short dynamics disfavoring seasonal processes over diurnal (or shorter) processes is consistent with our proposed mechanism.

797 Our model also suggests that slope streaks should preferentially form on east or equator facing slopes, where CO_2 frost sublimation at survise is assumed to be most vig-798 799 orous as they received more solar insolation. However, statistical work on slope streaks 800 orientation does not confirm this prediction (Baratoux et al., 2006; Heyer et al., 2019; 801 Schorghofer et al., 2002). Around Olympus Mons, slope streaks seem to form in a wide 802 range of directions (Heyer et al., 2019), without a clear preferential east or equator fac-803 ing orientation. This possible discrepancy between our model and the observations re-804 ported by Heyer et al. (2019) indicates that at a minimum, multiple factors control the 805 initiation of slope streaks.

806 If triggered by subliming CO_2 ice, slope streaks might preferentially form near the top of sloped terrains where radiative coupling with adjacent surfaces is minimized (and 807 therefore where CO_2 frost might preferentially form). Similarly, in our model, slope streaks 808 should preferentially form overnight during the colder seasons. Slope streak formation 809 810 seems to follow a mildly pronounced seasonality (Heyer et al., 2019; Schorghofer & King, 811 2011), but on the Olympus Mons Aureole, they form at all seasons, with a peak between $L_s = 140^{\circ}$ and $L_s = 220^{\circ}$ (Heyer et al., 2019), also corresponding to peak diurnal frost 812 thickness (Piqueux et al., 2016). Heyer et al. (2019) discuss the seasonality of slope streak 813 814 formation in terms of absolute peak surface temperatures, but not in terms of nighttime 815 temperatures. Peak daytime temperatures can be associated with some of the lowest night-816 time temperatures under low atmospheric aerosol opacity (Streeter et al., 2020; Wilson

817 et al., 2006) especially at high elevation on Tharsis. The seasonality of slope streak for-818 mation relative to the presence of CO_2 -ice-cold surfaces is not characterized and discussed 819 in the literature. Future investigations with repeated seasonal coverage at other loca-820 tions should be able to confirm and clarify any correlation.

Lastly, numerical simulations from Miyamoto et al. (2004) have shown that a flu-821 822 idization of the avalanching material is necessary to explain the length and width of the slope streaks given the angle of the slopes on which they develop. Similarly, work by Pilorget 823 824 and Forget (2015) on the formation of gullies at high latitudes, supported by Dundas et 825 al. (2019), showed that the sublimation of seasonal CO₂ frost originally present within 826 the pores of regolith at high latitudes created an upward gas flow leading gravity-driven 827 avalanches to behave like liquid-fluidized flows similar to pyroclastic flows on Earth. The 828 mechanism we propose here is comparable, as avalanching of material stirred from the warmer shallow regolith progressively covers subliming CO_2 ice downslope and promotes 829 sublimation/fluidization. Furthermore, gas velocities reached during sublimation (Fig. 830 831 6) can be high enough to exceed the velocity threshold computed in Cedillo-Flores et al. 832 (2011) required for fluidization of avalanching material.

833 This model involving carbon dioxide ice seems generally consistent with the spatio-834 temporal distribution of slope streaks versus diurnal CO_2 frost (Fig. 9). However, slope streaks initiate at well-defined discreet locations upslope (Schorghofer et al. (2007), Fig. 835 836 9), indicating that the conditions for their initiation are only rarely met. Otherwise, they 837 would ubiquitously cover dusty sloped terrains and probably frequently overlap each other. 838 For this reason, we conclude that subliming CO_2 frost alone does not seem to be suffi-839 cient to initiate slope streaks, and may need coupling with other seasonally-driven mechanisms. The interannual variability in slope streak formation rate identified by Heyer 840 841 et al. (2019) is consistent with the notion that other environmental conditions may be 842 required.

The presence of dirty diurnal CO_2 frost in the low thermal inertia terrains on Mars 843 844 at low latitude is consistent with visible and thermal infrared imagery acquired at dawn. Overnight crystal growth and sublimation-driven wind have the potential to instigate 845 846 grain movement. We hypothesize that diurnal CO₂ ice sublimation may be triggering and sustaining the growth of slope streaks, but other factors such as local winds may be 847 848 involved in this process. Our proposed model overcomes several limitations of compet-849 ing dry mechanisms presented in the literature, but it is not unequivocally validated by 850 observations constraining the seasonality or orientation of slopes highlighted by others. 851 Surface imaging at dawn and other local times by THEMIS and by the Colour and Stereo Surface Imaging System (CaSSIS) (Thomas et al., 2017) on the ExoMars Trace Gas Or-852 biter will certainly continue to provide new important constraints on the relationship be-853 854 tween the atmosphere and the surface.

855 5 Conclusions

We have conducted an analysis of THEMIS visible and thermal infrared data acquired at dawn. This work constrains the relationship between diurnal frost and the surficial regolith on Mars. It unveils the potential geomorphological impact of the diurnal
CO₂ cycle. Specifically:

- The distribution of THEMIS thermal infrared data acquired at dawn confirms the widespread nature of CO₂ frost on Mars previously reported (Piqueux et al., 2016;
 Khuller et al., 2021a), with surface temperature indicative of carbon dioxide ice presence observed at all latitudes, with a strong seasonal control;
- Multiband THEMIS visible wavelength images acquired at dawn frequently show
 blue/white hues interpreted as clean surface frost over a significant fraction of the
 surface. In the southern hemisphere, we find frosted surfaces in visible wavelength

867 868 869 870 871 872 873 874 875 876 877 878 879 880 881 882 883 884 885 886 887 888 889 890 891 892 893 894 895 896 897 898 899 900 901 902 903 904 905 906 907 908 909 9010	imagery at latitudes much lower than reported in the literature, i.e., 20°S vs. 33- 35°S previously. In the northern hemisphere, seasonal coverage is too sparse to precisely determine the limits of frost surface presence at THEMIS resolution; In the mid-to-low latitude low thermal inertia terrains (45°N-15°S), surface tem- peratures consistent with the presence of CO ₂ frost on the ground do not show any frost signature in visible wavelength imagery. The regolith in these low ther- mal inertia terrains consists of 1-2 µm dust grains conducive to the formation of diurnal CO ₂ ice, and we interpret the absence of frost in the visible wavelength imagery as an indication of the presence of dirty CO ₂ frost forming within the sur- ficial regolith. This conclusion is also supported by the high emissivity at 12.57 µm of these frosted surfaces; The dirty diurnal frost hypothesis is further supported by the notion that ther- mal infrared radiative cooling crucial to CO ₂ frost formation peaks at a wavelength significantly longer than the regolith dust grains (18 vs ~ 1 micron), thus within a layer of regolith several grains deep; The visible and thermal infrared observations presented in this paper could con- ceivably match other regolith/ice configurations and formation models than the one presented here (Fig. 1). However, the model detailed in this work is based on reasonable theoretical reasonings and is supported by our analysis; Within the regolith, frost sublimation-driven wind at sumrise reaches up to 3.2 cm s ⁻¹ , with a mean value between 1.5 and 2.0 cm s ⁻¹ (standard deviation of 0.8 cm s ⁻¹) as a function of season, latitude, and elevation. These values are upper limits and are generally larger than (but comparable to) other values reported in the liter- ature for seasonal ice sublimation and Knudsen pumping; The vertical drag exerted on individual grains can be larger and opposite in di- rection to cohesion plus gravity forces, suggesting that motion can be initiated by this wind. Local slope, and mor
910 911 912	tor over short time scales would represent another expression of the complex re- lationship between the atmosphere and the surface; While slope streaks are common in low thermal inertia regions, their formation may nonetheless remain relatively rare events, otherwise they would frequently over-
913 914 915 916	lap each other. This observation suggests that the conditions for their formation (either initiation or development along slopes) are not often met, despite the ubiquity of CO_2 frost at the low latitude low thermal inertia terrains on Mars. There- fore, the sublimation of CO_2 frost after dawn may not be the only necessary fac-
917	tor required for their formation.

918 Surface observations at dawn are generally limited by operational constraints. Fu-919 ture investigations able to provide local-time-dependent characterization of the regolith

will certainly help identify and characterize active processes such as the ones presentedin this paper.

922 Open Research

923 All THEMIS images, TES thermal inertia map, and MOLA data are publicly avail-924 able within the NASA Planetary Data System (Christensen et al., 2002; Putzig et al., 925 2005; Smith et al., 1999) at https://pds-geosciences.wustl.edu/missions/odyssey/ themis.html, https://pds-geosciences.wustl.edu/missions/mgs/tes-timap.html, 926 927 https://pds-geosciences.wustl.edu/missions/mgs/mola.html. The numerical model 928 used in this paper is available at the following address: http://krc.mars.asu.edu/. Data 929 files for figures and lists of THEMIS images used in this analysis are available in a pub-930 lic repository, see Lange et al. (2021).

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Supporting Information for

Gardening of the Martian Regolith by Diurnal CO₂ Frost and the Formation of Slope Streaks

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Contents of this file Figures S1

Introduction

This document presents a figure illustrating the distribution of THEMIS images indicative of the presence of CO_2 ice (i.e., within 5K of the local CO_2 frost point) and documents the frost identification confidence level (Class 1-5) following a classification described in section 2.2.



Figure S1. Distribution of THEMIS visible and infrared wavelength image pairs within 5 K of the local CO₂ frost point, near sun rise. The color of the dots indicates the confidence level of the frost identification in the THEMIS visible wavelengths images: red for Class 1 images (115 occurrences), blue for Class 2 images (534 occurrences), green for Class 3 images (708 occurrences), purple for Class 4 images (556 occurrences), black for Class 5 (848 occurrences). Grey background is a MOLA-shaded relief map. Maximum extent of the seasonal caps (Piqueux et al., 2015) shown as black curves.