# Observation of atmospheric and oceanic dynamics using ocean-bottom distributed acoustic sensing

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#### Abstract

Distributed Acoustic Sensing (DAS) leverages an ocean-bottom telecommunication fiber-optic cable into a densely-sampled massive array of strain sensors. We demonstrate DAS applications to Passive Acoustic Monitoring (PAM) through an experiment in Longyearbyen, Svalbard, Norway. We show that DAS can measure many types of signals generated by dynamics in the atmosphere, ocean, and solid earth. These include primary and secondary microseisms, Scholte waves, water-layer acoustic resonances, and seismic waves from earthquakes. In addition, we can trace the origin of primary microseisms back to distant storms a quarter of the way around the planet. We also find that the fjord acts as an amplifier for microseisms. Because DAS is capable of hydroacoustic monitoring with high spatial resolution over great distances, it can deliver great scientific value to ocean observation. We believe that DAS can and will become a valuable component of the Global Ocean Observing System.

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- 18 Key Points:
- We review the mechanisms of ocean-bottom seismo-acoustics recorded by Distributed
   Acoustic Sensing (DAS) of seafloor telecommunication fiber-optic cables.
- We are able to trace primary microseisms measured by DAS back to distant storms more than 10,000 km away.
- DAS could be a valuable and game-changing addition to the Global Ocean Observing
   System (GOOS), expanding coverage, spatial and temporal resolution.

# 25 Abstract

Distributed Acoustic Sensing (DAS) leverages an ocean-bottom telecommunication fiber-optic 26 cable into a densely-sampled massive array of strain sensors. We demonstrate DAS applications 27 to Passive Acoustic Monitoring (PAM) through an experiment in Longyearbyen, Svalbard, 28 Norway. We show that DAS can measure many types of signals generated by dynamics in the 29 30 atmosphere, ocean, and solid earth. These include primary and secondary microseisms, Scholte waves, water-layer acoustic resonances, and seismic waves from earthquakes. In addition, we can 31 trace the origin of primary microseisms back to distant storms a quarter of the way around the 32 planet. We also find that the fjord acts as an amplifier for microseisms. Because DAS is capable 33 of hydroacoustic monitoring with high spatial resolution over great distances, it can deliver great 34 scientific value to ocean observation. We believe that DAS can and will become a valuable 35 component of the Global Ocean Observing System. 36

# 37 Plain Language Summary

Over 1.3 million kilometers of submarine fiber-optic cables have been deployed around the Earth for telecommunications. In this study, we use one such cable in Svalbard, Norway, to measure vibrations at the seafloor. We describe the characteristics of these signals and deduce their origins, which include distant storms occurring in the South Atlantic Ocean more than 10,000 km away. We believe that this sensing technique will soon become a standard and powerful tool for the oceanographic community.

# 44 **1 Introduction**

The Earth's atmosphere and oceans are continuously in coupled motion. These complex 45 motions and interactions determine both weather and, over the longer term, the climate of the 46 planet. Oceans play a highly significant role in climate, because they can retain heat and distribute 47 it around the globe (Schmitt, 2018). Large-scale ocean currents, which are driven by variations in 48 water density caused by temperature and salinity gradients, influence the climate by exchanging 49 heat and water with the atmosphere. A change in ocean dynamics could induce major climate 50 variations over large areas of the Earth in the long term (Bigg & Hanna, 2016). Hence, ocean 51 52 surface winds, currents, and surface gravity waves are key climate variables that induce exchanges of momentum, energy, heat, salinity, gases, and other tracers between the ocean and atmosphere 53 (Villas Bôas et al., 2019). 54

Ocean surface gravity waves have random properties and evolve from complex 55 mechanisms. Their modern studies started in the 1940s (Mitsuyasu, 2002; Wunsch, 2021), with 56 57 seminal contributions from icons such as Sverdrup (Sverdrup, 1947), Stommel (Stommel, 1948) 58 and Munk (Munk, 1950). Ocean surface gravity waves are a primary source of turbulence in the upper ocean, and they are an important factor in the air-sea momentum transfer. In addition, they 59 60 directly affect navigation, offshore structure design, and coastal erosion (Abolfazli et al., 2020). However, they are not used explicitly in constraining most ocean-atmosphere models, because 61 62 high-spatial-resolution (scales under 25 km) two-dimensional (2D) measurements of waves are not normally available. Such measurements could significantly improve ocean models (Wu et al., 63 2019). 64

65 Many instruments have been developed to measure directional ocean surface gravity waves 66 (European cooperation in science and technology Action 714, Working Group 3, 2005). The 67 classical methods such as spatial arrays and pitch-and-roll buoys have been complemented by new

technologies such as the displacement and GPS buoys, acoustic Doppler current meters, 68 microwave and marine radars, coastal high-frequency radars, and real and synthetic aperture 69 radars. However, none of these instruments can provide all the data needed to make a complete 70 71 and robust estimate of the directional properties of ocean surface gravity waves. Data with high spatial resolution and extensive spatial coverage would be necessary to overcome this limit. In 72 principle, subsurface instruments that measure ocean-bottom pressure fluctuations due to surface 73 gravity waves could be deployed in spatially extended arrays for accurate estimation of swell 74 directional spectra, but this would be prohibitively expensive. Therefore, compact subsurface 75 instruments, whose dimensions are smaller than the typical wavelength, are more widely used by 76 the oceanographic community. 77

Distributed Acoustic Sensing (DAS) is a technology that is able to exploit the optical fiber 78 in standard telecommunication cables as an extended spatial array of acoustic sensors (Hartog, 79 2017). Over 1.3 million kilometers of submarine telecommunication cables have been deployed 80 around the Earth. Many optical fibers in these cables, often 'spares', are not currently used for 81 telecommunication. It is possible to repurpose these unused 'dark' fibers to serve as distributed 82 acoustic sensors to measure, among other signals, ocean-bottom pressure fluctuations. DAS 83 measures the strain fluctuations at each sensing element of an optical fiber. A DAS interrogator 84 can measure the strain data along the fiber with a length up to 171 km (Waagaard et al., 2021). 85 Therefore, DAS can form spatially extended arrays with very large dimension compared to the 86 typical length of ocean surface gravity waves. In addition, DAS measures data with a spatial 87 sampling interval of as little as 1 m, which creates arrays of many tens of thousands of sensors at 88 low cost. 89

90 DAS in submarine fiber-optic cables can measure pressure fluctuations at the ocean bottom originating from a variety of sources (Landrø et al., 2021). DAS in ocean-bottom 91 telecommunication cables can detect ocean surface gravity waves, microseisms and earthquakes 92 (Lindsey et al., 2019; Sladen et al., 2019). Furthermore, Williams et al. (2019) demonstrate that 93 94 DAS can record the seismic waves from a distant earthquake, ocean surface gravity waves, and Scholte waves. However, their spectral analyses were performed on a data record of only one-95 hour. DAS data with a longer recording length are necessary for studying the dynamics of ocean 96 surface gravity waves originating from distant storms. For example, Zhan et al. (2021) show 97 98 several dispersive signals from ocean swells from distant storms in a spectrogram computed over 11 days of the fiber-optic sensing data using the state of polarization technique. 99

In this article, we show that DAS can be employed as a valuable tool for studying ocean dynamics. First, we describe the DAS data used in our study and their acquisition parameters. Second, we review the mechanisms of the ocean-bottom vibrations that are recorded by DAS along an ocean-bottom telecommunication cable. We also review the characteristics of the DAS data corresponding to different mechanisms of the ocean-bottom vibrations. Then, we discuss the results of our analysis related to ocean surface gravity waves corresponding to distant storms. Finally, we address some potential applications of DAS in the oceanographic community.

# 107 **2 Method**

We used a dark fiber of SMF-28 single mode silica type in an existing submarine telecommunication cable, which was trenched into soft sediments at 0–2 m below the seafloor, from Longyearbyen to Ny-Ålesund in Svalbard, Norway (Figure 1). The cable is owned and operated by Uninett AS, which is the National Research and Education Network (NREN) in

Norway. We connected an OptoDAS interrogator, developed by Alcatel Submarine Networks, to 112 the cable end onshore Longyearbyen. The OptoDAS interrogator sends linear optical frequency-113 modulated swept pulses into the fiber and receives backscattered pulses from impurities in the 114 fiber (Waagaard et al., 2021). It calculates the time-differentiated phase changes of consecutive 115 backscattered pulses corresponding to every spatially sampled position along the fiber. These are 116 used to estimate longitudinal strains of the fiber at each sampling point. In this experiment, we use 117 light pulses with a free-space wavelength of 1,550 nm and a sampling period of  $1 \times 10^{-8}$  s at the 118 optical receiver. Defined by regions of interest, we extract 30,000 channels sampled every 4.08 m 119 along the fiber from 0 to 120 km from the interrogator. Figure 1 shows a map of the DAS array 120 used in our experiment. The DAS data was continuously recorded using 1.55 ms time sampling 121 interval throughout the survey. The spatial sampling interval is 1.02 m, while the gauge length is 122 8.16 m. The backscattered signal strength decays by  $\approx$  0.2 dB/km along the cable, amounting to 123 -40 dB over 100 km. The data were acquired over 44 days from 2020-06-23 to 2020-08-04. 124



Figure 1. Maps of the seabed DAS array. A World map showing the array location in Longyearbyen, the epicenter of the 2020-07-22  $M_{ww}$  7.8 earthquake on the Alaska Peninsula, and the approximate storm locations marked in Figure 4C. B Regional map showing the array and the KBS seismic station. C Local map showing the array annotated with the distance in km from the shore in Longyearbyen, and nearby weather stations.

The phenomena investigated in this study occur below 20 Hz. Therefore, we resample the DAS strain data from 1.55 to 20 ms with the antialiasing filter at 80% of the output Nyquist frequency. The resampled data with the Nyquist frequency of 25 Hz are used in our analysis. Data resampling also reduces the computational cost for analyzing data over a long-time window. We also attenuate interrogator noise that occurs in the whole DAS array, where the noise model is obtained by stacking all the DAS data traces from onshore channels in a calm environment. To
understand the characteristics of the data corresponding to different mechanisms of the oceanbottom vibrations, we compare the processed data with and without the excitation from seismic
waves. With this comparison, we can distinguish the microseisms from other ocean-bottom
pressure responses measured by DAS.

# 135 **2.1 Mechanisms of ocean-bottom vibrations**

The strain of a fiber section will be proportional to that component of the pressure gradient projected along the direction of the fiber. Pressure changes in space and time are therefore detectable by DAS at the seabed, providing they cause strains above the detection limit (due to noise) in the order of 1 nɛ. Ocean-bottom vibrations corresponding to pressure changes can be caused by four excitation mechanisms (Saito & Tsushima, 2016):

- Fluctuation of either sea-surface height or water density causing changes in the oceanbottom loading pressure. This hydrostatic pressure response is called a microseism, and it is associated with ocean surface gravity waves generated by winds and nonlinear wave-wave interaction mechanisms.
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   2. Fluctuation of the vertical seabed placement also causes changes in the ocean-bottom 146
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   3. Hydrodynamic responses associated with the interaction of propagating seismic waves at the seafloor interface between seawater and the solid earth.
- 4. Forces generated by the compressibility of seawater and the elasticity of the oceanbottom rock causes a hydroacoustic response associated with the acoustic resonance of the P-wave propagating within the water layer, resulting in different normal modes.

A primary microseism is driven by ocean surface gravity waves. Hence, the phase velocity ( $c_p$ ) of the primary microseism is given by  $c_p = \omega/k$  with the dispersive relation  $\omega^2 = gk \tanh(kH)$ , where  $\omega = 2\pi f$  is the angular frequency,  $k = 2\pi/\lambda$  is the angular wavenumber,  $g \approx 9.81 \text{ m/s}^2$  is the gravitational acceleration, and *H* is the water depth (Airy, 1841; Craik, 2004). According to linear wave theory (Dean & Dalrymple, 1991, sec. 3.4.4), the dispersive relation for deep water ( $H > 0.5 \lambda$ ) reduces to  $\omega^2 \approx gk$ , while the relation for shallow water ( $H < 0.05 \lambda$ ) reduces to  $\omega^2 \approx gk^2 H$ .

160 Matsumoto et al. (2012) comprehensively describe the frequency ranges of three pressure 161 responses at the seabed: hydrostatic, hydrodynamic, and hydroacoustic. First, for a given 162 wavelength, the frequency of an ocean surface gravity wave occurs in deep water at  $f \approx$ 163  $\sqrt{g/(2\pi\lambda)}$ . We may approximately derive the frequency  $(f_d)$  in deep water as a function of water 164 depth by assuming  $H = 0.5 \lambda$ , which is the lower limit of deep-water depth. Hence, the frequency 165 limit of ocean surface gravity waves in deep water as a function of water depth is approximately 166 defined by

$$f_d = \frac{1}{2} \sqrt{\frac{g}{\pi H}} \ . \tag{1}$$

167 Assuming  $H = 0.05 \lambda$  for the upper limit of shallow water depth, we can derive the frequency  $(f_s)$ 168 of an ocean surface gravity wave in shallow water from  $f \approx \sqrt{gH/\lambda^2}$ :

$$f_s = \frac{1}{20} \sqrt{\frac{g}{H}} . \tag{2}$$

A primary microseism is the direct hydrostatic pressure response onto the seafloor corresponding 169 to an ocean surface gravity wave; hence, it has the same frequency as the ocean surface wave. 170 When primary microseism wave trains (ocean surface gravity waves) propagate in opposite 171 directions (as occurs on reflection from topography, for example), secondary microseisms can be 172 generated by non-linear wave-wave interaction at double the frequency of the primary microseism 173 (Ardhuin & Herbers, 2013). Second, the lower frequency limit of the hydroacoustic responses is 174 governed by the fundamental (the 1<sup>st</sup> order) acoustic resonant frequency as formulated by  $f_1 =$ 175 c/(4H). Acoustic resonant frequencies are expressed as the cut-off (lower limit) frequency for 176 normal modes: 177

$$f_n = \frac{(2n-1)c}{4H} , \qquad (3)$$

where n is a positive integer indicating the order of the normal mode and c is the acoustic velocity in the water (Landrø & Hatchell, 2012). Last, the hydrodynamic responses associated with the seismic waves from an earthquake can be pronounced in a wide frequency range depending on the seismic source and the elastic properties of the subsurface. Unlike hydrostatic and hydroacoustic responses, the frequency of hydrodynamic responses from seismic waves is independent of water depth.

#### 184 **2.2 Origin of ocean swells**

Ocean surface gravity waves are generated by friction exerted by wind on the ocean surface. Propagating waves are generated when the restoration of the fluid to equilibrium is driven by gravity. Wave size depends on wind speed, wind duration and the area over which the wind is blowing (the fetch). Large ocean surface gravity waves generated by storms can propagate for a long distance. These waves are also called ocean swells.

Ocean-bottom seismic sensors can detect ocean swells generated from large storms 190 191 occurring several thousand kilometers away. Bromirski & Duennebier (2002) discuss the 192 amplitude characteristics and wave spectra of these microseisms. The dispersion relation for ocean 193 surface gravity waves in deep water predicts that low-frequency waves will arrive before higherfrequency waves. Also, it depicts the resulting linear up-sweep characteristics of ocean swells in 194 195 spectrograms (time-frequency representations) computed from ocean-bottom seismic data (Bromirski & Duennebier, 2002, fig. 11). Using the method described in Lin et al. (2018) based 196 197 on Munk et al. (1963), we can also trace ocean swells, measured by DAS, back to their originating distant storms. We use the time-frequency gradients measured in spectrograms to calculate the 198 great-circle distances and travel times of the storm-induced ocean swells traveling from the storm 199 centers to the DAS receiver. 200

Lin et al. (2018) derive the expression for the propagation distance of an ocean swell, based on Munk et al. (1963) as

$$x = \frac{g}{4\pi \left(\frac{df}{dt}\right)} , \tag{4}$$

where f is the frequency of the primary microseism associated with an ocean swell. Here, df/dtis the time-frequency gradient or slope of the linear up-sweep trend. Further, the group velocity  $(c_q)$  of an ocean surface gravity wave in deep water can be computed from

$$c_g = \frac{1}{2} \sqrt{\frac{g}{k}} \approx \frac{g}{4\pi f} , \qquad (5)$$

where f is the frequency of the wave. We can, then, estimate the travel time (t) of the ocean swell from the storm center to the DAS receiver from

$$t = \frac{x}{c_g} \ . \tag{6}$$

In short, we firstly estimate the slope of a linear up-sweep trend in the spectrogram and determine the propagation distance of the swell using equation (4). Next, we use the start frequency of the trend in the spectrogram to compute the group velocity and the travel time using equations (5) and (6), respectively. The estimated distance and travel time help to characterize the storms that produced the observed ocean swells.

#### 213 **3 Results and discussion**

#### 214 **3.1 Data characterization**

On 2020-07-22 at 06:12:44 (UTC), an earthquake with a moment W-phase magnitude 215 (M<sub>ww</sub>) of 7.8 occurred at a depth of 28 km approximately 100 km south of the Alaska Peninsula 216 (Figure 1A). The earthquake was detected by seismic stations worldwide and our DAS array near 217 Longyearbyen, which is approximately 5,100 km away from the epicenter on a great circle. At the 218 219 DAS channel at 36 km inline distance from the shore, the approximate arrival times of the P-, Sand SS-waves are respectively at 510, 950 and 1200 s from the earthquake's origin time. We 220 compare the data in the 1200-s time windows before and after 2020-07-22T06:20:02Z, which is 221 the timestamp at 72, 512 and 762 s before the first arrivals of P-, S- and SS-waves, respectively. 222

Figures 2A and 2B show the spectral analysis from 0.01 to 20 Hz of processed DAS strain 223 data from the 1200-s time window before the seismic event. Note that the processed data have a 224 time sampling interval of 20 ms, i.e., the Nyquist frequency is 25 Hz. A power spectrum is 225 computed by a discrete Fourier transform along the time axis of the processed data within the 226 whole 1200-s time window, in which the normalization factor is 1 (unscaled) for the forward 227 transform. The average power spectra over 251 channels (500 m radius) around three selected 228 locations are shown in Figure 2A. The locations are selected to represent the data recorded at 229 different water depths and distances from the shore. The power spectra of individual channels are 230 combined to produce the distance-frequency plot in Figure 2B. There are four energy peaks 231 (Events 1 to 4 marked in Figure 2A) in the spectra within 5 km from shore where the water is less 232 than 100 m deep (see water depth profile in Figure 3D). Only two of these energy peaks are present 233 where the water depth is greater than 100 m. The input strain data used in the spectral analysis are 234 given in Figure S1 for reference. 235

The 0.06-Hz peak (Event 1 in Figure 2A) present in all water depths corresponds to primary microseisms associated with ocean surface gravity waves, excited by either local winds or distant storms. The frequency of the primary microseisms visible in the figure has weak or no correlation with water depth and the waves are seen to propagate towards the shore, illustrated by the temporal-spatial correlations in Figure 3C. Figure 3C also shows that in shallow water near the shore, reflected waves propagate in the opposite direction. These two wave trains with the same frequency, propagating in opposite directions, generate the 0.12-Hz energy peak harmonic in the power spectra (Event 2 in Figure 2A) as discussed in Hasselmann (1963). Hence, the 0.12-Hz energy peak is a secondary microseism, generated only in shallow water, of twice the fundamental primary microseism frequency. These reflected waves are not apparent in water deeper than 100 m, so the secondary microseism is only observable near shore with water depth <100 m.</p>



Figure 2. Spectral analysis of DAS strain data. The average power spectra at selected locations along the cable (A) and the power spectral profile (B) immediately before the earthquake, without strong seismic energy, computed from the 1200-s time window before 2020-07-22T06:20:02Z. The average power spectra at the same locations (C) and the power spectral profile (D) with seismic energy from the earthquake (P-, S- and SS-waves), computed from the 1200-s time window after 2020-07-22T06:20:02Z. The average power spectra in A and C are computed over 251 recording channels (500 m radius) around each location. The numeric annotations in A and C highlight key events discussed in the text. The colored triangles in B and D mark the locations associated with the power spectra in A and C. The frequency limits corresponding to ocean surface gravity waves with water depth >200 m ( $f_d$ ) and with water depth <10 m ( $f_s$ ), and normal modes ( $f_1$  and  $f_2$ ) are also plotted in B and D.



**Figure 3. Band-limited DAS strain data.** The strain data filtered to three frequency bands: A 1.2–20 Hz, **B** 0.2–1.2 Hz, and **C** 0.005–0.2 Hz. **D** The water depth profile. The recording time starts at 2020-07-22T06:20:02Z, after which the first P-wave from the 2020-07-22  $M_{ww}$  7.8 earthquake on the Alaska Peninsula arrives at about 72 s. The colored triangles mark the locations associated with the spectra shown in Figure 2.

The energy peak between 1.0 and 1.2 Hz (Event 3 in Figure 2A) also only exists in shallow water. We believe that it represents Scholte or other seismic waves that are excited locally in shallow water by ocean-bottom pressure variations as a result of ocean surface gravity waves. The last energy peak (Event 4 in Figure 2A) is associated with the fundamental acoustic resonance ( $f_1$ ) of the water column, with frequency >1 Hz, varying with water depth. Note that this fundamental mode exists in all water depths for our DAS array.

In addition to the four features discussed above, there is a slight increase in energy around 253 0.36 Hz (Event 5 in Figure 2A) in the power spectra. An increase in energy around 0.36 Hz has 254 been reported as the secondary microseism associated with the 0.18-Hz opposing surface gravity 255 wave groups in a seabed DAS experiment in Belgium by Williams et al. (2019, fig. 2). In our data, 256 we observe no energy peak around 0.18 Hz; hence, it is unlikely that the 0.36-Hz energy observed 257 as Event 5 in Figure 2A is directly involved with ocean surface gravity waves. In addition, its 258 frequency is not close to the frequency limits of ocean swells or acoustic resonance in the water 259 column. Thus, we believe that Event 5 corresponds to hydrodynamic responses associated with 260 seismic waves, although we cannot identify their seismic origins. 261

Figures 2C and 2D show the spectral analysis of DAS strain data that contain strong seismic 262 waves (P-, S- and SS-waves) from the 2020-07-22 M<sub>ww</sub> 7.8 earthquake on the Alaska Peninsula. 263 Here, the responses caused by these seismic waves arriving at the seafloor significantly boost the 264 strain power in the frequency range below 4 Hz. This energy is superimposed on the initial ambient 265 levels shown in Figures 2A and 2B. Comparing the spectra with and without earthquake-related 266 energy, we see that the energy peaks corresponding to primary and secondary microseisms (Events 267 1 and 2) have similar characteristics. We also observe no significant change to the Scholte wave 268 response associated with ocean surface gravity waves (Event 3). In contrast, the water-layer 269 acoustic resonance (Event 4) is enhanced, and its 2<sup>nd</sup> order mode is visible, when excited by the 270 energy from the strong seismic waves. Therefore, we conclude that seismic waves from the 271 272 earthquake propagate into the seawater, causing stronger acoustic resonance in the water column. The earthquake-related responses enhance the power spectra between 0.01 and 4 Hz, and their 273 274 spectra have no correlation with water depth, as shown in Figures 2C and 2D. In addition, we observe that earthquake-related responses near the shore are stronger than in deeper water, which 275 276 we suspect to be due to an amplitude attenuation effect from different subsurface rock types. Accordingly, all the energy peaks outside the frequency ranges of primary microseisms (Event 1), 277 secondary microseisms (Event 2), Scholte waves (Event 3), and water-layer acoustic resonance 278 (Event 4) are related to hydrodynamic responses from earthquakes. The responses discussed above 279 have different characteristics in both temporal and spatial dimensions (see Figures S2 and S3 in 280 Supporting Information for detailed analyses in the frequency-wavenumber and frequency-281 velocity domains). 282

The theoretical cut-off frequencies for the water-layer acoustic resonance shown in Figures 283 2B and 2D do not match the power spectral energy distribution, especially in shallow water at 0-284 6 km from shore. In shallow water, the frequency of the acoustic resonance, which ranges between 285 1 and 4 Hz, is lower than the corresponding theoretical cut-off frequency. This energy must 286 therefore be evanescent or associated with acoustic resonance modes between the sea surface and 287 a strong reflector below the seafloor soft sediments. The amplitude of evanescent waves typically 288 decays rapidly. However, the energy peaks we observe have strong amplitudes. Therefore, they 289 are unlikely associated with evanescent modes. We believe that these energy peaks are more likely 290 related to propagating acoustic resonances, because their characteristics are similar to the acoustic 291 resonances in deeper water, e.g., the higher order modes are enhanced by seismic waves from an 292 earthquake. Thus, it is likely that these energy peaks in shallow water are caused by acoustic 293 resonance between the sea surface and a strong reflector below the seafloor. A possible candidate 294

is the Base of the Helvetiafjellet Formation which is approximately 200 m below the sea surface
(Bælum et al., 2012, fig. 6). To explore this possibility, we would need detailed knowledge of the
geological structure of this horizon, which is beyond the scope of this article.

Band-pass filtering the DAS strain data into three bands gives us the results illustrated in Figure 3, in which the first P-wave from the 2020-07-22 M<sub>ww</sub> 7.8 earthquake on the Alaska Peninsula arrives at about 72 s. Therefore, Figure 3 reveals the characteristics of the seabed DAS data with and without earthquake-related responses.

Figure 3A shows the data in the frequency band from 1.2 to 20 Hz, which mostly comprise signals from water-column acoustic resonance (Event 4 in Figure 2). These responses are further enhanced after 72 s by the arrival of strong P-waves from the Alaska earthquake. In addition, the times of the acoustic resonance events shown in Figure 3A vary with water depth. Therefore, these events are not direct responses to the seismic P-waves, that are generally independent of water depth.

Figure 3B shows the data in the frequency band from 0.2 to 1.2 Hz. Here we see a series 308 of strong P-waves from the Alaska earthquake arriving after 72 s. These P-waves are coherent and 309 almost flat on the data profile. The direct signals from P-waves are independent of water depth. 310 Note that we also observe weak P-waves before 72 s that represent seismic events from unknown 311 sources that form a slight increase in energy around 0.36 Hz (Event 5 in Figure 2). We also see 312 scattered events in shallow water (<100 m water depth) throughout the recording. These events 313 represent Scholte waves or other seismic waves that are excited locally by ocean-bottom pressure 314 315 variation due to ocean surface gravity waves (Event 3 in Figure 2).

In Figure 3C, the data in the frequency band from 0.005 to 0.2 Hz, we see right-dipping 316 events all along the cable. These are primary microseisms corresponding to ocean surface gravity 317 waves, propagating towards the shore (Event 1 in Figure 2). In addition, we observe left-dipping 318 events near the shore where the water depth is <100 m, creating a 'checkerboard' pattern. These 319 are ocean surface gravity waves that are reflected from the shoreface back to the ocean. The (non-320 linear) superposition of long-wavelength ocean surface gravity waves and their reflections in 321 shallow water near the shore creates secondary microseisms (Event 2 in Figure 2). Based on our 322 observation, the seismic waves from the earthquake do not change the characteristics of primary 323 and secondary microseisms. 324

# **325 3.2 Ocean wave monitoring**

326 Figure 4 shows spectrograms from selected DAS receivers at different water depths and 327 distances from the shore. Figures 4A-C show the linear up-sweep trends of different primary microseism events ranging from 0.04 to 0.1 Hz. Their frequencies monotonically increase with 328 time. These linear trends correspond to the ocean swells produced by distant storms. Over the 329 entire period of recording, we identify 12 linear trends in the spectrograms. Most of them last 330 between 50 and 100 hours, and can overlap in time and space. The amplitude level of these linear 331 332 trends increases towards the inner parts of the fjord, i.e., the shore in Longyearbyen. Hence, the ford appears to act as a 'narrowing amplifier' for the ocean swells produced by distant storms. In 333 shallow water (Figure 4A), we observe stronger amplitudes of primary microseism signals, 334 especially for the more locally-generated ones (with steeper gradients). While we do not expect to 335 see secondary microseisms (at double the frequency of the primary microseism) in deep water 336 (>100 m) at distances greater than 6 km on the DAS array, we do not see them even at 3.08 km 337

along the cable, with an average water depth of 71 m. This is unexpected and remains to be

339 understood. We believe that the secondary microseism should be strong enough to be seen in water

 $340 \quad \text{depths} < 100 \text{ m}.$ 

Figure 4D shows the maximum speed of local winds measured at the Isfjord Radio and 341 Svalbard Lufthavn weather stations near the DAS array. The Isfjord Radio station is located at the 342 entrance of the fjord and close to the 55 km distance point along the DAS array, whereas Svalbard 343 Lufthavn station is located at the Svalbard Airport in Longyearbyen and close to the start of the 344 DAS array (see Figure 1C). We found no correlation between the local wind speeds and the 345 primary microseisms associated with ocean swells. Therefore, we deduce that the primary 346 microseisms visible in the spectrograms are mostly generated from winds or storms outside the 347 fjord. 348

Four linear up-sweep trends of the primary microseisms corresponding to distant storms 349 are highlighted in the spectrogram in Figure 4C. Using the methods described in the previous 350 section, we can calculate the distance and time taken by the ocean swell to travel from each storm 351 center to the DAS array. Table 1 summarizes the calculation of the four storms as marked in Figure 352 4C. By applying geographical and topological constraints (there must be an open seaway between 353 our DAS array and the source) we can retrieve their approximate locations for comparison with 354 public records. The Arctic Ocean is isolated from other oceans by land. The Fram Strait, which 355 lies between Svalbard and Greenland, is the only deep passage into the Arctic Ocean. In addition, 356 the main orientation of our DAS array points towards the Atlantic Ocean. Therefore, the primary 357 358 microseisms detected by our DAS array are likely produced by storms in the Atlantic Ocean. It is unlikely that our DAS data are dominated by strong primary microseisms caused by storms in the 359 Pacific Ocean through the shallow Bering Strait. 360

From public records, we can trace all the four linear trends in Figure 4C back to their 361 corresponding storms in the Atlantic Ocean. Event 1 corresponds to the Tropical Storm Edouard 362 near Bermuda at about 4,100 km away from Longyearbyen, occurring from 2020-07-04 to 2020-363 07-06 (Pasch, 2021). Event 2 possibly corresponds to the bomb cyclone in offshore south Brazil 364 at about 13,000 km from Longyearbyen from 2020-06-30 to 2020-07-02 as reported in Gobato & 365 366 Heidari (2020) and Khalid et al. (2020). According to weather news in Iceland (Cirić, 2020), Event 3 should correspond to an extratropical depression between Iceland and Greenland at about 2,400 367 km away from the DAS array from 2020-07-15 to 2020-07-17. Lastly, Event 4 probably comes 368 from a storm in a remote region in offshore south Brazil at about 11,000 km from the DAS array 369 on 2020-07-12. 370



**Figure 4. Spectrograms for storm monitoring.** Spectrograms at 3.08 (A), 11.56 (B) and 66.71 (C) km along the DAS array from shore. Maximum wind speeds measured at Isfjord Radio and Svalbard Lufthavn weather stations (see Figure 1) are shown in **D**. Four storm events marked in C are discussed in the text. All the spectrograms are computed from the average power spectrum over 251 recording channels (500 m radius) around the selected locations within a 300-s time window on an hourly basis. In the spectrograms, the yellow vertical stripes are caused by dynamic range saturation, which is weakly correlated with the local storm noise from the winds illustrated in **D**, whereas the white vertical stripes indicate drop-out periods in the real-time data transfer.

Table 1. Estimated of gins of the four occan swens marked in figure 10.					
Parameters	Event 1	Event 2	Event 3	Event 4	
Start time at DAS $(t_0)$	2020-07-07	2020-07-10	2020-07-18	2020-07-19	
	T00:00:00Z	T00:00:00Z	T12:00:00Z	T00:00:00Z	
End time at DAS $(t_1)$	2020-07-09	2020-07-16	2020-07-20	2020-07-25	
	T12:00:00Z	T00:00:00Z	T00:00:00Z	T00:00:00Z	
Frequency at start time $(f_0)$	0.046 Hz	0.045 Hz	0.044 Hz	0.038 Hz	
Frequency at end time $(f_1)$	0.087 Hz	0.076 Hz	0.086 Hz	0.075 Hz	
Travel distance ( $x$ in equation (4))	4,113 km	13,055 km	2,409 km	10,938 km	
Group velocity for the lowest- frequency swell ( $c_g$ in equation (5) with $f = f_0$ )	16.97 m/s	17.35 m/s	17.74 m/s	20.54 m/s	
Travel time for the lowest- frequency swell ( <i>t</i> in equation (6))	67.32 hours	209.03 hours	37.71 hours	147.89 hours	
Estimated time at source $(t_0 - t)$	2020-07-04	2020-07-01	2020-07-16	2020-07-12	
	T04:40:00Z	T06:58:00Z	T22:17:00Z	T20:06:00Z	

Table 1. Estimated origins of the four ocean swells marked in Figure 4C.

# **371 3.3 Future oceanographic applications**

Functioning marine ecosystems are vital to healthy oceans on which a sustainable future 372 on Earth for all living beings ultimately depends (Danovaro et al., 2020). Marine acoustics plays 373 an important role in studying physical processes in the oceans and their interaction with the solid 374 earth, atmosphere and living organisms. Therefore, Passive Acoustic Monitoring (PAM) is 375 recognized as an important surveillance tool for the Earth's ecosystems, through the studies of 376 ocean ambient sound, marine mammal behavior, glacial/iceberg noise, anthropogenic ocean use, 377 unsanctioned nuclear or other polluting activity, earthquake and tsunami warning, in addition to 378 search and rescue. 379

We have shown that DAS, as a PAM system, can detect waves from various sources 380 through dynamic interactions between the atmosphere, ocean, and solid earth. DAS has many 381 valuable attributes to offer the oceanographic community, nicely complementing existing sensing 382 systems such as satellites (which are broadly limited to very near-surface observations), buoys, 383 moorings, and floats (which have limited spatial coverage and resolution). The advantages of DAS 384 include broadband and high-resolution spatial and temporal measurement capacities, with data 385 available in real-time to support active marine management and decision-making. The real-time 386 capability, bringing data from the seafloor, is unmatched by any other system other than fixed 387 installations cabled to shore or supporting long lines to surface buoys, both of which represent 388 expensive and complex engineering challenges. The potential for earthquake and tsunami warning 389 systems alone is therefore remarkable. This sensing network is also possible to create at low cost, 390 since we can use existing submarine telecommunication cables. These cables span more than a 391 392 million kilometers around all the oceans on the globe, potentially bringing a sensing capability to many less-sampled environments, and perhaps also able to support less developed countries in 393 responsibly managing their maritime resources. 394

Thus, DAS brings an innovative and game-changing new sensing modality to oceanography and planetary observation systems in general. Therefore, we believe that DAS will become a valuable new component of the Global Ocean Observing System (GOOS), of the Intergovernmental Oceanographic Commission (IOC) of UNESCO, as discussed in Howe et al. (2019).

#### 400 4 Conclusion

401 DAS in an ocean-bottom telecommunication cable can measure various types of oceanbottom pressure responses that are caused by dynamics in the atmosphere, ocean, and solid earth. 402 They comprise the responses from ocean surface gravity waves causing primary and secondary 403 microseisms, Scholte waves, water-layer acoustic resonances, and seismic waves (P-, S- and SS-404 waves) from earthquakes. We clearly describe and compare their characteristics in the DAS data. 405 Our interpretations are validated by redundant samples from the data acquired extensively in 406 407 spatial and temporal dimensions, over 44 days along 120 km of a fiber-optic cable, which extends along the ford across different water depths from 0 to 400 m. We observe primary microseisms 408 from distant storms, their reflections from the shore in shallow water and the resulting non-linear 409 wave-wave interaction, forming secondary microseisms. We also see an approximate correlation 410 between hydroacoustic first mode energy and the theoretical cutoff, but this is not supported in the 411 nearshore, shallow water, leading us to suspect that the energy may be associated with a mode 412 resonating between the sea surface and a deeper rigid structure, rather than soft unconsolidated 413 sediment. More detailed geological knowledge, beyond the scope of this paper, would be required 414 415 to explore this possible explanation. The DAS data do enable us to trace several primary microseisms associated with ocean swells back to their storm origins, which are significant ocean-416 atmosphere disruptions occurring up to 13,000 km away. We also find that the fjord acts as a 417 'narrowing amplifier' for microseisms, because their amplitudes increase towards the inner parts 418 of the fjord. Thus, it is possible to use DAS data acquired over 120 km to study dynamic 419 interactions between the atmosphere, ocean, and solid earth. Thanks to its high spatial and temporal 420 resolution, real-time data availability, broadband low frequency sensitivity and its ability to sense 421 what is happening close to the seabed, capturing both hydroacoustic and seismic events, DAS 422 423 offers great scientific value to Earth observation systems. We believe that DAS will become a key value sensing modality in the Global Ocean Observing System (GOOS). 424

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# 435 **Open research**

DAS data for this research are available in Taweesintananon & Landrø (2022) via
 <u>https://doi.org/10.18710/VPRD2H</u>. Seismic data from the KBS seismic station in Svalbard used

as our reference are available through IRIS web services: https://service.iris.edu/. Details on the 438 2020-07-22 Mww 7.8 earthquake on the Alaska Peninsula are available at USGS web site: 439 https://earthquake.usgs.gov/earthquakes/eventpage/us7000asvb/executive. The weather data are 440 available through the Norwegian Center for Climate Services (NCCS) at 441 https://seklima.met.no/observations/. 442

# 443 Author contributions

ML, SEJ, JKB, AH, OS and FS conceived and designed the experiment. AH and FS collected data. KT processed data and prepared the visualizations. KT and ML analyzed data with support from SEJ, JRP, RAR, LB and HJK. ML and JRP validated research outputs, acquired funding, and managed the project. KT wrote the original draft of the manuscript. All the authors conducted review & editing of the manuscript.

# 449 **Competing interests**

450 There is no competing interest related to this work.

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### Journal of Geophysical Research: Oceans

#### Supporting Information for

# Observation of atmospheric and oceanic dynamics using ocean-bottom distributed acoustic sensing

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# **Contents of this file**

Figures S1 to S3

# Additional Supporting Information (Files uploaded separately)

None

#### Introduction

The supporting information comprises the figures illustrating the data profiles and their additional spectral analyses to support the discussion in the manuscript.



**Figure S1. DAS strain data filtered to 0.005–20 Hz. A** The recording profile from 2020-07-22T06:00:02Z without strong seismic energy from known earthquake corresponding to the spectra in Figures 2A & 2B. **B** The recording profile from 2020-07-22T06:20:02Z with seismic energy (P-, S- and SS-waves arriving at 72, 512 and 762 s, respectively) from the 2020-07-22 M<sub>ww</sub> 7.8 earthquake on the Alaska Peninsula corresponding to the

spectra in Figures 2C & 2D. **C** The water depth profile. The colored triangles mark the locations associated with the spectra shown in Figure 2.



**Figure S2. Frequency-wavenumber spectra of DAS strain data. A & B** The power spectra in the *f-k* domain around two channels from the DAS recording profile from 2020-07-22T06:00:02Z without strong seismic energy from known earthquake corresponding to the data in Figure S1A. **C & D** The power spectra in the *f-k* domain around the same channels as in A & B from the DAS recording profile from 2020-07-22T06:20:02Z with seismic energy (P-, S- and SS-waves) from the 2020-07-22 M<sub>ww</sub> 7.8 earthquake on the Alaska Peninsula corresponding to the data in Figure S1B. All the *f-k* spectra are computed over 1001 recording channels (2 km radius) around the selected locations. The wavenumber is based on the distance along the fiber-optic cable from the shore in Longyearbyen—positive wavenumber for waves propagating to the ocean, and negative wavenumber for waves propagating toward the shore.



**Figure S3. Frequency-velocity spectra of DAS strain data. A & B** The power spectra in the *f*-*v* domain around two channels from the DAS recording profile from 2020-07-22T06:00:02Z without strong seismic energy from known earthquake corresponding to the data in Figure S1A. **C & D** The power spectra in the *f*-*v* domain around the same channels as in A & B from the DAS recording profile from 2020-07-22T06:20:02Z with seismic energy (P-, S- and SS-waves) from the 2020-07-22 M<sub>vwv</sub> 7.8 earthquake on the Alaska Peninsula corresponding to the data in Figure S1B. All the *f*-*v* spectra are computed over 1001 recording channels (2 km radius) around the selected locations. Positive velocity corresponds to the wave propagation to the ocean, whereas negative velocity corresponds to the wave propagation to spatial Nyquist sampling beyond which aliasing artifacts are shown ( $f_{Nyquist,k} = vk_{Nyquist}$ , where  $k_{Nyquist} = 1/(2 \Delta x)$ ).