

# Ionospheric Plasma Depletions at Mars Observed by the MAVEN spacecraft

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## Abstract

The Martian ionosphere, modulated by the solar wind from the topside and remnant crustal magnetic fields close to the surface, possesses unique structures different from Earth and Venus. Integrated observations by the plasma and magnetic field instruments onboard the Mars Atmosphere and Volatile Evolution (MAVEN) spacecraft show evidence of ionospheric plasma depletions, independent of seasonal variations at Mars. During such depletions, the density of all ionospheric ion species is reduced by more than an order of magnitude, and, at the same time, the electron temperature increases abruptly. An automated algorithm for the identification of such plasma depletions is developed. Altogether, as many as 1177 events are identified in 8618 orbits available from October 2014 to May 2021. A statistical investigation of these events reveals that they are more prominent on the nightside. Their higher occurrence in the southern hemisphere suggests a possible relation to the crustal magnetic fields. While the dayside events occur mainly at altitudes above about 200 km, nightside event altitudes are typically lower. Considering the relation between spacecraft velocity and observed event duration, we suggest that the depletions are bubble-like structures, more elongated horizontally than vertically. A possible mechanism of their formation is discussed.

1 **Ionospheric Plasma Depletions at Mars Observed by**  
2 **the MAVEN spacecraft**

3 **Praveen Basuvaraj, František Němec, Zdeněk Němeček, Jana Šafránková**

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5 **Key Points:**

- 6 • Density depletions larger than an order of magnitude are observed for all major  
7 ionospheric ion species and electrons.  
8 • The depletions are more prominent on the nightside, mainly in the southern hemi-  
9 sphere, suggesting a relation to crustal magnetic fields.  
10 • While most dayside depletions occur at altitudes above 200 km, their altitudes on  
11 the nightside are typically below 250 km.

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**Abstract**

The Martian ionosphere, modulated by the solar wind from the topside and remnant crustal magnetic fields close to the surface, possesses unique structures different from Earth and Venus. Integrated observations by the plasma and magnetic field instruments onboard the Mars Atmosphere and Volatile Evolution (MAVEN) spacecraft show evidence of ionospheric plasma depletions, independent of seasonal variations at Mars. During such depletions, the density of all ionospheric ion species is reduced by more than an order of magnitude, and, at the same time, the electron temperature increases abruptly. An automated algorithm for the identification of such plasma depletions is developed. Altogether, as many as 1177 events are identified in 8618 orbits available from October 2014 to May 2021. A statistical investigation of these events reveals that they are more prominent on the nightside. Their higher occurrence in the southern hemisphere suggests a possible relation to the crustal magnetic fields. While the dayside events occur mainly at altitudes above about 200 km, nightside event altitudes are typically lower. Considering the relation between spacecraft velocity and observed event duration, we suggest that the depletions are bubble-like structures, more elongated horizontally than vertically. A possible mechanism of their formation is discussed.

**Plain Language Summary**

Ions and electrons embedded in an atmosphere of a planet constitute the ionospheric plasma. At Mars, the ionospheric plasma has a well-defined altitude structure similar to Earth. However, we find that the normally smooth plasma density profiles exhibit at times significant sudden plasma depletions. We use nearly 7 years of data obtained by the Mars Atmosphere and Volatile Evolution (MAVEN) spacecraft to systematically analyze the occurrence and properties of such events. We show that, during the events, the density of all monitored ion species decreases at least ten times and, at the same time, the temperature increases. The events are more frequent and occur at lower altitudes on the nightside than on the dayside. We show that the event geometry corresponds to bubble-like structures rather than to layers and/or columns of decreased density. We discuss several possible mechanisms of their formation.

**1 Introduction**

Mars is an unmagnetized planet with a thin layer of atmosphere compared to Earth. Due to the precipitation of energetic solar charged particles (mostly electrons) and the absorption of incoming solar radiation (EUV and X-rays), the neutral atmosphere of Mars gets ionized to form the dayside ionosphere composed of ions and free electrons (Krasnopolsky, 2002; Němec et al., 2010). The dayside ionosphere is photochemically controlled below the exobase (about 220 km) (Mendillo et al., 2017). On the other hand, the nightside ionosphere is formed - apart from the impact ionization by precipitating energetic particles - due to the plasma transport from the dayside (Fox et al., 1993; Lillis et al., 2009; Lillis & Fang, 2015).

First in-situ observations of the Martian neutral atmosphere and charged ionosphere were performed by the Viking Landers 1 and 2 through their Neutral Mass Spectrometer (NMS) and Retarding Potential Analyzer (RPA) instruments (Owen & Biemann, 1976; Soffen, 1976; McElroy et al., 1976). Nier and McElroy (1976) used the NMS data to demonstrate that  $CO_2$  is the dominant neutral atmospheric species representing about 95% of the total atmosphere of Mars with few traces of  $N_2$ ,  $Ar$ , and  $O$ . The RPA data revealed that the Martian ionosphere is populated by  $O_2^+$  ion followed by  $CO_2^+$ ,  $O^+$ , and  $NO^+$  (Hanson et al., 1977). Under a photochemical equilibrium of the dayside ionosphere, the average peak ionization altitude is about 125 km with a peak electron density of about  $1.5 \times 10^5 \text{ cm}^{-3}$  (e.g., Withers, 2009; Němec, Morgan, Gurnett, Duru, & Truhlík, 2011; Andrews, Andersson, et al., 2015, and references therein).

62 The structure and composition of the ionosphere are highly variable due to the cou-  
 63 pling with the neutral atmosphere and plasma environment (Bougher, Cravens, et al.,  
 64 2015; Bougher, Pawlowski, et al., 2015). The absence of a magnetic field allows for a di-  
 65 rect interaction of the ionosphere with the solar wind plasma, thereby influencing the  
 66 ionospheric dynamics (Acuña et al., 1998; Mitchell et al., 2001). The localized remnant  
 67 crustal magnetic sources dispersed over the Martian crust are insignificant in balancing  
 68 the solar wind dynamic pressure on the dayside. Still, they considerably affect the for-  
 69 mation of the nightside ionosphere. They may, in particular, inhibit the plasma trans-  
 70 port from the dayside (Němec et al., 2010) and control the regions where energetic par-  
 71 ticles precipitate and form localized ionization patches (Němec, Morgan, Gurnett, & Brain,  
 72 2011; Lillis et al., 2011; Girazian et al., 2017). Ionospheric ion and electron densities are  
 73 related to the solar activity, solar zenith angle, precipitating energetic particle fluxes, and  
 74 neutral abundances (Withers et al., 2012; Lillis et al., 2015; Sánchez-Cano et al., 2016).  
 75 In addition, orbit-to-orbit ionospheric variability, as well as the variability of altitude-  
 76 density profiles within individual orbits, have been reported (Duru et al., 2008; Bougher,  
 77 Jakosky, et al., 2015).

78 Mitchell et al. (2001) noticed suprathermal electron flux depressions in the Mars  
 79 Global Surveyor (MGS) Electron Reflectometer data, which represent sudden drops in  
 80 the fluxes by at least two orders of magnitude, and they introduced the term “plasma  
 81 voids” for such depletions. These voids are seen within the nightside closed crustal mag-  
 82 netic field lines at altitudes below 200 km (Xu et al., 2017). The Mars Advanced Radar  
 83 for Subsurface and Ionospheric Sounding (MARSIS) instrument (Picardi et al., 2005)  
 84 on board the Mars Express spacecraft has revealed the existence of electron depressions  
 85 in the nightside ionosphere (Duru et al., 2011). However, the spacecraft orbit spatially  
 86 constrained this study to altitudes above 275 km. The average width of electron depres-  
 87 sions estimated from 66 full orbits is found to be 950 km.

88 Withers et al. (2012) utilized the electron density profiles, retrieved from the Mars  
 89 Express Radio Science Experiment (MaRS), to reveal peculiar features present at times  
 90 in the upper ionosphere, indicating major departures from conditions in models that pro-  
 91 duce a Chapman layer-like shape. Depressions in the suprathermal electrons that vary  
 92 with altitude and solar illumination conditions were studied using Mars Express, MGS,  
 93 and MAVEN probes (Steckiewicz et al., 2015; Hall et al., 2016; Steckiewicz et al., 2017).  
 94 Collinson et al. (2020) reported the existence of sporadic E-like structures (layers and  
 95 rifts) in the Martian dayside ionosphere, contributing to the sudden increase or deple-  
 96 tion of plasma density within the atmospheric dynamo region (all layers and rifts were  
 97 encountered below 223 km). It was suggested that they are, analogically to the situation  
 98 at Earth, formed in the dynamo region; the ions are dragged by neutral winds across the  
 99 magnetic field lines, resulting in electromagnetic forces and eventually, depending on the  
 100 configuration, in the layer/rift formation.

101 In the present study, we use nearly 7 years (about 3.5 Martian years) of MAVEN  
 102 spacecraft measurements to systematically check for the presence of ionospheric plasma  
 103 depletions and to analyze their occurrence and properties as a function of relevant pa-  
 104 rameters. The used data set is described in section 2. An automated method used for  
 105 the event identification is described in section 3. The results obtained are presented in  
 106 section 4 and discussed in section 5. Finally, section 6 contains a brief summary of the  
 107 main findings.

## 108 2 Data Set

109 The Mars Atmosphere and Volatile Evolution (MAVEN) probe reached Mars on  
 110 21 September 2014. Its primary objectives are to estimate the atmospheric loss rate, to  
 111 determine the composition of neutral and ionized constituents, and to reveal the major  
 112 parameters controlling their loss through time (Lillis et al., 2015; Jakosky et al., 2015).

113 The MAVEN spacecraft orbits Mars along an elliptical orbit with a nominal periapsis  
 114 of about 150 km to sample the upper atmospheric region and its orbital period is about  
 115 4.5 h. With an orbital inclination of  $75^\circ$ , the spacecraft observations can provide a quasi-  
 116 global coverage of Mars (Jakosky et al., 2015). We utilized the plasma and magnetic mea-  
 117 surements in the Martian ionosphere from three onboard instruments.

118 The Neutral Gas and Ion Mass Spectrometer (NGIMS) instrument is a quadrupole  
 119 mass spectrometer that samples the atmospheric neutrals and ions in the mass range of  
 120 2-150 Da (Mahaffy, Benna, King, et al., 2015). NGIMS operates at altitudes below 500 km.  
 121 Each inbound/outbound orbital part of the NGIMS measurements lasts for about 600 s  
 122 with a vertical resolution of about 5 km for targeted species (Bougher, Jakosky, et al.,  
 123 2015; Mahaffy, Benna, Elrod, et al., 2015). In our study, we use the NGIMS data from  
 124 18 October 2014 to 14 May 2021, comprising of 8618 orbits in total. Number densities  
 125 of the major ionospheric constituents  $O^+$ ,  $O_2^+$ ,  $NO^+$ , and  $CO_2^+$  are used.

126 The Langmuir Probe and Waves (LPW) instrument consists of two identical 40 cm  
 127 probe sensors mounted on a pair of about 7 m long booms (Andersson et al., 2015). These  
 128 sensors are used to measure current-voltage characteristics that are, in turn, used to cal-  
 129 culate the electron density and temperature of the local plasma environment. The time  
 130 resolution of the measurements is 4 s and the sensitivity of the instrument is  $10^2$ – $10^6$   $\text{cm}^{-3}$   
 131 for the electron density and 500–50000 K (*i.e.*,  $\approx 0.04$  to 4.3 eV) for the temperature. We  
 132 use the electron density ( $N_{e^-}$ ) and electron temperature ( $T_{e^-}$ ) measurements in concu-  
 133 rrence with the ion density measurements by NGIMS.

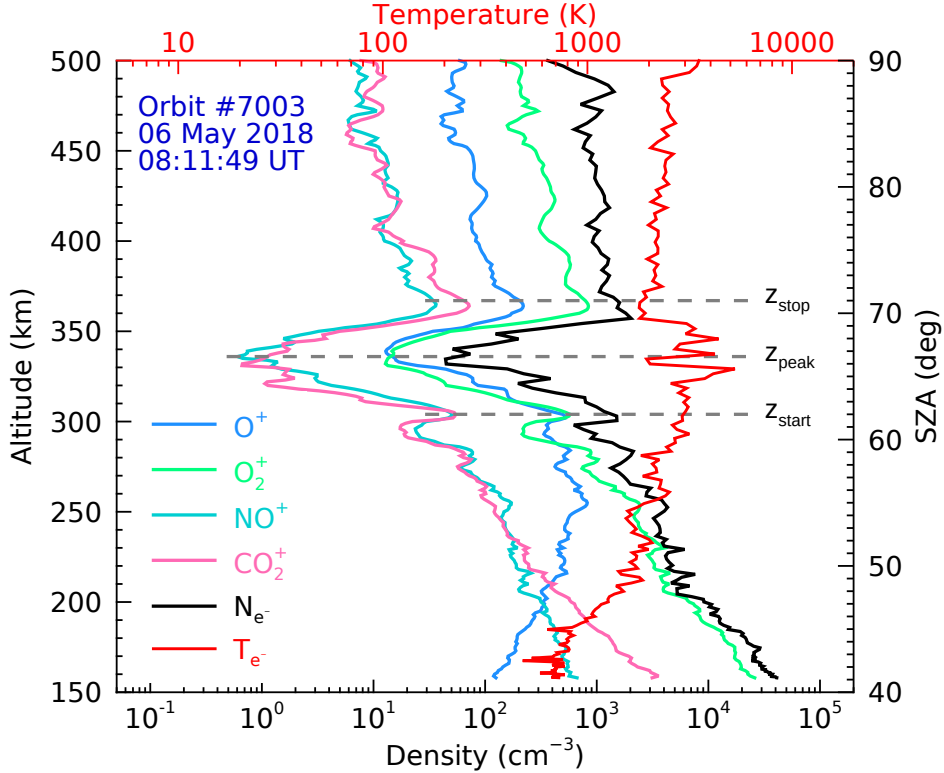
134 The Magnetometer (MAG) instrument comprises of twin three-axial flux-gate mag-  
 135 netometers mounted at the tip of the spacecraft solar panels. The MAG instrument pro-  
 136 vides vector measurements of the magnetic field along the spacecraft trajectory at an  
 137 intrinsic sampling rate of 32 vector samples per second (Connerney et al., 2015). The  
 138 dual magnetometers were designed to operate at a dynamic range of about 60,000 nT  
 139 with resolution up to 0.05 nT. The calibrated vector magnetic field measurements in the  
 140 planetocentric coordinate system with a time resolution of 4 s are used in our study.

### 141 3 Identification of Plasma Depletion Events

142 An example of a plasma depletion event (PDE) is shown in Figure 1. Altitudinal  
 143 profiles of ion densities ( $O^+$ ,  $O_2^+$ ,  $NO^+$ ,  $CO_2^+$ ) measured by the NGIMS instrument dur-  
 144 ing the MAVEN outbound orbit #7003 on 06 May 2018 are shown by the color curves.  
 145 Additionally, the electron density and temperature profiles measured by the LPW in-  
 146 strument are indicated by the black and red curves, respectively. The plasma depletion  
 147 event is seen at altitudes between about 304 and 367 km, with the peak depletion roughly  
 148 in the middle of the interval (336 km altitude). The start, peak, and stop altitudes of the  
 149 event are marked by horizontal dashed lines and labeled  $z_{start}$ ,  $z_{peak}$ , and  $z_{stop}$ , respec-  
 150 tively.

151 At altitudes around  $z_{peak}$ , number densities of all analyzed ionospheric ion species  
 152 are by more than an order of magnitude lower than at altitudes below/above. Similar  
 153 depletion is also observed in electron densities. This may be expected, as the ionospheric  
 154 plasma is quasi-neutral. However, given that the ion and electron densities are measured  
 155 by two independent instruments, this represents a nice confirmation that the effect is in-  
 156 deed real and not instrumental. Additionally, the electron temperature is found to in-  
 157 crease considerably during the event.

158 Showing an example PDE event and demonstrating its principal characteristics,  
 159 we aim to systematically identify such events in the MAVEN data. The events are typ-  
 160 ically visually well identifiable in individual orbital plots similar to that in Figure 1, and  
 161 their list might be thus, in principle, prepared manually. However, we develop an auto-



**Figure 1.** Ionospheric altitude profiles observed by the NGIMS and LPW instruments during orbit #7003 (outbound) on 06 May 2018. Plasma number densities of individual ion species and electron density are shown by the color and black curves, respectively. The altitudinal profile of the electron temperature is shown by the red curve using the scale at the top. A plasma depletion event, exhibiting a clear sudden decrease of plasma densities and an increase of the electron temperature, is seen at altitudes between about 304 km ( $z_{start}$ ) and 367 km ( $z_{stop}$ ). The peak depletion is observed at an altitude of about 336 km ( $z_{peak}$ ).

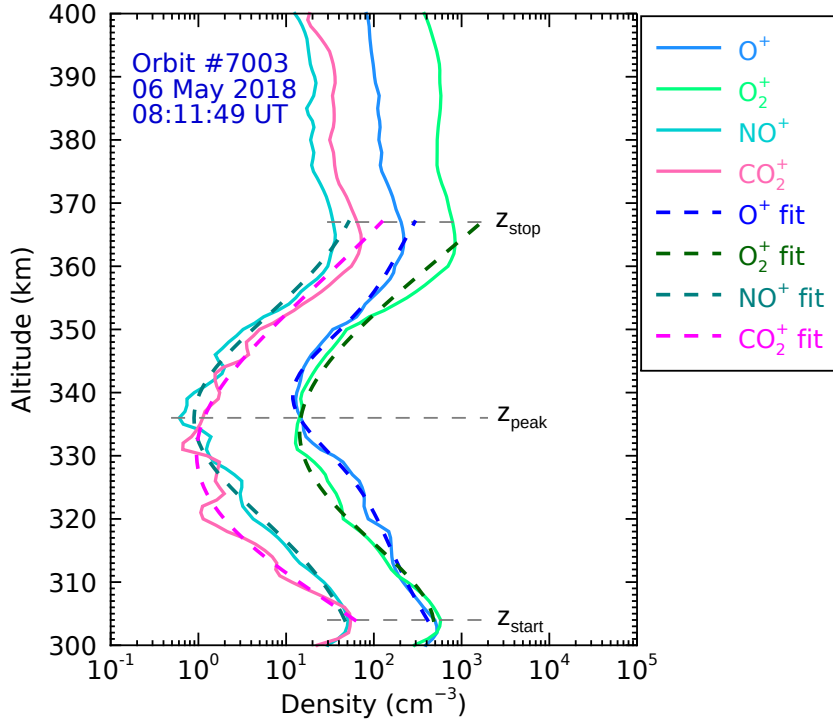
162 matic identification routine, which has the advantage of exact quantitative criteria and,  
163 at the same time, it determines the values  $z_{start}$ ,  $z_{peak}$ , and  $z_{stop}$ .

164 The procedure uses NGIMS measured density profiles of all major ionospheric ions  
165 ( $O^+$ ,  $O_2^+$ ,  $NO^+$ ,  $CO_2^+$ ). It assumes that the altitude range of PDE (depletion width here-  
166 after) is between 5 and 80 km, which safely encompasses the width ranges of manually  
167 pre-identified events used for the routine development.

168 An identified PDE is required to have an average (over the four major ion species)  
169 peak density depletion larger than 90%:

$$170 \quad \frac{1}{2N} \sum_{i=1}^N \left[ \frac{n_i(z_{peak})}{n_i(z_{start})} + \frac{n_i(z_{peak})}{n_i(z_{stop})} \right] \leq 0.1 \quad (1)$$

171 where  $n_i(z)$  is the number density of the ionospheric ion  $i$  at an altitude  $z$  and  $N$   
172 is the total number of ionospheric ion species analyzed. In our case,  $N=4$  corresponds  
173 to  $O^+$ ,  $O_2^+$ ,  $NO^+$ , and  $CO_2^+$  ions.



**Figure 2.** A zoomed view of ion density profiles from Figure 1, highlighting the depletion event as observed by the NGIMS instrument, is shown by the solid curves. The altitudes  $z_{start}$ ,  $z_{peak}$ , and  $z_{stop}$  are identified using an automatic routine (see text). Gaussian fits of individual ion density profiles in between these altitudes are shown by the dashed curves.

174 Altitudes of NGIMS measurements range from the periapsis up to about 500 km,  
 175 where the ionopause is periodically encountered (Benna et al., 2015). Considering this  
 176 altitude limitation, we apply the aforementioned PDE identification criteria to the data  
 177 measured at altitudes between 150 and 500 km during all the analyzed 8618 orbits. For  
 178 each identified event, the altitudes  $z_{start}$ ,  $z_{peak}$ , and  $z_{stop}$  are determined.

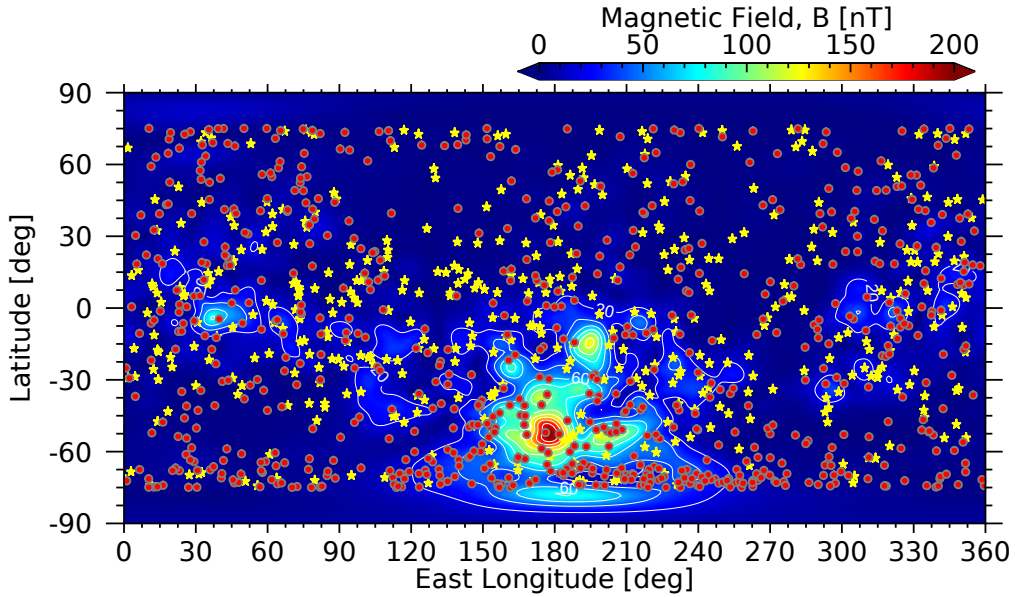
179 Figure 2 depicts a zoomed view of Figure 1, focusing on altitudes around the PDE  
 180 and demonstrating the performance of the developed identification routine. We notice  
 181 that the shapes of individual ion density profiles are similar over the depletion region,  
 182 regardless of the actual ion abundances. Gaussian fits of ion density profiles in the alti-  
 183 tude range between  $z_{start}$  and  $z_{stop}$  are over-plotted by dashed curves. The fits are based  
 184 on the logarithm of ion densities rather than on the densities themselves. They include  
 185 a linear background term to account for the gradual density decrease with the altitude.  
 186 This allows us to define (separately for each ion density profile) a density ratio correspond-  
 187 ing to the maximum plasma depletion as  $r_N = n_{obs}/n_{est}$ , where  $n_{obs}$  corresponds to  
 188 the peak of the Gaussian fit and  $n_{est}$  corresponds to the respective linear term.

189 We note that up to three distinct PDEs are identified at different altitudes in some  
 190 orbits. These are considered as separate events hereinafter. Altogether, as many as 1177  
 191 PDEs events are identified in 1060 half (inbound/outbound) orbits.

## 4 Results

### 4.1 Locations of Plasma Depletion Events

Figure 3 shows the locations of the detected PDEs over the Martian surface. The loci of PDEs depicted in the figure correspond to the altitude of peak depletions ( $z_{peak}$ ). The events observed during the day and night are plotted by yellow stars and red circles, respectively. Model crustal magnetic field magnitude at an altitude of 400 km based on the Cain et al. (2003) magnetic field model is color-coded using the scale at the top. The remnant crustal fields are typically larger in the southern hemisphere than in the northern hemisphere, but there is no apparent relation to the event locations.



**Figure 3.** Locations of identified plasma depletion events. The daytime and nighttime events are shown as yellow stars and red circles, respectively. Model crustal magnetic field magnitude at 400 km (Cain et al., 2003) is color-coded using the scale at the top.

The total number of events and corresponding normalized occurrence rates (*i.e.*, number of events per hour of measurement) in the northern and southern hemispheres are given in Table 1. It can be seen that the events are more frequently observed on the nightside than on the dayside. Additionally, while the event occurrence is roughly the same in the northern and southern hemispheres during the day, the nighttime occurrence rate is significantly higher in the southern hemisphere.

The normalized occurrence rate (*i.e.*, number of events per hour of measurement) of PDEs is investigated as a function of possible controlling parameters in Figure 4. The occurrence rates as a function of local solar time (LST) and solar zenith angle (SZA) are investigated in Figures 4a and 4c, respectively. It can be seen that the probability of observing a depletion event during the nighttime is about 2–3 times higher than during the daytime, in agreement with the results in Table 1. The normalized event occurrence rate as a function of the altitude is depicted in Figure 4b. Although the events occur all over the analyzed altitudinal range, their occurrence seems to be somewhat lower at low (< 170 km) and high (> 450 km) altitudes. The dependence of the event occurrence rate

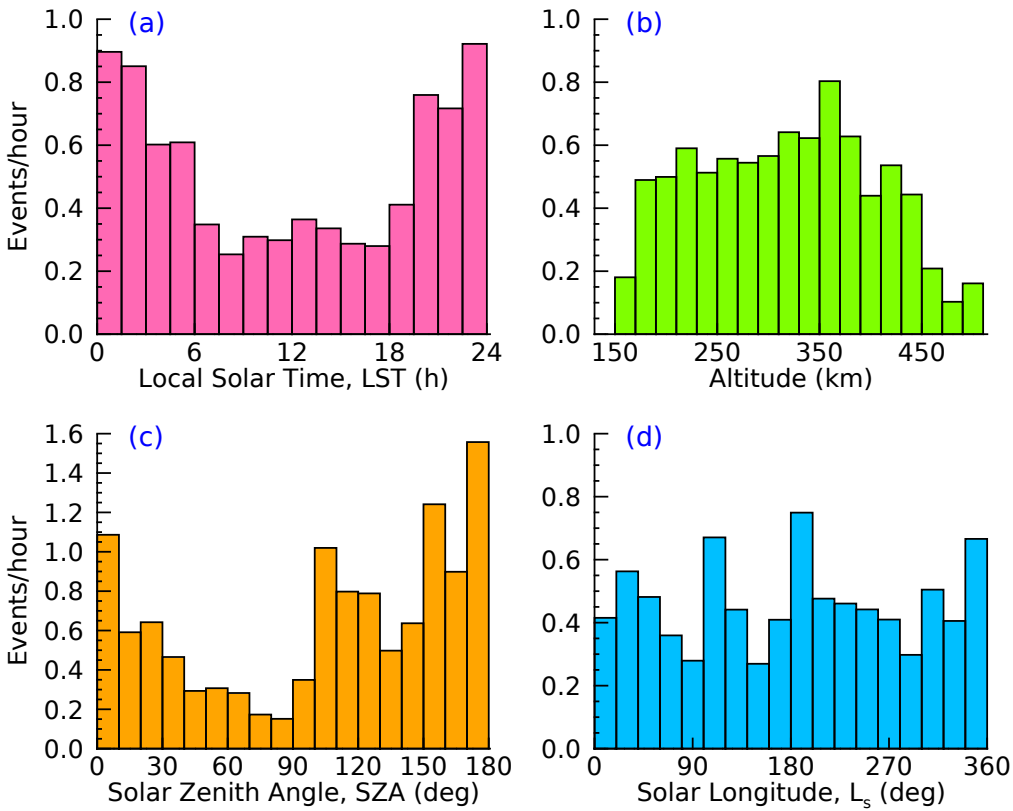


**Table 1.** Occurrence of plasma depletion events in northern/summer hemispheres

Parameter	Dayside A(B)	Nightside A(B)
Northern Hemisphere	205 (0.31)	275 (0.52)
Southern Hemisphere	224 (0.26)	473 (0.91)

A–number of events; B–normalized occurrence rate ( $\text{h}^{-1}$ )

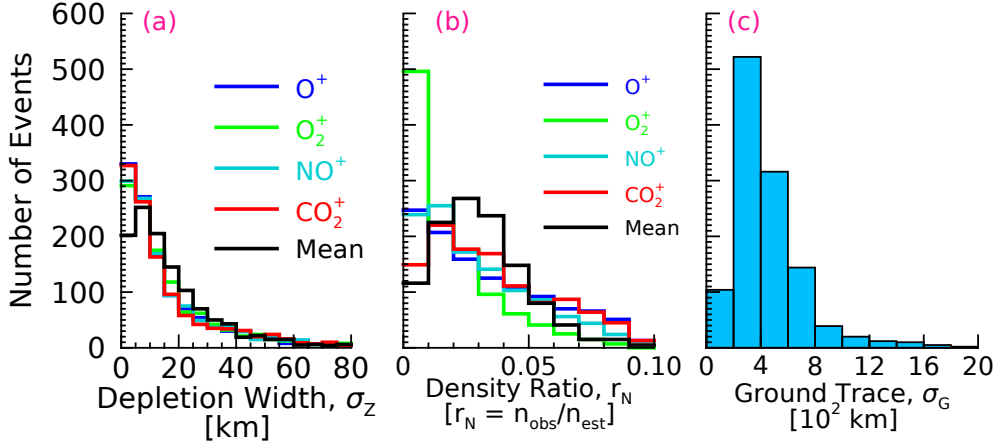
216 on the solar longitude shown in Figure 4d does not reveal any clear trend, demonstrat-  
 217 ing that the events occur relatively equally throughout all the seasons.



**Figure 4.** Normalized occurrence rates (number of events per hour) as a function of: (a) local solar time ( $LST$ ), (b) altitude, (c) solar zenith angle ( $SZA$ ), and (d) solar longitude ( $L_s$ ).

218 Histograms of basic parameters of PDEs are shown in Figure 5. The altitudinal range  
 219 of plasma depletion events ( $\sigma_Z$ ) is analyzed in Figure 5a, separately for each ion species.  
 220 The used bin size is 5 km. It is found that the depletion width is quite the same for all  
 221 the analyzed ion species. The mean depletion width is overplotted by the black line. About  
 222 69% of PDEs have a depletion width less than 20 km. The density ratios ( $r_N$ ) of ana-  
 223 lyzed ion species are shown in Figure 5b. A larger reduction in the ion number densi-  
 224 ties corresponds to lower values of  $r_N$ . The color lines show the results obtained for in-  
 225 dividual ion species, and the black line shows the mean ion density ratio. It can be seen

226 that while the distributions of density ratios obtained for  $O^+$ ,  $NO^+$ , and  $CO_2^+$  are roughly  
 227 the same, the density ratios obtained for the most abundant  $O_2^+$  ion tend to be often much  
 228 lower. A histogram of horizontal extents of PDEs ( $\sigma_G$ ) is shown in Figure 5c. It can be  
 229 seen that the ground trace of most events (about 84%) is between 200 and 800 km. For  
 230 all identified depletion events,  $\sigma_G > \sigma_Z$ . We note, however, that larger  $\sigma_Z$  does not al-  
 231 ways correspond to larger  $\sigma_G$  and vice versa.



**Figure 5.** Histograms of parameters of plasma depletion events: (a) depletion thickness ( $\sigma_Z$ ), (b) density ratios of ionospheric ions ( $r_N$ ), and (c) the ground trace ( $\sigma_G$ ).

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## 4.2 Event Properties

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Properties of PDEs as a function of solar zenith angle are shown in Figure 6. For ionospheric altitudes (above 120 km), the nightside ionosphere is illuminated up to SZA  $\approx 105^\circ$  and hence the actual nightside observations start at SZAs somewhat higher than  $90^\circ$  (Lillis et al., 2009; Dubinin et al., 2016). This threshold SZA can be clearly identified in all panels of Figure 6, as the number of identified PDEs increases significantly on the nightside.

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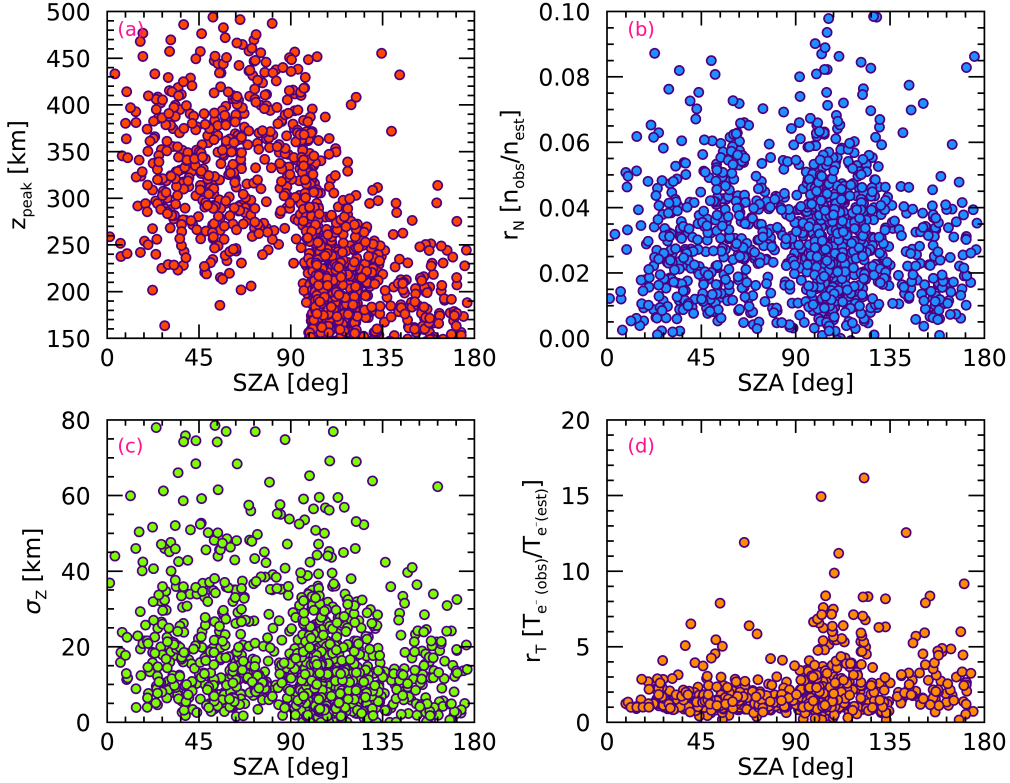
The peak depletion altitudes depicted in Figure 6a exhibit a sudden change at the dayside/nightside boundary, they are almost exclusively above 200 km on the dayside. On the other hand, the peak depletion altitudes on the nightside are generally lower, below about 250 km for the vast majority of the events. Solar zenith angle dependences of average number density ratios and depletion thicknesses depicted in Figures 6b and 6c, respectively, do not reveal any clear trend. The ratio between the electron temperature measured at the peak depletion altitude and the expected electron temperature at the same height without the depletion event is termed the electron temperature ratio ( $r_T$  in Figure 6d). For most dayside PDEs (66%), the electron temperature ratio is between 1 and 2, and it does not seem to depend on SZA. The electron temperature ratios obtained for the nightside PDEs exhibit considerably larger scatter. They are still larger than 1 for most events (88% of events), but the values larger than 2 are rather common (46% of events).

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Figure 7 shows the variations of peak depletion altitudes, magnetic field magnitudes at the times of the events, and electron temperature ratios with mean ion density ratio. The results obtained for the dayside and nightside events are plotted separately in the



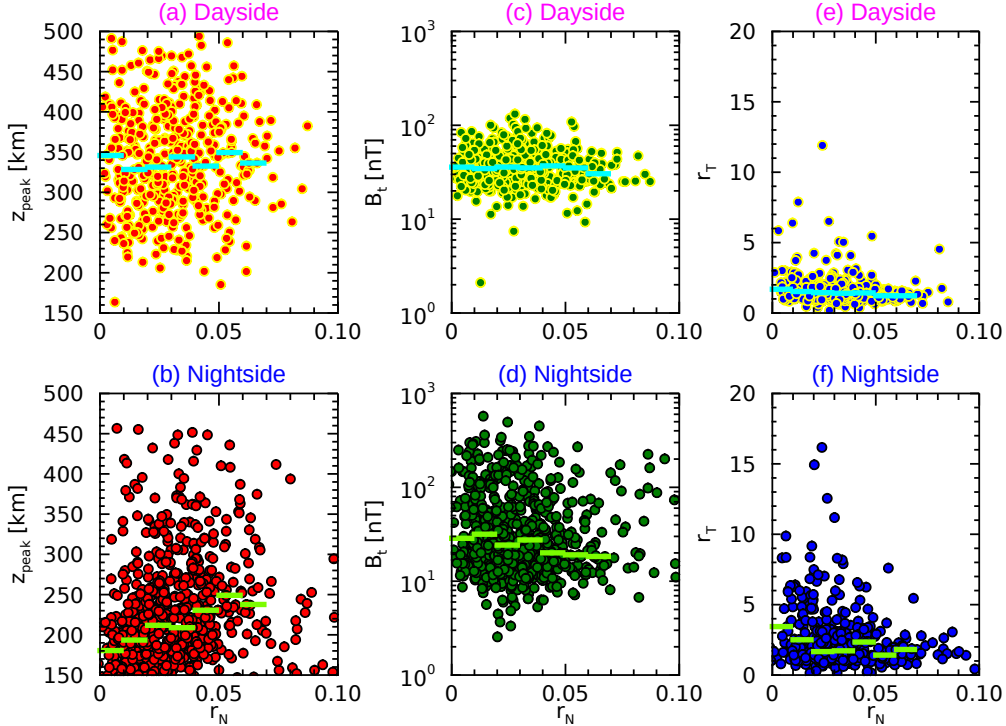
**Figure 6.** Parameters of individual identified plasma depletion events as a function of solar zenith angle: (a) peak depletion altitudes ( $z_{peak}$ ), (b) ion density ratio ( $r_N$ ), (c) depletion thickness ( $\sigma_z$ ), and (d) electron temperature ratio ( $r_T$ ).

255 top and bottom panels, respectively. Typical values of dayside peak depletion altitudes  
 256 depicted in Figure 7a are about 350 km, and they do not show any appreciable variations  
 257 with the plasma depletion ratio. In contrast, nightside peak depletion altitudes are cor-  
 258 related with the density ratio; more significant density depletions occur preferentially at  
 259 lower altitudes. We also note a considerable spread in  $z_{peak}$  at altitudes above about 250 km.

260 Figures 7c and 7d indicate that the magnetic field magnitude measured at the times  
 261 of the events is essentially independent of the peak density ratio, both on the dayside  
 262 and nightside. It is, however, noteworthy that the spread of magnetic field magnitudes  
 263 observed during the nightside events is somewhat larger than during the dayside events.  
 264 This is believed to be related to the crustal magnetic field influence and generally lower  
 265 altitudes of nightside events. We also note that the magnetic field measured at the times  
 266 of the events is typically almost horizontal both on the dayside and nightside. Although  
 267 the scatter of electron temperature ratios is quite large, in particular on the nightside,  
 268 Figures 7e and 7f indicate that they tend to be somewhat larger for more significant plasma  
 269 depletions.

### 270 4.3 Geometry of the Events

271 The detection of PDEs at all ionospheric altitudes might suggest that the deple-  
 272 tions are vertical column-like structures. Nevertheless, their observed horizontal extents  
 273 are generally larger than vertical extents, suggesting a possibility of a horizontal layer

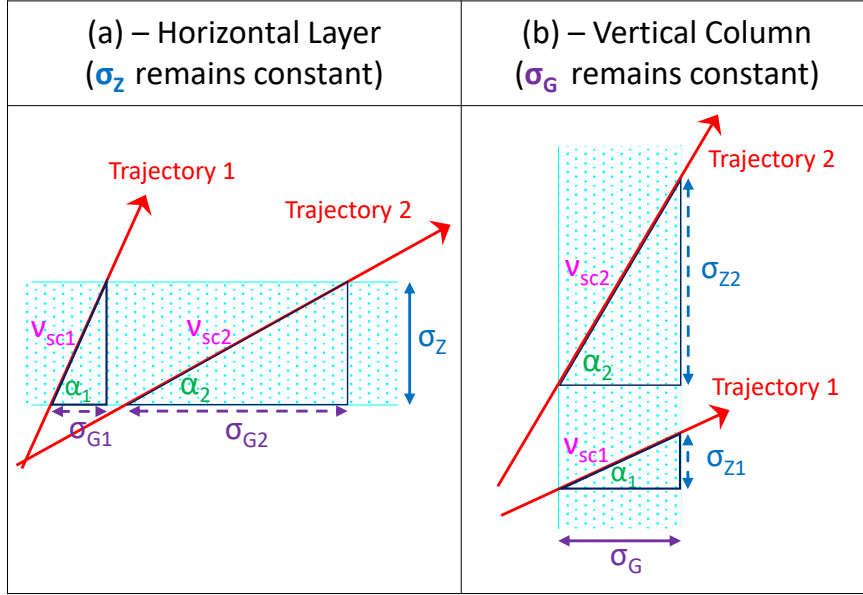


**Figure 7.** Parameters of individual identified PDEs as a function of ion density ratio. The results obtained for the dayside and nightside events are plotted in the top and bottom panels, respectively. (a)-(b) Peak depletion altitude ( $z_{peak}$ ). (c)-(d) Magnetic field magnitude ( $B_t$ ). (e)-(f) Electron temperature ratio ( $r_T$ ). The overplotted horizontal lines mark the median values in individual intervals.

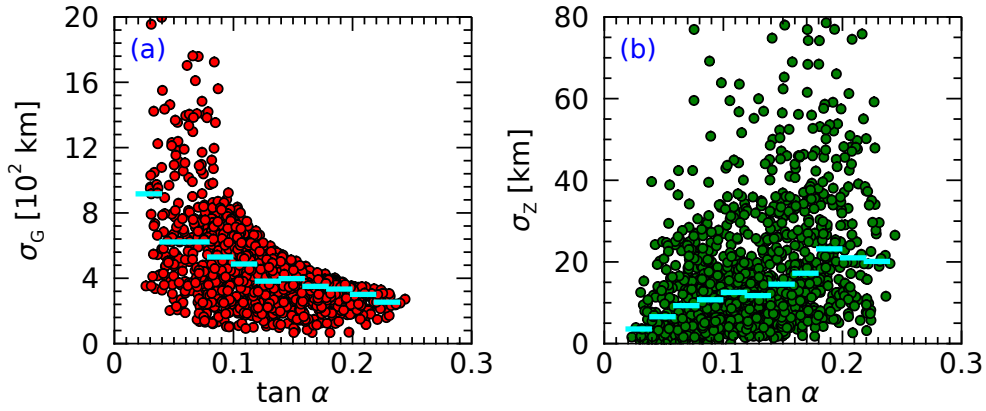
274 geometry. Given that only single-point measurements are available, it is difficult to dis-  
 275 tinguish these two scenarios experimentally. However, we attempt to use the fact that  
 276 the orbital inclination ( $\alpha = \arctan \nu_r / \nu_h$ , where  $\nu_r$  and  $\nu_h$  are radial and horizontal com-  
 277 ponents of the spacecraft velocity, respectively) varies from event to event to tackle this  
 278 issue.

279 Schematic views of the two extreme geometry scenarios corresponding, to a hor-  
 280 izontal layer depletion and a vertical column depletion are shown in Figures 8a and 8b  
 281 respectively. Two different spacecraft trajectories, with two different inclination angles  
 282  $\alpha$ , are shown in each figure. Assuming that the real PDEs would have a form of hori-  
 283 zontal layers (Figure 8a), the observed vertical extents of the events would be independ-  
 284 ent of the orbital inclination ( $\sigma_Z \approx const.$ ). Then, the observed horizontal extents would  
 285 depend on the orbital inclination and were proportional to  $1/\tan \alpha$ . On the other hand,  
 286 in a case of PDEs having a form of vertical columns (Figure 8b), the observed hori-  
 287 zontal extents of the events would be independent of the orbital inclination ( $\sigma_G \approx const.$ ),  
 288 and the vertical extents would be proportional to  $\tan \alpha$ .

289 Plotting the horizontal and vertical extents of the events as a function of the orb-  
 290 ital inclination should thus, in principle, allow distinguishing between the two geom-  
 291 etry scenarios. This is done in Figures 9a and 9b, respectively. Each point corresponds  
 292 to a single identified PDE, and the horizontal bars show the median values calculated  
 293 in individual  $\tan \alpha$  bins. It can be seen that the horizontal extent of the events  $\sigma_G$  is lower



**Figure 8.** Schematic plots of two possible extreme scenarios of the event geometry. (a) Horizontal layer. (b) Vertical column. Two spacecraft trajectories with different orbital inclinations ( $\alpha_1, \alpha_2$ ) are shown for each scenario, along with their respective ground ( $\sigma_G$ ) and vertical ( $\sigma_z$ ) projections.



**Figure 9.** (a) Horizontal and (b) vertical extents of individual events as a function of tangent of orbital inclination. The overplotted horizontal lines mark the median values in individual intervals.

294 for larger orbital inclinations  $\alpha$ , as would be expected for a horizontal layer geometry.  
 295 However, at the same time, the vertical extent  $\sigma_z$  increases with  $\tan \alpha$ , indicative of the  
 296 vertical column geometry. This suggests that neither of the two extreme geometry sce-  
 297 narios is correct, with the actual event geometry possibly anywhere in between.

298 Thus, PDEs appear to be bubble-like ionospheric structures. Also, the arguably  
 299 stronger dependence obtained for  $\sigma_G$  in Figure 9a than for  $\sigma_z$  in Figure 9b may suggest

that the real scenario is somewhat closer to the idealized case of a horizontal layer. This is in agreement with  $\sigma_G > \sigma_Z$  observed in Figure 4.

## 5 Discussion

Large-scale electron density depressions (ionospheric holes) are known to exist in the nightside ionosphere of Venus at altitudes above about 200 km already for decades (Brace et al., 1980, 1982; Luhmann & Russell, 1992). While the exact mechanism driving the formation of these holes remains unclear, it is inferred that strong radial magnetic fields observed within the depression regions obstruct the neighboring plasma from replenishing these holes. Such ionospheric holes are thus strongly controlled by the magnetic pressure and the orientation of the draped interplanetary magnetic field (Collinson et al., 2014). By contrast, the Martian PDEs analyzed in our study occur both during the daytime and nighttime, and, on the nightside, they are mostly seen at altitudes below 200 km. The characteristics disparate from the Venusian ionospheric holes suggest that Martian PDEs are likely governed by different mechanisms. Moreover PDEs seems to be connected with remnant crustal magnetic field that is virtually absent at Venus.

Suprathermal electron depressions (referred to as plasma voids or electron holes) in the induced magnetosphere and ionosphere of Mars have been studied using MGS (Mitchell et al., 2001), Mars Express (Soobiah et al., 2006; Duru et al., 2011; Hall et al., 2016), and MAVEN (Steckiewicz et al., 2015, 2019). Plasma voids are proposed to be the direct consequence of the presence of rotating stronger crustal magnetic field regions (Mitchell et al., 2001). Considering that the energies of suprathermal electrons are much higher than the energies of thermal electrons and plasma we analyze in the present study, it might be difficult to draw any direct analogies. However, plasma voids are expected to occur below 200 km in the nightside ionosphere with closed crustal field lines (Xu et al., 2016, 2017). On the other hand, PDEs are found to often occur at higher altitudes on the dayside, with virtually undetectable magnetic field changes as compared to plasma voids. Hall et al. (2016) surveyed the electron holes in the induced magnetosphere of Mars in the energy range of 20–200 eV using the data from the MGS spacecraft, reporting most events in the altitude bin of 300 to 400 km. Although this altitudinal range is arguably comparable to the altitudes of dayside PDEs, the analyzed energies are still way larger than those of the thermal plasma. Duru et al. (2011) proposed that electron density depression events at the nightside may be due to the solar wind perturbations and plasma instabilities in the ionosphere. This can cause turbulence and possibly the streamers and/or plasma clouds at high altitudes.

According to Collinson et al. (2020), sudden reductions of the plasma densities may be explained by the displacement of plasma resulting from the collisions with neutral winds ( $\vec{v}_{wind}$ ) and traversing across the magnetic field lines ( $\vec{B}$ ), resulting in the electromagnetic forces ( $\vec{v}_{wind} \times \vec{B}$ ). Their preliminary non-statistical study based on 34 events concluded that such a sporadic E-like structures occur primarily in the southern hemisphere with stronger crustal magnetic fields, near the magnetic field twist regions induced by the solar wind. Although the existence of sporadic E-like structures might explain small-scale increase/decrease in the electron density measurements (Withers et al., 2012), it does not seem to entirely account for the disparity in the dayside and nightside depletions, large-scale ground traces, and event occurrence also well above the Martian dynamo region. Keskinen (2018) argued that ionospheric plasma irregularities caused by the ion drifts (resulting from the neutral winds and magnetic field) are plausible only in the dynamo region at altitudes of about 125–250 km. In our case, almost all the dayside PDEs occur at altitudes well above.

The Martian dayside ionosphere is photochemically controlled below 200 km. In contrast, the plasma transport and diffusion dominate above these altitudes (Andrews, Edberg, et al., 2015). Additionally, the presence of strong crustal magnetic fields can af-



351 fect both electron densities and electron temperatures at altitudes above 200 km (Dubinin  
 352 et al., 2016). Nearly all identified daytime PDEs are in the plasma transport dominant  
 353 region. In a transport-controlled region, increased electron temperatures can result in  
 354 higher electron densities due to enhanced upward plasma flows (Flynn et al., 2017). Ergun  
 355 et al. (2016) demonstrated that a higher electron temperature at these altitudes could  
 356 develop strong enough ambipolar electric field to drift the major ions ( $O^+$ ,  $O_2^+$ ,  $CO_2^+$ ),  
 357 displacing them spatially and facilitating the ion escape. In addition, an increased plasma  
 358 temperature at the times of an enhanced solar wind dynamic pressure can result in plasma  
 359 clouds (Zhang et al., 2021), possibly removing the plasma from other regions. Neverthe-  
 360 less, the insignificant increase in electron temperature ratio eliminates the possibilities  
 361 of plasma cloud formation and ion escape processes.

362 At altitudes below 300 km, the most abundant ion is  $O_2^+$  (Benna et al., 2015). At  
 363 higher altitudes, its density decreases exponentially with the scale height similarly to other  
 364 species, except  $O^+$  (Withers, 2009). Due to its significant abundance,  $O_2^+$  ion has been  
 365 envisaged to be the principal mechanism of oxygen loss in the Martian atmosphere up  
 366 to 500 km. It has also been demonstrated that the loss due to a possible electric field-  
 367 induced drift is more effective than the dissociative recombination loss of  $O_2^+$  in com-  
 368 parison to other ions (Ergun et al., 2016). This suggests that the more profound deple-  
 369 tions noticed in the  $O_2^+$  ion densities could be due to additional loss processes through  
 370 different channels apart from the PDE event formation mechanism itself.

371 Most nightside PDEs are seen below 250 km with large-scale ground traces and con-  
 372 siderably increased electron temperatures. This may be due to the lack of a photochemically-  
 373 dominant region, which is characteristic for the dayside ionosphere. In the dayside pho-  
 374 tochemical region, any plasma possibly removed in any way would be quickly replenished  
 375 by photoionization, rendering the formation of depletion events very complicated and  
 376 explaining possibly the lack of PDEs at low altitudes on the dayside. We note that the  
 377 recombination coefficients are generally lower for higher electron temperatures (Schunk  
 378 & Nagy, 2009; Dubinin et al., 2016), eliminating the possibility that the PDEs could be  
 379 formed simply due to faster recombination. Since the nightside ionosphere of Mars is,  
 380 at least close to the terminator, formed mainly due to plasma transported from the day-  
 381 side (Němec et al., 2010; Lillis & Fang, 2015; Girazian et al., 2017), large number of PDEs  
 382 detected close to the terminator can be perhaps attributed to the related ionospheric abun-  
 383 dances. Additionally, the larger number of PDEs detected on the nightside than on the  
 384 dayside may be perhaps related to their comparatively longer time of life, as the plasma  
 385 sources on the nightside are extremely limited, and their eventual replenishing by the  
 386 plasma may thus take much longer.

387 Brain et al. (2007) used MGS observations at  $\approx 400$  km to reveal the magnetic topol-  
 388 ogy of Mars. The magnetic field lines are configured as draped (connected within the  
 389 solar wind), open (connecting solar wind and the crustal magnetic field), and closed (con-  
 390 nected within the crustal sources). Their study demonstrated that the northern hemi-  
 391 sphere of Mars is dominated by draped/open field lines. In contrast, the magnetic field  
 392 in the southern hemisphere is likely closed within the collisional atmosphere and some-  
 393 times connected to regions far from the strong crustal sources. Cusp-like configurations  
 394 are noticed in areas with large radial magnetic components. The draped fields are al-  
 395 most horizontal at low altitudes close to the terminator regions (Dubinin et al., 2016;  
 396 Weber et al., 2021). Below 200 km, the magnetic field lines are mostly closed and hor-  
 397 izontal at nightside (Xu et al., 2016) which can weaken the electron-impact ionization  
 398 process (Adams et al., 2018). At higher altitudes, draped and open field lines are observed  
 399 more often above the weakly magnetized regions, which can connect to the nightside iono-  
 400 sphere even below the exobase (Weber et al., 2021). While the dayside PDEs occur at  
 401 altitudes corresponding to draped and open field lines, most of nightside PDEs are iden-  
 402 tified within the closed field lines relating their origins to the crustal sources. This is con-  
 403 sistent with the dominant horizontal magnetic component observed during the events.

404 The PDEs may be in a sense analogous to equatorial plasma bubbles encountered  
 405 at Earth. These are sudden electron density depletions observed in the nightside equa-  
 406 torial ionosphere. Observations have confirmed the existence of plasma bubbles below  
 407 the peak ionization region (about 200 km) and their eventual raising to altitudes up to  
 408 about 1700 km (Oya et al., 1986; Otsuka et al., 2002, 2021). Rayleigh-Taylor instabil-  
 409 ity is considered to be the most likely cause of the bubble formation at Earth (Kil, 2015;  
 410 Oliveira et al., 2020). An uplifting plasma bubble can bifurcate due to an apparent shift  
 411 in its central part. The stability of plasma bubbles depends on the electric field gener-  
 412 ated within (Otsuka et al., 2002).

## 413 6 Conclusions

414 We used MAVEN spacecraft data to identify and analyze plasma depletion events  
 415 in the Martian ionosphere. These depletion events are distinguished by a sudden drop  
 416 in ion number densities by at least an order of magnitude. At the same time, the elec-  
 417 tron temperature typically increases, especially during the nightside events. The mag-  
 418 netic field during the event observations is mostly horizontal, and it does not exhibit any  
 419 significant variations.

420 Altogether, 1177 plasma depletion events are identified in the data measured from  
 421 18 October 2014 to 14 May 2021 using an automated identification algorithm. A sta-  
 422 tistical study of their occurrence reveals a strong solar zenith angle dependence, with more  
 423 events occurring in the nightside ionosphere. Additionally, the nightside events occur pre-  
 424 dominantly in the southern hemisphere, suggesting their relation to the crustal magnetic  
 425 fields. The dayside events appear mainly in the transport-dominant region at altitudes  
 426 above about 200 km, while nightside events occur mostly at altitudes lower than 250 km.

427 It has been statistically found that the most abundant  $O_2^+$  ion is depleted the most,  
 428 while the depletions in other major ionospheric ions are somewhat smaller and roughly  
 429 the same. The depletions likely have a form of bubble-like structures, with their ground  
 430 traces being much larger than the altitudinal extents.

431 Although the depletion events may appear similar to the equatorial ionospheric plasma  
 432 bubbles observed at Earth, it remains questionable whether they are formed by the same/similar  
 433 mechanism. The exact formation mechanism of the depletion events thus remains un-  
 434 known and will be addressed in the future.

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 439 archived in NASA's Planetary Atmospheres Node ([https://atmos.nmsu.edu/PDS/data/  
 440 PDS4/MAVEN/ngims\\_bundle/12/](https://atmos.nmsu.edu/PDS/data/PDS4/MAVEN/ngims_bundle/12/)). The MAVEN LPW and MAG data products (Level-  
 441 02, In situ Key Parameters) used in this study are publicly archived in NASA's Plane-  
 442 tary Plasma Interaction Node ([https://pds-ppi.igpp.ucla.edu/search/view/?f=  
 443 yes&id=pds://PPI/maven.insitu.calibrated/data](https://pds-ppi.igpp.ucla.edu/search/view/?f=yes&id=pds://PPI/maven.insitu.calibrated/data)).

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