

# Virtual height characteristics of ionospheric and ground scatter observed by mid-latitude SuperDARN HF radars

Evan G. Thomas<sup>1</sup> and Simon George Shepherd<sup>1</sup>

<sup>1</sup>Dartmouth College

November 26, 2022

## Abstract

Propagation of high-frequency (HF) radio signals is strongly dependent on the ionospheric electron density structure along a communications link. The ground-based, HF space weather radars of the Super Dual Auroral Radar Network (SuperDARN) utilize the ionospheric refraction of transmitted signals to monitor the global circulation of E- and F-region plasma irregularities. Previous studies have assessed the propagation characteristics of backscatter echoes from ionospheric irregularities in the auroral and polar regions of the Earth's ionosphere. By default, the geographic location of these echoes are found using empirical models which estimate the virtual backscattering height from the measured range along the radar signal path. However, the performance of these virtual height models has not yet been evaluated for mid-latitude SuperDARN radar observations or for ground scatter propagation modes. In this study, we derive a virtual height model suitable for mid-latitude SuperDARN observations using 5 years of data from the Christmas Valley East and West radars. This empirical model can be applied to both ionospheric and ground scatter observations and provides an improved estimate of the ground range to the backscatter location compared to existing high-latitude virtual height models. We also identify a region of overlapping half-hop F-region ionospheric scatter and one-hop E-region ground scatter where the measured radar parameters (e.g., velocity, spectral width, elevation angle) are insufficient to discriminate between the two scatter types. Further studies are required to determine whether these backscatter echoes of ambiguous origin are observed by other mid-latitude SuperDARN radars and their potential impact on scatter classification schemes.

1 **Virtual height characteristics of ionospheric and ground**  
2 **scatter observed by mid-latitude SuperDARN HF**  
3 **radars**

4 **E. G. Thomas<sup>1</sup>, S. G. Shepherd<sup>1</sup>**

5 <sup>1</sup>Thayer School of Engineering, Dartmouth College, Hanover, New Hampshire, USA

6 **Key Points:**

- 7 • We derive a new empirical virtual height model suitable for improved geolocation  
8 of mid-latitude SuperDARN HF radar observations  
9 • The new model provides the first characterization of ground scatter propagation  
10 modes  
11 • Characteristics of half-hop *F*-region ionospheric scatter and one-hop *E*-region ground  
12 scatter are examined

## Abstract

Propagation of high-frequency (HF) radio signals is strongly dependent on the ionospheric electron density structure along a communications link. The ground-based, HF space weather radars of the Super Dual Auroral Radar Network (SuperDARN) utilize the ionospheric refraction of transmitted signals to monitor the global circulation of  $E$ - and  $F$ -region plasma irregularities. Previous studies have assessed the propagation characteristics of backscatter echoes from ionospheric irregularities in the auroral and polar regions of the Earth's ionosphere. By default, the geographic location of these echoes are found using empirical models which estimate the virtual backscattering height from the measured range along the radar signal path. However, the performance of these virtual height models has not yet been evaluated for mid-latitude SuperDARN radar observations or for ground scatter propagation modes. In this study, we derive a virtual height model suitable for mid-latitude SuperDARN observations using 5 years of data from the Christmas Valley East and West radars. This empirical model can be applied to both ionospheric and ground scatter observations and provides an improved estimate of the ground range to the backscatter location compared to existing high-latitude virtual height models. We also identify a region of overlapping half-hop  $F$ -region ionospheric scatter and one-hop  $E$ -region ground scatter where the measured radar parameters (e.g., velocity, spectral width, elevation angle) are insufficient to discriminate between the two scatter types. Further studies are required to determine whether these backscatter echoes of ambiguous origin are observed by other mid-latitude SuperDARN radars and their potential impact on scatter classification schemes.

## 1 Introduction

The ground-based, high-frequency (HF) space weather radars of the Super Dual Auroral Radar Network (SuperDARN) utilize ionospheric refraction to routinely measure the line-of-sight (LOS) Doppler velocity of backscattered signals from  $E$ - and  $F$ -region plasma irregularities out to ranges of several thousand kilometers (Greenwald et al., 1995). The ability of SuperDARN radars to monitor ionospheric plasma convection therefore depends on two conditions: the presence of decameter-scale ionospheric irregularities, and suitable propagation conditions such that the transmitted HF radio waves can achieve perpendicularity to the magnetic field-aligned irregularities to satisfy the coherent Bragg scattering condition and return to the radar (Greenwald et al., 1985).

At auroral latitudes ( $60^\circ$ – $85^\circ$  magnetic latitude, or MLAT), where the first SuperDARN radars were located, both the occurrence of  $E$ - and  $F$ -region irregularities (e.g., Ruohoniemi & Greenwald, 1997; Ballatore et al., 2000; Koustov et al., 2004; Ghezlbash et al., 2014; Koustov et al., 2019; Marcucci et al., 2021) and HF propagation conditions (e.g., André et al., 1998; Yeoman et al., 2001; Gauld et al., 2002; Chisham et al., 2008; Yeoman et al., 2008; Ponomarenko et al., 2010) have been studied in great detail. More recently, new SuperDARN radars have been constructed at both mid-latitudes and in the polar cap for improved monitoring of global convection during periods of enhanced geomagnetic activity (Chisham et al., 2007; Nishitani et al., 2019). At midlatitudes, there is therefore a smaller body of work examining the irregularity and propagation characteristics (e.g., Nishitani & Ogawa, 2005; Ribeiro et al., 2012; de Larquier et al., 2013; Oinats et al., 2016; Shepherd et al., 2020; Wang et al., 2022).

It is important to note that most studies using SuperDARN data have focused on the occurrence and propagation modes of HF backscatter from ionospheric irregularities, or ionospheric scatter (IS). An important byproduct of the sky-wave propagation mode used by SuperDARN radars is the occurrence of ground scatter (GS) echoes from land and ocean surfaces along the signal path. While these GS returns are often treated as noise when producing global maps of ionospheric plasma motion (Chisham & Pinnock, 2002), they can be useful for monitoring different geophysical phenomena such as trav-

64 eling ionospheric disturbances (e.g., Bristow et al., 1996; He et al., 2004; Frissell et al.,  
 65 2016), HF absorption caused by solar flares (e.g., Hosokawa et al., 2000; Chakraborty  
 66 et al., 2018), or even land and ocean surface features (Shand et al., 1998; Ponomarenko  
 67 et al., 2010; Greenwood et al., 2011).

68 Many SuperDARN radars have a secondary interferometer antenna array to mea-  
 69 sure the angle of arrival, or elevation angle, of received signals. Because HF radio waves  
 70 undergo refraction as they traverse electron density gradients in the ionosphere, the ac-  
 71 tual height of the IS echo (or reflection height for GS) will always be lower in altitude  
 72 than for a signal traveling the same total distance along a straight-line path with the same  
 73 elevation angle. Breit and Tuve (1926) demonstrated how, for a flat Earth and planar  
 74 ionosphere, the propagation paths associated with these true and “virtual” heights have  
 75 the same ground range. SuperDARN radars can therefore use this virtual height infor-  
 76 mation to estimate the ground range to an IS or GS backscatter location as described  
 77 below.

78 The triangular virtual height geometry of Breit and Tuve (1926) has often been adapted  
 79 to describe  $\frac{1}{2}$ - and  $1\frac{1}{2}$ -hop IS propagation modes over a spherical Earth (e.g., Chisham  
 80 et al., 2008; Greenwald et al., 2017). One can extend this application of the law of cosines  
 81 to define a more general set of equations which describe both IS and GS propagation modes.  
 82 From the measured slant range  $r$  and elevation angle  $\alpha$  of the received radar signal, the  
 83 corresponding virtual height  $h_N$  for any  $N$ -hop propagation mode (assuming a spher-  
 84 ical Earth with radius  $R_E$ ) can be found using:

$$85 \quad h_N(r, \alpha) = \left[ R_E^2 + \left( \frac{r}{2N} \right)^2 + \left( \frac{r}{N} \right) R_E \sin(\alpha) \right]^{\frac{1}{2}} - R_E \quad (1)$$

86 where integer values of  $N$  (e.g., 1, 2, 3, etc.) correspond to GS propagation modes while  
 87 fractional values of  $N$  (e.g.,  $\frac{1}{2}$ ,  $1\frac{1}{2}$ ,  $2\frac{1}{2}$ , etc.) correspond to IS propagation modes. Note  
 88 that for the multi-hop case ( $N > 1$ ), the virtual height is assumed to be constant for  
 89 all ionospheric reflection and/or backscatter locations. The ground range  $G_N$  to each  
 90 IS or GS echo for any  $N$ -hop propagation mode can then be found using:

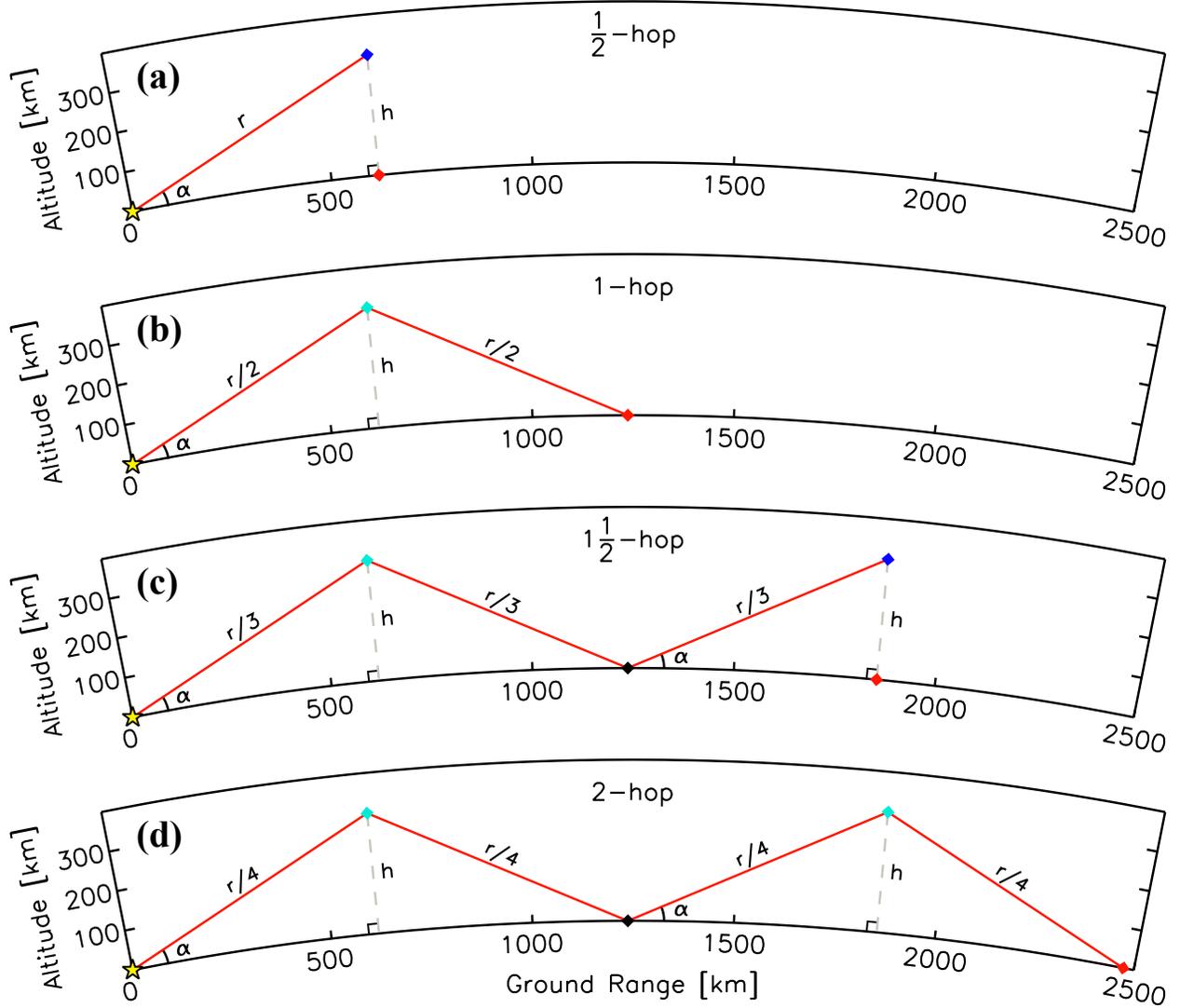
$$91 \quad G_N(r, \alpha, h_N) = 2NR_E \sin^{-1} \left[ \frac{\left( \frac{r}{2N} \right) \cos(\alpha)}{R_E + h_N} \right] \quad (2)$$

92 Alternatively, for cases where the elevation angle is not known (e.g., when using a vir-  
 93 tual height model) the ground range  $G_N$  can be found using:

$$94 \quad G_N(r, h_N) = 2NR_E \cos^{-1} \left[ \frac{R_E^2 + (R_E + h_N)^2 - \left( \frac{r}{2N} \right)^2}{2R_E(R_E + h_N)} \right] \quad (3)$$

95 Figure 1 illustrates sample HF propagation geometries found using Equations 1–3 for  
 96  $N \leq 2$ , where we have chosen input values of  $r$  and  $\alpha$  to obtain a representative  $F$ -region  
 97 virtual height of 300 km in each panel. In their consideration of the  $\frac{1}{2}$ -hop IS propaga-  
 98 tion mode, Greenwald et al. (2017) refer to this approach as the “two-parameter method”  
 99 due to the reliance on  $r$  and  $\alpha$  as input parameters. However, for the more general treat-  
 100 ment of either IS or GS, it is clear that a third input parameter specifying the number  
 101 of hops (i.e.,  $N$ ) is also required for an accurate ground range determination.

102 In practice, not all SuperDARN radars have a secondary interferometer array for  
 103 the measurement of elevation angles, or the time delays needed to accurately calculate  
 104 the elevation data ( $t_{\text{diff}}$ ) have not been properly calibrated (Chisham et al., 2021). For  
 105 this more common scenario, empirical models of virtual height are used for the geolo-  
 106 cation of line-of-sight (LOS) observations. The standard SuperDARN virtual height model  
 107 (hereafter referred to as the standard VHM) was derived from observations of IS mea-  
 108 sured by the original SuperDARN radar at Goose Bay (53.32° N, 60.46° W) overlook-  
 109 ing the auroral zone of the high-latitude ionosphere (Greenwald et al., 1985, 2017). This



**Figure 1.** Illustration of (a)  $\frac{1}{2}$ -hop, (b) 1-hop, (c)  $1\frac{1}{2}$ -hop, and (d) 2-hop  $F$ -region propagation geometries as a function of slant range  $r$ , elevation angle  $\alpha$ , and virtual height  $h$ , assuming a spherical Earth with radius  $R_E$ . Blue diamonds indicate ionospheric backscatter locations, while red diamonds indicate the ground range associated with each ionospheric or ground backscatter location. Cyan and black diamonds indicate ionospheric and ground reflection points, respectively, while the yellow star at zero ground range indicates the radar location.

110 model is divided into two segments, with the  $\frac{1}{2}$ -hop  $E$ -region and  $\frac{1}{2}$ -hop  $F$ -region prop-  
 111 agation modes connected by a simple linear transition:

$$112 \quad h(r) = \begin{cases} \frac{115r}{150} & 0 < r \leq 150 \text{ km} \\ \frac{115}{150} & 150 < r \leq 600 \text{ km} \\ \frac{r-600}{200}(h_i - 115) + 115 & 600 < r < 800 \text{ km} \\ h_i & r \geq 800 \text{ km} \end{cases} \quad (4)$$

113 where  $r$  is the measured slant range and  $h_i$  is the user-provided  $F$ -region virtual height  
 114 (typically either 300 or 400 km). Note that SuperDARN radars usually do not record  
 115 samples at ranges nearer than 180 km, although some non-standard operating modes de-

**Table 1.** The coefficients for Equation 5 for the Chisham et al. (2008) VHM.

Propagation Mode	A	B	C
$\frac{1}{2}$ -hop E-region	108.974	0.0191271	$6.68283 \times 10^{-5}$
$\frac{1}{2}$ -hop F-region	384.416	-0.178640	$1.81405 \times 10^{-4}$
$1\frac{1}{2}$ -hop F-region	1098.28	-0.354557	$9.39961 \times 10^{-5}$

116 signed for lower atmospheric measurements (such as mesospheric winds) may collect data  
 117 at these very near ranges (e.g., Yukimatu & Tsutsumi, 2002).

118 More recently, Chisham et al. (2008) derived a virtual height model (hereafter referred to as the Chisham VHM) using 5 years of IS measurements from the high-latitude  
 119 Saskatoon (SAS) SuperDARN radar (52.16° N, 106.53° W), fitting a low-order polynomial (or quadratic) of the form  
 120  
 121

$$122 \quad h(r) = A + Br + Cr^2 \quad (5)$$

123 to the  $\frac{1}{2}$ -hop *E*-region,  $\frac{1}{2}$ -hop *F*-region, and  $1\frac{1}{2}$ -hop *F*-region distributions; the coefficients for each model propagation mode are listed in Table 1. It should be noted that  
 124 the Chisham VHM was derived by first evaluating and then combining observations from  
 125 four equally-spaced azimuthal beam directions (beams 3, 6, 9, and 12) and all local times,  
 126 seasons, and radar operating frequencies. Greenwald et al. (2017) assessed the performance of the Chisham VHM using numerical ray-tracing simulations through the International Reference Ionosphere (IRI) model (Bilitza et al., 2011) at three different local times and a single frequency, finding the best agreement in terms of virtual height during nighttime conditions ( $\sim 21$  LT) when the SAS radar is most likely to observe IS.  
 127  
 128  
 129  
 130  
 131

132 Most recently, Liu et al. (2012) derived an alternative virtual height model using ray-tracing simulations and IS observations from the high-latitude Hankasalmi radar (62.32° N, 26.61° E) to fit a quadratic of the form  
 133  
 134

$$135 \quad h(r, \alpha) = Ar^2 + Br + C\alpha^2 + D\alpha + Er\alpha + F \quad (6)$$

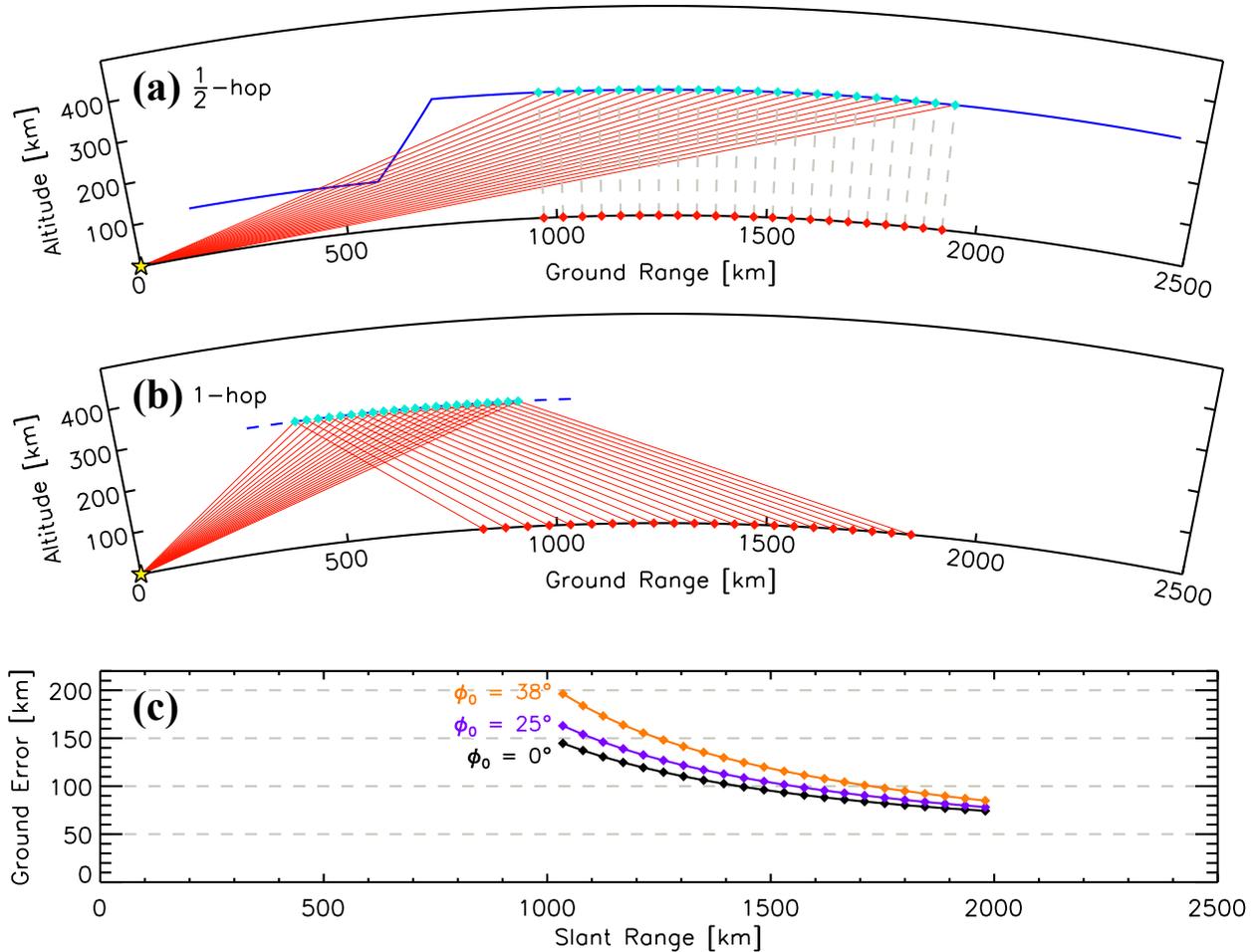
136 to the  $\frac{1}{2}$ -hop *F*-region distribution (630–1980 km slant range) only. Because this model is not relevant for  $\frac{1}{2}$ -hop *E*-region or  $1\frac{1}{2}$ -hop *F*-region scatter, it is of more limited use for geolocation purposes than either the standard or Chisham VHMs. Furthermore, this model relies upon both the measured range  $r$  and elevation angle  $\alpha$  as input, and is therefore not useful for radars without secondary interferometer arrays or calibrated  $t_{\text{diff}}$  values.  
 137  
 138  
 139  
 140  
 141

142 None of these virtual height models are appropriate for accurately mapping GS echoes to the Earth’s surface; neither is the current implementation of Equation 1 in the SuperDARN geolocation software, which can only support the  $\frac{1}{2}$ -hop propagation mode. As an example, the different propagation geometries of  $\frac{1}{2}$ -hop IS and 1-hop GS echoes at the same measured ranges ( $1000 < r < 2000$  km) are shown in Figures 2a and 2b for a spherical Earth. The cyan diamonds indicate the ionospheric reflection point of the IS or GS echoes predicted by the standard VHM (blue line), while the red diamonds indicate the associated ground range of the backscattered signals.  
 143  
 144  
 145  
 146  
 147  
 148  
 149

150 There is one additional factor which must be considered before comparing the ground ranges found for the IS and GS propagation modes in Figure 2. The azimuthal beam direction  $\phi$  relative to the radar boresight is given by  
 151  
 152

$$153 \quad \sin \phi = \frac{\sin \phi_0}{\cos \alpha} \quad (7)$$

154 where  $\phi_0$  is the direction at  $\alpha = 0^\circ$  (horizontal) set electronically by the radar hard-  
 155 ware. The final ground location of the backscattering target will, therefore, vary as a func-  
 156 tion of elevation angle not only in range but also in azimuth. The latter is often ignored  
 157 when determining ground range errors.



**Figure 2.** (a) Example of ground range mapping using the standard VHM (blue line) for measured slant ranges  $r$  between 1000–2000 km at 45 km resolution assuming a  $\frac{1}{2}$ -hop propagation mode (i.e., ionospheric backscatter) and a spherical Earth. (b) Ground range mapping for the same slant ranges  $r$  assuming a 1-hop propagation mode (i.e., ground backscatter) and an equivalent virtual reflection height at 300 km (blue dashed line). (c) Error in ground location obtained when applying the  $\frac{1}{2}$ -hop propagation assumption to 1-hop observations for three representative azimuthal beam directions; the black curve ( $\phi_0=0^\circ$ ) corresponds to the radar boresight direction, while the purple and orange curves correspond to the beams furthest from boresight for a nominal 16- or 24-beam SuperDARN radar, respectively.

158 The differences in the ground ranges obtained assuming  $\frac{1}{2}$ -hop IS (Figure 2a) versus  
 159 1-hop GS (Figure 2b) at the same measured ranges and virtual height along the radar  
 160 boresight direction ( $\phi_0 = 0^\circ$ ) can be seen in Figure 2c as indicated by the black curve.  
 161 Here we find that using a  $\frac{1}{2}$ -hop IS propagation model can result in positive ground range  
 162 offsets (i.e. away from the radar) of  $\sim 60$ – $150$  km. Along the beam directions furthest from  
 163 boresight for 16-beam ( $\phi_0 = 25^\circ$ ) or 24-beam ( $\phi_0 = 38^\circ$ ) SuperDARN radars, the dif-

ference in ground range increases to  $\sim 150$ – $200$  km for higher elevation angles as indicated by the purple and orange curves, respectively. Compared to the standard SuperDARN range resolution of 45 km, this ground range error is quite significant and must be accounted for when comparing radar measurements to land or sea features.

In this study, we examine how HF propagation characteristics of IS observed at mid-latitudes compare to the previously derived VHM for high-latitude propagation conditions. We also present the first statistical characterization of virtual height for different GS propagation modes. We use these results to derive a new VHM which can more accurately describe the propagation characteristics of both IS and GS measurements observed by the mid-latitude SuperDARN radars, thus allowing for improved geolocation of these LOS observations not only in ground range but also in azimuth.

## 2 Methodology and Data

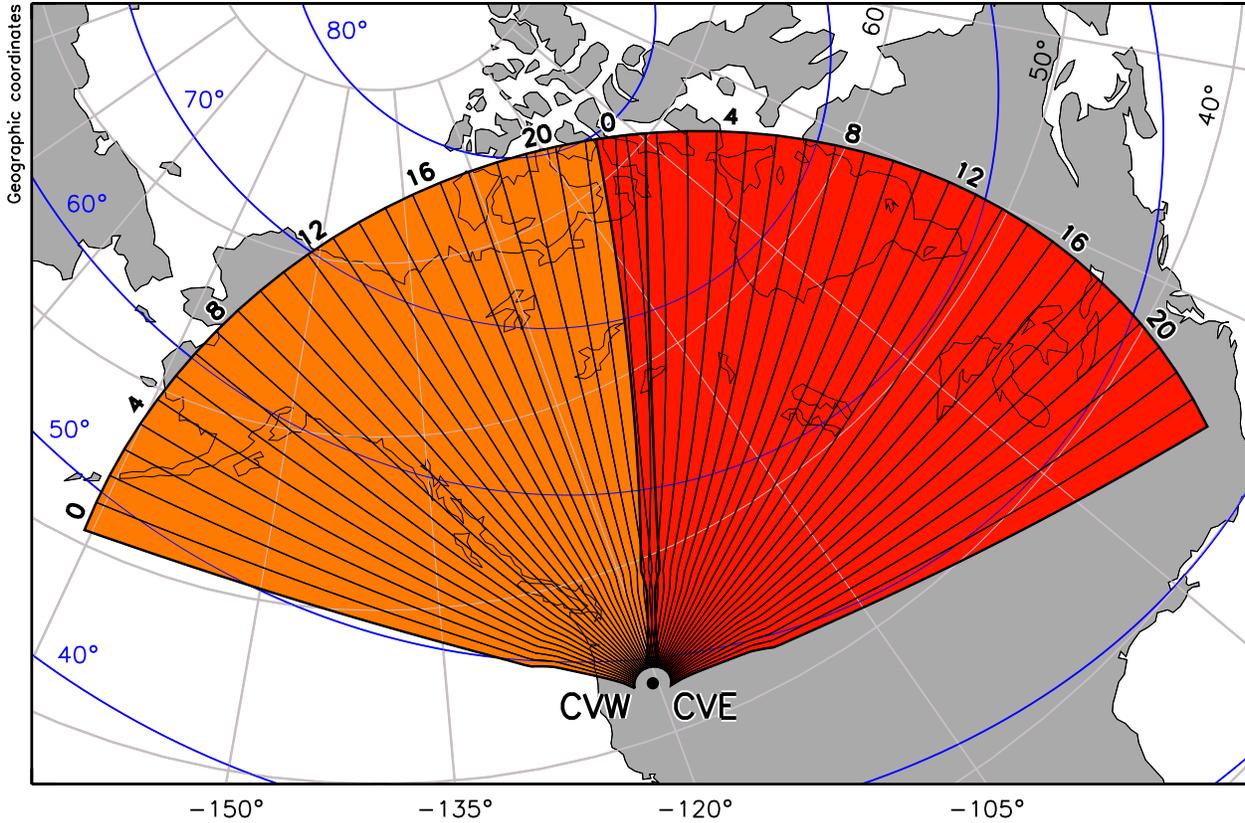
In this study we use 5 years of data (2014–2018) from the mid-latitude Christmas Valley East (CVE) and Christmas Valley West (CVW) pair of co-located SuperDARN radars ( $43.27^\circ$  N,  $120.36^\circ$  W). Figure 3 shows the nominal fields of view (FOVs) of each radar in geographic coordinates using the standard VHM, with contours of constant MLAT in Altitude-Adjusted Corrected Geomagnetic Coordinates (AACGM) (Shepherd, 2014) overlaid in blue. Both the CVE and CVW radars scan through up to 24 azimuthal beam directions across a sector of the mid- to high-latitude ionosphere spanning from  $50^\circ$ – $80^\circ$  MLAT. Each radar beam is separated by  $3.24^\circ$  in azimuth and sampled in 45 km range gates out to a maximum range of  $\sim 5000$  km at a cadence of 1–2 min. In practice however, only the 20 most-meridional beams of the CVE and CVW radars are typically sampled in order to synchronize scans to a 1 min boundary for standard radar operating modes.

LOS velocities, power, and spectral width are obtained from the raw data samples using the FITACF 2.5 library contained in version 4.3.1 of the Radar Software Toolkit (RST) (Thomas et al., 2020). Elevation angles are calculated using the generalized algorithm of Shepherd (2017) with fixed  $t_{\text{diff}}$  values of  $-398$  ns and  $-346$  ns for the CVE and CVW radars, respectively (Chisham et al., 2021). More than 450 million fitted LOS measurements with reliable elevation angles are available from each of the two radars during this 5-year interval, of which approximately 20% are identified as ionospheric scatter (IS) and 80% as ground scatter (GS) echoes using the default SuperDARN GS criterion:

$$|v| + \frac{w}{3} < 30 \text{ m/s} \quad (8)$$

where  $v$  is the fitted Doppler velocity and  $w$  is the spectral width. Note that echoes from meteor trails at near-ranges or slow-moving IS may be mis-identified as GS using the simple empirical criterion of Equation 8, particularly at mid-latitudes (Ribeiro et al., 2011). We will address the impact of potentially mis-identified scatter on our results in the following sections.

Figure 4 presents histograms of IS and GS echo occurrence for the CVE radar organized by six parameters: slant range, elevation angle, azimuthal beam number, radar operating frequency, Universal Time (UT), and month of year (all results for the CVW radar are shown in an equivalent set of figures in the supplementary material, and are generally similar to those shown for CVE). There is a large population of both IS and GS echoes found for slant ranges  $< 600$  km which is likely associated with  $\frac{1}{2}$ -hop backscatter from either meteor trails or  $E$ -region irregularities (Makarevich, 2010; Yakymenko et al., 2015). Secondary peaks in the IS and GS distributions in Figure 4a are located at slant ranges of  $\sim 1200$  and  $\sim 1500$  km, respectively. Both scatter types are observed across the full range of measurable elevation angles from  $0^\circ$  to  $50^\circ$  with a clear peak at  $18^\circ$  elevation (Figure 4b). The discontinuities seen above  $35^\circ$  elevation are related to the maximum observable elevation angle by the CVE radar, which for a given radar’s antenna configuration varies with both azimuthal beam direction and operating frequency

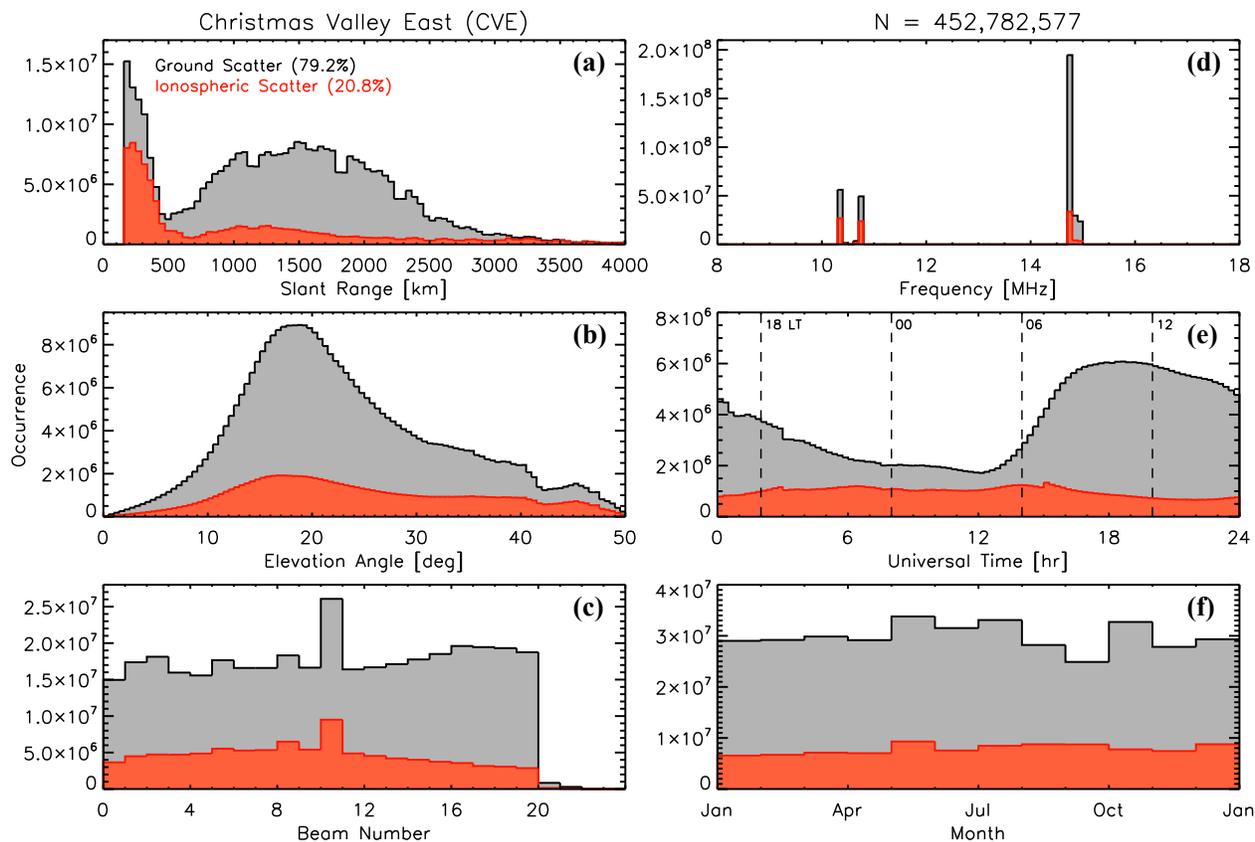


**Figure 3.** Nominal fields of view of the Christmas Valley East (CVE) and Christmas Valley West (CVW) radars in geographic coordinates, shaded red and orange respectively. Selected azimuthal beam numbers are labeled for each radar and contours of constant geomagnetic latitude at  $10^\circ$  intervals are overlaid in blue.

215 (Shepherd, 2017; Chisham, 2018). As previously described, the Christmas Valley radars  
 216 typically operate on only the 20 most-meridional beams; for the CVE radar this corre-  
 217 sponds to beam numbers 0–19 (Figure 4c). Beam number 10 of the CVE radar is the  
 218 designated “camping” beam used for special operating modes where finer temporal res-  
 219 olution (and thus increased occurrence rate) is obtained along a single azimuthal direc-  
 220 tion, at the expense of an increased scan duration across the full radar FOV.

221 Figure 4d shows that the CVE radar typically operates in one of two frequency bands:  
 222 10.3–10.8 MHz (during nighttime) and 14.7–15.0 MHz (during daytime). The lower fre-  
 223 quency band (10.3–10.8 MHz) has the appearance of being further divided into two bands  
 224 separated by only a few hundred kHz, which is due to an unresolved software issue in  
 225 the Christmas Valley radars’ clear frequency search algorithm. By combining the echo  
 226 occurrence from each of these lower (nighttime) frequency bands, approximately twice  
 227 as many IS echoes are observed than at the higher (daytime) frequency band. The op-  
 228 posite is true for the GS data, with significantly more echoes observed at the higher (day-  
 229 time) frequency band than for the lower (nighttime) band(s). Returning to the eleva-  
 230 tion histograms in Figure 4b, the maximum observable elevation angle at the lower fre-  
 231 quency band ranges from  $\sim 50^\circ$  (at boresight) to  $\sim 41^\circ$  (furthest from boresight), while  
 232 the elevation cutoff ranges from  $\sim 41^\circ$ – $35^\circ$  for the higher frequency band.

233 There is a clear diurnal variation in the GS occurrence seen in Figure 4e, with more  
 234 GS echoes observed during daytime hours (14–02 UT) than at nighttime. We find the



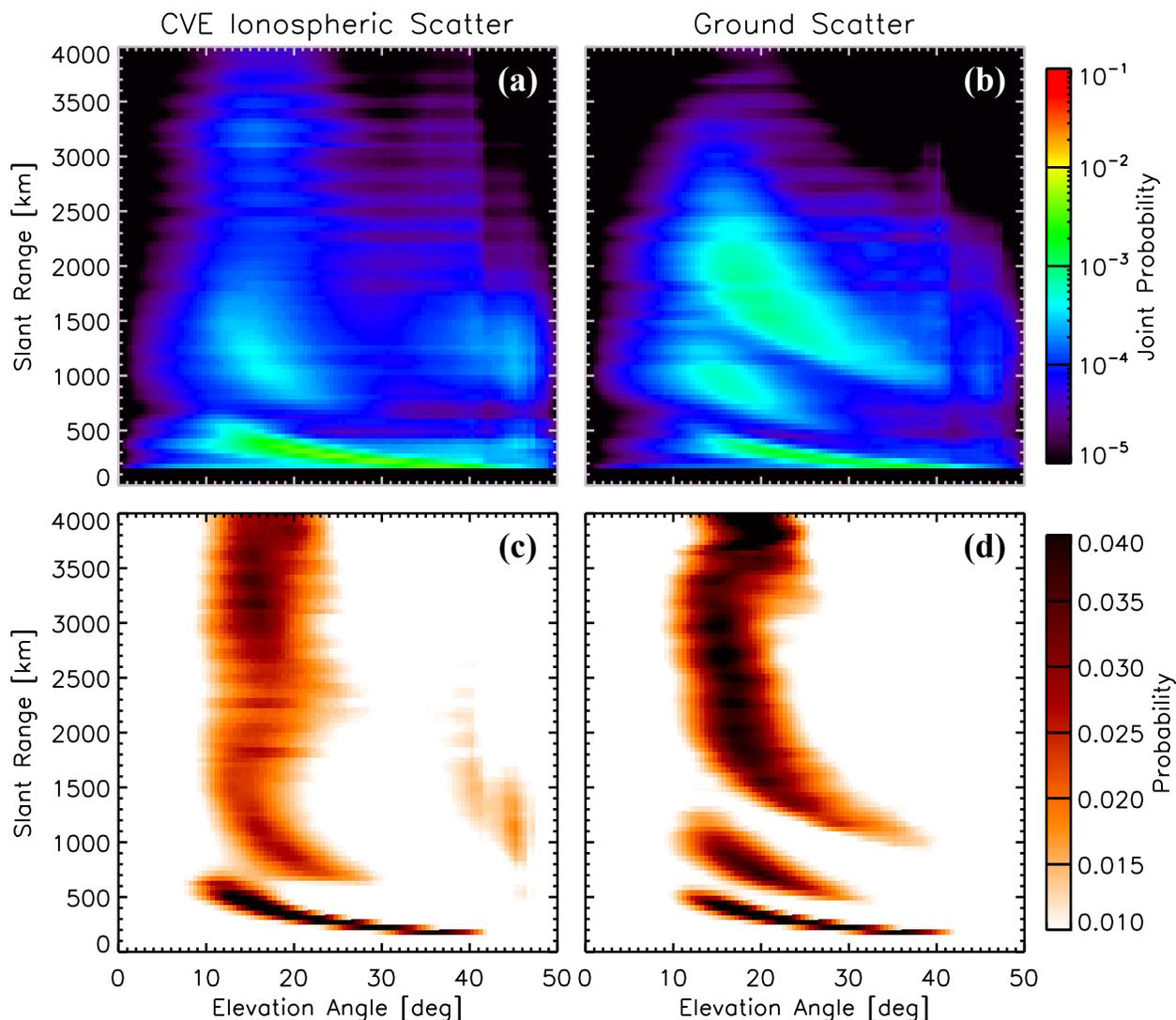
**Figure 4.** Statistical occurrence of CVE radar observations from 2014–2018 with (a) slant range, (b) elevation angle, (c) azimuthal beam number, (d) frequency, (e) Universal Time, and (f) month, sorted by ionospheric scatter (red) and ground scatter (black). Approximate local times at 6 hr intervals are indicated on panel (e) by vertical dashed lines, and the total number of measurements is given above panel (d).

235 opposite to be true for IS occurrence with slightly greater occurrence during nighttime  
 236 (02–14 UT) than daytime hours. No clear seasonal variations in the IS occurrence are  
 237 observed in Figure 4f and only a slight peak in the GS occurrence may be present dur-  
 238 ing summer months (May–July). These results are largely in agreement with previous  
 239 studies of SuperDARN backscatter occurrence rates at both mid- and high-latitudes (e.g.,  
 240 Hosokawa & Nishitani, 2010; Ribeiro et al., 2012; Ruohoniemi & Greenwald, 1997; Bal-  
 241 latore et al., 2000; Koustov et al., 2004; Ghezlbash et al., 2014; Koustov et al., 2019;  
 242 Marcucci et al., 2021).

### 243 3 Results

#### 244 3.1 Ionospheric and Ground Backscatter Distributions

245 We begin by considering the distribution of ionospheric and ground scatter observed  
 246 by the CVE radar in terms of the measured elevation angle versus slant range. The top  
 247 row of Figure 5 shows the joint probability distributions for both IS and GS, while the  
 248 bottom row shows the same distributions normalized by the maximum occurrence at each  
 249 range bin, after Chisham et al. (2008) and Chisham et al. (2021). The distributions in



**Figure 5.** (top) Joint probability distributions of elevation angle and slant range observed by the CVE radar for (a) ionospheric and (b) ground scatter, in  $0.5^\circ$  elevation and 45 km range bins. (bottom) The same probability distributions normalized by the maximum occurrence in each range bin, after Chisham et al. (2008) and Chisham et al. (2021); occurrence probabilities of less than 0.010 are not shown.

250 each panel of Figure 5 are divided into  $0.5^\circ$  elevation and 45 km range bins. Starting with  
 251 the IS distribution in Figures 5a and 5c, we observe three distinct populations:

- 252 1. At near ranges ( $\sim 180$ – $600$  km) across all elevation angles
- 253 2. Between  $\sim 600$ – $2000$  km range and  $10$ – $30^\circ$  elevation
- 254 3. At far ranges beyond  $\sim 2500$  km from  $10$ – $25^\circ$  elevation

255 There is one other population observed at higher elevation angles ( $35$ – $50^\circ$ ) between  
 256  $600$ – $2500$  km range (seen most clearly in Figure 5c). This region of range-elevation space  
 257 is sometimes associated with observations from the rear FOV (e.g., Milan et al., 1997;  
 258 André et al., 1998). However, we consider this scenario unlikely due to the Christmas

Valley radars' twin-terminated dipole (TTFD) wire antenna and corner reflector design which has an improved front-to-back ratio compared to the log-periodic antenna design of the original SuperDARN radars (Custovic et al., 2013). Instead, these anomalous measurements are almost certainly aliased from low elevation angles near zero degrees (McDonald et al., 2013) and are therefore excluded from further analysis.

Next we consider the GS distribution shown in Figures 5b and 5d, and again observe three distinct populations:

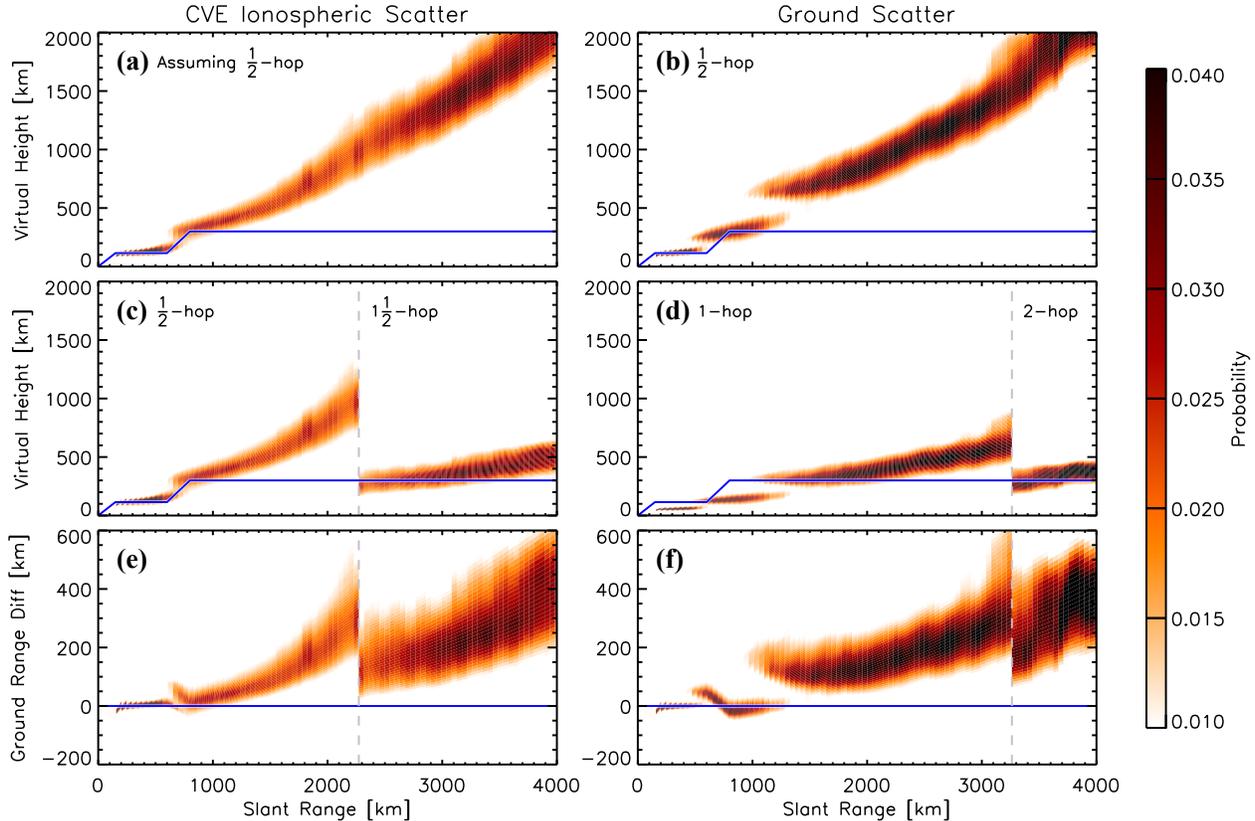
1. At near ranges ( $\sim 180$ – $500$  km) across all elevation angles
2. Between  $\sim 500$ – $1500$  km range and  $10$ – $30^\circ$  elevation
3. Between  $\sim 800$ – $3000$  km range and  $10$ – $40^\circ$  elevation

Unlike the IS distribution shown in the left column of Figure 5, there is considerable overlap between each of these three GS populations in the range dimension along the vertical axis. Elevation angle information is therefore critical for the identification of different GS propagation modes. From the normalized elevation-range distribution shown in Figure 5d, there appears to be a discontinuity near  $3500$  km slant range where the center of the elevation distribution shifts from  $\sim 14^\circ$  to  $\sim 19^\circ$ . This feature may indicate the presence of multi-hop GS echoes observed at extreme ranges. Note the two populations of what we assume to be aliased elevation angles near  $40^\circ$  and  $45^\circ$  are visible in Figure 5b but not the normalized representation in Figure 5d; these aliased data are also excluded from further analysis.

### 3.2 Comparison to Existing Virtual Height Models

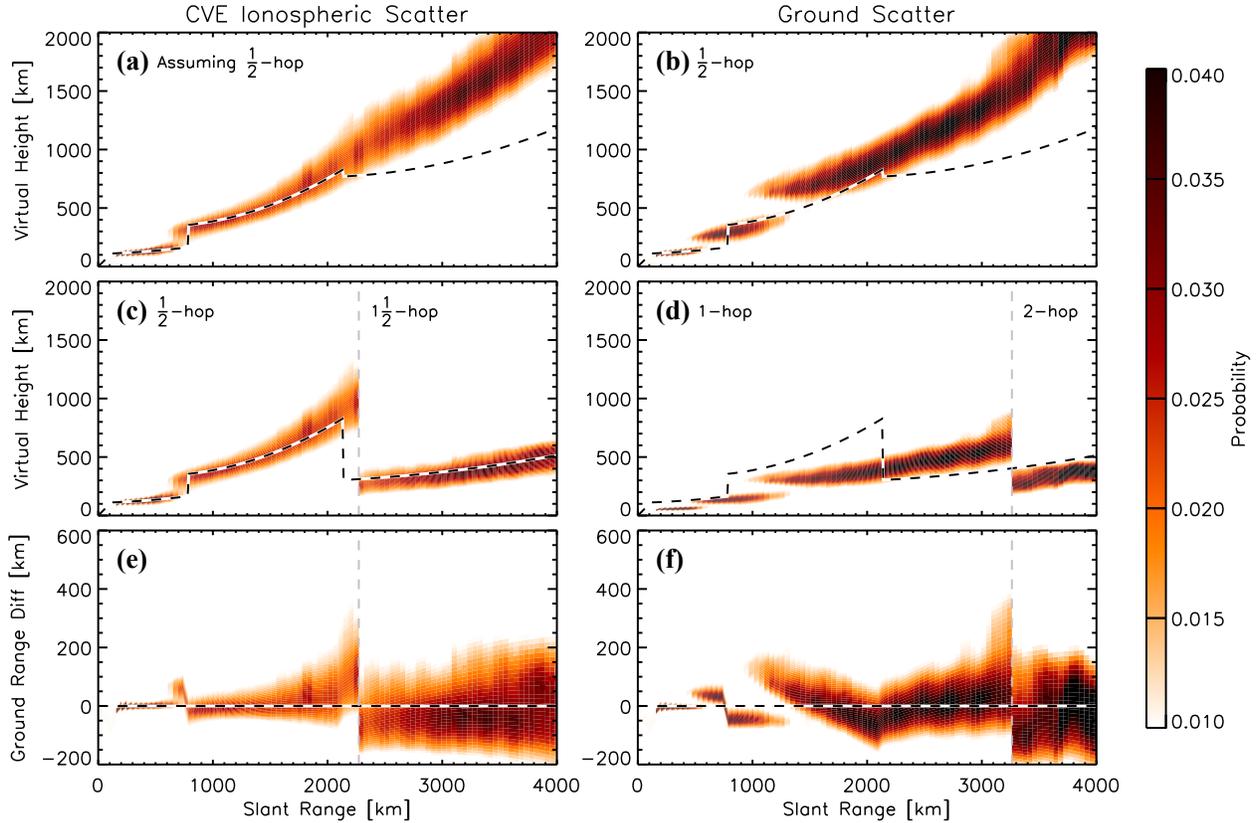
To aid in our physical interpretation of the results presented in Figure 5, we can use the measured elevation angle and slant range information to calculate the virtual height of the IS and GS probability distributions for any arbitrary number of hops with Equation 1. By doing so we may also evaluate the performance of existing SuperDARN VHMs when applied to the mid-latitude Christmas Valley radar observations. The resulting virtual height profiles, when assuming a  $\frac{1}{2}$ -hop propagation mode (e.g., Figure 1a), are shown versus range in Figures 6a and 6b for the normalized IS and GS probability distributions, respectively. The standard VHM with an  $F$ -region virtual height  $h_i$  of  $300$  km is overlaid on each panel in blue for reference (Equation 4). At near ranges in Figure 6a, there is excellent agreement between the observations and model prediction of  $115$  km virtual height, suggesting these echoes are in fact associated with  $\frac{1}{2}$ -hop backscatter from either meteor trails or  $E$ -region irregularities. Beyond  $\sim 600$  km range the IS distribution shifts from  $E$ - to  $F$ -region altitudes, and quickly curves upwards and away from the standard VHM  $h_i$  which remains constant at  $300$  km virtual height.

After careful examination of the elevation-range and range-height distributions, we estimate the transition between  $\frac{1}{2}$ - and  $1\frac{1}{2}$ -hop  $F$ -region IS propagation modes to be located near  $2270$  km slant range for the Christmas Valley radars. Figure 6c shows the same IS distribution as panel (a) but instead applying a  $1\frac{1}{2}$ -hop propagation assumption to the observations beyond  $2270$  km range (e.g., Figure 1c). This approach has the practical effect of lowering the virtual height by  $\sim 700$ – $1400$  km in our assumed  $1\frac{1}{2}$ -hop region. Figure 6e shows the difference in ground range obtained when using the  $\frac{1}{2}$ - and  $1\frac{1}{2}$ -hop assumptions applied in Figure 6c versus the predictions from the standard VHM. Here, a positive ground range difference indicates the standard VHM places the scatterer farther from the radar than Equations 1 and 2 would suggest. The ground range of the  $E$ -region IS is largely consistent whether the model or measured elevation angles are used, while the  $\frac{1}{2}$ -hop  $F$ -region IS can be located  $0$ – $400$  km closer to the radar than the standard VHM predicts. The  $1\frac{1}{2}$ -hop  $F$ -region IS is seen to always be in error with ground range differences of  $100$ – $600$  km found for all ranges, again with the model placing scatterer further from the radar than the measurements suggest.



**Figure 6.** Normalized probability distribution of elevation angle and slant range observed by the CVE radar for (a) ionospheric and (b) ground scatter from Figure 5 mapped to slant range versus virtual height assuming a  $\frac{1}{2}$ -hop propagation path, with the standard VHM overlaid in blue. (c) Same probability distribution for ionospheric scatter as panel (a) but assuming a  $1\frac{1}{2}$ -hop propagation mode for slant ranges beyond 2250 km (vertical dashed line). (d) Same probability distribution for ground scatter as panel (b) but assuming a 1-hop propagation mode for slant ranges less than 3240 km (vertical dashed line) and a 2-hop propagation mode for further ranges. The bottom row shows the difference in ground ranges for (e) ionospheric and (f) ground scatter from panels (c) and (d) compared to application of the standard VHM (which always assumes a  $\frac{1}{2}$ -hop propagation path); positive values indicate the true ground range is closer to the radar than suggested by the model.

309 The GS distribution shown in Figure 6b has similarly been mapped to virtual height  
 310 assuming a  $\frac{1}{2}$ -hop propagation mode using Equation 1. We see that again at near ranges  
 311 (180–500 km) the observations and standard VHM agree quite closely. This agreement  
 312 suggests the data have been mis-identified as GS and are instead associated with  $\frac{1}{2}$ -hop  
 313 backscatter from either meteor echoes or *E*-region irregularities, as the virtual height at  
 314 these ranges is unphysically low ( $\sim 50$  km) when assuming a 1-hop propagation mode (Fig-  
 315 ure 6d). The GS distribution at intermediate ranges (500–1300 km) in Figure 6b is lo-  
 316 cated between 250–400 km virtual height, straddling the standard VHM *F*-region  $h_i$  of  
 317 300 km. However, the virtual height of this population, when calculated assuming a 1-  
 318 hop propagation mode, also agrees with the standard VHM *E*-region virtual height of  
 319 115 km (Figure 6d), suggesting the echoes may be attributed to either mis-identified  $\frac{1}{2}$ -  
 320 hop IS from *F*-region irregularities or 1-hop GS reflected at *E*-region altitudes.



**Figure 7.** Normalized probability distributions of slant range and virtual height observed by the CVE radar for (left) ionospheric and (right) ground scatter in the same format as Figure 6, but instead compared against the Chisham VHM overlaid as a black and white dashed line. The “pseudo” virtual height predicted by the Chisham VHM for  $1\frac{1}{2}$ -hop scatter (beyond 2137 km slant range) is overlaid on panels (a–b), while the “true” virtual height predicted by the model for  $1\frac{1}{2}$ -hop scatter is overlaid on panels (c–d); see text for further details. The bottom row shows the difference in ground ranges for (e) ionospheric and (f) ground scatter from panels (c) and (d) compared to application of the Chisham VHM (which assumes either a  $\frac{1}{2}$ -hop or  $1\frac{1}{2}$ -hop propagation path); positive values indicate the true ground range is closer to the radar than suggested by the model.

321 At farther ranges (beyond  $\sim 1000$  km), the GS distribution in Figure 6b is offset  
 322 from the standard VHM prediction by at least 300 km and curves upwards to virtual heights  
 323 exceeding 2000 km. After applying a 1-hop propagation assumption to these data (e.g.,  
 324 Figure 1b) the distribution is brought downward to significantly lower virtual heights  
 325 spanning from 300–800 km. Again we have estimated a likely transition between 1- and  
 326 2-hop *F*-region GS to be located near 3260 km slant range, beyond which the observa-  
 327 tions in Figure 6d have been mapped to a virtual height assuming a 2-hop propagation  
 328 mode using Equation 1. Similar to Figure 6e, Figure 6f shows the difference in ground  
 329 range obtained when using the 1- and 2-hop assumptions applied in Figure 6d versus the  
 330 predictions from the standard VHM.

331 Figure 7 is in the same format as Figure 6, but instead we compare the CVE re-  
 332 sults to the Chisham VHM (black and white dashed line) rather than to the standard  
 333 VHM. In the top row, the IS and GS distributions are again mapped to range-virtual

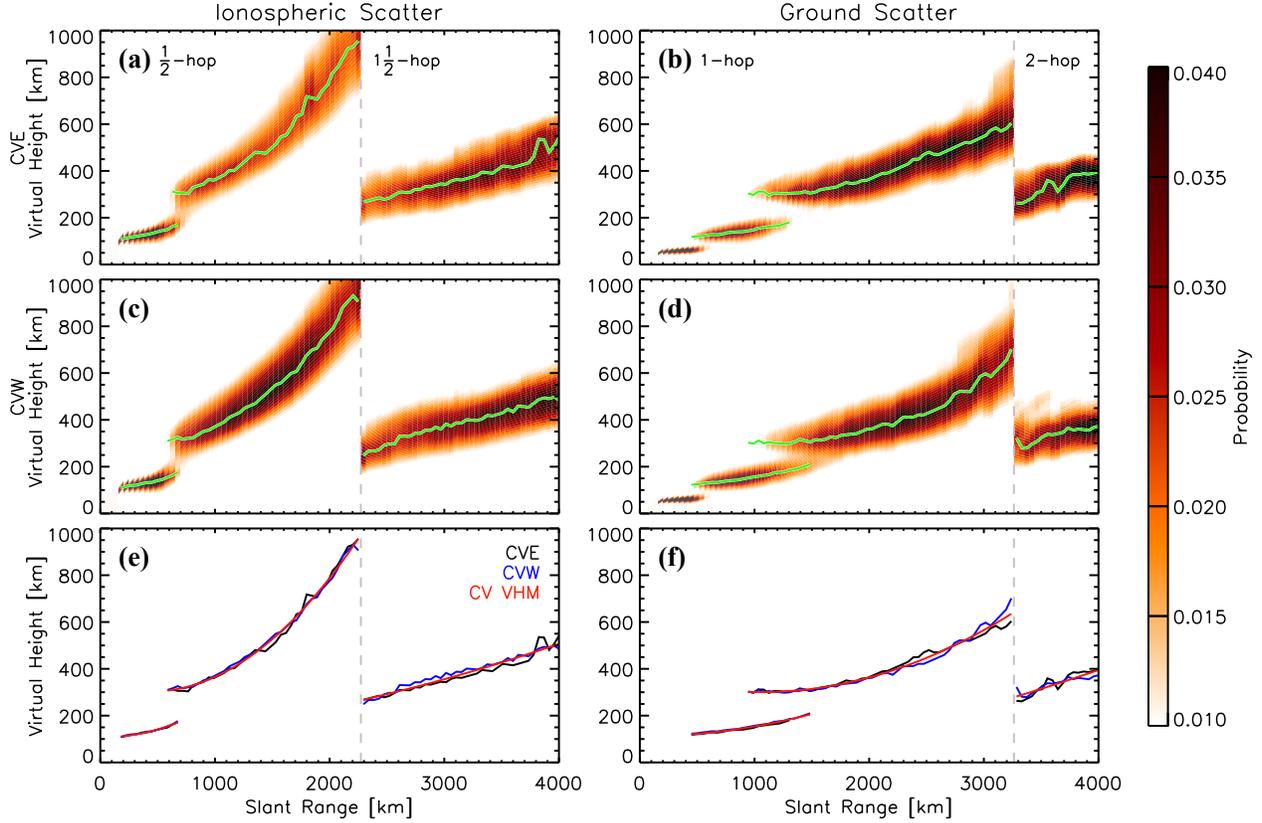
334 height space assuming a  $\frac{1}{2}$ -hop propagation mode. In these panels, the “pseudo” virtual  
 335 height of the Chisham VHM is shown for the  $1\frac{1}{2}$ -hop region beyond 2137 km range. Chisham  
 336 et al. (2008) derived this pseudo virtual height for compatibility with the existing Su-  
 337 perDARN range finding software, which can support only  $\frac{1}{2}$ -hop propagation modes. The  
 338 pseudo virtual height is therefore a  $\frac{1}{2}$ -hop virtual height with the same ground range as  
 339 the true  $1\frac{1}{2}$ -hop virtual height found by Chisham et al. (2008) in their statistical anal-  
 340 ysis of IS echoes. In the middle row of Figure 7 where the observed IS and GS distribu-  
 341 tions have been mapped using more appropriate assumptions, the true  $1\frac{1}{2}$ -hop virtual  
 342 height predicted by the Chisham VHM is shown. Here we see that the IS distribution  
 343 aligns quite well with the Chisham VHM in all three scatter regions in terms of virtual  
 344 height, except perhaps the  $1\frac{1}{2}$ -hop region where the observations are located slightly be-  
 345 low the model prediction. Where the IS observations and Chisham VHM differ most no-  
 346 ticeably are the ranges at which the peak of the virtual height distribution transitions  
 347 from one propagation mode to the next, i.e.  $\frac{1}{2}$ -hop  $E$ -region to  $\frac{1}{2}$ -hop  $F$ -region, and  $\frac{1}{2}$ -  
 348 to  $1\frac{1}{2}$ -hop  $F$ -region. This difference is further illustrated in Figure 7e, where the ground  
 349 range differences are largely centered about zero except at these propagation boundaries.

350 The GS distribution shown on the right side of Figure 7, on the other hand, does  
 351 not align particularly well with the Chisham VHM. The  $\frac{1}{2}$ -hop  $F$ -region portion of the  
 352 Chisham VHM is at significantly higher virtual heights than the 1-hop  $F$ -region height,  
 353 and the 2-hop GS is at a lower virtual height than the  $1\frac{1}{2}$ -hop model prediction. How-  
 354 ever, an incorrect virtual height and propagation mode can sometimes produce a real-  
 355 istic ground range estimate. For example, at 3500 km slant range the Chisham VHM pre-  
 356 dicts a  $1\frac{1}{2}$ -hop propagation mode with a virtual height of 435 km and ground range of  
 357 3146 km, while for a 2-hop propagation mode at that range the CVE GS measurements  
 358 suggest a virtual height of 315 km and ground range of 3189 km (only a  $\sim 1\%$  difference  
 359 of 43 km, or less than one standard range gate). At these ranges the elevation angle dif-  
 360 ference is less than  $1^\circ$  so the azimuthal errors will be within the beamwidth ( $3.24^\circ$ ), how-  
 361 ever at nearer ranges when the ground ranges agree but the virtual height is incorrect,  
 362 the errors in elevation angle, and therefore azimuth, will become more significant.

### 363 3.3 Christmas Valley Virtual Height Model

364 Figure 8 shows results for both the CVE and CVW radars in range-virtual height  
 365 space, with the IS and GS distributions again in the left and right columns, respectively.  
 366 The green lines overlaid on panels (a–d) indicate the virtual height of peak occurrence  
 367 at each slant range bin, i.e. the virtual height at which observations are most likely for  
 368 each range bin. Note there is significant overlap in range between the 1-hop  $E$ - and 1-  
 369 hop  $F$ -region GS distributions observed by both radars (Figures 8b and 8d), even moreso  
 370 for CVW. In this case we have attempted to find the peak occurrence for each virtual  
 371 height population regardless of whether there is some overlap in range.

372 The virtual height of maximum occurrence at each range bin is shown for both radars  
 373 in Figures 8e and 8f with results for CVE in black and CVW in blue. In the same man-  
 374 ner as Chisham et al. (2008), we have performed a least-squares fit to the average of the  
 375 black and blue curves using a quadratic function (Equation 5), which is overlaid on Fig-  
 376 ures 8e and 8f in red. The red curves therefore correspond to the Christmas Valley vir-  
 377 tual height model derived from both CVE and CVW observations, referred to hereafter  
 378 as the CV VHM, with 3 independent sets of coefficients for both the IS and GS prop-  
 379 agation modes. These model coefficients for Equation 5 are provided in Table 2. Note  
 380 that although we have shown an abrupt transition from  $\frac{1}{2}$ - to  $1\frac{1}{2}$ -hop  $F$ -region propa-  
 381 gation modes indicated by the vertical dashed line, this should in practice be a flexible  
 382 boundary where the transition may vary based on local time, season, or solar cycle con-  
 383 ditions (and similarly for the  $\frac{1}{2}$ -hop  $E$ - to  $F$ -region transition).



**Figure 8.** Normalized probability distributions of slant range and virtual height observed by the CVE (top row) and CVW (middle row) radars for ionospheric scatter (left column) and ground scatter (right column), with the virtual height of maximum occurrence at each range bin overlaid on panels (a–d) in green. The bottom row shows the virtual height of maximum occurrence from the CVE and CVW radars in black and blue, respectively, with a series of quadratic fits representing the new VHM overlaid in red. Note the different vertical axis scale used in these panels compared to Figures 6 and 7.

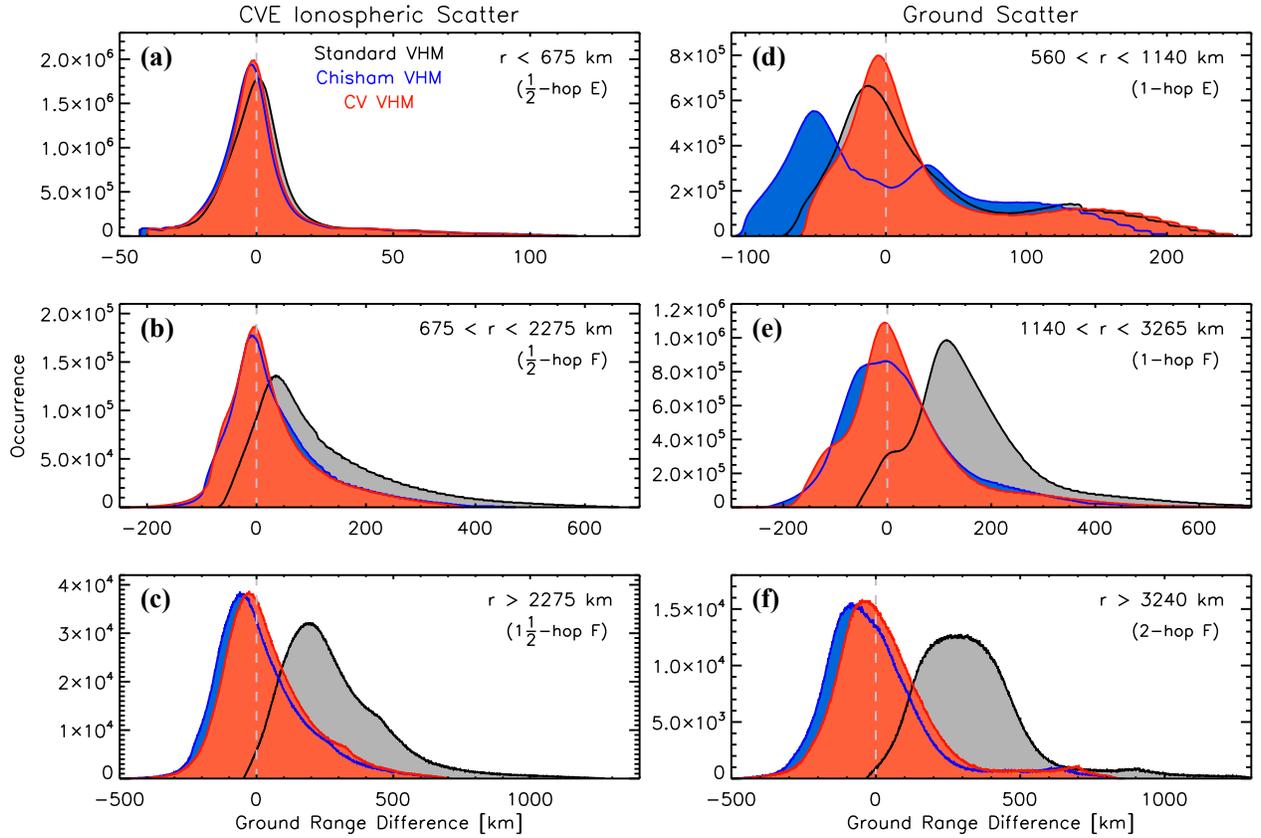
384 To assess the performance of the CV VHM relative to the standard and Chisham  
 385 VHMs, we calculate the ground range difference of all IS and GS observations when us-  
 386 ing each of the three VH models. These results are summarized in Figure 9, where each  
 387 panel corresponds to a separate IS or GS propagation mode; results for the standard VHM  
 388 are shown in gray, the Chisham VHM in blue, and the CV VHM in red. We can see that  
 389 for all six propagation modes, the ground range difference when using the CV VHM is  
 390 centered about zero, indicating there are no systematic biases (as is clearly seen for the  
 391 standard VHM with positive range offsets). The Chisham VHM performs surprisingly  
 392 well even for the 1- and 2-hop *F*-region GS propagation modes, although it is worth not-  
 393 ing that because the Chisham VHM virtual heights are incorrect, the inferred elevation  
 394 angle and thus coning angle correction to the beam azimuth will be incorrect (Equation 7).

#### 395 4 Discussion

396 The IS component of the CV model has large deviations from the SuperDARN com-  
 397 munity’s standard VHM which uses fixed *E*- and *F*-region virtual heights of 115 and 300 km  
 398 respectively (Equation 4). The IS component of the CV model agrees much more closely

**Table 2.** The coefficients for Equation 5 for the Christmas Valley ionospheric and ground scatter VHM.

Model	Propagation Mode	A	B	C
Ionospheric Scatter	$\frac{1}{2}$ -hop E-region	108.873	-0.01444	$1.57806 \times 10^{-4}$
	$\frac{1}{2}$ -hop F-region	341.005	-0.17484	$1.99144 \times 10^{-4}$
	$1\frac{1}{2}$ -hop F-region	92.9665	0.03967	$1.59501 \times 10^{-5}$
Ground Scatter	1-hop E-region	111.393	$-1.65773 \times 10^{-4}$	$4.26675 \times 10^{-5}$
	1-hop F-region	378.022	-0.14738	$6.99712 \times 10^{-5}$
	2-hop F-region	-76.2406	0.06854	$1.23078 \times 10^{-5}$


**Figure 9.** Differences in ground range when using the slant range and virtual height observed by the CVE radar for (left) ionospheric and (right) ground scatter compared to the standard (gray), Chisham (blue) and CV (red) virtual height models. Results are organized by (a)  $\frac{1}{2}$ -hop *E*-region, (b)  $\frac{1}{2}$ -hop *F*-region, (c)  $1\frac{1}{2}$ -hop *F*-region, (d) 1-hop *E*-region, (e) 1-hop *F*-region, and (f) 2-hop *F*-region scatter distributions.

399 with the Chisham VHM, suggesting that, in a statistical sense, the  $\frac{1}{2}$ - and  $1\frac{1}{2}$ -hop prop-  
 400 agation modes are consistent between mid- and high-latitude radar observations (or at  
 401 least for the CVE/CVW and SAS radars, from which the respective models were derived).  
 402 Where the CV and Chisham VHM differ the most are the ranges at which the models  
 403 predict a transition from  $\frac{1}{2}$ -hop *E*- to  $\frac{1}{2}$ -hop *F*-region modes, and from  $\frac{1}{2}$ -hop to  $1\frac{1}{2}$ -hop  
 404 *F*-region modes. While a fixed transition range is necessary for implementation in au-  
 405 tomated geolocation software, ideally these should be flexible boundaries which the user

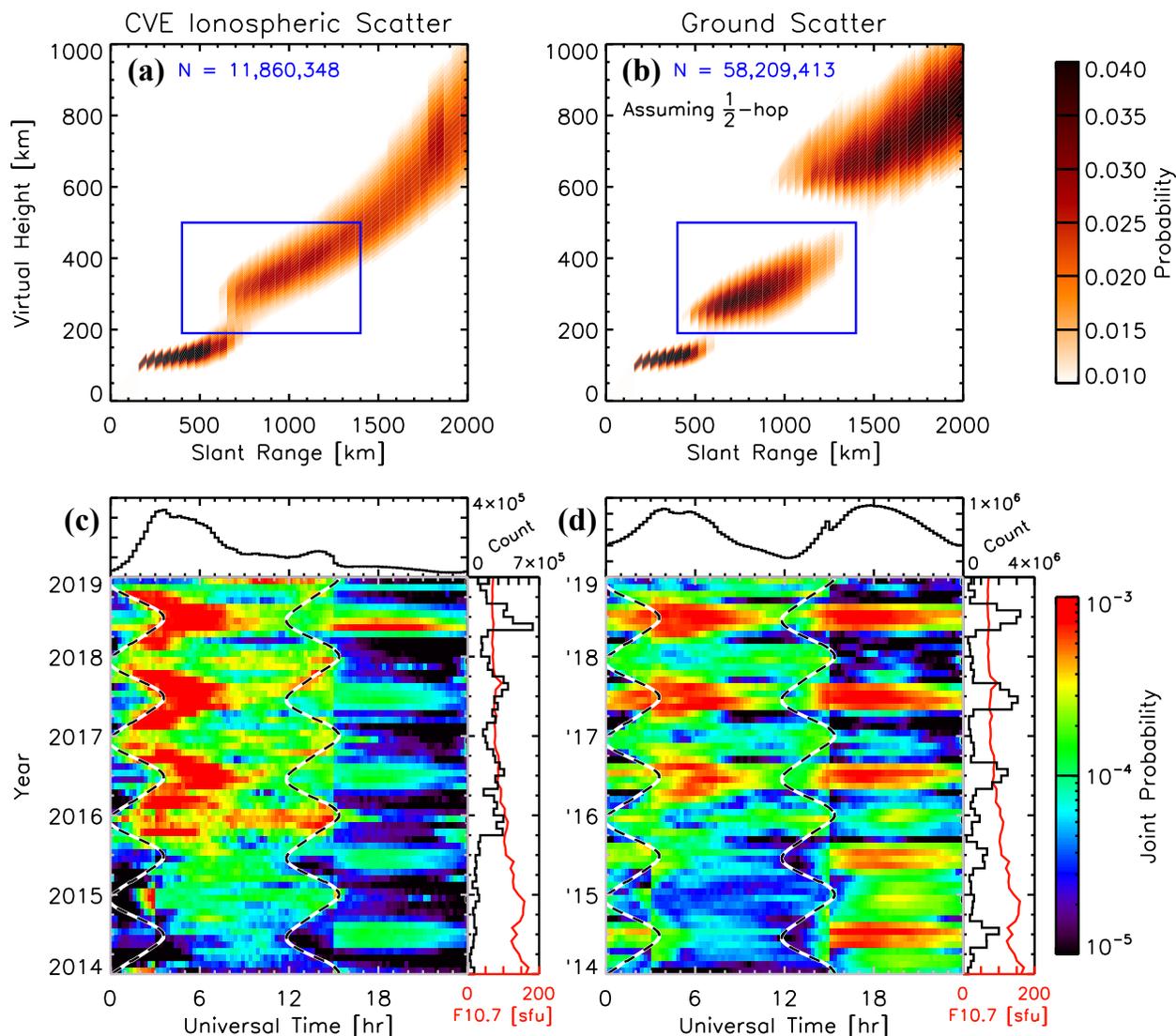
406 can shift to nearer or further ranges depending on instantaneous conditions. These tran-  
 407 sition ranges (as well as the virtual height) are also likely to vary with local time, sea-  
 408 son, operating frequency, etc., and will be examined in future work. We also emphasize  
 409 that despite the Chisham VHM providing reasonably accurate ground ranges when ap-  
 410 plied to 1- or 2-hop  $F$ -region GS propagation modes (Figure 9e–f), the resulting eleva-  
 411 tion angles and thus coning angle correction to the radar beam azimuth will be incor-  
 412 rect. This additional error in ground location due to an incorrect elevation angle pre-  
 413 dicted by any VHM is an often overlooked aspect of SuperDARN HF backscatter geolo-  
 414 cation.

415 The greatest source of uncertainty when applying the CV VHM described in this  
 416 study will likely arise from the initial determination of whether a backscatter echo be-  
 417 longs to either an IS or GS propagation mode. This issue was not relevant for previous  
 418 VHMs, as they did not consider GS propagation modes. Ribeiro et al. (2011) and oth-  
 419 ers have demonstrated how the default SuperDARN GS criteria can falsely identify slow-  
 420 moving IS as GS, particularly at subauroral latitudes in the nightside ionosphere. Ex-  
 421 isting techniques for identifying SuperDARN backscatter propagation modes (e.g., Bur-  
 422 rell et al., 2015; Bland et al., 2014) rely upon calibrated elevation angle measurements,  
 423 which are currently not available at many SuperDARN radar sites.

424 In our study, we have identified a backscatter region of ambiguous origin located  
 425 between  $\sim 500$ – $1400$  km slant range where the data could belong to either a  $\frac{1}{2}$ -hop  $F$ -  
 426 region or 1-hop  $E$ -region propagation mode. Figure 10 shows the CVE echo distribution  
 427 in range-virtual height space centered about this region for IS-flagged data in panel (a)  
 428 and GS-flagged data in panel (b); both distributions are mapped to virtual height as-  
 429 suming a  $\frac{1}{2}$ -hop propagation mode for easier comparison. The precise region of interest  
 430 is indicated by the blue rectangle: approximately 12 million IS-flagged echoes and 58 mil-  
 431 lion GS-flagged echoes are located within this space (a similar proportion to the over-  
 432 all IS-GS echo ratio, e.g. Figure 4a). In the lower panels of Figure 10 we have plotted  
 433 data from the echo region of “ambiguous” origin in terms of the joint UT versus month/year  
 434 probability distribution. The black and white dashed lines overlaid on each panel indi-  
 435 cate the local sunrise and sunset times at the approximate midpoint between the CVE  
 436 radar boresight direction and backscattering volume. The outer histograms along the  
 437 top and right edges of each panel show the 1-D occurrence distributions with UT and  
 438 month/year, respectively. Overlaid on the right-hand histogram in red is the monthly  
 439 average F10.7 solar radio flux (Tapping, 2013). We again refer the reader to the online  
 440 supplementary material for an equivalent figure of the region of “ambiguous” scatter ori-  
 441 gin observed by the CVW radar.

442 Starting with the IS-flagged echo occurrence in Figure 10c, there is a clear depen-  
 443 dence on the solar terminator with increased occurrence in the hours just after local sun-  
 444 set ( $\sim 0$ – $8$  UT depending on season). This dependence is superimposed on a larger trend  
 445 of increasing echo occurrence with decreasing F10.7 (i.e., the decline of solar cycle 24).  
 446 Ribeiro et al. (2012) found a similar relationship between the solar terminator and night-  
 447 time IS echo occurrence as observed by the mid-latitude Blackstone radar ( $37.10^\circ$  N,  $77.95^\circ$  W),  
 448 although they could not identify any seasonal or solar cycle trends because their study  
 449 was limited to only 2 years of data. There is an abrupt increase in the total occurrence  
 450 near the end of 2015, which we are unable to attribute to any operational changes at the  
 451 CVE radar (CVW did not observe a similar change in scatter occurrence). Throughout  
 452 all years in this interval, there is also a smaller population of daytime echoes observed  
 453 only during summer months. We therefore suggest that the IS-flagged distribution con-  
 454 tains primarily backscatter from nighttime ionospheric sources with some contamination  
 455 from daytime 1-hop  $E$ -region propagation modes.

456 Turning next to the GS-flagged echoes in Figure 10d, there is a much more even  
 457 distribution of daytime and nighttime echoes seen in the UT histogram along the top  
 458 of the figure. Again there are signs of increased echo occurrence following the local sun-



**Figure 10.** Normalized probability distributions of slant range and virtual height observed by the CVE radar for (a) ionospheric scatter and (b) ground scatter, centered on the region of possible “mixed” scatter indicated by the blue box. Both the ionospheric and ground scatter-flagged data in panels (a–b) have been mapped to virtual height assuming a  $\frac{1}{2}$ -hop propagation mode for easier comparison, and the number of observations falling within the “mixed” scatter region are given at the top of each panel in blue. (c) Joint probability distribution of Universal Time (UT) and month/year for the ionospheric-flagged scatter within the blue box in panel (a); along the top of the panel is a 1-D histogram of the same data as a function of UT only, while along the right side of the panel is a 1-D histogram as a function of month and year with the monthly average F10.7 solar radio flux overlaid in red. (d) Same as panel (c) but for the ground scatter-flagged data within the blue box in panel (b). Approximate sunrise and sunset times at the midpoint between the CVE radar and “mixed” scatter backscattering volume are overlaid on panels (c–d) as black and white dashed lines.

459 set terminator. The daytime summer population is present during all years, however, and  
 460 appears to be a dominant contributor to the 1-D histogram along the right edge of the

461 plot. This feature follows from the knowledge that mid-latitude  $E$ -region densities are  
 462 almost solely controlled by solar zenith angle, i.e. maximum during summer and min-  
 463 imum during winter (Chu et al., 2009), and thus more likely to support a 1-hop  $E$ -region  
 464 propagation mode. An increase in 1-hop  $E$ -region GS echo occurrence at mid-latitudes  
 465 during summer months was also predicted by the ray-tracing simulations for the Black-  
 466 stone radar by de Larquier et al. (2011) in their Figures 6 and 7. Unlike the IS-flagged  
 467 results in our Figure 10c, in the GS-flagged results there is another weaker population  
 468 of winter daytime echoes which disappears with decreasing F10.7 / declining solar cy-  
 469 cle phase. This feature also follows from the knowledge that  $N_mE$  has a secondary de-  
 470 pendence on F10.7 (Titheridge, 2000).

471 To summarize our observations of this measurement region, there is clear evidence  
 472 for contamination of GS echoes in the IS-flagged data and vice versa. However, there is  
 473 nearly an equal number of likely IS echoes in the GS-flagged data as there are true GS  
 474 echoes. Therefore, even the inclusion of measured elevation angles in an empirical GS  
 475 criteria will not be sufficient to accurately classify the measurements within this slant  
 476 range interval as either IS or GS for the Christmas Valley radars. Because the ionospheric  
 477  $E$ -region electron densities are almost solely controlled by the solar zenith angle at mid-  
 478 latitudes (with a secondary dependence on solar activity), consideration of local time,  
 479 season, and solar cycle factors may help with discrimination of IS versus GS sources. Fu-  
 480 ture work will determine whether echoes are observed by the other mid-latitude Super-  
 481 DARN radars which exhibit similar occurrence characteristics.

## 482 5 Conclusions

483 In this study we have examined 5 years of data from the mid-latitude Christmas  
 484 Valley East and West SuperDARN radars to derive an empirical virtual height model  
 485 with two sets of coefficients: one suitable for ionospheric scatter (IS) and, for the first  
 486 time, another exclusively for ground scatter (GS) echoes. Both components of the CV  
 487 model represent a significant advancement over the standard SuperDARN virtual height  
 488 model, which treats all backscatter echoes as belonging to either a  $\frac{1}{2}$ -hop  $E$ - or  $F$ -region  
 489 IS propagation mode. The IS component of our CV model performs similarly to the more  
 490 recent model of Chisham et al. (2008), suggesting that in a climatological sense, the HF  
 491 propagation modes for backscatter from ionospheric irregularities are similar at both au-  
 492 roral and mid-latitudes. We have also identified a measurement region ( $500 < r < 1400$  km)  
 493 where the LOS velocity, spectral width, slant range, and elevation angle are insufficient  
 494 for separation of IS and GS echoes. Local time, season, and solar cycle factors should  
 495 therefore be considered when analyzing scatter from this region. The CV IS and GS vir-  
 496 tual height models have been incorporated into the freely available SuperDARN RST  
 497 for use with the standard analysis routines, which will improve the geolocation accuracy  
 498 for scatter observed by all mid-latitude SuperDARN radars.

## 499 Open Research

500 The raw SuperDARN data used in this study are available from the British Antarc-  
 501 tic Survey (BAS) SuperDARN data mirror (<https://www.bas.ac.uk/project/superdarn>).  
 502 The Radar Software Toolkit (RST) to read and process the SuperDARN data can be down-  
 503 loaded from Zenodo (Thomas et al., 2020). The monthly average solar radio flux data  
 504 were obtained from Space Weather Canada at [https://spaceweather.gc.ca/solarflux/sx-](https://spaceweather.gc.ca/solarflux/sx-5-en.php)  
 505 [5-en.php](https://spaceweather.gc.ca/solarflux/sx-5-en.php).

## 506 Acknowledgments

507 This research was funded by the National Science Foundation under grants OPP-1836426  
 508 and AGS-1341925. The authors acknowledge the use of SuperDARN data. SuperDARN

509 is a collection of radars funded by the national scientific funding agencies of Australia,  
 510 Canada, China, France, Italy, Japan, Norway, South Africa, UK, and United States. The  
 511 authors also acknowledge the use of solar radio flux data from Space Weather Canada.

## 512 References

- 513 André, D., Sofko, G. J., Baker, K., & MacDougall, J. (1998). SuperDARN interfer-  
 514 ometry: Meteor echoes and electron densities from groundscatter. *J. Geophys.*  
 515 *Res.*, *103*(A4), 7003–7015. doi: 10.1029/97JA02923
- 516 Ballatore, P., Villain, J. P., Vilmer, N., & Pick, M. (2000). The influence of the in-  
 517 terplanetary medium on SuperDARN radar scattering occurrence. *Ann. Geo-*  
 518 *phys.*, *18*, 1576–1583. doi: 10.1007/s00585-001-1576-2
- 519 Bilitza, D., McKinnell, L.-A., Reinisch, B., & Fuller-Rowell, T. J. (2011). The inter-  
 520 national reference ionosphere today and in the future. *J. Geod.*, *85*, 909–920.  
 521 doi: 10.1007/s00190-010-0427-x
- 522 Bland, E. C., McDonald, J., de Larquier, S., & Devlin, J. C. (2014). Determination  
 523 of ionospheric parameters in real time using SuperDARN HF radars. *J. Geo-*  
 524 *phys. Res. Space Physics*, *119*, 5830–5846. doi: 10.1002/2014JA020076
- 525 Breit, G., & Tuve, M. (1926). A test of the existence of the conducting layer. *Phys.*  
 526 *Rev.*, *28*, 554–575.
- 527 Bristow, W. A., Greenwald, R. A., & Villain, J. P. (1996). On the seasonal de-  
 528 pendence of medium-scale atmospheric gravity waves in the upper atmo-  
 529 sphere at high latitudes. *J. Geophys. Res.*, *101*(A7), 15685–15699. doi:  
 530 10.1029/96JA01010
- 531 Burrell, A. G., Milan, S. E., Perry, G. W., Yeoman, T. K., & Lester, M. (2015).  
 532 Automatically determining the origin direction and propagation mode  
 533 of high-frequency radar backscatter. *Radio Sci.*, *50*, 1225–1245. doi:  
 534 10.1002/2015RS005808
- 535 Chakraborty, S., Ruohoniemi, J. M., Baker, J. B. H., & Nishitani, N. (2018). Char-  
 536 acterization of short-wave fadeout seen in daytime SuperDARN ground scatter  
 537 observations. *Radio Sci.*, *53*, 472–484. doi: 10.1002/2017RS006488
- 538 Chisham, G. (2018). Calibrating SuperDARN interferometers using meteor backscat-  
 539 ter. *Radio Sci.*, *53*, 761–774. doi: 10.1029/2017RS006492
- 540 Chisham, G., Burrell, A. G., Marchaudon, A., Shepherd, S. G., Thomas, E. G., &  
 541 Ponomarenko, P. (2021). Comparison of interferometer calibration techni-  
 542 ques for improved SuperDARN elevation angles. *Polar Sci.*, *28*, 100638. doi:  
 543 10.1016/j.polar.2021.100638
- 544 Chisham, G., Lester, M., Milan, S. E., Freeman, M. P., Bristow, W. A., Grocott,  
 545 A., . . . Walker, A. D. M. (2007). A decade of the Super Dual Auroral Radar  
 546 Network (SuperDARN): Scientific achievements, new techniques and future  
 547 directions. *Surv. Geophys.*, *28*(1), 33–109. doi: 10.1007/s10712-007-9017-8
- 548 Chisham, G., & Pinnock, M. (2002). Assessing the contamination of SuperDARN  
 549 global convection maps by non-F-region backscatter. *Ann. Geophys.*, *20*, 13–  
 550 28. doi: 10.5194/angeo-20-13-2002
- 551 Chisham, G., Yeoman, T. K., & Sofko, G. J. (2008). Mapping ionospheric backscat-  
 552 ter measured by the SuperDARN HF radars – Part 1: A new empirical virtual  
 553 height model. *Ann. Geophys.*, *26*, 823–841. doi: 10.5194/angeo-26-823-2008
- 554 Chu, Y.-H., Wu, K.-H., & Su, C.-L. (2009). A new aspect of ionospheric E region  
 555 electron density morphology. *J. Geophys. Res.*, *114*, A12314. doi: 10.1029/  
 556 2008JA014022
- 557 Custovic, E., McDonald, A. J., Whittington, J., Elton, D., Kane, T. A., & Devlin,  
 558 J. C. (2013). New antenna layout for a SuperDARN HF radar. *Radio Sci.*, *48*,  
 559 722–728. doi: 10.1002/2013RS005156
- 560 de Larquier, S., Ponomarenko, P., Ribeiro, A. J., Ruohoniemi, J. M., Baker,  
 561 J. B. H., Sterne, K. T., & Lester, M. (2013). On the spatial distribution

- 562 of decameter-scale subauroral ionospheric irregularities observed by Su-  
 563 perDARN radars. *J. Geophys. Res. Space Physics*, *118*, 5244–5254. doi:  
 564 10.1002/jgra.50475
- 565 de Larquier, S., Ruohoniemi, J. M., Baker, J. B. H., Ravindran Varrier, N., &  
 566 Lester, M. (2011). First observations of the midlatitude evening anomaly  
 567 using Super Dual Auroral Radar Network (SuperDARN) radars. *J. Geophys.*  
 568 *Res.*, *116*, A10321. doi: 10.1029/2011JA016787
- 569 Frissell, N. A., Baker, J. B. H., Ruohoniemi, J. M., Gerrard, A. J., Miller, E. S.,  
 570 West, M. L., & Bristow, W. A. (2016). Sources and characteristics of medium-  
 571 scale traveling ionospheric disturbances observed by high-frequency radars in  
 572 the North American sector. *J. Geophys. Res. Space Physics*, *121*, 3722–3739.  
 573 doi: 10.1002/2015JA022168
- 574 Gauld, J. K., Yeoman, T. K., Davies, J. A., & Milan, S. E. (2002). SuperDARN  
 575 radar HF propagation and absorption response to the substorm expansion  
 576 phase. *Ann. Geophys.*, *20*, 1631–1645. doi: 10.5194/angeo-20-1631-2002
- 577 Ghezelbash, M., Fiori, R. A. D., & Koustov, A. V. (2014). Variations in the occur-  
 578 rence of SuperDARN F region echoes. *Ann. Geophys.*, *32*, 147–156. doi: 10  
 579 .5194/angeo-32-147-2014
- 580 Greenwald, R. A., Baker, K. B., Dudeney, J. R., Pinnock, M., B., J. T., Thomas,  
 581 E. C., . . . Yamagishi, H. (1995). DARN/SuperDARN: A global view of the  
 582 dynamics of high-latitude convection. *Space Sci. Rev.*, *71*(1), 761–796. doi:  
 583 10.1007/BF00751350
- 584 Greenwald, R. A., Baker, K. B., Hutchins, R. A., & Hanuise, C. (1985). An HF  
 585 phased-array radar for studying small-scale structure in the high-latitude iono-  
 586 sphere. *Radio Sci.*, *20*(1), 63–79. doi: 10.1029/RS020i001p00063
- 587 Greenwald, R. A., Frissell, N., & de Larquier, S. (2017). The importance of elevation  
 588 angle measurements in HF radar investigations of the ionosphere. *Radio Sci.*,  
 589 *52*, 305–320. doi: 10.1002/2016RS006186
- 590 Greenwood, R. I., Parkinson, M. L., Dyson, P. L., & Schulz, E. W. (2011).  
 591 Dominant ocean wave direction measurements using the TIGER Super-  
 592 DARN systems. *J. Atmos. Sol.-Terr. Phys.*, *73*, 2379–2385. doi: 10.1016/  
 593 j.jastp.2011.08.006
- 594 He, L.-S., Dyson, P. L., Parkinson, M. L., & Wan, W. (2004). Studies of  
 595 medium scale travelling ionospheric disturbances using TIGER Super-  
 596 DARN radar sea echo observations. *Ann. Geophys.*, *22*, 4077–4088. doi:  
 597 10.5194/angeo-22-4077-2004
- 598 Hosokawa, K., Iyemori, T., Yukimatu, A. S., & Sato, N. (2000). Character-  
 599 istics of solar flare effect in the high-latitude ionosphere as observed by  
 600 the SuperDARN radars. *Adv. Polar Upper Atmos. Res.*, *14*, 66–75. doi:  
 601 10.15094/00006304
- 602 Hosokawa, K., & Nishitani, N. (2010). Plasma irregularities in the duskside subau-  
 603 roral ionosphere as observed with midlatitude SuperDARN radar in Hokkaido.  
 604 *Radio Sci.*, *45*, RS4003. doi: 10.1029/2009RS004244
- 605 Koustov, A. V., Sofko, G. J., André, D., Danskin, D. W., & Benkevitch, L. V.  
 606 (2004). Seasonal variation of HF radar F region echo occurrence in the mid-  
 607 night sector. *J. Geophys. Res.*, *109*, A06305. doi: 10.1029/2003JA010337
- 608 Koustov, A. V., Ullrich, S., Ponomarenko, P. V., Nishitani, N., Marcucci, F. M., &  
 609 Bristow, W. A. (2019). Occurrence of F region echoes for the polar cap Super-  
 610 DARN radars. *Earth Planet Sp*, *71*, 112. doi: 10.1186/s40623-019-1092-9
- 611 Liu, E. X., Hu, H. Q., Liu, R. Y., Wu, Z. S., & Lester, M. (2012). An adjusted loca-  
 612 tion model for SuperDARN backscatter echoes. *Ann. Geophys.*, *30*, 1769–1779.  
 613 doi: 10.5194/angeo-30-1769-2012
- 614 Makarevich, R. A. (2010). On the occurrence of high-velocity E-region echoes  
 615 in SuperDARN observations. *J. Geophys. Res.*, *115*, A07302. doi:  
 616 10.1029/2009JA014698

- 617 Marcucci, M. F., Coco, I., Massetti, S., Pignalberi, A., Forsythe, V., Pezzopane, M.,  
618 ... Salvati, A. (2021). Echo occurrence in the southern polar ionosphere for  
619 the SuperDARN Dome C East and Dome C North radars. *Polar Sci.*, *28*,  
620 100684. doi: 10.1016/j.polar.2021.100684
- 621 McDonald, A. J., Whittington, J., de Larquier, S., Custovic, E., Kane, T. A., &  
622 Devlin, J. (2013). Elevation angle-of-arrival determination for a standard  
623 and a modified SuperDARN HF radar layout. *Radio Sci.*, *48*, 709–721. doi:  
624 10.1002/2013RS005157
- 625 Milan, S. E., Jones, T. B., Robinson, T. R., Thomas, E. C., & Yeoman, T. K.  
626 (1997). Interferometric evidence for the observation of ground backscatter  
627 originating behind the CUTLASS coherent HF radars. *Ann. Geophys.*, *15*,  
628 29–39. doi: 10.1007/s00585-997-0029-y
- 629 Nishitani, N., & Ogawa, T. (2005). Model calculations of possible ionospheric  
630 backscatter echo area for a mid-latitude HF radar. *Adv. Polar Upper Atmos.*  
631 *Res.*, *19*, 55–62.
- 632 Nishitani, N., Ruohoniemi, J. M., Lester, M., Baker, J. B. H., Koustov, A. V., Shep-  
633 herd, S. G., ... Kikuchi, T. (2019). Review of the accomplishments of mid-  
634 latitude Super Dual Auroral Radar Network (SuperDARN) HF radars. *Prog*  
635 *Earth Planet Sci*, *6:27*. doi: 10.1186/s40645-019-0270-5
- 636 Oinats, A. V., Nishitani, N., Ponomarenko, P., & Ratovsky, K. G. (2016). Di-  
637 urnal and seasonal behavior of the Hokkaido East SuperDARN ground  
638 backscatter: simulation and observation. *Earth Planet Sp*, *68:18*. doi:  
639 10.1186/s40623-015-0378-9
- 640 Ponomarenko, P. V., St. Maurice, J.-P., Hussey, G. C., & Koustov, A. V. (2010).  
641 HF ground scatter from the polar cap: Ionospheric propagation and ground  
642 surface effects. *J. Geophys. Res.*, *115*, A10310. doi: 10.1029/2010JA015828
- 643 Ribeiro, A. J., Ruohoniemi, J. M., Baker, J. B. H., Clausen, L. B. N., Greenwald,  
644 R. A., & Lester, M. (2012). A survey of plasma irregularities as seen by the  
645 midlatitude Blackstone SuperDARN radar. *J. Geophys. Res.*, *117*, A02311.  
646 doi: 10.1029/2011JA017207
- 647 Ribeiro, A. J., Ruohoniemi, J. M., Baker, J. B. H., N., C. L. B., de Larquier, S.,  
648 & Greenwald, R. A. (2011). A new approach for identifying ionospheric  
649 backscatter in midlatitude SuperDARN HF radar observations. *Radio Sci.*, *46*,  
650 RS4011. doi: 10.1029/2011RS004676
- 651 Ruohoniemi, J. M., & Greenwald, R. A. (1997). Rates of scattering occurrence in  
652 routine HF radar observations during solar cycle maximum. *Radio Sci.*, *32(3)*,  
653 1051–1070. doi: 10.1029/97RS00116
- 654 Shand, B. A., Milan, S. E., Yeoman, T. K., Chapman, P. J., Wright, D. M.,  
655 Jones, T. B., & Pederson, L. T. (1998). CUTLASS HF radar obser-  
656 vations of the Odden ice tongue. *Ann. Geophys.*, *16*, 280–282. doi:  
657 10.1007/s00585-997-0063-9
- 658 Shepherd, S. G. (2014). Altitude-adjusted corrected geomagnetic coordinates: Def-  
659 inition and functional approximations. *J. Geophys. Res. Space Physics*, *119*,  
660 7501–7521. doi: 10.1002/2014JA020264
- 661 Shepherd, S. G. (2017). Elevation angle determination for SuperDARN HF radar  
662 layouts. *Radio Sci.*, *52*, 938–950. doi: 10.1002/2017RS006348
- 663 Shepherd, S. G., Sterne, K. T., Thomas, E. G., Ruohoniemi, J. M., Baker, J. B. H.,  
664 Parris, R. T., ... Holmes, J. M. (2020). Bistatic observations with Super-  
665 DARN HF radars: First results. *Radio Sci.*, *55*. doi: 10.1029/2020RS007121
- 666 Tapping, K. F. (2013). The 10.7 cm solar radio flux ( $F_{10.7}$ ). *Space Weather*, *11*,  
667 394–406. doi: doi:10.1002/swe.20064
- 668 Thomas, E. G., Shepherd, S. G., Sterne, K. T., Kotyk, K., Schmidt, M., Bland,  
669 E. C., ... Burrell, A. G. (2020). *SuperDARN Radar Software Toolkit (RST)*  
670 *4.3.1* [Software]. Zenodo. doi: 10.5281/zenodo.3634732
- 671 Titheridge, J. E. (2000). Modelling the peak of the ionospheric E-layer. *J. Atmos.*

- 672 *Sol. Terr. Phys.*, *62*, 93–114. doi: 10.1016/S1364-6826(99)00102-9  
673 Wang, W., Zhang, J. J., Wang, C., Nishitani, N., Yan, J. Y., Lan, A. L., ... Qui,  
674 H. B. (2022). Statistical characteristics of mid-latitude ionospheric ir-  
675 regularities at geomagnetic quiet time: Observations from the Jiamusi and  
676 Hokkaido East SuperDARN HF radars. *J. Geophys. Res. Space Physics*, *127*,  
677 e2021JA029502. doi: 10.1029/2021JA029502
- 678 Yakymenko, K. N., Koustov, A. V., & Nishitani, N. (2015). Statistical study of mid-  
679 latitude E region echoes observed by the Hokkaido SuperDARN HF radar. *J.*  
680 *Geophys. Res. Space Physics*, *120*, 9959–9976. doi: 10.1002/2015JA021685
- 681 Yeoman, T. K., Chisham, G., Baddeley, L. J., Dhillon, R. S., Karhunen, T. J. T.,  
682 Robinson, T. R., ... Wright, D. M. (2008). Mapping ionospheric backscat-  
683 ter measured by the SuperDARN HF radars – Part 2: Assessing Super-  
684 DARN virtual height models. *Ann. Geophys.*, *26*, 843–852. doi: 10.5194/  
685 angeo-26-843-2008
- 686 Yeoman, T. K., Wright, D. M., Stocker, A. J., & Jones, T. B. (2001). An evaluation  
687 of range accuracy in the Super Dual Auroral Radar Network over-the-horizon  
688 HF radar systems. *Radio Sci.*, *36*(4), 801–813. doi: 10.1029/2000RS002558
- 689 Yukimatu, A. S., & Tsutsumi, M. (2002). A new SuperDARN meteor wind mea-  
690 surement: Raw time series analysis method and its application to mesopause  
691 region dynamics. *Geophys. Res. Lett.*, *29*, 1981. doi: 10.1029/2002GL015210



*Radio Science*

Supporting Information for

**Virtual height characteristics of ionospheric and ground scatter observed by mid-latitude SuperDARN HF radars**

E. G. Thomas<sup>1</sup> and S. G. Shepherd<sup>1</sup>

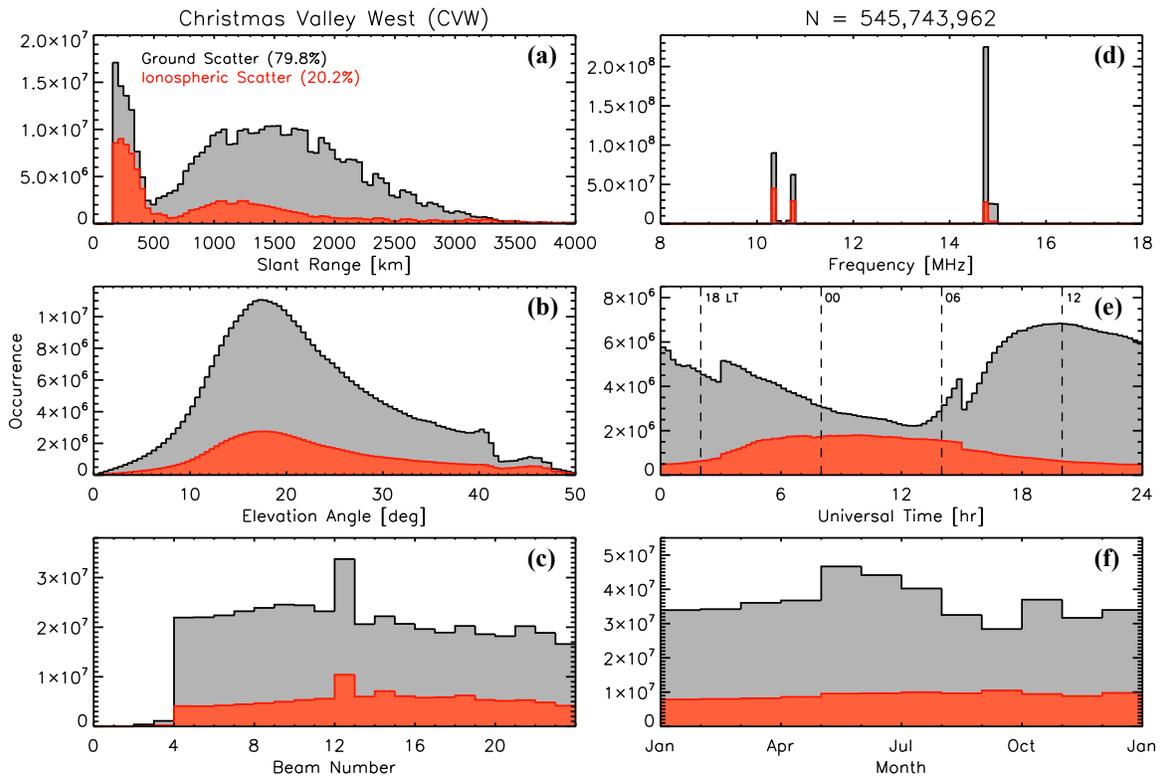
<sup>1</sup>Thayer School of Engineering, Dartmouth College, Hanover, New Hampshire, USA

**Contents of this file**

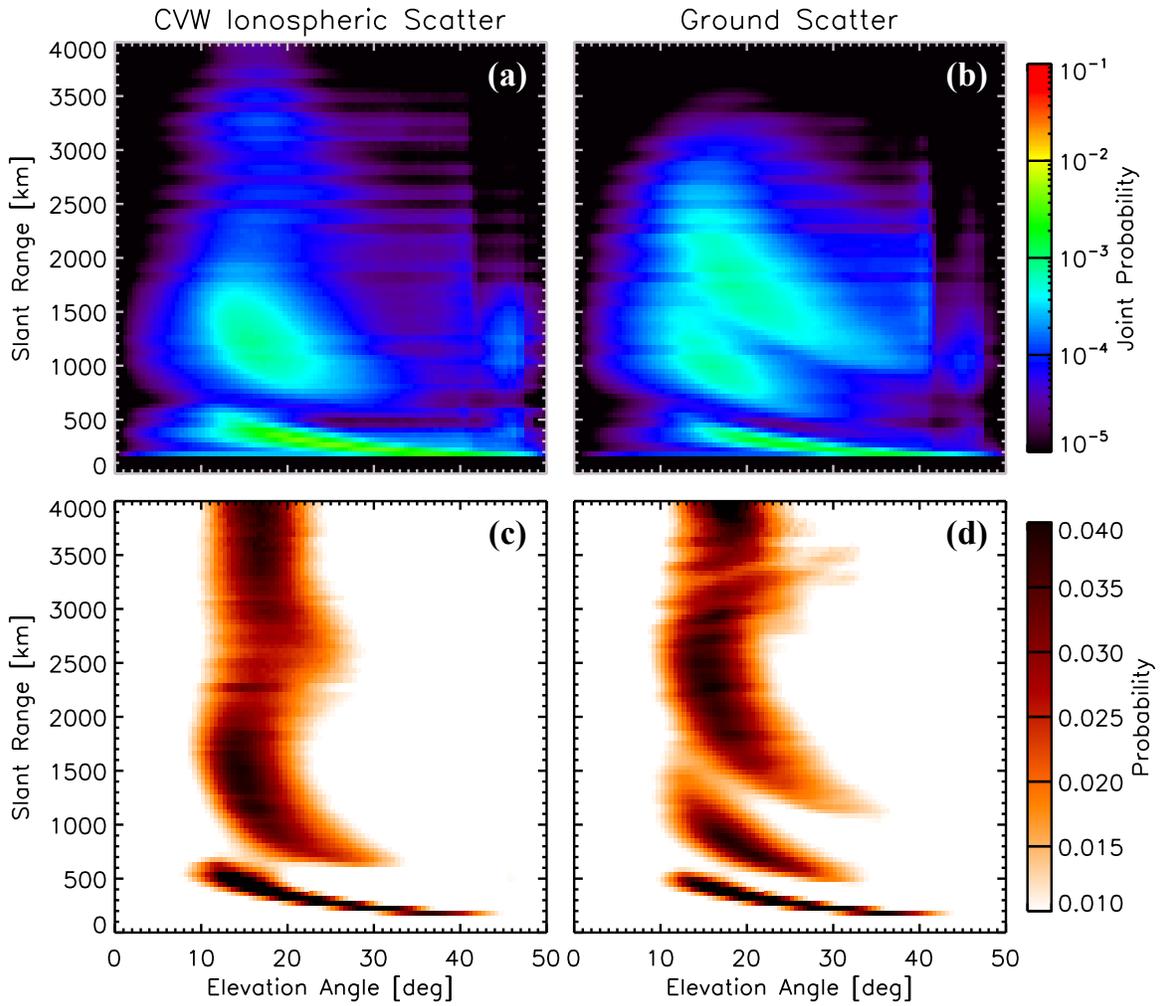
Figures S1 to S6

**Introduction**

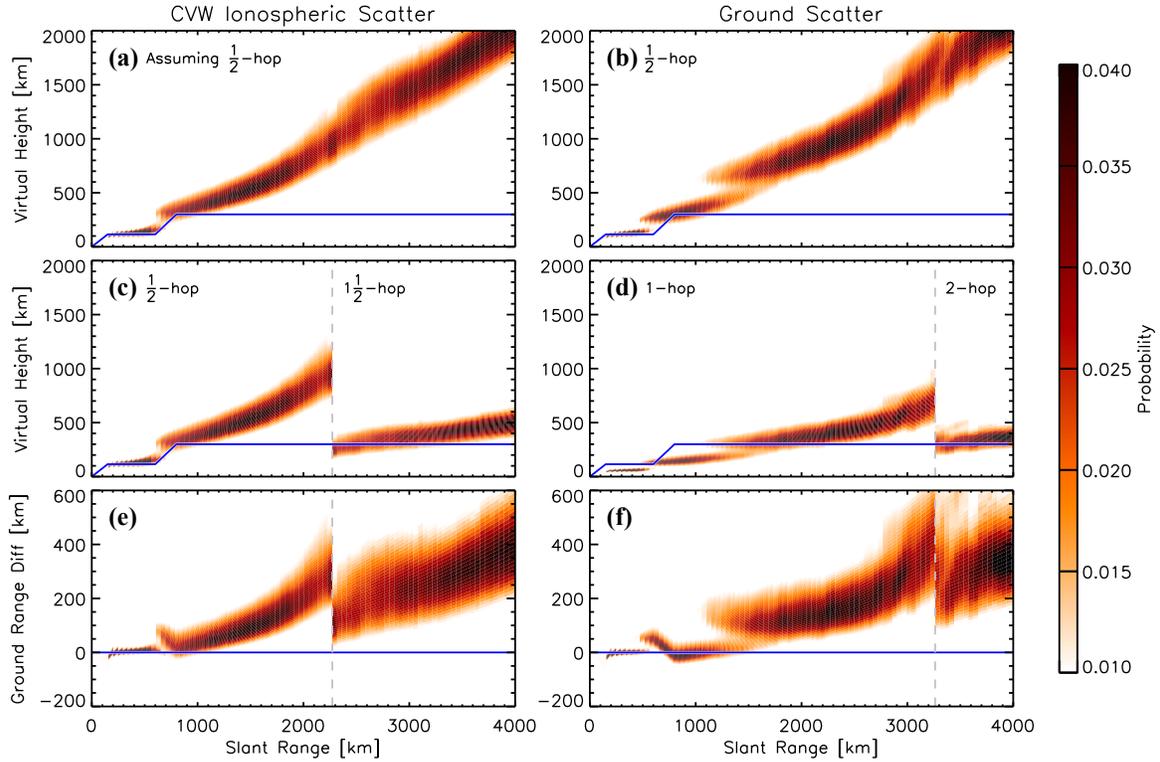
This supporting information provides a matching set of figures for the Christmas Valley West (CVW) SuperDARN radar (Figures S1-S6).



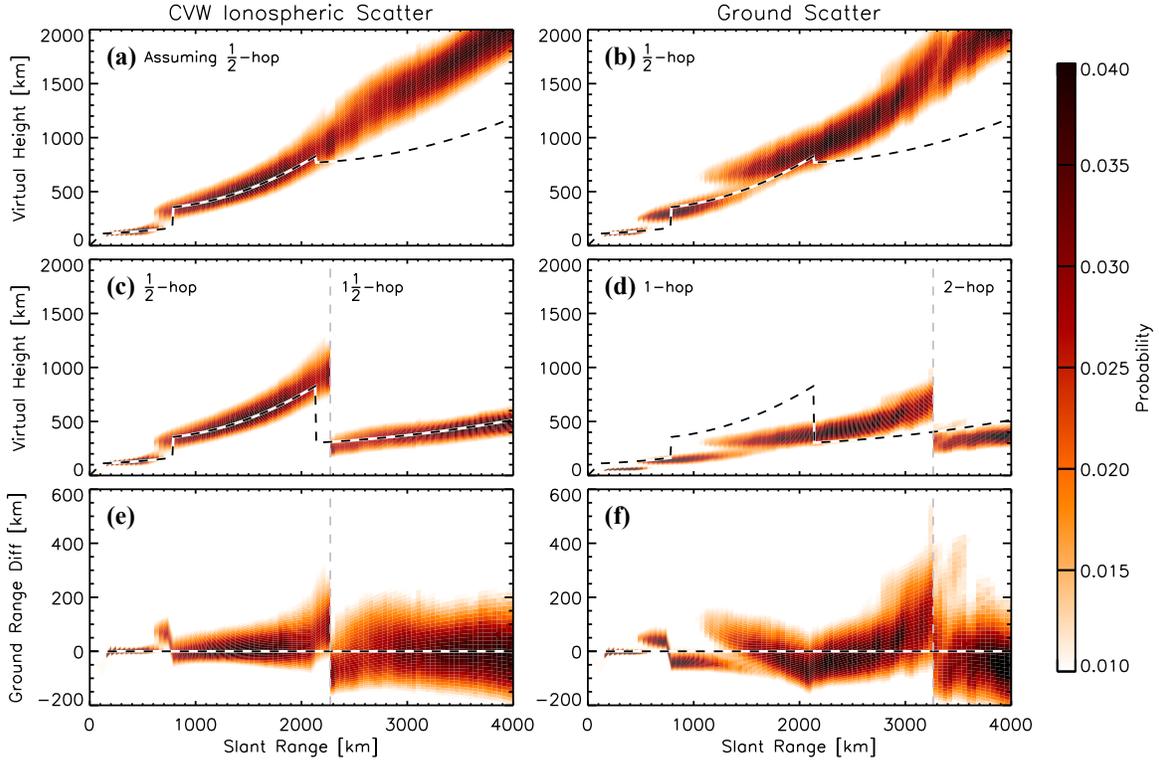
**Figure S1.** Statistical occurrence of Christmas Valley West (CVW) radar observations from 2014–2018 with (a) slant range, (b) elevation angle, (c) azimuthal beam number, (d) frequency, (e) Universal Time, and (f) month, sorted by ionospheric scatter (red) and ground scatter (black). Approximate local times at 6 hr intervals are indicated on panel (e) by vertical dashed lines, and the total number of measurements is given above panel (d).



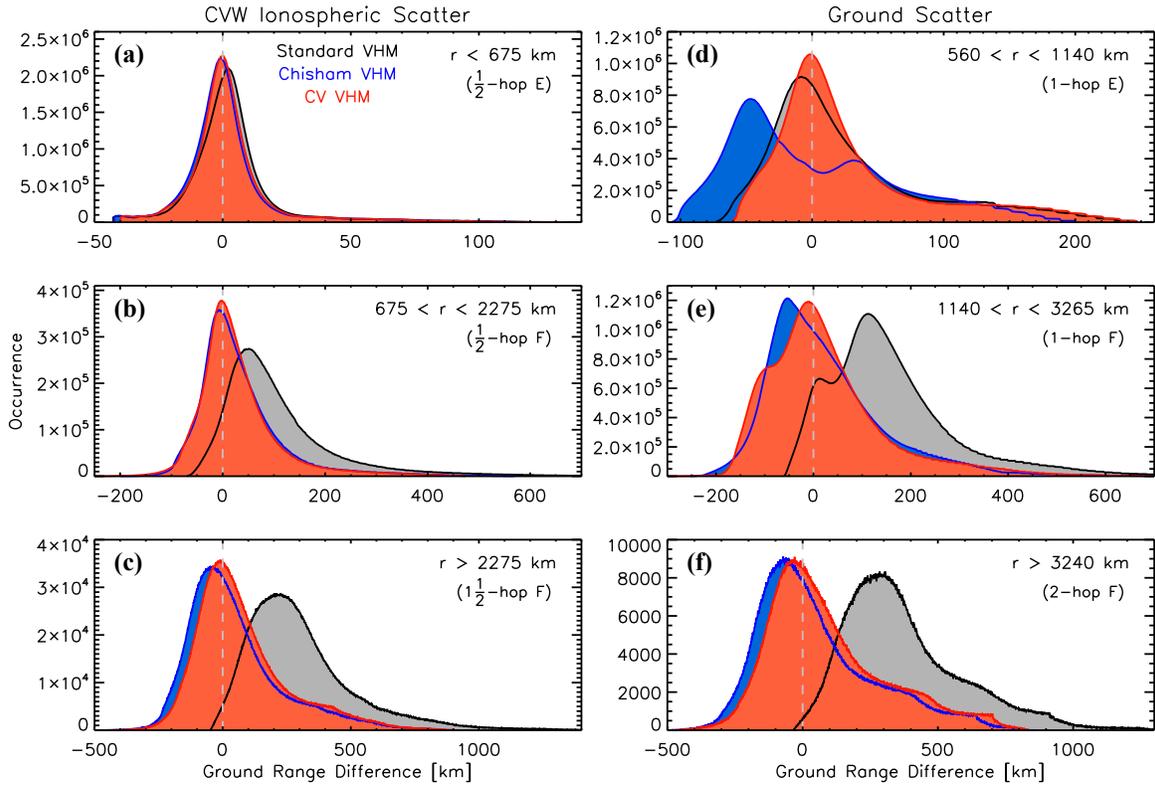
**Figure S2.** (top) Joint probability distributions of elevation angle and slant range observed by the CVW radar for (a) ionospheric and (b) ground scatter, in  $0.5^\circ$  elevation and 45 km range bins. (bottom) The same probability distributions normalized by the maximum occurrence in each range bin; occurrence probabilities of less than 0.010 are not shown.



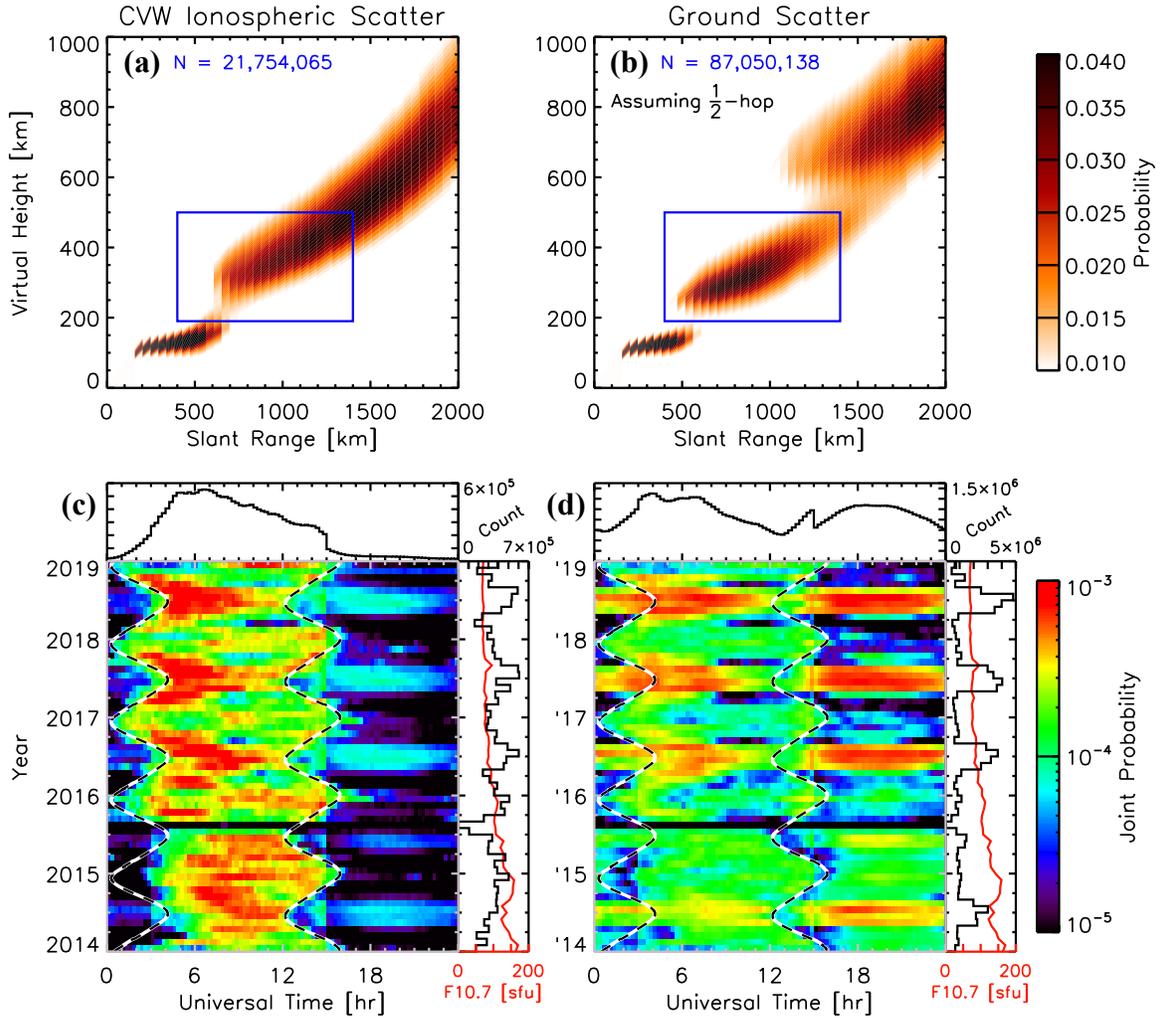
**Figure S3.** Normalized probability distribution of elevation angle and slant range observed by the CVW radar for (a) ionospheric and (b) ground scatter from Figure S2 mapped to slant range versus virtual height assuming a  $\frac{1}{2}$ -hop propagation path, with the standard virtual height model overlaid in blue. (c) Same probability distribution for ionospheric scatter as panel (a) but assuming a  $1\frac{1}{2}$ -hop propagation mode for slant ranges beyond 2250 km (vertical dashed line). (d) Same probability distribution for ground scatter as panel (b) but assuming a 1-hop propagation mode for slant ranges less than 3240 km (vertical dashed line) and a 2-hop propagation mode for further ranges. The bottom row shows the difference in ground ranges for (e) ionospheric and (f) ground scatter from panels (c) and (d) compared to application of the standard virtual height model (which always assumes a  $\frac{1}{2}$ -hop propagation path); positive values indicate the true ground range is closer to the radar than suggested by the model.



**Figure S4.** Normalized probability distributions of slant range and virtual height observed by the CVW radar for (left) ionospheric and (right) ground scatter in the same format as Figure S3, but instead compared against the Chisham virtual height model overlaid as a black and white dashed line. The “pseudo” virtual height predicted by the Chisham virtual height model for  $1\frac{1}{2}$ -hop scatter (beyond 2137 km slant range) is overlaid on panels (a–b), while the “true” virtual height predicted by the model for  $1\frac{1}{2}$ -hop scatter is overlaid on panels (c–d); see text for details. The bottom row shows the difference in ground ranges for (e) ionospheric and (f) ground scatter from panels (c) and (d) compared to application of the Chisham virtual height model (which assumes either a  $\frac{1}{2}$ -hop or  $1\frac{1}{2}$ -hop propagation path); positive values indicate the true ground range is closer to the radar than suggested by the model.



**Figure S5.** Differences in ground range when using the slant range and virtual height observed by the CVW radar for (left) ionospheric and (right) ground scatter compared to the standard (gray), Chisham (blue) and CV (red) virtual height models. Results are organized by (a)  $\frac{1}{2}$ -hop *E*-region, (b)  $\frac{1}{2}$ -hop *F*-region, (c)  $1\frac{1}{2}$ -hop *F*-region, (d) 1-hop *E*-region, (e) 1-hop *F*-region, and (f) 2-hop *F*-region scatter distributions.



**Figure S6.** Normalized probability distributions of slant range and virtual height observed by the CVW radar for (a) ionospheric scatter and (b) ground scatter, centered on the region of possible “mixed” scatter indicated by the blue box. Both the ionospheric and ground scatter-flagged data in panels (a–b) have been mapped to virtual height assuming a  $\frac{1}{2}$ -hop propagation mode for easier comparison, and the number of observations falling within the “mixed” scatter region are given at the top of each panel in blue. (c) Joint probability distribution of Universal Time (UT) and month/year for the ionospheric-flagged scatter within the blue box in panel (a); along the top of the panel is a 1-D histogram of the same data as a function of UT only, while along the right side of the panel is a 1-D histogram as a function of month and year with the monthly average F10.7 solar radio flux overlaid in red. (d) Same as panel (c) but for the ground scatter-flagged data within the blue box in panel (b). Approximate sunrise and sunset times at the midpoint between the CVW radar and “mixed” scatter backscattering volume are overlaid on panels (c–d) as black and white dashed lines.