

# N<sub>2</sub>O rate of change as a diagnostic of the Brewer-Dobson Circulation in the stratosphere

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## Abstract

The Brewer-Dobson Circulation (BDC) determines the distribution of long-lived tracers in the stratosphere; therefore, their changes can be used to diagnose changes in the BDC. We investigate decadal (2005-2018) trends of nitrous oxide (N<sub>2</sub>O) stratospheric columns (12-40 km) as measured by four Fourier transform infrared (FTIR) ground-based instruments and by the Atmospheric Chemistry Experiment Fourier Transform Spectrometer (ACE-FTS), and compare them with simulations by two models: a chemistry-transport model (CTM) driven by four different reanalyses, and the Whole Atmosphere Chemistry-Climate Model (WACCM). The limited sensitivity of the FTIR instruments can hide negative N<sub>2</sub>O trends in the mid-stratosphere because of the large increase in the lowermost stratosphere. When applying the ACE-FTS sampling on model datasets, the reanalyses by the European Centre for Medium Range Weather Forecast (ECMWF) compare best with ACE-FTS, but the N<sub>2</sub>O trends are consistently exaggerated. Model sensitivity tests show that while decadal N<sub>2</sub>O trends reflect changes in transport, these trends are less significant in the northern extratropics due to the larger variability of transport over timescales shorter than two years in that region. We further investigate the N<sub>2</sub>O Transformed Eulerian Mean (TEM) budget in three model datasets. The TEM analysis shows that enhanced advection affects the stratospheric N<sub>2</sub>O trends more than changes in mixing. While no ideal observational dataset currently exists, this model study of N<sub>2</sub>O trends still provides new insights about the BDC and its changes thanks to relevant sensitivity tests and the TEM analysis.

# Evaluation of the N<sub>2</sub>O rate of change to understand the stratospheric Brewer-Dobson Circulation in a Chemistry-Climate Model

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## Key Points:

- Sparse sampling of Atmospheric Chemistry Experiment Fourier Transform Spectrometer exaggerates the stratospheric nitrous oxide trends
- Transformed Eulerian Mean analysis shows that the residual mean advection contributes to the positive nitrous oxide trend in the Tropics
- The Whole Atmosphere Community-Climate Model simulates weaker hemispheric asymmetries of the nitrous oxide trends compared to reanalyses

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**Abstract**

The Brewer-Dobson Circulation (BDC) determines the distribution of long-lived tracers in the stratosphere; therefore, their changes can be used to diagnose changes in the BDC. We evaluate decadal (2005-2018) trends of nitrous oxide ( $\text{N}_2\text{O}$ ) in two versions of the Whole Atmosphere Chemistry-Climate Model (WACCM) by comparing them with measurements from four Fourier transform infrared (FTIR) ground-based instruments, the Atmospheric Chemistry Experiment Fourier Transform Spectrometer (ACE-FTS), and with a chemistry-transport model (CTM) driven by four different reanalyses. The limited sensitivity of the FTIR instruments can hide negative  $\text{N}_2\text{O}$  trends in the mid-stratosphere because of the large increase in the lowermost stratosphere. When applying ACE-FTS measurement sampling on model datasets, the reanalyses from the European Centre for Medium Range Weather Forecast (ECMWF) compare best with ACE-FTS, but the  $\text{N}_2\text{O}$  trends are consistently exaggerated. The  $\text{N}_2\text{O}$  trends obtained with WACCM disagree with those obtained from ACE-FTS, but the new WACCM version performs better than the previous above the Southern Hemisphere in the stratosphere. Model sensitivity tests show that the decadal  $\text{N}_2\text{O}$  trends reflect changes in the stratospheric transport. We further investigate the  $\text{N}_2\text{O}$  Transformed Eulerian Mean (TEM) budget in WACCM and in the CTM simulation driven by the latest ECMWF reanalysis. The TEM analysis shows that enhanced advection affects the stratospheric  $\text{N}_2\text{O}$  trends in the Tropics. While no ideal observational dataset currently exists, this model study of  $\text{N}_2\text{O}$  trends still provides new insights about the BDC and its changes because of the contribution from relevant sensitivity tests and the TEM analysis.

**Plain Language Summary**

The circulation in the stratosphere is characterized by upward motion above the Tropics, followed by poleward and downward motions above the high latitudes. Changes in the pattern of this stratospheric circulation are currently a challenging topic of research. We investigate the decennial changes of this stratospheric circulation using observations and numerical simulations of the long-lived tracer nitrous oxide. Observations are obtained from ground-based and satellite instruments. Numerical simulations include complex atmospheric models that reproduce the chemistry and dynamics of the stratosphere. Both observations and models show differences between the hemispheres in the nitrous oxide decennial changes. Unfortunately, the current observations of nitrous oxide are not perfect. The ground-based instruments cannot correctly measure the changes of nitrous oxide in the northern hemisphere. The satellite does not measure at all times, and it spatially covers more the high latitudes, which negatively affects the measurements of nitrous oxide. On the other hand, model simulations can provide valuable insights into the changes in the stratospheric circulation. They show that changes in the stratospheric circulation cause the differences between hemispheres in the nitrous oxide trends and show that the circulation changes can be associated with different physical processes.

**1 Introduction**

Nitrous oxide ( $\text{N}_2\text{O}$ ) is continuously emitted in the troposphere, with a nearly constant rate of change of 2% per decade, and transported into the stratosphere, where it is destroyed by photodissociation mainly in the Tropics above 35 km (Tian et al., 2020). The atmospheric lifetime of  $\text{N}_2\text{O}$  is approximately 120 years, which makes it an excellent tracer for stratospheric transport studies (Seinfeld & Pandis, 2016). Within the stratosphere, the lifetime of  $\text{N}_2\text{O}$  depends also on the solar activity because of its influence on the photolysis rates, with slightly decreased lifetime during solar maxima and increased lifetime during solar minima (Prather et al., 2015).

$\text{N}_2\text{O}$  enters the stratosphere in the Tropics, and is transported towards higher latitudes by the Brewer-Dobson Circulation (BDC, Dobson et al., 1929; Brewer, 1949; Dob-

76 son, 1956). The BDC is driven by the breaking of tropospheric waves that propagate into  
77 the stratosphere (e.g., Charney & Drazin, 1961) and is often separated into an advective  
78 component, the residual mean meridional circulation (hereafter residual circulation),  
79 and a mixing component (Garny et al., 2014). The residual circulation consists in up-  
80 welling in the Tropics, followed by poleward flow and downwelling over the middle and  
81 high latitudes (Plumb, 2002). The mixing is a two-way exchange of mass that, within  
82 the stratosphere, occurs mostly on isentropic surfaces, thus, it is mainly quasi-horizontal  
83 (Shepherd, 2007). The BDC has a significant impact in determining the stratospheric  
84 distribution of chemical tracers, like ozone and greenhouse gases (e.g., Butchart, 2014),  
85 and in maintaining the observed meridional and vertical temperature structure of the strato-  
86 sphere (Holton, 2004). Long-term changes in the BDC can have significant impacts on  
87 the climate system. One of the most important is the effect on the recovery of strato-  
88 spheric ozone, as a changing BDC would result in changes of its meridional distribution  
89 (e.g., Shepherd, 2008; Dhomse et al., 2018). Changes in the BDC also impact the life-  
90 time of Ozone Depleting Substances (ODS) in the stratosphere (Butchart & Scaife, 2001;  
91 Waugh & Hall, 2002), as well as the water vapor entering the stratosphere through the  
92 Tropics (e.g., Randel & Park, 2019). The troposphere is also affected by BDC changes  
93 because of the impact on the mass exchange with the stratosphere (e.g., ozone, Meul et  
94 al., 2018), and on the ultra-violet radiation reaching the surface (Meul et al., 2016).

95 Given the relevance of the BDC changes, understanding them is thus fundamen-  
96 tal to fully comprehend the past and future evolution of climate. Simulations by Chemistry-  
97 Climate Models (CCMs) robustly project an acceleration of the BDC throughout the strato-  
98 sphere in recent and coming decades due to the increase of greenhouse gases (e.g., Aba-  
99 los et al., 2021). On the other hand, Oberländer-Hayn et al. (2016) argue that the global  
100 BDC trends in the lower stratosphere in CCMs are caused to a large extent by a lift of  
101 the tropopause level in response to global warming rather than an actual speedup of the  
102 circulation. Another significant impact of the increase of greenhouse gases is the shrink-  
103 age of the stratosphere, i.e., the combination of the tropopause rise and the downward  
104 shift of the height of the pressure levels above 55 km, that results from its cooling over  
105 the last decades (Pisoft et al., 2021). Modelling studies have shown that this stratospheric  
106 shrinking can impact the BDC and modulate its changes over the past decades (Šácha  
107 et al., 2019; Eichinger & Šácha, 2020). Such modulation consists in a BDC acceleration  
108 similar to that resulting from the impact of the tropopause lift (Eichinger & Šácha, 2020).  
109 In addition, CCMs simulations show that also the vertical and meridional structure of  
110 the BDC has changed in the past decades in response to climate change (Hardiman et  
111 al., 2014). Other modeling studies have shown that mixing, both on resolved and un-  
112 resolved scales, also plays an important role in the simulated magnitudes of the BDC  
113 changes in addition to changes in the residual circulation among CCMs (e.g., Eichinger  
114 et al., 2019). Recent studies have also shown that ODS, through their impact on ozone,  
115 play a significant role in the modeled BDC changes (Abalos et al., 2019). In particular,  
116 the ODS decrease resulting from the Montreal Protocol, will reduce the global warming-  
117 induced acceleration of the BDC and potentially lead to hemispheric asymmetries in the  
118 BDC trends (Polvani et al., 2019).

119 The BDC and its changes cannot be measured directly (e.g., Minschwaner et al.,  
120 2016), but can be indirectly examined from measurements of stratospheric long-lived trac-  
121 ers (e.g., Engel et al., 2009; Hegglin et al., 2014) or temperature (Fu et al., 2015). Re-  
122 cently, Strahan et al. (2020) used ground-based observations of nitric acid and hydro-  
123 gen chloride to investigate hemispheric-dependent BDC changes in the stratosphere. Sim-  
124 ilarly, space-borne observations of stratospheric tracers are often used to investigate decadal  
125 changes in the BDC using, e.g., hydrogen fluoride (Harrison et al., 2016), ozone (Nedoluha  
126 et al., 2015) or  $N_2O$  (Han et al., 2019). Measurements of stratospheric tracers are often  
127 used to calculate the mean Age of Air (AoA, Hall & Plumb, 1994). The mean AoA is  
128 a widely used diagnostic for stratospheric transport and is defined as the transit time  
129 of an air parcel from the tropical tropopause (or the surface, depending on the defini-

tion) to a certain point of the stratosphere (Vaugh & Hall, 2002). Engel et al. (2017) used balloon-borne observations of carbon dioxide and methane to derive mean AoA trends above the northern mid-latitudes in the mid-lower stratosphere. Engel et al. (2017) found positive but not statistically significant mean AoA trends over about 40 years (corresponding to a possible slowdown of the BDC), which is in contrast with the modeling studies that simulate a significant acceleration of the BDC over the same region (e.g., Abalos et al., 2021). These discrepancies can be partly attributed to the temporal and spatial sparseness of the measurements and to uncertainties in the mean AoA trends derived from real tracers (Garcia et al., 2011; Fritsch et al., 2020). In addition to ground-based measurements, space-borne observations have been used to compute mean AoA trends as well (e.g., Stiller et al., 2012; Haenel et al., 2015). These observational studies using remote sensing measurements have shown a hemispheric asymmetry in the mean AoA trends over 2002-2012, with positive changes in the Northern Hemisphere (NH) and negative changes in the Southern Hemisphere (SH) (e.g., Mahieu et al., 2014; Stiller et al., 2017). The mean AoA indirectly obtained from satellite measurements in these studies does not allow the separation between residual circulation and mixing, which was proven to be important in CCMs (Dietmüller et al., 2018). However, Linz et al. (2021) showed that the effect of mixing can be explicitly calculated using AoA vertical gradients from both models and satellite measurements. In addition, von Clarmann and Grabowski (2021) (similarly to the early study of Holton & Choi, 1988) proposed an alternative method to infer the stratospheric circulation from satellite measurements of long-lived tracers by a direct inversion of the continuity equation.

Reanalysis datasets try to fill the gap between observations and free-running models, providing a global multi-decadal and continuous state of the past atmosphere by assimilating available observations. Dynamical fields from reanalyses can be used to drive Chemistry-Transport Models (CTMs) to simulate the distribution of real and synthetic tracers in the atmosphere. In the past decade, these CTM experiments have been used to investigate BDC changes in reanalyses using the AoA diagnostic (e.g., Monge-Sanz et al., 2012; Diallo et al., 2012; Ploeger et al., 2015). However, significant differences exist in the BDC changes obtained from different reanalyses, both over multi-decadal and decadal time scales (e.g., Abalos et al., 2015; Chabrilat et al., 2018). Furthermore, the computation of mean AoA largely depends on whether the kinematic velocities or the heating rates are used to drive the CTMs, leading to significant differences within the same reanalysis (Ploeger et al., 2019).

This study is based on the work performed by Minganti et al. (2020, hereafter M2020), who evaluated the climatological BDC in the Whole Atmosphere Community Climate Model (WACCM) version 4 (Garcia et al., 2017). The evaluation in M2020 consisted in studying the impact of the BDC on the climatologies of the stratospheric  $N_2O$  abundances and of the  $N_2O$  Transformed Eulerian Mean (TEM) budget (Andrews et al., 1987). This evaluation was performed by comparison with simulations of the Belgian Assimilation System for Chemical Observation Chemistry-Transport Model CTM (BASCOE CTM, Chabrilat et al., 2018) driven by dynamical reanalyses and with the BASCOE reanalysis of Aura Microwave Limb Sounder (MLS) version 2 (BRAM2, Errera et al., 2019). The TEM diagnostic was included in M2020 because it allows separating the effects of transport and chemistry on the rate of change of a stratospheric tracer such as  $N_2O$  (Randel et al., 1994). Within the TEM framework, the impact of transport can be further separated into the impact from the residual circulation and mixing, as was done for ozone and carbon monoxide in Abalos et al. (2013). It is important to note that the mixing obtained from the TEM analysis generally includes contributions from advective transport that are not represented by the residual circulation (Holton, 2004). After studying the climatologies in M2020, the present study aims to evaluate the changing BDC in WACCM in its versions 4 and 6 (Gettelman et al., 2019) by studying multi-decadal and decadal changes of  $N_2O$  in the stratosphere, comparing them with ground-based and space-borne observations and BASCOE CTM simulations. We also evaluate the changes

184 in TEM N<sub>2</sub>O budget in WACCM and in the BASCOE CTM. We compare the model  
 185 simulations with ground-based observations of N<sub>2</sub>O from Fourier transform infrared (FTIR)  
 186 spectrometers that are part of the Network for the Detection of Atmospheric Compo-  
 187 sition Change (NDACC) at four stations in the SH and NH subtropics as well as at mid-  
 188 latitudes (De Mazière et al., 2018, <http://www.ndaccdemo.org/>). We also use satellite  
 189 measurements from the Atmospheric Chemistry Experiment Fourier Transform Spec-  
 190 trometer (ACE-FTS, Bernath et al., 2021). Contrary to M2020, we cannot use N<sub>2</sub>O from  
 191 BRAM2 because of the unrealistic negative drift in the MLS N<sub>2</sub>O dataset (Livesey et  
 192 al., 2021). The BASCOE CTM is driven by four modern reanalyses that are part of the  
 193 SPARC (Stratosphere-troposphere Processes and their Role in Climate) Reanalysis In-  
 194 tercomparison Project (S-RIP, Fujiwara et al., 2017).

195 The present study is structured as follows. Section 2 describes the observational  
 196 and modeling datasets used in this study, as well as the TEM diagnostics and the regres-  
 197 sion model used to derive linear trends. In Section 3, we use FTIR observations to eval-  
 198 uate the trends in the stratospheric N<sub>2</sub>O columns obtained from WACCM and the CTM  
 199 simulations and from satellite measurements. In Section 4, using ACE-FTS as a refer-  
 200 ence, we study the global N<sub>2</sub>O trends in the stratosphere and focus on the differences  
 201 in the trend patterns among datasets. In Section 5, we investigate the N<sub>2</sub>O TEM bud-  
 202 get from WACCM version 6 and a BASCOE simulation in order to separate the impact  
 203 of the residual circulation and mixing on the N<sub>2</sub>O trends. Finally, Section 6 concludes  
 204 the study with a summary of the principal findings.

## 205 2 Data and Methods

206 This section describes the observational and model data as well as the methods used  
 207 in this study (see Tables 1 and 2). Throughout the study, we will refer to the CCMs and  
 208 the BASCOE CTM simulations as "models" to distinguish them from the observations  
 209 obtained from the FTIR and ACE-FTS. For the sake of brevity, we refer to M2020 for  
 210 a more detailed description of the dataset (BASCOE CTM, WACCM version 4, and S-  
 RIP reanalyses) and methods (TEM framework) already used there.

Dataset name	Full Name	Reference	Year range	Vert. resol. + top
WACCM-REFC1	Whole Atmosphere Commu- nity Climate Model	Marsh et al. (2013) Garcia et al. (2017)	1985-2018	L66, 5.96 10 <sup>-6</sup> hPa
WACCM-REFD1	Whole Atmosphere Commu- nity Climate Model	Gettelman et al. (2019)	1985-2018	L70, 5.96 10 <sup>-6</sup> hPa
CTM+ERA1	ECMWF Reanalysis Interim	Dee et al. (2011)	1985-2018	L60, 0.1 hPa
CTM+ERA5	ECMWF Reanalysis 5	Hersbach et al. (2020)	1985-2019	L86, 0.01 hPa
CTM+JRA55	Japanese 55-year Reanalysis	Kobayashi et al. (2015)	1985-2018	L60, 0.2 hPa
CTM+MERRA2	Modern-Era Retrospective analysis for Research and Applications	Gelaro et al. (2017)	1985-2018	L72, 0.01 hPa
ACE-FTS	Atmospheric Chemistry Ex- periment Fourier Transform Spectrometer	Bernath et al. (2021)	2005-present	L42, 150 km

**Table 1.** Overview of the models and satellite measurements used in this study.

Station name	Reference	Location (lat and lon)	Altitude	strato DOFS
Lauder	Zhou et al. (2019)	45.4°S and 169.68°E	370 m	2
Wollongong	Griffith et al. (2012)	34.45°S and 150.88°E	30 m	2
Izaña	García et al. (2021)	28.30°N and 16.48°E	2367 m	1.5
Jungfraujoch	Zander et al. (2008)	46.55°N and 7.98°E	3580 m	1.1

**Table 2.** Overview of FTIR stations considered in this study.

## 2.1 Ground-based FTIR Observations

We use ground-based measurements of stratospheric N<sub>2</sub>O columns obtained at four stations that are part of NDACC: Lauder (New Zealand, 45°S), Wollongong (Australia, 34°S), Izaña (Spain, 28°N) and Jungfraujoch (Switzerland, 46°N) (Zhou et al., 2019). The solar absorption spectra under clear-sky conditions with the ground-based FTIR measurements taken under the auspices of the NDACC allow the acquisition of long-term consistent data sets. The stations have been chosen at the mid-latitudes and subtropics where the observed BDC changes are the largest (e.g., Strahan et al., 2020).

At Jungfraujoch, measurements have been obtained from two spectrometers: an instrument developed at the University of Liège (1984-2008), and a Bruker IFS 120HR (early 1990's-present) (Zander et al., 2008; Prignon et al., 2019). In this study, we use the spectra taken by the Bruker spectrometer to investigate the most recent period. Ground-based measurements of N<sub>2</sub>O profiles at Lauder started in 2001 with a Bruker 120HR spectrometer, replaced in 2018 (with 6 months overlap) by a Bruker 125HR (Strong et al., 2008; Zhou et al., 2019). The Lauder station is particularly relevant as is the only FTIR site of NDACC located in the SH mid-latitudes. The Wollongong station has provided data for the SH subtropics since 1996. Solar spectra were measured with a Bomem instrument until 2007, which was then replaced by a Bruker 125HR (Griffith et al., 2012). N<sub>2</sub>O profiles are also measured at the Izaña Observatory since 1999. This high-altitude station is characterized by excellent conditions for FTIR spectroscopy, with clear sky conditions for most of the year. Observations started using a Bruker 120M spectrometer and continued, since 2005, with a Bruker 125HR (García et al., 2021). The retrieval code for the N<sub>2</sub>O profiles is the SFIT-v4 (v0.9.4.4) for the Jungfraujoch, Lauder and Wollongong stations, and PROFITT9 for the Izaña station (Zhou et al., 2019).

We consider stratospheric N<sub>2</sub>O columns between 12 and 40 km of altitude because the instruments at all stations are the most sensitive to the measured N<sub>2</sub>O profiles over this stratospheric region (not shown). The degrees of freedom for signal (DOFS), which quantify the vertical resolution of the measurement (Rodgers, 2000), vary largely between the stations. For N<sub>2</sub>O, the stratospheric DOFS between 12 and 40 km of the instruments in the SH are approximately 2, allowing the separation of two layers within the stratosphere. On the other hand, the stratospheric DOFS of the instruments in the NH are around 1.5 for Izaña, and 1 for Jungfraujoch, limiting the analysis to one stratospheric layer between 12 and 40 km. Thus, in order to perform a fair comparison, we compute one stratospheric N<sub>2</sub>O column between 12 and 40 km for all stations. In order to take into account the limited sensitivity of the FTIR measurements, we smooth the ACE-FTS data and the model output on the FTIR vertical grid using the FTIR averaging kernels as described in Langerock et al. (2015).

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## 2.2 Spaceborne Measurements - ACE-FTS

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ACE-FTS, onboard the SCISAT Canadian satellite, was launched in August 2003 on a high inclination ( $74^\circ$ ) low earth orbit (650 km) and is still in operation in 2022 (Bernath et al., 2005; Bernath, 2017). The ACE-FTS instrument measures the infrared absorptions from solar occultations between 2.2 and  $13.3 \mu\text{m}$  with a spectral resolution of  $0.02 \text{ cm}^{-1}$ . This allows the retrieval of vertically resolved mixing ratio profiles for 44 molecules and 24 isotopologues from each measurement (Bernath et al., 2020).

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In this study, we use version 4.1 of the ACE-FTS data. It differs from previous versions by the significantly better retrievals at low altitudes and led to substantially improved trends compared to the earlier version 3.5 (Bernath et al., 2021). For  $\text{N}_2\text{O}$ , previous comparisons of v3.6 with independent satellite instruments showed a good agreement below 35 km (within 10%) and larger biases above that level (within 20%, Sheese et al., 2017). In our study,  $\text{N}_2\text{O}$  profiles are filtered for outliers using the method described in Sheese et al. (2017) and are then vertically regridded to a constant pressure vertical grid using a mass-conservative scheme (Bader et al., 2017). For trend analysis, profiles are monthly averaged on latitude bins with  $5^\circ$  spacing from pole to pole.

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In order to compare the trend analysis of model simulations with those obtained by ACE-FTS, the model datasets are first re-sampled from their native temporal and spatial grids (model space) to match those of ACE-FTS (observational space). This is important in particular due to the low sampling of ACE-FTS - only 30 daily profiles due to the solar occultation method. The re-sampling is done by finding model output adjacent in time to each ACE-FTS profile (BASCOE and WACCM datasets used in this study have, respectively, 6 hourly and daily output) and then by linearly interpolating the model values in time and space at the profile geolocation. The re-sampled model datasets are then averaged over a month as done with ACE-FTS.

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## 2.3 BASCOE CTM and Driving Reanalyses

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In this study, we use the BASCOE CTM driven by four dynamical reanalyses: the European Centre for Medium-Range Weather Forecast Interim reanalysis (ERA-I, Dee et al., 2011), and its newer version ERA5 (Hersbach et al., 2020), the Modern-Era Retrospective analysis for Research and Applications version 2 (MERRA2, Gelaro et al., 2017), and the Japanese 55-year Reanalysis (JRA55, Kobayashi et al., 2015). In the following, we provide a brief overview of the BASCOE CTM and the ERA-I, MERRA2 and JRA55 reanalyses, as more detailed information can be found in such companion studies: Chabrillat et al. (2018); Prignon et al. (2019, 2021) and M2020. Since ERA5 is not detailed in these publications, we provide a more detailed description.

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The BASCOE CTM is built on a kinematic transport module (that takes as input the surface pressure and the horizontal winds) with a flux-form semi-Lagrangian (FFSL) advection scheme (Lin & Rood, 1996). The FFSL scheme is run on a common horizontal grid of  $2^\circ \times 2.5^\circ$  for all the reanalyses, while the vertical grid depends on the input reanalysis. The chemical scheme explicitly solves for stratospheric chemistry, and includes 65 chemical species and 243 reactions (Prignon et al., 2019). ERA-I and JRA55 have 60 levels up to 0.1 hPa, MERRA2 has 72 levels up to 0.01 hPa. The model setup, as well as the boundary conditions (including those for  $\text{N}_2\text{O}$ ), are the ones used in Prignon et al. (2019), M2020 and Prignon et al. (2021). Readers are directed towards Chabrillat et al. (2018) for a detailed description of the BASCOE CTM and its driving by the ERA-I, JRA55 and MERRA2 reanalyses.

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The ERA5 reanalysis is the fifth generation of reanalysis produced by the ECMWF and covers the 1979-present period, with a programmed extension back to 1950 (Hersbach et al., 2020). The horizontal resolution is 31 km, with hourly output frequency, and the vertical grid ranges from the surface to 0.01 hPa with 137 levels and with 300-600 m ver-

299 tical spacing in the troposphere and stratosphere, which increases to 1-3 km above 30  
 300 km. ERA5 suffers from a cold bias in the lower stratosphere from 2000 to 2006. For this  
 301 reason, a new analysis (ERA5.1) has been produced for that period to correct for that  
 302 bias (Simmons et al., 2020). In this study, the BASCOE CTM was driven by ERA5.1  
 303 for the 2000-2006 period. For computational reasons, the vertical resolution is reduced  
 304 to 86 levels from the original 137 keeping the original vertical spacing in the stratosphere,  
 305 and we used 6-hourly (0000, 0600, 1200, 1800 UTC) data. As done for the other reanal-  
 306 yses, the ERA5 data on the fine 31-km grid were truncated at wavenumber 47 to avoid  
 307 aliasing on the target  $2.5^\circ \times 2^\circ$  horizontal grid (Chabrillat et al., 2018).

308 In order to further investigate the contribution of transport in ERA5, we performed  
 309 two sensitivity tests with the BASCOE CTM driven by that reanalysis. To isolate the  
 310 contribution of transport, the first sensitivity test consists of a fixed  $\text{N}_2\text{O}$  run, i.e., a BAS-  
 311 COE CTM simulation where  $\text{N}_2\text{O}$  does not increase over time. We accomplished that  
 312 by performing a BASCOE CTM run exactly as the ERA5 simulation but keeping the  
 313  $\text{N}_2\text{O}$  volume mixing ratios at the surface fixed to their values at the beginning of the sim-  
 314 ulation (cst- $\text{N}_2\text{O}$ ). Any  $\text{N}_2\text{O}$  trend for the cst- $\text{N}_2\text{O}$  simulation is therefore due only to  
 315 the effect of transport. The second sensitivity test is a perpetual year simulation that  
 316 is complementary to cst- $\text{N}_2\text{O}$ , and consists of an experiment where the transport does  
 317 not change over time (cst-dyn). In order to include a complete Quasi Biennial Oscilla-  
 318 tion cycle (QBO, Baldwin et al., 2001), we used the years 2006 and 2007 from ERA5.1  
 319 and ERA5, respectively. Those years are unusual (but convenient) because the QBO lasted  
 320 exactly 24 months (see the zonal wind data at Singapore [https://www.geo.fu-berlin](https://www.geo.fu-berlin.de/met/ag/strat/produkte/qbo/singapore.dat)  
 321 [.de/met/ag/strat/produkte/qbo/singapore.dat](https://www.geo.fu-berlin.de/met/ag/strat/produkte/qbo/singapore.dat)). We used the dynamics of the year  
 322 2006 to simulate even years and from the year 2007 for odd years. All the  $\text{N}_2\text{O}$  changes  
 323 simulated by cst-dyn are due to its constant increase at the surface.

## 324 2.4 WACCM

325 In this study, we use two versions of WACCM: version 4 (Marsh et al., 2013; Gar-  
 326 cia et al., 2017) and version 6 (Gettelman et al., 2019). WACCM version 4 (WACCM4)  
 327 is the atmospheric component of the Community Earth System Model version 1.2.2 (CESM,  
 328 Hurrell et al., 2013), which has been developed by the U.S. National Center of Atmo-  
 329 spheric Research. It is the extended (whole atmosphere) version of the Community At-  
 330 mosphere Model version 4 (CAM4, Neale et al., 2013). WACCM4 has a longitude-latitude  
 331 grid of  $2.5^\circ \times 1.9^\circ$  and 66 vertical levels from the surface to about 140 km altitude, with  
 332 1.1-1.75 km vertical spacing in the stratosphere. The physics of WACCM4 is the same  
 333 as CAM4 and the dynamical core is a finite volume with a horizontal discretization based  
 334 on a conservative flux-form semi Lagrangian (FFSL) scheme (Lin, 2004). WACCM4 is  
 335 not able to internally generate the QBO; thus, it is nudged towards observations of strato-  
 336 spheric winds (Matthes et al., 2010). In this study, we use the WACCM4 version included  
 337 within the SPARC (Stratosphere-troposphere Processes And their Role in Climate) Chemistry-  
 338 Climate Model Intercomparison phase 1 (CCMI-1, Morgenstern et al., 2017). In partic-  
 339 ular, we use the REFC1 experiments (WACCM-REFC1), which consist of simulations  
 340 of the recent past (1960-2018) using state-of-the-art historical forcings and observed sea-  
 341 surface temperatures (Morgenstern et al., 2017). For  $\text{N}_2\text{O}$ , the boundary conditions are  
 342 prescribed using the forcing recommended by the CCMI (Eyring et al., 2013). Compared  
 343 to the default WACCM4 version, WACCM-REFC1 includes important modifications of  
 344 the treatment of heterogeneous chemistry and of the gravity waves parameterization, which  
 345 ultimately improve the simulation of ozone in the Southern Hemisphere (Garcia et al.,  
 346 2017). In this study, we use three realizations of the WACCM-REFC1 configuration for  
 347 the 1985-2018 period.

348 Version 6 of WACCM (WACCM6) is the extension to the whole atmosphere of ver-  
 349 sion 6 of CAM that is part of version 2 of CESM (Danabasoglu et al., 2020). The de-  
 350 fault horizontal resolution of WACCM6 is  $0.9^\circ \times 1.25^\circ$  latitude-longitude, with 70 levels

351 in the vertical from the ground to around 140 km, with vertical resolution similar to WACCM4.  
 352 The transition from WACCM4 to WACCM6 involved several changes in the physics and  
 353 chemistry that are described in Gettelman et al. (2019). WACCM6 is part of the Cou-  
 354 pled Model Intercomparison Project Phase 6 (CMIP6, Eyring et al., 2016), and is used  
 355 in the CCMI-2022 activity (i.e., the successor of CCMI-1, Plummer et al., 2021). Within  
 356 CCMI-2022, we use the REFD1 WACCM6 experiments (WACCM-REFD1), i.e., a suite  
 357 of hindcast experiments for the recent past (1960-2018) used to compare with observa-  
 358 tions. The REFD1 experiments use the databases for historical forcings and observed  
 359 sea surface temperatures developed for the CMIP6. The N<sub>2</sub>O emissions are specified fol-  
 360 lowing the CMIP6 recommendation for historical simulations, i.e., following Meinshausen  
 361 et al. (2020). Although WACCM6 can internally produce the QBO, the REFD1 exper-  
 362 iments require a nudged QBO towards observed winds to ensure synchronization with  
 363 historical variability. In this study, we use one realization of the WACCM-REFD1 ex-  
 364 periments for the 1985-2018 period.

## 365 2.5 TEM Diagnostics

For stratospheric tracers, the TEM diagnostics (Andrews et al., 1987) allows sep-  
 arating the impact of transport and chemistry on the zonal mean local rate of change  
 of a tracer with mixing ratio  $\chi$ :

$$\bar{\chi}_t = -v^* \bar{\chi}_y - w^* \bar{\chi}_z + e^{z/H} \nabla \cdot \mathbf{M} + \bar{S} + \bar{\epsilon}, \quad (1)$$

366 where  $\chi$  represents N<sub>2</sub>O,  $\mathbf{M} = -e^{-z/H} (\overline{v'\chi'} - \overline{v'\theta'} \bar{\chi}_z / \bar{\theta}_z, \overline{w'\chi'} + \overline{v'\theta'} \bar{\chi}_y / \bar{\theta}_z)$  is the eddy  
 367 flux vector, and  $(v^*, w^*)$  are the meridional and vertical components of the residual cir-  
 368 culation, respectively. Overbars denote zonal means and prime quantities indicate de-  
 369 viations from it, while subscripts indicate partial derivatives.  $H = 7$  km is the scale height,  
 370 and  $z \equiv -H \log_e(p/p_s)$  is the log-pressure altitude, with the surface pressure  $p_s = 10^5$   
 371 Pa. The  $S$  term is the net rate of change due to chemistry, defined as the difference be-  
 372 tween the production ( $\bar{P}$ ) and loss ( $\bar{L}$ ) rates  $\bar{S} = \bar{P} - \bar{L}$ . The  $\bar{\epsilon}$  contribution represents  
 373 the residual of the budget, i.e., the difference between the actual rate of change of  $\bar{\chi}$  and  
 374 the sum of the transport and chemistry terms on the right-side hand of Eq. 1.

The transport terms in Eq. 1 can be grouped as follows:

$$\bar{\chi}_t = ADV + MIX + \bar{S} + \bar{\epsilon}, \quad (2)$$

375 where  $ADV = (-v^* \bar{\chi}_y - w^* \bar{\chi}_z)$  and  $MIX = e^{z/H} \nabla \cdot \mathbf{M}$  represent the contribution of  
 376 the residual circulation and of the resolved mixing, respectively. We refer to M2020 for  
 377 a more detailed description of the TEM framework applied to the N<sub>2</sub>O mixing ratios in  
 378 the stratosphere and for a comprehensive discussion of the contribution of each term to  
 379 the N<sub>2</sub>O budget.

## 380 2.6 Derivation of Trends with the Dynamical Linear Modelling Tool

381 In this study, we investigate decadal trends using the Dynamical Linear Modeling  
 382 regression tool (DLM, Alsing, 2019). DLM is based on Bayesian inference and provides  
 383 a number of possible models to analyze time series. Each model is characterized by some  
 384 unknown parameters, and the DLM computes the posterior probability distribution of  
 385 those parameters using a combination of Kalman filtering and Markov chain Monte Carlo  
 386 method.

387 For a given atmospheric time-series  $y_t$ , a generic DLM model is composed of four  
 388 components: a linear background trend, a seasonal cycle with 12- and 6-months periods,  
 389 forcing terms described by a number of regressor variables and an auto-regressive com-

390 ponent:

$$\begin{aligned}
 391 \quad y_t &= \beta_{1,t}z_{1,t} + \beta_{2,t}z_{2,t} \dots + \beta_{n,t}z_{n,t} & (3) \\
 392 &+ \beta_{1,t}^{12} \sin(2\pi t/12) + \beta_{2,t}^{12} \cos(2\pi t/12) \\
 393 &+ \beta_{1,t}^6 \sin(2\pi t/6) + \beta_{2,t}^6 \cos(2\pi t/6) \\
 394 &+ \mu_t \\
 395 &+ z_t^{AR} \\
 396 &+ \epsilon_t.
 \end{aligned}$$

397 In Eq. 3, the terms  $\beta_{i,t}z_{i,t}$  represent the contribution to  $y_t$  from the regressors, where  
 398  $z_{i,t}$  is the corresponding time-series for each regressor. The 6- and 12-months seasonal  
 399 cycles are modeled respectively by  $\beta_{1,t}^6 \sin(2\pi t/6) + \beta_{2,t}^6 \cos(2\pi t/6)$  and  $\beta_{1,t}^{12} \sin(2\pi t/12) +$   
 400  $\beta_{2,t}^{12} \cos(2\pi t/12)$ . The  $\mu_t$  term denotes the linear fit term, and  $z_t^{AR}$  the auto-regressive  
 401 term, defined similarly to the Cochrane-Orcutt correction (Kyrölä et al., 2013), and  $\epsilon_t$   
 402 is the uncertainty.

403 Contrarily to a multi-linear regression (MLR) model, the background linear fit  $\mu_t$   
 404 and the amplitudes of the seasonal cycles  $\beta_{i,t}^{6,12}$  in DLM can vary with time (i.e., they  
 405 are non-parametric). Their degrees of time-dependence are the unknown model param-  
 406 eters and are initially set by the user and inferred from the data during the model run.  
 407 Furthermore, the auto-regressive process in the DLM is computed within the model run  
 408 together with the other parameters, not as a post-run correction as done in the MLR,  
 409 and its uncertainties are carefully taken into account within the error propagation. In  
 410 addition, the standard DLM implementation has time-varying (heteroscedastic) uncer-  
 411 tainty distribution, when time-varying uncertainties are available. DLM was recently used  
 412 to investigate stratospheric ozone trends in observations and models (Ball et al., 2017,  
 413 2018). A more detailed description of the DLM models and their implementation can  
 414 be found in Laine et al. (2014). For a more comprehensive review of time-series analy-  
 415 sis using DLM, refer to Durbin and Koopman (2012).

416 As regressor variables, we used the 30 cm radio flux as a solar proxy (de Wit et al.,  
 417 2014), an index for the El-Nino Southern Oscillation (Wolter & Timlin, 2011) from the  
 418 National Oceanic and Atmospheric Administration (<http://www.esrl.noaa.gov/psd/enso/mei/>), and two indices for the QBO at 30 and 50 hPa from the Freie Universität  
 419 Berlin (<http://www.geo.fu-berlin.de/en/met/ag/strat/produkte/qbo/index.html>).  
 420 We fed the DLM model with monthly data, running 3000 samples where the first 1000  
 421 were considered as a warmup and discarded. We also tried 10000 realizations and 3000  
 422 as warmup with very similar results (not shown). We performed several sensitivity tests  
 423 to determine the appropriate values of the initial model parameters, i.e., the degree of  
 424 time-dependence of the linear trend and seasonal cycles, in order to allow a reasonable  
 425 time-dependence without being unrealistic. The different combinations of these values  
 426 did not provide significant differences, so we kept the recommended values.  
 427

428 The linear trends are computed from the distribution of the fit samples  $\mu_t$  as the  
 429 difference between the end and start dates of the considered period ( $\text{delta} = \mu_t[\text{end}] -$   
 430  $\mu_t[\text{start}]$ ), weighted by the number of the years. From the resulting delta distribution,  
 431 the uncertainties associated with the trend are computed as the percentage of its posi-  
 432 tive (negative) values. This percentage can be interpreted as the posterior probability  
 433 that the trend is positive (negative) between the considered dates. In this way, we do  
 434 not make any assumption on the shape of the distribution of the trends.

### 435 3 Stratospheric N<sub>2</sub>O Columns and their Trends

436 Figure 1 shows the linear fits of the monthly stratospheric N<sub>2</sub>O columns (12-40 km)  
 437 at the four FTIR stations, together with the initial N<sub>2</sub>O columns from the observations

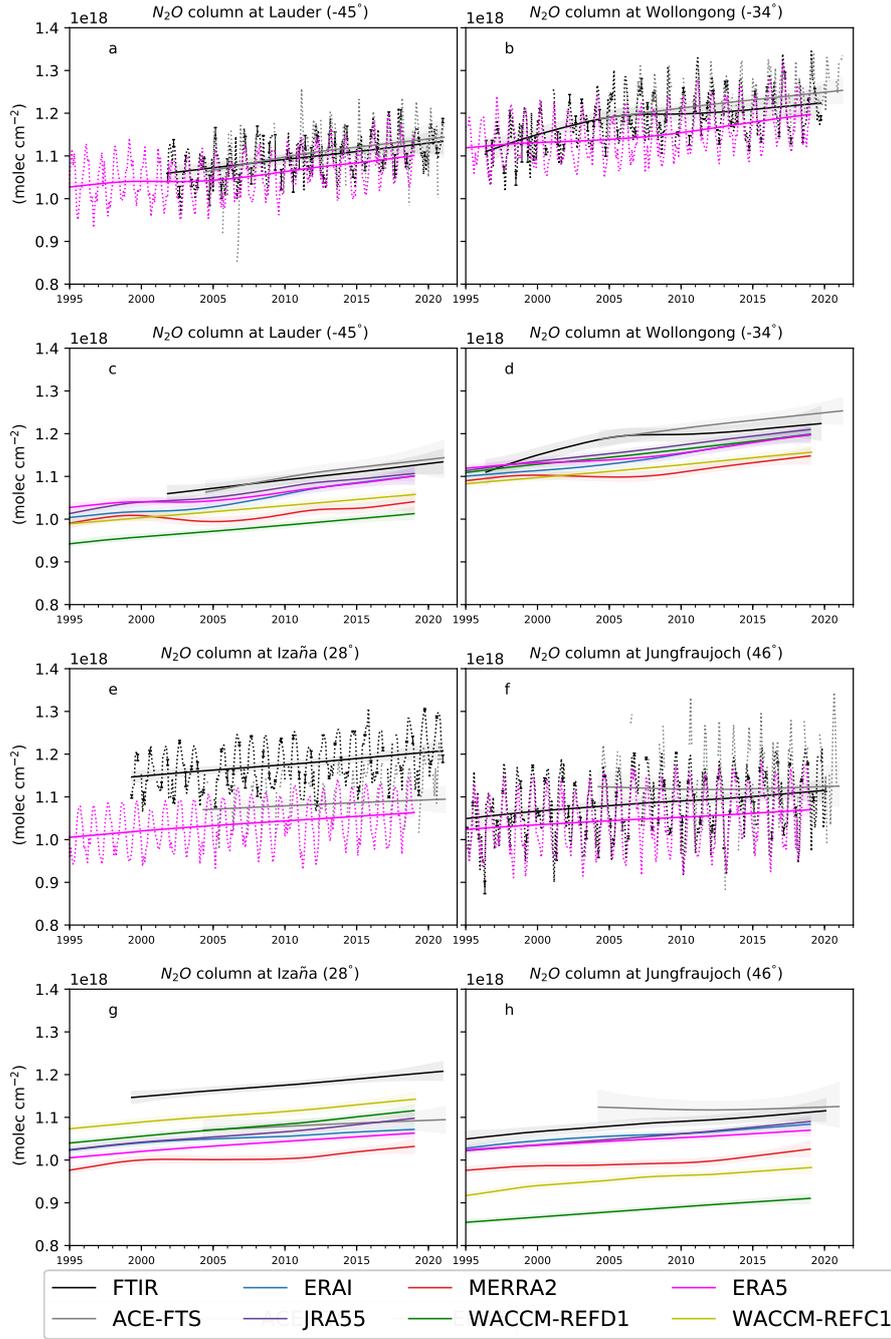
438 and the ERA5 simulation. In this analysis, we do not apply the FTIR time sampling to  
 439 the model output because sensitivity tests using WACCM-REFD1 at each station showed  
 440 no significant impact of the FTIR time sampling on the recovered trends of the N<sub>2</sub>O columns  
 441 (not shown). The stratospheric N<sub>2</sub>O columns computed between 12 and 40 km of alti-  
 442 tude are highly sensitive to the N<sub>2</sub>O increase in the lower stratosphere, which is mainly  
 443 the result of the continuous growth in the troposphere (Tian et al., 2020) and can also  
 444 be impacted by structural changes of the atmosphere (e.g., the global rise of the tropopause,  
 445 Xian & Homeyer, 2019). Consequently, all datasets exhibit an increase in the stratospheric  
 446 N<sub>2</sub>O columns over the last two decades.

447 Above Lauder, the linear fit of the stratospheric N<sub>2</sub>O columns from the ERA5 sim-  
 448 ulation is in agreement with the observations, similarly to JRA55 and ERAI. WACCM-  
 449 REFD1 underestimates the stratospheric N<sub>2</sub>O columns compared to the observations by  
 450 around 10%, and performs worse than its earlier version WACCM-REFC1, which dif-  
 451 fers from the observations by only 5%. At Wollongong, the slope of the linear fit of the  
 452 N<sub>2</sub>O columns measured by the FTIR, and to a lesser extent by ACE-FTS, is steeper be-  
 453 fore 2005 compared to the following period. This change of gradient is not visible in any  
 454 of the model simulations. On the contrary, some of the models show a slower increase  
 455 before 2005, followed by a more rapid increase (e.g., the ERA5 simulation). Contrarily  
 456 to Lauder, the WACCM-REFD1 simulation delivers more realistic stratospheric N<sub>2</sub>O columns  
 457 compared to its previous version WACCM-REFC1.

458 Above Izaña, all the models underestimate the stratospheric N<sub>2</sub>O columns with re-  
 459 spect to the FTIR observations, with the largest difference reaching 14% for MERRA2.  
 460 Concerning ACE-FTS, the bias with FTIR measurements is around 8%, which is qual-  
 461 itatively consistent with the results of Strong et al. (2008), even though they used v2.2  
 462 of ACE-FTS. However, García et al. (2021) showed good agreement above Izaña for tropo-  
 463 spheric N<sub>2</sub>O abundances and total N<sub>2</sub>O columns obtained from independent measure-  
 464 ments. The difference between the stratospheric N<sub>2</sub>O columns measured by FTIR and  
 465 ACE-FTS could be explained by the poor coverage of ACE-FTS over the tropical and  
 466 subtropical regions. Since the ACE-FTS measurements represent a latitude band, the  
 467 observed N<sub>2</sub>O results biased towards the values measured at higher latitudes, where more  
 468 occultations are available (Kolonjari et al., 2018). Since the N<sub>2</sub>O abundances decrease  
 469 poleward (Jin et al., 2009), this could explain the low bias in the stratospheric N<sub>2</sub>O columns  
 470 measured by ACE-FTS compared to those obtained from FTIR.

471 Above Jungfraujoch, there is the largest spread in the linear fits of the stratospheric  
 472 N<sub>2</sub>O columns, with the largest differences reaching around 25% between ACE-FTS and  
 473 WACCM-REFD1. Prignon et al. (2019) compared lower stratospheric columns of chlorod-  
 474 ifluoromethane (HCFC-22) between an earlier WACCM version and FTIR measurements,  
 475 and showed that WACCM consistently underestimates the HCFC-22 columns compared  
 476 to the FTIR measurements. Since both N<sub>2</sub>O and HCFC-22 (which has an atmospheric  
 477 lifetime of 12 years, Prignon et al., 2019) are produced at the surface and transported  
 478 into the stratosphere, this underestimation in WACCM could indicate a shortcoming in  
 479 simulating the accumulation of long-lived tracers in the stratosphere above the north-  
 480 ern mid-latitudes. Indeed, Angelbratt et al. (2011) already highlighted that the strato-  
 481 spheric transport has a large impact on the N<sub>2</sub>O columns above Jungfraujoch compared  
 482 to stations at higher latitudes. Regarding the observational datasets, there is a consid-  
 483 erable disagreement between the FTIR instrument and ACE-FTS before 2012, showing  
 484 increasing and decreasing N<sub>2</sub>O columns, respectively. This is in contrast with the remark-  
 485 ably good agreement in the SH between the two datasets. This difference between the  
 486 stratospheric N<sub>2</sub>O columns in ACE-FTS and FTIR measurements will be further addressed  
 487 in Sect. 4.

488 In the Tropics and above the lower stratospheric mid-latitudes, the N<sub>2</sub>O abundances  
 489 are inversely proportional to the mean AoA (Andrews et al., 2001; Strahan et al., 2011;  
 490 Galytska et al., 2019). The stratospheric N<sub>2</sub>O columns at mid-latitudes considered here



**Figure 1.** Time-series of stratospheric  $\text{N}_2\text{O}$  columns (12-40 km) from observations and models at four stations. Continuous lines show the linear fit obtained by the DLM regression, dashed lines depict the  $\text{N}_2\text{O}$  column data. The color code is shown in the legend. The vertical error bars in panels a,b,e,f represent the standard error of the monthly mean. Panels a,b show Lauder, panels b,d show Wollongong, panels e,g show Izaña and panels f,h show Jungfraujoch. Panels a,b,e,f: DLM fits and data for FTIR and ACE-FTS measurements and the BASCOE simulation driven by ERA5. Panels c,d,g,h: DLM fits for all the datasets considered. The model and satellite data are interpolated to the longitude and latitude of the station, and vertically regridded to match the retrieval layering schemes. After the regridding, the data were smoothed using the FTIR averaging kernels. The colored shadings represent the uncertainties from the 2.5 and 97.5 percentiles of the distributions from the DLM.

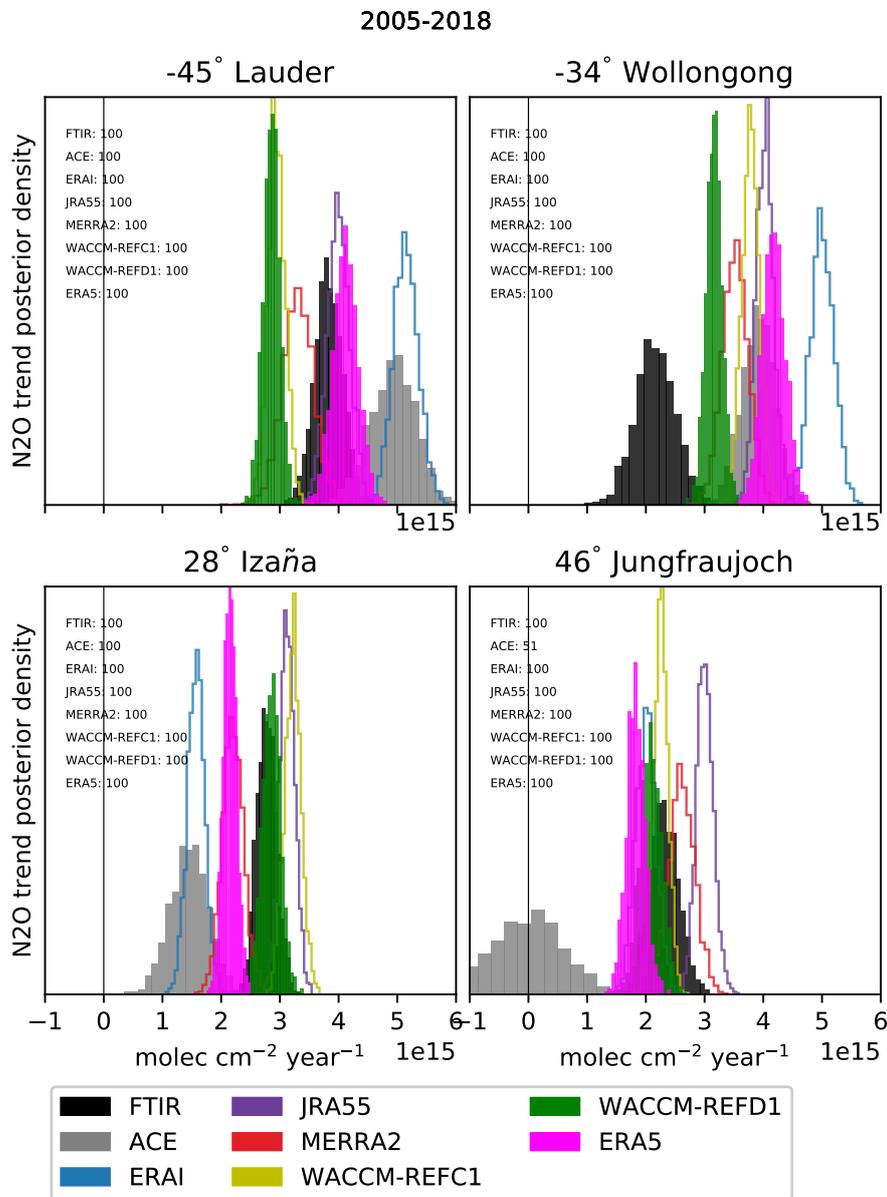
491 are highly sensitive to the N<sub>2</sub>O abundances in the lower stratosphere, hence the inverse  
 492 relationship also holds for the stratospheric N<sub>2</sub>O columns above the mid-latitudes. Thus,  
 493 the lower stratospheric N<sub>2</sub>O columns in MERRA2 compared to the other datasets across  
 494 the stations are consistent with the older mean AoA throughout the stratosphere found  
 495 using MERRA2 by Chabrillat et al. (2018). The N<sub>2</sub>O distribution in the stratosphere  
 496 is opposite also to the total inorganic fluorine F<sub>y</sub>. N<sub>2</sub>O is emitted in the troposphere while  
 497 F<sub>y</sub> is produced in the stratosphere, and, as a consequence of the poleward transport of  
 498 the BDC, N<sub>2</sub>O is removed and F<sub>y</sub> is increased in the stratospheric mid-latitudes. In the  
 499 light of this relationship between N<sub>2</sub>O and F<sub>y</sub>, the underestimated N<sub>2</sub>O columns above  
 500 Lauder and Jungfraujoch in MERRA2 are consistent with larger stratospheric F<sub>y</sub> columns  
 501 in MERRA2 compared to the other reanalyses above those stations (Prignon et al., 2021).

502 Figure 2 shows distributions of the trend of the stratospheric N<sub>2</sub>O columns obtained  
 503 from the respective linear fits over the common period 2005-2018. The N<sub>2</sub>O trends at  
 504 the surface have already been compared for a number of FTIR stations (including Lauder,  
 505 Wollongong and Izaña) against observations from flask samples, showing an excellent agree-  
 506 ment (Zhou et al., 2019).

507 Above Lauder, the N<sub>2</sub>O trends obtained with ERA5 and JRA55 are in good agree-  
 508 ment with the FTIR measurements, but are underestimated in WACCM-REFD1 (around  
 509 25%) with no particular improvement with respect to WACCM-REFC1. The ERAI sim-  
 510 ulation delivers the largest N<sub>2</sub>O trends, with more than 30% difference with respect to  
 511 the FTIR measurements. At Wollongong, the N<sub>2</sub>O trend obtained with the FTIR mea-  
 512 surements is the smallest because the N<sub>2</sub>O increase above that station is smoother com-  
 513 pared to the other datasets. Interestingly, the N<sub>2</sub>O trend simulated by WACCM-REFD1  
 514 is the closest to the trend obtained from the FTIR observations, while the trend obtained  
 515 with ERA5 is almost twice as large. As for Lauder, the N<sub>2</sub>O trends obtained from ERAI  
 516 are the largest at this station. Above Izaña, WACCM-REFD1 agrees remarkably well  
 517 with the FTIR (difference around 3%), while the trend from ERA5 lies between the trends  
 518 measured from FTIR and ACE-FTS, with around 20% difference compared to FTIR.  
 519 Above Jungfraujoch, the trend in the N<sub>2</sub>O columns from WACCM-REFD1 agrees with  
 520 the trend from the FTIR within 10% difference and is similar to what is obtained with  
 521 ERA5. The largest trends are obtained with MERRA2 and JRA55, reaching 13% and  
 522 30% difference compared to the FTIR, respectively. The decreasing N<sub>2</sub>O stratospheric  
 523 column in ACE-FTS before 2012 results in a near-zero trend, which is in contrast with  
 524 the trends obtained by the other datasets, which approximately range from 2 to 3 × 10<sup>15</sup>  
 525 molec cm<sup>-2</sup> year<sup>-1</sup>.

526 Considering decadal changes, the observations and the ERA5 and ERAI simula-  
 527 tions show larger trends of the stratospheric N<sub>2</sub>O columns in the SH than in the NH,  
 528 especially at mid-latitudes (respectively Lauder and Jungfraujoch). WACCM-REFD1  
 529 also shows this hemispheric difference at mid-latitudes, which is a clear improvement with  
 530 respect to WACCM-REFC1. Those asymmetries are consistent with the results of Strahan  
 531 et al. (2020), who found significantly negative mean AoA trends in the SH compared to  
 532 the NH using HCl and HNO<sub>3</sub> measured at several ground-based FTIR stations. In ad-  
 533 dition, the hemispheric differences of the N<sub>2</sub>O trends are also consistent with the results  
 534 of Prignon et al. (2021), who found larger and more significant F<sub>y</sub> trends from FTIR above  
 535 Jungfraujoch than above Lauder.

536 We conclude the section by providing a short description of the limits of using strato-  
 537 spheric columns of N<sub>2</sub>O from FTIR measurements. As mentioned earlier, the stratospheric  
 538 N<sub>2</sub>O columns between 12 and 40 km are primarily influenced by the steady increase in  
 539 the lowermost stratosphere below 15 km. The DOFS of the FTIR instrument at Jungfrau-  
 540 joch for the stratosphere (12-40 km) is close to 1.1. Thus, the FTIR measurements at  
 541 that station cannot resolve more than one partial column between 12 and 40 km, which  
 542 can hinder the detection of N<sub>2</sub>O trends in the middle and upper stratosphere (i.e., above  
 543 30 km) because of the influence of the increase in the lowermost stratosphere. Indeed,



**Figure 2.** Posterior probability of positive changes of the DLM linear trend of the stratospheric N<sub>2</sub>O columns (12-40 km) for the four FTIR stations (2005-2018). The color code is shown in the legend. For reference, the N<sub>2</sub>O trend in the troposphere (5.5-10.5 km) is approximated from the data in Bernath et al. (2020) as 4.3e15 molec cm<sup>-2</sup> year<sup>-1</sup>.

544 it was shown that stratospheric N<sub>2</sub>O trends over the last decades, obtained both from  
 545 satellite measurements and model simulations, do not consist of just a global increase,  
 546 but largely depend on latitude and height (e.g., Froidevaux et al., 2019). Therefore, we  
 547 will consider latitudinal- and vertical-dependent trends of N<sub>2</sub>O mixing ratios in the fol-  
 548 lowing section.

## 549 4 Global N<sub>2</sub>O Linear Trends

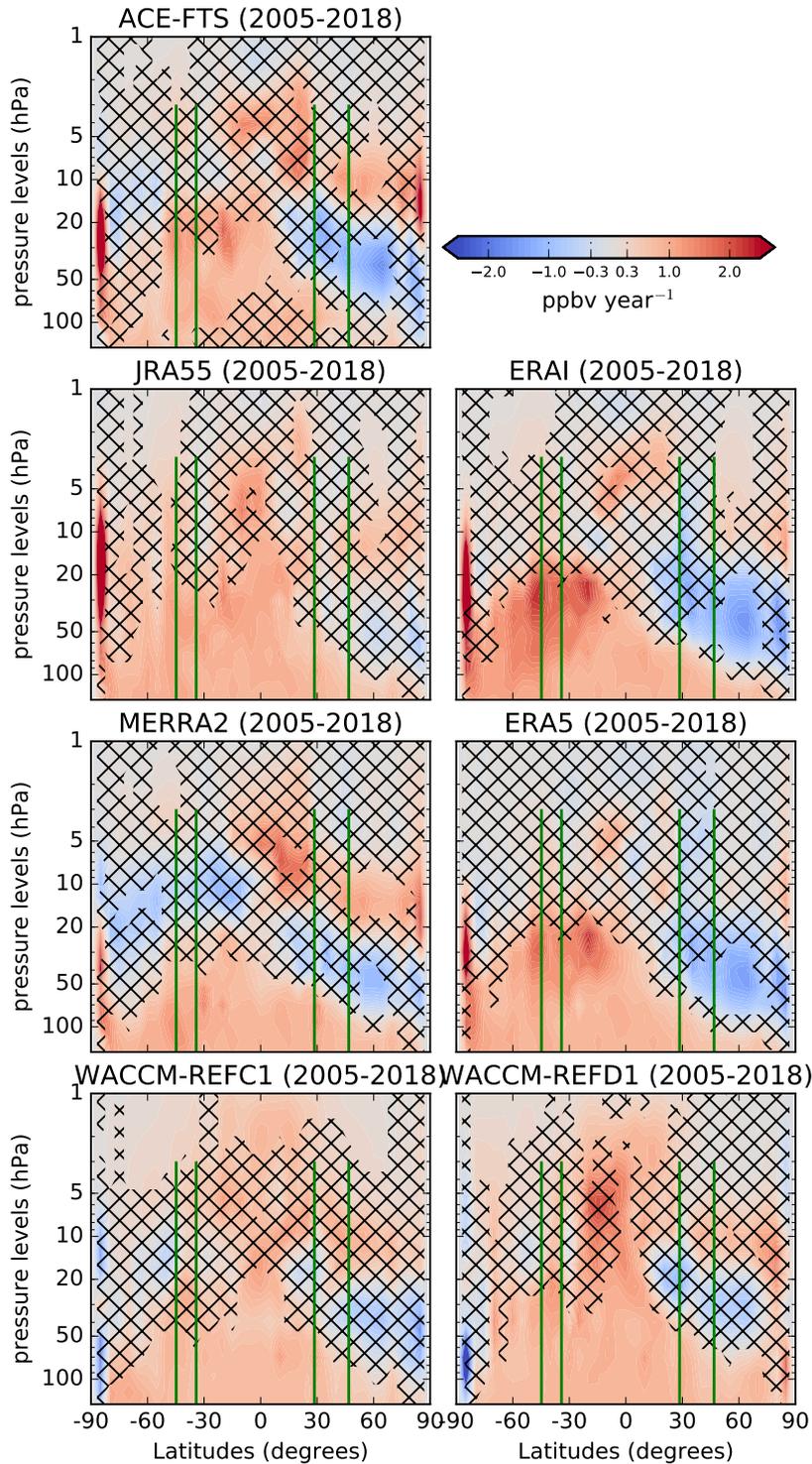
### 550 4.1 Trends in the ACE-FTS Observational Space

551 Figure 3 shows latitude-vertical cross sections of the linear trends of the N<sub>2</sub>O mix-  
 552 ing ratios for the various datasets, over the 2005-2018 period. In order to reduce the sam-  
 553 pling bias, the model datasets are sampled in space and time as the ACE-FTS measure-  
 554 ments before the computation of the trends. We use the ACE-FTS measurements as a  
 555 reference, because they encompass this period with global coverage and good stability  
 556 (Bernath et al., 2020, 2021).

557 In the upper stratosphere above 10 hPa, the N<sub>2</sub>O trends from ACE-FTS are pos-  
 558 itive, with larger trends in the NH that are found significant at lower levels than in the  
 559 SH. The ERAI-driven simulation qualitatively reproduces these patterns in the upper  
 560 stratosphere, while the other model datasets differ from ACE-FTS, especially ERA5. A  
 561 common feature among all datasets is an increase in N<sub>2</sub>O above the Equator in the up-  
 562 per stratosphere, around 5 hPa. At those altitudes of the tropical pipe, the upward trans-  
 563 port of N<sub>2</sub>O by the residual circulation reaches its maximum (see M2020).

564 In the mid-lower stratosphere below 20 hPa, ACE-FTS shows a clear hemispheri-  
 565 cal asymmetry (meridional dipole) in the N<sub>2</sub>O trends, with significantly negative val-  
 566 ues in the NH and significantly positive in the SH. Above the location of Jungfrauoch  
 567 (the most northern vertical green line), the negative N<sub>2</sub>O trend detected by ACE-FTS  
 568 in the mid-lower stratosphere is responsible for the disagreement with the FTIR obser-  
 569 vations discussed in the previous section, as the layer of the stratospheric N<sub>2</sub>O column  
 570 encompasses regions of both positive (lowermost and upper stratosphere) and negative  
 571 (mid-lower stratosphere) N<sub>2</sub>O trends. The meridional dipole is significant also over a shorter  
 572 period (2005-2012, not shown) and corroborates a number of previous findings over that  
 573 period using satellite measurements of HCl (Mahieu et al., 2014) and mean AoA derived  
 574 from space-borne measurements of SF<sub>6</sub> (Haenel et al., 2015). In regions where the N<sub>2</sub>O  
 575 abundances are larger than 100 ppbv, i.e., approximately below 10 hPa, the N<sub>2</sub>O linear  
 576 trends are opposite to those obtained with its product NO<sub>2</sub>, because the two tracers are  
 577 correlated by an inverse linear relationship (Plumb & Ko, 1992). Below 20 hPa, the N<sub>2</sub>O  
 578 meridional dipole from ACE-FTS is consistent with the pattern of the decadal trends  
 579 of NO<sub>2</sub> obtained from independent satellite measurements (Galytska et al., 2019; Dubé  
 580 et al., 2020).

581 The meridional dipole in the N<sub>2</sub>O trends derived from ACE-FTS is generally re-  
 582 produced by the CTM simulations, with ERAI and ERA5 delivering trends that are most  
 583 similar to the satellite measurements. Prignon et al. (2021) used the same simulations  
 584 as the present study to investigate global stratospheric trends of total inorganic fluorine  
 585 F<sub>y</sub>. The dipoles obtained here in the N<sub>2</sub>O trends from the ECMWF reanalyses are con-  
 586 sistent with the opposite trends of F<sub>y</sub> for almost the same period (Prignon et al., 2021).  
 587 For WACCM, the strength of the N<sub>2</sub>O meridional dipole is globally reduced compared  
 588 to ACE-FTS, with weaker and not significant negative N<sub>2</sub>O trends over the NH. How-  
 589 ever, WACCM-REFD1 performs better than WACCM-REFC1 over the SH, with stronger  
 590 and significant positive N<sub>2</sub>O trends that reach 30 hPa, similarly to those obtained with  
 591 ACE-FTS in the same region. This improvement is possibly related to the changes in  
 592 the parametrization of the gravity waves (i.e., small-scale tropospheric waves that drive  
 593 the BDC) in WACCM version 6 compared to version 4 that followed the increase of its



**Figure 3.** Latitude-pressure cross-sections of  $\text{N}_2\text{O}$  linear trends ( $\text{pptv year}^{-1}$ ) obtained from the DLM (2005-2018). The  $\text{N}_2\text{O}$  simulated by the model is interpolated to the location and timing of the observations, see text for details. The black crosses indicate grid-points where the probability of positive/negative  $\text{N}_2\text{O}$  changes is smaller than 95%. The green vertical lines identify the position of the FTIR stations together with their vertical coverage.

horizontal resolution (Gettelman et al., 2019). This new parametrization results in a good agreement between the gravity waves simulated by WACCM and the observations in the Tropics (Alexander et al., 2021). For tracers, the favorable effect of adjusting the parameterization of the gravity waves in WACCM was shown for ozone in the extratropical SH by Mills et al. (2017). Over the same region, the improved N<sub>2</sub>O trends in WACCM-REFD1 compared to WACCM-REFC1 could be attributed to the new parameterization of the gravity waves. This beneficial impact would be consistent with the results of M2020, which showed similar improvements in the N<sub>2</sub>O climatologies between two WACCM versions differing by the parametrization of gravity waves over the SH.

In the lowermost stratosphere (pressure greater than 100 hPa), all models and ACE-FTS show positive N<sub>2</sub>O trends, resulting from the constant increase in the troposphere. However, the N<sub>2</sub>O increase in the lowermost stratosphere (below 70 hPa) over the Tropics and the NH is not significant in ACE-FTS, contrary to the model simulations. This difference could be related to the stronger trends in the tropopause rise in the models: around 50 m/decade in CCMs (including WACCM) and ERA5 (Pissoft et al., 2021; Darrag et al., 2022) compared to the observations (around 35 m/decade, Darrag et al., 2022) when using the tropopause definition from the World Meteorological Organization.

## 4.2 Trends in the Model Space

Figure 4 shows the N<sub>2</sub>O trends as in Fig. 3, but without applying the ACE-FTS spatial and temporal sampling. A comparison between each model simulation in the observation and model space (respectively Fig. 3 and Fig. 4) reveals large differences in the N<sub>2</sub>O decadal trends. Generally, the sampling of the ACE-FTS observations enhances the trends simulated by the models, both in the negative and positive directions. For the ERA5 simulation, the significantly negative trend in the NH in observational space becomes insignificant in model space. In addition, one notes immediately that the N<sub>2</sub>O trends in the WACCM simulations change sign, with negative trends in the NH in the observational space becoming weakly positive in model space. In particular for WACCM-REFD1, the N<sub>2</sub>O trends over the northern mid-latitudes in the mid-low stratosphere substantially increase from -0.5 ppbv year<sup>-1</sup> in observational space to 0.3 ppbv year<sup>-1</sup> in native model space. However, this difference is not significant because neither of the N<sub>2</sub>O trends above that region is statistically significant with 95% probability.

For satellite measurements, the impact of the sampling in the detection of trends in long-lived species (including N<sub>2</sub>O) has been evaluated in Millán et al. (2016). They concluded that large errors may arise in the detected trends for coarse and non-uniform sampling obtained with occultation instruments (such as ACE-FTS), and that long time scales are required for a robust trend detection from these datasets. Such errors also occur in the models when they are sampled in space and time as the observations. In particular, within the DLM, the non-uniform time sampling of ACE-FTS considerably increases the standard deviation of the error in the N<sub>2</sub>O time series, which is zero for regular time sampling. This difference plays a role when deriving trends over these relatively short (decadal) time scales. For example, the non-uniform ACE-FTS sampling applied to the ERA5 output results in negative N<sub>2</sub>O trends that are 4 times stronger compared to the native grid above the northern mid-latitudes between 50 and 70 hPa. For WACCM, the issue of downsampling was also raised by Garcia et al. (2011) when comparing mean AoA trends obtained from balloon-borne observations and simulated by the model. Garcia et al. (2011) showed that sampling the model as the observations would deliver positive and non-significant mean AoA trends, similarly to the observations. We find consistent results with the WACCM simulations: sampling the WACCM output as the observations drives the N<sub>2</sub>O trends towards the observed values. In addition, the non-significant negative N<sub>2</sub>O trends simulated by WACCM are compatible with the non-significant positive mean AoA trends found by Garcia et al. (2011) when downsampling WACCM at

Figure 4. As in Figure 3, but in the model space.

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