Lithospheric Erosion in the Patagonian Slab Window, and Implications for Glacial Isostasy

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November 23, 2022

Abstract

The Patagonian slab window has been proposed to enhance the solid Earth response to ice mass load changes in the overlying Northern and Southern Patagonian Icefields (NPI and SPI, respectively). Here we present the first regional seismic velocity model covering the entire north-south extent of the slab window. A slow velocity anomaly in the uppermost mantle indicates warm mantle temperature, low viscosity, and possibly partial melt. Low velocities just below the Moho suggest that the lithospheric mantle has been thermally eroded over the youngest part of the slab window. The slowest part of the anomaly is north of 49°S, implying that the NPI and the northern SPI overlie lower viscosity mantle than the southern SPI. This comprehensive seismic mapping of the slab window provides key evidence supporting the previously hypothesized connection between post-Little Ice Age anthropogenic ice mass loss and rapid geodetically observed glacial isostatic uplift ([?] 4 cm/yr).

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13 Key Points:

- We observe anomalously low upper mantle shear wave velocities (less than 4.1 km/s) in
 the Patagonian slab window
- The lithospheric mantle has been thermally eroded over the youngest part of the slab
 window
- Low viscosities in the slab window link observed rapid geodetic uplift to geologically
 recent ice mass loss in Patagonia

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21 Abstract

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- load changes in the overlying Northern and Southern Patagonian Icefields (NPI and SPI,
- 24 respectively). Here we present the first regional seismic velocity model covering the entire north-
- 25 south extent of the slab window. A slow velocity anomaly in the uppermost mantle indicates
- 26 warm mantle temperature, low viscosity, and possibly partial melt. Low velocities just below the
- 27 Moho suggest that the lithospheric mantle has been thermally eroded over the youngest part of
- the slab window. The slowest part of the anomaly is north of 49°S, implying that the NPI and the
- 29 northern SPI overlie lower viscosity mantle than the southern SPI. This comprehensive seismic
- 30 mapping of the slab window provides key evidence supporting the previously hypothesized 31 connection between post-Little Ice Age anthropogenic ice mass loss and rapid geodetically
- 31 connection between post-Little ice Age anthropogenic ice mass loss a 32 observed glacial isostatic uplift (> 4 cm/yr).

33 Plain Language Summary

- A gap in the subducting plate beneath Patagonia has enabled hotter, less viscous mantle material
- to flow underneath South America. Icefields in the Austral Andes above the gap in the plate have
- 36 recently been shrinking, removing weight that had caused the continent to flex downward. We
- use seismic data to image the subsurface structure and find very low seismic velocity within and
- around the gap, as well as thinning of the rigid South American lithosphere overlying the gap.
- 39 The low mantle velocity implies that mantle viscosity is also low beneath the shrinking icefields,
- 40 and low viscosity enables the region to rebound upwards.

41 **1 Introduction**

Slab windows form when a spreading ridge subducts and the plates continue to diverge, 42 opening a gap in the subducting plate interface (Groome & Thorkelson, 2009; Thorkelson, 43 1996). Volcanic products associated with several Cenozoic slab windows can be found at 44 45 subduction margins around the Pacific Ocean, indicating that this phenomenon is widespread (McCrory et al., 2009). The Patagonian slab window began forming ~ 18 Ma, when the Chile 46 Ridge started subducting beneath South America near 54°S (Breitsprecher & Thorkelson, 2009). 47 The Chile Triple Junction (CTJ) has since migrated north to its present-day location offshore the 48 Península de Taitao near 46.5°S. Expressions of the slab window include gaps in arc volcanism 49 and subduction zone seismicity (Agurto-Detzel et al., 2014; DeLong et al., 1979); adakitic 50 volcanism near slab edges (Bourgois et al., 2016; Gorring et al., 1997; Stern & Kilian, 1996; 51 Thorkelson & Breitsprecher, 2005); near-trench volcanic activity (Forsythe et al., 1986; Guivel 52 et al., 2003; Marshak & Karig, 1977); anomalously high heat flow (Ávila & Dávila, 2018; Cande 53 et al., 1987) and low upper mantle seismic velocity (Gallego et al., 2010; Russo, VanDecar, et 54 al., 2010); positive dynamic topography (Georgieva et al., 2016; Guillaume et al., 2009) 55 associated with low-viscosity asthenospheric mantle upwelling (Boutonnet et al., 2010; Gorring 56 et al., 1997); and mantle flow patterns influenced by the slab window geometry (Murdie & 57 Russo, 1999; Russo, Gallego, et al., 2010; Russo, VanDecar, et al., 2010). Volcanic products 58 associated with several Cenozoic slab windows can be found at subduction margins around the 59 60 Pacific Ocean, indicating that this phenomenon is widespread (McCrory et al., 2009).

61 The extent of the Patagonian slab window has previously been estimated based on plate 62 kinematic reconstructions (Breitsprecher & Thorkelson, 2009) and has been mapped using body 63 wave tomography in the immediate vicinity of the CTJ (Russo, VanDecar, et al., 2010). These two methods are in good agreement near the CTJ (Figure 1), but the full extent of the slab

65 window remains poorly defined. Reconciling tectonic reconstructions with observations such as

the locations of slab-edge adakitic volcanism requires invoking ridge jumps and changes in

67 spreading rates, which are poorly constrained due to the subduction of seafloor magnetic

anomaly records (Bourgois et al., 2016).

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Figure 1: Map of the study region in Patagonia. Seismic stations and volcanoes are marked.
Data sources are described in Section 2. Previously estimated slab window extents are shown for
Russo, VanDecar, et al. (2010) and for Breitsprecher and Thorkelson (2009). The present-day
NPI and SPI are shaded in grey (Davies et al., 2020). Background map colors show bathymetry
and elevation data (Ryan et al., 2009). Red lines show present-day plate boundaries. Tierra del
Fuego (TdF) and the Península de Taitao are labeled.

The co-location of the Patagonian slab window with the Northern and Southern 77 Patagonian Icefields (NPI and SPI) in the Austral Andes makes Patagonia an excellent place to 78 study the effects of lateral variations in Earth structure on glacial isostatic adjustment (GIA). 79 Low-viscosity mantle in the slab window beneath the icefields is expected to speed the GIA 80 response to ice mass changes on decadal to centennial timescales, and ongoing glacial unloading 81 is thought to drive extremely high uplift rates (up to 40 mm/yr) measured in the NPI and SPI 82 (Dietrich et al., 2010; Ivins & James, 2004; Klemann et al., 2007; Lange et al., 2014; Richter et 83 al., 2016). Improved constraints on the extent of the slab window and on lateral variations in the 84 viscosity structure of the mantle beneath the icefields are necessary for improving GIA models 85 and interpreting geodetic observations in terms of changing ice mass. More robust GIA models 86 provide stronger constraints on past icefield mass and climate change (Oerlemans, 2005). 87

88 While surface volcanism provides some information on slab window formation and 89 geometry, reliably reconstructing slab window mantle dynamics is a long-standing challenge in 90 geodynamics (Dickinson, 1981; S. Lin, 2014). Seismic tomography is an essential tool for onnecting surface features to subsurface structure. In this study, we use seismic data recently

collected by the GUANACO broadband seismic deployment to derive a new seismic velocity

model for Patagonia, map the full extent of the slab window, and investigate how associated

mantle dynamics have affected the overriding lithosphere. We show that the dynamics of the slab

95 window are responsible for inferred low viscosities in the upper mantle and the unusually rapid

96 glacial isostatic response to ice mass loss in the region.

97 2 Data and methods

We obtained a shear velocity (V_{sv}) model for Patagonia by jointly inverting Rayleigh
wave dispersion curves from ambient noise and earthquake tomography with P receiver
functions in a Bayesian framework that enables us to quantify velocity uncertainties statistically.
Data were from the GUANACO (Magnani et al., 2020), SEPA (Wiens et al., 1998), and CRSP
(Russo, VanDecar, et al., 2010) temporary seismic networks; and from the Chilean National
Seismic Network, the GEOSCOPE Network, the Antarctic Seismographic Argentinian Italian

104 Network, and ENAP (Empresa Nacional del Petróleo) monitoring stations.

105 2.1 Rayleigh wave tomography from ambient noise and earthquake records

106 We used ambient noise tomography to obtain isotropic Rayleigh wave phase velocities at 8 to 40 seconds period, and group velocities from 10 to 30 seconds (Bensen et al., 2007) (Figure 107 S1). Temporal normalization was done with the running average method, and we incorporated 108 time-frequency phase weighted stacking of the daily cross-correlation records (Schimmel et al., 109 2011; Schimmel & Gallart, 2007). We then obtained dispersion curves from the stacked cross-110 correlations using Automated Frequency Time Analysis (Bensen et al., 2007) and performed 111 tomography using the method of Barmin et al. (2001). Uncertainties were estimated using a 112 scaling relationship with ray path coverage, with group velocity uncertainties taken to be double 113 the phase velocity uncertainties (Barmin et al., 2001; Shen et al., 2016). 114

Surface wave tomography was performed using shallow events (<50 km depth) within 115 20-150° of the study region (Figure S2). A balanced azimuthal distribution was constructed by 116 starting with all events with M_w>6 and adding non-overlapping events down to M_w 5.4 at 117 118 undersampled azimuths. We performed visual quality control on all waveforms. Helmholtzcorrected phase velocity maps were calculated from 20 to 100 seconds period using the 119 120 Automated Surface-Wave Measurement System (ASWMS) (Jin & Gaherty, 2015; F.-C. Lin et al., 2009; F.-C. Lin & Ritzwoller, 2011) (Figure S3). The minimum inter-station distance for 121 calculating cross-correlations was set to 50 km, and the maximum distance was varied with 122 period such that it did not exceed ~4 wavelengths. The bandwidth range for the Gaussian filters 123 applied to the waveform cross-correlations was 0.04-0.07 Hz. Automated quality controls based 124 on the fraction of good measurements per event were intentionally relaxed for stations in the 125 CRSP temporary network to obtain adequate data coverage near the Chile Triple Junction, and 126 the larger velocity uncertainties near the triple junction at short periods reflect this choice. 127

Phase velocity dispersion curves from 8 to 100 seconds were constructed by combining the results from ambient noise and earthquake tomography (Figure S4). We used linear weighting across the overlapping periods, with ambient noise velocities weighted more at shorter periods and earthquake results at longer periods.

132 2.2 P receiver functions

We selected events 30-90° away from our study area with M_w>5.1 and a signal-to-noise 133 ratio greater than 3 on the vertical component for the receiver function (RF) calculation. The 134 seismograms were filtered from 0.33-1 Hz, and P first arrival picks were refined using STA/LTA 135 (Withers et al., 1998) in a time window around the predicted onset time from the global model 136 IASP91 (Kennett & Engdahl, 1991). P-to-s RFs were then calculated using the multitaper 137 deconvolution method (Helffrich, 2006; Park et al., 1987; Park & Levin, 2000; Shibutani et al., 138 2008) and corrected for moveout using IASP91. For each station, a composite RF with 139 uncertainties was obtained by taking the zeroth order component from harmonic decomposition 140 (Bianchi et al., 2010). If azimuthal coverage was not sufficient to fit harmonics, the station 141 average RF was used instead. 142

143 2.3 Bayesian inversion for velocity-depth models

We inverted for 1D velocity-depth models using a Markov chain Monte Carlo (MCMC) method (Shen et al., 2013). Each velocity-depth model was described by 14 parameters: layer thicknesses for sediments and crust; top and bottom velocities for the sediment layer; four cubic basis spline coefficients for crustal velocities; and six cubic basis spline coefficients for mantle velocities (Table S1). The total model depth was fixed at 300 km, with the lowermost 100 km of the model gradually converging to the global model AK135 (Kennett et al., 1995) since the data provide no constraints at these depths.

Initial prior distributions for the sediment and crustal layer thicknesses were set based on 151 RFs. We used bootstrap stacking of H-k stacks to estimate an initial Moho depth for each station 152 (Sandvol et al., 1998; Zhu & Kanamori, 2000). For sediment thickness, we performed a K-means 153 154 clustering analysis (Pedregosa et al., 2011) on the first six seconds of the RF stacks for the stations, and used the clusters to divide the stations into those overlying "thick" sediments and 155 "thin" sediments. The "thick" sediment cluster agreed well with the mapped extent of the 156 Austral-Magallanes Basin (Cuitiño et al., 2019). For stations overlying "thick" sediment, the 157 prior distribution for sediment thickness was set to 4 ± 4 km, and for stations with "thin" 158 sediment the thickness prior distribution was set to 1 ± 1 km. 159

Prior distributions for velocities in all layers were set to typical values with large search ranges to allow for variation (Table S1). The sixth mantle spline coefficient at the base of our model was fixed at 4.7 km/s based on AK135, since the data have almost no sensitivity at 300 km depth.

We imposed some velocity constraints to ensure that accepted models were physically 164 reasonable: all velocity parameters were less than 4.9 km/s, crustal velocities were less than 4.2 165 km/s, the velocity jump across the Moho was less than 0.7 km/s, and velocities were not allowed 166 to decrease with depth through the sediments and crust. Dispersion curves and receiver functions 167 were weighted equally in the joint misfit function after normalizing their respective uncertainties. 168 169 We used 15 chains of 5000 steps each for the MCMC calculation. The posterior distributions for the 14 model parameters were calculated from the set of accepted models based on the misfit 170 function (Shen et al., 2013) (Figure S5). 171

172 Inversions were first done for station locations. Prior distributions for the sediment and 173 crustal layer thicknesses were then adjusted in places where the inversion failed to find a well-174 fitted model. This was particularly important for stations within the Austral-Magallanes Basin,

where the thick sediments violated the H-k stacking assumption of a constant-velocity crustal 175

layer. We then inverted dispersion curves alone for velocity-depth model at grid points set at 0.3° 176

intervals throughout the study area, using smoothed maps of crustal and sediment thicknesses 177

from the station inversions to set layer thickness prior distributions. 178

The grid point and station results were combined by averaging together the 14 model 179 parameters for each grid point with those for any stations within a 50 km radius. Weights for the 180 station parameters were calculated based on proximity to the grid point. The model was 181 smoothed laterally at each depth using a Gaussian filter with a standard deviation equal to the 182

grid spacing. 183

2.4 Mantle viscosity calculation 184

Although there is no direct relationship between seismic velocity and mantle viscosity, 185 velocity is commonly used to indirectly estimate viscosity since both are largely controlled by 186 temperature. Mantle viscosities were estimated based on differences between our velocity model 187 and V_{sv} from the global 1D model STW105 (Kustowski et al., 2008). The seismic anomalies 188 189 were used to estimate temperature anomalies relative to a global average temperature model, which were then used along with experimentally-derived flow laws to estimate deviations from a 190 global 1D viscosity model (Ivins et al., 2021; Wu et al., 2013). We used rheologic parameters for 191 dry diffusion creep of olivine (Hirth & Kohlstedt, 2003; Karato, 2008), a reference mantle 192 viscosity from IJ05-R2 (Ivins et al., 2013), and temperature derivatives that included both 193 anharmonic and anelastic contributions (Karato, 2008). The calculated viscosities would not be 194 significantly different for wet diffusion creep given parameter uncertainties (Hirth & Kohlstedt, 195 2003). The reference global average temperatures were calculated by proportionally weighting 196 continental average geotherms (Stacey & Davis, 2008) and adiabatic temperature gradients 197 beneath oceanic regions, giving 1486 K at 100 km, and 1582 K at 150 km. We set the fraction of 198 199 the velocity anomaly attributed to temperature to 0.65, as found in a geodetic study of North America and Fennoscandia (Wu et al., 2013). While this temperature fraction may be different in 200 Patagonia compared to stable cratonic regions, such variation would not change the pattern of 201 relative viscosity differences across Patagonia. The remaining velocity anomaly is attributed to 202 compositional variations in the mantle. 203

3 Results 204

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3.1 Extent of the Patagonian slab window

Mantle velocities are low throughout the inferred slab window region, with a minimum 206 velocity less than 4.1 km/s at 50 km depth, ~8% slower than the global average given by 207 STW105 (Kustowski et al., 2008). The most intense portion of the shallow slow anomaly is north 208 of 49°S, in the youngest part of the slab window (Figure 2). North of 51°S, the western edge of 209 the anomaly at 100 km depth aligns with estimates of the extent of the subducting Antarctic slab 210 from plate kinematic reconstructions (Breitsprecher & Thorkelson, 2009) and the trend of 211 adakitic volcanism along the Austral Andes Volcanic Arc (Figure 2). Increased velocities north 212 of Tierra del Fuego delineate the southeastern extent of the slab window effects and are 213 214 consistent with xenolith studies suggesting the presence of a continental lithospheric block with thicker lithosphere (Schilling et al., 2017). 215

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Figure 2: Velocity-depth slices. Maps show V_{sv} in km/s: (a) 50 km depth, contour at 4.15 km/s (solid black line); (b) 100 km depth; (c) 150 km depth. Other lines and symbols are as in Figure 1. Velocity uncertainties are ~0.1 km/s across most of the maps; velocity uncertainty maps are shown in Figure S6.

3.2 Thermal erosion of the South American lithosphere

Velocities directly beneath the crust near the CTJ are much lower than expected for continental lithosphere (4.1 km/s at 50 km, compared to 4.5 km/s for STW105; Kustowski et al., 2008), indicating that the lithospheric mantle is missing in the youngest part of the slab window (Figure 2a). Vertical cross sections through the slab window show that anomalously slow mantle velocities are present immediately below the Moho, with thin (<10 km thick) patches of faster

227 mantle material at the Moho in places (Figure 3a, 45 to 47°S).

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Figure 3: Velocity cross sections. Cross sections show V_{sv} along (a) 73.5°W, (b) 47.5°S, and (c) 50°S, with transect locations shown on a map (d). Other lines and symbols on the map are as in Figure 1. Topography over each transect is plotted above the velocities (Ryan et al., 2009).

232 Dashed black lines mark the base of the sediments and the Moho on each cross section, and

labeled lines show isovelocity contours. The V_{sv} color scale is saturated at 3.6 km/s on the low end to emphasize slow anomalies in the mantle.

The absence of lithospheric mantle near the CTJ suggests that mantle dynamics 235 associated with the slab window have eroded the base of the plate. Thermal erosion of the 236 overriding plate is predicted by thermo-mechanical models of ridge subduction (Groome & 237 Thorkelson, 2009), and shallow slow velocity anomalies beneath the Antarctic Peninsula have 238 been similarly interpreted (Lloyd et al., 2020). Previous studies have inferred that the Patagonian 239 lithosphere has thinned above the slab window based on regional heat flow data (Ávila & Dávila, 240 2018), crustal thickness measurements (Robertson Maurice et al., 2003), and GIA models fit to 241 observed uplift rates (Lange et al., 2014; Richter et al., 2016). 242

The process of thermal erosion likely requires the presence of melt or fluids, as purely 243 conductive heating is too slow relative to the age of the slab window. In the absence of a present-244 day slab and slab-derived volatiles, melt may be supplied by decompression of upwelling 245 asthenospheric mantle. There are no known active Holocene volcanos over the slowest part of 246 the velocity anomaly, but the presence of the icefields complicates the mapping of volcanic 247 activity. Shear wave splitting analyses show a strong E–W fast direction near the CTJ, indicating 248 vigorous mantle flow around the edge of the Nazca slab (Russo, Gallego, et al., 2010; Russo, 249 VanDecar, et al., 2010; Wiens et al., 2021), and fast mantle flow in the shallow asthenosphere 250 may assist in the removal of lithospheric material. 251

Lithospheric erosion near the CTJ contrasts with the structure farther south, where fast 252 253 velocities indicate that the lithospheric mantle is largely intact beneath the Austral-Magallanes Basin. Patagonia has a complex history of terrane accretion (Ramos & Ghiglione, 2008), and it is 254 possible that the lithosphere in the south was thicker prior to the opening of the slab window. 255 Alternatively, thermal erosion may have been less efficient during the earliest stages of ridge 256 subduction, becoming more efficient over time as the slab window thermally perturbed the 257 surrounding mantle. As ridge subduction initiated only 18 Ma and subsequent ridge segments 258 entered the trench at ~12 Ma, 6 Ma, and 3 Ma, we expect that timescales of conductive cooling 259 are too short relative to the age of the slab window to allow for significant re-formation of 260 261 mantle lithosphere even in the oldest parts of the window (Boutonnet et al., 2010).

The crust thins by >10 km from north to south over the slab window (Figure 3, S7, S8). 262 This trend is opposite the lithospheric erosion seen at the Moho, and is unlikely to be entirely due 263 to surface erosion since the thinning is not primarily in the upper crust. Along the west coast of 264 North America, the passage of the migrating Mendocino Triple Junction is thought to have 265 caused rapid, temporary crustal thickening followed by crustal thinning (Furlong & Govers, 266 1999). The same mechanism may be at work in Patagonia, but the trend of mean relief along the 267 Austral Andes does not match predictions for flexural downwarping associated with this model 268 269 for crustal modification (Georgieva et al., 2016). Preexisting structure, overthrusting of terranes, thermal erosion, and tectonic extension have also been proposed to explain crustal structure near 270 the CTJ (Rodriguez & Russo, 2020), and further measurements extending north of our study 271 region would help clarify the source of the variations in crustal thickness. 272

273 4 Implications for glacial isostatic adjustment

Mantle viscosity structure strongly controls GIA, which in turn responds to spatial and
 temporal ice-mass variations. High geodetic uplift rates around the NPI and SPI (≥4 cm/yr) have

been attributed to anomalous mantle viscosities lower than $2x10^{18}$ Pa s (Ivins & James, 2004; 276 Lange et al., 2014; Richter et al., 2016), and recent GIA models suggest that reproducing 277 observed uplift rates requires either mantle viscosities that are significantly lower beneath the 278 279 NPI compared to the SPI or more ice mass loss in the NPI than previously estimated (Lange et al., 2014; Russo et al., 2021). The observed location of the slowest part of the seismic velocity 280 anomaly north of 49°S is consistent with the former explanation, and viscosities estimated from 281 our velocity model also point to both low overall viscosity in the slab window and a difference in 282 structure beneath the two icefields (Figure 4). Our estimated viscosities are mostly higher than 283 the values obtained by previous geodetic studies (Lange et al., 2014; Richter et al., 2016), but we 284 emphasize that the absolute viscosity values we obtain are highly sensitive to uncertain 285 parameters such as the fraction of velocity variation due to temperature. The extent of the slow 286 anomaly also suggests that gradients in uplift rates along the SPI may reflect latitudinal variation 287 in mantle viscosity. In broader terms, the strong lateral heterogeneity in mantle viscosity 288 indicated by the velocity model implies that the geodetic response to glacial unloading in 289 Patagonia will be highly three-dimensional, and cannot be fully described by symmetric 290 deformation predicted for a radially layered mantle, particularly with respect to the prediction of 291 292 horizontal crustal motions (Klemann et al., 2007). 293



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Figure 4: Estimated mantle viscosity based on the seismic velocity model. Maps of $log_{10}(\eta)$ at (a) 100 km and (b) 150 km depth. Other lines and symbols on the map are as described in the

captions for Figures 1 and 2.

The patterns of lithospheric erosion and thinning observed in our velocity model are also expected to affect GIA. While the seismic lithosphere is not exactly equivalent to the elastic lithosphere relevant to geodetic models, both terms refer to a layer of colder, more rigid material

that acts as a lowpass filter on the response of the mantle to changes in surface loading. The thin 301 lithosphere observed near the CTJ enables shorter wavelength signals from GIA loading and 302 mantle viscosity variations to be observed in surface deformation and topography. In the south, 303 the fast velocities to the east indicate that thicker continental lithospheric mantle beneath the 304 Austral-Magallanes Basin may constrain mantle flow patterns in the older parts of the slab 305 window by blocking shallow latitudinal flow (Klemann et al., 2007). In this scenario, horizontal 306 surface motions on the eastern side of the SPI are predicted to be dominantly to the east, and 307 GNSS observations of horizontal surface displacements support this prediction where data are 308 available along the eastern side of the northern SPI (Richter et al., 2016). 309

Quantifying ice mass changes on the Patagonian Icefields is crucial for projecting future 310 water resources and informing models of global sea level rise. Temperate mountain glaciers 311 including the NPI and SPI currently contribute substantially to global mean sea level rise (Jacob 312 et al., 2012; Radić & Hock, 2011). Ice and hydrological mass changes are efficiently monitored 313 by satellite gravimetry, provided the GIA gravity effect is accurately removed (Ivins et al., 314 2011). However, there are large trade-offs in GIA modeling between ice mass history and mantle 315 viscosity structure (Lange et al., 2014), so uncertainties in mantle viscosity propagate forward 316 into highly uncertain ice mass change rates (Richter et al., 2019). Constraining the regional 3D 317 viscosity structure provides the space gravimetry community with key information on the solid 318 319 Earth contribution to the secular mass change signal in Patagonia.

The geologically rapid response of Patagonian topography to ice mass changes makes the 320 321 area above the slab window ideal for studying the connections between the solid Earth rheology and dynamics, surface processes, and climate. The heterogeneous mantle structure indicated by 322 seismic velocities implies that the isostatic response to ice mass changes will also be 323 heterogeneous, and the absence of mantle lithosphere over the slab window promotes the 324 observation of shorter wavelength variations in surface deformation. Lateral heterogeneity in 325 mantle viscosity and lithospheric thickness, guided by seismic models and other geophysical 326 327 observations, must be incorporated into GIA modeling in this region. The resulting higher quality models will help advance our understanding of the history and the future of the 328 cryosphere and hydrosphere in Patagonia. 329

330 Acknowledgments

We thank Patrick Shore, Gerardo Connon, Leticia Duca, Celeste Bollini, Nora Sabbione, 331 Gerd Sielfeld, Daniel Valladares, and many others for their work in planning, deploying, 332 servicing, and recovering the GUANACO seismic array. We also thank Maeva Pourpoint for her 333 work in the field and with initial data processing; Weisen Shen for sharing his Bayesian 334 inversion code; and Zhengyang Zhou and Zongshan Li for their help with several of the software 335 336 packages used for data processing. We thank reviewers R. M. Russo and Volker Klemann for their thoughtful comments. The OCCAP Foundation assisted in arranging boat transport to 337 remote stations in Chile. ENAP Magallanes provided additional seismic data. The seismic 338 instruments were provided by the IRIS-PASSCAL Instrument Center at New Mexico Tech. The 339 facilities of the IRIS Consortium are supported by the National Science Foundation's SAGE 340 Award under Cooperative Support Agreement EAR-1851048. The GUANACO project is funded 341 342 by the National Science Foundation under grants EAR-1714154 to WUSTL and EAR-1714662 to SMU, and E. R. I. was supported by NASA under grant NNH19ZDA001N-GRACEFO. 343

344 **Open Research**

345 Data used in this study is from the GUANACO, SEPA, and CRSP temporary seismic

- networks (network codes: 1P, 10/2018-03/2021; XB, 02/1997-10/1998; YJ, 12/2004-12/2006),
- 347 permanent stations from the Chile Network, GEOSCOPE, and the Antarctic Seismographic
- Argentinian Italian Network (network codes: C, C1, G, AI), and stations operated and
- maintained by ENAP. Data for all except the ENAP stations can be obtained from the IRIS DMC
- 350 (<u>https://ds.iris.edu/ds/nodes/dmc</u>). The earthquake records and ambient noise cross-correlations
- for ENAP stations used in this study are publicly available (Mark et al., 2021a,
- https://doi.org/10.5281/zenodo.5508198). The final velocity and viscosity models presented in
- this paper are also available (Mark et al., 2021b, https://doi.org/10.5281/zenodo.5794167).
- 354 Publicly released versions of the codes used for analysis can be found at:
- 355 <u>https://github.com/NoiseCIEI/Seed2Cor, https://github.com/NoiseCIEI/AFTAN,</u>
- 356 <u>https://github.com/NoiseCIEI/RayTomo</u>, <u>https://github.com/trichter/rf</u>,
- 357 <u>https://github.com/jinwar/matgsdf</u>, and
- 358 <u>http://diapiro.ictja.csic.es/gt/mschi/SCIENCE/tseries.html#software</u>. These codes are currently
- only available on github and researchers' personal websites. Color palettes for all figures are
- from Fabio Crameri's ScientificColourMaps7 (<u>https://doi.org/10.5281/zenodo.1243862</u>).

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Supporting Information for

Lithospheric Erosion in the Patagonian Slab Window, and Implications for Glacial Isostasy

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Introduction

This file contains supplemental figures and a table detailing parameters used for Bayesian inversion.





Figure S1. Phase velocity maps from ambient noise tomography. The left column shows phase velocities in km/s, and the right column shows path density smoothed over the map. Phase velocities are shown at periods of **(a)** 8 seconds, **(b)** 14 seconds, **(c)** 20 seconds, **(d)** 26 seconds, **(e)** 32 seconds, and **(f)** 40 seconds. Velocities are only shown where path density is greater than 10 and the resolution length scale is greater than 0.02.



Figure S2. Event distribution map. Blue dots mark the locations of earthquakes used for tomography. The study region is outlined by a red box.





- 2

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Figure S3. Phase velocity maps from earthquake tomography. The left column shows Helmholtz-corrected phase velocities in km/s, the middle column shows one standard deviation (in km/s) of the velocities calculated from stacking maps for each earthquake, and the right column shows log₁₀ of the number of paths used in each grid cell. Phase velocities are shown at periods of (a) 20 seconds, (b) 26 seconds, (c) 32 seconds, (d) 40 seconds, (e) 50 seconds, (f) 60 seconds, (g) 80 seconds, and (h) 100 seconds. Velocities are only shown for grid cells where velocity measurements were obtained for at least 10 events.

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Figure S4. Comparison between earthquake and ambient noise tomography. The left column shows earthquake phase velocities as in Figure S3, the middle column shows ambient noise phase velocities as in Figure S1, and the right column shows the difference between the two velocity maps at (a) 20 seconds, and **(b)** 40 seconds. All velocities and velocity differences are in km/s. For the maps of velocity differences, magenta contours outline regions where the difference between the two velocity maps is greater than 0.1 km/s.



TRLG: -71.57, -49.48

Figure S5. Inversion results for station TRLG. (a) Map showing station locations (black triangles) and the location of station TRLG (red inverted triangle). **(b)** Input station receiver function (grey line and shading) plotted with the forward-calculated receiver function for the average of all accepted velocity models for this station (dark blue line). **(c)** Input group and phase velocity measurements (pink and orange errorbars) with forward-calculated dispersion curves for the average of all accepted velocity models for this station (navy lines). **(d)** Inversion results to 200 km depth, showing the model for the centroid of the parameter prior distributions (orange line), the average of all accepted models (dark blue line), one standard deviation of all accepted models (green lines), and the range spanned by all accepted models (grey shading). **(e)** Shallow structure of accepted velocity models, with the same lines as in (d). **(f)** Histogram showing the distribution of Moho depths for all accepted models.



Figure S6. Velocity uncertainty maps. Maps show one standard deviation for Vsv (in km/s) at **(a)** 50 km, **(b)** 100 km, and **(c)** 150 km depth. The uncertainty maps correspond to the velocity depth slices shown in Figure 2 of the main text. On all maps, previously estimated slab window extents are shown from Russo et al. (2010) at depths of 50 km (dot-dashed line) and 100 km (dotted line), and from Breitsprecher and Thorkelson (2009) (dashed line)^{1,17}. The present-day NPI and SPI are shaded in grey, and the locations of volcanoes are marked with black triangles (adakitic) and black circles (basaltic).



Figure S7. Sediment thickness and Moho depth from the velocity model. Maps of **(a)** sediment layer thickness and **(b)** Moho depth (sediments plus crust) show thick sediments in the Austral-Magallanes basin to the south, and a north-to-south decrease in crustal thickness across the study region. Features marked on the maps, including slab window outlines from previous studies, locations of volcanoes, and present-day icefield extents, are the same as in Figure 1. Uncertainty maps for interface depths are shown in Figure S8.



Figure S8. Sediment thickness and crustal thickness uncertainty maps. Maps show one standard deviation of **(a)** the sediment layer thickness, and **(b)** the crustal thickness, excluding sediments. Features marked on the maps, including slab window outlines from previous studies, locations of volcanoes, and present-day icefield extents, are the same as in Figure 1.

	Layer type	Parameter	Value	Perturbation range
Sediments	Gradient	Thickness (T _{seds})	1 or 4 km	± 1 or 4 km
		Top velocity	1.5 km/s	± 0.5 km/s
		Bottom velocity	2.2 or 2.8 km/s	± 0.6 km/s
Crust	Cubic splines	Thickness (T _{crust})	From H-k stacks	± 9 km
		1 st spline coeff.	3.2 km/s	± 25%
		2 nd spline coeff.	3.4 km/s	± 25%
		3 rd spline coeff.	3.6 km/s	± 25%
		4 th spline coeff.	3.9 km/s	± 25%
Mantle	Cubic splines	Thickness	300 km - T _{crust} - T _{seds}	
		1 st spline coeff.	4.2 km/s	± 25%
		2 nd spline coeff.	4.3 km/s	± 25%
		3 rd spline coeff.	4.4 km/s	± 25%
		4 th spline coeff.	4.5 km/s	± 25%
		5 th spline coeff.	4.6 km/s	± 25%
		6 th spline coeff.	4.7 km/s	Held constant

Table S1. Prior parameter distributions for Bayesian inversion.