Properties of Ion-Inertial Scale Plasmoids Observed by the Juno Spacecraft in the Jovian Magnetotail

Yash Sarkango¹, James A. Slavin¹, Xianzhe Jia¹, Gina A. DiBraccio², George B. Clark³, Weijie Sun⁴, Barry H. Mauk⁵, William S. Kurth⁶, and George B. Hospodarsky⁶

¹University of Michigan-Ann Arbor ²NASA GSFC ³Johns Hopkins University Applied Physics Laboratory ⁴University of Michigan - Ann Arbor ⁵Johns Hopkins University ⁶University of Iowa

November 24, 2022

Abstract

We expand on previous observations of magnetic reconnection in Jupiter's magnetosphere by constructing a survey of ioninertial scale plasmoids in the Jovian magnetotail. We developed an automated detection algorithm to identify reversals in the component and performed the minimum variance analysis for each identified plasmoid to characterize its helical structure. The magnetic field observations were complemented by data collected by the Juno Waves instrument, which is used to estimate the total electron density, and the JEDI energetic particle detectors. We identified 87 plasmoids with 'peak-to-peak' durations between 10 s and 300 s. 31 plasmoids possessed a core field and were classified as flux-ropes. The other 56 plasmoids had minimum field strength at their centers and were termed O-lines. Out of the 87 plasmoids, 58 had in situ signatures shorter than 60 s, despite the algorithm's upper limit to be 300 s, suggesting that smaller plasmoids with shorter durations were more likely to be detected by Juno. We estimate the diameter of these plasmoids assuming a circular cross-section and a travel speed equal to the Alfven speed in the surrounding lobes. Using the electron density inferred by Waves, we contend that these plasmoid diameters were within an order of the local ion-inertial length. Our results demonstrate that magnetic reconnection in the Jovian magnetotail occurs at ion scales like in other space environments. We show that ion-scale plasmoids would need to be released every 0.1 s or less to match the canonical 1 ton/s rate of plasma production due to Io.

| 1 | Properties of Ion-Inertial Scale Plasmoids Observed by the Juno Spacecraft in the Iovian Magnetotail |
|----------|--|
| 2 | Jovian Magnetotan |
| 3 | |
| 4 | Enter authors here: |
| 5 6 | Yash Sarkango ¹ , James A. Slavin ¹ , Xianzhe Jia ¹ , Gina A. DiBraccio ² , George B. Clark ³ , Weijie Sun ¹ , Barry H. Mauk ³ , William S. Kurth ⁴ , and George B. Hospodarsky ⁴ |
| 7 | ¹ University of Michigan - Ann Arbor, USA. |
| 8 | ² NASA Goddard Space Flight Center, USA. |
| 9 | ³ John Hopkins University – Applied Physics Laboratory, USA |
| 10 | ⁴ University of Iowa, USA |
| 11 | |
| 12 | Corresponding author: Yash Sarkango (<u>sarkango@umich.edu</u>) |
| 13 | |
| 14 | Key Points: |
| 15 16 | • We identify and analyze 87 ion-inertial scale plasmoids (56 O-lines, 31 flux-ropes) in the Jovian magnetotail using an automated algorithm. |
| 17 18 | • North-South field reversals with peak-to-peak durations less than 60 s are more common than those with durations between 60 and 300 s. |
| 19 20 | Ion-inertial scale plasmoids alone cannot account for the >500 kg/s loss-rate deficit unless they are being produced every ~0.1 s or less. |

21 Abstract

We expand on previous observations of magnetic reconnection in Jupiter's magnetosphere by 22 constructing a survey of ion-inertial scale plasmoids in the Jovian magnetotail. We developed an 23 automated detection algorithm to identify reversals in the B_{θ} component and performed the 24 minimum variance analysis for each identified plasmoid to characterize its helical structure. The 25 magnetic field observations were complemented by data collected by the Juno Waves instrument, 26 which is used to estimate the total electron density, and the JEDI energetic particle detectors. We 27 identified 87 plasmoids with 'peak-to-peak' durations between 10 s and 300 s. 31 plasmoids 28 possessed a core field and were classified as flux-ropes. The other 56 plasmoids had minimum 29 field strength at their centers and were termed O-lines. Out of the 87 plasmoids, 58 had in situ 30 signatures shorter than 60 s, despite the algorithm's upper limit to be 300 s, suggesting that smaller 31 plasmoids with shorter durations were more likely to be detected by Juno. We estimate the 32 33 diameter of these plasmoids assuming a circular cross-section and a travel speed equal to the Alfven speed in the surrounding lobes. Using the electron density inferred by Waves, we contend 34 that these plasmoid diameters were within an order of the local ion-inertial length. Our results 35 demonstrate that magnetic reconnection in the Jovian magnetotail occurs at ion scales like in other 36 space environments. We show that ion-scale plasmoids would need to be released every 0.1 s or 37 less to match the canonical 1 ton/s rate of plasma production due to Io. 38

39 **1. Introduction**

Magnetic reconnection can be the primary mechanism through which the plasma created in Jupiter's inner magnetosphere from Io and Europa is ultimately lost to the external solar wind (Vasyliunas, 1983; Krupp et al., 2004). Many in situ observations support this hypothesis through different particle and field signatures, which are briefly summarized in the following paragraphs.

44 Recurring bursts of energetic particles occur in the Jovian magnetotail with flow velocities deviating from the corotation direction (Krimigis et al., 1981; Krupp et al., 1998; Woch et al., 45 2002; Kronberg et al., 2005, 2007; Kronberg, Woch, Krupp, & Lagg, 2008; Kasahara et al., 2013). 46 These flow bursts repeat on timescales between 1 to 4 days and can be directed either inward 47 (sunward) or outward (anti-sunward). More inward flow bursts are seen at radial locations closer 48 to Jupiter, whereas outward bursts are common farther away from the planet (Woch et al., 2002). 49 50 They are associated with an increase in the energetic particle fluxes and a decrease in the ion energy spectral index γ , or hardening of the ion energy spectra (Krupp et al., 1998; Woch et al., 51 1998). The peak energy of the ions also increases, which suggests acceleration associated with 52 these events (Woch et al., 1999). 53

54 Simultaneous magnetic field observations have shown that flow bursts occur during periods of reversals in the north-south component of the magnetic field (Nishida, 1983; Russell et 55 al., 1998; Woch et al., 1999; Kronberg et al., 2005; Vogt et al., 2010, 2014, 2020). Under typical 56 quiet conditions, the magnetic field in the Jovian magnetotail points predominantly southward at 57 58 the magnetic equator. Abrupt north-south reversals indicate a change in the topology through magnetic reconnection. Persistent northward magnetic fields may indicate open magnetic flux that 59 is 'disconnected' from Jupiter (Vogt et al., 2014). Multiple north-south reversals can be seen within 60 a single reconfiguration event lasting over a period of several days (Kronberg et al., 2007). The 61 sense of the reversal, i.e., from north to south or vice-versa, provides more information about the 62 placement of the measuring spacecraft with respect to the reconnection X-line, which can be 63 identified based on the meridional component for a given interval (Ge et al., 2010; Vogt et al., 64

65 2010). Some magnetic field reversals appear to have a helical or loop-like magnetic field structure 66 that is characteristic of plasmoid events (Vogt et al., 2014). Also seen are magnetic signatures of 67 'dipolarizations', which are the result of plasma compression due to fast planetward flows in the 68 so-called 'exhausts' emanating from the reconnection sites (Artemyev et al., 2013, 2020; Yao et 69 al., 2020).

The simultaneous occurrences of magnetic field reversals and flow bursts repeating on 70 similar timescales suggests a common origin. That such events recur over long timescales of 1-4 71 days, with no clear dependence on solar wind triggers, also suggests an internally driven 72 phenomenon. The dominant field component in the lobes parallel to the current sheet (B_r) 73 gradually increases between two consecutive active periods and decreases after the onset of the 74 events, while the opposite is true for the component normal to the current sheet (B_{θ}) at the equator. 75 This implies a gradual 'stretching' of the magnetodisc over the 1 to 4-day period (Ge et al., 2007; 76 Kronberg et al., 2007). Based on the above points, it is believed that the magnetosphere 77 experiences two states – a state of 'loading', characterized by an increase in magnetic stresses in 78 the magnetotail, and a state of 'unloading', when magnetic stresses and plasma are released from 79 80 the magnetosphere via magnetic reconnection and plasmoid release.

81 Based on these observations, which are also seen in Saturn's magnetosphere (Jackman et 82 al., 2011; Garton et al., 2021), it is now generally believed that reconnection occurs in the Jovian magnetotail and produces plasmoids. But the question remains if plasmoids can account for the 1 83 ton/s mass addition rate produced in the inner magnetosphere (Bagenal & Delamere, 2011). 84 85 Previous estimates vary on the size of and mass carried by Jovian plasmoids. Kronberg et al., (2008) assume plasmoid down-tail length pf 9 R_I , thickness of 2 R_I and azimuthal length of 200 86 R_I (1 R_J = 71492 km is the equatorial radius of Jupiter at 1 bar), and an oxygen ion density of 87 0.025 cm⁻³ and find that ~50 plasmoids should be released over the unloading period (~1 day) to 88 balance a 250 kg/s production, with each plasmoid contributing ~800 tons. Bagenal, (2007) 89 assumes a plasmoid diameter of 25 R_J, width of 2 R_J and density of 0.01 cm⁻³ and estimate the 90 mass of each plasmoid to be ~500 tons, which effectively translates to a loss rate of ~150 kg/s if 91 such plasmoids are released on an hourly basis. Vogt et al., (2014) consider larger plasmoid 92 dimensions with the higher estimate of the down-tail length to be 20 R_I , a width of 6 R_I , a cross-93 tail length of 70 R_I , and density of 0.01 cm⁻³ and calculate an upper estimate of the net loss rate to 94 be ~120 kg/s based on five such plasmoids released over one day. Similarly, Cowley et al., (2015) 95 find that the rate of 1 ton/s can only be achieved after including the large post-plasmoid plasma 96 sheet (PPPS), which may exist for the ~15 h recurrence time between consecutive reconnection 97 98 events and would increase effective plasmoid down-tail lengths up to $\sim 150 \text{ R}_{\text{J}}$. Hence, isolated, and infrequent plasmoids, which recur on a timescale of several hours or days, cannot match the 99 100 contribution due to the Galilean satellites without the inclusion of the PPPS. It can also be noted from the widely varying numbers shown above that the dimensions for Jovian plasmoids are not 101 well constrained because of the inherent limitations of single point measurements. Meanwhile, 102 other theories have also been proposed to explain the observed deficit which allow for mass loss 103 through other means e.g., through boundary interactions at the magnetopause (Delamere & 104 Bagenal, 2013; Masters, 2017) or through diffusive processes occurring at smaller scales than what 105 has been detected in the past (Kivelson & Southwood, 2005). 106

107 Observations of magnetotail reconnection noted above have analogs in the terrestrial 108 magnetosphere with the primary difference being that the driving mechanism to stress the 109 magnetotail in the terrestrial case is the external solar wind and interplanetary magnetic field

(Kronberg, Woch, Krupp, Lagg, et al., 2008). Our understanding of magnetic reconnection and 110 plasmoids has improved with multi-spacecraft observations in the terrestrial magnetosphere, high 111 cadence instrumentation and kinetic simulations. Plasmoids in the terrestrial magnetosphere and 112 other regions of the space environment are often accompanied with a strong core field within the 113 helical magnetic structure and are called magnetic flux-ropes (Slavin et al., 2003). The magnetic 114 pressure of the core region balances the magnetic tension force exerted by the outer regions and in 115 some cases, plasma pressure gradients are unnecessary to maintain this quasi-equilibrium. Flux-116 ropes in which pressure gradients are negligible and where the magnetic forces are self-balancing 117 are referred to as 'force-free'. It has been argued that the force-free configuration contains 118 minimum magnetic energy for helical structures (Taylor, 1974; Priest, 2011) toward which they 119 tend to evolve with time. Hence, knowledge about a particular flux-rope event's magnetic structure 120 could be used to determine its stage of evolution. Simultaneous energetic particle observations 121 have shown that flux-ropes which are produced on the ion-inertial scale can interact with or trap 122 electrons and ions, which get accelerated due to adiabatic processes such as due to the conservation 123 of the adiabatic invariants, or through non-adiabatic processes such as electromagnetic turbulence 124 (Grigorenko et al., 2015; Kronberg et al., 2019). Evidence exists for both Fermi and betatron 125 acceleration, which manifest as increases in the electron fluxes in the parallel and perpendicular 126 directions, respectively (Zhong et al., 2020; Vaivads et al., 2021). Similar results are found in 127 particle-in-cell simulations (Drake et al., 2006). 128

Plasmoids observed in the Jovian magnetotail have predominantly contained minimum 129 magnetic field strength in their interiors and are not force-free flux-ropes (Vogt et al., 2014). This 130 could be a result of the very large plasma β in the Jovian magnetotail, which produces plasmoids 131 containing dense plasma and large pressure gradients. The spatiotemporal scales over which these 132 plasmoids, which are also referred to as O-lines, evolve and possibly convert to force-free flux-133 ropes are also not known. It is not clear which parameter determines the direction and strength of 134 the core field for a plasmoid released in the magnetotail. For externally driven magnetospheres 135 such as Earth's and Mercury's, studies have found both strong (Moldwin & Hughes, 1992; Slavin 136 et al., 2003) or weak correlation (Smith, Slavin, Jackman, Poh, et al., 2017) between the IMF 137 orientation and the direction of the core field. In Jupiter's magnetosphere, the solar wind influence 138 and penetration of the IMF B_{y} into the plasma sheet is minimal. However, the 'bend-back' of the 139 magnetic field in the magnetodisc introduces a cross-tail magnetic field component in the mid-140 latitude regions (Khurana et al., 2004). 141

To understand the role of magnetic reconnection in facilitating mass loss, it is also 142 143 important to consider alternative theories such as smaller-scale reconnection in the Jovian magnetotail. However, plasmoids observed in the Jovian magnetotail based on the Galileo 144 measurements so far have been large, which could be due to the low temporal cadence of the 145 available instrumentation onboard the Galileo spacecraft. Kronberg, Woch, Krupp, & Lagg, (2008) 146 and Vogt et al., (2014) used data collected by Galileo to study the properties of tailward moving 147 Jovian plasmoids and found their average diameters to between ~9 to 10 RJ and between 2.6 to 20 148 R_J, respectively. Vogt et al., (2014) inferred the plasmoid size based on the time difference between 149 the two extrema in B_{θ} during the north-south reversal, whereas Kronberg, Woch, Krupp, & Lagg, 150 (2008) also included the period during which the magnetic field gradually returned to the 151 southward orientation, i.e., the post-plasmoid plasma sheet (PPPS). Kronberg, Woch, Krupp, & 152 Lagg, (2008) observed plasmoid durations to vary between 0 and 50 minutes, with most events 153 having durations between 10 and 20 minutes. They also calculated the plasma flow speeds during 154 the plasmoid events to be between 200 and 1200 km/s. Most plasmoids were associated with flows 155

of around 400 km/s. Similarly, in the Vogt et al., (2014) survey, the 2 to 20 R_J diameters correspond to an average in-situ duration of 6.8 minutes. The recent survey by (Vogt et al., 2020) used data collected by the Juno spacecraft to identify signatures of magnetic reconnection and found a similar result as for the Galileo data.

In a previous work, we reported on two Juno-based observations of flux-ropes in the Jovian 160 magnetotail with diameters comparable to or less than the local ion-inertial length (Sarkango et 161 al., 2021). This was made possible by the higher resolution magnetometer instrument aboard Juno, 162 as well as the presence of heavy ions in the magnetosphere which increased plasma length scales. 163 Our observations extended previous work on magnetic reconnection in the Jovian magnetosphere, 164 as most plasmoids observed previously by Galileo or Juno were found to have large diameters, 165 corresponding to in situ signatures lasting several minutes or longer. In contrast, the two events 166 discussed in Sarkango et al., (2021) had durations of 22 s and 62 s respectively. Under the 167 assumption that these plasmoids traveled at the Alfven speed corresponding to that in the lobes, 168 we had calculated their diameters to be ~11,000 and ~30,000 km, respectively. Based on these 169 observations, we hypothesized that magnetic reconnection in the Jovian magnetotail, like other 170 regions in the space environment, proceeds via the tearing instability in the magnetotail current 171 sheet. This had also been proposed by Kronberg et al., (2007), who demonstrated that the 172 magnetotail loading process would take roughly 2 days or more to create conditions conducive for 173 174 this instability to occur, which would lead to unloading.

In this work, we extend on previous works on plasmoids at Jupiter by conducting a survey 175 176 of all possible Jovian plasmoids with in situ signatures shorter than 5 minutes and corresponding diameters less than ~2 R_I . This is achieved by using an automated detection algorithm to identify 177 plasmoids using transient reversals in B_{θ} with 'peak-to-peak' signatures shorter than 5 minutes 178 and longer than 10 seconds. By using an empirical density profile and the local Alfven speed in 179 the surrounding lobes measured by Juno, we show that the events identified have diameters within 180 an order of magnitude of the local oxygen ion-inertial length. We also classify plasmoids based on 181 whether their core-region has magnetic fields that are stronger (flux-ropes) or weaker (O-lines) 182 than the surrounding magnetic fields in the outer layers of the plasmoid. A force-free flux rope 183 model is fitted to each flux-rope event, and it is found that out of the 31 plasmoids with strong core 184 fields, 6 events fit the force-free model well, i.e., they are self-balanced due to internal magnetic 185 stresses. However, 56 magnetic O-line-type plasmoids were identified and hence were more 186 187 commonly observed than flux-ropes. We also use data from the JEDI instrument to show properties of the energetic electrons and ions during two example plasmoid events. The fluxes of electrons 188 and ions are larger within these intervals and in the post-plasmoid plasma sheet. For one example, 189 the electron pitch-angles were isotropic during the interval but field-aligned before and after, 190 191 which could be due to betatron acceleration, as has been observed for electron distributions in the terrestrial magnetotail. Our results highlight that reconnection occurs in Jupiter's magnetosphere 192 193 over a wide range of scales and can accelerate plasma in the process. The frequent observations of plasmoids with small diameters, which are presumably easier to miss, by a single spacecraft in the 194 magnetotail also raises questions about the occurrence of ion scale magnetic reconnection in other 195 regions of the magnetosphere. Nevertheless, our estimates of the total mass carried by ion-inertial 196 197 scale plasmoids suggest that they do not directly contribute in a substantial manner to the loss of mass from the magnetosphere. 198

199 **2. Data and Methodology**

200 **2.1. Juno's Trajectory**



201

Figure 1. Juno's trajectory in the JSS coordinate system as seen in the equatorial and meridional projections. Dates and positions corresponding to every fifth apogee are highlighted. Also shown in panel (a) are the 75% percentile bow-shock (BS) and magnetopause (MP) model by Joy et al., (2002).

The Juno spacecraft was inserted into an elliptical orbit around Jupiter in June 2016 at around 06 LT (local time), or close to the dawn meridian. Each subsequent perijove pass was separated by a time of roughly 53 days. Over the years, Juno naturally precessed towards the nightside magnetotail, reaching 00 LT (midnight) in early 2020. Simultaneously, its apogee moved from near-equatorial to mid-latitudes in the southern hemisphere. Over the course of the highly

elliptical trajectory (shown in Figure 1), Juno spent a considerable amount of time in the central 211 plasma sheet, especially during the planet-bound portion of its orbit. During these periods, it 212 frequently crossed the oscillating magnetotail current sheet, which can be seen in the periodic 213 reversals of the radial component of the magnetic field (B_r) every 5 hours, or twice every Jovian 214 rotation period. Juno was most likely to encounter the magnetodisc current sheet at different radial 215 locations for different local times in the dawnside magnetotail. Initially (e.g., for years 2016-2018), 216 when the orbit was less inclined, the current sheet crossings were observed over a broad range of 217 radial distances, ranging from roughly 30 to 80 RJ. However, the increase in orbital inclination 218 during the later years (e.g., 2020) meant that current sheet crossings near midnight (00 LT) could 219 only be observed when the spacecraft was located at lower latitudes where the hinged oscillating 220 221 current sheet was expected to occur, i.e., at radial distances nearer to the planet and in the middle magnetosphere between roughly 20 and 50 RJ. This also had direct implications for the detection 222 of plasmoid signatures, as magnetic reconnection is also expected to occur close to the current 223 224 sheet location, which was not sampled uniformly by Juno.

225 **2.2. Data Description**

In this work, we used 1-s resolution vector magnetic field intensity data collected in situ 226 227 by Juno's onboard fluxgate magnetometers (Connerney et al., 2017). We used the magnetometer data to identify plasmoid signatures on the order of 10 s to 300 s, for which the 1 s cadence was 228 reasonable as it provides at least 10 magnetic field vectors per event. We also used data from the 229 230 Juno Waves instrument (Kurth et al., 2017), which measured the fluctuations in the electric field between frequencies of 50 Hz and 40 MHz, to identify the low-frequency cutoff for the continuum 231 radiation and estimate the local electron density (Gurnett et al., 1981; Barnhart et al., 2009). We 232 233 supported the fields observations using data from the JEDI energetic particle detectors (Mauk et al., 2013). Three JEDI instruments were located on the Juno spacecraft; two having a field-of-view 234 in the spacecraft equatorial plane (JEDI-90 and JEDI-270) and one looking perpendicular to the 235 spacecraft equatorial plane (JEDI-180). JEDI could measure electrons in the energy range between 236 18 keV to 740 keV. JEDI could also measure and distinguish between ion species based on the 237 time-of-flight channels, specifically protons (~37 keV - 2 MeV), oxygen (~130 keV - 10 MeV) 238 239 and sulfur (~130 keV - 11 MeV) ions. In favorable conditions, the rotation of Juno about its spin axis allowed for near-complete pitch-angle coverage with a typical collection time of 30 s. 240 However, for most times when Juno was in the middle and outer magnetosphere, the JEDI 241 instruments were operating at a lower data rate mode with reduced energy resolution to facilitate 242 data transfer during these periods. 243

244

2.3. Magnetic signatures of plasmoids

We used the Juno data to search for plasmoids in the Jovian magnetotail by identifying reversals in the north-south component of the magnetic field, i.e., B_z or B_θ in Cartesian JSS or spherical JSS coordinate systems respectively. An illustration of the expected magnetic signature of a tailward moving plasmoid containing helical magnetic fields is illustrated in Figure 2 (a).

The sense of the B_z reversal, i.e., from north-to-south or vice-versa, can be used to infer the direction of travel of the plasmoid. Observations in the terrestrial magnetosphere have shown that plasmoids which travel tailward usually correspond to a B_z reversal from positive to negative values (Slavin et al., 2003). While the Earth's magnetic moment points predominantly southward, Jupiter's internal magnetic moment points northward. This implies that tailward moving plasmoids

in the Jovian magnetotail would have the opposite sense of reversal as seen at Earth, i.e., B_z would 254

change from negative to positive values (assuming that a single plasmoid is released). On the other 255 hand, planetward moving plasmoids would result in B_z changing from positive to negative values.

256



257

Figure 2. (a) Magnetic signature and representative geometry of a tailward moving plasmoid 258 with a post-plasmoid plasma sheet. (b-c) Schematics showing the different magnetic field and 259 thermal pressure profiles of magnetic O-lines (b) and force-free flux-ropes (c) with circular 260 cross-sections. 261

262 To prevent ambiguity due to the oscillating Jovian current sheet and magnetodisc, analysis of magnetic field components in the Jovian system was conducted in the spherical JSS coordinate 263 system where B_r , B_{θ} and B_{ϕ} represented the radial, co-latitudinal and azimuthal components of 264 the magnetic field. In this system, the periodic motion of the current sheet is largely limited to the 265 radial and azimuthal components (B_r and B_{ϕ}), which are anti-correlated due to the bend-back 266 phenomena (Khurana et al., 2004). In a quiet time magnetotail and in the absence of reconnection, 267 B_{θ} is predominantly positive (In our work, like the previous studies, positive B_{θ} corresponds to 268

negative B_z at the equator). Hence, the magnetic signature of a tailward moving plasmoid would 269 be a positive-to-negative reversal in B_{θ} , and opposite (negative-to-positive) for a planetward 270 moving plasmoid (Figure 2). Additionally, plasmoids may or may not possess a core field inside 271 the outer helical magnetic structures (flux-ropes), which can be seen predominantly in the 272 azimuthal or radial components $(B_{\phi} \text{ or } B_r)$ or as a localized increase in the magnetic field strength 273 within the interval corresponding to the reversal of B_{θ} (Figure 2 (c)). 274

2.4. Minimum Variance Analysis 275

We used the magnetic field based minimum variance analysis (referred to as MVA or 276 277 BMVA in the literature) to characterize the helical magnetic structure. A local cartesian coordinate system was determined, whose orthogonal directions represent the directions of maximum, 278 intermediate and minimum variance of the magnetic field during the plasmoid interval. This was 279 done by first constructing the variance matrix M according to the following equation (Sonnerup & 280 Cahill, 1967), 281

 $M_{ii} = \langle B_i B_i \rangle - \langle B_i \rangle \langle B_i \rangle$, for $i, j \in \{x, y, z\}$ (1)

284 The eigenvectors of the variance matrix corresponding to the decreasing magnitude of the eigenvalues $(\lambda_L, \lambda_M, \lambda_N)$ provided the directions of maximum $(\vec{x_L})$, intermediate $(\vec{x_M})$ and 285 minimum variance $(\overrightarrow{x_N})$. 286

287 For plasmoid signatures in the magnetotail, the magnetic field varies most in the northsouth (Z or θ) direction. The maximally varying component (B_I) also reverses sign. In the case of 288 a flux-rope, the local increase in the core field direction can be seen in the component of 289 intermediate variance (B_M) . This leads to a near elliptical path when visualized as a hodogram of 290 the B_L and B_M components, also referred to as a rotation (Slavin et al., 1989). 291

MVA fails to determine the orientation of the variance directions if two or more 292 eigenvalues of the variance matrix are similar in magnitude. In other words, there are times when 293 the variance coordinate system is *degenerate*. This is verified in the present work by requiring that 294 the ratio of larger to smaller eigenvalues be greater than 3. Additionally, we also imposed the 295 condition by Rosa Oliveira et al., (2020) using the metric P (shown below), where P > P296 4.5 considered to be sufficient to validate the MVA eigensystem. 297

298
$$P = \frac{100}{\lambda_L^{1.5}} \times (\sqrt{\lambda_L} - \sqrt{\lambda_M})(\sqrt{\lambda_L} - \sqrt{\lambda_N})(\sqrt{\lambda_M} - \sqrt{\lambda_N})$$
299 (2)

299

2.5. Force-free Flux-rope Modeling 300

Plasmoids with strong axial core fields are referred to as 'flux-ropes'. A subset of flux-301 ropes within which pressure gradients are negligible and which are in force equilibrium due to the 302 self-balancing magnetic forces are termed 'force-free' (Kivelson & Khurana, 1995). The expected 303 pressure and magnetic field profile within O-line type plasmoids and magnetic flux ropes is shown 304 in Figure 2. In the O-line type plasmoids, the magnetic pressure of the helical wraps is balanced 305 by the thermal pressure gradient. One solution for axially symmetric force-free flux-ropes with 306 circular cross-sections takes the following form (Lepping et al., 1990; Slavin et al., 2003), 307

(3)

(4)

$$B_A = B_0 J_0(\alpha r)$$

$$B_T = B_0 H J_1(\alpha r)$$

311

Where B_A and B_T are the axial and tangential components of the magnetic field, H is the 312 handedness (either 1 or -1), J_0 and J_1 are the Bessel functions of the first kind, α is a constant 313 parameter and r here refers to the ratio between the impact parameter, which is the distance from 314 the center of the flux-rope at closest approach, and the radius of the flux-rope. In this work, $\alpha =$ 315 2.4048 was chosen as it results in the tangential field component being zero when r = 1 (at the 316 flux-rope edge). 317

For a given interval exhibiting a flux-rope-like signature, we used Eq. 3-4 to fit a constant-318 α flux-rope to the observations to determine whether the helical structures seen in the data were 319 force-free. This was achieved by varying B_0 , r, H and the spherical angles providing the axial 320 orientation of the flux-rope, θ_A and ϕ_A , such that the Chi-squared error between the observations 321 322 and the force-free model was minimized (Lepping et al., 1990).

323
324

$$\chi_r^2 = \frac{1}{N} \sum_{i=1}^{N} \left[\left(B_x - B_{x,m} \right)^2 + \left(B_y - B_{y,m} \right)^2 + \left(B_z - B_{z,m} \right)^2 \right]$$
(5)

524

This was achieved in two steps using an open-source nonlinear least-squares fitting 325 package (Newville et al., 2016). Firstly, the magnetic field components were normalized by the 326 local magnetic field strength within the interval. Next, the flux-rope model was fitted to the 327 observations for all parameters except for B_0 . The initial values for θ_A and ϕ_A , i.e., the flux-rope 328 orientation, were chosen based on the eigenvector for intermediate variance $(\vec{x_M})$ as provided by 329 330 the minimum-variance analysis. Then the minimization was repeated, keeping all parameters fixed to their optimized values, but varying B_0 to scale the modeled flux-rope's core field. 331

2.6. Automated detection of plasmoids 332

We used an automated algorithm to detect possible plasmoid signatures in the magnetic 333 field observations made by the Juno magnetometer. For identifying potential signatures, we used 334 a method like that used by Smith, Slavin, Jackman, Fear, et al., (2017) for the Kronian magnetotail 335 and by Vogt et al., (2014) for the Jovian magnetotail. Firstly, within an interval in the 1-s resolution 336 magnetometer data, all times corresponding to a reversal in B_{θ} , either from positive-to-negative or 337 vice-versa, were identified. Reversals which occurred beyond 05 LT on the dawnside and beyond 338 339 90 RJ in radial distance from the planet were discarded to prevent contamination due to proximity to the magnetopause, where the magnetic field is highly variable. 340

341 Next, for each reversal, the times corresponding to the extrema in B_{θ} for the event were identified. As there can be multiple local maxima and minima in the B_{θ} magnetic field 342 observations, we adopted the method used by Smith, Slavin, Jackman, Fear, et al., (2017) and 343 identified maxima-minima (or vice-versa) pairs within a period of +/- 10 min from the B_{θ} reversal. 344 Pairs in which the peak-to-peak B_{θ} extrema were less than 2 nT or the standard deviation of the 345 B_{θ} component during a 100 min interval centered on the reversal in question were considered 346

inconclusive and discarded. Pairs in which the excursion into negative B_{θ} values was a factor of 6 347 smaller than that in the positive direction (and vice-versa) were also discarded. Additionally, only 348 those extrema pairs were considered whose start and end times were separated by a duration of at 349 least 10 s and at most 5 min, which form the lower and upper limit for the events identified in this 350 study. The lower limit was chosen such that there are at least 10 vector measurements for 351 subsequent analysis. As the purpose of this work was to study small-scale plasmoid events, we 352 limited the algorithm to signatures lasting less than 5 min. The range in the present work overlaps 353 with the lower bins of a previous survey by Vogt et al., (2014) of plasmoids identified by the 354 Galileo magnetometer. 355

For each of the remaining pairs of extrema, which correspond to potential start and stop times for a plasmoid event, a linear function was fitted to the B_{θ} observations (Smith, Slavin, Jackman, Fear, et al., 2017). Those extrema pairs which showed low degree of correlation with the observations (quantified by the coefficient of determination $r^2 < 0.85$) were discarded.

Next, additional filters were applied based on the minimum-variance analysis, which was 360 conducted on all remaining extrema pairs. Events in which the eigenvector corresponding to the 361 direction of maximum variance did not have a predominantly Z_{ISS} component ($\hat{x}_L \cdot \hat{z} < 0.8$) were 362 discarded. Pairs for which the ratio of the maximum to intermediate and intermediate to minimum 363 eigenvalues were less than 3, or for which the P value was less than 4.5, were also removed from 364 consideration. Lastly, to only capture plasmoid events close to a magnetotail current sheet crossing 365 and prevent identification of traveling compression regions (TCRs) in the magnetotail lobes 366 (which was outside the scope of the present work), additional filters were applied to limit the 367 detection to events in which the minimally varying component of the magnetic field (B_N) as well 368 as the radial component in the JSS spherical coordinate system (B_r) , were less than 2 nT. 369

Of the remaining extrema pairs, the pair which fits best the B_{θ} observations using a linear function was identified (i.e., highest r^2) and chosen to be the start and end time for that event. In this procedure to detect B_{θ} reversals no distinction was made between positive-to-negative or negative-to-positive sense, as both tailward and planetward moving plasmoids (respectively) are likely to occur in the magnetotail. The conservative approach used by this algorithm ensured good candidates for plasmoids in the Jovian magnetotail. The procedure was repeated for each B_{θ} reversals detected by Juno, though only 87 reversals passed all criteria.

377 3. Results

378

3.1. Case Studies: Magnetic field and energetic particle signatures of plasmoids

In this section we discuss the magnetic field and energetic particle signatures of two examples of ion-inertial scale plasmoids identified by our automated algorithm.

381 3.1.1. Example 1: Flux-rope - DOY 76, 2017

In Figure 3 (a)-(d) we show the magnetic field data in the spherical JSS coordinate system for a plasmoid event occurring on DOY 75 in 2017 roughly between 09:55:25 and 09:56:26 UTC. Juno was in the dawnside magnetotail near ~05 LT at ~86.5 R_J from Jupiter. This interval was close to a current sheet crossing, as can be seen in the smooth reversals in the radial (B_r) and azimuthal (B_{ϕ}) components from positive to negative values, or vice-versa. The southward-to-





389

Figure 3. Magnetic field observations of a magnetic flux-rope event on DOY 75, 2017 by the 390 Juno spacecraft in the dawnside magnetotail. Panels (a)-(d) show the magnetic field intensity in 391 the spherical JSS coordinate system. The Waves electric field spectra is shown in Panel (e). 392 Panels (f)-(g) show the components of the magnetic field in the minimum-variance coordinate 393 system corresponding to the minimum, intermediate and maximum eigenvector. The magnetic 394 signature of the best-fit modeled force-free flux-rope is shown in blue. Panels (h)-(i) show the 395 hodograms between the field components in the MVA coordinates. Details of the minimum 396 variance analysis and the force-free flux-rope modeling are shown in the grey box. 397

Results of the MVA and force-free modeling for this interval are shown in Figure 3 (f)-(j), 398 where the magnetic field components are plotted in the variance coordinate system. The minimally 399 varying component (B_N) was predominantly in the X_{JSS} direction with values less than 0.3 nT, 400 compared to the field strength between 2 and 4 nT. A reversal could be seen in the maximally 401 402 varying component (B_L) which was accompanied by a moderate increase in the intermediate component (B_M) near the center of the interval, which can also be visualized as a rotation in the 403 $B_L - B_M$ hodogram (Figure 3 (j)). This transient increase in the core field was also seen in the 404 magnetic field magnitude and we classified this event as a magnetic flux rope. The ratios of the 405

eigenvalues of the variance matrix were 7.78 and 45.34, which were large enough to suggest that the MVA analysis unambiguously determined the variance coordinate system. Also shown in Figure 3 (f)-(j) is the modeled force-free flux rope (in blue) which results in the least χ_r^2 . The reduced Chi-squared error between the data and the modeled flux-rope was large (~2.39 nT²) compared to the mean field strength (~ 4 nT), which indicated that the observed flux-rope was not force-free. Nevertheless, the increase in the magnetic field strength and the intermediate component show that this event is a magnetic flux-rope with a strong axial core field.



413

2017-03-16, DOY=075

Figure 4. Energetic particle observations made by the JEDI instruments on board the Juno

spacecraft for the same event as shown in Figure 4. Panel a) shows the variation of B_{θ} during the

event for context. Each consecutive panel shows (b-c) the dynamic energy and pitch-angle

417 spectra for the electrons, (d) the omnidirectional energy spectral index γ for the electrons, the

energy spectra for (e) protons, (f) indistinguishable sulfur and oxygen ions O/S, (g) oxygen ions,
and (h) sulfur ions. All spectra have units of differential intensity i.e., counts/s/str/cm²/keV.

Figure 3 (e) shows the dynamic spectra for the high-frequency fluctuations in the electric field as measured by the Juno Waves instrument. The continuum radiation was observed throughout the interval with a low-frequency cut-off at roughly 1000 Hz. Assuming this to be the electron plasma frequency, we estimated the local electron density during this interval to be ~0.012 cm⁻³. Assuming quasi-neutrality and a singly charged ion mass of 16 amu, this density corresponds to an ion inertial length of 8178.15 km.

In Figure 4 we show the differential intensity measurements by the JEDI energetic particle 426 detector for the same event as shown in Figure 3 for the electrons, protons, oxygen, and sulfur 427 ions, averaged in bins of 30 s each. We show the energy and pitch-angle spectra for the electrons 428 (b)-(c) since higher resolution electron data was available. Also shown for the electrons in (d) is 429 the energy spectral index (γ) obtained by fitting the relation $I = I_0 (E/E_0)^{-\gamma}$ to the omnidirectional 430 differential intensities (1). Only the energy dynamic spectra are shown for the ions (e)-(h), and 431 432 their data was limited due to the lesser energy channels and pitch-angle coverage. JEDI was unable to distinguish between the heavy ions at relatively low energies and these are shown together in 433 panel (f). The JEDI data shown in this figure was resampled to a cadence of 30 s using data from 434 all operating JEDI detectors. There was a moderate increase in the electron flux near the plasmoid 435 event. The electron spectral index decreased after the passage of the plasmoid. Ion fluxes also 436 increased near the plasmoid. The increase in ion flux is most prominent for the sulfur and low 437 438 energy S/O ions.

439

3.1.2. Example 2: O-line - DOY 75, 2017

Juno observed another plasmoid event between 23:57:08 and 23:58:12 UTC on DOY 75, 440 441 2017. The magnetic field and Waves spectra for this example are shown in Figure 5. The magnetic field reversed from a southward to northward configuration before and after this event (B_{θ} changed 442 from positive to negative) and there was no increase seen in the radial or azimuthal component 443 within the reversal. This plasmoid did not have a core field signature and had its minimum field 444 strength at the center of the interval and was therefore classified as a magnetic O-line. The 445 minimum variance analysis showed a similar result, with a minimum in the intermediate variance 446 447 component near the center of the interval. The core field direction, as inferred based on the direction of minimum variance for O-lines, was skewed in the XY plane with a large out-of-plane 448 component. The ratios of the eigenvalues were very large $(\lambda_L/\lambda_M = 15.28 \text{ and } \lambda_M/\lambda_N = 11.46)$ 449 indicating that the MVA coordinate system was well defined. As this event did not have a core 450 451 field, no attempt was made to fit a force-free flux-rope. The Waves spectra, shown in panel (e) showed that the cutoff for the continuum radiation briefly increased during the plasmoid interval 452 from ~500 Hz to 700 Hz. If the cutoff occurs at the electron plasma frequency, this transient 453 increase indicated that the electron density also increased within the interval. This is consistent 454 with the low magnetic field in the center for the O-line type plasmoid. The ~700 Hz cutoff 455 frequency corresponds to an electron density close to 0.006 cm⁻³ and an ion inertial length of 456 ~11683 km assuming quasi-neutrality and a singly charged oxygen ion. 457

Data from the JEDI detectors for this interval is shown in Figure 6. A moderate increase in the electron flux was seen during the plasmoid interval (b). On the other hand, proton, oxygen and sulfur fluxes increased by almost two orders of magnitude compared to times before the plasmoid event (e)-(h). The large ion fluxes were seen consistently even after the B_{θ} reversal and during the

- 462 prolonged interval of negative B_{θ} , or the post-plasmoid plasma sheet. Increases in particle fluxes
- were also seen in data collected by the Galileo EPD (Kronberg et al., 2005) for the larger plasmoids
 discussed in previous surveys.



465

Figure 5. Magnetic field signatures (a)-(d) and electric field spectra (e) obtained by the Juno magnetometer and Waves instruments for a plasmoid event on DOY 75, 2017 in a similar format as Figure 3. No force-free flux rope was fitted for this interval as the field is weakest in the interior of the plasmoid and it is classified as a magnetic O-line.

There was near-complete pitch-angle coverage for the electrons, which were 470 predominantly field-aligned before and after the plasmoid. Before the event, electrons were seen 471 streaming mainly along the magnetic field. The distribution gradually became more isotropic near 472 the plasmoid interval and gradually returned to being field aligned after the passage of the 473 plasmoid, from 23:58 DOY 75 to ~02:00 DOY 76. It has been demonstrated that electron 474 distributions near flux-ropes are influenced by the Fermi and betatron acceleration processes 475 (Zhong et al., 2020; Vaivads et al., 2021) and it is plausible that similar processes are occurring in 476 the present situation. The observed pitch-angle dispersion is seen primarily for electrons with 477 pitch-angles less than 90°, which could also be because of the abundance of field-aligned electrons 478 before and during this interval. Field-aligned and anti-field-aligned electrons are observed in the 479 post-plasmoid plasma sheet (e.g., at DOY 76, 00:02). The electron spectral index γ , shown in 480





483

2017-03-16, DOY=075, 076

Figure 6. Energetic particle differential intensities measured by the JEDI instruments for the event shown in Figure 5 in a similar format as for Figure 4.

In Example 2, the magnetic field was predominantly southward before the event (positive *B*_{θ}) and was in a northward configuration (negative *B*_{θ}) for a brief period (>3 minutes) after. The gradual return to positive *B*_{θ} is consistent with the presence of a large post-plasmoid plasma sheet (PPPS). This is further supported by JEDI observations of higher ion fluxes lasting for the entire PPPS duration. This interpretation follows the schematic shown in Figure 2 (a) for a tailward moving plasmoid with a PPPS.

JEDI data for the electrons and ions was available at high cadence for the two examples discussed above. However, this was not the case for most events in our survey. For this reason, subsequent analysis of the plasmoids uses data gathered primarily by the Juno magnetometer and Waves instruments.



3.2. Survey results: Location and sense of magnetic field reversals

497

496



The automated algorithm searched for plasmoid signatures between DOY 49, 2017 and DOY 150, 2020. Although Juno was inserted into orbit in June 2016, earlier times were effectively not considered due to the LT < 5 filter used in the algorithm to prevent misidentification due to the magnetic field fluctuations near the magnetopause. In this period, corresponding to the stringent criteria described in Section 2.6, the algorithm detected 87 plasmoids with peak-to-peak durations less than 5 minutes. The list of the detected plasmoids is given in the Supporting Information. Out of the 87 events, 47 corresponded to a positive-to-negative reversal in B_{θ} and were thus likely tailward-moving plasmoids, while the remaining 40 events corresponded to a negative-to-positive reversal in B_{θ} and were planetward-moving plasmoids. The terms 'tailward' and 'planetward' are used interchangeably in this work with south-to-north and north-to-south reversals of the magnetic field, respectively, though this is strictly applicable only in the case of a single plasmoid.

511 The locations of the 87 events identified by the algorithm are shown in Figure 7. The vast 512 majority (N=86) of events were identified in the planet-bound portion of Juno's orbit, which lies 513 within the volume of the expected oscillating magnetotail current sheet. All events were located

roughly between 23 and 05 LT due to Juno's orbit being in the midnight to dawn quadrant of the

515 magnetotail between years 2016 and 2020.



516

Figure 8. Histograms showing Juno's dwell time in hours in different radial and local time bins (a-b), as well as the number of tailward moving (c-d) and planetward-moving (e-f) plasmoids identified by the algorithm in the different bins. The dashed line marks the limits of the detection algorithm. Here, 'tailward' and 'planetward' refers to the motion expected for a single plasmoid based on the sense of the B_{θ} reversal.

In Figure 8 we show using histograms the time spent by Juno in different radial and local time regions (a)-(b). Due to Juno's elliptical orbit, it spent more time in the outer magnetosphere as it slowed down near its apogee. On the other hand, each local time between 00 and 06 was sampled almost equally. In panels (c)-(d) we show the occurrence of the planetward-moving plasmoids different radial and local bins. The 47 tailward moving plasmoids were identified

between 30 and 90 R_J (the latter being specified by the automated algorithm). In general, tailward 527 events were observed more frequently at larger radial distances, with the most events (N=18) being 528 observed between 80 and 90 RJ. In terms of local time, >38 tailward events were seen between 03 529 and 05 LT, or close to the dawnside flank. The distribution for planetward moving plasmoids 530 (N=40) was not skewed towards larger radial distances like for the tailward events, though out of 531 a total of 40, 30 events were seen at distances beyond 50 R_I . Although they were observed at all 532 local times, the maximum number of planetward events (N=13) were seen between 03 and 04 LT. 533 4 planetward plasmoids were also seen between 00 and 02 LT, or close to midnight. 534

In interpreting Figure 8, we note again that Juno's trajectory and increasing inclination with time implied that current sheet crossings, and thus, small-scale plasmoids which occur close to the current sheet, were more likely to be seen for the earlier years (2017-2019 or between 03-05 LT) at larger radial distances, and for the later years (2019-2020 or between 00-03 LT) at smaller radial distances closer to the planet, respectively (see Figure 1).



3.3. Duration and size of plasmoids

541

Figure 9. Histogram showing the "peak-to-peak" durations of the identified plasmoids in the present study (a) based on the Juno data and by Vogt et al., (2014) and Kronberg et al. (2008)

using data from the Galileo magnetometer (b). The lower and upper limits specified in ouridentification algorithm are highlighted using red dashed lines.

In Figure 9 (a) we show a histogram of the peak-to-peak durations of the 87 plasmoids 546 identified by the algorithm. The minimum and maximum allowed event size were specified to be 547 10 s and 300 s respectively, as in this work we focus only on small-scale plasmoids with potential 548 diameters comparable to the ion-inertial length. A majority (N=50) of the identified plasmoid 549 events had durations less than 60 s. In general, plasmoids with shorter-duration in situ signatures 550 (and hence smaller diameters), were observed more frequently than plasmoids with longer-551 duration in situ signatures. The mean and median durations for the small-scale plasmoids were 552 found to be 66.44 s and 45 s, respectively. Small-scale events were most likely to be seen having 553 signatures lasting between 10 and 30 s, with 30 out of the 87 events in these two bins. For 554 comparison, the histogram of plasmoid durations by the Vogt et al., (2014) and Kronberg, Woch, 555 Krupp, & Lagg, (2008) surveys are shown in Figure 9 (b), with dashed lines showing the event 556 selection thresholds used by our algorithm. Vogt et al., (2014) found that plasmoids were most 557 likely to have signatures lasting for 5 minutes, with a decreasing trend toward the 0-2 min and 2-558 4 min bins. (Kronberg, Woch, Krupp, & Lagg, 2008) use a different definition for the plasmoid 559 duration, however, they found a similar result with the distribution being skewed towards smaller 560 durations which were less than 20 min. 561



562

Figure 10. a) The Bagenal and Delamere (2011) density profile as a function of radial distance. b) The average magnetotail lobe Alfven speed +/- 5 min from the corresponding B_{θ} reversal for each plasmoid event. c) Scatter plot showing the diameters of the identified events, calculated by assuming that the plasmoids travel at the lobe Alfven speed shown in panel b. The solid line shows the expected ion-inertial length calculated using the Bagenal and Delamere (2011) density profile and using an ion mass of 16 amu. The dashed line marks the limit of the detection algorithm.

We estimated the diameters of the identified plasmoids by assuming that they were 570 travelling at the Alfven speed in the magnetotail lobes, which we estimated based on the measured 571 in situ magnetic field strength and an empirical density profile (Bagenal & Delamere, 2011). 572 Figure 10 shows the typical Alfven speeds for each event (b) along with the empirical density 573 profile used (a)-(b). The Alfven speed was calculated by averaging the magnetic field strength 574 during an interval +/- 5 min from the reversal in B_{θ} . The mean and median Alfven speeds 575 calculated in this manner was 151.85 km/s and 150.27 km/s, respectively. In Figure 10 (c), we 576 show the calculated diameters of the identified plasmoids by multiplying the Alfven speed (b) with 577 the peak-to-peak duration (Δt) for each event. The spacecraft travel distance (Δx) during the Δt 578 579 time was also accounted for, although it is negligible compared to the plasmoid motion. Also shown in Figure 10 (c) is the ion-inertial length as a function of radial distance, calculated based 580 on the same density profile. An ion mass of 16 amu was assumed for the calculations due to the 581 dominance of heavy ions $(O^+/S^+/S^{++})$ in the Jovian magnetosphere. Figure 10 (c) illustrates that 582 the diameters of all events with in-situ signatures shorter than 5 minutes are within an order of 583 magnitude of the local ion inertial length, with some events having diameters even shorter than 584 585 this length scale.

However, we note that defining the plasmoid duration between the two extrema in B_{θ} may lead to underestimation of the size of the plasmoid as the reversal alone does not account for the post-reversal post-plasmoid plasma sheet (PPPS) (Vogt et al., 2014; Cowley et al., 2015).

589 **3.4. Relative occurrence of flux-ropes versus O-lines**

For some plasmoids identified by the algorithm, we observed that an increase in the 590 magnetic field strength near the center of the events i.e., near the reversal in B_{θ} , caused an 591 increased in the component of the magnetic field in the direction of intermediate variance (B_M) . 592 However, these increases in B_M did not always correspond to an increase the overall magnetic field 593 strength due to the near-zero values of the component of maximum variance (B_L) near the center 594 of the event as it reversed sign during the plasmoid interval. Hence, in our work, we classified 595 those events as flux-ropes in which the median of the intermediate-variance component (B_M) was 596 larger than that measured at the beginning and end of the event interval. In other cases, where there 597 was no localized increase in B_M , the reversal of the magnetic field usually resulted in a minimum 598 field strength at the center of the event interval; such events were classified as magnetic O-lines. 599 Out of 87 plasmoids events identified by the algorithm, 31 were classified as magnetic flux-ropes 600 and 56 were classified as magnetic O-lines. 601

Figure 11 shows the distributions of flux-ropes and O-lines identified by the algorithm in radial distance and local time. Out of the 31 flux-ropes, 28 were found at radial distances beyond $50 R_J$. Similarly, out of 56 O-lines, 43 were observed beyond $50 R_J$, which could be due to the longer time spent by Juno in distances beyond $50 R_J$. The greatest number of O-lines were observed between 80 and 90 R_J (N=18), whereas flux-ropes were found likely to occur between 50 and 90 R_J . More flux-ropes and O-lines were observed near the dawn-side magnetotail (03-05 LT) than near midnight (00-02 LT). This could be due to Juno's orbit, as between 00 and 02 LT it only crossed the current sheet at radial distances inward of 50 R_J, which could be planetward of the reconnection X-line in these local times.



611

Figure 11. Histograms of the locations of the identified flux-rope and O-line events in radial
distance and local time. Time spent by the Juno spacecraft in each bin is shown in panels a) and
b). The dashed lines mark the limits of the detection algorithm.

The direction of the core field for a magnetic flux-rope is closely associated with the MVA 615 eigenvector corresponding to the direction of intermediate variance (\vec{x}_M) , which is illustrated in 616 Figure 12 for all 31 flux-rope events. Also shown in Figure 12 are representative magnetic field 617 lines from the Sarkango et al., (2019) MHD model which show the bend-back of the magnetic 618 field lines due to the sub-corotation of the magnetospheric plasma. The flux-rope events were 619 observed at nearly all local times sampled by Juno, with a wide range of core field orientations. 620 Also shown in Figure 12 (b) are the eigenvectors corresponding to the direction of maximum 621 variance $(\vec{x_I})$, which should be predominantly in the north-south direction as it was the basis on 622 623 which these events were identified.

For an O-line, the direction of minimum variance of the field is used to infer the core-field direction as the intermediate variance is expected to be close to the radial direction. Figure 13 shows the MVA eigenvectors corresponding to the direction of minimum variance $(\overline{x_N})$ (top) and maximum variance $(\vec{x_L})$ for all 56 O-line events in a similar format as Figure 12. As for the fluxropes, O-lines were observed at nearly all local times between 00 and 05 and had a wide range of core field orientations compared to the local bend-back direction. For both flux-ropes and O-lines, the direction of maximum variance was close to the Z direction in the JSS coordinate system.



631

Figure 12. Eigenvectors of the intermediate and maximum variance shown in the equatorial and
 meridional projections for the 31 flux-rope events identified by the algorithm. Each event is
 colored according to its peak-to-peak duration. Also shown in blue are magnetic field lines from
 an MHD model.

636



637

Figure 13. Eigenvectors of the minimum and maximum variance shown in the equatorial and
 meridional projections for the 56 O-line events identified by the algorithm. Each event is colored
 according to its peak-to-peak duration. Also shown in blue are magnetic field lines from an
 MHD model.

We compared the core-field directions of the flux-ropes and the O-lines with the local bendback plane, which is presumably also the plane in which reconnection is occurring. The plane of the bend-back for each plasmoid events was determined by calculating the average of $\alpha =$ tan⁻¹(B_{ϕ}/B_r) for a 10-minute period starting 20 minutes prior to the observed B_{θ} reversal, during which time the spacecraft was typically sampling the off-equatorial / mid-latitudes, as seen in the

anti-correlated azimuthal (B_{ϕ}) and radial (B_r) components of the magnetic field. In Figure 14, we 647 show the angle between the core-field direction and the vector intersecting the local bend-back 648 649 plane and the JSS equatorial plane $(\vec{x_{\alpha}})$. The direction of intermediate variance was chosen as the core-field direction for the magnetic flux-rope events whereas the minimum variance direction 650 was chosen for the O-lines. Near 90° values of $\Delta \phi$ imply that the axial direction of the plasmoid 651 was perpendicular to the bend-back plane, which is predominantly the case for O-lines, which 652 show preferences for larger acute angles. Such a result is very consistent with the configuration of 653 O-lines formed in the midnight-to-dawn quadrant in the global simulation of Saturn's 654 magnetosphere by Jia et al., (2012). Although more flux-ropes were seen with $\Delta \phi > 45^{\circ}$ (N=19) 655 than vice-versa (N=13), the histograms show that angles between 20-30° were as likely to occur 656 as those between 70-80°. 657



658

Figure 14. Histograms showing the acute angle between the core-field direction and the local bend-back plane for a) flux-ropes and b) O-lines identified by the algorithm.

A constant- α force-free flux-rope model was fitted to each flux-rope event (N=31) based on the methodology described in Section 2.5. Out of the 31 which were originally classified as flux-ropes, 6 produced reduced-chi-squared errors less than 0.3 nT² for an average field strength between 2 and 6 nT, which suggests that they were close to being force-free.

665 4. Discussion

666 **4.1. Ion-inertial scale plasmoids in the Jovian magnetotail**

The magnetic flux rope event shown in Figure 3 lasted approximately 61 s between the two 667 extrema in B_{θ} , which roughly corresponds to a plasmoid diameter of 17141 km, assuming it 668 travelled at the Alfven speed of ~281 km/s corresponding to that in the surrounding magnetotail 669 lobes (B=4 nT, n_{lobe}=0.006 cm⁻³, based on Waves data). This diameter was roughly 2 times the 670 local ion inertial length, which we calculated to be ~8178 km (assuming quasi-neutrality and ion 671 mass of 16 amu). The second example, shown in Figure 5, had a 'peak-to-peak' duration of 63 s. 672 The magnetic field strength in the surrounding lobes was approximately 5 nT, which corresponds 673 to an Alfven speed of ~498 km/s (for $n_{lobe}=0.003$ cm⁻³) and a plasmoid diameter of ~31374 km. 674

The ion inertial length during this interval was also ~11683 km, making the plasmoid approximately 2.68 times larger than this length scale.

Figure 9 and Figure 10 illustrate the scale of the plasmoids detected by the algorithm using 677 similar calculations. A majority (N=50 out of 87) of plasmoids were seen to last less than 60 s, 678 measured between the peaks in B_{θ} . In contrast, previous surveys (Figure 9 (b)) have discussed 679 larger plasmoids with durations longer than 2 min. In Figure 10, we calculated the diameters of 680 the plasmoids based on the 'peak-to-peak' duration and the Alfven speed in the lobes and 681 demonstrate that the diameters of all events were within one order of magnitude of the local oxygen 682 ion-inertial length. Our survey demonstrates that reconnection occurs in the Jovian magnetotail at 683 kinetic scales and produces ion-inertial scale plasmoids, like in other regions in the space 684 environments such as the terrestrial magnetotail. Numerical simulations of ion-scale plasmoids 685 have shown that they evolve by interacting with other plasmoids, i.e., via coalescence, which 686 could ultimately produce the large plasmoids analyzed by previous surveys. 687

Our calculations of plasmoid diameters are subject to the assumption that the 'peak-to-688 689 peak' signature represents the entire plasmoid crossing. However, limiting the plasmoid duration between the two extrema in B_{θ} may underestimate the plasmoid size by a factor of 4 or more (Vogt 690 et al., 2014). Note that we have used a mass of 16 amu or the mass of O⁺ for the average mass of 691 singly charged ions. Due to the abundance of singly charged sulfur ions, the ion inertial length in 692 the Jovian magnetotail may be even larger than the numbers described above, which further 693 supports our conclusions that our identified small-scale plasmoids have spatial sizes comparable 694 695 to the local ion inertial length.

Our observations show that reconnection in Jupiter's magnetotail may be proceeding via 696 current-sheet instabilities, as found in other magnetospheres The tearing instability is most likely 697 to occur in a thin current sheet. In the magnetotail, this depends on the ratio of the radial field in 698 the lobes $(B_{r,lobe})$ and the meriodional field at the equator $(B_{\theta,equator})$ (Kronberg et al., 2007), 699 which changes over the course of the large-scale loading of the magnetosphere. Hence, on a global 700 scale, there are no different onset conditions for the small-scale reconnection events. However, as 701 702 smaller plasmoids with diameters roughly between $\sim 10,000$ to 50,000 km are easier to miss being detected by a single spacecraft in the magnetotail, the observations presented here raise new 703 questions about the frequency at which reconnection occurs and whether it could also be occurring 704 705 in other regions of the magnetosphere, such as the dusk side magnetotail.

706

4.2. Abundance of O-lines versus flux-ropes

On comparing the Waves electric field spectra for Example 1 (Figure 3 (e)), we observe 707 that the cutoff for the continuum radiation briefly decreased from 1.1kHz to 0.9 kHz within the 708 plasmoid event. This minor transient decrease was different from the larger increase in the cutoff 709 frequency (from 500 Hz to 1 kHz) over the broader interval containing the plasmoid event, and 710 the latter could a result of Juno entering the dense plasmasheet. The transient decrease in the cutoff 711 frequency indicates a decrease in the plasma density within the plasmoid. The magnetic field 712 713 strength increased within the event due to the core field, which led us to classify it as a magnetic flux-rope. Hence, the density decrease observed by Waves is consistent with the result that 714 Example 1 was a flux-rope where the magnetic tension of the outer wraps is balanced by the larger 715 magnetic pressure in the interior. However, this flux-rope was not entirely force-free, as discussed 716 717 in Section 3.1.1.

The cutoff briefly increased to higher frequencies for Example 2 (Figure 5 (e)) from 500 Hz to 700 Hz, indicating that the plasma frequency and electron density also increased. Simultaneously, the magnetic field strength decreased within the interval. This supports our interpretation that the second plasmoid (Figure 5 and Figure 6) was an O-line, and the magnetic pressure of the outer helical structure was balanced by the larger thermal pressure contained within the plasmoid core. The relatively high cutoff frequency of ~1 kHz and corresponding higher density were also seen after the reversal in the post-plasmoid plasma sheet.

Like the Vogt et al., (2014) survey for Jupiter and the Jackman et al., (2011) survey for 725 Saturn, we find that a majority of B_{θ} reversals associated with a plasmoid event do not have strong 726 core fields. Out of the total 87 plasmoid candidates only 31 can be classified as magnetic flux-727 ropes, and out of these only 6 were found to fit the force-free model well with χ_r^2 less than 0.3 nT². 728 This could be because the plasmoids are generated within a high β plasma, where plasma pressure 729 gradients are very large. Dense plasma inside plasmoids could increase the thermal pressure which 730 would oppose the magnetic tension of the outer helical wraps (Kivelson & Khurana, 1995). In 731 contrast, the abundance of flux-ropes seen in the terrestrial-like magnetospheres could also be 732 because of the stronger cross-magnetotail magnetic field introduced by the IMF (Slavin et al., 733 734 2003), which has a weaker, if not negligible, influence in the Jovian magnetosphere as the magnetopause reconnection is relatively weak. In the terrestrial case, the northern and southern 735 magnetotail lobes are sheared due to the solar wind interaction (Cowley, 1981), which can 736 contribute to a core-field for a plasmoid structure. (Here 'sheared' refers to the displacement of 737 the two lobes with respect to each other, rather than to the anti-parallel magnetic field across the 738 current sheet). Meanwhile, at Jupiter, the northern and southern lobes are not sheared, though the 739 740 entire reconnection plane is likely skewed due to the bend-back effect (Russell et al., 1998). This could also result in more O-lines being generated than flux-ropes at Jupiter and Saturn. 741

Results of the minimum variance analysis on all flux-rope and O-line events showed that 742 their core fields can be highly skewed in the XY plane (Figure 12 and Figure 13). In an ideal 743 scenario, the core field of a newly produced plasmoid would be perpendicular to the plane of 744 745 reconnection. In the Jovian magnetotail, this is determined by the 'bend-back' of the magnetic field lines due to sub-corotation of the plasma. So, it is interesting to compare the core field 746 directions with the orientation of the local bend-back plane. In Figure 14 we showed the angle 747 between these two vectors. The majority (~70%) of O-line events identified by our algorithm have 748 core fields which are oriented larger than 45° from the bend-back plane. A similar but less 749 prominent result is found for the flux-ropes, ~60% of which subtend angles larger than 45° with 750 respect to the bend-back plane. 751

In Figure 11 we showed the occurrence of flux-ropes and O-lines detected by our algorithm 752 at different radial distances and local times. Both flux-ropes and O-lines were seen more frequently 753 at larger radial distances due to Juno's orbital bias. However, very few flux-ropes were seen at 754 distances inward of 50 R_J, while O-lines were equally likely to occur at all bins between 30 and 755 70 R_J, though they are most likely between 70 and 90 R_I . Assuming that flux-ropes are plasmoids 756 757 at a later stage during their evolution than O-lines, it is plausible that their signatures would be seen more frequently for regions in the deep magnetotail. Conversely, O-lines, which contain 758 dense plasma and are possibly products of fresh reconnection, are more likely to occur closer to 759 the planet. 760

761 **4.3. Contribution of small-scale plasmoids to mass loss from the magnetosphere**

We can produce cursory estimates of the mass carried by a small-scale plasmoid by 762 assuming that it occupies a cylindrical volume with diameter equal to the average calculated in our 763 survey, ~10,000 km and possesses a uniform density of 0.05 cm⁻³. Consider three additional 764 parameters: f, a factor to scale the plasmoid diameter to obtain the cross-tail dimension, \dot{M} , the 765 net mass loss rate due to plasmoids, and δt , the time duration between two consecutive plasmoid 766 events. With f = 10 (plasmoid cross tail length equal to 10 times its diameter), one plasmoid with 767 cross-sectional diameter of 10000 km would need to be released every $\delta t = 0.1$ s to provide a 768 total mass loss rate $\dot{M} = mnV/\delta t$ (where V is the plasmoid volume) equal to 100 kg/s. Note that 769 the actual plasmoid-related mass loss deficit is >500 kg/s, which is even larger. Plasmoids which 770 were released at such high frequencies would have been detected more often by Juno, which spent 771 more than 150 hours at distances less than 10,000 km to the current sheet. 772

In the above calculation, the assumed plasmoid diameter is 0.13 R_J . In contrast, Cowley et al., (2015) assumed dimensions of ~150 R_J in the tailward direction, ~70 R_J in the cross-tail direction and ~7 R_J for the direction normal to the current sheet plane. They argue that the large tailward dimensions of the plasmoid are due to the post-plasmoid plasma sheet, which is present in the reconnection exhaust.

An alternative approach is to evaluate the effective outflow area needed to lose 1 ton/s of 778 mass continuously, e.g., via steady reconnection instead of sporadic plasmoid release. If we 779 assume the density at ~80 R_I downstream is ~ 0.01 cm⁻³ and the outflow speed is ~400 km/s, a 780 total area of ~1828 R_I^2 would be needed to lose 1 ton of mass per second. This translates to a square 781 region ~42.7 R_I in length on each side. The area and square length reduce to 366 R_I^2 and 19.1 R_I if 782 the density is increased by a factor of 5. Since observations have shown that reconnection in the 783 Jovian magnetotail does not appear to occur in a steady manner, and a period of several days passes 784 between consecutive large-scale 'unloading' events, the combined dimensions (in at least two 785 directions) of all plasmoids in an 'unloading' interval must be larger than the length scales needed 786 for continuous outflow of plasma, if they must account for all 1 ton/s of production. 787

Hence, the conclusion from this discussion is that small-scale magnetic reconnection does not contribute in a substantial manner to mass loss from the Jovian magnetosphere, and only large plasmoids with a prolonged PPPS may potentially account for the loss of 1 ton/s of mass from the magnetosphere. Other loss mechanisms apart from magnetic reconnection are out of scope of the present work and are not discussed here.

793

4.4. Orbit-by-orbit variability and comparison with other studies

The number of plasmoid events identified by our algorithm varied between perijove passes 794 795 or Juno orbits. It has been shown by Vogt et al., (2020) and Yao et al., (2019) that some orbits were more "active" than others, in the form of frequent in situ sightings of magnetic reconnection 796 or of more dynamic aurora. We found a similar result, and in Figure 15 we compare the occurrence 797 798 of the plasmoids presented in this survey with the reconnection events identified by Vogt et al., (2020) for different Juno perijoves. Both surveys show excellent agreement with each other. Orbits 799 5, 8, 9, 10, and 11 were especially active. On the other hand, fewer events were seen for orbit 7, 800 and for orbit numbers beyond 12. The latter could perhaps be explained by Juno's increasing 801 inclination and lack of sampling of larger radial distances (Section 2.1). The orbital bias is less 802

prominent for adjacent perijoves like 6 and 7 which nevertheless show different levels of activity. 803 This variation on time scales longer than fluctuations in the upstream solar wind conditions could 804 be explained by a different internal magnetospheric state, e.g., due to changes in Io's volcanic 805 activity. The two example plasmoid events discussed in Section 3.1 were observed by Juno on 806 DOY 75 and 76, 2017 during orbit 5 within a longer period of unloading of the magnetosphere, 807 which was discussed by Yao et al., (2019). They also noted that the UV aurora was dimmer during 808 orbit 7 than orbit 5. The long-term variability of the Jovian magnetosphere may be linked to other 809 visible changes such as in the strength of the magnetodisc current sheet and location of the Jovian 810 aurora, which could occur due to changes in mass loading from Io rather than external solar wind 811 conditions (Vogt et al., 2017). 812



813

Figure 15. Histograms showing the occurrence of reconnection events in the Vogt et al., (2020) survey and of plasmoid events in our present work for different Juno orbits or perijoves. Orbits 4 and below were excluded from the present work.

817 **5. Conclusions**

In this work, we used data from the Juno spacecraft to identify plasmoids in the Jovian 818 magnetotail with in situ durations on the order of 5 minutes or less. These small-scale events have 819 diameters comparable to the ion inertial length, which is an important length scale as it is related 820 to the size of the ion diffusion region in magnetic reconnection. We used data from the Juno 821 822 magnetometer to identify reversals in the north-south component of the magnetic field, which is expected to occur in a reconfiguration of the magnetotail during magnetic reconnection. An 823 automated algorithm is applied to detect B_{θ} reversals and identify corresponding extrema which 824 determine the start and end of the event. Various filters are applied during the detection procedure 825 based on the minimum variance analysis and other magnetic properties such as the proximity to 826 the current sheet and the magnitude of the perturbation associated with the event. Based on our 827 algorithm, we detected 87 plasmoids with 'peak-to-peak' durations between 10 s and 5 min, out 828 of which 31 were seen to have an increase in the field component normal to the reconnection plane 829 and were classified as magnetic flux-ropes, while the 56 events with minimum field strengths at 830 their centers were classified as magnetic O-lines. 831

We examined two such plasmoid events in more detail due to the availability of 832 simultaneous, higher cadence energetic particle intensities measured by the JEDI instruments. The 833 results of the minimum variance analysis show that the first event can be classified as a magnetic 834 flux-rope whereas the second had minimum field strength at its center and was an O-line. 835 Energetic particle fluxes were seen to maximize for the second plasmoid event for the electrons, 836 protons, oxygen and sulfur ions. Moreover, the electron pitch angle spectra indicate isotropic 837 distributions within the magnetic loop structure, which could be due to betatron acceleration either 838 at the front of, or within the plasmoid. For the second event, the isotropic pitch-angle distribution 839 gradually tends to become field-aligned in the proximity of the plasmoids. 840

We used the 'peak-to-peak' duration between the two extrema in B_{θ} to calculate the 841 duration for each plasmoid event identified by the algorithm. Within the interval of 10 s and 5 min 842 chosen for the algorithm, a majority of plasmoids (N=50 out of 87) were seen with durations lasting 843 844 less than 60 s. It is interesting to compare the distribution of plasmoid durations with previous surveys (Kronberg, Woch, Krupp, & Lagg, 2008; Vogt et al., 2014). Although the two previous 845 studies used different definitions to define a plasmoid and looked for signatures on different 846 timescales, their histograms also showed a similar behavior. In both studies, the histograms were 847 skewed toward smaller values, indicating that smaller plasmoids were more likely to be observed, 848 depending on the length scales under consideration. Similar results have also been reported for 849 flux-ropes observed in the solar wind (Hu et al., 2018), and for plasmoids seen in Saturn's 850 magnetosphere (Garton et al., 2021), which have shown that plasmoid diameters exhibit a power-851 law-like scaling. The duration of the 87 plasmoids observed in our survey is also used to estimate 852 the plasmoids' diameters using the lobe Alfven speed. We demonstrate that all events with 853 durations less than 5 minutes can have diameters within an order of magnitude larger or smaller 854 than the local ion-inertial length. These results demonstrate that magnetic reconnection occurs in 855 the Jovian magnetotail at ion kinetic scales, like in other regions in the space environments. This 856 is important as multiple ion-inertial scale plasmoids can coalesce to form larger plasmoids, such 857 858 as those analyzed by previous studies, and can also trap and accelerate ions and electrons.

The abundance of O-lines (N=56) versus flux-ropes (N=31) identified by the algorithm is 859 consistent with previous surveys of plasmoids in Jupiter's (Vogt et al., 2014) and Saturn's 860 (Jackman et al., 2011) magnetotail. Using the minimum variance analysis, we show that O-lines 861 detected by our algorithm were more likely to have an axial direction perpendicular to the 862 reconnection plane, which in the case of the Jovian magnetotail is assumed to be the local plane of 863 bent-back magnetic field. The core-fields for the flux-ropes do not show a clear relationship with 864 the bend-back plane, which could be because flux-ropes structures are at a later stage of plasmoid 865 evolution and are 'de-coupled' from the corotation dynamics that cause the bend-back. 866 Alternatively, the small number of flux ropes formed at Jupiter may simply be due to the weakness 867 of the IMF at 5.2 AU combined with vast dimensions of this huge magnetosphere. As a result, 868 reconnection deep in Jupiter's magnetosphere may not be aware of the direction of the IMF and 869 any shearing of the two tail lobes due to IMF stress may be extremely weak. Such a situation may 870 well favor the development of O-lines as opposed to flux-ropes. In contrast, in terrestrial-like 871 magnetospheres like those of Earth and Mercury, the IMF B_{y} component is an important factor to 872 produce the core field of the plasmoid. 873

⁸⁷⁴ Despite being smaller, 87 ion-inertial scale plasmoids were detected by Juno and captured ⁸⁷⁵ by our algorithm. We demonstrate that a plasmoid with a cross-sectional diameter of 10,000 km ⁸⁷⁶ (= $0.13 R_I$) and density of 0.5 cm⁻³ would need to be released at least once every 0.1 s to result in

a 100 kg/s loss rate. We argue that this release frequency is unlikely since Juno would have 877 detected more such events had it been the case. Based on cursory calculations, we argue also that 878 the total dimensions of plasmoids in at least two dimensions must be larger than several tens of R_{I} 879 to match the dimensions of an effective outflow area needed to lose 1 ton/s of mass. This may 880 become possible if the post-plasmoid plasma sheet is included in the calculations, as argued by 881 Cowley et al., (2015). Hence, we suggest that ion-inertial scale plasmoid release, by itself, is an 882 insignificant loss mechanism, unless several hundred such events are produced simultaneously 883 every 1 s or less, which is unlikely according to the current observations. 884

Lastly, we compare the events detected by our algorithm with the survey of reconnection signatures observed by Vogt et al., (2020). The relative occurrence of plasmoid and reconnection events show a very good agreement and both studies find different magnetospheric behavior for different yet consecutive Juno orbits. For example, in both studies, the number of reconnection signatures seen during Juno orbit #8 were less than half of the total number seen during orbit #9. The mechanisms which can lead to such variability over the long timescale associated with each orbit (~53 days), need to be examined further.

892 Acknowledgments, Samples, and Data

893 This work was supported by NASA through the NASA Earth and Space Science Fellowship

674 Grant 80NSSC17K0604 and Early Career Fellow Startup Grant 80NSSC20K1286. The research

- at the University of Iowa is supported by NASA through contract 699041X with the Southwest
- 896 Research Institute. The Juno magnetometer (MAG), Waves and JEDI data used in this study is
- ⁸⁹⁷ publicly available from the Planetary Plasma Interactions node of the Planetary Data System at
- 898 https://pds-ppi.igpp.ucla.edu/ .

899 **References**

- Artemyev, A. V., Kasahara, S., Ukhorskiy, A. Y., & Fujimoto, M. (2013). Acceleration of ions
 in the Jupiter magnetotail: Particle resonant interaction with dipolarization fronts. *Planetary and Space Science*, 82–83, 134–148. https://doi.org/10.1016/j.pss.2013.04.013
- Artemyev, A. V., Clark, G., Mauk, B., Vogt, M. F., & Zhang, X. J. (2020). Juno Observations of
 Heavy Ion Energization During Transient Dipolarizations in Jupiter Magnetotail. *Journal of Geophysical Research: Space Physics*, 125(5), 1–14. https://doi.org/10.1029/2020JA027933
- Bagenal, F. (2007). The magnetosphere of Jupiter: Coupling the equator to the poles. *Journal of Atmospheric and Solar-Terrestrial Physics*, 69(3), 387–402.
 https://doi.org/10.1016/j.jastp.2006.08.012
- Bagenal, F., & Delamere, P. A. (2011). Flow of mass and energy in the magnetospheres of
 Jupiter and Saturn. *Journal of Geophysical Research: Space Physics*, *116*(5).
 https://doi.org/10.1029/2010JA016294
- Barnhart, B. L., Kurth, W. S., Groene, J. B., Faden, J. B., Santolik, O., & Gurnett, D. A. (2009).
- Electron densities in Jupiter's outer magnetosphere determined from Voyager 1 and 2
 plasma wave spectra. *Journal of Geophysical Research: Space Physics*, *114*(5), 5218.
 https://doi.org/10.1029/2009JA014069
- Connerney, J. E. P., Benn, M., Bjarno, J. B., Denver, T., Espley, J., Jorgensen, J. L., et al. (2017).
 The Juno Magnetic Field Investigation. *Space Science Reviews*.

- 918 https://doi.org/10.1007/s11214-017-0334-z
- Cowley, S. W. H. (1981). Magnetospheric asymmetries associated with the y-component of the
 IMF. *Planetary and Space Science*, 29(1), 79–96. https://doi.org/10.1016/00320633(81)90141-0
- Cowley, S. W. H., Nichols, J. D., & Jackman, C. M. (2015). Down-tail mass loss by plasmoids in
 Jupiter's and Saturn's magnetospheres. *Journal of Geophysical Research A: Space Physics*,
 120(8), 6347–6356. https://doi.org/10.1002/2015JA021500
- Delamere, P. A., & Bagenal, F. (2013). Magnetotail structure of the giant magnetospheres:
 Implications of the viscous interaction with the solar wind. *Journal of Geophysical Research: Space Physics*, *118*(11), 7045–7053. https://doi.org/10.1002/2013JA019179
- Drake, J. F., Swisdak, M., Che, H., & Shay, M. A. (2006). Electron acceleration from contracting
 magnetic islands during reconnection. *Nature*, 443(7111), 553–556.
 https://doi.org/10.1038/nature05116
- Garton, T. M., Jackman, C. M., & Smith, A. W. (2021). Kronian Magnetospheric Reconnection
 Statistics Across Cassini's Lifetime. *Journal of Geophysical Research: Space Physics*,
 126(8), e2021JA029361. https://doi.org/10.1029/2021JA029361
- Ge, Y. S., Jian, L. K., & Russell, C. T. (2007). Growth phase of Jovian substorms. *Geophysical Research Letters*, *34*(23), 1–6. https://doi.org/10.1029/2007GL031987
- Ge, Y. S., Russell, C. T., & Khurana, K. K. (2010). Reconnection sites in Jupiter's magnetotail
 and relation to Jovian auroras. *Planetary and Space Science*, 58(11), 1455–1469.
 https://doi.org/10.1016/j.pss.2010.06.013
- Grigorenko, E. E., Malykhin, A. Y., Kronberg, E. A., Malova, K. V., & Daly, P. W. (2015).
 Acceleration of ions to suprathermal energies by turbulence in the plasmoid-like magnetic
 structures. *Journal of Geophysical Research A: Space Physics*, *120*(8), 6541–6558.
 https://doi.org/10.1002/2015JA021314
- Gurnett, D. A., Scarf, F. L., Kurth, W. S., Shaw, R. R., & Poynter, R. L. (1981). Determination
 of Jupiter's electron density profile from plasma wave observations. *Journal of Geophysical Research: Space Physics*, 86(A10), 8199–8212. https://doi.org/10.1029/JA086IA10P08199
- Hu, Q., Zheng, J., Chen, Y., le Roux, J., & Zhao, L. (2018). Automated Detection of Small-scale
 Magnetic Flux Ropes in the Solar Wind: First Results from the Wind Spacecraft
 Measurements . *The Astrophysical Journal Supplement Series*, 239(1), 12.
 https://doi.org/10.3847/1538-4365/aae57d
- Jackman, C. M., Slavin, J. A., & Cowley, S. W. H. (2011). Cassini observations of plasmoid
 structure and dynamics: Implications for the role of magnetic reconnection in
 magnetospheric circulation at Saturn. *Journal of Geophysical Research: Space Physics*.
 https://doi.org/10.1029/2011JA016682
- Jia, X., Hansen, K. C., Gombosi, T. I., Kivelson, M. G., Tóth, G., Dezeeuw, D. L., & Ridley, A.
 J. (2012). Magnetospheric configuration and dynamics of Saturn's magnetosphere: A global
 MHD simulation. *Journal of Geophysical Research: Space Physics*, *117*(5).
 https://doi.org/10.1029/2012JA017575
- 958 Joy, S. P., Kivelson, M. G., Walker, R. J., Khurana, K. K., Russell, C. T., & Ogino, T. (2002).

Probabilistic models of the Jovian magnetopause and bow shock locations. Journal of 959 Geophysical Research: Space Physics, 107(A10). https://doi.org/10.1029/2001JA009146 960 Kasahara, S., Kronberg, E. A., Kimura, T., Tao, C., Badman, S. V., Masters, A., et al. (2013). 961 Asymmetric distribution of reconnection jet fronts in the Jovian nightside magnetosphere. 962 Journal of Geophysical Research: Space Physics, 118(1), 375–384. 963 https://doi.org/10.1029/2012JA018130 964 Khurana, K. K., Kivelson, M. G., Vasyliunas, V. M., Krupp, N., Woch, J., Lagg, A., et al. 965 (2004). The Configuration of Jupiter's Magnetosphere (Chapter 24). Jupiter: The Planet, 966 Satellites and Magnetosphere. 967 Kivelson, M. G., & Khurana, K. K. (1995). Models of flux ropes embedded in a harris neutral 968 sheet: Force-free solutions in low and high beta plasmas. Journal of Geophysical Research, 969 100(A12), 23637. https://doi.org/10.1029/95ja01548 970 971 Kivelson, M. G., & Southwood, D. J. (2005). Dynamical consequences of two modes of centrifugal instability in Jupiter's outer magnetosphere. Journal of Geophysical Research: 972 973 Space Physics, 110(A12). https://doi.org/10.1029/2005JA011176 Krimigis, S. M., Carbary, J. F., Keath, E. P., Bostrom, C. O., Axford, W. I., Gloeckler, G., et al. 974 975 (1981). Characteristics of hot plasma in the Jovian magnetosphere: Results from the 976 Voyager spacecraft. Journal of Geophysical Research: Space Physics, 86(A10), 8227-8257. https://doi.org/10.1029/JA086IA10P08227 977 Kronberg, E. A., Woch, J., Krupp, N., Lagg, A., Khurana, K. K., & Glassmeier, K. H. (2005). 978 Mass release at Jupiter: Substorm-like processes in the Jovian magnetotail. Journal of 979 Geophysical Research: Space Physics, 110(A3), 1–10. 980 https://doi.org/10.1029/2004JA010777 981 982 Kronberg, E. A., Glassmeier, K. H., Woch, J., Krupp, N., Lagg, A., & Dougherty, M. K. (2007). A possible intrinsic mechanism for the quasi-periodic dynamics of the Jovian 983 magnetosphere. Journal of Geophysical Research: Space Physics, 112(5). 984 https://doi.org/10.1029/2006JA011994 985 Kronberg, E. A., Woch, J., Krupp, N., Lagg, A., Daly, P. W., & Korth, A. (2008). Comparison of 986 987 periodic substorms at Jupiter and Earth. Journal of Geophysical Research: Space Physics, 113(4), 1-11. https://doi.org/10.1029/2007JA012880 988 989 Kronberg, E. A., Woch, J., Krupp, N., & Lagg, A. (2008). Mass release process in the Jovian magnetosphere: Statistics on particle burst parameters. Journal of Geophysical Research: 990 Space Physics, 113(10). https://doi.org/10.1029/2008JA013332 991 Kronberg, E. A., Grigorenko, E. E., Malykhin, A., Kozak, L., Petrenko, B., Vogt, M. F., et al. 992 (2019). Acceleration of Ions in Jovian Plasmoids: Does Turbulence Play a Role? Journal of 993 Geophysical Research: Space Physics, 124(7), 5056–5069. 994 https://doi.org/10.1029/2019JA026553 995 Krupp, N., Woch, J., Lagg, A., Wilken, B., Livi, S., & Williams, D. J. (1998). Energetic particle 996 bursts in the predawn Jovian magnetotail. Geophysical Research Letters, 25(8), 1249–1252. 997 https://doi.org/10.1029/98GL00863 998 999 Krupp, N., Vasyliunas, V. M., Woch, J., Lagg, A., Khurana, K. K., Kivelson, M. G., et al.

- 1000 (2004). Dynamics of the Jovian Magnetosphere. In *Jupiter. The planet, satellites and*1001 *magnetosphere.*
- Kurth, W. S., Hospodarsky, G. B., Kirchner, D. L., Mokrzycki, B. T., Averkamp, T. F., Robison,
 W. T., et al. (2017). The Juno Waves Investigation. *Space Science Reviews*.
 https://doi.org/10.1007/s11214-017-0396-y
- Lepping, R. P., Jones, J. A., & Burlaga, L. F. (1990). Magnetic field structure of interplanetary
 magnetic clouds at 1 AU. *Journal of Geophysical Research*, 95(A8), 11957.
 https://doi.org/10.1029/ja095ia08p11957
- Masters, A. (2017). Model-Based Assessments of Magnetic Reconnection and Kelvin-Helmholtz
 Instability at Jupiter's Magnetopause. *Journal of Geophysical Research: Space Physics*,
 122(11), 11,154-11,174. https://doi.org/10.1002/2017JA024736
- Mauk, B. H., Haggerty, D. K., Jaskulek, S. E., Schlemm, C. E., Brown, L. E., Cooper, S. A., et
 al. (2013). The Jupiter Energetic Particle Detector Instrument (JEDI) Investigation for the
 Juno Mission. *Space Science Reviews 2013 213:1, 213*(1), 289–346.
 https://doi.org/10.1007/S11214-013-0025-3
- Moldwin, M. B., & Hughes, W. J. (1992). On the formation and evolution of plasmoids: A
 survey of ISEE 3 geotail data. *Journal of Geophysical Research: Space Physics*, 97(A12),
 19259–19282. https://doi.org/10.1029/92JA01598
- Nishida, A. (1983). Reconnection in the Jovian magnetosphere. *Geophysical Research Letters*,
 10(6), 451–454. https://doi.org/10.1029/GL010i006p00451
- Priest, E. R. (2011). The equilibrium of magnetic flux ropes (tutorial lecture). *Physics of Magnetic Flux Ropes (Eds C.T. Russell, E.R. Priest and L.C. Lee)*, 1–22.
 https://doi.org/10.1029/gm058p0001
- Rosa Oliveira, R. A., da Silva Oliveira, M. W., Ojeda-González, A., & De La Luz, V. (2020).
 New Metric for Minimum Variance Analysis Validation in the Study of Interplanetary
 Magnetic Clouds. *Solar Physics*, 295(3), 45. https://doi.org/10.1007/s11207-020-01610-6
- Russell, C. T., Khurana, K. K., Huddleston, D. E., & Kivelson, M. G. (1998). Localized
 reconnection in the near jovian magnetotail. *Science*, 280(5366), 1061–1064.
 https://doi.org/10.1126/science.280.5366.1061
- Sarkango, Y., Jia, X., & Toth, G. (2019). Global MHD simulations of the Response of Jupiter's
 Magnetosphere and Ionosphere to Changes in the Solar Wind and IMF. *Journal of Geophysical Research: Space Physics*, 124(7), 5317–5341.
- 1032 https://doi.org/10.1029/2019JA026787
- Sarkango, Y., Slavin, J. A., Jia, X., DiBraccio, G. A., Gershman, D. J., Connerney, J. E. P., et al.
 (2021). Juno Observations of Ion-Inertial Scale Flux Ropes in the Jovian Magnetotail.
 Geophysical Research Letters, 48(2). https://doi.org/10.1029/2020GL089721
- Slavin, J. A., Baker, D. N., Craven, J. D., Elphic, R. C., Fairfield, D. H., Frank, L. A., et al.
 (1989). CDAW 8 observations of plasmoid signatures in the geomagnetic tail: An
 assessment. *Journal of Geophysical Research*, 94(A11), 15153.
- 1039 https://doi.org/10.1029/ja094ia11p15153
- 1040 Slavin, J. A., Lepping, R. P., Gjerloev, J., Fairfield, D. H., Hesse, M., Owen, C. J., et al. (2003).

- 1041Geotail observations of magnetic flux ropes in the plasma sheet. Journal of Geophysical1042Research: Space Physics, 108(A1). https://doi.org/10.1029/2002JA009557
- Smith, A. W., Slavin, J. A., Jackman, C. M., Fear, R. C., Poh, G. K., DiBraccio, G. A., et al.
 (2017). Automated force-free flux rope identification. *Journal of Geophysical Research: Space Physics*, *122*(1), 780–791. https://doi.org/10.1002/2016JA022994
- Smith, A. W., Slavin, J. A., Jackman, C. M., Poh, G. K., & Fear, R. C. (2017). Flux ropes in the
 Hermean magnetotail: Distribution, properties, and formation. *Journal of Geophysical Research: Space Physics*, *122*(8), 8136–8153. https://doi.org/10.1002/2017JA024295
- Sonnerup, B. U. Ö., & Cahill, L. J. (1967). Magnetopause structure and attitude from Explorer
 12 observations. *Journal of Geophysical Research*, 72(1), 171.
 https://doi.org/10.1029/jz072i001p00171
- Taylor, J. B. (1974). Relaxation of toroidal plasma and generation of reverse magnetic fields.
 Physical Review Letters, 33(19), 1139–1141. https://doi.org/10.1103/PhysRevLett.33.1139
- Vaivads, A., Khotyaintsev, Y. V., Retinò, A., Fu, H. S., Kronberg, E. A., & Daly, P. W. (2021).
 Cluster Observations of Energetic Electron Acceleration Within Earthward Reconnection
 Jet and Associated Magnetic Flux Rope. *Journal of Geophysical Research: Space Physics*,
 126(8), 1–13. https://doi.org/10.1029/2021ja029545
- Vasyliunas, V. M. (1983). Plasma distribution and flow. In *Physics of the Jovian Magnetosphere* (pp. 395–453). https://doi.org/10.1017/cbo9780511564574.013
- Vogt, M. F., Kivelson, M. G., Khurana, K. K., Joy, S. P., & Walker, R. J. (2010). Reconnection
 and flows in the Jovian magnetotail as inferred from magnetometer observations. *Journal of Geophysical Research: Space Physics*, *115*(6). https://doi.org/10.1029/2009JA015098
- Vogt, M. F., Jackman, C. M., Slavin, J. A., Bunce, E. J., Cowley, S. W. H., Kivelson, M. G., &
 Khurana, K. K. (2014). Structure and statistical properties of plasmoids in Jupiter's
 magnetotail. *Journal of Geophysical Research A: Space Physics*, *119*(2), 821–843.
 https://doi.org/10.1002/2013JA019393
- 1067 Vogt, M. F., Bunce, E. J., Nichols, J. D., Clarke, J. T., & Kurth, W. S. (2017). Long-Term
 1068 Variability of Jupiter's Magnetodisk and Implications for the Aurora. *Journal of*1069 *Geophysical Research: Space Physics*, *122*(12), 12,090-12,110.
 1070 https://doi.org/10.1002/2017JA024066
- 1071 Vogt, M. F., Connerney, J. E. P., DiBraccio, G. A., Wilson, R. J., Thomsen, M. F., Ebert, R. W.,
 1072 et al. (2020). Magnetotail Reconnection at Jupiter: A Survey of Juno Magnetic Field
 1073 Observations. *Journal of Geophysical Research: Space Physics*, *125*(3).
 1074 https://doi.org/10.1029/2019JA027486
- Woch, J., Krupp, N., Lagg, A., Wilken, B., Livi, S., & Williams, D. J. (1998). Quasi-periodic
 modulations of the Jovian magnetotail. *Geophysical Research Letters*, 25(8), 1253–1256.
 https://doi.org/10.1029/98GL00861
- Woch, J., Krupp, N., Khurana, K. K., Kivelson, M. G., Roux, A., Perraut, S., et al. (1999).
 Plasma sheet dynamics in the Jovian magnetotail: Signatures for substorm-like processes ?
 Geophysical Research Letters, 26(14), 2137–2140. https://doi.org/10.1029/1999GL900493
- 1081 Woch, J., Krupp, N., & Lagg, A. (2002). Particle bursts in the Jovian magnetosphere: Evidence
- 1082for a near-Jupiter neutral line. Geophysical Research Letters, 29(7), 42-1-42-4.1083https://doi.org/10.1029/2001GL014080
- Yao, Z. H., Grodent, D., Kurth, W. S., Clark, G., Mauk, B. H., Kimura, T., et al. (2019). On the
 Relation Between Jovian Aurorae and the Loading/Unloading of the Magnetic Flux:
 Simultaneous Measurements From Juno, Hubble Space Telescope, and Hisaki. *Geophysical Research Letters*, 46(21), 11632–11641. https://doi.org/10.1029/2019GL084201
- Yao, Z. H., Bonfond, B., Clark, G., Grodent, D., Dunn, W. R., Vogt, M. F., et al. (2020).
 Reconnection- and Dipolarization-Driven Auroral Dawn Storms and Injections. *Journal of Geophysical Research: Space Physics*, *125*(8). https://doi.org/10.1029/2019JA027663
- Zhong, Z. H., Zhou, M., Tang, R. X., Deng, X. H., Turner, D. L., Cohen, I. J., et al. (2020).
 Direct Evidence for Electron Acceleration Within Ion-Scale Flux Rope. *Geophysical Research Letters*, 47(1), e2019GL085141. https://doi.org/10.1029/2019GL085141

1094



Journal of Geophysical Research – Space Physics

Supporting Information for

Properties of Ion-Inertial Scale Plasmoids Observed by the Juno Spacecraft in the Jovian Magnetotail

Yash Sarkango¹, James A. Slavin¹, Xianzhe Jia¹, Gina A. DiBraccio², George B. Clark³, Weijie Sun¹, Barry H. Mauk³, William S. Kurth⁴, and George B. Hospodarsky⁴

¹University of Michigan - Ann Arbor, USA., ²NASA Goddard Space Flight Center, USA., ³John Hopkins University – Applied Physics Laboratory, USA, ⁴University of Iowa, USA

Contents of this file

Table S1 Figures S1-S120

Introduction

- This file contains a list of 87 plasmoids which were identified in the Jovian magnetotail. (Table S1)
- Magnetic field and Waves spectra, along with the results of MVA analysis are shown for each event. (Figures S1 S87)
- Energetic particle spectra is shown for selective events for which high cadence JEDI data was available. (Figures S88 S120)

| | | | Reversal | Plasmoid | Duration | < B > | ΔB_{θ} |
|---------|-------------------|-------------------|------------|-----------|----------|---------|---------------------|
| Event # | Start | End | Sense | Туре | (s) | (nT) | (nT) |
| 1 | 2017-075T23:57:09 | 2017-075T23:58:14 | tailward | o-line | 65 | 1.16 | 2.07 |
| 2 | 2017-075T09:55:25 | 2017-075T09:56:27 | tailward | flux-rope | 62 | 3.85 | 3.84 |
| 3 | 2017-075T09:35:23 | 2017-075T09:36:08 | planetward | o-line | 45 | 2.13 | 4.79 |
| 4 | 2017-076T06:44:50 | 2017-076T06:46:05 | planetward | flux-rope | 75 | 4.13 | 6.82 |
| 5 | 2017-076T15:56:33 | 2017-076T15:56:58 | planetward | o-line | 25 | 2.34 | 2.93 |
| 6 | 2017-076T00:09:54 | 2017-076T00:11:45 | planetward | flux-rope | 111 | 2.58 | 3.64 |
| 7 | 2017-077T02:57:16 | 2017-077T02:59:55 | tailward | flux-rope | 159 | 2.22 | 3.16 |
| 8 | 2017-078T18:31:53 | 2017-078T18:32:15 | tailward | o-line | 22 | 1.63 | 3.48 |
| 9 | 2017-083T13:01:29 | 2017-083T13:02:02 | tailward | o-line | 33 | 2.87 | 8.45 |
| 10 | 2017-096T03:20:06 | 2017-096T03:21:31 | planetward | o-line | 85 | 2.86 | 5.19 |
| 11 | 2017-096T03:16:48 | 2017-096T03:17:12 | tailward | o-line | 24 | 1.62 | 2.05 |
| 12 | 2017-129T23:24:18 | 2017-129T23:26:02 | planetward | flux-rope | 104 | 2.10 | 2.92 |
| 13 | 2017-132T05:23:39 | 2017-132T05:25:09 | tailward | flux-rope | 90 | 3.26 | 2.09 |
| 14 | 2017-136T02:37:05 | 2017-136T02:41:32 | tailward | o-line | 267 | 1.93 | 4.58 |
| 15 | 2017-150T18:31:35 | 2017-150T18:31:51 | tailward | o-line | 16 | 1.63 | 3.22 |
| 16 | 2017-186T04:22:10 | 2017-186T04:24:32 | tailward | flux-rope | 142 | 2.25 | 2.49 |
| 17 | 2017-189T17:56:00 | 2017-189T17:56:15 | tailward | flux-rope | 15 | 1.06 | 2.31 |
| 18 | 2017-234T07:44:32 | 2017-234T07:45:17 | tailward | o-line | 45 | 2.30 | 3.68 |
| 19 | 2017-234T16:41:50 | 2017-234T16:42:51 | planetward | o-line | 61 | 1.95 | 5.67 |
| 20 | 2017-235T10:11:28 | 2017-235T10:13:45 | tailward | o-line | 137 | 1.18 | 3.00 |
| 21 | 2017-235T02:26:06 | 2017-235T02:26:41 | tailward | o-line | 35 | 2.17 | 2.49 |
| 22 | 2017-236T20:21:14 | 2017-236T20:21:34 | planetward | flux-rope | 20 | 3.67 | 3.50 |
| 23 | 2017-239T12:46:13 | 2017-239T12:46:40 | tailward | flux-rope | 27 | 2.00 | 2.31 |
| 24 | 2017-239T08:31:58 | 2017-239T08:32:27 | tailward | flux-rope | 29 | 2.67 | 3.23 |
| 25 | 2017-239T22:54:38 | 2017-239T22:55:07 | tailward | o-line | 29 | 1.63 | 2.42 |
| 26 | 2017-239T08:31:05 | 2017-239T08:31:43 | planetward | o-line | 38 | 1.73 | 3.33 |
| 27 | 2017-241T03:48:43 | 2017-241T03:49:08 | tailward | o-line | 25 | 5.23 | 11.12 |
| 28 | 2017-242T08:56:02 | 2017-242T08:56:15 | planetward | o-line | 13 | 2.19 | 2.48 |
| 29 | 2017-287T13:20:33 | 2017-287T13:23:05 | tailward | o-line | 152 | 2.97 | 7.91 |

| 30 | 2017-287T07:07:40 | 2017-287T07:09:11 | tailward | o-line | 91 | 1.71 | 4.29 |
|----|-------------------|-------------------|------------|-----------|-----|------|------|
| 31 | 2017-289T09:00:42 | 2017-289T09:02:16 | tailward | o-line | 94 | 4.96 | 2.22 |
| 32 | 2017-289T09:11:43 | 2017-289T09:12:46 | planetward | o-line | 63 | 2.28 | 4.18 |
| 33 | 2017-291T16:09:44 | 2017-291T16:10:27 | tailward | o-line | 43 | 1.06 | 2.18 |
| 34 | 2017-291T15:48:32 | 2017-291T15:49:04 | tailward | flux-rope | 32 | 3.58 | 4.85 |
| 35 | 2017-291T19:27:19 | 2017-291T19:27:35 | planetward | o-line | 16 | 1.93 | 2.20 |
| 36 | 2017-291T10:22:01 | 2017-291T10:23:45 | tailward | o-line | 104 | 1.87 | 3.29 |
| 37 | 2017-293T16:36:31 | 2017-293T16:37:37 | planetward | o-line | 66 | 1.88 | 4.29 |
| 38 | 2017-294T18:13:43 | 2017-294T18:14:09 | planetward | o-line | 26 | 2.31 | 3.68 |
| 39 | 2017-338T20:58:39 | 2017-338T20:59:11 | planetward | o-line | 32 | 1.67 | 2.67 |
| 40 | 2017-338T21:49:25 | 2017-338T21:51:52 | tailward | flux-rope | 147 | 4.54 | 5.87 |
| 41 | 2017-338T21:33:00 | 2017-338T21:33:18 | planetward | flux-rope | 18 | 1.35 | 3.53 |
| 42 | 2017-338T20:55:15 | 2017-338T20:56:04 | tailward | o-line | 49 | 3.51 | 5.32 |
| 43 | 2017-338T21:43:31 | 2017-338T21:43:47 | tailward | o-line | 16 | 2.89 | 2.12 |
| 44 | 2017-340T13:13:23 | 2017-340T13:13:36 | planetward | o-line | 13 | 0.94 | 2.75 |
| 45 | 2017-342T03:00:19 | 2017-342T03:02:21 | tailward | o-line | 122 | 1.64 | 2.13 |
| 46 | 2017-345T00:02:43 | 2017-345T00:04:13 | planetward | o-line | 90 | 1.47 | 2.30 |
| 47 | 2018-026T20:15:53 | 2018-026T20:17:51 | tailward | flux-rope | 118 | 3.76 | 3.82 |
| 48 | 2018-027T15:12:45 | 2018-027T15:16:54 | tailward | o-line | 249 | 2.31 | 2.56 |
| 49 | 2018-027T16:26:19 | 2018-027T16:26:51 | tailward | o-line | 32 | 1.26 | 2.19 |
| 50 | 2018-027T16:16:08 | 2018-027T16:16:30 | tailward | o-line | 22 | 2.19 | 2.25 |
| 51 | 2018-028T12:26:21 | 2018-028T12:30:13 | tailward | flux-rope | 232 | 2.04 | 2.05 |
| 52 | 2018-029T18:17:07 | 2018-029T18:20:05 | planetward | o-line | 178 | 1.26 | 3.42 |
| 53 | 2018-030T13:38:18 | 2018-030T13:39:12 | planetward | o-line | 54 | 1.67 | 3.27 |
| 54 | 2018-030T21:26:31 | 2018-030T21:27:11 | tailward | flux-rope | 40 | 1.84 | 2.16 |
| 55 | 2018-030T01:27:08 | 2018-030T01:28:41 | tailward | o-line | 93 | 1.38 | 3.12 |
| 56 | 2018-032T22:47:36 | 2018-032T22:48:27 | planetward | flux-rope | 51 | 2.68 | 2.62 |
| 57 | 2018-033T11:30:05 | 2018-033T11:31:06 | tailward | o-line | 61 | 1.69 | 3.09 |
| 58 | 2018-036T10:36:44 | 2018-036T10:36:59 | planetward | o-line | 15 | 1.94 | 4.82 |
| 59 | 2018-083T11:21:16 | 2018-083T11:21:44 | tailward | o-line | 28 | 1.34 | 3.33 |
| 60 | 2018-085T11:18:14 | 2018-085T11:19:09 | tailward | flux-rope | 55 | 2.59 | 4.27 |
| 61 | 2018-139T18:52:52 | 2018-139T18:53:30 | planetward | o-line | 38 | 4.15 | 6.16 |

| 62 | 2018-140T14:27:52 | 2018-140T14:28:15 | planetward | o-line | 23 | 2.42 | 3.25 |
|----|-------------------|-------------------|------------|-----------|-----|------|------|
| 63 | 2018-142T01:27:49 | 2018-142T01:28:18 | planetward | o-line | 29 | 4.73 | 6.55 |
| 64 | 2018-188T23:12:51 | 2018-188T23:14:44 | tailward | flux-rope | 113 | 1.86 | 2.27 |
| 65 | 2018-190T14:48:13 | 2018-190T14:48:41 | planetward | flux-rope | 28 | 1.54 | 3.01 |
| 66 | 2018-191T10:23:12 | 2018-191T10:23:35 | planetward | o-line | 23 | 5.36 | 4.88 |
| 67 | 2018-245T15:02:56 | 2018-245T15:03:20 | planetward | o-line | 24 | 0.89 | 2.21 |
| 68 | 2018-297T14:55:14 | 2018-297T14:58:31 | tailward | o-line | 197 | 2.34 | 3.60 |
| 69 | 2018-297T17:05:54 | 2018-297T17:06:24 | planetward | o-line | 30 | 1.11 | 2.58 |
| 70 | 2018-298T01:02:35 | 2018-298T01:03:48 | planetward | flux-rope | 73 | 3.59 | 5.47 |
| 71 | 2018-298T10:48:22 | 2018-298T10:48:53 | planetward | o-line | 31 | 1.84 | 4.51 |
| 72 | 2018-350T04:57:17 | 2018-350T04:57:30 | tailward | flux-rope | 13 | 2.41 | 2.72 |
| 73 | 2018-351T10:04:44 | 2018-351T10:05:03 | planetward | flux-rope | 19 | 1.93 | 2.08 |
| 74 | 2018-352T04:27:41 | 2018-352T04:28:18 | planetward | flux-rope | 37 | 4.27 | 9.31 |
| 75 | 2018-353T14:00:16 | 2018-353T14:02:33 | tailward | o-line | 137 | 2.71 | 4.46 |
| 76 | 2018-353T13:59:12 | 2018-353T13:59:40 | tailward | o-line | 28 | 1.95 | 5.25 |
| 77 | 2019-089T20:35:44 | 2019-089T20:36:34 | tailward | flux-rope | 50 | 4.46 | 2.59 |
| 78 | 2019-090T05:24:56 | 2019-090T05:27:20 | tailward | flux-rope | 144 | 6.35 | 3.12 |
| 79 | 2019-091T21:33:28 | 2019-091T21:35:36 | tailward | flux-rope | 128 | 3.56 | 2.81 |
| 80 | 2019-091T00:47:51 | 2019-091T00:50:08 | tailward | flux-rope | 137 | 2.02 | 3.63 |
| 81 | 2019-092T17:30:38 | 2019-092T17:32:03 | planetward | o-line | 85 | 2.33 | 4.62 |
| 82 | 2019-093T20:50:14 | 2019-093T20:50:31 | planetward | o-line | 17 | 2.04 | 5.95 |
| 83 | 2019-093T01:38:18 | 2019-093T01:39:10 | planetward | o-line | 52 | 2.72 | 3.23 |
| 84 | 2019-144T19:03:24 | 2019-144T19:03:37 | planetward | flux-rope | 13 | 2.63 | 4.10 |
| 85 | 2019-199T17:28:05 | 2019-199T17:28:22 | planetward | flux-rope | 17 | 1.72 | 3.52 |
| 86 | 2019-249T19:23:52 | 2019-249T19:24:35 | planetward | o-line | 43 | 1.09 | 2.30 |
| 87 | 2019-301T22:13:15 | 2019-301T22:14:05 | planetward | o-line | 50 | 2.29 | 4.75 |

Table S1. List of plasmoids identified by the algorithm.



Figure S1



Figure S2



Figure S3



Figure S4



Figure S5



Figure S6



Figure S7



Figure S8



Figure S9



Figure S10



Figure S11



Figure S12



Figure S13



Figure S14



Figure S15



Figure S16



Figure S17



Figure S18



Figure S19



Figure S20



Figure S21



Figure S22



Figure S23



Figure S24



Figure S25



Figure S26



Figure S27



Figure S28



Figure S29



Figure S30



Figure S31


Figure S32



Figure S33



Figure S34



Figure S35



Figure S36



Figure S37



Figure S38



Figure S39



Figure S40



Figure S41



Figure S42



Figure S43



Figure S44



Figure S45



Figure S46



Figure S47



Figure S48



Figure S49



Figure S50



Figure S51



Figure S52



Figure S53



Figure S54



Figure S55



Figure S56



Figure S57



Figure S58



Figure S59



Figure S60



Figure S61



Figure S62



Figure S63



Figure S64



Figure S65



Figure S66



Figure S67


Figure S68



Figure S69



Figure S70



Figure S71



Figure S72



Figure S73



Figure S74



Figure S75



Figure S76



Figure S77



Figure S78



Figure S79



Figure S80



Figure S81



Figure S82



Figure S83



Figure S84



Figure S85



Figure S86



Figure S87



2017-03-16, DOY=075



2017-03-16, DOY=075



2017-03-16, DOY=075



2017-03-17, DOY=076



2017-03-17, DOY=076



2017-03-17, DOY=076



2017-03-18, DOY=077



2017-03-19, DOY=078



2017-03-24, DOY=083



2017-04-06, DOY=096



2017-04-06, DOY=096



2017-05-09, DOY=129



2017-05-12, DOY=132



2017-08-22, DOY=234



2017-08-22, DOY=234



2017-08-24, DOY=236


2017-08-27, DOY=239



2017-08-27, DOY=239



2017-08-30, DOY=242



2018-05-20, DOY=140



2018-10-25, DOY=298

Figure S108



2018-10-25, DOY=298



2018-12-17, DOY=351



2018-12-18, DOY=352

Figure S111



2018-12-19, DOY=353



2018-12-19, DOY=353



2019-03-31, DOY=090



2019-04-01, DOY=091



2019-04-02, DOY=092

Figure S116



2019-04-03, DOY=093



2019-05-24, DOY=144



2019-09-06, DOY=249



2019-10-28, DOY=301