Examining Parameterizations of Potential Temperature Variance Across Varied Landscapes for use in Earth System Models

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Abstract

Earth system models (ESMs) and mesoscale models have come to employ increasingly complex parameterization schemes for the atmospheric boundary layer (ABL), requiring surface boundary conditions for numerous higher order turbulence statistics. Of particular interest is the potential temperature variance (PTV), which is used not only as a boundary condition itself but also to close boundary conditions of other statistics. The existing schemes in ESMs largely rely on the assumptions of Monin-Obukhov similarity theory (MOST), and are not necessarily applicable over complex and heterogeneous surfaces where large scale circulations and roughness sub-layer effects may cause deviations from MOST. The National Ecological Network (NEON) is used here to evaluate existing parameterizations for the surface boundary of PTV, note key deficiencies, and explore possible remedies. The results indicate that existing schemes are acceptable over a variety of surface conditions provided the analysis of a priori filters out low frequency variability not associated with turbulent time scales. There was, however, significant inter-site variability in observed similarity constants and a significant bias when compared to the textbook values of these parameters. Existing models displayed the poorest performance over heterogeneous sites, and rough landscapes. Attempts to use canopy structure and surface roughness characteristics to improve the results confirmed a relation between these variables and PTV, but failed to significantly improve the predictive power of the models. The results did not find strong evidence indicating that large scale circulations caused substantial deviations from textbook models, although additional analysis is required to assess their full impacts.

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Key Points:

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8	•	Models of potential temperature variance in the surface layer based on similar-
9		ity theory were evaluated using data from 39 varied sites
10	•	Existing schemes perform well across most surfaces, although the data shows a
11		significant bias in the values of the similarity constants
12	•	Canopy structure and surface heterogeneity drive a large portion of inter-site vari
13		ability in model performance

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14 Abstract

Earth system models (ESMs) and mesoscale models have come to employ increas-15 ingly complex parameterization schemes for the atmospheric boundary layer (ABL), re-16 quiring surface boundary conditions for numerous higher order turbulence statistics. Of 17 particular interest is the potential temperature variance (PTV), which is used not only 18 as a boundary condition itself but also to close boundary conditions of other statistics. 19 The existing schemes in ESMs largely rely on the assumptions of Monin-Obukhov sim-20 ilarity theory (MOST), and are not necessarily applicable over complex and heteroge-21 22 neous surfaces where large scale circulations and roughness sub-layer effects may cause deviations from MOST. The National Ecological Network (NEON) is used here to eval-23 uate existing parameterizations for the surface boundary of PTV, note key deficiencies, 24 and explore possible remedies. The results indicate that existing schemes are acceptable 25 over a variety of surface conditions provided the analysis of a priori filters out low fre-26 quency variability not associated with turbulent time scales. There was, however, sig-27 nificant inter-site variability in observed similarity constants and a significant bias when 28 compared to the textbook values of these parameters. Existing models displayed the poor-29 est performance over heterogeneous sites, and rough landscapes. Attempts to use canopy 30 structure and surface roughness characteristics to improve the results confirmed a rela-31 tion between these variables and PTV, but failed to significantly improve the predictive 32 power of the models. The results did not find strong evidence indicating that large scale 33 circulations caused substantial deviations from textbook models, although additional anal-34 ysis is required to assess their full impacts. 35

³⁶ Plain Language Summary

Modern models of the lower atmosphere, which are used to analyze climate change 37 and weather, resolve increasingly complex characteristics of the turbulence in the atmo-38 sphere. An estimate for the value of many of these characteristics at the land surface is 39 required to set boundary conditions for these models. An important boundary condi-40 tion is the variance of very small temperature fluctuations that occur in the atmosphere 41 due to turbulence. Currently, model estimates for these values assume the surface is flat 42 and its characteristics do not change in space, which doesn't represent many of the con-43 ditions we wish to model over the earth. In addition, existing studies tend to only an-44 alyze data from a small number of locations. We analyzed data from a network of 39 sites 45 and found that the current estimates work fairly well across a large variety of conditions, 46 but that there is a bias in the constants often used and there are notable differences over 47 forests, complex surfaces, and heterogeneous terrain. There is a clear relationship be-48 tween surface characteristics such as tree canopy height and performance of the model, 49 however it was not clear enough to improve our ability to predict the surface boundary 50 condition. 51

52 1 Introduction

The atmospheric boundary layer (ABL) plays a fundamental role in the climate 53 system due to its significance in bridging land surface fluxes of heat and water vapor to 54 convection and cloud formation (Siqueira et al., 2009; Huang & Margulis, 2010; Garratt, 55 1992). The ABL is characterized by the coexistence of mechanically and thermally gen-56 erated turbulence, which regulate mixing and transport properties and exchanges between 57 the land surface and the lower atmosphere. The variances of turbulent quantities are of 58 particular interest due to their emerging role in state-of-the-science Earth System Mod-59 els (ESMs) and numerical weather prediction. They have accordingly received attention 60 in the literature, although most of these studies have focused on the velocity variances. 61 Comparatively few examine the potential temperature variance (PTV) and those that 62 do often focus on flat homogeneous terrain (Albertson et al., 1995; Asanuma & Brut-63

saert, 1999; G. G. Katul & Hsieh, 1999; Mironov & Sullivan, 2016; van de Boer et al., 64 2014; Maronga & Reuder, 2017; Otić et al., 2005; Antonia et al., 1981; D. Li et al., 2016; 65 Monji, 1973; Champagne et al., 1977; Kiely et al., 1996). Traditional boundary layer schemes 66 in ESMs employed first-order or 1.5-order closure schemes (Cohen et al., 2015; Lock et al., 2000), although increasingly many higher order schemes that resolve PTV prognos-68 tically throughout the ABL are now in use, such as the Cloud Layers Unified by Bino-69 mials (CLUBB) scheme in the Community Earth System Model (CESM) and the En-70 ergy Exascale Earth System Model (E3SM) (Larson, 2017), the Mellor-Yamada-Nakanishi-71 Niino model (MYNN) implemented in the meso-scale Weather Research and Forecast-72 ing model and the Model for Interdisciplinary Research on Climate (MIROC) (Nakanishi 73 & Niino, 2009), and the intermediately prognostic higher-order turbulencce closure (IPHOC) 74 implemented in the Community Atmosphere Model, version 5 (Cheng & Xu, 2015). How-75 ever, less attention has been placed on the surface boundary condition of PTV of these 76 schemes despite their use in the aforementioned models and the fact that many higher 77 order terms are closed based this temperature variance. 78

The specification of the lower boundary conditions in such schemes utilize Monin-79 Obukhov Similarity Theory (MOST) that rests on the assumptions of stationary and pla-80 nar homogeneous, high Reynolds number flow in the absence of subsidence (Monin & 81 Obukhov, 1954). For these idealized conditions, the turbulent fluxes are assumed to be 82 invariant with distance from the boundary and all flow statistics can be reduced to a set 83 of universal curves that vary with the atmospheric stability parameter (Foken, 2006). 84 Currently, one of two parameterization schemes, both consistent with MOST (Tillman, 85 1972; J. Wyngaard & Coté, 1971) for unstable atmospheric conditions are used in ESMs 86 and are given by 87

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 $\frac{\overline{\theta'^2}}{T_{+}^2} = a(1 - b\zeta)^{-2/3},$

$$\frac{\overline{\theta'^2}}{\pi^2} = C_1(-\zeta)^{-2/3},\tag{2}$$

(1)

over

where
$$\theta'$$
 is the fluctuating potential temperature, overline indicates time-averaging over
a period that is sufficiently long to reliably capture the ensemble statistics of turbulence

but short enough relative to variations in the mean state of the ABL, a, b, and C_1 are 93 similarity constants, ζ is the atmospheric stability parameter defined as 94

$$\zeta = \frac{z - z_d}{L},\tag{3}$$

with z_d being the zero-plane displacement height, z is the measurement height and L is 96 the Obukhov length (Obukhov, 1946) given by 97

$$L = -\frac{u_*^3 \overline{\theta_v}}{kg \overline{w'\theta'}},\tag{4}$$

where k = 0.4 is the von Kármán constant, q is the gravitational acceleration, u_* is the 99 friction velocity, $\overline{\theta_v}$ is the mean virtual potential temperature, $\overline{w'\theta'}$ is the kinematic tur-100 bulent sensible heat flux, and w' is the turbulent vertical velocity. Unstable atmospheric 101 stability conditions is defined by $\zeta < 0$ whereas near-neutral atmospheric stability con-102 ditions occurs when $|\zeta| < 0.05$. The T_* is the non-dimensional temperature scale de-103 fined as 104

$$T_* = \frac{\overline{w'\theta'}}{u_*}.$$
(5)

Equations (1) and (2) converge as near-convective conditions ($-\zeta \gg 1$) are ap-106 proached resulting in $ab^{-2/3} = C_1$. For these conditions, the turbulent heat flux can 107

be linked to $\sigma_T = \sqrt{\overline{\theta'^2}}$ through the well known flux-variance expression (Tillman, 1972)

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$$\overline{w'\theta'} = C_1^{-3} [kg(z - z_d)]^{1/2} \overline{\theta_v}^{-1/2} \sigma_T^{3/2}.$$

(6)

This expression suggests that sensible heat only depends on σ_T and $(z - z_d)$ independent of u_* as expected when convective conditions are approached. For near neutral conditions with $u_* > 0$ and $\zeta \to 0$, equation (1) ensures $\frac{\sigma_T}{T_*} \to \sqrt{a}$ whereas equation (2) suggests that σ_T is indeterminate by MOST. The two-third scaling is fixed for the purposes of this study - a reasonable assumption as it matches the logical, dimensional limits of free convection.

The 'textbook' similarity constants estimated in the literature are a = 4, b = 8.3, 116 and $C_1 = 0.95$. These values were initially derived from experiments over flat, homo-117 geneous wheat stubble in Kansas and confirmed by other studies over similarly homo-118 geneous and largely flat terrain (Tillman, 1972; J. Wyngaard & Coté, 1971; J. C. Wyn-119 gaard & Coté, 1974; Andre et al., 1978; Albertson et al., 1995; Haugen et al., 1971; Monji, 120 1973). This lends some support to their supposed universal character. However, MOST 121 is not readily generalizable for application in ESMs over more realistic landscapes, tall 122 forests and a variety of atmospheric conditions such as those associated with significant 123 entrainment and mesoscale phenomenon (Kroon & de Bruin, 1995; Asanuma & Brut-124 saert, 1999; Lloyd et al., 1991; Hang et al., 2018; Mcnaughton, 2006; van de Boer et al., 125 2014; Wilson, 2008; Harman, 2012; Brunet, 2020; Q. Li et al., 2018). Previous literature 126 examining the scaling relation between ζ and non-dimensional flow statistics has focused 127 on conditions that satisfy the assumptions of flat uniform surfaces so that the univer-128 sal character suggested by MOST can be readily tested (Kader & Yaglom, 1990). How-129 ever, comparatively less research has been carried out over non-idealized terrain. These 130 few studies have found that MOST derived functions may not hold over surfaces such 131 as sparse and open canopies and heterogenous surfaces (Lee, 2009; Kroon & de Bruin, 132 1995; van de Boer et al., 2014; Hang et al., 2018; Detto et al., 2008). Few studies have 133 consistently examined PTV across a wide variety of land cover types (G. Katul et al., 134 1995). The latter study suggested that local similarity may still hold (i.e. a local T_* and 135 L can explain the mathematical form of PTV) provided the similarity coefficients (e.g. 136 C_1) are allowed to vary with land cover type. Despite these issues raised, the use of MOST 137 scaling over various landscapes is widespread in ESMs that require it (Nakanishi & Ni-138 ino, 2009; Larson, 2017; Zhao et al., 2018; Golaz et al., 2019, 2002; Cheng & Xu, 2015). 139 To explore PTV in the atmospheric surface layer across differing landscapes and a wide 140 range of atmospheric conditions, observations covering many ecosystems and canopy struc-141 tures with appropriate parameterizations are becoming necessary and motivates the present 142 work. 143

Since these parameterizations were developed, there has been a significant growth 144 in the availability of data across differing surfaces that can be used to re-examine MOST 145 parameterizations. One example is the National Ecological Observation Network (NEON). 146 NEON is a continent-scale network where high frequency (20 Hz) velocity and air tem-147 perature fluctuations are sampled in a consistent manner (i.e. same instrumentation, rel-148 ative heights, pre- and post-processing algorithms, etc..) over 39 sites that vary in cli-149 mate and land-cover across the United States. Hence, the NEON high frequency data 150 set offers a unique opportunity to explore these similarity relations over many land cover 151 types (ideal and non-ideal) and ζ conditions. Using this information, it is possible to ex-152 plore validity and modifications to the traditional MOST PTV parameterizations. The 153 initial focus spans near-neutral to unstable stratification ($\zeta < 0$), where the turbulence 154 is fully developed. Stably stratified conditions are characterized by a shallow boundary 155 layer depth and are infected with numerous non-turbulent phenomena that will require 156 a separate investigation that is better kept for a future study. 157

¹⁵⁸ With this large data set, the time is ripe to revisit and reevaluate traditional schemes ¹⁵⁹ for PTV in light of these contemporary needs of ESM. In doing so, the focus is on two

deviations from the assumptions of MOST. The first is mesoscale phenomenon and outer-160 layer eddies that impinge onto the atmospheric surface layer, potentially introducing ad-161 ditional length scales not captured by ζ . The second is roughness sublayer effects, es-162 pecially over forests or other forms of structured heterogeneity, which is not included as 163 part of MOST. This study seeks to quantify the significance of the distortions from both 164 mesoscale and roughness sublayer effects on equations (1) and (2), and examine if such 165 distortions can be partly absorbed in the parameters a and b (or C_1). The approach that 166 follows takes advantage of the wealth of data provided by NEON as well as remotely sensed 167 sources, and the Random Forest (RF) method, which is a machine learning method able 168 to classify the significance of surrogate terms such as boundary layer height, land cover 169 type, canopy height, and other ancillary variables on $\theta^{\prime 2}/T_*^2$. 170

171 **2 Data**

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The core turbulence data are publicly available from NEON and includes the turbulence statistics ($\overline{w'\theta'}$, θ' , u_*), meteorological variables (mean temperature, humidity and wind speed), as well as site specific information (h_c , tower characteristics). Additional information from remotely sensed datasets and reanalysis data (see Table 1) colocated with the NEON site data are used in predictive models that seek to link environmental variables and land surface features to PTV.

2.1 National Ecological Observation Network

The NEON sites are located within the continental United States (CONUS), Alaska, 179 Hawaii and Puerto Rico. Sites are centrally managed and designed, which means that 180 sampling and post-processing high frequency data are consistent, and differences can be 181 attributed to site characteristics rather than differences in management, methods and 182 instrumentation as is the case for other locally managed flux tower networks such as FLUXNET 183 or AmeriFlux (Novick et al., 2018). Moreover, the high frequency time series spanning 184 several years are publicly available thus enabling the determination of variances and heat 185 fluxes in a coherent manner when post-processed. Sites are also spread across different 186 ecological domains to ensure coverage of the different landscapes and ecosystems in North 187 America. 188

Each site includes a full suite of meteorological instrumentation, eddy covariance 189 measurements from a CSAT-3 sonic anemometer recording a 20 Hz and time averaged 190 to 30 min, and mean wind profiles throughout the canopy and above it, compiled into 191 one dataset (National Ecological Observatory Network (NEON), 2021). Data is exam-192 ined from first availability at each site, which varies by tower but is generally in mid 2017, 193 to May 2020. Towers at sites with a canopy less than three meters are designed to be 194 8m tall, whereas towers at sites with a canopy greater than three meters are designed 195 to have a height corresponding to $z_d + 4(h_c - z_d)$, with canopy height h_c , to ensure that 196 the turbulence exchange assembly samples largely above the momentum roughness layer 197 (Metzger et al., 2019). In addition, detailed canopy structure at each site is acquired through 198 near-annual airborne remote sensing surveys with discrete and full waveform LiDAR. Soil, 199 vegetative and meteorological characteristics are described and continuously collected 200 when appropriate at each site. Only the 39 CONUS sites are included in this analysis. 201

For illustrative purposes, eight representative sites were selected as examples of site 202 level differences throughout the study. Wind River Experimental Forest (WREF) - a tall 203 evergreen forest in the Pacific Northwest, Northern Great Plains Research Laboratory 204 (NOGP) - a flat grassland site in North Dakota, Bartlett Experimental Forest (BART) 205 - a mixed deciduous evergreen forest in New England, Soaproot Saddle (SOAP) - a conifer 206 forest with complex terrain in the Sierra Nevada mountains, Oak Ridge National Lab 207 (ORNL) - a deciduous forest with some pine in Appalachia, Santa Rita Experimental 208 Range (SRER) - a semiarid scrub environment site in the Sonoran Desert, Konza Prairie 209

Variable	Spatial Resolution	Temporal Resolution	Source
LAI	250 m	8 days	MODIS
f_{veg}	$250 \mathrm{~m}$	1 year	MODIS
f_{tree}	$250 \mathrm{~m}$	1 year	MODIS
f_{bare}	$250 \mathrm{~m}$	1 year	MODIS
Land Cover	30 m	N/A	NLCD (Landsat)
BLH	$30 \mathrm{km}$	1 hour	ERA 5 Reanalysis
f_{cloud}	$30 \mathrm{km}$	1 hour	ERA 5 Reanalysis
CAPE	30 km	1 hour	ERA 5 Reanalysis

 Table 1.
 Summary table of the remotely sensed data and reanalysis products used in this

 project with their native spatial and temporal resolution as well as the source of the data

Biological Station (KONZ) - a pristine prairie site in Kansas, and Disney Wilderness Preserve (DSNY) - a wetland site in the headwaters of the everglades.

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2.2 Remotely Sensed Data

Three satellite remote sensing datasets are collocated with the NEON site data to complete a full coverage of vegetation and land cover at each site. This data, as well as reanalysis data discussed in section 2.3, are summarized in Table 1.

Two MODIS derived satellite products are used including MODIS Leaf Area In-216 dex (LAI) (Myneni et al., 2015), which is reported at about 250m resolution every 8 days. 217 The site is assigned the LAI of the grid cell in which the tower is contained, and linear 218 interpolation is used to determine LAI for days in which MODIS LAI is unavailable. MODIS 219 Vegetative Continuous Fields (VCF) (DiMiceli et al., 2015) includes measurements of veg-220 etation cover at about 250m resolution on a yearly basis, with linear interpolation used 221 to fill in gaps. Similar to MODIS LAI, MODIS VCF at each site for each point is assigned 222 based on the VCF of the grid cell in which the tower is contained. The VCF product de-223 tails low lying vegetation cover (f_{veg}) , tree cover (f_{tree}) , and bare soil coverage (f_{bare}) 224 percentages around each site. These products provide basic information about the veg-225 etation structure. 226

The National Land Cover Dataset (NLCD), a Landsat derived product defining the land cover at 30m pixels over CONUS (Jin et al., 2019), is the third remote dataset employed. Fractional coverage of each landcover type within a 250m radius from the tower location is computed for each site, as well as the dominant NLCD land cover type.

2.3 Reanalysis Data

ERA5 (Hersbach et al., 2018) is a reanalysis dataset that combines historical ob-232 servations and modelling results to generate hourly data of a variety of land surface and 233 atmospheric characteristics. For this analysis, the boundary layer height (BLH), total 234 cloud cover (f_{cloud}) and Convective Available Potential Energy (CAPE) are used to in-235 clude the impacts on mesoscale phenomenon using commonly reported variables in the 236 meteorological community. BLH is selected as the depth of the boundary layer is closely 237 related to thermal convection strength and the size of some circulations are closely re-238 lated to this value. f_{cloud} is used as it may also serve as a proxy for deep convection and 239 identification of cloudy conditions that could impact temperature statistics. CAPE is 240 chosen as it is related to updraft and general convection strength in the atmosphere. 241

$\mathbf{3}$ Methods

Turbulence statistics as directly acquired from NEON includes variance informa-243 tion from non-turbulence sources whereas models such as CLUBB focus on variances pro-244 duced by turbulent eddies. A filtering process is required to remove non-turbulent events 245 (and lack of stationarity) before they can be used for analysis. In addition, computed 246 z_d is required for the tower area, as values reported by NEON are suspect and represent 247 the physical characteristics of the entire ecological site rather than the local tower foot-248 print. These values are needed to assess the influence of surface roughness on the devel-249 250 opment of turbulence. One method of analysis to be used is the Random Forest (RF) method, which is employed to determine what physical and environmental characteris-251 tics are most significant for the development of variance without constraints imposed by 252 similarity theory and concomitant dimensional analysis. 253

3.1 Filtering

One of the key assumptions for MOST parameterizations is the stationarity of the 255 temperature time series. The data for the majority of atmospheric conditions at each 256 site are not strictly stationary. Any computed temperature variance value captures vari-257 ance associated with turbulent eddies and meso-scale disturbances as well as non-stationarity 258 found at transitions from night to day and vice-a-versa. To fulfill the requirement of solely 259 including PTV caused by turbulence as required by ESMs, a high pass filter with a cut-260 off time scale of 5 min is applied to the high frequency air temperature time series in the 261 Fourier domain. An example application of the high pass filter is featured in figure 1. 262 Time scales exceeding 5 minutes in the air temperature spectra are assumed to be not 263 associated with turbulent eddies produced by mechanical or buoyant production near 264 the surface. In fact, the choice of 5 minutes exceeds by at least one to two orders of mag-265 nitude measured peaks in the co-spectra of w' and θ' or the shear time scale $k(z-z_d)/u_*$ 266 linked with MOST. These events do not significantly impact turbulent sensible heat flux 267 but contribute appreciably to temperature variance. For some points the filtering resulted 268 in reductions of variance of up to 50%, however for the majority of the points the change 269 in variance was near negligible. Alternative filter cutoffs greater and less than 5 minutes 270 were examined and the variances were found to have only a very small sensitivity to the 271 exact cutoff value. The remainder of the analysis presented herein uses the filtered tem-272 perature variance. 273

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3.2 Canopy Structure Determination

While NEON does report site level h_c and z_d , which is directly estimated from h_c 275 and required for the analysis, these values appear to be reported as averages for the whole 276 ecological site and not the direct tower area, windshed, or source weight function. Re-277 ported z_d in particular appear to deviate significantly from experimentally derived val-278 ues at a number of sites, and since they are constant they do not reflect seasonal changes 279 in canopy, in particular over deciduous forest. z_d also is estimated simply as $z_d = (2/3)h_c$, 280 which while a good rule of thumb is not always accurate. As such, the z_d is estimated 281 here from measured mean wind profile data (u(z)) assuming a log wind profile and that 282 u^* is approximately constant with z (as required by MOST). We apply the following us-283 ing the top three points in the mean wind profile under near-neutral conditions so that 284 the stability correction terms can be ignored (Oke, 1987) 285

$$\frac{du}{dz} = \frac{u_*}{k(z-z_d)}.$$
(7)

The resulting heights are seasonally averaged at each site and then interpolated for use in the computation of ζ via (3). The resulting z_d , seen in figure 2, approximately follow the 2/3 relation reported in the literature (Garratt, 1992) over most sites. This relation, however, is less clear at a number of sites with short vegetation. This deviation may not



Figure 1. Illustration of filtering to reduce the effects of non-stationarity. Top row shows the raw unfiltered time series for one 30-minute run (left) and the same data after the high pass filter was employed where the nonlinear trend is removed (right). The second row shows the raw unfiltered data for another 30-minute run (left) and the same data after the high pass filter where the approximate linear trend is removed (right). Dotted lines are the actual data, and the solid line is the 1 minute average value. The bottom plot illustrates the change in over a one-month period at the ABBY site from the unfiltered NEON product and the filtered data



Figure 2. Comparison of NEON reported canopy height (h_c) and mean calculated zero plane displacement height from equation (7) (z_d) . Colored according to the dominant land cover type at each site. Dotted line represents $z_d = (2/3)h_c$ relations.

be surprising. A basis for the $z_d = (2/3)h_c$ relation stems from an exponential mean velocity profile characterized by an extinction coefficient $a_c > 1$ inside the canopy as derived from a constant mixing length hypothesis for the turbulent eddy diffusivity (Raupach & Thom, 1981). These arguments, when combined with the drag-force centroid method to estimate z_d for (i) constant drag coefficient and leaf area density and (ii) rigid, tall and dense canopy yield

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$$\frac{z_d}{h_c} = 1 - \frac{1}{2a_c} = 1 - \frac{1}{2} \frac{L_s}{h_c},\tag{8}$$

where $L_s = u/(du/dz)$ evaluated at $z = h_c$ is known as the canopy shear length scale. 298 For the flow near the canopy top to behave as 'mixing layers' requires an inflection point 299 in the mean velocity profile at $z/h_c = 1$ (Raupach et al., 1996). This condition leads 300 to a constraint on $1/2 < L_s/h_c < 1$ thereby bounding z_d/h_c to be between 1/2 and 301 3/4. All these assumptions (i.e. rigid, tall and dense canopy, constant mixing length within 302 the canopy, etc..) break down for short and sparse canopies (Poggi, Porporato, et al., 303 2004) as evidenced by the near independence between z_d and h_c in Figure 2 for short 304 h_c . 305

3.3 Quality Assurance and Quality Control

To ensure that the data are both of high quality and readily applicable, a number of quality assurance steps are applied: (1) All points that fail NEON quality assurance for air temperature are removed, (2) data where the reported energy balance has a residual greater than 20% are removed, as large residuals indicate high likelihood of significant advective fluxes and thus complicate the analysis (Mauder et al., 2020). The 20% threshold was selected to preserve as much data as possible for site-by-site analysis, and ³¹³ no significant difference in data quality or observed trends was noted when tightening ³¹⁴ this threshold further. (3) All data points with $\zeta > 0$ are removed as uncertainties in ³¹⁵ this range are high, data availability is relatively low, and this is not the intended focus ³¹⁶ of the study. (4) Periods with non-negligible precipitation are removed. (5) Any site which, ³¹⁷ after all previous quality control is applied, retains less than 100 half-hourly runs are re-³¹⁸ moved. Quality control retained just over 32,000 half-hourly runs across 39 NEON sites, ³¹⁹ roughly equivalent to about 2 site-years at 30-min averaging.

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3.4 Random Forest (RF) Method

The RF method is used to generate an initial data-driven alternative model to equa-321 tions 1 and 2 with no regards to dimensional constraints as required by similarity the-322 ory. This allows us to examine empirically impacts that various environmental predic-323 tors not present in current MOST based formulations might have on the development 324 of PTV. The RF is a machine learning method that uses ensemble decision trees for pro-325 ducing a regression, with each decision tree run using a random subsample of the data 326 to generate the model (Breiman, 2001). Data are randomly split into testing (10%) and 327 training (90%) datasets for RF as well as other fitting analysis. Accuracy is evaluated 328 primarily using normalized root mean squared error (nRMSE). From the results of the 329 model, we extract feature importance, a measure of which predictors play the largest role 330 in the model fit. In this case, high feature importance indicates that the value of a given 331 predictor is essential for describing and predicting PTV using the RF method. We have 332 elected to use sensible heat, ζ , and u_* due to their role in MOST. Tree cover fraction, 333 bare soil fraction, vegetative fraction, LAI, and effective drag $C_d = [u_*/u(z)]^2$ are used 334 to potentially represent canopy structure and roughness effects. BLH, CAPE, and cloud 335 cover fraction are selected due to their relation with mesoscale phenomenon and large 336 scale eddies. 337

338 4 Results

The analysis begins with basic examination of the data across all sites. These are 339 summarized as comparisons between σ_T and the environmental predictors presented in 340 section 3.4, followed by analysis of the diurnal cycle of sensible heat and σ_T . The data 341 are then compared to the curves of (1) and (2). Analysis continues focused on explor-342 ing site level differences, first with RF over the entire dataset as well as individually for 343 each site. A bar plot showing the relation between predicted and observed at each site 344 is then featured to illustrate differences between land cover types. The final section of 345 the analysis focuses on evaluating potential model improvements leveraging the results 346 in the previous sections. The observations are compared to Equations 1 and 2 with up-347 dated parameter values selected through curve fitting. This comparison is shown over 348 both the overall dataset and a select few sites. Finally, select parameterizations of the 349 b parameter in (1) based on a variety of metrics that represent canopy structure are pre-350 sented, evaluated, and compared to traditional formulations. 351

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4.1 Holistic Exploration

The data from the remotely sensed products and NEON were merged and then qual-353 ity controlled as described in section 3.3. Figure 3 presents a comparison between σ_T and 354 collocated environmental and meteorological data. The results show a clear relation be-355 tween PTV and sensible heat flux $H = \rho C_p \overline{w'\theta'}$ where ρ is the mean air density and 356 C_p is the specific heat capacity of dry air at constant pressure as well as ζ and to a lesser 357 extent effective drag C_d . In addition, some patterns seem apparent with LAI and BLH. 358 Other environmental variables not included in figure 3 have no significant relation with 359 PTV. For H, there appears to be a family of curves rather than one defined shape, im-360 plying some additional parameter is influencing that relation. Effective drag, similarly, 361



Figure 3. Relation between environmental variables and σ_T across all sites. The resulting scatterplots are binned into small hexagons; the colors illustrate the concentration of points in each hexagon where blue is low and yellow/brown is high.

has two families of curves, with the larger effective drag values arising primarily fromforested sites.

The diurnal cycles of PTV and H are plotted in figure 4. In the four forested sites, the raio of H relative to σ_T is higher when compared to the four low lying sites, already suggestive of the importance of site level difference. Figure 4 is also illustrative of the differences between unfiltered and filtered PTV, with the change being most significant in the mornings when sensible heat flux is small but rapid changes in mean air temperature would artificially inflate the apparent PTV caused by turbulence only.

The data covers a range of stability conditions in the near neutral and unstable range, 370 as indicated in figure 5a. The shape of the data generally follows expectations from MOST 371 with an extensive $\zeta^{-1/3}$ scaling (Tillman, 1972) in figure 5b in the unstable range, al-372 though in the near neutral range this is less clear. Similarly, in figure 5a, there is some 373 deviation from the established formulation in equation (1), especially as ζ increases in 374 magnitude. Comparing equations (1) and (2) directly to the data show significant er-375 rors. Equation (1) has an nRMSE of 21.5% and a 1% bias, although the bias in equa-376 tion (1) is deceptive as the model has significant negative bias at low values and a pos-377 itive bias at larger values. Equation (2) performs significantly worse over that range, with 378 an nRMSE of 27.6% and a bias of 15%. 379



Figure 4. The diurnal cycle of sensible heat flux (*H*), filtered (red-dashed) and unfiltered (orange-dashed) σ_T for 8 selected sites, with their locations indicated on the central map of CONUS.



Figure 5. $\mathbf{a}(\text{left})$: Relation between the dimensionless standard deviation of potential temperature σ_T/T_* and the stability parameter ζ for the data with the modeled values from Equation (1) in black and Equation (2) in red. \mathbf{b} (right): Relation between the dimensionless standard deviation of potential temperature and $(-\zeta)^{-1/3}$. The resulting scatterplots in both panels are binned into small hexagons; the colors illustrate the number of points in each hexagon where blue is low and yellow/brown is high.



Figure 6. a (left): The feature importance from the random forest on the aggregate dataset sorted by overall importance. b (right): Results of the site level random forest feature importance. Violin plot shows distribution of site level feature importance for each predictor

4.2 Site by Site Comparison

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Random Forest provides a first pass at the potential to improve upon the model 381 when noting deviations from the data in figure 5. RF does perform significantly better 382 than either model with nRMSE of 13.4% and a bias of less than 0.1%, although com-383 putational constraints prevent its use in ESMs. The feature importance provides dynam-384 ically interesting results as can be seen in figure 6. Sensible heat flux dominates the de-385 termination of PTV as expected from flux variance literature. The relative unimportance 386 of friction velocity is also consistent with equation (2) and with equation (1) when the 387 magnitude of ζ is large. The high importance of f_{tree} may be consistent with the results 388 in figure 4 as well, further indicating that tree cover has a significant impact on the re-389 lation between sensible heat flux and PTV. Somewhat surprising is the relatively low im-390 portance for ζ . Although it is notable that since ζ is a function of H and z_d , which is 391 related to canopy height, a significant portion of the stability effect may be captured by 392 these two aforementioned variables. 393

The results shown in figures 3, 4, and 6 all indicate the possibility of variable curves 394 for each site in the network. RF was run again, separately, for each individual site to ex-395 amine these possible relations and remove any attempts by the algorithm to use a pre-396 dictor as a proxy for the site. The violin plot in figure 6 shows the distribution of the 397 feature importance of each predictor across sites. When examined site by site, H is an 398 even more dominant predictor for PTV. The stability parameter becomes the second most 399 important indicator, consistent with preexisting MOST formulations, although again it 400 is small when compared to sensible heat flux. 401

When further exploring site level differences, key patterns begin to emerge. When 402 comparing the PTV predicted by equation 1 and the PTV observed at each site, there 403 is a significant variability in the slope of the best fit line of the data, which would ide-404 ally sit at 1 indicating close agreement between the observations and the model. Figure 405 7 illustrates how that slope changes site to site and with land cover type. Sites with slopes 406 close to 1 are generally flat, homogeneous, and dominated by low lying vegetation, which 407 is the ideal landscape for MOST, and matches the landscapes where the values of the 408 parameters a, b and C_1 were originally derived. Forested sites however, especially those 409 with significant heterogeneities, have slopes significantly higher than 1, indicating that 410 the pre-existing model underpredicts PTV at low values and overpredicts PTV at high 411 values. LENO, SJER and SOAP in particular are all sparsely forested sites with signif-412 icant open water at LENO, an oak savannah at SJER, and sparse evergreens with vary-413



Figure 7. Bar plot showing the best fit slope between the observed and predicted temperature variance at each site using equation 1. Normalized RMSE is also listed in red above each bar. Site bars are colored by the dominant NLCD land cover

ing topography at SOAP. In addition, it is notable that ABBY is a logging site, so while
it is classified as evergreen, the actual canopy is quite short, and the tower is located in
a clearing. Site sensible heat and Bowen ratio were also examined, and can be seen in
Figures S1 and S2 respectively in the supplementary material, however few trends were
found outside of poor performance under very low Bowen ratios or sensible heat fluxes.

4.3 Adjusting Existing Models

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The results in figure (5) show that there is good predictive value in the existing schema, 420 but also imply that operational adjustments could yield improvements. An iterative fit-421 ting process was used to determine the optimal values for the constants a, b, and C_1 over 422 the aggregate data. This global fit resulted in small but non-trivial model improvements 423 for equation (1) and equation (2) both in error and bias. Equation (1) after a global fit 424 to the data nRMSE changes from 20.5% to 17% while bias remains constant and for equa-425 tion (2) the error is reduced from 27.6% to 18.4% and bias shifts from 15% to -1.4%. These 426 values are compared with additional models in Table 2. Figure 8 illustrates how the newly 427 fitted curves describe the entire data in two ways: dimensional and dimensionless forms. 428 The dimensional form of the comparison was selected because it does not suffer from any 429 self-correlation. Self correlation arises here because H impacts both T_* (ordinate) and 430 ζ (abscissa) in the stability correction function, which can lead to spurious agreement 431 (especially in the exponent). In dimensional form (left panels 8), the σ_T is computed from 432 measured u_* , H, z, and inferred z_d and compared to independently measured σ_T obtained 433 after filtering the high frequency air temperature series. Fitted equation (1) performs 434 better than fitted equation (2) though in dimensional form, this difference appears mi-435 nor. This difference becomes clear when the two formulations are assessed by stability 436 class and appear to diverge in the near neutral range. In this range, there is much greater 437 uncertainty in the values of σ_T/T_* (though the variances themselves are small). As $T_* \rightarrow$ 438 0 but σ_T remains finite due to entertainment of heat and due to finite signal-to-noise ra-439 tio in the measurements, σ_T/T_* becomes ill-defined or suspect in equation (1). Interest-440 ingly, equation (2) suggests that both the left-hand and right-hand side becomes unbounded 441 as $T_* \to 0$, and thus predicts rapid increase in σ_T/T_* as $|\zeta| \to 0$. While this increase 442 in σ_T/T_* appears to be consistent with some data sets, it is simply a statement that σ_T/T_* 443 may be ill-defined. As such, we will be focusing our analysis on equation (1) in which 444 σ_T/T_* is forced to approach a constant and σ_T maintains its scaling with local sensible 445 heat flux. Last, a realizability constraint was also developed (described later) so as to 446 illustrate a theoretical lower limit for the applicability of equations (1) and (2). Figure 447 8 demonstrates that the majority of the observations (in dimensionless form) as well as 448



Figure 8. (upper left): Comparison between observed and predicted temperature variance by equation (1) after a global fit (a=7.5, b=34.0) over all sites. (bottom left): Comparison between observed and predicted temperature variance by equation (2) after a global fit (=.812) over all sites. Note the comparisons in the left panels do not suffer from self-correlation. (right): the stability correction function for the non-dimensional temperature variance σ_T^2/T_*^2 . The original forms of equation (1) (a=4, b=8.3) and equation (2) (=.95) are shown as well as the fitted versions of both equations. In addition, the limit imposed by the realizability constraint is featured. The resulting scatterplots in all three are binned into small hexagons; the colors illustrate the number of points in each bin where blue is low and yellow/brown is high.



Figure 9. Stability correction function for the temperature variance at selected sites. In red, the modeled results are plotted for both the original Equation 1 (a=4, b=8.3) and the global fit Equation 1 (a=7.5, b=34) as well.

the changes in parameter values remain above the line and satisfy this realizability constraint.

The global fit does not perform universally well at each site, although most sites 451 realize some improvements. Figure 9 shows the global fit and original equation (2) over 452 select sites as well as scatterplots of the data, similar to the hexbin scatterplot for the 453 aggregate data in figure 8c. At three of the sites, ORNL, WREF, and SOAP, the data 454 lies largely below both the original equation (1) curve as well as the fitted curve. These 455 sites all have significant forest cover, especially compared to the 4 sites where the data 456 lies largely above the fitted curve, NOGP, SRER, KONZ and DSNY, which are all flat 457 sites with only bare soil or low-lying vegetation. The performance of these site by site 458 fits are presented in Table 2. 459

These site level differences can also be examined more quantitatively. The fitting 460 exercise was repeated, this time doing a separate fit of the b parameter for each site while 461 holding a constant. The a is held constant and b is adjusted because under highly con-462 vective conditions, which are the main conditions of interest, b is the dominant param-463 eter whereas a dominates in the more uncertain near-neutral range. After comparing b464 and other environmental predictors, it became clear that there is a close relation between 465 a variety of measurements of canopy structure around the tower and the best fit value 466 of the b parameter. As shown in Figure 10, there is a clear linear relation between the 467 cube root of b and the different measurements of vegetative structure: canopy height, 468 leaf area density (LAI/h_c) , effective drag C_d , z/z_d , LAI/z_d , and $h_c^2/[(z-z_d)^2 LAI]$. 469



Figure 10. Scatterplot of the selected predictors at each site on a log scale compared to the best fit value for the b parameter from equation (1). The points are colored according to the dominant NLCD landcover type

Equation 1				
Model	a	b	nRMSE	nBias
Standard	4	8.3	20.5%	1.4%
Global Fit	7.5	34	17.0%	-2.7%
Site by Site Fit	Site by Site Fit 7.5 20-80		15.5%	-3.0%
Equation 2				
Model	C_1		nRMSE	nBias
Standard	0.95		27.6%	15%
Global Fit	0.812		18.4%	1.4%
Random Forest				
Global			13.4%	< 0.1%

Table 2. Summary table of the results of various selections for the parameters of equations 1 and 2 as well as the random forest model. The table includes the normalized RMSEs and normalized biases of the different models for PTV.

Table 3. Summary of a variety of possible parameterizations of the b parameter in equation (1) following the form of equation (9)

χ	α	β	nRMSE	nBias	R^2
LAI/h_c	0.036	0.289	16.2%	-5.0%	0.55
LAI/z_d	0.036	0.284	16.3%	-4.8%	0.57
z/z_d	0.042	0.277	16.7%	-3.6%	0.54
C_d	-0.047	0.228	16.8%	-2.9%	0.46
$h_c^2/[(z-z_d)^2 LAI]$	-0.024	0.258	16.4%	-4.7%	0.52
h_c	-0.025	0.323	16.5%	-3.4%	0.5

Taking advantage of this relation between b and the various proxies for canopy structure around the tower, a linear model was developed for b based on a linear regression. Applying the model in Equation (6) to update the parameters in Equation (1) with the existing data yields only marginal improvements on the updated Equation (1) based on globally fit parameters, as is clear in Table 3. The relation

$$b^{-1/3} = \alpha \log(\chi) + \beta \tag{9}$$

does suggest that as canopy height increases, *b* also increases thereby amplifying the modulations introduced by ζ that act to reduce the dimensionless temperature variance. Hence, it appears that tall canopies make the dimensionless temperature variance more sensitive to ζ .

$_{480}$ 5 Discussion

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5.1 Results Summary and Synthesis

Overall, the filtered data indicates agreement with MOST based formulations (Tillman, 482 1972). While this study emphasizes that there is room for improvement, it is also im-483 portant to note that equation (1) holds even over non-idealized landscapes and atmo-484 spheric conditions despite the fact that the formulation and parameter values were de-485 rived over highly idealized flows. The evaluation of equation (1) over these wide-ranging 486 landscapes offers one of the clearest pictures in the literature of its broad applicability 487 for understanding PTV at the bottom of the surface layer (provided non-turbulent phe-488 nomenon are filtered). From the results of the random forest, and inspection of figure 489 2, it is evident that other local physical and meteorological characteristics that were thought 490 to have some influence on the development of PTV in the surface boundary are largely 491 unimportant. Heat flux and local stability continue to be the driving factors and can yield 492 good predictions for PTV over flat landscapes using the standard parameter values. Based 493 on model error analysis, there is additional uncertainty to be captured. 494

Numerous studies have shown how parameter values for various local sites can de-495 viate from the global values described in the early literature, however few have proposed 496 updates to models used in ESMs as these studies often include only a very small num-497 ber of sites and therefore painted a limited picture of the variety that one can find in the 498 field. Site by site fitting to the parameter values indicate that most sites have param-499 eter values larger than those defined in the literature, and only one site was found to have 500 parameter values smaller. Since the best fit values of the parameters across landscapes 501 do not oscillate around these 'ideal' values, but rather are all greater than or equal to 502 them, ESMs can benefit from alternative global parameters to cover regions with var-503 ied and heterogeneous canopies. The inter-site best fit parameter variation is quite sig-504 nificant, with best fit values of b ranging from 20 to 80. This suggests that, while global 505 parameter values may be useful for broad application, localized studies will benefit most 506 from a local, site based empirical fit, especially if they deviate from ideal (i.e. flat, ho-507 mogeneous, short vegetation) surfaces. An important note with respect to these model 508 fit values is the role of the filtering process. The filtering process described yielded closer 509 agreement to MOST for all of these model fits; unfiltered data overall yields slightly more 510 noise and a greater deviation from traditional MOST relations with a larger magnitude 511 of bias, but maintains inter site trends with unfiltered data. 512

Attempts to use environmental predictors to capture the local variation were only marginally successful outside of a random forest model. The RF method detailed significant improvement, and was able to capture most of the inter-site variability based on the tree cover fraction. This implies that canopy structure and surface roughness characteristics are partly responsible for a significant portion of the deviations from ideal conditions. The RF method, however, is too computationally intensive for application in ESMs. As such, there was an attempt to generate a compact model for the *b* parameter based on environmental variables related to surface roughness and canopy. Of a long list of possible predictors to model values of b, the most successful are shown in figure 10. A plausibility argument for the inclusion of some of these variables to parameterize b may be obtained by examining qualitatively the potential temperature variance budget (PTV).

524 5.2 PTV Budget: A scaling analysis

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Using index notation, the PTV above the canopy is given by

$$\underbrace{\frac{1}{2}\frac{\partial\overline{\theta'^2}}{\partial t}}_{\mathrm{I}} + \underbrace{\frac{1}{2}\overline{U_j}\frac{\partial\overline{\theta'^2}}{\partial x_j}}_{\mathrm{II}} = \underbrace{-\overline{u'_j\theta'}\frac{\partial\overline{\theta}}{\partial x_j}}_{\mathrm{III}} - \underbrace{\frac{1}{2}\frac{\partial\overline{u'_j\theta'^2}}{\partial x_j}}_{\mathrm{IV}} - \underbrace{\frac{\epsilon_{\theta}}{\mathrm{V}}}_{\mathrm{V}},\tag{10}$$

where i = 1, 2, 3 indicates longitudinal $(x_1 \text{ or } x)$, lateral $(x_2 \text{ or } y)$, and vertical $(x_3 \text{ or } x)$ 527 z) directions, respectively, repeated indicies imply summation, t is time, and u'_i are ve-528 locity fluctuations along direction x_i from the mean value $\overline{U_i}$. In this budget, term I is 529 the local variance storage, which as an outcome of the filtering exercise, can be neglected. 530 Term II is advection of PTV by the mean flow. In the analysis here, it is assumed that 531 subsidence $(\overline{U}_3 = 0)$ is small and due to the choice of coordinate systems, the mean lat-532 eral velocity is zero ($\overline{U_2} = 0$). Term III is the production due to the turbulent heat fluxes 533 and mean potential temperature gradients, which we assume to be driven by the ver-534 tical direction components. Term IV is the turbulent transport term. This term is sig-535 nificant in some cases, but the filtering process used here will remove some portion of 536 the non-local heat transport effects captured by this term. Hence, for simplicity, this term 537 is momentarily ignored. Term V is the molecular dissipation of PTV, which, along with 538 term III, tends to comprise the largest portion of the budget (Champagne et al., 1977; 539 Monji, 1973). There are additional terms, not presented here, that represent radiative 540 destruction, conductive diffusion and another dissipation term. These terms are gener-541 ally considered negligible near the surface. With these simplifications, 542

$$\frac{1}{2}\overline{U}\frac{\partial\overline{\theta'^2}}{\partial x} = -\overline{w'\theta'}\frac{\partial\overline{\theta}}{\partial z} - \epsilon_{\theta}.$$
(11)

To proceed further, closure schemes and scaling analyses are needed for the mean advection and variance dissipation terms, which are given as

$$\epsilon_{\theta} = C_{\epsilon,\theta} \frac{\overline{\theta'^2}}{\tau}; \frac{1}{2} \overline{U} \frac{\partial \overline{\theta'^2}}{\partial x} = -C_{adv,\theta} \frac{\overline{\theta'^2}}{\tau_{adv}}; \tau_{adv} = \frac{L_x}{\overline{U}}, \tag{12}$$

where $C_{\epsilon,\theta}$ and $C_{adv,\theta}$ are closure constants, τ is a relaxation time scale describing how 547 long a potential temperature excursion lasts before it gets dissipated by molecular pro-548 cesses, τ_{adv} is an advection time scale formed by the local mean velocity at z and a po-549 tential temperature spatial variability integral length scale L_x . This length scale reflects 550 imprints of 'near-field' heat sources from tree crowns or vegetation upper layers not 'blended 551 out' by turbulence within the roughness sublayer above the canopy. When the spatial 552 imprint of these heterogeneous heat sources is entirely blended out by turbulence mix-553 ing, $L_x \to \infty$. Inserting these closure schemes into the simplified PTV budget and af-554 ter some algebra results in 555

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$$\frac{\theta^{\prime 2}}{T_*^2} = -\frac{u_*\tau}{\kappa(z-z_d)} \frac{\phi_h(\xi)}{C_{\epsilon,\theta}} \frac{1}{\left(1 - \frac{C_{adv,\theta}}{C_{\epsilon,\theta}} \frac{\tau}{\tau_{adv}}\right)}; \phi_h(\xi) = -\frac{d\theta}{dz} \frac{\kappa(z-z_d)}{T_*} \phi_h^*(\xi), \tag{13}$$

where $\phi_h(\xi)$ is the stability correction function for temperature defined in the roughness sublayer defined using a surface layer representation adjusted by $\phi_h^*(\xi)$, a roughness sublayer modification. The $\phi_h^*(\xi) \to 1$ as $|\xi| \to \infty$. As discussed previously, over forested terrain and those with high heterogeneity in canopy structure, equations (1) and (2) have

a tendency to over predict temperature variance unless the fit is adjusted (in the case 561 of C_1 , a lower value). If the simplified version of the PTV budget is examined, then this 562 underestimation emerges from competition between three terms: (i) the finite ratio of 563 τ/τ_{adv} increasing the normalized PTV, (ii) the role of $\phi_h^*(\xi) \subset [0,1]$ reducing the nor-564 malized PTV, and (iii) the reduced τ relative to its surface layer value. The L_x is finite 565 in the roughness sublayer and may be commensurate with crown size, crown-to-crown 566 spacing, and measurement height z (i.e. more blending with increasing z), which may 567 explain why z/z_d emerges as an explanatory variable to b. Moreover, tree spacing and 568 canopy height are indirectly captured by canopy LAI, another variable impacting b. The 569 finite L_x here is shown to increase normalized PTV above its surface layer similarity value, 570 not reduce it. The two processes that act to reduce the normalized PTV are $\phi_h^*(\xi)$ and 571 τ . The $\tau = TKE/\epsilon$ may be smaller than predicted by its surface layer value ($\propto (z - \tau)$ 572 $(z_d)/(u_*)$, where TKE and ϵ are the turbulent kinetic energy and its dissipation rate. This 573 underestimation of τ arises because of an enhancement of ϵ near the canopy top rela-574 tive to predictions from extrapolations of the surface layer similarity value (Poggi, Katul, 575 & Albertson, 2004). Likewise, roughness layer corrections to $\phi_h(\xi)$ are usually smaller 576 than unity, meaning that $\phi_h(\xi)$ in the roughness sublayer must be smaller than its sur-577 face layer counterpart (Garratt & Segal, 1988; Harman, 2012). 578

In addition to possible reductions in C_1 due to τ , and $\phi_h(\xi)$, deviations from equa-579 tion (1) due to large scale eddies was also proposed as a source of uncertainty. This ef-580 fect is primarily captured by term (IV) in the PTV budget. This term was neglected in 581 the analysis here primarily because of the filtering of temperature time series employed 582 to minimize the effect of term I (the storage term) and likely have filtered some of the 583 very large scale inactive eddies that can contribute to PTV but not T_* . Thus, the effect 584 of large-eddies in the ABL is to increase the normalized PTV, not reduce it. Using the 585 5-min filtered series in PTV calculations, the RF algorithm here paints an unclear pic-586 ture that neither confirms this hypothesis nor offers a clear refute. Figure 3 shows that 587 there is a weak linear relation between BLH and the temperature variance, suggesting 588 that when the ABL is large, additional external sources of variance could exist, intro-589 ducing new length scales not represented in MOST for the PTV budget. The relation 590 in Figure 3, however, is weak. Figure 6a and 6b also support this weak relation. The three 591 predictors included in the RF to represent these effects, BLH, cloud cover fraction, and 592 CAPE, all have a feature importance of less than 1%. These three predictors, however, 593 do have their limitations and the weak relation shown here does not clearly refute the 594 hypothesis. The data comes from a reanalysis product with relatively poor resolution, 595 meaning that measurements of these values, particularly BLH are not necessarily reli-596 able. Another important note is the specific selection of time periods where the surface 597 energy balance is largely closed. Previous literature indicates that sub-mesoscale circu-598 lations (usually of time scales longer than several minutes) may cause the non-closure 599 of the surface energy balance (Mauder et al., 2020), which means that by virtue of con-600 straining the study to a closed energy balance we may be excluding the study periods 601 where these circulations would have an impact. Likewise, the removal of time scales longer 602 than 5 min may also ameliorate sub-mesoscale circulations. Examining the PTV bud-603 get equations, one could see how advection may increase PTV and the scatter of PTV 604 as well if included. Upon initial examination, the data appears to support this hypoth-605 esis with greater observed scatter of PTV, although a more comprehensive analysis is 606 outside the scope of this manuscript. Additional studies, with more reliable and locally 607 relevant measurements of BLH such as through surface to air LIDAR as well as consid-608 eration of the surface energy balance, are required to adequately assess this hypothesis. 609

5.3 Realizability Constraint

Equations 1 and 2 must satisfy the realizability constraint requiring that θ' and w'must not be perfectly correlated resulting in the inequality

$$\sigma_w^2 \sigma_T^2 > (\overline{w'\theta'})^2 = (T_* u_*)^2 \tag{14}$$

614 and

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$$\frac{\sigma_w^2}{u_*^2} \frac{\sigma_T^2}{T_*^2} > 1 \tag{15}$$

Paired with the original formulation of equation (1) are equivalent MOST consistent forms for σ_w^2/u_*^2 (Andre et al., 1978)

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 $\frac{\sigma_w^2}{u_*^2} = 1.75 + 2(-\zeta)^{2/3} \tag{16}$

 $_{619}$ Thus, combining equation (15) and equation (16) we obtain

$$\frac{\sigma_T^2}{T_*^2} > \frac{1}{\frac{\sigma_w^2}{u_*^2}} = \frac{1}{(1.75 + 2(-\zeta)^{2/3})}$$
(17)

Equation (17) is plotted as part of figure 8c, which clearly illustrates that the model adjustments proposed do not violate any realizability constraint.

5.4 Roughness Sub-Layer Effects

A significant source of concern when analyzing above canopy PTV and how it com-624 pares between grasslands and forested areas is the thickness of the roughness layer for 625 heat. It is difficult to ensure that the towers are reporting flow statistics outside of the 626 roughness layer and in the inertial layer, where MOST scaling is intended to apply. NEON 627 towers are designed to lie above the roughness layer (i.e. in the surface layer), and mean 628 wind profiles at the sites indicate that most sites are within the surface layer based on 629 momentum considerations. In addition, past studies have indicated that the thickness 630 of the momentum roughness layer and a roughness layer for scalar quantities such as po-631 tential temperature and water vapor are not necessarily the same. For scalar quantities, 632 the roughness layer can be significantly thicker than those for momentum upon which 633 the tower design is based (Raupach & Thom, 1981). If points are indeed interrogated 634 inside the scalar roughness sublayer, this could yield significant changes in the values of 635 temperature variance as the canopy and surface elements play a greater role in introduc-636 ing variable heat sources and sinks. Related to this is a concern inherent in the design; 637 towers over forested sites in NEON are designed as a factor of the canopy height, but 638 over low-lying vegetation it is defined as simply a constant 8m. This means that the in-639 strumentation may lie further up in a normalized profile for some of the flatter sites than 640 the forested ones, potentially explaining some of the differences between these two cat-641 egories. Initial exploration does show a poor but persistent relation between the ratio 642 of tower height to canopy height and the model parameter values. These challenges with 643 defining where the instrumentation lies above the canopy, however, is of lesser concern 644 for the primary intended application of this study in earth system models, where the rough-645 ness layer is inconsistently defined across different models. 646

⁶⁴⁷ 5.5 Future Work

A robust evaluation of the primary models of PTV at the surface layer was undertaken and avenues for improvements proposed. In addition to the importance of refining models over sparse canopies, as discussed in the previous section, exploration of these

models and PTV more generally in stable and, to a lesser extent, near neutral atmospheric 651 regimes is needed. The model currently assumes that the non-dimensional variance re-652 mains constant with stability, although when exploring the data this became less clear. 653 Previous studies have suggested high errors in the near neutral range is a consequence 654 of the non-stationarity (Kroon & de Bruin, 1995; J. Wyngaard & Coté, 1971). This ef-655 fect may be partly ameliorated by spectral filtering as shown. Yet, the data scatter at 656 the near neutral limit is undisputed. Significant scatter also exists in the mildly unsta-657 ble range, implying that there are issues with the application of MOST under these con-658 ditions that require more work. 659

An important limitation to note is that this analysis is all based on one point within 660 a tower reading from a flux footprint covering areas on the order of a few square kilo-661 meter. In ESMs, PTV is computed at the tile level which is intended to be representa-662 tive of the landscape level but, depending upon the tower height, can be on a larger scale 663 and is an area rather than point measurement. This scaling issue, which to the best of 664 the authors knowledge has not been examined previously, may have an impact especially 665 when considered in light of circulations and advection of PTV, where tiling schemes may 666 fail to be representative of the landscape. 667

There are three other avenues of future research that require further exploration. 668 First, while this study focused on the case where the energy balance is closed and there 669 is no significant advection, unbalanced conditions where sensible heat, latent heat and 670 ground heat flux fail to account for the energy balance constitute a large fraction of the 671 data. Initial work shows a clear shift in the fit of the data to Equation 1 under these con-672 ditions, with lower best fit parameter values and larger scatter. Unfortunately, explor-673 ing the potential effect of significant advection require model simulations or data not avail-674 able through NEON. The second avenue of future research is examining the analogous 675 models for the other primary atmospheric scalar, water vapor. The model assumes that 676 temperature and moisture behave similarly, with the same parameter values. Numerous 677 studies, as well as initial examination of the NEON data, illustrate that water vapor and 678 temperature do not behave identically (G. G. Katul & Hsieh, 1999; Asanuma & Brut-679 saert, 1999; De Bruin et al., 1993; Liu et al., 2021) in the surface layer as previously the-680 orized, and as such an alternative model, or at the very least, alternative values for the 681 a and b parameters are needed. Finally, results and previous literature have indicated 682 that surface heterogeneity, especially in heating on scales large enough to induce circu-683 lations, can have a significant impact on MOST derived parameterizations such as the 684 ones discussed here. A brief examination, not presented here, implies there is a complex 685 relation between heterogeneity and temperature variance statistics, and as such addi-686 tional work considering different length scales of surface heterogeneity may indicate new 687 directions for improvement and model analysis. 688

689 6 Conclusion

High frequency time series across 39 similarly instrumented sites covering varied 690 landscapes across CONUS were analyzed to assess the validity of existing models for tem-691 perature variance in the surface layer, note key deficiencies and recommend avenues for 692 improvement. Results indicated that conventional flux-variance similarity formulations 693 are largely corroborated by data in both dynamic-convective and nearly convective cases 694 provided non-turbulent features are spectrally filtered out. This filtering reduced the tem-695 perature variance by factors of up to 2 to 3 in some cases when compared to the unfil-696 tered runs. The most significant deviations from standard MOST formulations were ob-697 served over heterogeneous and forested sites. Site by site analysis also revealed bias to-698 wards similarity constants larger than the traditional parameter values used in the lit-699 erature and ESMs. A random forest model illustrated that there is variability not cap-700 tured by the traditional formulations. Results generally indicate that canopy structure, 701 surface heterogeneity, and roughness characteristics drive a portion of the inter-site vari-702

ability, although a dimensional approach was unable to illustrate superior predictive value. 703 Future studies will expand this analysis to include situations with non-local energy bal-704 ance closure as well as landscapes with sparse canopies or large surface heterogeneities. 705 Water vapor and carbon dioxide concentration, the other primary atmospheric scalars, 706 use the same formulation in CLUBB and other models as PTV, although the literature 707 shows a difference in behavior. As such, any updated parameter values for temperature 708 cannot be applied to other scalars and additional work is required to make similar im-709 provements to their variance fluctuations. 710

711 Open Research

The ERA 5 Reanalysis data (Hersbach et al., 2018) was downloaded from the Coper-712 nicus Climate Change Service (C3S) Climate Data Store. The results contain modified 713 Copernicus Climate Change Service information 2020. Neither the European Commis-714 sion nor ECMWF is responsible for any use that may be made of the Copernicus infor-715 mation or data it contains. MODIS Vegetative cover data (DiMiceli et al., 2015) and Leaf 716 Area Index (Myneni et al., 2015) is available through https://lpdaac.usgs.gov/. NEON 717 turbulence data (National Ecological Observatory Network (NEON), 2021) is available 718 through https://data.neonscience.org/data-products/DP4.00200.001. Finally, Land Cover 719 types (2016 Version) from the National Land Cover Database (Dewitz, 2019) are ava-720 ialble from the Multi-Resolution Land Characteristics Consortium database https://www.mrlc.gov/data. 721 Software used to process this data (Waterman, 2021) and generate results can be found 722 here: https://tinyurl.com/tswneon. 723

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