

Strongly eddying ocean simulations required to resolve Eocene model-data mismatch

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Abstract

Model simulations of past climates are increasingly found to compare well with proxy data at a global scale, but regional discrepancies remain. A persistent issue in modeling past greenhouse climates has been the temperature difference between equatorial and (sub-)polar regions, which is typically much larger in simulations than proxy data suggest. Particularly in the Eocene, multiple temperature proxies suggest extreme warmth in the southwest Pacific Ocean, where model simulations consistently suggest temperate conditions. Here we present new global ocean model simulations at 0.1° horizontal resolution for the middle-late Eocene. The eddies in the high-resolution model affect poleward heat transport and local time-mean flow in critical regions compared to the non-eddy flow in the standard low-resolution simulations. As a result, the high-resolution simulations produce higher surface temperatures near Antarctica and lower surface temperatures near the equator compared to the low-resolution simulations, leading to better correspondence with proxy reconstructions. Crucially, the high-resolution simulations are also much more consistent with biogeographic patterns in endemic-Antarctic and low-latitude-derived plankton, and thus resolve the long-standing discrepancy of warm subpolar ocean temperatures and isolating polar gyre circulation. The results imply that strongly eddying model simulations are required to reconcile discrepancies between regional proxy data and models, and demonstrate the importance of accurate regional paleobathymetry for proxy-model comparisons.

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Key Points:

- Eddy ocean simulations provide a profoundly different local flow compared to non-eddy simulations
- Heat transport is enhanced in eddy simulations leading to reduced equator-to-pole sea surface temperature gradients
- Eddy simulations reduce model-data mismatches for sea surface temperature and ocean flow

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Abstract

Model simulations of past climates are increasingly found to compare well with proxy data at a global scale, but regional discrepancies remain. A persistent issue in modeling past greenhouse climates has been the temperature difference between equatorial and (sub-)polar regions, which is typically much larger in simulations than proxy data suggest. Particularly in the Eocene, multiple temperature proxies suggest extreme warmth in the southwest Pacific Ocean, where model simulations consistently suggest temperate conditions. Here we present new global ocean model simulations at 0.1° horizontal resolution for the middle-late Eocene. The eddies in the high-resolution model affect poleward heat transport and local time-mean flow in critical regions compared to the non-eddy flow in the standard low-resolution simulations. As a result, the high-resolution simulations produce higher surface temperatures near Antarctica and lower surface temperatures near the equator compared to the low-resolution simulations, leading to better correspondence with proxy reconstructions. Crucially, the high-resolution simulations are also much more consistent with biogeographic patterns in endemic-Antarctic and low-latitude-derived plankton, and thus resolve the long-standing discrepancy of warm sub-polar ocean temperatures and isolating polar gyre circulation. The results imply that strongly eddy model simulations are required to reconcile discrepancies between regional proxy data and models, and demonstrate the importance of accurate regional paleobathymetry for proxy-model comparisons.

Plain Language Summary

Climate models are widely used to understand warm climates in the geologic past such as the late Eocene (38 million years ago; $\sim 8^\circ\text{C}$ warmer than today). To determine the quality of these models, simulations are often compared to measured proxies representing the regional environment. Here we show that a finer-than-typical detail in the ocean model causes a profoundly different regional ocean flow and environmental conditions. The improved correspondence to proxy data implies that high resolution simulations are required for a meaningful point-by-point data-model comparison.

1 Introduction

Model-data comparisons for warm periods in the geological past can be used to test the performance of climate models under greenhouse conditions (Tierney et al., 2020;

48 Tabor et al., 2016; Hutchinson et al., 2021; Lunt et al., 2021; Kennedy-Asser et al., 2020;
49 Braconnot et al., 2012; Zhu et al., 2020; Liu et al., 2009; Schmidt et al., 2014; Dowsett
50 et al., 2013; Cramwinckel et al., 2018). Some fully coupled climate models using state-
51 of-the-art Eocene geographic boundary conditions (Baatsen et al., 2016) and greenhouse
52 gas forcing simulate climates that correspond well to reconstructions of tropical sea sur-
53 face temperature (SST) and deep ocean temperature (Cramwinckel et al., 2018). How-
54 ever, in such simulations these models regionally simulate much cooler conditions in ex-
55 tratropical regions than proxy data suggest, particularly in the southwest Pacific (Lunt
56 et al., 2021; Cramwinckel et al., 2018; Baatsen et al., 2020; Lunt et al., 2012; Huber &
57 Caballero, 2011). Consequently, depending on the radiative forcing, models either pro-
58 duce SSTs near the equator that are higher than proxy data indicate or SSTs at mid-
59 to-high latitudes that are much lower than proxy reconstructions, leading to stronger merid-
60 ional SST gradients.

61 One challenge in paleoclimate model-data comparisons is the scale difference be-
62 tween proxies and models. The proxies capture a regional environment and effects of small-
63 scale regional setting (e.g. geography, bathymetry, and oceanography), while general cir-
64 culation models have difficulties capturing regional climate correctly due to the coarse
65 resolution that is typically used (1° horizontally or coarser for the ocean) (Nooteboom
66 et al., 2020; Dowsett et al., 2013; Harrison et al., 2016; Eyring et al., 2019; Kennedy-Asser
67 et al., 2020; Tabor et al., 2016). The quality of ocean models improves considerably at
68 a higher horizontal resolution (0.1°) (Griffies et al., 2015; Dong et al., 2014; Sun et al.,
69 2019; Müller et al., 2019; Dong et al., 2014; Viebahn et al., 2016; Hewitt et al., 2016; Mc-
70 Clean et al., 2006), especially their regional flow (Delworth et al., 2012; Marzocchi et al.,
71 2015; Nooteboom et al., 2020). This is not only due to higher level of detail, but also
72 because of the smaller scale interactions resolved (including mesoscale eddies of 10-30km
73 size) that influence the large-scale flow properties (Porta Mana & Zanna, 2014) and in-
74 crease the importance of the local setting (i.e. the paleogeography and bathymetry) in
75 the resulting regional ocean flow.

76 Biogeographic patterns of microplankton (e.g. dinoflagellate cysts; dinocysts) in
77 Southern Ocean marine sediments have been used as tracer of past surface oceanogra-
78 phy (Huber et al., 2004). For instance, Eocene sediments deposited near Antarctica con-
79 tain dinocyst species that are endemic to circum-Antarctic locations (Bijl et al., 2011).
80 Hence, Southern Ocean regions with many of these endemic species, as opposed to those

81 with abundant cosmopolitan species, must be oceanographically connected. This implies
82 that these biogeographic patterns of dinocysts provide a direct proxy of the flow direc-
83 tion itself (Bijl et al., 2011). So far, climate models were broadly able to match the cir-
84 culation patterns deduced from microplankton endemism in the Southern Ocean, some-
85 times after adaptations of the model paleobathymetry (Houben et al., 2019; Bijl et al.,
86 2013; Huber et al., 2004) or details of the configuration of critical Southern Ocean gate-
87 ways (Sijp et al., 2016). However, these model simulations cannot explain the occurrence
88 or absence of endemic dinocysts at some sites. In addition, state-of-the-art fully coupled
89 climate model simulations did come close to the proxy-based warmth in the southwest
90 Pacific Ocean, but this required a flow through the Tasmanian Gateway which was in-
91 compatible with microplankton-based evidence of surface ocean flow (Stickley et al., 2004;
92 Cramwinckel et al., 2020). Consequently, no model simulation exists that can reconcile
93 southwest Pacific Ocean warmth with ocean flow that is compatible with the plankton
94 records (Baatsen et al., 2020).

95 Here we show that high resolution ocean model simulations partly solve this mis-
96 match, using sinking Lagrangian particles to represent biogeographic patterns of microplank-
97 ton in the ocean model simulations (Nooteboom et al., 2019; Huber et al., 2004). We present
98 the first simulations of a global eddying Eocene ocean model with a 0.1° horizontal res-
99 olution (HR2 and HR4; Table 1). These simulations are initialized and forced with at-
100 mospheric fields from an equilibrium state of a coarser (1°) resolution model with a fully
101 coupled ocean and atmosphere (LR2 and LR4; Table 1) (Baatsen et al., 2020). Hence,
102 the high- and low-resolution simulations have a similar atmospheric forcing and bathymetry.
103 The new high-resolution simulations are run for a few decades (42 and 27 years for HR2
104 and HR4 respectively), sufficient for upper-ocean circulation to equilibrate.

105 **2 Effect of model resolution on Eocene flow**

106 The resulting ocean circulation is different between the eddying and non-eddying
107 configurations (Fig. 1). In the eddying simulations, the time-mean flow strength has a
108 higher spatial variability, the bathymetry has a larger influence on the flow strength and
109 direction (especially in the Southern Ocean; see Supporting Information Fig. S3 for the
110 bathymetry), and local scale features are much more pronounced, compared to the low-
111 resolution model. All western boundary and equatorial currents increase in strength, ex-
112 cept in the North Atlantic. The spatial structure and separation locations of the west-
113 ern boundary currents are also shifted. For instance, the eastward Agulhas separation

Table 1. The ocean model simulations of the middle-late Eocene (38Ma) in this paper

Run	resolution	layers	type	forcing ^a	years run
LR2 ^a	1°	60	fully coupled with atmosphere (CESM)	2×pre-industrial CO ₂	3000
LR4 ^a	1°	60	fully coupled with atmosphere (CESM)	4×pre-industrial CO ₂	4000
HR2	0.1°	42	ocean only (POP), forced by LR2 atmosphere	2×pre-industrial CO ₂	42
HR4	0.1°	42	ocean only (POP), forced by LR4 atmosphere	4×pre-industrial CO ₂	27

^aFrom (Baatsen et al., 2020)

114 (near South-Africa) is only present in the eddying simulations (it retroreflects more east-
 115 ward compared to the present day). Moreover, east of Australia, the East Australia Cur-
 116 rent (EAC) extends further southeastwards in the eddying compared to the non-eddying
 117 simulation, while there is a narrow but strong northward current east of Tasmania that
 118 is not present in the low-resolution simulations.

119 The EAC flow provides an example of the stronger influence of the paleobathymetry
 120 on the flow in HR4 compared to LR4, even though the bathymetry is the same in both
 121 configurations. Eddies are responsible for the downward transfer of momentum input at
 122 the ocean surface by winds that is eventually balanced by bottom form stresses (Munday
 123 et al., 2015). As a consequence, the flow is strongly determined by isobaths (i.e. lines
 124 of constant bathymetry) (Rintoul, 2018; Marshall, 1994). Hence, the bathymetry has a
 125 much larger influence on the flow if the ocean is eddying (in HR4 and HR2) than if it
 126 is not (LR4). In HR4, the EAC is steered further southeastward than in LR4 along the
 127 submerged continental block of Lord Howe Rise (see Supporting Information Fig. S3 for
 128 the bathymetry). Moreover, jets like the EAC have a narrower structure in the eddy-
 129 ing flow, due to interactions between eddies and the time-mean flow (Waterman et al.,
 130 2011), which has profound impacts on the regional oceanography.

131 **2.1 Model-data comparison: plankton biogeography**

132 The new Eocene ocean model velocity fields enable us to use sinking Lagrangian
 133 particles (Nootboom et al., 2020) to reveal biogeographic provinces of endemic microplank-

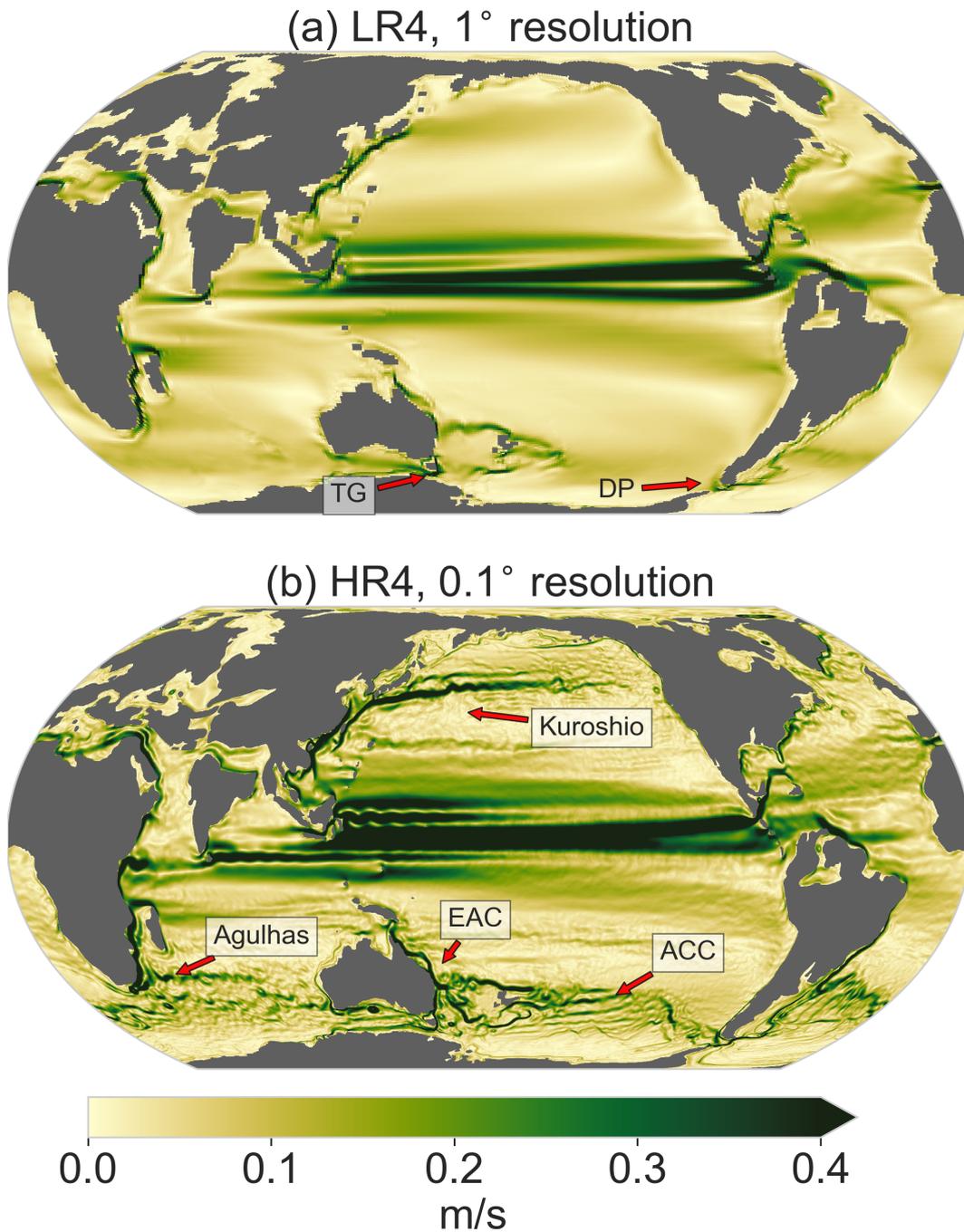


Figure 1. Magnitude of the time-mean surface horizontal flow velocity in the model of (a) 1° (mean over years 3995-4000) and (b) 0.1° horizontal resolution (mean over years 23-27). Both with 4×pre-industrial atmospheric CO₂ (LR4 and HR4). The Drake Passage (DP), Tasman Gateway (TG), East Australian Current (EAC), Kuroshio current and proto-Antarctic Circumpolar Current (ACC) are labeled.

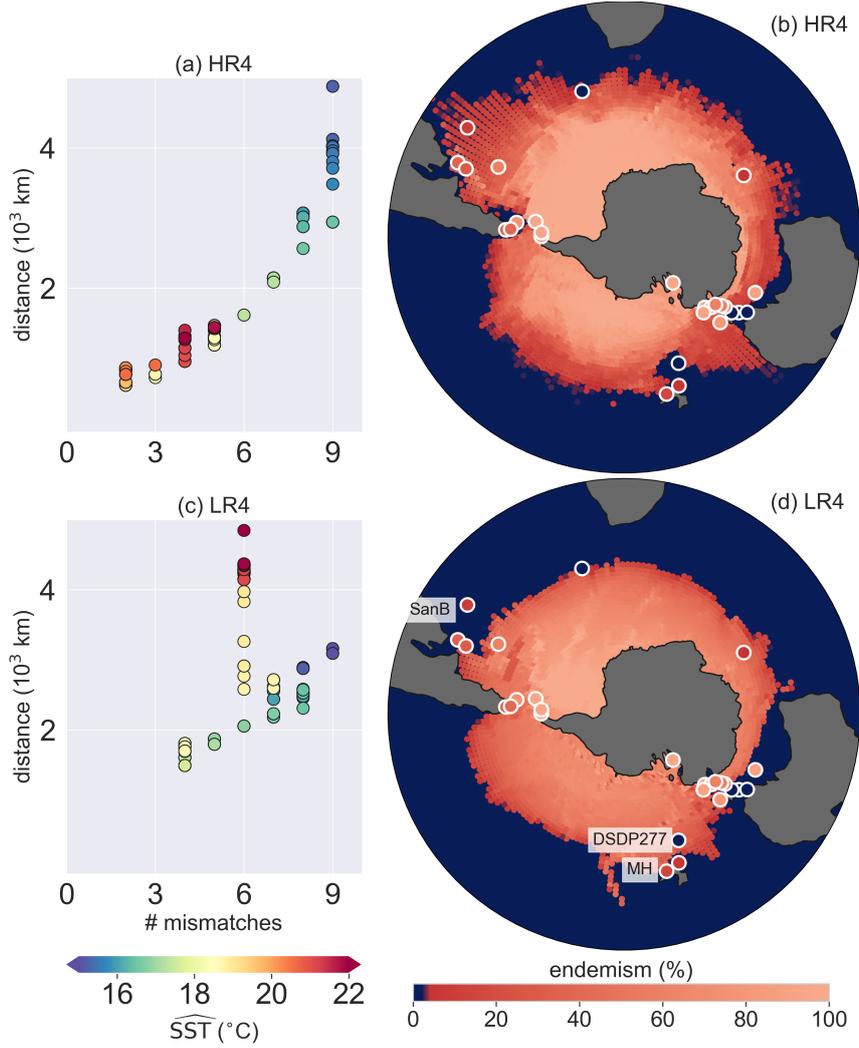


Figure 2. Model-data comparison: Antarctic endemism of sedimentary dinocysts in configurations HR4 and LR4. The model dinocyst endemism at the ocean bottom is determined by the percentage of virtual particles that started sinking (with 6 m day^{-1} sinking speed) in a surface environment with temperature below \widehat{SST} (see Supporting Information Fig. S4 for an illustration). (a), (c) Model-data fit for HR4 and LR4 respectively, for different values of \widehat{SST} (given by the dot colors). Model and data compare better if the following two measures of fit are lower: 1) the number of sites with a point-to-point model-data mismatch in terms of endemic dinocyst species occurrence and 2) shortest cumulative distance of these sites to a location in the model that does match in terms of endemic dinocyst occurrence (i.e. $\sum_i D_i$, where D_i is the distance between a site i and a location in the model that does match with site i in terms of the endemic dinocyst occurrence). (b), (d) Model-data comparison of dinocyst endemism at the \widehat{SST} value that minimizes the measures of fit in (a), (c). The sedimentary endemism of the data is the percentage of measured endemic species at the site (Bijl et al., 2011), representative of 41-39 Ma. Labeled sites are named in the main text.

134 ton in the Eocene Southern Ocean. In this way, we can test how representative the mod-
135 eled flow is compared to the reconstructed ocean flow from sediment records. In this ap-
136 proach, it is determined where sedimentary particles originated from at the ocean sur-
137 face, while taking into account how the particles were advected by ocean currents dur-
138 ing their sinking journey. If these virtual particles originate from an environment with
139 a temperature below a threshold value indicated by \widehat{SST} , the particle is assumed to orig-
140 inate close to Antarctica, and flagged as representing Antarctic-endemic dinocyst species
141 (see Materials and Methods section and Supporting Information Fig. S4 for an illustra-
142 tion).

143 Due to the circulation differences between eddying and non-eddy simulations,
144 the model-derived occurrence of Antarctic-endemic sedimentary dinocysts is clearly dif-
145 ferent between both configurations (Fig. 2). While the endemism is more strongly de-
146 pendent on latitude and a sharper boundary exists between low-endemism and high-endemism
147 in LR4, sinking particles are transported further away from Antarctica in specific areas
148 (especially near western boundary currents) in HR4. As a consequence, the occurrence
149 of several recorded endemic species can be explained in HR4, while it cannot in LR4 (see
150 e.g. site SanB). Moreover, the modeled endemism in the non-eddy LR4 cannot match
151 with both DSDP277 and MH at the same time, because these sites contain an opposite
152 signal (i.e. MH contains endemic species and DSDP277 does not) while being located
153 closely to each other. In HR4 on the other hand, the sedimentary particles in site DSDP277
154 (Fig. 2) originate only from the warm waters of the southeastward flowing EAC, while
155 the closely located site MH also contains particles originating from cold waters in the
156 east, in agreement with the occurrence of endemic species at MH.

157 Overall we find that only the eddying simulations produce circulation patterns con-
158 sistent with plankton biogeographic patterns. As a result, the model-data comparison
159 has a better overall fit in HR4 compared to LR4 (Fig. 2a,c). The model-data fit improve-
160 ment in HR4 compared to LR4 highlights the need for accurate reconstructions of the
161 geographic boundary conditions (Baatsen et al., 2016) to optimize model-data matches
162 as in Fig. 2a,b: It are the details in the ocean flow that induce a better model-data fit
163 in HR4 compared to LR4.

164 The modeled dinocyst endemisms in the $2\times$ and $4\times$ pre-industrial atmospheric CO_2
165 configurations are similar (see Supporting Information Fig S1), even though HR2 and
166 HR4 are forced by a different atmosphere and respond differently after initialisation (Fig.

167 3). However, the transient response of the upper ocean equilibrates similarly in the $2\times$
 168 and $4\times$ pre-industrial CO_2 cases in a few decades, which also results in a similar time-
 169 mean surface flow (Supporting Information Fig. S2). This implies that plankton biogeo-
 170 graphic patterns and surface ocean circulation are to a large extent affected by bathymetry,
 171 rather than the climate boundary conditions (e.g. atmospheric CO_2) of the model.

172 At the beginning of the HR2 and HR4 simulations, much of the energy input at
 173 the surface is used to set up the circulation and the development of eddies, as can be seen
 174 from a reduction of Southern Ocean gateway transports in the first 5 years (similar in
 175 both HR2 and HR4), after which they recover (Fig. 3c-f). After 9 years, the Drake Pas-
 176 sage transport (through the gateway between South America and Antarctica) exceeds
 177 the transport in the low-resolution simulations and equilibrates at a higher level. The
 178 increased Drake Passage transport is mainly caused by the lower (more realistic) viscos-
 179 ity that the high-resolution models allows compared to the low-resolution model (which
 180 becomes numerically unstable at this low viscosity value). Interestingly, the volume trans-
 181 port through the Tasman Gateway in HR2 and HR4 does not exceed the volume trans-
 182 port in LR2 and LR4. Instead, a larger fraction of the water is transported north of Aus-
 183 tralia, resulting in the stronger southeastward East Australian Current (EAC) in the South
 184 Pacific (Fig. 1).

185 **2.2 Model-data comparison: sea surface temperatures**

186 Now that the high-resolution Parallel Ocean Program (POP) model simulates an
 187 Eocene ocean flow, which is consistent with proxy data for ocean circulation, we com-
 188 pare the results of these simulations to proxy data for SST. SST distributions however,
 189 are also influenced by the model background state and sensitive to their global-scale equi-
 190 libration. Moreover, the background flow affects the distribution of heat differently in
 191 the eddying versus non-eddying simulations. Meso-scale eddies are important for the dis-
 192 tribution of heat, and eddying ocean models do a better job in representing heat trans-
 193 port compared to non-eddying models that use parameterizations for eddy-induced heat
 194 transport (Viebahn et al., 2016; Griffies et al., 2015; Dong et al., 2014).

195 Indeed, heat is distributed differently in the top km of the eddying compared to
 196 the non-eddying simulations (Fig. 3a and 3b). Eddies efficiently transport heat to the
 197 subsurface (Delworth et al., 2012), which leads to subsurface warming in both eddying
 198 simulations (HR2 and HR4) and a lower vertical temperature gradient compared to LR2

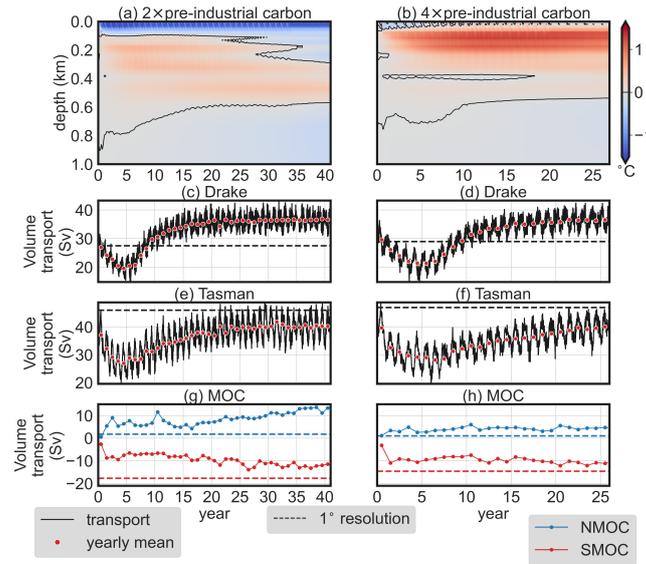


Figure 3. Response of the ocean model after initialisation, HR2 (left) and HR4 (right). Note that the initial state of HR2 (HR4) corresponds to LR2 (LR4). (a), (b) Depth-dependent evolution of the horizontal mean temperature increase compared to the initialisation state (upper 1km only). Water volume transport through the (c), (d) Drake Passage (65°W) and (e),(f) Tasman Gateway (150°E). (g), (h) Northern and southern maximum meridional overturning. MOC=Meridional Overturning Circulation, NMOC=Northern MOC, SMOC=Southern MOC, Sv=Sverdrup.

199 and LR4. However, in HR2 the surface cools more, while the subsurface warms less com-
200 pared to HR4.

201 Much of the heat transport change from LR to HR is related to the Southern and
202 Northern Meridional Overturning Circulation (SMOC and NMOC respectively). In both
203 HR2 and HR4, North Pacific sinking develops (in a few decades) next to existing South
204 Pacific sinking, while in the low-resolution simulations there is only Southern Hemisphere
205 sinking (see Supporting Information Fig. S8). Overall, the North Pacific sinking leads
206 to an increase in the NMOC and a decrease in the SMOC. These changes in the MOC
207 are stronger in HR2 compared to HR4, and both the NMOC and SMOC are still increas-
208 ing in magnitude at the end of the HR2 simulation.

209 The SMOC also differs in structure between the high- and low-resolution simula-
210 tions (see the mixed layer depth in Supporting Information Figure S8). In HR2 and HR4,
211 more volume transport through the Drake Passage increases the surface salinity in the
212 South Atlantic resulting in denser surface water in the Weddell Sea (Tumoulin et al., 2020).
213 Therefore, the main deepwater formation location is the South Atlantic in HR2 and HR4,
214 while it is the South Pacific in LR2 and LR4.

215 These results imply that HR2 and HR4 are run long enough for the upper-ocean
216 circulation to equilibrate, while the deep ocean is not in equilibrium yet, as can be seen
217 from the MOC in HR2 (Fig. 3g). Although the transient evolution of the deep ocean cir-
218 culation differs between HR2 and HR4, we can nevertheless investigate their impact on
219 SST distributions and compare those to proxy-data.

220 Both the tropical and Arctic Ocean cool significantly in HR2 compared to LR2,
221 while in HR4 the equatorial regions cool less and high-latitude (north and south) regions
222 warm more as compared to LR4 (see Fig. 4). For both atmospheric CO₂ levels, local SST
223 differences between the high- and low-resolution simulations mostly occur near western
224 boundary currents of which the location shifts in the eddy simulation (Fig. 4a and
225 d). These shifts have an effect on the model-data comparison at sites near western bound-
226 ary currents. In fact, the EAC transports warm waters southeastwards in the southwest
227 Pacific, which (partly) explains why sites in the southwest Pacific are found to be warmer
228 compared to model simulations with a coarse resolution. Notably, similar SST changes
229 occur near the Kuroshio and Agulhas currents. The Weddell Sea warms up in HR2 and
230 HR4 compared to LR2 and LR4 respectively, which is related to the South Atlantic sink-
231 ing that occurs in HR2 and HR4.

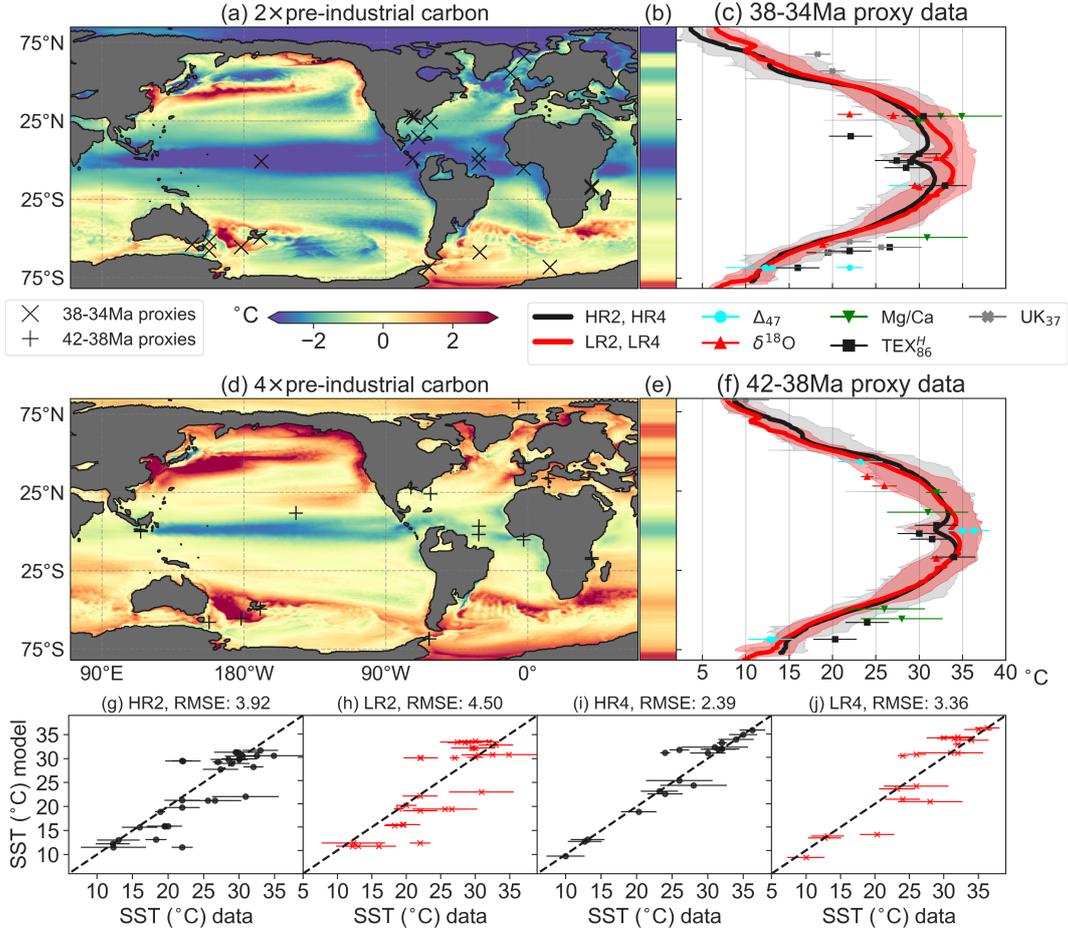


Figure 4. Model-proxy data comparison: sea surface temperature (SST). The $2\times$ and $4\times$ pre-industrial case are compared to SST proxy data of 38-34Ma and 42-38Ma respectively. (a), (d) SST difference of the high- compared to the low-resolution model with the site locations of the SST proxies for $2\times$ and $4\times$ pre-industrial carbon configuration, respectively and (b), (e) their zonal mean. (c), (f) the zonally averaged annual mean SST in the high resolution (black) and the low resolution (red) model for $2\times$ and $4\times$ pre-industrial carbon configuration respectively. The shaded areas show zonal spread (i.e. minimum and maximum) of the annual mean SST. Markers indicate SST proxy estimates with their uncertainty. (g-j) Scatter plots between proxy-derived and model-derived SST for all four configurations, with Root Mean Squared Errors (RMSE). Error bars represent proxy calibration errors. To consider the paleolocation uncertainty of sites (van Hinsbergen et al., 2015), each site is compared to the model SST value from up to 3° distance of the site that minimizes the RMSE of the scatter plot (similar to (Baatsen et al., 2020); see Supporting Information Fig. S7 for a point-to-point comparison). The dashed black line is the one-to-one line representing the perfect match between model and proxy data.

232 Climate models generally do not produce the low meridional temperature gradi-
233 ents of warm climates as inferred from proxy data (Huber & Caballero, 2011; Sijp et al.,
234 2014). While the simulations LR2 and LR4 were found to generate a lower meridional
235 SST gradient compared to other models of 1° horizontal resolution or coarser (Baatsen
236 et al., 2020), this gradient reduces further in HR2 and HR4. The tropics are cooler in
237 HR2 and HR4 compared to LR2 and LR4, while in the zonal-mean the southern high-
238 latitudes are only slightly warmer in HR4 (Fig. 4d-f). Regionally, there is, however, sig-
239 nificant warming of Southern Ocean SSTs in HR4. Overall, this improves consistency
240 between the high-resolution model results and SST proxies in the tropics, while the mod-
241 eled high-latitude SST values are often still lower than the proxy-derived SST values.
242 The eddying simulations show stronger horizontal gradients in the time-mean SST field
243 compared to the non-eddying simulations, which results in a higher time-mean SST vari-
244 ation in the model around the sediment sample sites. The model-data fit greatly improves
245 in the eddying compared to non-eddying simulations (Fig. 4g-j), although a mismatch
246 with some sites remains (especially for the $2\times$ pre-industrial CO_2 case) and the high-latitude
247 temperatures are overall lower compared to the proxy data.

248 Overall, the eddying ocean model improves the SST model-data match from the
249 non-eddying model, because it alters the local transport of heat. However, the SST model-
250 data comparison is also sensitive to the model background state (i.e. the state of the ocean
251 at a global scale), which depends on the used atmospheric forcing, paleogeography and
252 long time scales phenomena, such as the deep meridional overturning circulation. Hence,
253 the SST model-data mismatch could be reduced even further if better model boundary
254 conditions are used which lead to a more realistic background state of the late Eocene.

255 **3 Conclusion and outlook**

256 We have shown that an eddying Eocene ocean simulation provides a more detailed
257 ocean flow compared to a non-eddying version of the same model. As a result, model-
258 data mismatches in the geologic past (Lunt et al., 2021; Hutchinson et al., 2021; Baat-
259 sen et al., 2020; Houben et al., 2019; Bijl et al., 2011; Huber et al., 2004) can at least
260 partly be explained by the lack of eddies in the ocean models used. Our eddying sim-
261 ulations of the late Eocene are better able to explain the occurrence or absence of en-
262 demic dinocyst species near Antarctica compared to non-eddying simulations. The SST

263 model-data comparison also improved in the eddying compared to non-eddying simu-
264 lations.

265 The explicit representation of eddies in ocean models may have implications for com-
266 parison of models with other proxy types than considered here. For instance, pollen-based
267 temperature reconstructions imply that it did not freeze at the Antarctic coast during
268 winter in the early Eocene (globally $\sim 6^\circ\text{C}$ warmer than the late Eocene), despite po-
269 lar darkness (Pross et al., 2012). Eddy-induced flow, and its impact on ocean heat trans-
270 port, could in part explain such conditions.

271 The simulations in this paper are computationally expensive. However, other types
272 of model set-ups may be interesting if not limited by computational capabilities. First,
273 the strong influence of bathymetry on the eddying flow implies that the uncertainty of
274 paleogeography reconstructions has a major impact on model-data comparisons. Future
275 studies could make adaptations on the bathymetry within uncertainty of paleogeographic
276 reconstructions, to find its impact on the modeled ocean circulation and model-data com-
277 parison. Moreover, since the eddying flow has a direct response to bottom topography,
278 it seems suitable for a downscaling, or eddy parameterization type of approach to ob-
279 tain this influence of bathymetry on the flow with reduced computational costs. How-
280 ever, these type of approaches are found to be challenging in present-day configurations
281 (Fox-Kemper et al., 2019; Nooteboom et al., 2020; Lanzante et al., 2018).

282 Second, we used the model equilibrium of the non-eddying climate model simula-
283 tions (which are in radiative equilibrium (Baatsen et al., 2020)) to start and force the
284 eddying model. However, this switch induces a drift of the deep ocean circulation, which
285 is not equilibrated yet in the high-resolution simulations of this paper. Hence, the back-
286 ground state of the model will change further if the model is run for longer time peri-
287 ods (a few millennia). Future simulations may have the capabilities to perform longer
288 simulations. These changes of the model background state on long time scales might have
289 implications for the regional flow and the quality of the model-data comparisons.

290 Finally, atmospheric feedbacks greatly influence the ocean model background state
291 on long time scales, such as the meridional overturning circulation (den Toom et al., 2012;
292 Rahmstorf & Willebrand, 1995; Arzel et al., 2011; Zhang et al., 2010). Hence, the high-
293 resolution ocean should be coupled to a high-resolution atmosphere, which could further
294 enhance the meridional transport of heat and lead to an improved model-data compar-
295 ison.

4 Materials and Methods

4.1 Data

We used two datasets in this paper. The first includes the SST proxies from U_{37}^k , TEX_{86}^H , Mg/Ca, Δ_{47} and $\delta^{18}O$, which are described in detail in (Baatsen et al., 2020). Proxy-based SST reconstructions come with uncertainties, limitations and biases (Hollis et al., 2019), related to the depth, or season they represent. The second dataset are sediment samples with dinocysts from (Bijl et al., 2011), combined with the samples described in (Houben et al., 2019; Cramwinckel et al., 2020; Bijl et al., 2021). We averaged dinocyst abundance of Endemic-Antarctic, cosmopolitan and low-latitude-derived for the respective time slices.

4.2 Model set-up

We used the Parallel Ocean Program (POP) (Viebahn et al., 2016; den Toom et al., 2014; Smith et al., 2010) to perform eddy ocean model simulations for the middle-late Eocene (38Ma). To derive the forcing of this model, we made use of the fully-coupled (ocean and atmosphere) simulations with Community Earth System Model v1.0.5 (CESM) from (Baatsen et al., 2020), with a non-eddy ocean. We used both CESM simulations with $2\times$ pre-industrial atmospheric CO_2 (LR2) and $4\times$ pre-industrial CO_2 (LR4) configuration.

The high-resolution POP is forced at the surface by a fixed atmosphere of the CESM simulation. To construct the surface forcing, we interpolated the average (over the last 50 model years of LR2 and LR4) Sea Surface Temperature (SST), Sea Surface Salinity (SSS) and wind stress (zonal and meridional) of the CESM simulation for every month of the year (such that a seasonal cycle is included in the surface forcing). These SST and SSS fields were used as restoring boundary conditions at the surface. The restoring boundary conditions imply that POP is ‘pushed’ towards the SST and SSS output of the CESM at the surface with a specific timescale (30 and 10^{20} days respectively). This means that differences between the SST and SSS at different model resolutions arise due to the internal transport (vertical and horizontal) of heat and salt in the ocean, not due to the surface forcing. The bathymetry that CESM uses was interpolated linearly on the high-resolution grid that POP uses, making both bathymetries similar (see the code at <https://github.com/pdnootboom/MCEocene>).

327 For initialisation of the eddying model, the three-dimensional ocean output at the
328 end of the CESM simulations (LR2 and LR4) is interpolated on the higher resolution
329 grid that the POP (HR2 and HR4) uses. We simulated 42 and 27 years in total for HR2
330 and HR4, respectively. Since we investigate the response of the simulations to an increase
331 in horizontal resolution, the same 5 model years of both HR2 and HR4 are used in most
332 analyses in this paper: year 23 to 27. For the same analyses of the low-resolution sim-
333 ulations (LR2 and LR4), we used the last 5 years of these simulations.

334 Using this setup of POP, we can investigate the sensitivity of simulations to the
335 studied resolution difference only, because the model is forced by the same atmosphere
336 and their geographic boundary conditions are based on the same reconstruction of (Baatsen
337 et al., 2016), and the three-dimensional eddying ocean is initialized by the equilibrated
338 output of the CESM. As a result, the atmosphere is representative of the middle-late Eocene
339 climate, but does not respond to changes in the ocean. We hence cannot investigate the
340 effect of atmospheric feedbacks on the results (den Toom et al., 2012; Rahmstorf & Wille-
341 brand, 1995; Arzel et al., 2011; Zhang et al., 2010).

342 The model set-up is suited to study the effects of model resolution on Eocene ocean
343 flows, but it is not suitable to study dynamics which involve atmospheric coupling, such
344 as the El Niño Southern Oscillation. The model set-up can best be used to investigate
345 the upper ocean circulation, as the deep ocean is not in equilibrium yet. Therefore, we
346 can only use this setup to obtain a transient response of the deep meridional overturn-
347 ing, not its equilibrium.

348 **4.3 Sinking Lagrangian particles**

349 To quantify sedimentary dinocyst endemism in the model, we applied a similar back-
350 tracking analysis of virtual sinking Lagrangian particles as in (Nooteboom et al., 2019)
351 (Supporting Information Fig. S4). This implies that we released these particles at the
352 ocean bottom and tracked them back in time while sinking and being advected by the
353 three-dimensional flow from POP, until they reached 10m depth. We released particles
354 on a $2^\circ \times 1^\circ$ grid of locations between 32-80°S every day for a year and waited until all
355 of the particles reached the near-surface (i.e. 17,520 particles in total). This analysis re-
356 quires a higher than monthly temporal resolution of model output (Nooteboom et al.,
357 2020; Qin et al., 2014). Therefore we used daily fields for the years 35-42 (HR2) and years
358 20-27 (HR4) to perform this backtracking analysis.

359 The used particle sinking speed of the Lagrangian particles in this paper is 6 m day^{-1} .
360 This represents a low sinking speed for single dinocysts (Anderson et al., 1985). We choose
361 this low sinking speed, because it is considered as a lower bound of the realistic sinking
362 speeds where most lateral transport occurs, which makes it easier to explain low abun-
363 dances of dinocyst species. However, this sinking speed could in reality be different due
364 to e.g. aggregation with other particles. We also applied a sinking speed of 25 m day^{-1}
365 (see Supporting Information Figure S6), which represent small aggregates (Nootboom
366 et al., 2019). The main conclusions on the model-data comparison do not change if 25
367 m day^{-1} instead of 6 m day^{-1} sinking speed is used.

368 The percentage of dinocyst endemism in the model is determined by the percent-
369 age of particles that originated from an environment with a temperature below \widehat{SST} (which
370 must be close to Antarctica; see Supporting Information Figure S4; similar approach as
371 in (Huber et al., 2004)). The percentage of modeled dinocyst endemism is not expected
372 to compare well with the percentage of measured endemic dinocyst, because this match
373 is also sensitive to the species-specific susceptibility of dissolution during the sinking jour-
374 ney and their productivity at the ocean surface (Nootboom et al., 2019). Therefore, we
375 compare whether any endemic species occur in sites (0% or $\geq 0\%$) between model and data
376 instead of the exact percentage.

377 We assume that the sinking Lagrangian particles are not greatly influenced by the
378 fact that the deep circulation is not in full equilibrium yet in the eddying simulations.
379 Most of the lateral particle displacement occurs near the surface which is in equilibrium
380 and where the currents are the strongest. Moreover, the eddying simulations are initialised
381 with output from the non-eddying simulations, which are in reasonable equilibrium. The
382 mechanistic development of the flow, given the heat and salt distribution from the ini-
383 tialisation, occurs in a few years (see also Fig. 3c-h). Hereafter, the flow changes slowly
384 and may equilibrate after ~ 1000 years due to the flow response to changing density dis-
385 tributions. The assumption that sinking Lagrangian particles are not greatly affected
386 by the deep ocean equilibration, is supported by the results that use sinking Lagrangian
387 particles in HR2 and HR4: These results are similar, even though the deep ocean cir-
388 culation is different in HR2 and HR4.

Acknowledgments

The code used for this work and the results are distributed under the MIT license and can be found at the website <https://github.com/pdnootboom/MCEocene>. The model data used to generate the main figures in this paper are publicly available on the Utrecht University Yoda platform: <https://doi.org/10.24416/UU01-AYNLZP>.

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Supporting Information for ”Strongly eddying ocean simulations required to resolve Eocene model-data mismatch”

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Contents of this file

1. Figures S1 to S10

Additional Supporting Information (Files uploaded separately)

1. Caption for Movie S1

Introduction These Supplementary Materials include 11 figures and 1 animation that support the results described in the main article.

Movie S1. Animation of sea surface temperature during the spin-up of the LR2 simulation.

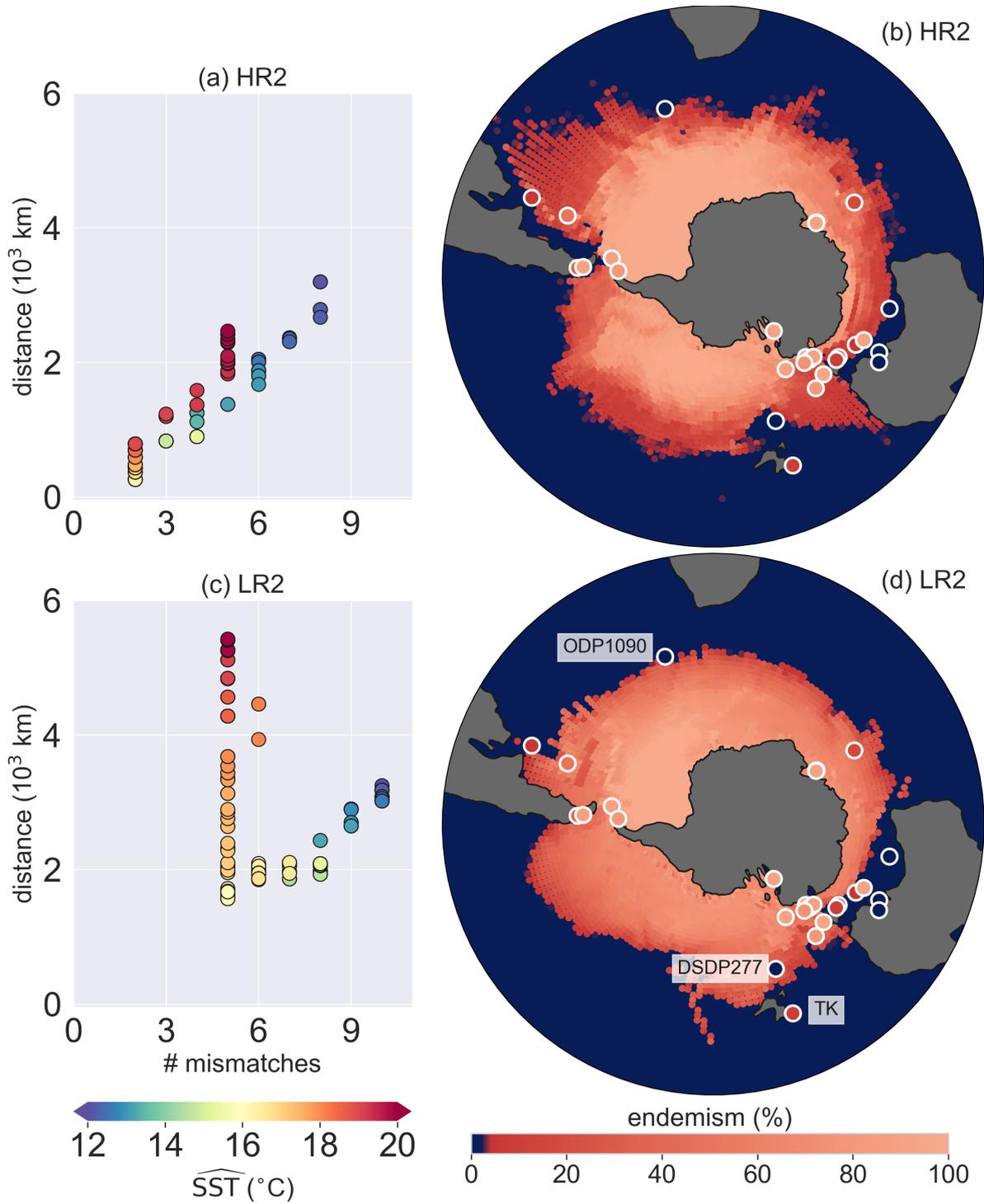


Figure S1. Same as figure 2, but for the $2\times$ pre-industrial case (LR2 and HR2).

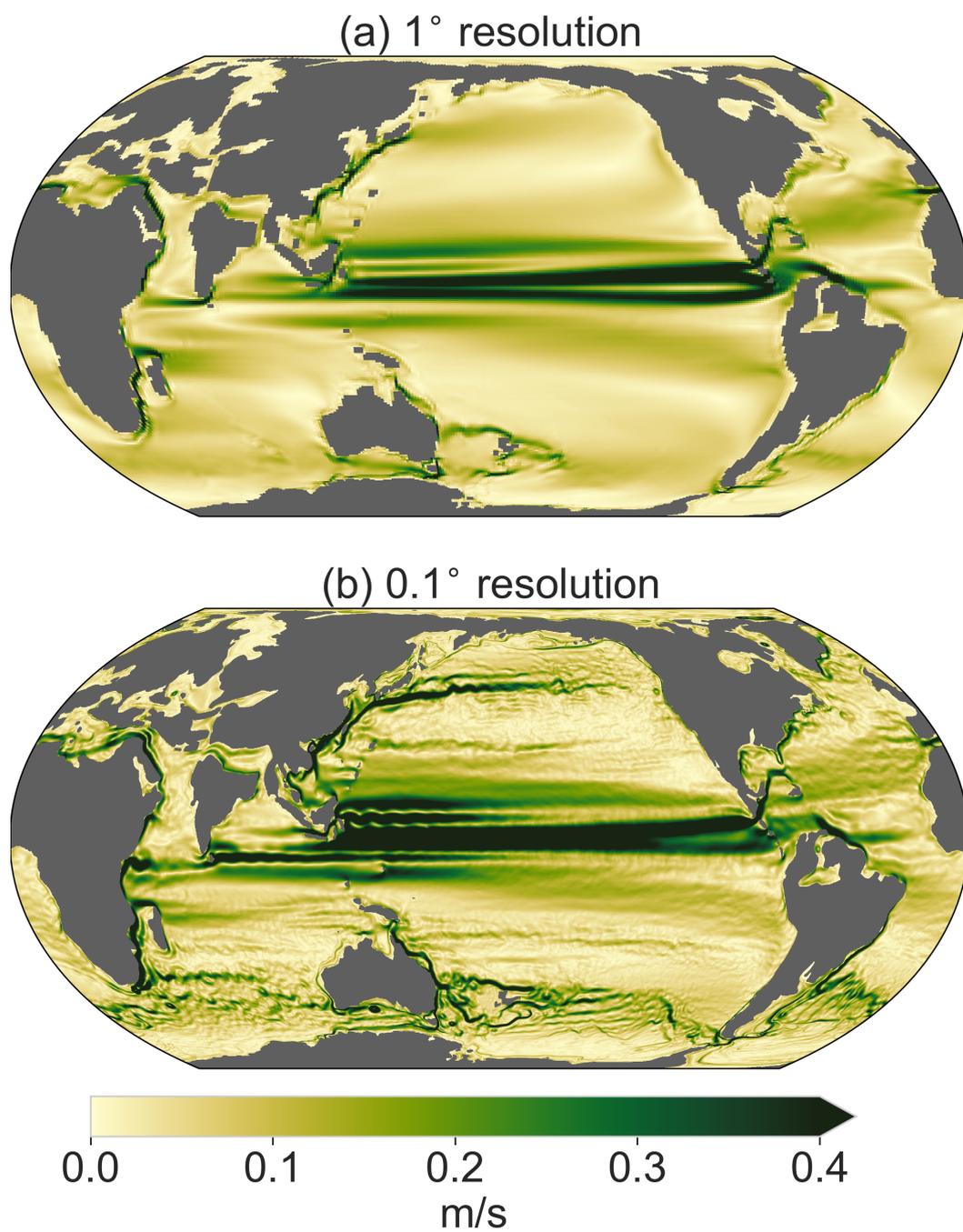


Figure S2. Same as figure 1, but for the 2×pre-industrial case (LR2 and HR2).

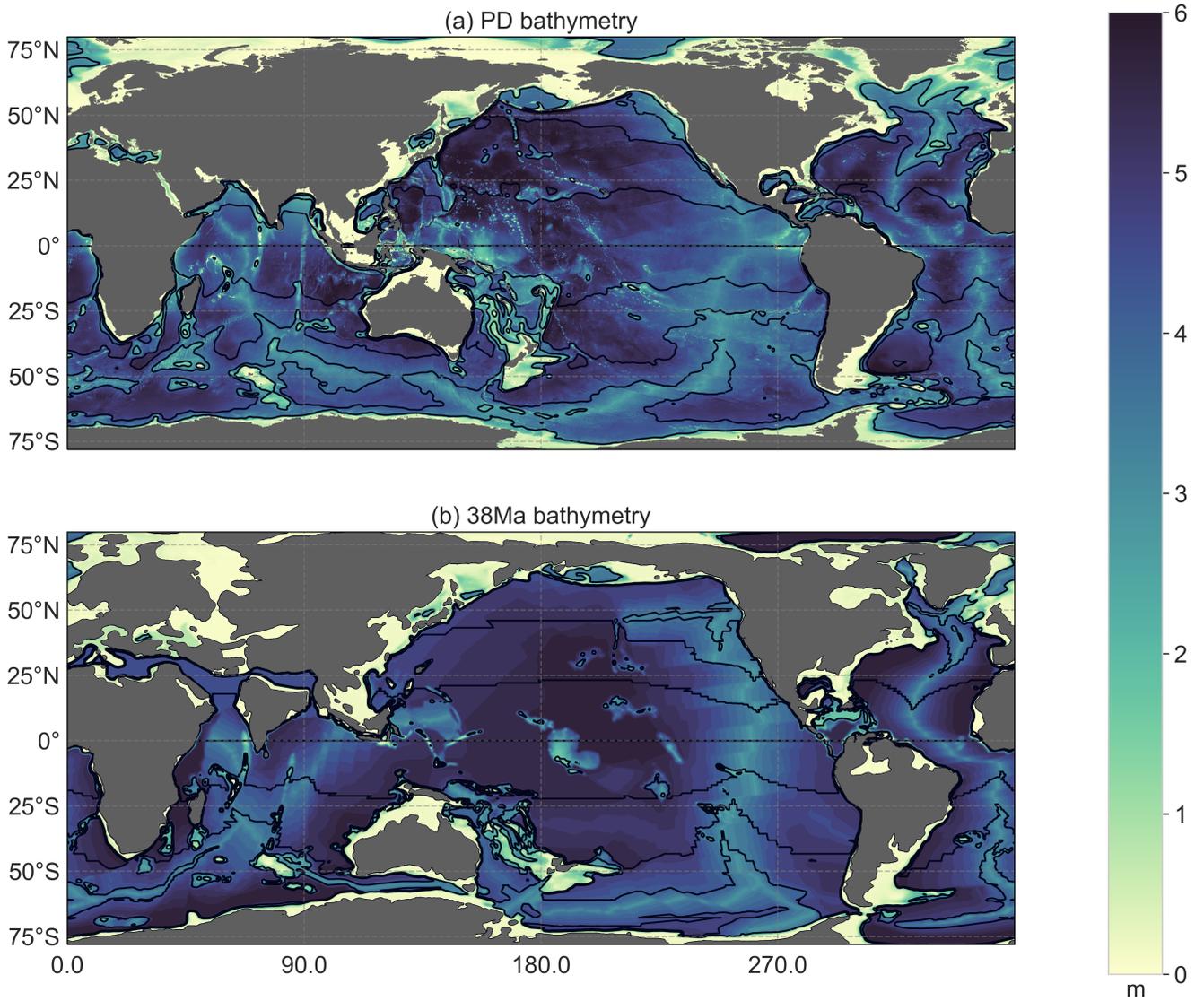


Figure S3. Global bathymetry in (a) the present-day (PD) and (b) the middle-late Eocene (38Ma). Black contours are lines of constant $\frac{f}{H}$ that the flow tends to follow in eddying simulations to conserve potential vorticity, with $f = 2\Omega \sin(\phi)$ the coriolis parameter (Ω is the rotation rate of the Earth and ϕ the latitude) and H the bathymetry.

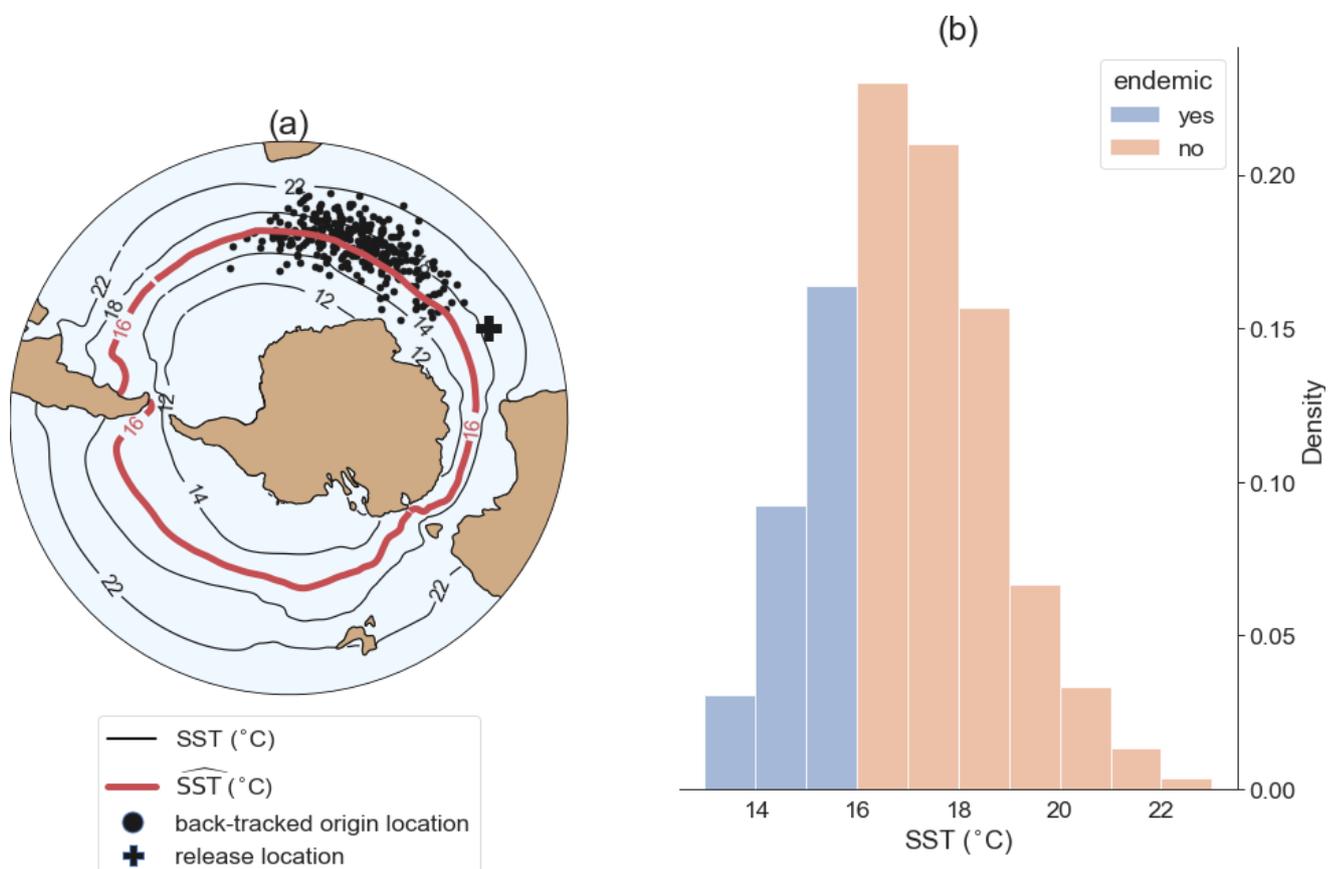


Figure S4. Illustration of the modelled dinocyst endemism near Antarctica. (a) Virtual particles are released at the bottom release location and tracked back in time with some sinking speed to determine their surface origin location. If the SST at the back-tracked origin location is lower than the threshold SST ($\widehat{SST} = 16^{\circ}\text{C}$ in this illustration), it is assumed to originate close to Antarctica, hence it is flagged as endemic. (b) A histogram of SSTs at the surface origin locations.

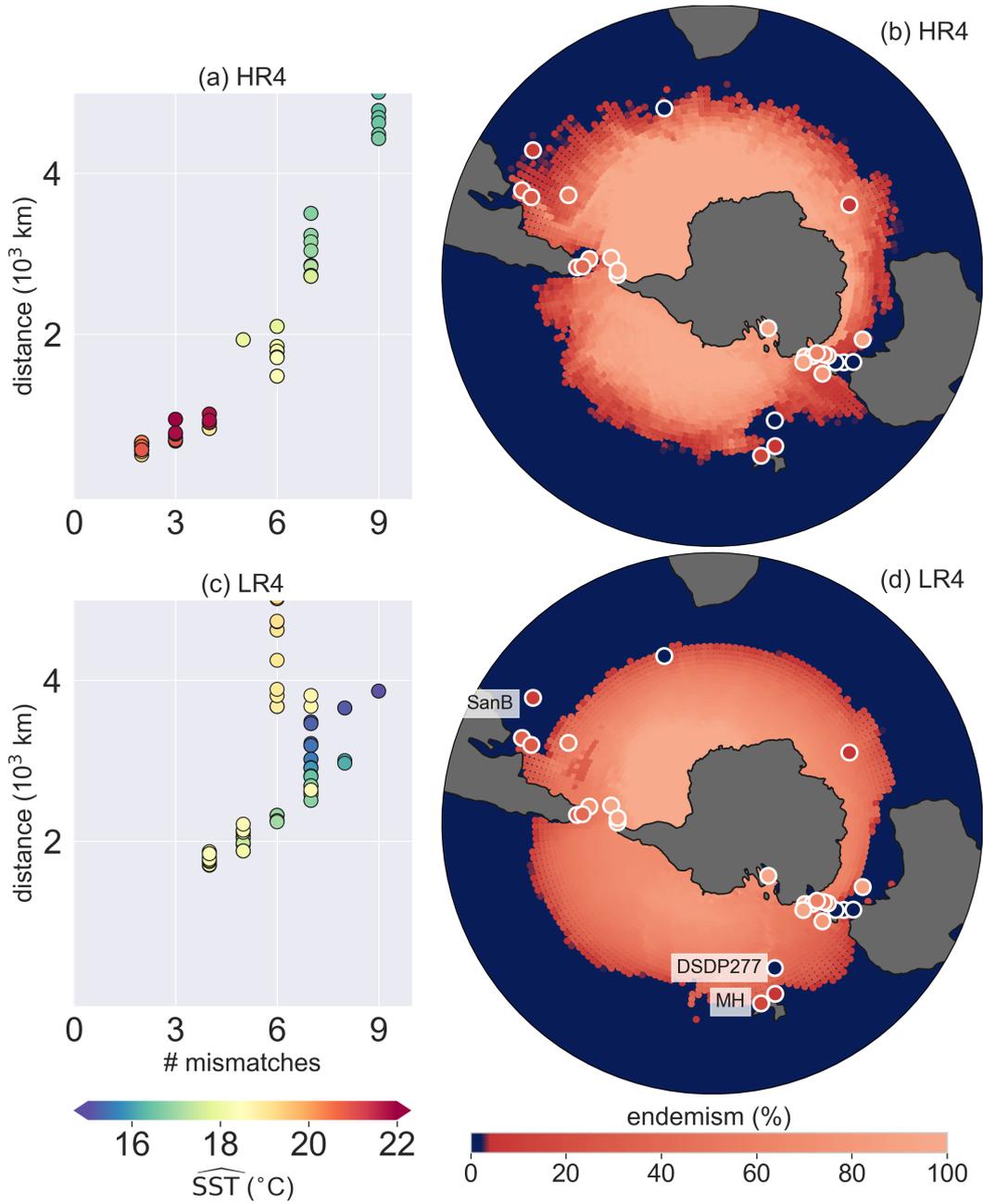


Figure S5. Same as figure 2, but with 25 m day^{-1} sinking speed.

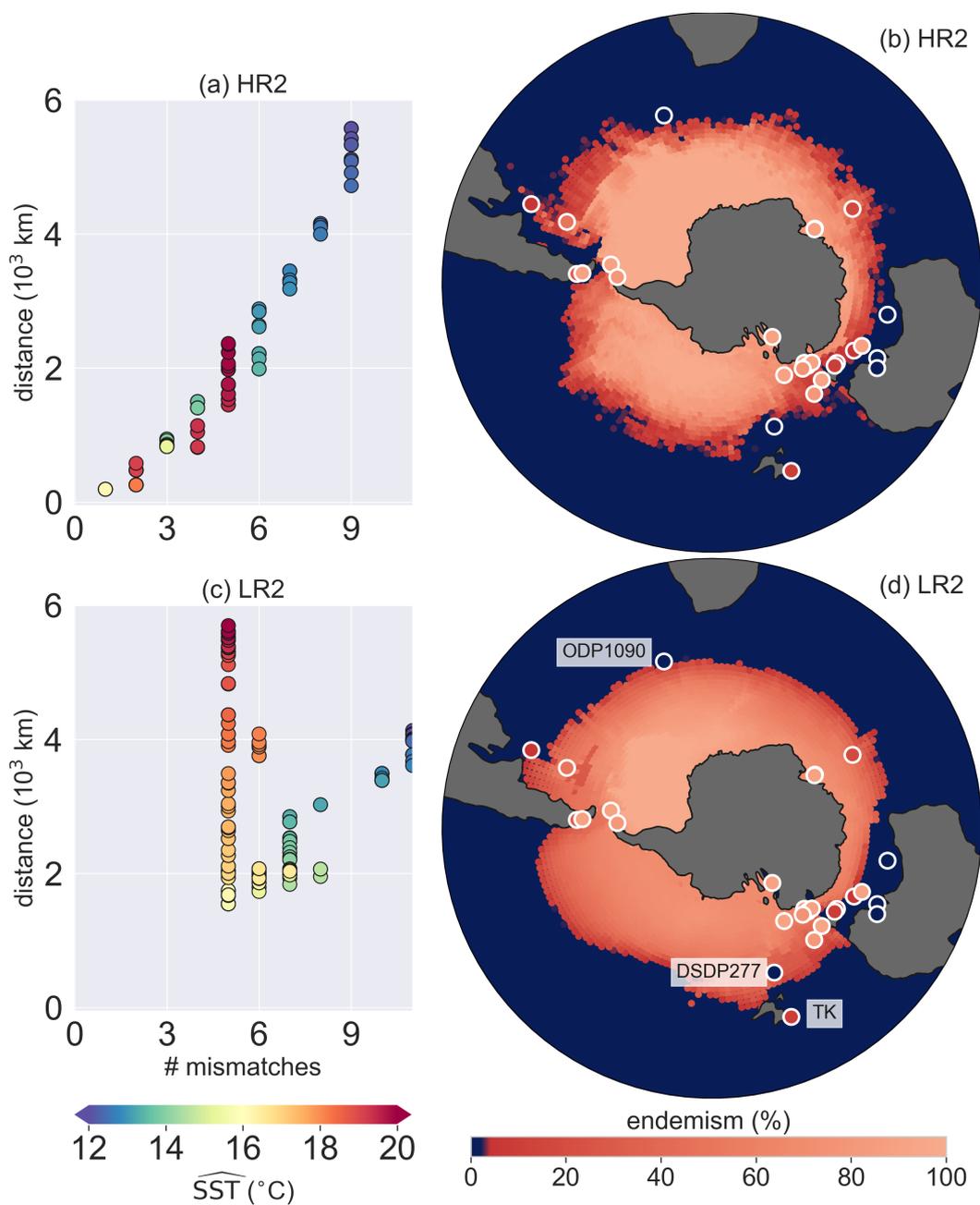


Figure S6. Same as figure 2, but with 25 m day^{-1} sinking speed and $2\times$ pre-industrial case (LR2 and HR2).

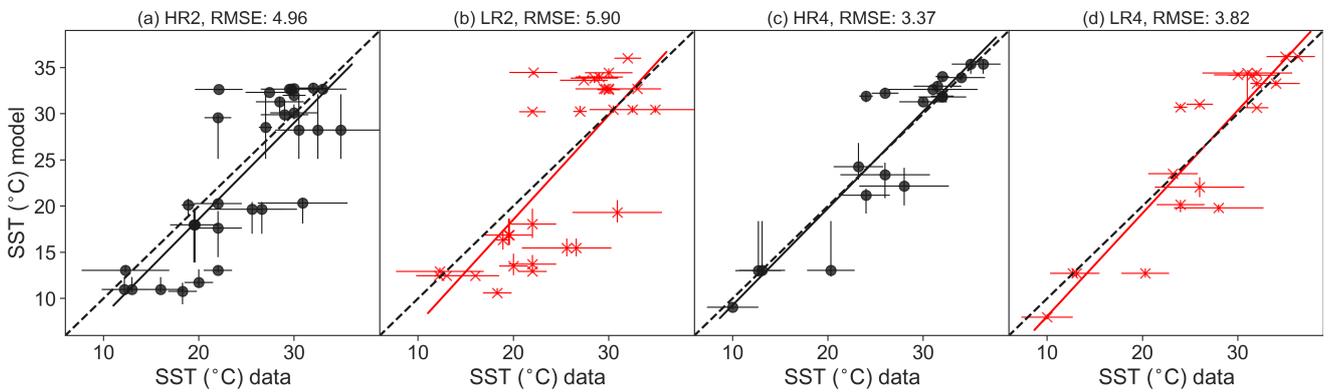


Figure S7. Same as figure 4g-j, but with a point-to-point comparison of model and data. The vertical uncertainty bars show the SST spread (minima and maxima) within a $4 \times 4^\circ$ box around the sites.

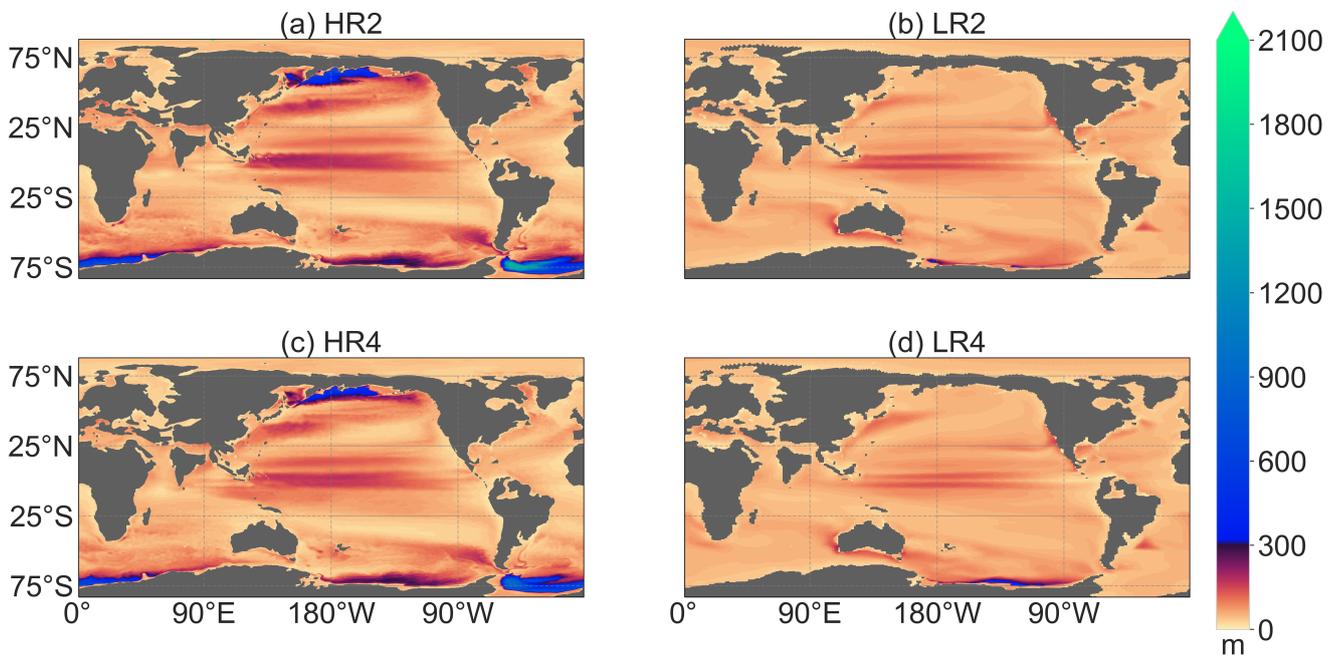


Figure S8. Maximum monthly mean of the mixed layer depth.

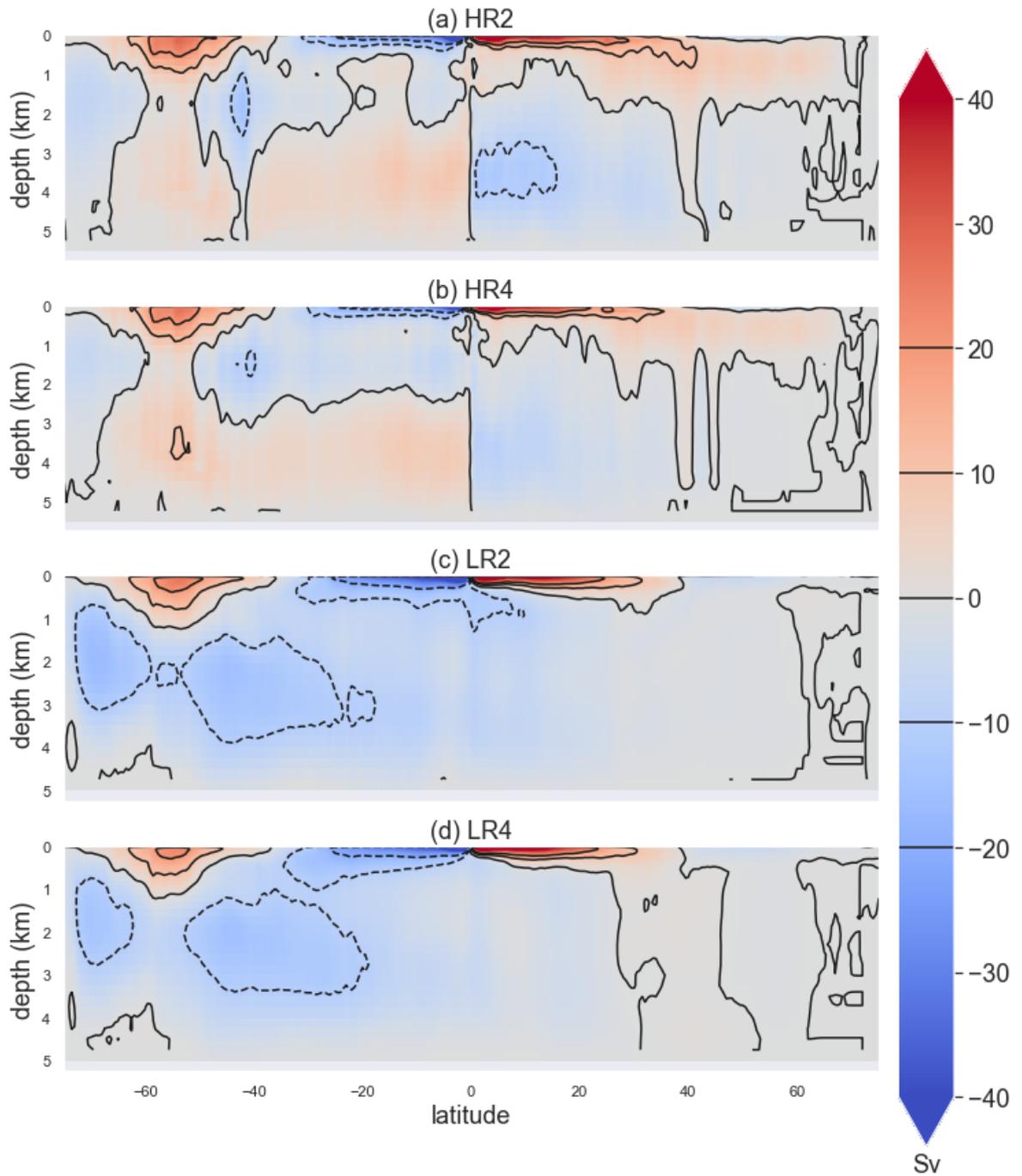


Figure S9. Meridional overturning stream function of the time-mean flow (over the same years as figure 1) in configuration (a) HR2, (b) HR4, (c) LR2 and (d) LR4.

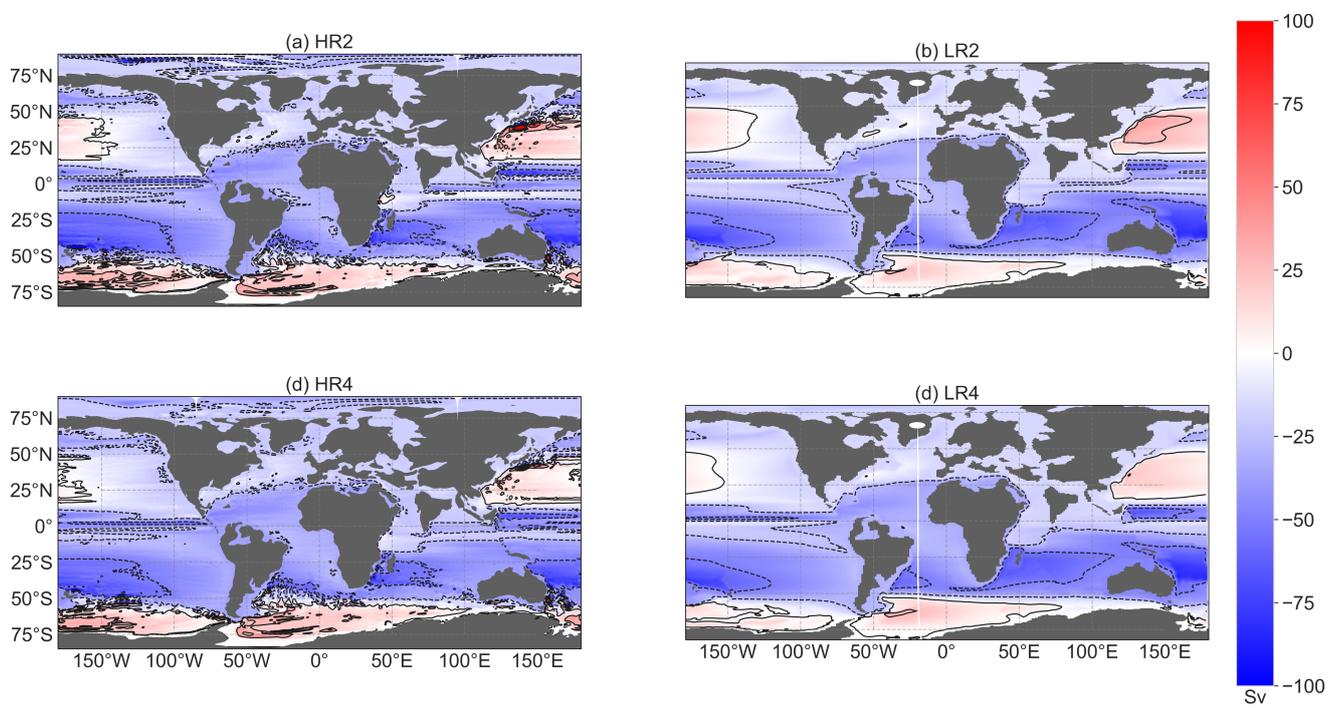


Figure S10. Barotropic stream functions of the annual mean circulation.