Labrador Sea winter heat and freshwater content observations from glider and Argo data

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Abstract

The Labrador Sea undergoes deep mixing in the wintertime, with mixed layer depths frequently reaching down to 2000 m. The resulting water mass that is formed - Labrador Sea Water (LSW) - has long been thought to be important for the deep Western Boundary Current (dWBC) and the upper limb of the AMOC. Direct observations of the overturning have, however, been rather limited. Limited Argo profiles and moorings in key locations offered winter measurements in a region challenged by severe weather conditions. Here we discuss observations of a winter-spring glider deployment in the Labrador Sea, but more specifically where deep convection occurs, from December 2019- to June 2020. Using the glider data, we describe the evolution of the mixed layer, changes in heat and freshwater content for surface (0-500 m) and intermediate depth (500-1000 m) layers for the central Labrador Sea convection region inside a box 200 by 100 km wide and spatial scales of T and S. We compare the observations with reanalysis data (air-sea heat fluxes and winds) and Argo profiles to better understand the variability missed by existing datasets. These observations highlight the role played by eddies in the overall variability of heat and salt in this region, something that is missed by Argo observations. They also show changes in spatial scales of T-S over the months from January to May, pointing towards the modulating effect of eddies on LSW winter formation.

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Key Points:

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9	• We use a glider survey and Argo profiles in the Central Labrador Sea during the
10	winter season 2020 to estimate mixed layer depths, heat and freshwater content
11	and the role of eddies.
12	• We find that overall convection during the 2019-2020 winter is weakened compared
13	to previous years, but that despite this trend individual Argo profile that convec-
14	tion reached to almost 2000 m .
15	• From the glider data we calculate correlation scales and find that scales, possibly
16	also due to the eddies, exhibit strong heterogeneity rendering local small-scale pro-
17	cesses important.
18	• We explore the discrepancy between Argo and the glider observations in terms of
19	the impact on the Labrador Sea freshwater budget, suggesting that the glider data
20	offers useful information not captured by the existing Argo float observations.

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21 Abstract

The Labrador Sea undergoes deep mixing in the wintertime, with mixed layer depths fre-22 quently reaching down to 2000 m. The resulting water mass that is formed - Labrador 23 Sea Water (LSW) - has long been thought to be important for the deep Western Bound-24 ary Current (dWBC) and the upper limb of the AMOC. Direct observations of the over-25 turning have, however, been rather limited. Limited Argo profiles and moorings in key 26 locations offered winter measurements in a region challenged by severe weather condi-27 tions. Here we discuss observations of a winter-spring glider deployment in the Labrador 28 Sea, but more specifically where deep convection occurs, from December 2019- to June 29 2020. Using the glider data, we describe the evolution of the mixed layer, changes in heat 30 and freshwater content for surface (0-500 m) and intermediate depth (500-1000 m) lay-31 ers for the central Labrador Sea convection region inside a box 200 by 100 km wide and 32 spatial scales of T and S. We compare the observations with reanalysis data (air-sea heat 33 fluxes and winds) and Argo profiles to better understand the variability missed by ex-34 isting datasets. These observations highlight the role played by eddies in the overall vari-35 ability of heat and salt in this region, something that is missed by Argo observations. 36 They also show changes in spatial scales of T-S over the months from January to May, 37 pointing towards the modulating effect of eddies on LSW winter formation. 38

³⁹ Plain Language Summary

In this paper we describe the Labrador Sea winter-time convection period based on ob-40 servations from an ocean glider and Argo floats. The data reveal that mixed layer depths 41 reached nearly 2000 m in the winter of 2020. There is good agreement between the Argo 42 and glider data, although the glider resolves spatial-temporal dynamics not captured by 43 Argo. Such variability is important to resolve the amount of freshwater in the Labrador 44 Sea. Freshwater content is important for different reasons, one of which is its role in con-45 straining convection. Because the Labrador Sea is a key place to understanding the global 46 overturning circulation, our study suggests that glider data provide useful information 47 not captured by existing Argo float observations. 48

49 **1** Introduction

The Labrador Sea plays an outsized role in influencing the large-scale climate cir-50 culation as one of the few regions where water convects to depths of 2000 m. Winter con-51 vection driven by strong winds and storms drives oceanic mixing frequently exceeding 52 depths of 1000 m and large water mass formation of Labrador Sea Water (LSW). Through 53 subduction and outflow within the the deep-Western Boundary Current (dWBC), LSW 54 is present in much of the North Atlantic and beyond. Historically, convection in the Labrador 55 Sea was thought to serve as the downwelling limb of the Atlantic meridional overturn-56 ing circulation (AMOC), and that its intensity would be related to the strength of the 57 overturning (Clarke & Gascard, 1983; Aagaard & Carmack, 1989). However, OSNAP measurements (Lozier et al., 2019; Li et al., 2021) appear to show only a minor role for 59 Labrador Sea convection, and comparisons between overturning transport at 45N and 60 air-sea transformation to the north of 45N (Desbruyères et al., 2019) suggest that the 61 reason convection and LSW formation do not imprint on transports is because the wa-62 ters entering the Labrador Sea - prior to convection - are already denser than the wa-63 ters of density of maximum overturning. While the relationship between deep water for-64 mation and the MOC remains uncertain, it is clear that the ventilation that occurs in 65 regions of deep water formation plays a role of hotspots for the storage of anthropogenic 66 carbon (Khatiwala et al., 2013) and the supply of oxygen to the deep ocean (Körtzinger 67 et al., 2008; Koelling et al., 2017). 68

The large scale processes connected to deep convection and formation of LSW have been described in various studies (Clarke & Gascard, 1983; Lilly et al., 1999; Yashayaev

& Loder, 2016). These studies are in general agreement over the location and mecha-71 nisms behind LSW water formation in the months from February to April when heat loss 72 reaches its maximum. These studies in also point to significant variability on annual and 73 decadal scales in the properties and volume of LSW. It is also accepted that significant 74 water property (T-S, gasses, nutrients) modification occurs at sub-mesoscale length scales 75 (Tagklis et al., 2020). Winter storms on 5–10 day timescales (Sathiyamoorthy & Moore, 76 2002) can trigger heat loss exceeding 1500 W m⁻². Boundary currents generate fresh 77 and salty eddies shedding in regions of steep topography (Lilly et al., 2003; Lilly & Rhines, 78 2002). These eddies can be sources of salt or heat that can alter stratification and ei-79 ther add or remove buoyancy, decreasing or enhancing convection. Not all eddies form 80 at the shelf break with some formed as a result of local instability during convective mix-81 ing. These still could be responsible for significant contribution to winter heat loss (Gelderloos 82 et al., 2011). Beyond the smaller scale features, the oceanographic conditions in the Labrador 83 Sea change on annual and multi-annual time-scales, with convection intensifying in cer-84 tain phases and weakening in others (Clarke & Gascard, 1983; Gascard & Clarke, 1983; 85 Yashayaev, 2007; Yashayaev & Loder, 2016). The important role and influence of this 86 ocean region supports and ongoing observing presence. 87

The Labrador Sea is well-known for its extreme winter weather and sea states 88 with significant wave heights regularly exceeding 10 m, winter winds stronger than 30 89 m s⁻¹ and atmospheric temperatures below -20 °C (Renfrew & Moore, 1999; Moore et 90 al., 2008). These conditions pose challenges for any observational programs. Prior to the 91 Argo program, most of the data for the Labrador Sea came from moorings and hydro-92 graphic ship cruises. These programs, in particular the WOCE, provided seasonally-biased 93 (towards non-winter seasons) snapshots or point-measurement time series (BRAVO) but it was not until the Argo program that a more holistic view of the Labrador Sea emerged. 95 The presence of floats in all of the Atlantic, allowed exploration of the larger-scale pic-96 ture of LSW spreading (Fischer et al., 2018), however, Argo does not close all gaps in 97 tracking LSW. For one, Argo is designed to capture seasonal variability at a 3° by 3° spa-98 tial scale, and to this end, each float spends approximately 10 days at 1000 m between 99 profiles. These scales are much larger than eddies that are thought to be important for 100 the heat exchange (Hátún et al., 2007) and larger than the winter mixing plumes at these 101 latitudes (Mertens, 2000). Modelling studies such as Bailey et al. (2005) suggest that ac-102 curate knowledge of heat budgets in the Labrador Sea is vital to properly constrain deep 103 water formation. Floats are also excluded from the continental shelf, ice covered and most 104 of the shelf areas. Moored observations provide sampling that is more frequent in time. 105 Several long term moorings exist in the Labrador Sea with the German K1 mooring in 106 the central Labrador Sea (2004-present) e.g. Avsic et al. (2006) and moorings nearer to 107 the shelf break by Fisheries and Oceans Canada (e.g. Yashayaev & Loder, 2016). Ship 108 cruises can provide sampling in higher spatial resolution, but they can be rather treated 109 as snapshots in time and they are typically limited to summer season, meaning that the 110 winter periods are undersampled. Satellites can provide higher spatial resolution of e.g. 111 sea surface temperature at daily or higher frequencies, but cloud coverage (pervasive dur-112 ing the cold winter months) limits visibility in the infrared and visible range of the spec-113 trum. In addition, subsurface data are required to assess the strength of wintertime mix-114 ing and to understand the export of newly convected water. 115

Gliders can help to close this observational gap because they can operate in win-116 ter conditions and can be directed to sample in particular areas of interest. Glider data 117 quality is improving but to achieve more confidence in the data, glider observations should 118 be coordinated with other platforms, as glider sensor payloads are limited by the strict 119 power considerations of the platform. To target a particular oceanic process it is also im-120 portant to separate the inherent smearing of time and space signals in glider data, given 121 that they move at most at 25 km day⁻¹ (e.g. Rudnick, 2016). In the Labrador Sea, glid-122 ers have been successfully deployed on several occasions (Hátún et al., 2007; Eriksen & 123 Rhines, 2008; Frajka-Williams et al., 2009, 2014; von Oppeln-Bronikowski et al., 2021) 124

and they have shed new light on the dynamics in this region, particularly on winter cooling, mixing, density stratification and role of eddies, to name a few.

In this study, we deployed a glider in the wintertime into the Labrador Sea (Figure 1). The focus of the sampling was on understanding the wintertime exchange of heat and salt, in particular the role of transient processes such as eddies and storms for the mixed layer depth formation and ocean heat storage. This glider was stationed in the deep convection zone in the western part of the Labrador Sea (see Figure 1b). In this paper we address three questions:

- To what extent do small scale processes contribute to the cycle of heat and salt
 exchange in this region and therefore the formation of LSW?
 - 2. What are the dominant scales of water property evolution in the Labrador Sea?
- In particular, do gliders offer critical information on scales below that offered by
 Argo floats?
- ¹³⁸ 2 Data and Methods

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139 2.1 Glider Data

We deployed a Slocum 1000 m glider into the Labrador Sea during the Winter 2019-140 2020. The glider sampled the deep convection area highlighted by the <800 m mixed layer 141 depth contours (Figure 1b). We use a similar geographic definition as in other studies 142 in the area (see Yashayaev (2007); Yashayaev and Loder (2017)). In this study we present 143 observations from January 15 to May 20, 2020. The glider sampled with a mean spa-144 tial resolution of 1.5 km and 2-hr between profiles (Figure 1a). The details of this mis-145 sion are described in de Young et al. (2020), including the challenges and many lessons 146 learned as part of this long-duration (7-month) mission in a harsh and remote environ-147 ment. 148

The Slocum glider data were processed following basic quality control procedures 149 recommended by the Australian National Facility for Ocean Gliders (ANFOG) Integrated 150 Marine Observing System (IMOS) best practices document (Woo, 2019). The processed 151 T-S profiles are shown in Figure 2. We converted raw glider data to level-1 and level2 152 data sets using the SOCIB toolbox (Troupin et al., 2015) following ANFOG/IMOS best 153 practices. We used the ANFOG/IMOS QC manual for data flagging, linear interpola-154 tion of longitude and latitude, time vectors filling gaps, outlier detection, and spike re-155 movals. Profile identification and splitting was done using the pressure inversion method. 156 Product profiles of T and S were bin-averaged on a 1 m depth grid going from 0 m to 157 max profile depth (1020 m). Thermal lag was corrected with respect to T in pumped 158 CTD cell following the standard correction from Morison et al. (1994) using 8 ml s⁻¹ 159 flow speed inside the conductivity cell. The exact procedure for glider pumped CTDs 160 is described in the ANFOG/IMOS manual. Absolute salinity, conservative temperature, 161 potential density are calculated with the TEOS10 toolbox (McDougall & Barker, 2011). 162 Gridded profiles were corrected for up/downcast mismatch using RMS minimization de-163 scribed in the ANFOG/IMOS manual, median delay $(T=2.39s^{-1}, C=2.05s^{-1})$ was used 164 for all variables. For the final data set we used a Savitzky-Golay filter for salinity (7 steps, 165 1st order) to remove occasional spikes in the lag corrected and up/downcast mismatch 166 corrected profiles. 167

168 2.2 Ancillary Data

We used ERA5 global reanalysis (Hersbach et al., 2020) winds and surface forcing with 1-hr temporal resolution and $1/12^{\circ}$ grid spacing resolution. Data were downloaded from the Copernicus Climate Data Store (CDS) website (Hersbach et al., 2018). We estimated the surface heat fluxes from ERA5 by including latent (Q_l) and sensible



Figure 1. (a) shows the spatial and temporal resolutions between glider profiles with the pink dot indicating the mean. (b) Location of glider profiles (green from January 15 - May 15, 2020) and of Argo profiles (grey stars) taken over the same period and within 60 km of the glider profiles. Background color is the winter (January 1 - March 31) mixed layer depth extracted from the extended Roemmich-Gilson (Roemmich & Gilson, 2009) Argo climatology (2004-2020) based on the 0.05 kg m⁻³ density change criterion. Isobars (1000 m spacing) are superimposed in magenta.



Figure 2. (a) T and (b) S profiles (x-axis is a profile number and the time) from the glider corresponding to the highlighted green track (Figure 1b) from January 15 to May 15, 2020. Gaps are periods when the glider was not sampling due to problems with the onboard computer. Contours of the 27.71 and 27.73 kg m⁻³ potential density surfaces are shown in black. Towards the end of April - beginning of May a strong freshening and warming is due to an eddy that trapped the glider for nearly 2 weeks.

 (Q_s) turbulent heat fluxes, as well as the net longwave (Q_{lw}) and shortwave radiation (Q_{sw}) . We calculated the net surface heat flux (Q_{net}) in W m⁻², where negative values indicate a net heat loss from the ocean to the atmosphere and positive values indicate a net heat gain by the ocean from the atmosphere:

$$Q_{net} = Q_l + Q_s + Q_{sw} + Q_{lw} \tag{1}$$

Argo data were extracted from Argo Global Data Assembly Centre (GDAC) (Argo, 177 GDAC, 2000). We extracted all Argo profiles (total of 2678) from 2002 to 2020 between 178 44°W and 65°W and 44°N and 66°N. Data were checked for QC using the Argo DAC 179 guide (Wong et al., 2020). Outliers were removed for T and S outside the ranges -2 < -2180 $T < 25^{\circ}$ C and 30 < S < 35.5. For delayed time mode data, only data with quality 181 flags of 1 or better were selected. In winter-spring 2019-2020, there were seven separate 182 Argo floats present (WMO: 3901668, 4902395, 4902469, 4902471, 4902478, 4902481, 6902684) 183 in this region within 60 km of the glider profiles during most of the glider sampling pe-184 riod. These floats collected a total of 48 profiles. The mean separation between profiles 185 for the Argo floats (profiling every 10 days) was 55 km, compared to 1.5 km for the glider. 186

2.3 Mixed Layer Depth Estimation

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We estimated the mixed layer depth using a threshold criterion of d = 0.01 kg m-188 3 density change for glider and Argo profiles, finding the depth Z_{MLD} , where $\Delta \rho = \rho(z(i)) - \rho(z(i))$ 189 $\rho(z=10) < 0.01 \text{ kg m}^{-3}$. We ignored profiles shallower than 10 m, because the glider 190 was inflecting 12 m below the surface and because mixed layer depths during Jan-May 191 were all greater than 10 m. The shallowest mixed layer depth we computed for the glider 192 using this method was 30 m in April and the deepest was 1020 m, the maximum dive 193 depth of the glider. The glider probably underestimated the MLD from Feb 20 to April 194 5 (Figure 3) as in this time the MLD likely exceeded 1000 m. Using this density thresh-195 old, we find MLD comparable with those reported in the literature (Yashayaev, 2007; 196 Körtzinger et al., 2008; Yashayaev & Loder, 2016) for the Argo observing period . The 197 choice of criterion or method has a large effect on calculated values of MLD as discussed 198 and compared in Holte and Talley (2009). As an example, we found that MLD could vary 199 between 373 m ($\Delta \rho = 0.01$ kg m⁻³) and 848 m ($\Delta \rho = 0.05$ kg m⁻³) for an identical T-S 200 profile. Applying the hybrid method from Holte and Talley (2009) usually returned val-201 ues close to those estimated with the d=0.01 kg m⁻³ criterion. A detailed overview and 202 sensitivity analysis is beyond the scope of this study. We justified our choice of MLD den-203 sity criterion based on how well our estimates aligned with those in the literature to pro-204 vide a better comparison. 205

2.4 Vertical Heat and Freshwater Content

To estimate the Ocean Heat Content (OHC) per unit area ($a = 1 m^{-2}$), we followed Boyer et al. (2007) (their Equation 2),

$$OHC = a \int_{z_1}^{z_2} \rho_m C_p (T_c - T_m) dz$$
 (2)

Here OHC is the ocean heat content (J / m⁻²), ρ_m and T_m are the reference den-209 sity (1027.3 kg m⁻³) and temperature (3.2°C) for the upper Labrador Sea Water (uLSW) 210 (Rhein et al. (2007) their Figure 4c), respectively. C_p is the heat capacity of seawater 211 $(4000 \text{ J kg}^{-1} \text{ C}^{-1})$. We integrated temperature profiles for two layers 0-500 m and 500-212 1000 m, setting z_1 and z_2 appropriately for both Argo and glider profiles. In Boyer et 213 al. (2007) they integrated in-situ temperature but we used conservative temperature T_c 214 calculated from absolute salinity, in-situ temperature and pressure (McDougall & Barker, 215 2011). 216

To estimate the Freshwater Content (FWC), we followed the method of Florindo-López et al. (2020) (Equation 2), except that we are switching around the terms to interpret positive values as an addition of FWC (fresher) and negative removal of FWC (saltier). As before we integrate independently from z_1 to z_2 , for both the top 0–500 m and intermediate 500–1000 m layer.

$$FWC = a \int_{z_1}^{z_2} \frac{\rho(S, T, P)}{\rho(0, T, P)} \frac{(S_m - S)}{S_m} dz$$
(3)

As above, ρ is the seawater density. We used (S) absolute salinity and not practical salinity as in the original method, as well as the reference Salinity (S_m) of 35 g kg⁻¹ which is the mean salinity of the upper 1000 m of the study area from Jan 15-May 15, 2020 to look at the anomaly in salinity over the water column.

226 **2.5** Correlation Scale Calculations

We use the glider data to estimate the correlation length scales of the observed T 227 and S fields for each glider transect to investigate the spatial structure of dynamic events 228 (eddies, convection, storms) during this period. This is difficult with the sparse Argo pro-229 files, but we can use the higher resolution repeated glider tracks in each month to char-230 acterize T and S correlation lengths scales over the potential density layer where the new 231 Labrador Sea water forms through the mixing of intermediate-depth and surface cooled 232 waters (top 1000 m). To quantify the spatial scales, we used the definition of the auto-233 correlation function r(k) following the results in von Oppeln–Bronikowski et al. (2021) 234 for any isopycnal layer of interest. 235

$$r(k) = \frac{\sum_{t=1}^{N-k} (x_t - \bar{x})(x_{t+k} - \bar{x})}{\sum_{t=1}^{N} (x_t - \bar{x})^2}$$
(4)

Here x_t denotes measurements of T and S along with the isopycnal layer at step 236 t, and \bar{x} is the mean value along the spatial dimension k, N is the total number of sam-237 ples. Before calculating the spatial correlation functions we average T-S data along the 238 transect inside a particular isopycnal layer (± 0.01 kg m⁻³) and detrended the T-S data, 239 to remove non-stationary spatial trends. Equation 4 then gives us the autocorrelation function as a function of lags. We use the lag at the first zero-crossing as the correla-241 tion length scale (see Supplement Figures S2 and S3). In our calculations we spatially 242 reference the glider data to the westernmost measurements of the glider so that distances 243 are consistent between transects. 244

245 3 Results

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3.1 Argo-Glider Heat and Freshwater Content

We calculated the OHC and FWC from Argo and glider profiles following Equations 1 and 2, for two depth layers from 0-500 m and 0-1000 m. We did this for every glider and co-located Argo profile in the region. We chose these two layers to determine the difference and relative contribution of each layer to heat and freshwater change in the water column. ERA5 data have been extracted closest to the glider track positions (1-hr time steps).

The ERA5 surface heat fluxes (Figure 3a bars, black line) track the general cycle of cooling (-200 W m⁻²) in winter (January - early April) with a gradual warming (200 W m⁻²) towards spring (April-May) reported in other studies (Straneo (2006), their Figure 1a). The shift from predominantly cooling to warming is noticeable as well as a rather abrupt shift near the end of April after which cooling events subside, except a few events



Figure 3. (a) Net heat flux from ERA5 with red shading for wind speeds from a northern direction in excess of 50 km hr-1 and (b) glider and Argo mixed layer depths. (c) and (d) OHC and FWC content calculated from glider and Argo profiles for the joint sampling period. OHC and FWC are shown for two depth layers from the surface to 500 m (orange) and 500-1000 m (bluish) with a glider (line) and Argo profiles (markers and dashed)

lasting only a few hours. Inside these phases of cooling and warming there is a lot of vari-258 ability with sudden spikes of warming in excess of 700 W m^{-2} noticeable as deviation 259 from the daily average (black line). Storm systems frequently track through the Labrador 260 Sea from a Northern direction that import cold air and enhance cooling or so-called cold 261 air outbreaks (Moore et al., 2014). We color-coded negative heat fluxes when the air was 262 coming from a northern direction and winds were in excess of 50 km hr^{-1} and appear 263 to be influenced by a 2 week time-scale as described by Sathiyamoorthy and Moore (2002). 264 Overall the lowest heat fluxes in our data do not exceed -500 W m^{-2} . This is less than 265 previously observed winter-minimums (-1500 W m⁻²) in the vicinity (Lilly et al., 1999). 266

The time series of measured mixed layer depths (Figure 3b) reveal a deepening of the mixed layer during this period of heat loss, with estimated MLD from the glider and Argo floats agreeing well until March. After this time, the glider turnaround depth of 1020 m is too shallow to capture the full extent of mixing. During March, mixed layer depths from Argo profiles exceeded 1500 m, with the deepest measured mixed layer depth of nearly 1900 m deep near the end of March (Figure 3b). The onset of shoaling MLD is rather sudden (within 1 week) after which the MLD is generally less than 200 m. The



Figure 4. (a) Implied Argo, glider and ERA5 13-day averaged heat fluxes. Glider and Argo heat flux is the derivative of the OHC (0-500 m layer). Background shading is the upper and lower bounds based on the standard deviation from the bin averaging. (b) Correlation results Argo and Glider vs. ERA5 for different bin-sizes with 13-day bin size highlighted by vertical dashed black line. (c) Correlation of glider, Argo and ERA5 fluxes for the 13-day bin size. Line in the background shows 1:1 fit.

higher sampling resolution of the glider captures some spatio-temporal variability of the
mixed layer depth not observed by the sparse Argo profiles. Filtering the glider data with
a low-pass filter for time periods shorter than 3-days (blue in Figure 3b) removes those
features from the glider record.

OHC provides a different measure of the intensity of watermass transformation than 278 mixed layer depth (Figure 3c). In the near surface layer (0-500 m), the OHC started de-279 creasing from the start of the record (January) reaching a minimum in the first week of 280 February. In the intermediate depth 500-1000 m a decrease in OHC was only noticeable 281 after Feb 10, reaching a minimum around March 1. OHC estimates from glider and Argo profiles coincide well, however the high temporal and spatial sampling of the glider cap-283 tures much greater episodic variability compared to the Argo profiles. OHC in both lay-284 ers becomes nearly equal after March. OHC at different depths are quite different in Febru-285 ary when convection intensifies. This suggests that convection may start in the upper 286 layer before affecting the deeper layer as expected from convection forced by surface heat 287 flux. 288

On the other hand, FWC shows distinct differences between the two integrated layers. The top 500 m (orange, Figure 3d) is saltier compared to the mean reference salinity (35 g kg⁻¹), while the 500-1000 m layer (blue, Figure 3d) is initially fresher than the surface layer. As the mixed layer deepens, and the two layers mix, the FWC in the surface layer decreases (water becomes saltier) while the FWC in the deeper layer increases (water becomes fresher). The FWC of the two layers merges in March during the period of the deepest mixing. Some short-duration events appear $(1 \text{ m}^{-1} \text{ drop in FWC for}$ in the surface layer Feb 3-11) in the glider time series with the removal of FW on the order of 1 m³ in the timespan of a day. These apparent salinifying events are associated with the glider passing through eddies with different water masses and FWC than the surrounding water. These salinity altering events extend to the bottom layer, pointing to their energetic nature.

We are interested in how well the sub-surface glider and Argo data resolve surface driven processes such as wind-driven surface mixing. We expect that by calculating the correlation between surface fluxes (ERA5) and inferred surface fluxes from the changes in OHC (Argo and glider) the residuals of this correlation may be partly explained by wind-driven surface mixing. To estimate implied fluxes from the change in observed ocean heat content $Q_{net,obs}$, we take in Figure 4a the derivative in time

$$Q_{net,obs} = dOHC/dt \tag{5}$$

where dOHC is the difference of OHC over the time interval dt for both observed 307 profiles. Here we take the surface layer (0-500 m) as our best estimate with surface fluxes 308 from ERA5. We bin average results to a common time scale between the Argo and glider 309 observations. Argo profiles are spaced approximately every 10–11 days. The Argo data 310 is sparse and 10-days is the practical limit over which we can compare surface processes 311 with Argo data. Indeed, the correlation between ERA5 and implied heat fluxes from Argo 312 reveal a spike in correlation for 10 day averaging and a steep drop for values smaller than 313 that (Figure 4b). However, this story is more complicated for correlation between ERA5 314 and glider data. For small averaging windows (<10 days) the data do not agree well with 315 $r^2 < 0.2$. On 13-day time scales (Figure 4c) we find that Argo ($r^2 = 0.82$) and glider 316 data $(r^2 = 0.68)$ are in overall agreement. During this averaging interval the glider cov-317 ered approximately 154 km. Thereafter, the correlation between ERA5 and Argo improves 318 steadily the larger the window size and on monthly scales (30-day) they agree over 90%. 319 This makes sense given that increasing window size filters out small scale variability and 320 that ERA5 data assimilates data from Argo. In contrast, larger averaging windows do 321 not necessarily produce better correlation between implied heat fluxes from the glider 322 data and ERA5. There are spikes and drops for different windows with 25-day rising $r^2 >$ 323 0.8 and 31-day dropping $r^2 < 0.2$. 324

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3.2 Winter 2020 Convection and Correlation Scales

The Argo time series goes back to 2002, but significantly more profiles are avail-326 able from 2012. Yashayaev and Loder (2016) (their Figure 4d) show a trend of increas-327 ing density for the thickest density layer from 2012-2016 with the 2016 Labrador Sea wa-328 ter being denser and colder than previous convection classes such as the LSW 1994, 2008 329 classes (Yashayaev, 2007; Yashayaev & Loder, 2016). We extend in Figure 5 the anal-330 ysis from (Yashayaev & Loder, 2016) to the time of the glider observations to provide 331 context for the high resolution glider observations and put them into the context of other 332 studies. 333

We use the Argo profiles in the region near the glider observations with a box sim-334 ilar to Yashayaev and Loder (2016) (their Figure 1) from 2011 to 2020. We bin average 335 all Argo profiles inside this spatial box in 14 day and 0.005 kg m⁻³ density bins (Fig-336 ure 5a). A figure showing all Argo floats selected for the average is given in the supple-337 ment (see Supplement, Figure S1). Different from Yashayaev and Loder (2016) we use 338 the potential density to 0 dbar (σ_0 instead of σ_1). The result in Figure 5a shows an in-339 creasing trend in LSW density from 2012 to 2017, similar to the results in Yashayaev and 340 Loder (2016). The period of 2017 to 2018 shows density of LSW in the convective zone 341



Figure 5. (a) Isopycnal layer thickness diagram is broken down by year from Argo data (2012-2020) in 14-day by 0.005 kg m⁻³ bins. Box indicates range of glider observations. (b) Layer thickness for each month for glider observations in the Labrador Sea in 2020 in bins of 12-hr and 0.005 kg m⁻³.

nearly constant and then decreasing in 2019-2020. The onset of convection in 2019-2020
 appears to be also later compared to the previous five years.

Convection in 2020 from Argo data in the glider observing period (pink box in Fig-344 ure 5a) shows similar convection intensity to 2013 (Figure 5a) albeit with less thickness 345 observed in the density layers. The glider observations for the period of winter 2019-2020 346 (Figure 5b) show that the volume of density groups exhibit variability not seen by Argo 347 in terms of thickness of density layers and the variability of the density layer thickness. 348 There is an event of lighter water occupying a larger depth range from a salinifying but 349 warm feature (see Figure 3c-d), which later analysis (see Figure 8) shows to be an eddy. 350 The glider record stops at 1000 m and misses the densest layers of convected water in 351 March ($\sigma_0 > 27.75$) but shows well mixed layers for the entire water column occupy-352 ing the same density. The glider data also shows that events like an eddy, that was de-353 tected in February, results in water with different density than otherwise observed in the 354 record. The onset of spring and end of convection shows gradual spread of density layer 355 thickness as the water column re-stratifies. 356

One of the key questions around the heat and freshwater calculations (Section 3.1) 357 concerns the spatial homogeneity in the obtained results. Here we break the glider data 358 into individual transects and calculate the spatial correlation length scale as per meth-359 ods (Section 2.5). A sample autocorrelation function for an isopycnal is given in the sup-360 plement (Figures S2,S3). We repeat the analysis for every isopycnal group observed by 361 the glider range from 27.65 to 27.75 kg m^{-3} to find the correlation scales. The glider man-362 aged to do 16 transects during the 5-month period in the central Labrador Sea, each at 363 least 90 km long and lasting a week, with the majority 150-200 km long and lasting two 364



Figure 6. (a), (b), (c) Glider sections with eddy or ventilation events (left side) and corresponding correlation scales (right hand side) for T(blue) and S(red). Note the potential y-axis is reversed to align with the glider transects. White contour indicates the mixed layer depth, and black contours the potential density contours spaced every 0.01 kg m⁻³. (d) Inset map shows the individual tracks used for the glider section analysis.



Figure 7. Same as Figure 6 but for the glider sampling transects without eddies.

weeks. There were sections containing several eddy signatures from either the same or 365 different eddies (Figure 6) and there were sections that contained no apparent eddies (Fig-366 ure 7). The sections were picked for the stages of convection with the upper panel cor-367 responding to the preconditioning, middle panel with active convection and lower panel 368 with re-stratifying post convection. The presence of eddies influences the correlation length 369 scale, though it is not possible to accurately estimate the eddy size because of the slow 370 speed of the glider, causing spatial aliasing. For example the uppermost panel (Figure 371 (6) shows an early cyclonic eddy event around 110 km, the middle panel is the edge (150 372 km) of an eddy smearing the glider signal and the bottom panel is an energetic anti-cyclonic 373 eddy centred at 60 km that trapped the glider for nearly two weeks from end of April 374 to middle of May. The corresponding correlation scales of these sections containing ed-375 dies appear to show that the density layers corresponding to the lens of the eddy appear 376 to increase. To properly determine the size of the eddy it would be better to attempt 377 to estimate the dynamic height from the glider data and attempt to correspond these 378 observations with another independent reference like single track satellite altimetry data. 379 The temperature and salinity scales (Figure 6, 7) generally coincide with each other. 380

We summarize the correlation scale analysis (Figure 8) over the entire record with bars identifying sections sampling an eddy. The mean mixed layer depth for each section is the green horizontal line. The time bar is scaled to fit the sections into equal bins and is slightly distorting the timing of the events. Overall, the correlation scales within the mixed layer (above the green line, Figure 8) are smaller than 15 km unless eddies are detected. This confirms a result of von Oppeln–Bronikowski et al. (2021) where the glider sections were uniform in T and S, sampled in October and November, showing monthly



Figure 8. a) Length scales of temperature averaged for all glider sections over the entire record, broken down by isopycnal group observed in the glider record. Shading is ± 1 STD (b) Length scales for each section in each density bin (0.01 kg m^{-3}) . The green line is the mean mixed layer depth isopycnal in the period the sections were surveyed by the glider. Vertical lines indicate when an eddy was present in the surveyed sections. Note the y-axis (σ_0) is oriented from shallower (lighter) to deeper (denser) water.

averaged scales around 10–15 km. In general, our data point to a non-heterogenous state 388 of mixing and T and S variability in the Labrador Sea before, during and after convec-389 tion. 390

These results suggest that there is significant spatial variability at scales much shorter 391 than those sampled by the Argo floats. It is noteworthy that the scales observed are smaller 392 than what would be captured in CMIP style models (Hasumi, 2014) used in climate pre-393 diction scenarios (IPCC). Given the importance of convection and ventilation to longer 394 climate timescales (storage of carbon in the deep ocean Khatiwala et al. (2013)) our re-395 sults suggest that important processes happening on these scales are also not captured 396 in these simulations. 397

4 Discussion 398

399

4.1 Difference in OHC and FWC between Glider-Argo Observations

Based on the correlation scale analysis of the previous section, it becomes appar-400 ent that the processes that drive T and S variability are not heterogeneous. The extent 401 to which this variability is captured by Argo floats is not entirely clear given that the sampling does not allow us to repeat the same analysis as used on the glider data. In-403 stead, we focus again on the OHC and FWC calculations from Section 3.1, to investi-404 gate the role eddies play on the overall OHC and FWC measured from the glider and 405 Argo floats. 406

For a robust comparison between Argo and the glider, we bin averaged the profile 407 data for each month in Figure 9. The OHC and FWC averages \pm their standard errors 408 (STD/\sqrt{n}) with STD meaning the standard deviation and n the number of indepen-409 dent samples) are displayed in Figure 9. The glider covered a total distance ranging from 410 405-493 km per month. The correlation length scale analysis (Figure 8) showed that the 411 glider data (monthly averages) are correlated within 19-26 km of the zero-crossings and 412 therefore are not independent within/below this distance. This suggests the presence of 413 15-23 independent samples, n, in each section. We estimate the number of independent 414



Figure 9. Comparison of monthly averaged OHC (b) and FWC (c) calculated with Argo (dark purple) and glider (orange) profiles in the subsurface 500-1000 m layer, with error bars indicating +/- 1 standard errors. Bars in (b) show the number of Argo profiles used for the monthly average.

Argo profiles with 6 to 14 profiles sampled each month. Given that the distance between 415 float profiles averaged around 55 km, compared to just a few km for the gliders (Figure 416 1 1, and the low mean correlation scales (from the glider), we assume that each profile 417 is independent (n = 6-14). On average between two profiles over 10 days a float would 418 travel 55 km. We focus on the deeper layer 500-1000 m to look at subsurface (Jan-Feb, 419 Apr-May) effects below the mixed layer. From the correlation scale analysis, we know 420 that the subsurface correlation scales are larger than at the surface, providing a more 421 robust comparison between the two platforms. 422

Comparing the Argo and glider data (Figure 9), we find that the deep-layer (500-423 1000 m) cooled and freshened during the period of convection. While Argo and glider 424 monthly OHC are in agreement, FWC are not, with the exception of January. What is 425 causing this difference of up to a maximum of 0.3 m^{-1} in FWC in February and min-426 imum 0.035 m^{-1} in January in the subsurface intermediate depth layer between Argo 427 and glider observations? February is the month of the lowest number of Argo float ob-428 servations (6). There appears to be a decrease in difference between FWC estimation 429 with an increase in monthly Argo observations (maximum of 14 in April, FWC differ-430 ence 0.0939 m^{-1}). Accurate knowledge on surface and subsurface salt fluxes in this re-431 gion is critical to properly represent mixed layer dynamics. From the previous section 432 on correlation scales we also know that it was during these months that the glider en-433 countered the majority of eddy events though eddies were also present in January. It could 434 be that these observed eddies have an influence on the overall FWC and that the Argo 435 data may not capture all of this variability. 436

Given the disagreements in the monthly FWC data between Argo and glider, what is the importance of the small scale variability in the data vs data averaged over a time period like 10 days? If we subsample the glider data to Argo spacing of 10 days and compare this with the glider data averaged in 10 day bins we can investigate how much variability is missed by averaging vs an instantaneous representation given by the denser sampling of the glider (Figure 10). This gives a sense of how important dense sampling could



Figure 10. (a) Comparison of 10 day averaged glider OHC and (b) FWC vs 10 day subsampled data (approximately Argo sample spacing).

be. Looking at the result we see that the glider data subsampled and averaged are overall in good agreement, but that in months with increased eddy activity (February and
May) there are differences of up to 50% between FWC and OHC subsampled vs averaged data.

447

4.2 Importance of Eddies on Heat and Salt Budget

From the previous section, it appears that Argo floats and glider data estimate sim-448 ilar monthly heat content means in the lower-layer. In contrast, there are differences 449 in FWC between the two data sets, possibly due to the presence of eddies that are only 450 detected with the spatio-temporal resolutions of glider measurements whereas Argo floats 451 can sample within an eddy but their resolutions are too sparse to fully resolve the scale 452 and energy of these events. The glider data show a non-heterogenous state with regards 453 to the spatial structure of T and S important for our understanding of air-sea gas ex-454 change, subduction and mixing. 455

We describe our idea of the observations and dynamics in Figure 11. The diagram 456 shows Argo profiles and the glider trajectory inside a sampling box 200 by 100 km wide 457 corresponding with the spatial coverage of the observation (see Figure 1b). There are 458 eddies representing small-scale features throughout the sampling box as well as convec-459 tion driven by both surface and eddy heat and freshwater fluxes. The glider encounters the eddies on several passes, while the Argo floats measure the large-scale process well, 461 but do not resolve the smaller-scale eddy dynamics associated with the eddy. In Section 462 4.1 we estimate that the difference in glider and Argo OHC is small (within a standard 463 error) but that differences in FWC can reach up to 0.3 m^{-1} . For some of the glider tran-464 sects it is clear that it was the same eddy that was measured as the glider passed through 465 the same T-S signature when it turned around. A total of 6 eddies were sampled by the 466 glider. Our result of increased eddy activity in May aligns with other studies (Körtzinger 467 et al., 2008; Fischer et al., 2018) which point to increased eddy kinetic energy in spring 468 time following the tail of convection. Other studies (Chanut et al., 2008; Rieck et al., 2019) 469 have also shown eddy events in the period during fall and winter. In addition, the glider 470



Figure 11. Schematic diagram of example glider sampling in the Labrador Sea together with Argo float profiles and eddies.

likely sampled the signal of convective mixing cells but their scales of meters to a few
km are likely too small to be reliably sampled by the glider.

An open question remains around the relative contribution of freshwater due to ed-473 dies prior and after convection in terms of their impact on the onset of convection and 474 restratification. A few studies have pointed at eddies and their role in modulating OHC 475 and FWC in the Labrador Sea (Chanut et al., 2008; Hátún et al., 2007). We use our com-476 parison between glider and Argo FWC and our eddy passes to address the importance 477 of eddies over the whole basin using a simple scaling approach with the dimensions pro-478 vided by the glider sampling. The approximate area of the glider sampling with Argo 479 profiles is a box of 200 by 100 km, with 48 Argo profiles over the time from January to 480 May (Figure 3). The glider did 16 passes along the 200 km corridor with 8 eddy events 481 in all the transects completing on average 420 km per month. If we assume that eddies 482 (approximately 40 km wide) are equally distributed and that the glider was managing 483 the same section, the glider should have sampled an eddy approximately 50% of the time 484 or every second pass through the 200 km corridor. We recognize that this is a crude as-485 sumption given that the eddy detection is not distributed equally every month. For the 486 months with eddies, 20% (80 km / 420 km = 20%) of the total distance sampled by the glider spanned through at least one eddy. If we ignore issues around the glider sampling, 488 we can approximate that 20% of the box area could have been occupied by eddies at any 489 given time. In the previous section, we discussed the possibility that the start in FWC 490 subsurface discrepancy between glider and Argo could partly be due to eddies. We there-491 fore apply this area scaling directly to the difference in observed monthly FWC to work 492 out an approximate contribution of eddies to FWC in the Labrador Sea using 300 km 493 radius of convection region (Lilly et al., 2003). Using February as an example of strong 494 eddy activity, we found a difference of 0.3 m^{-1} FWC between Argo and glider. We con-495 clude that eddies could possibly contribute up to 4.2 m^{-1} FWC for the entire basin in 496 that month. It would be interesting to see how our results compare against modelling 497 studies. This could help to pinpoint discrepancies between models and observations, as 498 well as help to understand the sources and sinks of freshwater in the Labrador Sea, which 499 is an ongoing question (Aagaard & Carmack, 1989; Zhang et al., 2021). We note that 500 the number of Argo profiles in winter is lower compared with the summer (see Supple-501

ment, Figure S1), and that the glider moves slowly and likely not sampling the area well
 enough to confidently isolate the eddy contribution. Therefore, this exercise is more of
 a thought experiment but could help inspire the sampling strategies of future glider mis sions in the region.

506 5 Conclusions

In this study we presented winter time observations of heat and freshwater content 507 and mixing in the Labrador Sea from Argo floats and a glider. In the Introduction, we 508 posed three main questions for this paper: (1) To what extent do small scale processes 509 contribute to the cycle of heat and salt exchange in this region and therefore the forma-510 tion of LSW? (2) what are the dominant scales of water property evolution in the Labrador 511 Sea? and (3) do gliders offer critical information on scales below that offered by Argo 512 floats? Here we summarize the results and discussion as well as offer a future outlook 513 on glider observation needs in the Labrador Sea. 514

We examined OHC and FWC for Argo and glider observations and found that un-515 surprisingly the glider exhibits a lot of small-scale energetic structure that is not seen 516 by the coarser sampling Argo floats. These distinct events occur due to eddies and hence 517 the question on the importance of the small scale structure is in part a question on the 518 role of eddies for FWC and OHC in this region. The glider encountered a number of salin-519 ifying events which are associated with the glider passing through eddies with different 520 water masses and FWC than the surrounding water. These salinity altering events ex-521 tend to the deep (500-1000 m) layer, pointing to their energetic nature. Filtering the glider 522 data by a 3-day low-pass filter removes the small–scale variability and most of the eddy 523 signal. To answer the question to what extent these small scale processes contribute to 524 the overall OHC and FWC in this region is difficult by the fact that we do not have enough 525 glider data for a full year and not enough independent measurements of a field with ed-526 dies and without. From our discussion in Section 4.2 we think that eddies may play a 527 significant role in redistributing OHC and FWC in the Labrador Sea. Similar conclu-528 sions are offered by eddy-resolving modelling studies that consider these contributions 529 (Yeager et al., 2021). We further underline this point, comparing the three heat flux datasets 530 (ERA5, Argo and glider derived implied heat fluxes) we find that the correlation between 531 glider and Argo implied surface heat fluxes and ERA5 are dependent on window aver-532 aging. For the glider it appears that the spatial scales may impact the correlation scales. 533 The more widely distributed Argo floats on the other hand show good agreement with 534 ERA5. 535

Convection in the Labrador Sea in winter 2019–2020 was weaker than the previ-536 ous year. The maximum MLD from Argo profiles reached nearly 2 km based on indi-537 vidual profiles. Though overall the volume of LSW formed during this year convection 538 season was less than in the previous year's winter season (2018–2019). The deepening 539 of MLD was rather gradual (Jan-March) but the end of convection was very abrupt (March 540 31–April 1). For the FWC and OHC calculations we broke the data into two separate 541 layers: a surface (0-500 m) and a deep (500-1000 m) layer. There is a distinct separa-542 tion in the two layers in January–February between FWC and OHC prior to convection 543 (March). The merging of OHC and FWC for the top and deep layer is rather sudden, 544 even though MLD changes are gradual. The end of convection and change in MLD is 545 very dramatic occurring in a matter of days as evidenced by both data-sets. The ERA5 546 surface heat flux show that the end of convection is preceded by a week of weak warm-547 ing (100 W m^2). These observations suggest convection is largely driven by surface heat 548 fluxes. From the glider high-resolution transects we estimated the correlation length scales 549 and find that the T and S structures exhibit a lot of heterogeneity due to the aforemen-550 tioned eddy events. On average, spatial scales are small on the order of 15-20 km in agree-551 ment with previous glider experiments in the region (von Oppeln–Bronikowski et al., 2021). 552 However, eddies have a significant influence on the spatial structure increasing length 553

scales to 40-50 km. We note several limitations to our approach. The glider only managed to sample one transect about twice a month. The slow speed likely aliased the space
and time variability in the record. We limit the influence of aliasing in our assumption
of spatial snapshots of the sections by detrending the data and working on isopycnal coordinates.

Monthly averages in FWC show significant differences between platforms during 559 February, when eddy sampling by the glider was highest. We estimate, based on the dis-560 crepancy in FWC between Argo and glider data, that eddies could potentially play a sig-561 nificant role. Ignoring limitations of the glider data, we estimate that up to 4 m^{-1} of FWC change could be attributed to eddies for the entire Labrador Sea. The lack of compa-563 rable datasets limits how far we can push this result or allow us to further test our re-564 sults. The Argo floats probably then sample the field in sufficient spatial density and cov-565 erage that they adequately sample the mean state of the Labrador Sea and so for bulk 566 estimates do a good job of how has the Labrador Sea changed but do not give further 567 insight into why. Gliders reveal a more detailed look at eddies which is important for 568 understanding their role towards the bigger system of the Labrador Sea. Future glider 569 missions could sample the region more densely to improve our understanding of the last-570 ing influence of these eddies on the region. This is not possible to do with a single glider 571 since the gliders are too slow to repeatedly and effectively sample such a large region. 572

Our results do not have a definitive answer as to the role of eddies, but rather of-573 fer another glimpse (see e.g. Hátún et al., 2007) that eddies might play an important role 574 and that the role of the interior of the Labrador Sea towards pathways of watermasses is not yet sufficiently well understood. A recent paper by (MacGilchrist et al., 2021) point 576 577 to the boundary as important for driving subduction and hence export and not the interior, however those results do not consider eddies. An observing strategy targeting both 578 the export at boundaries and a way to track eddies is needed. We are planning to do a 579 follow up long duration glider sampling experiment in the Labrador Sea with several glider 580 sampling near the boundary current, flying along the boundary current to observe sub-581 duction as well as gliders offshore from the boundary current tacking snapshots of ed-582 dies that propagate into the interior of the Labrador Sea. The goal of this study would 583 be to quantify the ideas in MacGilchrist et al. (2021) as well further analyze the contri-584 bution of eddies to FWC and OHC in the Labrador Sea. To properly disentangle the role 585 of eddies, eddy tracking is needed which would relay on novel tools for glider navigation 586 and remote sensing from satellites. 587

588 Data Availability

Processed glider data have been archived on SEANOE (https://www.seanoe.org/data/ 589 00681/79349/) as part of the Memorial University glider data repository (von Oppeln-590 Bronikowski et al., 2021). Argo data were collected and made freely available by the In-591 ternational Argo Program (Argo, GDAC, 2000) and the national programs that contribute 592 to it (https://argo.ucsd.edu, https://www.ocean-ops.org). The Argo Program is 593 part of the Global Ocean Observing System. ERA5 (Hersbach et al., 2018) 1-hr single 594 level data were accessed from Copernicus Climate Data Store (CDS) (https://cds.climate 595 .copernicus.eu/cdsapp#!/dataset/reanalysis-era5-single-levels?tab=form). 596

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Figure 1.





Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.



Figure 9.





Figure 10.



Figure 11.

