

# Observations of Small Large-Amplitude Magnetic Structures (SLAMS) at Mars by MAVEN

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## Abstract

According to the different orientations of the interplanetary magnetic field (IMF), the planetary shock can be either quasi-parallel or quasi-perpendicular. Under quasi-parallel conditions a significant number of solar wind suprathermal particles are reflected from the shock and drift along IMF, forming an extended and highly turbulent region called the foreshock where various nonlinear plasma phenomena are observed. In this research, we perform a case study of the structures in the foreshock region at Mars observed by Mars Atmosphere and Volatile Evolution (MAVEN). We use data from plasma analyzer STATIC and magnetometer MAG to analyze ion beams angular spectrum and magnetic field dynamics. We show that the observed structures are consistent with Short Large-Amplitude Magnetic Structures (SLAMS), commonly detected in foreshock regions of magnetized and unmagnetized bodies throughout the Solar system. Finally, we calculate the magnetic Mach number to analyze the characteristics of the observed foreshock structures. The analysis shows, that SLAMS are formed by the resonance between plasma waves propagating along the IMF and the backstreaming scattered solar wind H<sup>+</sup> and exospheric O<sup>+</sup> and O<sub>2</sub><sup>+</sup> ions, with the dominant impact of O<sub>2</sub><sup>+</sup> ions.

1           **Observations of Small Large-Amplitude Magnetic**  
2           **Structures (SLAMS) at Mars by MAVEN**

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5           **Key Points:**

- 6           • The observed plasma structures are associated with Small Large-Amplitude Mag-  
7           netic Structures  
8           • The acceleration of  $O^+$  and  $O_2^+$  ions is the result of wave-particle interaction via  
9           Landau damping of ULF waves with  $H^+$  ions

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## Abstract

According to the different orientations of the interplanetary magnetic field (IMF), the planetary shock can be either quasi-parallel or quasi-perpendicular. Under quasi-parallel conditions a significant number of solar wind suprathermal particles are reflected from the shock and drift along IMF, forming an extended and highly turbulent region called the foreshock where various nonlinear plasma phenomena are observed. In this research, we perform a case study of the structures in the foreshock region at Mars observed by Mars Atmosphere and Volatile Evolution (MAVEN). We use data from plasma analyzer STATIC and magnetometer MAG to analyze ion beams angular spectrum and magnetic field dynamics. We show that the observed structures are consistent with Short Large-Amplitude Magnetic Structures (SLAMS), commonly detected in foreshock regions of magnetized and unmagnetized bodies throughout the Solar system. Finally, we calculate the magnetic Mach number to analyze the characteristics of the observed foreshock structures. The analysis shows, that SLAMS are formed by the resonance between plasma waves propagating along the IMF and the backstreaming scattered solar wind  $H^+$  and exospheric  $O^+$  and  $O_2^+$  ions, with the dominant impact of  $O_2^+$  ions.

## 1 Introduction

The solar wind interaction with the Martian plasma environment has been actively investigated for the past few decades. One of the mostly discussed research areas is the solar wind interactions with planetary plasma environment. As the supersonic solar wind flow becomes subsonic at closer distances to Mars, a bow shock is formed at which solar wind is decelerated, deflected and thermalized. The observation of the Martian bow shock suggests the existence of the region upstream of the bow shock filled with ULF waves, diffusive ions and electrons. This region is known as foreshock region, which is forming under quasi-parallel shock conditions. In foreshock numerous plasma phenomena occur. The largest structures observed in the foreshock are foreshock cavities (Sibeck et al., 2002), foreshock bubbles (Turner et al., 2013) and hot flow anomalies (Schwartz et al., 1985; Thomsen et al., 1986; Paschmann et al., 1988).

The wave-particle interaction of ULF waves and ions under quasi-parallel shock conditions may lead to the formation of Short Large-Amplitude Magnetic Structures (SLAMS). The observations of SLAMS at terrestrial foreshock are described as long pulsations on a short time interval (Schwartz & Burgess, 1991; Schwartz et al., 1992; Wilson III et al., 2013). To date, SLAMS have been already observed at Venus (Omidi et al., 2017), Saturn (Bebesí et al., 2019) and Jupiter (Tsurutani et al., 1993), comets (Tsurutani et al., 2013). On Mars there are evident observations of SLAMS presence (Halekas et al., 2017; Collinson et al., 2018). However, the impact on the modulation of energetic neutral atoms flux by foreshock structures like SLAMS on Mars was described by (Fowler et al., 2019), but no analysis of SLAMS themselves was conducted. Also, there is an insight on the physical model of SLAMS formation at terrestrial foreshock as observed by Magnetospheric Multiscale (MMS) Mission (Chen et al., 2021).

This article provides the results of the SLAMS analysis during near-radial interplanetary magnetic field (IMF) conditions observed by Mars Atmosphere and Volatile Evolution (MAVEN) mission. We analyze plasma properties of foreshock ions upstream SLAMS formation region and solar wind ions both upstream and downstream SLAMS. We use minimum variance analysis technique to investigate wave nature of the observed process.

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## 2 Instrumentation

This research is based on the data obtained from MAVEN spacecraft, launched on 13 November 2013. MAVEN is inserted on an elliptical orbit around Mars with an orbital period of 4.5 hours. Its periapsis is at 150 km and apoapsis at 6200 km with 75° inclination. Data from magnetometer MAG and Suprathermal And Thermal Ion Composition (STATIC) instrument from the Particle and Fields package is used in this study.

STATIC is a top-hat ion energy-mass analyzer (McFadden et al., 2015). The instrument measures the ion energy distribution in a wide energy range from 0.1 eV to 30 keV and can resolve  $H^+$ ,  $O^+$  and  $O_2^+$  ions, covering ionospheric, magnetospheric and tail plasma. Its field of view (FOV) is  $360^\circ \times 90^\circ$  which is decreasing at high energies, with angular resolution  $22.5^\circ \times 22.5^\circ$ . The energy resolution of the instrument is  $dE/E \sim 15\%$  and mass resolution is  $M/dM \geq 4$ . During the data analysis routine, the contamination of  $H^+$  mass channel is considered. The data sampling rate can be switched from 4 to 16 sec, and was 4 sec during the analyzed time interval.

The MAG measures 3 components of local magnetic field in the solar wind, magnetosheath and crustal magnetic field with 32 Hz time cadance (Connerney et al., 2015). Its dynamic range is 60000 nT with a resolution is 0.05 nT.

## 3 Observation

The event was detected in a time interval within 02.00 – 02.30 UT 23 October 2019. The observed orbit of MAVEN lies through the dayside of Mars, crossing the subsolar region from the northern to the southern hemisphere of Mars. The altitude varies from 500 to 2000 km. The observational period is divided into two regions: oscillation region from 02.00 to 02.25 UT and solar wind region from 02.25 to 02.30 UT. The solar wind conditions are characterized by narrow energy spectra of  $H^+$  ions with the maximum energy flux approximately at 1 keV and a weakly disturbed magnetic field. The (IMF) vector averaged over the solar wind region has components [4.1020 -1.1556 0.6414] nT in Mars Solar Orbital (MSO) frame, in which x-axis is pointed to the Sun, y-axis is directed against the orbital motion of the planet, and z-axis completes the system to the right-handed basis. Considering orbital characteristics and IMF conditions, the observed oscillation region is consistent with the foreshock region.

The data from STATIC and MAG is demonstrated in form of time series in Figure 1. Starting from 02.00 UT quasi-periodic pulsations of the magnetic field, accompanied by deceleration of  $H^+$  ions, are observed. The pulsations of the magnetic field have a period of  $66 \pm 36.6$  sec and a time width of approximately 12 sec. The magnetic field in the structures increases by factor of 4 to 8 compared to the total value of IMF.

No significant correlation between ion density variations and magnetic field pulsations are observed, despite the time interval from 02.12 to 02.15 UT, where peaks of magnetic field pulsations coincide with minimums of light-to-heavy ion density ratio.

The Figure 2 demonstrates an example of a typical angular distribution function during SLAMS crossing from 02.22 to 02.24 UT. Hammer projection is used to show the measured part of velocity space by STATIC instrument. Each bin corresponds to one angular cell with sizes  $25^\circ \times 25^\circ$  and the color of the bin shows the differential energy flux. We also consider a feature of STATIC FOV by which it is narrowing as the measured energy increases. The direction of the local magnetic field is shown as a red cross. The MSO basis vectors of frame  $\{X_{MSO}, Y_{MSO}, Z_{MSO}\}$  in STATIC frame are shown by red, green and black dots and circles respectively, highlighting the positive and negative directions. Empty angular distribution functions are neglected.

106 The detailed analysis of the  $H^+$  ion angular distribution function (ADF) in Fig-  
 107 ure 2 shows the observation of two ion beams: the solar wind  $H^+$  ions beam and fore-  
 108 shock  $H^+$  ions. The solar wind is seen as a narrow ion beam with high differential en-  
 109 ergy flux in the sunward ( $X_{MSO+}$ ) direction. At the same time foreshock ions are char-  
 110 acterized by wide angular distribution detected near the direction of IMF. The typical  
 111 ADF of  $H^+$  ions is seen from 02.25.11 to 02.25.19 UT which corresponds to the upstream  
 112 region of SLAMS in Figure 1. At 02.25.07 UT the process of deceleration and heating  
 113 of the solar wind is seen. The energy of ions on average decreases by 4-5 times and the  
 114 angular coverage of FOV by the solar wind ions is significantly higher, compared to the  
 115 upstream region. From 02.24.45 to 02.25.07 UT the ions beam population of low energy  
 116 scattered  $H^+$  ions is observed propagating in the direction of IMF. At the same time  $O^+$   
 117 and  $O_2^+$  ADF shows the appearance of narrow ion beams in the sunward and pick-up  
 118 ions in anti-sunward directions.

119 The presence of the scattered/deflected ion beam is the consequence of the solar  
 120 wind interaction with the bow shock-like sharp front of SLAMS. As result, the interac-  
 121 tion of the deflected ion beams with Alfvén waves is described by Landau damping. Com-  
 122 paring phase velocity of Alfvén wave  $v_A$  and peak velocity  $v$  of the solar wind  $H^+$  ions  
 123 EDA, we see that  $v > v_A$ . Thus energetic ions transfer their energy to the Alfvén wave  
 124 modulating the amplitude of SLAMS.

## 125 4 Analysis

126 In further chapters, the results of the case study will be performed and discussed.  
 127 In a time interval from 02.22 to 02.26 UT, only several events have clear observations  
 128 of the magnetic field oscillations and plasma properties.

### 129 4.1 Minimum Variance Analysis of magnetic field oscillations (MVAB)

130 We apply MVAB to calculate the wave vector  $k$  of magnetic field  $B$  oscillations. We  
 131 also assume that the observed oscillations are more temporal rather than spatial due to  
 132 the specific process of SLAMS formation. To estimate the orientation of the wave vec-  
 133 tor  $k$  the condition of  $(k, B) = 0$  should be considered. According to MVAB, solving the  
 134 eigenvalue and eigenvector problem for matrix  $M = \langle B_i B_j \rangle - \langle B_i \rangle \langle B_j \rangle$ , where  $i, j =$   
 135  $\{x, y, z\}$ , results as three vectors  $B_{max}, B_{int}, B_{min}$  of maximum, intermediate and min-  
 136 imum variance of magnetic field respectively. In  $B_{max}, B_{int}, B_{min}$  frame wave activity  
 137 in a given time interval looks like rotation in one of the planes which is clearly seen in  
 138 Figure 3. If the rotation is in the  $B_{max}$  and  $B_{int}$  plane, then it gives a rough estimation  
 139 of the line, which contains  $k$  and  $B_{min} \parallel k$ . The direction of the wave vector is chosen  
 140 based on the physical conditions of the processes that occurred in foreshock.

141 The major results of MVAB are listed in Table 1. It is seen that the observed waves  
 142 propagate along with the IMF into the Sun direction. The polarization of waves is both  
 143 left-handed and right-handed. Though data can't provide accurate information about  
 144  $H^+$  temperature due to different problems, we assume plasma is cold. In cold magne-  
 145 toactive plasma, several types of wave modes exist. As the observed waves have  $k \parallel B$ ,  
 146 they are possibly Alfvén or magnetosonic waves. The relation of  $B_{\parallel}/B_{\perp}$  shows that the  
 147 waves have major oscillations perpendicular to magnetic field direction, thus we assume  
 148 that observed waves have Alfvén nature.

### 149 4.2 Wavelet analysis

150 To go deeper with the wave analysis, we apply continuous wavelet transform (CWT)  
 151 on the magnetic field with the Morlet wavelet. In Figure 4a the CWT is demonstrated  
 152 for the whole period of time. The colorbar is indicating the square module of the am-  
 153 plitude. In the time interval from 02.12 to 02.16 UT, strong oscillations near  $O^+$  and  $O_2^+$

154 ions cyclotron frequency are observed. A normalized general wavelet spectrum (GWS)  
 155 with overlaying mean value and standard deviation of cyclotron frequencies of  $H^+$ ,  $O^+$   
 156 and  $O_2^+$  ions is shown in Figure 4b. It is seen that  $O_2^+$  cyclotron frequency corresponds  
 157 to the global maximum of GWS. At the same time,  $H^+$  and  $O^+$  cyclotron frequencies  
 158 are located in the vicinity of GWS local maximums. This pattern proves the hypothe-  
 159 sis of ultra-low frequency Alfvén waves observation, which originates from the solar wind  
 160 interaction with Martian quasi-parallel bow shock. We also observe an intensive inter-  
 161 action between Alfvén waves and  $O_2^+$  ions in foreshock region.

### 162 4.3 SLAMS shock characteristics

163 One of the most interesting features of SLAMS is their shock characteristics. To  
 164 investigate the shock parameters of SLAMS, shock normal and the angle between the  
 165 IMF and shock normal for SLAMS and bow shock are calculated. The formula for cal-  
 166 culation of the normal vector to shock surface of rotational discontinuity requires the val-  
 167 ues of the magnetic field upstream and downstream of the shock. The upstream mag-  
 168 netic field was averaged in the vicinity of SLAMS, and the downstream magnetic field  
 169 was averaged in the core of SLAMS. The results of the calculation are listed in Table 2.  
 170 The average duration of the magnetic field amplification is  $12.6 \pm 3.72$  sec, the angle be-  
 171 tween the bow shock and IMF on average is  $\Theta_{Bn} = 14.9 \pm 0.9$ , which is consistent with  
 172 quasi-parallel bow shock. The SLAMS shock parameters are varying drastically compared  
 173 to bow shock, however, the average angle between the normal vector of SLAMS and IMF  
 174 is less than  $45^\circ$ . This means SLAMS may inherit the configuration of the bow shock.

175 The dynamics of solar wind during the interaction with planetary bow shock can  
 176 be described by magnetic Mach number  $M_A = v_{sw}/v_A$ , with  $v_{sw}$  – solar wind veloc-  
 177 ity,  $v_A = B/\sqrt{4\pi np}$  – Alfvén velocity of  $H^+$  ions. The Alfvén velocity is important phys-  
 178 ical parameter in space plasma, which closely related to the wave activity. If  $M_A > 1$ ,  
 179 then the solar wind velocity has supersonic values; if  $M_A < 1$  – the solar wind veloc-  
 180 ity is subsonic. High Mach numbers ( $M_A > 1$ ) are typical for the solar wind in the up-  
 181 stream region. As the solar wind interacts with the bow shock, the solar wind deceler-  
 182 ates, the total magnetic field is increasing by 2-3 times according to the Rankine-Hugoniot  
 183 conditions. All this factors cause the decrease of magnetic Mach number to subsonic val-  
 184 ues. Previously, we showed that SLAMS can have shock nature.

185 In Figure 5 the scatterplot of Mach number and density of  $H^+$  ions is shown. The  
 186 colorbar indicates the ratio of the measured magnetic field to the total value of IMF. Two  
 187 populations of  $H^+$  ions are seen on the scatterplot with different density and Mach num-  
 188 ber. The labeled with red color distribution of density and Mach number indicates the  
 189 solar wind  $H^+$  ions, and labeled with blue indicates the shocked  $H^+$  ions. Pretty log-  
 190 ical to assume less dense and faster ion population as an upstream  $H^+$  ions of the so-  
 191 lar wind.

### 192 4.4 Wave-particle interaction of ULF waves and foreshock ions

193 Considering MAVEN altitude during the analyzed time interval, the observed back-  
 194 streaming  $O^+$  and  $O_2^+$  ions are originated from the Martian exosphere. The energies of  
 195 these ions far exceed the thermal energies of exospheric ions. One of the possible mech-  
 196 anisms for the growth of the ion energies might be Landau damping. As shown in Sec-  
 197 tion 4.1, the observed ULF waves are propagating along the IMF and interact with ex-  
 198 ospheric ions. The Figure 6 represents the comparison of calculated average velocities  
 199 of backstreaming  $H^+$ ,  $O^+$  and  $O_2^+$  ions with Alfvén velocity, which is considered as phase  
 200 velocity of the observed waves.

201 The peak velocities of  $O^+$  and  $O_2^+$  ions correspond to 17.5 and 24.3 km/s, respec-  
 202 tively, which is lower than mean Alfvén velocity (around 34.9 km/s). At the same time,

203  $H^+$  maximum velocity is around 126.9 km/s, which is higher compared to Alfvén veloc-  
 204 ity. In terms of wave-particle interaction this observation can be interpreted as a com-  
 205 plex multicomponent plasma interaction. The backstreaming  $H^+$  ions transfer their ki-  
 206 netic energy to ULF waves, amplifying them via Landau damping mechanism. Then, am-  
 207 plified ULF waves transfer energy to  $O^+$  and  $O_2^+$ , accelerating them to suprathermal ve-  
 208 locities in the sunward direction.

## 209 5 Conclusion and discussion

210 In conclusion, this article performs the case study of foreshock structures, commonly  
 211 known as Short Large-Amplitude Magnetic Structures, in the time interval 02.20 to 02.30  
 212 UT of 23 October 2019. On the time scale of roughly tens of seconds, the magnetic field  
 213 amplifies by a factor of 4-5 times compared to IMF. The MVAB applied on the time in-  
 214 terval of SLAMS observation shows the presence of ULF Alfvén and magnetosonic waves,  
 215 which are originated from the interaction of the solar wind ions and backstreaming ions.  
 216 It was found by the wavelet analysis that  $O_2^+$  ions cyclotron frequency is the dominant  
 217 observed frequency in the oscillations of the magnetic field. The possible explanation is  
 218 the high inertia of  $O_2^+$  ions, compared to  $O^+$  and  $H^+$  ions. The shock parameters of SLAMS  
 219 are inherited from the planetary bowshock, having the same quasi-parallel structure.

220 The observed process of  $O^+$  and  $O_2^+$  ions acceleration from thermal to suprather-  
 221 mal velocities is not analyzed in this article. Though, the proposed mechanism of Lan-  
 222 dau damping during the wave-particle interaction seems logical, more analysis should  
 223 be done.

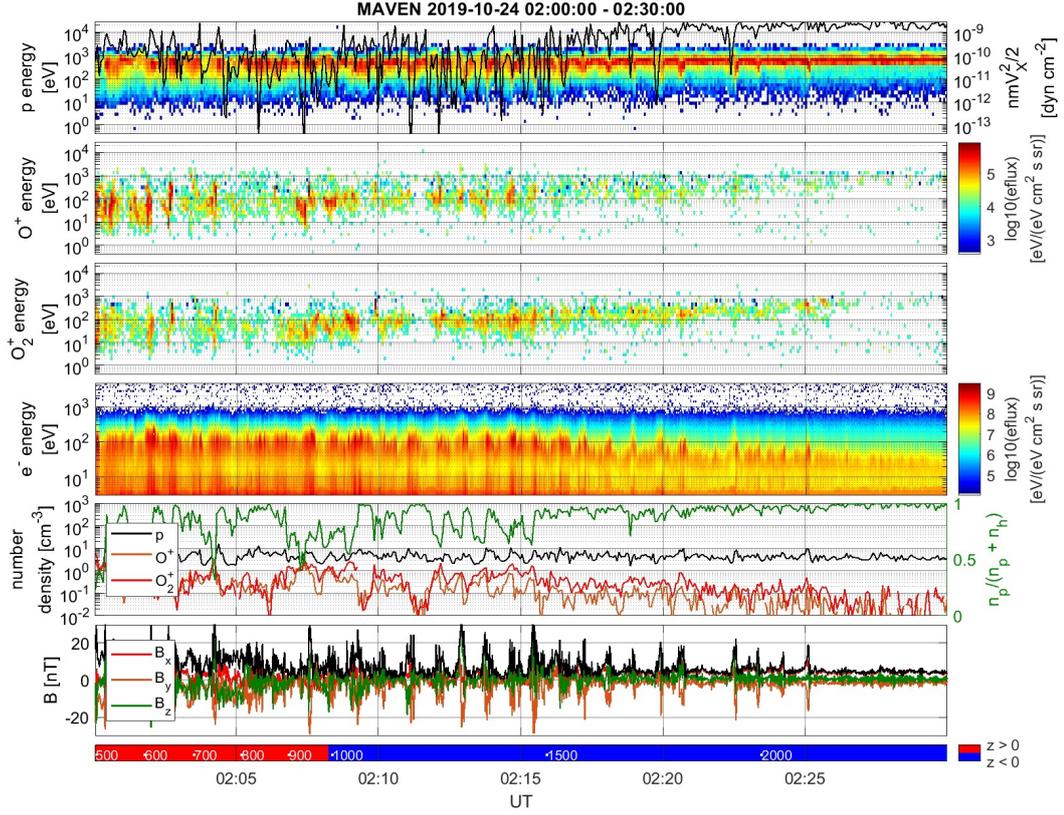
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 228 sions.

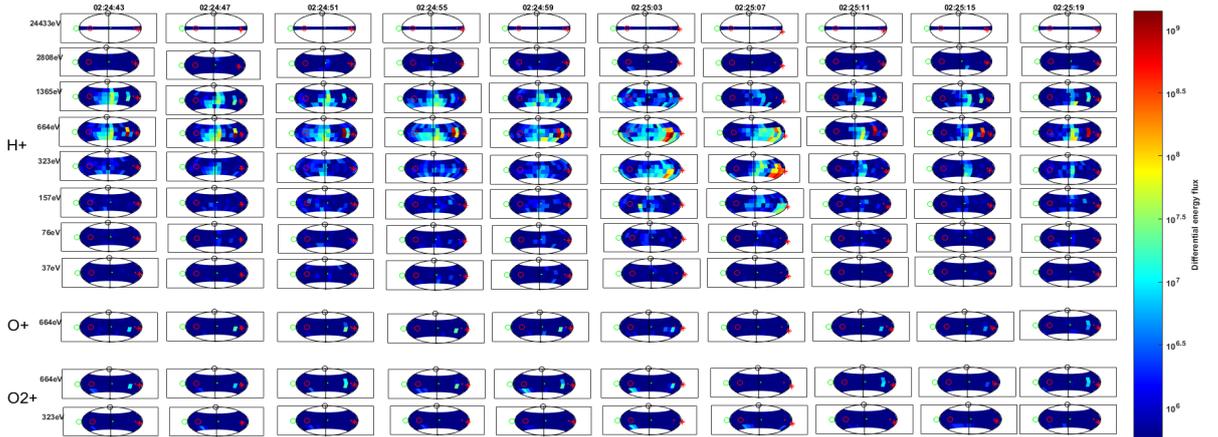
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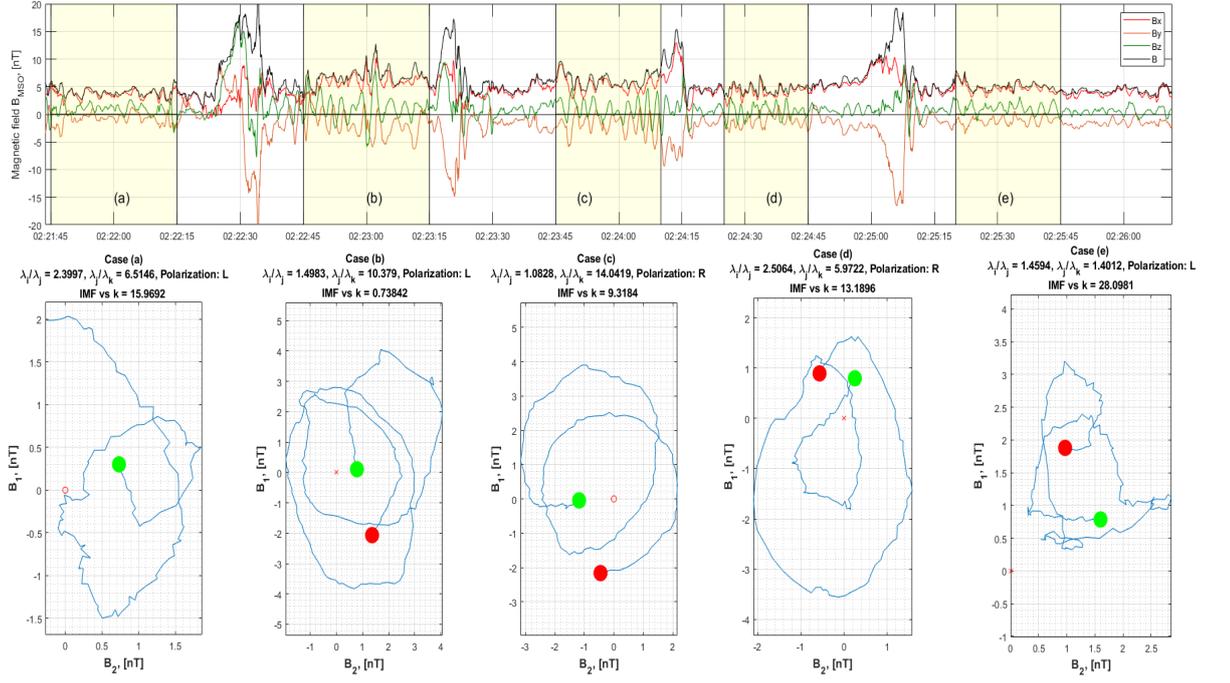
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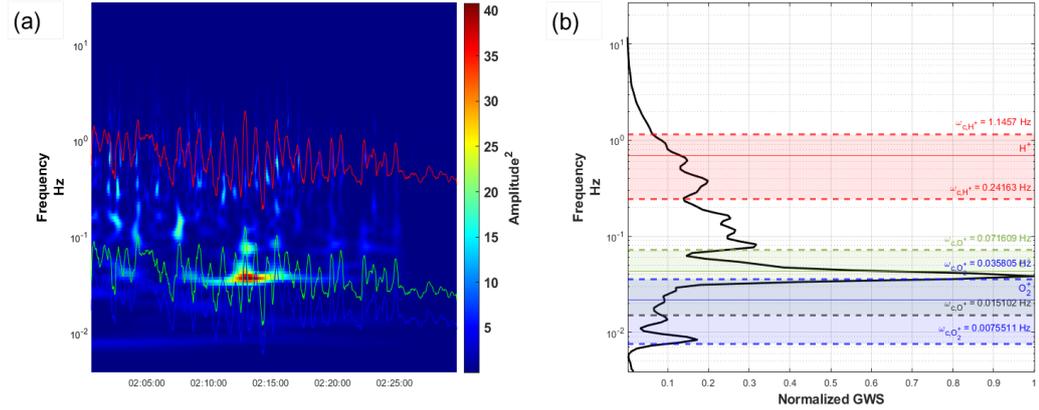
**Figure 1.** Observation of SLAMS at 24 October 2019 in a time interval from 02:00:00 to 02:30:00 (from up to the bottom): energy-time spectrograms of  $H^+$ ,  $O^+$  and  $O_2^+$  ions and electrons; number density of  $H^+$ ,  $O^+$  and  $O_2^+$  ions, and light to heavy ions ratio overlaid; vector of magnetic field



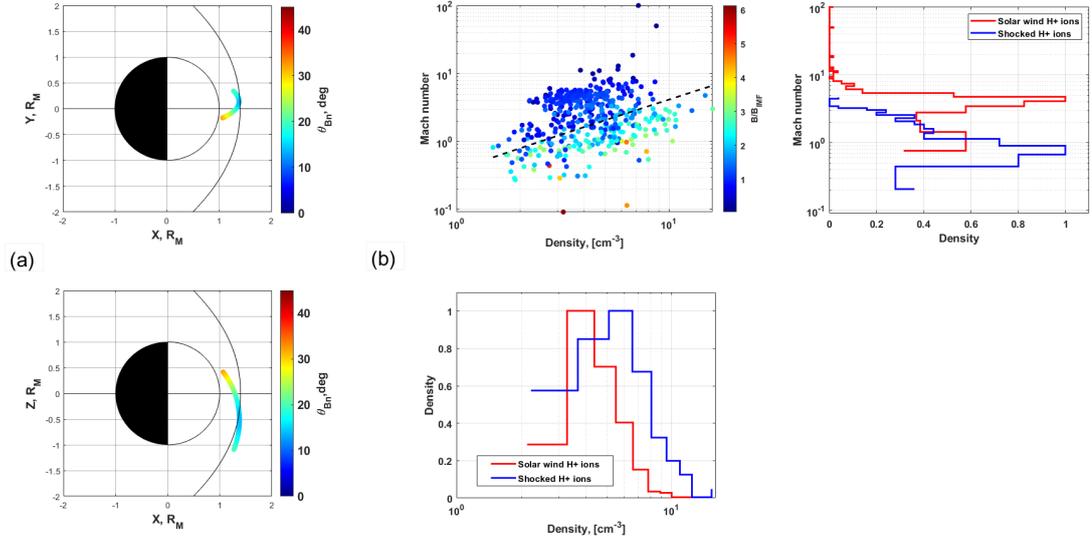
**Figure 2.** An example of ion angular distribution function of  $H^+$ ,  $O^+$  and  $O_2^+$  ions. Red cross corresponds to the IMF direction in STATIC frame. Red, green and blue dot and circles corresponds to XMSO, YMSO and ZMSO in STATIC frame, where dots and circles indicate positive and negative direction of each axis, respectively



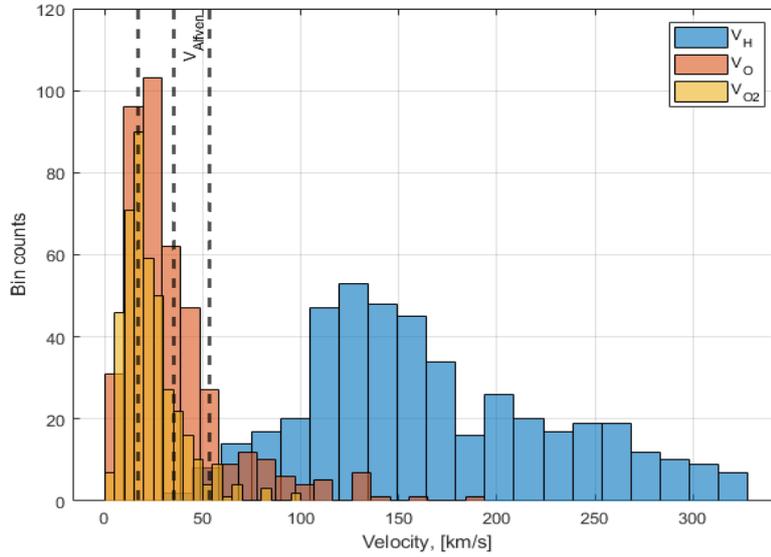
**Figure 3.** The results of MVAB for the time interval from 02.22 to 02.26 UT. Green and red dots correspond to the beginning and end of the time interval. With dot and cross the direction of vector  $k$  is demonstrated



**Figure 4.** (a) Wavelet spectrum of the magnetic field and (b) General Wavelet Spectrum (GWS). Colored areas indicate confidence interval for cyclotron frequencies of  $H^+$ ,  $O^+$  and  $O_2^+$  ions



**Figure 5.** (a) MAVEN orbit projections (colorbar shows the angle between the magnetic field and normal to the bow shock  $\Theta_{Bn}$ ), (b) scatterplot of  $H^+$  ions density and magnetic Mach number



**Figure 6.** Comparison of the backstreaming  $H^+$ ,  $O^+$  and  $O_2^+$  ions velocities with the Alfvén velocity

**Table 1.** The calculated characteristics of the observed waves

Event	Time, UT	$B_{max}/B_{int}$	$B_{int}/B_{min}$	$k$	IMF vs $k$ , deg	Polarization	$B_{\parallel}/B_{\perp}$
a	02.22.00 - 02.22.04	2.4	6.5	[0.98 0.01 0.18]	15.9	L	3.6
b	02.22.48 - 02.22.58	1.5	10.4	[0.96 -0.26 0.14]	0.7	L	2.2
c	02.24.00 - 02.24.05	1.1	14.1	[0.93 -0.22 0.30]	9.3	R	2.2
d	02.24.35 - 02.24.40	2.5	5.9	[0.91 -0.19 0.36]	13.2	R	2.3
e	02.25.30 - 02.25.35	1.4	1.4	[0.75 -0.29 0.59]	28	L	4.3

**Table 2.** Model shock parameters compared with SLAMS shock parameters in the time period of 02:00:00 - 02:30:00 24 October 2019.

Observation time, UT	Duration, s	$n_{BS}$	$\Theta_{Bn}$ , deg	$n_{SLAMS}$	$\Theta'_{Bn}$ , deg
02:24:47 - 02:25:47	10	[0.92 -0.11 0.38]	16.31	[0.88 0.47 0.11]	43.32
02:23:55 - 02:24:39	5	[0.93 -0.09 0.36]	15.79	[0.63 0.75 0.23]	64.57
02:23:11 - 02:23:39	8	[0.93 -0.09 0.35]	15.28	[0.93 0.26 -0.27]	39.41
02:22:15 - 02:22:47	10	[0.94 -0.09 0.33]	14.86	[0.99 0.09 -0.08]	24.64
02:20:19 - 02:21:03	14	[0.95 -0.07 0.29]	14.16	[0.87 0.48 0.14]	43.90
02:19:39 - 02:20:19	14	[0.96 -0.06 0.28]	13.96	[0.97 -0.14 0.19]	7.82
02:18:39 - 02:19:27	12	[0.96 -0.05 0.26]	13.84	[0.93 0.35 0.15]	35.84
02:17:47 - 02:18:35	12	[0.97 -0.05 0.24]	13.83	[0.99 0.13 0.04]	23.99
02:13:27 - 02:14:19	22	[0.94 0 0.13]	15.55	[0.94 -0.02 0.3]	17.76
02:12:35 - 02:13:27	19	[0.99 0.01 0.11]	16.27	[0.75 -0.59 -0.28]	33.18
Mean value	$12.60 \pm 3.72$	-	$14.99 \pm 0.85$	-	$33.43 \pm 11.96$