

Impact of ocean heat transport on the natural and forced variability of Arctic sea-ice in the GFDL CM2-O model suite

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November 30, 2022

Abstract

The impact of horizontal resolution on meridional Ocean Heat Transport (OHT) and sea ice in the Arctic is investigated using the GFDL CM2-O climate model suite (1° , $1/4^\circ$, and $1/10^\circ$) in both preindustrial control and climate change simulations. Results show an increase in OHT associated to a decrease in sea ice extent (SIE) in the Arctic on inter-annual and decadal time scales. This link, however, is not monotonic with spatial resolution. While OHT increases and SIE decreases from the Low to the Medium resolution, the reverse is true from the Medium to the High resolution. Differences in OHT and SIE between the three model configurations mostly arise from the preindustrial state. As the spatial resolution increases, the Irminger Current is favored at the expense of the North Atlantic Drift. This rerouting of water to the Western side of Greenland results in less heat delivered to the Arctic in the High resolution configuration than in its Medium counterpart. As a result, the Medium resolution configuration is in best agreement with observed SIE and Atlantic OHT. Concurrent with the change in the partitioning in volume is a change in deep convection centers from the Greenland-Irminger-Norwegian Seas in the Low resolution to the Labrador Sea in the Medium and High resolutions. Results suggest a coupling between OHT into the Arctic and deep convection in the North Atlantic.

1 **Impact of ocean heat transport on Arctic sea ice**
2 **variability in the GFDL CM2-O model suite**

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5 **Key Points:**

- 6 • Ocean heat transport into the Arctic does not systematically increase with hor-
7 zontal resolution in the GFDL CM2-O model suite.
8 • The eddy-permitting and eddy-rich configurations show a stronger response to cli-
9 mate change than the eddy-parameterized configuration.
10 • Flow partitioning in the northern North Atlantic and location of deep convection
11 centers are key to the heat transport into the Arctic.

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Abstract

The impact of horizontal resolution on meridional Ocean Heat Transport (OHT) and sea ice in the Arctic is investigated using the GFDL CM2-O climate model suite (1° , $1/4^\circ$, and $1/10^\circ$) in both preindustrial control and climate change simulations. Results show an increase in OHT associated to a decrease in sea ice extent (SIE) in the Arctic on inter-annual and decadal time scales. This link, however, is not monotonic with spatial resolution. While OHT increases and SIE decreases from the Low to the Medium resolution, the reverse is true from the Medium to the High resolution. Differences in OHT and SIE between the three model configurations mostly arise from the preindustrial state. As the spatial resolution increases, the Irminger Current is favored at the expense of the North Atlantic Drift. This rerouting of water to the Western side of Greenland results in less heat delivered to the Arctic in the High resolution configuration than in its Medium counterpart. As a result, the Medium resolution configuration is in best agreement with observed SIE and Atlantic OHT. Concurrent with the change in the partitioning in volume is a change in deep convection centers from the Greenland-Irminger-Norwegian Seas in the Low resolution to the Labrador Sea in the Medium and High resolutions. Results suggest a coupling between OHT into the Arctic and deep convection in the North Atlantic.

Plain Language Summary

The Arctic has experienced a dramatic decrease in its sea ice cover over the past four decades. One of the main drivers of this intense melting is ocean heat transport from lower latitudes into the Arctic. This transport takes place at three main gates linking the North Pacific and Atlantic oceans to the Arctic. Thus, proper representation of ocean currents and the associated heat transport is necessary to make accurate projections of the Arctic pack ice in climate models. Here, we study the response of the Arctic sea ice to an increase in atmospheric carbon dioxide concentration using three configurations of a climate model that differ in their horizontal resolution of the ocean. Changing resolution can affect the strength, pattern and amount of heat carried by the currents. Our results confirm that the greater the ocean heat transport into the Arctic, the lower the sea ice extent. In contrast with previous studies, however, the ocean heat transport does not systematically increase when refining the ocean horizontal resolution. This result points to the fact that not only the currents strength, but also the pathways are influenced by the ocean horizontal resolution, impacting the penetration of warm Atlantic waters into the Arctic.

1 Introduction

Three different ways of improving climate projections are increasing the complexity of climate processes, refining spatial resolution or advancing parameterizations. Refining spatial resolution is costly numerically, as the total integration time increases by a factor of at least 8 for each doubling of horizontal spatial resolution (Flato, 2011). It is also costly in terms of workforce since most parameterizations are still required and must be recalibrated as a function of newly resolved spatial scales (Molinari & Dudek, 1992). Human and computational resources in the last decade have been invested in the development of new or improved parameterizations of sub-grid scale processes (e.g. Fox-Kemper et al., 2011; Brankart, 2013; Jansen et al., 2015), increased ensemble size and number of scenarios, as well as on increasing spatial resolution of all the components of the climate system. Still, the majority of the Earth System Models participating to the Coupled Model Intercomparison Project version 6 (CMIP6) DECK use a 1° ocean component that require to employ eddy parameterizations Hewitt2020. In the context of Arctic climate, the new parameterizations include surface melt pond (M. M. Holland et al., 2012), ice thickness distribution (Bitz et al., 2001; Ungermann et al., 2017), lateral melt

62 (Tsamados et al., 2015; Smith et al., 2021) and ice-ocean heat exchange (Shi et al., 2020),
63 among others. These developments have led to significant improvements in the simula-
64 tion of the mean state and variability (forced and natural) of the ice-ocean system, in-
65 cluding the sea ice thickness distribution (Bitz et al., 2002; Bitz & Roe, 2004; Shi et al.,
66 2020), and sensitivity of the sea ice cover to increased carbon dioxide (CO₂) concentra-
67 tion (M. M. Holland et al., 2006; Stroeve et al., 2014; Jahn et al., 2016; Auclair & Trem-
68 blay, 2018).

69 Recently, climate groups have started to explore the sensitivity of the climate sys-
70 tem to an eddying ocean. For instance, the High Resolution Model Intercomparison Project
71 (HighResMIP) proposed a common protocol for low (1°) and high (1/4° to 1/12°) res-
72 olution model simulations under the umbrella of the World Climate Research Program
73 (WCRP; Haarsma et al., 2016). Studies using global climate models and ocean-only mod-
74 els have investigated the effect of refining spatial resolution on the sub-polar gyre and
75 Atlantic water pathways in the northern North Atlantic, Irminger Sea, Labrador Sea and
76 Baffin Bay in the context of ice shelf-ocean interactions and increased rate of advance
77 of tidewater glaciers (Myers et al., 2007; Straneo & Heimbach, 2013). Marzocchi et al.
78 (2015) find that a high resolution model (1/12° resolution) leads to an improved repre-
79 sentation of the subpolar gyre and a better representation of Labrador Sea Water for-
80 mation and variability compared to the 1° and 1/4° versions of the same model. Koenigk
81 et al. (2021) find that increasing the ocean model resolution from 1° to 1/4° leads to an
82 increase in deep mixing in the Labrador Sea and draw a direct link between the subpo-
83 lar gyre strength, surface ocean salinity and depth of convection. García-Quintana et al.
84 (2019) find less formation of Labrador Sea Water in a 1/12° model compared to a 1/4°
85 model, due to a shallowing of the mixed layer and a smaller area of deep convection. Pennelly
86 and Myers (2020) study the impact of resolution (from 1/4° to 1/12° to 1/60°) on Labrador
87 Sea circulation, and find that the mixed layer depth in the Labrador sea is shallower as
88 the resolution increases thanks to an increase in eddy kinetic energy, and that Labrador
89 Sea Waters density is better represented in the 1/60° model.

90 Several studies showed that an increase in resolution leads to an increase in mid-
91 latitude meridional ocean heat transport (OHT) in general (Griffies et al., 2015; Hewitt
92 et al., 2016) and in the Atlantic Ocean in particular (Grist et al., 2018). A better rep-
93 resentation of OHT is needed to improve projections of sea ice extent (SIE), as the ocean
94 is one of the main drivers of sea ice loss and variability in the Arctic (Bitz et al., 2005).
95 Indeed, in recent years, an increase in the Barents Sea Opening OHT led Atlantic Wa-
96 ters to penetrate deeper into the Eurasian Basin (Smedsrud et al., 2010), a process known
97 as the Atlantification of the Arctic (Årthun et al., 2012; Polyakov et al., 2017). This was
98 accompanied by a weakening of the stratification in the Eurasian Basin and enhanced
99 vertical heat fluxes from Atlantic Waters (Polyakov et al., 2017), and a limited winter
100 sea ice growth in the Barents Sea (Barton et al., 2018). Variability in Atlantic OHT is
101 responsible for the interannual variability SIE in the Barents Sea (Årthun et al., 2012,
102 2019). The impact of the Atlantic multidecadal variability on the Arctic SIE has been
103 highlighted especially for Barents Sea ocean surface temperature and ice extent (Drinkwater
104 et al., 2014; Årthun et al., 2019; Mette et al., 2021) and the Greenland Ice Sheet (Drinkwater
105 et al., 2014). Pacific Waters also play a key role in sea ice loss : for instance, Woodgate
106 et al. (2010) argued that a doubling of ocean heat flux through the Bering Strait between
107 2001 and 2007 was responsible for a third of the 2007 seasonal sea ice loss. Finally, cor-
108 relation between OHT and SIE is shown at interannual and decadal time scales during
109 rapid decline events in the Community Earth System Model - LE (Auclair & Tremblay,
110 2018; Li et al., 2017).

111 While the impact of spatial resolution on global scale circulation patterns has been
112 widely discussed, relatively fewer studies focus on the impact of resolution on OHT and
113 SIE variability in the Arctic Ocean. Griffies et al. (2015) find a lower poleward OHT in
114 the coarse resolution model configuration (1° resolution) than in the finer resolution model

115 configurations ($1/4^\circ$ and $1/10^\circ$ resolution), due to weaker sub-tropical and sub-polar gyre
 116 transports. Furthermore, increased ocean and atmosphere resolutions in the HadGEM3-
 117 GC2 model (from $1/4^\circ$ and 60 km to $1/12^\circ$ and 25 km, respectively), together with higher
 118 coupling frequency lead to stronger boundary currents, increased OHT, warmer surface
 119 ocean in the North Atlantic, and lower SIE (Hewitt et al., 2016; M. J. Roberts et al., 2016).
 120 Similarly, the ocean processes in the ECMWF-IFS are very sensitive to changes in ocean
 121 resolution from 1° to $1/4^\circ$, especially North Atlantic and Arctic Ocean, with improved
 122 representations of the Atlantic Meridional Overtuning Circulation (AMOC), OHT, and
 123 sea ice cover (C. D. Roberts et al., 2018). A recent study by Docquier et al. (2019) shows
 124 that, in the CMIP6 models participating in HighResMIP, the increase of spatial reso-
 125 lution from 1° to $1/4^\circ$ yields a larger Atlantic OHT and lower sea ice extent and volume.
 126 Furthermore, while the models exhibit strong correlations between the Atlantic OHT and
 127 the SIE variability in the Barents, Kara and Greenland Seas, the correlations do not in-
 128 crease uniformly with resolution across the models studied.

129 In the early 2010s, both the Geophysical Fluid Dynamics Laboratory (GFDL) and
 130 the National Center for Atmospheric Research (NCAR) have developed a climate model
 131 with an ocean component at $1/10^\circ$ for century scale simulations of the past, present and
 132 future climate (Delworth et al., 2012; Kirtman et al., 2012). Using the GFDL $1/10^\circ$ model,
 133 Griffies et al. (2015) find that mesoscale eddies play a significant role in the upward ver-
 134 tical heat transport and ocean heat uptake, and that this model yields a generally more
 135 accurate representation of global ocean temperature and heat budget. Using the same
 136 model, Saba et al. (2016) show that a refined resolution provides a more realistic rep-
 137 resentation of the Northwest Atlantic Shelf circulation, and a higher warming rate to in-
 138 creased CO_2 forcing. Dufour et al. (2017) show that this same model enables the for-
 139 mation of polynyas in the Weddell Sea compared to a coarser resolution, thanks to a stronger
 140 stratification in the Southern ocean and a better representation of transient eddies and
 141 topographical features. Drake et al. (2018) find that this fine resolution model leads to
 142 a significantly shorter advective upwelling time scale of Circumpolar Deep Waters in the
 143 Southern Ocean compared to the coarser resolution configurations, because of eddy vari-
 144 ability, thus highlighting the role of mesoscale eddies in large scale circulation time scale.

145 In this paper, we use the GFDL CM2-O model suite which comprises three con-
 146 figurations of different horizontal resolutions of the ocean component. We investigate the
 147 impact of refining the horizontal grid spacing of the ocean component on OHT in the
 148 Arctic, SIE and their relationship. We find that the magnitude of OHT and sea ice are
 149 strongly correlated on (multi) decadal time scales; however the links between OHT and
 150 SIE at interannual scale differ between model configurations. While the increase from
 151 the 1° resolution to the $1/4^\circ$ resolution does lead to an increase in OHT and decrease
 152 in SIE, the increase from the $1/4^\circ$ resolution to the $1/10^\circ$ leads to an opposite response.
 153 In addition, the change in resolution impacts the partitioning of North Atlantic heat trans-
 154 port thus resulting in different sea ice conditions.

155 The paper is structured as follows. In section 2, we present the GFDL CM2-O model
 156 suite and the simulations, and we describe the methods used to analyse the model out-
 157 put. In section 3, we present the SIE and OHT mean states, their response to an ide-
 158 alised climate change simulation as well as the impact of OHT on SIE. In section 4, we
 159 discuss the differences in the ocean circulation in the North Atlantic across the model
 160 suite and their potential impact on the OHT and sea ice.

161 **2 Model Description and Simulations**

162 **2.1 The CM2-O Model Suite**

163 In this study, we use the GFDL CM2-O model suite which comprises three con-
 164 figurations of the same climate model differing by the horizontal resolution of the ocean

165 component: CM2-1deg (1°; eddy-parameterized), CM2.5 (1/4°; eddy-permitting), and
 166 CM2.6 (1/10°; eddy-rich) (Delworth et al., 2012; Griffies et al., 2015). In the following,
 167 we refer to the three configurations as Low, Medium and High, respectively.

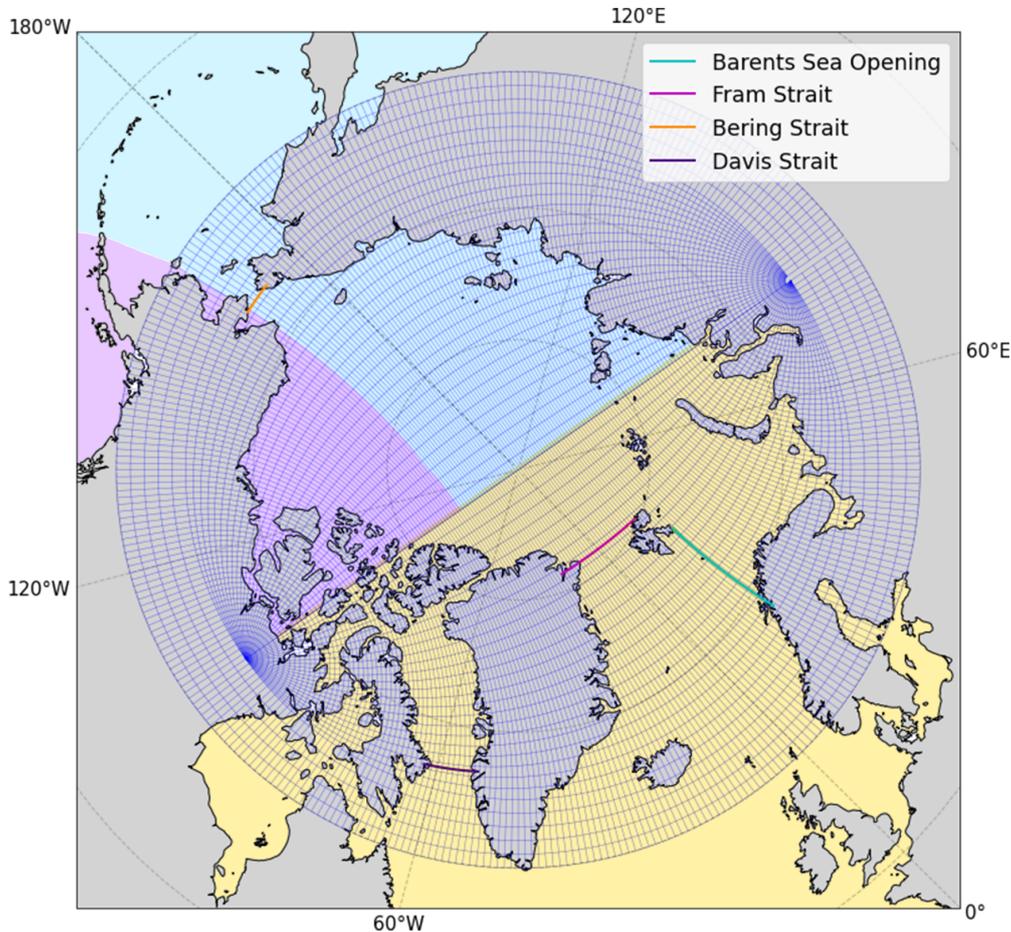


Figure 1. Arctic model domain and tripolar grid in CM2-1deg (Low resolution model configuration). The main gates used in the study are: the Fram Strait (pink), the Barents Sea Opening (cyan) the Bering Strait (orange) and the Davis Strait (purple). The coastlines are drawn from observations. The three sectors are : the Atlantic sector (yellow), the Pacific sector (purple) and the Eurasian sector (blue).

168 The ocean component is the version 5 of the Modular Ocean Model (MOM5; Griffies
 169 et al., 2015) run with volume-conserving Boussinesq kinematics. The model uses a tripolar
 170 grid, with one pole at the South Pole, and two poles placed over northern Canada
 171 and Russia (Figure 1; Murray, 1996). The ocean model is run with a z^* geopotential
 172 vertical coordinate (meaning that grid cell thickness is time dependent) and 50 layers
 173 in the vertical. At rest, the thickness of the layers ranges from 10 m in the first 250 m
 174 to 210 m at the bottom. The thickness of bottom cells is adjusted to match topography
 175 using the partial cell method (Pacanowski & Gnanadesikan, 1998). The model uses the
 176 piecewise parabolic method for the advection scheme (Delworth et al., 2012), and the
 177 non-local K-profile parameterization for vertical mixing (Large et al., 1994). The Low
 178 resolution model configuration includes the Ferrari et al. (2010) modified version of the

Table 1. Summary of key differences between the Low, Medium and High resolution model configurations of the CM2-O suite. The SIE trends are calculated over the 80 years of the CC simulation, and the observed trends are computed over the equivalent years of CO₂ concentrations (1979-2019; Fetterer et al., 2017). Note that the simulated SIE trends are linear over the 80 year period. Interannual variability is the standard deviation relative to a five-year running mean. OHT into the Arctic Ocean is defined as positive. The observed OHTs for the Fram Strait, Bering Strait and Barents Sea Opening are from Beszczynska-Möller et al. (2011), and the observational periods are 1997-2009, 1998-2007, and 1997-2007 respectively. The observed OHT for Davis Strait is from Cuny et al. (2004), and the observational period is 1987-1990. The observed Atlantic OHT at 26.5°N is from Johns et al. (2011). The model OHT is the average over the years with equivalent CO₂ concentration to the observation periods.

	Low	Medium	High	Observations
Nominal horizontal resolution (°)	1	1/4	1/10	-
Horizontal resolution at 65°N (km)	46 x 111	11 x 11	4 x 4	-
Mesoscale eddy parameterization	Yes	No	No	-
March SIE trend (10 ⁶ km ² / decade)	-0.1	-0.5	-0.5	-0.9
September SIE trend (10 ⁶ km ² / decade)	-0.3	-0.6	-0.6	-1.6
March SIE interannual variability (10 ⁶ km ²)	0.24	0.18	0.23	0.23
September SIE interannual variability (10 ⁶ km ²)	0.41	0.36	0.32	0.44
Fram Strait OHT (TW)	17	37	23	30 – 42
Bering Strait OHT (TW)	1	5	3	10 – 20
Barents Sea Opening OHT (TW)	3	76	38	50 – 70
Davis Strait OHT (TW)	5.1	9.5	18	1 – 35
Total Arctic OHT (TW)	26.1	127.5	82	91 – 167
Atlantic OHT at 26.5°N (TW)	670	560	800	1350

179 Gent and McWilliams mesoscale eddy parameterization (Gent et al., 1995) with a maximum
 180 diffusivity of 1200 m²s⁻¹ (Griffies et al., 2015) compared with 800 m²s⁻¹ in the
 181 ESM2M Earth System Model (Dunne et al., 2012), a model similar to the Low resolution
 182 in many aspects. The Medium and High resolution model configurations enable some
 183 explicit representation of the mesoscale, though incomplete, and do not use a mesoscale
 184 eddy parameterization (Griffies et al., 2015). The resolution needed to resolve the baro-
 185 clinic deformation radius in the Arctic ranges from 1/12° in the Central Arctic to 1/50°
 186 in the shallow waters near the coast (see Figure 2 of Hallberg, 2013). All three model
 187 configurations use the submesoscale mixed layer eddy parameterization of Fox-Kemper
 188 et al. (2011). Key characteristics of the model configurations are summarized in Table
 189 1.

190 In the High resolution configuration, the refined horizontal resolution allows for a
 191 better representation of the Gulf of Ob in the Kara Sea and of the Canadian Arctic Archipelago.
 192 Key differences between the High resolution and the Medium and Low resolutions also
 193 include the resolution of the Alpha and Lomonosov ridges, the Barents Sea and the steep-
 194 ness of the continental slopes. In the Medium resolution, the Victoria Strait, the Cor-
 195 onation Gulf, the Prince Regent Inlet and the Foxe Basin are closed. In the Low resolu-
 196 tion, the Fury and Hecla Strait connecting the Gulf of Boothia and Foxe Basin is closed.
 197 In contrast, all these basins and straits are open in the High resolution configuration.

198 The sea ice component is the GFDL Sea Ice Simulator (SIS) which uses a three-
 199 layer Semtner thermodynamic model (one layer of snow, two layers of ice) with five ice-
 200 thickness categories (Semtner, 1976; Winton, 2000; Delworth et al., 2006) and a brine
 201 pocket parameterization (Bitz & Lipscomb, 1999). The model uses the same tripolar grid
 202 as the ocean component (Dunne et al., 2012). The dynamic component of the sea ice model
 203 uses the elastic-viscous-plastic rheology of Hunke and Dukowicz (1997). The maximum
 204 value for albedos are set to 0.85 for snow on ice and 0.68 for bare sea ice (Delworth et
 205 al., 2012).

206 The atmospheric component is the GFDL AM2.1 (Atmospheric Model 2.1). AM2.1
 207 is run on a "cubed-sphere" grid with a horizontal resolution of 50 km and 32 vertical lev-
 208 els (Delworth et al., 2012), compared with 200 km and 24 levels in the GFDL CM2.1 de-
 209 scribed in Delworth et al. (2006). The advective terms are calculated with a modified
 210 Euler backward scheme (Kurihara & Tripoli, 1976). The atmospheric physics module
 211 is the GFDL AM2-LM2 model (Anderson et al., 2004) that includes three prognostic trac-
 212 ers for clouds: cloud liquid, cloud ice and cloud fraction. Finally, the suite uses the land
 213 component LM3 (Land Model 3) with a drainage route from Milly et al. (2014). More
 214 details about the suite or individual configurations' performance can be found in Delworth
 215 et al. (2012) and Griffies et al. (2015).

216 In the following sections, we will be discussing the model versions in this order :
 217 Medium, High and Low, as the Medium resolution model configuration is the closest to
 218 the observed SIE and OHT and the Low resolution model configuration is the farthest.

219 2.2 Simulations

220 We analyse a pre-industrial control run and a climate change run for each config-
 221 uration, hereafter referred to as *CTRL* and *CC*, respectively. The *CTRL* simulation is
 222 run for 200 years with constant globally averaged CO₂ concentration of 286 ppmv cor-
 223 responding to 1860. All model configurations started from the same initial conditions.
 224 The *CC* run branches off from the control run at year 121 with an atmospheric CO₂ con-
 225 centration increasing at 1% per year over 80 years leading to a doubling of CO₂ levels
 226 after 70 years (year 190 of the simulation). For the sake of clarity, we refer to the 80 years
 227 of the *CC* run as years 0 to 79 (not 121 to 200) in figures and text. Only one ensemble
 228 member was run for each of the configuration of the suite due to the high computational
 229 and storage cost of the high-resolution configuration.

230 In order to compare model output and observations, we use the annual mean CO₂
 231 atmospheric concentration from the Mauna Loa record (Keeling & Keeling, 2017). Note
 232 that the actual increase in CO₂ concentration is slower than the 1% CO₂ increase per
 233 year of the model. For this reason, the 41 years of satellite era from 1979 to 2020, cor-
 234 responding to CO₂ concentrations between 336.84 and 414.24 ppm, are compared to 21
 235 years in the *CC* run (years 16 to 37). In the following sections, the years between 1930
 236 and 1979, and between years 7 and 16 in the model, are referred to as the "pre-satellite"
 237 period.

238 2.3 Method

239 The total Ocean Heat Transport diagnostic (hereafter referred to as OHT) in the
 240 CM2-O suite is calculated online at each time step as $\int_{section} \rho_0 c_p U \Theta dS$ where ρ_0 is the
 241 constant Boussinesq reference density ($=1035 \text{ kg m}^{-3}$), c_p is the ocean heat capacity ($=3992.1$
 242 $\text{J kg}^{-1} \text{K}^{-1}$), U is the ocean velocity perpendicular to the section, Θ is the potential tem-
 243 perature, and dS is the surface of the grid cell normal to the flow. The OHT at each gate
 244 is calculated by integrating the monthly or yearly averaged OHT across the gate and the
 245 full water column. Each gate is located on the same constant latitude or longitude grid
 246 points in all the configurations, and is defined from the Low resolution for simplicity (Fig-

247 ure 1). We find that the positioning of the gates can have a minor impact on the mag-
 248 nitude of the OHT, but the changes are uniform across the configurations and within
 249 the ranges of observation errors at the gates (not shown). Furthermore, the positioning
 250 has a negligible impact on the variability (not shown). We analyse monthly mean out-
 251 put from the last 80 years of each simulation, except for the mass and heat transports
 252 of the High resolution where we use yearly means due to storage constraints. The inter-
 253 annual variability is defined as the variability around the five-year running mean.

254 The Arctic is divided into three sectors in the analysis presented in section 3.3 :
 255 the Atlantic sector, the Pacific sector and the Eurasian sector. The delimitations of those
 256 regions are shown in Figure 1.

257 **3 Results**

258 **3.1 Mean Arctic Ocean Climate over the Historical Period**

259 *3.1.1 Sea Ice Extent and Thickness*

260 Over the historical record, all three configurations reproduce the pan-Arctic win-
 261 ter sea ice thickness distribution with thicker ice on the Canadian side and thinner ice
 262 on the Eurasian side of the Arctic, and an east-west asymmetry north of the Canadian
 263 Arctic Archipelago (Figure 2 a-c). The winter sea ice thickness in the Low and High res-
 264 olutions is in general agreement with submarine observations from 1960 to 1982 (Bourke
 265 & Garrett, 1987), except along the Alaskan coastline where thicker ice is present in model
 266 configurations (2.5 to 3 meters-thick ice as opposed to 1 to 2 meters-thick ice in obser-
 267 vations), indicative of a small bias in the location of the Arctic High. In the Medium res-
 268 olution, the sea ice is too thin by a few meters in the Central Arctic and Canada Basin
 269 (2 meters-thick ice in the winter as opposed to 3-6 meters-thick ice in observations), and
 270 has a thick bias along the Alaskan coastline that is similar to the other configurations.
 271 In the High and Low resolutions, the thicker ice in the East Siberian Sea is typical of cli-
 272 mate models, where easterly winds interact with Wrangle Island and the New Siberian
 273 Islands (DeWeaver & Bitz, 2006). In the summer, the sea ice thickness is again in gen-
 274 eral agreement with observations in the Low and High resolutions, and too thin in the
 275 Medium resolution (Figure 2 d-f).

276 In the Medium resolution, the winter and summer SIE are in very good agreement
 277 with early satellite observations (Figure 3a). In the winter, in the Low and High reso-
 278 lutions, more sea ice is found in the Bering and Greenland Seas, suggesting a weaker sub-
 279 polar gyre in both the northern North Pacific and Atlantic (Figure 2a,c). The overes-
 280 timation of sea-ice in those regions leads to a larger March SIE in the Low and High res-
 281 olutions compared to the Medium resolution (Figure 3a). The thick bias in summer SIE
 282 is associated with an absence of sea ice melt in all peripheral seas (Figure 2d,f). This
 283 bias could be due to winter sea ice thickness anomalies in the western Arctic (Figures
 284 2a,c), or a smaller summer melt. We will see in section 3.3 that the sea ice in periph-
 285 eral seas is strongly correlated with the OHT into the Arctic, which is weaker in the Low
 286 and High resolutions. Despite a similar SIE bias, the Low and High resolutions strongly
 287 differ in their response to climate change, as the Low resolution has a much weaker trend
 288 than the High resolution (see Section 3.2 for details).

289 The September and March SIE of the Medium resolution are also in very good agree-
 290 ment with observations over the satellite era (1979-2019), with a small underestimation
 291 for September SIE, mostly in the Greenland and Barents Seas (Figures 3a and 4). Con-
 292 versely, in the Low and High resolutions, the September and March SIE (~ 9 and ~ 19
 293 million km^2) are too large by about ~ 1 to 3 million km^2 in September and 3 million
 294 km^2 in March. While the September total SIE is realistic in the Medium resolution, the
 295 spatial extent is too extensive in the East Siberian sea and too retreated in the Atlantic
 296 sector when compared to the satellite record (Figure 4). In the Low resolution, the Septem-

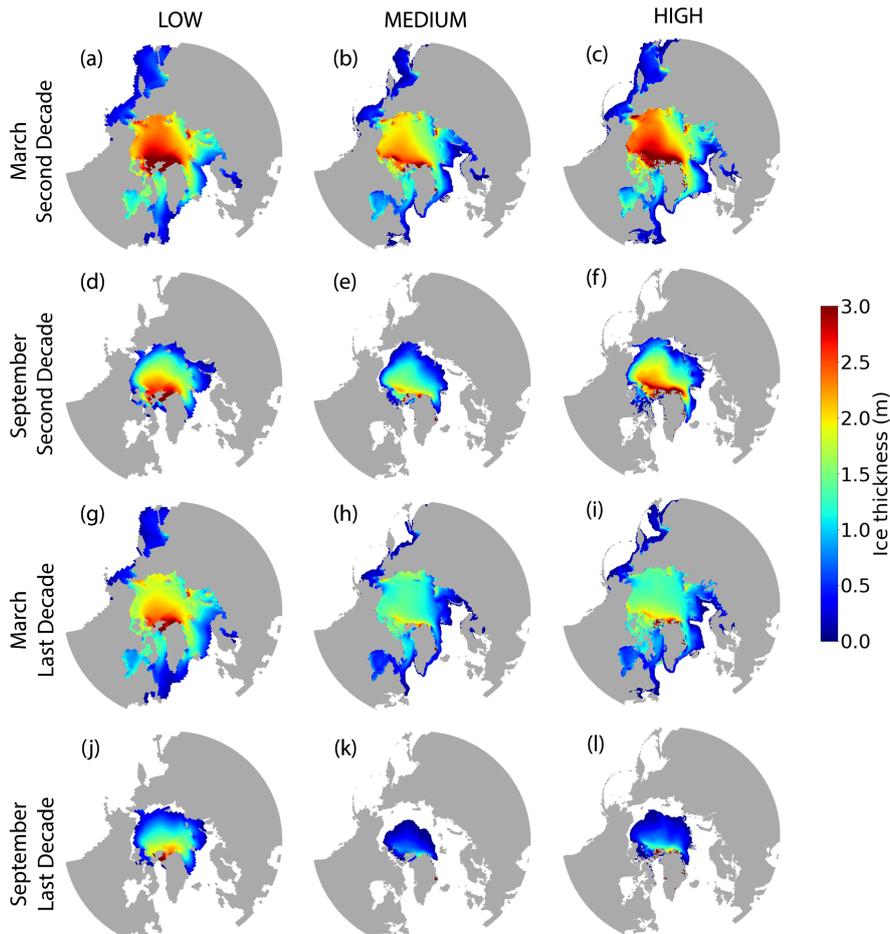


Figure 2. Mean sea ice thickness in the CC simulation averaged over the second decade (a-f) and the last decade (g-l) in March (a-c and g-i) and September (d-f and j-l) for the Low, Medium and High resolutions. The thicker ice reaches 4.5 m which is within realistic values, and some areas have an accumulation of anomalously thick ice due to the ice being trapped in the simulations (3 km thickness on the coast of Greenland for instance).

297 ber SIE is too large in all three sectors of the Arctic (Figure 4). The High resolution sea
 298 ice is too extensive in the Pacific and Eurasian sectors and in good agreement with the
 299 observations in the Atlantic sector (Figure 4). While the Medium resolution simulates
 300 the correct SIE, it does so with a much thinner ice cover throughout the simulation (as
 301 the initial sea ice thickness is thinner compared to observations from 1960 to 1982, Bourke
 302 & Garrett, 1987).

303 The interannual variability of SIE is in good agreement with observations in all three
 304 model configurations (see Table 1), though it is slightly underestimated in September.
 305 The increase in interannual variability observed during the transition to a seasonally ice-
 306 free Arctic is entirely missing in all the model configurations (not shown, Desmarais &
 307 Tremblay, 2021). The decadal variability of SIE is larger than observations in Septem-
 308 ber across the model suite, and in March for the Low resolution (see Fig 3a). Interan-
 309 nual variability and the correlations between SIE and OHT variability is discussed fur-
 310 ther in section 3.2.

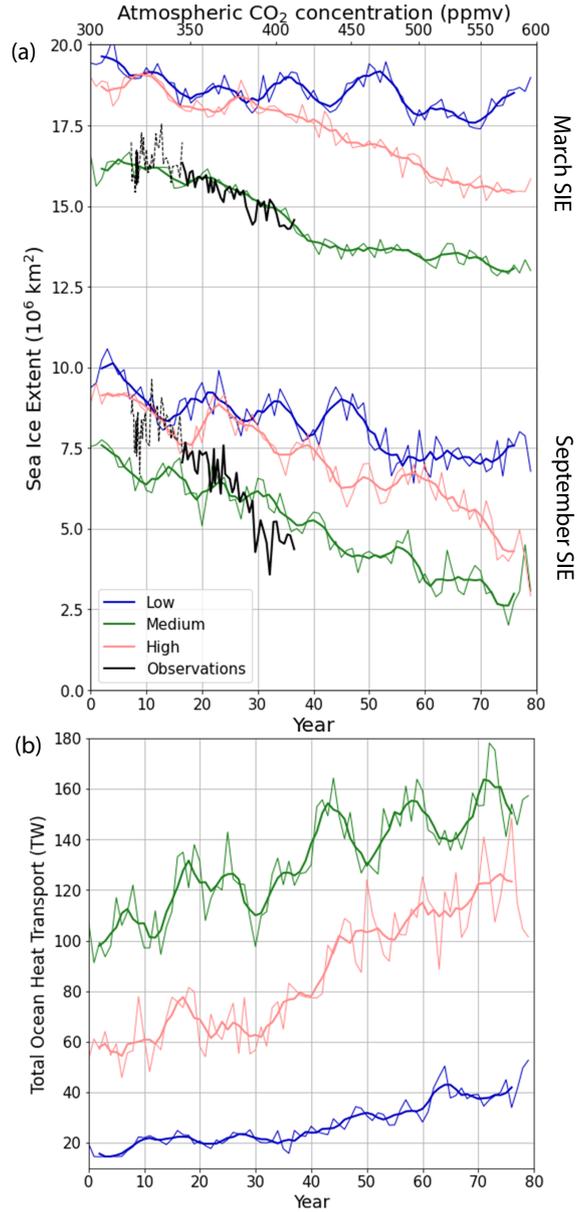


Figure 3. a) Observed March and September SIE between 1930 and 1979 from the historical record (dashed black line, Walsh et al., 2019) and between 1979 and 2019 from the satellite record (thick black line, Fetterer et al., 2017), and simulated March and September SIE (thin lines) and five-year running mean (thick lines) and b) Yearly mean total OHT into the Arctic as the sum of Barents Sea Opening, Fram Strait and Bering Strait OHT (thin lines) and five-year running mean (thick lines) as a function of time (model years; bottom axis) and CO_2 concentration (top axis) in the CC run for the Low, Medium and High resolutions. Note that the observations are plotted with respect to the CO_2 concentration for comparison with the model. The SIE is calculated as the area of grid cells where the sea ice concentration exceeds 15%.

311

3.1.2 Ocean Heat Transport

312

During the observational period (see Table 1, corresponding to the end of the third decade and the beginning of the fourth decade in the model), the Medium resolution has

313

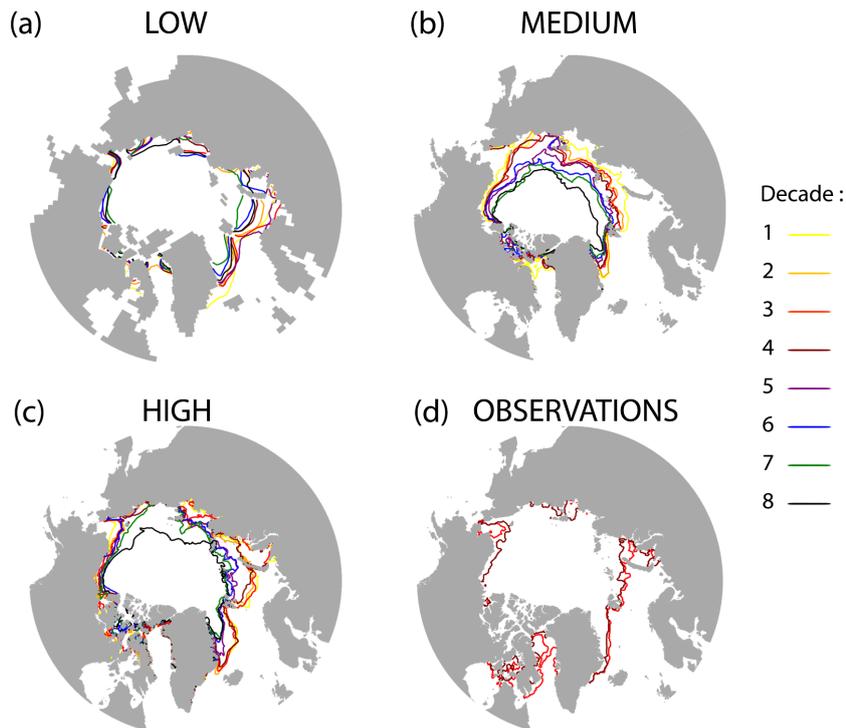


Figure 4. September sea ice edge averaged over each decade of the CC simulation for the (a) Low, (b) Medium and (c) High resolutions, and for (d) observations over the satellite record. The first decade begins with model year 121 while the last decade ends with model year 200. The satellite era in the model corresponds to the third and fourth decades according to equivalent CO_2 levels.

314 a total OHT of 112 TW into the Arctic and is in good agreement with observations in
 315 the Fram Strait. In the Barents Sea Opening, the simulated OHT is slightly overesti-
 316 mated compared to estimates from Beszczynska-Möller et al. (2011), but very close to
 317 an observational estimate of 73 TW (Smedsrud et al., 2010). The Low resolution greatly
 318 underestimates the total OHT, with little heat entering the Arctic through the Barents
 319 Sea Opening and Bering Strait (3 TW and 1 TW respectively; Table 1). The OHT through
 320 the Fram Strait is also underestimated, by at least 13 TW. This lack of heat transport
 321 is the result of low Atlantic waters intrusion onto the Barents Sea shelf in the Low Res-
 322 olution compared to the other two configurations (Figure 8a-c). This is presumably due
 323 to discrepancies in the large scale atmospheric circulation, since OHT variability is mostly
 324 driven by volume transport variability (Madonna & Sandø, 2022), rather than its low
 325 spatial resolution, given that other climate models with similar spatial resolution sim-
 326 ulates much higher ocean heat transport in the Barents Sea Opening (e.g. the Commu-
 327 nity Earth System Model, Auclair & Tremblay, 2018). In the High resolution, the OHTs
 328 in the Fram Strait and Barents Sea Opening are underestimated by at least 7 TW and
 329 12 TW respectively (Table 1). All model configurations strongly underestimate the OHT
 330 across the Bering Strait with the modelled OHTs reaching at most 50% of the observa-
 331 tional estimates. The OHT at 26.5°N in the Low (0.67 PW), Medium (0.56 PW) and
 332 High (0.80 PW) resolutions for the third decade are comparable, although somewhat lower,
 333 compared to that of the RAPID array (1.35 PW, Johns et al., 2011).

334 Over the observational period, the Medium resolution is the closest to the obser-
 335 vational estimates of total OHT and SIE (Table 1 and Figure 3). Of the three resolu-
 336 tions, that model configuration also carries the most heat into the Arctic ($\sim 50\%$ more
 337 heat than the High resolution). Both the Low and High resolutions underestimate the
 338 OHT and overestimate SIE over the observational period, with the High resolution show-
 339 ing significantly greater OHT but only slightly lower SIE than its lower resolution coun-
 340 terpart. Hence, in the CM2-O model suite, the greater the OHT, the lower the SIE, which
 341 suggests a major impact of OHT on SIE, in agreement with several studies (Mahlstein
 342 & Knutti, 2011; Sandø et al., 2014; Li et al., 2017; Muilwijk et al., 2019; Docquier et al.,
 343 2019).

344 3.2 Impact of OHT on SIE at a Pan-Arctic Scale

345 In response to the CO₂ forcing, all configurations show a linear decline in SIE with
 346 a clear decadal to multidecadal signal super-imposed (Figure 3a, Table 1). The trends
 347 in the September SIE in the Medium and High resolutions are around -0.6×10^6 km²/
 348 model decade (significant at the 95% confidence level), much smaller in absolute value
 349 than the observed trend of -1.6 million km²/model decade in the satellite era. We note
 350 that, even without adjusting the observed trend to the CO₂ concentration in the model
 351 simulation, the trend in observations is still higher than in the model configurations ($-$
 352 0.8 million km²/decade, Onarheim et al., 2018). The underestimation of September sea
 353 ice decline in the CM2-O suite is common among climate models; for instance, the CMIP6
 354 multi-model mean trend is -0.7 million km² /decade (Shu et al., 2020). The trends in
 355 the March SIE are $\sim 50\%$ of the observed trend in the Medium and High resolution mod-
 356 els (significantly different than zero), and comparable to that of the 1980-1999 observa-
 357 tional record in the Low resolution model (non significant, not shown). Note that the
 358 Medium resolution is in very good agreement with observations over the satellite era.

359 All three simulations have a weak trend in sea ice extent compared to observations
 360 and do not reach an ice-free Arctic (defined as SIE < 1 million km²; IPCC, 2013) af-
 361 ter a doubling of CO₂ concentration. Whether this is caused by too weak OHT in the
 362 Arctic or other processes (e.g. atmospheric circulation, cloud phase, etc.) is unclear and
 363 beyond the scope of the paper. The minima of SIE reached by the CM2-O suite at the
 364 end of the CC simulation are generally higher than in the other models participating in
 365 CMIP6. Indeed, the majority of climate model simulations reach a sea ice free Arctic
 366 in the summer by the year 2050 with a CO₂ concentration ranging between 500 and 550
 367 ppm depending on the emission scenario (Figure 3 and Table S4 of Notz & SIMIP Com-
 368 munity, 2020).

369 We note that the High resolution loses significantly more sea ice under climate change
 370 than the Low resolution (Figures 2 and 3a, and Table 1), though both have very sim-
 371 ilar initial conditions throughout the preindustrial era (not shown). Conversely, the Medium
 372 and High resolutions display the same trends under climate change in both seasons de-
 373 spite starting from very different SIE preindustrial conditions (Figure 3a and Table 1).
 374 Hence, the lower SIE at the end of the CC run in the Medium resolution is mostly due
 375 to the preindustrial mean state (low initial sea ice cover), rather than to a strong response
 376 to the CO₂ increase.

377 The OHT is sensitive to the CO₂ increase in all three model configurations, but
 378 the intensity of the response varies across the configurations (Figure 3b). By the end of
 379 the simulation, the total OHT has increased by $\sim 50\%$ in the Medium resolution while
 380 it has doubled in the High and Low resolutions. In the Medium resolution, the OHT in-
 381 crease is mostly linear, with a strong decadal variability. In the Low and High resolu-
 382 tions, a significant multi-decadal signal is super imposed on the linear increase in OHT,
 383 resulting in two "apparent" stable periods without OHT trends (in the first three decades

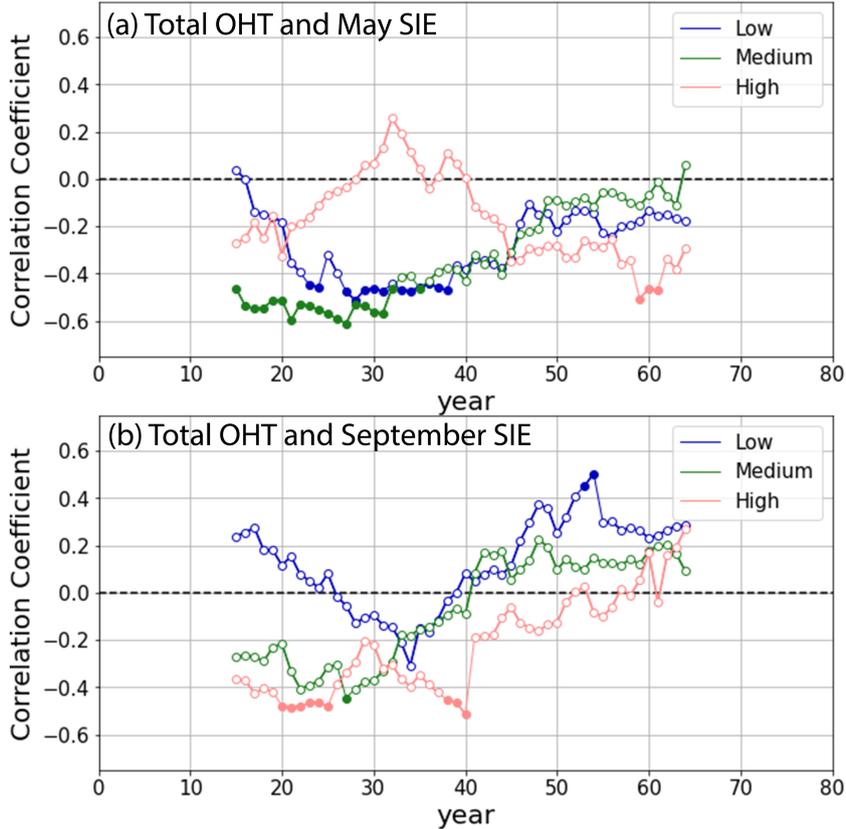


Figure 5. Twenty-year moving window correlation between the detrended annual (January–December) total OHT and the detrended (a) May SIE and (b) September SIE in the CC run. Full circles indicate instances where the correlation exceeds the 95% confidence level.

384 and last two-three decades) and a relatively rapid increase between the fourth and fifth
 385 decades (see Figure 3b).

386 In the Medium resolution, we see a weak signal at decadal time scale in March SIE
 387 in the first half of the record, and a stronger decadal signal in September SIE that per-
 388 sists until the end of the simulation (see Figure 3a). We will see in Section 3.3 that this
 389 signal is driven mostly by the OHT from the Atlantic driving sea ice loss in the Green-
 390 land and Barents Seas. We note that the signal is not as strong as for the High resolu-
 391 tion. Presumably, this is due to the fact that the sea ice cover retreats north of the Bar-
 392 ents Sea continental shelf in the middle of the simulation (\sim year 30, i.e. between the
 393 third and fourth decade; see Figure 4), at which point the ocean heat is not in direct con-
 394 tact with the sea ice anymore (Auclair & Tremblay, 2018). Similarly, at an interannual
 395 time scale, the total OHT in the Medium resolution is negatively correlated with the May
 396 SIE until \sim year 30, after which the correlation reduces (Figure 5) when the sea ice has
 397 completely retreated in the Barents Sea.

398 In the Low resolution, the decadal variability in SIE and OHT are the largest and
 399 smallest (respectively) of the CM2-O suite (Figure 3). Hence, the decadal variability in
 400 the pan-Arctic SIE is not dominated by OHT variability in that configuration. We will
 401 see in Section 3.3 that the OHT and SIE are linked at regional scale (i.e. in the Atlantic
 402 and Pacific sectors), but that the two regional signals are out of phase and not appar-

ent in the total SIE and OHT. At the interannual time scale, the total OHT is significantly correlated with May SIE in the Low resolution from year 23 until year 38, with higher OHT leading to lower SIE. From year 38 onward, no significant correlation is found (Figure 5). Several studies highlight that the atmosphere-ocean coupling is generally weaker and poorly represented when the ocean component is at a low, non-eddying spatial resolution (Bryan et al., 2010) resulting in less air-sea fluxes especially in the North Atlantic (M. J. Roberts et al., 2016). This weak atmosphere-ocean coupling should in principle lead to stronger correlation between ice edge location and OHT variability, as seen in other low resolution GCMs (e.g. Auclair & Tremblay, 2018). The absence of SIE - OHT coupling in the Low resolution is instead attributed to the very small OHT (non differentiable from noise) through the Barents Sea Opening in this configuration.

In the High resolution, the variability in OHT at decadal time scale is linked with variability in September SIE (correlation coefficient of -0.51 significant at the 95% confidence level; Figure 3). At interannual time scale, the total OHT is correlated with September SIE until \sim year 160, and with May SIE in the last 20 years of the simulation, although the signal is not robust (i.e. the correlation is only significant for the last few years; Figure 5). Again, the shift in correlations at year 160 corresponds to a significant retreat of sea ice in the Barents Sea (Figure 4 and discussion in section 3.3).

The links between the interannual variability in SIE and total OHT in the CM2-O model suite do not persist throughout the CC simulation, and are not always present in the CTRL simulation (not shown). Hence, OHT variability is not the only driver of SIE variability for any of the model configurations on a global and Pan-Arctic scale, where atmospheric processes also play a key role. However, all configurations show correlations between SIE and OHT at decadal or interannual scale at the beginning of the simulation (except for the September SIE and OHT in the Low resolution model configuration), which corresponds to the period when sea ice cover is larger, especially in the Barents Sea where a strong influence of the ocean on sea ice is expected (Auclair & Tremblay, 2018; Årthun et al., 2012). This suggests that OHT variability is a major driver of sea ice variability at regional scale, especially when the sea ice extends to the Barents Sea where ocean-ice interactions are more important.

3.3 Impact of OHT on SIE at Regional Scale

3.3.1 Temporal Scales of Correlations between OHT and SIE

All model versions show an increase in OHT at the three gates concurrent with a decrease in SIE in the three main Arctic sectors in the CC simulations (Figure 6). There is an exception in the March SIE for the Low resolution (Atlantic sector) which shows an increase in SIE (years 165-175) despite the increase in OHT, indicating that the natural variability at decadal time scale in this version is larger than the forced change associated with the CO₂ increase. This partly explains the very weak March SIE trend on the Pan-Arctic scale discussed in section 3.2. At the multi-decadal time scale, the Medium and High resolutions show an abrupt increase in OHT in the Barents Sea Opening at the mid-simulation that is concurrent with an abrupt decline in SIE mostly visible in March in the Atlantic (Figures 6 a-c and 4). The September SIE does not react to the abrupt change in Barents Sea Opening OHT, however, as the September sea ice covers only a small part of the Barents Sea shelf. Furthermore, the weak reaction to OHT changes is indicative that summer processes (e.g. ice-albedo feedback) have more impact than later winter preconditioning in the model suite.

In the Medium resolution, the Fram Strait OHT increases in the second decade by about 15 TW, which is concurrent with a very slight local minimum in SIE. The Fram Strait OHT sees another sharp increase of 10 TW in the fourth decade, which is followed by an abrupt increase of 20 TW in the Barents Sea Opening in the fifth decade. Those increases match a sudden decrease in March SIE in the Atlantic sector that is sustained

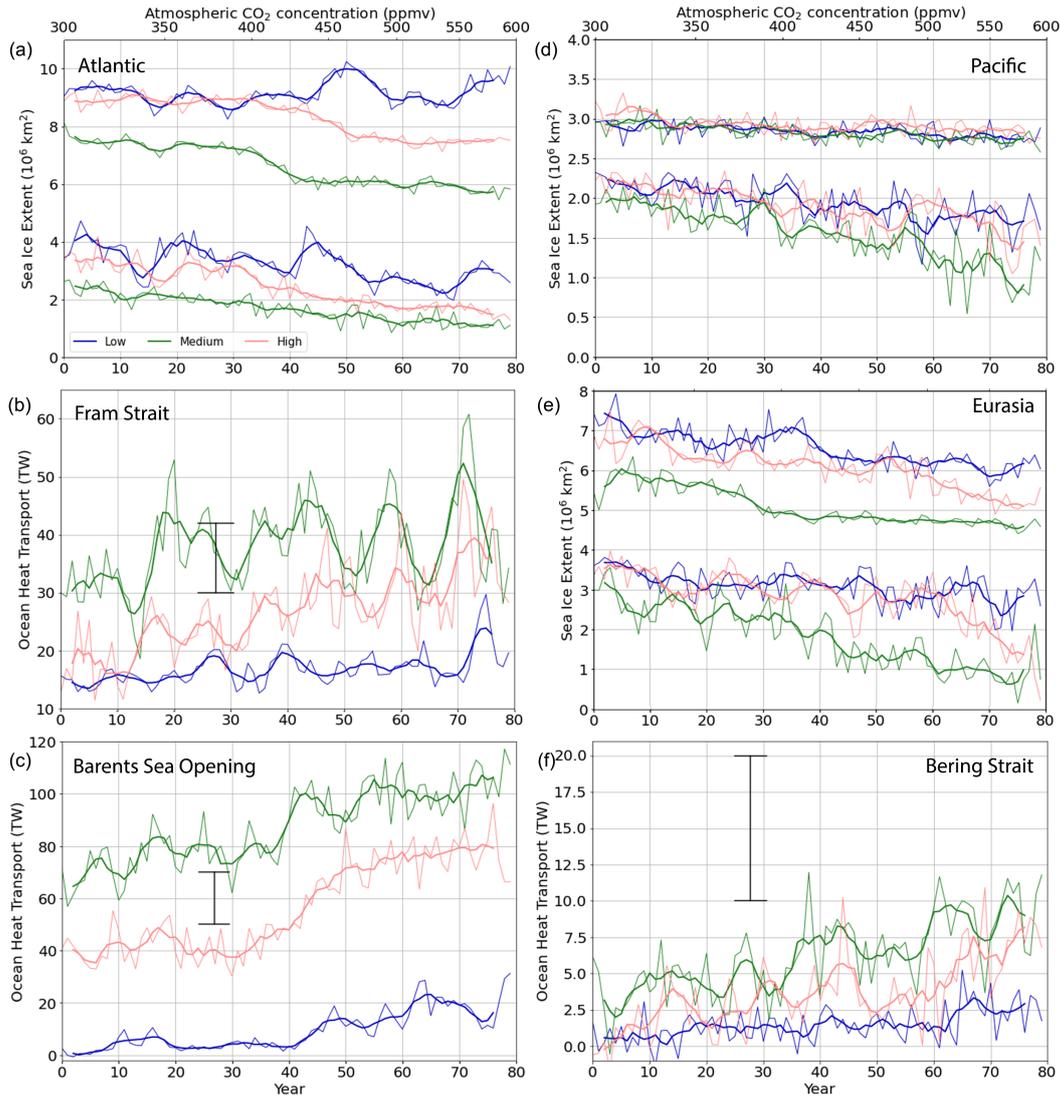


Figure 6. March and September SIE (thin lines) and five year running mean (thick lines) in (a) the Atlantic sector, (d) the Pacific sector and (e) the Eurasian sector, and annual OHT (thin lines) and five year running mean (thick lines) through (b) the Fram Strait, (c) the Barents Sea Opening and (f) the Bering Strait in the CC run for the Low, Medium and High resolutions as a function of time (top axis) and CO₂ concentration (bottom axis). Observational estimates are indicated as vertical bars with the horizontal line corresponding to the time period of observations: (b) 1997-2009 (Schauer & Beszczynska-Möller, 2009), (c) 1997-2007 (Beszczynska-Möller et al., 2011), and (f) 1998-2007 (Woodgate et al., 2010).

454 until the end of the simulation (Figure 6 a-c). The Bering Strait OHT increases through-
 455 out the simulation, with a sharper increase in the fourth decade that also matches sig-
 456 nificant sea ice loss in the Eurasian sector (Figures 4 b and 6 e-f). By the end of the sim-
 457 ulation, the OHT in the Bering Strait reaches the lower range of current observations
 458 (10 TW).

459 In the High resolution, the OHT remains fairly constant in the Barents Sea Open-
 460 ing until the fourth decade (equivalent CO₂ concentration around 400 ppmv) when an
 461 OHT increase of 30 TW occurs, after which the OHT stabilizes again until the end of
 462 the simulation (Figure 6 c). These changes match well the pattern of sea ice melt in the
 463 Atlantic Sector in March (Figures 4 c and 6 a). We note that while the September sea
 464 ice loss is concurrent with the Barents Sea Opening OHT increase in the High resolu-
 465 tion, the March sea ice loss is delayed by ~ 10 years. In the Atlantic sector, the decadal
 466 variability in the September SIE is driven by Fram Strait and Barents Sea Opening vari-
 467 ability in the first half of the simulation (with a significant correlation coefficients of -
 468 0.92 between the September SIE and the Fram Strait and Barents Sea Opening OHT).
 469 The decadal variability in the Bering Strait OHT is also well correlated with the decadal
 470 variability in September SIE over the whole simulation in the Eurasian Sector (with a
 471 significant correlation coefficients of -0.66 between the September SIE and the Bering
 472 Strait OHT), as a sharp increase in Bering Strait OHT in the last 25 years of the sim-
 473 ulation is concurrent with a decline in March and September SIE in the Eurasian sec-
 474 tor (Figure 6 e-f).

475 In the Low resolution, the significant increase in Barents Sea Opening OHT hap-
 476 pens around the fourth decade when the OHT goes from near zero to about 20 TW by
 477 the end of the simulation. This increase in OHT is concurrent with the retreat of sea ice
 478 in the Atlantic sector (Figure 6 a and c) and especially the Barents Sea (Figure 4a) af-
 479 ter the fourth decade. The decadal variability in the Fram Strait and Barents Sea Open-
 480 ing OHT are well correlated with decadal variability in the Atlantic sector September
 481 SIE during the first half of the simulation (with a significant correlation coefficient of -
 482 0.69 between the September SIE and the Fram Strait and Barents Sea Opening OHT).
 483 In the Eurasian sector, the decrease in SIE at the end of the simulation is concurrent with
 484 an OHT increase in the Bering Strait.

485 *3.3.2 Spatial Patterns of Correlations between OHT and SIE*

486 We now turn to spatial correlations between OHT and Sea Ice Concentration (SIC)
 487 anomalies to unravel some major modes of variability at the Pan-Arctic scale and the
 488 impact of OHT on sea ice decline at the regional scale.

489 A tripole between the three sectors defined in Figure 1 appears in the Low and High
 490 resolutions, with the Bering Strait OHT and SIC anomalies having opposite sign cor-
 491 relations in the Eurasian sector and the Atlantic/Pacific sector (Figure 7 a-c). In the Medium
 492 resolution, the Bering Strait OHT is still positively correlated with the SIC in the At-
 493 lantic sector, but negatively correlated in the Pacific sector and on the shelf in the Eurasian
 494 sector (we also note an anticorrelation in the Eurasian sector away from the shelf, although
 495 it is not significant). This is in accord with results from the CESM-LE (Auclair & Trem-
 496 blay, 2018), and follows from the fact that, to first order, the volume of water in the Arc-
 497 tic is conserved, hence there is a compensation of ocean volume transport (OVT) between
 498 the two sectors (Timmermans & Marshall, 2020). In the Medium and High resolutions,
 499 we also find a consistent dipole with opposite sign correlations between SIC variability
 500 in the Barents/Greenland seas, and the Labrador Sea. This is a standard signal in the
 501 observational record linked with the North Atlantic Oscillation (NAO) variability (Venegas
 502 & Mysak, 2000). In the Medium resolution, the correlations are weaker in the Barents
 503 Sea Opening because the sea ice edge is retreated northward compared to the Low and
 504 High resolutions (see Figure 2, 4).

505 In the Medium resolution, the Bering Strait OHT is correlated negatively with most
 506 of the Pacific side of the Arctic, even well into the Kara Sea, and is positively correlated
 507 with SIC in the Barents Sea and Greenland Sea (Figure 7 b). This is in accord with the
 508 three major pathways of Pacific Waters into the Arctic : the Alaskan current branch,
 509 the branch that spills over the Chukchi shelf and enters the Canada/Makarov Basin, and

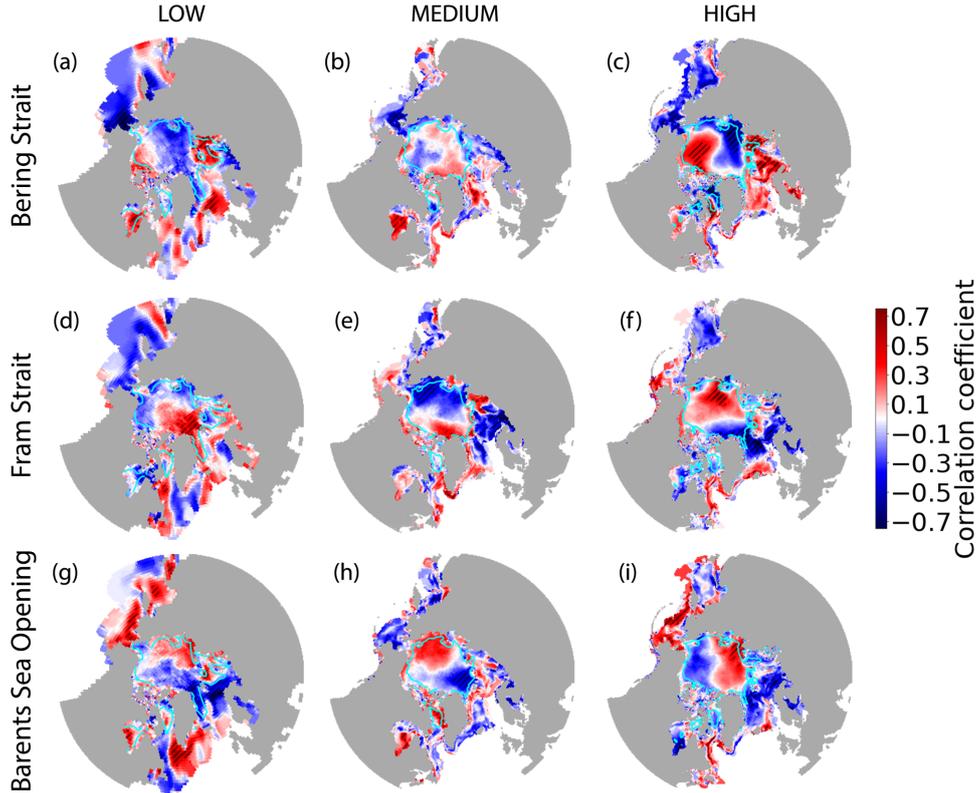


Figure 7. Correlation maps between the detrended annual (January-December) OHT in the Bering Strait (a-c), Fram Strait (d-f) and Barents Sea Opening (g-i) and the detrended May sea ice concentration (SIC), in the Low (left), Medium (center) and High (right) resolutions in the CC experiments. Inside the blue contour lines are areas where the SIC varies by less than 5%. The dashed areas is the 95% significance level. We note that the intensity of the correlation may vary depending on the month used for SIC, but the correlation patterns are similar.

510 the branch that stays on the Eurasian shelf (Pickart, 2004; Yamamoto-Kawai et al., 2008).
 511 The Fram Strait OHT is strongly linked with sea ice melt in the Greenland Sea, Bar-
 512 rents Sea and even in the Chuchki Sea (Figure 7 e). The Barents Sea Opening OHT is
 513 significantly anti-correlated with SIC in the Central Arctic, and a weak but widespread
 514 negative correlation pattern appears in the Barents Sea and in the Eurasian Basin (Fig-
 515 ure 7 h). The weakness of this negative correlation (non-significant at the 95% level)
 516 in the Barents Sea is surprising, but could be due to several factors such as the lack of sea
 517 ice in that area or the importance of summer processes unrelated to OHT (e.g. surface
 518 albedo etc.). OHTs across the three main gates are shown to be mostly positively cor-
 519 related with SIC into the Baffin Bay, Hudson Bay and Labrador Sea, which reflects a
 520 partitioning of the heat transport between the Arctic and the Irminger current, associ-
 521 ated with the NAO variability Straneo and Heimbach (2013) as we will see in section 4.

522 In the Low resolution, we see a significant negative correlation between Bering Strait
 523 OHT and SIC in the East Siberian sector (Figure 7 a). In this configuration, the branch
 524 of Pacific Waters that stays on the Eurasian shelf is dominant for the sea ice variabil-
 525 ity, in accord with the CESM-LE (Auclair & Tremblay, 2018). The Fram Strait OHT
 526 is significantly correlated to sea ice loss in the Greenland Sea and the Labrador Sea, as

527 well as on the Barents Sea Shelf, but positively correlated with SIC around the Fram Strait
 528 itself (Figure 7 d). Finally, the Barents Sea Opening OHT is strongly correlated with
 529 sea ice loss both in the Barents Sea and the Fram Strait (Figure 7 g).

530 In the High resolution, the Bering Strait OHT is negatively correlated with SIC
 531 in the Bering Sea and Chuchki Sea, as well as the Baffin Bay. For both Fram Strait and
 532 Barents Sea Opening OHT, the negative correlations with SIC are significant in the Green-
 533 land Sea, the Fram Strait and the Barents Sea and Kara Sea (Figure 7 f,i). We also note
 534 that Fram Strait and Barents Sea Opening OHTs are positively correlated with SIC in
 535 the Baffin Bay and Labrador sea, although the correlations are less significant (again,
 536 this is the typical dipole in SIC in the Barents and Labrador Seas).

537 This analysis reveals more robust coupling between OHT and SIC at the regional
 538 scale, especially in the Atlantic sector where Atlantic sea ice loss is driven by OHT in-
 539 crease, in particular in the Barents Sea as shown in Figure 4. We note that the corre-
 540 lations can weaken depending on the month used for the calculation, as atmospheric pro-
 541 cesses play a more important role in late summer SIE, however the patterns of negative
 542 correlations mostly remain consistent (not shown). Significant correlations at interan-
 543 nual and decadal time scales are exhibited between Bering Strait OHT and SIE in the
 544 Eurasian Sector, and between the Fram Strait and Barents Sea Opening OHT in the Green-
 545 land and Barents Seas.

546 4 Discussion

547 Of the three model configurations, the Medium resolution has the largest OHT into
 548 the Arctic and smallest winter and summer SIE, both of which are in good agreement
 549 with observations. The correct seasonal cycle in SIE is achieved at the expense of a thin
 550 bias in sea ice thickness. We find that in the CM2-O suite, the OHT increases with in-
 551 creasing resolution from the Low to the Medium resolution, in agreement with results
 552 from Docquier et al. (2019), but decreases as the resolution increases further from the
 553 Medium to the High resolution. We find that in the CM2-O suite, the OHT increases
 554 with increasing resolution from the Low to the Medium resolution model configuration,
 555 but decreases as the resolution increases further from the Medium to the High resolu-
 556 tion model configuration. This non-monotonic behaviour with spatial resolution is in con-
 557 trast with other studies, which show a systematic increase in OHT with a finer resolu-
 558 tion (M. J. Roberts et al., 2016; Grist et al., 2018; Docquier et al., 2019). The possible
 559 explanations for this specificity are discussed in this section, with the main candidate
 560 being the different partitioning of Atlantic Waters between the Barents Sea Opening, Fram
 561 Strait and Irminger Current between the model configurations. Apart from the non-monotonous
 562 increase of OHT with resolution, the other results found in the CM2-O model suite are
 563 robust across model families, including weaker ocean heat transport into the Arctic for
 564 low resolution (1°) models, and an increase in ocean heat transport northward (North
 565 Atlantic Drift or Irminger branch) as the spatial resolution increases (M. J. Roberts et
 566 al., 2016; Grist et al., 2018; Docquier et al., 2019). This conclusion is robust with respect
 567 to the exact location where OHT is calculated: i.e., along latitudinal transect at 60°N
 568 and 65°N as in M. J. Roberts et al. (2016); Grist et al. (2018); Docquier et al. (2019)
 569 or at Arctic gates (results not shown).

570 The increase in OHT in response to the CO_2 increase is slightly larger in the High
 571 resolution than in the Medium resolution, so that the higher OHT and lower SIE in the
 572 Medium resolution at the end of the CC simulation are primarily due to the preindus-
 573 trial mean state. The High resolution OHT is larger than that of the Low resolution, yet
 574 the mean sea ice states in the preindustrial and early CC simulations are similar. This
 575 is in contrast with the study by Kirtman et al. (2012) who also find a larger OHT when
 576 increasing the resolution in their analysis of the NCAR Community Climate System Model
 577 version 3.5 (CCSM3.5) from 1° to $1/10^\circ$ but a smaller sea ice extent in the High reso-

578 lution. The low OHT in the NCAR Low resolution model is mostly attributed to the poor
 579 representation of the Norwegian Coastal Current in the model, in accord with results from
 580 the CM2-O Low resolution (Figure 8 g). The decrease in OHT from the Medium res-
 581 olution to the High resolution is also in contrast with the results from Hewitt et al. (2016)
 582 though the resolution of the atmosphere and the frequency of the ocean/atmosphere cou-
 583 pling is also increased between their two model versions. We note that OHT and SIE
 584 correlations are not sensitive to an increase in spatial resolution of the atmosphere com-
 585 ponent (Docquier et al., 2019). The stronger OHT in the CM2-O Medium resolution oc-
 586 curs despite a weaker AMOC (not shown), in agreement with Oldenburg et al. (2018)
 587 and in contrast with results by Jackson et al. (2020). This suggests that the higher OHT
 588 in the Medium resolution is linked with the surface ocean circulation (gyre transport)
 589 rather than the meridional circulation (Griffies et al., 2015). We argue that differences
 590 in current pathways could explain the changes in Arctic OHT in the model versions.

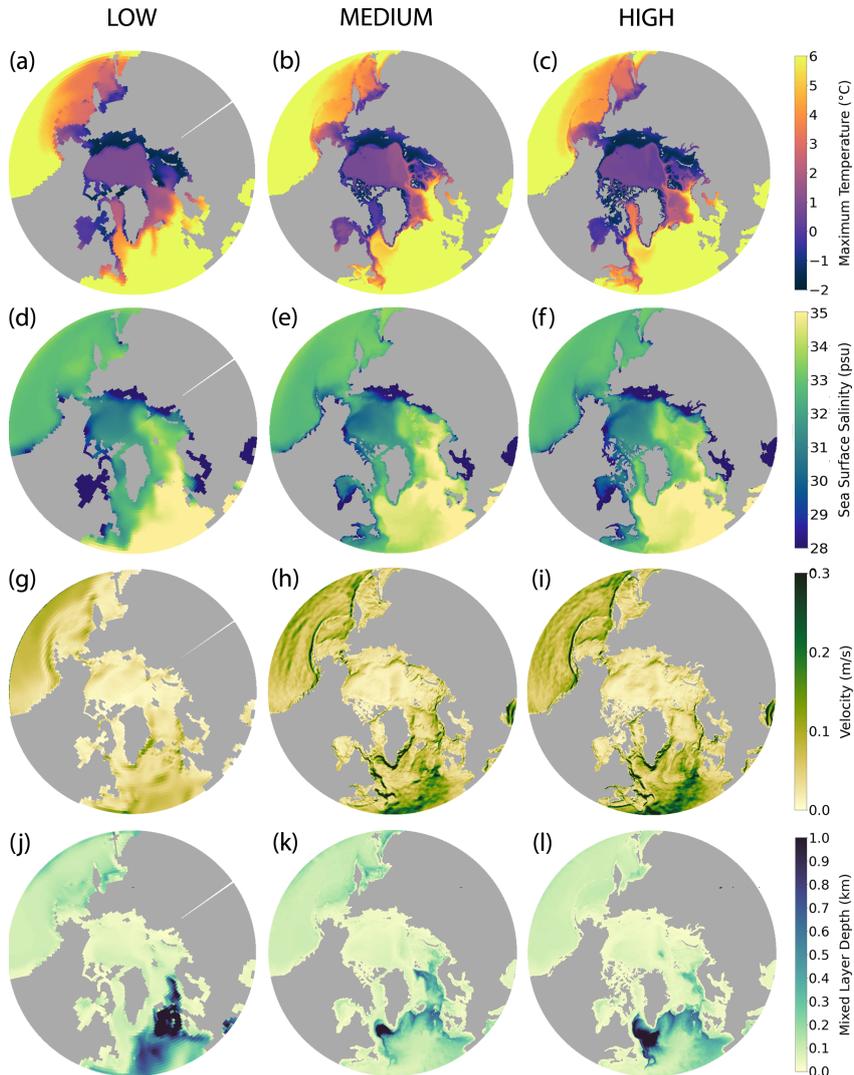


Figure 8. Mean temperature maximum (a-c), sea surface salinity (d-f), surface velocities (g-i) and winter mixed layer depth (h-l) averaged over the first decade (years 120-129) of the CC experiment for the Low (left column), Medium (middle column) and High (right column) resolutions. Note that the colorbars are always the same between the configurations.

591 All three model versions agree broadly in the structure of the currents in north-
 592 ern North Atlantic (Figure 8 g-i). The Low resolution however has broader and signif-
 593 icantly weaker currents than the Medium and High resolutions over the Arctic. This re-
 594 sult is in agreement with Docquier et al. (2019). The most striking difference with the
 595 other two resolutions is the absence of the West Greenland Current and Labrador cur-
 596 rent at the surface (Figure 8 g). In the Low resolution, Atlantic Waters enter the Labrador
 597 Sea and Baffin Bay at depth (Figure 8 a, d) and fresh cold Arctic Waters - entering from
 598 Lancaster Sound and the Nares Strait - flow southward at the surface. The same top/bottom
 599 structure of ocean current is present in the Fram Strait, where Arctic Waters flow south-
 600 ward along the East Greenland coastline and Atlantic Waters flow northward at depth
 601 (West Spitsbergen current; results not shown). In the High resolution, Atlantic Waters
 602 penetrate far north into the Baffin Bay. The Medium resolution contrasts with the other
 603 two resolutions in the Baffin Bay, where very little Atlantic Water enters (Figures 8b and
 604 9b). Instead, Atlantic Waters flow cyclonically around the Labrador Sea along the con-
 605 tinental shelf (Figure 8 h).

606 The path of the Atlantic Waters and penetration of heat into the Baffin Bay is known
 607 to be influenced by the atmospheric forcing (D. Holland et al., 2008). In particular, the
 608 partitioning of OHT between the North Atlantic Drift and the Irminger Current (south
 609 of Iceland) is sensitive to the state of the NAO, with positive phase of the NAO favor-
 610 ing the eastern branch of the circulation, which is then associated with a reduced ice cover
 611 in the Greenland and Barents Seas (Myers et al., 2007; Strong et al., 2009; Straneo &
 612 Heimbach, 2013). In climate models, the NAO has been shown to influence Labrador Sea
 613 Water formation on decadal time scales, which in turn affects the subpolar gyre (Langehaug
 614 et al., 2012). During the spin up of our model (years 1 to 120), the mean state of the
 615 atmosphere changes to a more positive NAO state in the Low and High resolutions com-
 616 pared to the Medium resolution (not shown). This state persists throughout the CC sim-
 617 ulation (see Figure 10), and should promote deeper penetration of Atlantic Waters in
 618 the Fram Strait and Barents Sea Opening in the Low and High resolutions (Langehaug
 619 et al., 2012). Instead, we see more recirculation of Atlantic Waters in the Irminger Sea
 620 in the High resolution compared to the Medium resolution, indicating the NAO variabil-
 621 ity is not the leading factor in determining the current pathways in the Arctic.

622 The path of warmer Atlantic Waters into the Baffin Bay is also sensitive to spa-
 623 tial resolution in models, with high resolutions (up to $1/60^\circ$) favoring the Irminger branch
 624 (Pennelly & Myers, 2020). Although all three model configurations fall within the range
 625 of observations for OHT through the Davis Strait, the High resolution is the closest to
 626 the mean and has the largest interannual and decadal variability of the suite, yet still
 627 smaller than observations (Figure 9). Importantly, the OHT across the Davis Strait in
 628 the High resolution is the highest across the suite, about twice as large as the OHT in
 629 the Medium resolution, and four times that of the Low resolution (Figure 9a). The OVT
 630 in the Medium and High resolutions is close to observations, but is much weaker in ab-
 631 solute value in the Low resolution (Figure 9c). Very little poleward volume transport is
 632 found in the Medium and Low resolutions compared to the High resolution, and the south-
 633 wards volume transport averages in the Medium and High resolutions are of a similar
 634 order (Figure 9b). In the Low and Medium resolution configurations, the currents do not
 635 penetrate the Baffin Bay and continue along the western boundary towards the Labrador
 636 shelf (Figure 8 g-h), whereas in the High resolution configuration, the current penetrates
 637 into the Baffin Bay (Figure 8 si). The interannual variability of OVT in Davis Strait is
 638 significantly anti-correlated (at the 95% confidence level) with the sum of the transport
 639 through the Fram Strait and Barents Sea Opening, with correlation coefficients of -0.82,
 640 -0.96 and -0.82 in the Low, Medium and High resolutions, respectively. This suggests that
 641 the Irminger Branch dominates the variability in the Davis Strait as opposed to the East
 642 Greenland Current branch bringing polar surface waters southward. This anticorrela-
 643 tion also illustrates the partitioning of the transport of Atlantic Waters between the Arc-
 644 tic and the Labrador Sea and Davis Strait. Hence, in the Medium resolution, the weaker

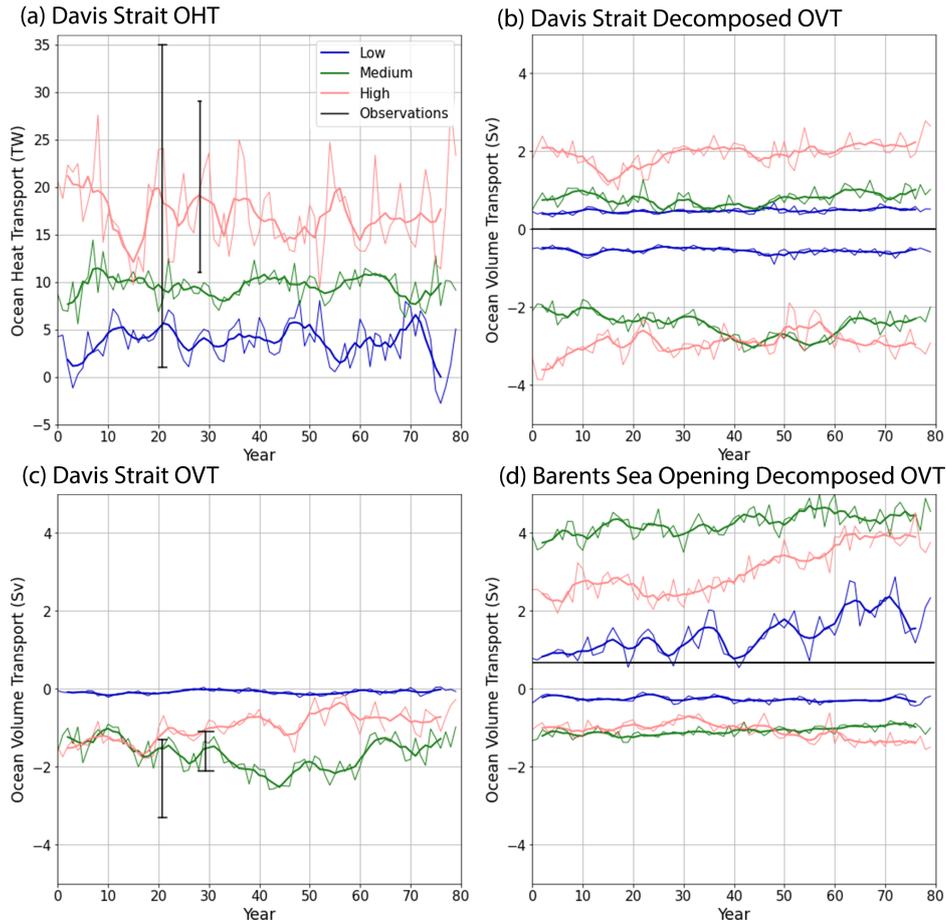


Figure 9. Timeseries of (a) total OHT and (c) total OVT across Davis Strait in the CC run for the Low, Medium and High resolutions. The total OVT is decomposed into its northward (positive) and southward (negative) component in (b). The same decomposition is made for the Barents Sea Opening OVT in (d). Thin lines correspond to annual averages and thick lines to five-year running mean. Observational estimates are indicated as vertical bars with the horizontal extent corresponding to the period of observations: 1987–1990 (Cuny et al., 2004) and 2004–2005 (Curry et al., 2011) for OHT, 1987–1990 (Cuny et al., 2004) and 2004–2010 (Curry et al., 2014) for OVT.

645 OVT into the Davis Strait is tied to the higher OHT into the Arctic through the Fram
 646 Strait and Barents Sea Opening. In particular, the poleward OVT through the Barents
 647 Sea Opening is twice as large in the Medium resolution as in the High resolution (see
 648 Figure 9d). Furthermore, OHT variability is strongly driven by OVT variability at the
 649 interannual and decadal scale, resulting in warmer waters in the Norwegian and Barents
 650 Seas. This suggests that the partitioning of OVT into the Arctic is a key component of
 651 the representation of the sea ice in the model suite.

652 In the model suite, the difference in the partitioning of Atlantic waters between the
 653 Irminger branch and Norwegian branch can be partly related to the difference in con-
 654 vection centers. In the Medium resolution, mixed layer in the Labrador Sea is slightly
 655 deeper but more localized than in the High resolution (Figure 8 j-l). The maximum win-

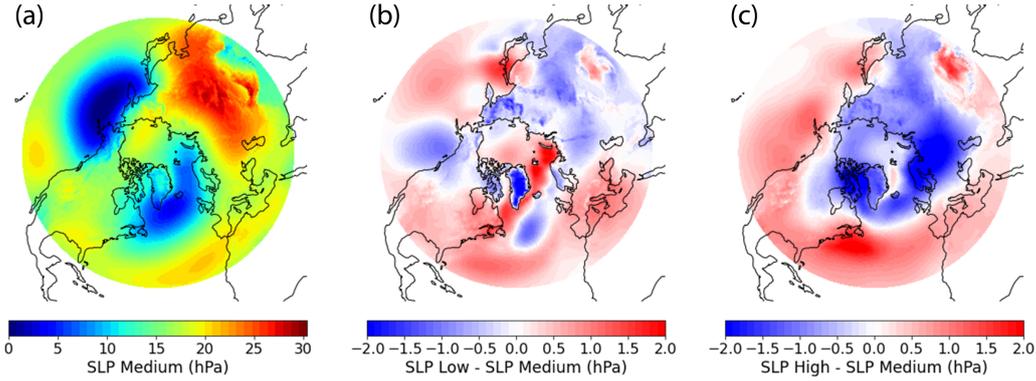


Figure 10. Winter SLP (JFM) in the Medium resolution (a), winter SLP difference between the Low and Medium resolutions (b) and winter SLP difference between the High and Medium resolutions (c) averaged over the CC simulation

656 ter mixed layer depth (MLD) in the Medium resolution is 1.7 km in the first decade of
 657 the CC run, around 300 m deeper than in the High resolution. The area of deep con-
 658 vection in the High resolution extends to the western boundary of the Labrador Sea, with
 659 MLD of around 1 km. Similarly in the HighResMIP models, $1/4^\circ$ models show deeper
 660 convection than 1° models, and overestimate MLD compared to observations (Koenigk
 661 et al., 2021). Pennelly and Myers (2020) also find that increasing the ocean resolution
 662 from $1/4^\circ$ to $1/12^\circ$ (and even $1/60^\circ$) leads to a shallower mixed layer thanks to more
 663 representation of eddy fluxes; however, they also find that the area of deep convection
 664 is less extensive. Conversely, in the Icelandic and Norwegian Seas, the depth and area
 665 of deep mixed layer are greater in the Medium resolution than in the High resolution.
 666 From Low to High resolutions, we see a south-westward transfer of deep convection re-
 667 gions from the Greenland-Icelandic-Norwegian (GIN) Seas towards the Labrador Sea (see
 668 Figure 8 j-l), in agreement with results from Jackson et al. (2020). These results are also
 669 in agreement with those of HadGEM3 and ECMWF, with a shift in convection centers
 670 towards the Labrador Sea in $1/4^\circ$ and $1/12^\circ$ model configurations, compared to the 1°
 671 model (Koenigk et al., 2021, Figure 1).

672 In the northern North Atlantic, where deep convection is present, we find a strong
 673 negative correlation at interannual scale between the OVT across the Barents Sea Open-
 674 ing and winter MLD in both the Medium and High resolutions (Figures 11 b-c and 8
 675 h-l), indicating that deep penetration of Atlantic Waters into the Barents Sea Opening
 676 is associated with weak convection in the GIN Seas. Similarly, the OVT across Davis Strait
 677 is negatively correlated with winter MLD in the Labrador Sea in all three model con-
 678 figurations (Figure 11), indicating that deep penetration of Atlantic Waters into Baffin
 679 Bay through the Davis Strait is associated with weak convection in the Labrador Sea.
 680 This negative correlation suggests that a weak subpolar gyre circulation is associated with
 681 strong deep convection and meridional circulation, in agreement with results from Drijfhout
 682 and Hazeleger (2006).

683 5 Conclusion

684 In this study, we investigated the impact of ocean heat transport on Arctic sea ice
 685 under climate change in the GFDL CM2-O model suite. The model suite only differs in
 686 the horizontal spatial resolution of the ocean component : from 1° (Low) to $1/4^\circ$ (Medium)
 687 to $1/10^\circ$ (High), with a mesoscale eddy parameterization for the Low resolution. We in-

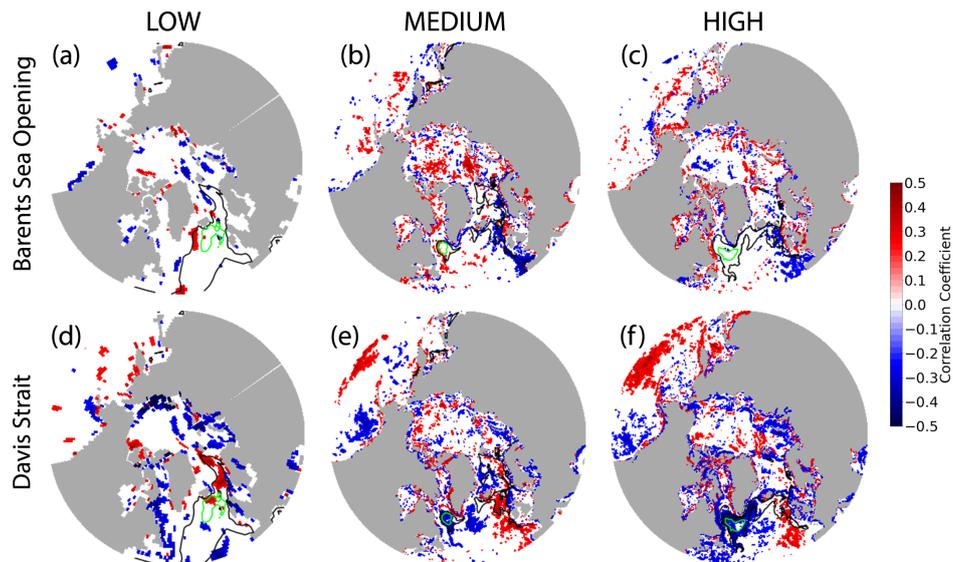


Figure 11. Correlation maps between the detrended annual (January-December) OVT in the Barents Sea Opening (a-c) and Davis Strait (d-f) and the detrended Winter MLD, in the Low (left), Medium (center) and High (right) resolutions in the CC experiments. Only correlations that are significant at the 95% level are shown. Black contours indicate the CC simulation average winter MLD at 300 m (black line) and 1 km (green line).

688 investigated the potential impact of resolution on the mean ocean and sea ice states, and
 689 the relationship between Arctic ocean heat transport and sea ice, on the Pan-Arctic and
 690 regional scale, at annual and decadal time scales. We found that :

- 691 • Models with a higher total ocean heat transport into the Arctic have a smaller sea
 692 ice extent in all seasons, in agreement with previous studies (Hewitt et al., 2016;
 693 M. J. Roberts et al., 2016; Docquier et al., 2019).
- 694 • Decadal variability in ocean heat transport explains a large fraction of decadal vari-
 695 ability in sea ice extent.
- 696 • At interannual time scale, the impact of ocean heat transport on sea ice extent
 697 is limited to the shelf regions.
- 698 • The SIE in the Medium resolution model configuration is in best agreement with
 699 the observational record at the beginning of the satellite era.
- 700 • In the CM2-O model, the refining of spatial resolution does not induce a system-
 701 atic increase in OHT, as opposed to other model suites that show a monotonous
 702 decrease in sea ice extent with increasing ocean heat transport (Hewitt et al., 2016;
 703 M. J. Roberts et al., 2016).
- 704 • The shift from non-eddying to eddy-permitting resolutions tends to improve the
 705 representation of currents and heat transport, particularly in the North Atlantic,
 706 in agreement with other studies (Docquier et al., 2019; M. J. Roberts et al., 2016).
- 707 • Though the Medium resolution has a higher ocean heat transport and lower sea
 708 ice extent when compared with those of the High resolution in the pre-industrial
 709 mean state, the trends in sea ice loss and ocean heat transport in the two model
 710 configurations under increasing CO₂ forcing are similar.

- 711 • The Low and High resolutions have the same pre-industrial sea ice extent and thick-
712 ness distribution, but very different response in sea ice extent to CO₂ forcing, with
713 the High resolution being more sensitive than its coarser resolution counterpart.
- 714 • As the spatial resolution of the model increases from medium to high, greater heat
715 transport is found into Davis Strait at the expense of the Atlantic-Arctic gates sug-
716 gesting that the Irminger branch is favored over the Faroe-Scotland branch. The
717 differences in deep convection between these two model configurations might partly
718 explain the difference in heat partitioning.

719 While the increase in OHT and shift of convection centers from east to west of the
720 basin as resolution increases are robust findings across different climate and ocean model
721 families, the lack of sensitivity of sea ice to OHT in Low (eddy-parameterized) is pre-
722 sumably due to the very low OHT in this model. A more complete analysis of different
723 GCMs with eddy-rich ocean component would be required to determine whether the in-
724 crease in OHT going from eddy-permitting (Medium) to eddy-rich (High) is a robust fea-
725 ture.

726 In the GFDL CM2-O suite, the poleward heat transport does increase with reso-
727 lution, with stronger narrower currents, until about 50°N. However, the partitioning of
728 the currents in the high latitudes greatly impacts the penetration of heat into the Arct-
729 ic and in turn the projections of Arctic sea ice. This highlights the need for a realistic
730 representation of said partitioning on top of that of temperature and current strength.
731 The Overturning in the Subpolar North Atlantic Program (OSNAP) provides contin-
732 uous records of mass and heat transports in the eastern and western subpolar regions
733 against which models' partitioning can be assessed.

734 Acknowledgments

735 M.D. acknowledges support from the Fonds de recherche du Québec Nature et tech-
736 nologies (FRQNT) through the Bourse de maîtrise en recherche. M.D. and C.O.D. ac-
737 knowledge the support of the Natural Sciences and Engineering Research Council of Canada
738 (NSERC) - Accelerator program. M.D. is also grateful for the support of McGill Uni-
739 versity, Québec-Océan and Arctrain Canada. This work is a contribution to the NSERC
740 - Discovery Program, Marine Environmental Observation, Prediction and Response (MEOPAR),
741 NASA grant #80NSSC20K1259, NSF – Office of Polar Program grants #1504023, #1603350
742 and #1928126, awarded to BT.

743 All output variables from the the NSIDC are available at <https://nsidc.org/data/G02135/versions/3>
744 for monthly sea ice extent and <https://nsidc.org/data/G10010> for the SIBT1850 (Sea
745 Ice Back To 1850) data. Output of the GFDL CM2-O suite that were used to make the
746 figures of the paper will be available from the Polar Data Catalogue (<https://www.polardata.ca/>)
747 by acceptance.

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