

High-resolution model simulation of the interannual and seasonal variability of the Weddell Gyre during 1958 - 2018

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Abstract

The Weddell Gyre's variability on seasonal and interannual timescales is investigated using an ocean-sea ice model at three different horizontal resolutions. The model is evaluated against available observations to demonstrate that the highest resolution configuration (0.1° in the horizontal) best reproduces observed features of the region. The simulations suggest that the gyre is subject to large variability in its circulation that is not captured by summer-biased or short-term observations. The Weddell Gyre's seasonal cycle consists of a summer minimum and a winter maximum and accounts for changes that are between one third and a half of its mean transport. On interannual time scales we find that the gyre's strength is correlated with the local Antarctic easterlies and that extreme events of gyre circulation are associated with changes in sea ice concentration and the characteristics of warm inflow at the eastern boundary.

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3 **1958 - 2018**

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10 **Key Points:**

- 11 • The Weddell Gyre, as simulated by a global coupled ocean-sea ice model, displays
12 large variability on seasonal and interannual time scales.
- 13 • There is evidence of low frequency (decadal) variability in the Weddell Gyre's strength,
14 but no significant trends during study period.
- 15 • Years of extreme strong/weak Weddell Gyre flow are shown to be linked to anoma-
16 lous winds, ocean temperatures and sea ice concentration.

Abstract

The Weddell Gyre’s variability on seasonal and interannual timescales is investigated using an ocean-sea ice model at three different horizontal resolutions. The model is evaluated against available observations to demonstrate that the highest resolution configuration (0.1° in the horizontal) best reproduces observed features of the region. The simulations suggest that the gyre is subject to large variability in its circulation that is not captured by summer-biased or short-term observations. The Weddell Gyre’s seasonal cycle consists of a summer minimum and a winter maximum and accounts for changes that are between one third and a half of its mean transport. On interannual time scales we find that the gyre’s strength is correlated with the local Antarctic easterlies and that extreme events of gyre circulation are associated with changes in sea ice concentration and the characteristics of warm inflow at the eastern boundary.

Plain Language Summary

The Weddell Gyre, located east of the Antarctic Peninsula, is one of the largest features of the ocean circulation of the Southern Hemisphere. It is adjacent to an important site of bottom water formation, a process that sequesters carbon and heat from the atmosphere and sets the density of the deep ocean, therefore making the region important for global climate. However, extensive sea ice cover throughout the year has historically prevented continuous observations. Several unique features of the gyre, such as open boundaries and intense surface buoyancy fluxes, make the identification of its forcing mechanisms difficult. A deeper understanding of the dynamics in this remote region will shed light on the role of the gyre in our present climate, and help us understand its potential evolution with climate change. We use a high resolution numerical model which shows that the Weddell Gyre undergoes large seasonal and interannual changes. We find that the gyre spins up during winter and slows down during summer, and that strong/weak events in our model simulation are correlated with the strength of the regional easterly winds close to the Antarctic continent. These strong/weak events affect sea ice cover, water mass characteristics and bottom water production.

1 Introduction

The lateral circulation in the Weddell Sea is dominated by the Weddell Gyre, one of the largest features of the ocean circulation south of the Antarctic Circumpolar Current (ACC). The Weddell Gyre is characterized by a broad cyclonic circulation spanning from the Antarctic Peninsula until approximately 30°E (Deacon, 1979), schematized in Figure 1. The gyre is one of the southernmost open ocean reaches in the world, with several features that make it an intriguing and relevant component of the Southern Ocean circulation. A strong interaction between ocean and sea ice favours large surface buoyancy fluxes that, in combination with a weak stratification, create a connection between the atmosphere, the ocean surface and the ocean bottom. For example, located in the southwestern Weddell Sea is one of the major formation sites of Antarctic Bottom Water (AABW) (Meredith, 2013), a process which supports the deepest limb of the global overturning circulation and involves waters that are circulated and transformed within the gyre. Despite the Weddell Gyre’s relevance to global climate, the extensive sea ice cover throughout most part of the year has historically hampered long-term, continuous observational efforts and the present knowledge of the gyre’s circulation and variability is mostly limited to the summer months.

The Weddell Gyre’s circulation has been traditionally associated with the negative stress curl given by a large scale surface wind field consisting of westerlies to the north, a circumpolar band of easterly winds surrounding Antarctica and a low pressure system embedded at around 30°E , 65°S (Gordon et al., 1981; Deacon, 1979). The surface wind pattern makes the center of the gyre a region of divergence of Ekman transport, char-

acterized by a depression of sea level (Armitage et al., 2018), wind driven Ekman upwelling and a doming of subsurface isopycnals (Klatt et al., 2005; Schröder & Fahrbach, 1999). The Antarctic continent to the south and west, together with a series of submarine ridges to the north provide a clear topographic barrier to the flow. There is no such barrier on the eastern boundary of the gyre, which has been suggested to be a dynamic feature located between 30°E and 70°E. It is at this boundary where relatively warm and saline Circumpolar Deep Water (CDW) enters the gyre, partly following an advective and partly following an eddy-driven pathway (Ryan et al., 2016; Cisewski et al., 2011; Leach et al., 2011; Schröder & Fahrbach, 1999), hereafter called Warm Deep Water (WDW). The properties of WDW that enter the gyre are modified along its westward path by mixing and upwelling until some of this modified Warm Deep Water (mWDW) crosses the shelf break at the southern Weddell Sea, where it mixes with High Salinity Shelf Water (HSSW), formed by cooling and brine rejection, to produce Weddell Sea Deep Water (WSDW). Some of this WSDW is able to escape the gyre towards the Scotia Sea and becomes AABW (R. A. Locarnini et al., 1993).

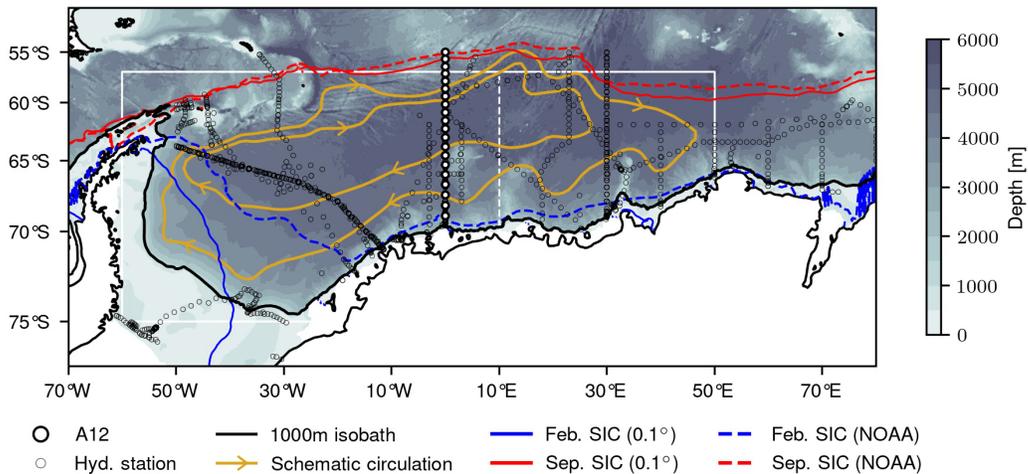


Figure 1. Bathymetry of the Weddell Sea with the 1000m isobath (black contour) and February (blue) and September (red) sea ice extent defined by 15% sea ice concentration from ACCESS-OM2-01 (solid line) and NOAA/NSIDC Climate Data Record of Passive Microwave Sea Ice Concentration, Version 3 (dashed line). The arrows show a schematic of the Weddell Gyre circulation and the white solid box marks the area within which surface stress and buoyancy fluxes are averaged to build a time series in Figures 8 and 10. The Weddell Gyre strength is calculated west of the dashed orange line at 10°E.

Observational and modelling studies have estimated the Weddell Gyre strength to be between 30 and 100 Sv, a wide range of values that reflect the strong dependence of gyre strength on the methodology used. Hydrography-based estimates may underestimate gyre strength because they are biased towards summer conditions, thus excluding a possible winter intensification. Furthermore, the weak stratification that characterizes the region means that the assumption of a level of no motion, required to derive transports from the thermal wind equations, may not be appropriate (Cisewski et al., 2011; Park & Gambéroni, 1995; Fahrbach et al., 1991). Numerical models are not subject to such experimental difficulties but their evaluation is limited by the scarcity of observations in the region. The gyre exhibits a persistent double-cell circulation structure whose origin remains unclear, but further contributes to the uncertainties in gyre transport (Reeve

93 et al., 2019; Mazloff et al., 2010; Wang & Meredith, 2008; Beckmann et al., 1999; Orsi
94 et al., 1990).

95 The present knowledge of the variability of the Weddell Gyre, its magnitude, drivers
96 and time scales is greatly limited by the lack of continuous observations. Repeat hydro-
97 graphic sections have been used to determine that the properties of the Weddell Gyre’s
98 water masses display significant variations associated with the gyre’s circulation (Fahrbach
99 et al., 2011; Behrendt et al., 2011; Fahrbach et al., 2004). For example, changes in the
100 properties of AABW exported from the Weddell Sea have been linked to wind-forced changes
101 in the gyre’s baroclinic structure (Jullion et al., 2010; Meredith et al., 2011, 2008). The
102 northern and southern limbs of the gyre have been suggested to vary independently and
103 not necessarily in phase, responding to different forcing mechanisms (Fahrbach et al.,
104 2011). The northern boundary could be driven by the westerly winds, while the south-
105 ern limb of the gyre, comprised of a westward flowing current system referred to as the
106 Antarctic Slope Current (ASC), partly forced by the easterly winds close to the Antarc-
107 tic continent. Easterlies force an Ekman transport towards the coast, raising sea level
108 and driving a geostrophic flow that displays significant variability on seasonal and in-
109 terannual time scales (A. Naveira Garabato et al., 2019; Armitage et al., 2018; Math-
110 iot et al., 2011). Interannual variability in gyre strength has also been related to climate
111 modes, such as the Southern Annular Mode (SAM) or El Nio oscillation (Armitage et
112 al., 2018; Martinson & Iannuzzi, 2003). Few studies take into consideration the influence
113 of sea ice and surface buoyancy fluxes as possible driving mechanisms, despite their rel-
114 evance to the dynamics of the region.

115 This study analyzes numerical simulations derived from a coupled numerical model
116 (described in Section 2.1) configured over three different horizontal resolutions (1° , 0.25°
117 and 0.1°) to study the Weddell Gyre’s variability on seasonal to interannual timescales,
118 and its connection to possible forcing mechanisms including wind stress and buoyancy
119 fluxes. In Section 3 the model is evaluated against available observations to assess its
120 strengths and weaknesses in the study region. For this evaluation we use the Armitage
121 et al. (2018) dynamic ocean topography product and temperature and salinity profiles
122 from available hydrographic stations in the region. Section 4 characterizes the Weddell
123 Gyre mean state in model simulations and Section 5 and 6 the gyre’s variability on sea-
124 sonal and interannual timescales respectively. Section 7 summarizes the results and dis-
125 cusses their relevance for our current knowledge of the Weddell Gyre circulation.

126 1.1 Ocean/Sea Ice Model

127 This study uses the Australian Community Climate and Earth System Simulator
128 (ACCESS-OM2) (Kiss et al., 2020), a global ocean-sea ice model. This model is avail-
129 able at three different horizontal resolutions, namely 1° and 0.25° with 50 vertical lev-
130 els and 0.1° with 75 vertical levels. The ocean component of the model is the Modular
131 Ocean Model (MOM) version 5.1, developed by the Geophysical Fluid Dynamics Lab-
132 oratory (<https://mom-ocean.github.io/>) and the sea ice component is the Los Alamos
133 sea ice model (CICE) version 5.1.2 from Los Alamos National Laboratories ([https://
134 github.com/CICE-Consortium/CICE-svn-trunk/tree/cice-5.1.2](https://github.com/CICE-Consortium/CICE-svn-trunk/tree/cice-5.1.2)). All three config-
135 urations are initialized from a rest state with zero sea level, temperature and salinity given
136 by the World Ocean Atlas 2013 v2 0.25° decav product (M. Locarnini et al., 2018; Zweng
137 et al., 2019) and are forced by a prescribed atmosphere from JRA55-do reanalysis v1.4
138 for the period 1958 to 2018 (Tsujino et al., 2018). The model is continuously cycled through
139 this 61-year period, from which we select the third forcing cycle since it is the latest cy-
140 cle available for all three resolutions. The three model configurations are consistent with
141 each other which allows to infer the importance of resolution in the study region. There
142 is a general improvement at the highest 0.1° configuration when solving several features
143 of the ocean: those relevant for this study are the representation Southern Ocean wa-

ter masses, the overturning circulation and the characteristics of the circulation on the Antarctic continental shelf (Kiss et al., 2020; Moorman et al., 2020; Morrison et al., 2020).

1.2 Satellite Observations

We evaluate ACCESS-OM2 in the Weddell Gyre region against the dynamic topography of the satellite product developed by Armitage et al. (2018), a unique data set in the Southern Ocean where extensive sea ice cover has historically hampered long term continuous observations using traditional radar altimetry. The product consists of monthly composites of dynamic ocean topography spanning a five year period (2011 to 2016) at 50km horizontal resolution, referenced to the GOCO05c combined gravity field model (Fecher et al., 2017). A limitation of this satellite product is the dependence of its mean dynamic topography on the geoid model, which has larger errors and uncertainties toward the Antarctic continent due to the sparsity of the data being assimilated. These uncertainties do not affect sea level anomalies, i.e. observed variability, because the geoid is time invariant. The product also displays some north-south striping artifacts generated by the satellite’s orbital precession (see Fig. 5 of Armitage et al. (2018)). To compare ACCESS-OM2 sea level with satellite observations, we interpolated the model to the satellite product’s grid and applied an offset to remove the spatially uniform footprint of geoid uncertainties in the satellite product. The offset is defined as the average difference between the mean sea level fields of each model simulation and observations for the entire region covered by the satellite product.

1.3 Hydrographic Data

The model’s temperature and salinity is evaluated against hydrographic observations in the Weddell Gyre region downloaded from the Clivar and Carbon Hydrographic Data Office (CCHDO). We selected all publicly available CTD profiles in the region, a total number of 1576. Table 1 has information on the hydrographic cruises and their locations are shown in Figure 1. To compare against ACCESS-OM2, we select the nearest grid point and corresponding monthly composite within the model output and we interpolate it to the same vertical levels. With this approach we obtain synthetic profiles within the ACCESS-OM2 model simulations to evaluate the temperature-salinity structure of the gyre. To give further insight into the gyres vertical structure, we selected one of the most repeated WOCE transects in the region, A12 nominally at the Greenwich Meridian (see Figure 1), to calculate an average potential density cross-section. Analogous to the procedure for individual hydrographic stations, we select the monthly composites of model data that correspond to the repeat A12 cruises, interpolate the vertical level, and then calculate the anomalies of the model with respect to the observations.

1.4 Barotropic streamfunction definition

The barotropic streamfunction, ψ , is used to study the gyre’s transport and is defined as the meridional integral of the depth-integrated zonal mass transport M_x :

$$\psi(x, y) = \int_{y_0}^y \frac{M_x(x, y')}{\rho(y - y_0)} dy' \quad (1)$$

where we take ρ as an average density of 1035 kg m^{-3} and the integration goes from south to north starting at the Antarctic continent (y_0). By this definition, ψ takes negative values for the cyclonic circulation of the Weddell Gyre. The gyre strength (GS) is derived from ψ by calculating its minimum in the region bounded by the white box in Figure 1 west of 10°E (dashed white line) and taking the absolute value:

Table 1. Summary of the hydrographic cruises used to evaluate the model. Cruise locations are indicated in Figure 1. The expocode is the identifier for the data sets archived at the Carbon and Climate Hydrographic Data Office (CCHDO; <http://cchdo.ucsd.edu>).

Expocode	Line	Start date	End date	Principal Investigator
06AQANTVIII_2	SR02, SR04	1989-09	1989-10	Eberhard Fahrbach
06AQANTIX_2	SR04	1990-11	1990-12	Eberhard Fahrbach
06AQANTX_4	A12, SR04	1992-05	1992-08	Peter Lemke
35MFCIVA_1	I06S	1993-01	1993-03	Alain Poisson
74DI200_1	S04	1993-02	1993-03	Robert R. Dickson
74JC10_1	A23	1995-03	1995-05	Karen J. Heywood and Brian A. King
320696_3	S04, S04I	1996-03	1996-04	Thomas Whitworth
35MF103_1	I06S	1996-02	1996-03	Alain Poisson
06AQANTXIII_4	S04A, SR04	1996-03	1996-05	Eberhard Fahrbach
06AQANTXV_4	SR04	1998-03	1998-05	Eberhard Fahrbach
06AQ199901_2	A12	1999-01	1999-03	Eberhard Fahrbach
06AQ200012_3	A12	2000-12	2001-01	Eberhard Fahrbach
06AQ200211_2	A12	2002-11	2003-01	D.K. Ffterer
06AQ20050102	A12	2005-01	2005-04	Eberhard Fahrbach
09AR20060102	S04I	2006-01	2006-03	Mark Rosenberg
06AQ20071128	A12	2007-11	2008-02	Ulrich Bathmann
33RR20080204	I06S	2008-02	2008-03	Kevin Speer
06AQ20080210	A12	2008-02	2008-04	Eberhard Fahrbach
06AQ20101128	A12	2010-11	2011-02	Eberhard Fahrbach
06AQ20141202	PS89	2014-12	2015-02	Olaf Boebel

$$GS = |\min\{\psi\}| \quad (2)$$

187 The gyre strength definition is limited to 10°E to exclude the unstable, eddy-rich flow
 188 of the eastern region, but nonetheless encompasses the main structure of the mean gyre
 189 (see Section 3.2). The barotropic streamfunction is also used to define the boundary of
 190 the Weddell Gyre as the 12 Sv contour with the purpose of tracking changes in the gyre’s
 191 area over time.

192 1.5 Surface stress and buoyancy fluxes

193 To assess the role of stress over the ocean’s surface as a possible driver of the cir-
 194 culation, we calculate the total surface stress from model output taking into account the
 195 relative contributions of air/ocean and ice/ocean stresses ($\tau_{air/ocean}$ and $\tau_{ice/ocean}$ re-
 196 spectively) weighted by sea ice concentration. The curl of the total surface stress, $\nabla \times$
 197 τ , is defined as:

$$\nabla \times \tau = \nabla \times ((1 - \alpha_{ice})\tau_{air/ocean}) + \nabla \times (\alpha_{ice}\tau_{ice/ocean}) \quad (3)$$

198 where α_{ice} is sea ice concentration. To build a time series, $\nabla \times \tau$ is averaged in the re-
 199 gion bounded by the white box in Figure 1 excluding grid cells in which the bathymetry
 200 is shallower than 1000m. We show the absolute value of the surface stress curl so that
 201 stronger curl anomalies are represented by positive values.

202 Surface buoyancy fluxes are also considered as a possible driver of the gyre’s vari-
 203 ability, taking into account contributions due to both heat and freshwater exchanges. The
 204 surface buoyancy flux is defined as:

$$\mathcal{B} = \frac{g\alpha Q}{c_w\rho} + g\beta F_w s \quad (4)$$

205 where ρ is surface density, g is the acceleration due to gravity, c_w the specific heat of sea
 206 water, α the thermal expansion coefficient, β the saline contraction coefficient, s surface
 207 salinity and Q and F_w are the heat and freshwater fluxes, positive representing a buoy-
 208 ancy gain by the ocean. The Python implementation of the Gibbs Sea Water Oceanog-
 209 raphic Toolbox of TEOS-10 (<https://teos-10.github.io/GSW-Python/>) was used
 210 to compute the quantities derived from temperature and salinity. The largest buoyancy
 211 fluxes in the Weddell Sea region are found over the continental shelf, so \mathcal{B} is averaged
 212 at depths shallower than 1000m to build a time series.

213 1.6 Climate indices

214 We use the Southern Annular Mode (SAM) index as a proxy for the meridional ex-
 215 pansion/contraction of the band of westerlies surrounding the Antarctica to evaluate the
 216 connection between the gyre and large-scale atmospheric modes of variability. This in-
 217 dex is calculated as the difference between the zonal anomalies of sea level pressure at
 218 40°S and 65°S, (Marshall, 2003), from the JRA55 reanalysis, the atmospheric data set
 219 used to force the model. To verify that local changes in the westerlies are not significantly
 220 different from the circumpolar average we re-evaluate the SAM index in a limited do-
 221 main, calculating the zonal anomalies with respect to the mean of the Weddell Sea and
 222 find no significant differences. By analogy we define a different index to track the inten-
 223 sification/weakening of the local easterly winds, hereafter denoted EAS, defined as the
 224 difference in the zonal anomalies of sea level pressure at 65°S and 72°S. Contrary to the
 225 SAM index, since the easterlies display a large zonal variation, the EAS is calculated be-
 226 tween 30°W and 70°E with the intention of focusing only on regional changes.

227 1.7 Analysis period

228 Long term climatological fields are calculated for the period 1958 to 2018 for ψ , $\nabla \times$
 229 τ and \mathcal{B} , from which we calculate the annual cycles of gyre strength, surface stress curl
 230 and buoyancy fluxes as described in the previous sections. To study interannual variabil-
 231 ity, these metrics together with the SAM and EAS indices are calculated using monthly
 232 composite fields for the full 61-year period, and subsequently the annual cycle is removed
 233 to obtain anomalies. We also apply a 12-month running mean to filter out high frequency
 234 variability. The 0.1° experiment is used to define strong and weak events of Weddell Gyre
 235 flow as periods longer than 6 months during which gyre strength is beyond $\pm 0.8SD$ of
 236 the gyre strength time series smoothed with a 10-year running mean. Once the events
 237 are identified we composite the streamfunction, sea level pressure, surface buoyancy fluxes,
 238 subsurface temperature maximum and sea ice concentration anomaly fields to charac-
 239 terise the differences between anomalously strong and weak gyre periods and the impact
 240 on the region. Additionally, we calculate the linear correlation coefficients between the
 241 gyre strength time series and the anomaly fields of surface buoyancy fluxes, subsurface
 242 temperature maximum and sea ice concentration in order to identify regions where they
 243 are significantly correlated to, and thus likely affected by, gyre strength.

244 2 Model Evaluation

245 Satellite observations (Armitage et al., 2018) show the Weddell Gyre as an elon-
 246 gated depression of sea level, indicative of a cyclonic geostrophic current, with its main
 247 axis oriented in the northeast-southwest direction (Figure 2a). Figures 2b, c and d show
 248 the difference between the model and satellite observed mean dynamic topography for
 249 the 1°, 0.25° and 0.1° configurations respectively. If this difference were spatially con-
 250 stant, then the derived surface geostrophic circulation would be the same since the ve-

251 locality in between two grid points is proportional to the difference in elevation between
 252 those two points. The 0.1° simulation exhibits the smallest horizontal gradients in the
 253 sea level difference with observations (Figure 2d), not only within the gyre’s boundary,
 254 but also in the southwestern continental shelf, the eastern region and within the ACC,
 255 which indicates that the highest resolution configuration better represents the observed
 256 Weddell Gyre geostrophic flow.

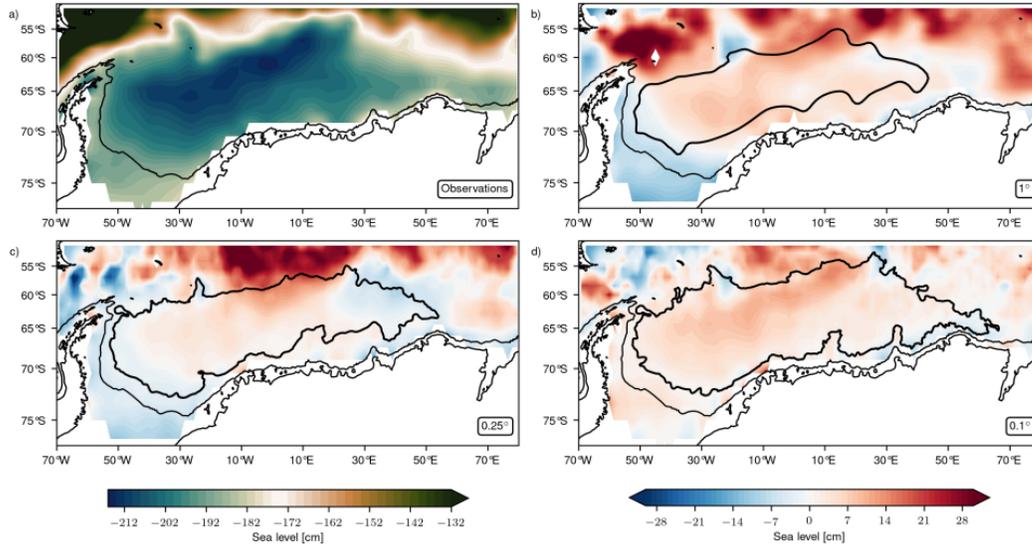


Figure 2. (a) Mean sea level from satellite observations for the period 2011 to 2016 (Armitage et al., 2018). Model’s mean sea level for the same period minus satellite observations for the (b) 1° , (c) 0.25° and (d) 0.1° resolution configurations, with gyre’s mean boundary (defined as the 12 Sv barotropic streamfunction contour) in thick black. Thin black contour in all panels marks the 1000m isobath.

257 The Armitage et al. (2018) data set can also be used to compare variability. We
 258 compute the correlation coefficient between the satellite and model sea level at every grid
 259 point for the period 2011 to 2016 and use them to build correlation maps, shown in Fig-
 260 ures 3a, b and c with significant correlations hatched. A north-south striping pattern
 261 is visible in the correlation coefficients that reflects an artifact from the satellite’s orbit.
 262 Taking this artifact into consideration, we infer that within the gyre, model and obser-
 263 vations are significantly correlated, with no large differences between model configura-
 264 tions. However, the correspondence between model and observations breaks down within
 265 the ACC, which we attribute to the presence of eddies that the model is not expected
 266 to reproduce at the precise location and time as they appear in observations. Unlike the
 267 comparison between model and observations within the gyre, in the southwestern contin-
 268 tental shelf the ability of the model to reproduce observed variability improves signif-
 269 icantly with resolution, with higher, significant correlations for the 0.1° configuration.
 270 However, the correlation maps do not fully illustrate differences in the model’s variabil-
 271 ity with respect to observations, which is why we calculate the root mean square error
 272 (RMSE) within the gyre over the observational record, Figure 3d. We observe that the
 273 RMSE within the Weddell Gyre is in phase between model configurations and of simi-
 274 lar magnitude, indicating that the departure from observations is consistent across res-
 275 olution and therefore likely due to errors in model forcing compared to observations. There
 276 is also no seasonality in the error (e.g. the error is not consistently larger in winter), which

277 further suggests that the error is due to biases in the model's forcing rather than biases
 278 in satellite observations associated with sea ice cover.

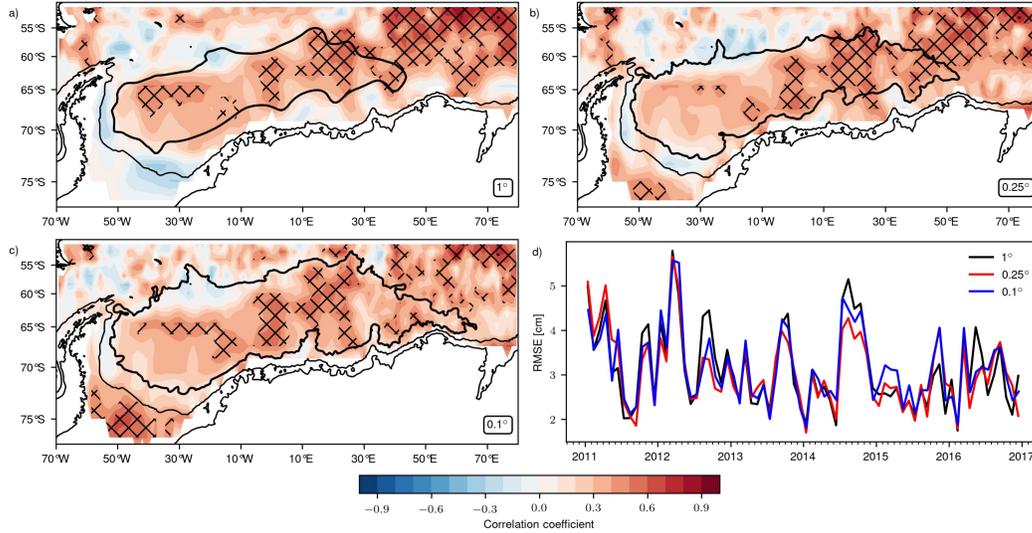


Figure 3. Correlation coefficients with satellite observations for models at (a) 1° , (b) 0.25° and (c) 0.1° resolution with hatching for significant correlations ($p < 0.05$). Mean gyre's boundary in thick black contour and 1000m isobath in thin black contour. (d) RMSE between model and observations calculated for a region encompassing the gyre $[30^\circ\text{W}, 30^\circ\text{E}] \times [70^\circ\text{S}, 60^\circ\text{S}]$.

279 Apart from evaluating the model's sea level, we can evaluate the model's temper-
 280 ature and salinity structure against observations. Figure 4 shows modeled and observed
 281 temperature-salinity diagrams comparing all available hydrographic stations in the Wed-
 282 dell Gyre region at locations indicated in Figure 1. It is possible to identify different wa-
 283 ter masses in observations (Figure 4a) according to their potential temperature and salin-
 284 ity: namely warm and saline CDW, colder WDW and a distinct tail of colder Weddell
 285 Sea Deep and Bottom Waters (WSDW and WSBW respectively). There are also some
 286 stations on the continental shelf that indicate the presence of Ice Shelf Water (ISW) with
 287 temperatures lower than the surface freezing point (-1.9°C). The ability of the model
 288 to reproduce the distinct characteristics of these water masses is highly dependent on
 289 resolution. The coarser 1° resolution does not show the presence of CDW in the region,
 290 has a colder, fresher than observed WDW and the characteristics of its bottom waters
 291 are less distinct from shallower waters than observed. There is a slight improvement in
 292 the 0.25° configuration which shows a trace of CDW and a better reproduction of WDW
 293 characteristics. The improvement is clearest for the 0.1° resolution configuration which
 294 captures the presence of CDW in the region, although warmer than observed, as well as
 295 achieving a better representation of WDW than the two coarser resolutions. The 0.1°
 296 simulation also shows some trace of the bottom waters, albeit less distinct than observed,
 297 which could be related to the coarse vertical resolution of the model relative to the sta-
 298 tion observations (50 levels for the 1° and 0.25° and 75 levels for the 0.1° configuration).
 299 Moreover, none of the model configurations capture the presence of supercold ISW be-
 300 cause the model lacks the ice shelf cavities where this water mass is formed. Since ISW
 301 participates in deep and bottom water production, the model's misrepresentation of these
 302 dense waters could be related to its inability to generate ISW.

303 The average cross section of potential density at the Greenwich Meridian, Figure
 304 5, shows the doming of isopycnals that characterizes the center of the gyre and the in-

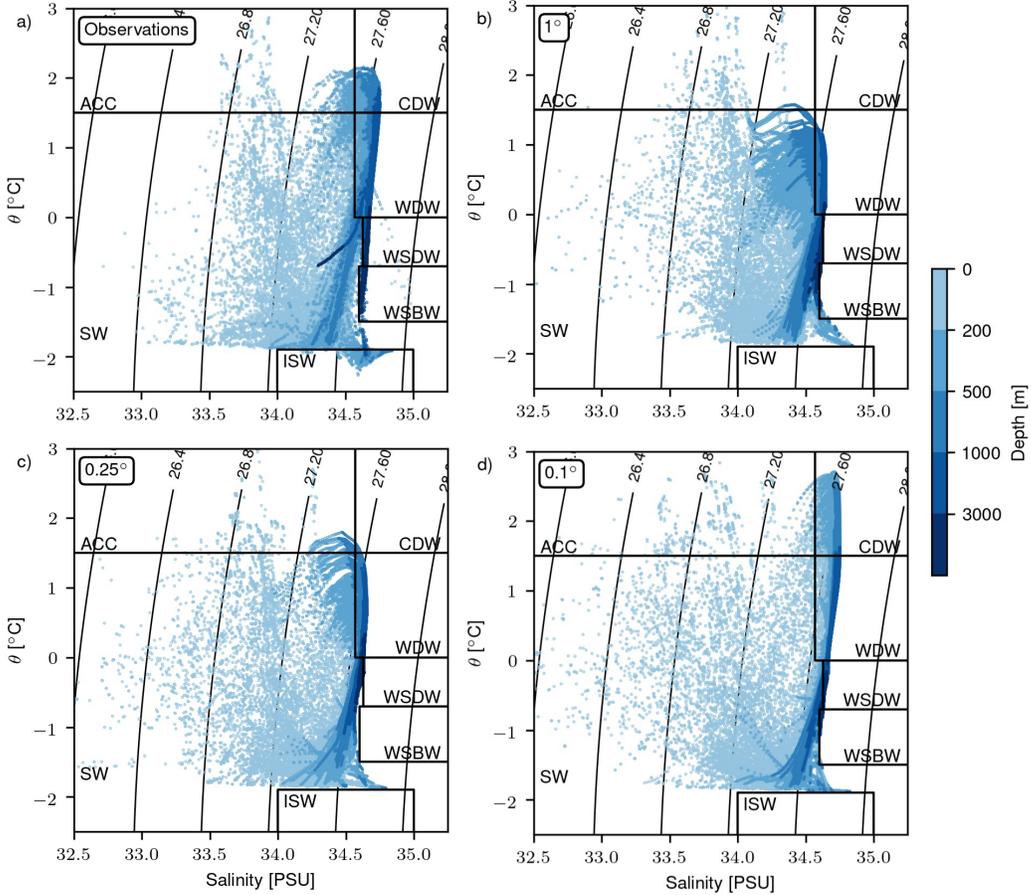


Figure 4. Temperature-salinity diagrams for all stations in Figure 1 for (a) hydrographic observations and for model runs at (b) 1° , (c) 0.25° and (d) 0.1° resolution, color coded by their depth. Antarctic Circumpolar Current = ACC, Circumpolar Deep Water = CDW, Warm Deep Water = WDW, Weddell Sea Deep Water = WSDW, Weddell Sea Bottom Water = WSBW, Surface Water = SW, Ice Shelf Water = ISW.

305 fluence of Maud Rise, north of 67°S , on top of which sits a cold, fresh Taylor column.
 306 At around 58°S the sloping of the isopycnals shows the front that marks the transition
 307 from Weddell Gyre to the ACC. The differences between model and observations (Fig-
 308 ures 5b, c and d) show that the largest discrepancies are located in the upper 250m of
 309 the water column. The most striking difference is the clear contrast between Weddell Gyre
 310 waters (denser than observed) and ACC waters (lighter than observed) for the 1° and
 311 0.25° configurations, which means that for these resolutions the oceanic front at the north-
 312 ern boundary of the gyre is steeper than observed. This Weddell Gyre/ACC contrast
 313 is not apparent in the 0.1° resolution. On the other hand, all three configurations show
 314 anomalies of alternate sign within the gyre in a thin surface layer shallower than 50m
 315 and the underlying subsurface layers. Since the cross sections are representative of sum-
 316 mer conditions, there is a shallow layer of light waters associated with sea ice melt that
 317 sits on top of the Winter Water (WW), so-called because it surfaces during winter when
 318 there is a weakening of the summer pycnocline associated with ice production. The model
 319 cross sections show for all resolutions a lighter than observed summer surface layer and
 320 a denser than observed winter subsurface layer, which means that the stratification is

321 larger than observed in the model during the summer months. Below the upper 250m,
 322 the 0.1° configuration exhibits an improved representation of the observed potential density
 323 relative to the 1° and 0.25° simulations.

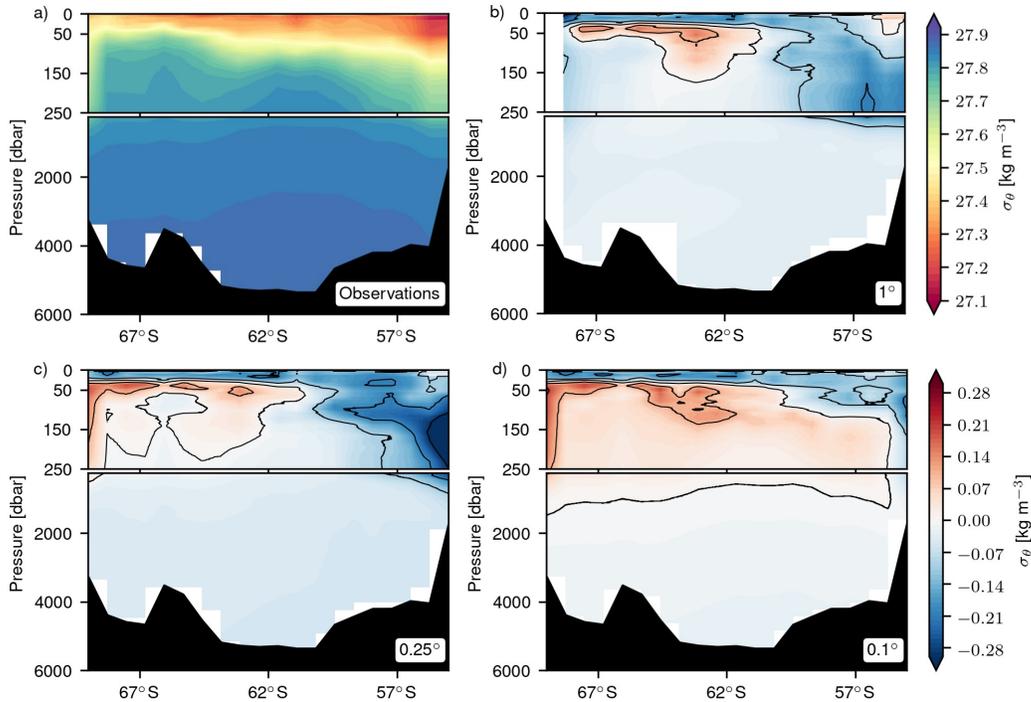


Figure 5. Average potential density (kg m^{-3}) section for A12 transect shown in Figure 1 for (a) hydrographic observations and the difference with observations for (b) 1° , (c) 0.25° and (d) 0.1° resolutions with 0.1 kg m^{-3} spacing between contours in black.

324 **3 Weddell Gyre Mean State**

325 The mean barotropic streamfunction for the entire 61 year period of the model run
 326 (1958 to 2018) in Figures 6a, c, e, shows an elongated cyclonic gyre with mean strengths
 327 of 34, 33 and 41 Sv and a seasonal range of 9, 14 and 15 Sv for the 1° , 0.25° and 0.1°
 328 resolutions respectively. The magnitude and spatial pattern of the Weddell Gyre as de-
 329 picted by the model is in agreement with both observational and past modeling stud-
 330 ies. All resolutions display a double lobed circulation structure, with a larger western
 331 cell and a recirculation close to the Greenwich Meridian. As the resolution increases, the
 332 region east of approximately 10°E becomes increasingly unstable and the cyclonic cir-
 333 culation expands towards the east along the coastline.

334 The barotropic circulation generally follows potential vorticity contours (approx-
 335 imated by f/H) at the south, west and northern boundaries; these contours are largely
 336 controlled by topographic features, indicating that the motion is steered by bathymetry.
 337 The larger spacing between f/H contours in the southern limb is associated with a broad,
 338 slow circulation that contrasts with a more intense northern boundary flow. At the north-
 339 eastern boundary the flow crosses potential vorticity contours, where there is an abrupt
 340 deflection of f/H contours owing to a gap in the Southwestern Indian Ridge at around
 341 30°E , 55°S . The crossing of f/H contours indicates that the flow experiences changes
 342 in its potential vorticity due to its interaction with topography, an interaction that is

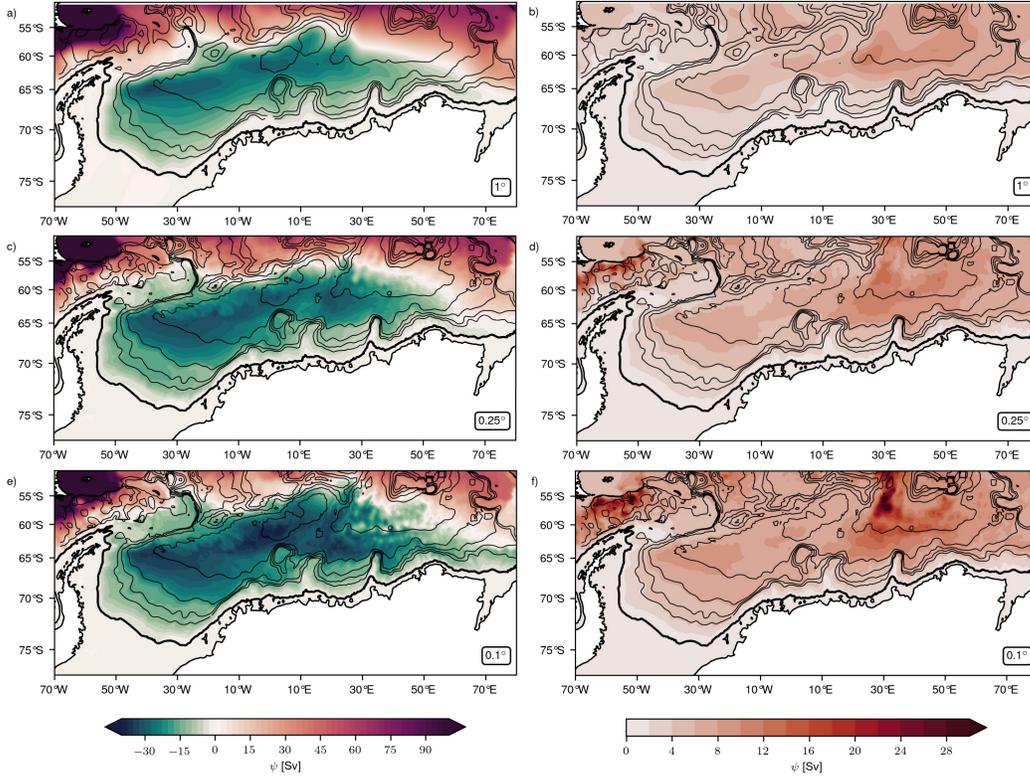


Figure 6. Mean barotropic streamfunction (Sv, $1 \text{ Sv} = 10^6 \text{ m}^3/\text{sec}$) for the period 1958 - 2018 for (a) 1° , (c) 0.25° and (e) 0.1° and its standard deviation for the (b) 1° , (d) 0.25° and (f) 0.1° resolution models with contours of potential vorticity (f/H) and 1000m isobath indicated by the thin and thick black contours respectively.

343 highly dependent on resolution. The deflection of f/H contours, together with the strong
 344 flow of the ACC, give rise to instabilities at the tail of the ridge particularly at 0.1° res-
 345 olution, that leave a trace in the mean circulation (Figure 6e), as highlighted by the large
 346 values in the standard deviation of ψ (Figure 6f).

347 The subsurface temperature maximum can be used to trace the path CDW takes
 348 after entering the gyre at its eastern boundary. Observations show this warm inflow and
 349 the subsequent lowering of its temperature due to mixing as CDW flows to the west (see
 350 Figure 7a using Reeve et al. (2019) gridded data set as an example). The 0.1° simu-
 351 lation is the only configuration that displays such warm inflow. The improved accuracy
 352 of the 0.1° model in reproducing this feature of the gyre could be crucial, since the in-
 353 flow affects the characteristics of bottom waters formed at the continental shelf (Couldrey
 354 et al., 2013; Jullion et al., 2014; Fahrback et al., 2011).

355 4 Weddell Gyre Seasonal Cycle

356 The three model simulations show a seasonal cycle of gyre strength consisting of
 357 a winter intensification and a summer weakening (Figure 8a) that increases in amplitude
 358 with resolution and accounts for 36%, 46% and 44% of the variability in gyre strength
 359 for the 1° , 0.25° and 0.1° configurations respectively. There is a clear increase in trans-
 360 port from the 1° and 0.25° resolutions to the 0.1° resolution. The seasonal cycle of sur-
 361 face stress curl is strongly modulated by changes in sea ice concentration (α_{ice} in Equa-

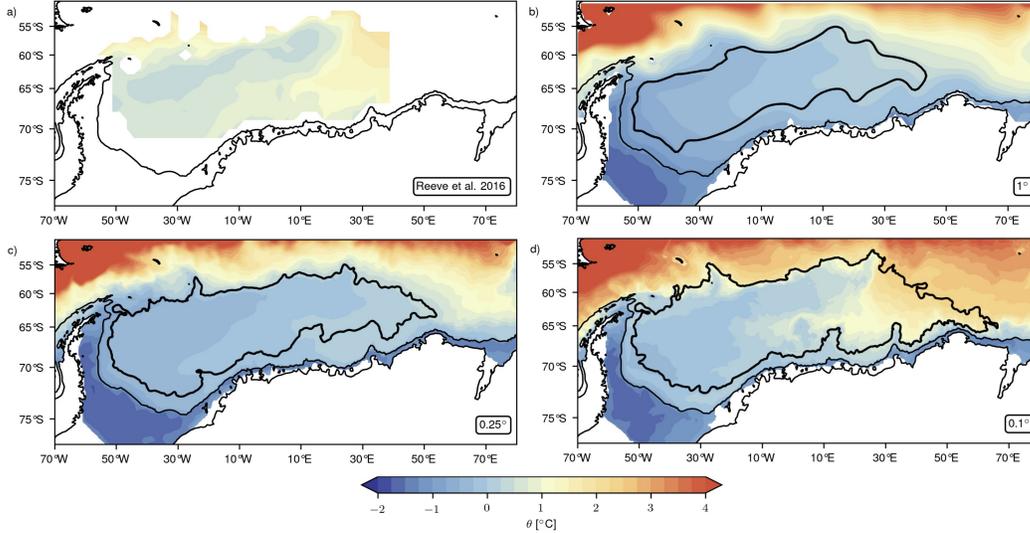


Figure 7. Subsurface mean potential temperature maximum for the period 2001 to 2014 from (a) the Reeve et al. (2016) product and for the period 1958 - 2018 from the (b) 1° , (c) 0.25° and (d) 0.1° resolution simulations. The thick black contour marks the gyre’s boundary and the thin black contour the 1000m isobath.

362 tion 3): during the summer months, the minimum in sea ice concentration makes wind
 363 stress the primary contributor to the total stress curl. As the ice pack begins to develop,
 364 it damps the transfer of momentum from the atmosphere to the ocean, generating a winter
 365 minimum in total stress curl relative to the autumn peak. After reaching its max-
 366 imum during August-September, α_{ice} begins to decrease, allowing for a second maximum
 367 in total stress curl during spring. The seasonal cycle of surface buoyancy fluxes, Figure
 368 8c, is also strongly modulated by the presence of sea ice. During the summer, months
 369 sea ice melt, along with warmer air temperatures, result in a net surface buoyancy gain,
 370 with the opposite occurring during winter via sea-ice formation, brine rejection, and at-
 371 mospheric cooling.

372 Correlations between gyre strength, surface stress curl and surface buoyancy fluxes
 373 are shown in Table 2. Gyre strength is significantly correlated with both forcing mech-
 374 anisms considered, indicating that buoyancy losses and surface stress curl intensification
 375 are concurrent with a stronger gyre, with the exception of the 0.1° resolution simulation
 376 whose annual cycle of gyre strength is not correlated with the semi-annual cycle of sur-
 377 face stress curl. There is a clear increase in transport from the 1° and 0.25° simulations
 378 to the 0.1° that is not explained by differences between simulations in either surface stress
 379 curl or surface buoyancy fluxes. Moreover, the surface stress curl (Figure 8b) is weak-
 380 est for the 0.1° simulation and there are no appreciable differences in the seasonal cycle
 381 of surface buoyancy fluxes (Figure 8c).

382 There are spatial differences in the amplitude of seasonal variability. The summer/winter
 383 anomalies with respect to the mean for the barotropic streamfunction in Figure 9a and
 384 c show that the largest seasonal variations are located in two separate regions around
 385 30°W and 30°E respectively. The first region is located adjacent to a region of signif-
 386 icant buoyancy gain during summer and the latter is located beneath the regional low
 387 pressure system (Figures 9b and d) which deepens and expands during the winter months
 388 and is shallower during summer, modifying the local wind pattern. The gyre boundary,

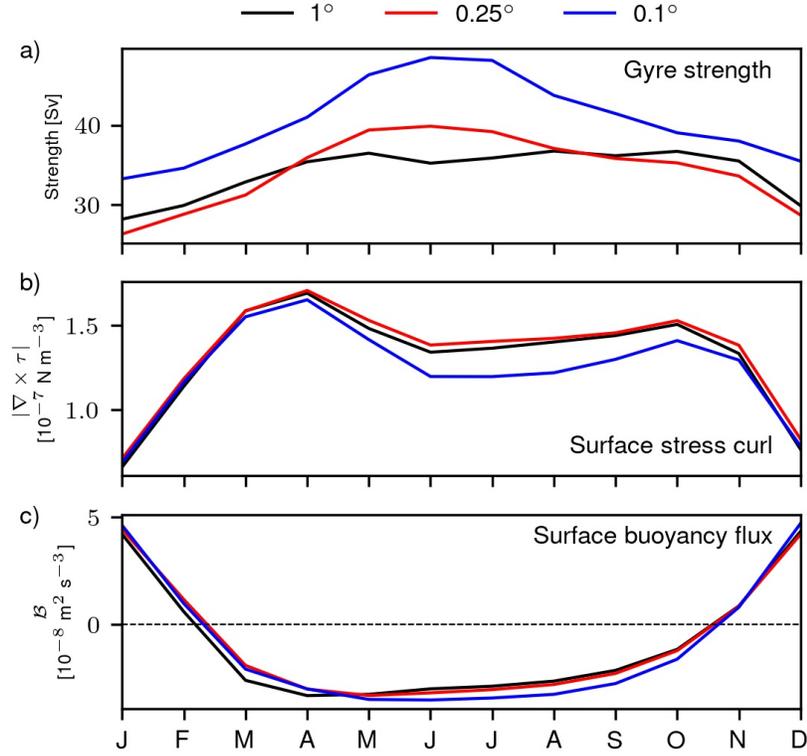


Figure 8. Annual cycles of (a) gyre strength (Sv), (b) surface stress curl ($10^{-7}Nm^{-3}$) averaged over the region bounded by the solid orange box in Figure 1 and (c) surface buoyancy flux ($10^{-8}m^2s^{-3}$) averaged at depths shallower than 1000 m.

Table 2. Correlation coefficients between gyre strength and surface stress curl, buoyancy flux, SAM index and EAS index for seasonal and interannual time scales. Bold values indicate significant correlations with $p < 0.05$.

		Seasonal			
		$ \nabla \times \tau $		B	
Gyre strength	1°	0.82		-0.81	
	0.25°	0.71		-0.89	
	0.1°	0.38		-0.83	
		Interannual			
		$ \nabla \times \tau $	B	SAM	EAS
Gyre strength	1°	0.04	-0.11	-0.36	0.26
	0.25°	0.21	-0.32	-0.24	0.34
	0.1°	0.21	-0.29	-0.19	0.51

389 defined by the 12 Sv contour, shows a general expansion during winter in all directions,
 390 with a particularly large excursion towards the east.

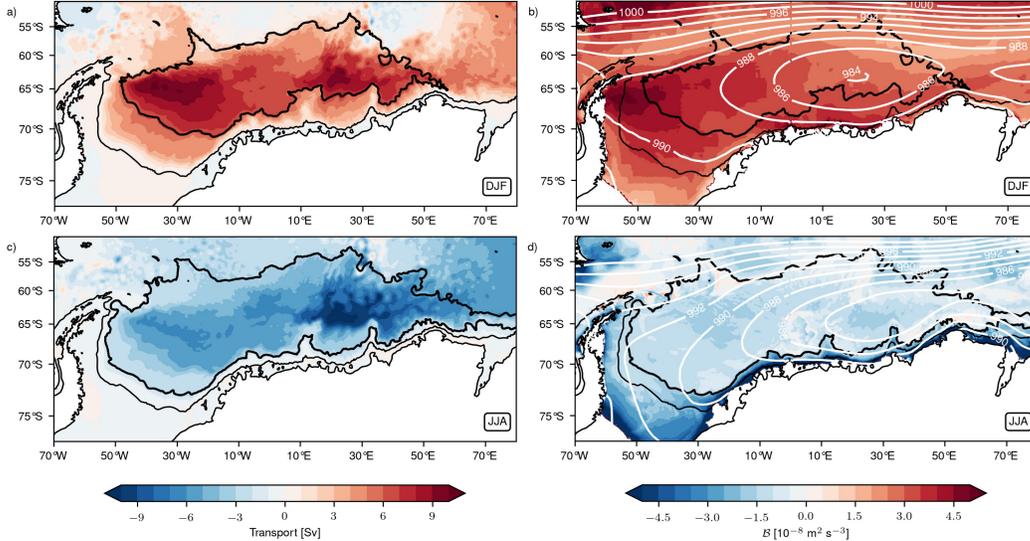


Figure 9. (a) Summer and (c) winter anomalies of ψ and (b) summer and (d) winter climatological fields of net surface buoyancy fluxes with and sea level pressure (white contours) for the 0.1° resolution case. Thin black contour indicates the 1000m isobath and thick black contour the Weddell Gyre’s boundary in the corresponding season.

5 Weddell Gyre Interannual Variability

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In order to explore the interannual variability of the Weddell Gyre and its connection to regional and large scale climate, we analyse the full 61-year period time series of anomalies with respect to the seasonal cycle of buoyancy fluxes, surface stress curl, gyre strength and the SAM and EAS index calculated as described in Section 2. These time series are shown in Figure 10. Gyre strength displays significant interannual variability that becomes greater in magnitude with resolution, with standard deviations of 4.5, 5 and 6 Sv for the 1° , 0.25° and 0.1° configurations respectively. There are some years which show an intensification of around 10 Sv, which is between a third and a quarter of the gyre’s mean strength depending on the resolution. The strengths of the three configurations are mostly in phase and significantly correlated with each other, indicating the predominance of external forcing in driving the gyre’s interannual variability. Similar to what was observed for the seasonal cycles, there is a clear influence of model resolution in setting the magnitude of gyre strength and the increase in magnitude from the 1° and 0.25° simulations to the 0.1° is not related to differences between resolutions in surface stress curl or buoyancy fluxes. Correlations between gyre strength and the other time series in Figure 10 are shown in Table 2. From the climate metrics considered here, only EAS shows a significant correlation with gyre strength for the 0.1° resolution configuration.

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Next we identify strong events and weak events of gyre strength in the 0.1° resolution simulation, highlighted in red and blue shading respectively in Figure 10. There is no periodicity or seasonality in the occurrence of these events, i.e. extreme strong and weak events do not have a preference for particular months or seasons. Figure 11a shows the composites of strong events minus the composites of weak events for the barotropic streamfunction and sea level pressure. Since the cyclonic circulation of the Weddell Gyre is represented by $\psi < 0$ (Figure 6), negative anomalies indicate an intensified circulation. Sea level pressure shows a deepening and displacement to the west of a low pressure system that deepens the trough in the southern portion of the gyre, thus intensi-

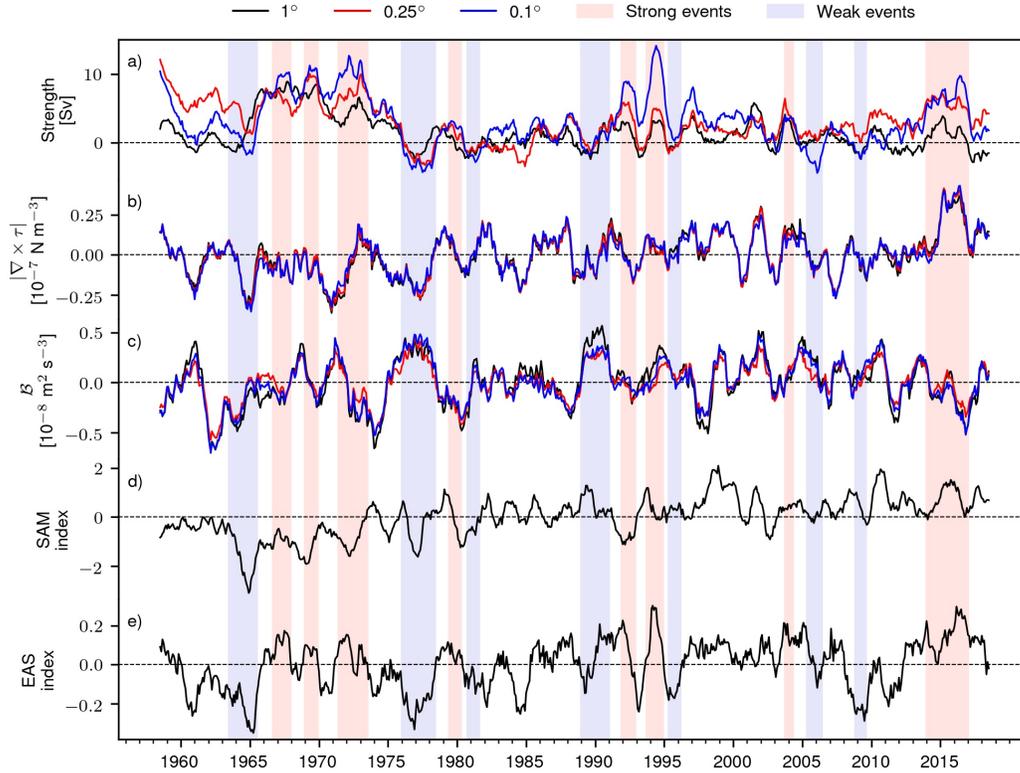


Figure 10. Time series of monthly anomalies with respect to the long term seasonal cycle of (a) gyre strength (Sv), (b) surface buoyancy flux ($10^{-8}m^2s^{-3}$), (c) surface stress curl ($10^{-7}Nm^{-3}$), (d) SAM index and (e) EAS index for the 1° , 0.25° and 0.1° resolutions. A 12-month running filter was applied to all time series. Red and blue shading indicate strong and weak events respectively for the 0.1° simulation.

419 flying the local easterly winds. The composites of surface buoyancy flux, Figure 11b show
 420 that during strong years, the southwestern region of the gyre experiences buoyancy loss
 421 in a region where buoyancy fluxes are significantly correlated with gyre strength, as marked
 422 by the hatching. These changes in buoyancy fluxes are consistent with the composite for
 423 sea ice concentration in Figure 11d: i.e., a stronger circulation advects sea ice from the
 424 southwestern region, decreasing its concentration during strong years, exposing the ocean's
 425 surface to the atmosphere, thus enabling a larger buoyancy loss. The ice is advected to
 426 the north and then to the east, generating an increase of almost 10% in sea ice concen-
 427 tration at the northern rim of the gyre. Another region where strength impacts sea ice
 428 concentration is the eastern boundary, where we see a decrease in concentration during
 429 stronger years consistent with a stronger warm inflow. Stronger years also display warmer
 430 subsurface temperature maxima along the southern region of the gyre that translate into
 431 a warmer continental shelf, Figure 11c. At around $30^\circ E$ there are significant correlations
 432 between the subsurface temperature maximum, indicating synchronous changes in gyre
 433 strength and the characteristics of the warm inflow.

434 6 Discussion and Summary

435 In this study we have used a coupled ocean-sea ice model at three different hor-
 436 zontal resolutions (1° , 0.25° and 0.1°) to diagnose the circulation of the Weddell Gyre

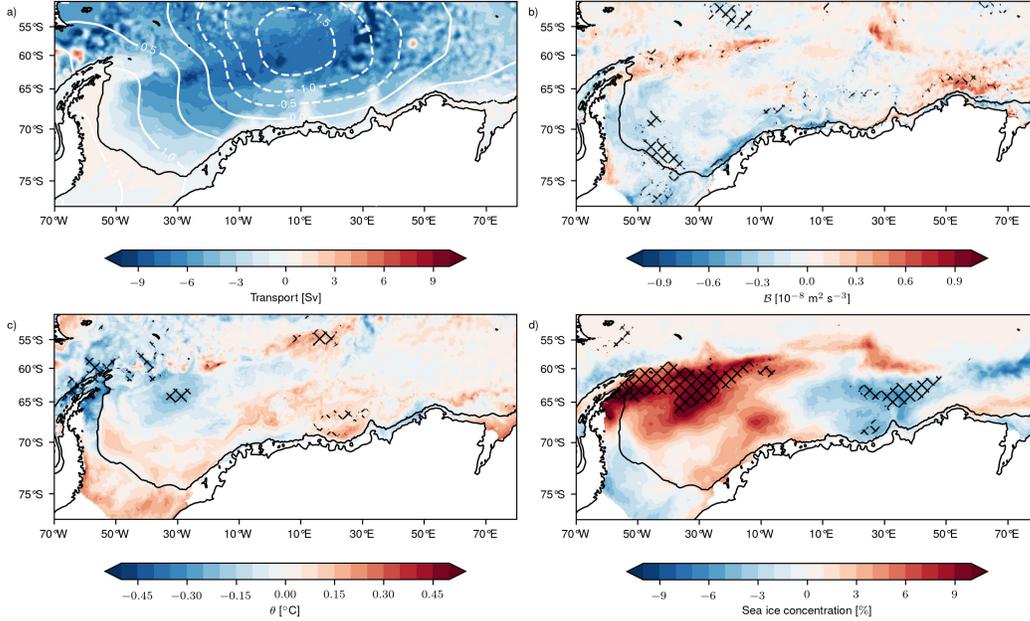


Figure 11. Strong minus weak Weddell Gyre composites of monthly anomalies for the events highlighted in Figure 10. Composites of (a) barotropic streamfunction (Sv) with sea level pressure contours overlaid (HPa), (b) surface buoyancy flux ($10^{-8} Nm^{-3}$), (c) subsurface temperature maximum ($^{\circ}C$) and (d) sea ice concentration (%) for the 0.1° simulation. Hatching indicates significant correlations ($p < 0.1$) between gyre strength time series and the corresponding field..

437 and its connection to possible forcing mechanisms, namely surface stress forcing (including
 438 the contributions of wind and sea ice) and surface buoyancy fluxes. The inclusion
 439 of sea ice and the additional consideration of surface buoyancy fluxes expands on past
 440 studies that solely consider wind forcing due to the lack of appropriate observations. We
 441 evaluate the model against available satellite observations and hydrographic data and
 442 find a distinct improvement with resolution, particularly in the representation of the char-
 443 acteristics of regional water masses, which we suggest is a direct consequence of the abil-
 444 ity of the 0.1° configuration to resolve instabilities along the eastern boundary of the gyre.
 445 The model displays significant seasonal and interannual variability in gyre strength, the
 446 magnitude of which increases with resolution. We find that extreme events of gyre cir-
 447 culation are climatically distinct, with significant differences in sea ice concentration and
 448 water mass characteristics during strong and weak phases of Weddell Gyre flow.

449 Our analysis indicates that most of the variability in gyre strength is concentrated
 450 in the seasonal cycle, which explains between 35% and 45% of the variance depending
 451 on the model's horizontal resolution. This seasonal cycle consists of a winter intensifi-
 452 cation and a summer weakening, in agreement with past studies (Dellnitz et al., 2009;
 453 Beckmann et al., 1999), that increases in magnitude with resolution. Consequently, es-
 454 timates obtained from summer-biased observations will underestimate the gyre's trans-
 455 port, stressing the importance of sustained, continuous observations in the Weddell re-
 456 gion. We further find that the seasonal changes in the barotropic circulation are not spa-
 457 tially uniform, but are concentrated in an eastern and a western region at approximately
 458 $30^{\circ}W$ and $30^{\circ}E$ respectively. Surface stress curl displays a semi-annual seasonal cycle
 459 peaking in April and October, highlighting the importance of sea ice in modulating the
 460 transfer of momentum from the atmosphere to the ocean, as emphasized by A. Naveira Gara-
 461 bato et al. (2019). Sea ice is also relevant for the seasonal cycle of surface buoyancy fluxes,

462 where the ocean gains buoyancy during summer due to sea ice melt and loses surface buoy-
 463 ancy during winter due to sea ice formation.

464 The gyre also displays significant variability on interannual timescales, with gyre
 465 strength in phase across the three model resolutions, indicating the primacy of atmo-
 466 spheric forcing in driving the gyre’s variability. However, the variations in gyre strength
 467 become more pronounced with resolution, the consequences of which we explored by means
 468 of composites of strong and weak events on interannual timescales using the 0.1° con-
 469 figuration. The composites show that the gyre circulation has an impact on relevant pro-
 470 cesses within the region: a stronger gyre coincides with a warmer inflow at the eastern
 471 boundary that warms up the continental shelf and induces large changes in sea ice con-
 472 centration by advection towards the northern rim. Anomalous gyre events over inter-
 473 annual timescales are associated with a deepening and expansion to the west of the re-
 474 gional low pressure system that deepens the trough and accelerates the coastal easter-
 475 lies. This acceleration of the easterlies increases the onshore Ekman transport, raising
 476 the slope in sea level and accelerating the southern limb of the Weddell Gyre, a mech-
 477 anism supported by the significant correlation between gyre strength and the easterlies
 478 at 0.1° model resolution.

479 For the purpose of this study we have considered two possible forcing mechanisms
 480 for the Weddell Gyre’s variability: namely surface stress curl, traditionally supposed to
 481 drive the circulation in the depth-integrated circulation of the ocean via a linear vortic-
 482 ity relationship with the ocean’s meridional transport (Munk, 1950; Sverdrup, 1947), and
 483 buoyancy fluxes, which some studies suggest can maintain a gyre-like circulation (Hogg
 484 & Gayen, 2020; Wang & Meredith, 2008) and even drive a mean ACC (Howard et al.,
 485 2015; Hogg, 2010). The weak stratification that characterizes the region, together with
 486 intense surface buoyancy fluxes, drive an overturning circulation (Jullion et al., 2014; A. C. Naveira Gara-
 487 bato et al., 2016) that could affect the gyre’s circulation via changes in its stratification.
 488 The coupling between ocean and sea ice in the model used for this study does not al-
 489 low us to consider the independent contributions from buoyancy fluxes and surface stress.
 490 For example, on seasonal timescales, the lower resolution model’s gyre strength is cor-
 491 related with both surface stress curl and surface buoyancy fluxes, yet on interannual timescales
 492 neither show a clear correlation. Thus, the simulations presented here cannot ascertain
 493 which factors, or non-linear interactions between them, are more relevant in driving the
 494 gyre’s variability. We conclude that more targeted numerical experiments are needed to
 495 address this question and to separate out stress from buoyancy forcing.

496 The importance of model resolution in the region is emphasized by its evaluation
 497 against hydrographic observations. Because of the Weddell Gyre’s open configuration
 498 at its eastern boundary, the properties of the water masses in the region are highly in-
 499 fluenced by the characteristics of the inflow (Kerr et al., 2018; Jullion et al., 2014; Coul-
 500 drey et al., 2013). Two possible pathways have been suggested for this inflow: an eddy-
 501 driven path in the northeastern gyre boundary and an advective path further south to-
 502 wards the Antarctic continent (Ryan et al., 2016; Cisewski et al., 2011; Leach et al., 2011;
 503 Gouretski & Danilov, 1993). The eddy-driven inflow can be identified by a subsurface
 504 temperature maxima that is only reproduced by the 0.1° configuration, indicating that
 505 finer scale resolution processes are important for the exchange of waters between the gyre
 506 and its surroundings. The characteristics of this warm inflow are likely responsible for
 507 the improvement with resolution of the temperature-salinity structure of the Weddell Gyre,
 508 potentially also affecting the characteristics of bottom waters in the model.

509 The variability of the Weddell Gyre has the potential to impact regional processes
 510 of relevance to global climate, but there is still not a complete understanding of the gyre’s
 511 driving mechanisms and timescales of interaction. We have considered surface stress and
 512 buoyancy fluxes as possible drivers of variability, but in a coupled ocean-sea ice model
 513 their intertwined nature does not allow us to consider them as independent mechanisms.
 514 It has also been suggested that the northern and southern limbs of the gyre vary inde-

515 pendently, forced by the westerlies and easterlies respectively, with the difference bal-
 516 anced by inflows/outflows enabled by the gyre’s open boundaries (Fahrbach et al., 2011).
 517 Our definition of gyre strength using the barotropic streamfunction is not intended to
 518 capture this independent variation. Remote forcing via the propagation of anomalies along
 519 the Antarctic continent could also affect the Weddell Gyre’s circulation, but assessing
 520 this hypothesis is beyond the scope of this study. The 0.1° resolution model has the abil-
 521 ity to reproduce observed key features of this polar region and provides a 61-year long
 522 data set that can prove useful for expanding our current knowledge of the Weddell Gyre.

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 531 model output for the simulations presented in this paper are stored in the COSIMA data
 532 collection, available from <http://dx.doi.org/10.4225/41/5a2dc8543105a>. The an-
 533 alytical scripts used to generate the figures are available at [https://github.com/julia](https://github.com/julia-neme)
 534 `-neme`

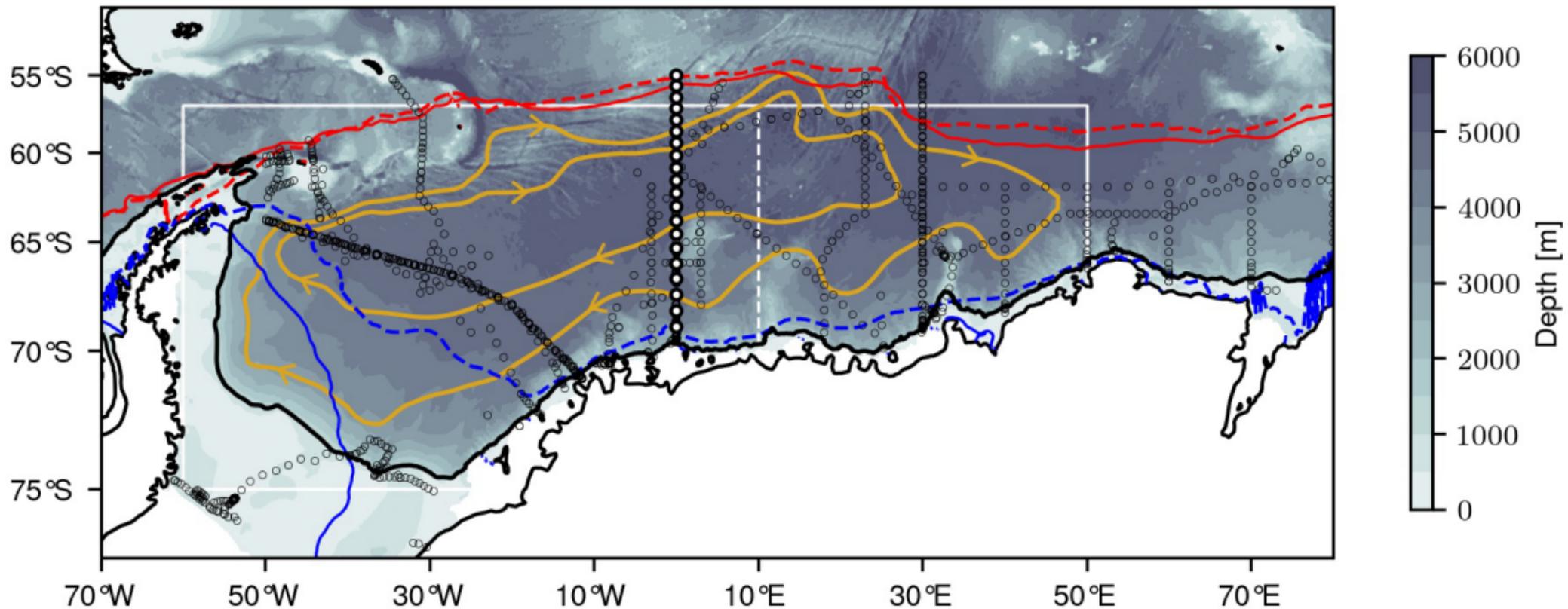
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Figure 1.



- | | | | |
|----------------|-------------------------|-------------------|-------------------|
| ○ A12 | — 1000m isobath | — Feb. SIC (0.1°) | — Feb. SIC (NOAA) |
| ○ Hyd. station | → Schematic circulation | — Sep. SIC (0.1°) | — Sep. SIC (NOAA) |

Figure 2.

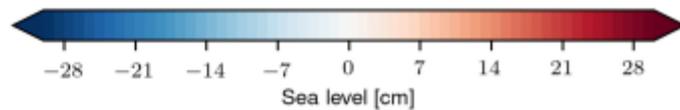
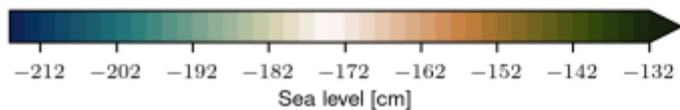
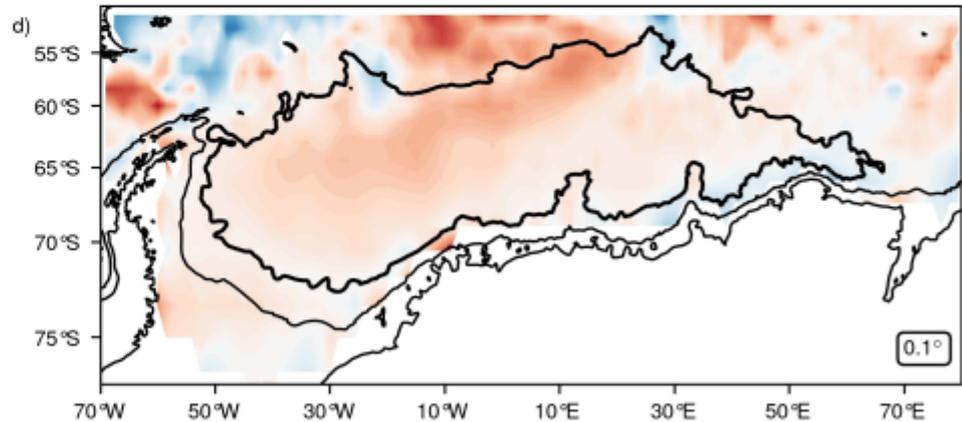
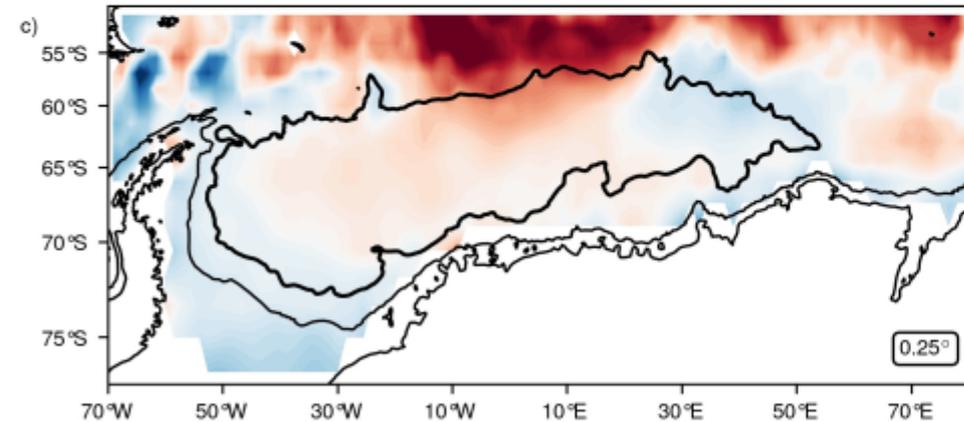
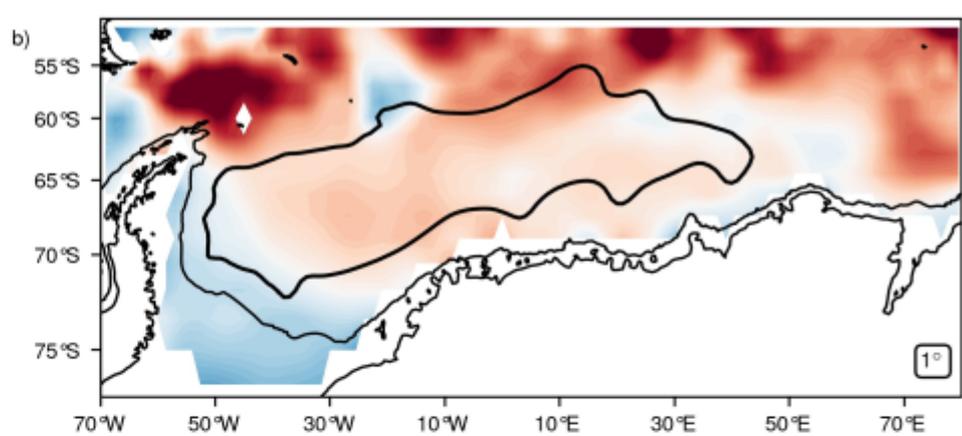
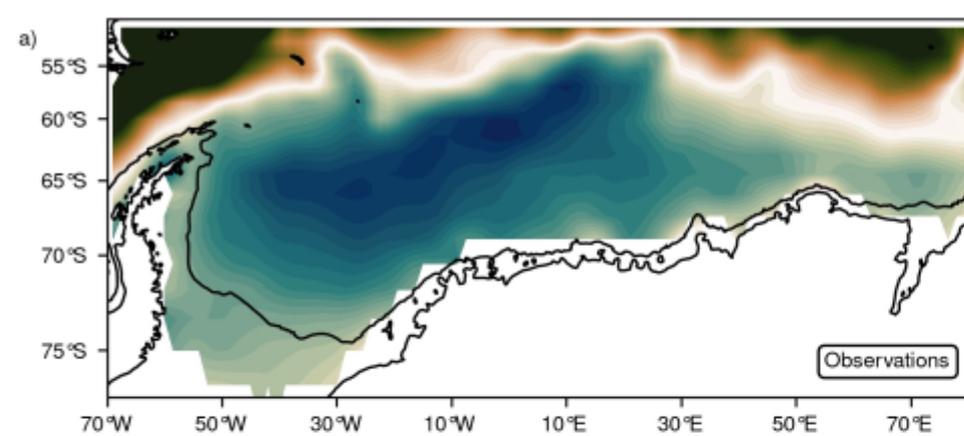


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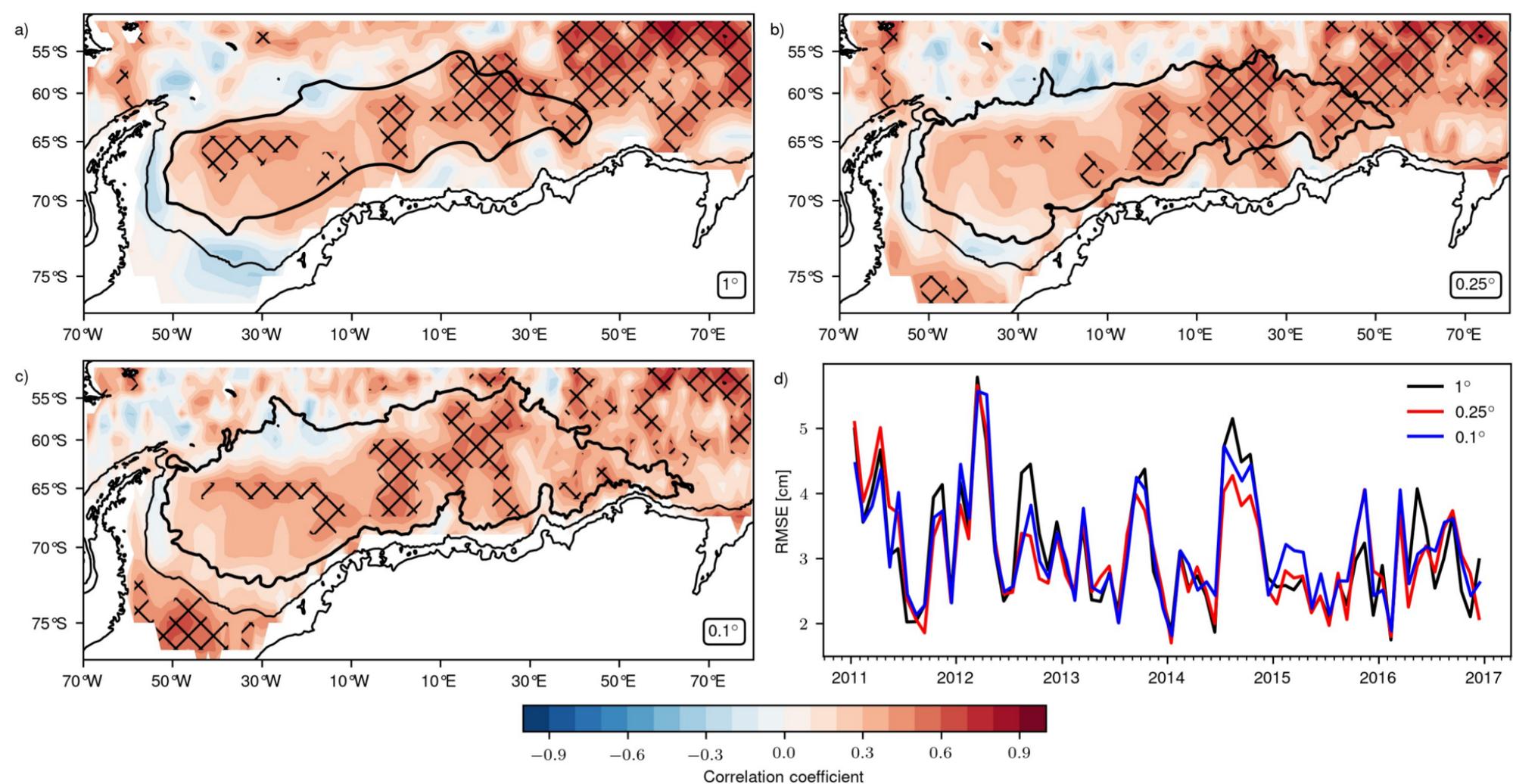


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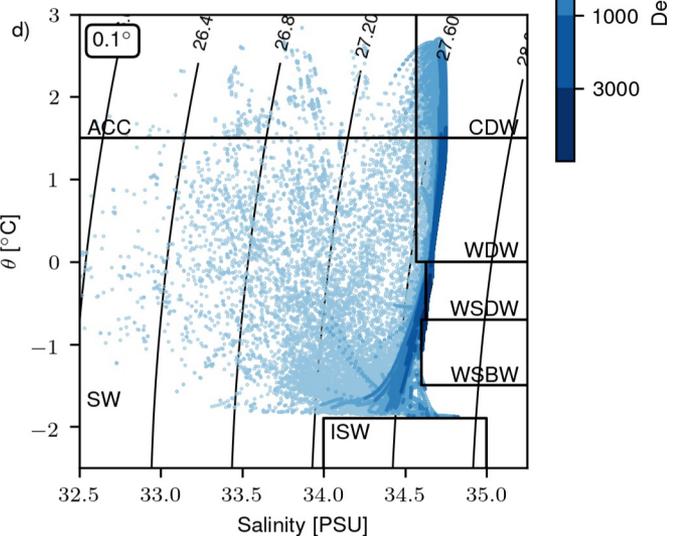
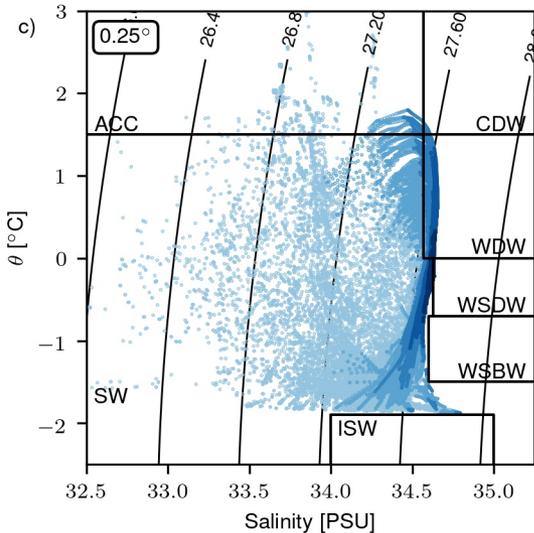
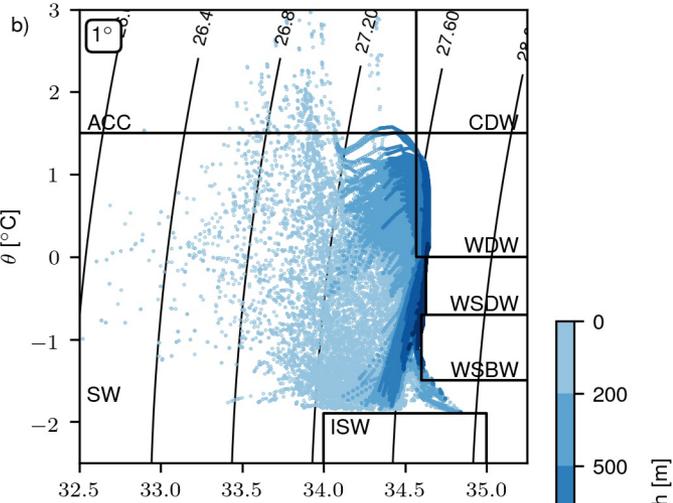
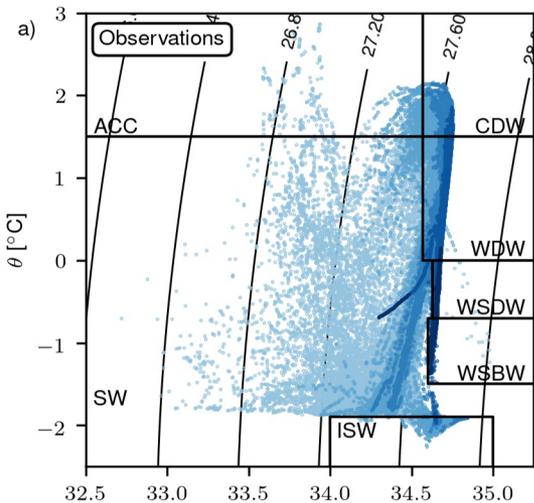


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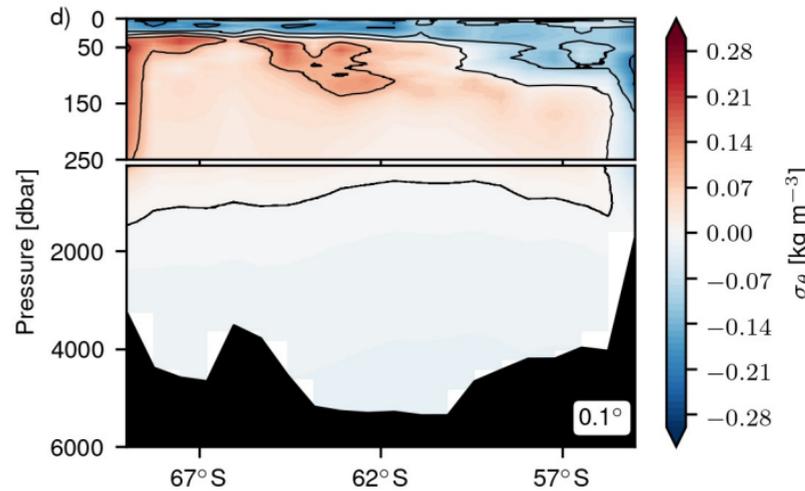
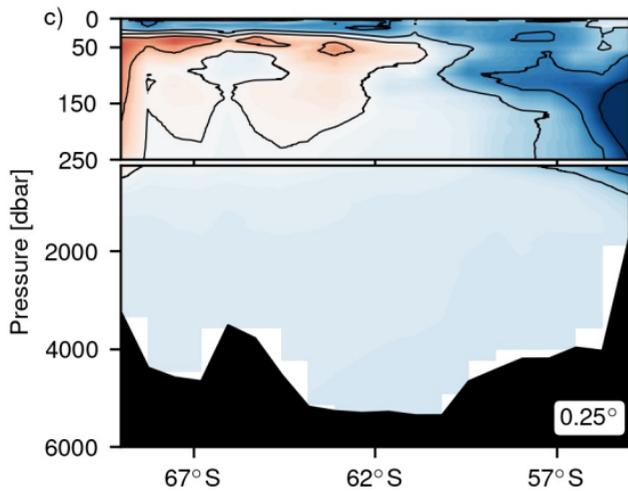
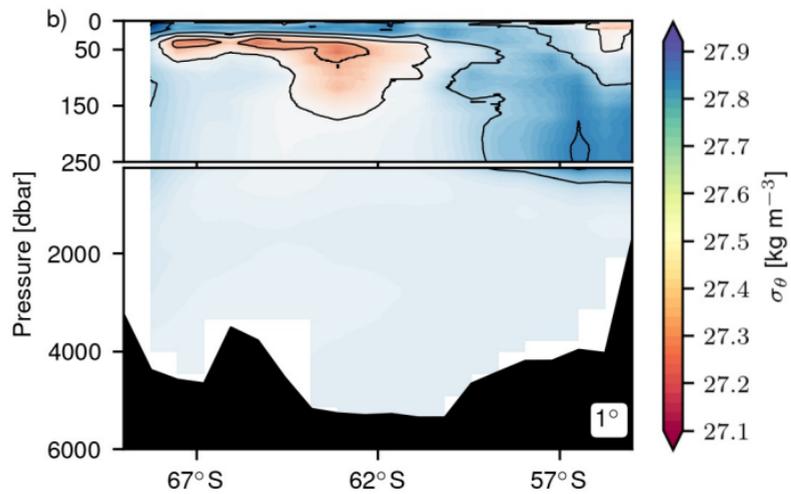
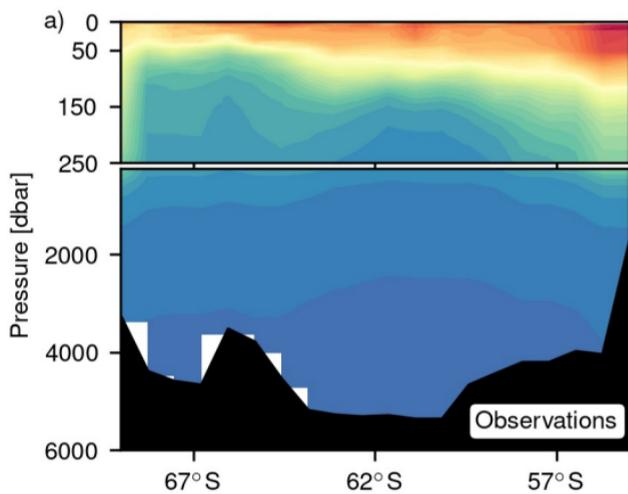


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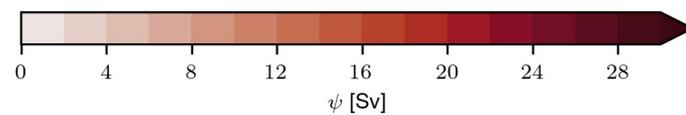
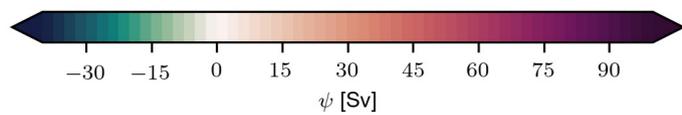
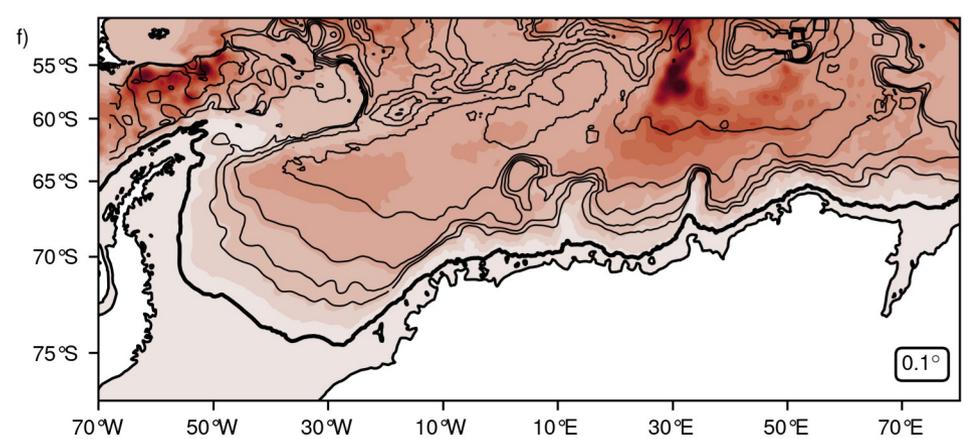
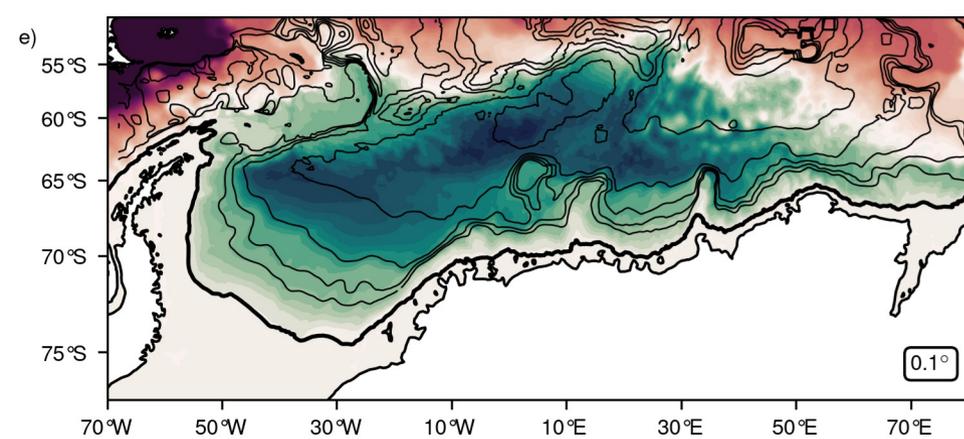
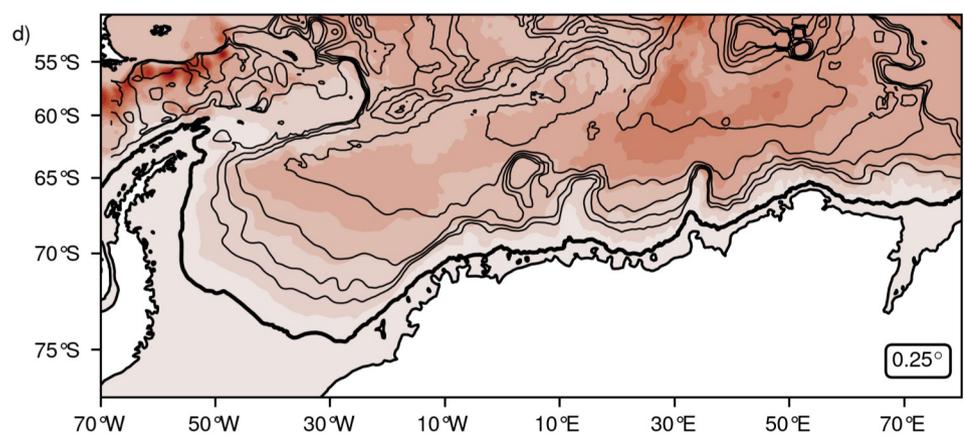
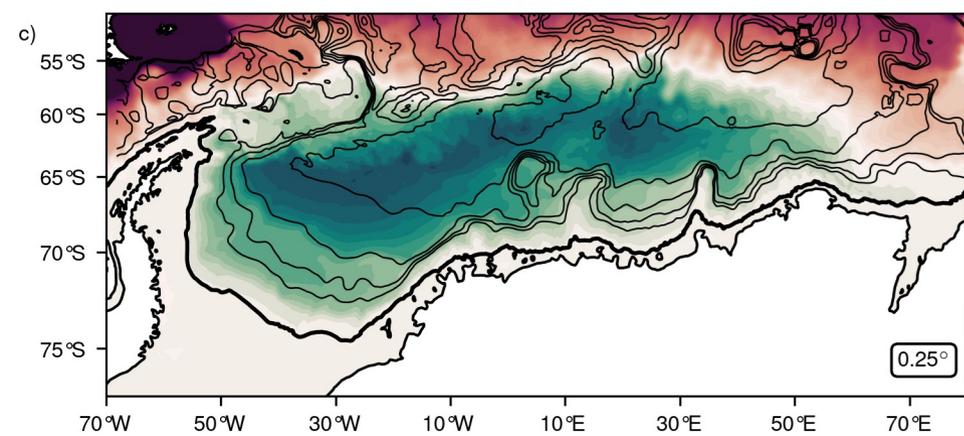
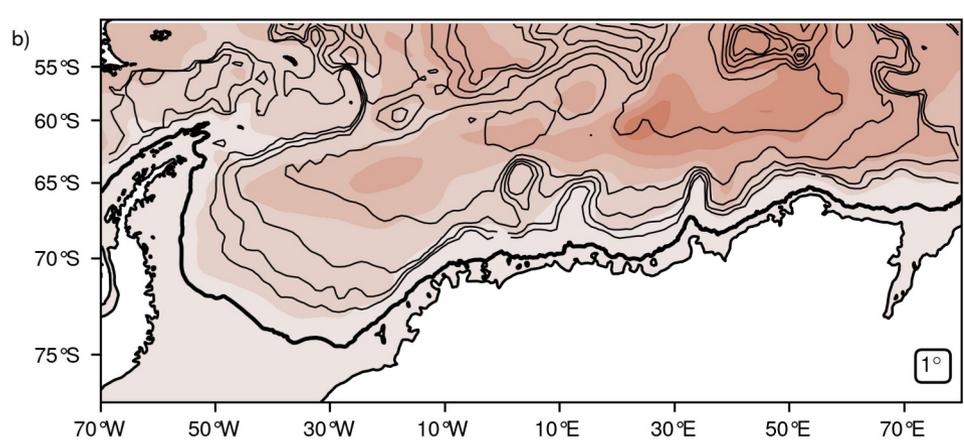
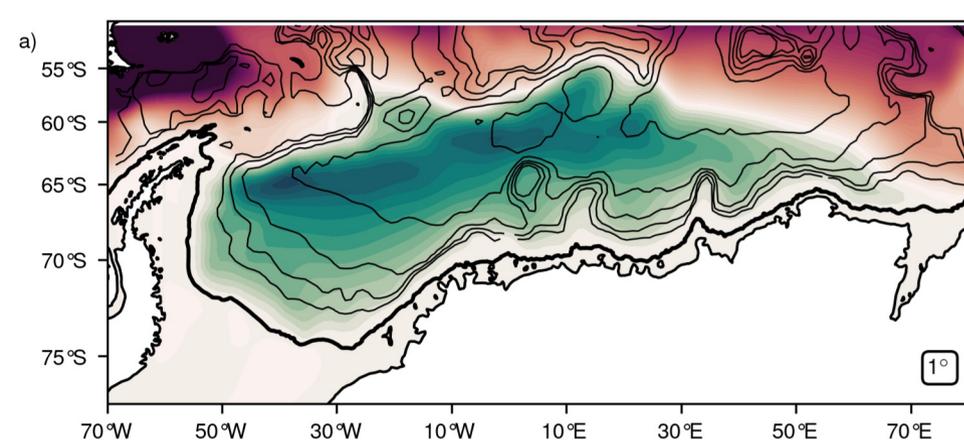


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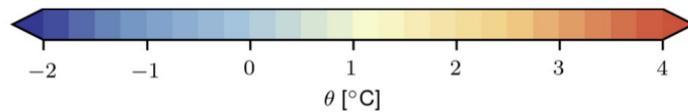
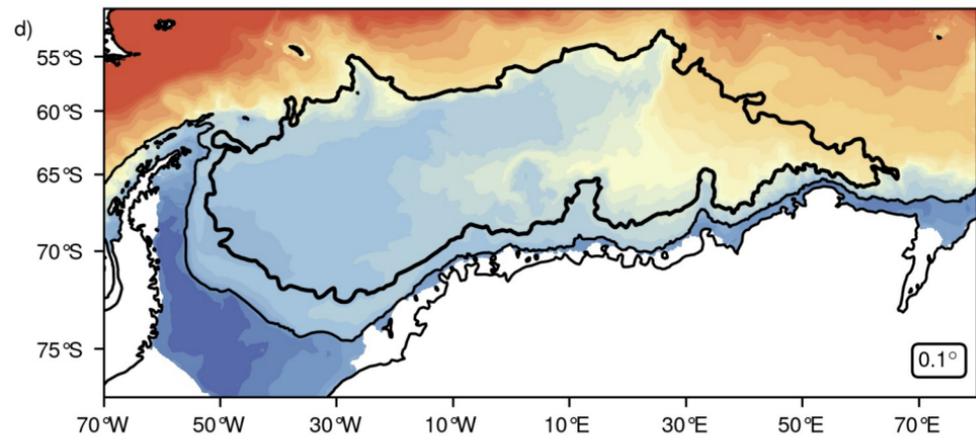
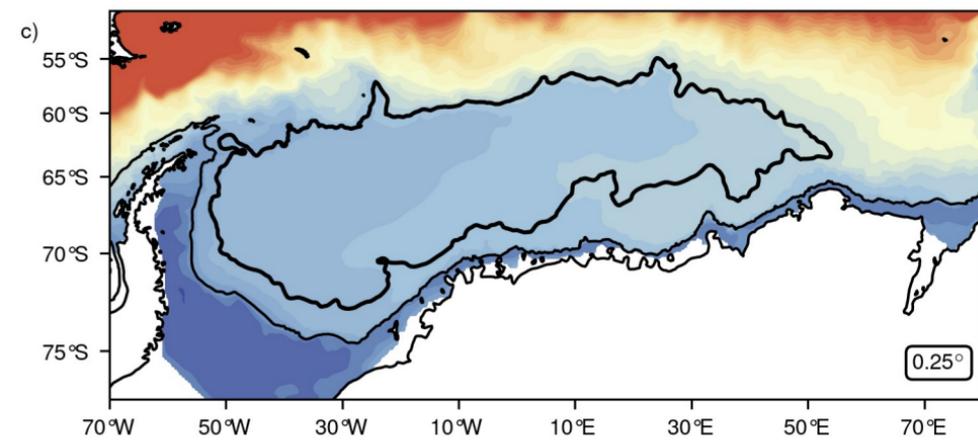
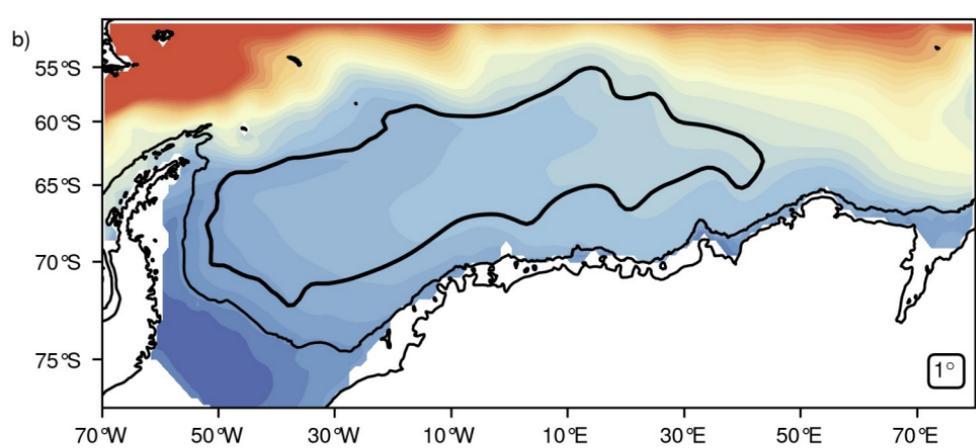
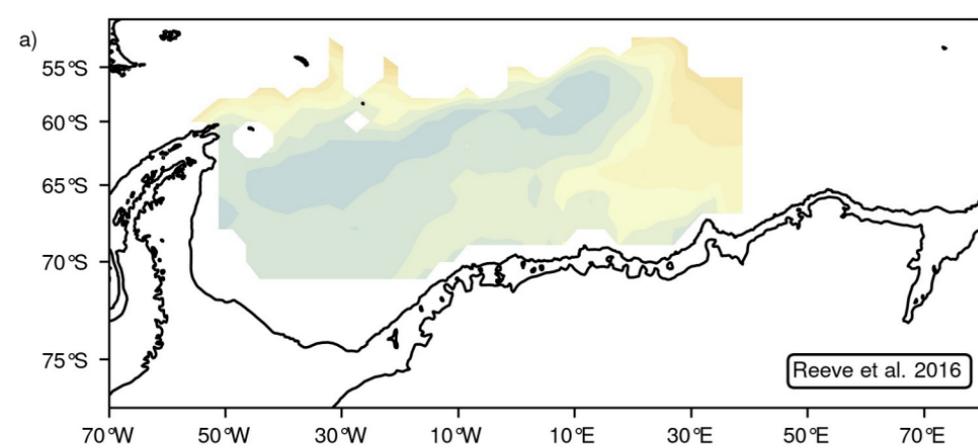


Figure 8.

— 1° — 0.25° — 0.1°

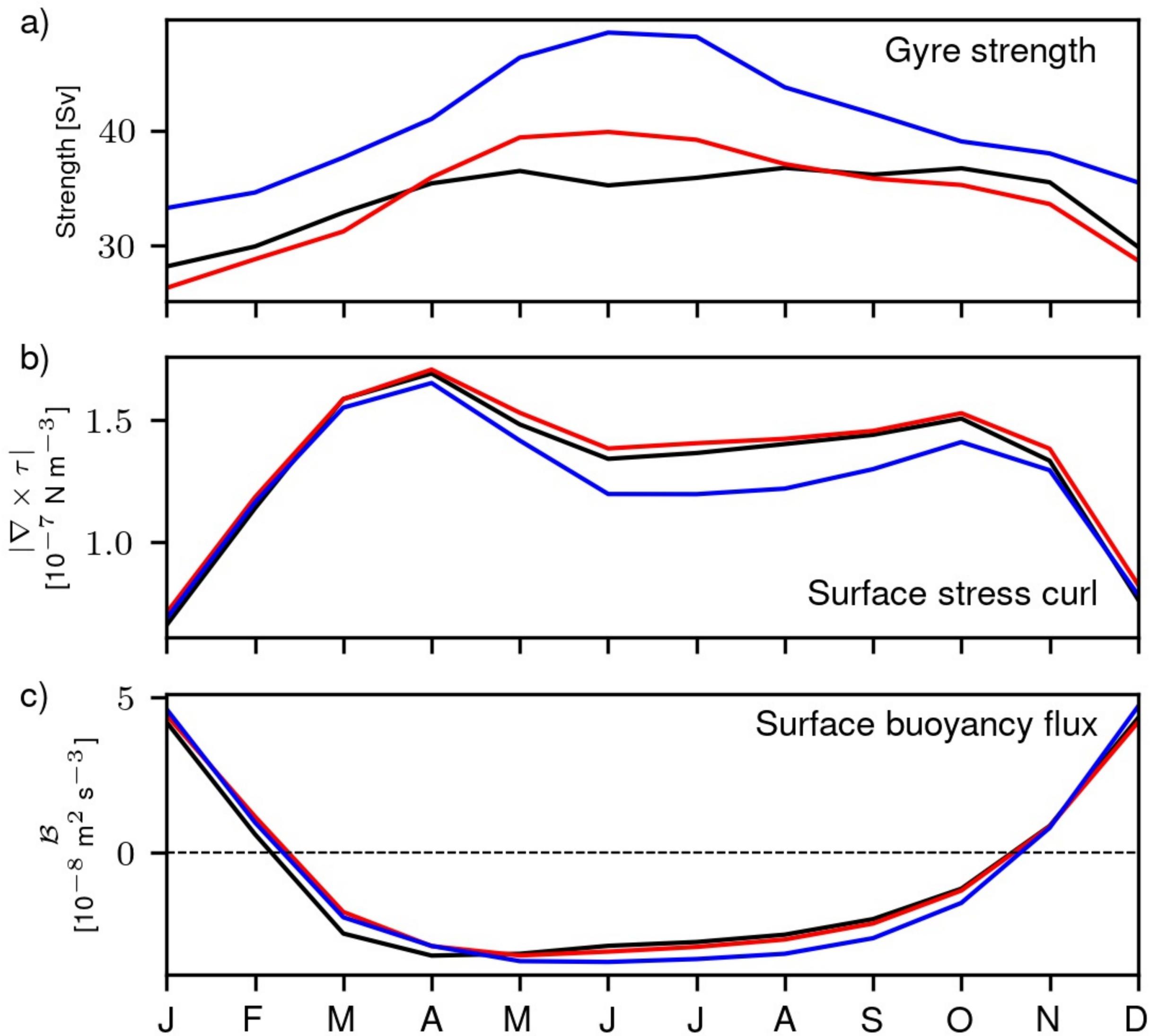


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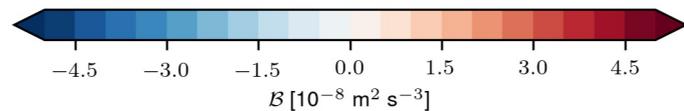
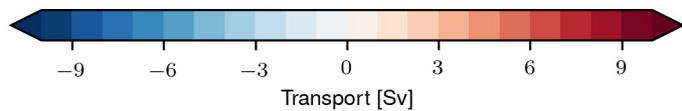
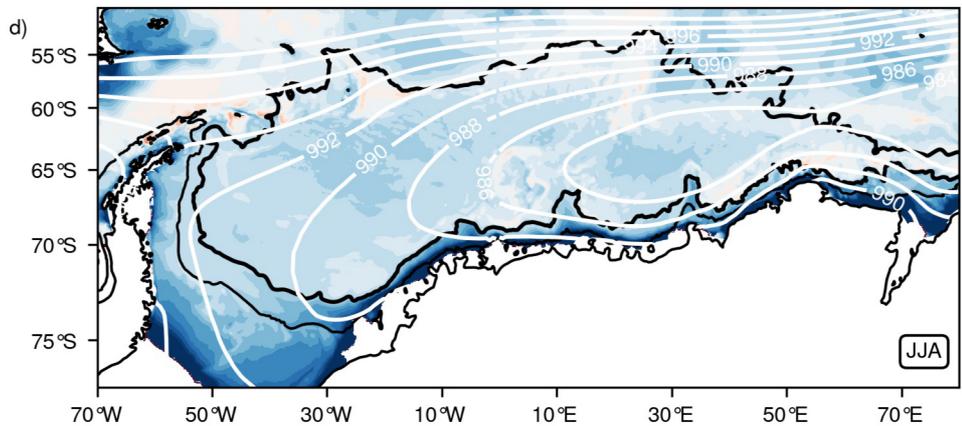
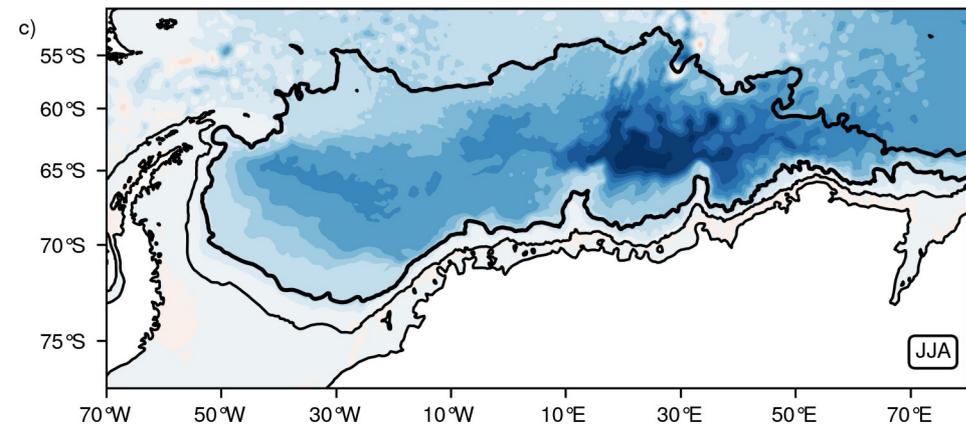
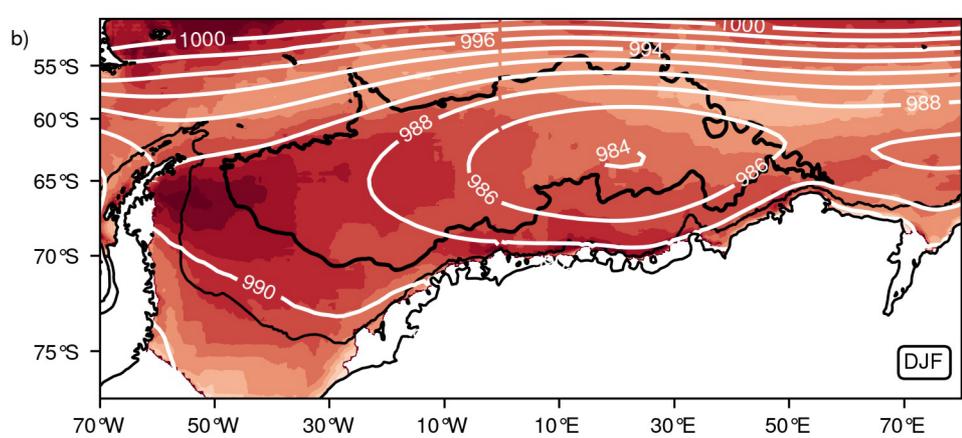
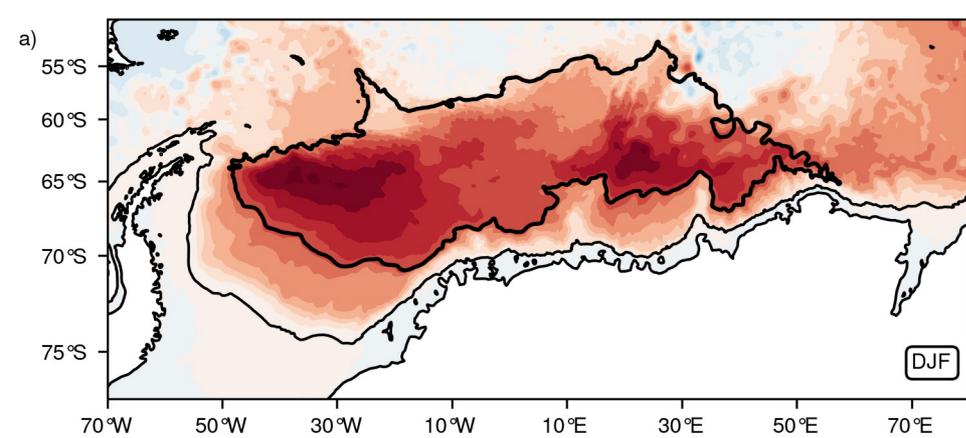


Figure 10.

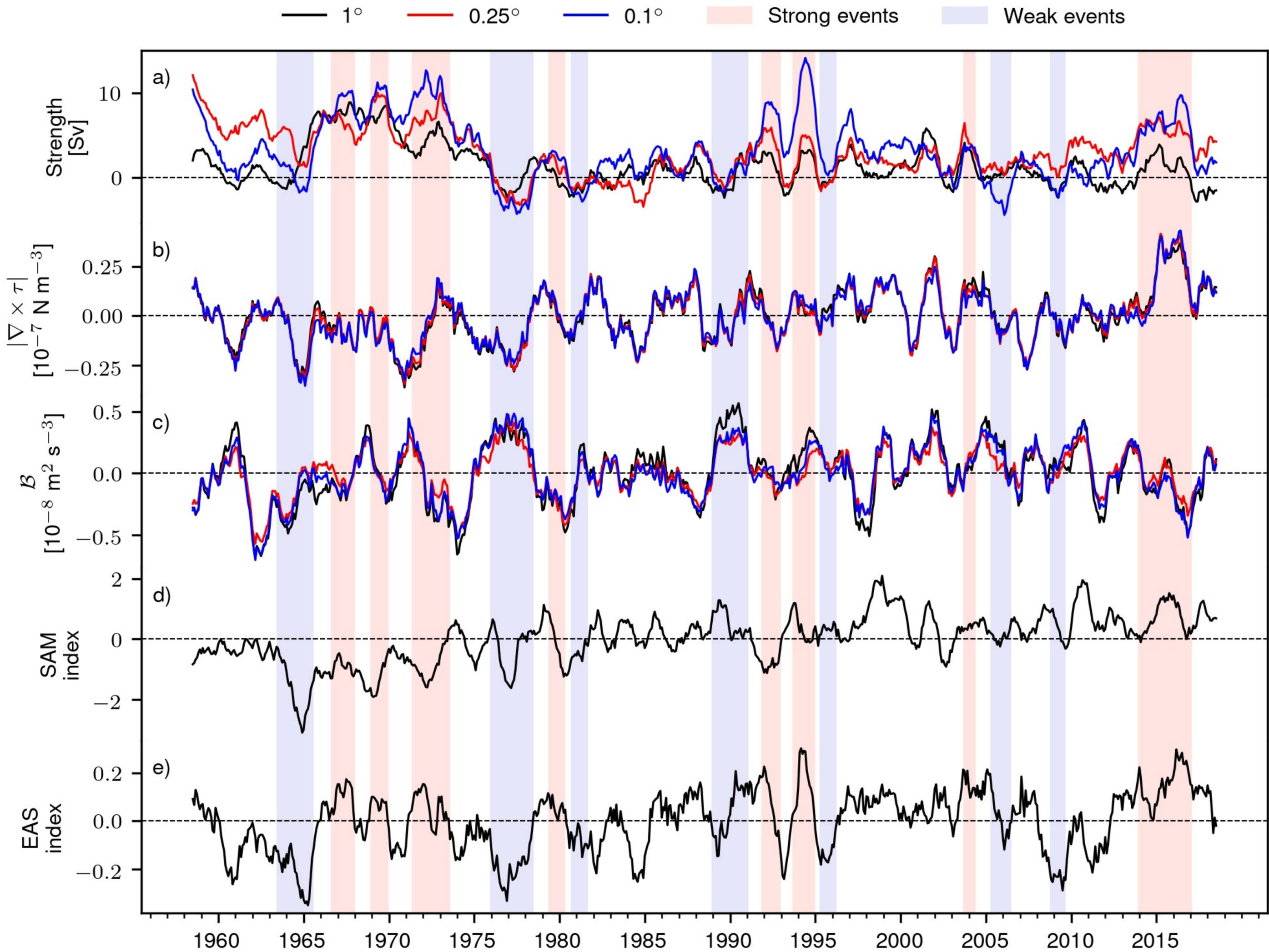
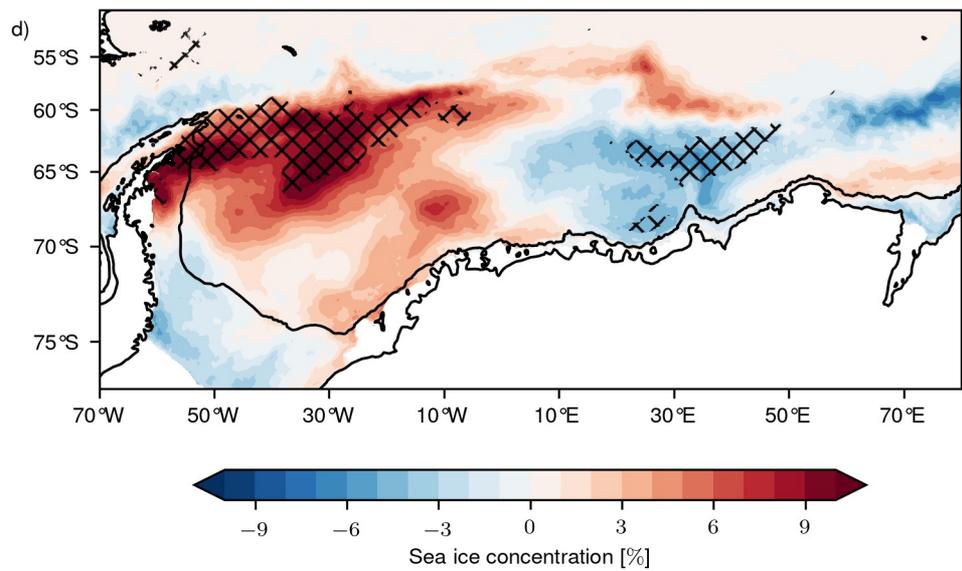
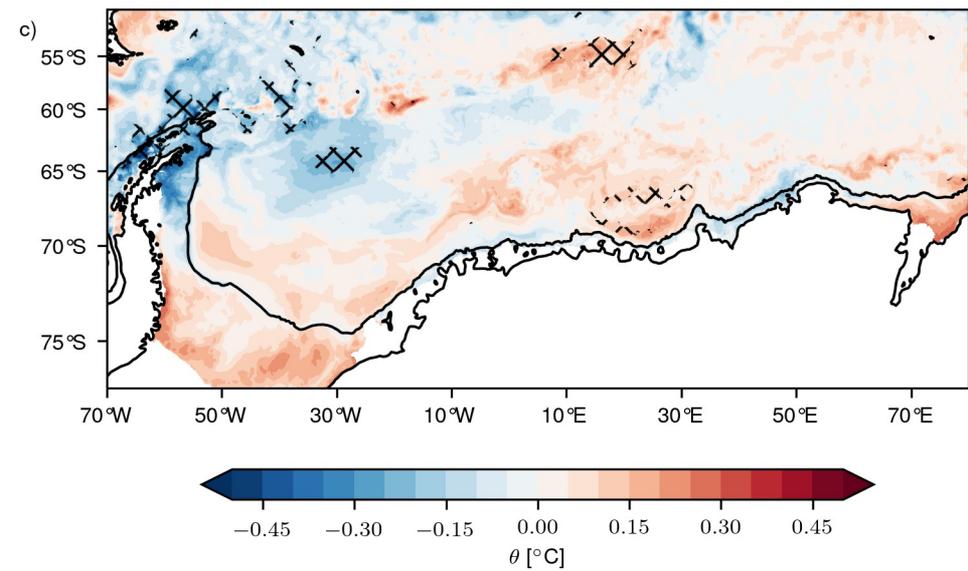
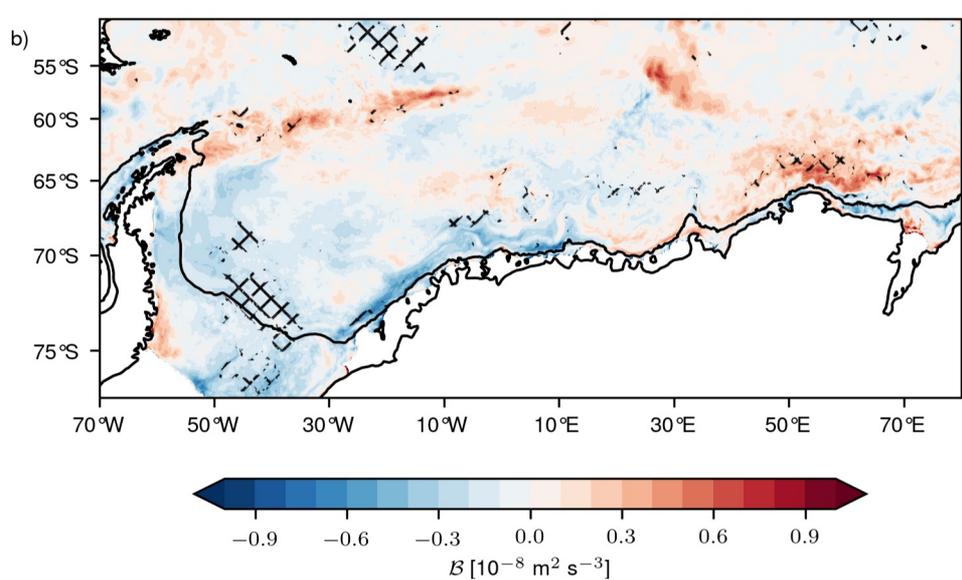
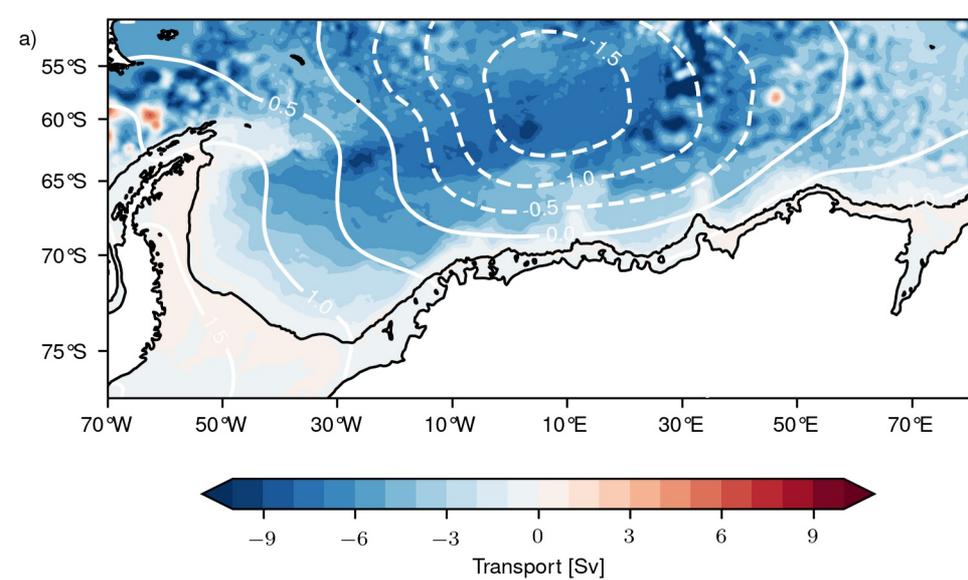


Figure 11.



1 **Seasonal and interannual variability of the Weddell**
2 **Gyre from a high-resolution global ocean-sea ice**
3 **simulation during 1958 - 2018**

4 **J. Neme^{1,2}, M. H. England^{1,2}, A. McC. Hogg³**

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6 New South Wales, Sydney, Australia.

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8 ³Research School of Earth Science and ARC Centre of Excellence for Climate Extremes, Australian
9 National University, Canberra, Australia.

10 **Key Points:**

- 11 • The Weddell Gyre, as simulated by a global coupled ocean-sea ice model, displays
12 large variability on seasonal and interannual time scales.
13 • There is evidence of low frequency (decadal) variability in the Weddell Gyre's strength,
14 but no significant trends during study period.
15 • Years of extreme strong/weak Weddell Gyre flow are shown to be linked to anoma-
16 lous winds, ocean temperatures and sea ice concentration.

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Abstract

The Weddell Gyre’s variability on seasonal and interannual timescales is investigated using an ocean-sea ice model at three different horizontal resolutions. The model is evaluated against available observations to demonstrate that the highest resolution configuration (0.1° in the horizontal) best reproduces observed features of the region. The simulations suggest that the gyre is subject to large variability in its circulation that is not captured by summer-biased or short-term observations. The Weddell Gyre’s seasonal cycle consists of a summer minimum and a winter maximum and accounts for changes that are between one third and a half of its mean transport. On interannual time scales we find that the gyre’s strength is correlated with the local Antarctic easterlies and that extreme events of gyre circulation are associated with changes in the characteristics of the warm inflow at the eastern boundary, that in turn drives changes in sea ice concentration.

Plain Language Summary

The Weddell Gyre, located east of the Antarctic Peninsula, is one of the largest features of the ocean circulation of the Southern Hemisphere. It is adjacent to an important site of bottom water formation, a process that sequesters carbon and heat from the atmosphere and sets the density of the deep ocean, therefore making the region important for global climate. However, extensive sea ice cover throughout the year has historically prevented continuous observations. Several unique features of the gyre, such as open boundaries and intense surface buoyancy fluxes, make the identification of its forcing mechanisms difficult. A deeper understanding of the dynamics in this remote region will shed light on the role of the gyre in our present climate, and help us understand its potential evolution with climate change. We use a high resolution numerical model which shows that the Weddell Gyre undergoes large seasonal and interannual changes. We find that the gyre spins up during winter and slows down during summer, and that strong/weak events in our model simulation are correlated with the strength of the regional easterly winds close to the Antarctic continent. These strong/weak events affect sea ice cover, water mass characteristics and bottom water production.

1 Introduction

The lateral circulation in the Weddell Sea is dominated by the Weddell Gyre, one of the largest features of the ocean circulation south of the Antarctic Circumpolar Current (ACC). The Weddell Gyre is characterized by a broad cyclonic circulation spanning from the Antarctic Peninsula until approximately 30°E (Deacon, 1979), Figure 1. The gyre is one of the southernmost open ocean reaches in the world, with several features that make it a relevant component of the Southern Ocean circulation (Vernet et al., 2019). A strong interaction between ocean and sea ice favours large surface buoyancy fluxes that, in combination with a weak stratification, create a connection between the atmosphere, the ocean surface and the ocean bottom (Abrahamsen et al., 2019; Nicholls et al., 2009; Robertson et al., 2002). For example, located over the continental shelf in the southwestern Weddell Sea, adjacent to the Filchner-Ronne Ice Shelf, is one of the major formation sites of Antarctic Bottom Water (AABW) (Meredith, 2013), a process which drives the deepest limb of the global overturning circulation and involves waters that are circulated and transformed within the gyre (Orsi et al., 2002). Part of the dense waters formed at the continental shelf cascade down the slope and enter the deep circulation of the Weddell Gyre (Foldvik et al., 2004). Despite the Weddell Gyre’s relevance to global climate, the extensive sea ice cover throughout most part of the year has historically hampered long-term, continuous observational efforts and the present knowledge of the gyre’s circulation and variability is mostly limited to the summer months (Vernet et al., 2019).

66 The Weddell Gyre’s circulation has been traditionally associated with the nega-
 67 tive stress curl given by a large scale surface wind field consisting of westerlies to the north,
 68 a circumpolar band of easterly winds surrounding Antarctica and a low pressure system
 69 embedded at around 30°E, 65°S (Gordon et al., 1981; Deacon, 1979). The surface wind
 70 pattern makes the center of the gyre a region of divergence of Ekman transport, char-
 71 acterized by a depression of sea level (Armitage et al., 2018), wind driven Ekman up-
 72 welling and a doming of subsurface isopycnals (Klatt et al., 2005; Schröder & Fahrbach,
 73 1999). The Antarctic continent to the south and west, together with a series of subma-
 74 rine ridges to the north provide a clear topographic barrier to the flow. There is no such
 75 barrier on the eastern boundary of the gyre, which has been suggested to be a dynamic
 76 feature located between 30°E and 60°E (Park & Gambéroni, 1995; Deacon, 1979). It is
 77 at this boundary where relatively warm and saline Circumpolar Deep Water (CDW) en-
 78 ters the gyre, partly following an advective and partly following an eddy-driven path-
 79 way (Ryan et al., 2016; Cisewski et al., 2011; Leach et al., 2011; Schröder & Fahrbach,
 80 1999). The properties of CDW, locally termed Warm Deep Water (WDW), are modi-
 81 fied along its path to the west by mixing and upwelling until some of this WDW crosses
 82 the shelf break at the southern Weddell Sea. On the continental shelf, modified WDW
 83 mixes with High Salinity Shelf Water (HSSW) and Ice Shelf Water (ISW) which are formed
 84 by cooling and brine rejection, to produce Weddell Sea Bottom Water (WSBW) (Nicholls
 85 et al., 2009; Foldvik et al., 2004; Jacobs, 1991; Foster & Carmack, 1976). Some of this
 86 WSBW cascades down the continental slope to fill the abyssal Weddell Sea and mixes
 87 to form a less dense variety termed Weddell Sea Deep Water (WSDW) that is able to
 88 escape the gyre towards the Scotia Sea and becomes AABW (R. A. Locarnini et al., 1993).

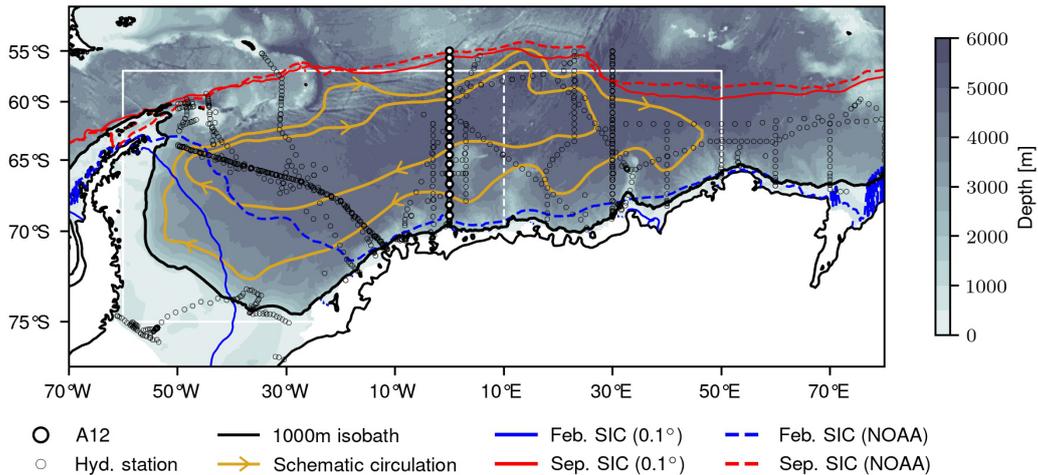


Figure 1. Bathymetry of the Weddell Sea with the 1000m isobath (black contour) and February (blue) and September (red) sea ice extent defined by 15% sea ice concentration from ACCESS-OM2-01 (solid line) and NOAA/NSIDC Climate Data Record of Passive Microwave Sea Ice Concentration, Version 3 (dashed line). The yellow streamlines shown indicate two arbitrary selected contours of the barotropic streamfunction in our 1° simulation that illustrate the schematic circulation of the Weddell Gyre. The white solid box marks the area within which surface stress and buoyancy fluxes are averaged to build a time series, Figures 8 and 10. The Weddell Gyre strength is calculated west of the dashed white line at 10°E.

89 Observational and modelling studies have estimated the Weddell Gyre strength to
 90 be between 30 and 100 Sv, a wide range of values that reflects the strong dependence

91 of gyre strength on the methodology used. Hydrography-based estimates may underes-
92 timate gyre strength because they are biased towards summer conditions, thus exclud-
93 ing a possible winter intensification. Furthermore, the weak stratification that charac-
94 terizes the region means that the assumption of a level of no motion, required to derive
95 transports from the thermal wind equations, may not be appropriate (Cisewski et al.,
96 2011; Park & Gambéroni, 1995; Fahrbach et al., 1991). Numerical models are not sub-
97 ject to such experimental difficulties but their evaluation is limited by the scarcity of ob-
98 servations in the region. The gyre exhibits a persistent double-cell circulation structure
99 whose origin remains unclear, but further contributes to the uncertainties in gyre trans-
100 port (Reeve et al., 2019; Mazloff et al., 2010; Wang & Meredith, 2008; Beckmann et al.,
101 1999).

102 The present knowledge of the variability of the Weddell Gyre, its magnitude, drivers
103 and time scales is greatly limited by the lack of continuous observations and lack of ocean
104 models that simulate the gyre realistically. Repeat hydrographic sections have been used
105 to determine that the properties of the Weddell Gyre’s water masses display significant
106 variations associated with the gyre’s circulation (Fahrbach et al., 2011; Behrendt et al.,
107 2011; Fahrbach et al., 2004). For example, changes in the properties of AABW exported
108 from the Weddell Sea have been linked to wind-forced changes in the gyre’s baroclinic
109 structure (Jullion et al., 2010; Meredith et al., 2011, 2008). Interannual variability in gyre
110 strength has also been related to climate modes, such as the Southern Annular Mode
111 (SAM) and the El Niño - Southern Oscillation (Armitage et al., 2018; Martinson & Ian-
112 nuzzi, 2003). Few studies take into consideration the influence of sea ice and surface buoy-
113 ancy fluxes as possible driving mechanisms, despite their relevance to the dynamics of
114 the region (A. Naveira Garabato et al., 2019; Meneghello et al., 2018).

115 The northern and southern limbs of the gyre have been suggested to vary indepen-
116 dently and not necessarily in phase, responding to different forcing mechanisms (Fahrbach
117 et al., 2011). For example, the northern boundary could be driven by the westerly winds,
118 while the southern limb of the gyre, comprised of a westward flowing current system re-
119 ferred to as the Antarctic Slope Current (ASC), could be partly forced by the easterly
120 winds close to the Antarctic continent. The ASC is a complex current system with a barotropic
121 component driven by the easterly winds along the Antarctic continent that drive an on-
122 shore Ekman transport that raises the sea level towards the coast, driving a westward
123 geostrophic flow that varies according to changes in these winds (Le Paih et al., 2020;
124 A. Naveira Garabato et al., 2019; Armitage et al., 2018; Mathiot et al., 2011; Núñez-Riboni
125 & Fahrbach, 2009). The ASC also comprises an important baroclinic component asso-
126 ciated with the Antarctic Slope Front, characterized in the Weddell Sea by a sloping down
127 of the isopycnals towards the coast due to a fresh continental shelf (Thompson et al., 2018).
128 The ASC and ASF act as a dynamic barrier, mediating the exchange of waters between
129 the continental shelf and the open ocean (Daae et al., 2017; Smedsrud et al., 2006).

130 This study analyzes numerical simulations derived from a coupled numerical model
131 (described in Section 2.1) configured over three different horizontal resolutions (1° , 0.25°
132 and 0.1°) to study the Weddell Gyre’s variability on seasonal to interannual timescales,
133 and its connection to possible forcing mechanisms including wind stress and buoyancy
134 fluxes. In Section 3 the model is evaluated against available observations to assess its
135 strengths and weaknesses in the study region. For this evaluation we use the Armitage
136 et al. (2018) dynamic ocean topography product and temperature and salinity profiles
137 from available hydrographic stations in the region. Section 4 characterizes the Weddell
138 Gyre mean state in model simulations and Section 5 and 6 the gyre’s variability on sea-
139 sonal and interannual timescales respectively. Section 7 summarizes the results and dis-
140 cusses their relevance for our current knowledge of the Weddell Gyre circulation.

2 Model Description, Data and Methods

2.1 Ocean/Sea Ice Model

This study uses the Australian Community Climate and Earth System Simulator (ACCESS-OM2) (Kiss et al., 2020), a global ocean-sea ice model. This model is available at three different horizontal resolutions, namely 1° and 0.25° with 50 vertical levels and 0.1° with 75 vertical levels. The ocean component of the model is the Modular Ocean Model (MOM) version 5.1, developed by the Geophysical Fluid Dynamics Laboratory (<https://mom-ocean.github.io/>) and the sea ice component is the Los Alamos sea ice model (CICE) version 5.1.2 from Los Alamos National Laboratories (<https://github.com/CICE-Consortium/CICE-svn-trunk/tree/cice-5.1.2>). All three configurations are initialized from a rest state with zero sea level, temperature and salinity given by the World Ocean Atlas 2013 v2 0.25° “decav” product (M. Locarnini et al., 2018; Zweng et al., 2019) and are forced by a prescribed atmosphere from JRA55-do reanalysis v1.4 for the period 1958 to 2018 (Tsujino et al., 2018). The model is continuously cycled through this 61-year period, from which we select the third forcing cycle since it is the latest cycle available for all three model resolutions. We further note that the Weddell Gyre takes two full 61-year forcing cycles to properly equilibrate, with longer-term adjustments and drift only stabilizing toward the end of the second cycle. The three model configurations are consistent with each other which allows to infer the importance of resolution in the study region. There is a general improvement at the highest 0.1° configuration when solving several features of the ocean: those relevant for this study are the representation Southern Ocean water masses, the overturning circulation and the characteristics of the circulation on the Antarctic continental shelf and slope (Kiss et al., 2020; Moorman et al., 2020; Morrison et al., 2020). The Rossby radius of deformation in the Weddell Gyre region is less than 10km, so only the 0.1° resolution simulation is close to eddy-resolving (Hallberg, 2013; Chelton et al., 1998).

2.2 Satellite Observations

We evaluate ACCESS-OM2 in the Weddell Gyre region against the dynamic topography of the satellite product developed by Armitage et al. (2018), a unique data set in the Southern Ocean where extensive sea ice cover has historically hampered long term continuous observations using traditional radar altimetry. The product consists of monthly composites of dynamic ocean topography spanning a five year period (2011 to 2016) at 50km horizontal resolution, referenced to the GOCO05c combined gravity field model (Fecher et al., 2017). A limitation of this satellite product is the dependence of its mean dynamic topography on the geoid model, which has larger errors and uncertainties toward the Antarctic continent due to the sparsity of the data being assimilated. These uncertainties do not affect sea level anomalies, i.e. observed variability, because the geoid is time invariant. The product also displays some north-south striping artifacts generated by the satellite’s orbital precession (see Fig. 5 of Armitage et al. (2018)). To compare ACCESS-OM2 sea level with satellite observations, we interpolated the model to the satellite product’s grid and applied an offset to remove the spatially uniform footprint of geoid uncertainties in the satellite product. The offset is defined as the average difference between the mean sea level fields of each model simulation and observations for the entire region covered by the satellite product.

2.3 Hydrographic Data

The model’s temperature and salinity is evaluated against hydrographic observations in the Weddell Gyre region downloaded from the Clivar and Carbon Hydrographic Data Office (CCHDO, <https://cchdo.ucsd.edu/>), a total number of 1576 CTD profiles from 20 oceanographic cruises in the region. Table 1 has information on the hydrographic cruises and their locations are shown in Figure 1. To compare against ACCESS-

Table 1. Summary of the hydrographic cruises used to evaluate the model. Cruise locations are indicated in Figure 1. The expocode is the identifier for the data sets archived at the Carbon and Climate Hydrographic Data Office (CCHDO; <http://cchdo.ucsd.edu>), the line identifies the World Ocean Circulation Experiment (WOCE) section of the cruise, and start and end dates with corresponding principal investigator

Expocode	Line	Start date	End date	Principal Investigator
06AQANTVIII_2	SR02, SR04	1989-09	1989-10	Eberhard Fahrbach
06AQANTIX_2	SR04	1990-11	1990-12	Eberhard Fahrbach
06AQANTX_4	A12, SR04	1992-05	1992-08	Peter Lemke
35MFCIVA_1	I06S	1993-01	1993-03	Alain Poisson
74DI200_1	S04	1993-02	1993-03	Robert R. Dickson
74JC10_1	A23	1995-03	1995-05	Karen J. Heywood and Brian A. King
320696_3	S04, S04I	1996-03	1996-04	Thomas Whitworth
35MF103_1	I06S	1996-02	1996-03	Alain Poisson
06AQANTXIII_4	S04A, SR04	1996-03	1996-05	Eberhard Fahrbach
06AQANTXV_4	SR04	1998-03	1998-05	Eberhard Fahrbach
06AQ199901_2	A12	1999-01	1999-03	Eberhard Fahrbach
06AQ200012_3	A12	2000-12	2001-01	Eberhard Fahrbach
06AQ200211_2	A12	2002-11	2003-01	D.K. Fütterer
06AQ20050102	A12	2005-01	2005-04	Eberhard Fahrbach
09AR20060102	S04I	2006-01	2006-03	Mark Rosenberg
06AQ20071128	A12	2007-11	2008-02	Ulrich Bathmann
33RR20080204	I06S	2008-02	2008-03	Kevin Speer
06AQ20080210	A12	2008-02	2008-04	Eberhard Fahrbach
06AQ20101128	A12	2010-11	2011-02	Eberhard Fahrbach
06AQ20141202	A12	2014-12	2015-02	Olaf Boebel

191 OM2, we select the nearest grid point and corresponding monthly composite within the
 192 model output and we interpolate it to the same vertical levels. With this approach we
 193 obtain synthetic profiles within the ACCESS-OM2 model simulations to evaluate the temperature-
 194 salinity structure of the gyre. To give further insight into the gyre’s vertical structure,
 195 we selected one of the most repeated WOCE transects in the region, A12 nominally at
 196 the Greenwich Meridian (see Figure 1), to calculate an average potential density cross-
 197 section. Analogous to the procedure for individual hydrographic stations, we select the
 198 monthly composites of model data that correspond to the repeat A12 cruises, interpo-
 199 late the vertical level, and then calculate the anomalies of the model with respect to the
 200 observations.

201 2.4 Barotropic Streamfunction

202 The barotropic streamfunction, ψ , is used to study the gyre’s transport and is de-
 203 fined as the meridional integral of the depth-integrated zonal mass transport M_x :

$$\psi(x, y) = \int_{y_0}^y \frac{M_x(x, y')}{\rho(y - y_0)} dy' \quad (1)$$

204 where we take ρ as an average density of 1035 kg m^{-3} and the integration goes from
 205 south to north starting at the Antarctic continent (y_0). By this definition, ψ takes neg-
 206 ative values for the cyclonic circulation of the Weddell Gyre. The gyre strength (GS) is

207 derived from ψ by calculating its minimum in the region bounded by the white box in
 208 Figure 1 west of 10°E ($[75^\circ S : 57^\circ S] \times [60^\circ W : 10^\circ E]$) and taking the absolute value:

$$GS = |\min\{\psi\}| \quad (2)$$

209 The gyre strength definition is limited to 10°E to exclude the unstable, eddy-rich flow
 210 of the eastern region, but nonetheless encompasses the main structure of the mean gyre
 211 (see Section 3.2). The barotropic streamfunction is also used to define the boundary of
 212 the Weddell Gyre as the 12 Sv contour. The choice of a specific streamfunction contour,
 213 while somewhat arbitrary, allows us to track changes in the gyre’s area over time in or-
 214 der to identify expansions or contractions of its boundary.

215 2.5 Surface Stress and Buoyancy Fluxes

216 To assess the role of stress over the ocean’s surface as a possible driver of the cir-
 217 culation, we calculate the total surface stress from model output taking into account the
 218 relative contributions of air/ocean and ice/ocean stresses ($\tau_{air/ocean}$ and $\tau_{ice/ocean}$ re-
 219 spectively) weighted by sea ice concentration. The curl of the total surface stress, $\nabla \times$
 220 τ , is defined as:

$$\nabla \times \tau = \nabla \times ((1 - \alpha_{ice})\tau_{air/ocean}) + \nabla \times (\alpha_{ice}\tau_{ice/ocean}) \quad (3)$$

221 where α_{ice} is sea ice concentration. To build a time series, $\nabla \times \tau$ is averaged in the re-
 222 gion bounded by the white box in Figure 1 excluding grid cells in which the bathymetry
 223 is shallower than 1000m. We show the absolute value of the surface stress curl so that
 224 stronger curl anomalies are represented by positive values.

225 Surface buoyancy fluxes are also considered as a possible driver of the gyre’s vari-
 226 ability, taking into account contributions due to both heat and freshwater exchanges. The
 227 surface buoyancy flux is defined as:

$$\mathcal{B} = \frac{g\alpha Q}{c_w \rho} + g\beta F_w s \quad (4)$$

228 where ρ is surface density, g is the acceleration due to gravity, c_w the specific heat of sea
 229 water, α the thermal expansion coefficient, β the saline contraction coefficient, s surface
 230 absolute salinity and Q and F_w are the heat and freshwater fluxes, positive represent-
 231 ing a buoyancy gain by the ocean. The Python implementation of the Gibbs Sea Wa-
 232 ter Oceanographic Toolbox of TEOS-10 (<https://teos-10.github.io/GSW-Python/>)
 233 was used to compute the quantities derived from conservative temperature and absolute
 234 salinity, such as α , β and ρ . The largest buoyancy fluxes in the Weddell Sea region are
 235 found over the continental shelf, so \mathcal{B} is averaged at depths shallower than 1000m to build
 236 a time series.

237 2.6 Climate Indices

238 We use the Southern Annular Mode (SAM) index as a proxy for the meridional ex-
 239 pansion/contraction of the band of westerlies surrounding the Antarctica to evaluate the
 240 connection between the gyre and large-scale atmospheric modes of variability. This in-
 241 dex is calculated as the difference between the zonal anomalies of sea level pressure at
 242 40°S and 65°S, (Marshall, 2003), from the JRA55 reanalysis, the atmospheric data set
 243 used to force the model. To verify that local changes in the westerlies are not significantly
 244 different from the circumpolar average we re-evaluate the SAM index in a limited do-
 245 main, calculating the zonal anomalies with respect to the mean of the Weddell Sea and

246 find no significant differences. By analogy we define a different index to track the inten-
 247 sification/weakening of the local easterly winds, hereafter denoted EAS, defined as the
 248 difference in the zonal anomalies of sea level pressure at 65°S and 72°S. Contrary to the
 249 SAM index, since the easterlies display a large zonal variation, the EAS is calculated be-
 250 tween 30°W and 70°E with the intention of focusing only on regional changes.

251 2.7 Analysis Period

252 Long term climatological fields are calculated for the period 1958 to 2018 for ψ , $\nabla \times$
 253 τ and \mathcal{B} , from which we calculate the annual cycles of gyre strength, surface stress curl
 254 and buoyancy fluxes as described in the previous sections. To study interannual variabil-
 255 ity, these metrics together with the SAM and EAS indices are calculated using monthly
 256 composite fields for the full 61-year period, and subsequently the annual cycle is removed
 257 to obtain anomalies. We also apply a 12-month running mean to filter out high frequency
 258 variability. The 0.1° experiment is used to define strong and weak events of Weddell Gyre
 259 flow as periods longer than 6 months during which gyre strength is beyond $\pm 0.8SD$ of
 260 the gyre strength time series smoothed with a 10-year running mean. Once the events
 261 are identified we composite the streamfunction, sea level pressure, surface buoyancy fluxes,
 262 subsurface temperature maximum and sea ice concentration anomaly fields to charac-
 263 terize the differences in these variables between anomalously strong and weak gyre pe-
 264 riods. Additionally, we calculate the linear correlation coefficients between the gyre strength
 265 time series and the anomaly fields of surface buoyancy fluxes, subsurface temperature
 266 maximum and sea ice concentration in order to identify regions where they are signif-
 267 icantly correlated to, and thus likely affected by, gyre strength.

268 3 Model Evaluation

269 Satellite observations (Armitage et al., 2018) show the Weddell Gyre as an elon-
 270 gated depression of sea level, indicative of a cyclonic geostrophic current, with its main
 271 axis oriented in the northeast-southwest direction, Figure 2a. Figures 2b, c and d show
 272 the difference between the model and satellite observed mean dynamic topography for
 273 the 1°, 0.25° and 0.1° configurations respectively. If this difference were spatially con-
 274 stant, then the derived surface geostrophic circulation would be the same since the ve-
 275 locity in between two grid points is proportional to the difference in elevation between
 276 those two points. The 0.1° simulation exhibits the smallest horizontal gradients in the
 277 sea level difference with observations, not only within the gyre’s boundary, but also in
 278 the southwestern continental shelf, the eastern region and within the ACC, which indi-
 279 cates that the highest resolution configuration better represents the observed Weddell
 280 Gyre geostrophic flow, Figure 2d.

281 Sea level variability in the model simulations is also evaluated using the Armitage
 282 et al. (2018) data set. We compute the correlation coefficient between the satellite and
 283 model sea level at every grid point for the period 2011 to 2016 and use them to build
 284 correlation maps, Figures 3b, c and d with significant correlations hatched. A north-south
 285 striping pattern is visible in the correlation coefficients that reflects an artifact from the
 286 satellite’s orbit. Taking this artifact into consideration, we infer that within the gyre,
 287 model and observations are significantly correlated, with no large differences between model
 288 configurations. However, the correspondence between model and observations breaks down
 289 within the ACC, which we attribute to the presence of eddies that the model is not ex-
 290 pected to reproduce at the precise location and time as they appear in observations. Un-
 291 like the comparison between model and observations within the gyre, in the southwest-
 292 ern continental shelf the ability of the model to reproduce observed variability improves
 293 significantly with resolution, with higher, significant correlations for the 0.1° configura-
 294 tion. However, the correlation maps do not fully illustrate differences in the model’s vari-
 295 ability with respect to observations, which is why we calculate the root mean square er-

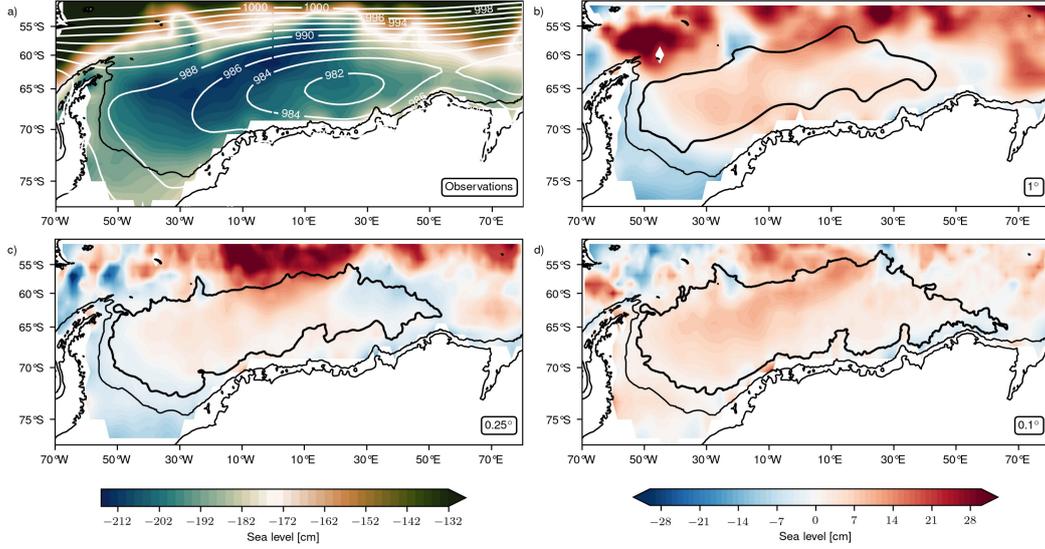


Figure 2. (a) Mean sea level from satellite observations (Armitage et al., 2018) with mean sea level pressure from JRA55-do v1.3 for the period 2011 to 2016. Model’s mean sea level for the same period minus satellite observations for the (b) 1° , (c) 0.25° and (d) 0.1° resolution configurations, with gyre’s mean boundary (defined as the 12 Sv barotropic streamfunction contour) in thick black. Thin black contour in all panels marks the 1000m isobath.

296 ror, RMSE, within the gyre between 30°W and 30°E , 70°S and 60°S over the observa-
 297 tional record, Figure 3a. We observe that the RMSE within the Weddell Gyre is in phase
 298 between model configurations and of similar magnitude, indicating that the departure
 299 from observations is consistent across resolution and therefore likely due to errors in model
 300 forcing compared to observations. There is also no seasonality in the error (e.g. the er-
 301 ror is not consistently larger in winter), which further suggests that the error is due to
 302 biases in the model’s forcing rather than biases in satellite observations associated with
 303 sea ice cover.

304 Apart from evaluating the model’s sea level, we can evaluate the model’s temper-
 305 ature and salinity structure against observations. Figure 4 shows modeled and observed
 306 temperature-salinity diagrams comparing hydrographic stations in the Weddell Gyre re-
 307 gion at locations indicated in Figure 1. It is possible to identify different water masses
 308 in observations (Figure 4a) according to their potential temperature and salinity: namely
 309 warm and saline CDW, colder WDW and a distinct tail of colder WSDW and WSBW.
 310 There are also some stations on the continental shelf that indicate the presence of ISW
 311 with temperatures lower than the surface freezing point (-1.9°C), which is formed through
 312 the interaction with the ice shelf base. The ability of the model to reproduce the distinct
 313 characteristics of these water masses is highly dependent on resolution. The coarser 1°
 314 resolution does not show the presence of CDW in the region, has a colder, fresher than
 315 observed WDW and the characteristics of its bottom waters are less distinct from shall-
 316 lower waters than observed. There is a slight improvement in the 0.25° configuration which
 317 shows a trace of CDW and a better reproduction of WDW characteristics. The improve-
 318 ment is clearest for the 0.1° resolution configuration which captures the presence of CDW
 319 in the region, although warmer than observed, as well as achieving a better representa-
 320 tion of WDW than the two coarser resolutions. The 0.1° simulation also shows some trace
 321 of the bottom waters, albeit less distinct than observed, which could be related to the
 322 coarse vertical resolution of the model relative to the station observations (50 levels for

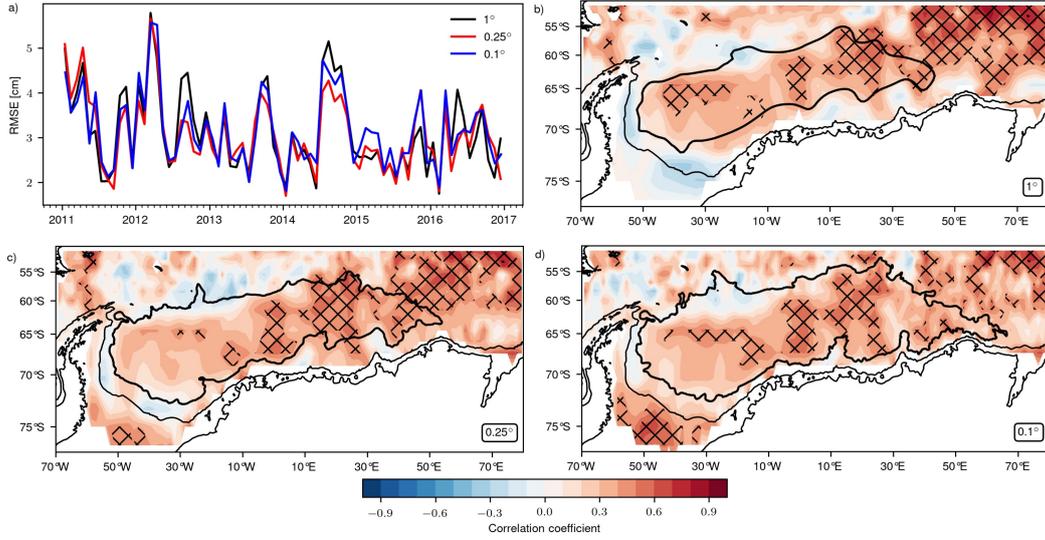


Figure 3. (a) RMSE between model and observations calculated for a region encompassing the gyre [30°W, 30°E]x[70°S, 60°S]. Correlation coefficients of sea surface height with satellite observations (Armitage et al., 2018) for models at (b) 1°, (c) 0.25° and (d) 0.1° resolution with hatching for significant correlations ($p < 0.05$). Mean gyre’s boundary in thick black contour and 1000m isobath in thin black contour.

323 the 1° and 0.25° and 75 levels for the 0.1° configuration). Moreover, none of the model
 324 configurations capture the presence of supercold ISW because the model lacks the ice
 325 shelf cavities where this water mass is formed. Since ISW participates in deep and bot-
 326 tom water production, the model’s misrepresentation of these dense waters could be re-
 327 lated to its inability to generate ISW.

328 The average cross section of potential density at the Greenwich Meridian shows the
 329 doming of isopycnals that characterizes the center of the gyre roughly south of 60°S and
 330 the influence of Maud Rise, north of 67°S, on top of which sits a cold, fresh Taylor col-
 331 umn, Figure 5. At around 58°S the sloping of the isopycnals shows the Weddell Front
 332 that marks the transition from Weddell Gyre to the ACC. The differences between model
 333 and observations show that the largest discrepancies are located in the upper 250m of
 334 the water column, Figures 5b, c and d. The most striking difference is the clear contrast
 335 between Weddell Gyre waters (denser than observed) and ACC waters (lighter than ob-
 336 served) for the 1° and 0.25° configurations, which means that for these resolutions the
 337 oceanic front at the northern boundary of the gyre is steeper than observed. This Wed-
 338 dell Gyre/ACC contrast is not apparent in the 0.1° resolution. On the other hand, all
 339 three configurations show anomalies of alternate sign within the gyre in a thin surface
 340 layer shallower than 50m and the underlying subsurface layers. Since the cross sections
 341 are representative of summer conditions, there is a shallow layer of light waters associ-
 342 ated with sea ice melt that sits on top of the Winter Water (WW), so-called because it
 343 surfaces during winter when there is a weakening of the summer pycnocline associated
 344 with ice production (Fahrbach et al., 2004). The model cross sections show for all res-
 345 olutions a lighter than observed summer surface layer and a denser than observed win-
 346 ter subsurface layer, which means that the stratification is larger than observed in the
 347 model during the summer months. Below the upper 250m, the 0.1° configuration exhibits
 348 an improved representation of the observed potential density relative to the 1° and 0.25°
 349 simulations.

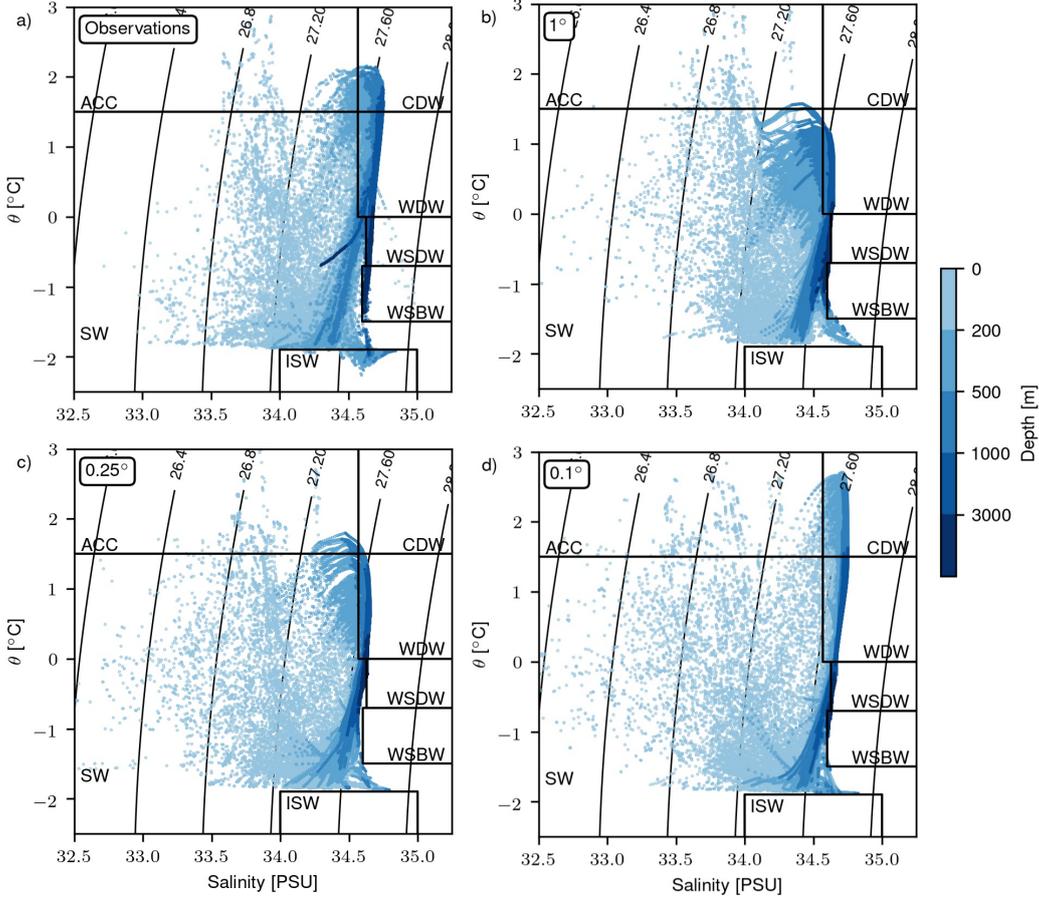


Figure 4. Temperature-salinity diagrams with potential density isolines for all stations in Figure 1 for (a) hydrographic observations and for model runs at (b) 1° , (c) 0.25° and (d) 0.1° resolution, color coded by their depth. Antarctic Circumpolar Current = ACC, Circumpolar Deep Water = CDW, Warm Deep Water = WDW, Weddell Sea Deep Water = WSDW, Weddell Sea Bottom Water = WSBW, Surface Water = SW from Schröder and Fahrback (1999), Ice Shelf Water = ISW from Foldvik et al. (2004).

350 On the continental slope, observations show the ASF with sloping isopycnals that
 351 separate lighter waters on the continental shelf from denser waters off-shore. The 1°
 352 simulation does not resolve the continental slope at this longitude, and is therefore not able
 353 to properly resolve the ASC associated with the front. On the other hand, the 0.25° and
 354 0.1° simulations are able to resolve the continental slope, and both show denser than ob-
 355 served waters below the upper 50m of the water column, less so for the 0.1° simulation.
 356 This means that the downward sloping isopycnals in the model are less steep than ob-
 357 served, and consequently the model underestimates the baroclinic transport of the ASC
 358 at this location.

359 4 Weddell Gyre Mean State

360 The mean barotropic streamfunction for the entire 61 year period of the model run
 361 (1958 to 2018) shows an elongated cyclonic gyre with mean strengths of 34, 33 and 41
 362 Sv and a seasonal range of 9, 14 and 15 Sv for the 1° , 0.25° and 0.1° resolutions respec-

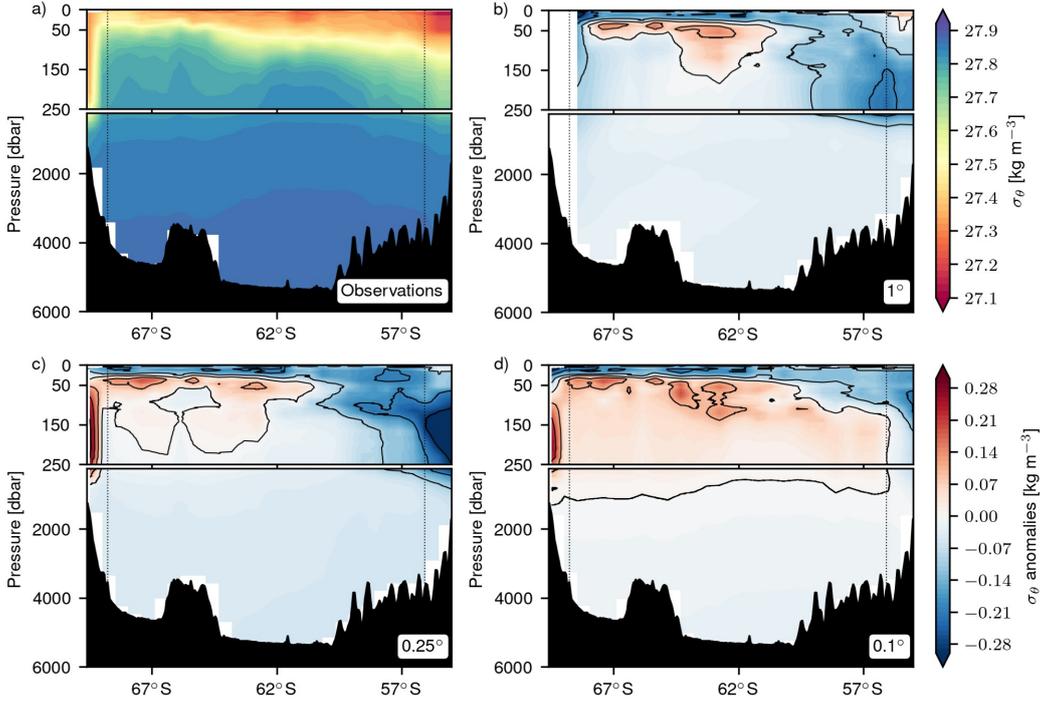


Figure 5. Average potential density (kg m^{-3}) section for WOCE's A12 transect shown in Figure 1 for (a) hydrographic observations and the difference with observations for (b) 1° , (c) 0.25° and (d) 0.1° resolutions with 0.1 kg m^{-3} spacing between contours in black. Vertical dotted lines mark the gyre's boundary at the Greenwich Meridian from the 0.1° simulation for reference.

363 tively, Figures 6a, c and e. The magnitude and spatial pattern of the Weddell Gyre as
 364 as depicted by the model is in agreement with both observational and past modeling stud-
 365 ies. All resolutions display a double lobed circulation structure, with a larger western
 366 cell and a recirculation close to the Greenwich Meridian. As the resolution increases, the
 367 region east of approximately 10°E becomes increasingly unstable and the cyclonic cir-
 368 culation expands towards the east along the coastline.

369 The barotropic circulation generally follows potential vorticity contours (approx-
 370 imated by f/H) at the south, west and northern boundaries; these contours are largely
 371 controlled by topographic features, indicating that the motion is steered by bathymetry.
 372 The larger spacing between f/H contours in the southern limb is associated with a broad,
 373 slow circulation that contrasts with a more intense northern boundary flow. At the north-
 374 eastern boundary the flow crosses potential vorticity contours, where there is an abrupt
 375 deflection of f/H contours owing to a gap in the Southwestern Indian Ridge at around
 376 30°E , 55°S . The crossing of f/H contours indicates that the flow experiences changes
 377 in its potential vorticity due to its interaction with topography, an interaction that is
 378 highly dependent on resolution. The deflection of f/H contours, together with the strong
 379 flow of the ACC, give rise to instabilities at the tail of the ridge particularly at 0.1° res-
 380 olution, that leave a trace in the mean circulation (Figure 6e), as highlighted by the large
 381 values in the standard deviation of ψ (Figure 6f).

382 The subsurface temperature maximum can be used to trace the path CDW takes
 383 after entering the gyre at its eastern boundary. Observations show this warm inflow and
 384 the subsequent lowering of its temperature due to mixing as CDW flows to the west (see
 385 Figure 7a using Reeve et al. (2019) gridded data set as an example). The 0.1° simula-

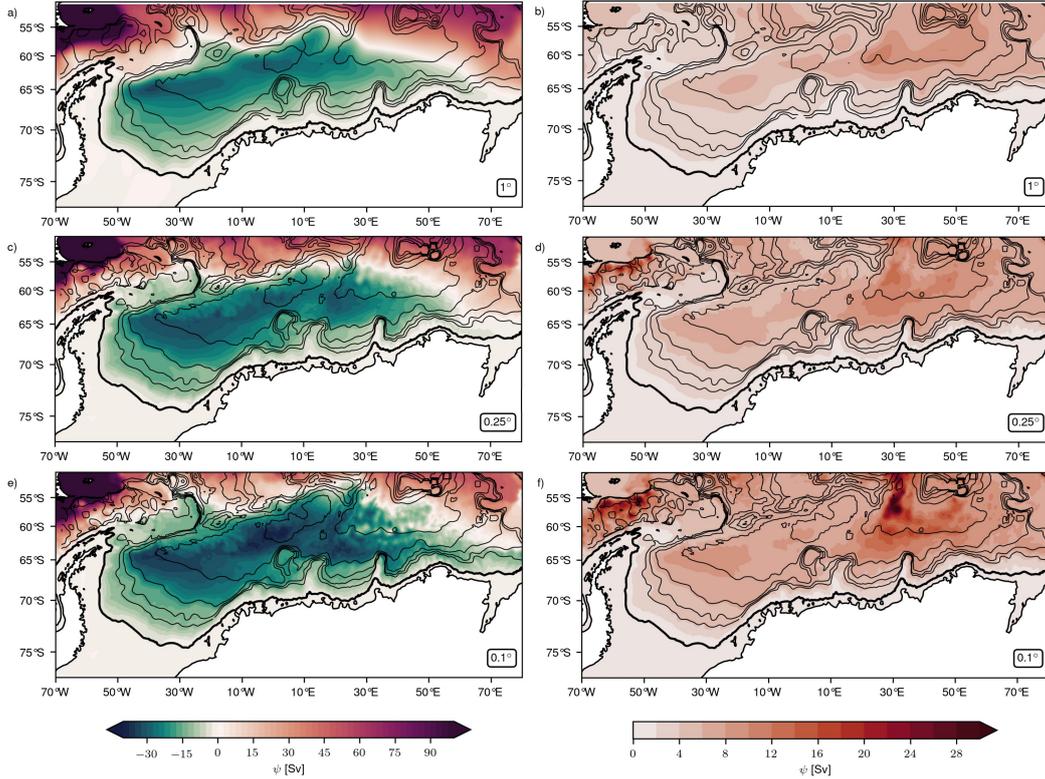


Figure 6. Mean barotropic streamfunction (Sv, $1 \text{ Sv} = 10^6 \text{ m}^3/\text{s}$) for the period 1958 - 2018 for (a) 1° , (c) 0.25° and (e) 0.1° and its standard deviation for the (b) 1° , (d) 0.25° and (f) 0.1° resolution models with contours of potential vorticity (f/H) and 1000m isobath indicated by the thin and thick black contours respectively.

386 tion is the only configuration that displays such warm inflow. The improved accuracy
 387 of the 0.1° model in reproducing this feature of the gyre could be crucial, since the in-
 388 flow affects the characteristics of bottom waters formed at the continental shelf (Couldrey
 389 et al., 2013; Jullion et al., 2014; Fahrback et al., 2011).

390 5 Weddell Gyre Seasonal Cycle

391 The three model simulations show a seasonal cycle of gyre strength consisting of
 392 a winter intensification and a summer weakening that increases in amplitude with res-
 393 olution and accounts for 36%, 46% and 44% of the variability in gyre strength for the
 394 1° , 0.25° and 0.1° configurations respectively, Figure 8a. There is a clear increase in trans-
 395 port from the 1° and 0.25° resolutions to the 0.1° resolution. The seasonal cycle of sur-
 396 face stress curl is strongly modulated by changes in sea ice concentration (α_{ice} in Equa-
 397 tion 3): during the summer months, the minimum in sea ice concentration makes wind
 398 stress the primary contributor to the total stress curl. As the ice pack begins to develop,
 399 after exceeding 50% concentration on average over the region in May, it damps the trans-
 400 fer of momentum from the atmosphere to the ocean, generating a winter minimum in
 401 total stress curl relative to the autumn peak. After reaching its maximum during August-
 402 September, α_{ice} begins to decrease, allowing for a second maximum in total stress curl
 403 during spring. The seasonal cycle of surface buoyancy fluxes is also strongly modulated
 404 by the presence of sea ice, Figure 8c. During the summer, months sea ice melt, along with

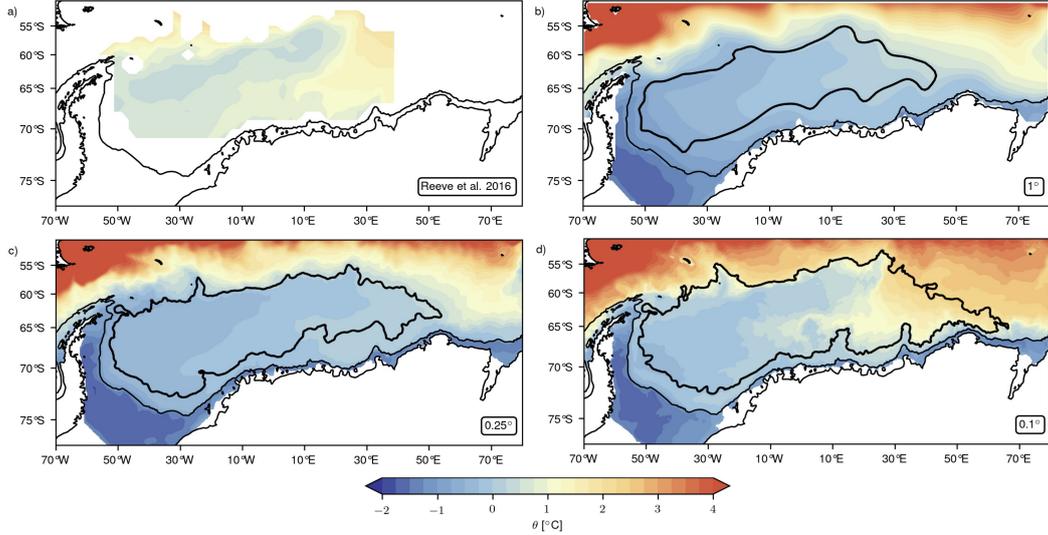


Figure 7. Mean subsurface potential temperature maximum for the period 2001 to 2014 from (a) the Reeve et al. (2016) product and for the period 1958 - 2018 from the (b) 1° , (c) 0.25° and (d) 0.1° resolution simulations. The thick black contour marks the gyre's boundary and the thin black contour the 1000m isobath.

405 warmer air temperatures, result in a net surface buoyancy gain, with the opposite oc-
 406 ccurring during winter via sea-ice formation, brine rejection, and atmospheric cooling.

407 Correlations between gyre strength, surface stress curl and surface buoyancy fluxes
 408 are shown in Table 2. Gyre strength is significantly correlated with both forcing mech-
 409 anisms considered, indicating that buoyancy losses and surface stress curl intensification
 410 are concurrent with a stronger gyre, with the exception of the 0.1° resolution simu-
 411 lation whose annual cycle of gyre strength is not correlated with the semi-annual cycle of sur-
 412 face stress curl. There is a clear increase in transport from the 1° and 0.25° simu-
 413 lations to the 0.1° that is not explained by differences between simulations in either surface stress
 414 curl or surface buoyancy fluxes. Moreover, the surface stress curl is weakest for the 0.1°
 415 simulation and there are no appreciable differences in the seasonal cycle of surface buoy-
 416 ancy fluxes, Figures 8b and c.

417 There are spatial differences in the amplitude of seasonal variability. The summer/winter
 418 anomalies with respect to the mean for the barotropic streamfunction show that the largest
 419 seasonal variations are located in two separate regions around 30°W and 30°E respec-
 420 tively, Figure 9a and c. The first region is located adjacent to a region of significant buoy-
 421 ancy gain during summer and the latter is located beneath the regional low pressure sys-
 422 tem which deepens and expands during the winter months and is shallower during sum-
 423 mer, modifying the local wind pattern, Figures 9b and d. The gyre boundary, defined
 424 by the 12 Sv contour, shows a general expansion during winter in all directions, with a
 425 particularly large excursion towards the east.

426 6 Weddell Gyre Interannual Variability

427 In order to explore the interannual variability of the Weddell Gyre and its connec-
 428 tion to regional and large scale climate, we analyse the full 61-year period time series of
 429 anomalies with respect to the seasonal cycle of buoyancy fluxes, surface stress curl, gyre
 430 strength and the SAM and EAS index calculated as described in Section 2, Figure 10.

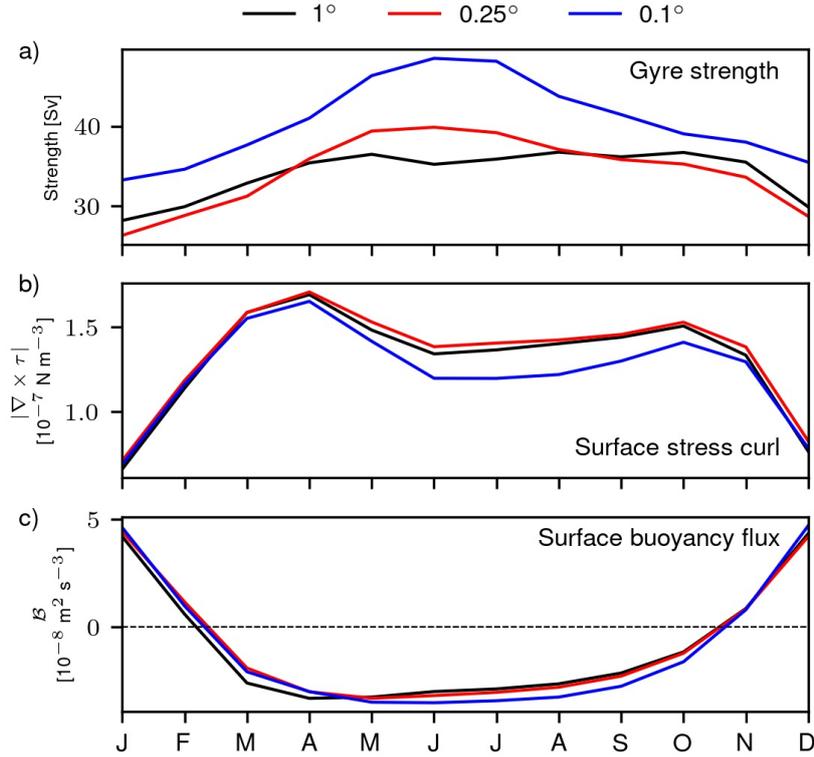


Figure 8. Annual cycles of (a) gyre strength (Sv), (b) surface stress curl (10^{-7} Nm^{-3}) averaged over the region bounded by the solid white box in Figure 1 and (c) surface buoyancy flux ($10^{-8} \text{ m}^2 \text{ s}^{-3}$) averaged in the same regions as the surface stress curl for depths shallower than 1000 m.

Table 2. Correlation coefficients between gyre strength and surface stress curl, buoyancy flux, SAM index and EAS index for seasonal and interannual time scales. Bold values indicate significant correlations with $p < 0.05$.

Seasonal					
	Model resolution	$ \nabla \times \tau $	\mathcal{B}		
Gyre strength	1°	0.82	-0.81		
	0.25°	0.71	-0.89		
	0.1°	0.38	-0.83		
Interannual					
	Model resolution	$ \nabla \times \tau $	\mathcal{B}	SAM	EAS
Gyre strength	1°	0.04	-0.11	-0.36	0.26
	0.25°	0.21	-0.32	-0.24	0.34
	0.1°	0.21	-0.29	-0.19	0.51

431 Gyre strength displays significant interannual variability that becomes greater in mag-
 432 nitude with resolution, with standard deviations of 4.5, 5 and 6 Sv for the 1°, 0.25° and

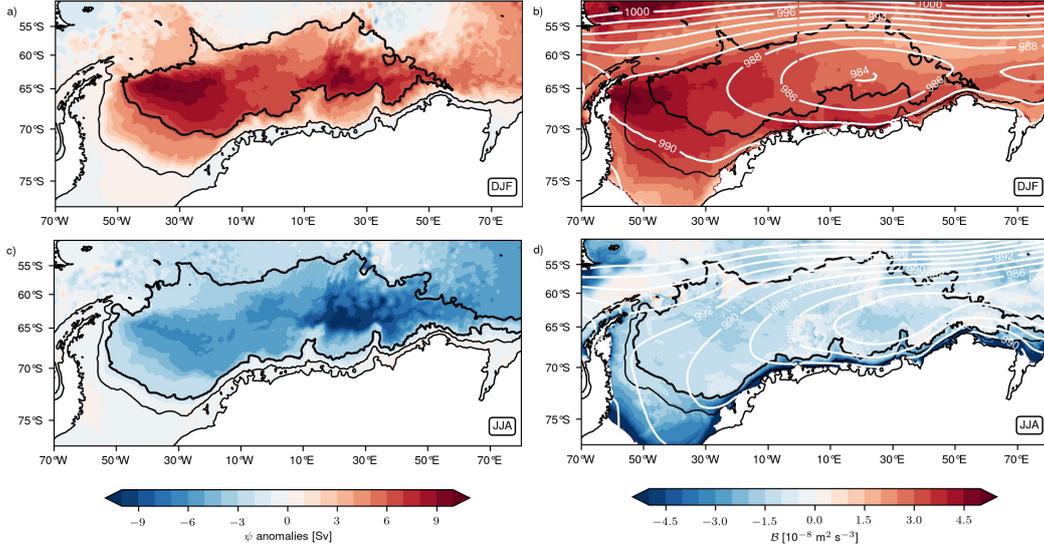


Figure 9. (a) Summer and (c) winter climatological anomalies of ψ and (b) summer and (d) winter climatological fields of net surface buoyancy fluxes with and sea level pressure (white contours) for the 0.1° resolution case calculated for the period 1958 - 2018. Thin black contour indicates the 1000m isobath and thick black contour the Weddell Gyre's boundary in the corresponding season.

433 0.1° configurations respectively. There are some years which show an intensification of
 434 around 10 Sv, which is between a third and a quarter of the gyre's mean strength de-
 435 pending on the resolution. The strengths of the three configurations are mostly in phase
 436 and significantly correlated with each other, indicating the predominance of external forc-
 437 ing in driving the gyre's interannual variability. Similar to what was observed for the sea-
 438 sonal cycles, there is a clear influence of model resolution in setting the magnitude of gyre
 439 strength and the increase in magnitude from the 1° and 0.25° simulations to the 0.1°
 440 is not related to differences between resolutions in surface stress curl or buoyancy fluxes.
 441 Correlations between gyre strength and the other time series in Figure 10 are shown in
 442 Table 2. From the climate metrics considered here, only EAS shows a significant correla-
 443 tion with gyre strength for the 0.1° resolution configuration.

444 Next we identify strong events and weak events of gyre strength in the 0.1° reso-
 445 lution simulation, highlighted in red and blue shading respectively in Figure 10. There
 446 is no periodicity or seasonality in the occurrence of these events, i.e. extreme strong and
 447 weak events do not have a preference for particular months or seasons. Figure 11a shows
 448 the composites of strong events minus the composites of weak events for the barotropic
 449 streamfunction and sea level pressure. Since the cyclonic circulation of the Weddell Gyre
 450 is represented by $\psi < 0$ (Figure 6), negative anomalies indicate an intensified circula-
 451 tion. Sea level pressure shows a deepening and displacement to the west of a low pres-
 452 sure system that deepens the trough in the southern portion of the gyre, thus intensify-
 453 ing the local easterly winds. The composites of surface buoyancy flux show that during
 454 strong years, the southwestern region of the gyre experiences buoyancy loss in a re-
 455 gion where buoyancy fluxes are significantly correlated with gyre strength, as marked
 456 by the hatching, Figure 11b. These changes in buoyancy fluxes are consistent with the
 457 composite for sea ice concentration, Figure 11d: i.e., a stronger circulation advects sea
 458 ice from the southwestern region, decreasing its concentration during strong years, ex-
 459 posing the ocean's surface to the atmosphere, thus enabling a larger buoyancy loss. The

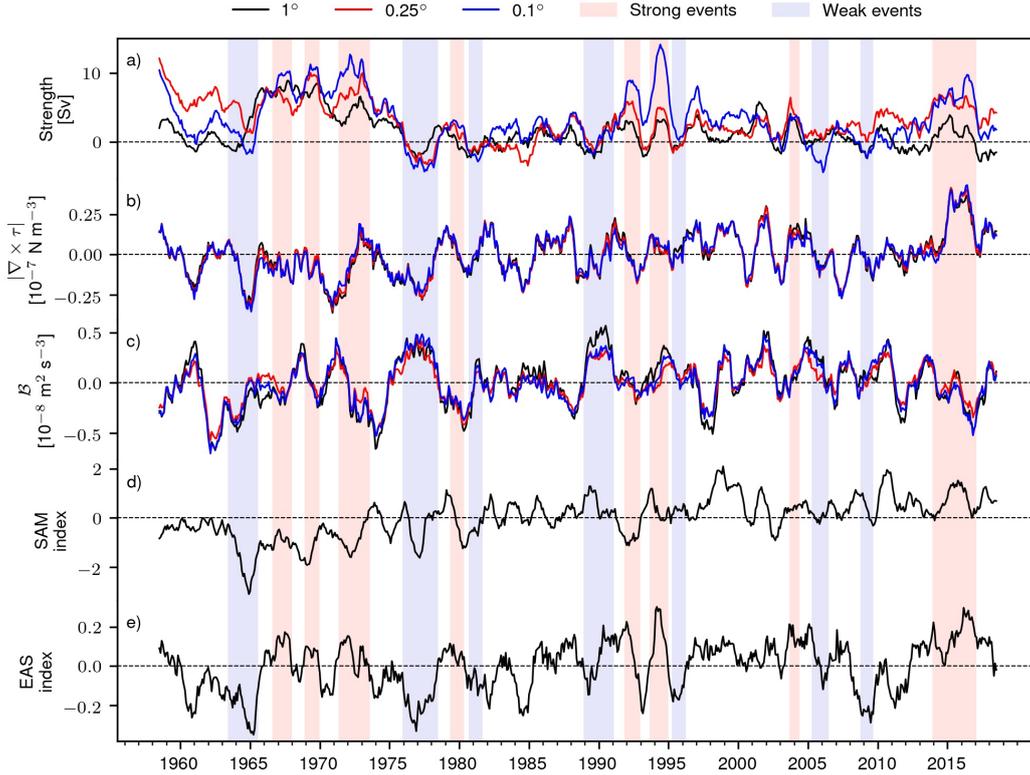


Figure 10. Time series of monthly anomalies with respect to the long term seasonal cycle of (a) gyre strength (Sv), (b) surface buoyancy flux ($10^{-8}m^2s^{-3}$), (c) surface stress curl ($10^{-7}Nm^{-3}$), (d) SAM index and (e) EAS index for the 1° , 0.25° and 0.1° resolutions. A 12-month running filter was applied to all time series. Red and blue shading indicate strong and weak events respectively for the 0.1° simulation, calculated with the method described in Section 2.7.

460 ice is advected to the north and then to the east, generating an increase of almost 10%
 461 in sea ice concentration at the northern rim of the gyre. Another region where strength
 462 impacts sea ice concentration is the eastern boundary, where we see a decrease in con-
 463 centration during stronger years consistent with a stronger warm inflow. Stronger years
 464 also display warmer subsurface temperature maximum along the southern region of the
 465 gyre that translate into a warmer continental shelf, Figure 11c. At around $30^\circ E$ there
 466 are significant correlations between the subsurface temperature maximum, indicating syn-
 467 chronous changes in gyre strength and the characteristics of the warm inflow.

468 7 Discussion and Summary

469 In this study we have used a coupled ocean-sea ice model at three different hori-
 470 zontal resolutions (1° , 0.25° and 0.1°) to diagnose the circulation of the Weddell Gyre
 471 and its connection to possible forcing mechanisms, namely surface stress forcing (includ-
 472 ing the contributions of wind and sea ice) and surface buoyancy fluxes. The inclusion
 473 of sea ice and the additional consideration of surface buoyancy fluxes expands on past
 474 studies that solely consider wind forcing due to the lack of appropriate observations (A. Naveira Gara-
 475 bato et al., 2019). We evaluate the model against available satellite observations and hy-
 476 drographic data and find a distinct improvement with resolution, particularly in the rep-

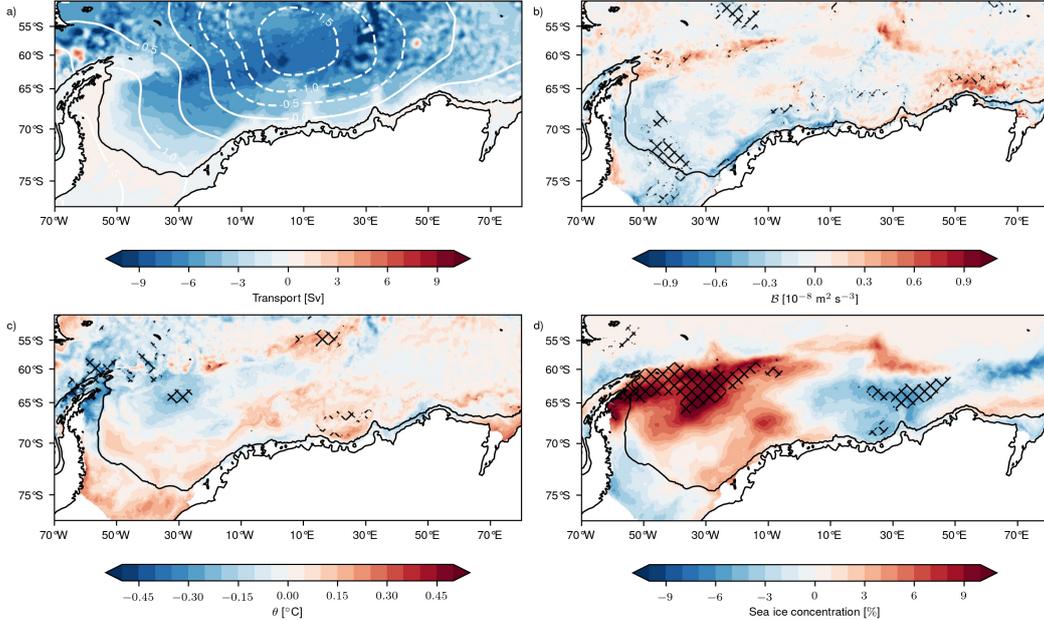


Figure 11. Strong minus weak Weddell Gyre composites of monthly anomalies for the events highlighted in Figure 10. Composites of (a) barotropic streamfunction (Sv) with sea level pressure contours overlaid (HPa), (b) surface buoyancy flux ($10^{-8} Nm^{-3}$), (c) subsurface temperature maximum ($^{\circ}C$) and (d) sea ice concentration (%) for the 0.1° simulation. Hatching indicates significant correlations ($p < 0.1$) between gyre strength time series and the corresponding field.

477 representation of the characteristics of regional water masses, which we suggest is a direct
 478 consequence of the ability of the 0.1° configuration to resolve instabilities along the east-
 479 ern boundary of the gyre. The model displays significant seasonal and interannual vari-
 480 ability in gyre strength, the magnitude of which increases with resolution. We find that
 481 extreme events of gyre circulation are climatically distinct, with significant differences
 482 in sea ice concentration and water mass characteristics during strong and weak phases
 483 of Weddell Gyre flow.

484 **7.1 Seasonal Cycle**

485 Our analysis indicates that most of the variability in gyre strength is concentrated
 486 in the seasonal cycle, which explains between 35% and 45% of the variance depending
 487 on the model’s horizontal resolution. This seasonal cycle consists of a winter intensifi-
 488 cation and a summer weakening, in agreement with past studies (Dellnitz et al., 2009;
 489 Beckmann et al., 1999), that increases in magnitude with resolution. Consequently, es-
 490 timates obtained from summer-biased observations will underestimate the gyre’s trans-
 491 port, stressing the importance of sustained, continuous observations in the Weddell re-
 492 gion. We further find that the seasonal changes in the barotropic circulation are not spa-
 493 tially uniform, but are concentrated in an eastern and a western region at approximately
 494 $30^{\circ}W$ and $30^{\circ}E$ respectively. Surface stress curl displays a semi-annual seasonal cycle
 495 peaking in April and October, highlighting the importance of sea ice in modulating the
 496 transfer of momentum from the atmosphere to the ocean, as emphasized by A. Naveira Gara-
 497 bato et al. (2019). Sea ice is also relevant for the seasonal cycle of surface buoyancy fluxes,
 498 where the ocean gains buoyancy during summer due to sea ice melt and loses surface buoy-
 499 ancy during winter due to sea ice formation.

7.2 Interannual Variability

The gyre also displays significant variability on interannual timescales, with gyre strength in phase across the three model resolutions, indicating the primacy of atmospheric forcing in driving the gyre’s variability. However, the variations in gyre strength become more pronounced with resolution, the consequences of which we explored by means of composites of strong and weak events on interannual timescales using the 0.1° configuration. The composites show that the gyre circulation has an impact on relevant processes within the region: a stronger gyre coincides with a warmer inflow at the eastern boundary that warms up the continental shelf and induces large changes in sea ice concentration by advection towards the northern rim. Anomalous gyre events over interannual timescales are associated with a deepening and expansion to the west of the regional low pressure system that deepens the trough and accelerates the coastal easterlies. This acceleration of the easterlies increases the onshore Ekman transport, raising the slope in sea level and accelerating the southern limb of the Weddell Gyre, a mechanism supported by the significant correlation between gyre strength and the easterlies at 0.1° model resolution, noting that this simulation is the model run that most accurately reproduces the ASC.

For the purpose of this study we have considered two possible forcing mechanisms for the Weddell Gyre’s variability: namely surface stress curl and surface buoyancy fluxes. Surface stress curl is traditionally understood to drive the depth-integrated circulation of the ocean via a linear vorticity relationship with the ocean’s meridional transport (Munk, 1950; Sverdrup, 1947). Meridional gradients of surface buoyancy fluxes have also been shown to be capable of maintaining a gyre-like circulation, and can even drive a mean ACC (Hogg & Gayen, 2020; Howard et al., 2015; Hogg, 2010; Wang & Meredith, 2008; England, 1993). In the Weddell Gyre region, surface buoyancy fluxes vary with largest amplitude on the continental shelf, where sea ice production and melt lead to large fresh-water fluxes, and therefore dominate the variability of the meridional buoyancy flux gradient. On the other hand, surface buoyancy fluxes have an impact on SSH via a steric effect, which in turn can drive changes in the barotropic circulation. However, changes in SSH over the continental shelf are rapidly propagated around the Antarctic continent via coastal-trapped barotropic Kelvin waves in a matter of days (Spence et al., 2017), which is why Weddell Gyre strength changes in SSH due to surface buoyancy fluxes are more likely to be found in the gyre’s interior following the cascading of dense waters down the continental slope following an overturning circulation (Jullion et al., 2014; A. C. Naveira Garabato et al., 2016). The coupling between ocean and sea ice in the model used for this study does not allow us to consider the independent contributions from buoyancy fluxes and surface stress. This is because surface winds not only impart momentum to the ocean’s surface, either directly or mediated by sea ice, they also affect surface heat and fresh-water fluxes. Thus, the simulations presented here cannot ascertain which factors, or non-linear interactions between them, are more relevant in driving the gyre’s variability. We conclude that more targeted numerical experiments are needed to address this question and to separate out stress from buoyancy forcing.

The importance of model resolution in the region is emphasized by its evaluation against hydrographic observations. Because of the Weddell Gyre’s open configuration at its eastern boundary, the properties of the water masses in the region are highly influenced by the characteristics of the inflow (Kerr et al., 2018; Jullion et al., 2014; Coul-drey et al., 2013). Two possible pathways have been suggested for this inflow: an eddy-driven path in the northeastern gyre boundary and an advective path further south towards the Antarctic continent (Ryan et al., 2016; Cisewski et al., 2011; Leach et al., 2011; Gouretski & Danilov, 1993). The warm inflow can be identified by a subsurface temperature maximum that is only reproduced by the 0.1° configuration, indicating that finer scale resolution processes are important for the exchange of waters between the gyre and its surroundings. The improved eastern boundary in the 0.1° simulation is not only re-

lated to the ability of this configuration to resolve instabilities, but also to a better representation of the ASF and the ASC at this resolution. And in turn, the characteristics of this warm inflow are likely responsible for the improvement with resolution of the temperature-salinity structure of the Weddell Gyre, potentially also affecting the characteristics of bottom waters in the model.

The variability of the Weddell Gyre has the potential to impact regional processes of relevance to global climate, but there is still not a complete understanding of the gyre's driving mechanisms and timescales of interaction (Vernet et al., 2019). We have considered surface stress and buoyancy fluxes as possible drivers of variability, but in a coupled ocean-sea ice model their intertwined nature does not allow us to consider them as independent mechanisms. It has also been suggested that the northern and southern limbs of the gyre vary independently, forced by the westerlies and easterlies respectively, with the difference balanced by inflows/outflows enabled by the gyre's open boundaries (Fahrbach et al., 2011). Our definition of gyre strength using the barotropic streamfunction is not intended to capture this independent variation. Remote forcing via the propagation of anomalies along the Antarctic continent could also affect the Weddell Gyre's circulation, but assessing this hypothesis is beyond the scope of this study. We have shown that the 0.1° resolution model has the ability to reproduce observed key features of this polar region and provides a 61-year long hindcast simulation that can prove useful for expanding our current knowledge of the Weddell Gyre. The model can also be used to explore the implications of the gyre for relevant processes such as warm water intrusions onto the Antarctic continental shelf and bottom water production. This is the subject of ongoing research into the dynamics and climatic relevance of the Weddell Gyre.

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