A new estimate of Global Ocean Carbon Flux from In Situ Optical Observations and Supervised Learning.

Daniel J Clements¹, Simon Yang¹, Thomas S Weber², Andrew M. P. McDonnell³, Rainer Kiko⁴, Lars Stemmann⁵, and Daniele Bianchi⁶

¹UCLA ²University of Rochester ³University of Alaska Fairbanks ⁴Sorbonne University ⁵U. Paris VI/ Villefranche ⁶University of California Los Angeles

November 23, 2022

Abstract

Export of sinking particles from the surface ocean is critical for carbon sequestration and for providing energy to the deepocean biosphere. The magnitude and spatial patterns of this flux have been estimated in the past by satellite-based algorithms and ocean biogeochemical models; however, these estimates remain uncertain. Here, we present a novel analysis of a global compilation of \textit{in situ} ocean particle size spectra from Underwater Vision Profiler 5 (UVP5) measurements, from which we determine particulate carbon fluxes. Using a machine learning algorithm, we extrapolate sparse observations of particle abundance by size to the global ocean from oceanographic variables that are more commonly observed. We reconstruct global maps of particle size distribution parameters for large sinking particles (80 \textmu{}m to 2.6 cm), and combine them with empirical relationships to calculate the sinking carbon flux from the euphotic zone and the wintertime mixed layer depth. Our flux reconstructions are comparable to other estimates, but suggest a less variable seasonal cycle in the tropical ocean, and a more continuous export in the Southern Ocean than previously thought. Because our estimates are not bounded by a specific depth horizon, we reconstruct export at multiple depths, and find that export from the wintertime mixed layer globally exceeds that from the euphotic zone. Our estimates provide a baseline for more accurate understanding of particle cycles in the ocean, and open the way to fully three-dimensional global reconstructions of particle size spectra and fluxes in the ocean, supported by the growing database of UVP5 observations.

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D.J. Clements¹, S. Yang¹, T. Weber², A.M.P. McDonnell³, R.Kiko⁴, L.Stemmann⁴, D.Bianchi¹

¹ Department of Atmospheric and Oceanic Sciences, University of California Los Angeles, Los Angeles,
CA, USA.
² Department of Earth and Environmental Sciences, University of Rochester, Rochester, New York, USA
3 College of Fisheries and Ocean Sciences, University of Alaska Fairbanks, Fairbanks, Alaska 99775-7220,
USA.
4 Sorbonne Université, CNRS, UMR 7093, Institut de la Mer de Villefranche sur mer, Laboratoire
d'Océanographie de Villefranche, Villefranche-sur-Mer, France.

Key Points:

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13	•	We use optical observations of particle size distribution to reconstruct global par-
14		ticulate carbon fluxes.
15	•	We quantify the importance of particle size distribution slope and biovolume on
16		the reconstructed fluxes.
17	•	Fluxes estimated at two depths show that more carbon is sequestered from the
18		wintertime mixed layer than the euphotic zone.

Corresponding author: D.J Clements, dclements@atmos.ucla.edu

 $Corresponding \ author: \ D. \ Bianchi, \ \texttt{dbianchi@atmos.ucla.edu}$

19 Abstract

Export of sinking particles from the surface ocean is critical for carbon sequestration and 20 for providing energy to the deep-ocean biosphere. The magnitude and spatial patterns 21 of this flux have been estimated in the past by satellite-based algorithms and ocean biogeochemical models; however, these estimates remain uncertain. Here, we present a novel 23 analysis of a global compilation of *in situ* ocean particle size spectra from Underwater 24 Vision Profiler 5 (UVP5) measurements, from which we determine particulate carbon 25 fluxes. Using a machine learning algorithm, we extrapolate sparse observations of par-26 ticle abundance by size to the global ocean from oceanographic variables that are more 27 commonly observed. We reconstruct global maps of particle size distribution parame-28 ters for large sinking particles (80 μ m to 2.6 cm), and combine them with empirical re-29 lationships to calculate the sinking carbon flux from the euphotic zone and the winter-30 time mixed layer depth. Our flux reconstructions are comparable to other estimates, but 31 suggest a less variable seasonal cycle in the tropical ocean, and a more continuous ex-32 port in the Southern Ocean than previously thought. Because our estimates are not bounded 33 by a specific depth horizon, we reconstruct export at multiple depths, and find that ex-34 port from the wintertime mixed layer globally exceeds that from the euphotic zone. Our 35 estimates provide a baseline for more accurate understanding of particle cycles in the ocean, 36 and open the way to fully three-dimensional global reconstructions of particle size spec-37 tra and fluxes in the ocean, supported by the growing database of UVP5 observations. 38

³⁹ 1 Introduction

In the ocean, primary production and other complex biogeochemical processes in-40 teract to form the ocean's biological pump. Aggregation of particulate organic matter 41 into particles denser than seawater leads to gravitational settling (Alldredge & Gotschalk, 42 1988), eventually storing inorganic carbon and nutrients in the deep ocean for timescales 43 of decades to centuries (Boyd et al., 1999, 2019). The export of particulate organic car-44 bon provides energy to the deep ocean ecosystem (Siegel et al., 2014), influences atmo-45 spheric carbon dioxide and climate (Kwon et al., 2009; Palevsky & Doney, 2018), and 46 indirectly affects on the ocean's microbiome (Karl et al., 1984; Bianchi et al., 2018). Sev-47 eral studies have estimated this global particle flux from the euphotic zone, resulting in 48 substantially variable estimates ranging from 3 to 10 PgC/y (Henson et al., 2011; Siegel 49 et al., 2014; DeVries & Weber, 2017; Dunne et al., 2007). 50

Current reconstructions of the sinking particle flux tend to vary in the total global 51 export, depending on the methods used (Quay et al., 2020). Biogeochemical models yield 52 a global export of 4-6 PgC/y when tuned to match particle observations, but could reach 53 up to 10 PgC/y when tuned to match in situ profiles of nutrients and other biogeochem-54 ical tracers (Siegel et al., 2014; DeVries et al., 2017). Models that use satellite inputs and 55 empirically derived export ratios tend to result in a larger flux (Dunne et al., 2007; Laws 56 et al., 2011; Guidi et al., 2015), near 10 PgC/y. This is similar to annual net commu-57 nity production estimates at the base of the mixed layer, which include additional export of dissolved organic carbon (Emerson, 2013; Quay et al., 2020). 59

Although the globally integrated particle flux is similar when comparing model, geo-60 chemical, and satellite-based estimates, the regional patterns predicted by these meth-61 ods differ substantially. Differences in regional flux estimates have been attributed to method-62 ological limitations, including scarcity and variability of in situ data used to constrain 63 models, variability in satellite-based primary production algorithms, and models not able 64 to fully capture underlying physical and biological processes. Based on in situ geochem-65 ical observations, Quay et al. (2020) suggest a weaker meridional variability in export 66 flux than other estimates, stressing the need for expanding and combining observational 67 approaches and models to fully constrain particle export.

Recent studies have highlighted the importance of standardized methods and met-69 rics used to quantify particle fluxes (Buesseler et al., 2020). In particular, the depth hori-70 zon of export has been identified as a leading cause of diverging estimates (Palevsky & 71 Doney, 2018). Two choices of export horizon have been commonly adopted: the base of the euphotic zone, either as a variable depth or global average (Buesseler & Boyd, 2009; 73 Siegel et al., 2016, 2014; Bisson et al., 2018; Dunne et al., 2007; DeVries et al., 2017; Hen-74 son et al., 2011), and the mixed layer depth, both as seasonally varying and maximum 75 depth (Emerson, 2013; Quay et al., 2020). These choices underlie different interpreta-76 tions of export fluxes: euphotic export takes an ecosystem view, while mixed layer ex-77 port takes a carbon storage view. 78

Gravitational settling is thought to be the primary export mechanism globally, contributing to about 60% of the total carbon export, and about half of the carbon storage in the deep ocean (Boyd et al., 2019). Using both a euphotic viewpoint, and considering only gravitational settling, particle flux estimates have begun to converge on a value of 5-6 PgC/y (Palevsky & Doney, 2018; Boyd et al., 2019).

Advances in ocean optical observations have begun to provide a three-dimensional 84 view into the life of particulate matter in the ocean (Stemmann & Boss, 2012; Kiko et 85 al., 2017; Guidi et al., 2009). The Underwater Vision Profiler 5 (UVP5) is an optical par-86 ticle counter which provides the *in situ* particle abundance for large particles ($80\mu m$ -87 2.6cm) in a given sampled volume (Picheral et al., 2010). The UVP5 consists of a cam-88 era attached to the CTD rosette, and quantifies the particle abundances at high frequency as it is lowered in the water column. Vertical profiles of particle size distribution (PSD) 90 from the UVP5 are commonly taken at up to 20 images per second with downward speeds 91 of 1m/s, with observations as deep at 6 km (Picheral et al., 2010). Since 2008, UVP5s 92 have been routinely deployed on ocean expeditions, resulting in over 9,000 profiles to date, 93 with observations from all ocean basins. 94

Although the UVP5 cannot directly determine carbon flux, because particle sink-95 ing speed and carbon content are not measured, empirical relationships have been used 96 to define these as a function of size, making flux estimates possible (e.g., as compiled in 97 Kriest (2002) and Stemmann et al. (2004)). The vertical resolution of the UVP5, cou-98 pled with these empirical relationships, enables a unique high-resolution view into the 99 three-dimensional ocean particle flux. Observations from UVP5 have been used to quan-100 tify particulate flux to the mesopelagic ocean on a regional basis (Guidi et al., 2008, 2009; 101 Kiko et al., 2017), and a smaller dataset has been used to reconstruct global fluxes by 102 large-scale biomes (Guidi et al., 2015). However, the expanded dataset has not yet been 103 used to quantify fluxes from the surface ocean yet. 104

In this study, we reconstruct global particle carbon export by training a supervised 105 machine learning algorithm to extrapolate PSD from a rapidly growing dataset of UVP5 106 observations and well-sampled oceanographic variables. We combine these estimates with 107 in situ sediment trap and thorium-derived particle flux observations to better constrain 108 empirical relationships between particle size, sinking speed and carbon content, produc-109 ing robust estimates of regional and seasonal flux variability. By comparing patterns in 110 PSD and flux with observations of environmental drivers, we further gain insight into 111 the mechanisms responsible for particle export and its variability. Finally, we exploit the 112 high vertical resolution of UVP5 measurements to estimate particle fluxes at both the climatological euphotic zone and the maximum mixed layer depth, revealing the impor-114 tance of the export horizon for this quantity. 115

The rest of the paper is organized as follows. Section 2 describes the machine-learning approach used to globally extrapolate PSD and reconstruct particle fluxes. Section 3 presents the reconstructions of particle distributions and export fluxes, and compares our results to previous studies at global and regional scales. Section 4 summarizes the main findings, discussing the uncertainties and caveats inherent to our approach, and future di-rections.

¹²² 2 Methods

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The flux of particulate carbon (ϕ) at any given depth can be expressed as a function of three size-dependent quantities: PSD $(n(s), \frac{\#}{L \cdot cm})$, sinking speed $(w(s), \frac{m}{s})$, and particle carbon content (c(s), mgC), according to the following equation (Stemmann & Boss, 2012):

$$\phi = \int_{s_{min}}^{s_{max}} n(s) \cdot w(s) \cdot c(s) \, ds, \tag{1}$$

Here, s indicates the particle equivalent spherical diameter, or size (cm), s_{min} and s_{max} are respectively the minimum and maximum size of particles considered for export. We further assume that the quantities in Equation 1 can be approximated by power laws that depend on particle size, each characterized by an intercept (the size-independent coefficient) and a slope (the exponent for size-dependence) (Stemmann & Boss, 2012):

$$n(s) = n_0 \cdot s^{-\beta} \tag{2}$$

$$w(s) = w_0 \cdot s^\eta \tag{3}$$

$$c(s) = c_0 \cdot s^{\zeta},\tag{4}$$

Thus, by using Equations 2-4, the total particle flux can be expressed as:

$$\phi = \int_{s_{min}}^{s_{max}} n_0 \cdot w_0 \cdot c_0 \cdot s^{-\beta + \eta + \zeta} \, ds = \int_{s_{min}}^{s_{max}} n_0 \cdot m_0 \cdot s^{-\beta + \mu} \, ds \tag{5}$$

where we combined the intercepts and exponents of the sinking speed and carbon content relationships by setting $m_0 = w_0 \cdot c_0$ and $\mu = \eta + \zeta$. We further approximate m_0 and μ with globally constant values.

We use UVP5 observations to reconstruct PSDs (i.e., n_0 and β) at the chosen ex-141 port horizon, by fitting Equation 2 to the observed particle abundance. We use a 20 me-142 ter depth bin around the export horizon to average the observations and smooth out small-143 scale noise and variability. We then extrapolate sparse UVP5 observations to a global 144 grid, by training a supervised learning algorithm to predict spatially-varying PSD from 145 well-sampled environmental predictors. To completely reconstruct fluxes based on Equation (1), we constrain the parameters of the combined sinking speed and carbon content 147 relationships, i.e., m_0 and μ , by optimizing predicted particle fluxes against *in situ* flux 148 estimates from sediment traps and thorium-uranium disequilibrium measurements. 149

We exploit the three-dimensional nature of UVP5 observations to perform these calculation at two different export horizons: the base of the euphotic zone (here defined by the 1% light level according to Morel et al. (2007)) and the annual maximum mixed layer depth. The steps used to solve Equation (1) are illustrated in the workflow schematic in Fig. 1, and are detailed in the following sections.



Figure 1. Schematic diagram illustrating the general workflow of processing UVP5 observations into a flux reconstruction. Observations are ensembled onto a normal 1 degree grid, with observation representing an average of a 20 meter vertical bin about the export horizon. All observations are the calculated values for the 105 μm to 5 mm size bins. Specific parameters are globally extrapolated using the Random Forest algorithm. These new global data are used to construct the flux, using equation 1.

2.1 Reconstructions of particle size spectra from UVP5 data

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156 Observations from UVP5 provide particle counts for the 80 μ m - 2.6 cm size range 157 at any specific location and depth. Under the power law assumption, the two parame-158 ters n_0 and β are needed to fully capture the PSD (Equation (2)) (Brun-Cottan, 1971; 159 Stemmann et al., 2004; Stemmann & Boss, 2012; Devries et al., 2014).

The slope β quantifies the relative abundance of large vs. small particles, while the intercept n_0 is a measure of the total abundance of particles for an arbitrary reference size. We determine the slope for each UVP5 measurement by fitting a linear least-squares regression through the log-transformed particle abundance and size. We determine the intercept by using the particle biovolume (BV), which can be directly derived from UVP5 observations, and applying the following definition:

$$BV = \int_{s_{min}}^{s_{max}} \frac{\pi}{6} \cdot n_0 \cdot s^{3-\beta} \, ds. \tag{6}$$

By fixing the size range, we solve Equation 6 for the intercept as a function of slope and biovolume:

$$n_0 = \frac{6 \cdot BV}{\pi} \cdot \left(\frac{s_{max}^{4-\beta}}{4-\beta} - \frac{s_{min}^{4-\beta}}{4-\beta}\right)^{-1}.$$
 (7)

We set the minimum and maximum size for this equation to the same values used to estimate the slope and biovolume from UVP5 observations. We use a minimum size of 105 μ m to avoid a potential slight instrument bias in the lowest size classes. We constrain the maximum size to 5 mm, which corresponds to the size where zooplankton start to dominate the biovolume at a variety of locations sampled by UVP5 (Forest et al., 2012; Stemmann et al., 2008; Stemmann & Boss, 2012). The sensitivity of our results to these choices is discussed in Section 2.2.

We coarsen the temporal and spatial resolution of the 6856 UVP5 profiles by binning them onto the standard monthly 1 degree-resolution grid of the World Ocean Atlas (H. Garcia et al., 2018; H. E. Garcia et al., 2019). To this end, we combine profiles in a given grid cell and month together, thus reducing the noisy and episodic nature of particle observations. To reconstruct global PSDs, we calculate slope and biovolume for each grid cell, at a given export horizon, using the gridded dataset, and assume that these averages are representative of the climatological monthly PSD in each grid cell.

Although the gridded observations reduce data patchiness in well sampled regions, many grid cells only contain one observation. For each grid cell with observations, we place an objective goodness of fit threshold to determine the robustness of the power law fit. If a power law fit has a correlation coefficient of less than 0.9, we remove the data point, as it likely does not follow a power law distribution. This quality control step removes less than 1% of data (Supplementary Information Fig. S1). The final processed UVP5 observation dataset contains 2,034 gridded observations at the export horizon, which together cover less than 10% of the ocean surface.

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2.1.1 Training and evaluating a Random Forest model

Monthly flux reconstructions require extrapolation of PSD parameters to the whole 193 ocean on monthly time scales. We use a Random Forest (RF) algorithm to reconstruct 194 PSD slope and biovolume globally, following an approach similar to Yang et al. (2020). A RF deploys a decision tree learning scheme to solve a regression equation iteratively, 196 and reports the ensembled average. Using a RF, each individual decision tree is trained 197 on a subset of the available data, with a subset of predictors, but the power of the method 198 emerges when considering the ensemble average. The RF is able to learn statistical re-199 lationships between target variables (here, UVP5-derived slope and biovolume) and a 200 series of predictors (here, environmental variables), to make reconstructions that min-201 imize the error between predicted and observed data. Because a RF is highly non-linear, 202 it runs the risk of overfitting the data, producing solutions with low error, but also limited extrapolation power outside of the training dataset. To mitigate the risk of over-204 fitting, the RF does not use all data points for training. Instead, a bootstrapped sam-205 ple of the data is selected for each tree in the forest. The degree of overfitting is deter-206 mined by finding the error between the model and the data not used for training, i.e., 207 the "out-of-bag" data. 208

The rank of predictors is given by the out-of-bag error coupled with an internally 209 derived measure of importance, using a "recursive feature elimination" approach. A re-210 cursive feature elimination systematically removes the least important predictor and records 211 the out-of-bag error to describe the contribution of each predictor to the final solution. 212 When there is relatively no change in the out-of-bag error for every additional predic-213 tor, these predictors are considered not important for this RF (Supplementary Fig. S2). 214 We determine statistical importance in order to establish a reduced set of predictors, reducing the risk of over-fitting while not losing predictive power. When interpreting the 216 results, we apply qualitative understanding of the predictors, combined with the recur-217 sive feature elimination, to determine if a predictor is should be included in the final re-218 gression. 219

220 2.1.2 Environmental Predictors

The RF algorithm relies on a set of predictors and target data at the resolution of the desired reconstruction. In our case, we use climatological monthly predictors at 1degree spatial resolution. We include a variety of predictors that could be mechanistically related to particle production and export in the surface ocean, ranging from physical variables (e.g., temperature and salinity) to ecosystem-level quantities (e.g., primary production, euphotic zone depth). A list of all predictors used is shown in Table 1.

Some of these predictors are obtained from satellite products at high spatial and 227 temporal resolution (e.g., surface chlorophyll and net primary production), and include 228 missing values caused by the presence of clouds or sea-ice. For these variables, we first 229 average observations into monthly climatologies, then replace missing data by using a 230 spherical interpolation algorithm (D'Errico, 2016). To avoid excessive extrapolation in 231 high latitude regions in wintertime, only points with at least 8 months of satellite observations are used for the final reconstruction, following the approach of Siegel et al. 233 (2014). To process net primary production, we also calculate the critical depth, where 234 light becomes too limiting to support photosynthesis, based on climatological chlorophyll 235 concentration and incident shortwave radiation (Siegel et al., 2002). Net primary pro-236 duction is then set to zero at all points where the mixed layer depth exceeds the crit-237 ical depth, before interpolating. We also include the standard deviation of the primary 238 production, as a proxy for intermittency and sub-seasonal variability. Similarly, we re-220 strict chlorophyll and net primary production based on climatological sea ice cover.

We use three-dimensional variables (e.g., temperature, nutrients) to generate two-241 dimensional surface predictors based on mixed layer averages. We also include predic-242 tors that quantify the change of variables with depth, by calculating the average of the 243 variable from the base of the mixed layer to 100m below it. For surface-only variables (e.g., chlorophyll, net primary production) and nutrients we also include predictors that 245 quantify change in time, calculated by applying a finite-difference time derivative, be-246 cause change over time might be more indicative of export flux than the actual variable. 247 We refer to these depth- and time-change variables as "variations" in Table 1. We test 248 the significance of each predictor, including vertical and time variations, with the recur-249 sive feature elimination. Finally, we group predictors into different categories, with vari-250 ations for selected variables (Table 1). If a predictor is in the universal category, it is al-251 ways included in all RF realizations. For all other categories, only one predictor is chosen for each realization, but if a predictor is chosen, all variations are included too. 253

Based on the categories listed in Table 1, we use a total of 29 predictors for each 254 RF realization (Table 1). We generate 100 realizations, with variable hyper-parameters 255 (the number of trees and their complexity) and randomly chosen predictors from each category, and take the ensemble average as the final product, with the inter-model spread 257 representing the error. Generating an ensemble of 100 RFs, with varying hyper-parameters 258 and predictors, reduces biases and overfitting, making the results robust with respect to 259 parameter tuning, and the choice of different data products. Thus, our reconstructions 260 are not the result of tuning the hyper-parameters, or choosing only the best predictors. 261 We evaluate the overall robustness of the predictions by considering goodness-of-fit statis-262 tics that include the correlation coefficient, the root mean square error (RMSE), and the 263 average bias, calculated by comparing predictions to in situ data.

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2.2 Sinking Speed and Carbon Content

Particle sinking speed and carbon content have been empirically estimated using
power law relationships analogous to Equations (3) and (4). Most of these studies measured a range of particles that does not wholly encompass the sizes detected by the UVP5.
Also, these relationships are defined for specific particles types, which are not distinguished
in the dataset used. Since estimates of total flux are sensitive to the sinking speed and

Category	Variable	Variations	Source
Universal			
	Topography		N.G.D.C (2006)
	Temperature below mixed layer	Time Derivative	Locarnini et al. (2019)
	Chlorophyll	Time Derivative	NASA G.S.F.C (2014)
	Oxygen	ML/ ML+100m Time Derivative	H. E. Garcia et al. (2019)
	Shortwave Radiation	Time Derivative	C.C.C.S (2017)
	Nitrate	$\begin{array}{c} \mathrm{ML}/~\mathrm{ML}{+100\mathrm{m}}\\ \mathrm{Time~Derivative} \end{array}$	H. Garcia et al. (2018)
	Phosphate	$\begin{array}{c} \mathrm{ML}/~\mathrm{ML}{+}100\mathrm{m}\\ \mathrm{Time~Derivative} \end{array}$	H. Garcia et al. (2018)
	Salinity	$\rm ML/~ML{+}100m$	Zweng et al. (2019)
Mixed Layer			
	Mixed Layer depth	Time Derivative	Johnson et al. (2012)
	Mixed Layer depth	Time Derivative	de Boyer Montégut et al. (2004)
Primary Production			
	Eppley VGPM	Time Derivative	Antoine and Morel (1996)
	VGPM	Time Derivative	Behrenfeld and Falkowski (1997)
	CBPM	Time Derivative	Westberry et al. (2008)
	CAFE	Time Derivative	Silsbe et al. (2016)
NPP Standard Deviation			
	Eppley VGPM		Antoine and Morel (1996)
	VGPM		Behrenfeld and Falkowski (1997)
	CBPM		Westberry et al. (2008)
Euphotic Zone Depth			
	VGPM		Morel et al. (2007)
	CBPM		Morel et al. (2007)
Iron			
	Soluble Iron Fraction	Time Derivative	Hamilton et al. (2019)
	Labile Iron Fraction	Time Derivative	Myriokefalitakis et al. (2018)

Table 1. Variables used to predict PSD parameters, their source and variations. The categories are organized based on their predictor type where the universal predictors are used in every random forest.

carbon content relationships, here encapsulated by the parameters m_0 and μ , we apply an optimization procedure to keep our results consistent with *in situ* particle flux measurements. Specifically, we find the values of m_0 and μ that minimize the sum of the square errors between particle flux predictions (Equation 5) and co-located *in situ* carbon flux measurements (Bisson et al., 2018).

Because the size distribution of particles that contribute to the flux is poorly con-276 strained, we perform this optimization for a range of plausible minimum and maximum 277 sizes for Equation (5), selecting reasonable combination for the final estimate. Ultimately, when optimizing the sinking carbon parameters, the global export flux is insensitive to the size range; however the resulting empirical relationships are (see Section 3.2 and Sup-280 plementary Fig. S5). The insensitivity of the carbon flux to the size range of that flux 281 indicates a compensatory effect between the sinking carbon parameters and the size range 282 selected. Thus, choosing different size combinations would result in a similar total flux, 283 although it may slightly alter spatial or temporal patterns. 284

Our final choice of size range is informed by average sinking speeds and carbon con-285 tent previously reported (Kriest, 2002). We assume the slope and intercept (calculated 286 from the biovolume; Equation 7) of the particle size distribution to be constant for sink-287 ing particulate matter, and expand the minimum size to include all sinking particle sizes. 288 Based on this analysis, we set the minimum size class to be 35 μ m, where the average 289 sinking speed is near 1 m d⁻¹ (Smayda, 1970; Kriest, 2002), a size that is likely rapidly remineralized, making its contribution to the sinking flux negligible (Riley et al., 2012). 291 Similar to the observation limitations, we choose 5 mm as the maximum size of sinking 292 particles. 293

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2.3 Flux reconstruction and evaluation

Using the PSD reconstructions and the optimized sinking carbon parameters, we 295 calculate particle export fluxes following Equation (5). We evaluate these reconstructions by comparing them to *in situ* flux observations and previous global reconstructions. Specif-297 ically, we compare total fluxes, meridional averages, and seasonal cycles. For these com-298 parisons, we divide the ocean into 14 biogeochemically-consistent regions based on the 299 boundaries identified by Weber et al. (2016), with an additional boundary along the equa-300 tor to separate Northern and Southern Hemispheres. We evaluate seasonal cycles by an-301 alyzing temporal correlations between reconstructions, and by introducing a seasonal-302 ity index defined by the ratio between the seasonal range and annual mean flux in each region. 304

We first present results for fluxes estimated at the climatological euphotic zone depth, and then repeat the calculation at the maximum mixed layer depth. For the latter, we keep the same sinking speed and carbon content parameters as determined for the euphotic zone depth. Thus, the only methodological difference between the two estimates is the depth used to calculate the export flux.

- 310 3 Results
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3.1 Particle Size Distribution reconstructions

Our global reconstructions capture most of the variability of UVP5-based PSD slope and biovolume data (Figs. 2 and 3), and robustly reproduce observations, with global average values of 0.6 ppm for biovolume $(r^2 = 0.91)$ and 3.9 for slope $(r^2 = 0.86)$. Observations that are not used in the training (out-of-bag) are also robustly predicted, with a RMSE of 2.1 ppm for biovolume $(r^2 = 0.74)$ and 0.33 for slope $(r^2 = 0.68)$. Some of the remaining uncertainty in our reconstructions can be attributed to the episodic nature of particle production and export. Our method operates under the assumption that



Figure 2. Observed and reconstructed Particle biovolume at the base of the euphotic zone. (A) Map of the observations of intercept, for locations with observations in multiple months the average is shown (B) Map of the biovolume reconstructions. (C) Performance of the RF reconstruction shown as density scatter plots of predicted vs. observed Biovolume (colors indicate the normalized density of grid points surrounding the given cell). (D) Same as B, but using out-of-bag (OOB) predictions, i.e., predictions vs. observations withheld from training. Annotations in B and C show the square of the correlation coefficient (r2), the RMSE and the global bias.

the input data (i.e., the UVP5 measurements) are monthly climatological averages, rather than snapshots. By ensembling these snapshot measurements into 2,034 monthly observational data points, we reduce part of the episodic nature of these observations; however some patchy behavior may still exist in the gridded data. Overall, the reconstructions show slight underestimates of extreme values (i.e., a reduced range), but negligible mean biases for both variables (Figs. 2 and 3).

We find high biovolume in productive regions such as high latitudes, coastal wa-325 ters, and upwelling systems, and low biovolume in the oligotrophic suptropical gyres. These 326 patterns mirror satellite-derived primary productivity and chlorophyll measurements (Sup-327 plementary Fig. S3), suggesting that phytoplankton and photosynthesis ultimately con-328 trol the total abundance of particles in any given region. Reconstructions of PSD slope 329 show a similar but negative correlation with primary production and chlorophyll, with 330 smaller slopes (i.e., "flatter" PSD) in more productive regions, and larger slopes (i.e., "steeper" 331 PSD) in oligotrophic waters. 332



Figure 3. Observed and reconstructed PSD slope at the base of the euphotic zone. (A) Map of the observations of PSD slope, for locations with observations in multiple months the average is shown (B) Map of the PSD slope reconstructions, colored dots show the observations from A. (C) Performance of the RF reconstruction shown as density scatter plots of predicted vs. observed particulate slope (colors indicate the normalized density of grid points surrounding a given cell). (D) Same as B, but using out-of-bag (OOB) predictions, i.e., predictions vs. observations withheld from training. Annotations in B and C show the square of the correlation coefficient (r2), the RMSE and the global bias.

Consistently, we find that slope and biovolume are negatively correlated ($r^2 = 0.4, p < 0.01$, Supplementary Fig. S3), roughly indicating that particle-rich regions (higher biovolume) are also characterized by an excess of large particles over small particles (i.e., flatter slope), relative to average oceanic conditions. Since large particles contribute proportionally more than smaller particles to export fluxes, because of their faster sinking speed, this relationship suggests that biovolume and slope will synergistically enhance export fluxes in particle-rich regions, and depress them in particle-poor regions.

While this pattern of correlations holds true for most regions, we find few significant exceptions where the PSD slope and biovolume do not co-vary as closely as expected. For example, in the North Pacific Subpolar Gyre, flatter slopes are found in the open ocean (Fig. 3), in particular close to the subpolar-subtropical transition, while the highest biovolumes are found closer to the coast and in marginal seas. Similarly, slopes in coastal upwelling system, such as the California Current and the Arabian Sea upwelling, are not as flat as the high biovolumes would suggest. We also find relatively flatter slopesin the North Pacific subtropical gyre as compared to other oligotrophic regions.

The seasonal dynamics of biovolume and slope confirms the general anti-correlation 348 of these two variables, and reveals significant seasonal cycles, with maximum biovolume 349 and minimum slope generally found in spring, and minimum biovolume and maximum 350 slope in late fall to winter (Supplementary Fig S4). Similar to the spatial distribution, 351 we find significant deviations from the general anti-correlation between biovolume and 352 slope. For example, in the North Atlantic, the peak in biovolume (May) precedes the min-353 imum in slope (July). In some of the tropical regions (e.g., in the North Pacific and North Atlantic) the anti-correlation is also less robust, with periods of several months where 355 biovolume and slope increase or decrease simultaneously. As discussed above, spatial and 356 temporal decoupling of the biovolume-slope relationship could have important consequences 357 for the patterns of particle export flux. 358

A recursive feature elimination suggests that multiple variables are required for a 359 robust reconstruction of PSD (Supplementary Fig. S2). Among these, we highlight chloro-360 phyll, mixed layer depth, and oxygen, each with different importance for explaining bio-361 volume and slope variability. Interpretation of these rankings should be done with care 362 because of the statistical nature of the RF algorithm. However, while a mechanistic un-363 derstanding of PSD patterns can not be directly tied to these rankings, highlighted pre-364 dictors can provide insights into the role of different processes that may be affecting PSDs. We note that certain predictors with mechanistic links to particle export, for example silicate, which serves as a proxy for diatom production, were deemed unimportant and 367 are not included in the final RF regressions. It is likely that information contained in 368 these predictors is shared by correlations with other variables, and is picked up by the 369 main predictors used by the method. 370

Multiple variables are significantly correlated with biovolume and slope (Supple-371 mentary Fig. S6). In particular, we find that biovolume correlates positively and robustly 372 with chlorophyll ($r^2 = 0.25, p < 0.01$, Supplementary Fig. S3). This is not surpris-373 ing, because phytoplankton are ultimately the main source of organic matter and sink-374 ing particles in the ocean. However, we find that chlorophyll is not as strong a predic-375 tor of slope, when the whole ocean is considered $(r^2 = 0.04, p < 0.01, \text{Supplementary})$ 376 Fig. S3), and that additional predictors are needed for robust slope reconstructions. This 277 result contrasts with previous findings based on UVP5 observations along a meridional section in the Pacific Ocean (Cram et al., 2018). Slope reconstructions also reveal a sig-379 nificant predictive power for subsurface oxygen. Although likely not directly related to 380 the PSD, oxygen is a proxy of respiration in the water column, which in turn reflects the 381 characteristics of both the surface community that drives export, and of the mesopelagic 382 community responsible for this respiration. We note that the PSD slope is an emergent 383 property that reflects the interaction of physical and biological processes that are still 384 poorly understood. Not surprisingly, slope is harder to reconstruct than biovolume, and shows overall weaker correlations with other individual predictors (Supplementary Fig. S6). 387

Spatial patterns in slope share several features with estimates of phytoplankton size 388 spectra from observations and models (Kostadinov et al., 2009; Roy et al., 2013; Barton et al., 2013; Ward et al., 2014), reflecting the importance of phytoplankton size structure and composition for particle export. Satellite-based reconstructions of phytoplank-391 ton functional groups (e.g., Mouw et al., 2017) could be included as predictors. How-392 ever, methodological shortcomings and disagreement between different approaches cur-393 394 rently limit their applicability — something that may be mitigated by future advances. It is also likely that information related to phytoplankton composition implicitly enters 395 the RF regression via relationships with predictors such as surface chlorophyll and tem-396 perature. 397

3.2 Particle export flux



Figure 4. Density scatter plots showing the relationships between *in situ* flux observations and global flux reconstructions from (A) This study, (B) Siegel et al. (2014)l, (C) Henson et al. (2011), (D) (Dunne et al., 2007). Colored dots represent the relative density of grid points surrounding the data point, dashed line indicates a 1:1 ratio. Annotations show the correlation coefficient (r^2) RMSE, and average bias.

Based on the PSD reconstructions and Equation 5, we optimize for the particle sinking speed and carbon content parameters $(m_0 \text{ and } \mu)$ that produce export fluxes in best agreement with *in situ* observations (see Section 2.2). This approach results in a value of 27.65 mgC m s⁻¹ for m_0 , and of 2.90 for μ , both in the range suggested by *in situ* observations (Kriest, 2002), and comparable to values adopted by previous studies (Kriest, 2002; Stemmann et al., 2004; Guidi et al., 2008; Kiko et al., 2017; Bianchi et al., 2018).

The resulting carbon fluxes compare well with sediment trap and thorium-based observations (Fig. 4), performing in a similar way or better than previous estimates (Henson et al., 2011; Dunne et al., 2007; Siegel et al., 2014). Our estimate reduces the overall uncertainty (here expressed by the RMSE) compared to previous work, and shows negligible bias. However, our method also reduces the overall range of reconstructed fluxes, i.e., it overestimates the flux at low values and underestimates it a high values compared to observations. This bias may be related to the tendence of the RF algorithm to underpredict extremes in both biovolume and slope (Figs. 2 and 3).

Extrapolated to the whole ocean, our method reveals spatial patterns of export fluxes 413 in broad agreement with previous studies, with some notable differences (Fig. 5). Sim-414 ilar to other estimates, particle fluxes tend to decrease from high to low latitudes, and 415 from coastal regions to the open ocean. A local maximum of export is reproduced along 416 the equator, and is particularly evident in the Pacific Ocean. Compared to previous work, 417 our method produces somewhat weaker gradients between coastal and offshore waters, 418 with relatively high fluxes even near the centers of subtropical gyres, and suggests an 419 asymmetry between the subpolar Atlantic and Pacific Oceans, with more intense par-420 ticle export along the gulf of Alaska than in the North Atlantic (see also Section 3.2.1). 421 We also reconstruct substantially stronger export than previously found in the South-422 ern Ocean, in particular south of 50S (see discussion in Section 3.3). 423

Globally integrated, we estimate a particle export flux of $6.7 \pm 0.4 \text{ PgC/y}$, in excellent agreement with the range of observational and model-based estimates of the biological gravitational pump (4-9 PgC/y, Boyd et al. (2019)). Compared to other spatiallyresolved reconstructions, our global flux sits between the low-value estimate of Henson et al. (2011) ($3.0 \pm 0.3 \text{ PgC/y}$) and the high-value estimate of Dunne et al. (2007) (9.8 $\pm 0.4 \text{ PgC/y}$).



Figure 5. Annual average particle export from the euphotic zone for the (a) Random forest derived compared to the *in situ* data of (b) Bisson et al. (2018), the steady state satellite driven model (c) SIMPLE-TRIM of DeVries et al. (2017), empirical models of (d) Dunne et al. (2007) and (f) Henson et al. (2011) and the satellite-driven euphotic zone food web model of (e) Siegel et al. (2014). Annotated in each figure is the calculated total export and uncertainty reported by each study.

430 3.2.1 Meridional variability

We illustrate the main spatial differences between our and other reconstructions 431 by considering zonally averaged export fluxes (Fig. 6). The largest export rates are ob-432 served around the equator, in the subpolar Pacific Ocean, and in the mid- to high-latitudes 433 of the South Atlantic Ocean, while more uniform exports are observed in the Indian Ocean. 434 In all basins, the minimum export rates are generally located at the latitude of the sub-435 tropical gyres. While export is nearly symmetrical around the equator in the Pacific Ocean 436 (Fig. 6a), in the Atlantic Ocean it dramatically increases moving from the Northern to 437 the Southern Hemisphere (Fig. 6b). These patterns reflect a combination of open-ocean and shelf enhanced particle export. Specifically, high exports in the Northern Pacific and 439 Southern Atlantic Oceans are partly driven by large fluxes in the Bering Sea, the Sea 440 of Okhotsk, and the Patagonian shelf. At lower latitudes, coastal upwelling systems sus-441 tain high export in the northern Indian Ocean and the tropical to subtropical Atlantic 442 Ocean. 443

Variations in export patterns reflect a combination of varying particle biovolume 444 (Fig. 2) and PSD slope (Fig. 3). These two quantities generally correlate in such a way 445 as to increase export fluxes in particle-rich waters, where large, fast-sinking particles tend 446 to be relatively more abundant than small particles, and decrease them in particle-poor 447 waters where small particles dominate (see Section 3.2, and Supplementary Fig. S3). High 448 export in the eastern equatorial and tropical Pacific can be attributed to a substantial 449 increase in biovolume, with a minor contribution from PSD slope, which appears to be more uniform across the region. The picture is somewhat different in the equatorial At-451 lantic Ocean, where a more substantial flattening of the PSD suggests a more important 452 role of large particles in driving elevated export fluxes. A similar interaction of particle 453 abundance and size-structure dramatically intensify fluxes in the subpolar North Pacific 454 and Southern Ocean, and to a lesser extent the subpolar Atlantic, where a relative in-455 crease in particle abundance is followed by a shift of the PSD toward large particles. In 456 contrast, along many coastal regions, including eastern boundary upwelling systems and the Arabian Sea upwelling, increase in particle biovolume, rather than substantial changes 458 in size structure, appears to drive enhanced export fluxes. 459

Our reconstruction shows broad meridional patterns similar to previous estimates 460 (Fig. 6); however, significant regional-level discrepancies remain. For example, in the low 461 latitudes, we predict somewhat less intense equatorial export peaks and subtropical lows, compared to the estimates of Dunne et al. (2007) and Siegel et al. (2014). In this respect, 463 our reconstruction is more in line with the results of DeVries and Weber (2017). In the 464 subpolar Pacific, our estimate shows a northward shift of maximum export that is com-465 parable to the results of Dunne et al. (2007). This is likely caused by intensification of 466 particle fluxes in coastal waters and marginal seas, which may be related to regional pro-467 cesses such as iron leakage from shelves and marginal seas (Nishioka et al., 2020). In the 468 Atlantic Ocean, the gradual increase of export from northern to southern latitudes (mostly driven by high export near the coast), and the rapid increase in the Southern Ocean (caused 470 by high export near the Patagonian shelf), are similar to the reconstruction of Henson 471 et al. (2011), although the magnitude is larger. In the Indian Ocean, our reconstruction 472 matches other studies at low latitudes; however, it shows a more dramatic increase in 473 export towards the Southern Ocean sector (see also Section 3.3). 474

3.2.2 Seasonal cycle

475

The seasonal cycle of particle export is comparable to previous studies (Fig. 7), when averaged over large-scale coherent biomes (Weber et al., 2016). However, significant discrepancies are also revealed. In general, our seasonal cycle is more muted than previous work, suggesting weaker month-to-month variability in some regions, while other regions match previous reconstructions more closely.



Figure 6. Zonally integrated annual export for the (a) Global Ocean, (b) Pacific Ocean, (c) Atlantic Ocean, and (d) Indian Ocean. Each color represents a different study, as listed in the legend (bottom).

Similar to other estimates, we capture well-known seasonal export pulses associated with spring phytoplankton blooms in the North Atlantic and North Pacific Oceans. Over most of the tropics, our reconstruction reveals nearly constant seasonal cycles, and a slight asymmetry about the equator with more pronounced seasonality in the Northern Hemisphere compared to the Southern Hemisphere. The most significant discrepancy is observed in the Southern Ocean, in particular in the Antarctic Zone, where our reconstruction is substantially higher than previous estimates, with sustained export throughout winter months. We discuss this deviation in detail in Section 3.3.

To better quantify seasonality and provide a more objective comparison to other 489 studies, we define a seasonality index as the range of monthly export divided by the an-490 nual mean flux in each region (Supplementary Fig. S10). A higher seasonality index is 491 indicative of a more dynamic export cycle, with more dramatic variations between low 492 and high export periods. As expected, seasonality is larger in high latitudes, and decreases 493 toward the tropics; the highest values are reached in the mid-latitude to subpolar North 494 Pacific and Atlantic, and higher variability is confirmed in the tropics and subtropics of 495 the Southern Hemisphere compared to the Northern Hemisphere. 496

The relatively muted seasonality of our reconstruction, compared to previous work, is consistent with the weaker spatial gradients discussed in the previous sections, and suggests less dramatic gradients in net community production and export than previously assumed. The machine learning approach used in this study relies on non-linear relationships with multiple ocean variables to reconstruct particle export fluxes, which may reveal compensatory relationships between different predictors. Surface chlorophyll, temperature, and net primary production have all been used in previous global reconstructions (Dunne et al., 2007; Henson et al., 2011; Siegel et al., 2014), but rarely together



Figure 7. Annual seasonal cycle of particle flux from the euphotic zone for the regions specified in the map (top). Each line corresponds to a study shown in the legend (top right). The same seasonal spatial mask was applied to each study.

with other variables that may be important in modulating spatial and seasonal export

estimating variability in particle export fluxes derived from these quantities. When com-

patterns. It is also possible that our method somewhat underestimates variability com-

pared to previous work. As previously noted (see Section 3.1), RF ensemble reconstruc-

tions appear to reduce extremes in both biovolume and PSD slope, potentially under-

pared with other studies, our results show overall similar patterns in seasonal variabil-

ity, but lower seasonality in most regions, particularly at higher latitudes (Fig. 7 and Supplementary Fig. S10).

⁵¹³ 3.3 Southern Ocean Export

Export flux in the Antarctic Zone are substantially larger than other global recon-514 structions, especially during winter (Fig. 7). A regional study based on 10 years of bio-515 Argo measurements from 2006-2014, combined with satellite-based net primary produc-516 tion and export algorithms, similarly suggests higher than previously reported particle 517 fluxes throughout the region (Arteaga et al., 2018), in better agreement with our results 518 (Fig. 8). This similarity is mostly evident in the open ocean, and varies depending on 519 the primary production algorithm chosen for the comparison. However, our estimate also reveals substantially higher export near landmasses, for example the Patagonian Shelf, 521 South Georgia and the South Sandwich Islands, and the Kerguelen Plateau. Although 522 estimates from Arteaga et al. (2018) do not show the same high flux in winter as our re-523 construction, they do demonstrate that export fluxes from the Antarctic Zone of the South-524 ern Ocean likely never decrease to the nearly negligible levels shown by other global es-525 timates in winter (Fig. 7). 526

This discrepancy with prior global estimates in Antarctic Zone export could be due 527 to a variety of factors. First, UVP5 observations in the Southern Ocean, particularly in 528 winter, are extremely scarce. Similarly, satellite-based observations of predictors based 529 on ocean color dwindle in wintertime. Other climatological variables, such nutrients and 530 oxygen, are also the results of interpolation of fewer *in situ* observations relative to the 531 rest of the ocean. The scarcity of observations to train our model result in strong inter-532 model variability, highlighting the spread between different RF realizations. Second, our 533 reconstruction reveals significant export primarily next to land masses in the Atlantic 534 and Indian sectors of the Southern Ocean. Proximity to land masses has been shown to 535 increase productivity and carbon flux (Jouandet et al., 2014), presumably via iron fer-536 tilization from terrestrial and sedimentary sources in otherwise high-nutrient low-chlorophyll 537 waters. It is possible that other flux reconstructions underestimate this rapid aggrega-538 tion of particles and increased particle export, in particular during winter, when obser-530 vations are scarce. Increasing the number of particle flux and size distribution observations from the Antarctic Zone, in particular downstream of major land masses and in 541 wintertime, together with a better regional understanding of export processes, could help 542 shed more light on the patterns of export in the region. 543

3.4 Mixed layer versus euphotic zone export

544

By exploiting the high vertical resolution of UVP5 measurements, our approach 545 allows reconstruction of particle export at arbitrary horizons. We illustrate this capability by estimating particle fluxes at the depth of the maximum climatological winter-547 time mixed layer, and compare them to export from the climatologically-varying euphotic 548 zone, shown in Fig. 5. We find that, globally integrated, the particulate carbon export 549 from the mixed layer is $9.4 \pm 1.1 \text{ PgC/y}$, i.e., about 3 Pg/year larger than the global ex-550 port from the euphotic zone. This estimate is slightly lower than observational estimates 551 of carbon export and sequestration from the same depth horizon based on ANCP and 552 POC export estimates (Emerson, 2013). 553

Overall, export at the wintertime mixed layer follows broad spatial patterns sim-554 ilar to the export from the euphotic zone (Fig. 9a). However, tropics and subtropics show 555 larger export fluxes from the mixed layer (locally, up to a few times), while high latitudes 556 show overall weaker export fluxes (Fig. 9b). The low-latitude intensification of mixed 557 layer particle fluxes is similar in all ocean basins, and more than compensates for the reduction at high latitudes (Supplementary Fig. S12), thus producing an overall larger ex-559 port from this horizon. Because of this low-latitude intensification, export from the mixed 560 layer shows stronger gradients between the tropics and high latitudes. Gradients between 561 the equatorial export peak and the subtropical export low are also intensified. Finally, 562



Figure 8. Southern Ocean specific export for (a) This study, and (b-e) the various models from Arteaga et al. (2018), and (f) the mean from that study. Each model from Arteaga et al. (2018) represents a different net primary productivity algorithm used to derive export. (g) Seasonal cycle of export for each model in the Antarctic zone (shown in figure 7).

export from the mixed layer in the Southern Ocean is substantially depressed compared to export from the euphotic zone.

Differences between euphotic zone and mixed layer export can be best interpreted 565 by considering the different depth of these horizons (Palevsky & Doney, 2018). The depth 566 of the maximum mixed layer is shallower than the euphotic zone in the tropics and sub-567 tropics, and is deeper in high latitudes (Fig. 9c). This suggests that shallower export hori-568 zons are generally characterized by higher fluxes than deeper export horizons, likely be-569 cause of remineralization of particles in the upper layers of the ocean. Specifically, we 570 identify three main latitudinal bands with different horizon depths and export patterns, 571 roughly corresponding to tropics and subtropics, mid-latitudes, and subpolar regions. 572 Over most of the tropics and the subtropics, the maximum wintertime mixed layer is shal-573 lower on average than the climatological euphotic zone (blue colors in Fig. 9c). Here, 574 particle remineralization between the wintertime mixed layer and the euphotic zone depth 575 reduces export from the latter horizon, indicating net heterotrophy in the deeper lay-576 ers of the euphotic zone. Over subpolar regions, the wintertime mixed layer is deeper 577 on average than the climatological euphotic zone. Here, export fluxes reach maximum values within the euphotic zone, and decrease below it following remineralization. Fi-579 nally, over most of mid-latitudes, the wintertime mixed layer is deeper on average than 580 the climatological euphotic zone; however, export fluxes from the mixed layer and eu-581 photic zone depth are very similar in magnitude, suggesting a close seasonal compen-582



Figure 9. (a) Annual mean particle export from the maximum mixed layer depth. (b) Ratio of the export from the maximum mixed layer depth to the export from the euphotic zone. (c) Ratio of the maximum mixed layer depth to the euphotic zone depth.

sations between enhanced euphotic zone fluxes when this horizon is found above the win tertime mixed layer, and reduced euphotic zone fluxes when it is found below it.

Ultimately, differences in export between the euphotic zone and the wintertime mixed layer are important when considering the role of the biological pump for carbon sequestration (Palevsky & Doney, 2018). Our results suggest that more carbon is sequestered below the wintertime mixed layer than leaves the euphotic zone. The sensitivity of these fluxes to climate variability and change is also likely to differ (Palevsky & Doney, 2021).

⁵⁹⁰ 3.5 Caveats to our approach

There are multiple sources of uncertainty and inherent limitations that could affect our estimates and call for further work. First, expanding the coverage of observations with UVP5 and similar instruments, in particular in under-sampled regions characterized by large variability, such as the Southern Ocean, would improve the robustness of our estimates, and shed additional light on regional export patterns not captured
by previous work. Regional correlations between environmental properties and PSD may
not be well captured by extrapolation with a RF algorithm trained on data from different regions, especially when non-linear relationships between variables are important.

Second, supervised learning methods are only as reliable as the data used for training; therefore, continued work on improving satellite reconstructions of surface chlorophyll, net primary production, and other remotely-sensed variables, in particular at high latitudes, would help improve the robustness of these methods.

Third, different machine learning approaches are likely characterized by different biases. Here, we note a slight underestimate of extreme values in reconstructed PSD properties, which may affect the reconstructed variability in particle fluxes. Different machine learning methods that are better suited to capture extreme values should be tested to address this limitation, in conjunction with more detailed analyses of particle flux observations, including at time-series stations.

Finally, we use globally-averaged relationships between particle size, sinking speed, and carbon content, which we optimize against observations. However, these quantities are likely to depend on region and time of the year, reflecting variable underlying ecological processes. More work combining *in situ* and optical measurements should focus on constraining these quantities and their regional and temporal variability.

614 4 Conclusions

In this paper, we provide a new, data-constrained estimate of particle fluxes based on global UVP5 observations of PSDs. It captures regional and seasonal variability in observed PSD properties and export fluxes, and demonstrates the ability of statistical machine learning methods to extrapolate these quantities globally. Our approach also allows reconstruction of export fluxes from multiple depths, highlighting the importance of the choice of export horizon (Palevsky & Doney, 2018), and paving the way to fully three-dimensional particle flux reconstructions.

We obtain a global particle export flux of 6.7 ± 0.4 PgC/year from the euphotic 622 zone, in line with previous work, although with significant regional and temporal differ-623 ences. Our results suggest weaker spatial and seasonal variability in particle fluxes com-624 pared to previous studies, in particular in the open ocean, while highlighting the impor-625 tance of coastal waters and marginal seas for high latitude export. We also capture sim-626 ilar patterns of high latitude seasonal blooms in the Northern Hemisphere as previous studies, but less variable flux in the tropical to subtropical ocean, and substantially higher 628 year-round export in the Southern Ocean, in better agreement with regional estimates 629 (Arteaga et al., 2018). Results from the Southern Ocean suggest that processes that sus-630 tain elevated fluxes, in particular in wintertime, may not be completely captured by other 631 global reconstructions, and that waters downstream of coasts and islands may harbor 632 a significant source of carbon export to the deep ocean, which is only partially captured 633 in one other model (Dunne et al., 2007). 634

The statistical nature of our machine learning approach does not directly reveal 635 mechanisms behind PSD and export fluxes. However, we are able to highlight globally 636 coherent patterns, and the relative roles of particle abundance and size structure in driv-637 ing export fluxes. Specifically, we show that the total particle biovolume and the PSD 638 slope show similar patterns, and are in fact correlated in such a way to act synergisti-630 cally on particle fluxes (Supplementary Fig. S3). We also suggest distinct deviation from 640 these patterns, for example in the tropical and northern subtropical Pacific Ocean, where 641 high availability of particles, rather than dominance of large vs. small particles, tends 642 to drive elevated export. 643

We illustrate the ability of our method to reconstruct particle fluxes at multiple depth horizons by reconstructing and comparing carbon export from the euphotic zone and the wintertime mixed layer depth. Export from the base of the wintertime mixed layer is overall stronger than export from the euphotic zone in low and mid latitudes, and weaker in high latitudes, driving a significantly larger global particle export of 9.4 ± 1.1 PgC/year. This suggests that more carbon is sequestered for decadal to centennial timescales than is available for remineralization in the twilight zone of the ocean.

We attribute the disparity in flux from the euphotic zone and mixed layer to differences in their depths: maximum mixed layers are typically shallower than the euphotic zone over large swathes of the low latitude ocean, but deeper in high latitudes. Future three-dimensional reconstructions of particle fluxes would allow a closer investigation of the processes controlling transfer of carbon between the surface and the deep ocean, providing further insights into particle transformation processes, transfer efficiency, and ocean carbon sequestration.

658 Acknowledgments

This material is based upon work supported by the U.S. National Science Foundation 659 under Grant No. OCE-1635632. D.B. acknowledges support from the Alfred P. Sloan 660 Foundation. A.M.P.M acknowledges support from NSF Award No. 1654663. T.W. was supported by NSF award OCE-1635414. RK acknowledges support via the BMBF funded 662 project CUSCO, the EU project TRIATLAS (European Union's Horizon 2020 programme, 663 grant agreement No 817578) and a "Make Our Planet Great Again" grant of the ANR 664 within the "Programme d'Investissements d'Avenir"; reference "ANR-19-MPGA-0012". 665 Data generated by this analysis has been uploaded to BCO-DMO, with a DOI pending. 666 The individual UVP5 profiles used to generate the reconstructions can be obtained on 667 the EcoTaxa website https://ecotaxa.obs-vlfr.fr/part/. 668

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