Solar Wind Interaction and Pressure Balance at the Dayside Ionopause of Mars

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Abstract

Due to the lower ionospheric thermal pressure and existence of the crustal magnetism at Mars, the Martian ionopause is expected to behave differently from the ionopause at Venus. We study the solar wind interaction and pressure balance at the ionopause of Mars using both in situ and remote sounding measurements from the MARSIS (Mars Advanced Radar for Subsurface and Ionosphere Sounding) instrument on the Mars Express orbiter. We show that the magnetic pressure usually dominates the thermal pressure to hold off the solar wind in the ionopause at Mars, with only 13% unmagnetized ionopauses observed over a 11-year period. We also find that the ionopause altitude decreases as the normal component of the solar wind dynamic pressure increases. Moreover, our results show that the ionopause thickness at Mars is mainly determined by the ion gyromotion and equivalent to about 5.7 ion gyroradii.

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9	Key Points:
10	+ 13% of the ionopauses at Mars are unmagnetized; the unmagnetized ionopauses
11	occur 65% at Venus at solar maximum and 52% at solar minimum
12	• The ionopause altitude decreases with the solar wind dynamic pressure at Mars,
13	similar to the altitude variation of the ionopauses at Venus
14	• The ionopause thickness at Mars is mainly determined by the ion gyromotion and
15	equivalent to about 5.7 ion gyroradii

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16 Abstract

Due to the lower ionospheric thermal pressure and existence of the crustal magnetism 17 at Mars, the Martian ionopause is expected to behave differently from the ionopause at 18 Venus. We study the solar wind interaction and pressure balance at the ionopause of Mars 19 using both in situ and remote sounding measurements from the MARSIS (Mars Advanced 20 Radar for Subsurface and Ionosphere Sounding) instrument on the Mars Express orbiter. 21 We show that the magnetic pressure usually dominates the thermal pressure to hold off 22 the solar wind in the ionopause at Mars, with only 13% unmagnetized ionopauses ob-23 served over a 11-year period. We also find that the ionopause altitude decreases as the 24 normal component of the solar wind dynamic pressure increases. Moreover, our results 25 show that the ionopause thickness at Mars is mainly determined by the ion gyromotion 26 and equivalent to about 5.7 ion gyroradii. 27

²⁸ Plain Language Summary

An ionopause is a sharp decrease in the plasma density at the top of the ionosphere. 29 It was found to be a common feature at Venus, where its variability is well constrained 30 by observations from the Pioneer Venus Orbiter (PVO). Past studies have shown that 31 there are many similarities between the ionopauses at Mars and Venus. However, because 32 the thermal pressure of the ionosphere at Mars is lower than that of Venus, the Martian 33 ionopause is also thought to behave differently from the ionopause at Venus. We study 34 the pressure configuration inside the ionopause at Mars and find that most of the time 35 the magnetic pressure is greater than the thermal pressure. We also find that higher so-36 lar wind pressure pushes the ionopause downward at Mars. Moreover, we show that the 37 thickness of the ionopauses at Mars equals a few radii of ion circular motion in the mag-38 netic field. Our results provide insight to the process that controls the formation of the 39 ionopause at Mars. 40

41 **1** Introduction

Since Mars does not possess a strong global intrinsic magnetic field, the incident solar wind plasma and interplanetary magnetic field interacts inductively with its upper atmosphere and highly conductive ionosphere. This interaction induces currents that produce a magnetic barrier to prevent the solar wind from further penetrating into the atmosphere (Ramstad et al., 2020), resulting in the formation of several plasma boundaries around Mars, such as the magnetic pileup boundary (Crider et al., 2002; Bertucci
et al., 2004, 2005), photoelectron boundary (Garnier et al., 2017; Q. Han et al., 2019),

and the ionopause (Nagy et al., 2004; X. Han et al., 2014).

The Martian ionopause is a feature identified as a steep gradient in electron den-50 sity at the topside of the ionosphere. It is a tangential discontinuity that marks the tran-51 sition from the hot plasma in the induced magnetosphere to the cold, dense ionospheric 52 plasma. When the solar wind flows around Mars, it exerts its dynamic pressure indirectly 53 on the ionosphere through the magnetosheath and magnetic pileup region. Therefore, 54 as the ionospheric thermal pressure decreases sharply across the ionopause, a magnetic 55 pressure from an intrinsic or induced field is required to maintain a steady state of the 56 ionosphere (Holmberg et al., 2019; Sánchez-Cano et al., 2020). 57

Since Venus also lacks a global-scale magnetic field and its upper atmosphere in-58 teracts directly with the solar wind, the ionopauses at Venus and Mars are similar in many 59 aspects. For example, solar wind conditions can heavily influence the dynamics of the 60 ionopauses at both planets (Phillips et al., 1985; Sánchez-Cano et al., 2020). Moreover, 61 the altitudes of Venusian and Martian ionopauses both vary over the solar cycle as a re-62 sult of the periodic effects of solar EUV on the ionospheric thermal pressure (Kliore & 63 Luhmann, 1991; Duru et al., 2020). Thanks to the Pioneer Venus Orbiter (PVO) mis-64 sion, the studies of the ionopause at Venus have provided many deep insights that can 65 be applied to understanding the formation of the ionopause at Mars. 66

On the other hand, because of the lower ionospheric thermal pressure and presence 67 of the highly-localized crustal magnetic fields at Mars, the behavior of the Martian ionopause 68 is also quite different from that at Venus. As the peak thermal pressure in the ionosphere 69 at Mars rarely exceeds the solar wind dynamic pressure, unlike Venus, the Martian iono-70 sphere is often found to be magnetized (Nagy et al., 2004). The magnetic pressure, there-71 fore, plays a more important role in standing off the solar wind ram pressure in the ionopause 72 at Mars. In addition, Chu et al. (2019) found that the strong crustal fields act as min-73 imagnetospheres that prevent the ionopause from forming. 74

In the past, a number of studies have been dedicated to the investigation of the dependence of the ionopause altitude on solar zenith angle (SZA) and solar extreme ultraviolet (EUV) flux, as well as the effects of the crustal magnetic fields on the formation of the ionopause at Mars (Vogt et al., 2015; Chu et al., 2019; Duru et al., 2020). In this

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paper, we take advantage of both in situ and remote sounding measurements from the 79 Mars Advanced Radar for Subsurface and Ionosphere Sounding (MARSIS) instrument 80 on board the Mars Express (MEX) spacecraft to study the pressure configuration and 81 balance in the Martian ionopause. We also report for the first time on the mechanisms 82 that control the thickness of the ionopause at Mars. The paper is organized as follows: 83 section 2 gives a description of the ionopause observations from the MARSIS instrument, 84 section 3 explains the pressure terms used in the analysis, section 4 presents the results, 85 and section 5 gives conclusions of the paper. 86

⁸⁷ 2 Ionopause Observations

The MARSIS instrument on board the Mars Express spacecraft is a low-frequency radar sounder designed to perform both subsurface and ionospheric soundings (Picardi et al., 2004). In this study, the ionopause is detected by MARSIS using two different techniques – remote radar sounding and in situ measurements.

To obtain the plasma density profile of the ionosphere, MARSIS sends a short ra-92 dio pulse from 0.1 to 5.4 MHz and detects any return echoes that are reflected off the 93 ionosphere (Gurnett et al., 2005). The reflection occurs at the point where the frequency 94 of the electromagnetic wave is below the local electron plasma frequency $f_{\rm p} = 8980 \sqrt{n_{\rm e}}$ Hz, 95 where $n_{\rm e}$ is the electron density in cm⁻³ (Gurnett & Bhattacharjee, 2005). By measur-96 ing the time delay Δt between the transmission of the pulse and the time that the echo 97 is received, the apparent altitude of the reflection point can be expressed as $h=h_{\rm MEX}-$ 98 $c\Delta t/2$, where $h_{\rm MEX}$ is the spacecraft altitude and c is the speed of light. The term ap-99 parent altitude refers to the scale that has not been corrected for dispersion of the radar 100 pulses that propagate in a plasma; however, in our situation where the dispersion effects 101 are small, the apparent altitude is very close to the real altitude. The detailed discus-102 sion about the apparent altitude can be found in Chu et al. (2019). MARSIS remote sound-103 ing data are displayed in ionograms, the intensity of the return echoes as a function of 104 frequency and time delay. An example of an ionospheric sounding ionogram is presented 105 in Figure 1a, where the ionopause can be seen as a horizontal line at frequencies below 106 0.4 MHz, representing a steep density change over a short vertical distance (Chu et al., 107 2019). 108

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The ionopause can also be identified through MARSIS in situ density measurements 109 (Duru et al., 2009). When MARSIS sounds the ionosphere, intense electrostatic electron 110 plasma oscillations can be excited at the local electron plasma frequency surrounding 111 the spacecraft (Gurnett et al., 2008). When these oscillations are picked up by MAR-112 SIS, the received waveforms are often severely clipped, resulting in closely spaced ver-113 tical harmonic lines in the low frequency region of the ionogram as shown in Figure 1b. 114 The local electron density can then be calculated from the harmonic spacing Δf (in the 115 units of Hz) using $n_{\rm e} = (\Delta f/8980)^2$ cm⁻³. As MEX enters or exits the topside iono-116 sphere, one can identify the ionopause by looking for the signature of steep density gra-117 dient in the local electron density profile. 118

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3 Pressure Terms at the Ionopause

To investigate the solar wind interaction and pressure balance at the ionopause, 120 we need to evaluate the pressures that are exerted on the ionopause. In this study specif-121 ically, we consider three different pressure terms, the thermal pressure of the ionosphere, 122 magnetic pressure, and the solar wind dynamic pressure. 123

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3.1 Ionospheric Thermal Pressure

Based on MARSIS in situ electron density measurements, the thermal pressure of 125 the ionosphere can be estimated as 126

$$P_{\rm th} = n_{\rm i} k T_{\rm i} + n_{\rm e} k T_{\rm e} \approx 2 n_{\rm e} k T_{\rm e},\tag{1}$$

where T_i and T_e are the ion and electron temperature, respectively, n_i is the ion density 127 $(n_{
m i}~pprox~n_{
m e}),$ and k is the Boltzmann constant. Here we assume equal ion and electron 128 temperature $T_{\rm i} \approx T_{\rm e}$ as the first order approximation for at least up to 350 km (Hanson 129 & Mantas, 1988; Matta et al., 2014). We also use a fixed representation of the $T_{\rm e}$ pro-130 file that does not account for potential SZA, latitude, seasonal, or solar activity varia-131 tions (Ergun et al., 2015; Pilinski et al., 2019). 132

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3.2 Magnetic Pressure in the Ionosphere

In addition to the vertical electron plasma oscillation harmonics discussed in sec-134 tion 2, another commonly found feature in many MARSIS ionograms is a series of equally 135 spaced horizontal echoes in time at frequencies below 1 MHz, as shown in Figure 1b. When 136



Figure 1. (a) An example of the color-coded MARSIS ionogram from the orbit 2360 on 14 November 2005. The ionopause can be seen as a horizontal line at frequencies below 0.4 MHz. (b) Another ionogram from the same orbit showing the features of the electron plasma oscillation harmonics and electron cyclotron echoes. (c) Schematic illustration of the pressure terms at the ionopause. $P_{\rm th}$ is the ionospheric thermal pressure, $P_{\rm Bi}$ is the magnetic pressure in the ionosphere, $P_{\rm Bp}$ is the magnetic pressure in the magnetic pileup region, and $P_{\rm sw}$ is the solar wind dynamic pressure.

electrons near the antenna are accelerated by the strong electric fields during each transmission cycle, they go through cyclotron motions in the local magnetic field and periodically return to the vicinity of the antenna, causing the electron cyclotron echoes to
appear in the ionogram (Gurnett et al., 2005). The repetition rate of these echoes is equal
to the local electron cyclotron frequency

$$f_{\rm c} = \frac{Be}{2\pi m_{\rm e}},\tag{2}$$

where B is the magnetic field strength, e is the electron charge, and $m_{\rm e}$ is the electron mass. Since MEX is not equipped with a magnetometer, these electron cyclotron echoes provide the only method to measure the local magnetic field.

As the peak thermal pressure in the Mars ionosphere rarely exceeds the solar wind dynamic pressure, the dayside ionosphere is often found to be magnetized in order to stand off the solar wind (Zhang et al., 1990; Nagy et al., 2004). Over the regions away from the strong crustal magnetism, we can assume that the magnetic field is approximately tangential to the ionopause. Thus, the magnetic pressure in the ionosphere normal to the ionopause can be estimated to the first order as

$$P_{\rm Bi} = \frac{B^2}{2\mu_0},\tag{3}$$

where μ_0 is the permeability of free space.

3.3 Solar Wind Dynamic Pressure

Since MEX does not directly measure the properties of the solar wind, the solar wind dynamic pressure P_{sw} is estimated based on the ASPERA (Analyzer of Space Plasma and Energetic Atoms) solar wind moments, which are calculated from averaged proton distributions collected over the inbound/outbound segments of MEX outside the bow shock (Barabash et al., 2006; Ramstad et al., 2015, 2017)

$$P_{\rm sw} = n_{\rm p} m_{\rm p} v_{\rm p}^2, \tag{4}$$

where $n_{\rm p}$ is the proton density, $m_{\rm p}$ is the proton mass, and $v_{\rm p}$ is the speed of the solar wind.

Note that the normal component of the solar wind ram pressure is not directly exerted on the ionopause; rather it is first converted to thermal pressure in the magnetosheath, then to magnetic pressure in the pileup region. The normal component of the solar wind ¹⁶³ dynamic pressure can be written as

$$P_{\rm sw\perp} = \alpha P_{\rm sw} \cos^2 \theta, \tag{5}$$

where θ is the angle between the magnetic pileup boundary normal and the flow direction of the upstream solar wind, and $\alpha \approx 0.88$ is the proportionality constant (Crider et al., 2003). For lower solar zenith angles, θ can be approximately replaced by SZA. At higher SZAs, however, this approximation breaks down because the curvature of the obstacle must be accounted. The pressure values used in the study are the closest ones to the time when the ionopause is observed within 6 hours.

¹⁷⁰ A schematic illustration of the three pressure terms $P_{\rm th}$, $P_{\rm Bi}$, and $P_{\rm sw}$ is shown in ¹⁷¹ Figure 1c. In theory, a pressure balance across the ionopause requires

$$P_{\rm th} + P_{\rm Bi} = P_{\rm sw\perp}.\tag{6}$$

In other words, the ionospheric thermal pressure and magnetic pressure should stand off the normal component of the solar wind dynamic pressure. A detailed examination of this pressure balance relation will be presented in section 4.

175 4 Results

In this study, we utilize both MARSIS in situ (79 detections) and remote sound-176 ing (1,791 detections) measurements, excluding crustal magnetic field regions, to study 177 the pressure balance in Martian ionopauses and their interactions with the solar wind. 178 The descriptions of these datasets can be found in Chu et al. (2019) and Duru et al. (2020). 179 We first compare the pressure configuration in the ionopause at Venus and Mars, and 180 then investigate the role of the solar wind in the formation of the Martian ionopause. 181 Finally, we study the dependence of the ionopause thickness at Mars on altitude and mag-182 netic field strength. 183

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4.1 Comparison of Pressure Configuration in Ionopauses at Venus and Mars

The maximum ionospheric thermal pressure at Venus often exceeds the solar wind dynamic pressure (Zhang & Luhmann, 1992). The thermal pressure of the ionosphere, therefore, plays a dominant role in the pressure balance underneath the ionopause. Figure 2a shows a typical example of the pressure configuration in the ionopause at Venus



Figure 2. (a)–(b) Thermal pressure and magnetic pressure profiles in the Venusian ionopause as a function of the altitude. The data shown here are based on the in situ measurements from the Pioneer Venus Orbiter's Electron Temperature Probe. A typical example of the pressure configuration is shown in (a) and a rare case is shown in (b). (c)–(d) Thermal pressure and magnetic pressure profiles in the Martian ionopause as a function of the altitude. A rare, Venusian-like example of the pressure configuration is shown in (c) and a typical case is shown in (d). As an example, the red arrows in (d) mark the ionopause thickness and the red star represents the location where the total pressure $P_{\text{tot}} = P_{\text{th}} + P_{\text{Bi}}$ is measured.

(Orbit 168 on 21 May 1979) during solar maximum. Inside the ionopause, as the thermal pressure decreases, the magnetic pressure is forced to increase to maintain the pressure balance. There are some rare occasions shown in Figure 2b (Orbit 190 on 12 June
1979), where the magnetic pressure underneath the ionopause is around the same order
of magnitude as the thermal pressure. In such cases, both of the pressure terms are equally
important in balancing the solar wind dynamic pressure above the ionopause.

Due to the lower ionospheric thermal pressure at Mars, the behavior of the Mar-196 tian ionopause is expected to be different from the ionopause at Venus. Figure 2c shows 197 a rare, Venusian-like (typical pressure configuration at Venus) pressure configuration in 198 the Martian ionopause, only found in 13% of all the cases over a 11-year period. In con-199 trast, the similar pressure configuration at Venus shown in Figure 2a is observed at least 200 65% at solar maximum (Luhmann et al., 1980; Elphic et al., 1981) and ~ 52% at so-201 lar minimum (Angsmann et al., 2011). Figure 2d shows a typical pressure configuration 202 in the ionopause at Mars, where the magnetic pressure is clearly seen to play a domi-203 nant roll in the pressure balance anywhere in the ionopause. 204

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4.2 Influence of Solar Wind on Martian Ionopause

At Venus, the dynamic behavior of the ionopause is strongly affected by the up-206 stream solar wind conditions. In our previous study, we showed how crustal magnetism 207 and solar extreme ultraviolet flux control the ionopause formation at Mars (Chu et al., 208 2019). Here, we examine the influence of the solar wind dynamic pressure on ionopause 209 apparent altitude at Mars based on 1,791 ionopause detections (SZA $< 65^{\circ}$) obtained 210 using the MARSIS remote sounding technique (Figure 3a). Despite a strong scattering 211 of the data points in Figure 3a, likely due to the variations of seasons and solar extreme 212 ultraviolet (EUV) flux (more plots showing the distribution in season, EUV flux, lati-213 tude, and SZA can be found in the supporting information), we find that the ionopause 214 altitude decreases with the normal component of the solar wind dynamic pressure by $131\pm$ 215 50 km/nPa. Similar ionopause altitude variation with the solar wind dynamic pressure 216 was also observed at Venus by PVO (Brace et al., 1980). The trend shown in Figure 3a, 217 however, was not found in Vogt et al. (2015), possibly because they had a smaller data 218 set or a different detection method. These solar wind effects are also observed in other 219 plasma boundaries at Mars (Edberg et al., 2009; Withers et al., 2016; Garnier et al., 2017; 220 Halekas et al., 2018; Girazian, Halekas, et al., 2019). 221

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Figure 3. (a) Scatter plot of the ionopause apparent altitude (SZA < 65°) as a function of the normal component of the solar wind dynamic pressure at Mars. The ionopause data are collected using the MARSIS remote sounding technique. The mean ionopause altitude averaged over each 0.1 nPa bin is shown in blue dots. (b) Correlation between the total ionopause pressure (SZA < 65°) and normal component of the solar wind dynamic pressure at Mars. The mean ionopause pressure for each 0.05 nPa bin is shown in blue dots. Error bars in (a)–(b) represent the standard deviation of the mean (σ/\sqrt{N} , σ standard deviation and N number of data points in each bin). SZA=solar zenith angle.



Figure 4. (a) Ionopause thickness as a function of altitude at Venus, reproduced from Figure 6 in Elphic et al. (1981). The orange and green dots mark the ionopauses forming at low and high altitudes, respectively. (b) Ionopause thickness as a function of altitude at Mars, based on the MARSIS in situ measurements.

To test the theory for pressure balance in the ionopause, we plot the total pressure 222 inside the ionopause $(P_{\text{tot}} = P_{\text{th}} + P_{\text{Bi}})$, the location of which is measured is shown as 223 an example in Figure 2d) as a function of the normal component of the solar wind dy-224 namic pressure for SZA $< 65^{\circ}$ in Figure 3b. We then perform a linear fit (least squares) 225 for the average pressure in each pressure bin using the formula $P_{\text{tot}} = aP_{\text{sw}\perp} + b$, and 226 find that $a = 1.36 \pm 0.44$. The slope close to unity indicates that the total pressure 227 $P_{\rm tot}$ inside the ionopause on average balances the normal component of the solar wind 228 dynamic pressure $P_{\mathrm{sw}\perp}$, consistent with the similar test based on the simultaneous mea-229 surements by ASPERA and MAVEN (Mars Atmosphere and Volatile Evolution) in Sánchez-230 Cano et al. (2020). Additionally, we find that $b = 0.19 \pm 0.06$ nPa, a small offset be-231 tween P_{tot} and $P_{\text{sw}\perp}$, which suggests that P_{tot} may slightly exceed $P_{\text{sw}\perp}$ at the location 232 where the ionopause forms, in agreement with the results shown in Holmberg et al. (2019). 233

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4.3 Dependence of Ionopause Thickness on Altitude

We identify the thickness of a Martian ionopause as the length scale of the steep change in ionospheric plasma density greater than $\Delta n/n > 0.1$ within a single MAR-SIS frequency sweep period (7.54 s), equivalent to a minimum density gradient of ~ 2 cm⁻³/km. An example illustrating the ionopause thickness is shown in Figure 2d. Since the ionopause is essentially a current sheet induced to shield the solar wind magnetic field from penetrating into the upper ionosphere, the thickness of this boundary layer is expected to
be on the order of the ion gyroradius scale (Cravens & Shinagawa, 1991). However, there
are also other factors that can affect the ionopause thickness, such as diffusion (Elphic
et al., 1981).

At Venus, past observations from PVO showed that the ionopauses can be mainly 244 grouped into two classes based on the mechanisms that determine their thickness – the 245 thick ionopauses (thickness > 60 km) at low altitudes and thinner ones at high altitudes 246 (Elphic et al., 1981). The PVO measurements of the ionopause thickness as a function 247 of altitude are shown in Figure 4a (reproduced from Figure 6 in Elphic et al., 1981). When 248 the ionopause forms at low altitudes (colored in orange), due to relatively high ionospheric 249 density, the ion coulomb collision rate usually dominates the ion gyrofrequency in this 250 boundary layer, making diffusive broadening an important process in determining the 251 ionopause thickness. However, as the altitude increases, the coulomb collision rate falls 252 off rapidly, causing the thickness of the ionopause forming in high altitudes (colored in 253 green) simply proportional to the ion gyroradius. Since the variation of the ion temper-254 ature is small above 300 km, thus the thickness of these ionopauses is also inversely pro-255 portional to the magnetic field strength (Miller et al., 1980). Figure 5a shows the PVO 256 measurements of the ionopause thickness as a function of the field strength (reproduced 257 from Figure 5 in Elphic et al., 1981). If we rewrite the ion gyroradius as $\rho_{\rm i} = \sqrt{m_{\rm i} k T_{\rm i}}/Be$, 258 where m_i is the ion mass, by fitting the thickness of the ionopauses at high altitudes (col-259 ored in green) with formula $d = \gamma \rho_i$ and assuming the temperature of O⁺ ions $T_i \approx$ 260 3000 K, we find that on average the ionopause thickness is about $\gamma = 5.3 \pm 0.9$ ion gy-261 roradii at Venus (Elphic et al., 1981). 262

At Mars, however, the ionospheric density is much smaller than that at Venus. Pre-263 vious studies have shown that the ion coulomb collision rate is only comparable to the 264 ion gyrofrequency in the ionospheric dynamo region that lies between 100 and 250 km 265 in altitude, well below that of the ionopause (Withers, 2008; Opgenoorth et al., 2010). 266 Therefore, in contrast to the two main mechanisms that affect the thickness of Venusian 267 ionopauses at low and high altitudes, the ion gyroradius scale becomes the most impor-268 tant factor that determines the ionopause thickness at Mars. Figure 4b shows the ionopause 269 thickness as a function of altitude based on the MARSIS in situ measurements. In gen-270 eral, we find that the ionopause thickness increases with altitude at Mars due to lower 271

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Figure 5. (a) Ionopause thickness as a function of magnetic field strength at Venus, reproduced from Figure 5 in Elphic et al. (1981). Same as Figure 4a, the orange and green dots represent the ionopauses forming at low and high altitudes, respectively. (b) Ionopause thickness as a function of magnetic field strength at Mars, based on the MARSIS in situ measurements. The red curves in (a) and (b) represent the best fit of the green dots with the formula $d = \gamma \sqrt{m_{i}kT_{i}}/Be \propto 1/B$.

magnetic field strength at higher altitudes (Holmberg et al., 2019). Figure 5b shows the ionopause thickness as a function of the field strength. If we assume the topside ionospheric composition is about 50% O⁺ ions and 50% O⁺₂ ions, and $T_i \approx T_e \approx 3000$ K, by repeating the same procedure as in Figure 5a, we find that the ionopause at Mars has a thickness of 5.7±0.8 ion gyroradii, comparable to 5.3 ion gyroradii at Venus (Girazian, Mahaffy, et al., 2019).

²⁷⁸ 5 Conclusions

In conclusion, we have investigated the solar wind interaction and pressure balance 279 at the dayside ionopause of Mars using both in situ and remote sounding measurements 280 from the MARSIS instrument. We have found that most of the time the magnetic pres-281 sure dominates the thermal pressure to hold off the solar wind in the ionopause at Mars. 282 Only about 13% of the ionopauses that we examined over a 11-year period are unmag-283 netized, whereas the unmagnetized ionopauses account for at least 52% at Venus. Ad-284 ditionally, our analysis has shown that the ionopause altitude decreases as the solar wind 285 dynamic pressure increases at Mars, similar to the altitude variation of the ionopauses 286

at Venus. Finally, we have shown for the first time that the thickness of the ionopauses

at Mars is mainly determined by the ion gyromotion, much alike the ionopauses form-

ing in high altitudes at Venus. The ionopauses at Mars are found to have a thickness of

about 5.7 ion gyroradii, surprisingly close to the ionopause thickness of 5.3 ion gyroradii
 at Venus.

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299 References

- Angsmann, A., Fränz, M., Dubinin, E., Woch, J., Barabash, S., Zhang, T. L., &
- Motschmann, U. (2011, March). Magnetic states of the ionosphere of Venus observed by Venus Express. *Planetary and Space Science*, 59(4), 327–337. doi: 10.1016/j.pss.2010.12.004
- Barabash, S., Lundin, R., Andersson, H., Brinkfeldt, K., Grigoriev, A., Gunell, H., ... Thocaven, J.-J. (2006, October). The Analyzer of Space Plasmas and
- Energetic Atoms (ASPERA-3) for the Mars Express Mission. Space Sci Rev,
 126(1), 113–164. doi: 10.1007/s11214-006-9124-8
- Bertucci, C., Mazelle, C., Acuña, M. H., Russell, C. T., & Slavin, J. A. (2005).
 Structure of the magnetic pileup boundary at Mars and Venus. J. Geophys.
 Res. Space Phys., 110(A1). doi: 10.1029/2004JA010592
- Bertucci, C., Mazelle, C., Crider, D. H., Mitchell, D. L., Sauer, K., Acuña, M. H.,
- ... Winterhalter, D. (2004, January). MGS MAG/ER observations at
 the magnetic pileup boundary of Mars: Draping enhancement and low
 frequency waves. Advances in Space Research, 33(11), 1938–1944. doi:
- ³¹⁵ 10.1016/j.asr.2003.04.054
- Brace, L. H., Theis, R. F., Hoegy, W. R., Wolfe, J. H., Mihalov, J. D., Russell,
 C. T., ... Nagy, A. F. (1980). The dynamic behavior of the Venus ionosphere

318	in response to solar wind interactions. J. Geophys. Res. Space Phys., $85(A13)$,
319	7663–7678. doi: $10.1029/JA085iA13p07663$
320	Chu, F., Girazian, Z., Gurnett, D. A., Morgan, D. D., Halekas, J., Kopf, A. J.,
321	Duru, F. (2019). The Effects of Crustal Magnetic Fields and Solar EUV Flux
322	on Ionopause Formation at Mars. Geophys. Res. Lett., $46(17-18)$, 10257–10266.
323	doi: 10.1029/2019GL083499
324	Cravens, T. E., & Shinagawa, H. (1991). The ionopause current layer at Venus. J .
325	Geophys. Res. Space Phys., $96(A7)$, 11119–11131. doi: 10.1029/91JA00674
326	Crider, D. H., Acuña, M. H., Connerney, J. E. P., Vignes, D., Ness, N. F., Krymskii,
327	A. M., Winterhalter, D. (2002). Observations of the latitude dependence
328	of the location of the martian magnetic pileup boundary. Geophys. Res. Lett.,
329	29(8), 11-1-11-4. doi: 10.1029/2001GL013860
330	Crider, D. H., Vignes, D., Krymskii, A. M., Breus, T. K., Ness, N. F., Mitchell,
331	D. L., Acuña, M. H. (2003). A proxy for determining solar wind dynamic
332	pressure at Mars using Mars Global Surveyor data. J. Geophys. Res. Space
333	<i>Phys.</i> , 108 (A12). doi: 10.1029/2003JA009875
334	Duru, F., Baker, N., De Boer, M., Chamberlain, A., Verchimak, R., Morgan, D. D.,
335	Kopf, A. (2020, May). Martian Ionopause Boundary: Coincidence With
336	Photoelectron Boundary and Response to Internal and External Drivers. Jour-
337	nal of Geophysical Research: Space Physics, 125(5), e2019JA027409. doi:
338	10.1029/2019JA027409
339	Duru, F., Gurnett, D. A., Frahm, R. A., Winningham, J. D., Morgan, D. D., &
340	Howes, G. G. (2009). Steep, transient density gradients in the Martian iono-
341	sphere similar to the ionopause at Venus. J. Geophys. Res. Space Phys.,
342	114(A12). doi: 10.1029/2009JA014711
343	Edberg, N. J. T., Brain, D. A., Lester, M., Cowley, S. W. H., Modolo, R., Fränz, M.,
344	& Barabash, S. (2009, September). Plasma boundary variability at Mars as
345	observed by Mars Global Surveyor and Mars Express. Ann. Geophys., 27(9),
346	3537–3550. doi: 10.5194/angeo-27-3537-2009
347	Elphic, R. C., Russell, C. T., Luhmann, J. G., Scarf, F. L., & Brace, L. H. (1981).
348	The Venus ionopause current sheet: Thickness length scale and control-
349	ling factors. J. Geophys. Res. Space Phys., 86(A13), 11430–11438. doi:
350	10.1029/JA086iA13p11430

351	Ergun, R. E., Morooka, M. W., Andersson, L. A., Fowler, C. M., Delory, G. T.,
352	Andrews, D. J., Jakosky, B. M. (2015). Dayside electron tempera-
353	ture and density profiles at Mars: First results from the MAVEN Langmuir
354	probe and waves instrument. Geophys. Res. Lett., $42(21)$, 8846–8853. doi:
355	10.1002/2015GL065280
356	Garnier, P., Steckiewicz, M., Mazelle, C., Xu, S., Mitchell, D., Holmberg, M. K. G.,
357	\dots Jakosky, B. M. (2017). The Martian Photoelectron Boundary as Seen
358	by MAVEN. J. Geophys. Res. Space Phys., 122(10), 10,472–10,485. doi:
359	10.1002/2017JA024497
360	Girazian, Z., Halekas, J., Morgan, D. D., Kopf, A. J., Gurnett, D. A., & Chu, F.
361	(2019). The Effects of Solar Wind Dynamic Pressure on the Structure of the
362	Topside Ionosphere of Mars. Geophys. Res. Lett., $46(15)$, 8652–8662. doi:
363	10.1029/2019GL083643
364	Girazian, Z., Mahaffy, P., Lee, Y., & Thiemann, E. M. B. (2019, April). Seasonal,
365	Solar Zenith Angle, and Solar Flux Variations of O+ in the Topside Ionosphere
366	of Mars. Journal of Geophysical Research: Space Physics, 124(4), 3125–3138.
367	doi: 10.1029/2018JA026086
368	Gurnett, D. A., & Bhattacharjee, A. (2005). Introduction to Plasma Physics: With
369	Space and Laboratory Applications. Cambridge University Press.
370	Gurnett, D. A., Huff, R. L., Morgan, D. D., Persoon, A. M., Averkamp, T. F.,
371	Kirchner, D. L., Picardi, G. (2008, January). An overview of radar sound-
372	ings of the martian ionosphere from the Mars Express spacecraft. $Advances$ in
373	Space Research, 41(9), 1335–1346. doi: 10.1016/j.asr.2007.01.062
374	Gurnett, D. A., Kirchner, D. L., Huff, R. L., Morgan, D. D., Persoon, A. M.,
375	Averkamp, T. F., Picardi, G. (2005, December). Radar Sound-
376	ings of the Ionosphere of Mars. Science, $310(5756)$, 1929–1933. doi:
377	10.1126/science.1121868
378	Halekas, J. S., McFadden, J. P., Brain, D. A., Luhmann, J. G., DiBraccio, G. A.,
379	Connerney, J. E. P., Jakosky, B. M. (2018, October). Structure and Vari-
380	ability of the Martian Ion Composition Boundary Layer. Journal of Geophysi-
381	cal Research: Space Physics, 123(10), 8439–8458. doi: 10.1029/2018JA025866
382	Han, Q., Fan, K., Cui, J., Wei, Y., Fraenz, M., Dubinin, E., Connerney, J. (2019,
383	October). The Relationship Between Photoelectron Boundary and Steep Elec-

384	tron Density Gradient on Mars: MAVEN Observations. J. Geophys. Res. Space
385	<i>Physics</i> , $124(10)$, 8015–8022. doi: 10.1029/2019JA026739
386	Han, X., Fraenz, M., Dubinin, E., Wei, Y., Andrews, D. J., Wan, W., Barabash,
387	S. (2014). Discrepancy between ionopause and photoelectron boundary de-
388	termined from Mars Express measurements. Geophys. Res. Lett., 41(23),
389	8221–8227. doi: 10.1002/2014GL062287
390	Hanson, W. B., & Mantas, G. P. (1988, July). Viking electron temperature
391	measurements: Evidence for a magnetic field in the Martian ionosphere.
392	Journal of Geophysical Research: Space Physics, 93(A7), 7538–7544. doi:
393	10.1029/JA093iA07p07538
394	Holmberg, M. K. G., André, N., Garnier, P., Modolo, R., Andersson, L., Halekas,
395	J., Mitchell, D. L. (2019). MAVEN and MEX Multi-instrument Study
396	of the Dayside of the Martian Induced Magnetospheric Structure Revealed by
397	Pressure Analyses. J. Geophys. Res. Space Phys., 124(11), 8564–8589. doi:
398	10.1029/2019JA026954
399	Kliore, A. J., & Luhmann, J. G. (1991). Solar cycle effects on the structure of the
400	electron density profiles in the dayside ionosphere of Venus. J. Geophys. Res.
401	Space Phys., 96(A12), 21281–21289. doi: 10.1029/91JA01829
402	Luhmann, J. G., Elphic, R. C., Russell, C. T., Mihalov, J. D., & Wolfe, J. H. (1980,
403	November). Observations of large scale steady magnetic fields in the dayside
404	Venus ionosphere. Geophysical Research Letters, 7(11), 917–920. doi: 10.1029/
405	GL007i011p00917
406	Matta, M., Galand, M., Moore, L., Mendillo, M., & Withers, P. (2014, January).
407	Numerical simulations of ion and electron temperatures in the ionosphere
408	of Mars: Multiple ions and diurnal variations. <i>Icarus</i> , 227, 78–88. doi:
409	10.1016/j.icarus.2013.09.006
410	Miller, K. L., Knudsen, W. C., Spenner, K., Whitten, R. C., & Novak, V. (1980,
411	December). Solar zenith angle dependence of ionospheric ion and electron
412	temperatures and density on Venus. Journal of Geophysical Research: Space
413	Physics, 85 (A13), 7759–7764. doi: 10.1029/JA085iA13p07759
414	Nagy, A. F., Winterhalter, D., Sauer, K., Cravens, T. E., Brecht, S., Mazelle, C.,
415	Trotignon, J. G. (2004, March). The plasma Environment of Mars. Space
416	Science Reviews, 111, 33–114. doi: 10.1023/B:SPAC.0000032718.47512.92

417	Opgenoorth, H., Dhillon, R., Rosenqvist, L., Lester, M., Edberg, N., Milan, S.,
418	Brain, D. (2010, August). Day-side ionospheric conductivities at Mars. Plane-
419	tary and Space Science, $58(10)$, 1139–1151. doi: 10.1016/j.pss.2010.04.004
420	Phillips, J., Luhmann, J., & Russell, C. (1985, January). Dependence of
421	Venus ionopause altitude and ionospheric magnetic field on solar wind
422	dynamic pressure. Advances in Space Research, 5(9), 173–176. doi:
423	10.1016/0273-1177(85)90286-8
424	Picardi, G., Biccari, D., Seu, R., Plaut, J., Johnson, W. T. K., Jordan, R. L.,
425	Zampolini, E. (2004, August). MARSIS: Mars Advanced Radar for Subsurface
426	and Ionosphere Sounding. In Mars Express: The Scientific Payload (Vol. 1240,
427	pp. 51–69).
428	Pilinski, M., Andersson, L., Fowler, C., Peterson, W. K., Thiemann, E., & Elrod,
429	M. K. (2019, November). Electron Temperature Response to Solar Forcing
430	in the Low-Latitude Martian Ionosphere. Journal of Geophysical Research:
431	Planets, 124(11), 3082-3094.doi: 10.1029/2019JE006090
432	Ramstad, R., Barabash, S., Futaana, Y., & Holmström, M. (2017). Solar wind-
433	and EUV-dependent models for the shapes of the Martian plasma boundaries
434	based on Mars Express measurements. J. Geophys. Res. Space Phys., 122(7),
435	7279–7290. doi: 10.1002/2017JA024098
436	Ramstad, R., Barabash, S., Futaana, Y., Nilsson, H., Wang, XD., & Holmström,
437	M. (2015). The Martian atmospheric ion escape rate dependence on solar wind
438	and solar EUV conditions: 1. Seven years of Mars Express observations. J .
439	Geophys. Res. Planets, $120(7)$, 1298–1309. doi: 10.1002/2015 JE004816
440	Ramstad, R., Brain, D. A., Dong, Y., Espley, J., Halekas, J., & Jakosky, B. (2020,
441	October). The global current systems of the Martian induced magnetosphere.
442	<i>Nat. Astron.</i> , 4(10), 979–985. doi: 10.1038/s41550-020-1099-y
443	Sánchez-Cano, B., Narvaez, C., Lester, M., Mendillo, M., Mayyasi, M., Holmstrom,
444	M., Durward, S. (2020). Mars' Ionopause: A Matter of Pressures. J. Geo-
445	phys. Res. Space Phys., $125(9)$, e2020JA028145. doi: 10.1029/2020JA028145
446	Vogt, M. F., Withers, P., Mahaffy, P. R., Benna, M., Elrod, M. K., Halekas, J. S.,
447	Jakosky, B. M. (2015). Ionopause-like density gradients in the Martian
448	ionosphere: A first look with MAVEN. Geophys. Res. Lett., 42(21), 8885–8893.
449	doi: 10.1002/2015GL065269

450	Withers, P. (2008, July). Theoretical models of ionospheric electrodynamics and
451	plasma transport. Journal of Geophysical Research: Space Physics, 113(A7).
452	doi: 10.1029/2007JA012918
453	Withers, P., Matta, M., Lester, M., Andrews, D., Edberg, N. J. T., Nilsson, H.,
454	Witasse, O. (2016, January). The morphology of the topside ionosphere of
455	Mars under different solar wind conditions: Results of a multi-instrument ob-
456	serving campaign by Mars Express in 2010. Planetary and Space Science, 120,
457	24–34. doi: 10.1016/j.pss.2015.10.013
458	Zhang, M. H. G., & Luhmann, J. G. (1992). Comparisons of peak ionosphere pres-
459	sures at Mars and Venus with incident solar wind dynamic Pressure. $J.$ Geo-
460	phys. Res. Planets, 97(E1), 1017–1025. doi: 10.1029/91JE02721
461	Zhang, M. H. G., Luhmann, J. G., Kliore, A. J., & Kim, J. (1990). A post-Pioneer
462	Venus reassessment of the Martian dayside ionosphere as observed by radio
463	occultation methods. J. Geophys. Res. Solid Earth, 95(B9), 14829–14839. doi:
464	10.1029/JB095iB09p14829