

# Crustal-Scale Thermal Models: Revisiting the Influence of Deep Boundary Conditions

Denise Degen<sup>1,1</sup>, Karen Veroy<sup>2,2</sup>, Magdalena Scheck-Wenderoth<sup>3,3</sup>, and Florian Wellmann<sup>1,1</sup>

<sup>1</sup>RWTH Aachen University

<sup>2</sup>Eindhoven University of Technology

<sup>3</sup>Helmholtz Centre Potsdam, GFZ German Research Science for Geosciences

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## Abstract

The societal importance of geothermal energy is significantly increasing because of its low carbon-dioxide footprint. However, geothermal exploration is also subject to high risks. For a better assessment of these risks, extensive parameter studies are required that improve the understanding of the subsurface. This yields computationally demanding analyses. Often this is compensated by constructing models with a small vertical extent. This paper demonstrates that this leads to entirely boundary-dominated and hence uninformative models. It demonstrates the indispensable requirement to construct models with a large vertical extent to obtain informative models with respect to the model parameters. For this quantitative investigation, global sensitivity studies are essential since they also consider parameter correlations. To compensate for the computationally demanding nature of the analyses, a physics-based machine learning approach is employed, namely the reduced basis method, instead of reducing the physical dimensionality of the model. The reduced basis method yields a significant cost reduction while preserving the physics and a high accuracy, thus providing a more efficient alternative to considering, for instance, a small vertical extent. The reduction of the mathematical instead of physical space leads to less restrictive models and, hence, maintains the model prediction capabilities. The combination of methods is used for a detailed investigation of the influence of model boundary settings in typical regional-scale geothermal simulations and highlights potential problems.

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D. Degen  
Computational Geoscience and Reservoir Engineering (CGRE)  
RWTH Aachen University  
Wüllnerstraße 2, 52062 Aachen, Germany  
E-mail: denise.degen@cgre.rwth-aachen.de

K. Veroy  
Centre for Analysis, Scientific Computing and Applications  
Department of Mathematics & Computer Science  
Eindhoven University of Technology (TU/e)  
Groene Loper 5, Eindhoven, The Netherlands  
Faculty of Civil Engineering  
RWTH Aachen University  
Schinkelstraße 2, 52062 Aachen, Germany

M. Scheck-Wenderoth  
GFZ German Research Centre for Geosciences  
Telegrafenberg, 14473 Potsdam, Germany  
Department of Geology, Geochemistry of Petroleum and Coal  
RWTH Aachen University  
Aachen, Germany

F. Wellmann  
Computational Geoscience and Reservoir Engineering (CGRE)  
RWTH Aachen University  
Wüllnerstraße 2, 52062 Aachen, Germany

is employed, namely the reduced basis method, instead of reducing the physical dimensionality of the model. The reduced basis method yields a significant cost reduction while preserving the physics and a high accuracy, thus providing a more efficient alternative to considering, for instance, a small vertical extent. The reduction of the mathematical instead of physical space leads to less restrictive models and, hence, maintains the model prediction capabilities. The combination of methods is used for a detailed investigation of the influence of model boundary settings in typical regional-scale geothermal simulations and highlights potential problems.

**Keywords** : boundary conditions · global sensitivity analysis · sensitivity-driven model calibration · reduced basis method

## 1 Introduction

Geothermal energy is an important part of the future energy mix on the path to a more sustainable use of resources. Many aspects influence the potential use of a geothermal resource, with one prime parameter being the temperature in the subsurface. In order to determine expected temperatures on a regional scale, geothermal simulations are often performed (Gelet et al., 2012; Kohl et al., 1995; O’Sullivan et al., 2001; Taron et al., 2009; Watanabe et al., 2010). A common procedure is to start with a geological model, representing the main geological sequences, grouped by similar thermal properties, and to use this information for the parameterization of a geothermal simulation (Cacace et al., 2010; Mottaghy et al., 2011; Sippel et al., 2015). However, the (effective) thermal parameters of subsurface geological units (e.g. thermal conductivity, heat production rate) are generally uncertain and the material parameters are therefore often calibrated on the basis of temperature observations.

Extensive parameter studies or full uncertainty quantification studies are non-trivial since basin-scale models tend to be computationally demanding. To overcome this issue, a common approach is to generate models that have a large horizontal extension but a very small vertical extend red(often only a dew kilometers) that can be up to 40 times smaller than the horizontal extend (Freyemark et al., 2019; Kastner et al., 2015; Noack et al., 2013; Pribnow and Clauser, 2000; Pujol et al., 2015). The boundary conditions for these models are either based on best estimates or retrieved from larger models (Noack et al., 2013). This work investigates in detail how these typical approaches to treat boundary conditions influence all subsequent analyses, leading partly to fully boundary-dominated models. In this paper, it is demonstrated that they only have very limited capabilities for the analysis and understanding of the physical processes. During the model calibration, a compensation for possible boundary errors through an adjustment of the thermal properties is possible. Consequently, this has no direct impact on the temperature distribution but a significant impact on the physical plausibility of our model. Hence, for scenarios that lay outside of the calibrated regime, any prediction capabilities are lost. This is a major restriction when considering the sparse nature of observations. The models with a small vertical extent are commonly used, although it is well known that diffusion problems a majorly impacted by the boundary conditions. Therefore, this paper illustrates the consequences of this model choice and demonstrates that crustal-scale models are crucial for basin-scale applications.

37 In order to investigate the influence of thermal boundaries, full global sensitiv-  
38 ity analyses (SA) are employed for several case studies. These types of global SA  
39 approaches are usually not performed due to the high associated computational  
40 cost. To address these computational challenges, the full finite element solution  
41 of the forward solves is replaced with the reduced basis solution. This approach  
42 aims to reduce the complexity of the mathematical instead of physical space, yield-  
43 ing fast, accurate, and physics-preserving surrogate models. With these surrogate  
44 models, global sensitivity analyses are performed on several model realizations of  
45 a regional-scale geothermal basin model in northern Germany (around Berlin and  
46 the state of Brandenburg) to demonstrate the influence of the lower boundary  
47 condition on the simulation.

48 Additionally, an automated model calibration is executed to provide an ob-  
49 jective and reproducible way to compensate for the errors of both the physical  
50 and geological model. Sensitivity analyses for basin-scale models have been per-  
51 formed before in Noack et al. (2012) and also been combined with automated  
52 model calibrations (Wellmann and Reid, 2014). Also, Fuchs and Balling (2016)  
53 consider model calibrations but in their case without sensitivity analyses. Fur-  
54 thermore, local sensitivity studies are presented in Ebigbo et al. (2016). However,  
55 none of these can address the computationally demanding nature of the problem.  
56 Therefore, they are limited in the number of parameters, sensitivity analyses, and  
57 model calibrations they can perform. By using a physics-based machine learning  
58 approach instead of the finite element method, the computation time of the for-  
59 ward solve is reduced by several orders of magnitude. It allows, in turn, to perform  
60 global sensitivity analysis and full flexibility in the model calibration.

61 Global sensitivity analyses have been performed for hydrological problems (Ba-  
62 roni and Tarantola, 2014; Cloke et al., 2008; Song et al., 2015; Tang et al., 2007;  
63 van Griensven et al., 2006; Zhan et al., 2013), for volcanic source modeling (Can-  
64 navó, 2012), and for geothermal heat exchangers (Fernández et al., 2017). In Degen  
65 et al. (2020a), the authors have investigated the influence of both local and global  
66 sensitivity studies for the Upper Rhine Graben. In this paper, the combination of  
67 the global sensitivity study and model calibration, as presented in Degen et al.  
68 (2020a), is used to investigate the influence of the placement of the boundaries on  
69 the model predictions.

70 The paper is structured as follows: The methodologies and the governing equa-  
71 tions are presented in Section 2 and in Section 3, the problem of the lower boundary  
72 condition is conceptually introduced using a simple 1D model. Section 4 presents  
73 the impact of the lower boundary conditions, by focusing on a real-case basin-scale  
74 application. Therefore, the results of both global sensitivity analyses and model  
75 calibrations are presented and discussed.

## 76 **2 Materials and Methods**

77 In the following, the geothermal conduction problem used for the forward simu-  
78 lations of the temperature is briefly described. Furthermore, the concept of sensi-  
79 tivity analyses is introduced.

## 80 2.1 Physical Model

For the simulation of the temperature field, a geothermal conduction problem with the radiogenic heat production  $S$  as the source term (Bayer et al., 1997) is considered:

$$\lambda \nabla^2 T + S = 0, \quad (1)$$

81 where  $\lambda$  is the thermal conductivity, and  $T$  the temperature. For efficiency reasons  
 82 and to investigate the relative importance, the nondimensional form of eq. 1 is  
 83 considered. Therefore, the nondimensional properties  $T^* = \frac{T - T_{\text{ref}}}{T_{\text{ref}}}$ ,  $\lambda^* = \frac{\lambda}{\lambda_{\text{ref}}}$ ,  
 84  $S^* = \frac{S}{S_{\text{ref}}}$ , and  $\nabla^* = \frac{\nabla}{l_{\text{ref}}}$ , where the asterisk denotes the nondimensional quantity,  
 85 are required. Inserting them into eq. 1, leads to eq. 2:

$$\frac{\lambda}{\lambda_{\text{ref}}} \frac{S}{S_{\text{ref}}} \frac{\nabla^2}{l_{\text{ref}}^2} \left( \frac{T - T_{\text{ref}}}{T_{\text{ref}}} \right) + \frac{S}{S_{\text{ref}}} \frac{1}{T_{\text{ref}}} \frac{1}{\lambda_{\text{ref}}} = 0. \quad (2)$$

86 Here,  $\lambda_{\text{ref}}$  is the reference thermal conductivity,  $T_{\text{ref}}$  the reference temperature,  
 87  $S_{\text{ref}}$  the reference radiogenic heat production, and  $l_{\text{ref}}$  the reference length. Note  
 88 that the equation operates on the nondimensional space. For the motivational  
 89 study, the radiogenic heat production is neglected to focus the analysis on the  
 90 heat diffusion and the boundary condition. Furthermore, for all models Dirichlet  
 91 boundary conditions are applied at the top and bottom of the model domain.

## 92 2.2 Sensitivity Analysis

93 Sensitivity analyses aim to determine which model parameters influence the model  
 94 response to what extent. So, these studies investigate, which thermal conductivities  
 95 and radiogenic heat productions have a significant impact on the temperature  
 96 distribution. One distinguishes two types of sensitivity analyses: local and global  
 97 ones. Local sensitivity analyses consider that all parameters are independent of  
 98 each other. In contrast, global sensitivity studies investigate also the parameter  
 99 correlations. A detailed comparison of both methods for hydro-geological problems  
 100 is presented in Wainwright et al. (2014) and for basin-scale geothermal application  
 101 in Degen et al. (2020a).

102 For the sensitivity analysis (SA), a quantity of interest needs to be defined.  
 103 Here, the L2-norm of the temperature misfit to the reference model is used as the  
 104 quantity of interest, for the motivational example. The quantity of interest for the  
 105 real-case model is the L2-norm of the temperature misfit between the simulated  
 106 and observed temperature values.

107 For the global sensitivity analysis, the Sobol method with the Saltelli sampler  
 108 is used, this is a variance-based sensitivity analysis operating in a probabilistic  
 109 framework. From the variances, the sensitivity indices are derived as the ratio  
 110 between the partial and total variance. In this work, the main interest is on the  
 111 first- and total-order indices. The first-order index is the ratio between the variance  
 112 of the  $p$  parameter and the total variance and defines the impact of the parameter  
 113 itself. In addition, the total-order index captures all parameter correlations. This  
 114 includes second-order but also any higher-order terms. Second-order terms describe  
 115 the correlation between two parameters only, whereas higher-order terms define  
 116 the correlation between multiple parameters. Further information regarding the

117 Sobol method can be found in Sobol (2001); Saltelli (2002); Saltelli et al. (2010).  
118 For the sensitivity analyses the python library SALib (Herman and Usher, 2017)  
119 is used.

## 120 2.3 Model Calibration

121 The main aim of this paper is to investigate the influence of the lower boundary  
122 condition on the physical interpretation through an evaluation of the temperature  
123 distribution. This is the reason why global sensitivity analyses are used. In practical  
124 applications, it is often desired to calibrate the model against existing temperature  
125 measurements to ensure the correctness of the model.

126 For this, model calibrations are required, which aim to compensate for existing  
127 model errors by an adjustment of the model parameters. For deep geothermal  
128 applications calibrations are challenging since one usually has only a few shallow  
129 data points (Degen et al., 2020a). As the real-case study will show, it is possible to  
130 adjust a given model to the observed temperatures. However, larger model errors  
131 yield unphysical model parameters, imposing the danger of losing the predictability  
132 for observation points that have not been included in the calibration. This aspect  
133 will be discussed in detail later on.

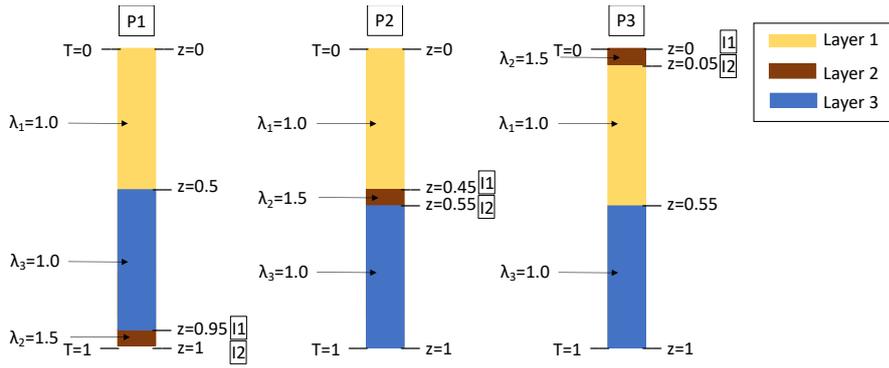
134 In this work, a trust region reflective algorithm is employed as the calibration  
135 method, which is a suitable choice for constrained problems, meaning that the  
136 thermal parameters have defined ranges (Branch et al., 1999). During the calibra-  
137 tion, the L1 norm of the misfit between the simulated and observed temperature  
138 measurements is minimized. The L1 norm is considered to put less weight on out-  
139 liers. The analysis is performed through the python library SciPy (Virtanen et al.,  
140 2020). For more details regarding the method, refer to Branch et al. (1999) and  
141 more details regarding the application to basin-scale models refer to Degen et al.  
142 (2020a).

## 143 3 Motivational Example

144 This paper investigates the influence of the impact of the lower boundary condition  
145 on the temperature distribution. This is an issue concerning geological models in  
146 general. For this reason, the problem is first demonstrated using a highly simplified  
147 motivational model. The motivational study aims to illustrate the general problems  
148 and not to represent a realistic geothermal application. To demonstrate that the  
149 issue has a major impact on real-case geothermal applications, the investigation  
150 is extended to the real-case study of Berlin-Brandenburg (a sedimentary basin in  
151 north-eastern Germany which is introduced in Section 4).

### 152 3.1 Forward Model

153 First, the forward problem used for the motivational study is introduced, for which  
154 a simplified 1D model is considered. The 3-layer model, schematically shown in  
155 Fig. 1, consists of three layers, where the middle layer is thinner than both adja-  
156 cent layers. A thermal conductivity of 1.0 is chosen for the top and bottom layer



**Fig. 1** Schematic representation of the 3-layer 1D model used for the motivational study of the boundary condition problem. Shown are the three different positions of the thin layer (P1-P3) for which the sensitivity analysis (Fig. 2) is conducted. Additionally, the interfaces of the thin layer are indicated with I1 and I2. The depth is denoted with  $z$ , the temperature with  $T$ , and the thermal conductivity with  $\lambda$ . Please note that a 1D model is considered, the 2D model representation of this figure was only chosen for an improved visibility.

157 and a thermal conductivity of 1.5 for the thin layer. To recall, throughout the  
 158 entire section, the dimensionless formulation is used. Consequently, the thermal  
 159 conductivity has no unit. At the top of the model, a Dirichlet boundary condition  
 160 of zero is applied for the temperature, and at the bottom a Dirichlet boundary  
 161 condition of one. The model is solved analytically. Note that the nondimensional  
 162 form is considered to focus the analysis on the relative difference.

163 In the following analyses, the influence of the thermal conductivity of the thin  
 164 middle layer (Layer 2 in Fig. 1) with respect to its distance from the boundary  
 165 conditions is analyzed. Therefore, the position of the thin layer changes. Three  
 166 different positions of the thin layer are considered: i) the thin layer adjacent to the  
 167 base boundary condition (position P1 in Fig. 2), ii) the thin layer in the center  
 168 of the model (position P2 in Fig. 2), and iii) the thin layer adjacent to the top  
 169 boundary condition (position P3 in Fig. 2). For the sensitivity analysis, scenario  
 170 P2 is defined as the reference model, where the thin layer is located around the  
 171 center (see Fig. 1). Consequently, the reference model represents the case of the  
 172 lowest possible boundary influence.

### 173 3.2 Impact of the Boundary Condition

174 To determine the influence of the lower boundary condition, a global sensitivity  
 175 analysis with 100 equally spaced temperature measurements in depth ranging from  
 176 zero to one is performed. Equally spaced measurements are chosen to avoid any  
 177 bias induced by the spatial distribution of the measurements in the sensitivity  
 178 analysis. Furthermore, the thermal conductivities of all three layers have an allowed  
 179 variation range of  $\pm 50\%$ .

180 The results of the global SA are shown in Fig. 2. Before discussing the results  
 181 for this SA, the terminology needs to be specified. From Fig. 2 first- and total-order  
 182 terms are obtained. The first-order terms describe the influence from the parameter  
 183 itself, whereas the total-order term describes the influence from the parameter

184 plus any parameter correlations. Consequently, the correlation is defined as the  
185 difference between the total- and first-order contributions. This motivational study  
186 investigates the influence of both boundary conditions on the model. Therefore, it  
187 needs to take the scenario, where the thin layer is in the center of the model (P2) as  
188 the reference case. This means that high influences of the parameters correspond  
189 to a high boundary dominance.

190 For the simple model, all thermal conductivities are dominated by total-order  
191 contributions for all three scenarios (P1-P3). This means that the parameters have  
192 high correlations. The high correlations are induced by the set-up of the model,  
193 where the temperature distribution is only determined by the two Dirichlet bound-  
194 ary conditions and by the ratio of the thermal conductivities between adjacent  
195 layers. Furthermore, the influence of  $\lambda_2$  is at all three positions the lowest, which  
196 is an effect of the lower thickness of this layer. Also note that for  $\lambda_2$ , nearly no  
197 first-order influences are observed.

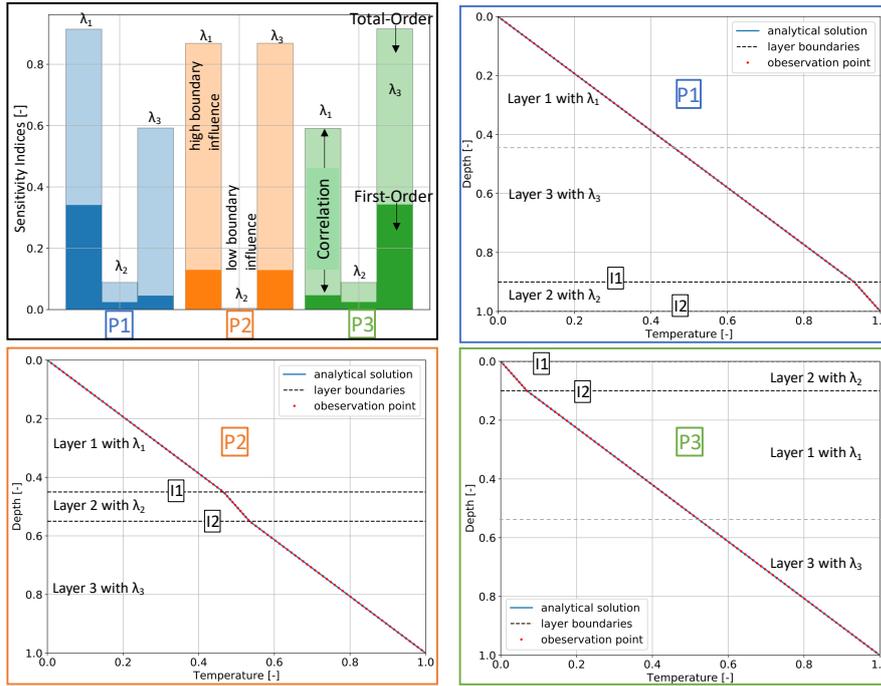
198 Focusing on scenario P1, the highest boundary dominance is achieved for  $\lambda_1$ ,  
199 which is situated at the upper boundary condition. The lowest influence is obtained  
200 for  $\lambda_2$  because of the above-described reason.  $\lambda_3$  has a significantly lower influence  
201 of the boundary than  $\lambda_1$ , which is logical since it is further away from the boundary.  
202 Interesting is that the decrease in the first-order contributions is more pronounced  
203 than the decrease in the total-order contributions. This shows that the remaining  
204 boundary influences are mainly arising from parameter correlations. By having  
205 a detailed look at the SA, one observes that the main correlations are arising  
206 from the correlation between  $\lambda_1$  and  $\lambda_3$ . For scenario P3, the same behavior with  
207 reversed roles for  $\lambda_1$  and  $\lambda_3$  is observed. In contrast for scenario P2, a boundary  
208 dominance of  $\lambda_1$  and  $\lambda_3$ , which are both adjacent to the boundaries, is obtained.  
209  $\lambda_2$  is situated in the center of the model, resulting in negligible contributions.

210 The results for all three scenarios are following the expectations since the  
211 smallest boundary influences are observed if the layers are further away from the  
212 boundaries. Note that these results can only be returned by a global SA. A local  
213 SA would assume that the influence is coming from the parameter itself. As an  
214 example, in P1 this would lead to a significant overestimation of the influence of  
215  $\lambda_3$ . In the worst case, this yields the misleading conclusion that  $\lambda_3$  is still greatly  
216 influenced by the boundary.

217 To conclude, for the motivational example the information about the thin layer  
218 is lost when it approaches the boundary condition. Or, as an alternative view-  
219 point, these two examples highlight the strong influence of boundary conditions  
220 on the simulation results. In a typical geothermal simulation setting, the position  
221 of the top boundary condition is usually defined as the land surface and cannot  
222 be changed. Its impact and possible ways to solve the issue have been discussed  
223 in Degen et al. (2020b). In contrast, the position of the lower boundary condition  
224 is usually adjustable.

## 225 4 Case Study Berlin-Brandenburg

226 After the demonstration of the general problem of the placement of the bound-  
227 ary for geological models, the consequences for real-case studies are illustrated.  
228 Therefore, the simplified 1D example is exchanged with various representations of



**Fig. 2** Black Box: Global sensitivity analysis to determine the impact of the boundary condition. Shown are the first- and total-order Sobol sensitivity indices of the thermal conductivities for the 3-layer model with respect to the distance from the boundaries. Blue Box: Scenarios P1, where the thin layer is adjacent to the bottom model boundary. Orange Box: Scenarios P2, where the thin layer is in the middle of the model boundary. Green Box: Scenarios P3, where the thin layer is adjacent to the top model boundary. Note that the interfaces of the thin layer are denoted with  $I$ . For a further illustration of the positions of the layers for P1 to P3, refer to Fig. 1

229 the Berlin-Brandenburg model, which cover a sedimentary basin in north-eastern  
 230 Germany (see Fig. 3).

#### 231 4.1 Berlin-Brandenburg Models

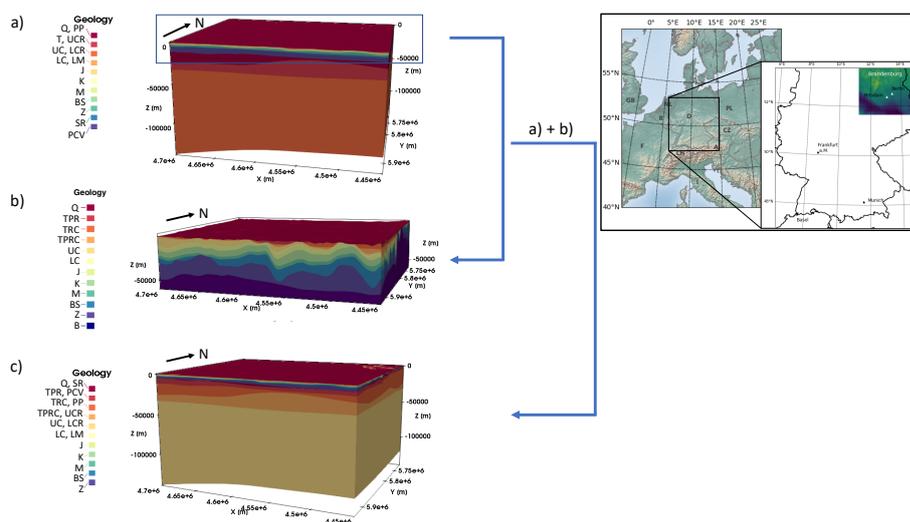
232 This paper uses three different versions of the Berlin-Brandenburg (BB) model.  
 233 The model is located in the southeastern part of the Northeast German Basin,  
 234 which is part of the Central European Basin System. The formation of the basin  
 235 started in the Late Carboniferous / Early Permian with a period of extensive  
 236 volcanism (Benek et al., 1996; Noack et al., 2012). Permian and Cenozoic sediments  
 237 are deposited above the volcanic rocks (Noack et al., 2012). The model is mainly  
 238 characterized by mobilized Upper Permian Zechstein salt, which forms salt pillows  
 239 and diapirs due to halokinetic movements (Noack et al., 2012; Scheck et al., 2003).  
 240 Also, the deeper crustal domains of the model are further differentiated to account  
 241 for the different consolidation ages (Noack et al., 2012). For further information  
 242 regarding the geological background, refer to Noack et al. (2012, 2013). The area is  
 243 of interest for geothermal studies due to a temperature anomaly consisting of high

244 heat flow values. This anomaly stretches from Poland to the river Elbe (Noack  
245 et al., 2012).

246 In the following, the numerical discretizations of the Berlin-Brandenburg mod-  
247 els are presented.

248 The first version of Berlin-Brandenburg, from now on denoted as the Berlin-  
249 Brandenburg LAB model (BB-LAB), has already been presented in Noack et al.  
250 (2012) and can be seen in Figure 3a. It has an extension of 250 km in the  $x$ - and  
251 of 210 km in the  $y$ -direction and extends down to the lithosphere–asthenosphere  
252 boundary (LAB). The model consists of 15 lithological units and the mesh consists  
253 of deformed eight-noded prisms. The grid resolution is one km in the horizontal  
254 directions, whereas the vertical length of the layers corresponds to the vertical  
255 element length, resulting in a mesh with 840,000 degrees of freedom.

256 The second model, in the following, referred to as the Berlin-Brandenburg 6  
257 km model, or BB-6km (Figure 3b), has the same horizontal extent but extends to a  
258 depth of 6 km instead of down to the LAB. It is presented in Noack et al. (2013) and  
259 consists of 12 lithological units. The model is discretized into a tetrahedral mesh.  
260 In comparison to the Brandenburg LAB model, it is refined in both geological and  
261 grid resolution terms. The horizontal element resolution is  $0.22 \text{ km}^2$  and vertical  
262 resolution is interpolated from the  $z$ -evaluations of the geological layers with a  
263 minimum thickness of 0.1 m, resulting in a mesh of 1,546,675 degrees of freedom.



**Fig. 3** Geology of the a) Berlin-Brandenburg LAB model, b) Berlin-Brandenburg 6 km model, and the c) Berlin-Brandenburg combined model. Please refer to Tab. 1, for the acronyms of the geological layers.

264 Combining the Berlin-Brandenburg 6 km model, the Berlin-Brandenburg LAB  
265 model, and removing the minimal vertical thickness of 0.1 m results in the third  
266 version of the Brandenburg model, denoted as the Berlin-Brandenburg combined  
267 model, or BB-combined (Figure 3c). Consequently, this model consists of 17 geo-  
268 logical layers, where the upper 11 layers have the same resolution as in the BB-6km  
269 model. The lower six layers have the same vertical resolution as the BB-LAB model

270 and the same horizontal resolution as the Berlin-Brandenburg 6 km model. As the  
 271 Berlin-Brandenburg LAB model this model extends to the LAB (the LAB depth  
 272 varies between about 100 to 140 km). This results in a tetrahedral mesh with  
 273 2,141,550 degrees of freedom.

274 For both the BB-LAB and BB-combined model, a Dirichlet boundary condi-  
 275 tion of 8 °C, corresponding to the average annual temperature, is applied at the  
 276 top of the model. Moreover, the Dirichlet boundary condition at the base of the  
 277 LAB is set to 1300 °C corresponding to the melting temperature of the mantle  
 278 rocks (Turcotte and Schubert, 2002). Additionally, a variation of the temperature  
 279 at this boundary condition of  $\pm 10\%$  is allowed to account for errors in the ge-  
 280 ometrical description of the LAB. The Berlin-Brandenburg 6 km model has the  
 281 same upper boundary condition, but at the base, various Dirichlet boundary con-  
 282 ditions directly taken from the Berlin-Brandenburg LAB model are considered.  
 283 Furthermore, a lower boundary conditions derived by Kriging is taken into ac-  
 284 count. For this interpolation, 900 equally spaced temperature observation from  
 285 the BB-LAB model in a depth of 6 km are considered and the interpolated bound-  
 286 ary is derived with a spherical variogram. All thermal properties are summarized  
 287 in Table 3 in the Supplementary Material. The forward simulations are performed  
 288 using the DwarfElephant package (Degen et al., 2020c) with a linear and nonlin-  
 289 ear solver tolerance of  $10^{-10}$ . Due to the nondimensional nature of the problem,  
 290 no preconditioners are needed for the finite element evaluations.

291 The reference thermal conductivity  $\lambda_{\text{ref}}$  is equal to the maximum thermal con-  
 292 ductivity of the BB-LAB model of  $3.95 \text{ W m}^{-1} \text{ K}^{-1}$ . For the BB-LAB and the BB-  
 293 combined model, the maximum temperature of 1300 °C is the reference tempera-  
 294 ture  $T_{\text{ref}}$ , whereas for the BB-6km model a reference temperature of 8 °C is chosen.  
 295 Homogeneous Dirichlet boundary conditions are used to achieve a better perfor-  
 296 mance of the numerical methods (Degen et al., 2020c). The Berlin-Brandenburg 6  
 297 km model has a constant Dirichlet boundary condition at the top. At the base, the  
 298 model has a Dirichlet boundary condition with a different temperature value for  
 299 each element. The top boundary condition is normalized to zero by using the value  
 300 of the top boundary as the reference parameter. The bottom boundary condition  
 301 is set to zero via a lifting function. In case of the Berlin-Brandenburg LAB and  
 302 combined model, the models have constant Dirichlet boundary condition values  
 303 for both upper and lower boundary, and hence one can use both of them as the re-  
 304 ference parameter. The value of the lower boundary condition was chosen to better  
 305 reduce the magnitude of the temperatures, which yields a better performance. The  
 306 maximum radiogenic heat production of the BB-LAB model of  $2.5 \mu\text{W m}^3$  is the  
 307 reference radiogenic heat production  $S_{\text{ref}}$ . The reference length  $l_{\text{ref}}$  corresponds to  
 308 the maximum x-extent of all models (250,000 m).

309 For the validation of the models the temperature measurements presented in  
 310 Noack et al. (2012, 2013) and based on Förster (2001) were used. The observations  
 311 consist of 81 temperature measurements from 44 wells in the area of Brandenburg.  
 312 It has been measured at various depth and stratigraphic levels.

#### 313 4.1.1 Reduced Models

314 The reduced basis (RB) method is a model order reduction technique that aims to  
 315 significantly reduce the dimensionality of problems resulting from a discretization  
 316 (e.g. via finite elements) of parameterized partial differential equations (PDE).

317 The method is decomposed into an offline and online stage, where the offline  
 318 stage, being a one time cost, constructs a reduced basis, and therefore comprises  
 319 all expensive pre-computations.

320 The online stage uses this reduced basis to allow very fast forward evalua-  
 321 tions, typically in the range of a few milliseconds (Degen et al., 2020c). In contrast  
 322 to other surrogate models, the RB method has the advantage that the physics  
 323 is preserved. Other surrogate model techniques build their models upon observa-  
 324 tions (Miao et al., 2019), without explicitly considering the PDE. The RB method  
 325 maintains the input-output relationship, meaning that the structure of the origi-  
 326 nal finite element problem (and consequently the PDE) is preserved (Hesthaven  
 327 et al., 2016). Hence, the method allows an extraction of the entire state vector  
 328 (e.g. the temperatures at every node of the model). Furthermore, for geothermal  
 329 conduction problems, it provides an error bound, enabling an objective evaluation  
 330 of the approximation quality. For further information regarding the RB method  
 331 refer to Hesthaven et al. (2016); Prud’homme et al. (2002); Veroy et al. (2003) and  
 332 for further information in the context of geosciences refer to Degen et al. (2020c).

333 For using the RB method, the geothermal problem is decomposed into a  
 334 parameter-dependent and -independent part. In the following, the affine decom-  
 335 positions of the integral formulation of the PDE for the various scenarios of the  
 336 Brandenburg model are defined. Note that this paper uses the operator represen-  
 337 tation. Therefore, it presents the bilinear form instead of the stiffness matrix, and  
 338 the linear form instead of the load vector.

For all Berlin-Brandenburg models, the bilinear form  $a$  has the following de-  
 composition:

$$a(w, v; \lambda) = - \sum_{q=0}^n \lambda_q \int_{\Omega} \nabla w \nabla v \, d\Omega, \quad \forall v, w \in X, \forall \lambda \in \mathcal{D}, \quad (3)$$

339 where  $w \in X$  is the trial function,  $v \in X$  the test function, “ $q$ ” denotes the index  
 340 of the training parameter (for more information see Tab. 3 in the Supplementary  
 341 Material),  $X$  the function space ( $H_0^1(\Omega) \subset X \subset H_1(\Omega)$ ),  $\Omega$  the spatial domain  
 342 in  $\mathbb{R}^3$ ,  $\lambda \in \mathcal{D}$  the parameter, and  $\mathcal{D}$  the parameter domain in  $\mathbb{R}^n$ . The number of  
 343 thermal conductivities in the training sample is denoted with  $n$ . Consequently,  $n$  is  
 344 equal to thirteen, nine, and fourteen for the BB-LAB, BB-6km, and BB-combined  
 345 model, respectively.

For all Berlin-Brandenburg models, except the BB-6km model with a lower  
 boundary condition derived via Kriging, the linear form  $f$  is decomposed in the  
 following way:

$$f(v; \lambda, s) = - \sum_{q=0}^n \lambda_q s \int_{\Gamma} \nabla v g(x, y, z) \, d\Gamma + s \int_{\Gamma} \nabla v S \, d\Gamma, \quad \forall v \in X, \forall \lambda \in \mathcal{D},$$

$$\text{with } g(x, y, z) = T_{\text{top}} \frac{h(x, y, z) - z_{\text{bottom}}(x, y)}{d(x, y)}.$$

(4)

346 Here,  $\Gamma$  is the boundary in  $\mathbb{R}^3$ ,  $s$  the scaling parameter for the lower boundary  
 347 condition,  $g(x, y, z)$  the lifting function,  $T_{\text{top}}$  the temperature at the top of the  
 348 model,  $h(x, y, z)$  the location in the model,  $z_{\text{bottom}}(x, y)$  the depth of the bottom  
 349 surface, and  $d(x, y)$  the distance between the bottom and top surface.

For the BB-6km with a Kriging lower boundary condition, the linear form slightly changes to the following:

$$f(v; \lambda, s) = - \sum_{q=0}^8 \sum_{i=0}^3 \lambda_q s_i \int_{\Gamma} \nabla v g_i(x, y, z) d\Gamma + s_2 \int_{\Gamma} \nabla v S d\Gamma, \quad \forall v \in X, \forall \lambda \in \mathcal{D},$$

with  $g_1(x, y, z) = g_3(x, y, z) = 1 - \frac{h(x, y, z) - z_{\text{bottom}}(x, y)}{d(x, y)}$ ,

$$g_2(x, y, z) = \left( \frac{3 d(x, y)}{2a} - \frac{1}{2} \left( \frac{d(x, y)}{a} \right)^3 \right) \left( 1 - \frac{h(x, y, z) - z_{\text{bottom}}(x, y)}{d(x, y)} \right). \quad (5)$$

Here  $g_1$ ,  $g_2$ , and  $g_3$  are again the lifting functions, with  $s_1$  being the nugget,  $s_2$  the partial sill,  $s_3$  the scaling parameter for the mean temperature, and  $a$  the range.

#### 4.1.2 Parameterization and Set-Up of the Sensitivity Analysis

The sensitivity analyses are performed with 13 (BB-LAB model – Fig. 3a), 11 (BB-6km model – Fig. 3b), 14 parameters (BB-combined model – Fig. 3c) and with 10,000 realizations for each parameter to reduce the statistical error. Note that for the Berlin-Brandenburg 6 km model exemplarily the results using the Kriging lower boundary condition are shown. The results of the sensitivity analyses using the other boundary conditions are analog to the one shown in this manuscript. In this paper, only the thermal conductivities are varied and the radiogenic heat productions are kept constant, to reduce the number of parameters within the reduction and all further analyses. The radiogenic heat productions are fixed and not the thermal conductivities because their influence on the overall temperature distribution is smaller. In Tab. 1, a list of all rock properties is provided. A variation of  $\pm 50$  % from the initial thermal conductivities is allowed for all thermal conductivities. Also, for the nugget and the partial sill, a variation of  $\pm 50$  % is enabled. For the scaling parameter of the lower boundary of both the Berlin-Brandenburg LAB model and Berlin-Brandenburg combined model a variation  $\pm 10$  % and for the scaling parameter of the mean temperature at the lower boundary condition of the BB-6km model  $\pm 20$  % is used, in order to account for the uncertainties related to those boundary conditions.

Table 1: Initial thermal properties Noack et al. (2012, 2013) of all models for and after the automated model calibration. The radiogenic heat production is denoted with  $S$ , and the initial thermal conductivity with  $\lambda_{\text{init}}$ .

ID	Layer	$S$ [ $\mu W m^{-3}$ ]	$\lambda_{\text{init}}$ [ $W m^{-1} K^{-1}$ ]
Q	Quaternary	0.7	1.50
T	Tertiary	0.7	1.50
TPR	Tertiary-post-Rupelian clay	0.7	1.50
TRC	Tertiary Rupelian-clay	0.45	1.00
TPRC	Tertiary-pre-Rupelian-clay	0.3	1.90
UC	Upper Cretaceous	0.3	1.90
LC	Lower Cretaceous	1.4	2.00

J	Jurassic	1.4	2.00
K	Keuper	1.4	2.30
M	Muschelkalk	0.3	1.85
BS	Buntsandstein	1.0	2.0
Z	Zechstein	0.09	3.5
B	Basement	1.5	2.50
SR	Sedimentary Rotliegend	1.0	2.16
PCV	Permo-Carboniferous Volcanics	2.0	2.50
PP	Pre-permian	1.5	2.65
UCR	Upper crust	2.5	3.10
LCR	Lower crust	0.8	2.70
LM	Lithospheric Mantle	0.03	3.95

## 371 4.2 Results

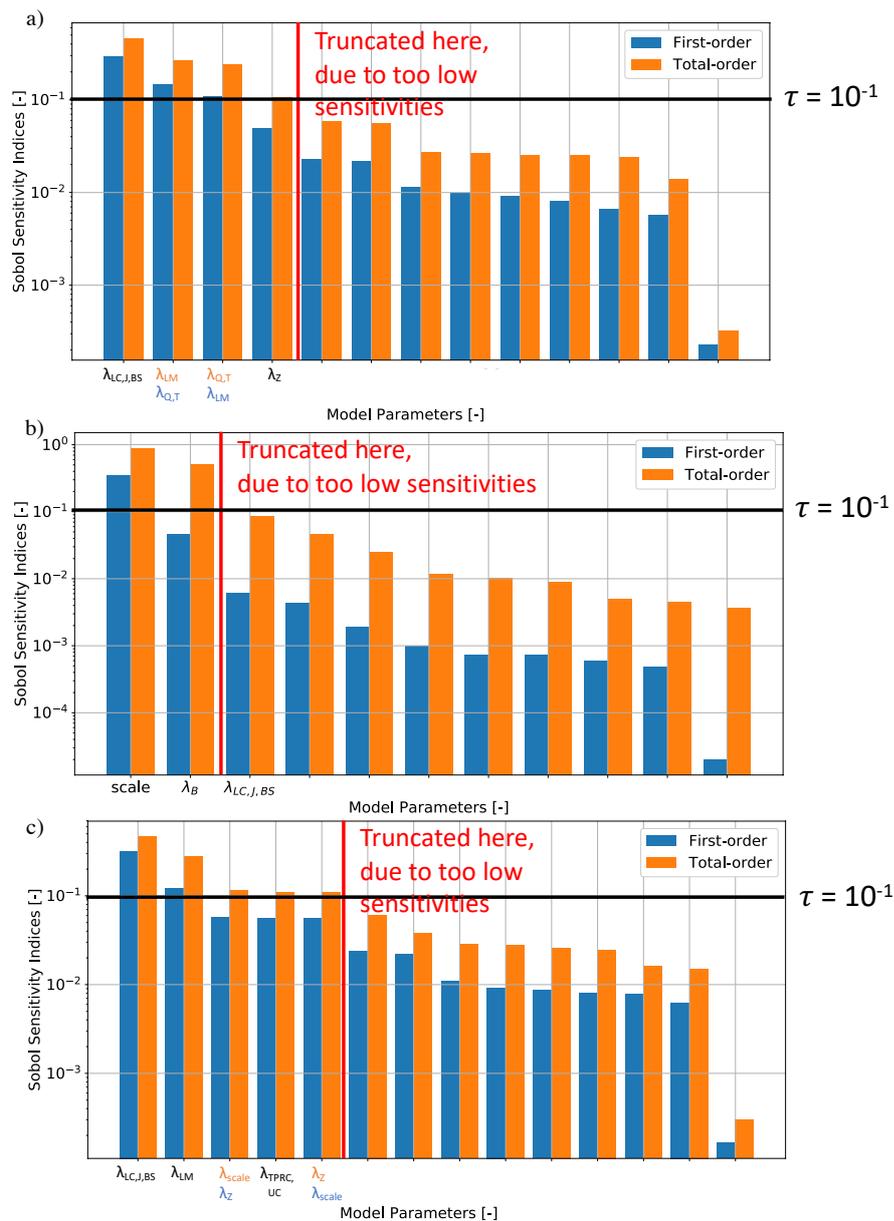
372 As for the conceptual study, this work demonstrates the influence of the lower  
 373 boundary condition. Therefore, first the results from the sensitivity analysis and  
 374 then the results from the model calibration are presented.

### 375 4.2.1 Sensitivity Analysis

376 Before presenting the results of the sensitivity analyses, note that all analyses  
 377 were performed with the aim to investigate the influence of the lower boundary  
 378 condition. The paper does not aim to characterize the influences of every single  
 379 thermal parameter in the model. Nevertheless, some geological impacts can be  
 380 derived and are presented in the following.

381 Regarding the sensitivities, the Berlin-Brandenburg LAB (Fig. 4a) is mostly  
 382 influenced by the Lower Cretaceous/Jurassic/Buntsandstein layer. The first-order  
 383 sensitivity index is dominant over the higher-order indices. Furthermore, the model  
 384 is sensitive to the Quaternary/Tertiary layer and the Lithospheric Mantle. For  
 385 the Quaternary/Tertiary layer, one again has predominantly first-order influences,  
 386 whereas the Lithospheric Mantle mostly impacts through higher-order contribu-  
 387 tions. Less pronounced is the influence from the Zechstein layer. The observed in-  
 388 fluence has similar first- and higher-order contributions. This is counter-intuitive  
 389 since one would expect a high influence of the Zechstein layer due to its high ther-  
 390 mal conductivity and highly variable thickness resulting in significant property  
 391 contrast. To explain this discrepancy, a closer look at the set-up of the sensitivity  
 392 analysis is required. In the analysis, layers with equal thermal conductivities were  
 393 combined. Therefore, the thermal conductivities of the Lower Cretaceous, Juras-  
 394 sic, and Buntsandstein layer are combined. Consequently, the high influence of this  
 395 layer is originating from this high combined sediment thickness. Keep in mind that  
 396 the aim of this analysis is to determine the influence of the boundary condition.  
 397 For determining which individual thermal conductivity has the highest influence  
 398 a separate analysis is required. The remaining thermal conductivities have minor  
 399 influences and are therefore disregarded in further analyses.

400 The Berlin-Brandenburg 6 km model is only influenced by the Basement layer  
 401 and by the variability of the lower boundary condition (Fig. 4b). The influence of



**Fig. 4** Global Sensitivity analysis for a) the Berlin-Brandenburg LAB, b) Berlin-Brandenburg 6 km model, and c) Berlin-Brandenburg combined model. Shown are the first- (blue) and total-order contributions (orange). The black line denotes the threshold value  $\tau$  for the truncation. Please refer to Tab. 1, for the acronyms of the thermal conductivities.

402 the scaling parameter of the mean temperature is significantly higher than the one  
 403 from the Basement layer. Higher-order contributions dominate both parameters.

404 Note that the Basement layer has nearly no first-order contributions, whereas the  
405 scaling parameter has non-dominant first-order contributions.

406 For the Berlin-Brandenburg combined model (Fig. 4c), one observes a similar  
407 pattern. The highest influences, dominated by first-order contributions, are arising  
408 from the Lower Cretaceous/Jurassic/Buntsandstein layer. The influence of both  
409 the Lithospheric Mantle and the scaling parameter of the lower boundary condi-  
410 tion increased, but higher-order contributions still dominate both parameters.  
411 The Tertiary-pre-Rupelian-clay/Upper Cretaceous, and the Zechstein layers are  
412 also influencing on the model and comparable first- and higher-order contribu-  
413 tions to each other.

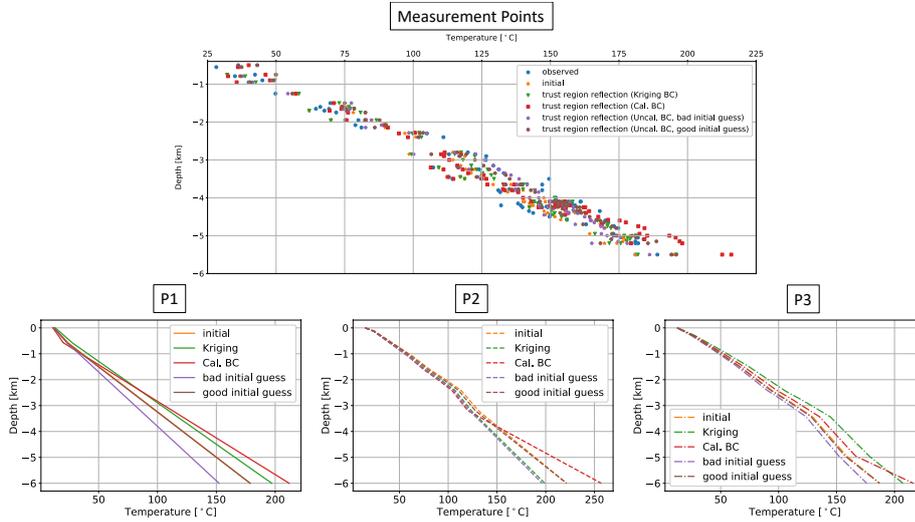
#### 414 *4.2.2 Model Calibration – Temperature Distribution*

415 The results from the global sensitivity analysis are taken as an input for the follow-  
416 ing model calibration. Therefore, only the influencing model parameters are consid-  
417 ered as shown in Fig. 4. Hence, four model parameters for the Berlin-Brandenburg  
418 LAB, two parameters for the Berlin-Brandenburg 6 km, and five parameters for  
419 the Berlin-Brandenburg combined model have to be taken into account for the  
420 model calibration. The remaining parameters are kept constant within the cali-  
421 bration since the sensitivity analysis identified them as having no impact on the  
422 temperature response. Model calibration is necessary to account for model errors  
423 of the Berlin-Brandenburg model.

424 The calibration of the Berlin-Brandenburg 6 km model is challenging because  
425 of the lower boundary condition. The conventional way to define this boundary  
426 condition is to extract it from the calibrated BB-LAB model and apply it to  
427 the BB-6km model, although it is generally not clear that the calibration for the  
428 larger model is also valid for the shallower model. To evaluate the influence of  
429 different calibration results, the model calibration for the shallow model using the  
430 boundary condition from two uncalibrated Brandenburg LAB model versions and  
431 various hierarchical model calibrations are compared. For the hierarchical models,  
432 either the boundary condition from the calibrated BC or a boundary condition  
433 obtained via Kriging as the lower boundary condition are chosen.

434 Therefore, Figure 5 compares the model calibrations using various lower bound-  
435 ary conditions of the Berlin-Brandenburg 6 km model. At the top panel, it shows  
436 the difference at the observation points. The differences between the various meth-  
437 ods are comparably small, which is not surprising since the calibration aims to  
438 minimize the difference between the simulated and observed temperatures at these  
439 locations. However, if one looks at the three points (P1 to P3, positions shown in  
440 Fig. 5), one observes differences between the various calibrations that can exceed  
441 50 °C. This means that for temperature prediction for points included inside the  
442 calibration data set good fits are obtained (regardless of the chosen boundary  
443 conditions). This changes once the points outside the calibration data set (P1 to  
444 P3) are considered, here significant differences for the different boundary condi-  
445 tions are obtained. This is of great importance for geoscientific applications since  
446 many studies face the problem of data sparsity. The model has many regions,  
447 where no data is available. Still, these regions might be of major importance. Con-  
448 sequently, it is desired to obtain models that are physically plausible to maintain  
449 the predictability of the models. To conclude, one can fit every model to a given  
450 temperature data set, with the consequences that the thermal conductivities get

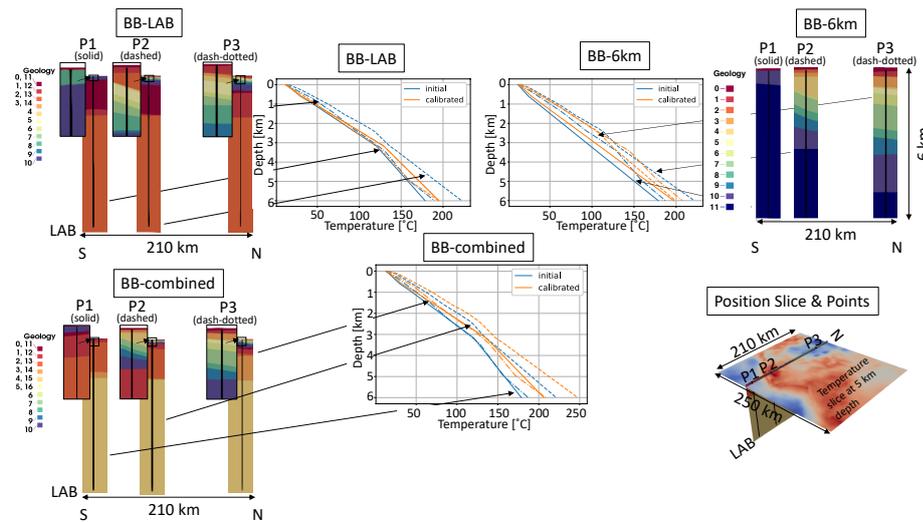
451 partly unphysical. This is less important if the target area coincides with a high  
 452 data density. However, this is often not the case. Therefore, the need to ensure  
 453 that the generality of the model is preserved remains.



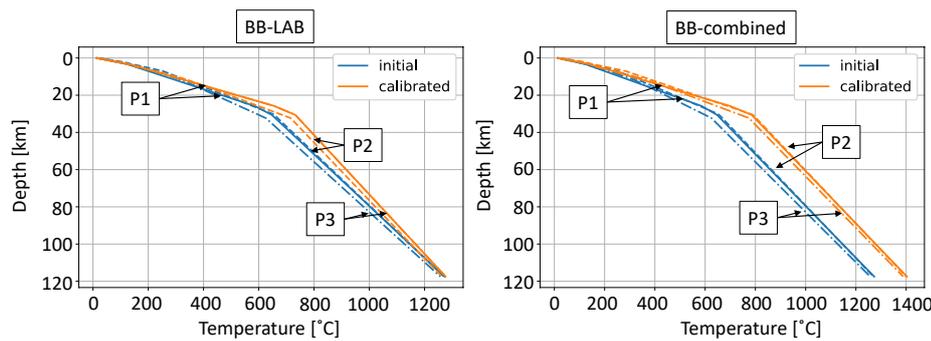
**Fig. 5** Comparison of the different calibration versions of the Berlin-Brandenburg-6 km model for the observed temperatures at all temperature measurements within the model (top panel) and at three points in the model (bottom panels) The position of the three points P1-P3 are shown in Fig. 6. They were chosen to cover the low temperature, the high temperature, and the by salt structures influenced temperature regions.

454 Fig. 6 compares the temperature distributions for the interval of the uppermost  
 455 6 kilometers of all three versions of the Berlin-Brandenburg model. For the BB-6km  
 456 model, exemplarily the hierarchical model calibration is shown. The differences  
 457 for all three points (P1 to P3) are comparable among the models. Note that the  
 458 possible variation range of the BB-6km is much larger since the determination  
 459 of the lower boundary condition is uncertain (see Fig. 5). The BB-LAB and BB-  
 460 combined model already show the maximum possible variation, whereas the BB-  
 461 6km model shows only the maximum variation range of the good-fit model.

462 Lastly, Fig. 7 shows the differences in the temperature distributions at the three  
 463 points (P1 to P3) for the entire depth of the BB-LAB and BB-combined models.  
 464 The major difference between both models is induced by the different treatments  
 465 of the boundary condition. During the sensitivity analysis of the BB-LAB model,  
 466 the scaling parameter of the lower boundary condition did not significantly in-  
 467 fluence the model response, contrary to the analysis of the BB-combined model.  
 468 Therefore, in the latter model the scaling parameter in the calibration is consid-  
 469 ered, whereas the value is kept constant for the former model. Although, with a  
 470 maximum temperature increase of 10 % a great amount of variation is allowed, the  
 471 possible variations at a depth of 6 km are comparable to those of the Brandenburg  
 472 6 km model.



**Fig. 6** Comparison of the temperature distribution over an interval of 6 km depth for all three versions of the Berlin-Brandenburg model at three different points in the models. The top left panels show the initial and calibrated temperature values or BB-LAB model and the stratigraphic columns for the points P1-P3. The top right panels show the same for the BB-6km model and the bottom panels for the BB-combined model. The bottom right panel shows the spatial position of the three points P1-P3.



**Fig. 7** The left panel shows the calibrated and initial temperature distributions at the points P1-P3 for the BB-LAB model over the entire model depth. The right panel displays the initial and calibrated temperature distributions at the points P1-P3 for the BB-combined model over the entire model depth. For the positions of P1-P3 refer to Fig. 6.

473 4.3 Discussion

474 In the following, the dangers of constructing models with a small vertical depth  
 475 are demonstrated. To further illustrate the importance of the placement of the  
 476 lower boundary condition, first its impact is demonstrated by using the results  
 477 of the global sensitivity study. Afterwards, the consequences for inverse processes  
 478 are emphasized, by using a deterministic model calibration. Both analyses are  
 479 presented for the case study of the Berlin-Brandenburg model.

#### 4.3.1 Sensitivity Analysis

The impact of the lower boundary condition is apparent by focusing on the difference between the BB-6km, and the BB-LAB and combined models. For the Berlin-Brandenburg 6 km model, the boundary condition is fixed at 6 km depth, resulting in an entirely boundary dominated model. This is observable due to the enormous sensitivity of the model to the:

- Basement layer,
- scaling parameter of the respective boundary condition, and
- correlation between both parameters.

Consequently, all information that is obtained from the Brandenburg 6 km model is coming from the boundary condition. Hence, the model is uninformative concerning the upper layers. However, these are the layers that are of interest since the target region is within these layers. Loosing the information about the thermal conductivities means that only the boundary is determining the solution. Hence, any errors of the boundary conditions have a possible huge impact on the temperature distribution at the target depth. This demonstrates that generating diffusive models with an extremely small vertical to horizontal length ratio is to be avoided at any cost.

The results of the global sensitivity analysis of the BB-LAB and combined model are matching the expectations. A high sensitivity is observed for the upper layers, which is caused by the shallow measurements (500 m to 6,820 m). First-order contributions of the Lower Cretaceous/Jurassic/Buntsandstein layers mostly impact the model. That means that the thermal conductivities of these layers are influencing the model themselves and not through a correlation with other layers. For the BB-LAB model, the thermal conductivity of the Quaternary and the Tertiary layer were combined into one training parameter. For the Brandenburg combined model, the thermal conductivities of the Quaternary and Tertiary-post-Rupelian, and the Tertiary-pre-Rupelian-clay and Upper Cretaceous were combined. Comparing the sensitivity analysis of both the BB-LAB model and combined model, one can conclude that the Tertiary-pre-Rupelian-clay is the layer that the model is sensitive to. The Quaternary, and the Tertiary-post-Rupelian layer can be ruled out because the Berlin-Brandenburg combined model is insensitive to it. Furthermore, also the Upper Cretaceous can be eliminated because the Berlin-Brandenburg LAB model is insensitive to it. Also, the influence of the thermal conductivity of the Tertiary-pre-Rupelian-clay is mainly originating from the parameter itself and not from interactions between various parameters. Again, the influence of the Tertiary-post-Rupelian-clay seems counter-intuitive due to its low thickness. This influence is a combination of the shallow measurements, which lead to higher influences for the upper layers and the Dirichlet boundary condition at the top. This boundary conditions fixes the temperature for each evaluation to the same value, yielding a reduced influence of the Quaternary and therefore a relatively higher influence of the Tertiary layers.

Additionally, for both models a significant influence of the Lithospheric Mantle is retrieved. Higher-order contributions dominate this parameter, and the second-order sensitivity indices show the parameter is correlated to the scaling parameter of the lower boundary condition. The Zechstein layer has similar influences in both model versions and is less significant in comparison to the overall influences.

527 To conclude, the only meaningful way to construct the model is by inserting  
528 the refined model into the original Berlin-Brandenburg LAB model. This results in  
529 the BB-combined model, which again shows the expected sensitivity distribution.  
530 One needs to keep in mind that this means an increase in degrees of freedom  
531 from 1,546,675 to 2,141,550. Nonetheless, both the finite element and the online  
532 execution time for both models are comparable since the complexity in these two  
533 models remains similar. This demonstrates that a reduction in the mathematical  
534 and not in the physical space is advantageous since it is much less restrictive.

#### 535 4.3.2 Model Calibration

536 At first hierarchical model calibrations seem to be a way to transfer the knowledge  
537 from large-scale coarse models to smaller-scale fine discretized models. However,  
538 the sensitivities clearly show that the smaller model becomes uninformative to-  
539 wards the upper layers. That is especially dangerous because it is not noticeable  
540 looking at the temperature distributions at the observation points only. Hence, at  
541 a first glance, one would get to the conclusion that cutting-of the model at 6 km is  
542 a valid approach. However, this would only be possible if our sole interests are the  
543 temperatures at the measurement points used within the calibration. Naturally,  
544 a calibration will match the simulation to the observed temperatures. However,  
545 that comes at a cost. For the various model calibrations of the BB-6km model  
546 one obtain thermal conductivities ranging between  $1.49 \text{ W m}^{-1} \text{ K}^{-1}$  and  $2.83 \text{ W}$   
547  $\text{m}^{-1} \text{ K}^{-1}$  for the Basement layer. Meaning that no longer physical thermal conduc-  
548 tivities but effective ones are retrieved. These effective thermal conductivities are  
549 tailored to our measurements. However, if a different location (e.g. new drill-hole  
550 location) is of interest, one can no longer derive reliable temperatures since the  
551 model calibration is not valid for this point and the model lost the information  
552 about the physical system.

553 This reveals the next important point. The above-described procedure is valid  
554 in a limited application field. However, one should be aware that the model is  
555 no longer representative of the physical processes. In contrast, both the BB-LAB  
556 and combined model have significant influences from various thermal conductivi-  
557 ties. The lower boundary condition is further away from the target area, reducing  
558 possible effects from this condition.

559 In general, one wants to improve through global SA the understanding of the  
560 physical model. In this specific case study, it a way to determine the most in-  
561 fluencing parameters allowing a back correlation to the geoscientific context was  
562 demonstrated. Note that both the SA and the calibration focus on the observa-  
563 tion locations. Hence, higher influences of shallower layers are observed. A study  
564 focusing solely on the temperatures at certain locations is applicable for some geo-  
565 physical studies but if the interest goes beyond fitting the temperatures it is not  
566 advisable to use models that are cut-off at a shallow depth.

567 Note that the changes for the thermal conductivities were not discussed in de-  
568 tail here. The reason is that the discussion of this paper focuses on the influence of  
569 the boundary condition. For further information about the thermal conductivities,  
570 refer to the Supplementary Material.

### 571 4.3.3 Outlook

572 Through this study, the path to subsequent tasks is opened. It would be interesting  
573 to further investigate the lower boundary condition. For some of the calibrations,  
574 very high thermal conductivities of the Lithospheric Mantle were obtained, which  
575 might be caused by the geometrical inaccuracies of the LAB. These inaccura-  
576 cies would impact the lower boundary condition and the calibration would try  
577 to compensate for this by adjusting the thermal conductivity of the Lithospheric  
578 Mantle. A scaling factor to the temperature value of this boundary to account  
579 for these inaccuracies was applied, which slightly improved the results. However,  
580 a single parameter is not enough to compensate for the model errors. Therefore,  
581 it would be interesting to replace the scaling factor by a function, which could  
582 be, for instance, determined through data assimilation. For this reason, a promis-  
583 ing next step to take would be to investigate if 3D-Var data assimilation yields  
584 improved results. In contrast to classical sequential data assimilation techniques,  
585 such as the Ensemble Kalman Filter (Burgers et al., 1998; Evensen, 1994), vari-  
586 ational data assimilation is a continuous approach, where the entire time frame  
587 is considered. Variational data assimilation methods minimize a cost function to  
588 obtain an estimate of the state variable. Three dimensional variational data assim-  
589 ilation has been studied intensively in numerical weather forecast by, for instance,  
590 Barker et al. (2004); Lorenc et al. (2000) but is fairly unknown for geothermal  
591 simulations. It has been studied in combination with the RB method already by  
592 Aretz-Nellesen et al. (2019). However, so far, the study is using benchmark prob-  
593 lems only. Therefore, it would be interesting to investigate the performance of the  
594 method for complex geophysical problems.

## 595 5 Conclusion

596 Throughout the entire paper, the high impact of the lower boundary conditions for  
597 conductive crustal-scale applications was demonstrated. Using a novel combination  
598 of reduced-order modeling techniques and global sensitivity analysis, the paper  
599 illustrated that cutting-off models at a shallow depth has severe consequences. For  
600 these models, the information content of the geological structures is entirely lost.  
601 This is of utmost importance if one aims to derive physical knowledge from the  
602 model and or want to perform predictions with the given model. These findings  
603 should be well known, still, it is a common procedure to construct models with a  
604 small vertical extent. Therefore, this work aims to explicitly show the consequences  
605 of this approach. The clear visualization of the boundary problem becomes only  
606 apparent through the utilization of a global sensitivity analysis since this method  
607 allows also the investigation of parameter correlations. Note that the value of a  
608 “too” small vertical extend differs for each model since it is dependent on various  
609 factors such as the type of boundary condition, the geological structure, and the  
610 governing physical principles. This further highlights the importance of sensitivity  
611 analyses in order to reliably determine whether a model is boundary-dominated.  
612 In order to construct informative models with a smaller vertical extend one could  
613 use, for instance, the Moho as the base boundary condition and apply a Neumann  
614 boundary condition, which is less restrictive than a Dirichlet boundary condition.

615 Another possibility is to use optimal experimental design techniques to determine  
616 a feasible depth of the model.

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## 622 **Declarations**

### 623 Funding

624 This work has been founded by the DFG through DFG Project GSC111.

### 625 Conflict of interest

626 The authors declare that they have no conflict of interest.

### 627 Availability of data

628 The temperature data used throughout this paper is available in Noack et al.  
629 (2012, 2013) and based on Förster (2001).

### 630 Code availability

631 For the construction of the reduced models, the software package DwarfElephant  
632 (Degen et al., 2020c) has been used. The software, which is based on the finite ele-  
633 ment solver MOOSE (Permann et al., 2020), is freely available on GitHub (<https://github.com/cgre-aachen/DwarfElephant>). The sensitivity analyses are per-  
634 formed with the Python library SALib (Herman and Usher, 2017).  
635

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# Supporting Information for “Crustal-Scale Thermal Models: Revisiting the Influence of Deep Boundary Conditions”

Denise Degen<sup>1</sup>, Karen Veroy<sup>2,3</sup>, Magdalena Scheck-Wenderoth<sup>4,5</sup>, Florian

Wellmann<sup>1</sup>

<sup>1</sup>Computational Geoscience and Reservoir Engineering (CGRE), RWTH Aachen University, Wüllnerstraße 2, 52072 Aachen,

Germany

<sup>2</sup>Centre for Analysis, Scientific Computing and Applications, Department of Mathematics & Computer Science, Eindhoven

University of Technology (TU/e), Groene Loper 5, Eindhoven, The Netherlands

<sup>3</sup>Faculty of Civil Engineering, RWTH Aachen University, Schinkelstraße 2, Aachen, Germany

<sup>4</sup>Faculty of Georesources and Materials Engineering, RWTH Aachen University, Aachen, Germany

<sup>5</sup>GFZ German Research Centre for Geosciences, Telegrafenberg, 14473 Potsdam, Germany

## Contents of this file

1. Text S1
2. Figure S1

## Additional Supporting Information (Files uploaded separately)

1. Caption for large Tables S1: Thermal Properties of all models for and after the automated model calibration. The initial thermal properties are from Noack, Scheck-Wenderoth, and Cacace (2012); Noack, Scheck-Wenderoth, Cacace, and Schneider (2013). We denote all parameters that are not involved in the model calibration, due to too low

sensitivities or that are not applicable for the specific model version, with n/a. Additionally, the IDs of the training parameters  $\mu$  are provided. We denote the radiogenic heat production with  $S$ , the initial thermal conductivity with  $\lambda_{\text{init}}$ , and the calibrated thermal conductivity with  $\lambda_{\text{cal}}$ .

## Introduction

This supporting material provides additional information regarding the results of the model calibration for the thermal conductivities (Text S1 and Table S1). Figure S1 shows the convergence for the maximum relative error for all three versions of the Berlin-Brandenburg model.

## Text S1: Model Calibration – Thermal Conductivities

Table S1 presents the changes in the thermal conductivity between the initial and the best-calibrated values for all models. Regarding the Berlin-Brandenburg LAB model, we observe an increase from  $2.0 \text{ W m}^{-1} \text{ K}^{-1}$  to  $2.45 \text{ W m}^{-1} \text{ K}^{-1}$  for the Lower Cretaceous-/Jurassic/Buntsandstein layers. For the Berlin-Brandenburg combined model, we see a more pronounced increase, resulting in a value of  $2.53 \text{ W m}^{-1} \text{ K}^{-1}$ . Furthermore, a substantial increase in thermal conductivity is observed for the Lithospheric Mantle, resulting in conductivities of  $5.93 \text{ W m}^{-1} \text{ K}^{-1}$  for the BB-LAB and combined model. Additionally, the scaling parameter for the lower boundary condition shows an increase to 1.10 for BB-combined model. Also, we obtain decreased thermal conductivities for the Zechstein layer of  $3.45 \text{ W m}^{-1} \text{ K}^{-1}$  (BB-LAB model), and an increased value of  $3.73 \text{ W m}^{-1} \text{ K}^{-1}$  (BB-combined model). The calibration of the Berlin-Brandenburg LAB model leads to an increased thermal conductivity of  $1.99 \text{ W m}^{-1} \text{ K}^{-1}$  for the Quaternary/Tertiary layer, and

the calibration of the Berlin-Brandenburg combined model to an increased thermal conductivity of the Tertiary-pre-Rupelian-clay/Upper Cretaceous layer to  $1.93 \text{ W m}^{-1} \text{ K}^{-1}$ . The parameter distribution for the BB-6km model shows an increased thermal conductivity of  $2.78 \text{ W m}^{-1} \text{ K}^{-1}$  for the Basement layer. The scaling factor for the mean temperature is 21.54 after the calibration resulting in a mean temperature of  $180 \text{ }^\circ\text{C}$ .

For the discussion of the thermal conductivities of the calibration results, we talk about the results from the Berlin-Brandenburg LAB and combined model because of the uninformative nature of the Berlin-Brandenburg 6 km model. We observe similar trends for both the original and the refined model. Although, the parameter distribution of the BB-combined model after the calibration is closer to the initial parameter distribution than the one from the BB-LAB model. This demonstrates the need for model calibration. It is incredibly challenging and time-consuming to construct a model that accounts for all structural effects. Taking the lack of data into account, it becomes a somehow impossible task. Therefore, we follow a different approach in this work. We compensate, for the model errors, by replacing the physical by effective thermal conductivities. In that way, we obtain a representative model. Also, keep in mind that the major shortcoming of the BB-6km model could only be revealed using a global sensitivity analysis. However, this requires so many forward simulations that it is not realizable, even with state-of-the-art finite element solvers, on a basin-scale without using surrogate models.

Considering the geological setting, the significant increase in thermal conductivity for the Lower Cretaceous/Jurassic/Buntsandstein layers (BB-LAB model and BB-combined model) is most likely caused by unresolved salt structures. Further investigations are required to analyze whether all layers (Lower Cretaceous/Jurassic/Buntsandstein) or only

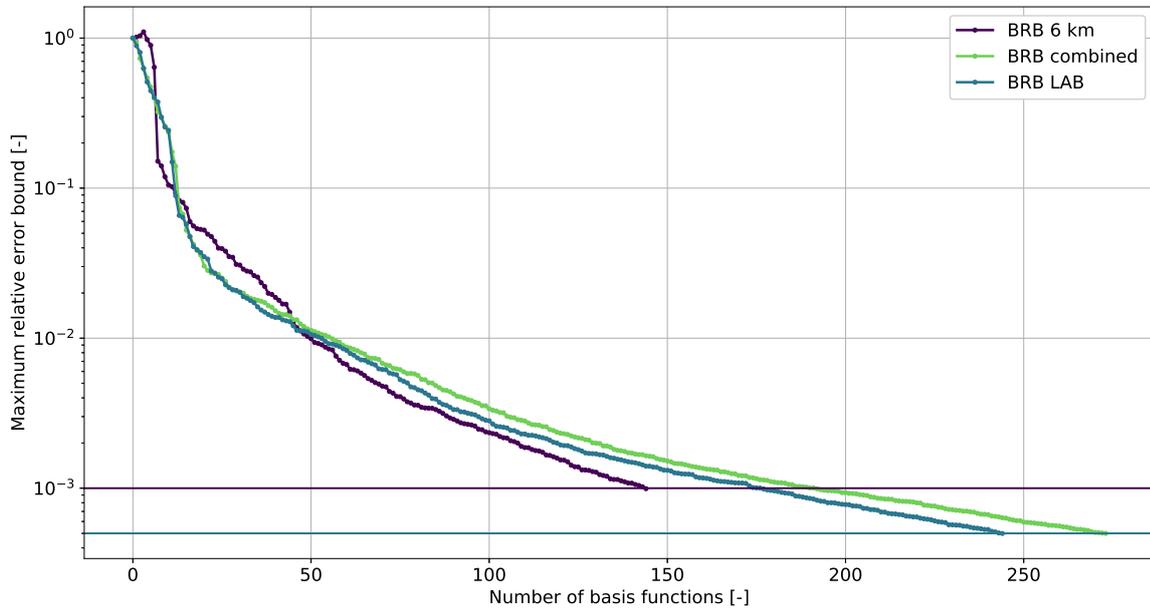
one layer contain unaccounted structures. The increase in thermal conductivity for the Zechstein layer is also most likely caused by unaccounted salt structures. Note that the increase in thermal conductivity for the Zechstein layer is significantly lower for the Berlin-Brandenburg combined model than for the Berlin-Brandenburg LAB model leading to the conclusion that the geological refinement captured successfully missing salt structures. The Keuper layer has after the calibration nearly the same value as the Muschelkalk layer leading to the assumption that we underestimated the sediment thickness of the Muschelkalk.

Combining the increase in thermal conductivity of the Lithospheric Mantle and the variation of the lower boundary conditions leads to the conclusion that a wrong geometrical parameterization of the LAB causes this increase. The correlation between both parameters confirms this. The fact that fixing the boundary condition to 1300 °C leads to an even higher increase in thermal conductivity further emphasizes this. For further studies, one could either allow a larger variation at the lower boundary condition or use a boundary condition that is either derived by data assimilation or considers tomography-derived temperatures.

Regarding the BB-6km model, we already stated that it is completely boundary dominated. Therefore, we focus the discussion on the Berlin-Brandenburg LAB and combined model. From the sensitivity analysis, we know that the main influence from the upper layers is arising from the Tertiary-pre-Rupelian-clay. Again, the increase is less dominant for the BB-combined model than for the BB-LAB model. This leads us to the conclusion that structural effects mainly cause this mismatch. The refined model version resolves this much better.

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**Figure S1.** Convergence of the maximum relative error bound for the Brandenburg 6 km model (denoted in purple), the Brandenburg combined model (denoted in green), and the Brandenburg LAB model (denoted in blue). We are using an error tolerance of  $1 \cdot 10^{-3}$  for the Brandenburg 6 km model, and an error tolerance of  $5 \cdot 10^{-4}$  for the Brandenburg combined and LAB model.