# The Migrating Speed of Alternate Bars

Michihide Ishihara<sup>1</sup> and Hiroyasu Yasuda<sup>1</sup>

<sup>1</sup>Niigata University

November 22, 2022

#### Abstract

Alternate bars can spontaneously occur and develop in rivers. They are considered to be a wave phenomenon due to their geometrical features and propagation characteristics. Presently, there is insufficient knowledge about their propagation, which is an important wave phenomenon property. In this study, a flume experiment was conducted under the condition that alternate bars occur and develop. This investigation aims to understand the existence and the scale of migrating speed of these alternate bars. The bed and water levels during the occurrence and development of the alternate bars were measured frequently with a high spatial resolution. By comparing the geometrical changes in the bed shape, the migrating speed of the alternate bars has a spatial distribution that changes with time. To quantify the spatial distribution of the migrating speed of the alternate bars, a hyperbolic partial differential equation for the bed level and migrating speed formula were derived. A comparison of the measured values for the flume experiment showed that the derived formula is applicable. Using the formula of the migrating speed in this hyperbolic partial differential equation, the migrating speed was verified to have a spatial distribution. In addition, the distribution changes with the development of the alternate bars over time. This study demonstrates that the dominant physical quantity of the migrating speed is the energy slope from the experimental results and the migrating speed formula.

# The Migrating Speed of Alternate Bars

# Michihide Ishihara<sup>1</sup>, Hiroyasu Yasuda<sup>2</sup>

<sup>1</sup>Graduate School of Science and Technology, Niigata University, Niigata, Japan <sup>2</sup>Research Institute for Natural Hazards & Disaster Recovery Niigata University, Niigata, Japan

# Key Points:

1

2

3 4

5

6	•	The spatial distribution of the migrating speed of alternate bars that occur in
7		rivers was determined.
8	•	A hyperbolic partial differential equation for the bed level and migrating speed
9		formula were derived.
10	•	The most dominant physical quantity of migrating speed of alternate bars is
11		the energy slope.

Corresponding author: Hiroyasu Yasuda, hiro@gs.niigata-u.ac.jp

#### 12 Abstract

Alternate bars can spontaneously occur and develop in rivers. They are considered 13 to be a wave phenomenon due to their geometrical features and propagation char-14 acteristics. Presently, there is insufficient knowledge about their propagation, which 15 is an important wave phenomenon property. In this study, a flume experiment was 16 conducted under the condition that alternate bars occur and develop. This investi-17 gation aims to understand the existence and the scale of migrating speed of these 18 alternate bars. The bed and water levels during the occurrence and development 19 of the alternate bars were measured frequently with a high spatial resolution. By 20 comparing the geometrical changes in the bed shape, the migrating speed of the al-21 ternate bars has a spatial distribution that changes with time. To quantify the spa-22 tial distribution of the migrating speed of the alternate bars, a hyperbolic partial 23 differential equation for the bed level and migrating speed formula were derived. A 24 comparison of the measured values for the flume experiment showed that the derived 25 formula is applicable. Using the formula of the migrating speed in this hyperbolic 26 partial differential equation, the migrating speed was verified to have a spatial distri-27 bution. In addition, the distribution changes with the development of the alternate 28 bars over time. This study demonstrates that the dominant physical quantity of the 29 migrating speed is the energy slope from the experimental results and the migrating 30 speed formula. 31

# 32 1 Introduction

Periodic forms can spontaneously form along a river channel's bed surface. 33 These forms are called riverbed waves because of their geometrical shapes and physi-34 cal properties. Riverbed waves can be classified as small-scale, mesoscale, and large-35 scale based on spatial scales, which include the wavelength and wave height (Seminara, 36 2010). Small-scale riverbed waves have wavelengths on the scale of the water depth; 37 meso-scale riverbed waves have wavelengths on the river width scale and wave heights 38 on the water depth scale. Large-scale riverbed waves have larger scales. The target 39 of this study is alternate bars that correspond to meso-scale riverbed waves. Alter-40 nate bars are riverbed waves that spontaneously form in rivers. They are located in 41 sites from the alluvial fan to the natural embankment. When observing alternate 42 bars from the sky with aerial photographs, the tip part is diagonally connected to 43 the left and right riverbanks; a deep-water pool is located on the downstream side of 44 this tip. In addition, it is known that the phase of the alternate bars propagates in 45 the same way as the water surface waves during flooding. 46

Over the years, many studies have been conducted on alternate bars. One of 47 the initial studies consisted of the flume experiments that were performed by Ki-48 noshita (Ryosaku, 1961). Kinoshita conducted long-term flume experiment to un-49 derstand the dynamics of the alternate bars that can produce meandering streams. 50 According to this experiment experiments, he reported that 1) alternate bars have 51 a globally uniform migrating speed and wavelength, 2) alternate bars in the early 52 stages of development have short wavelengths and fast migrating speeds, and 3) the 53 migrating speed becomes slower with the development of wavelengths. These results 54 have been confirmed in subsequent studies (Ikeda, 1983; Fujita & Muramoto, 1985; 55 Nobuhisa et al., 1999). In addition to the aforementioned conclusions, he proposed 56 a formula to calculate the migration speed of the alternate bars based on the experi-57 mental results, with the Froude number and shear velocity as the dominant physical 58 quantities. However, the validity of this formula has not been demonstrated in the 59 same study. 60

<sup>61</sup> Besides studies using flume experiments, several studies have applied math-<sup>62</sup> ematical analyses to understand the alternate bar phenomenon. Perhaps the first

mathematical study on alternate bars was that performed by Callander (Callander, 63 1969). He extended the stability analysis of Kennedy (Kennedy, 1963) for small-64 scale bed waves to a two-dimensional plane problem and theoretically discussed the 65 physical quantities that govern the generation of meso-scale riverbed waves. This 66 study was the starting point for the research that aimed at predicting the conditions 67 under which alternate bars occur and the wavelength and wave height of the alter-68 nate bars after development(Kuroki & Kishi, 1984; Colombini et al., 1987; Colombini 69 & Tubino, 1991; Tubino, 1991; Doelman et al., 1993). When considering the liter-70 ature that used these stability analyses, the studies by Callander (Callander, 1969) 71 and Kuroki (Kuroki & Kishi, 1984) are important. They derived a formula to cal-72 culate the migrating speed corresponding to the wave number that maximizes the 73 time amplification factor. The dominant physical quantities in the formula were 74 the Froude number, Shields number (shear velocity), bed slope, and wave number. 75 Moreover, in both the aforementioned studies, the value of the formula is compared 76 with the measured value. It has been reported that the reproducibility of the for-77 mula to calculate the migrating speed is good. 78

With the rise of stability analysis, numerical analyses of the riverbed fluctuations during the occurrence and development of alternate bars began to be carried out. Shimizu et al. (Shimizu & Itakura, 1989) reported for the first time that numerical analysis can satisfactorily reproduce each process of the occurrence and the development of alternate bars. In recent years, Federici et al. (Federici & Seminara, 2003) reported the propagation direction of the riverbed waves by performing stability and numerical analyses.

Recent studies that have used flume experiments (Lanzoni, 2000a, 2000b; Miwa 86 et al., 2007; Crosato et al., 2011, 2012; Venditti et al., 2012; Podolak & Wilcock, 87 2013) have investigated the effects of external factors such as the amount of sedi-88 ment supply on the dynamics of the alternate bars. Crosato et al. (Crosato et al., 89 2011, 2012) reported that alternate bars eventually shift from being migrating bars 90 to steady bars; they performed flume experiments and a numerical analysis to ver-91 ify this. Next, Venditti et al. (Venditti et al., 2012) reported that when the sedi-92 ment supply was interrupted after alternate bars occurred, the bed slope and shear 93 stress decreased, and the bars disappeared accordingly. Podolak et al. (Podolak & 94 Wilcock, 2013) also studied the response of alternate bars to sediment supply by in-95 creasing the sediment supply during the occurrence and development of alternate 96 bars. It was demonstrated that a non-migrating bar changed to a migrating bar with 97 an increase in the bed slope and shear stress owing to increase in the sediment supply. In addition, Eekhout et al. (Eekhout et al., 2013) investigated the dynamics of 99 alternate bars in rivers for nearly three years and reported that the migrating speed 100 decreased as the wavelength and wave height of the alternate bars increased and the 101 bed slope decreased. 102

Thus far, the geometrical shape and physical properties of the alternate bars 103 have been investigated. Based on the previous studies, it is possible to predict the 104 presence or absence of the alternate bars and their geometric shapes to some extent. 105 However, there is very little understanding of the nature of the migration speed of 106 the alternate bars. Therefore, in this study, we focused on the migration speed while 107 focusing on the physics of alternate bars. As this is not well understood, we carried 108 out the following to clarify the existence and scale of the spatial distribution. Sec-109 tion 2 describes the outline of the flume experiment that uses the Stream Tomog-110 raphy (ST) method, which can simultaneously measure the geometric shapes of the 111 water surface and the bed surface with a high spatial resolution; the results are also 112 described. In Section 3, we assumed that the alternate bars can be regarded as a 113 wave phenomenon, and we derived a hyperbolic partial differential equation (HPDE) 114 for the bed level. In this study, the advection velocity that is given to the advection 115

term of the HPDE was used to calculate the migration speed of the alternate bars.
In Section 4, the validity of the calculation formula that was derived in Section 3
was verified based on the characteristics of the HPDE and the measured values of
the bed level that was obtained in Section 2. In Section 5, the spatial distribution of
the migration speed of the alternate bars is quantified using the formula to calculate
the migration speed. Section 6 describes the results that were obtained in Section 5,
and Section 7 summarizes the research results.

# 2 Quantification of the Propagation Phenomenon in Alternate Bars based on the Flume Experiment

# 2.1 Experimental Setup

125

138

Figure 2 shows the plan view of the experiment flume. The experimental chan-126 nel consisted of a flume channel with a straight rectangular cross section. The flume 127 had a length of 12.0 m, a width of 0.45 m, and a depth of 0.15 m. Fixed weirs with 128 the same width as the flume were located 2.0 m from the upstream and downstream 129 ends of the flume. Over the section between 2.0 m and 10.0 m from the upstream 130 end that was sandwiched by these weirs, the initial bed of the channel for the exper-131 iment was a set flat bed. The bed was composed of a non-cohesive material with a 132 mean diameter of 0.76 mm and had a thickness of 5.0 cm. 133

For the water supply to the channel, circulation-type pumping from a water tank at the downstream end to a water tank at the upstream end was adopted; the water was steadily supplied. The accuracy of the water discharge was confirmed using an electromagnetic flowmeter.

#### 2.2 Experimental Condition

The purpose of this study is to understand the existence and scale of the spatial distribution of the migration speed of the alternate bars. An alternate bar is a
typical bed wave in a river that is attributed from an alluvial fan to a natural embankment.

Therefore, in the following experiments, we set the hydraulic conditions in 143 which the alternate bars developed. It has been theoretically shown that the occur-144 rence of alternate bars can be estimated using the river width depth ratio. There-145 fore, in this study, we set  $BI_0^{0.2}/h_0$  to 13.5, which corresponds to the occurrence area 146 of the alternate bars in the area division map as shown by Kuroki and Kichi(Kuroki 147 & Kishi, 1984). B is channel width,  $I_0$  is the initial uniform bed slope,  $h_o$  is the ini-148 tial uniform water depth. The bed slope of the flume is 1/160, the water discharge is 149 1.5 L/s, the flow velocity and water depth are 0.28 m/s and 0.012 m, respectively, on 150 the initial flat bed. The Shields number during the initial condition is 0.06, which is 151 higher than the critical Shields number (0.034) that was obtained from Iwagaki's for-152 mula (Yuichi, 1956). The sediment supply along the upstream end was not provided. 153 This is because by comparing the effect of the sediment supply with and without a 154 preliminary experiment, it was observed that the spatial distribution of the migrat-155 ing speed of alternate bars and its temporal changes occurred without the sediment 156 supply. 157

The water flow was carried out for 4 h during this experiment with the aforementioned conditions. At this time, alternate bars developed, and its propagation and shape change became slow.

#### 2.3 Measurement Method for the Bed Surface and Water Surface

In this study, we used ST, which was developed by Hoshino et al., to measure the bed level and water level in a plane while the water is flowing. For details on the principles of the ST measurement, refer to Appendix A. In this study, the aforementioned measurements were performed with a spatial resolution of 2 cm<sup>2</sup> for every minute. The water depth was calculated from the difference between the water level and bed level. The water surface slope was calculated from the central spatial difference between the water levels.

#### 2.4 Measurement Results

161

169

This section explains the propagation phenomenon of alternate bars using Fig. 3 and Fig. 4 based on the results of the high spatial resolution that was measured by ST.

Figure 3 shows the plan view of the deviation of the bed level by ST. The ori-173 gin of the vertical coordinate of the ST is the flume bottom. Therefore, the water 174 level and bed level represent the height from the bed of the flume. In this study, 175 we measured the bed level at 1-minute intervals using ST. However, the results at 176 20-minute intervals indicate that a clear change can be easily confirmed under the 177 set of hydraulic conditions that are shown. In this study, the initial bed surface was 178 created so that it was as flat as possible. However, it was difficult to obtain a per-179 fectly flat bed due to the accuracy limit of the bed surface that shapes the setup. It 180 has been confirmed that the alternate bars, which occurred and developed under the 181 aforementioned initial conditions, are almost the same as that developed in previous 182 studies (Ryosaku, 1958; Federici & Seminara, 2003; Crosato et al., 2011; Venditti et 183 al., 2012; Podolak & Wilcock, 2013). 184

First, the bed shape did not change from the flat bed as the initial condition 185 (Fig. 3(a),(b)). Second, it was possible to see the bed topography in which the 186 deposition and scouring are alternately repeated in the downstream direction, that 187 is 1.0 m, 2.5 m, and 4.0 m from the upstream end; thus, it was possible to confirm 188 that the alternate bars occurred (Fig. 3(c)). In this study, we defined 40 min, in 189 which the geometric features of the alternate bars were confirmed from the mea-190 sured result by the ST, as the occurrence time of the alternate bars. The alternate 191 bars developed the topography over time; they were deposited more on the riffle side 192 and scoured more on the pool. Subsequently, the entire alternate bars moved gradu-193 ally in the downstream direction. The development and propagation of the alternate 194 bars was significant from 40 min to 140 min (Fig. 3(c) to (h)). However, there was 195 minimal development or propagation after 140 min (Fig. 3 (h) to (m)). From this 196 result, comparing the migrating speed of the alternate bars during the early stage of 197 development with the migrating speed of the fully developed alternate bars, it can 198 be observed that the migrating speed in the former state is faster and the in latter is 199 slower. Figure 4 shows the longitudinal distribution of the deviation in the bed level 200 on the green dotted line in Fig. 3. Figure 4 shows (a) the initial stage of the exper-201 iment, (b) the occurrence of the alternate bars, (c) the intermediate stage of the ex-202 periment, and (d) the final stage of the experiment. Figure 4 shows three results, 203 where each one is 20 min apart. First, the deviation of the bed level was confirmed 204 to maintain a nearly flat bed from 1 min to 20 min (Fig. 4(a)). After (b) 40 min, 205 two bed undulations developed that were 1.5 m and 4.0 m from the upstream end. 206 The bed undulations developed their amplitudes and propagated in the downstream 207 direction. As a result, ST was confirmed to observe the wave nature of the alternate 208 bars. 209

The linear wave theory indicates that the phase propagates without deforming the waveform if a wave propagates with a spatial and temporal constant migrating

speed. Conversely, in the nonlinear wave theory, in which the migrating speed has a 212 spatial distribution and temporal changes, the wave propagates with deformation of 213 the waveform. From the viewpoint of the aforementioned wave theories, it is inferred 214 that the migrating speed of the alternate bars at the time of the occurrence of the 215 alternate bars in (b) has a spatial distribution; it changes with time and has nonlin-216 ear wave properties. Conversely, in (c) the intermediate stage of the experiment and 217 (d) the final stage of the experiment, there was no deformation of the waveform and 218 propagation. Thus, it was inferred that the migrating speed was almost zero; more-219 over, it is presumed that the wave nature was lost (Fig. 4(c),(d)). 220

#### 221 222

# 3 Derivation of the Calculation Formula for the Migrating Speed of the Alternate Bars

As shown in the previous section, the measurement results of this study show 223 the wave nature in the process of the occurrence and development of the alternate 224 bars. These findings are similar to what has been reported in the literature (Ryosaku, 225 1958; Federici & Seminara, 2003; Crosato et al., 2011; Venditti et al., 2012; Podolak 226 & Wilcock, 2013). In other words, there is scope for quantifying the spatial distribu-227 tion of the migration speed by an indirect method that uses a mathematical model 228 such as the HPDE(Fujita et al., 1985), which is suitable for describing the wave phe-229 nomena. The formula for calculating the migration speed is also derived from the 230 stability analysis (Callander, 1969; Kuroki & Kishi, 1984). However, because the for-231 mula calculates the migrating speed for each wave number, the spatial distribution 232 of the migrating speed cannot be quantified. Therefore, in this study, we can use 233 the HPDE for the bed level z and quantify the spatial distribution of the migration 234 speed of the alternate bars using the advection velocity of the advection term that 235 has the same formula. 236

This section describes the derivation process of the HPDE for the bed level 237 z. In addition, four different formulas can be obtained depending on the physical as-238 sumptions. This includes whether the dimension is one-dimensional or two-dimensional, 239 and whether the flow is stationary or unsteady. First, regarding the stationarity of 240 the flow, as it was confirmed that the non-stationary state in the phenomenon tar-241 geted by this study is very small from the verification results that are described in 242 Appendix B, we decided to deal only with the stationary state. In terms of the di-243 mensions, the geometric shape of the alternate bars and the flow there each has two-244 dimensional plane characteristics. Therefore, we decided to derive a two-dimensional 245 stationary equation. The derivation of the HPDE for the bed level can be used for 246 the continuous equation of the sediment, sediment functions, and the equation of the 247 water surface profile. For the derivation, the Exner equation was used as the contin-248 uous equation of the sediment, and the Meyer–Peter and Müller (M.P.M) formula 249 was used as the sediment function and the two-dimensional equation of the water 250 surface profile. The application of the HPDE to the various sediment functions was 251 examined using a method that is described in the next section. In this study, the 252 M.P.M formula, which is simple and has good applicability, was adopted. Vectors for 253 the longitudinal Eq. (2) and transverse Eq. (3) for the sediment flux are assumed 254 to match the flows. Equation (6) was used to calculate the Shields number. Above 255 all, we derived the steady two-dimensional equation of the water surface profile (Eq. 256 (4), Eq. (5)) to derive the HPDE for the bed level. For the details on the derivation 257 process of the steady two-dimensional equation for the water surface profile, please 258 refer to Appendix C. 259

$$\frac{\partial z}{\partial t} + \frac{1}{1 - \lambda} \left( \frac{\partial q_{Bx}}{\partial x} + \frac{\partial q_{By}}{\partial y} \right) = 0 \tag{1}$$

260 261

262

$$q_{Bx} = 8(\tau_* - \tau_{*c})^{3/2} \sqrt{sgd^3} \frac{u}{V}$$
(2)

$$q_{By} = 8(\tau_* - \tau_{*c})^{3/2} \sqrt{sgd^3} \frac{v}{v}$$
(3)

$$\frac{\partial h}{\partial x} = -\frac{\partial z}{\partial x} - I_{ex} - \frac{3}{10} \frac{u^2}{gI_{ex}} \frac{\partial I_{ex}}{\partial x}$$
(4)

$$+rac{1}{5}rac{uv}{gI_{ey}}rac{\partial I_{ey}}{\partial y}-rac{1}{2}rac{uv}{gI_{ex}}rac{\partial I_{ex}}{\partial y}$$

$$\frac{\partial h}{\partial y} = -\frac{\partial z}{\partial y} - I_{ey} - \frac{3}{10} \frac{v^2}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} -\frac{1}{2} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial x} + \frac{1}{5} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial x}$$
(5)

(6)

$$-\overline{2} \overline{gI_{ey}} \overline{d}$$

$$au_{*}^{269}=rac{hI_e}{sd}$$

where: z is the bed level, t is the time,  $\lambda$  is the porosity of the bed,  $q_{Bx}$  is the lon-271 gitudinal sediment flux, x is the distance of the longitudinal direction,  $q_{By}$  is the 272 transverse sediment flux, y is the distance of the transverse direction,  $\tau_*$  is the com-273 posite Shields number,  $\tau_{*c}$  is the critical Shields number, s is the specific gravity of 274 the sediments in water, g is the gravity acceleration, d is the sediment size, u is the 275 longitudinal flow velocity, V is the composite flow velocity, v is the transverse of the 276 flow velocity, and h is the depth. In addition,  $I_{bx} = -\partial z/\partial x$  is the longitudinal bed 277 slope,  $I_{ex}$  is the longitudinal energy slope,  $I_{by} = -\partial z/\partial y$  is the transverse bed slope, 278  $I_{ey}$  is the transverse energy slope, and n is the coefficient of roughness. 279

First, by applying the chain rule of differentiation to  $\partial q_{Bx}/\partial x$  in Eq. (1), we 280 obtain the following. 281

$$\frac{\partial q_{Bx}}{\partial x} = \frac{\partial q_{Bx}}{\partial \tau_*} \left( \frac{\partial \tau_*}{\partial h} \frac{\partial h}{\partial x} + \frac{\partial \tau_*}{\partial I_e} \frac{\partial I_e}{\partial x} \right) 
= \frac{\partial q_{Bx}}{\partial \tau_*} \left( \frac{I_e}{sd} \frac{\partial h}{\partial x} + \frac{h}{sd} \frac{\partial I_e}{\partial x} \right) 
= \frac{\partial q_{Bx}}{\partial \tau_*} \frac{I_e}{sd} \left( \frac{\partial h}{\partial x} + \frac{h}{I_e} \frac{\partial I_e}{\partial x} \right)$$
(7)

282

285

290

292

265

266

267

268

In addition,  $\partial I_e/\partial x$  in Eq. (7) becomes the following due to the application of the 283 chain rule to differentiate the Manning flow velocity Eq. (8). 284

$$V = \frac{1}{n} I_e^{1/2} h^{2/3} \tag{8}$$

$$\frac{\partial I_e}{\partial x} = \frac{\partial I_e}{\partial h}\frac{\partial h}{\partial x} + \frac{\partial I_e}{\partial V}\frac{\partial V}{\partial x} = -\frac{4}{3}\frac{I_e}{h}\frac{\partial h}{\partial x} + 2\frac{I_e}{V}\frac{\partial V}{\partial x}$$
(9)

Substituting Eq. (9) in Eq. (7) and rearranging this, we can obtain the following 288 equation. 289

$$\frac{\partial q_{Bx}}{\partial x} = \frac{\partial q_{Bx}}{\partial \tau_*} \frac{I_e}{sd} \left( -\frac{1}{3} \frac{\partial h}{\partial x} + 2 \frac{h}{V} \frac{\partial V}{\partial x} \right) \tag{10}$$

 $\partial q_{Bx}/\partial \tau_*$  in the aforementioned equation is as follows. 291

$$\frac{\partial q_{Bx}}{\partial \tau_*} = 12(\tau_* - \tau_{*c})^{1/2} \sqrt{sgd^3} \left(\frac{u}{V}\right) \tag{11}$$

Equation (4) is used for  $\partial h/\partial x$ . Substituting Eq. (4) and Eq. (11) in Eq. (10), Eq. 293 (10) is as follows. 294

$$\frac{\partial q_{Bx}}{\partial x} = 4(\tau_* - \tau_{*c})^{1/2} \sqrt{sgd^3} \left(\frac{u}{V}\right) \frac{I_e}{sd} \left(\frac{\partial z}{\partial x} + I_{ex}\right) \\ + \frac{3}{10} \frac{u^2}{gI_{ex}} \frac{\partial I_{ex}}{\partial x} - \frac{1}{5} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} + \frac{1}{2} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial y} + 6\frac{h}{V} \frac{\partial V}{\partial x}$$
(12)

In addition,  $\partial q_{By}/\partial y$  is arranged in the same process as Eq. (12), and the following equation is derived.

$$\frac{\partial q_{By}}{\partial y} = 4(\tau_* - \tau_{*c})^{1/2} \sqrt{sgd^3} \left(\frac{v}{V}\right) \frac{I_e}{sd} \left(\frac{\partial z}{\partial y} + I_{ey}\right) + \frac{3}{10} \frac{v^2}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} - \frac{1}{5} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial x} + \frac{1}{2} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial x} + 6\frac{h}{V} \frac{\partial V}{\partial y})$$
(13)

By substituting Eq. (12) and Eq. (13) in Eq. (1), the following HPDE for the bed level z is derived.

$$\frac{\partial z}{\partial t} + M_x \frac{\partial z}{\partial x} + M_x I_{ex} + M_x F_x + M_y \frac{\partial z}{\partial y} + M_y I_{ey} + M_y F_y = 0$$
(14)

In the aforementioned equation,  $M_x$  is the advection velocity of the longitudinal component of the bed level z. It is assumed to be closely related to the migration speed of the longitudinal component of the alternate bars, which is the subject of this study.  $M_y$  is the transverse migration speed of the alternate bars.  $M_x$  and  $M_y$ are not velocities of the sediments; they are supposed to be the propagation velocities of the bed level z.  $M_x$  and  $M_y$  are as follows.

$$M_x = \frac{4(\tau_* - \tau_{*c})^{1/2} \sqrt{sgd^3} \left(\frac{u}{V}\right) I_e}{sd(1-\lambda)}$$
(15)

309

308

$$M_y = \frac{4(\tau_* - \tau_{*c})^{1/2} \sqrt{sgd^3} \left(\frac{v}{V}\right) I_e}{sd(1-\lambda)}$$
(16)

From Eq. (15) and Eq. (16), it can be observed that the dominant physical quantities of the migrating speed are  $I_e$  and  $\tau_*$ . In addition,  $F_x$  and  $F_y$  are as follows.

$$F_x = \frac{3}{10} \frac{u^2}{gI_{ex}} \frac{\partial I_{ex}}{\partial x} - \frac{1}{5} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} + \frac{1}{2} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial y} + 6 \frac{h}{V} \frac{\partial V}{\partial x}$$
(17)

$$F_y = \frac{3}{10} \frac{v^2}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} - \frac{1}{5} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial x} + \frac{1}{2} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial x} + 6 \frac{h}{V} \frac{\partial V}{\partial y}$$
(18)

# <sup>316</sup> 4 Verifying the Applications of the HPDE for the Bed Level z and <sup>317</sup> the Migrating Speed Formula based on the Measured Values

This section verifies the applicability of the HPDE for the bed level z and the formula for the migrating speed, which was described in the previous section. Verification is achieved using the values of the ST and hydraulic analysis.

321

#### 4.1 Hydraulics required to verify applicability

This section describes the hydraulic quantities that are required to verify the 322 applicability of the HPDE and the calculation formula for the migrating speed, as 323 explained in the next section. As demonstrated from the HPDE and the calculation 324 formula of the migrating speed shown in the previous section, the hydraulic quan-325 tities that are required for the verification of the applicability are the water depth, 326 energy slope, and flow velocity. The water depth can be obtained from the bed level 327 and water level that is measured by the ST. However, the flow velocity and energy 328 slope that are paired with the water depth have not been measured; this measure-329 ment is generally difficult. Therefore, we determined the flow velocity and energy 330 slope by performing a numerical analysis. 331

For the numerical analysis, Nays2D included in iRIC (http://www.i-ric.org) can solve the two-dimensional plane hydraulic analysis. It was conducted with a bed level that was measured by the ST as a fixed bed.

The arrangement interval of the calculation points was 2 cm, which is the same 335 as the resolution of the ST in the transverse and longitudinal directions. The up-336 stream end boundary condition had a flow rate of 1.5 L /s; the downstream end 337 boundary condition provided the measured water depth. In addition, the roughness 338 coefficient was adjusted each time so that the calculated and measured values of the 339 water depth matched and they were uniform over the entire section. Figure 5 shows 340 the measured water depth. Figure 6 shows the difference between the measured wa-341 ter depth and the calculated value, which is dimensionless  $\Delta h_*$ . In addition, Fig. 342 7 illustrates the calculated flow velocity. Of these,  $\Delta h_*$  represents the calculation 343 accuracy of the numerical analysis. By focusing on  $\Delta h_*$  in Fig. 6,  $\Delta h_*$  is approxi-344 mately 10% for the entire channel at any time, regardless of the developmental state 345 of the alternate bars. In all the areas where  $\Delta h_*$  is 60% or more, the water depth is 346 very shallow. Most of these regions are not included when calculating the migrating 347 speed. This is demonstrated from the plan view of the migrating speed in Fig. 10 348 that is shown in the next section. Therefore, we decided to use the calculated value 349 of this part as well. In the next section, we verify the applicability of the derived 350 calculation formulas using these hydraulic quantities. 351

352 353

# 4.2 Verifying the Application of the Time Waveform for the Bed Level and the Riverbed Fluctuation Amount

We verified the applicability of the calculation formula that is derived in the previous section from two points of view. First, can the time waveform of the measured bed level be reproduced? Second, can the riverbed fluctuation amount that is measured in the entire section be reproduced? The verification results are described in this section.

359

# 4.2.1 Bed-level Time Waveform

The verification method that uses the time waveform at the bed level is described here. The HPDE from Eq. (14) that was derived in the previous section is numerically calculated as follows, and the amount of the riverbed fluctuation is between  $\Delta t$ .

$$\Delta z = \left(-M_x \frac{\partial z}{\partial x} - M_x I_{ex} - M_x F_x - M_y \frac{\partial z}{\partial y} - M_y I_{ey} - M_y F_y\right) \Delta t \tag{19}$$

This calculation used the bed level and water depth that is measured with ST, as well as the calculated values of the energy slope and flow velocity based on the hydraulic analysis that was described in the previous section. A time waveform at the bed level was obtained by repeating this numerical integration during each ST measurement time.

The applicability of the HPDE that was obtained in the previous section was investigated by comparing the time waveform of the bed level. This was calculated using the aforementioned method with the time waveform of the bed level that was observed by the ST. In this study, because the ST measurements were performed at 1-min intervals,  $\Delta t$  in the aforementioned calculation was set to 1 min.

Figure 8 shows the time waveform at the bed level. Figure 8 shows the time 375 waveforms of (a) the left bank side, (b) central part, and (c) right bank side at 6.0376 m from the upstream end; the red line shows the bed level of the measured value. In 377 the figure, the blue line shows the bed level that is calculated from the calculation 378 formula. By focusing on the measured values that are shown by the red line in Fig. 379 8, (a) the bed level decreases on the left bank side, and (b) the bed level increases 380 on the central part and (c) right bank side from the start of the water flow to 50 381 min. In addition, it can be observed that the bed level of (a) and (b) settled down 382

and the bed level of (c) increased from 50 to 150 min for the water flow. After 150 min, the bed level increased slightly for (a), (b), and (c).

Next, by looking at the time waveform of the bed level with the calculation 385 formula, the time waveform of the bed level is well reproduced from the start of the 386 water flow to 100–120 min in Figures (a), (b), and (c). It can be confirmed that af-387 ter 100 min, the reproducibility decreases on (a) the left bank side; after 150 min, 388 the reproducibility decreases on (b) the center and (c) the right bank side. By over-389 looking the process from the occurrence to the development of alternate bars using 390 Fig. 4 and Fig. 8, it was confirmed that the reproducibility of the time waveform, 391 when the propagation of the alternate bars is remarkable, is good. Conversely, it was 392 determined that the reproducibility becomes poor, especially in the sedimentary part 393 when the propagation of the alternate bars is slow. 394

As mentioned earlier, the time waveform was obtained by setting the time integration interval to 1 min. Although this time interval cannot be simply compared, it is much larger than the time interval in a general numerical analysis. From this result, it was determined that the verification method that uses the aforementioned numerical integration and the applicability of the calculation formula that was derived in the previous section are excellent.

#### 4.2.2 Riverbed Variation Amount

401

408

The verification in the previous section showed that the HPDE for Eq. (14) has sufficient applicability, but its applicability decreases as the alternate bars develop. In this section, we discuss how much of this reduced applicability occupies the entire waterway and where it occurs. This is achieved using the riverbed variation amount. The verification of the riverbed variation was performed using the following equation.

$$\Delta z_* = |\Delta z_{obs} - \Delta z_{cal}| / d \times 100 \tag{20}$$

where  $\Delta z_{obs}$  is the riverbed variation that was obtained from the bed level between 409 the two times that were measured by the ST. In addition,  $\Delta z_{cal}$  is the amount of the 410 riverbed variation by the HPDE and the calculation formula of the migration speed. 411  $\Delta z_*$  in the aforementioned equation is a dimensionless quantity that is obtained by 412 dividing the difference between the measured value of the riverbed variation amount 413 and the calculated value using the equation based on the particle size. In addition, 414 the difference between the two shows how much the divergence is based on the parti-415 cle size. 416

Figure 9 shows the plan view for the calculation accuracy of the riverbed vari-417 ation  $\Delta z_*$ . Figure 9 shows the bed level,  $\Delta z_*$  from the top. In addition, (a) to (i) 418 indicates the time zone in which the change is remarkable from the occurrence to 419 the development at 10-min intervals. In addition, (i) to (l) indicates the time zone 420 in which the change is slow until the end of the water flow at 40-min intervals. (a) 421 Looking at the results for 1 min of water flow,  $\Delta z_*$  is generally within 100%, and 422 the estimation accuracy of the waveform after 1 min at this time is the same as the 423 particle size. 424

From (a) 1 min of water flow to (h) 70 min, it can be observed that  $\Delta z_*$  is 425 generally within 100% of the entire channel. However, it can be observed that  $\Delta z_*$ 426 increases at the front edge of the bar after 80 min of water flow in comparison to 427 the other areas. From this, it is inferred that the applicability of this formula de-428 creases at the front edge of the alternate bars where the water depth becomes shal-429 low. Presently, the factors that reduce the applicability of the calculation formula 430 have not been sufficiently identified; however, it is suggested that careful handling is 431 required in places where the water depth becomes shallow. Conversely, it is unlikely 432

that alternate bars with a high wave height such that the water depth becomes ex-433 tremely shallow as considered in this experiment actually exists in an actual river; 434 therefore, it is considered that the decrease in applicability at the aforementioned 435 points is acceptable. It should be noted that there is a striped area where  $\Delta z_*$  is close to 500% on the side wall of the figure. This area is considered to occur owing 437 to a decrease in the ST measurement accuracy. In addition, (b) an increase in  $\Delta z_*$ 438 at the upstream end for 10 min of water flow is considered to occur owing to the 439 same aforementioned reason. Therefore, in the following discussion, the striped pat-440 tern of  $\Delta z_*$  on the right bank side and  $\Delta z_*$  at the upstream end were excluded from 441 consideration. 442

# 5 Quantification of the Migrating Speed for the Alternate Bars

The previous section confirmed that HPDE and the calculation formula for the migrating speed can reproduce the propagation phenomenon of the alternate bars. In this section, the migration speed of the alternate bars in each process during the occurrence and development is quantified using the calculation formula of the migrating speed.

449 450

# 5.1 Spatial Distribution of the Migrating Speed of the Alternate Bars

Figure 10 shows a plan view of the dimensionless migrating speed that is obtained by dividing the migrating speed that is obtained from the calculation formula for the bed level by the initial uniform flow velocity. The dimensionless migrating speed was used to understand the magnitude of the running water velocity and bed velocity. In a stability analysis(Callander, 1969; Kuroki & Kishi, 1984) that was conducted in the past, by applying the dimensionless governing equation, the migrating speed that was made dimensionless during the uniform velocity flow is often used.

The figure shows the bed level and  $M/u_0$  from the top. M is the magnitude 458 migrating speed,  $u_0$  is the uniform flow velocity. The area surrounded by the hatch 459 in the figure is the area in which the Shields number does not exceed the critical 460 Shields number (hereinafter referred to as the effective Shields number); in this area, 461 the migrating speed is 0. It is demonstrated that (a) to (i) indicates the time zone in 462 which the change is significant from the occurrence to the development of the alter-463 nate bars at 10-min intervals. In addition, (i) to (l) indicates the time zone in which 464 the change is slow until the end of the water flow at 40-min intervals. It should be 465 noted that the colors of the case law regarding the bed level and  $M/u_0$  are different 466 between (a) to (h) and (i) to (l). 467

First, by focusing on (a) 1 min of water flow in the figure,  $M/u_0$  has almost no 468 spatial distribution on a floor with an almost flat bed. It is also confirmed that the 469 bed surface propagates uniformly at a speed of approximately 0.002. After the bed 470 changes slightly from (b) 10 min to (c) 20 min,  $M/u_0$  begins to have a spatial dis-471 tribution. Subsequently, the spatial distribution of  $M/u_0$  changes significantly from 472 (d) 30 min of water flow to (e) 40 min when the alternate bars occurred. Looking at 473 this change with a spatial distribution from place to place, it can be observed that 474  $M/u_0$  increases at the sedimentary part and the front edge of the alternate bars, and 475 it decreases at other places. Subsequently, the spatial distribution of  $M/u_0$  became 476 clearer from (f) 50 min to (i) 80 min. In addition, before this time, it can be demon-477 strated that the area, where migrating speed is zero, is expanded on the downstream 478 side of the front edge. In addition, (i)  $M/u_0$  at 80 min of water flow decreased to 479 approximately 0.001 or less, while excluding the sedimentary part. (1) The spatial 480 distribution of the migrating speed at 240 min of the final time is not significantly 481 different from the distribution of (i) 80 min, but the area, where migrating speed 482

is zero, is further expanded. In addition,  $M/u_0$  also decreases slightly in the entire channel.

Next, Fig. 11 illustrates a histogram that quantitatively shows the spatial dis-485 tribution degree of  $M/u_0$  at each time. The red and blue vertical lines in the fig-486 ure represent the mean and mean  $\pm$  standard deviation of  $M/u_0$  at each time, and 487 each value is shown at the top of the figure. First, (a) the shape of the histogram 488 after 1 min of water flow was concentrated around the average value of 0.157. In ad-489 dition, because the standard deviation is 0.029, which is small with respect to the 490 mean value, it can be observed that the spatial distribution of  $M/u_0$  at this time is 491 small. Then, from (b) 10 min of water flow to (e) 40 min of water flow when the al-492 ternate bars occurred, the shape of the histogram became flat, and the mean value 493 of  $M/u_0$  was 0.147, and the standard deviation was 0.051. Comparing (a) 1 min and 494 (e) 40 min of water flow, although the mean value decreased by approximately 4 %, 495 the standard deviation increased to nearly 40 % of the mean value. From this, it is 496 demonstrated that the spatial distribution of the migrating speed greatly expanded 497 from the flat bed to the occurrence of the alternate bars. After that, from (e) 40 min to (i) 80 min of water flow, the flattening of the histogram, the increase in the 499 standard deviation, and the decrease in the mean value of  $M/u_0$  became more sig-500 nificant. From this, the increase in the standard deviation is particularly significant, 501 and the standard deviation value during 80 min of water flow reaches 70 % of the 502 average value of  $M/u_0$ . (i) After 80 min of water flow, there is no significant change 503 from (e) to (i). However, the average value of  $M/u_0$  gradually increases over 240 504 min of the final time, and it decreases and the standard deviation increases. Com-505 paring (a) 1 min of water flow and (k) 240 min, which was the final time, the mean 506 value of  $M/u_0$  is 0.72 times, and the standard deviation is 4.8 times. 507

From these results, it was demonstrated that the migration speed of the alternate bars has a spatial distribution, and this spatial distribution expands from the stage of occurrence to the development of the alternate bars.

511

# 5.2 Scale of the Migrating Speed of the Alternate Bars

This section discusses the scale of the migration speed of the alternate bars. As 512 shown in the previous section, from Fig. 11, it can be confirmed that the migrating 513 speed has a spatial distribution and it gradually expands from 1 min of water flow 514 to 240 min, which was the final time. The non-dimensional migrating speed in the 515 figure is divided by the uniform flow velocity (0.28 m/s) on the flat floor. However, 516 the scale of the migrating speed is on the order of  $10^{-4}$  to  $10^{-3}$  of the uniform flow 517 velocity at any place, regardless of the developmental state of the alternate bars. 518 Therefore, it is inferred that the deformation rate of the bed surface is sufficiently 519 smaller than the deformation rate of the running water. This tendency is similar 520 to the result of the non-dimensional migrating speed that corresponds to the wave 521 number of the maximum development rate as described by Callander (Callander, 522 1969).523

# 524 6 Discussion

In the previous sections, it was clarified that the migration speed of the alternate bars has a spatial distribution, and this spatial distribution changes with time. This section discusses the following three aspects of the migration speed of the alternate bars.

# 6.1 Approximate Description of the most Dominant Physical Quantity of the Migrating Speed

This study derived a formula to calculate the migration speed and its appli-531 cability is confirmed. Therefore, by considering the mathematical structure of the 532 same formula, it is possible to determine the most dominant physical quantity that 533 controls the migration speed. As demonstrated in Eq. (15) and Eq. (16), the dom-534 inant physical quantities of the migrating speed are the Shields number and energy 535 slope, which excludes the component decomposition of the migrating speed. Because 536 537 the dominant physical quantities of the Shields number consist of the water depth and energy slope, it can be concluded that the dominant physical quantities of the 538 migrating speed are the water depth and energy slope. In addition, the migrating 539 speed is calculated by multiplying the depth and the energy slope as shown in Eq. 540 (15) and Eq. (16). The depth in a river with alternate bars is generally on the order 541 of  $10^{\circ}$ , and the energy slope in the rivers with alternate bars is on the order of  $10^{-2}$ 542 to  $10^{-4}$ . Therefore, because the migration speed is the product of the water depth 543 and energy slope, it is inferred that the energy slope is dominant in terms of regulat-544 ing the magnitude of the migrating speed. 545

Figure 12 shows the relationship between the depth, energy slope, and migrating speed at each time. As demonstrated in Fig. 12, the relationship between the energy slope and migrating speed is nearly linear. Regarding the relationship between the depth and migrating speed, the migrating speed decreased with an increasing depth. From these results, it can be demonstrated that the migrating speed of the alternate bars is defined according to the energy slope.

It was previously demonstrated that the dominant physical quantity of the mi-552 grating speed is the energy slope. It is believed that an approximate description of 553 the migrating speed is possible using the energy slope and a water surface slope, 554 which is very similar to the energy slope. We verified whether the above approxi-555 mate description is possible from Fig. 13. The red points show the relationship be-556 tween  $M/u_0$  and  $0.4 \times I_e$ , and the green point shows  $M/u_0$  and  $0.4 \times I_w$  in Fig. 13. 557  $I_w$  is the value that is obtained by the central spatial difference of the water level 558 that is measured by the ST. Each correlation coefficient is shown on the upper side 559 of Fig. 13. Furthermore, 0.4, which is multiplied by  $I_e$  and  $I_w$ , is a coefficient that is 560 determined by the particle size. This is one of the variables on the denominator side 561 of Eq. (15) and Eq. (16). 562

First, it can be observed that the relationship between  $M/u_0$  and  $0.4 \times I_e$  is a nearly one-to-one relationship at all times. Furthermore, both relationships have a highly positive correlation because the correlation coefficients are over 0.99, on average. Second, it can be demonstrated that there is some variation in the relationship between  $M/u_0$  and  $0.4 \times I_w$ , but it is also a nearly one-one relationship at all times. Furthermore, both relationships have a highly positive correlation because the correlation coefficients are over 0.96 on average.

From these results, it was determined that the dominant physical quantity of the alternate bars is the energy slope, and the migration speed can be approximated using the energy slope and water surface slope.

573

529

530

### 6.2 Comparison of the Migrating Speed and the Stability Analysis

The migration speed based on the stability analysis(Callander, 1969; Kuroki & Kishi, 1984) has four dominant physical quantities: the Froude number, Shields number, bed slope, and wavenumber. In addition, the migration speed that was obtained from the HPDE, which was discussed in the previous section, is dominated by the Shields number and energy slope. The migration speed that was obtained from each stability analysis and the HPDE both have the water depth and energy slope as
the dominant physical quantities. However, the migration speed that was obtained
from the stability analysis was limited to those that correspond to any wavenumber.

582

# 6.3 Decreasing Factor for the Migrating Speed of the Alternate Bars

This subsection discusses the decreasing factor for the migrating speed of the 583 alternate bars. Figure 14 shows the average longitudinal distributions of the migrat-584 ing speed, energy line, hydraulic grade line, and bed line over time. The sediment 585 condition for the flume experiment in this study is that there is no sediment sup-586 ply. Therefore, this study confirmed that the bed level decreases with the passage of 587 time along the upstream side of the movable bed. At the final time, the bed level de-588 creased significantly from the upstream end to a point that is 3.5 m away. It can be 589 demonstrated that (b) the water level and energy head in this section have decreased 590 compared to the initial stage, and the riverbed slope and energy slope become more 591 gentle. In addition, (a) the migrating speed, which was calculated from Eq. (15) at 592 the final time in this section, is lower than that during the beginning of the water 593 flow. 594

A previous study by Eekhout et al. (Eekhout et al., 2013), which observed the occurrence and development process of the alternate bars in rivers, reported that the bed slope decreased while reducing the migration speed of the alternate bars. The results shown in Fig. 14 are consistent with those reported by Eekhout et al.

As mentioned earlier, the dominant physical quantity of the migrating speed of the alternate bars is the energy slope. The decrease in the migrating speed, which was confirmed by the flume experiments in this study and the observations in the actual river by Eekhout et al., should be interpreted as a decrease in the migration speed of the alternate bars due to the decrease in the riverbed slope rather than the decrease in the migrating speed with the development of the alternate bars.

# 605 7 Conclusion

In this study, we first conducted a flume experiment under the condition that alternate bars can occur and develop. We measured the hydraulic quantity and bed shape with a high spatial resolution. Next, we quantified the migrating speed of the alternate bars using the measured values that were obtained in the flume experiment and the calculation formula. This study determined that the migration speed of the alternate bars has a spatial distribution and it changes with time. The results of this study are presented below.

- 1) We were able to measure the water level and bed level with a high resolution 613 while continuing the water flow. In addition, the water level and bed level of 614 the occurrence and development process of the alternate bars are measured, 615 and it is demonstrated that the migrating speed of the alternate bars has a 616 spatial distribution from the measured geometric shape of the bed surface. 617 The HPDE for the bed level z and the formula of the migrating speed were 618 derived to quantitatively determine the migrating speed of the alternate bars. 619 By comparing the measured values of the flume experiment, it was demon-620 strated that the formula can appropriately describe the propagation phe-621 nomenon of the alternate bars. 622
- By calculating the migrating speed of the alternate bars based on the afore mentioned formula, it was clarified that the migrating speed of the alternate
   bars has a spatial distribution. In addition, the spatial distribution changes
   with the development of the bars over time, which was unconfirmed in the lit erature.

- 4) The dominant physical quantity of the migrating speed is the energy slope based on the results of the experiment and calculation formula. In addition, by comparing the scales of  $I_e$  and  $I_w$  with the non-dimensional migrating speed, it is suggested that the non-dimensional migrating speed can be described using  $I_e$  and  $I_w$ ; the behavior of the bars in the rivers can also be explained by  $I_e$  and  $I_w$ .
  - 5) It was observed that the migrating speed of the alternate bars is about three to four orders of magnitude smaller than the initial uniform flow velocity, regardless of the developmental state and the location of the bars.

It was found that the decrease in the applicability of the HPDE for the bed level and calculation formula of migration speed for the alternate bars, which was derived in this study, occurs in the sediment parts with extremely shallow water depth. Therefore, it is difficult to apply the aforementioned formula. However, it is unlikely that alternate bars, which has sediment parts with extremely shallow water depth, will occur in actual rivers, so this problem is not considered to be a practical problem.

Previous studies have revealed that the bank failure and the channel evolution
are closely related to alternate bars (Ryosaku, 1957; Callander, 1969). In the future,
we will quantify the spatial distribution of the migration speed of alternate bars in
actual rivers, and consider the above relationship.

#### 648 Acknowledgments

<sup>649</sup> The data used in this study can be accessed at https://doi.org/10.4121/14384999.

For details of the data, please refer to the enclosed README.md. We would like to thank Editage (www.editage.com) for English language editing.

#### 652 References

634

635

636

- Callander, R. A. (1969). Instability and river channels. Journal of Fluid Mechanics, 36(3), 465-480. doi: 10.1017/S0022112069001765
- <sup>655</sup> Colombini, M., Seminara, G., & Tubino, M. (1987). Finite-amplitude alternate bars.
   <sup>656</sup> Journal of Fluid Mechanics, 181, 213-232. doi: 10.1017/S0022112087002064
- <sup>657</sup> Colombini, M., & Tubino, M. (1991). Finite-amplitude free bars: A fully nonlinear
   <sup>658</sup> spectral solution. in Sand Transport in Rivers, Estuaries and the Sea, edited by
   <sup>659</sup> R. Soulsby and R. Bettes, A. A. Balkema, Brookfield, Vt., 163-169.
- Crosato, A., Desta, F. B., Cornelisse, J., Schuurman, F., & Uijttewaal, W. S. J.
   (2012). Experimental and numerical findings on the long-term evolution of migrating alternate bars in alluvial channels. *Water Resources Research*, 48(6).
   doi: 10.1029/2011WR011320
- Crosato, A., Mosselman, E., Beidmariam Desta, F., & Uijttewaal, W. S. J. (2011).
   Experimental and numerical evidence for intrinsic nonmigrating bars in alluvial channels. *Water Resources Research*, 47(3). doi: 10.1029/2010WR009714
- Doelman, A., de Swart, H., & Schielen, R. (1993, 07). On the non-linear dynamics of
   free bars in straight channel. Journal of Fluid Mechanics, 252, 325 356. doi:
   10.1017/S0022112093003787
- Eekhout, J. P. C., Hoitink, A. J. F., & Mosselman, E. (2013). Field experiment
   on alternate bar development in a straight sand-bed stream. Water Resources
   *Research*, 49(12), 8357-8369. doi: 10.1002/2013WR014259
- Federici, B., & Seminara, G. (2003). On the convective nature of bar instability.
   Journal of Fluid Mechanics, 487, 125-145. doi: 10.1017/S0022112003004737
- Fujita, Y., Koike, T., Furukawa, R., & Y., M. (1985). Experiments on the initial
  stage of alternate bar formation. *Disaster Prevention Research Institute Annu- (in Japanese)*, 28(B-2), 379-398.

Fujita, Y., & Muramoto, Y. (1985). Studies on the process of development of al-678 ternate bars. Bulletin of the Disaster Prevention Research Institute, 30(3), 679 55-86. 680 Ikeda, H. (1983). Experiments on bedload transport, bed forms, and sedimentary 681 structures using fine gravel in the 4-meter-wide flume.. 682 Kennedy, J. F. (1963). The mechanics of dunes and antidunes in erodible-683 bed channels. Journal of Fluid Mechanics, 16(4), 521-544. doi: 10.1017/ 684 S0022112063000975 685 Kuroki, M., & Kishi, T. (1984). Regime criteria on bars and braids in allu-686 vial straight channels. Proceedings of the Japan Society of Civil Engineers, 687 1984 (342), 87-96. doi: 10.2208/jscej1969.1984.342\_87 688 Lanzoni, S. (2000a). Experiments on bar formation in a straight flume: 1. uni-689 form sediment. Water Resources Research, 36(11), 3337-3349. doi: 10.1029/ 690 2000WR900160 Lanzoni, S. (2000b). Experiments on bar formation in a straight flume: 2. graded 692 sediment. Water Resources Research, 36(11), 3351-3363. doi: 10.1029/ 693 2000WR900161 694 Miwa, H., Daisdo, A., & Katayama, T. (2007). Effects of water and sediment dis-695 charge conditions on variation in alternate bar morphology. Proceedings of hy-696 draulic engineering (in japanese), 51, 1051-1056. doi: 10.2208/prohe.51.1051 697 Nobuhisa, N., Yoshio, M., Yoshihiko, U., Takashi, H., Masayuki, Y., Yasuhiko, T., & 698 Michiaki, I. (1999). On the behaviour of alternate bars under several kinds of 699 channel conditions. PROCEEDINGS OF HYDRAULIC ENGINEERING (in 700 *japanese*), 43, 743-748. doi: https://doi.org/10.2208/prohe.43.743 701 Podolak, C. J. P., & Wilcock, P. R. (2013). Experimental study of the response of 702 a gravel streambed to increased sediment supply. Earth Surface Processes and 703 Landforms, 38(14), 1748-1764. doi: 10.1002/esp.3468 704 Ryosaku, K. (1957). Formation of "dunes" on river-bed. Transactions of the 705 Japan Society of Civil Engineers in Japan(in Japanese), 1957(42), 1-21. doi: 706 10.2208/jscej1949.1957.1707 Ryosaku, K. (1958). Experiment on dune length in straight channel. Journal of 708 the Japan Society of Erosion Control Engineering (in japanese), 1958(30), 1-8. 709 doi: 10.11475/sabo1948.1958.30\_1 710 Ryosaku, K. (1961). Investigation of channel deformation in ishikari river. Rep. 711 Bureau of Resources, Dept. Science & Technology, Japan. (in japanese). 712 Seminara, G. (2010). Fluvial sedimentary patterns. Annual Review of Fluid Mechan-713 ics, 42(1), 43-66. doi: 10.1146/annurev-fluid-121108-145612 714 Shimizu, Y., & Itakura, T. (1989). Calculation of bed variation in alluvial channels. 715 Journal of Hydraulic Engineering, 115(3), 367-384. doi: 10.1061/(ASCE)0733 716 -9429(1989)115:3(367) 717 Tubino, M. (1991). Growth of alternate bars in unsteady flow. Water Resources 718 Research, 27(1), 37-52. doi: 10.1029/90WR01699 719 Venditti, J. G., Nelson, P. A., Minear, J. T., Wooster, J., & Dietrich, W. E. (2012). 720 Alternate bar response to sediment supply termination. Journal of Geophysical 721 Research: Earth Surface, 117(F2). doi: 10.1029/2011JF002254 722 Yuichi, I. (1956). Hydrodynamical study on critical tractive force. Transactions 723 of the Japan Society of Civil Engineers (in Japanese), 1956(41), 1-21. doi: 724 10.2208/jscej1949.1956.41\_1 725 Zhang, Z. (2000). A flexible new technique for camera calibration. *IEEE Transac*-726 tions on Pattern Analysis and Machine Intelligence, 22(11), 1330-1334. doi: 727 10.1109/34.888718 728



Figure 1. Aerial photos in the Naka river of Japan. (c) Google Earth



Figure 2. Plan view of the experiment flume

# 729 Appendix A Stream Tomography

732

733

Here, we describe the measurement principle of the stream tomography thatwas used in the flume experiment.

# A1 Outline of the Measurement Device and Measurement Procedure

Figure A1 and Fig. A2 show the overall plan view of the measurement device 734 and the layout of the equipment. The overall configuration of the measurement de-735 vice consists of a laser sheet light source and a traveling platform that has two dig-736 ital cameras installed. The laser sheet light source that was used in this study is a 737 yttrium aluminum garnet (YAG) laser with a wavelength of 532 nm. In addition, in 738 order to promote the emission of the laser light in water, the water that was used in 739 the flume experiment was green from dissolving sodium fluorescein. As shown in Fig. 740 A1 and Fig. A2, the two digital cameras sandwich the laser sheet light source so it 741



Figure 3. Temporal changes of the plan view in the observed bed topography

Distance from upstream [m]

Figure 4. Longitudinal view of the measured bed shape. (a) Initial stage of the experiment, (b) occurrence of the alternate bars, (c) intermediate stage of the experiment, and (d) the final stage of the experiment

---- 40 [min]

.....

6

80 [min]

180 [min]

240 [min]

742 was upstream and downstream on the traveling platform. The camera was installed so that it looks diagonally downward toward the center of the stream. The three-743 dimensional coordinates of the water level and bed level by the ST can be obtained 744 based on the intersection of the origin coordinates (lens center point) for each of the 745 two aforementioned cameras and the geometric vector that connects the water level 746 and bed position that will be measured. 747

ST is a non-contact measurement method that is based on triangulation and 748 it is used for photogrammetry. The geometric relationship in this method is shown 749 in the figures. The water level is obtained as the intersection of the two geometric 750 vectors that connect the origin coordinates for each of the two cameras and the wa-751 ter surface. The bed level can be determined as the intersection of the two geometric 752 vectors for the aforementioned water level and bed level. From these, the calculation 753

of the three-dimensional coordinates of the bed level requires consideration of the
 refraction of the irradiation light on the water surface. The three-dimensional coor dinates of the bed level can be obtained using a geometric vector that considers the
 refraction of the irradiation light. This refraction is based on the surface water level
 that was obtained by this method.

The measurement procedures consist of the following three stages. 1) Movie 759 shooting: a camera was installed on the traveling platform while running it in the 760 longitudinal direction. 2) Image analysis: after decomposing the shot video into still 761 photos, the pixel number that corresponds to the intersection of the bed and water 762 with laser light is calculated. 3) The water level and bed level are obtained by trian-763 gulation. The following sections show the geometric calculations based on the image 764 analysis and triangulation. In addition, it can be obtained either before or after the 765 movie shooting that was described in 1). Zhang's calibration method (Zhang, 2000) 766 was used to calculate the internal and external parameters of the camera. The ori-767 gin coordinates of the cameras that were installed on the upstream and downstream 768 sides of the waterway are  $c_u$  and  $c_d$ , respectively.  $c_u$  and  $c_d$  are the number vectors 769 whose components are three-dimensional spatial coordinates, and  $c_{u} = (x_{c_{u}}, y_{c_{u}}, z_{c_{u}})$ 770 and  $c_d = (x_{c_d}, y_{c_d}, z_{c_d}).$ 771

772 773

# A2 Geometric Calculation of the Water Level and Bed Level based on the Image Analysis and Triangulation

The method for detecting the water surface and bed surface in the ST is de-774 scribed as follows. From the still photograph that was created from the moving im-775 age captured by the camera, the pixel number corresponding to the intersection of 776 the laser beam and the water surface, and the intersection of the laser beam and the 777 water bed was calculated. First, the pixel number  $(i_w, j_w)$  corresponding to the in-778 tersection of the laser beam and the water surface is detected using the green bright-779 ness of the photograph as the threshold value. Similarly, the pixel numbers  $(i_b, j_b)$ 780 corresponding to the intersection of the laser beam and the bed surface are distin-781 guished from the position of the maximum green brightness. The reflection inten-782 sity on the water surface and bed surface varies depending on the environment, laser 783 light intensity, and the riverbed material. Therefore, it is necessary to adjust the 784 thresholds for detecting the water and bed surfaces according to the measurement 785 conditions. The threshold value for the green lightness in this analysis was set to 786 40 - 210.787

Next, the geometric calculation of the water level is explained. The water level 788 was calculated from the geometric relationship that is shown in Fig. A4. The three-789 dimensional coordinates of the water level h are obtained as the intersection of  $c_{wu}$ 790 and  $c_{wd}$ .  $c_{wu}$  is a geometric vector connecting  $c_u$  and the water surface level to be 791 measured  $c_u$ .  $c_{wd}$  is a geometric vector connecting  $c_d$  and the water surface level 792 to be measured  $c_d$ . In this method, the nearest points  $h_u$  and  $h_d$  for these geomet-793 ric vectors were calculated. The threshold values are set, which can be regarded as 794 the closest distance between the two points; the intersections are identified h when 795 it is below the threshold value. The latest points are calculated using the following 796 equations. 797

$$\boldsymbol{h}_{\boldsymbol{u}} = \boldsymbol{c}_{\boldsymbol{u}} + \boldsymbol{d}_{\boldsymbol{c}\boldsymbol{u}} \hat{\boldsymbol{c}}_{\boldsymbol{w}\boldsymbol{u}} \tag{A1}$$

798 799 800

801

$$\boldsymbol{h_d} = \boldsymbol{c_d} + \boldsymbol{d_{cd}} \boldsymbol{\hat{c}_{wd}} \tag{A2}$$

$$\boldsymbol{d_{cu}} = \frac{(\hat{\boldsymbol{c}}_{wd} \cdot \hat{\boldsymbol{c}}_{wu})(\hat{\boldsymbol{c}}_{wd} \cdot \boldsymbol{c}_{d} \vec{c}_{u}) - \hat{\boldsymbol{c}}_{wu} \cdot \boldsymbol{c}_{d} \vec{c}_{u}}{1 - (\hat{\boldsymbol{c}}_{wd} \cdot \hat{\boldsymbol{c}}_{wu})(\hat{\boldsymbol{c}}_{md} \cdot \hat{\boldsymbol{c}}_{wu})}$$
(A3)

802 803 804

$$\boldsymbol{d_{cu}} = \frac{(\hat{\boldsymbol{c}}_{\boldsymbol{wu}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{wd}})(\hat{\boldsymbol{c}}_{\boldsymbol{wu}} \cdot \boldsymbol{c}_{\boldsymbol{u}} \vec{c}_{d}) - \hat{\boldsymbol{c}}_{\boldsymbol{wd}} \cdot \boldsymbol{c}_{\boldsymbol{u}} \vec{c}_{d}}{1 - (\hat{\boldsymbol{c}}_{\boldsymbol{wu}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{wd}})(\hat{\boldsymbol{c}}_{\boldsymbol{wu}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{wd}})}$$
(A4)

805 806

828

842

843

849

85 852 853

$$\boldsymbol{h} = \frac{1}{2}(\boldsymbol{h}_{\boldsymbol{u}} + \boldsymbol{h}_{\boldsymbol{d}}) \tag{A5}$$

where h is a number vector whose components are the three-dimensional coordinates 807 of the calculated water level, and  $h_u$  and  $h_d$  are the number vectors of the bed level. 808 These are calculated based on the origin coordinates of the upstream and down-809 stream cameras. In addition,  $\hat{c_{wu}}$  and  $\hat{c_{wd}}$  are the unit vectors of  $c_{wu}$  and  $c_{wd}$ . 810

The water level in the transverse direction is calculated using the aforemen-811 tioned calculation. This calculation is repeated for the image that is taken by the 812 traveling platform that is moving at a constant speed. As a result, transverse wa-813 ter surface shapes in multiple longitudinal directions were obtained, and the surface 814 shape of the water level was obtained by combining these. During the final process-815 ing,  $H_{(i,j)}$  is obtained. This is a structural discrete function that is rearranged in a 816 grid pattern at arbitrary intervals that is based on h. 817

Next, the geometric calculation of the bed level is explained. The bed level was 818 calculated from the geometric relationship as shown in Fig. A5. As illustrated in 819 the figure,  $c_{biu}$  and  $c_{bid}$ , which are the geometric vectors that are incident in water, 820 are refracted on the water surface. The intersections of the geometric vector after 821 refracting from  $c_{bru}$  and  $c_{brd}$  are the bed pixels  $(i_b, j_b)$  on the image. 822

Therefore, to calculate the bed level, it is necessary to obtain the geometric 823 vectors  $c_{bru}$  and  $c_{brd}$ . Here, to prepare for the aforementioned calculation, the in-824 tersections  $c_{hu}$  and  $c_{hd}$  for the water surface and both of the geometric vectors 825  $c_{biu}$  and  $c_{bid}$  can be calculated.  $c_{bru}$  and  $c_{brd}$  can be calculated from the follow-826 ing equations while using  $H_{(i,j)}$ , which is arranged in a grid, and  $c_{biu}$  and  $c_{bid}$ . 827

$$\mathbf{\acute{c}_{hu}} = \mathbf{c_u} + \mathbf{c_{biu}} \underbrace{\overrightarrow{H_{(i,j)}\mathbf{c_u}} \cdot \mathbf{n_u}}_{\overrightarrow{H_{(i,j)}\mathbf{c_u}} \cdot \mathbf{n_u} + \overrightarrow{H_{(i,j)}\mathbf{c_{eu}}} \cdot \mathbf{n_u}}$$
(A6)

$$a_{u1} = \overrightarrow{H_{(i,j)}H_{(i+1,j)}} \times \overrightarrow{H_{(i+1,j)}} \xrightarrow{\leftarrow} (A7)$$

$$a_{u3} = H_{(i,j+1)}H_{(i,j)} \times H_{(i,j)}c_{hu}$$
(A9)
$$b_{u1} = \overrightarrow{H_{(i+1,j+1)}H_{(i,j+1)}} \times \overrightarrow{H_{(i,j+1)}c_{hu}}$$
(A10)

$$\boldsymbol{b_{u2}} = \overrightarrow{H_{(i,j+1)}H_{(i+1,j)}} \times \overrightarrow{H_{(i+1,j)}} \overrightarrow{\boldsymbol{c}_{hu}}$$
(A11)

$$\boldsymbol{b_{u3}} = \overrightarrow{H_{(i+1,j)}H_{(i+1,j+1)}} \times \overrightarrow{H_{(i+1,j+1)}} \overrightarrow{\boldsymbol{c}_{hu}}$$
(A12)

if  $(\hat{a_{u1}} = \hat{a_{u2}} = \hat{a_{u3}})$  or  $(\hat{b_{u1}} = \hat{b_{u2}} = \hat{b_{u3}})$  then, 841

$$\boldsymbol{c_{hu}} = \boldsymbol{\acute{c}}_{hu} \tag{A13}$$

$$\dot{\boldsymbol{c}}_{\boldsymbol{h}\boldsymbol{d}} = \boldsymbol{c}_{\boldsymbol{d}} + \boldsymbol{c}_{\boldsymbol{b}\boldsymbol{i}\boldsymbol{d}} \frac{\overrightarrow{H_{(i,j)}\boldsymbol{c}_{\boldsymbol{d}}} \cdot \boldsymbol{n}_{\boldsymbol{d}}}{\overrightarrow{H_{(i,j)}\boldsymbol{c}_{\boldsymbol{d}}} \cdot \boldsymbol{n}_{\boldsymbol{d}} + \overrightarrow{H_{(i,j)}\boldsymbol{c}_{\boldsymbol{e}\boldsymbol{d}}} \cdot \boldsymbol{n}_{\boldsymbol{d}}}$$
(A14)

$$a_{d1} = H_{(i,j)}H_{(i+1,j)} \times H_{(i,j+1)}c_{hd}$$

$$a_{d2} = \overrightarrow{H_{(i+1,j)}H_{(i,j+1)}} \times \overrightarrow{H_{(i,j+1)}c_{hd}}$$
(A15)
(A16)

$$a_{d2} = H_{(i+1,j)}H_{(i,j+1)} \times H_{(i,j+1)}\dot{c}_{hd}$$
(A16)

$$\boldsymbol{a_{d3}} = \overline{H_{(i,j+1)}H_{(i,j)}} \times H_{(i,j)}\boldsymbol{\acute{c}_{hd}}$$
(A17)

$$\boldsymbol{b_{d1}} = \overrightarrow{H_{(i+1,j+1)}H_{(i,j+1)}} \times \overrightarrow{H_{(i,j+1)}} \overrightarrow{\boldsymbol{b_{hd}}}$$
(A18)

$$\boldsymbol{b_{d2}} = \overrightarrow{H_{(i,j+1)}H_{(i+1,j)}} \times \overrightarrow{H_{(i+1,j)}} \overrightarrow{\boldsymbol{c}_{hd}}$$
(A19)

$$\boldsymbol{b_{d3}} = \overrightarrow{H_{(i+1,j)}H_{(i+1,j+1)}} \times \overrightarrow{H_{(i+1,j+1)}} \overrightarrow{\boldsymbol{c}_{hd}}$$
(A20)

if 
$$(\hat{a_{d1}} = \hat{a_{d2}} = \hat{a_{d3}})$$
 or  $(b_{d1} = b_{d2} = b_{d3})$  then,

$$\boldsymbol{c_{hd}} = \boldsymbol{\acute{c}}_{hd} \tag{A21}$$

868

874

878 879 880

$$\boldsymbol{c_{hu}} = \boldsymbol{c_{hu}} + \hat{\boldsymbol{c}}_{bu} \left( \frac{\overline{p_{u1}c_u} \cdot \boldsymbol{n_u}}{\overline{p_{u1}c_u} \cdot \boldsymbol{n_u} + \overline{p_{u1}c_{eu}} \cdot \boldsymbol{n_u}} \right)$$
(A22)

$$c_{hd} = c_{hd} + \hat{c}_{bd} \left( \frac{\overrightarrow{p_{d1}c_d} \cdot n_d}{\overrightarrow{p_{d1}c_d} \cdot n_d + \overrightarrow{p_{d1}c_{ed}} \cdot n_d} \right)$$
(A23)

where  $c_{hu}$ ,  $\dot{c}_{hu}$ ,  $c_{hd}$ , and  $\dot{c}_{hd}$  represent a number vector with three-dimensional coordinates. These are the starting points of the geometric vectors  $c_{bru}$  and  $c_{brd}$ , and  $p_1$ ,  $p_2$ , and  $p_3$  are the three-dimensional coordinates of the structured water surface that defines the function of each water surface.

In the second step, using Snell's law, we can calculate  $c_{bru}$  and  $c_{brd}$  with the following equations.

$$\boldsymbol{c_{bru}} = \frac{1}{n_0} \left\{ \boldsymbol{c_{biu}} + (e_u - m_u) \boldsymbol{n_u} \right\}$$
(A24)

$$e_u^{869} = -(c_{biu} \cdot n_u) \tag{A25}$$

$$m_u = \sqrt{n_0^2 + e_u^2 - 1} \tag{A26}$$

$$\boldsymbol{c_{brd}} = \frac{1}{n_0} \left\{ \boldsymbol{c_{bid}} + (e_d - m_d) \boldsymbol{n_d} \right\}$$
(A27)

$$e_d = -(\boldsymbol{c_{bid}} \cdot \boldsymbol{n_d}) \tag{A28}$$

$$m_d = \sqrt{n_0^2 + e_d^2 - 1} \tag{A29}$$

$$n_0 = \frac{n_2}{n_1} \tag{A30}$$

 $n_1$  and  $n_2$  are  $n_1 = 1.0$  and  $n_2 = 1.33$ , which represent the refractive indices of air and water, respectively.

Furthermore, in the third step,  $b_u$  and  $b_d$ , which is the closest distance between  $c_{bru}$  and  $c_{brd}$ , is calculated using the following equation.

$$\boldsymbol{b}_{\boldsymbol{u}} = \boldsymbol{c}_{\boldsymbol{h}\boldsymbol{u}} + \boldsymbol{d}_{\boldsymbol{u}\boldsymbol{2}} \hat{\boldsymbol{c}}_{\boldsymbol{b}\boldsymbol{r}\boldsymbol{u}} \tag{A31}$$

$$\boldsymbol{b_d} = \boldsymbol{c_{hd}} + \boldsymbol{d_{d2}} \hat{\boldsymbol{c}_{brd}} \tag{A32}$$

$$d_{u2} = \frac{(\hat{c}_{bru} \cdot \hat{c}_{brd})(\hat{c}_{brd} \cdot c_{hd}\vec{c}_{hu}) - \hat{c}_{bru} \cdot c_{hd}\vec{c}_{hu})}{1 - (\hat{c}_{bru} \cdot \hat{c}_{brd})(\hat{c}_{bru} \cdot \hat{c}_{brd})}$$
(A33)

$$\boldsymbol{d_{d2}} = \frac{(\hat{\boldsymbol{c}}_{\boldsymbol{bru}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{brd}})(\hat{\boldsymbol{c}}_{\boldsymbol{bru}} \cdot \boldsymbol{c}_{\boldsymbol{hu}} \hat{\boldsymbol{c}}_{\boldsymbol{hd}}) - \hat{\boldsymbol{c}}_{\boldsymbol{brd}} \cdot \boldsymbol{c}_{\boldsymbol{hu}} \hat{\boldsymbol{c}}_{\boldsymbol{hd}})}{1 - (\hat{\boldsymbol{c}}_{\boldsymbol{bru}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{brd}})(\hat{\boldsymbol{c}}_{\boldsymbol{bru}} \cdot \hat{\boldsymbol{c}}_{\boldsymbol{brd}})}$$
(A34)

890

$$\boldsymbol{b} = \frac{1}{2}(\boldsymbol{b}_{\boldsymbol{u}} + \boldsymbol{b}_{\boldsymbol{d}}) \tag{A35}$$

where  $b_u$ ,  $b_d$ , and b are the number vectors whose components are the three-dimensional coordinates of the bed level,  $\hat{b}_{bru}$ , and  $\hat{b}_{brd}$  represents the unit vectors  $b_{bru}$  and  $b_{brd}$ .

In this method, the nearest points  $b_u$  and  $b_d$  for both the vectors were calculated. If the distance between the two points below the threshold values, which can be regarded as the distance between the two points, is set sufficiently, b is identified as an intersection for both the points.

From the aforementioned, the ST acquired the three-dimensional coordinates of the water surface and bed surface.

#### A3 Verification of the Measurement Accuracy

#### A31 Experiment Outline

The measurement accuracy of the water level and bed level of the ST for the alternate bars was verified using the following procedure. The channel that was used in the experiment was the same that was used in this study. In the movable bed experiment, in which the alternate bars were formed prior to the verification of the measurement accuracy, the channel slope was set to 1/200, the flow discharge was set to 1.7 L/s, and the riverbed material with an average particle size of 0.76 mm was used.

The bed shape that was measured consisted of a fixed bed was created by sprinkling thinly the cement powder. This bed shape is the alternate bars after 120 min when the water flow started from a flat floor. Immobilization with cement powder was carried out in a section that was 2 m in the central part of the waterway used in this experiment.

After confirming the adhesion of the cement powder, the water flow was restarted, 917 and the water level and bed level were measured using the ST and point gauge. The 918 measurement with the point gauge was set at intervals of 10 cm in the longitudinal 919 direction and 2 cm in the transverse direction to capture the geometrical features of 920 the alternate bars. The measurement section in the longitudinal direction was 200 921 cm for one wavelength of the alternate bar, which formed in the experimental chan-922 nel. In addition, the measurement section in the transverse direction was 40 cm at 923 the center of the channel. This was not affected by the reflection of the irradiation 924 laser from the upper surface of the channel by the side wall. The ST measurement 925 interval was set to 1 cm in the longitudinal and transverse directions. 926

927

904

#### A32 Measurement Results

Figure A6 shows the measurement results of the water surface and bed surface. In the figure, (a) and (b) are the shapes of the photographed water surface in the flow and the bed surface after drainage, and (c) and (d) display the water surface and bed surface levels that was measured by ST. In addition, (e) and (f) present the water surface and bed surface levels that were measured by the point gauge. Finally, (g) and (h) display the difference between the point gauge and ST for the water surface and bed surface levels.

As shown in Fig. A6, a wave sequence of the standing waves with a wavelength of approximately 7—10 cm is formed on the water surface in the water flow in the longitudinal direction. In addition, from (b), it can be observed that the alternate bars with a sedimentary height of approximately 0.5 cm are formed on the left bank side. It should be noted that the characteristic shape of the front edge and small undulations formed on the bars.

Next, by comparing (c), (d), (e), and (f) in the figure, the undulations of the 941 high waves that are several centimeters in height cannot be reproduced at the wa-942 ter level and bed level at the measurement point intervals of the point gauge. Con-943 versely, it can be reproduced in ST. For the measurement targets in this study, the 944 wave with the shortest wavelength is the standing wave on the water surface. It can 945 be observed that this wave can be reproduced well with a resolution of 1 cm, which 946 can provide 10 measurement points in this wave. Figures A7 (a) and (b) show the 947 histograms that were created from Fig. A6 (g) and (h), respectively. In this exper-948 iment, the water surface was constantly changing; it did not have a constant shape. 949 Because a vertical fluctuation of approximately 2 mm was observed at the maximum 950 point, the point gauge was measured to capture the center of the fluctuation range. 951

For this reason, Fig. A7 (a) shows the effect of the time fluctuations on the water surface. Therefore, even if there is no measurement error, it should be noted that the difference between the point gauge and the ST measurement value does not become 0.

By looking at the difference between ST and the point gauge,  $\mu$  is -0.005 cm,  $\sigma$ 956 is 0.063 cm for the water level measurement,  $\mu$  is -0.203 cm, and  $\sigma$  is 0.093 cm when 957 the bed level measurement is performed. When the measured value of the point 958 gauge is used as the standard of the ST measurement accuracy, the measurement 959 accuracy of the bed level is lower than the water level from the aforementioned re-960 sults. However, it can be observed that the measurement accuracy is approximately 961 10~% of the maximum wave height of the alternate bars. From this, it can be con-962 cluded that this method has a sufficient measurement accuracy to measure alternate 963 bars. 964

# Appendix B Validity of the Pseudo-steady Flow Assumption that is applied to the Bars-Scale Riverbed Waves

This section describes the validity of the pseudo-steady flow assumption that 967 is applied to the bars-scale riverbed waves. In this study, we introduced the assumption of a pseudo-steady flow when deriving the HPDE for the bed level z. This as-969 sumption is often introduced in stability analyses of bars-scale riverbed waves(Callander, 970 1969; Kuroki & Kishi, 1984). In the aforementioned stability analysis, it is assumed 971 that the migration speed of the bed is sufficiently slower than the propagation ve-972 locity of the flow, and the flow can be treated as a pseudo-steady flow if the flow 973 rate is constant. Based on this assumption, the stability analysis ignores the term 974 of the time gradient in the continuity equation of flow and the equation of motion 975 of flow among the governing equations that are used in the analysis. The aforemen-976 tioned assumptions are considered to be valid. This is because the stability analysis 977 explains the occurrence and developmental mechanisms of the alternate bars. How-978 ever, to the best of our knowledge, whether the term of the time gradient of the flow 979 can actually be ignored cannot be confirmed from the actual phenomenon. There-980 fore, we verified whether the term of the flow time gradient can be ignored with the 981 ST measurement values and hydraulic analysis. 982

The aforementioned verification was performed by comparing the contributions of each term in the equation of motion for the flow.

985

 $\frac{1}{g}\frac{\partial u}{\partial t} + \frac{u}{g}\frac{\partial u}{\partial x} + \frac{\partial H}{\partial x} + I_{ex} = 0$ (B1)

H is the water level. As an explanation of the various physical quantities has al ready been mentioned, it is omitted here. The contribution of each term in the afore mentioned equation was calculated for each ST measurement time, and the magni tudes were compared.

 $\partial H/\partial x$  was obtained with the measured value of the water level of the ST. 990 Other terms was obtained with the results of the hydraulic analysis, which is de-991 scribed in Section 4.1 in the main text. The time interval and spatial interval of the 992 calculation were 1 min and 2 cm, which are the time resolutions and spatial reso-993 lutions of ST. The flow velocity and migration speed of the y component under the 994 experimental conditions are  $10^{-4}$  to  $10^{1}$  of the x components at any location regard-995 less of the developmental state of the alternate bars. For simplicity, the y component 996 is ignored in this section. 997

Figure B1 shows the time change of the box-beard diagram that displays the contribution of each term. This figure shows the (a) local term, (b) advection term, (c) pressure term, and (d) friction term; they correspond to the order of each term

in Eq. (B1). By looking at the figure, although the (b) advection term, (c) pressure 1001 term, and (d) friction term dominate the flow at any time, it can be confirmed that 1002 (a) the local term can be ignored because the local term is smaller than the afore-1003 mentioned three terms. Even if the advection term with the smallest contribution in 1004 (b),(c), and(d) is compared with the local term, the contribution of the local term is 1005  $10^{-4}$  to  $10^{-2}$  of the (b) advection term. In addition, it can be observed that the lo-1006 cal term is extremely small. From this, it is inferred that it is physically appropriate 1007 to ignore the time gradient of the flow in the alternate bars. 1008

# Appendix C Derivation of the Two-Dimensional Equation of the Water Surface Profile

Appendix C presents the derivation processes of the two-dimensional equation of the water surface profile to derive the HPDE for the bed level. The governing equations that were used for the derivation consist of the following continuous equations and the equations of motion. When deriving the equation, the flow can be treated as a pseudo-steady-state flow based on the verification results in Appendix B. Therefore, the following continuous equations and equations of motion are used for the derivation.

$$\frac{\partial [hu]}{\partial x} + \frac{\partial [hv]}{\partial y} = 0 \tag{C1}$$

$$u \partial u = u \partial u = \partial z = \partial b$$

$$\frac{u}{g}\frac{\partial u}{\partial x} + \frac{v}{g}\frac{\partial u}{\partial y} + \frac{\partial z}{\partial x} + \frac{\partial h}{\partial x} + I_{ex} = 0$$
(C2)

$$\frac{u}{g}\frac{\partial v}{\partial x} + \frac{v}{g}\frac{\partial v}{\partial y} + \frac{\partial z}{\partial y} + \frac{\partial h}{\partial y} + I_{ey} = 0$$
(C3)

As an explanation of the various physical quantities has already been mentioned, it is omitted here.

The derivation of  $\partial h/\partial x$  is described as follows. First, applying the product rule to Eq. (C1) results in the following equation.

$$h\frac{\partial u}{\partial x} + u\frac{\partial h}{\partial x} + h\frac{\partial v}{\partial y} + v\frac{\partial h}{\partial y} = 0 \tag{C4}$$

<sup>1028</sup> Next, for the first and third terms on the left side of Eq. (C4),

1018 1019

1020 1021 1022

1027

1029

1

$$u = \frac{1}{n} I_{ex}{}^{1/2} h^{2/3} \tag{C5}$$

$$v = \frac{1}{n} I_{ey}^{1/2} h^{2/3}$$
(C6)

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial h}\frac{\partial h}{\partial x} + \frac{\partial u}{\partial I_{ex}}\frac{\partial I_{ex}}{\partial x} = \frac{2}{3}\frac{u}{h}\frac{\partial h}{\partial x} + \frac{1}{2}\frac{u}{I_{ex}}\frac{\partial I_{ex}}{\partial x}$$
(C7)

$$\frac{\partial v}{\partial y} = \frac{\partial v}{\partial h}\frac{\partial h}{\partial y} + \frac{\partial v}{\partial I_{ey}}\frac{\partial I_{ey}}{\partial y} = \frac{2}{3}\frac{v}{h}\frac{\partial h}{\partial y} + \frac{1}{2}\frac{v}{I_{ey}}\frac{\partial I_{ey}}{\partial y}$$
(C8)

After the differentiation of the composite function (Eq. (C7) and Eq. (C8)) using Manning's flow velocity formula (Eq. (C5), Eq. (C6)), substituting it into Eq. (C4), and rearranging  $\partial h/\partial x$ , the following equation is obtained.

$$\frac{\partial h}{\partial x} = -\frac{v}{u}\frac{\partial h}{\partial y} - \frac{3}{10}\frac{h}{I_{ex}}\frac{\partial I_{ex}}{\partial x} - \frac{3}{10}\frac{vh}{uI_{ey}}\frac{\partial I_{ey}}{\partial y}$$
(C9)

Next, after substituting Eq. (C7) and the following Eq. (C10) into the first  
and second terms of the equation of motion in the 
$$x$$
 direction for Eq. (C2),

$$\frac{\partial u}{\partial y} = \frac{\partial u}{\partial h}\frac{\partial h}{\partial y} + \frac{\partial u}{\partial I_{ex}}\frac{\partial I_{ex}}{\partial y} = \frac{2}{3}\frac{u}{h}\frac{\partial h}{\partial y} + \frac{1}{2}\frac{u}{I_{ex}}\frac{\partial I_{ex}}{\partial y}$$
(C10)

the following equation is obtained. After substituting Eq. (C9), which was organized earlier into Eq. (C11),

$$\frac{3}{10}\frac{u^2}{gI_{ex}}\frac{\partial I_{ex}}{\partial x} - \frac{1}{5}\frac{uv}{gI_{ey}}\frac{\partial I_{ey}}{\partial y} + \frac{1}{2}\frac{uv}{gI_{ex}}\frac{\partial I_{ex}}{\partial y} + \frac{\partial z}{\partial x} - \frac{v}{u}\frac{\partial h}{\partial y} - \frac{3}{10}\frac{h}{I_{ex}}\frac{\partial I_{ex}}{\partial x} - \frac{3}{10}\frac{vh}{uI_{ey}}\frac{\partial I_{ey}}{\partial y} + I_{ex} = 0$$
(C11)

1046 The following equation can be obtained by rearranging  $\partial h/\partial y$ .

$$\frac{\partial h}{\partial y} = \frac{3}{10} \frac{u^3}{vgI_{ex}} \frac{\partial I_{ex}}{\partial x} - \frac{1}{5} \frac{u^2}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} + \frac{1}{2} \frac{u^2}{gI_{ex}} \frac{\partial I_{ex}}{\partial y} 
+ \frac{u}{v} \frac{\partial z}{\partial x} - \frac{3}{10} \frac{uh}{vI_{ex}} \frac{\partial I_{ex}}{\partial x} - \frac{3}{10} \frac{h}{I_{ey}} \frac{\partial I_{ey}}{\partial y} + \frac{u}{v} I_{ex}$$
(C12)

<sup>1048</sup> After substituting Eq. (C12) into Eq. (C9) and rearranging it, the following  $\partial h/\partial x$ <sup>1049</sup> is derived.

$$\frac{\partial h}{\partial x} = -\frac{\partial z}{\partial x} - I_{ex} - \frac{3}{10} \frac{u^2}{gI_{ex}} \frac{\partial I_{ex}}{\partial x} + \frac{1}{5} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} - \frac{1}{2} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial y}$$
(C13)

<sup>1051</sup> By rearranging  $\partial h/\partial y$  using the same process as earlier, the following equation <sup>1052</sup> for  $\partial h/\partial y$  is obtained.

$$\frac{\partial h}{\partial y} = -\frac{\partial z}{\partial y} - I_{ey} - \frac{3}{10} \frac{v^2}{gI_{ey}} \frac{\partial I_{ey}}{\partial y} 
-\frac{1}{2} \frac{uv}{gI_{ey}} \frac{\partial I_{ey}}{\partial x} + \frac{1}{5} \frac{uv}{gI_{ex}} \frac{\partial I_{ex}}{\partial x}$$
(C14)

1053

1047

1050



Figure 5. Temporal changes of the plan view for the observed water depth

Figure 6. Difference between the measured value and the calculated value of the water depth that is made dimensionless by the measured value.



Figure 7. Temporal changes of the plan view for the calculated flow velocity



Figure 8. Bed-level time waveform. (a) Left bank side, (b) center, (c) right bank side



Figure 9. Temporal changes of the plan view for the observed bed topography and  $\Delta z_*$ 



Figure 10. Temporal changes of the plan view for the observed bed topography and the calculated migrating speed



Figure 11. Histograms of the migrating speed



Figure 12. Relationship between the migrating speed, energy slope, and water depth



Figure 13. Relationship between the migrating speed, energy slope, and water surface slope



Figure 14. Longitudinal view of the (a) cross-sectional averaged migrating speed, (b) and cross-sectional averaged bed level







Figure A2. Equipment layout of the measurement equipment



Figure A3. Outline of the measurement principle of the water surface and bed surface by triangulation



Figure A4. Relationship between h and H that is rearranged in a grid by triangulation



Figure A5. Outline of the bed surface measurement method while considering the refraction on the water surface



Figure A6. Measurement results of the water surface and bed surface



Figure A7. Histogram of the difference for the ST and point gauge measurements



Figure B1. Temporal changes of the box plots for the (a) local term, (b) advection term, (c) pressure term, (d) and friction term