Revealing a High Water Abundance in the Upper Mesosphere of Mars with ACS onboard TGO

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Abstract

We present the first water vapor profiles encompassing the upper mesosphere of Mars, 100–120 km, far exceeding the maximum altitudes where remote sensing has been able to observe water to date. Our results are based on solar occultation measurements by Atmospheric Chemistry Suite (ACS) onboard the ExoMars Trace Gas Orbiter (TGO). The observed wavelength range around 2.7 μ m possesses strong CO2 and H2O absorption lines allowing sensitive temperature and density retrievals. We report a maximum H2O mixing ratio varying from 10 to 50 ppmv at 100–120 km during the global dust storm (GDS) of Martian Year (MY) 34 and around southern summer solstice of MY 34 and 35. During other seasons water remains persistently below ~2 ppmv. We claim that contributions of the MY34 GDS and perihelion periods into the projected hydrogen escape from Mars are nearly equivalent.

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10 Key Points:

- Water abundances are reported in a previously unexplored altitude range: from 100 to
 120 km
- The observed GDS (MY34) and two perihelion seasons (MY34, 35) reveal the H2O content around 10-50 parts per million by volume at 100-120 km
- Contributions of the MY34 GDS and perihelion periods into the projected hydrogen
 escape from Mars are nearly equivalent

17

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- 20 km, far exceeding the maximum altitudes where remote sensing has been able to observe water
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- around 2.7 μ m possesses strong CO₂ and H₂O absorption lines allowing sensitive temperature
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- summer solstice of MY 34 and 35. During other seasons water remains persistently below ~ 2
- 27 ppmv. We claim that contributions of the MY34 GDS and perihelion periods into the projected
- hydrogen escape from Mars are nearly equivalent.

29 Plain Language Summary

- 30 We report regular events of high abundances of the water vapor (H_2O) in the upper atmosphere
- of Mars (100–120 km). So far, any water enrichment has not been revealed by remote sensing at
- 32 such high altitudes. Higher than 80 km, solar light breaks water vapor molecules into H and O
- atoms, which may reach the exosphere and escape the planet. When Mars is closer to the Sun
- 34 (the perihelion season), the atmosphere's circulation intensifies, causing increased dust activity
- 35 with global dust storms (GDS), occurring every 3–4 Mars years. We observed during the second
- halves of Martian years 34 and 35 (2018–2020), including one GDS and two perihelion seasons.
- We report that the maximum water relative abundance reaches 10-50 parts per million in volume
- (ppmv) at 100-120 km during the GDS and every perihelion season. These high values indicate that the Martian atmosphere above 100 km regularly hosts large amounts of water, facilitating
- that the long terms accore of water, facilitat
- 40 the long-term escape of water from the planet.

41 **1 Introduction**

- 42 The vertical distribution of water vapor (H_2O) on Mars is an indicator of the intricate coupling of distinct phenomena: temperature variations, cloud formation, sublimation, turbulent 43 and convective mixing, as well as general circulation and wave/eddy transport. H₂O has long 44 been thought to remain confined below the hygropause, which is the level where the saturation 45 condition is met and where water ice clouds may form, as occurs on Earth. The existence of this 46 layer on Mars was established for the first time by ground-based microwave soundings of Clancy 47 et al. (1996) with a saturation level between 10-20 km around the aphelion, i.e., Solar 48 Longitudes (L_s) 70°, and 40–60 km around perihelion (L_s 250°). In parallel, Rodin et al. (1997) 49 reported water vapor profiles retrieved from the solar occultations made by Auguste on Phobos-2 50 in 1989. The existence of a hygropause at 30–35 km (with a mixing ratio of 3 ppm) in the 51 northern spring ($L_s=0^{\circ}-20^{\circ}$) near the equator was subsequently claimed. The first climatology of 52 water vapor profiles was derived from SPICAM-IR solar occultations on Mars Express (MEx) 53 (Fedorova et al., 2009; 2018; 2021; Maltagliati et al., 2013), covering eight Martian years. The 54 hygropause level was found to vary from 40 to 80 km depending on season, latitude, and dust 55 events. Hygropause is also indirectly sensed in CRISM limb profiles of $O_2(^1\Delta g)$ emission, a 56 confident indicator of O₃, from which water vapor mixing ratios were inferred by Clancy et al. 57 58 (2017).
- The observation of large amounts of water vapor in and above the middle atmosphere (>40 km, Maltagliati et al., 2013) was then complemented by the discovery of short-term decline

of the hydrogen corona brightness over several weeks (Chaffin et al., 2014; Clarke et al., 2014).

62 This variability exposed a new paradigm in our perception of how water escapes from Mars

63 (Chaffin et al., 2017). So far, water escape was thought to be controlled by a slow conversion

64 process involving H_2 , formed from the catalytic recombination of carbon dioxide with odd

hydrogen (McElroy and Donahue, 1972; Krasnopolsky, 2002). The non-condensable H_2 can overcome the hygropause and reach the mesosphere (80–120 km), while transported by turbulen

overcome the hygropause and reach the mesosphere (80–120 km), while transported by turbulent
 mixing or circulation. There, it can dissociate and release H atoms that will escape the planet

68 once above the exobase.

Observations have revealed that water vapor transport from the troposphere to the lower 69 mesosphere of Mars occurs during the dusty season and is enhanced at times of major dust 70 storms. In particular, a significant H₂O enhancement in the middle atmosphere was observed 71 72 during the global dust storm (GDS) in 2007 (MY28) with a rise of the hygropause altitude to >60 km (Fedorova et al., 2018; Heavens et al., 2018; 2019). Sensitive solar occultation measurements 73 by NOMAD and ACS NIR instruments onboard the ExoMars Trace Gas Orbiter (TGO) have 74 showed that water vapor reached 80-100 km (Aoki et al., 2019; Fedorova et al., 2020) during 75 two storms in 2018 and 2019 (a global one at L_s 190°–220° and a regional one at L_s 330° in 76 MY34; Montabone et al., 2020). Fedorova et al. (2020) revealed the water supersaturation at 70-77 90 km even in the presence of H₂O ice clouds not only during the GDS but also near the 78 Southern summer solstice ($L_{s}\sim 270^{\circ}$) when water reached 90-100 km as well. Altogether, these 79 studies promote a new mechanism for controlling H escape through direct delivery at above 80 80 km and further photodissociation of H₂O molecules (Chaffin et al., 2017; Krasnopolsky et al., 81 2019). General circulation models predict an upward water flux into the thermosphere (>120 km) 82 during the GDS and perihelion periods (Shaposhnikov et al., 2019; Neary et al., 2020; Rossi et 83 al. 2021). 84

The discussion regarding a relative contribution of perihelion or GDS to the mesospheric 85 water enrichment was recently stimulated by SPICAM/MEx long-term observations covering 86 87 Martian Years 28 through 35. Here, Fedorova et al. (2021) claimed an annual rise of water abundance up to ~90 km in perihelion, which is compatible with GDS enhancements. The new 88 ACS/TGO dataset confirms those conclusions for altitudes below 100 km in MY34-MY35 89 (Fedorova et al. 2020, Alday et al., 2021). In parallel, during the perihelion season, the D/H ratio 90 in water decreases with altitude from 4-6 times SMOW (Standard Mean Ocean Water) in the 91 lower atmosphere to 2–3 times in the mesosphere (50-70 km) as measured by ACS MIR (Alday 92 93 et al., 2021) and NOMAD (Villanueva et al., 2021) spectrometers. Alday et al. (2021) show that ultraviolet H₂O photolysis dominates the production of H relative to D atoms in the upper 94 95 atmosphere.

From above, ion chemistry in the thermosphere has been characterized by the NGIMS 96 mass-spectrometer on MAVEN (Benna et al., 2015) and interpreted by the ionospheric model of 97 Fox et al. (2015). Using NGIMS data, Stone et al. (2020) measured H₂O ion concentrations 98 99 around ~150 km for the 2014–2018 period (MY32–MY34). With the help of the model by Fox et al. (2015), Stone et al. (2020) found the relative abundance of water at this altitude on the 100 dayside varying seasonally on average from 2 to 5 ppm. Several enhanced dusty episodes disrupt 101 this seasonal signal: 3–9 ppm during the regional storm of MY32, 10–20 ppm during the storm 102 of MY33, and up to 60 ppm in the GDS of MY34. Stone et al. (2020) concluded that water 103 104 transport into the ionosphere and its destruction are the main mechanisms in the overall hydrogen escape from Mars. 105

We used the data of the middle infrared spectrometer of the Atmospheric Chemistry Suite 106 107 (ACS MIR) onboard the ExoMars TGO, which measures water vapor VMR and atmospheric density in a wide range of altitudes, from the troposphere to the lower thermosphere, using the 108 109 strong absorption bands of H₂O and CO₂ around 2.66–2.70 µm. The high spectral resolution and the good signal-to-noise ratio of ACS MIR allow the measurements of water profiles up to 120 110 km, inaccessible altitudes for the ACS NIR and SPICAM measurements, sensing the 1.38 µm 111 absorption band (Fedorova et al., 2020; 2021). The strong H₂O absorption around 2.6 µm is also 112 used by NOMAD, yielding water profiles up to ~90 km (Aoki et al., 2019; Villanueva et al., 113 2021). 114

Here we report the first water vapor abundance measurements in the upper mesosphere (up to 120 km) of Mars. The goal of our paper is to compare the mesospheric water behavior between the second halves of MY34 and MY35 when the high H₂O content is observed. We aim to clarify the principal mechanism of H₂O delivery to the upper mesosphere: it is sporadic dust events, or the result of seasonal variability in the Martian circulation that peaks each year, around Southern summer solstice. For that, we analyze seasonal and latitudinal variations of H₂O VMR

121 vertical profiles retrieved from the ACS MIR solar occultation experiment.

122 2 Measurements and dataset overview

123 2.1 ACS MIR spectroscopy and retrievals

ACS MIR, a solar occultation cross-dispersion echelle spectrometer, records spectra from 124 a set of adjacent diffraction orders (from 10 to 20 per occultation) projected onto a 2D detector 125 array (Korablev et al., 2018). To retrieve high altitude water vapor abundances together with the 126 127 atmospheric temperature and pressure, we use MIR spectra from the diffraction order #223. They cover a narrow wavelength interval of 2.66–2.68 μ m (3733–3753 cm⁻¹), including a part of the 128 2.7-µm CO₂ absorption band and a few strong H₂O lines near 2.66 µm (Fig. 1a, 1b, Fig. S1). The 129 instrument's spectral resolution is ~ 0.15 cm⁻¹, while the signal-to-noise ratio ranges from 2,000 130 to 4,000, which provides high sensitivity for detections in the upper atmosphere where 131 atmospheric constituent densities are low. Temperature (Fig. 1c) is retrieved by fitting a 132 synthetic model to the CO₂ rotational band taking advantage of its temperature dependence as 133 seen in Fig. S1 of the Supplementary Material (SM). This procedure was applied iteratively, with 134 the pressure calculated from the retrieved temperature profile under the assumption of 135 hydrostatic equilibrium. The temperature measurements were then validated against those made 136 by MIR near the 2.6 μ m CO₂ band (Alday et al., 2019) and by ACS NIR around the 1.58 μ m 137 band (Fedorova et al., 2020). As a result, one occultation session allows us to simultaneously 138 retrieve profiles of pressure and temperature (from CO₂ absorption bands) and the H₂O number 139 density (Fig. 2c). The water abundance can then be expressed relative to the total atmospheric 140 density, that is, in VMR (in ppmv). Specific details of the algorithms pertaining to this work can 141

142 be found in SM.

The dataset analyzed here consists of a series of transmission spectra obtained during a solar occultation while the line of sight of the instrument progressively penetrates from the upper into deeper layers of the atmosphere, or vice versa (see examples in Fig. 1a, 1b). The transmission is determined as the solar spectrum ratio measured through the atmosphere to the reference one, taken from the data above a tangent height of 200 km. This altitude level is negligibly attenuated by the atmosphere even within the very strong CO₂ band system at 2.7 µm.

149 The typical integration time is 2 seconds, which provides an altitude resolution ranging from 0.5

- to 2.5 km, depending on the occultation duration. It gives sufficiently fine vertical sampling for an atmosphere whose scale height ranges from 5 to 10 km depending on temperature. The
- instrument field of view projected at the limb is around 1-3 km in altitude equivalent.





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with the best-fit models, including both CO_2 and H_2O absorptions (black), and only CO_2

- absorption (green). Blue arrows indicate water absorption lines. Zoom in (a) shows a part with
- the strongest H₂O absorption detected at 115 km. (c) Retrieved vertical profiles of temperature
- 159 (left), number densities (center), and H₂O volume mixing ratio (VMR) (right). The atmospheric

density (grey squares) is scaled by the factor of 10^{-6} . Black triangles mark H₂O upper limits (see SM for the description of uncertainties).

162 2.2 Data selection

Our measurements of the mesospheric water focus on the second halves of MY34 and 163 MY35, which correspond to ACS MIR observations from May 2018 to March 2019 and from 164 April 2020 to January 2021. The selected dataset comprises 187 occultation sessions in the 165 Northern Hemisphere and 156 sessions in the Southern Hemisphere, encompassing seasonal 166 periods from L_S 180° to 355° in MY34 and from L_S 185° to 356° in MY35 (Fig. 2a). Figure 2a 167 shows the latitude coverage with the corresponding aerosol activity, which was defined for each 168 occultation at the altitude level where the slant opacity equals 0.3 (~0.75 of the atmospheric 169 transmittance in the continuum). Measurements in the Northern Hemisphere occurred mostly in 170 the high latitude range, between 40°N and 70°N. In the South, the perihelion observations (Ls 171 172 270°) were made in mid-latitudes, while the rest of occultations occurred close to the polar region (60°S–90°S). Only a few sessions were localized nearby the equator: at $L_S \sim 240^\circ$ and 173 L_{s} ~300° of MY34 and at L_{s} ~210° and L_{s} ~280° of MY35. These observations are accompanied 174

by a higher aerosol loading than for high latitude and polar regions (Fig. 2a).

176 **3 Seasonal variation of altitude profiles**

Observations in the second halves of MY34 and MY35 uncover events, which drastically 177 perturbed the temperature and water vapor vertical distributions. The peculiar pattern to compare 178 with is the MY34 GDS and perihelion periods in MY34 versus MY35, which had no GDS but a 179 regional dust activity in its second half. The seasonal variation of the processed altitude profiles 180 is presented in Figure 2(b, c). We binned the profiles into intervals of 2° in solar longitude and 2 181 km in altitude. Depending on the L_s and altitude sampling, the value in each bin is calculated as 182 the weighted mean of one to five individual points. We excluded all points with 1-sigma 183 uncertainties exceeding 20 K in temperature and 100% of the H₂O mixing ratio. The second 184 rejection criterion corresponds to the detection limit ($\sim 10^7$ cm⁻³) of water number density (see in 185

Fig. 1c) that defines the seasonal variations of the uppermost detectable points in Figure 2c.

We observe seasonal temperature (Fig. 2b) and H₂O (Fig. 2c) peaks in the middle 187 atmosphere (40–80 km) during the GDS of MY34, L_s 190°–220°, and an additional smaller peak 188 189 at $L_{s} 320^{\circ} - 330^{\circ}$, corresponding to a regional storm. Moreover, the rise of water vapor to higher altitudes, up to the mesopause at 110–130 km where temperature encounters a minimum, is 190 observed during the two perihelion intervals (Ls 250°-290°) of MY34 and MY35. Here, the 191 192 Southern summer (Fig. 2c, right panel) is accompanied by a more humid mesosphere (40–60 ppm of H₂O) than the Northern Winter (Fig. 2c, left panel) where the mean mesospheric water 193 194 reaches 20–30 ppmv on average between 80 and 120 km. In contrast, out of the perihelion peak 195 or dust events, i.e. for the selected data at the beginning of the MY34 GDS and at the very end of MY 34, 35, water content above 80 km never exceeds 2-3 ppmv. 196





205 4 H₂O variations around perihelion

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To quantify seasonal trends of water content in the mesosphere, we selected three altitude layers corresponding to 80 km, 100 km, and 110–120 km. The first layer, which corresponds to the middle mesosphere, is accessible in all profiles (Fig. 2c) when the vapor concentration exceeds the detection limit of ~ 10^7 cm⁻³, even in low water loading periods. Water at 100–120 km shows up only in stormy periods and around perihelion (Fig. 2c).

Observed variations during perihelion for the three selected levels are presented in Figure 211 3 for both Martian Years. The number of MIR observations at the considered spectral range is 212 low during the dust storm activity of MY34. Nevertheless, a comparison with MY35 reveals 213 significant increases of H₂O mixing ratios during the GDS: by a factor of 6–8 at 80 km (Fig. 3e, 214 3f) and by a factor of 3–5 at 100-120 km from L_S 190° to 220° (Fig. 3a-3d). Increases at L_S 215 320° -330° follow annually repeatable dust storm activity at this season, although injecting far 216 217 less water into the mesosphere than the GDS in MY34. Around Mars perihelion ($L_s=240^{\circ}-300^{\circ}$) water behaves almost identically between MY34 and MY35. For both Martian Years, the 218 maximum H₂O mixing ratio was observed near the Southern summer solstice ($L_{S} \sim 270^{\circ}$), 219 reaching values of 40-80 ppm at 80 km, 30-60 ppm at 100 km, and 20-50 ppm at 110-120 km. 220 In the Northern winter solstice, it varied from 20 to 40 ppm at all levels, 80-120 km. There are 221 groups of points out of general behaviour, i.e. at $L_s=270^{\circ}-280^{\circ}$ in Fig. 3, that results from 222 223 latitudinal variations of the water content (see Fig. S3 in SM). We compare our results at 80 km

with the corresponding ACS NIR dataset derived from the MY34 profiles of Fedorova et al.

225 (2020, grey points in Fig. 3e, 3f). The NIR dataset is five times denser than used in the present

work, and it observed the H₂O seasonal variations in greater detail, especially during the dust events of MY34.



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Figure 3. Seasonal trends of H₂O volume mixing ratio (VMR) at three altitude levels: 80, 100,
110-120 km. The season is the second halves of MY34 (in blue) and of MY35 (in red). Each
point corresponds to an individual vertical profile: weighted mean value obtained in between
110-120 km (a, b), and interpolated value for the levels of 100 km (c, d) and 80 km (e, f). Data at

80 km in grey (**e**, **f**) are taken from the ACS NIR profiles (Fedorova et al., 2020).

234 4 Discussion and Conclusions

For the first time, we report observations of H₂O abundances in a previously unexplored 235 altitude range (from 100 to 120 km). There we find 10-30 ppm of water vapor during the MY34 236 GDS and 20–50 ppm around Mars perihelions ($L_s=250^\circ-290^\circ$) of MY34 and MY35 in both 237 hemispheres. Our GDS retrievals at 100-120 km are of the same order of magnitude as MAVEN 238 NGIMS results reported by Stone et al. (2020) at ~150 km. Surprisingly, NGIMS water 239 240 abundances reveal a 2014-2018 mission-wise maximum of H₂O at 150 km, only during the MY34 GDS, whereas we repeatedly observe the annual maximum around the Southern summer 241 solstice both in MY34 and MY35. 242

NGIMS measures $[H_2O^+]$ ions, from which neutral H_2O abundances at 150 km were derived on the basis of 1D photochemical modeling. The model was adjusted to reproduce the H_2O VMR at 150 km inferred from the $[H_2O^+]$ ions measured under two scenarios: low water, corresponding to 2 ppm prescribed at 80 km in a non-GDS case; and high water of 40 ppm in a GDS case (Stone et al., 2020; as corrected in March 2021). Notably, all the solar occultation observations performed by TGO and MEX to date (Fedorova et al., 2018, 2020, 2021; Aoki et al., 2019; Villanueva et al., 2021), including the present dataset (Fig. 3e, 3f), report even higher
water vapor VMRs at 80 km, of 50–80 ppm during the GDS. Stone et al. (2020) indicate a
systematic uncertainty of 69% on their neutral H₂O inference, which is consistent with observed
MAVEN, TGO, and MEX values within such error bars.

It is important to consider how ACS's high altitude water vapor abundances combine 253 254 with photolysis since this process has been hypothesized to be essential, if not the dominant, source for the H atoms observed in the exosphere (Chaffin et al., 2017). The conclusion of Stone 255 et al. (2020) argues for the GDS's predominance and related ion chemistry in the H atoms' 256 production. Our observations suggest that while the GDS period corresponds to the maximum of 257 water abundance at 80 km, H₂O at 120 km peaks only later, at the Southern summer solstice, 258 when it is twice as large as during the GDS. This enhanced solstice maximum suggests that 259 relative water abundance declines more rapidly above 80 km during the GDS than after, during 260 perihelion (Fig. 4). 261



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Figure 4. Altitude profiles of average H₂O volume mixing ratio (VMR) during the GDS (MY 263 34) and the perihelion season (MY 34, 35). The dataset used includes all ACS MIR observations 264 highlighted in Fig. 2: GDS of MY34 (Ls=195°-220°) (brown); the bin around the perihelion 265 point ($L_s=250^{\circ}-295^{\circ}$) (blue). Panels (**a**, **b**) for MY34 and (**c**, **d**) for MY35 from the Northern (**a**, 266 c) and Southern (b, d) hemispheres. Profiles are presented with 1-sigma dispersion over 4-km 267 altitude bins. Light grey curves are the averages of the NIR data during GDS of MY34 268 (Fedorova et al., 2020). The GDS curves are also indicated for MY35 to facilitate comparison. 269 For the data points used in the averaging see Fig. S4 of SM. 270

In Figure 4, we combined altitude profiles from GDS-only (Ls $195^{\circ}-220^{\circ}$) and perihelion (Ls $250^{\circ}-295^{\circ}$) intervals to compare averaged vertical trends between them. Here, we also see a coincidence between MIR and NIR GDS profiles in frames of dispersions, which reflect high variability of the observed GDS points (Fig. S4 of SM). The considered H₂O distributions allow

estimating an integral escape flux of the atomic hydrogen in each case. For that, we applied the 275 model of Chaffin et al. (2017), which predicts the atmospheric escape rate depending on the 276 water injection into different altitudes (see Figure 3 of their paper). Our rough calculations show 277 that the H escape flux is about $\sim 5 \times 10^9$ cm⁻¹s⁻¹ during the considered intervals of MY34's GDS 278 and Southern summer solstices of MY 34 and 35. Thus, we claim nearly equivalent contributions 279 from a single GDS and the perihelion period into the hydrogen escape by the high water 280 enrichment in the middle/upper atmosphere. Fedorova et al. (2021) come to a similar conclusion 281 based on SPICAM/MEx water profiles up to 80 km. A GDS occurs every 3-4 martian years on 282 average (Zurek and Martin, 1993; Wang and Richardson, 2015), making the yearly perihelion 283 contribution to the hydrogen escape reasonable. The water enhancements are tied to the 284 circulation regime (Clancy et al., 1996; Richardson and Wilson, 2002; Montmessin et al., 2005). 285 More measurements and modeling would be needed to decide whether the southern summer 286 solstitial transport, currently near perihelion, or the GDS equinoctial circulation prevailed during 287 the history of Mars. 288

Overall, our results cannot be easily reconciled with water values (up to 60 ppm at 150 km) inferred from NGIMS ion measurements, which suggested that the thermosphere hosted much more water during the GDS than during the rest of the year. However, we note that the only time when Stone et al. (2020) reported measurements around perihelion concerned the Ls interval between 240° and 265° of MY33 (Figure 4 of Stone et al., 2020). It showed the same rough trend as during the onset of the MY34 GDS, with values far exceeding those reported for the regional dust storm of MY33, still a factor of 3 smaller than during the GDS.

Our results remain in line with the conclusion of Fedorova et al. (2020, 2021) that the perihelion season is the primary conveyor of water to high altitudes on a long-term basis. The high values above 100 km fill the gap between the water observed below 100 km and water ions measured by NGIMS at 150 km. Both measurements bring unique constraints in our attempt to understand how the water in the lower atmosphere connects with the escaping hydrogen in the exosphere, an essential step before confidently extrapolating the water escape back in time.

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307 available from ESA's Planetary Science Archive at

308 https://archives.esac.esa.int/psa/#!Table%20View/ACS=instrument. The retrieved data with

altitude profiles of H₂O VMR for the considered seasons are available at

- 310 https://data.mendeley.com/datasets/995y7ymdgm/draft?a=daa72362-898d-4c86-8a13-
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Supporting Information for

Revealing a High Water Abundance in the Upper Mesosphere of Mars with ACS onboard TGO

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Introduction

The present Supporting Information describes details of the ACS MIR instrument and the concept of altitude profiles retrievals from the measured transmission spectra. The retrieval algorithm is supplemented by Figure S1, which demonstrates instrumental and temperature peculiarities of the CO₂ and H₂O absorption spectroscopy around 2.7 μ m band (echelle order #223 of ACS MIR). Figure S2 shows validation of the derived temperature, H₂O number density and mixing ratio profiles for different seasonal scenarios in MY34. Latitudinal variation of the high altitude water around perihelion is presented in Figure S3 for both MY34 and MY35. Figure S4 presents vertical distribution of all data points for H₂O VMR in the GDS and perihelion seasons. Uncertainties of the obtained data are described separately in section Text S3.

Text S1. Instrument description

The middle infrared channel of the Atmospheric Chemistry Suite is a cross-dispersion echelle spectrometer dedicated to solar occultation measurements in the 2.3-4.3 µm wavelength range (Korablev et al., 2018). Each occultation session is devoted to one of ten angular positions of the MIR secondary grating that disperses about 10-20 echelle orders spatially separated and recorded simultaneously at the focal plane array (FPA) by 640x512 pixels. The spectral range of our interest, 2.66-2.68 μm, lies within the diffraction order #223 at the grating position #4. One order covers a spectral interval of about 30 cm⁻¹ (~25 nm); a spectrum is dispersed along 640 FPA elements with the sampling of 0.05 cm⁻¹ and the resolving power $\lambda/\delta\lambda$ reaching ~25000 $(\delta\lambda \sim 0.15$ cm⁻¹). In the occultation field of view (FOV), the instrumental rectangular slit cuts a part from the solar disk, so that one order occupies a stripe with about 20 FPA rows. In such a manner, the consequences of orders are located on the matrix as stripe-by-stripe, one above the other, with some dark rows between them. For our analysis, we selected a central row from the stripe correspondent to order #223 that provides the maximal signal-to-noise ratio. The calibration procedure for the analyzed transmissions encompassed a pixel-to-wavenumber assignment and a determination of the instrument line shape (ILS). The procedure was applied for every occultation and in various spectral intervals separately since we observed variability of the light spreading over the FPA from orbit to orbit. Using strong CO₂ and H₂O absorption lines in order #223, we derived an optimal wavenumber assignment by a parabolic law and an asymmetric ILS including a superposition of two Gaussian functions (Fig. S1a). Previously, Alday et al. (2019) implemented an analogous approach for MIR spectra but in the grating position #5.

Text S2. Retrieval concept

A scheme for the water VMR retrieval consists of several iterations with a fitting of a forwardly modeled transmission spectrum to a measured one at each observed altitude. The forward modeling includes contributions from H₂O, HDO, and CO₂ molecular absorption cross-sections that are abundant in the 2.66-2.68 μm spectral range, i.e. 3730-3755 cm⁻¹ (Fig. S1b). They were calculated using line-by-line modeling for specified temperatures and pressures on a basis of the HITRAN2016 database (Gordon et al., 2017) taking into account H₂O and HDO lines broadening in the CO₂-rich atmosphere (Gamache et al., 2016; Malathy Devi et al., 2017) and self-broadening in the case of CO_2 . A transmission spectrum J(v, T, p) is defined by the crosssections $\sigma(v, T, p)$, gaseous concentrations $N_{mol} = f_{mol} \bullet p/k_B T$, and aerosol slant opacity τ_a . Here, v is wavenumber; T and p are the atmospheric temperature and pressure respectively; f_{mol} volume mixing ratio (VMR) of a molecule; k_B – Boltzmann constant. In order to fit such a synthetic spectrum J_{mod} to the measured one J_{mes} , we convolved the model by the previously derived ILS using wavenumber calibrations. The aerosol opacity $\tau_a = -log(J_a)$ was fitted from the measured transmission as a continuum level J_a out of molecular absorption. The fitting procedure includes minimization of the function $\chi^2 = \sum_i A^2(v_i)$, which is a residuals sum over all considered spectral points *i* (pixels) divided by the transmission errors $\delta J(v_i)$, where $A(v_i) = (J_{mod}(v_i) - J_{mes}(v_i))/\delta J(v_i)$. An optimal search for the minimum was based on Levenberg-Marquardt algorithm (Marquardt, 1963) with the Jacobian matrix containing transmission derivatives $[\partial A/\partial X]$ on free parameters X, i.e. T, N_{CO2} , f_{H2O} or τ_a . Here, the clue contribution comes from the partial cross-section derivative on temperature $\partial \sigma / \partial T$. It differs for each molecule from line to line with a change of sign as seen in the Figure S1c. Thanks to that, an independent and simultaneous retrieval of temperature and molecular concentrations is possible.

In our multi iteration scheme, the first step is devoted to temperature and pressure retrievals with the rotational structure of CO_2 absorption bands. They are selected out from strong H_2O and HDO lines, e.g. in intervals 3733-3735 cm⁻¹, 3738-3743 cm⁻¹ of the order #223. Total number of spectral points exceeds 150 that makes the χ^2 -minimization confidential. Altitude profiles $T_1(z)$ and $N_{CO_2}(z)$ are directly retrieved, while the pressure is expressed on the first stage through the ideal gas law as $p_1(z) = {\binom{N_{CO2}(z)}{f_{CO2}(z)}} k_B T_1(z)$. The CO₂ VMR vertical distribution $f_{CO_2}(z)$ is taken from the Mars Climate Database (MCD) (Millour et al., 2018). Then, in order to meet the hydrostatic law, we constrain the pressure by a formula for the hydrostatic equilibrium, $p_{hyd}(z) = p_o exp \left[-\int_{z_o}^{z} \frac{g(h)M(h)}{k_BT_1(h)} dh \right]$, using the retrieved $T_1(z)$. Here, g is the gravity acceleration, *M* is the molecular weight (taken from MCD). A reference pressure $p_0 = p(z_0)$ is selected at an altitude z_o , where the retrieval uncertainties are minimal. In all derived profiles, this level lies between 30 and 60 km corresponding to the deepest, but not saturated, absorption lines (see Fig. 1b). Once the hydrostatic pressure profile is calculated, we keep it fixed when a new temperature and CO₂ density are retrieved on the next step. We repeat this procedure about 4-5 iterations reaching the profiles convergence. Values on an (i-1)th step are used as a priori for the ith one, while for the 1st stage we take the values from MCD. Analogous hydrostatic approach has been recently verified by other temperature and pressure retrievals from the ACS data (Alday et al., 2019, 2021; Fedorova et al., 2020). On the last stage, we apply the temperature and density profiles as a reference atmosphere for the H₂O mixing ratio retrieval (Fig. S2c) in the wavenumber interval from 3743 to 3753 cm⁻¹. We do not fit the HDO concentration separately and keep it as 5 times of the terrestrial SMOW value (Standard Mean Ocean Water), which is 1.56•10⁻⁴ for this isotope (Owen et al., 1988; Krasnopolsky, 2015). This assumption simplifies the fit, decreasing the number of free parameters, and it makes sense for the tiny detection of high altitude water. In parallel, each retrieval is accompanied by the derived aerosol opacity τ_{a} , which may characterize the dust activity during an occultation (see Figure 2a).

Examples of the derived temperature, H_2O number density and mixing ratio profiles are shown on Figure S2 in comparison with collocated and simultaneous occultations by the ACS near-IR channel (see in Fedorova et al. (2020)), as well as with the MCD predictions for particular occultations. The validation is demonstrated for three orbits that correspond to different atmospheric scenarios in MY34: before GDS ($L_S 164^\circ$), during GDS ($L_S 196.6^\circ$) and close to perihelion ($L_S 260.9^\circ$). Sensitivity of the NIR spectroscopy does not allow measuring temperature and densities above 100 km. In the MIR case, one can potentially retrieve the temperature altitude profiles spreading from 10-40 km (depending on aerosol opacity) up to 150 km in the considered spectral range (Fig. S2a). The upper altitude of the water detection is 20-25 km higher for the 2.66-2.67 µm band in MIR than 1.38 µm in NIR (Fig. S2b, S2c).

Text S₃. Estimation of uncertainties

The transmission errors δJ , included in the Jacobian matrix, determine uncertainties of the retrievals for each of free parameters. We estimated them for any observed altitude as a multiplication of the derivatives array by the transposed one, $\delta X = \sqrt{\left[\frac{\partial A}{\partial X}\right] \times \left[\frac{\partial A}{\partial X}\right]'}$, over the spectral points. The error bars δJ are composed of a statistical noise and some systematic uncertainties. We derived the statistical noise analyzing transmission fluctuations near 1 at altitudes above 200 km. It gave the standard deviation of about 5•10⁻⁴ for a signal recorded in the considered FPA row and in the center of the range, 3735-3750 cm⁻¹. The systematics comes

from the obtained calibrations (wavenumber, ILS) and from the hydrostatic pressure and temperature uncertainties when modeling transmission spectra on the final step of H_2O VMR retrievals. On average, it increases the δJ value up to 0.001-0.0015. The described uncertainties establish a detection limit of 5•10⁶-1•10⁷ cm⁻³ for H_2O number density above 100 km. When the water enrichment exceeds this order of magnitude, we observe a positive detection of the molecular abundance at high altitudes. In the opposite case, an upper limit is determined (see Fig. 1c).



Supplementing Figures:

Figure S1. Absorption spectroscopy in the echelle order #223 of ACS MIR. **a**: Example of measured transmission spectrum at a target altitude of 100 km (black axis) and the instrumental line shape (ILS) for both edges of the considered wavenumber range (blue axis). **b**: Absorption cross-sections for molecules CO₂ (black), H₂O (red), and HDO (green) calculated for T=150 K and p=10⁻⁵ mbar and convolved by ILS. **c**: Partial temperature derivatives of the considered cross-sections.



Figure S2. Validation of vertical profiles retrieved by ACS-MIR (red) with analogous profiles of ACS-NIR (from Fedorova et. al, 2020) (blue) and with predictions of MCD (Millour et al., 2018) (black): (a) temperature; (b) H₂O number density; (c) H₂O volume mixing ratio (VMR).



Figure S3. Latitudinal trend of H2O volume mixing ratio (VMR) in the perihelion seasons (Ls 250°–295°) of MY34 (blue) and MY35 (red). Each individual point is a weighted mean value revealed over altitude range of 100-120 km from a vertical profile of water VMR.



Figure S4. Altitude distribution of H_2O volume mixing ratio (VMR) during the GDS (MY 34) and the perihelion season (MY 34, 35). The dataset used includes all ACS MIR observations at GDS of MY34 ($L_s=195^{\circ}-220^{\circ}$) (brown) and around the perihelion point ($L_s=250^{\circ}-295^{\circ}$) (blue). Panels (a, b) for MY34 and (c, d) for MY35 from the Northern (a, c) and Southern (b, d) hemispheres. Light grey points are the NIR data during GDS of MY34 (Fedorova et al., 2020).