

# Revealing a High Water Abundance in the Upper Mesosphere of Mars with ACS onboard TGO

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## Abstract

We present the first water vapor profiles encompassing the upper mesosphere of Mars, 100–120 km, far exceeding the maximum altitudes where remote sensing has been able to observe water to date. Our results are based on solar occultation measurements by Atmospheric Chemistry Suite (ACS) onboard the ExoMars Trace Gas Orbiter (TGO). The observed wavelength range around 2.7  $\mu\text{m}$  possesses strong CO<sub>2</sub> and H<sub>2</sub>O absorption lines allowing sensitive temperature and density retrievals. We report a maximum H<sub>2</sub>O mixing ratio varying from 10 to 50 ppmv at 100–120 km during the global dust storm (GDS) of Martian Year (MY) 34 and around southern summer solstice of MY 34 and 35. During other seasons water remains persistently below  $\sim$ 2 ppmv. We claim that contributions of the MY34 GDS and perihelion periods into the projected hydrogen escape from Mars are nearly equivalent.



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## 29 **Plain Language Summary**

30 We report regular events of high abundances of the water vapor ( $\text{H}_2\text{O}$ ) in the upper atmosphere  
31 of Mars (100–120 km). So far, any water enrichment has not been revealed by remote sensing at  
32 such high altitudes. Higher than 80 km, solar light breaks water vapor molecules into H and O  
33 atoms, which may reach the exosphere and escape the planet. When Mars is closer to the Sun  
34 (the perihelion season), the atmosphere's circulation intensifies, causing increased dust activity  
35 with global dust storms (GDS), occurring every 3–4 Mars years. We observed during the second  
36 halves of Martian years 34 and 35 (2018–2020), including one GDS and two perihelion seasons.  
37 We report that the maximum water relative abundance reaches 10–50 parts per million in volume  
38 (ppmv) at 100–120 km during the GDS and every perihelion season. These high values indicate  
39 that the Martian atmosphere above 100 km regularly hosts large amounts of water, facilitating  
40 the long-term escape of water from the planet.

## 41 **1 Introduction**

42 The vertical distribution of water vapor ( $\text{H}_2\text{O}$ ) on Mars is an indicator of the intricate  
43 coupling of distinct phenomena: temperature variations, cloud formation, sublimation, turbulent  
44 and convective mixing, as well as general circulation and wave/eddy transport.  $\text{H}_2\text{O}$  has long  
45 been thought to remain confined below the hygropause, which is the level where the saturation  
46 condition is met and where water ice clouds may form, as occurs on Earth. The existence of this  
47 layer on Mars was established for the first time by ground-based microwave soundings of Clancy  
48 et al. (1996) with a saturation level between 10–20 km around the aphelion, i.e., Solar  
49 Longitudes ( $L_S$ )  $70^\circ$ , and 40–60 km around perihelion ( $L_S$   $250^\circ$ ). In parallel, Rodin et al. (1997)  
50 reported water vapor profiles retrieved from the solar occultations made by Auguste on Phobos-2  
51 in 1989. The existence of a hygropause at 30–35 km (with a mixing ratio of 3 ppm) in the  
52 northern spring ( $L_S=0^\circ$ – $20^\circ$ ) near the equator was subsequently claimed. The first climatology of  
53 water vapor profiles was derived from SPICAM-IR solar occultations on Mars Express (MEx)  
54 (Fedorova et al., 2009; 2018; 2021; Maltagliati et al., 2013), covering eight Martian years. The  
55 hygropause level was found to vary from 40 to 80 km depending on season, latitude, and dust  
56 events. Hygropause is also indirectly sensed in CRISM limb profiles of  $\text{O}_2(^1\Delta\text{g})$  emission, a  
57 confident indicator of  $\text{O}_3$ , from which water vapor mixing ratios were inferred by Clancy et al.  
58 (2017).

59 The observation of large amounts of water vapor in and above the middle atmosphere  
60 ( $>40$  km, Maltagliati et al., 2013) was then complemented by the discovery of short-term decline

61 of the hydrogen corona brightness over several weeks (Chaffin et al., 2014; Clarke et al., 2014).  
62 This variability exposed a new paradigm in our perception of how water escapes from Mars  
63 (Chaffin et al., 2017). So far, water escape was thought to be controlled by a slow conversion  
64 process involving H<sub>2</sub>, formed from the catalytic recombination of carbon dioxide with odd  
65 hydrogen (McElroy and Donahue, 1972; Krasnopolsky, 2002). The non-condensable H<sub>2</sub> can  
66 overcome the hygropause and reach the mesosphere (80–120 km), while transported by turbulent  
67 mixing or circulation. There, it can dissociate and release H atoms that will escape the planet  
68 once above the exobase.

69 Observations have revealed that water vapor transport from the troposphere to the lower  
70 mesosphere of Mars occurs during the dusty season and is enhanced at times of major dust  
71 storms. In particular, a significant H<sub>2</sub>O enhancement in the middle atmosphere was observed  
72 during the global dust storm (GDS) in 2007 (MY28) with a rise of the hygropause altitude to >60  
73 km (Fedorova et al., 2018; Heavens et al., 2018; 2019). Sensitive solar occultation measurements  
74 by NOMAD and ACS NIR instruments onboard the ExoMars Trace Gas Orbiter (TGO) have  
75 showed that water vapor reached 80–100 km (Aoki et al., 2019; Fedorova et al., 2020) during  
76 two storms in 2018 and 2019 (a global one at L<sub>S</sub> 190°–220° and a regional one at L<sub>S</sub> 330° in  
77 MY34; Montabone et al., 2020). Fedorova et al. (2020) revealed the water supersaturation at 70–  
78 90 km even in the presence of H<sub>2</sub>O ice clouds not only during the GDS but also near the  
79 Southern summer solstice (L<sub>S</sub>~270°) when water reached 90-100 km as well. Altogether, these  
80 studies promote a new mechanism for controlling H escape through direct delivery at above 80  
81 km and further photodissociation of H<sub>2</sub>O molecules (Chaffin et al., 2017; Krasnopolsky et al.,  
82 2019). General circulation models predict an upward water flux into the thermosphere (>120 km)  
83 during the GDS and perihelion periods (Shaposhnikov et al., 2019; Neary et al., 2020; Rossi et  
84 al. 2021).

85 The discussion regarding a relative contribution of perihelion or GDS to the mesospheric  
86 water enrichment was recently stimulated by SPICAM/ME<sub>x</sub> long-term observations covering  
87 Martian Years 28 through 35. Here, Fedorova et al. (2021) claimed an annual rise of water  
88 abundance up to ~90 km in perihelion, which is compatible with GDS enhancements. The new  
89 ACS/TGO dataset confirms those conclusions for altitudes below 100 km in MY34–MY35  
90 (Fedorova et al. 2020, Alday et al., 2021). In parallel, during the perihelion season, the D/H ratio  
91 in water decreases with altitude from 4–6 times SMOW (Standard Mean Ocean Water) in the  
92 lower atmosphere to 2–3 times in the mesosphere (50-70 km) as measured by ACS MIR (Alday  
93 et al., 2021) and NOMAD (Villanueva et al., 2021) spectrometers. Alday et al. (2021) show that  
94 ultraviolet H<sub>2</sub>O photolysis dominates the production of H relative to D atoms in the upper  
95 atmosphere.

96 From above, ion chemistry in the thermosphere has been characterized by the NGIMS  
97 mass-spectrometer on MAVEN (Benna et al., 2015) and interpreted by the ionospheric model of  
98 Fox et al. (2015). Using NGIMS data, Stone et al. (2020) measured H<sub>2</sub>O ion concentrations  
99 around ~150 km for the 2014–2018 period (MY32–MY34). With the help of the model by Fox et  
100 al. (2015), Stone et al. (2020) found the relative abundance of water at this altitude on the  
101 dayside varying seasonally on average from 2 to 5 ppm. Several enhanced dusty episodes disrupt  
102 this seasonal signal: 3–9 ppm during the regional storm of MY32, 10–20 ppm during the storm  
103 of MY33, and up to 60 ppm in the GDS of MY34. Stone et al. (2020) concluded that water  
104 transport into the ionosphere and its destruction are the main mechanisms in the overall hydrogen  
105 escape from Mars.

106 We used the data of the middle infrared spectrometer of the Atmospheric Chemistry Suite  
107 (ACS MIR) onboard the ExoMars TGO, which measures water vapor VMR and atmospheric  
108 density in a wide range of altitudes, from the troposphere to the lower thermosphere, using the  
109 strong absorption bands of H<sub>2</sub>O and CO<sub>2</sub> around 2.66–2.70 μm. The high spectral resolution and  
110 the good signal-to-noise ratio of ACS MIR allow the measurements of water profiles up to 120  
111 km, inaccessible altitudes for the ACS NIR and SPICAM measurements, sensing the 1.38 μm  
112 absorption band (Fedorova et al., 2020; 2021). The strong H<sub>2</sub>O absorption around 2.6 μm is also  
113 used by NOMAD, yielding water profiles up to ~90 km (Aoki et al., 2019; Villanueva et al.,  
114 2021).

115 Here we report the first water vapor abundance measurements in the upper mesosphere  
116 (up to 120 km) of Mars. The goal of our paper is to compare the mesospheric water behavior  
117 between the second halves of MY34 and MY35 when the high H<sub>2</sub>O content is observed. We aim  
118 to clarify the principal mechanism of H<sub>2</sub>O delivery to the upper mesosphere: it is sporadic dust  
119 events, or the result of seasonal variability in the Martian circulation that peaks each year, around  
120 Southern summer solstice. For that, we analyze seasonal and latitudinal variations of H<sub>2</sub>O VMR  
121 vertical profiles retrieved from the ACS MIR solar occultation experiment.

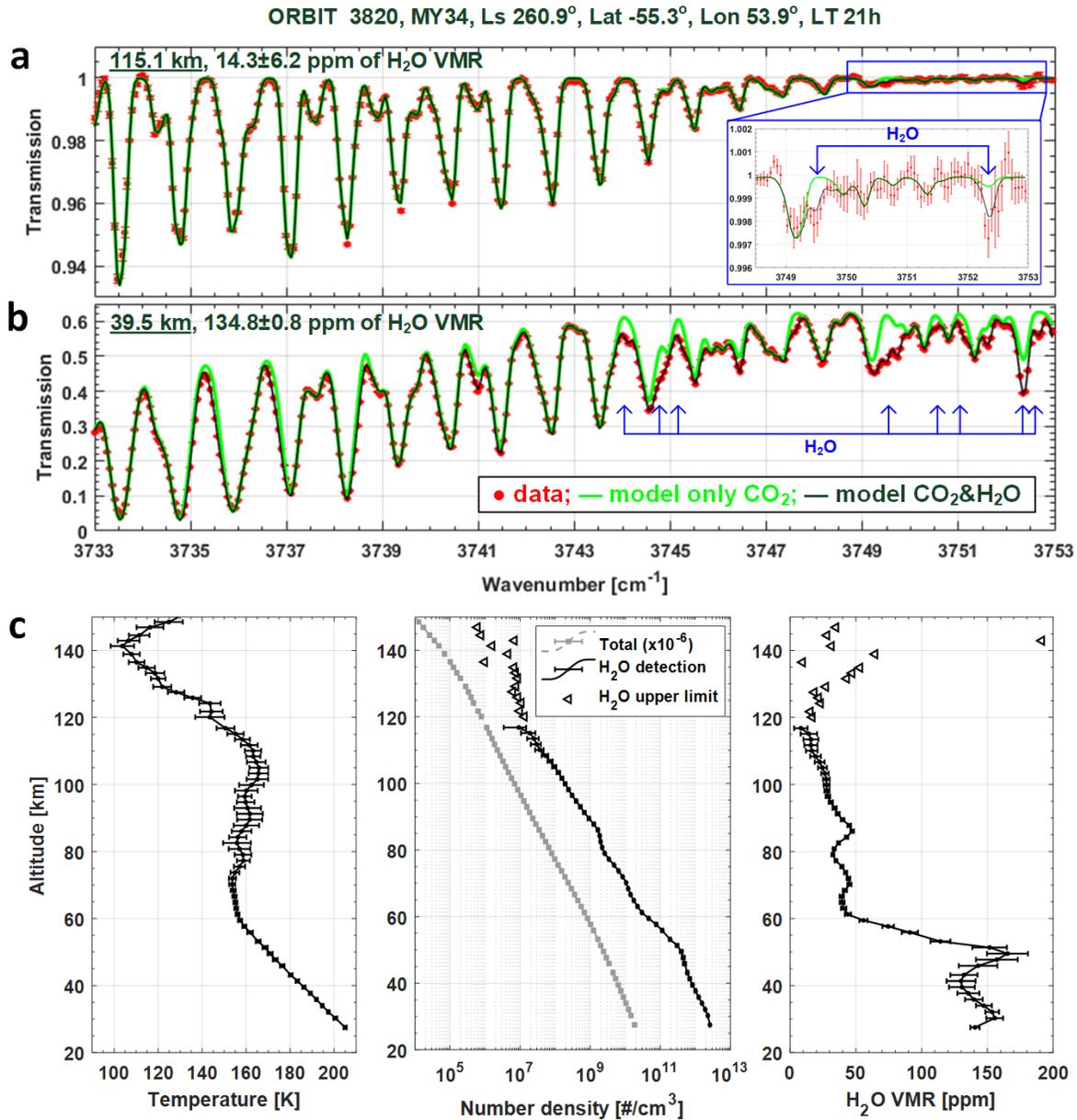
## 122 **2 Measurements and dataset overview**

### 123 **2.1 ACS MIR spectroscopy and retrievals**

124 ACS MIR, a solar occultation cross-dispersion echelle spectrometer, records spectra from  
125 a set of adjacent diffraction orders (from 10 to 20 per occultation) projected onto a 2D detector  
126 array (Korablev et al., 2018). To retrieve high altitude water vapor abundances together with the  
127 atmospheric temperature and pressure, we use MIR spectra from the diffraction order #223. They  
128 cover a narrow wavelength interval of 2.66–2.68 μm (3733–3753 cm<sup>-1</sup>), including a part of the  
129 2.7-μm CO<sub>2</sub> absorption band and a few strong H<sub>2</sub>O lines near 2.66 μm (Fig. 1a, 1b, Fig. S1). The  
130 instrument's spectral resolution is ~0.15 cm<sup>-1</sup>, while the signal-to-noise ratio ranges from 2,000  
131 to 4,000, which provides high sensitivity for detections in the upper atmosphere where  
132 atmospheric constituent densities are low. Temperature (Fig. 1c) is retrieved by fitting a  
133 synthetic model to the CO<sub>2</sub> rotational band taking advantage of its temperature dependence as  
134 seen in Fig. S1 of the Supplementary Material (SM). This procedure was applied iteratively, with  
135 the pressure calculated from the retrieved temperature profile under the assumption of  
136 hydrostatic equilibrium. The temperature measurements were then validated against those made  
137 by MIR near the 2.6 μm CO<sub>2</sub> band (Alday et al., 2019) and by ACS NIR around the 1.58 μm  
138 band (Fedorova et al., 2020). As a result, one occultation session allows us to simultaneously  
139 retrieve profiles of pressure and temperature (from CO<sub>2</sub> absorption bands) and the H<sub>2</sub>O number  
140 density (Fig. 2c). The water abundance can then be expressed relative to the total atmospheric  
141 density, that is, in VMR (in ppmv). Specific details of the algorithms pertaining to this work can  
142 be found in SM.

143 The dataset analyzed here consists of a series of transmission spectra obtained during a  
144 solar occultation while the line of sight of the instrument progressively penetrates from the upper  
145 into deeper layers of the atmosphere, or vice versa (see examples in Fig. 1a, 1b). The  
146 transmission is determined as the solar spectrum ratio measured through the atmosphere to the  
147 reference one, taken from the data above a tangent height of 200 km. This altitude level is  
148 negligibly attenuated by the atmosphere even within the very strong CO<sub>2</sub> band system at 2.7 μm.

149 The typical integration time is 2 seconds, which provides an altitude resolution ranging from 0.5  
 150 to 2.5 km, depending on the occultation duration. It gives sufficiently fine vertical sampling for  
 151 an atmosphere whose scale height ranges from 5 to 10 km depending on temperature. The  
 152 instrument field of view projected at the limb is around 1-3 km in altitude equivalent.



153

154 **Figure 1.** Example ACS MIR spectra and profiles of the retrieved quantities. Measured  
 155 transmission spectra (red) at tangent altitudes of 115.1 km (**a**) and 39.5 km (**b**) are compared  
 156 with the best-fit models, including both CO<sub>2</sub> and H<sub>2</sub>O absorptions (black), and only CO<sub>2</sub>  
 157 absorption (green). Blue arrows indicate water absorption lines. Zoom in (**a**) shows a part with  
 158 the strongest H<sub>2</sub>O absorption detected at 115 km. (**c**) Retrieved vertical profiles of temperature  
 159 (left), number densities (center), and H<sub>2</sub>O volume mixing ratio (VMR) (right). The atmospheric

160 density (grey squares) is scaled by the factor of  $10^{-6}$ . Black triangles mark H<sub>2</sub>O upper limits (see  
161 SM for the description of uncertainties).

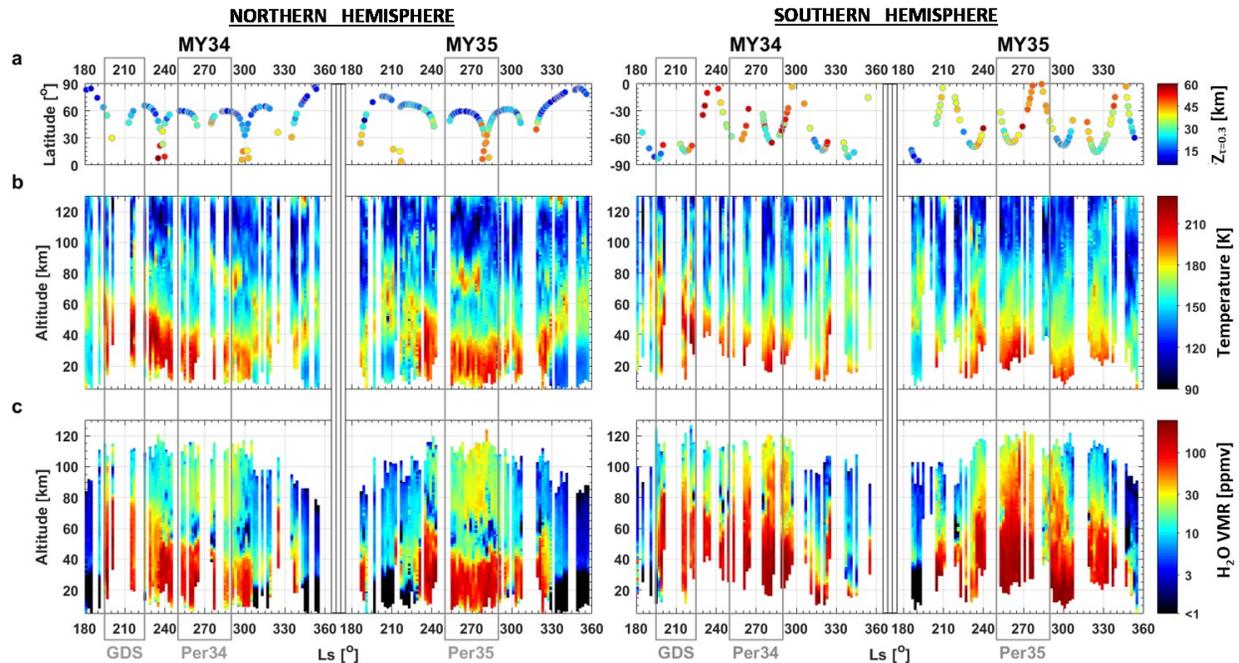
## 162 2.2 Data selection

163 Our measurements of the mesospheric water focus on the second halves of MY34 and  
164 MY35, which correspond to ACS MIR observations from May 2018 to March 2019 and from  
165 April 2020 to January 2021. The selected dataset comprises 187 occultation sessions in the  
166 Northern Hemisphere and 156 sessions in the Southern Hemisphere, encompassing seasonal  
167 periods from L<sub>S</sub> 180° to 355° in MY34 and from L<sub>S</sub> 185° to 356° in MY35 (Fig. 2a). Figure 2a  
168 shows the latitude coverage with the corresponding aerosol activity, which was defined for each  
169 occultation at the altitude level where the slant opacity equals 0.3 (~0.75 of the atmospheric  
170 transmittance in the continuum). Measurements in the Northern Hemisphere occurred mostly in  
171 the high latitude range, between 40°N and 70°N. In the South, the perihelion observations (L<sub>S</sub>  
172 270°) were made in mid-latitudes, while the rest of occultations occurred close to the polar  
173 region (60°S–90°S). Only a few sessions were localized nearby the equator: at L<sub>S</sub>~240° and  
174 L<sub>S</sub>~300° of MY34 and at L<sub>S</sub>~210° and L<sub>S</sub>~280° of MY35. These observations are accompanied  
175 by a higher aerosol loading than for high latitude and polar regions (Fig. 2a).

## 176 3 Seasonal variation of altitude profiles

177 Observations in the second halves of MY34 and MY35 uncover events, which drastically  
178 perturbed the temperature and water vapor vertical distributions. The peculiar pattern to compare  
179 with is the MY34 GDS and perihelion periods in MY34 versus MY35, which had no GDS but a  
180 regional dust activity in its second half. The seasonal variation of the processed altitude profiles  
181 is presented in Figure 2(b, c). We binned the profiles into intervals of 2° in solar longitude and 2  
182 km in altitude. Depending on the L<sub>S</sub> and altitude sampling, the value in each bin is calculated as  
183 the weighted mean of one to five individual points. We excluded all points with 1-sigma  
184 uncertainties exceeding 20 K in temperature and 100% of the H<sub>2</sub>O mixing ratio. The second  
185 rejection criterion corresponds to the detection limit (~10<sup>7</sup> cm<sup>-3</sup>) of water number density (see in  
186 Fig. 1c) that defines the seasonal variations of the uppermost detectable points in Figure 2c.

187 We observe seasonal temperature (Fig. 2b) and H<sub>2</sub>O (Fig. 2c) peaks in the middle  
188 atmosphere (40–80 km) during the GDS of MY34, L<sub>S</sub> 190°–220°, and an additional smaller peak  
189 at L<sub>S</sub> 320°–330°, corresponding to a regional storm. Moreover, the rise of water vapor to higher  
190 altitudes, up to the mesopause at 110–130 km where temperature encounters a minimum, is  
191 observed during the two perihelion intervals (L<sub>S</sub> 250°–290°) of MY34 and MY35. Here, the  
192 Southern summer (Fig. 2c, right panel) is accompanied by a more humid mesosphere (40–60  
193 ppm of H<sub>2</sub>O) than the Northern Winter (Fig. 2c, left panel) where the mean mesospheric water  
194 reaches 20–30 ppmv on average between 80 and 120 km. In contrast, out of the perihelion peak  
195 or dust events, i.e. for the selected data at the beginning of the MY34 GDS and at the very end of  
196 MY 34, 35, water content above 80 km never exceeds 2-3 ppmv.



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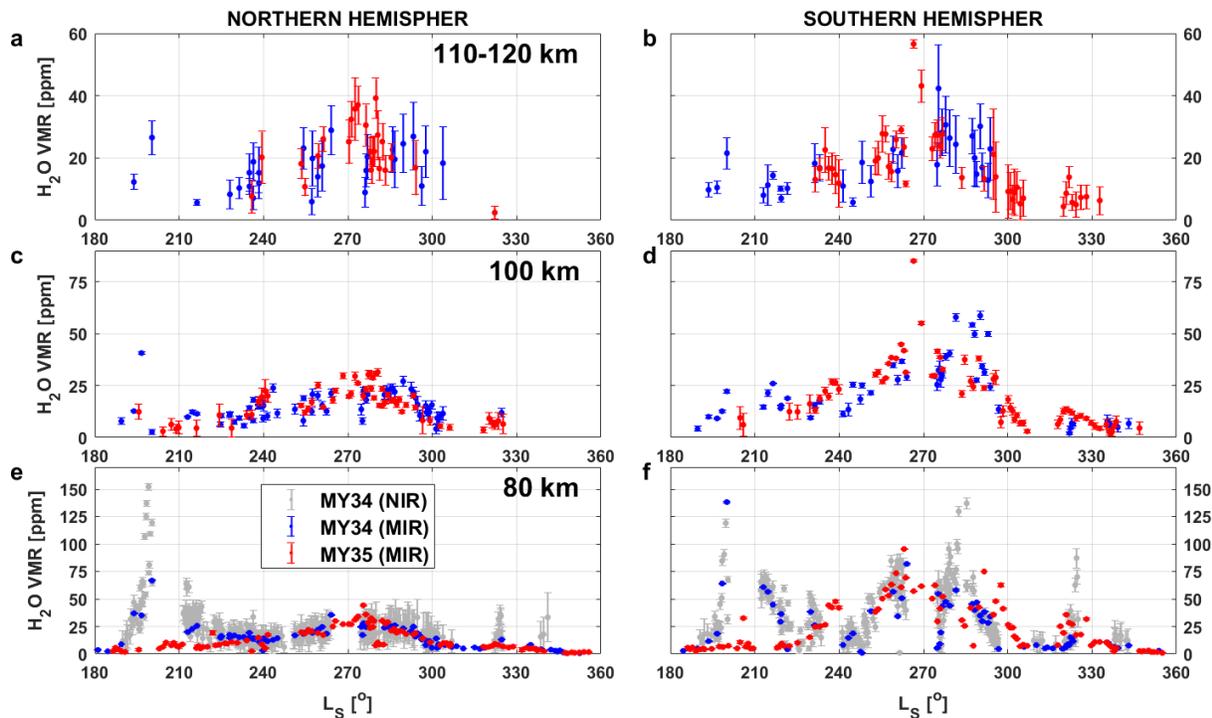
198 **Figure 2.** Seasonal map of atmospheric temperature and H<sub>2</sub>O mixing ratio during the second half  
 199 of MY34 and MY35. The data are plotted in function of L<sub>S</sub> and altitude for the Northern (left)  
 200 and the Southern (right) hemispheres. (a) Latitudinal distribution of the ACS-MIR solar  
 201 occultations, depending on an altitude level where the aerosol slant opacity ( $\tau$ ) equals 0.3. (b)  
 202 Temperature. (c) Volume mixing ratio (VMR) of water vapor. Grey frames outline time intervals  
 203 of the global dust storm (GDS) in MY34 and the two perihelions in MY34 (Per34) and MY35  
 204 (Per35).

#### 205 4 H<sub>2</sub>O variations around perihelion

206 To quantify seasonal trends of water content in the mesosphere, we selected three altitude  
 207 layers corresponding to 80 km, 100 km, and 110–120 km. The first layer, which corresponds to  
 208 the middle mesosphere, is accessible in all profiles (Fig. 2c) when the vapor concentration  
 209 exceeds the detection limit of  $\sim 10^7 \text{ cm}^{-3}$ , even in low water loading periods. Water at 100–120  
 210 km shows up only in stormy periods and around perihelion (Fig. 2c).

211 Observed variations during perihelion for the three selected levels are presented in Figure  
 212 3 for both Martian Years. The number of MIR observations at the considered spectral range is  
 213 low during the dust storm activity of MY34. Nevertheless, a comparison with MY35 reveals  
 214 significant increases of H<sub>2</sub>O mixing ratios during the GDS: by a factor of 6–8 at 80 km (Fig. 3e,  
 215 3f) and by a factor of 3–5 at 100–120 km from L<sub>S</sub> 190° to 220° (Fig. 3a–3d). Increases at L<sub>S</sub>  
 216 320°–330° follow annually repeatable dust storm activity at this season, although injecting far  
 217 less water into the mesosphere than the GDS in MY34. Around Mars perihelion (L<sub>S</sub>=240°–300°)  
 218 water behaves almost identically between MY34 and MY35. For both Martian Years, the  
 219 maximum H<sub>2</sub>O mixing ratio was observed near the Southern summer solstice (L<sub>S</sub>~270°),  
 220 reaching values of 40–80 ppm at 80 km, 30–60 ppm at 100 km, and 20–50 ppm at 110–120 km.  
 221 In the Northern winter solstice, it varied from 20 to 40 ppm at all levels, 80–120 km. There are  
 222 groups of points out of general behaviour, i.e. at L<sub>S</sub>=270°–280° in Fig. 3, that results from  
 223 latitudinal variations of the water content (see Fig. S3 in SM). We compare our results at 80 km

224 with the corresponding ACS NIR dataset derived from the MY34 profiles of Fedorova et al.  
 225 (2020, grey points in Fig. 3e, 3f). The NIR dataset is five times denser than used in the present  
 226 work, and it observed the H<sub>2</sub>O seasonal variations in greater detail, especially during the dust  
 227 events of MY34.



228  
 229 **Figure 3.** Seasonal trends of H<sub>2</sub>O volume mixing ratio (VMR) at three altitude levels: 80, 100,  
 230 110–120 km. The season is the second halves of MY34 (in blue) and of MY35 (in red). Each  
 231 point corresponds to an individual vertical profile: weighted mean value obtained in between  
 232 110–120 km (a, b), and interpolated value for the levels of 100 km (c, d) and 80 km (e, f). Data at  
 233 80 km in grey (e, f) are taken from the ACS NIR profiles (Fedorova et al., 2020).

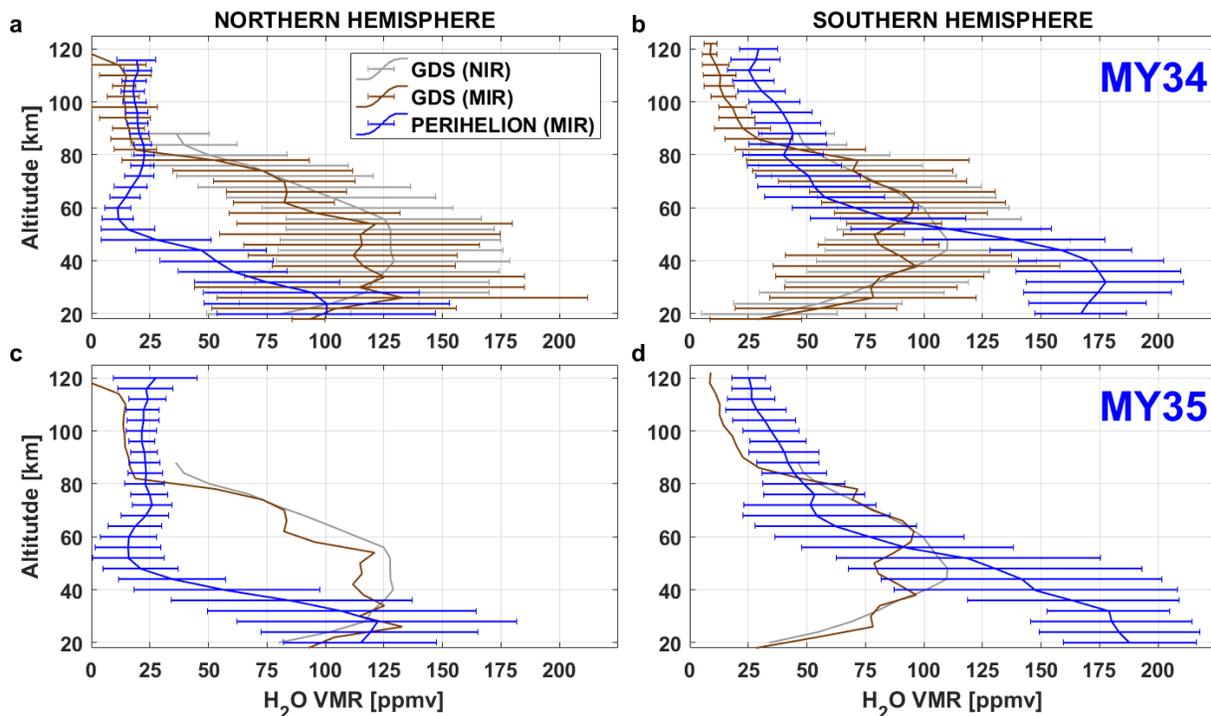
#### 234 4 Discussion and Conclusions

235 For the first time, we report observations of H<sub>2</sub>O abundances in a previously unexplored  
 236 altitude range (from 100 to 120 km). There we find 10–30 ppm of water vapor during the MY34  
 237 GDS and 20–50 ppm around Mars perihelions ( $L_s=250^\circ$ – $290^\circ$ ) of MY34 and MY35 in both  
 238 hemispheres. Our GDS retrievals at 100–120 km are of the same order of magnitude as MAVEN  
 239 NGIMS results reported by Stone et al. (2020) at  $\sim 150$  km. Surprisingly, NGIMS water  
 240 abundances reveal a 2014–2018 mission-wise maximum of H<sub>2</sub>O at 150 km, only during the  
 241 MY34 GDS, whereas we repeatedly observe the annual maximum around the Southern summer  
 242 solstice both in MY34 and MY35.

243 NGIMS measures  $[H_2O^+]$  ions, from which neutral H<sub>2</sub>O abundances at 150 km were  
 244 derived on the basis of 1D photochemical modeling. The model was adjusted to reproduce the  
 245 H<sub>2</sub>O VMR at 150 km inferred from the  $[H_2O^+]$  ions measured under two scenarios: low water,  
 246 corresponding to 2 ppm prescribed at 80 km in a non-GDS case; and high water of 40 ppm in a  
 247 GDS case (Stone et al., 2020; as corrected in March 2021). Notably, all the solar occultation  
 248 observations performed by TGO and MEX to date (Fedorova et al., 2018, 2020, 2021; Aoki et

249 al., 2019; Villanueva et al., 2021), including the present dataset (Fig. 3e, 3f), report even higher  
 250 water vapor VMRs at 80 km, of 50–80 ppm during the GDS. Stone et al. (2020) indicate a  
 251 systematic uncertainty of 69% on their neutral H<sub>2</sub>O inference, which is consistent with observed  
 252 MAVEN, TGO, and MEX values within such error bars.

253 It is important to consider how ACS's high altitude water vapor abundances combine  
 254 with photolysis since this process has been hypothesized to be essential, if not the dominant,  
 255 source for the H atoms observed in the exosphere (Chaffin et al., 2017). The conclusion of Stone  
 256 et al. (2020) argues for the GDS's predominance and related ion chemistry in the H atoms'  
 257 production. Our observations suggest that while the GDS period corresponds to the maximum of  
 258 water abundance at 80 km, H<sub>2</sub>O at 120 km peaks only later, at the Southern summer solstice,  
 259 when it is twice as large as during the GDS. This enhanced solstice maximum suggests that  
 260 relative water abundance declines more rapidly above 80 km during the GDS than after, during  
 261 perihelion (Fig. 4).



262  
 263 **Figure 4.** Altitude profiles of average H<sub>2</sub>O volume mixing ratio (VMR) during the GDS (MY  
 264 34) and the perihelion season (MY 34, 35). The dataset used includes all ACS MIR observations  
 265 highlighted in Fig. 2: GDS of MY34 ( $L_S=195^\circ-220^\circ$ ) (brown); the bin around the perihelion  
 266 point ( $L_S=250^\circ-295^\circ$ ) (blue). Panels (a, b) for MY34 and (c, d) for MY35 from the Northern (a,  
 267 c) and Southern (b, d) hemispheres. Profiles are presented with 1-sigma dispersion over 4-km  
 268 altitude bins. Light grey curves are the averages of the NIR data during GDS of MY34  
 269 (Fedorova et al., 2020). The GDS curves are also indicated for MY35 to facilitate comparison.  
 270 For the data points used in the averaging see Fig. S4 of SM.

271 In Figure 4, we combined altitude profiles from GDS-only ( $L_S 195^\circ-220^\circ$ ) and perihelion  
 272 ( $L_S 250^\circ-295^\circ$ ) intervals to compare averaged vertical trends between them. Here, we also see a  
 273 coincidence between MIR and NIR GDS profiles in frames of dispersions, which reflect high  
 274 variability of the observed GDS points (Fig. S4 of SM). The considered H<sub>2</sub>O distributions allow

275 estimating an integral escape flux of the atomic hydrogen in each case. For that, we applied the  
276 model of Chaffin et al. (2017), which predicts the atmospheric escape rate depending on the  
277 water injection into different altitudes (see Figure 3 of their paper). Our rough calculations show  
278 that the H escape flux is about  $\sim 5 \times 10^9 \text{ cm}^{-1} \text{ s}^{-1}$  during the considered intervals of MY34's GDS  
279 and Southern summer solstices of MY 34 and 35. Thus, we claim nearly equivalent contributions  
280 from a single GDS and the perihelion period into the hydrogen escape by the high water  
281 enrichment in the middle/upper atmosphere. Fedorova et al. (2021) come to a similar conclusion  
282 based on SPICAM/MEx water profiles up to 80 km. A GDS occurs every 3–4 martian years on  
283 average (Zurek and Martin, 1993; Wang and Richardson, 2015), making the yearly perihelion  
284 contribution to the hydrogen escape reasonable. The water enhancements are tied to the  
285 circulation regime (Clancy et al., 1996; Richardson and Wilson, 2002; Montmessin et al., 2005).  
286 More measurements and modeling would be needed to decide whether the southern summer  
287 solstitial transport, currently near perihelion, or the GDS equinoctial circulation prevailed during  
288 the history of Mars.

289 Overall, our results cannot be easily reconciled with water values (up to 60 ppm at 150  
290 km) inferred from NGIMS ion measurements, which suggested that the thermosphere hosted  
291 much more water during the GDS than during the rest of the year. However, we note that the  
292 only time when Stone et al. (2020) reported measurements around perihelion concerned the L<sub>s</sub>  
293 interval between 240° and 265° of MY33 (Figure 4 of Stone et al., 2020). It showed the same  
294 rough trend as during the onset of the MY34 GDS, with values far exceeding those reported for  
295 the regional dust storm of MY33, still a factor of 3 smaller than during the GDS.

296 Our results remain in line with the conclusion of Fedorova et al. (2020, 2021) that the  
297 perihelion season is the primary conveyor of water to high altitudes on a long-term basis. The  
298 high values above 100 km fill the gap between the water observed below 100 km and water ions  
299 measured by NGIMS at 150 km. Both measurements bring unique constraints in our attempt to  
300 understand how the water in the lower atmosphere connects with the escaping hydrogen in the  
301 exosphere, an essential step before confidently extrapolating the water escape back in time.

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307 available from ESA's Planetary Science Archive at  
308 <https://archives.esac.esa.int/psa/#!Table%20View/ACS=instrument>. The retrieved data with  
309 altitude profiles of H<sub>2</sub>O VMR for the considered seasons are available at  
310 [https://data.mendeley.com/datasets/995y7ymdgm/draft?a=daa72362-898d-4c86-8a13-](https://data.mendeley.com/datasets/995y7ymdgm/draft?a=daa72362-898d-4c86-8a13-023b4b59134c)  
311 [023b4b59134c](https://data.mendeley.com/datasets/995y7ymdgm/draft?a=daa72362-898d-4c86-8a13-023b4b59134c).

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*Geophysical Research Letters*

Supporting Information for

**Revealing a High Water Abundance in the Upper Mesosphere of Mars with ACS onboard TGO**

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## Introduction

The present Supporting Information describes details of the ACS MIR instrument and the concept of altitude profiles retrievals from the measured transmission spectra. The retrieval algorithm is supplemented by Figure S1, which demonstrates instrumental and temperature peculiarities of the CO<sub>2</sub> and H<sub>2</sub>O absorption spectroscopy around 2.7 μm band (echelle order #223 of ACS MIR). Figure S2 shows validation of the derived temperature, H<sub>2</sub>O number density and mixing ratio profiles for different seasonal scenarios in MY<sub>34</sub>. Latitudinal variation of the high altitude water around perihelion is presented in Figure S3 for both MY<sub>34</sub> and MY<sub>35</sub>. Figure S4 presents vertical distribution of all data points for H<sub>2</sub>O VMR in the GDS and perihelion seasons. Uncertainties of the obtained data are described separately in section Text S3.

### Text S1. Instrument description

The middle infrared channel of the Atmospheric Chemistry Suite is a cross-dispersion echelle spectrometer dedicated to solar occultation measurements in the 2.3–4.3  $\mu\text{m}$  wavelength range (Korablev et al., 2018). Each occultation session is devoted to one of ten angular positions of the MIR secondary grating that disperses about 10–20 echelle orders spatially separated and recorded simultaneously at the focal plane array (FPA) by 640x512 pixels. The spectral range of our interest, 2.66–2.68  $\mu\text{m}$ , lies within the diffraction order #223 at the grating position #4. One order covers a spectral interval of about 30  $\text{cm}^{-1}$  (~25 nm); a spectrum is dispersed along 640 FPA elements with the sampling of 0.05  $\text{cm}^{-1}$  and the resolving power  $\lambda/\delta\lambda$  reaching ~25000 ( $\delta\lambda \sim 0.15 \text{ cm}^{-1}$ ). In the occultation field of view (FOV), the instrumental rectangular slit cuts a part from the solar disk, so that one order occupies a stripe with about 20 FPA rows. In such a manner, the consequences of orders are located on the matrix as stripe-by-stripe, one above the other, with some dark rows between them. For our analysis, we selected a central row from the stripe correspondent to order #223 that provides the maximal signal-to-noise ratio. The calibration procedure for the analyzed transmissions encompassed a pixel-to-wavenumber assignment and a determination of the instrument line shape (ILS). The procedure was applied for every occultation and in various spectral intervals separately since we observed variability of the light spreading over the FPA from orbit to orbit. Using strong  $\text{CO}_2$  and  $\text{H}_2\text{O}$  absorption lines in order #223, we derived an optimal wavenumber assignment by a parabolic law and an asymmetric ILS including a superposition of two Gaussian functions (Fig. S1a). Previously, Alday et al. (2019) implemented an analogous approach for MIR spectra but in the grating position #5.

### Text S2. Retrieval concept

A scheme for the water VMR retrieval consists of several iterations with a fitting of a forwardly modeled transmission spectrum to a measured one at each observed altitude. The forward modeling includes contributions from  $\text{H}_2\text{O}$ , HDO, and  $\text{CO}_2$  molecular absorption cross-sections that are abundant in the 2.66–2.68  $\mu\text{m}$  spectral range, i.e. 3730–3755  $\text{cm}^{-1}$  (Fig. S1b). They were calculated using line-by-line modeling for specified temperatures and pressures on a basis of the HITRAN2016 database (Gordon et al., 2017) taking into account  $\text{H}_2\text{O}$  and HDO lines broadening in the  $\text{CO}_2$ -rich atmosphere (Gamache et al., 2016; Malathy Devi et al., 2017) and self-broadening in the case of  $\text{CO}_2$ . A transmission spectrum  $J(\nu, T, p)$  is defined by the cross-sections  $\sigma(\nu, T, p)$ , gaseous concentrations  $N_{mol} = f_{mol} \cdot p / k_B T$ , and aerosol slant opacity  $\tau_a$ . Here,  $\nu$  is wavenumber;  $T$  and  $p$  are the atmospheric temperature and pressure respectively;  $f_{mol}$  – volume mixing ratio (VMR) of a molecule;  $k_B$  – Boltzmann constant. In order to fit such a synthetic spectrum  $J_{mod}$  to the measured one  $J_{mes}$ , we convolved the model by the previously derived ILS using wavenumber calibrations. The aerosol opacity  $\tau_a = -\log(J_a)$  was fitted from the measured transmission as a continuum level  $J_a$  out of molecular absorption. The fitting procedure includes minimization of the function  $\chi^2 = \sum_i A^2(\nu_i)$ , which is a residuals sum over all considered spectral points  $i$  (pixels) divided by the transmission errors  $\delta J(\nu_i)$ , where  $A(\nu_i) = (J_{mod}(\nu_i) - J_{mes}(\nu_i)) / \delta J(\nu_i)$ . An optimal search for the minimum was based on Levenberg-Marquardt algorithm (Marquardt, 1963) with the Jacobian matrix containing transmission derivatives  $[\partial A / \partial X]$  on free parameters  $X$ , i.e.  $T$ ,  $N_{\text{CO}_2}$ ,  $f_{\text{H}_2\text{O}}$  or  $\tau_a$ . Here, the clue contribution comes from the partial cross-section derivative on temperature  $\partial\sigma/\partial T$ . It differs for each molecule from line to line with a change of sign as seen in the Figure S1c. Thanks to that, an independent and simultaneous retrieval of temperature and molecular concentrations is possible.

In our multi iteration scheme, the first step is devoted to temperature and pressure retrievals with the rotational structure of CO<sub>2</sub> absorption bands. They are selected out from strong H<sub>2</sub>O and HDO lines, e.g. in intervals 3733-3735 cm<sup>-1</sup>, 3738-3743 cm<sup>-1</sup> of the order #223. Total number of spectral points exceeds 150 that makes the  $\chi^2$ -minimization confidential. Altitude profiles  $T_1(z)$  and  $N_{CO_2}(z)$  are directly retrieved, while the pressure is expressed on the first stage through the ideal gas law as  $p_1(z) = \left( \frac{N_{CO_2}(z)}{f_{CO_2}(z)} \right) k_B T_1(z)$ . The CO<sub>2</sub> VMR vertical distribution  $f_{CO_2}(z)$  is taken from the Mars Climate Database (MCD) (Millour et al., 2018). Then, in order to meet the hydrostatic law, we constrain the pressure by a formula for the hydrostatic equilibrium,  $p_{hyd}(z) = p_o \exp \left[ - \int_{z_o}^z \frac{g(h)M(h)}{k_B T_1(h)} dh \right]$ , using the retrieved  $T_1(z)$ . Here,  $g$  is the gravity acceleration,  $M$  is the molecular weight (taken from MCD). A reference pressure  $p_o = p(z_o)$  is selected at an altitude  $z_o$ , where the retrieval uncertainties are minimal. In all derived profiles, this level lies between 30 and 60 km corresponding to the deepest, but not saturated, absorption lines (see Fig. 1b). Once the hydrostatic pressure profile is calculated, we keep it fixed when a new temperature and CO<sub>2</sub> density are retrieved on the next step. We repeat this procedure about 4-5 iterations reaching the profiles convergence. Values on an (i-1)<sup>th</sup> step are used as a priori for the i<sup>th</sup> one, while for the 1<sup>st</sup> stage we take the values from MCD. Analogous hydrostatic approach has been recently verified by other temperature and pressure retrievals from the ACS data (Alday et al., 2019, 2021; Fedorova et al., 2020). On the last stage, we apply the temperature and density profiles as a reference atmosphere for the H<sub>2</sub>O mixing ratio retrieval (Fig. S2c) in the wavenumber interval from 3743 to 3753 cm<sup>-1</sup>. We do not fit the HDO concentration separately and keep it as 5 times of the terrestrial SMOW value (Standard Mean Ocean Water), which is  $1.56 \cdot 10^{-4}$  for this isotope (Owen et al., 1988; Krasnopolsky, 2015). This assumption simplifies the fit, decreasing the number of free parameters, and it makes sense for the tiny detection of high altitude water. In parallel, each retrieval is accompanied by the derived aerosol opacity  $\tau_a$ , which may characterize the dust activity during an occultation (see Figure 2a).

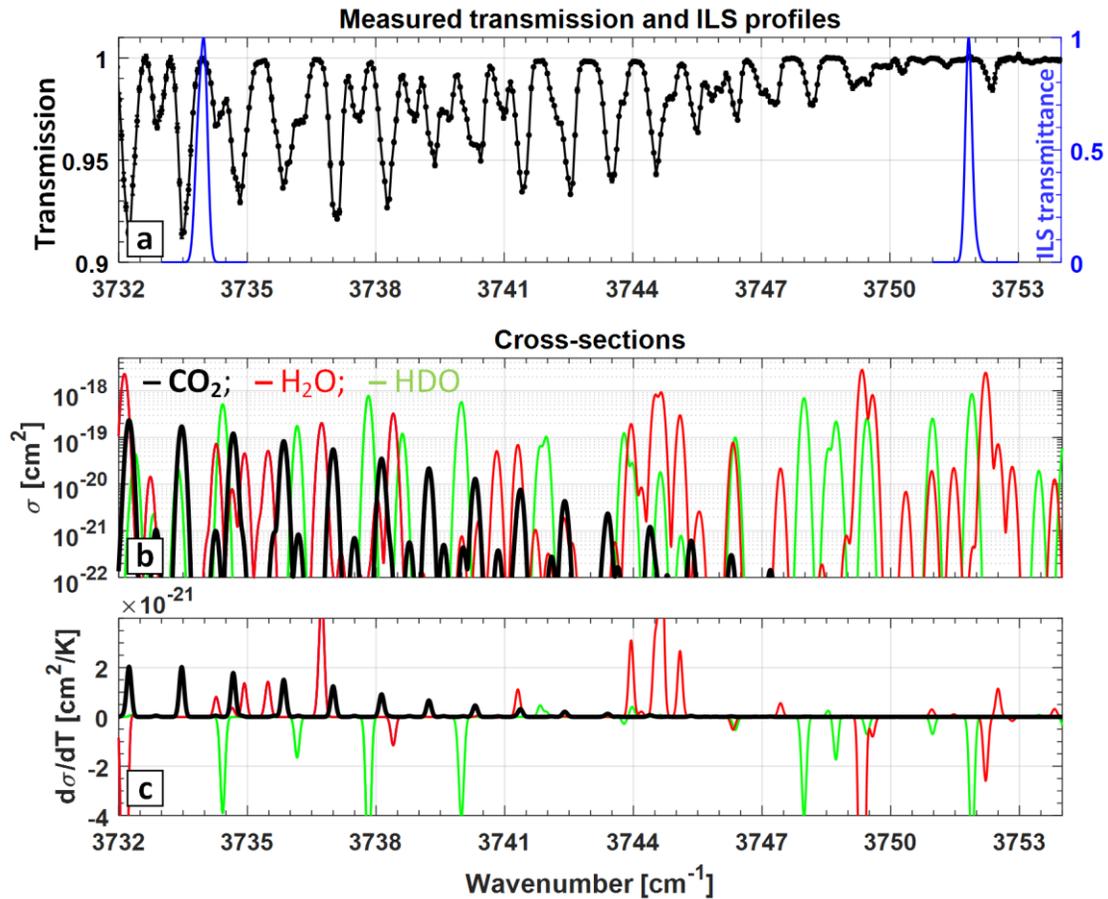
Examples of the derived temperature, H<sub>2</sub>O number density and mixing ratio profiles are shown on Figure S2 in comparison with collocated and simultaneous occultations by the ACS near-IR channel (see in Fedorova et al. (2020)), as well as with the MCD predictions for particular occultations. The validation is demonstrated for three orbits that correspond to different atmospheric scenarios in MY34: before GDS (L<sub>S</sub> 164°), during GDS (L<sub>S</sub> 196.6°) and close to perihelion (L<sub>S</sub> 260.9°). Sensitivity of the NIR spectroscopy does not allow measuring temperature and densities above 100 km. In the MIR case, one can potentially retrieve the temperature altitude profiles spreading from 10-40 km (depending on aerosol opacity) up to 150 km in the considered spectral range (Fig. S2a). The upper altitude of the water detection is 20-25 km higher for the 2.66-2.67  $\mu$ m band in MIR than 1.38  $\mu$ m in NIR (Fig. S2b, S2c).

### Text S3. Estimation of uncertainties

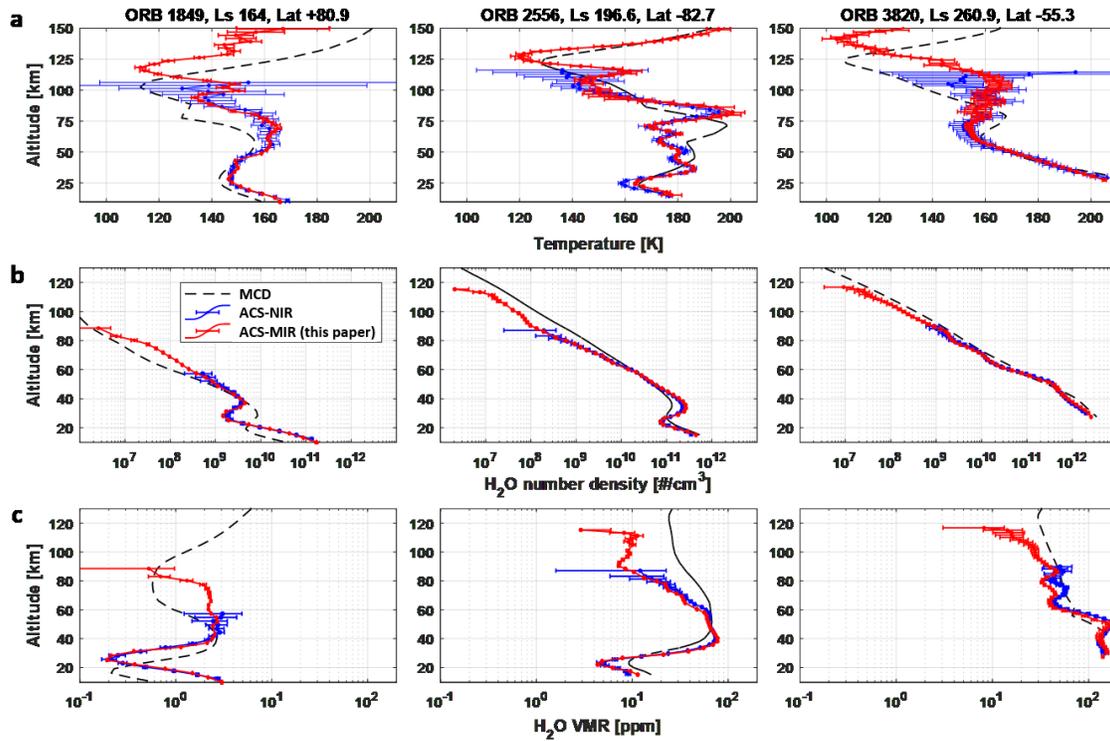
The transmission errors  $\delta J$ , included in the Jacobian matrix, determine uncertainties of the retrievals for each of free parameters. We estimated them for any observed altitude as a multiplication of the derivatives array by the transposed one,  $\delta X = \sqrt{\left[ \frac{\partial A}{\partial X} \right] \times \left[ \frac{\partial A}{\partial X} \right]^T}$ , over the spectral points. The error bars  $\delta J$  are composed of a statistical noise and some systematic uncertainties. We derived the statistical noise analyzing transmission fluctuations near 1 at altitudes above 200 km. It gave the standard deviation of about  $5 \cdot 10^{-4}$  for a signal recorded in the considered FPA row and in the center of the range, 3735-3750 cm<sup>-1</sup>. The systematics comes

from the obtained calibrations (wavenumber, ILS) and from the hydrostatic pressure and temperature uncertainties when modeling transmission spectra on the final step of H<sub>2</sub>O VMR retrievals. On average, it increases the  $\delta J$  value up to 0.001-0.0015. The described uncertainties establish a detection limit of  $5 \cdot 10^6 - 1 \cdot 10^7 \text{ cm}^{-3}$  for H<sub>2</sub>O number density above 100 km. When the water enrichment exceeds this order of magnitude, we observe a positive detection of the molecular abundance at high altitudes. In the opposite case, an upper limit is determined (see Fig. 1c).

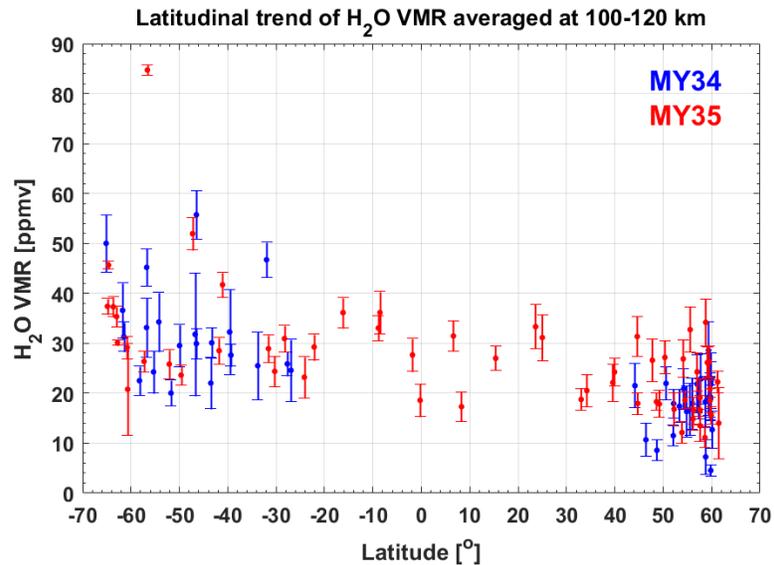
Supplementing Figures:



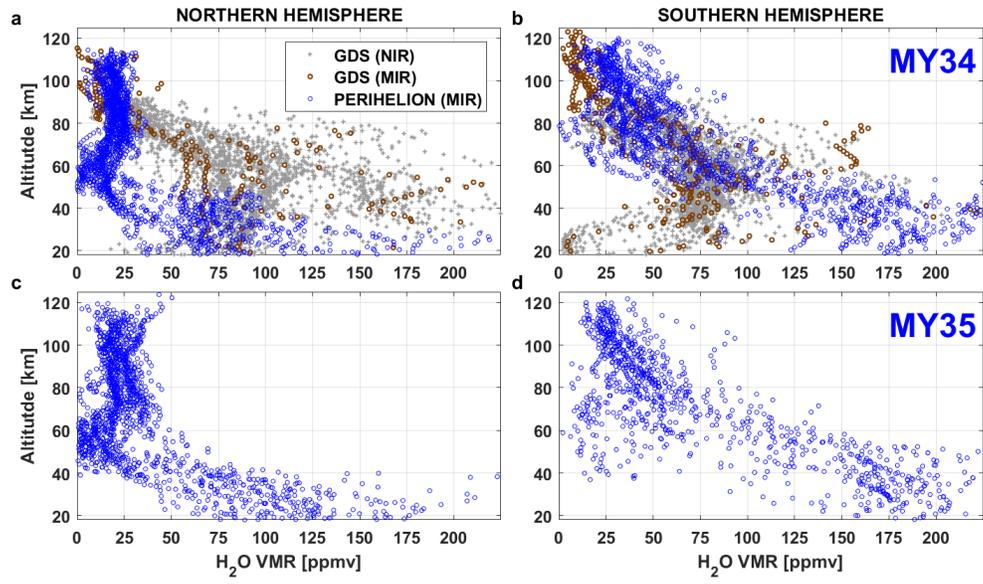
**Figure S1.** Absorption spectroscopy in the echelle order #223 of ACS MIR. **a:** Example of measured transmission spectrum at a target altitude of 100 km (black axis) and the instrumental line shape (ILS) for both edges of the considered wavenumber range (blue axis). **b:** Absorption cross-sections for molecules CO<sub>2</sub> (black), H<sub>2</sub>O (red), and HDO (green) calculated for  $T=150 \text{ K}$  and  $p=10^{-5} \text{ mbar}$  and convolved by ILS. **c:** Partial temperature derivatives of the considered cross-sections.



**Figure S2.** Validation of vertical profiles retrieved by ACS-MIR (red) with analogous profiles of ACS-NIR (from Fedorova et. al, 2020) (blue) and with predictions of MCD (Millour et al., 2018) (black): (a) temperature; (b) H<sub>2</sub>O number density; (c) H<sub>2</sub>O volume mixing ratio (VMR).



**Figure S3.** Latitudinal trend of H<sub>2</sub>O volume mixing ratio (VMR) in the perihelion seasons (Ls 250°–295°) of MY<sub>34</sub> (blue) and MY<sub>35</sub> (red). Each individual point is a weighted mean value revealed over altitude range of 100-120 km from a vertical profile of water VMR.



**Figure S4.** Altitude distribution of H<sub>2</sub>O volume mixing ratio (VMR) during the GDS (MY 34) and the perihelion season (MY 34, 35). The dataset used includes all ACS MIR observations at GDS of MY34 ( $L_S=195^\circ-220^\circ$ ) (brown) and around the perihelion point ( $L_S=250^\circ-295^\circ$ ) (blue). Panels (a, b) for MY34 and (c, d) for MY35 from the Northern (a, c) and Southern (b, d) hemispheres. Light grey points are the NIR data during GDS of MY34 (Fedorova et al., 2020).