

# Free infragravity waves in the North Sea

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## Abstract

Infragravity waves are low-frequency surface waves that can impact a variety of nearshore and oceanic processes. Recent measurements in the North Sea showed that significant bursts of infragravity energy occurred during storm events. Using a spectral wave model, we show that a substantial part of this energy was radiated from distant shorelines where it was generated by the incident sea-swell waves. These radiated infragravity waves can cross the sea basin and reach distant shorelines where they add to locally generated infragravity waves that are trapped by refraction. During storms, the shoreward directed component of the infragravity waves can reach up to  $O(0.5)$  m in height along the coastline, suggesting that they can potentially impact the coastal environment of the North Sea.

# Free infragravity waves in the North Sea

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## Key Points:

- Significant bursts of free infragravity energy in the North Sea can be explained by radiation from distant shorelines
- The origin of the free infragravity waves depends on storm intensity and track, and particularly where largest sea-swell waves make landfall
- The shoreward directed free infragravity waves can reach heights up to 0.6 m, suggesting that they may impact the coastal environment

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**Abstract**

Infragravity waves are low-frequency surface waves that can impact a variety of nearshore and oceanic processes. Recent measurements in the North Sea showed that significant bursts of infragravity energy occurred during storm events. Using a spectral wave model, we show that a substantial part of this energy was radiated from distant shorelines where it was generated by the incident sea-swell waves. These radiated infragravity waves can cross the sea basin and reach distant shorelines where they add to locally generated infragravity waves that are trapped by refraction. During storms, the shoreward directed component of the infragravity waves can reach up to  $\mathcal{O}(0.5)$  m in height along the coastline, suggesting that they can potentially impact the coastal environment of the North Sea.

**Plain Language Summary**

Infragravity waves are long period surface waves that are small in the ocean but much larger in the shallow waters closer to the coast. In coastal regions, they are known to affect coastal safety assessments as they can, for example, impact the overtopping of coastal structures, enhance the erosion of dunes, and trigger the resonance of harbours. To date, such assessments have typically assumed that these infragravity waves were locally generated by sea-swell waves (waves generated by the wind). Recent observations in the North Sea showed that large infragravity waves occurred during storm events. By analyzing the measurements and using a wave model, we have shown that a substantial part of these waves was not generated locally but originated from distant shorelines. At such a distant shoreline, infragravity waves were generated by the breaking of sea-swell waves at the beach. After generation, the infragravity waves radiate into the sea, and our modelling shows that they are able to cross the North Sea and reach distant shorelines. Here, they combine with locally radiated infragravity waves. Our study suggests that these radiated infragravity waves can be of substantial height during storm events, and may impact the coastal environment of the North Sea.

**1 Introduction**

Infragravity (IG) waves are longer period surface gravity waves with typical frequencies ranging between 0.005-0.05 Hz. They are typically considered small in oceanic waters with heights of  $\mathcal{O}$  (cm) (e.g., [Webb et al., 1991](#); [Aucan & Ardhuin, 2013](#)), but can reach heights up to  $\mathcal{O}$  (m) in shallow water during severe weather events (e.g., [Sheremet et al., 2014](#); [Matsuba et al., 2020](#)). In the past decades it has been well established that IG waves contribute to various nearshore processes, such as nearshore hydrodynamics (e.g., [Guza & Thornton, 1985](#); [Henderson & Bowen, 2002](#); [Pomeroy et al., 2012](#)), sediment transport (e.g., [Aagaard & Greenwood, 1994](#); [de Bakker et al., 2016](#)), and erosion of beaches and dunes (e.g., [Russell, 1993](#); [Van Thiel de Vries et al., 2008](#)). Associated with their longer periods, they can also trigger harbour seiches (e.g., [Bowers, 1977](#); [Okiihiro et al., 1993](#); [Thotagamuwage & Pattiaratchi, 2014](#); [Cuomo & Guza, 2017](#)) and excite large motions of moored vessels (e.g., [Van der Molen et al., 2006](#); [van der Molen et al., 2016](#)). Despite their relative small amplitude in deeper water, they have been found to be the source of seismic hum (e.g., [Rhie & Romanowicz, 2006](#); [Webb, 2007](#); [Ardhuin et al., 2015](#)) and may impact the integrity of ice-shelves in polar regions (e.g., [Bromirski et al., 2010, 2015](#)).

IG waves are generally considered to result from the interactions among wind-generated (sea-swell) surface gravity waves (e.g., [Longuet-Higgins & Stewart, 1962](#); [Hasselmann, 1962](#)), see [Bertin et al. \(2018\)](#) for a recent review. The strength of these interactions is depth dependant and only become significant (i.e., approach resonance) in shallow water, resulting in negligible forced IG waves in oceanic water. In the surf-zone, the (near) resonant interactions combined with other generation mechanisms

(e.g., Symonds et al., 1982) result in a substantial transfer of energy from the sea-swell to the IG frequencies. After (partial) reflection at the shoreline (e.g., Battjes et al., 2004; Van Dongeren et al., 2007; De Bakker et al., 2014), free infragravity (FIG) waves that are no longer bound to their forcing radiate seaward into oceanic basins and onto the shelf (e.g., Herbers, Elgar, Guza, & O’Reilly, 1995; Herbers, Elgar, & Guza, 1995). FIG energy levels reduce in deeper water due to refractive trapping (e.g., Gallagher, 1971; Okihiro et al., 1992; Herbers, Elgar, & Guza, 1995), limiting their radiation into ocean basins (e.g., Smit et al., 2018). Nonetheless, part of the FIG energy can reach and cross ocean basins (e.g. Rawat et al., 2014; Crawford et al., 2015; Neale et al., 2015; Bogiatzis et al., 2020), where it can explain the majority of the energy at IG frequencies due to weak (non-resonant) local forcing by sea-swell waves. Other generation mechanics in deep water have also been proposed, such as atmospheric forcing by wind speed fluctuations (de Jong & Battjes, 2004; Vrećica et al., 2019) and IG-tidal interactions (Sugioka et al., 2010), but – to date – most of the IG energy in intermediate and deep water has been explained by FIG radiation from distant shorelines.

Recent analysis of measurements in intermediate water depths (approx. 30 m) of the southern North Sea showed significant bursts of IG energy during storm events (Reniers et al., 2021). Only part of this energy could be attributed to local forcing from sea-swell (ranging between 10-100% depending on geographic location and timing relative to the peak of the storm). As a result, the source and origin of a substantial part of the IG energy remains unclear. The objective of this paper is to understand the dominant source and origin of this FIG energy for four of the most severe storm events in the observational record. We use the spectral wave model SWAN (Booij et al., 1999) extended with an empirical source of FIG energy along the shoreline (Ardhuin et al., 2014) to determine the contribution from FIG radiated from the coastlines bordering the North Sea (described in Section 2). Model-data comparisons at the three available measurement stations show that most of the FIG energy can be explained by radiation from distant shorelines (Section 3). In section 3.4, the model results are further analysed to gain insight into the onshore component of the FIG energy along the coastlines of the southern North Sea. This is followed by a discussion of the main conclusions of this work (Section 4).

## 2 Methods

### 2.1 Observations

Between 2010-2018, significant bursts of IG energy were observed during storm events at three measurement stations located in the intermediate water depths (approximately 30 m) of the southern North Sea (Reniers et al., 2021). At such intermediate water depths, both FIG and bound IG waves (locally forced by the sea-swell waves) may contribute to energy at the IG frequencies. Decomposing the total IG energy into bound and free components indicated that a substantial part of the energy could not be attributed to local forcing from sea-swell waves.

The contribution by bound IG waves was estimated using second-order equilibrium theory (Hasselmann, 1962) based on measured directional sea-swell spectra. The contribution from FIG was subsequently estimated by subtracting the predicted bound IG spectrum  $E_b(f)$  from the measured IG spectrum  $E(f)$ . The FIG wave height  $H_{FIG}$  was computed by integrating the resulting spectrum over the IG frequency band (due to measurement limitations only lower IG frequencies were available,  $0.005 \leq f \leq 0.01$  Hz),

$$H_{FIG} = 4 \sqrt{\int_{0.005}^{0.01} (E(f) - E_b(f)) df} \quad (1)$$

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## 2.2 Model

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The spectral wave model SWAN was used to simulate the temporal and spatial evolution of FIG waves that were radiated from the shorelines bordering the North Sea. To account for the radiation of FIG waves from the shoreline, the SWAN model was extended with an empirical source of FIG energy following the approach of [Ardhuin et al. \(2014\)](#). The empirical source is based on a parametrization that prescribes the bulk IG wave height based on local sea-swell wave parameters (taken seaward from the surf-zone). This parametrization provided a good correlation at several locations in moderate to deep water ([Ardhuin et al., 2014](#)). Combined with an empirical spectral shape and isotropic directional distribution, the FIG source is given by,

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$$E(f, \theta) = 1.2\alpha_1^2 \frac{kg^2}{c_g 2\pi f} \left( \frac{1}{4} H_{m0} T_{m0,-2} \right)^2 \times \frac{1}{\Delta_f} (\min(1., 0.015/f))^{1.5} \times \frac{1}{2\pi}, \quad (2)$$

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in which  $f$  is the wave frequency,  $k$  is the wave number,  $c_g$  is the group velocity, and  $\Delta_f$  ensures that the frequency distribution integrates to 1.  $H_{m0}$  is the sea-swell significant wave height,  $T_{m0,-2}$  is a sea-swell mean wave period, and  $\alpha_1$  is a (dimensional) calibration parameter. The source term can be imposed along waters of variable depth as the term  $kg^2/c_g 2\pi f$  accounts for the shoaling of a directional broad wave spectrum ([Ardhuin et al., 2014](#)).

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The IG source term was implemented as part of the obstacle functionality in SWAN, by which a line can be specified along which FIG energy should be radiated. In this work, we impose the IG source at intermediate water depths ( $\approx 15$  m) along all coastlines that border the North Sea. In regions with complex shorelines, such as the Scheldt estuary and the Wadden Sea, the source was occasionally located in waters of 10 m depth. In regions with steep and irregular coastlines (e.g., along the Norwegian coast), the source term was located in water depths of  $\approx 20$  m.

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The SWAN model was run in nonstationary mode with a spatial resolution of  $0.025^\circ$  and  $0.0165^\circ$  in longitudinal and latitudinal direction, respectively, a directional resolution of  $8^\circ$  and a time step of 1 hr. Twenty-five discrete frequencies with default logarithmic spacing were used to discretize the IG frequency band ( $0.005 \leq f \leq 0.05$  Hz). The spatial and directional resolutions were found to be sufficient based on sensitivity tests (refer to the Supporting Information for more details). Reducing the time-step did not affect the model results (not shown). No additional source terms were included in the simulations, except for dissipation due to bottom friction using the JONSWAP formulation of [Hasselmann et al. \(1973\)](#). To understand the origin of the FIG waves in the North Sea, additional simulations were run with the same model-setup, but with a subset of coastlines radiating IG energy. Individual simulations were conducted with IG waves radiating from only the Belgium, Dutch, German, Danish, UK, or Norwegian coast.

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The sea-swell wave height  $H_{m0}$  and mean wave period  $T_{m0,-1}$  (due to unavailability of  $T_{m0,-2}$ ) of Eq. 2 were obtained from the global ECMWF ERA5 reanalysis ([Hersbach et al., 2020](#)). These bulk wave parameter were available every hour at a  $0.5^\circ$  resolution, and were interpolated to the depths at which the IG source was specified. The  $\alpha_1$  parameter and the bottom friction coefficient were set based on a calibration study for a single storm event (storm Friedhelm) during which the majority of radiated FIG energy originated from the Danish coast (refer to the Supporting Information for more details). Satisfactory results were obtained for  $\alpha_1 = 18 \times 10^{-4} \text{ s}^{-1}$ , which is of the same order of magnitude as the values used in [Ardhuin et al. \(2014\)](#) and [Rawat et al. \(2014\)](#), in combination with a bottom friction coefficient of  $\chi = 0.01 \text{ m}^2 \text{ s}^{-3}$ . This bottom friction coefficient is lower than the default value  $\chi = 0.038 \text{ m}^2 \text{ s}^{-3}$  typically used for sea-swell waves (e.g., [Zijlema et al., 2012](#)). To study the influence of bottom friction, additional simulations without bottom friction were conducted, for which the calibration study indicated optimal results for a slightly smaller  $\alpha_1$  value

161 of  $14.4 \times 10^{-4} \text{ s}^{-1}$ . In this work, a constant  $\alpha_1$  was used for all shorelines, and no  
 162 attempt was made to optimise  $\alpha_1$  by varying it for different shorelines to account for  
 163 differences between geographic regions (e.g., steep versus mild bottom slopes) that  
 164 may affect FIG radiation.

165 The SWAN model was used to hindcast four storm events that resulted in the  
 166 largest observed  $H_{FIG}$  at the three measurement stations (storm Friedhelm, 8-11 Dec  
 167 2011; Xaver, 5-8 Dec 2013; Axel, 4-5 Jan 2017; and Egon, 13-15 Jan 2017). Directional  
 168 spectra were outputted at the three measurement stations and along the shorelines  
 169 bordering the southern North Sea (at approximately 20 m depth). The predicted  $H_{FIG}$   
 170 was computed by integrating the directional spectra from SWAN over all directions and  
 171 the IG frequency band. When comparing model results with measurements, spectra  
 172 were integrated over the measured IG frequency band ( $0.005 \text{ Hz} \leq f \leq 0.01 \text{ Hz}$ ).

### 173 3 Results

#### 174 3.1 Storm Xaver

175 In Dec 2013, a severe winter storm tracked from north of the UK towards the  
 176 south of Norway and north of Denmark. Significant sea-swell wave heights  $H_{SS}$   
 177 in the central North Sea at station A12 reached 8 m with mean sea swell periods  $T_{m0,-1}$   
 178 exceeding 10 s (Fig. 1c), with largest sea-swell waves occurring in the northern part  
 179 of the North Sea (e.g., Fig. 1a). Storm Xaver produced the largest IG response in the  
 180 observational record at A12, with FIG heights  $H_{FIG} > 0.2 \text{ m}$  for approximately 24 h  
 181 (Fig. 1d).  $H_{FIG}$  was lower but still significant at stations Q1 and EUR, which are  
 182 located further southward and in closer vicinity to the Dutch coast. During this storm  
 183 event, FIG waves contributed significantly to the total IG variance (compare full and  
 184 dashed lines in Fig. 1d, and see also Fig. 4 in Reniers et al. (2021)).

185 In accordance with the measurements (Fig. 1d), the predicted FIG wave heights  
 186 showed great spatial variability in the North Sea (e.g., during the peak of the storm  
 187 (06:00 UTC 6 Dec, Fig. 1b). Largest  $H_{FIG}$  typically occurred near the Danish coast  
 188 due to local radiation, and  $H_{FIG}$  decreased for an increasing distance away from the  
 189 Danish coast.  $H_{FIG}$  was amplified by shoaling in regions of relatively shallow water,  
 190 such as near the Dogger and Norfolk banks.

191 Predicted  $H_{FIG}$  captured the typical magnitude and trend of the observations  
 192 at the three stations (Fig. 1d), which indicates that the observed FIG levels can be  
 193 partly explained by IG wave radiation from neighbouring shorelines. The model failed  
 194 to capture relatively large  $H_{FIG}$  at A12 prior to the peak of the storm (00:00 UTC  
 195 6 Dec), indicating that FIG wave radiation from surrounding beaches cannot explain  
 196 these FIG energy levels. We will return to this in Section 3.3.

197 The model simulations further allow us to understand the temporal and spatial  
 198 variability in  $H_{FIG}$  by considering the origin of the IG energy. For the majority of the  
 199 storm, the modelled  $H_{FIG}$  at A12 primarily originated from the Danish coast (Fig. 1e),  
 200 which can be explained by the occurrence of the largest sea-swell waves in the northern  
 201 part of the North Sea (Fig. 1a). At Q1 and EUR however, most of the modelled IG  
 202 variance originated from the Dutch coast (Fig. 1f-g), with a smaller but non-negligible  
 203 contribution from the Danish coast ( $\approx 10 - 20\%$  of the FIG variance). This further  
 204 illustrates how FIG waves radiated from the Danish coast were attenuated by the  
 205 combined effect of refraction and dissipation by bottom friction before they reached  
 206 these stations in the southern North Sea (we will return to the effect of bottom friction  
 207 in Section 3.3).

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### 3.2 Storm Egon

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During storm Egon, large sea-swell waves occurred between 13 and 15 Jan, with  $H_{SS}$  up to 8 m at A12, and with  $H_{SS}$  peaking at 5 m at EUR in the southern North Sea. Storm Egon was preceded by another storm (11-13 Jan), with weaker but still significant waves ( $H_{SS} > 4$  m) at A12 and Q1. Between 11-13 Jan, FIG patterns were comparable to storm Xaver with largest  $H_{FIG}$  at A12 and progressively smaller  $H_{FIG}$  towards the south at Q1 and EUR. During storm Egon (13-15 Jan) on the other hand,  $H_{FIG}$  was generally largest at Q1 and smallest at A12.

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The model captured the typical magnitude of  $H_{FIG}$  at all three stations, except for large  $H_{FIG}$  at A12 prior to the peak of both storms (similar to the findings for storm Xaver). Notably, the model captured the trends in  $H_{FIG}$  that occurred during the first (largest  $H_{FIG}$  at A12) and second storm event (largest  $H_{FIG}$  at Q1). The general agreement between the model and the observations indicates that a significant part of the IG energy can be attributed to the arrival of FIG waves that were radiated from nearby shorelines.

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During the first storm (11-13 Jan) modelled  $H_{FIG}$  at A12 mainly originated from Denmark, whereas between 13-15 Jan,  $H_{FIG}$  was primarily explained by contributions from Denmark and the UK. The modelling results thus suggest that the variability in  $H_{FIG}$  (largest  $H_{FIG}$  at A12 between 11-13 Jan and largest  $H_{FIG}$  at Q1 between 13-15 Jan) is related to the storm trajectory and the spatial variability of the sea-swell waves, with stronger IG radiation from the Danish coast between 11-13 Jan due to a more northerly storm track, and stronger IG radiation along the shorelines of the southern North Sea (UK and NL) between 13-15 Jan due to a more southerly storm track.

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### 3.3 All storms and the influence of bottom friction

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Comparing the observed and modelled  $H_{FIG}$  at all stations for all considered storm events shows that the model (including bottom friction) generally reproduced the observations (Fig. 3). The agreement was best at Q1 and EURO, where the model Skill indicates that the model explained 75% of the variability in the FIG wave height (Fig. 3 b-c). At these stations, the model suggests that the FIG variance was typically dominated by radiation from the nearby Dutch coast (Fig. 3 e-f).

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The agreement at A12 was poorer (Fig. 3 a), with a lower Skill and larger Relative Bias (RB) consistent with a typical under prediction of  $H_{FIG}$ . This under prediction at A12 occurred consistently at times prior to the peak of all four storms (blue markers in Fig. 3a), as was observed previously for storm Xaver and Egon (Fig. 1-2). Near and following the peak of the storm (gray to red markers in Fig. 3a), the model did capture  $H_{FIG}$ , suggesting that at these time the observed FIG waves originated from surrounding shorelines. At station A12, the modelled FIG variance was typically dominated by radiation from the Danish coast (Fig. 3d).

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The model-data mismatch at A12 prior to the peak of the storm indicates that radiation from surrounding shorelines cannot explain these FIG energy levels. This suggests that other physical mechanisms could be at play. At times of the unexplained FIG energy, the sea-swell waves typically originated from WNW to NNW and passed the Dogger Bank (a shallow shoal located in the central North Sea) prior to reaching station A12 (which is located on the south-eastern side of the Dogger Bank). Although the precise generation mechanism remains unclear, this indicates that the unexplained energy in the IG band could originate from local excitation of FIG waves by sea-swell waves as they pass over the shallow shoal (e.g., Molin, 1982; Mei & Benmoussa, 1984; Li et al., 2020; Contardo et al., 2021).

257 To study the influence by bottom friction, the SWAN simulations were repeated  
 258 excluding this source term and a slightly lower  $\alpha_1$  value (to compensate for the ab-  
 259 sence of dissipation by bottom friction). The influence of bottom friction on  $H_{FIG}$   
 260 is particularly notable at stations Q1 and EUR, where it reduced the Danish con-  
 261 tribution at times that significant IG variance originated from Denmark (Fig. 3e-f  
 262 versus Fig. 3k-i), resulting in a minor improvement of the overall model skill. These  
 263 results indicate that bottom friction can provide some additional attenuation of FIG  
 264 waves as they propagate through the relatively shallow North Sea basin. For the four  
 265 storms considered, the bottom friction only led to noticeable attenuation when most  
 266 FIG energy originated from the Danish coast.

### 267 3.4 Shoreward directed FIG energy

268 The results in the previous subsections did not consider the full IG frequency  
 269 band (due to measurement limitations) and focused on three stations in intermediate  
 270 water depths of the southern North Sea. To gain insight into the potential magnitude of  
 271 FIG energy that is incident to the coastlines, two-dimensional spectra were outputted  
 272 at approximately 20 m depth along the coastlines bordering the southern North Sea.  
 273 Incident FIG wave heights were computed by integrating the SWAN spectra over the  
 274 full IG frequency band and the shoreward directed directional bins,

$$H_{FIG}^+ = \int_{0.005}^{0.05} \int_{\theta_p-90^\circ}^{\theta_p+90^\circ} E(f, \theta) d\theta df, \quad (3)$$

275 in which  $\theta_p$  is the angle perpendicular to the shore.

276 Fig. 4 shows the maximum  $H_{FIG}^+$  that occurred during each of the four storms  
 277 (top to bottom panels) along the UK coastline (left panels) and the coastline between  
 278 Belgium and Denmark (right panels). The maximum  $H_{FIG}^+$  shows a strong spatial  
 279 variation along the coast of western Europe (Fig. 4), with typically largest  $H_{FIG}^+$   
 280 of up to 0.6 m along the Danish coast. Here,  $H_{FIG}^+$  was primarily explained by IG energy  
 281 from local origin (refractive trapping), associated with significant local radiation due  
 282 to large sea-swell waves along the Danish coast (especially during storm Friedhelm and  
 283 Xaver).

284 The maximum  $H_{FIG}^+$  decreased towards the south (from the German to the Bel-  
 285 gian coast), with a decreasing contribution from the Danish coast (associated with  
 286 increased attenuation due to refraction and bottom friction). Along the Belgian to  
 287 German coast, the maximum  $H_{FIG}^+$  was of mixed origin, with different major con-  
 288 tributing shorelines depending on where significant sea-swell waves made landfall dur-  
 289 ing the particular storm. Along the UK coast, the maximum  $H_{FIG}^+$  was of similar  
 290 magnitude as along the Belgian and Dutch coast, with small contribution from local  
 291 radiation as the UK coast lacked exposure to significant sea-swell waves during these  
 292 four storms (associated with a typical north to east storm track). Nonetheless, the  
 293 maximum  $H_{FIG}^+$  along the UK coast reached up to 0.2-0.3 m due to arrival of FIG  
 294 waves from remote shorelines (e.g., from Denmark during storm Friedhelm Fig. 4a,  
 295 and from the Netherlands during storm Egon at lower latitudes, Fig. 4g).

## 296 4 Discussion and conclusions

297 In-situ observations at three measurement stations in the southern North Sea  
 298 revealed the occurrence of substantial bursts of free infragravity wave energy during  
 299 four significant storm events. A spectral wave model that accounts for the radiation  
 300 of FIG waves from adjacent beaches was able to explain up to 75% of the observed  
 301 variability in the FIG wave height. The model captured the typical magnitude and

302 temporal variation of the FIG wave height at the three measurement stations. Al-  
 303 though the model failed to explain significant FIG energy levels that occurred prior to  
 304 the peak of storms at the measurement site in the central North Sea (suggesting that  
 305 other processes also contributed to FIG energy at this location), the overall model-data  
 306 agreement suggests that a significant fraction of the FIG energy levels in the southern  
 307 North Sea originated from distant shorelines. These results are in accordance with  
 308 previous studies on the shelf (Smit et al., 2018) and in the deep ocean (e.g., Rawat et  
 309 al., 2014; Crawford et al., 2015; Neale et al., 2015; Bogiatzis et al., 2020) that linked  
 310 significant energy at IG frequencies to radiation of FIG from distant shorelines.

311 The origin of the FIG energy in the southern North Sea was found to depend  
 312 on the storm characteristics, and in particular on where large sea-swell waves made  
 313 landfall. The model suggests that radiated FIG energy was able to cross the North  
 314 Sea basin and reach neighbouring shorelines, where it adds to FIG energy with a local  
 315 origin (regionally trapped due to refraction). The FIG energy levels from a distant  
 316 source were attenuated by refraction and bottom friction, and the predicted nearshore  
 317 FIG energy levels were typically dominated by waves with a local origin. Along the  
 318 coastlines at approximately 20 m depth, the model showed that during storms the  
 319 shoreward directed component of the FIG waves can reach up to 0.6 m in height,  
 320 which suggests that FIG waves can potentially impact the coastal environment in the  
 321 North Sea and influence processes like flooding, dune erosion, and seiching of harbours.

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 328 sured infragravity wave data set and the source code of the adapted SWAN model are  
 329 available from <https://surfdrive.surf.nl/files/index.php/s/dITA6eD1r4iFxDj>.  
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 331 the 4TU.ResearchData repository (<https://data.4tu.nl/repository>). We thank  
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 333 (SWAN-DCSM-j15-v1).

### 334 References

- 335 Aagaard, T., & Greenwood, B. (1994). Suspended sediment transport and the role  
 336 of infragravity waves in a barred surf zone. *Marine Geology*, *118*(1-2), 23–48.  
 337 doi: 10.1016/0025-3227(94)90111-2
- 338 Ardhuin, F., Gualtieri, L., & Stutzmann, E. (2015). How ocean waves rock the  
 339 Earth: Two mechanisms explain microseisms with periods 3 to 300s. *Geophys-  
 340 ical Research Letters*, *42*(3), 765–772. doi: 10.1002/2014GL062782
- 341 Ardhuin, F., Rawat, A., & Aucan, J. (2014, 5). A numerical model for free infra-  
 342 gravity waves: Definition and validation at regional and global scales. *Ocean  
 343 Modelling*, *77*, 20–32. Retrieved from [http://linkinghub.elsevier.com/  
 344 retrieve/pii/S1463500314000274](http://linkinghub.elsevier.com/retrieve/pii/S1463500314000274) [http://www.sciencedirect.com/  
 345 science/article/pii/S1463500314000274](http://www.sciencedirect.com/science/article/pii/S1463500314000274) doi: 10.1016/j.ocemod.2014  
 346 .02.006
- 347 Aucan, J., & Ardhuin, F. (2013, 7). Infragravity waves in the deep ocean: An up-  
 348 ward revision. *Geophysical Research Letters*, *40*(13), 3435–3439. Retrieved  
 349 from <http://doi.wiley.com/10.1002/grl.50321> doi: 10.1002/grl.50321
- 350 Battjes, J. A., Bakkenes, H. J., Janssen, T. T., & Van Dongeren, A. R. (2004).  
 351 Shoaling of subharmonic gravity waves. *Journal of Geophysical Research*,  
 352 *109*(C2). Retrieved from <http://doi.wiley.com/10.1029/2003JC001863>

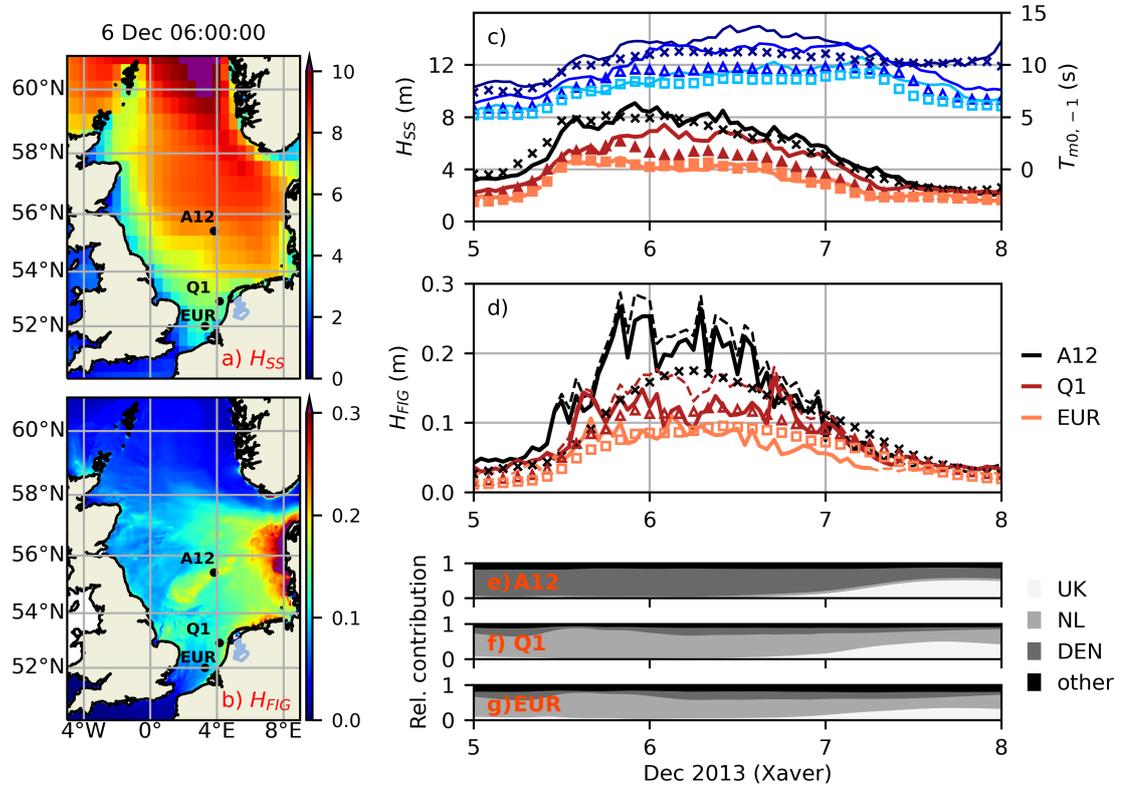
- 353 doi: 10.1029/2003JC001863
- 354 Bertin, X., Bakker, A. d., Dongeren, A. v., Coco, G., André, G., Ardhuin, F., ...  
 355 Tissier, M. (2018, 1). Infragravity waves: from driving mechanisms to impacts.  
 356 *Earth-Science Reviews*. Retrieved from [http://linkinghub.elsevier.com/  
 357 retrieve/pii/S0012825217303239](http://linkinghub.elsevier.com/retrieve/pii/S0012825217303239) doi: 10.1016/j.earscirev.2018.01.002
- 358 Bogiatzis, P., Karamitrou, A., Ward Neale, J., Harmon, N., Rychert, C. A., &  
 359 Srokosz, M. (2020, 6). Source Regions of Infragravity Waves Recorded at  
 360 the Bottom of the Equatorial Atlantic Ocean, Using OBS of the PI-LAB Ex-  
 361 periment. *Journal of Geophysical Research: Oceans*, 125(6), 1–17. Retrieved  
 362 from <https://onlinelibrary.wiley.com/doi/abs/10.1029/2019JC015430>  
 363 doi: 10.1029/2019JC015430
- 364 Booij, N., Ris, R. C., & Holthuijsen, L. H. (1999). A third-generation wave model  
 365 for coastal regions: 1. Model description and validation. *Journal of Geophysi-  
 366 cal Research*, 104(C4), 7649–7666. Retrieved from [http://doi.wiley.com/10  
 367 .1029/98JC02622](http://doi.wiley.com/10.1029/98JC02622) doi: 10.1029/98JC02622
- 368 Bowers, E. C. (1977, 4). Harbour resonance due to set-down beneath wave  
 369 groups. *Journal of Fluid Mechanics*, 79(1), 71–92. Retrieved from  
 370 [http://www.journals.cambridge.org/abstract\\_S0022112077000044](http://www.journals.cambridge.org/abstract_S0022112077000044) doi:  
 371 10.1017/S0022112077000044
- 372 Bromirski, P. D., Diez, A., Gerstoft, P., Stephen, R. A., Bolmer, T., Wiens, D. A.,  
 373 ... Nyblade, A. (2015). Ross ice shelf vibrations. *Geophysical Research  
 374 Letters*, 42(18), 7589–7597. doi: 10.1002/2015GL065284
- 375 Bromirski, P. D., Sergienko, O. V., & MacAyeal, D. R. (2010, 1). Transoceanic in-  
 376 fragravity waves impacting Antarctic ice shelves. *Geophysical Research Letters*,  
 377 37(2). Retrieved from <http://doi.wiley.com/10.1029/2009GL041488> doi:  
 378 10.1029/2009GL041488
- 379 Contardo, S., Lowe, R. J., Hansen, J. E., Rijnsdorp, D. P., Dufois, F., & Symonds,  
 380 G. (2021, 2). Free and forced components of shoaling long waves in the  
 381 absence of short wave breaking. *Journal of Physical Oceanography*. Re-  
 382 trieved from [https://journals.ametsoc.org/view/journals/phoc/aop/  
 383 JPO-D-20-0214.1/JPO-D-20-0214.1.xml](https://journals.ametsoc.org/view/journals/phoc/aop/JPO-D-20-0214.1/JPO-D-20-0214.1.xml) doi: 10.1175/JPO-D-20-0214.1
- 384 Crawford, W., Ballu, V., Bertin, X., & Karpytchev, M. (2015, 7). The sources of  
 385 deep ocean infragravity waves observed in the North Atlantic Ocean. *Jour-  
 386 nal of Geophysical Research: Oceans*, 120(7), 5120–5133. Retrieved from  
 387 [http://www.nature.com/articles/175238c0http://doi.wiley.com/  
 388 10.1002/2014JC010657](http://www.nature.com/articles/175238c0http://doi.wiley.com/10.1002/2014JC010657) doi: 10.1002/2014JC010657
- 389 Cuomo, G., & Guza, R. T. (2017). Infragravity Seiches in a Small Harbor. *Journal  
 390 of Waterway, Port, Coastal, and Ocean Engineering*, 143(5), 04017032. doi: 10  
 391 .1061/(asce)ww.1943-5460.0000392
- 392 de Bakker, A. T., Brinkkemper, J. A., van der Steen, F., Tissier, M. F., & Ruessink,  
 393 B. G. (2016). Cross-shore sand transport by infragravity waves as a function  
 394 of beach steepness. *Journal of Geophysical Research: Earth Surface*, 121(10),  
 395 1786–1799. doi: 10.1002/2016JF003878
- 396 De Bakker, A. T. M., Tissier, M. F. S., & Ruessink, B. G. (2014, 1). Shoreline dissi-  
 397 pation of infragravity waves. *Continental Shelf Research*, 72, 73–82. Retrieved  
 398 from <http://linkinghub.elsevier.com/retrieve/pii/S0278434313003786>  
 399 doi: 10.1016/j.csr.2013.11.013
- 400 de Jong, M. P., & Battjes, J. A. (2004). Low-frequency sea waves generated by at-  
 401 mospheric convection cells. *Journal of Geophysical Research: Oceans*, 109(1),  
 402 1–18. doi: 10.1029/2003jc001931
- 403 Gallagher, B. (1971). Generation of surf beat by non-linear wave interactions. *Jour-  
 404 nal of Fluid Mechanics*, 49(1), 1–20. doi: 10.1017/S0022112071001897
- 405 Guza, R. T., & Thornton, E. B. (1985). Observations of surf beat. *Journal of  
 406 Geophysical Research*, 90(C2), 3161–3172. Retrieved from [http://doi.wiley  
 407 .com/10.1029/JC090iC02p03161](http://doi.wiley.com/10.1029/JC090iC02p03161) doi: 10.1029/JC090iC02p03161

- 408 Hasselmann, K. (1962, 4). On the non-linear energy transfer in a gravity-wave  
 409 spectrum: Part 1. General theory. *Journal of Fluid Mechanics*, 12(4), 481–  
 410 500. Retrieved from [http://journals.cambridge.org/production/action/](http://journals.cambridge.org/production/action/cjoGetFulltext?fulltextid=368683)  
 411 [cjoGetFulltext?fulltextid=368683](http://www.journals.cambridge.org/abstract_S0022112062000373)[http://www.journals.cambridge.org/](http://www.journals.cambridge.org/abstract_S0022112062000373)  
 412 [abstract\\_S0022112062000373](http://www.journals.cambridge.org/abstract_S0022112062000373) doi: 10.1017/S0022112062000373
- 413 Hasselmann, K., Barnett, T. P., Bouws, E., Carlson, H., Cartwright, D. E., Enke,  
 414 K., ... Walden, H. (1973). Measurements of wind-wave growth and swell  
 415 decay during the Joint North Sea Wave Project (JONSWAP). *Erganzungsheft*  
 416 *zur Deutschen Hydrographischen Zeitschrift*, A8(12).
- 417 Henderson, S. M., & Bowen, A. J. (2002). Observations of surf beat forcing and dis-  
 418 sipation. *Journal of Geophysical Research*, 107(C11). Retrieved from [http://](http://doi.wiley.com/10.1029/2000JC000498)  
 419 [doi.wiley.com/10.1029/2000JC000498](http://doi.wiley.com/10.1029/2000JC000498) doi: 10.1029/2000JC000498
- 420 Herbers, T. H. C., Elgar, S., & Guza, R. T. (1995). Generation and propagation of  
 421 infragravity waves. *Journal of Geophysical Research*, 100(C12), 24863–24872.  
 422 Retrieved from [http://onlinelibrary.wiley.com/doi/10.1029/95JC02680/](http://onlinelibrary.wiley.com/doi/10.1029/95JC02680/full)  
 423 [full](http://onlinelibrary.wiley.com/doi/10.1029/95JC02680/full)[http://doi.wiley.com/10.1029/95JC02680](http://onlinelibrary.wiley.com/doi/10.1029/95JC02680/full) doi: 10.1029/95JC02680
- 424 Herbers, T. H. C., Elgar, S., Guza, R. T., & O'Reilly, W. C. (1995, 6). Infragravity-  
 425 Frequency (0.005–0.05 Hz) Motions on the Shelf. Part II: Free Waves. *Jour-*  
 426 *nal of Physical Oceanography*, 25(6), 1063–1079. Retrieved from [http://](http://journals.ametsoc.org/doi/abs/10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2)  
 427 [journals.ametsoc.org/doi/abs/10.1175/1520-0485\(1995\)025<1063:](http://journals.ametsoc.org/doi/abs/10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2)  
 428 [IFHMOT>2.0.CO;2](http://journals.ametsoc.org/doi/abs/10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2);  
 429 [2# .U\\_xzxkbnwr4.mendeleymhttp://journals.ametsoc.org/](http://journals.ametsoc.org/doi/abs/10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2)  
 430 [doi/abs/10.1175/1520-0485\(1995\)025<1063:IFHMOT>2.0.CO;2](http://journals.ametsoc.org/doi/abs/10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2) doi:  
 431 10.1175/1520-0485(1995)025<1063:IFHMOT>2.0.CO;2
- 431 Hersbach, H., Bell, B., Berrisford, P., Hirahara, S., Horányi, A., Muñoz-Sabater, J.,  
 432 ... Thépaut, J. N. (2020). The ERA5 global reanalysis. *Quarterly Journal of*  
 433 *the Royal Meteorological Society*, 146(730), 1999–2049. doi: 10.1002/qj.3803
- 434 Li, S., Liao, Z., Liu, Y., & Zou, Q. (2020). Evolution of Infragravity Waves Over  
 435 a Shoal Under Nonbreaking Conditions. *Journal of Geophysical Research:*  
 436 *Oceans*, 125(8), 1–15. doi: 10.1029/2019JC015864
- 437 Longuet-Higgins, M. S., & Stewart, R. W. (1962, 3). Radiation stress and mass  
 438 transport in gravity waves, with application to ‘surf beats’. *Journal of Fluid*  
 439 *Mechanics*, 13(4), 481–504. Retrieved from [http://www.journals.cambridge](http://www.journals.cambridge.org/abstract_S0022112062000877)  
 440 [.org/abstract\\_S0022112062000877](http://www.journals.cambridge.org/abstract_S0022112062000877)[http://www.journals.cambridge.org/](http://www.journals.cambridge.org/abstract_S0022112062000877)  
 441 [abstract\\_S0022112062000877](http://www.journals.cambridge.org/abstract_S0022112062000877) doi: 10.1017/S0022112062000877
- 442 Matsuba, Y., Shimozono, T., & Sato, S. (2020). Wave-breaking modulation by in-  
 443 fragravity waves during an extreme typhoon. *PLoS ONE*, 15(4), 1–14. Re-  
 444 trieved from <http://dx.doi.org/10.1371/journal.pone.0231242> doi: 10  
 445 .1371/journal.pone.0231242
- 446 Mei, C. C., & Benmoussa, C. (1984, 4). Long waves induced by short-wave groups  
 447 over an uneven bottom. *Journal of Fluid Mechanics*, 139, 219. Retrieved from  
 448 [http://journals.cambridge.org/abstract\\_S0022112084000331](http://journals.cambridge.org/abstract_S0022112084000331)[http://](http://journals.cambridge.org/abstract_S0022112084000331)  
 449 [www.journals.cambridge.org/abstract\\_S0022112084000331](http://www.journals.cambridge.org/abstract_S0022112084000331) doi:  
 450 10.1017/S0022112084000331
- 451 Molin, B. (1982). *On the generation of long-period second order free waves due to*  
 452 *changes in the bottom profile* (Tech. Rep.). Tokyo, Japan.
- 453 Neale, J., Harmon, N., & Srokosz, M. (2015, 9). Source regions and reflection of  
 454 infragravity waves offshore of the Pacific Northwest. *Journal of Geophysical Research: Oceans*, 120(9),  
 455 6474–6491. Retrieved from [https://onlinelibrary.wiley.com/doi/abs/](https://onlinelibrary.wiley.com/doi/abs/10.1002/2015JC010891)  
 456 [10.1002/2015JC010891](https://onlinelibrary.wiley.com/doi/abs/10.1002/2015JC010891) doi: 10.1002/2015JC010891
- 457 Okihiro, M., Guza, R. T., & Seymour, R. J. (1992). Bound infragravity waves.  
 458 *Journal of Geophysical Research*, 97(C7), 11453–11469. Retrieved from  
 459 <http://onlinelibrary.wiley.com/doi/10.1029/92JC00270/full>[http://](http://onlinelibrary.wiley.com/doi/10.1029/92JC00270/full)  
 460 [doi.wiley.com/10.1029/92JC00270](http://onlinelibrary.wiley.com/doi/10.1029/92JC00270/full) doi: 10.1029/92JC00270
- 461 Okihiro, M., Guza, R. T., & Seymour, R. J. (1993). Excitation of seiche ob-

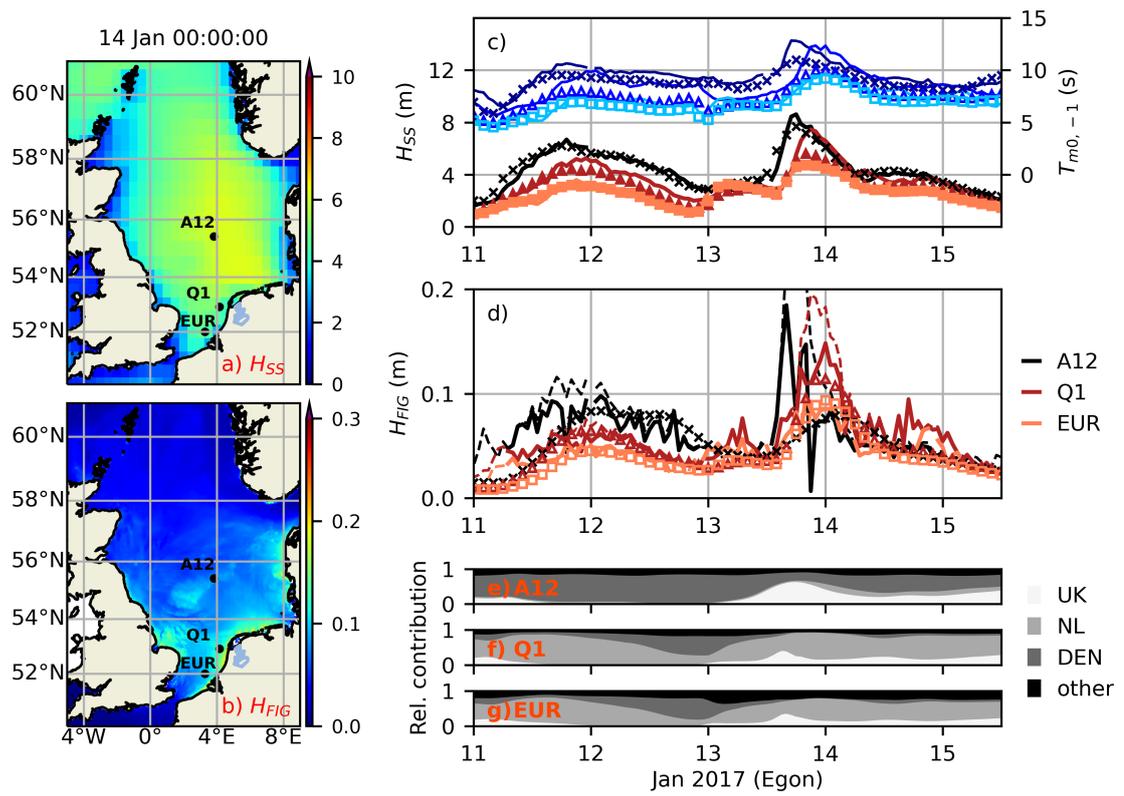
- 463 served in a small harbor. *Journal of Geophysical Research*, 98(C10), 18201–  
 464 18211. Retrieved from <http://doi.wiley.com/10.1029/93JC01760> doi:  
 465 10.1029/93JC01760
- 466 Pomeroy, A., Lowe, R., Symonds, G., Van Dongeren, A., & Moore, C. (2012, 11).  
 467 The dynamics of infragravity wave transformation over a fringing reef. *Journal*  
 468 *of Geophysical Research*, 117(C11). Retrieved from [http://doi.wiley.com/](http://doi.wiley.com/10.1029/2012JC008310)  
 469 10.1029/2012JC008310 doi: 10.1029/2012JC008310
- 470 Rawat, A., Arduin, F., Ballu, V., Crawford, W., Corela, C., & Aucan, J. (2014,  
 471 11). Infragravity waves across the oceans. *Geophysical Research Letters*,  
 472 41(22), 7957–7963. Retrieved from [http://doi.wiley.com/10.1002/](http://doi.wiley.com/10.1002/2014GL061604)  
 473 2014GL061604 doi: 10.1002/2014GL061604
- 474 Reniers, A. J. H. M., Naporowski, R., Tissier, M. F. S., de Schipper, M. A., Akrish,  
 475 G., & Rijnsdorp, D. P. (2021, 1). North Sea Infragravity Wave Observa-  
 476 tions. *Journal of Marine Science and Engineering*, 9(2), 141. Retrieved from  
 477 <https://www.mdpi.com/2077-1312/9/2/141> doi: 10.3390/jmse9020141
- 478 Rhie, J., & Romanowicz, B. (2006). A study of the relation between ocean storms  
 479 and the Earth's hum. *Geochemistry, Geophysics, Geosystems*, 7(10). doi: 10  
 480 .1029/2006GC001274
- 481 Russell, P. E. (1993, 11). Mechanisms for beach erosion during storms. *Con-*  
 482 *tinental Shelf Research*, 13(11), 1243–1265. Retrieved from [http://](http://www.sciencedirect.com/science/article/pii/027843439390051X)  
 483 [www.sciencedirect.com/science/article/pii/027843439390051X](http://www.sciencedirect.com/science/article/pii/027843439390051X)  
 484 <http://linkinghub.elsevier.com/retrieve/pii/027843439390051X> doi:  
 485 10.1016/0278-4343(93)90051-X
- 486 Sheremet, A., Staples, T., Arduin, F., Suanez, S., & Fichaut, B. (2014, 1). Ob-  
 487 servations of large infragravity-wave run-up at Banneg Island, France. *Geo-*  
 488 *physical Research Letters*. Retrieved from [http://dx.doi.org/10.1002/](http://dx.doi.org/10.1002/2013GL058880)  
 489 2013GL058880 doi: 10.1002/2013GL058880
- 490 Smit, P. B., Janssen, T. T., Herbers, T. H. C., Taira, T., & Romanowicz, B. A.  
 491 (2018, 7). Infragravity Wave Radiation Across the Shelf Break. *Journal of*  
 492 *Geophysical Research: Oceans*, 123(7), 4483–4490. Retrieved from [http://](http://doi.wiley.com/10.1029/2018JC013986)  
 493 [doi.wiley.com/10.1029/2018JC013986](http://doi.wiley.com/10.1029/2018JC013986) doi: 10.1029/2018JC013986
- 494 Sugioka, H., Fukao, Y., & Kanazawa, T. (2010). Evidence for infragravity wave-tide  
 495 resonance in deep oceans. *Nature Communications*, 1(7), 1–7. doi: 10.1038/  
 496 ncomms1083
- 497 Symonds, G., Huntley, D. A., & Bowen, A. J. (1982). Two-dimensional surf beat:  
 498 Long wave generation by a time-varying breakpoint. *Journal of Geophysical*  
 499 *Research*, 87(C1), 492–498. Retrieved from [http://onlinelibrary.wiley](http://onlinelibrary.wiley.com/doi/10.1029/JC087iC01p00492/full)  
 500 [.com/doi/10.1029/JC087iC01p00492/full](http://onlinelibrary.wiley.com/doi/10.1029/JC087iC01p00492/full)  
 501 [http://doi.wiley.com/10.1029/](http://doi.wiley.com/10.1029/JC087iC01p00492)  
 502 JC087iC01p00492 doi: 10.1029/JC087iC01p00492
- 503 Thotagamuwage, D. T., & Pattiaratchi, C. B. (2014, 9). Influence of off-  
 504 shore topography on infragravity period oscillations in Two Rocks Marina,  
 505 Western Australia. *Coastal Engineering*, 91, 220–230. Retrieved from  
 506 <http://www.sciencedirect.com/science/article/pii/S0378383914001197>  
 507 doi: 10.1016/j.coastaleng.2014.05.011
- 508 Van der Molen, W., Monardez, P., & Van Dongeren, A. (2006, 3). Numeri-  
 509 cal simulation of long-period waves and ship motions in Tomakomai port,  
 510 Japan. *Coastal Engineering Journal*, 48(1), 59–79. Retrieved from  
 511 <http://www.worldscientific.com/doi/abs/10.1142/S0578563406001301>  
 512 doi: 10.1142/S0578563406001301
- 513 van der Molen, W., Scott, D., Taylor, D., & Elliott, T. (2016). Improvement of  
 514 mooring configurations in Geraldton harbour. *Journal of Marine Science and*  
 515 *Engineering*, 4(1), 1–20. doi: 10.3390/jmse4010003
- 516 Van Dongeren, A., Battjes, J., Janssen, T., van Noorloos, J., Steenhauer, K.,  
 517 Steenbergen, G., & Reniers, A. (2007, 2). Shoaling and shoreline dissipa-  
 tion of low-frequency waves. *Journal of Geophysical Research*, 112(C2).

- 518 Retrieved from <http://doi.wiley.com/10.1029/2006JC003701> doi:  
519 10.1029/2006JC003701
- 520 Van Thiel de Vries, J. S. M., Van Gent, M. R. A., Walstra, D. J. R., & Reniers,  
521 A. J. H. M. (2008, 12). Analysis of dune erosion processes in large-scale  
522 flume experiments. *Coastal Engineering*, *55*(12), 1028–1040. Retrieved from  
523 <http://linkinghub.elsevier.com/retrieve/pii/S0378383908000860> doi:  
524 10.1016/j.coastaleng.2008.04.004
- 525 Vrećica, T., Soffer, R., & Toledo, Y. (2019, 8). Infragravity Wave Gener-  
526 ation by Wind Gusts. *Geophysical Research Letters*, *46*(16), 9728–  
527 9738. Retrieved from [https://onlinelibrary.wiley.com/doi/abs/](https://onlinelibrary.wiley.com/doi/abs/10.1029/2019GL084241)  
528 [10.1029/2019GL084241](https://onlinelibrary.wiley.com/doi/10.1029/2019GL084241)  
529 [https://onlinelibrary.wiley.com/doi/10.1029/](https://onlinelibrary.wiley.com/doi/10.1029/2019GL084241)  
530 [2019GL084241](https://onlinelibrary.wiley.com/doi/10.1029/2019GL084241) doi: 10.1029/2019GL084241
- 531 Webb, S. C. (2007). The Earth's 'hum' is driven by ocean waves over the continental  
532 shelves. *Nature*, *445*(7129), 754–756. doi: 10.1038/nature05536
- 533 Webb, S. C., Zhang, X., & Crawford, W. (1991). Infragravity waves in the deep  
534 ocean. *Journal of Geophysical Research*, *96*(C2), 2723–2736. Retrieved from  
535 <http://doi.wiley.com/10.1029/90JC02212> doi: 10.1029/90JC02212
- 536 Zijlema, M., Van Vledder, G. P., & Holthuijsen, L. H. (2012). Bottom friction  
537 and wind drag for wave models. *Coastal Engineering*, *65*, 19–26. Retrieved  
538 from <http://dx.doi.org/10.1016/j.coastaleng.2012.03.002> doi:  
10.1016/j.coastaleng.2012.03.002

Figures

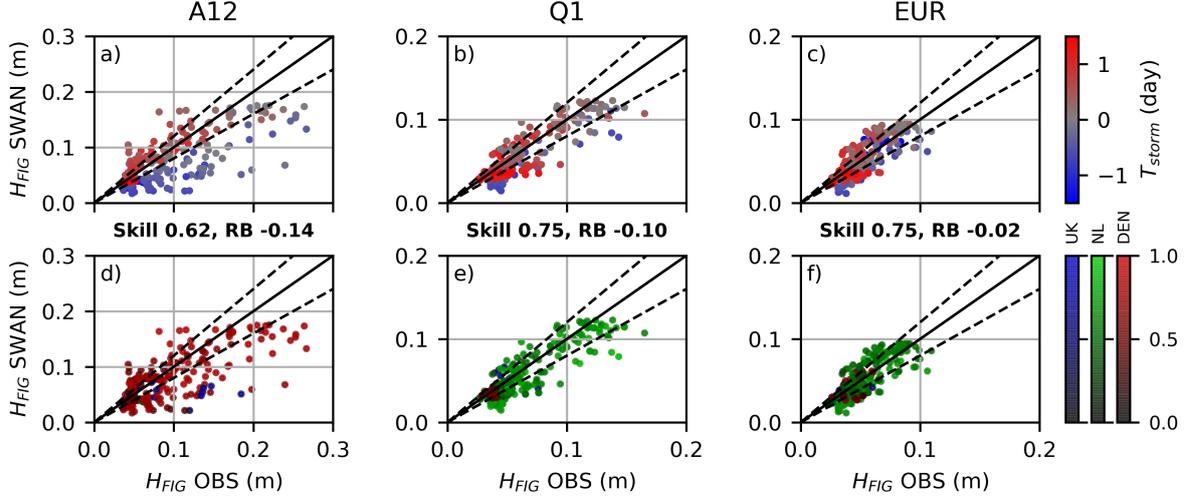


**Figure 1.** Infragravity wave conditions in the North Sea during storm Xaver (5-8 Dec 2013). Snapshot of the instantaneous significant SS wave height  $H_{SS}$  from the ERA5 reanalysis (panel a) and FIG wave height  $H_{FIG}$  from the SWAN model (panel b) at 6 Dec 06:00:00 UTC. Observed (lines) and ERA5 (markers) time-series of  $H_{SS}$  (full line and filled markers, left axis) and  $T_{m0,-1}$  (dashed lines and open markers, right axis) at three stations in the southern North Sea (A12, Q1 and EUR) (panel c). Observed (full lines) and SWAN (markers) FIG wave heights  $H_{FIG}$  at A12, Q1 and EUR (panel d). The dashed lines in panel d represent the observations of the total IG wave height (including both bound and free contributions). Relative contribution of different shorelines to the SWAN predicted FIG variance ( $\propto H_{FIG}^2$ ) at the three stations (panel e-g).

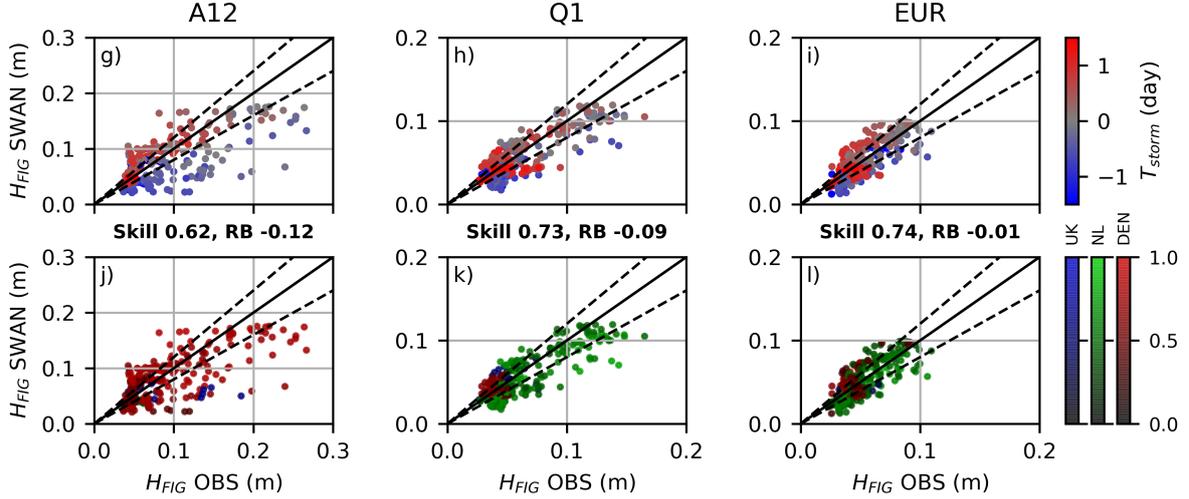


**Figure 2.** Infragravity wave conditions in the North Sea prior to and during storm Egon (13-15 Jan 2017). Refer to the caption of Fig. 1 for further details.

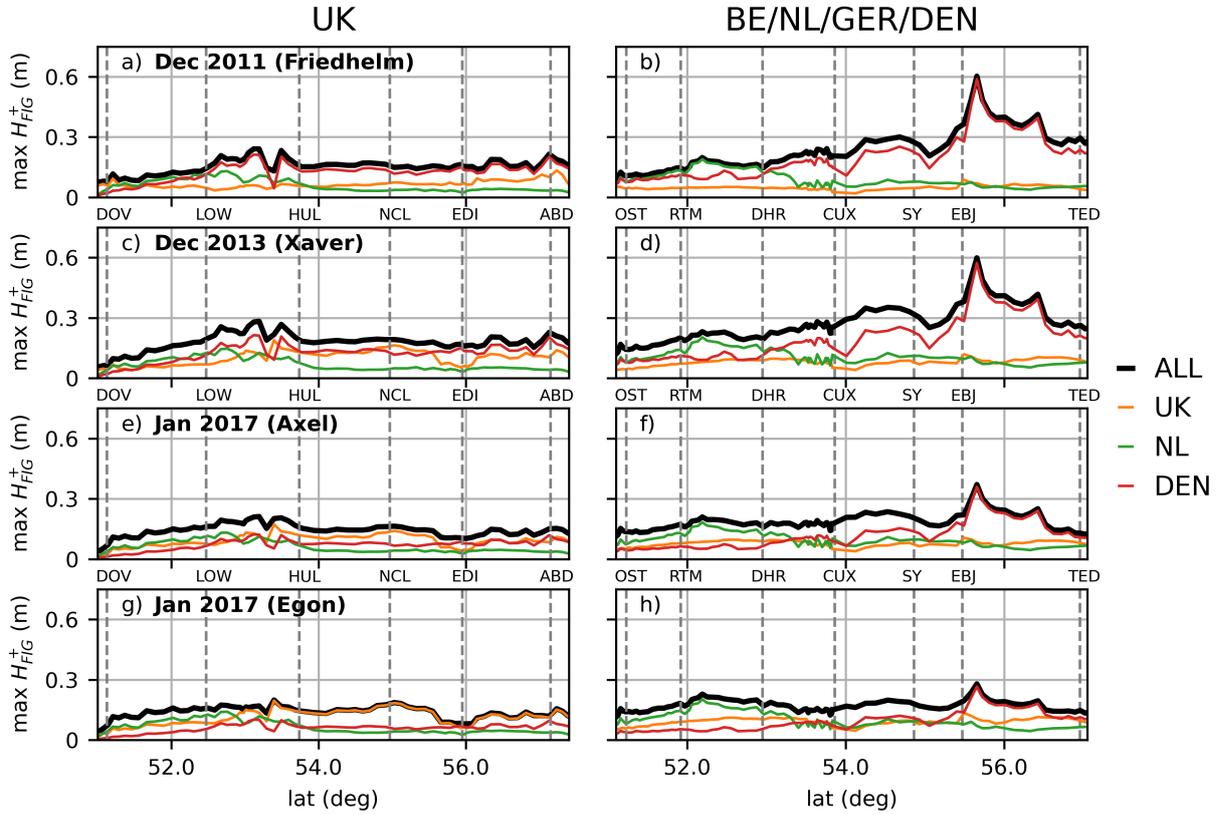
**including bottom friction** ( $\alpha_1 = 18.0 \times 10^{-4} \text{ s}^{-1}$ ,  $\chi = 0.01 \text{ m}^2$ )



**no bottom friction** ( $\alpha_1 = 14.4 \times 10^{-4} \text{ s}^{-1}$ ,  $\chi = 0. \text{ m}^2$ )



**Figure 3.** Scatter plot of the instantaneous observed (OBS) versus modelled (SWAN) FIG wave height  $H_{FIG}$  at station A12 (left panels), Q1 (middle panels) and EUR (right panels). The top plots (panel a-f) show the modelling results excluding the effect of bottom friction, and the bottom plots (panel g-l) show the results including bottom friction. In panels (a-c) and (g-i), the marker colour indicates the time relative to the approximate peak in the observed  $H_{FIG}$  at A12. In panels (d-f) and (j-l), the marker colour indicates the region from which the majority of the IG variance ( $\propto H_{FIG}^2$ ) originated. The colour shading indicates the relative contribution of this region to the total IG variance, with darker colors indicative for significant contributions from multiple regions. Two statistical parameters to quantify the model-data comparison are shown for all stations. The model skill is computed as  $\text{Skill} = 1 - \frac{\sqrt{\frac{1}{N} \sum (H_{FIG}^o - H_{FIG}^c)^2}}{\sqrt{\frac{1}{N} \sum (H_{FIG}^o)^2}}$ , and the relative bias is computed as  $\text{RB} = \frac{\sum (H_{FIG}^c - H_{FIG}^o)}{\sum H_{FIG}^o}$ , where  $N$  is the total number of samples and superscripts  $o$  and  $c$  indicate an observed and computed variable, respectively.



**Figure 4.** Maximum of the SWAN modelled (including bottom friction) shoreward directed FIG wave height  $H_{FIG}^+$  (black line) during individual storm events along the Eastern UK coastline (left panels) and the coastlines between the Belgium and Denmark (right panels) at approximately 20 m water depth. The coloured lines indicate the  $H_{FIG}^+$  that was radiated from the UK, Dutch (NL) or Danish (DEN) coast (as indicated by the legend). The vertical dashed lines indicate geographical locations along the coastline. Left panels: DOV (Dover), LOW (Lowestoft), Hul (Hull), NCL (Newcastle), EDI (Edinburgh), ABD (Aberdeen); Right panels: OST (Ostend), RTM (Rotterdam), DHR (Den Helder), CUX (Cuxhaven), SY (Sylt), EBJ (Esbjerg), TED (Thisted).

# Supporting Information for ”Free infragravity waves in the North Sea”

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## Contents of this file

Figures S1 to S2

## Introduction

This supporting information provides the details and results of the model sensitivity study (spatial and directional resolution) and the model calibration study.

## Model sensitivity study (spatial and directional resolution)

A sensitivity study was conducted to determine the optimal spatial and directional resolution of the SWAN simulations. A stationary simulation (with the 1<sup>st</sup> order BSBT propagation scheme) was ran with all shorelines along the North Sea radiating an equal amount of IG energy. The grid and directional resolution were varied and compared with a reference simulation. The reference simulation was run with the 2<sup>nd</sup> order SORDUP propagation scheme and with the finest considered spatial and directional resolution. The

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directional resolution was varied between  $2^\circ - 15^\circ$  and the longitudinal grid resolution was varied between  $0.01 - 0.05^\circ$ . The latitudinal resolution was kept as a constant fraction ( $\approx 1.5$ ) of the longitudinal resolution. For each simulation, we compute the model skill relative to the reference simulation for the free infragravity wave height  $H_{FIG}$  at the three measurement stations ( $N = 3$ ),

$$Skill = 1 - \frac{\sqrt{\frac{1}{N} \sum (H_{FIG} - H_{FIG,R})^2}}{\sqrt{\frac{1}{N} \sum H_{FIG,R}^2}}, \quad (1)$$

where subscript  $R$  is indicative for the reference simulation.

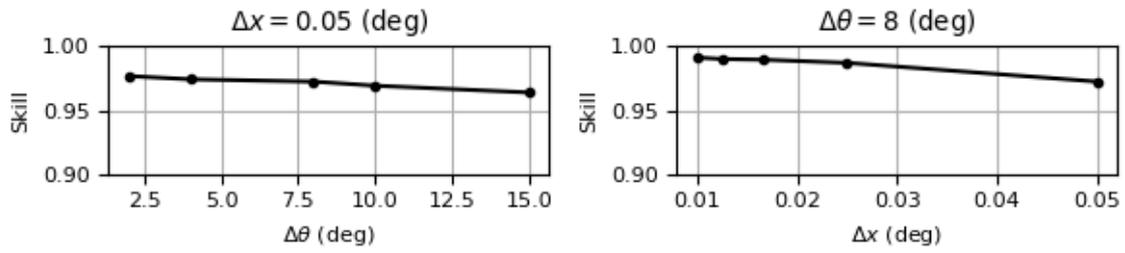
This sensitivity study showed that the model results were only weakly affected by the directional and grid resolution (Fig. S1). Based on these results, the model simulations were conducted with a resolution of  $0.025^\circ$  in longitudinal and  $0.0165^\circ$  in latitudinal direction, and a directional resolution  $\Delta\theta = 8^\circ$ .

### Model calibration study ( $\alpha_1$ and $\chi$ )

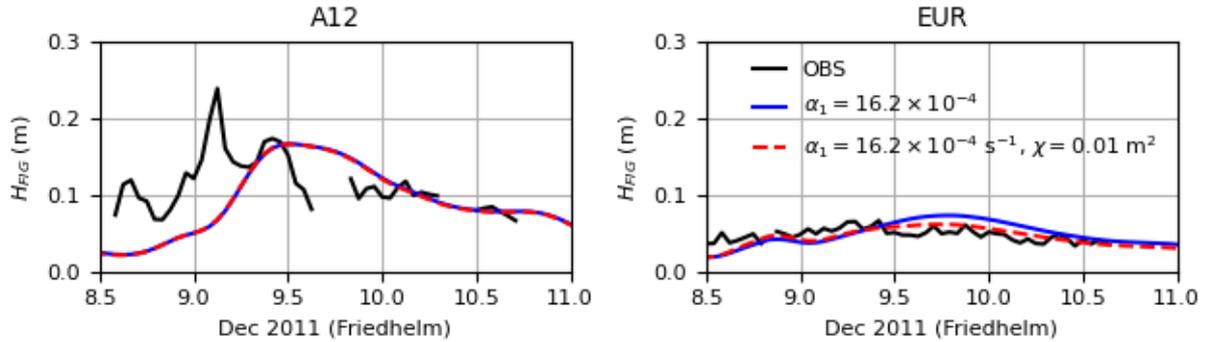
Following the model sensitivity study, a calibration study was conducted to decide the  $\alpha_1$  parameter of the infragravity source term, and the  $\chi$  parameter of the JONSWAP friction formulation. For this purpose, we conducted a non-stationary simulation of storm Friedhelm. During storm Friedhelm, significant sea-swell waves mainly made landfall in Denmark, and modelled  $H_{FIG}$  in the North Sea were dominated by infragravity waves radiated from the Danish coast. This allowed for a more straightforward model calibration compared to a storm during which  $H_{FIG}$  originated from multiple coastal sections.

First, we selected the  $\alpha_1$  parameter based on the model-data agreement at station A12 for the simulation without bottom friction. This resulted in  $\alpha_1 = 14.4 \times 10^{-4} \text{ s}^{-1}$ . Without friction,  $H_{FIG}$  is slightly over predicted at station EUR (Fig. S2). Subsequently,

we varied  $\alpha_1$  and  $\chi$  to 1) match the predictions without bottom friction at A12, and 2) improve the model-data agreement at station EUR. This resulted in  $\alpha_1 = 18 \times 10^{-4} \text{ s}^{-1}$  and  $\chi = 0.01 \text{ m}^2$ . As the model-data agreement with these model parameters was satisfactory for the other storms, we made no attempt to optimize the model results by calibrating  $\alpha_1$  for separate coastal sections.



**Figure S1.** Model skill for a varying directional resolution ( $\Delta\theta$ ) with a fixed grid resolution (with  $\Delta x$  the longitudinal resolution) (left panel), and for a varying grid resolution with a fixed directional resolution (right panel).



**Figure S2.** Observed (black line) and predicted (excluding friction, blue line; and including friction, red line) infragravity wave height  $H_{FIG}$  at station A12 in the central North Sea (left panel) and station EUR in the southern North Sea (right panel) during storm Friedhelm.