

High-resolution ambient noise imaging of geothermal reservoir using 3C dense seismic nodal array and ultra-short observation

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November 22, 2022

Abstract

Tomographic imaging based on long-term ambient seismic noise measurements, mainly the phase information from surface waves, has been shown to be a powerful tool for geothermal reservoir imaging and monitoring. In this study, we utilize seismic noise data from a dense nodal array (192 3C nodes within 20km²) over a ultra-short observation period (4.7 days) to reconstruct surface waves and determine the high-resolution (0.2km) three-dimensional (3-D) S wave velocity structure beneath a rural town in Zhejiang, China. We report the advantage of cross-coherence over cross-correlation in suppressing pseudo-arrivals caused by persistent sources. We use ambient noise interferometry to retrieve high quality Rayleigh waves and Love waves. Body waves are also observed on the R-R component interferograms. We apply phase velocity dispersion measurements on both Rayleigh waves and Love waves and automatically pick more than 23,000 dispersion curves by using a Machine Learning technique. 3-D surface wave tomographic results after depth inversion indicate low-velocity anomalies (between -1% and -4%) from the surface to 2km depth in the central area. Combined with the conductive characteristics observed on resistivity profile, the low-velocity anomalies are inferred to be a fluid saturated zone of highly fractured rock. Joint interpretation based on HVSR measurements, and existing temperature and fluid resistivity records observed in a nearby well, suggests the existence of the high-temperature geothermal field through the fracture channel. Strong correlation between HVSR measurements and S wave velocity model sheds light on the potential of extraction of both amplitude and phase information from ambient noise.

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Key Points:

- Ambient noise data have been recorded using a dense nodal array (192 3C nodes within 20km^2) over ultra-short observation period (4.7 days)
- Both surface waves (Rayleigh and Love waves) and P waves are identified in the cross-coherence functions
- S wave velocity model is consistent with existing geophysical data and suggests the existence of high-temperature geothermal resources at depth

Abstract

Tomographic imaging based on long-term ambient seismic noise measurements, mainly the phase information from surface waves, has been shown to be a powerful tool for geothermal reservoir imaging and monitoring. In this study, we utilize seismic noise data from a dense nodal array (192 3C nodes within 20km^2) over a ultra-short observation period (4.7 days) to reconstruct surface waves and determine the high-resolution (0.2km) three-dimensional (3-D) S wave velocity structure beneath a rural town in Zhejiang, China. We report the advantage of cross-coherence over cross-correlation in suppressing pseudo-arrivals caused by persistent sources. We use ambient noise interferometry to retrieve high quality Rayleigh waves and Love waves. Body waves are also observed on the R-R component interferograms. We apply phase velocity dispersion measurements on both Rayleigh waves and Love waves and automatically pick more than 23,000 dispersion curves by using a Machine Learning technique. 3-D surface wave tomographic results after depth inversion indicate low-velocity anomalies (between -1% and -4%) from the surface to 2 km depth in the central area. Combined with the conductive characteristics observed on resistivity profile, the low-velocity anomalies are inferred to be a fluid saturated zone of highly fractured rock. Joint interpretation based on HVSR measurements, and existing temperature and fluid resistivity records observed in a nearby well, suggests the existence of the high-temperature geothermal field through the fracture channel. Strong correlation between HVSR measurements and S wave velocity model sheds light on the potential of extraction of both amplitude and phase information from ambient noise.

1 Introduction

Geothermal energy is one of the most promising renewable energy sources, particularly within the context of China’s energy structure optimization, environmental protection measures, energy conservation, and rising pressure on emission reduction. By the end of 2020, renewable energy facilities, including solar, wind, geothermal and other types of energy, in China will supply 27% of total power generation, according to the government’s 2016-2020 plan for renewable energy. However, geothermal resources accounted for only 0.6% of the total energy consumption in 2019 (Liu et al., 2019). Therefore, significant work is required for the development of national geothermal resources. Geothermal energy production converts heat energy stored in the Earth into energy forms useful for humans; in most implementations, geothermal energy production is clean, sustainable, and can provide baseload capacity to regional power grids (Tomac & Sauter, 2018). Geothermal energy resources can be classified into two types: shallow geothermal and deep geothermal resources (Ganguly & Kumar, 2012). Shallow geothermal energy is often tapped in the form of hot water or steam (e.g. hydrothermal production), while deep geothermal energy often takes the form of ”hot dry rock” resources that usually exists at depths greater than 3–5 km beneath the Earth’s surface (Rubio-Maya et al., 2015; Xie et al., 2020). China has enormous geothermal resource potential, however, low-temperature geothermal resources are more common than high-temperature ones. The high-temperature geothermal resources are located in the marginal zone of the plate with an abnormal tectonic activity, e.g., Himalayan and Taiwan geothermal belts, (Zhang et al., 2019). The low- and medium-temperature geothermal resources are mostly located in uplifted mountain-type and sedimentary basin areas within the plates (Long et al., 2015). Geothermal resources distributed in mountain fault zones are generally quite small in scale (Wang et al., 2017). Therefore, evaluation and utilization of low/medium-temperature geothermal energy dependent on the high-resolution geothermal reservoir imaging techniques.

Geothermal systems often give distinctive and fairly easily measured discontinuities in physical properties (e.g., high heat flow, low electrical resistivity, attenuation of high frequency elastic waves), and geophysical methods play a key role in geothermal reservoir exploration (Combs, 1978). For examples, a gravity survey can be used to study the

74 depth of fill in intermontaine valleys, and to locate intrusive masses of rock (e.g., San-
 75 tos & Rivas, 2009; Atef et al., 2016); magnetic surveys can be used to identify the bound-
 76 aries to the flows in volcanic areas (e.g., Hochstein & Soengkono, 1997; Zaher et al., 2018);
 77 a combination of resistivity studies and heat flow determinations is advisable to search
 78 for zones of fracture permeability in the reservoir (e.g., Wright et al., 1985; Thanassoulas,
 79 1991; Munoz, 2014); a seismic reflection survey can be used where there is a bedded struc-
 80 ture to the subsurface to allow the recognition of faults by the disruption of the conti-
 81 nuity of the bedding (e.g., Lüschen et al., 2011); a microseismic survey is also a widely
 82 used tool for studying activity on fracture zones in a prospect area since high temper-
 83 ature hydrothermal areas are characterized by a relatively high level of microearthquake
 84 activity (e.g., Ward, 1972; Combs & Hadley, 1977; Obermann et al., 2015). However, no
 85 one exploration technique is likely to be universally effective in defining a geothermal reser-
 86 voir. Some methods lack the maturity of development to be used effectively under dif-
 87 ficult conditions, while others become less useful for deep exploration because of lack of
 88 sensitivity (Keller, 1981). Considering the limitations of various methods, it is proba-
 89 bly necessary to employ a wide variety of techniques.

90 Over the last decade, ambient noise interferometry techniques have found a vari-
 91 ety of applications for geothermal reservoir imaging (e.g., Tibuleac et al., 2009; Tibuleac
 92 & Eneva, 2011; Obermann et al., 2015; Lehujeur et al., 2016, 2018; Spica et al., 2018;
 93 Martins et al., 2019, 2020; Planès et al., 2020). Compared to relatively expensive active
 94 seismic imaging methods, ambient noise imaging is a passive and low-cost approach. Fol-
 95 lowing the pioneering work of Campillo and Paul (2003), ambient noise interferometry
 96 can be used to estimate an approximate Green’s function between two receivers by cross-
 97 correlating the ambient seismic wave field (Shapiro & Campillo, 2004; Snieder, 2004; Wape-
 98 naar, 2004a; Bensen et al., 2007; Snieder et al., 2009; Nakata et al., 2015). This approach
 99 has been applied to characterize multiple scales of earth structure: from the global scale
 100 or continental scale deep-structure imaging in seismology (e.g., Yang et al., 2007; Lin et
 101 al., 2008; Yao & van der Hilst, 2009; Lin et al., 2009; Strobbia & Cassiani, 2011) to lo-
 102 cal scale exploration (e.g., Bakulin & Calvert, 2006; Wapenaar et al., 2008; Draganov
 103 et al., 2009; Nakata et al., 2011; Ali et al., 2013; Behm et al., 2014; Cheng et al., 2015,
 104 2016; Nakata et al., 2016; Cheng et al., 2018; Behm et al., 2019; Castellanos et al., 2020).

105 To date, ambient noise interferometry is almost exclusively performed with surface
 106 waves tomography based on multiple-station (tens or more) networks and long-term (months
 107 or years) continuous observations (Lin et al., 2008; Martins et al., 2019; Planès et al.,
 108 2020). Here we investigate the potential of high-resolution (i.e., 0.2km) ambient noise
 109 imaging of geothermal reservoir using a dense seismic nodal array (i.e., 192 nodes within
 110 $20km^2$) over an ultra-short observation period (i.e., 4.7 days). In the following study, we
 111 present the acquisition and the main characteristics of the ambient seismic noise records
 112 obtained from a dense network deployed in a rural town in Zhejiang, China. We extract
 113 high quality Rayleigh waves and Love waves based on ambient noise interferometry, and
 114 automatically pick more than 23k phase velocity dispersion curves to allow three-dimensional
 115 (3-D) S wave velocity model construction. The model is finally discussed in the light of
 116 HVSR measurements and existing borehole records and resistivity surveys.

117 2 Area and Data

118 The area of investigation (Fig.1) stands on the eastern margin of the Jinqu basin
 119 in southeastern China, where the deep NE-SW Jiangshan-Shaoxing fault crosses the basin.
 120 The Jiangshan-Shaoxing fault is a major structural feature which traverses Zhejiang Province
 121 and divides it into two distinct geological zones, the northwest Yangtze paraplatform,
 122 a relatively stable tectonic area dominantly composed of sedimentary rocks ranging from
 123 the Sinian (Pre-Cambrian) System to the Lower Triassic Series, and the southeast South
 124 China fold system, which is commonly overlain by Yanshanian (Mesozoic to Cenozoic)
 125 volcanic and acid igneous rocks (Zhejiang, 1989). The fault itself has been active since

126 Proterozoic times when it was initiated (Ren, 1987). The pattern of heat flow in south-
 127 eastern China has been investigated by Hu and Wang (2000) and Yuan et al. (2006). A
 128 high heat flow of $75\text{-}80\text{mW}/\text{m}^2$ has been found in the Jinqu basin. Our survey region
 129 is centered in a rural town (Fig.1a), Andi, where surface hot water has been founded by
 130 residents in recent years.

131 Based on the 1:50000 geological map released by Zhejiang Geophysical and Geo-
 132 chemical Prospecting Academy (ZGGPA) in 2006 (Fig.1b), the survey zone mainly con-
 133 sists of the younger upper Jurassic system (J_3x , J_3d , J_3g) at the northern, central and
 134 southwestern region, and the older PreSinian system ($AnZch^2$) at the south. The allu-
 135 vial deposits from the Quaternary system (Q_4) split the area. The Plum Creek River
 136 (the north blue line on Fig.1a) starts from north and crosses through the town area be-
 137 fore it reaches a water reservoir (outlined by the blue lines on Fig.1a) at the southwest.
 138 Coarse-grained adamellites ($\eta\gamma$) are widely distributed in the south mountain areas and
 139 adamellite dykes are in unconformable contact with the banded biotite plagiogneiss of the
 140 PreSinian system; granite (ν) dykes are intruded in the fracture system by the rifting
 141 unconformity.

142 A total of 192 nodal seismic stations (Fairfield ZLand 3C 5Hz), as indicated by the
 143 triangles on Figure.1a, were deployed over Andi town with an average aperture around
 144 4.8km. The nodes recorded continuously from 12:30 pm, May 9th 2019 to 7:45 am, May
 145 14th 2019 (about 4.7 days) with the sampling frequency of 500Hz. The nodes were buried
 146 at 30 cm and coupled to the ground with 15-cm metallic spikes. The interstation dis-
 147 tances vary from the nearest 100m to the farthest 4.7km. In addition to the nodal seis-
 148 mic survey, one 2-km-long CSAMT (controlled-source audio-frequency magnetotellurics)
 149 profile was available (the red line in Fig.1) for reference, and one test well (the white
 150 cross in Fig.1) was drilled in 2016 by ZGGPA.

151 3 Methods

152 3.1 Noise characteristics

153 To understand the temporal and spectral characteristics of the observed ambient
 154 noise records, we employed spectral analysis on the raw waveforms. For each station, we
 155 first split the continuous noise records into 1-min-long time segments without overlaps.
 156 We computer the power spectral density (PSD, McNamara & Buland, 2004) of the raw
 157 waveforms, and stack all segments along the time direction to build a time-frequency spec-
 158 tra image. The PSD spectrograms within each one-hour epoch are averaged together as
 159 the spectrogram at the corresponding epoch. Note that we did not correct the absolute
 160 amplitude of the PSD with the instrument response. The obtained spectrograms present
 161 strong temporal and spatial amplitude variations. The PSD difference between daytime
 162 and nighttime is around 10dB in spatial average. Figure.2 displays three examples of the
 163 vertical component spectrograms at different locations (indicated by the magenta tri-
 164 angles on Fig.1a), north (a), central (b) and south (c). Compared with the central sta-
 165 tion in the rural town area, the north station (Fig.2a) shows stronger noise energy as well
 166 as higher dominant frequencies (around 20-30 Hz) because of the existence of several busy
 167 highways and express roads that connect the northern urban city, Jinhua, with the sur-
 168 rounding rural towns. As for the central station (Fig.2b), it shows dominant frequen-
 169 cies around 10Hz which is similar as that in urban area, and the distinct daily pattern
 170 which reflects much regular human activities in the peaceful rural town compared with
 171 that on the north station. Several long duration and very narrow-band signals, visible
 172 as horizontal lines or spikes (as indicated by the double arrow around 4.2 Hz), were also
 173 observed; these seismic waves are most probably excited by rotating machinery operat-
 174 ing at fixed frequencies, like electrical motors and gearboxes of industrial machinery (Plesinger
 175 & Wielandt, 1974; Groos & Ritter, 2009; Cheng et al., 2019). As for the south station
 176 located in the mountain area (Fig.2c), the PSD energy is generally 10dB lower than that

177 in the central town area, and the weak daily pattern indicates the observed noise energy
 178 from the distant human activities. Note that the strong energy around 35Hz in the day-
 179 time of May 13th (highlighted by the gray box on Fig.2c) is supposed to be the signa-
 180 ture from weather associated with rain and potentially thunder (Dean, 2017), and it is
 181 also consistent with the relative weak energy between 1 and 20Hz observed at the same
 182 duration on the central station (Fig.2b) which indicates less human activities affected
 183 by the rain weather. The similar seismic signature from weather has been successfully
 184 reported by Zhu and Stensrud (2019) by using a fiber-optic distributed acoustic sens-
 185 ing array.

186 We also apply beamforming analysis (Lacoss et al., 1969; Rost & Thomas, 2002;
 187 Gerstoft & Tanimoto, 2007) on the raw waveforms, to figure out the spatial distribution
 188 of the seismic noise sources, which is necessary for our further ambient noise interferom-
 189 etry work. Beamforming analysis presents the constructive summation of all signals shifted
 190 appropriately in the time or frequency domain for the matching azimuth (clockwise from
 191 the north) and slowness. Figure.3 displays the averaged beam energy at different frequency
 192 bands based on vertical component (the upper panels) and the north horizontal compo-
 193 nent (the bottom panels). Beam energy plots below 1.0Hz (Fig.3a1 and Fig.3b1) show
 194 distinct source energy from the southeastern direction with apparent velocity $>3\text{km/s}$,
 195 and we infer the source to be primary and secondary microseism i.e. nonlinear inter-
 196 actions of ocean waves (Ardhuin et al., 2015) with the southeastern coast of Zhejiang
 197 province (as indicated by the China map on Fig.1). For the frequency band between 1Hz
 198 and 5Hz, we observe distinct spectral energy peak from the north (Fig.3a2 and Fig.3b2)
 199 with an apparent velocity of $\sim 2.5\text{km/s}$; the noise source is likely surface waves gener-
 200 ated by activity 18 km to the north in the urban city of Jinhua as well as the northern
 201 traffic lines. For the higher frequency band between 5 and 10Hz, we observe almost an
 202 isotropic noise distribution in the beam domain (Fig.3a3 and Fig.3b3), except for the
 203 southeastern direction where the mountain area located. In general, it indicates a rela-
 204 tive homogeneous source distribution which is advantageous for ambient noise inter-
 205 ferometry (Weaver & Lobkis, 2001; Wapenaar, 2004b). It is worth noting that the weak
 206 beam energy with apparent velocity $>4\text{km/s}$ can be observed on the horizontal compo-
 207 nent with frequency $>1\text{Hz}$ (Fig.3b2 and Fig.3b3), as well as the vertical component with
 208 frequency $>5\text{Hz}$ (Fig.3a3). Beam energy with higher frequencies and higher velocities
 209 is likely associated with body waves which will be further discussed. Accounting for the
 210 body wave energy, we observed the horizontal component (Fig.3b3) presents the rela-
 211 tive stronger energy than the vertical component (Fig.3a3).

212 3.2 Ambient Noise Interferometry

213 We follow the workflow of Bensen et al. (2007) to preprocess the recorded noise prior
 214 to retrieval of surface waves by ambient noise interferometry. We first downsample the
 215 raw data to 100Hz; next we split the continuous noise waveforms into a series of 1-min-
 216 long time segments without overlaps, and remove means and trends of the short noise
 217 segment followed by tapering; next we utilize temporal normalization to attenuate ar-
 218 tifacts, e.g., near-field interferences and earthquakes, by using a running absolute mean
 219 filter (e.g. Bensen et al., 2007); finally, spectral normalization is utilized to extend the
 220 frequency band with a frequency-domain whitening approach, which computes the run-
 221 ning smoothed amplitude of complex Fourier spectrum as the whiten weights with a mov-
 222 ing window of 0.5% length of the frequency series.

223 Two main algorithms exist for empirical Green’s function (EGF) extraction from
 224 ambient noise, cross-correlation (Shapiro & Campillo, 2004; Stehly et al., 2006) and cross-
 225 coherence (Aki, 1957; Combs, 1978; Schuster et al., 2003; Nakata et al., 2011). Alter-
 226 native approaches include deconvolution (Vasconcelos & Snieder, 2008a, 2008b; Snieder
 227 et al., 2009) and multi-dimensional deconvolution (Wapenaar et al., 2008, 2011; Van Dalen
 228 et al., 2015; Weemstra et al., 2016; Cheng et al., 2017), both of which have been utilized

for seismic interferometry. In general, generation of EGF's using cross-correlation is the simplest and currently most popular technique with numerous examples of successful field application. The cross-coherence algorithm is also referred to as whitened cross-correlation; Prieto et al. (2009) demonstrates performing cross-correlation with spectral whitening is equivalent to calculating the cross-coherence. However, the choice of spectral whitening approach and the corresponding parameters can yield differences in the extracted EGFs. Figure.4 presents a comparison between cross-correlation and cross-coherence with the same preprocessed noise waveforms. We observe distinct pseudo-arrivals existing on extracted cross-correlation functions (Fig.4a and Fig.4b). These may be caused by insufficient spectral normalization during the data preprocessing procedure. Figure.4c shows the spectral difference between the averaged cross-correlation functions and the cross-coherence functions; several distinct spikes observed in the cross-correlation functions have been significantly attenuated after the further spectral normalization included in the cross-coherence algorithm. These kinds of pseudo-arrivals are almost inevitable since selection of the appropriate data preprocessing workflow requires substantial manual tuning. However, the existence of the pseudo-arrivals could mislead interpreters, particularly for coda wave interferometry, since they could be mistaken for coda waves while not encoding any subsurface information. In fact, they only reflect the seismic signatures associated with some specific sources, for example, the narrow-band persistent source signatures observed on the spectrograms (indicated by the double arrow on Fig.2b) with peak frequencies around 3, 4.2, 5.5, 7.5 Hz which are consistent with the spikes presenting on the cross-correlation spectrum (Fig.4c). Therefore, we recommend the use of cross-coherence for ambient noise interferometry. With pseudo-arrivals removal, the cross-coherence functions show much cleaner virtual-source gather (see Fig.S1 in the supporting information) with higher signal to noise ratio (SNR) (see Fig.S2 in the supporting information).

3.3 Surface Waves and Body Waves

Due to limitation from the single-component instrument as well as the poor data quality on horizontal components, most ambient noise interferometry studies focus primarily on the vertical component, accounting for Rayleigh waves retrieval, even with months- or years- duration time observations (Lehuteur et al., 2018; Martins et al., 2019; Planès et al., 2020). In order to go beyond the retrieval of Rayleigh waves from ambient noise interferometry, we apply cross-coherence on both vertical component (Z-Z) and horizontal components (NN, EE, NE, EN). After cross-coherence, we employ phase-weighted stacking on the 4.7 days cross-coherence functions to further improve the coherence signals (Schimmel & Paulssen, 1997; Schimmel et al., 2011; Ventosa et al., 2017), and rotate the north and east horizontal components into radial (R) and transverse (T) components (Lin et al., 2008). Finally, we obtain $3 * C_N^2 = 3 * 18336$ high quality cross-coherence functions for Z-Z, R-R and T-T components.

Figure.5 presents the bin-stacked cross-coherence gathers from Z-Z (a), R-R (b) and T-T (c) components by stacking all available interstation cross-coherence pairs in a 70m offset bin. A bandpass filter between 1 and 10Hz has been applied for better display. Clear Rayleigh waves with apparent velocity around 2.5 km/s can be observed on both Z-Z and R-R components; high quality Love waves with slightly higher apparent velocity around 2.7 km/s are also distinct on the T-T component. Moreover, we can also distinguish body waves with apparent velocity around 4.2 km/s on the whole offset range of R-R component (highlighted on Fig.5b), as well as on the near-offset (<2 km) sections of Z-Z and T-T components.

We applied dispersion analysis on the bin-stacked virtual-sources gathers by using a frequency-domain slant-stacking technique that has been frequently utilized for multi-channel analysis of surface wave (MASW Park et al., 1998). For clearer presentation, all dispersion images in this work have been normalized along the frequency direction. Figure.6

281 displays the obtained dispersion spectra from Z-Z (a), R-R (b) and T-T (c) components.
 282 A clear dispersive energy trend can be observed for Rayleigh waves on Figure.6a and Figure.6b,
 283 and Love waves on Figure.6c. We are able to pick continuous dispersion curves from 1Hz
 284 to 10Hz for both Rayleigh and Love wave, and the picked dispersion curves (see Fig.S3
 285 in the supporting information) can be taken as reference for later two-station surface wave
 286 dispersion analysis.

287 Compared with the surface wave dispersion energy, the non-dispersive energy trend
 288 with higher frequencies and higher velocities suggests the presence of body waves. It is
 289 in good agreement with the observation of the high frequency higher velocity (>4 km/s)
 290 beam energy seen in our previous beamforming analysis (Fig.3). Both virtual-sources
 291 gathers (Fig.5) and dispersion spectra (Fig.6) illustrate that the body-wave energy is dom-
 292 inant on the radial (R-R) component. For typical velocity structures (e.g. velocity in-
 293 creasing with depth), P-waves at larger offsets should be stronger on the vertical com-
 294 ponent due to the bending of the upcoming waves towards the vertical. However, a sonic
 295 log from the center of the area (see Fig.S4 in the supporting information) indicates the
 296 existence of a thin hard (high-velocity) surface layer, resulting in a significant velocity
 297 decrease with depth. Consequently, upcoming waves will be bent away from the verti-
 298 cal, and P-wave energy can be strong on the radial component. We use the ray tracing
 299 code ANRAY (Gajewski & Pšencík, 1987; Gajewski & Psencik, 1989) to model travel
 300 times and amplitudes based on the 1D velocity model from the sonic log. It shows that
 301 the shallow high-velocity layer leads to significant bending of the raypath (Fig.7a), and
 302 that the presumed body-wave moveout in virtual source gathers fits well with the cal-
 303 culated travel times (Fig.7b). The radial component of the modeled P-wave amplitudes
 304 is significantly stronger than the vertical component (Fig.7c), confirming the assump-
 305 tion of observation of P-wave energy on the radial component of the interferograms. These
 306 observations suggest that double-beamforming techniques might be useful for isolating
 307 the body wave energy from the ambient noise field and enhancing P first arrivals for body
 308 wave tomography (Nakata et al., 2016; Castellanos et al., 2020).

309 3.4 Phase Velocity Dispersion Measurement

310 Recent ambient noise tomography applications for geothermal reservoir imaging
 311 focus on measurement of Rayleigh wave group velocities (Lehuteur et al., 2016, 2018; Planès
 312 et al., 2020), probably because of the directivity bias on phase velocity estimation from
 313 the inhomogeneous source distribution (Lin et al., 2008). However, the phase velocity
 314 measurements have the advantages of less uncertainty and higher depth sensitivity over
 315 the group velocity measurements. Based on the beamforming analysis and the perfect
 316 symmetry between the negative and positive time lags of the obtained interferograms,
 317 we believe the source distribution in Andi town is able to provide sufficient illumination
 318 for complete EGFs retrieval as well as dispersion measurements, at least for the frequency
 319 band between 1Hz and 10Hz.

320 We employ the an image transformation technique introduced by Yao et al. (2006)
 321 for phase velocity estimation based on the extracted EGFs. Considering the higher qual-
 322 ity of the retrieved Rayleigh waves on Z-Z component over that on R-R component (see
 323 Fig.S5 in the supporting information), we choose the Z-Z component EGFs for Rayleigh
 324 waves phase velocity estimation. The T-T component EGFs are used for Love waves phase
 325 velocity estimation.

326 In order to ensure the quality of dispersion measurements, we set a series of cri-
 327 teria for quality control:

- 328 1. we reject interstation pairs with distance <0.6 km to avoid the potential near-field
 329 effects on the dispersion measurement (Yoon & Rix, 2009; Foti et al., 2018);

- 330 2. we reject EGFs with $\text{SNR} < 7$ although most of EGFs show $\text{SNR} > 10$ (see Fig.S5
 331 in the supporting information);
 332 3. we apply velocity filter on extracted EGFs and mute arrivals beyond the veloc-
 333 ity band from 1.5 km/s to 3.5 km/s considering the local velocity range;
 334 4. we set the interstation distance has to be longer than 1.5 times wavelength;
 335 5. we limit the frequency band of dispersion curves above 1Hz to ensure the appro-
 336 priate illumination.

337 As for the wavelength criterion, it usually varies with the data as well as the nec-
 338 essaries change. Bensen et al. (2007) suggests a strict criterion with 3 times wavelength
 339 accounting for the far-field approximation; others choose a criterion with 1.5 times wave-
 340 length (e.g., Mordret et al., 2015; Obermann et al., 2016; Fallahi et al., 2017); Luo et al.
 341 (2015) demonstrates that one wavelength is still consistent with and also reliable as these
 342 with stricter wavelength criterion. In this work, we choose the 1.5 times wavelength cri-
 343 terion in order to get rid of the potential directional noise effects although our high qual-
 344 ity EGFs allow us to go beyond the 1.5 times wavelength.

345 Figure.8 displays examples of dispersion analysis on extracted Rayleigh and Love
 346 wave by using the an image transformation technique. The red waveforms present the
 347 velocity filtered EGFs used for Rayleigh and Love wave dispersion measurements, sep-
 348 arately. Clear fundamental modes can be observed on the obtained dispersion spectra.
 349 The spectral energy besides the fundamental modes indicates the 2π ambiguity caused
 350 by phase velocity measurement (Yao et al., 2006). We overly the dispersion spectra with
 351 the averaged dispersion curves picked from the bin-stacked virtual-source gathers for ref-
 352 erence. The blue dash lines indicate the 1.5 times wavelength criterion.

353 3.5 Dispersion Curves Picking using Machine Learning

354 After dispersion measurements for both Rayleigh and Love waves, noisy dispersion
 355 images are manually rejected by visual inspection. Finally, a task pool with more than
 356 23k dispersion measurements is prepared for dispersion curves picking. We utilize a deep
 357 learning model, named DCNet developed by Dai et al. (2020), for full automatic disper-
 358 sion curve picking by regrading dispersion curves extracted as an instance segmentation
 359 task. To help the machine to distinguish the target dispersion curves in this work, we
 360 set a confidence region based on the reference dispersion curves picked from the bin-stacked
 361 virtual-source gathers. First, we smooth the reference dispersion curves v_{ref} by linear
 362 regression; next, we build the upper and bottom boundaries of the confidence region with
 363 an extreme 25% velocity variation, $v_{upper} = 1.25 * v_{ref}$ and $v_{bottom} = 0.75 * v_{ref}$. The
 364 dispersion spectra beyond the confidence region has masked. We manually pick 1% dis-
 365 persion curves which are randomly selected from the task pool, and the high cross-correlation
 366 (95.32%) between the manually picked dispersion curves and the ML picked dispersion
 367 curves indicates the high quality of the automatically dispersion curves picking.

368 Figure.9 shows examples of dispersion curves picking by using the Machine Learn-
 369 ing (ML) technique. The excellent match between the ML picked dispersion curves (the
 370 cyan curves) and the manually picked dispersion curves (the magenta curves) demon-
 371 strates the accuracy of the ML picks. Figure.10 presents all the automatically picked dis-
 372 persion curves for Rayleigh waves (a) and Love waves (b) between 1 and 10Hz. The ma-
 373 genta curves indicate the smoothed reference dispersion curves, and the green dash lines
 374 highlight the confidence region defined by the smoothed reference dispersion curves. Fil-
 375 tered by a series of quality control criteria, we obtain 12,593 fundamental dispersion curves
 376 for Rayleigh waves and 11,105 fundamental dispersion curves for Love waves from $2 * C_N^2 = 2 * 18336$ EGFs, with a data utilization coefficient of 64.6%. See Fig.S6 for the
 377 distribution of the interstation distance as well as SNR for the picked dispersion curves
 378 in the supporting information.
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3.6 Surface Wave Tomography

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Based on the picked dispersion curves for both Rayleigh waves and Love waves, we construct two-dimensional (2-D) phase velocity distribution maps for a series of frequencies from 1Hz to 10Hz using a non-linear 2-D tomographic inversion technique (Rawlinson, 2005; Rawlinson & Sambridge, 2005). The inversion algorithm includes an eikonal solver based on the fast marching method (FMM, Rawlinson & Sambridge, 2005) for ray tracking and a subspace inversion scheme for the iterative inversion steps (Kennett et al., 1988). It accounts for propagation effects caused by rapid changes in the velocity field, and allows both smoothing and damping regularization to be imposed in order to address the problem of solution non-uniqueness. To tune the value of the two regularization parameters, we apply the L-curve approach (Hansen, 1992) to coordinate the traveltimes misfit and model variance as well as model roughness and define the optimal parameters.

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An adequate model resolution can help to identify subsurface anomalies' geometries, which is relevant for subsurface characteristics and geothermal purposes. We employ checkerboard sensitivity tests (Lévéque et al., 1993) to check the ability of the inversion algorithm to reconstruct structure at different locations in model space. Based on the picked dispersion curves, a background velocity of 2.5 km/s with $\pm 10\%$ perturbations has been targeted for forward modeling. Since in our case Rayleigh waves picks have a better raypath distribution compared with Love waves (see Fig.S7 and Fig.S8 in the supporting information), we check the geometry limitation as well as spatial resolution based on the Love wave observations. A series of checkerboard with different spatial resolutions ranging from 0.1 km to 0.5 km have been reproduced for each frequency because the numbers and spatial densities of the raypaths vary with frequencies. Figure.11 presents the simulated checkerboard models and tomographic results with two different grid sizes, 0.3 km (left panels) and 0.2 km (right panels). A series of simulated tests demonstrate that the inversion should allow us to estimate a spatial resolution of around $0.2\text{km} \sim 0.3\text{km}$ with higher resolution in the center of our seismic network and lower resolution towards the border of our investigation area. We define a resolved zone with the raypath density of Love wave (see Fig.S7 in the supporting information) greater than 60 per $0.3 \times 0.3 \text{km}^2$ cell grid at median frequency 5.0Hz where the simulated tests with grid sizes of both 0.3 km and 0.2 km can be well recovered, as indicated by the black curves on Figure.11. Note that the resolved zone has been smoothed for spatial consistency. Pixels outside the resolved zone are masked in the final maps.

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Figure.12 presents the tomographic inversion results at 1.4Hz, 6.0Hz and 9.8Hz for both Rayleigh waves (the left panels) and Love waves (the right panels). The black contour plotted on each map delineates the resolved zone defined by the raypath density. Any features outside this contour should be interpreted with caution. Broadly, the average velocity obtained decreases with the frequency (see the reference velocity in the subplot titles); the relative phase velocity variations of Rayleigh and Love waves exhibit similar patterns for all frequencies with lower velocities in the north and higher velocities at south. The histograms of traveltimes residuals have small standard deviations indicating good coherence between the measurements, on average (see Fig.S9 in the supporting information).

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4 Results

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4.1 Three-dimensional S wave velocity

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To build a high-resolution 3-D S wave velocity model of the area, we jointly invert the Rayleigh and Love wave dispersion curves obtained in each pixel of the tomographic inversions (Fig.12) by using a neighborhood algorithm (NA) with a Monte Carlo solver, Geopsy (Wathelet et al., 2004). Compared with single wave type dispersion inversion, joint inversion of both Rayleigh and Love wave has the advantages of reducing non-uniqueness

430 inherent in surface-wave methods and improving the accuracy of the inverted S wave ve-
 431 locity model (Joh et al., 2006; Chmiel et al., 2019; Yin et al., 2020). The 1-D depth model
 432 is parameterized with 13 layers including a half-space at the bottom (Table.1). The den-
 433 sity parameter is gradually increased with the depth based on borehole observations; V_p
 434 is linked to V_s during the inversion with a dynamic Poisson ratio range from 0.2 to 0.5;
 435 a loose prior constraint, based on the empirical formula on Xia et al. (1999), is applied
 436 to V_s . For each location, we invert the obtained dispersion curves with 3 runs of the in-
 437 version process. Each run retains 2500 models and we end up with the best 500 mod-
 438 els from all 3 runs. Figure.13 presents an example of the 1-D depth inversion. The for-
 439 ward modeled dispersion curves simultaneously converge to the measured Rayleigh (Fig.13a)
 440 and Love (Fig.13b) data with very small misfits. Although surface wave phase velocity
 441 is less sensitive to V_p compared to V_s , the coherence of both inverted V_s (Fig.13d) and
 442 V_p models (Fig.13d) still indicates a stable inversion processing.

443 We run 1-D depth inversions independently for each pixel (151x139) on the tomo-
 444 graphic maps, and combine the optimum 1-D V_s model obtained at each pixel to form
 445 a 3-D V_s model. Figure.14 presents the inverted V_s maps at depths of 0.10 km, 0.48 km,
 446 0.90 km, 1.22 km, 1.64 km, and 2.20 km. The primary pattern that emerges from both
 447 the tomographic phase velocity maps and the inverted V_s maps is the negative veloc-
 448 ity variations at north, which corresponds to the young Jurassic sediments, and the pos-
 449 itive velocity variations at south, which corresponds to the coarse-grained adamellites
 450 in the mountain areas. The negative velocity anomalies along the north river channel
 451 (the gray lines on Fig.14) also coincide with the alluvial deposits from the Quaternary
 452 system; we can observe similar anomalies across the CSAMT line (the black lines on Fig.14)
 453 at shallower depths. It is worth noting that the small negative velocity anomalies (< 3%)
 454 in the southwestern water reservoir area (outlined by the green lines on Fig.14) could
 455 be artifacts caused by the influence from the water reservoir on the surface wave ray-
 456 paths. Figure.15 shows variable performances of the sensitivity kernels between Rayleigh
 457 and Love wave for the 1-D velocity model at the well location (the magenta crosses on
 458 Fig.14). Joint analysis of Rayleigh and Love wave offers a multiple-scale vertical reso-
 459 lution investigation result from surface to depth 2 km.

460 4.2 HVSR Measurement

461 Most ambient noise studies only focus on extraction of the phase information for
 462 travel time tomography but abandon the amplitude information. Ambient noise (microtremors)
 463 horizontal-to-vertical spectral ratio (HVSR) provides the opportunity to extend the am-
 464 bient noise studies beyond the phase extraction scope, since multiple-component sens-
 465 ing has been more and more regular for seismic data acquisition. HVSR method has been
 466 widely used for estimation of predominant vibration frequency of soils, mainly for mi-
 467 crozonation and site effect purposes (Acerra et al., 2004; Gosar et al., 2010; Leyton et
 468 al., 2013; García-Jerez et al., 2016, 2019). Although the theoretical basis of the HVSR
 469 method is still debated, HVSR has been widely accepted as related to the ellipticity of
 470 Rayleigh waves and frequency dependent (Bard et al., 1999; Sylvestre et al., 2006). There-
 471 fore, HVSR exhibits a sharp peak at the fundamental frequency of the sediments, when
 472 there is a high impedance contrast between the sediments and underlying bedrock.

473 Following spectral analysis on the raw waveforms as described above, we apply Konno-
 474 Ohmachi smoothing (Konno & Ohmachi, 1998) with a b value of 40 on each one-hour
 475 averaged PSD spectrogram for 3 components of all available stations. Next, we calcu-
 476 late the HVSR as the square root of the ratio of the spectral energy components:

$$477 \frac{H}{V}(x, w) = \sqrt{\frac{E_1(x, w) + E_2(x, w)}{E_3(x, w)}}, \quad (1)$$

478 where, E_1 and E_2 stand for the spectral energy for the horizontal components; E_3 stands
 479 for the spectral energy for the vertical component; x indicates the station location; w

480 is the angular frequency. The spectral energy can also be computed from the average au-
 481 tuncorrelation of the ambient noise wavefield (Pertou et al., 2018). Finally, we obtain over
 482 100 one-hour averaged HVSR functions from 4.7 days observation, and the average HVSR
 483 function of all windows is taken as the final HVSR measurement of the corresponding
 484 location. We estimate the measurement precision for each frequency by 0.6 times of the
 485 standard deviation.

486 Figure.16 presents two examples of the obtained HVSR measurements from stations
 487 located at north (a) and south (b), separately. A clear peak with large amplitude
 488 (> 3) (Fig.16a) is related to a high impedance contrast between the sedimentary cover
 489 and the basement while a low amplitude (< 3) (Fig.16b) usually indicates a lower con-
 490 trast, for example of the presence of a hard soil at rock sites (Bard et al., 1999; Wool-
 491 ery & Street, 2002; Bonnefoy-Claudet et al., 2006, 2008). The peak frequency (f_0), or
 492 natural frequency, from HVSR measurement also reflects the sediment depth (h) with
 493 a general relationship $f_0 = \frac{V_s}{4h}$ (Castellaro & Mulargia, 2009; Pazzi et al., 2017). We
 494 reject the HVSR measurements (24/192) with flat HVSR curves and amplitudes smaller
 495 than 1 (see Fig.S10 in the supporting information) according to Acerra et al. (2004).

496 Figure.17a displays all available HVSR measurements overlaying on the surface ge-
 497 ology map (Fig.1b). We cluster the HVSR peaks based on the peak frequencies coded
 498 by the colors as well as the peak values coded by the scales. We observe four primary
 499 units that strongly correlate with the background geology:

- 500 1. to the north part of the area, large (>4) and dark blue ($<8\text{Hz}$) HVSR peaks in-
 501 dicate the strong impedance contrast between the thick sediment and the base-
 502 ment, which is collocated with the younger Jurassic sediment;
- 503 2. to the south of the area, small (<2) and dark red ($>11\text{Hz}$) HVSR peaks indicate
 504 the weak contrast between the hard rock surface, where coarse-grained adamel-
 505 lites are widely distributed, and the basement;
- 506 3. along the river crossing the town, medium size HVSR peaks with peak frequen-
 507 cies around 9Hz coincide with the fluvial deposits from the Quaternary system;
- 508 4. in the central area, a transition zone with various peak frequencies and peak val-
 509 ues HVSR measurements is highlighted by the magenta shadow belt.

510 Figure.17b presents the iso-surface of the median velocity of the inverted Vs model
 511 with $V_s = 3.0\text{km/s}$. The surface colors are coded by depths. Although it does not strictly
 512 reflect the basement surface, the northern cavern and the southern hump indicate the
 513 rather deeper basement in the north than that in the south. These results are consis-
 514 tent with HVSR observations with the lower peak frequencies distributed in the north
 515 and higher peak frequencies distributed in the south. The lower surface velocities in the
 516 north (Fig.14a) coincide with the stronger impedance contrast inferred from the larger
 517 HVSR peaks.

518 5 Discussion

519 We successfully resolve 3-D S wave velocity model from the surface to 2 km depth
 520 with spatial resolutions of 0.2 km~0.3 km, and observe the negative velocity variation
 521 around 3% to depths of up to 1 km in the fluvial deposit area. These low-velocity anoma-
 522 lies are consistently present at depth in both frequency (Fig.12) and depth (Fig.14) do-
 523 mains. The transition belt observed on HVSR measurements also covers this area. These
 524 lower velocity anomalies may be associated with the high-temperature geothermal field
 525 from deep to surface. The cross sections of the S wave velocity structure as well as the
 526 electric resistivity measurement in Figure.18 provide additional insight on the spatial dis-
 527 tribution of the potential geothermal field.

Vertical cross sections along the CSAMT profile line present clear velocity anomalies from both the absolute velocity profile (Fig.18a) and the velocity variation (Fig.18b) profile. In particular, distinct low velocity anomalies can be observed on the velocity variation profile from location 1.0 km to location 1.6 km. We are able to infer three steep faults, Fw^1 , Fw^2 and Fw^3 , based on the boundary of the positive and negative velocity variation. Fw^3 has been proven by the surface geology observation. The low velocity anomalies among Fw^1 and Fw^2 appear as a channel from deep to surface, which is consistent with the colocated conductive region on the resistivity profile, which is suggested as hot flow signature. Note that we mask the two sides of the cross sections that beyond the resolvable zone.

Figure.19 presents the 3-D view of the iso-surfaces of the velocity variations at $\Delta V = -3\%$ (a) and $\Delta V = -1\%$ (b). The surface colors are coded by depths. The horizontal slice shows the plane Vs variation at depth 2.0 km. The slice colors are coded by the velocity variations as Figure.14. The resolvable zone is indicated by the black line. We observe a clear reservoir bounded by the negative velocity anomalies surface on Figure.19a, which is colocated with the observed transition belt on HVSR measurements. However, we are not going to further discuss the anomalies beyond the resolvable zone since they could be potentially stretched during the tomography due the poor resolve resolution. Within the resolvable zone, a funnel-shaped low-velocity zone (LVZ) (indicated by the green arrow) is visible on Figure.19a. It biases from south to north with the depth increasing, and the root turns to be broader with velocity variation decreasing to $\Delta V = -1\%$ (Fig.19b). The CSAMT profile (indicated by the white line) intersects this LVZ particularly the smaller velocity variation surface on Figure.19b. Combined with the conductive characteristics observed on CSAMT profile, we interpret this LVZ as a zone of more intense fracturing with conductive fluids (Guéguen & Palciauskas, 1994; Paterson & Wong, 2005; Lehujeur et al., 2018). We also observe the drilling well (indicated by the thick black stick on Fig.19b) crossing into the LVZ from depth around 1 km, which has been proven by the sudden decrease around 1 km in fluid resistivity well records (the black curve on Fig.20). The sudden decrease in fluid resistivity also supports our interpretation that the fluid filled fracture channel presents as a more conductive zone than the surrounding rocks. The gradually increasing borehole temperature logs (the magenta curve on Fig.20a) show a gradient of around $3^\circ C/100m$ indicating the existence of the higher-temperature geothermal resources at depth.

We also observe a columnar LVZ (indicated by the blue arrow) on Figure.19b which is located beneath the water reservoir. As described above, this abnormal body could be artifacts caused by the influence from the water reservoir on the surface wave ray-paths. Another shallow (above 0.9 km) abnormal body is indicated by the red arrow on Figure.19b. We cannot rule out the possibility that the observed shallow anomalies are related with the geothermal activity. Unlike the LVZ indicated by the green arrow, however, this one loses the surface water resource from the nearby Plum Creek River which might limit the condition for the generation of a good geothermal field since the underground water layer in this area is usually deeper than the maximum depth of the abnormal body. Both questions could be addressed if we can include more constraints, for example extraction body wave from interferograms for traveltimes tomography or application of additional magnetotelluric (MT) surveys.

6 Conclusions

We successfully retrieve surface waves, both Rayleigh and Love waves, from ambient noise over an ultra-short observation period using a dense nodal array, and apply tomographic imaging of the subsurface 2 km S wave velocity structure beneath a rural town. For the first time, we demonstrate the advantage of cross-coherence over cross-correlation on suppressing pseudo-arrivals caused by persistent sources. Body waves are also observed on the cross-coherence functions which offer the possibility for the further

580 body wave tomography study. We investigate spatial horizontal resolutions for the to-
 581 morphic inversion and present a resolvable zone with the highest resolution of 0.2 km.
 582 Strong correlation between HVSR measurements and S wave velocity model indicates
 583 the potential of extraction of both amplitude and phase information from 3C ambient
 584 noise data, which will increase the data utilization coefficient and provide more constraints
 585 for ambient noise imaging. Given the continually increasing demands for the develop-
 586 ment of local geothermal resources, particularly in China, our work demonstrates the util-
 587 ity of high spatial-resolution geothermal characterization with affordable seismic nodal
 588 array observation, as well as high temporal-resolution geothermal monitoring due to the
 589 ultra-short observation period.

590 We detect low-velocity anomalies (between -1% and -4%) from surface to depth in
 591 the central area, which is inferred as a fracture channel filled in with the fluid contents
 592 in the light of observation of the colocated conductive zone on resistivity profile. Joint
 593 interpretation based on HVSR measurements, the temperature and fluid resistivity records
 594 observed in a nearby well suggests the existence of the high-temperature geothermal field
 595 through the fracture channel.

596 Acknowledgments

597 This study is supported by the National Natural Science Foundation of China under grant
 598 No. 41830103 and Zhejiang Geophysical and Geochemical Prospecting Academy. The
 599 seismic time-series datasets utilized in this paper were archived in the Mendeley Data
 600 (<http://dx.doi.org/10.17632/382frvcy6r.3>). We appreciate Dr. Shucheng Wu for
 601 many useful suggestions on surface wave tomography.

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Table 1. Prior boundaries of uniform probability distributions used for each parameter of the depth model.

Layer number	Thickness(<i>km</i>)	Poisson ratio	V_s (<i>km/s</i>)	Density(<i>g/cm³</i>)
01	0.01 ~ 0.03	0.2 ~ 0.5	1.3 ~ 2.3	1.9 ~ 2.6
02	0.02 ~ 0.04	0.2 ~ 0.5	1.5 ~ 2.8	2.0 ~ 2.7
03	0.02 ~ 0.06	0.2 ~ 0.5	1.7 ~ 3.2	2.0 ~ 2.8
04	0.03 ~ 0.08	0.2 ~ 0.5	1.8 ~ 3.4	2.0 ~ 2.8
05	0.04 ~ 0.10	0.2 ~ 0.5	1.9 ~ 3.5	2.1 ~ 2.8
06	0.05 ~ 0.13	0.2 ~ 0.5	2.0 ~ 3.6	2.1 ~ 2.8
07	0.07 ~ 0.18	0.2 ~ 0.5	2.0 ~ 3.7	2.1 ~ 2.9
08	0.09 ~ 0.23	0.2 ~ 0.5	2.0 ~ 3.8	2.1 ~ 2.9
09	0.12 ~ 0.31	0.2 ~ 0.5	2.1 ~ 3.9	2.1 ~ 2.9
10	0.16 ~ 0.41	0.2 ~ 0.5	2.2 ~ 4.0	2.1 ~ 2.9
11	0.21 ~ 0.55	0.2 ~ 0.5	2.3 ~ 4.2	2.2 ~ 3.0
12	0.28 ~ 0.73	0.2 ~ 0.5	2.4 ~ 4.4	2.2 ~ 3.0
half-space	0.28 ~ 0.73	0.2 ~ 0.5	2.4 ~ 4.5	2.2 ~ 3.0

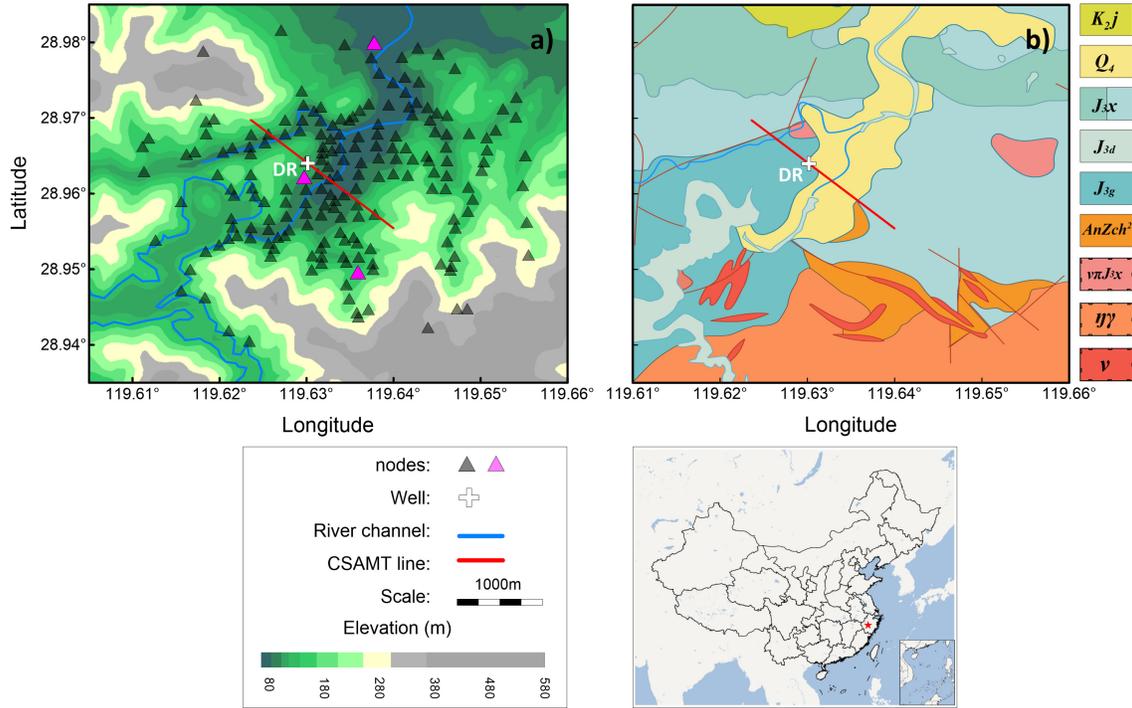


Figure 1. Maps of the geothermal site and the Andi network. (a). Topography map of the survey region and available seismic nodal network. The triangles denote the Zland nodes; three magenta triangles denote three stations used in spectral analysis; the white cross denotes the well location; the red line indicates the CSAMT profile line; the blue curves indicate the river channel as well as the water reservoir outlines located in the southwest. (b). Geology map of the survey region. K_2j denotes the Cretaceous system; Q_4 denotes the Quaternary system; J_{3x} , J_{3d} , and J_{3g} denote three different groups of the upper Jurassic system; $AnZch^2$ denotes the PreSinian system; $\nu\pi J_{3x}$ indicates the felsophyre; $\eta\gamma$ denotes the coarse-grained adamellites; ν denotes the granite dykes. The red star on China map indicates the location of the Andi town.

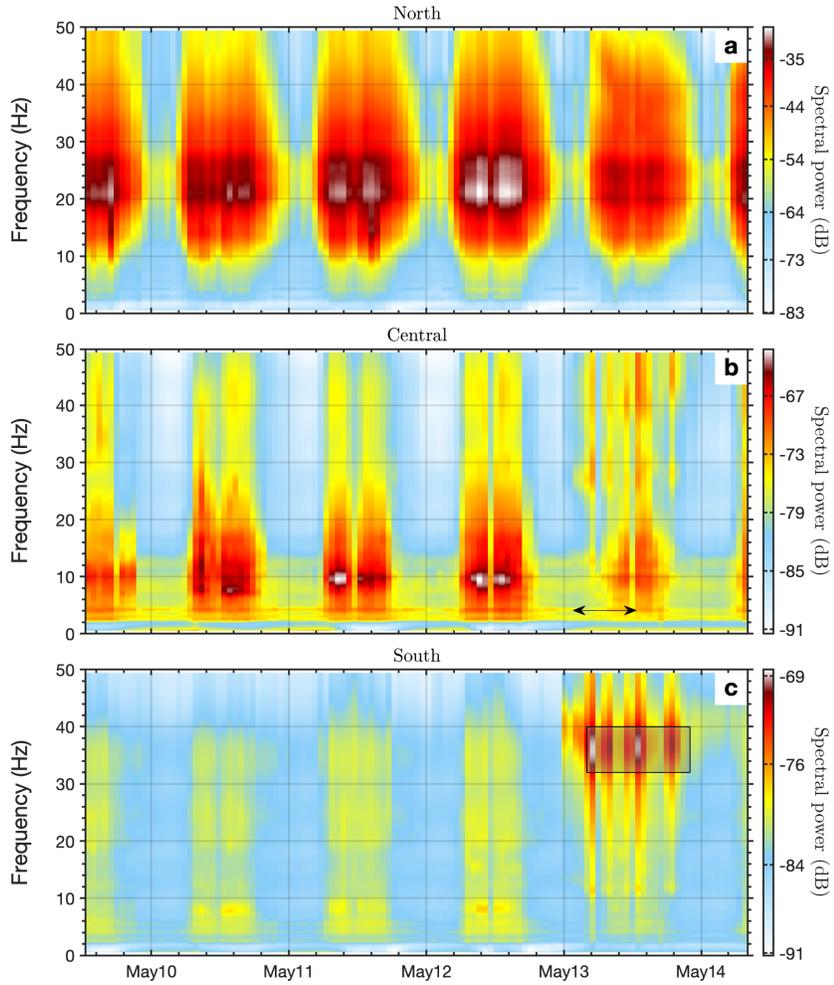


Figure 2. Vertical component spectrograms of over 4 days ambient noise data at three different stations, north (a), central (b) and south (c). The three stations are denoted by the three magenta triangles on Fig.1a. The black double-arrow on b indicates the spectrum of consistent sources from anthropogenic activities. The gray box on c highlights the source spectrum from the rain- and thunder-induced ground motions.

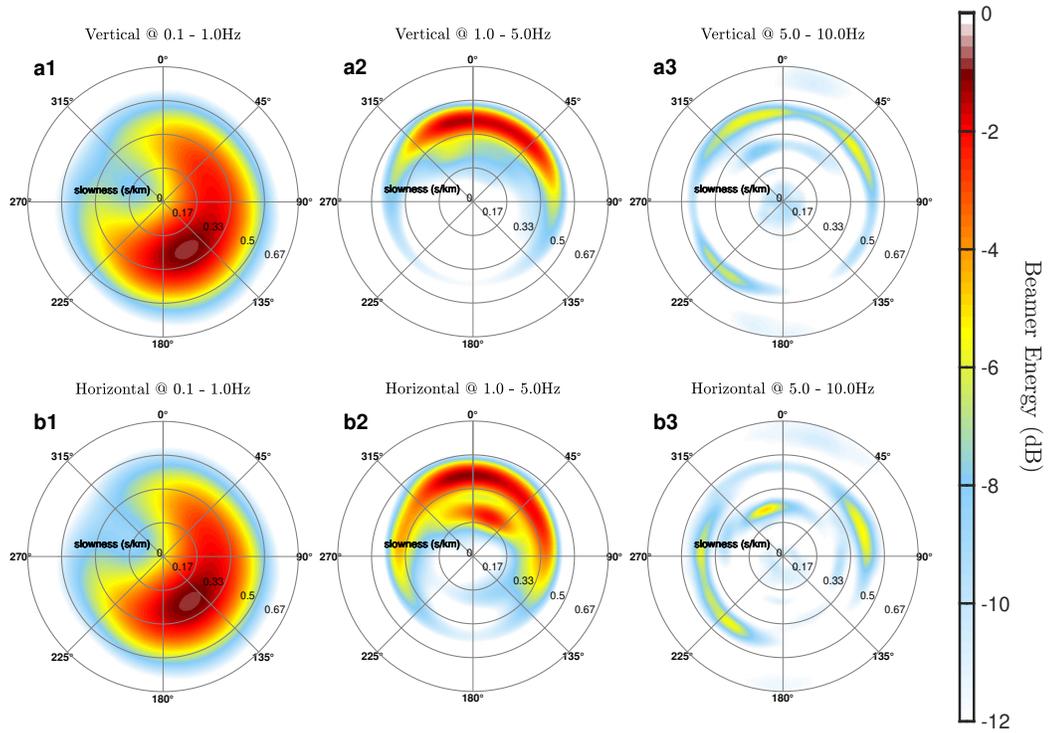


Figure 3. Beamforming analysis performed on the vertical component (the upper panels, a1, a2, a3) and the north horizontal component (the bottom panels, b1, b2, b3) of Andi network at different frequency bands 0.1-1.0Hz (a1, b1), 1.0-5.0Hz (a2, b2) and 5.0-10Hz (a3,b3).

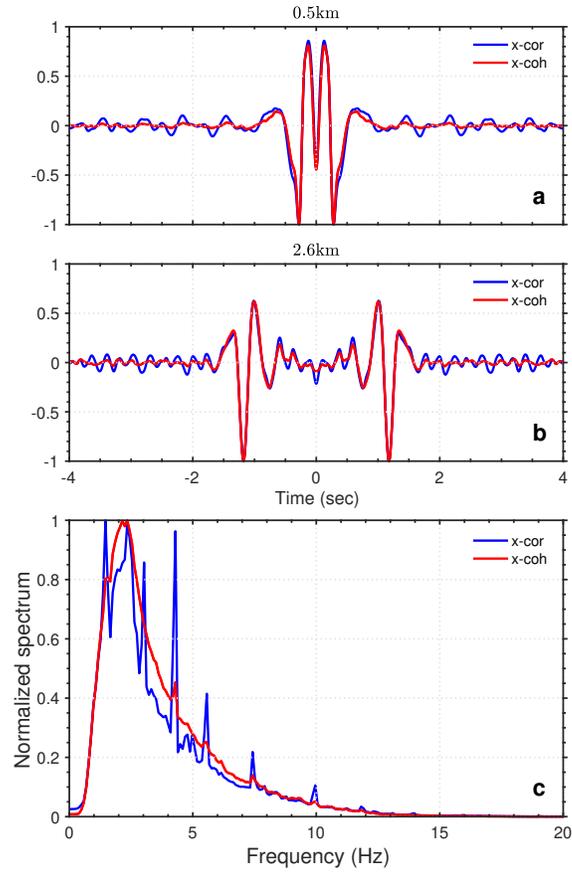


Figure 4. Comparisons between the bin-stacked cross-correlation (the blue curves) and cross-coherence (the red curves) functions with interstation distance at 0.5km (a) and 2.6km (b). (c) Comparison between the averaged spectrum from the bin-stacked virtual-source gathers from cross-correlation and cross-coherence (see Fig.S1 in the supporting information).

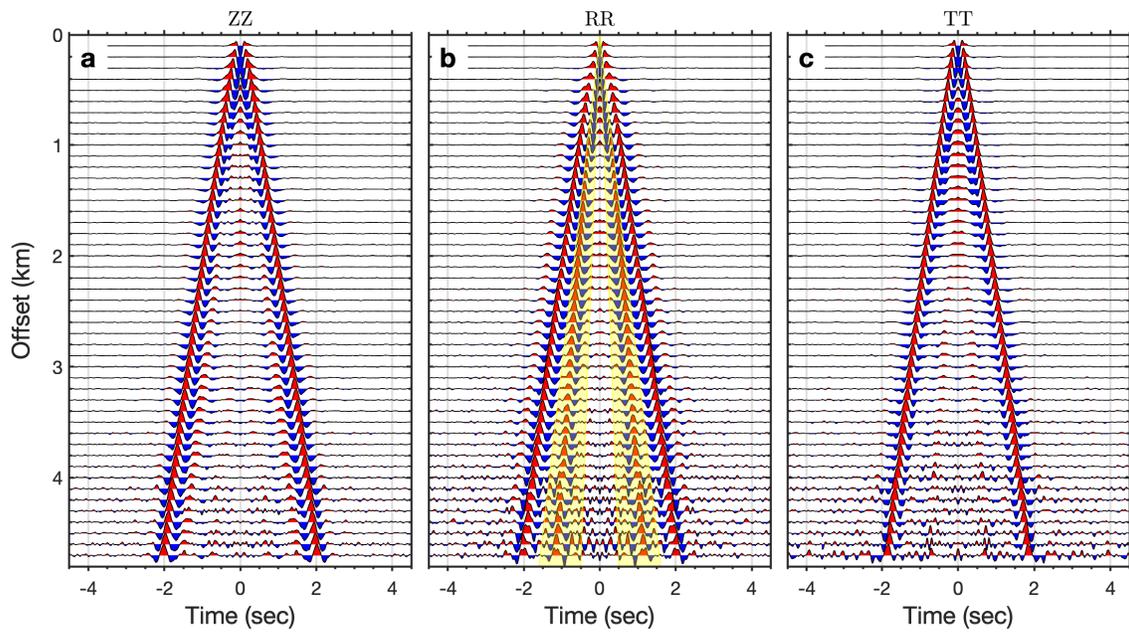


Figure 5. Bin-stacked virtual-source gathers from Z-Z (a), R-R (b), and T-T (c) cross-coherence functions. Bandpass filter between 1 and 10Hz has been applied. The body waves are highlighted by yellow color on b.

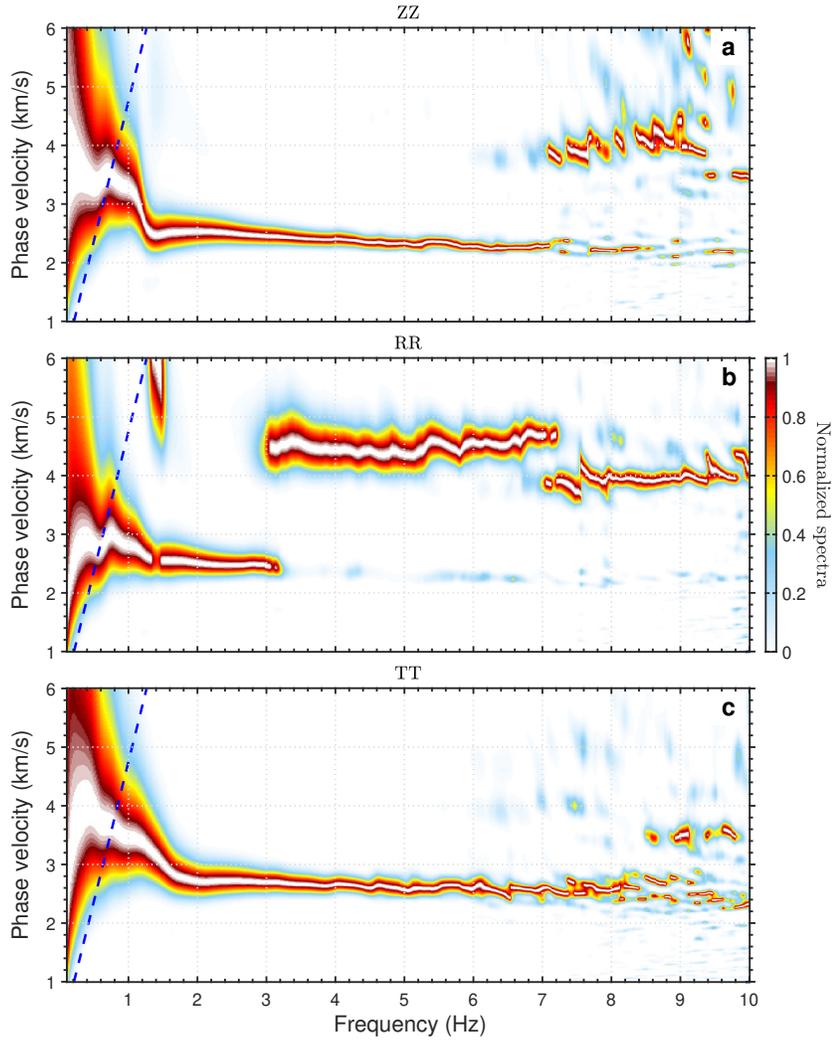


Figure 6. Dispersion measurements from the bin-stacked virtual-source gathers at Z-Z component (a), and R-R component (b), and T-T component (c). The blue dash lines indicate the minimum wavenumber defined by $k_{min} = \frac{1}{Array\ length}$.

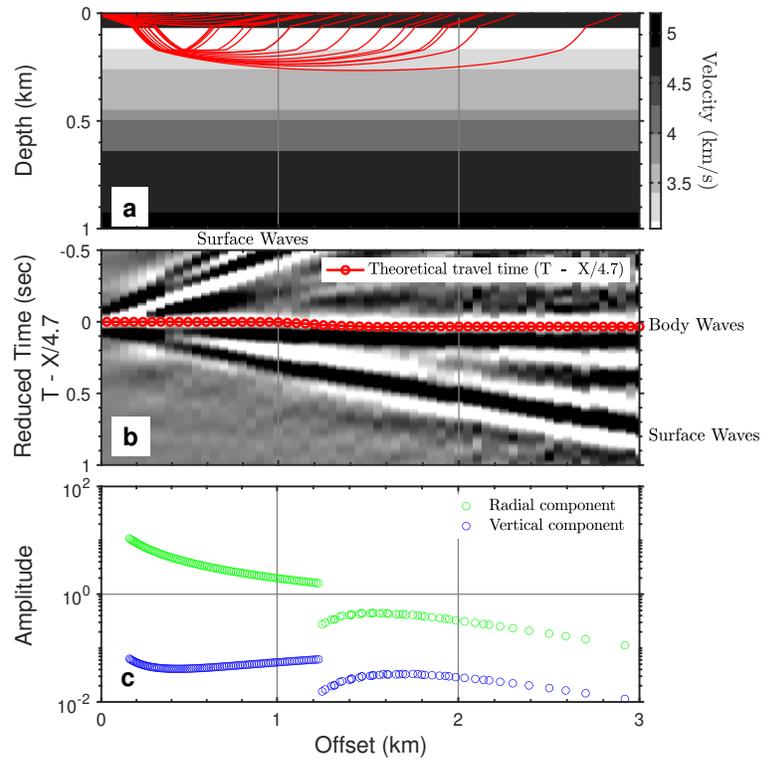


Figure 7. (a) 1D velocity model used for ray tracing as derived from smoothed sonic log and resulting ray geometry of the P-wave. (b) Superposition of the theoretical raytracing travel times on the interferometric wave field of the bin-stacked R-R component. Traveltimes are corrected with a linear move-out velocity of 4.7km/s. Bandpass filter between 4 and 10Hz has been applied. (c) Ray tracing amplitudes of the radial and vertical components.

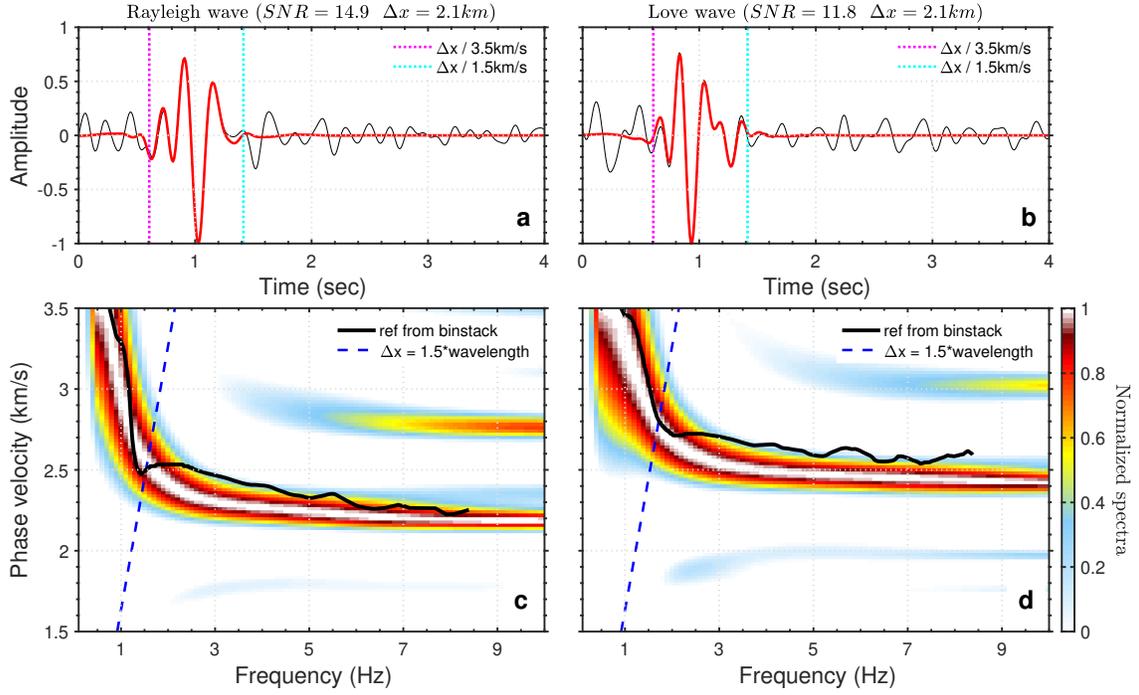


Figure 8. Examples of dispersion analysis for Rayleigh wave (left panels) and Love wave (right panels) using the image transformation technique by Yao et al. (2006). The thin black and red curves on a and b present the empirical Green's functions before and after velocity filter (or mute) with velocity range from 1.5km/s to 3.5km/s. Colored dashed lines indicate the time window estimated from the corresponding velocity window. For better presentation, all EGFs have been bandpass filtered (1~10Hz). The thick black curves on c and d present the reference average dispersion curves picked from Fig.6. The blue dashed lines indicate the 1.5 times wavelength criterion.

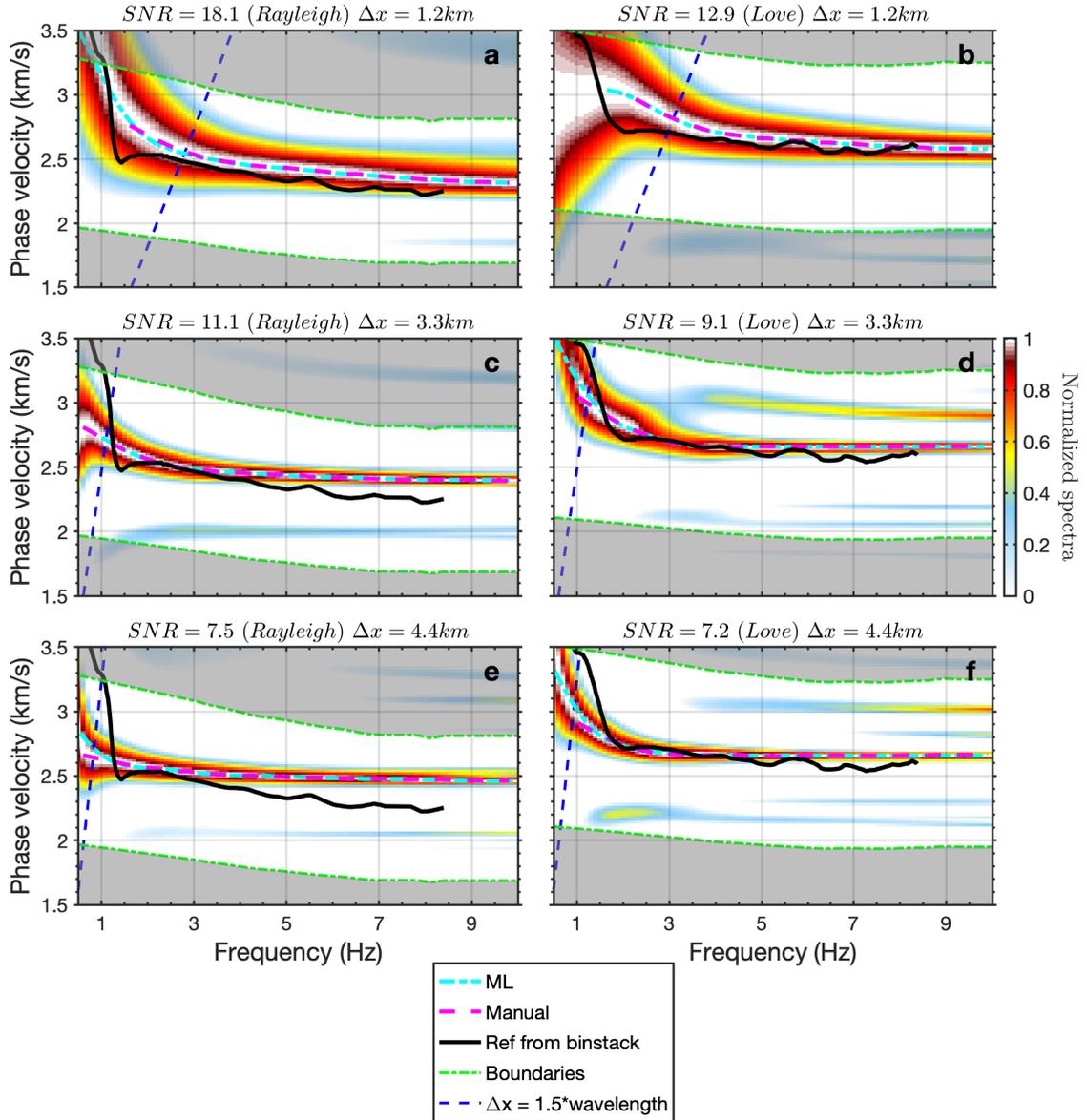


Figure 9. Examples of dispersion curves picking for Rayleigh waves (left panels) and Love waves (right panels) using Machine Learning. The cyan curves indicate the picked dispersion curves using ML technique; the magenta curves indicate the manually picked dispersion curves; the black curves indicate the reference average dispersion curves picked from Fig.6; the green curves indicate the upper and bottom boundaries defined by the smoothed reference dispersion curves; the blue lines indicate the 1.5 times wavelength criterion. We mask the dispersion spectra beyond the confidence region. We label each sub-figure with the corresponding SNR, wave type and the interstation distance.

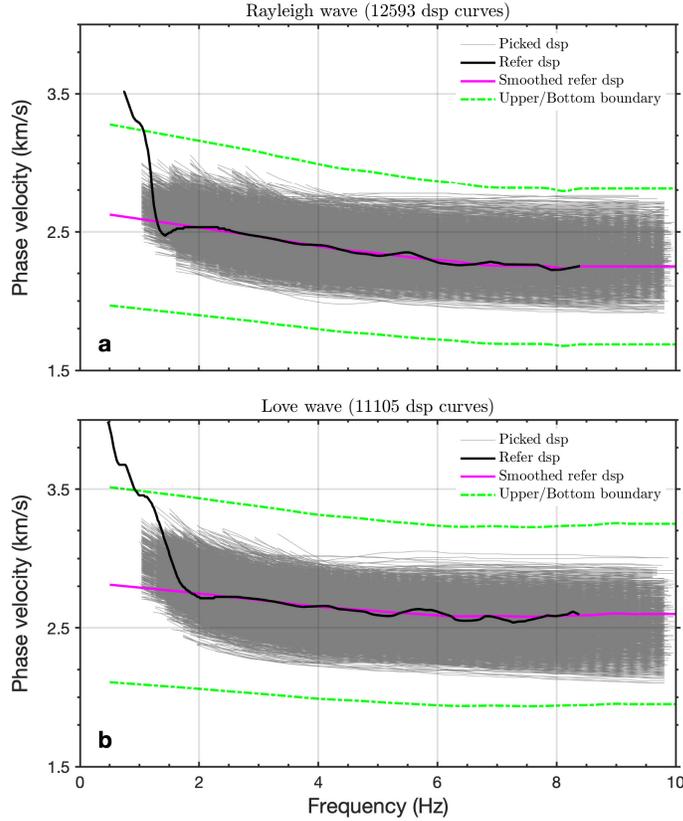


Figure 10. Picked dispersion curves for Rayleigh waves (a) and Love waves (b) by using Machine Learning. The thick gray curves indicate the final picked dispersion curves using ML technique; the black curves indicate the reference average dispersion curves picked from Fig.6; the magenta curves indicate the smoothed reference dispersion curves; the green curves indicate the upper and bottom boundaries defined by the smoothed reference dispersion curves. We label each sub-figure with the corresponding wave type and the total number of the picked dispersion curves.

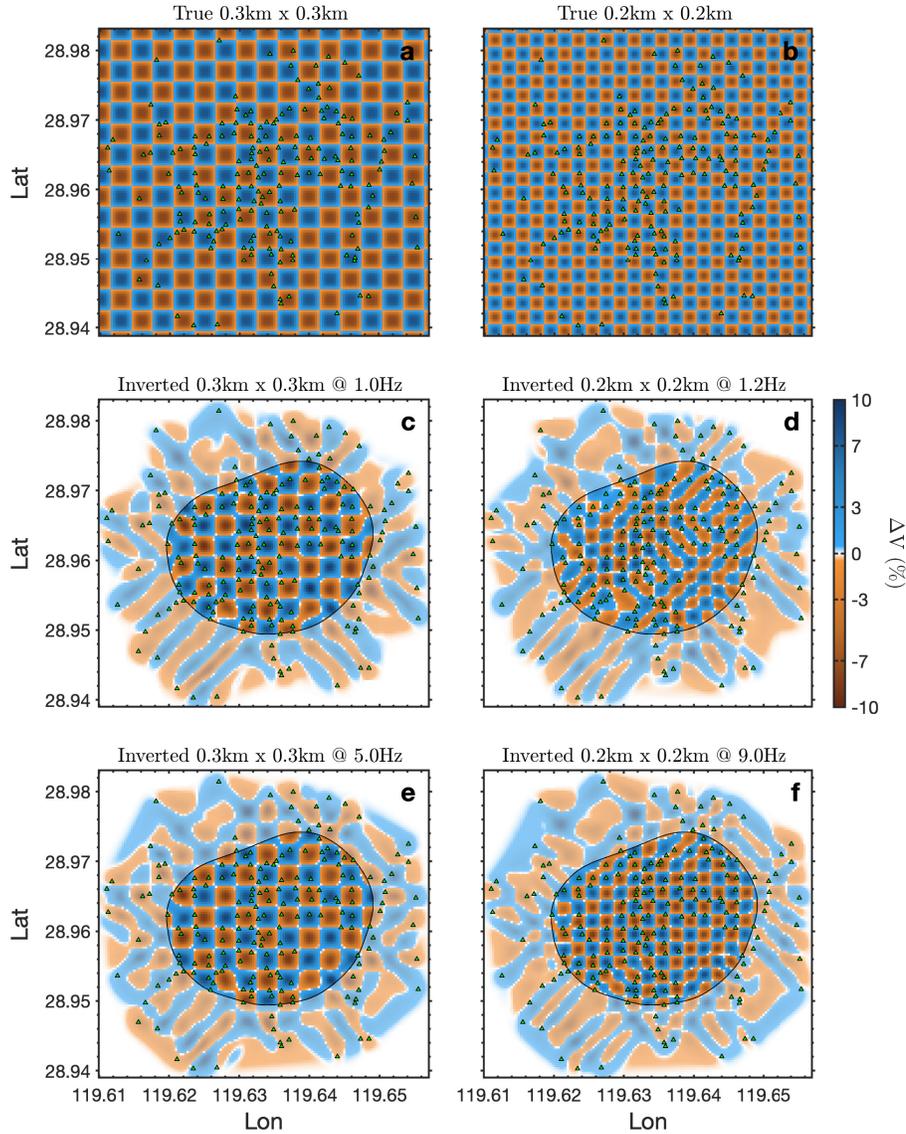


Figure 11. Checkerboard tests for surface wave tomography with two spatial resolutions, 0.3km (the left panels) and 0.2km (the right panels). (a) and (b) present the simulated models; (c-f) present the recovered models at different frequencies, 1.0Hz, 1.2Hz, 5.0Hz, and 9.0Hz. The green triangles denote the seismic network; the black contours indicate the resolvable zone defined with raypath density. We mask the area beyond the resolvable zone. Here we consider the Love waves as an example.

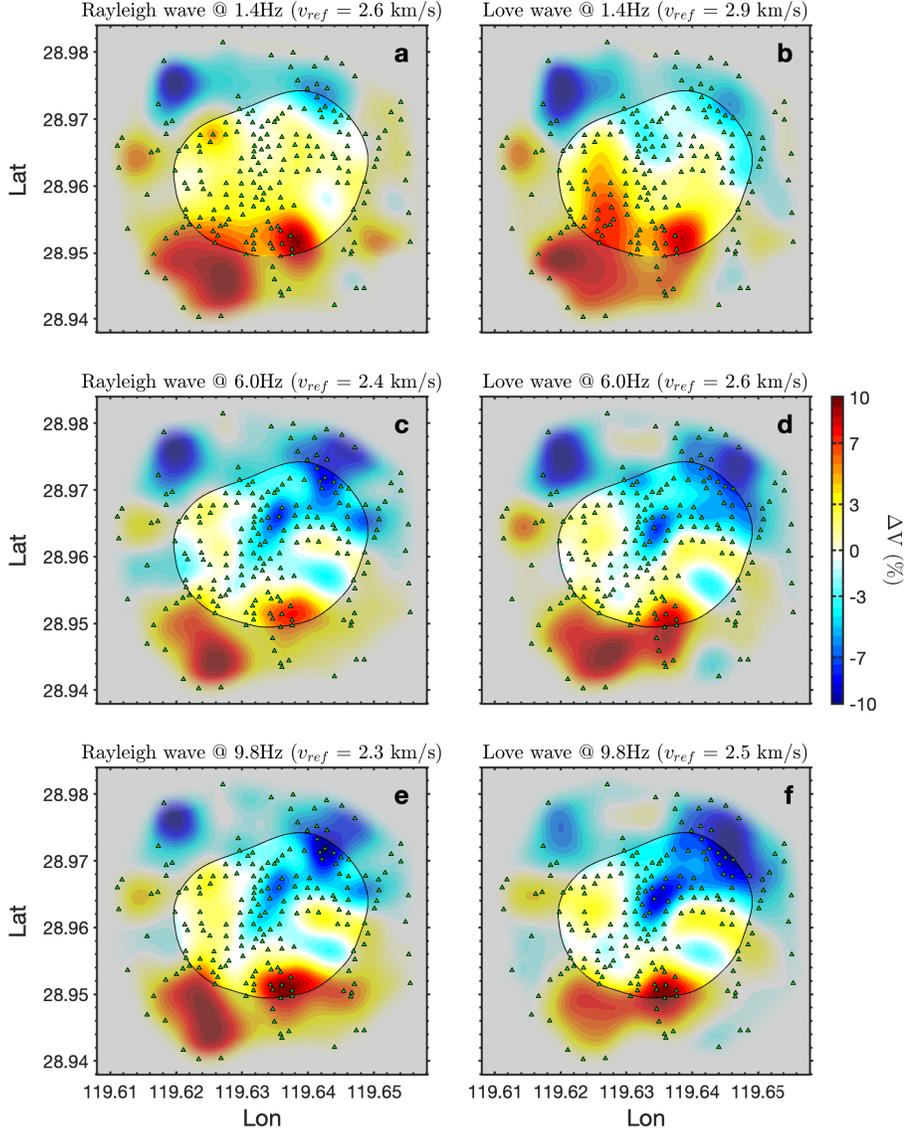


Figure 12. Phase velocity tomographic results for Rayleigh waves (the left panels) and Love waves (the right panels) at different frequencies, 1.4Hz (a and b), 6.0Hz (c and d), 9.8Hz (e and f). The green triangles denote the seismic network; the black contours indicate the resolvable zone defined with raypath density. We mask the area beyond the resolvable zone.

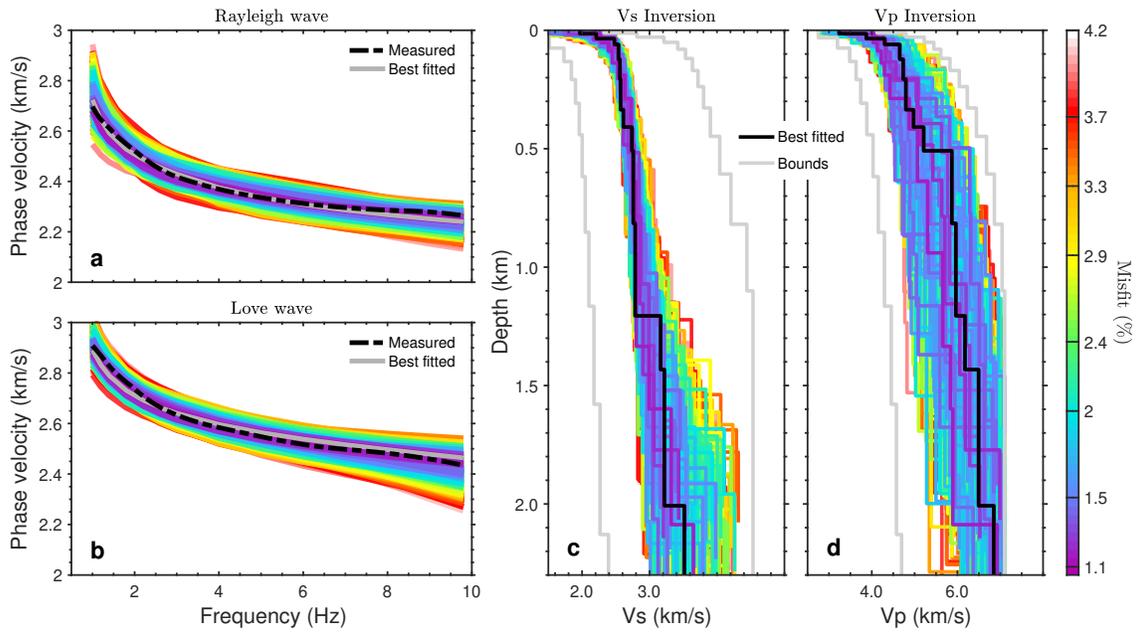


Figure 13. Joint inversion of Rayleigh wave and Love wave dispersion curves. (a) and (b) present examples of the measured (the black dashed curves) and the best 500 forwarded (the colored curves) dispersion curves; the gray curves indicate the best fitted dispersion curves. (c) and (d) present the best 500 Vs and Vp models; the black curves indicates the best fitted model; the gray curves indicate the upper and bottom velocity boundaries. Colors are coded by misfits as shown on the color map.

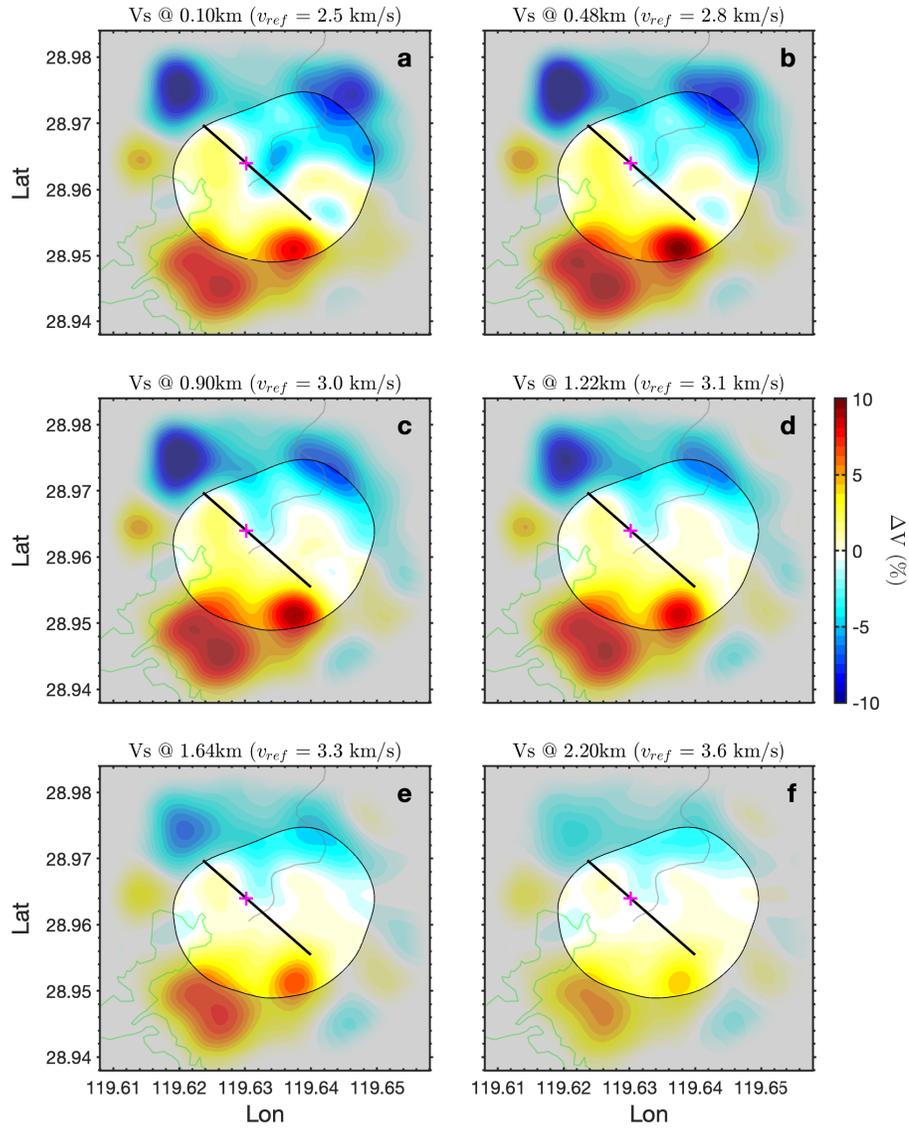


Figure 14. Horizontal slices of the obtained S wave velocity model at different depths. The thick black line indicates the CSAMT profile line; the magenta cross indicates the well location; the thin contour indicates the resolvable zone. We mask the area beyond the resolvable zone. We label each sub-figure with the corresponding depth and the reference velocity.

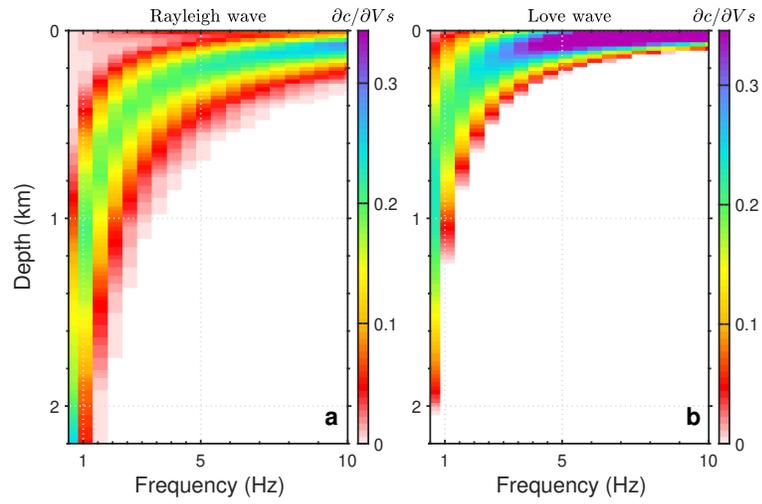


Figure 15. Sensitivity kernels of Rayleigh (a) and Love wave (b) based on the 1D velocity model at the well location.

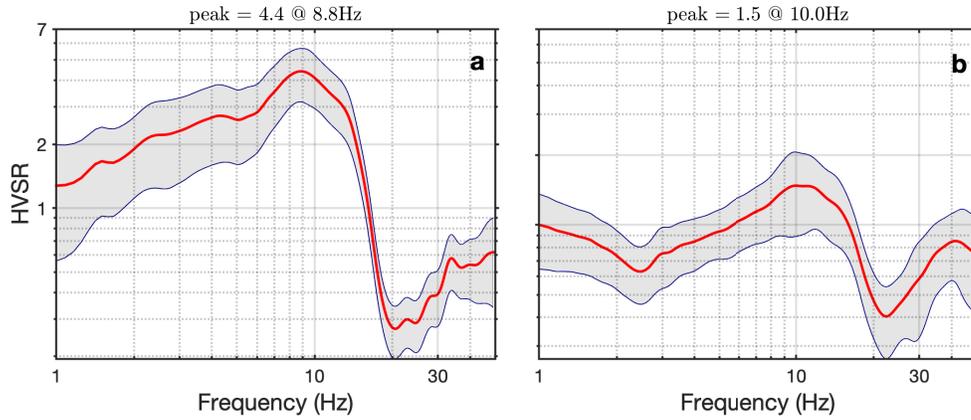


Figure 16. Examples of the obtained HVSR curves at north (a) and south (b). The red curves denote the measured HVSR curves; the blue curves indicate the measurement precisions defined by 0.6 times of the standard deviation.

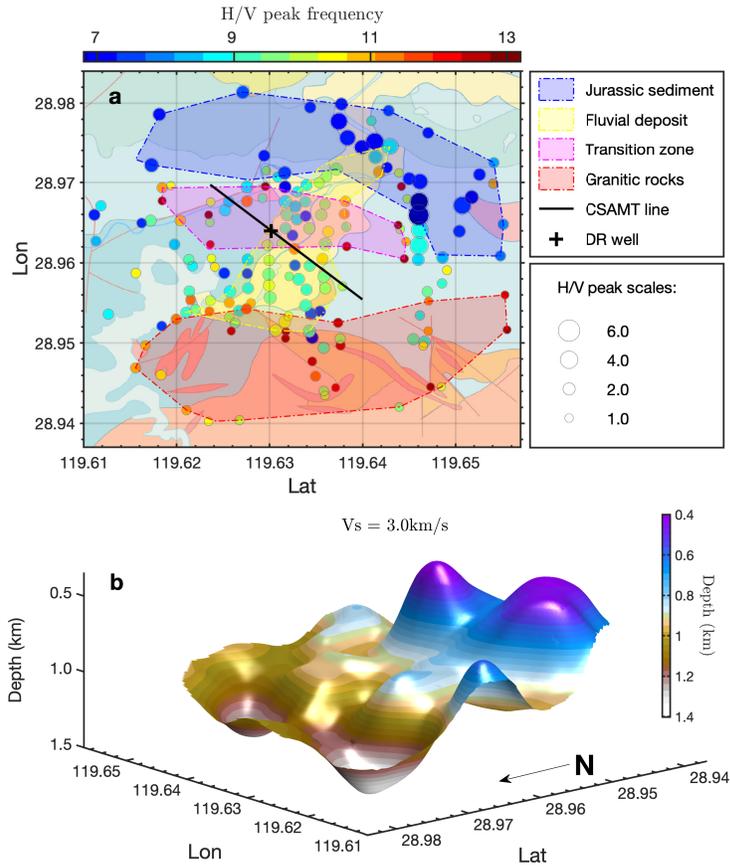


Figure 17. (a). The distribution map of the measured HVSR peaks overlaying on the surface geology map. The scatters denote the HVSR peaks from different stations. The scatter colors are coded by the HVSR peak frequencies; the scatter sizes are scaled by the HVSR peak values. The black line indicates the CSAMT profile line; the cross indicates the well location. Four colored shadows present four main clusters of HVSR measurements as indicated on the legend box. (b). The iso-surface of the median velocity of the inverted V_s model with $V_s = 3.0 \text{ km/s}$. The surface colors are coded by depths.

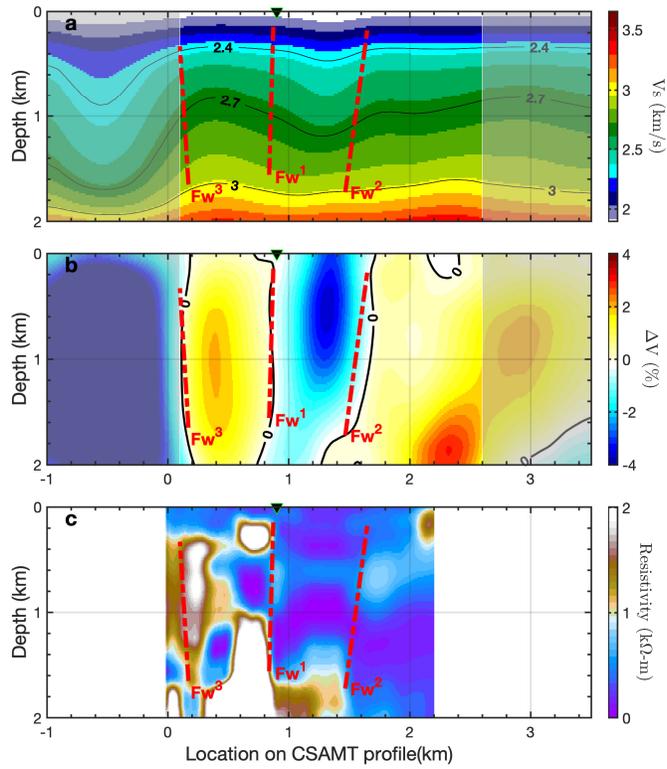


Figure 18. Vertical sections of the obtained S wave velocity model, (a) the absolute velocity model and (b) the velocity variation model, along the CSAMT profile (c). The red dashed lines indicate the inferred fault. We mask the sections beyond the resolvable zone on a and b.

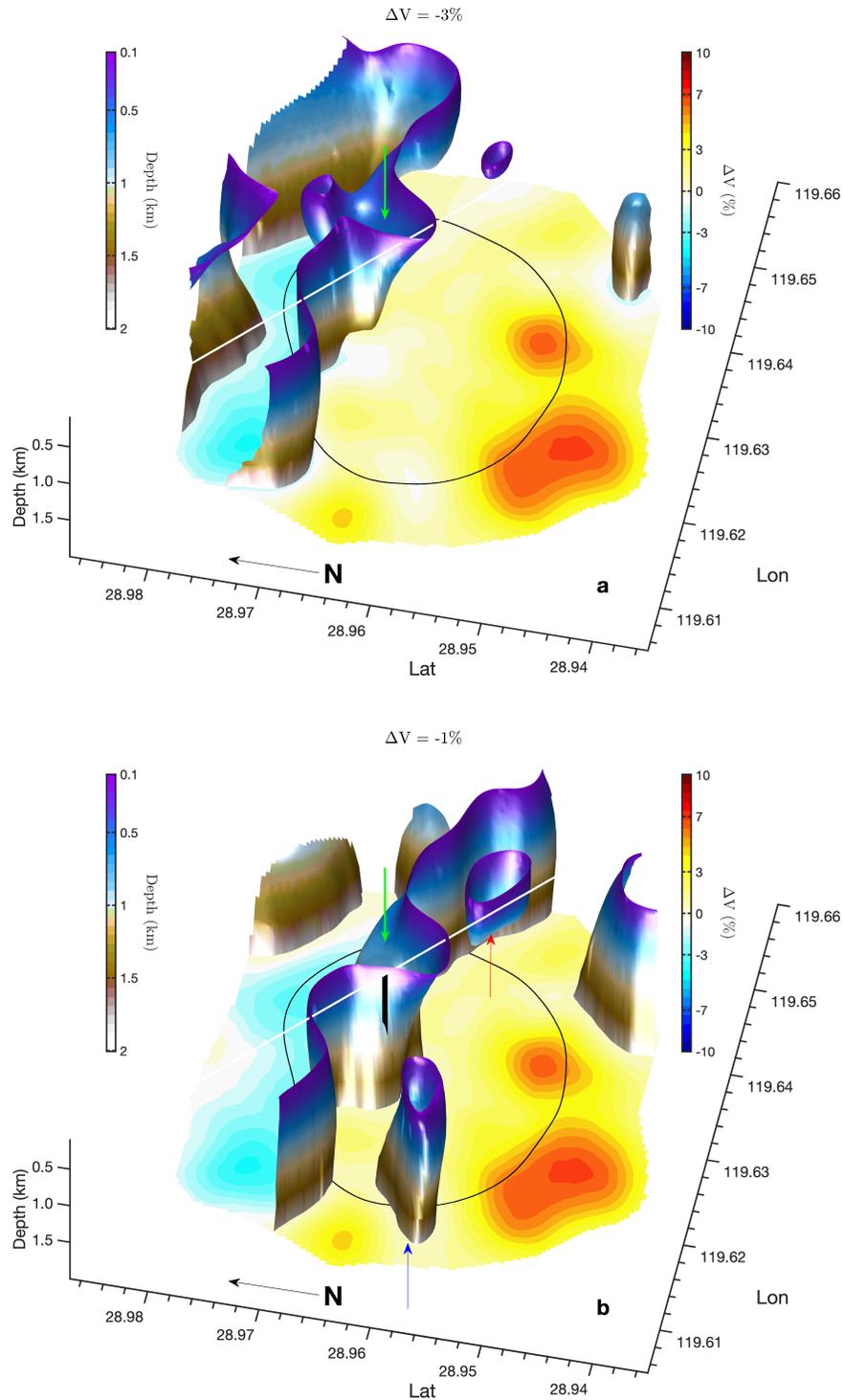


Figure 19. 3-D iso-surface of the obtained Vs variations at $\Delta V = -3\%$ (a) and $\Delta V = -1\%$ (b). The horizontal slice on a and b present the plane Vs variation at depth 2.0km. The black contour on a and b indicate the resolvable zone; the white line on a and b indicate the CSAMT profile line; the green arrow on a and b indicate the interpreted geothermal reservoir channel; the black arrow on b indicates the possible artifacts caused by water reservoir; the red arrow on b indicates the shallow low-velocity anomalies. Colors of the iso-surfaces are coded by depths; colors on the horizontal slices are coded by the velocity variations.

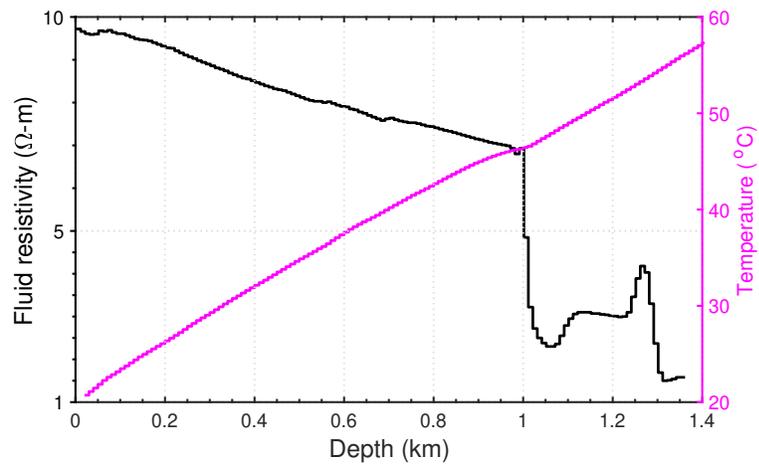


Figure 20. Well log of fluid resistivity (the black curve) and temperature (the magenta curve).

Supporting Information for "Cheng et al., High-resolution ambient noise imaging of geothermal reservoir using dense seismic nodal array and ultra-short observation"

Contents of Supporting Information

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9. Figure.S9 Histograms of the final residuals of surface wave traveltime tomography.
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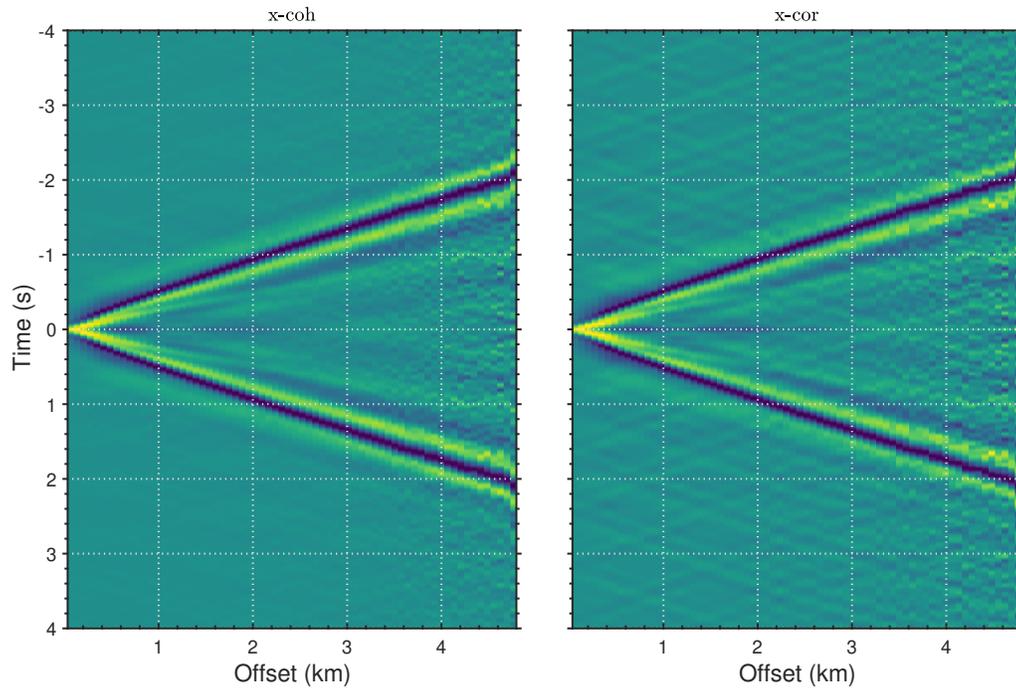


Figure S1. Comparison of the bin-stacked virtual-source gathers between cross-coherence (left) and cross-correlation (right) at vertical component.

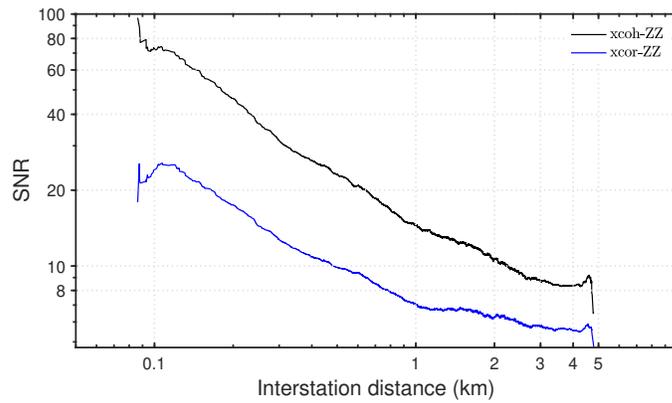


Figure S2. Comparison of SNR between extracted cross-coherence functions (black) and cross-correlation functions (blue). SNR curves have been smoothed for better display.

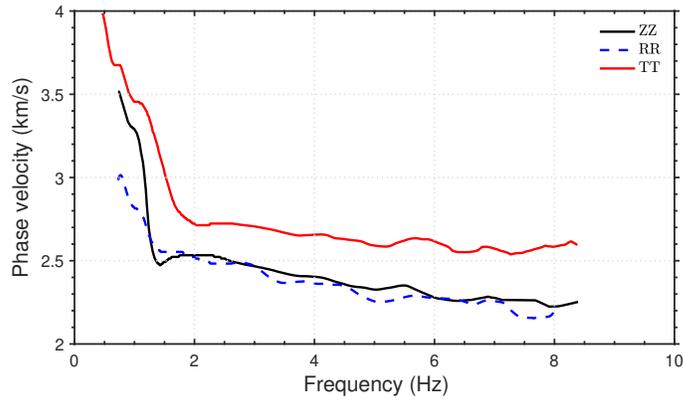


Figure S3. The reference dispersion curves picked from Fig.6 for Z-Z component (black), R-R component (blue) and T-T component (red).

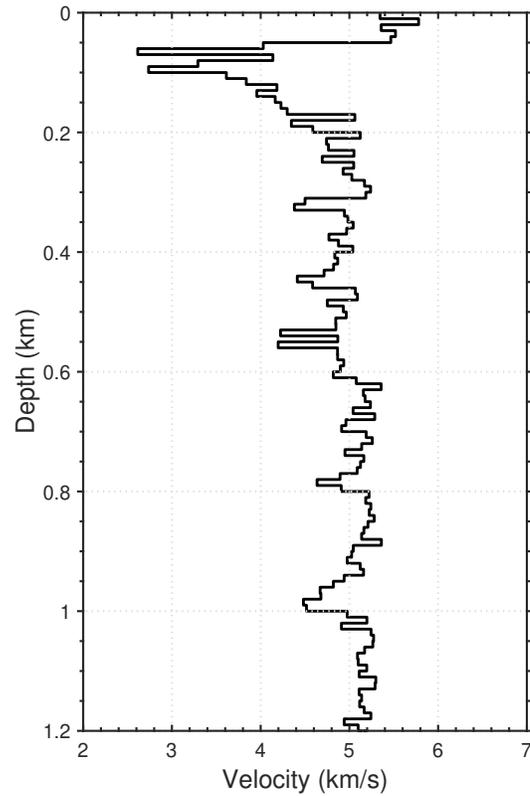


Figure S4. Sonic log from the center of the area (white cross on Fig.1).

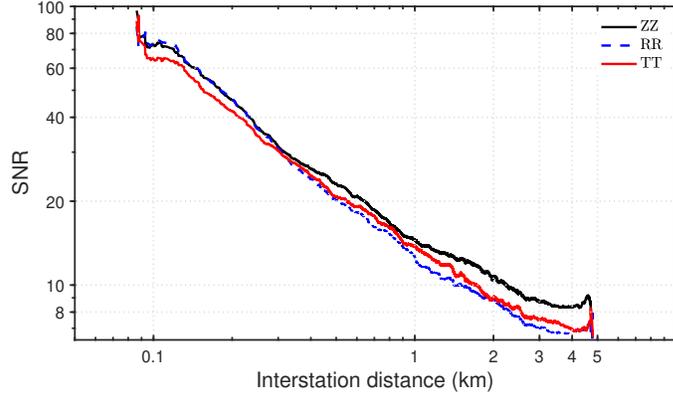


Figure S5. Comparison of SNR of cross-coherence functions between Z-Z component (black), R-R component (blue) and T-T component (red). SNR curves have been smoothed for better display.

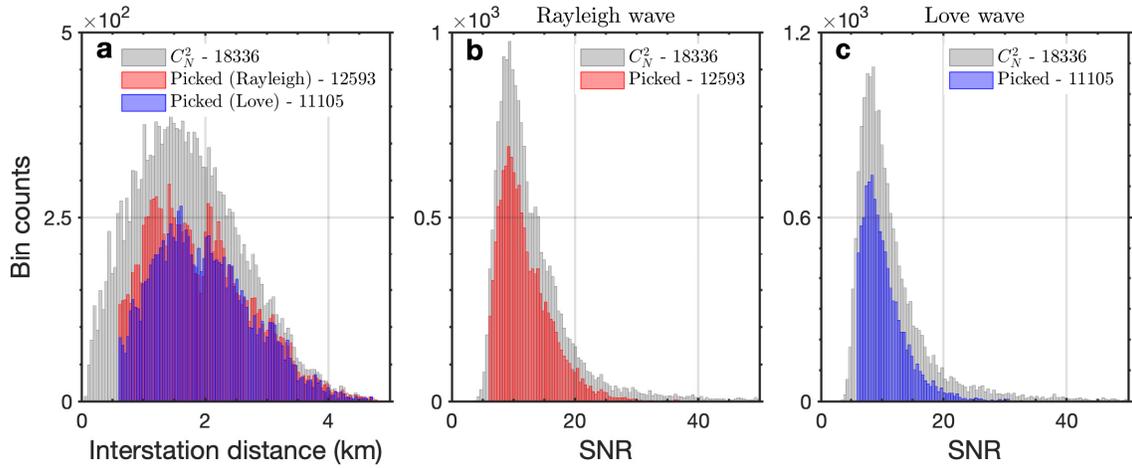


Figure S6. Histograms of interstation distances and SNRs of the picked dispersion curves. (a) Histograms of the interstation distances of all C_N^2 interstation pairs (gray), all picked Rayleigh waves (red), and all picked Love waves (blue). (b) Histograms of SNRs of all C_N^2 interstation pairs (gray) and all picked Rayleigh waves (red). (c) Histograms of SNRs of all C_N^2 interstation pairs (gray) and all picked Love waves (blue).

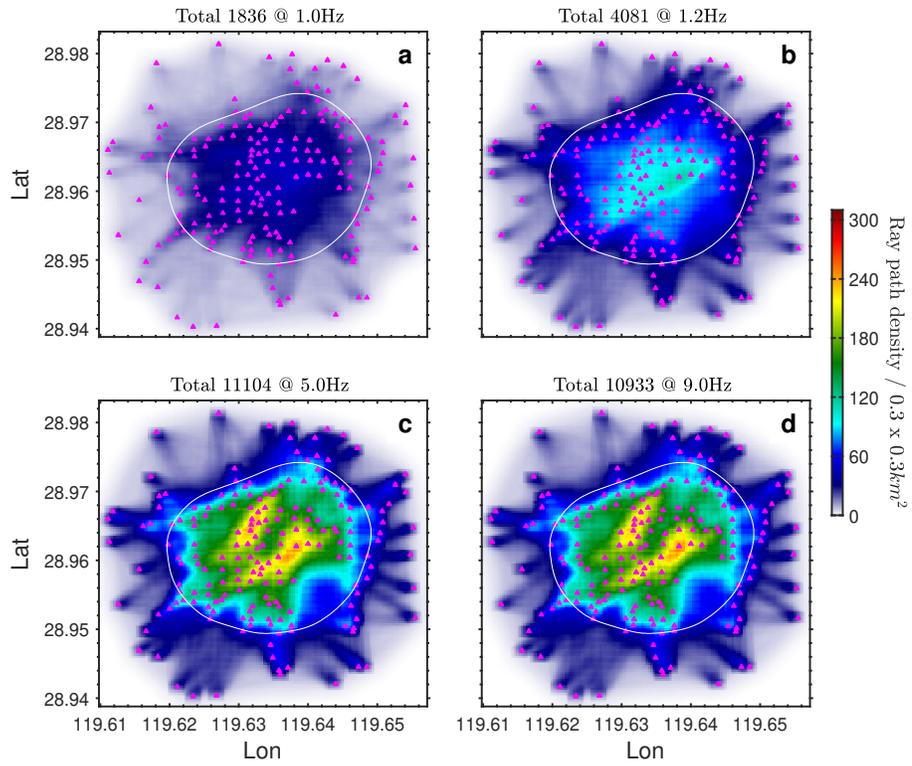


Figure S7. Raypath density maps of Love wave at different frequencies. The magenta triangles denote the seismic network. The white contour indicates the resolvable zone defined by raypath density map of Love wave at 5.0Hz. We label each sub-figure with the corresponding raypath number and frequency.

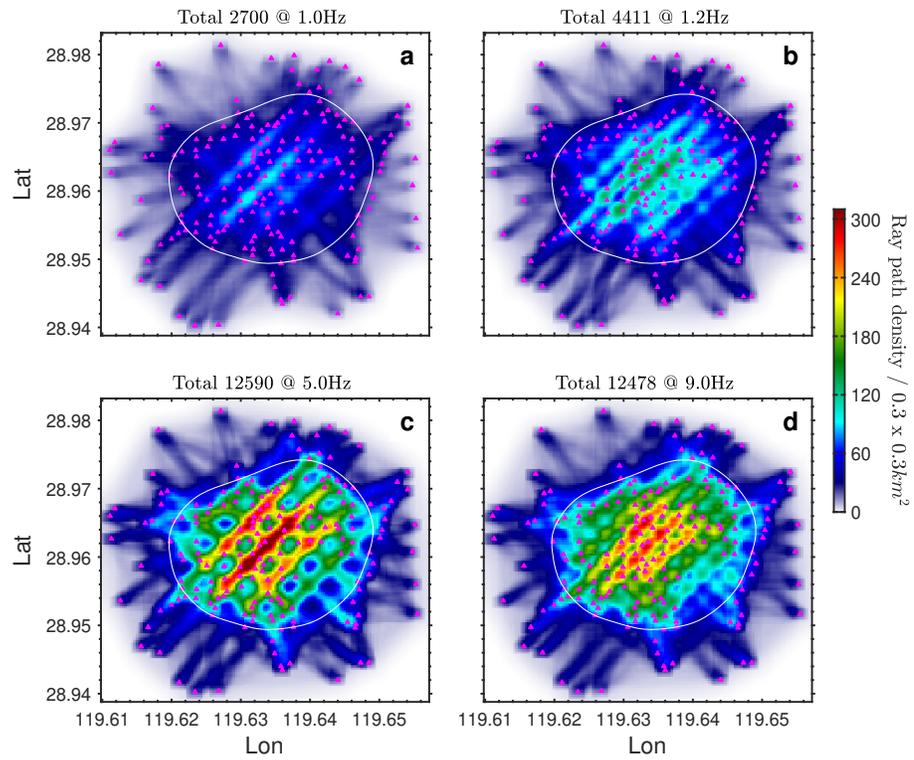


Figure S8. Raypath density maps of Rayleigh wave at different frequencies. The magenta triangles denote the seismic network. The white contour indicates the resolvable zone defined by raypath density map of Love wave at 5.0Hz. We label each sub-figure with the corresponding raypath number and frequency.

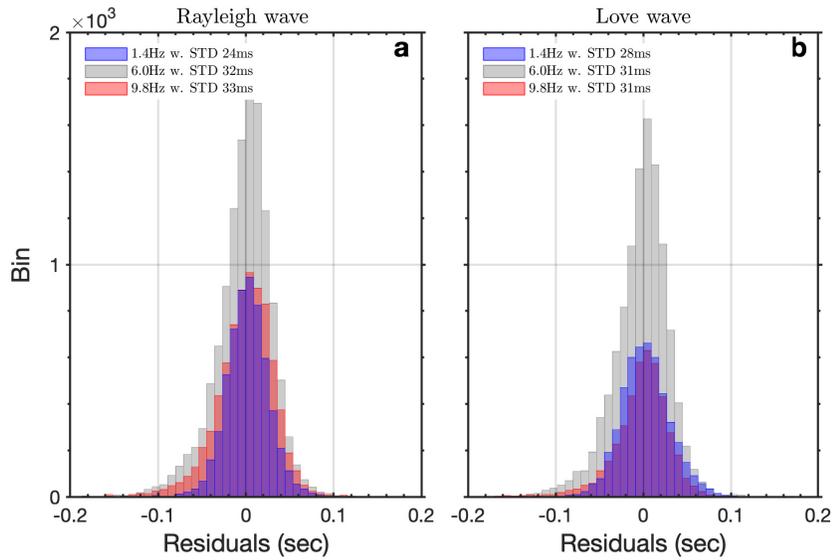


Figure S9. Histograms of the final residuals of surface wave traveltime tomography. (a) and (b) present the residuals for Rayleigh wave tomography and Love wave tomography. The different colors, blue, gray and red, indicate three different frequencies presented on Fig.12. The standard deviations are indicated on the legends.

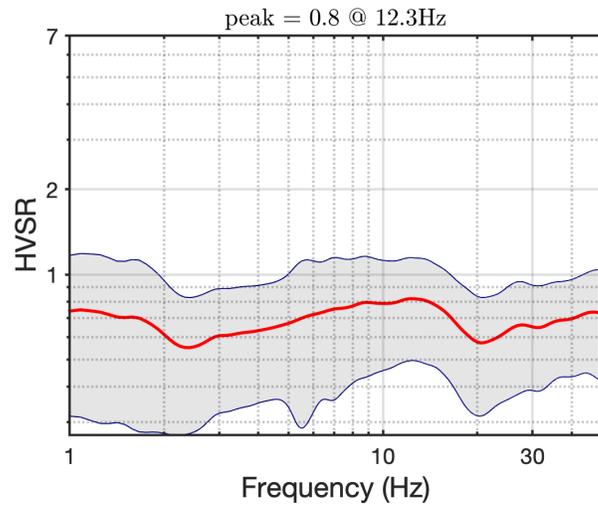


Figure S10. Example of the rejected HVSR curve.