

# Measurements of atmospheric HDO/H<sub>2</sub>O in southern California from CLARS-FTS

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## Abstract

Atmospheric isotopologues of water vapor (e.g., HDO) are important tracers for understanding Earth's hydrological cycles. Most remote sensing measurements of these isotopologues, however, are column averaged values and sparse in space and time. Measurements targeting the planetary boundary layer (PBL) are much rarer. In this study, we retrieved HDO and H from CLARS-FTS observations (2011-2019). The isotopological abundance  $\delta D$ , which represents the relative difference of the HDO/H<sub>2</sub>O ratio to a standard abundance ratio, is also calculated. The averaged  $\delta D$  retrievals are  $(-156.1 \pm 60.0)$  uncertainty of  $(6.1 + -10.2)$  an uncertainty of  $(42.4 + -31.6)$  In LA, the  $\delta D$  shows a seasonal cycle that is primarily driven by the change of atmospheric humidity. The temporal variabilities in  $\delta D$  data between CLARS-FTS and a collocated Total Carbon Column Observing Network (TCCON) observatory are highly correlated. The difference between CLARS and TCCON  $\delta D$  retrievals can primarily be attributed to the difference in their observation geometries. We conclude that the HDO and  $\delta D$  measurements from CLARS-FTS provide high spatial and temporal resolution datasets for further study of hydrological processes in the LA megacity.



32 **Abstract**

33 Atmospheric isotopologues of water vapor (e.g., HDO) are important tracers for understanding  
34 Earth's hydrological cycles. Most remote sensing measurements of these isotopologues,  
35 however, are column averaged values and sparse in space and time. Measurements targeting the  
36 planetary boundary layer (PBL) are much rarer. In this study, we retrieved HDO and H<sub>2</sub>O  
37 columns from observations by the California Laboratory for Atmospheric Remote Sensing  
38 Fourier Transform Spectrometer (CLARS-FTS), a mountaintop observatory on Mt. Wilson (1.67  
39 km a.s.l.) overlooking the Los Angeles (LA) basin in southern California. CLARS-FTS  
40 observations are highly sensitive to the lower atmosphere due to the long light path along the  
41 PBL. Retrievals were conducted using spectral windows between 6000-7000 cm<sup>-1</sup> from CLARS-  
42 FTS observations (2011-2019). The isotopological abundance  $\delta D$ , which represents the relative  
43 difference of the HDO/H<sub>2</sub>O ratio to a standard abundance ratio, is also calculated. The averaged  
44  $\delta D$  retrievals are  $(-156.1 \pm 60.0)\text{‰}$  with an uncertainty of  $(6.1 \pm 10.2)\text{‰}$  for LA Basin Survey m  
45 and  $(-344.7 \pm 95.0)\text{‰}$  with an uncertainty of  $(42.4 \pm 31.6)\text{‰}$  for Spectralon Viewing Observation  
46 mode. In LA, the  $\delta D$  shows a seasonal cycle that is primarily driven by the change of  
47 atmospheric humidity. The temporal variabilities in  $\delta D$  data between CLARS-FTS and a  
48 collocated Total Carbon Column Observing Network (TCCON) observatory are highly  
49 correlated. The difference between CLARS and TCCON  $\delta D$  retrievals can primarily be  
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51 measurements from CLARS-FTS provide high spatial and temporal resolution datasets for  
52 further study of hydrological processes in the LA megacity.

53

## 54 1. Introduction

55 Water vapor is both the most abundant greenhouse gas and an important component of  
56 the global hydrological cycle. Uncertainties in global concentrations of water vapor are a major  
57 challenge for global climate modeling, but observations of water vapor isotopologues provide  
58 additional information in constraining meteorological models and improving understandings of  
59 current weather processes and past climate events (Galewsky et al., 2016). Given the differences  
60 in binding energies for different molecular masses, concentrations of water isotopologues are  
61 influenced by phase changes in a process referred to as fractionation, in which heavier  
62 isotopologues are more likely to condense compared to lighter isotopologues. Fractionation  
63 allows measurements of water isotopologue concentrations to serve as a proxy for observing  
64 water vapor transport through the global hydrological cycle.

65 There are several naturally occurring stable oxygen ( $^{16}\text{O}$ ,  $^{17}\text{O}$ , and  $^{18}\text{O}$ ) and hydrogen ( $^1\text{H}$   
66 and  $^2\text{H}$  or D) isotopes, all of which can combine to form various stable water isotopologues. The  
67 lightest isotopologue ( $\text{H}_2^{16}\text{O}$ ) is the most abundant, but heavier isotopologues (e.g., HDO) are  
68 still observed commonly on Earth (Yoshimura, 2015). In liquid water, heavier isotopologues  
69 have higher binding energies and lower diffusive velocities, making them less likely to evaporate  
70 compared to lighter isotopologues (Craig, 1961; Craig et al., 1965). Therefore, the resulting  
71 water vapor tends to contain a smaller proportion of heavy isotopologues compared to the  
72 remaining liquid water, i.e. is more isotopically “depleted” than the liquid water (Galewsky et  
73 al., 2016). Similarly, when precipitation occurs, heavier isotopologues are more likely to  
74 condense, again leaving the remaining gaseous water vapor more isotopically depleted  
75 (Yoshimura, 2015). Such fractionation processes allow measurements of isotopologue  
76 concentrations to serve as a proxy for observing water vapor movement. Generally, a ratio, R, of  
77 heavy to light isotopologues, is defined to quantify relative concentrations. Using delta-notation  
78 “ $\delta$ ”, one can measure the deviation of a given isotopologue concentration (here D) from the  
79 standard composition of ocean water:

$$80 \quad \delta D = \left( \frac{R_D}{R_{D,\text{VSMOW}}} - 1 \right) * 1000 \text{ [‰]} \quad (1)$$

81 where,  $R_D$  and  $R_{D,\text{VSMOW}}$  are the ratios of the heavy to light isotopologue (HDO/ $\text{H}_2\text{O}$ ) in the  
82 sample and standard, respectively. The standard is defined using Vienna Standard Mean Ocean  
83 Water (VSMOW).  $R_{D,\text{VSMOW}} = 3.1152 \times 10^{-4}$  (Hagemann et al., 1970). The units of  $\delta D$  for a  
84 given isotopologue are in parts per thousand, or per mil (‰).  $\delta D$  values are generally negative,  
85 such that lower values, i.e. more negative, imply greater depletion of the sample in the heavy  
86 isotopologue, whereas higher  $\delta D$  values, i.e., closer to zero, imply greater enrichment compared  
87 to the sample.

88 In recent years, important advancements have been made in the field of remote sensing  
89 observations of stable water isotopologues, especially HDO. The growth in available global data  
90 sets from remote-sensing instruments along with the continual improvements of retrievals has in  
91 turn generated interest in data set validation and use of these data sets in modeling. One of the  
92 first instruments to demonstrate the potential of remote sensing to study water vapor  
93 isotopologues in the stratosphere was the Atmospheric Trace Molecule Spectroscopy (ATMOS)  
94 mission on the Space Shuttle (e.g., Kuang et al., 2003). Following this, instruments were  
95 developed that were sensitive to water vapor isotopologues in the troposphere. Greater spatial  
96 coverage of  $\delta D$  was provided by the Interferometric Monitor for Greenhouse gases (IMG) sensor  
97 on the Advanced Earth Observing Satellite 1 (ADEOS-1) (Zhakalov et al., 2004). Since then,

98 subsequent satellite instruments have improved the temporal and spatial resolution of  $\delta D$  data  
99 retrievals and increased the number of data sets (e.g., **Worden et al., 2007; Lacour et al., 2012**).  
100 For example, the Aura Tropospheric Emission Spectrometer (TES) and Atmospheric Infrared  
101 Sounder (AIRS) use thermal IR radiances to measure the HDO on a global scale (e.g., **Worden**  
102 **et al., 2007; Worden et al., 2019**). The Greenhouse Gas Observing Satellite (GOSAT) launched  
103 in 2009, houses a Fourier Transform spectrometer, Thermal and Near Infrared Sensor for Carbon  
104 Observations (TANSO-FTS), with the capability of retrieving global HDO and H<sub>2</sub>O  
105 concentrations (**Frankenberg et al., 2013; Boesch et al., 2013**). The Scanning Imaging  
106 Absorption Spectrometer for Atmospheric Chartography instrument (SCIAMACHY) aboard the  
107 ESA's environmental research satellite ENVISAT used near infrared (NIR) spectra to retrieve  
108 global  $\delta D$  values with high sensitivity in the lower troposphere, where most of the water vapor  
109 resides (**Frankenberg et al., 2009; Schneider et al., 2018**). Most recently, The Tropospheric  
110 Monitoring Instrument (TROPOMI) instrument on board Sentinel-5 also uses short-wave  
111 infrared spectra to make global total column measurements of HDO and  $\delta D$ , reporting  
112 improvements in the signal-to-noise ratio of observations compared to SCIAMACHY  
113 (**Schneider et al 2020**).

114 Ground-based remote sensing instruments, such as global Total Carbon Column  
115 Observing Network (TCCON; **Wunch et al., 2011**) which operates in the NIR, have also  
116 computed  $\delta D$  values from water vapor isotopologue retrievals (**Rokotyan et al., 2014**).  
117 Comparable measurements taken from Network for the Detection of Atmospheric Composition  
118 Change (NDACC) global tower network, which are similar to TCCON measurements in viewing  
119 geometry but use spectra from middle infrared, are incorporated into project MUSICA (Multi-  
120 platform remote Sensing of isotopologues for investigating the Cycle of Atmospheric Water),  
121 which includes measurements from ground-based, space-based, and in-situ instruments.  
122 MUSICA has performed valuable validation work of water vapor isotopologue measurements  
123 (**Schneider et al., 2016**). Specifically, MUSICA applies a bias correction to remote sensing data  
124 using vertical isotopologue profiles measured by well-calibrated in-situ instruments with low  
125 instrumental uncertainty (**Gonzales et al., 2016**). The MUSICA data product has been used for  
126 validation and bias correction of other remote sensing data sets (**Scheepmaker et al., 2015**).  
127 Overall, the availability of new data sets has allowed for the improvement of general circulation  
128 models (GCM) and prompted many studies aimed at better understanding complicated  
129 meteorological processes such as convection, cloud formation, and stratospheric-tropospheric  
130 exchange processes, along with the relative contribution of different global sources to  
131 atmospheric water vapor (**Yoshimura, 2015; Galewsky et al., 2016**).

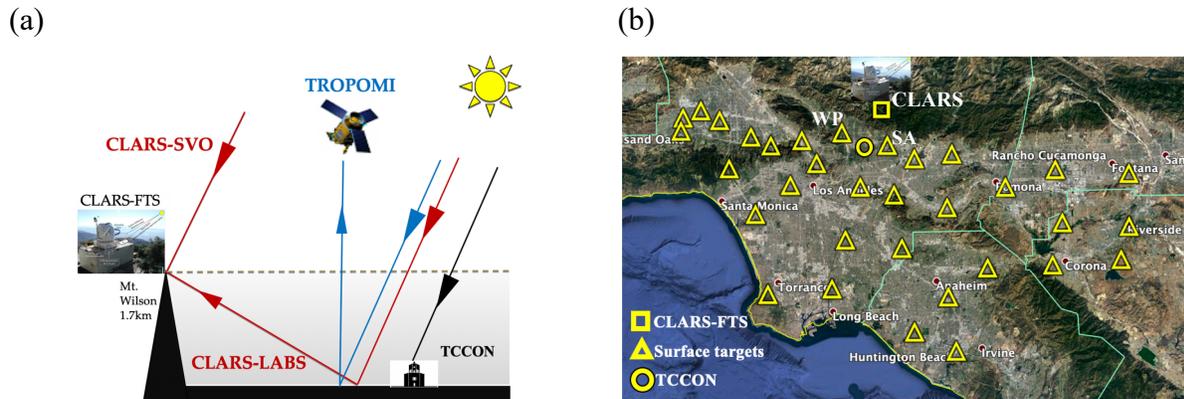
132 However, most measurements of these isotopologues from space-borne and ground-based  
133 remote sensing instruments are column averaged values with contributions from all altitudes. For  
134 thermal IR based satellites, e.g., TES or AIRS, the measurements have low sensitivity in the  
135 PBL. The measurements are also sparse in space and time on an urban scale like the Los Angeles  
136 (LA) basin in southern California. Measurements targeting the planetary boundary layer (PBL),  
137 the layer that couples the Earth's surface and the atmosphere above, are much rarer. In this study,  
138 we study the PBL-targeted measurements of HDO and  $\delta D$  in LA from observations by the  
139 California Laboratory for Atmospheric Remote Sensing - Fourier Transform Spectrometer  
140 (CLARS-FTS). Compared to conventional remote sensing observation networks, CLARS-FTS  
141 observations are highly sensitive to the lower atmosphere due to the long light path along the  
142 PBL.

143           In this paper, we first demonstrate that HDO and  $\delta D$  can be retrieved from CLARS-FTS  
144 spectra with sufficiently small fitting error and retrieval uncertainty. The entire spectral record  
145 from 2011 to 2019 was processed to provide a novel HDO, H<sub>2</sub>O, and  $\delta D$  dataset. Secondly, an  
146 examination of the temporal variability of XHDO, XH<sub>2</sub>O, and  $\delta D$  on annual and interannual  
147 timescales is performed using these resulting datasets. Thirdly, we compare CLARS-FTS  
148 retrievals to TCCON and TROPOMI retrievals at Caltech to demonstrate their consistency and  
149 discrepancy. Finally, we show that the discrepancies between CLARS-FTS and TCCON can be  
150 explained by the differences between their observation geometries.

151 **2. Data and methods**

152 **2.1 CLARS-FTS**

153



154 **Figure 1.** (a) Schematic figure showing the observations of CLARS-FTS, TROPOMI, and TCCON in the  
155 Los Angeles Basin; (b) The spatial distribution of CLARS-FTS surface reflection targets. The location of  
156 TCCON at Caltech, and the surface targets of West Pasadena (WP) and Santa Anita (SA) are also indicated.  
157

158 The CLARS-FTS instrument is located on Mt. Wilson at 1.67 km a.s.l. and makes daily  
159 observations of solar spectra reflected from 33 different surface targets distributed around the LA  
160 basin (**Figure 1(b)**). CLARS-FTS is a first-of-its-kind mountaintop observation system to  
161 monitor urban emissions by collecting surface reflected light from a top-down perspective.  
162 CLARS-FTS has two observing modes: Los Angeles Basin Survey (LABS) and Spectralon  
163 Viewing Observation (SVO) (**Figure 1(a)**). In the SVO mode, the spectrometer is pointed  
164 toward a Spectralon plate directly below the instrument. It receives the reflected sunlight to  
165 retrieve the column abundance of the atmosphere above the CLARS level. The SVO observation  
166 enables CLARS-FTS to measure background concentration during the day. In the LABS mode,  
167 the instrument is pointed towards one of the 33 surface reflection targets. The LABS observation  
168 mode has a longer light path in the PBL relative to satellite measurements and therefore higher  
169 sensitivity to urban emissions. The observation time for each surface target is about 3 minutes,  
170 which means high temporal resolution retrievals can be achieved. The spatial coverage spanned  
171 by the collection of reflection points provides a mapping capability over the entire LA basin.  
172 CLARS-FTS can perform one basin-wide scan in approximately 90 minutes and cycles through  
173 the entire measurement cycle around 5-8 times a day. Given its location above the top of the  
174 PBL, CLARS-FTS is a unique instrument which is both ground-based but employs a retrieval  
175 geometry similar to that of a geo-stationary satellite. The result of this instrument configuration,  
176 coupled with the frequency with which measurements are taken, yields a record of spectral  
177 observations possessing high spatial and temporal resolution, compared to other remote sensing  
178 instruments. Furthermore, as CLARS-FTS began taking daily measurements in 2011, the spectral  
179 record represents the longest available data record of atmospheric gases for the entire LA basin.  
180 A detailed description of the FTS and the surface reflection targets can be found in **Fu et al.**  
181 **(2014)** and **Wong et al. (2015)**.

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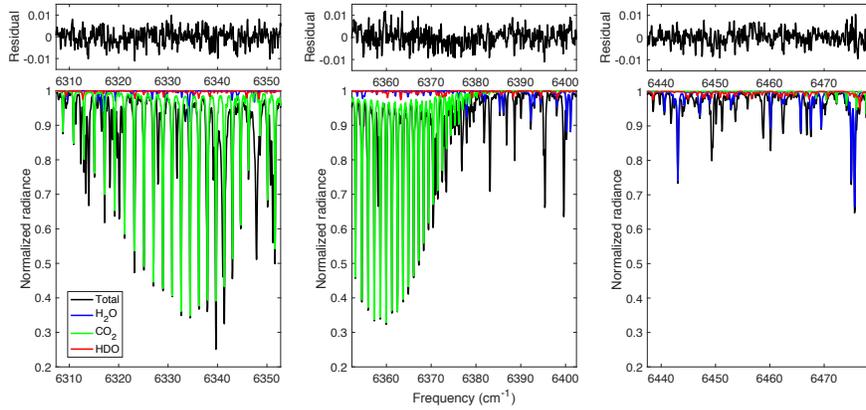
## 2.2 Retrievals of HDO, H<sub>2</sub>O, and $\delta D$

CLARS-FTS operates in the NIR from 4000 cm<sup>-1</sup> to 13500 cm<sup>-1</sup> with a spectral resolution of 0.06 cm<sup>-1</sup>. Recorded solar spectra are converted to slant-column densities (SCD), or total numbers of absorbing molecules per unit area along a Sun-Earth-instrument optical path, using a modified GFIT algorithm developed at JPL (Fu et al. 2014). The GFIT algorithm, within the GGG 2014 Software Suite, is employed by other ground-based remote sensing instruments, including TCCON FTS, for the retrieval of greenhouse gases (Wunch et al., 2011). GFIT provides a recommended list of spectral windows for HDO and H<sub>2</sub>O retrieval, along with associated input parameters. The broad spectral interval for HDO windows cover multiple HDO features, which leads to better and more consistent retrievals. **Figure A1** in **Appendix A** shows a comparison of the nine spectral window candidates by their contribution from HDO absorption to the overall gas absorption and the spectral fitting error for each spectral window using a set of ~4000 observations from 4 distinct days in 2013. We rejected the six spectral windows between 4000 cm<sup>-1</sup> to 6000 cm<sup>-1</sup> because of their large fitting error and used the three spectra windows between 6000 cm<sup>-1</sup> and 7000 cm<sup>-1</sup> that are more robust, as shown in **Table 1**. For H<sub>2</sub>O, we selected the five spectral windows between 6200 cm<sup>-1</sup> and 6500 cm<sup>-1</sup> from the TCCON list (Wunch et al., 2015) that are close to the HDO windows, as shown in Table 1. These spectral windows have been fully tested for TCCON observations. In this study, we further evaluate these spectral windows for CLARS observations by their fitting errors and retrieval uncertainties. Examples of spectral fit, including their fitting residuals and contributions from target gases as well as interfering gases, are shown in **Figure 2** for HDO windows and in **Appendix B** for H<sub>2</sub>O windows.

Table 1. Spectral Windows for HDO and H<sub>2</sub>O with Associated Parameters

| Gas              | Center (cm <sup>-1</sup> ) | Width (cm <sup>-1</sup> ) | Gases to fit                            | Continuum Basis Functions        |
|------------------|----------------------------|---------------------------|---|----------------------------------|
| HDO              | 6330.05                    | 45.50                     | HDO, H <sub>2</sub> O, CO <sub>2</sub>  | 2-order polynomial for continuum |
|                  | 6377.40                    | 50.20                     | HDO, H <sub>2</sub> O, CO <sub>2</sub>  |                                  |
|                  | 6458.10                    | 41.40                     | HDO, H <sub>2</sub> O, CO <sub>2</sub>  |                                  |
| H <sub>2</sub> O | 6255.95                    | 3.60                      | H <sub>2</sub> O, CO <sub>2</sub> , HDO | 2-order polynomial for continuum |
|                  | 6301.35                    | 7.90                      | H <sub>2</sub> O, CO <sub>2</sub> , HDO |                                  |
|                  | 6392.45                    | 3.10                      | H <sub>2</sub> O, HDO                   |                                  |
|                  | 6401.15                    | 1.15                      | H <sub>2</sub> O, HDO, CO <sub>2</sub>  |                                  |
|                  | 6469.60                    | 3.50                      | H <sub>2</sub> O, CO <sub>2</sub> , HDO |                                  |

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209 **Figure 2.** Examples of CLARS-FTS spectral windows of (left)  $6330.05 \text{ cm}^{-1}$ , (center)  $6377.40 \text{ cm}^{-1}$ , and  
 210 (right)  $6458.10 \text{ cm}^{-1}$  for retrieving HDO in this study. These samples of normalized spectra are taken from  
 211 a mid-day observation on 7/14/2013 over the West Pasadena surface target. The lower panel shows the full  
 212 spectral fit with total contribution in black, contribution from HDO in red, from H<sub>2</sub>O in blue, and from CO<sub>2</sub>  
 213 in green. The upper label shows the residuals of the spectral fits, defined as the difference in total measured  
 214 and total calculated radiance. Similar examples of spectral fit for H<sub>2</sub>O are shown in the Appendix **Figure**  
 215 **B1**.  
 216

217 Using these optimized spectral windows, we separately retrieved HDO and H<sub>2</sub>O SCDs  
 218 using the CLARS GFIT algorithm. For each spectral observation, CLARS GFIT outputs four  
 219 numerical results with which to calculate a SCD for HDO and H<sub>2</sub>O: an original vertical column  
 220 density (OVC) in unit of molecules/cm<sup>2</sup>, an air mass value (AM), a volume mixing ratio (VMR)  
 221 scale factor (VSF), and the error in the VSF (VSF error). The air mass value represents the  
 222 number of vertical columns the light travels through in its slant column path. The first three  
 223 results are multiplied together to determine the SCD:

$$224 \quad \text{HDO}_{\text{SCD}} = \text{AM} \times \text{OVC} \times \text{VSF} \quad (2)$$

225 The same calculation is conducted to produce H<sub>2</sub>O SCD. The uncertainty in the SCD is  
 226 determined using the same formula, but using VSF error in place of VSF, according to  
 227 conventional rules of error propagation. Dry-air column averaged mixing ratios of HDO  
 228 (XHDO) and H<sub>2</sub>O (XH<sub>2</sub>O) are computed from the retrieved SCDs by normalizing the SCD  
 229 measurements to the dry-air total column, which can be derived from the measured SCD for O<sub>2</sub>  
 230 and the dry-air O<sub>2</sub> mole fraction:

$$231 \quad \text{XHDO} = 0.2095 \times \frac{\text{SCD}_{\text{HDO}}}{\text{SCD}_{\text{O}_2}} \quad (3)$$

232 Using this method improves mixing ratio measurements since any existing systematic  
 233 errors in retrievals of both HDO or H<sub>2</sub>O and O<sub>2</sub> SCDs will be minimized in computing the ratio  
 234 (Fu et al., 2014). In this study, O<sub>2</sub> is retrieved using a spectral window centered on  $7885 \text{ cm}^{-1}$ ,  
 235 whose retrieval results yield very low spectral fitting residual and VSF error values (Zeng et al.,  
 236 2020). XH<sub>2</sub>O is retrieved using the spectral windows in the same wavelength range as XHDO.  
 237 This method reduces overall uncertainty in the final  $\delta D$  result by avoiding the possible  
 238 complication of wavelength dependent noise such as that from the aerosol scattering effect  
 239 (further discussion in Section 4.1). Ratioing of XHDO and XH<sub>2</sub>O values to some extent cancels  
 240 out possible errors existing in both XHDO and XH<sub>2</sub>O retrievals. Finally, XHDO and XH<sub>2</sub>O

241 measurements can be used to compute  $\delta D$  (**Equation 1**) for each CLARS-FTS observation,  
242 yielding a data set extending from 2011-2019 and spanning the LA basin.

### 243 **2.3 Calibration of XH<sub>2</sub>O, XHDO, and $\delta D$ from CLARS-GFIT algorithm**

244 Since CLARS-FTS retrievals employ the same GFIT algorithm as TCCON, the  
245 calibration developed by TCCON can also be applied to CLARS-FTS retrievals. As shown in  
246 **Wunch et al. (2015)**, XH<sub>2</sub>O retrievals from TCCON observations have been compared against  
247 radiosonde measurements, resulting in a bias correction factor of 1/1.0183 being applied to  
248 TCCON XH<sub>2</sub>O values. For XHDO retrievals, **Schneider et al. (2020)** derived a correction  
249 factor by scaling the TCCON XHDO to match the calculated  $\delta D$  between TCCON and the  
250 MUSICA dataset, whose  $\delta D$  profiles have been validated against aircraft measurements  
251 (**Schneider et al., 2016**). As a result, an error-weighted average correction factor of 1/1.0778 for  
252 XHDO was derived based on multiple TCCON sites. These two scale factors are used to  
253 calibrate XH<sub>2</sub>O and XHDO retrievals, respectively, from CLARS-FTS, and the resulting  
254 calibrated  $\delta D$  is re-generated.

### 255 **2.4 HDO and $\delta D$ observations from TCCON and TROPOMI**

#### 256 **2.4.1 TCCON**

257 The TCCON FTS measures direct solar spectra in the NIR and retrieves the column-  
258 averaged abundances of many atmospheric gases, including H<sub>2</sub>O and HDO, using the GFIT  
259 algorithm. XH<sub>2</sub>O reported in the TCCON data product is the mean of retrievals from fifteen  
260 spectral windows and for XHDO it is from six spectral windows. A detailed introduction of the  
261 configuration of TCCON FTS, the characteristics of the observed spectra, and the calibration of  
262 the GFIT retrievals using aircraft measured profiles can be found in **Wunch et al. (2011)**. As  
263 discussed in **Section 2.3**, for TCCON XH<sub>2</sub>O, a correction factor of 1/1.0183 is applied, and for  
264 XHDO, a correction factor 1/1.0778 is applied. The XH<sub>2</sub>O, XHDO, and  $\delta D$  data were collected  
265 from 2012 to 2019 by the TCCON-Caltech site (**Figure 1(a)**; **Wennberg et al., 2015**) on the  
266 Caltech campus.

#### 267 **2.4.2 TROPOMI**

268 The TROPOMI instrument on board Sentinel-5 uses short-wave infrared spectra at 4225  
269 cm<sup>-1</sup> (i.e., 2.3  $\mu\text{m}$ ) to retrieval global total column measurements of HDO, H<sub>2</sub>O, and  $\delta D$   
270 (**Schneider et al., 2020**). The measurements have a daily coverage (overpass at around 1:30 pm  
271 local time) and a spatial resolution of up to 3.5 km $\times$ 7 km at nadir. The retrievals are filtered by  
272 strict criteria to exclude measurements contaminated by clouds and aerosols. A detailed  
273 description of the instruments, retrieval algorithm, and data screening can be found in **Schneider**  
274 **et al. (2020)**. From a comparison with collocated TCCON measurements, the TROPOMI  
275 retrievals have a mean bias of  $(-1.1\pm 7.3)\%$  for XHDO and  $(-14\pm 17)\%$  for  $\delta D$ . The  $\delta D$  dataset  
276 available from late 2017 to 2019 in the LA region is used in this study. To allow more  
277 measurements for comparison, we relaxed the cloud filter slightly from 1% to 5% in the filters.  
278 Averaged  $\delta D$  before and after relaxing the filter are the same (about -236‰), indicating no bias  
279 caused. In total, there are 635 valid observations over the study area shown in **Figure 1(b)**.

### 280 3. Results

#### 281 3.1 Retrieval and spectra fitting errors from CLARS-FTS

282 Based on preliminary investigation of the quality of spectral fittings and the retrieval  
283 uncertainties in HDO and H<sub>2</sub>O mixing ratios, the entire spectral record of CLARS-FTS from  
284 2011 to 2019, was reprocessed using the spectral windows with central wavenumbers in the  
285 range 6000-7000 cm<sup>-1</sup>. The CLARS-GFIT parameters used are the same as those shown **Table 1**.  
286 Before additional analysis was performed on HDO and H<sub>2</sub>O observations, the data were passed  
287 through a series of filters, as summarized in **Table 2**. Data with poor spectral fitting, identified as  
288 instances with large solar zenith angles (SZA), low signal-to-noise (SNR) ratios, and large root-  
289 mean-square-error (RMSE) values from the spectral fitting, are removed. Additionally, the ratio  
290 between retrieved and geometric O<sub>2</sub> slant column densities (O<sub>2</sub> ratio) are used to remove  
291 retrievals affected by cloud and aerosol scattering. The geometric O<sub>2</sub> SCD is calculated  
292 assuming no scattering occurs, along with additional assumptions outlined in **Fu et al. (2014)**.  
293 Because oxygen is well-mixed in the atmosphere, deviations in the retrieved O<sub>2</sub> SCD from the  
294 geometric O<sub>2</sub> SCD implies variations in the light path due to clouds and/or aerosols (**Zeng et al.,**  
295 **2018; Zeng et al., 2020**), and can therefore be used to identify observations that represent  
296 especially cloudy or hazy days. Along with the above criteria, retrievals with high uncertainty  
297 values, defined VSF error values (one of the outputs from CLARS-GFIT) as more than two times  
298 the standard deviation from the mean VSF error, are also removed. This filter helps to remove  
299  $\delta D$  results which would necessarily have very high uncertainties, since VSF error is propagated  
300 through calculation of SCDs, mixing ratios, and eventually  $\delta D$  values. Note that the VSF error is  
301 the uncertainty of the profile scaling factor while the fitting RMSE is the error of residuals from  
302 the spectral fit.

**Table 2. Data filters for HDO and H<sub>2</sub>O retrievals from CLARS-GFIT**

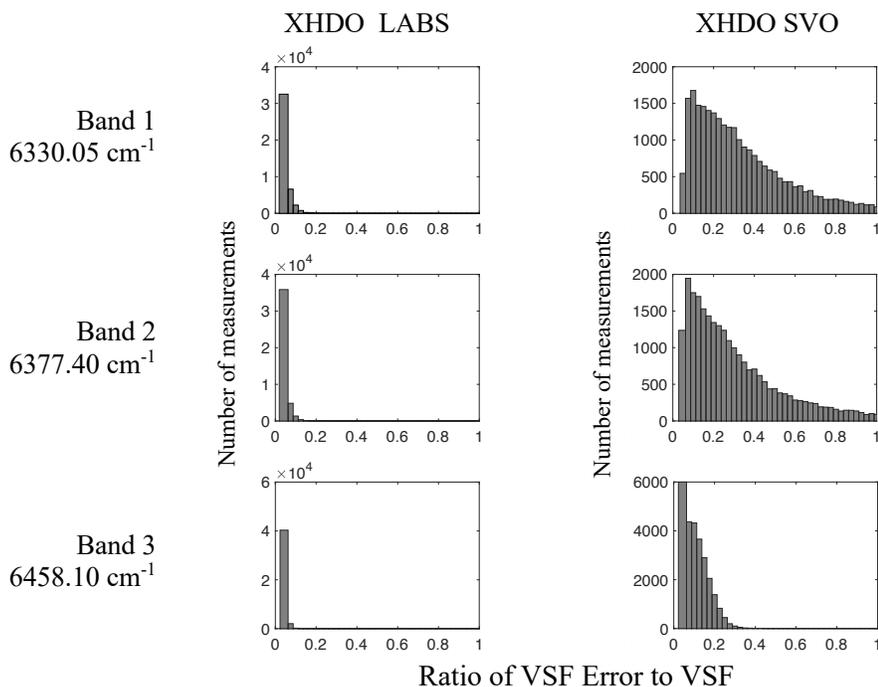
| <b>Filters</b>             | <b>Selection Criterion</b>                                      |
|----------------------------|---|
| Low clouds and/or aerosols | LABS O <sub>2</sub> ratio between 0.9 and 1.04                  |
| High clouds                | SVO O <sub>2</sub> ratio between 1.0 and 1.08                   |
| Large SZA                  | SZA less than 70 degrees  |
| Low SNR                    | SNR larger than 100   |
| Poor spectral fit          | Spectral fitting RMSE less than 1 standard deviation above mean |
| High VSF Error             | VSF error less than 2 standard deviation above mean             |

303

304 **Figure 3** shows relevant statistics for the entire data set once it has been filtered  
305 according to the above criteria. The figure shows the fraction VSF error in retrievals, a  
306 quantification of retrieval uncertainty, for HDO, where the data is separated into histograms  
307 according to observation modes (LABS and SVO) and spectral windows (6330.05 cm<sup>-1</sup>, 6377.40  
308 cm<sup>-1</sup>, and 6458.10 cm<sup>-1</sup>). For LABS retrievals, one can observe that the fraction VSF error in  
309 both HDO and H<sub>2</sub>O yield similar distributions for the three spectral windows. In other words,  
310 retrievals from the three spectral windows have similar retrieval uncertainties for LABS  
311 measurements. In addition, the majority of retrievals have VSF error values less than 10% for  
312 HDO. However, this consistency is not observed for SVO measurements, where the distributions  
313 for the two higher wavenumber spectral windows (6377.40 cm<sup>-1</sup> and 6458.10 cm<sup>-1</sup>) have very  
314 long tails. The reason maybe at these two windows are contributions from interference due to  
315 other gases, mainly CO<sub>2</sub> and H<sub>2</sub>O as shown in **Figure 2**, are much stronger than HDO for the

316 portion of the atmosphere above CLARS. As a result, retrieved VSFs for HDO are associated  
 317 with large uncertainty even as the fitting errors are small, as shown in the following **Figure 4**.  
 318 For H<sub>2</sub>O as shown in **Figure B2**, most retrievals, except for the 6255.95 cm<sup>-1</sup> spectral window,  
 319 have VSF error less than 10% for LABS and 20% for SVO. The 6255.95 cm<sup>-1</sup> spectral window,  
 320 however, shows larger VSF error, especially for SVO. We therefore conclude that this window is  
 321 not robust for H<sub>2</sub>O retrievals and are not included in the following analysis.

322



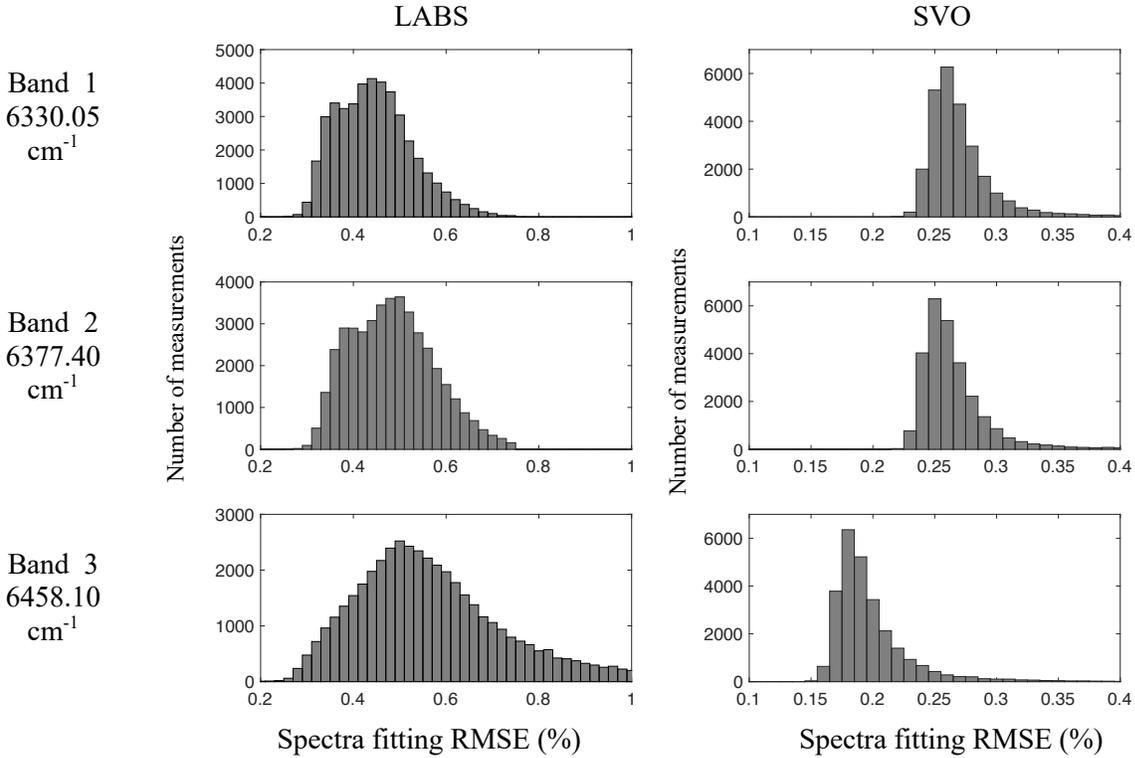
323 **Figure 3.** Retrieval error, in ratio of VSF error to VSF, for the entire filtered XHDO datasets from 2011 to  
 324 2019. Both VSF and VSF error values are calculated from CLARS-GFIT. The data are separated into  
 325 histograms according to observation modes (LABS and SVO) and three spectral windows (6330.05 cm<sup>-1</sup>,  
 326 6377.40 cm<sup>-1</sup>, and 6458.10 cm<sup>-1</sup>). The retrieval errors for H<sub>2</sub>O data are shown in the Appendix **Figure B2**.

327

328 **Figure 4** shows the histograms of RMSE from spectral fitting for the entire HDO data  
 329 set, again with SVO and LABS observations separated. In the LABS histograms, we do see  
 330 slightly different distributions when comparing the two lower frequency (6330.05 cm<sup>-1</sup> and  
 331 6377.40 cm<sup>-1</sup>) spectral windows to the 6458.10 cm<sup>-1</sup> spectral window, which tends to have  
 332 greater counts of higher RMSE. For SVO, the fitting errors are in general less than those for  
 333 LABS, because SVO mode is measuring the background concentration above PBL with small  
 334 perturbations such as impacts from aerosol and surface scattering. However, the 6330.05 cm<sup>-1</sup>  
 335 and 6377.40 cm<sup>-1</sup> retrievals have significantly larger fitting errors compared to the 6458 cm<sup>-1</sup>  
 336 spectral window. These discrepancies are consistent with what is shown in **Figure 3** for the  
 337 retrieval uncertainty. Therefore, we conclude that the 6330.05 cm<sup>-1</sup> and 6377.40 cm<sup>-1</sup> retrievals  
 338 for SVO mode are not robust for HDO retrievals. For the following analysis, we used the  
 339 6458.10 cm<sup>-1</sup> spectral window only for SVO HDO retrievals. For H<sub>2</sub>O spectral windows as  
 340 shown in **Figure B3**, all windows show consistent RMSE for both LABS and SVO.

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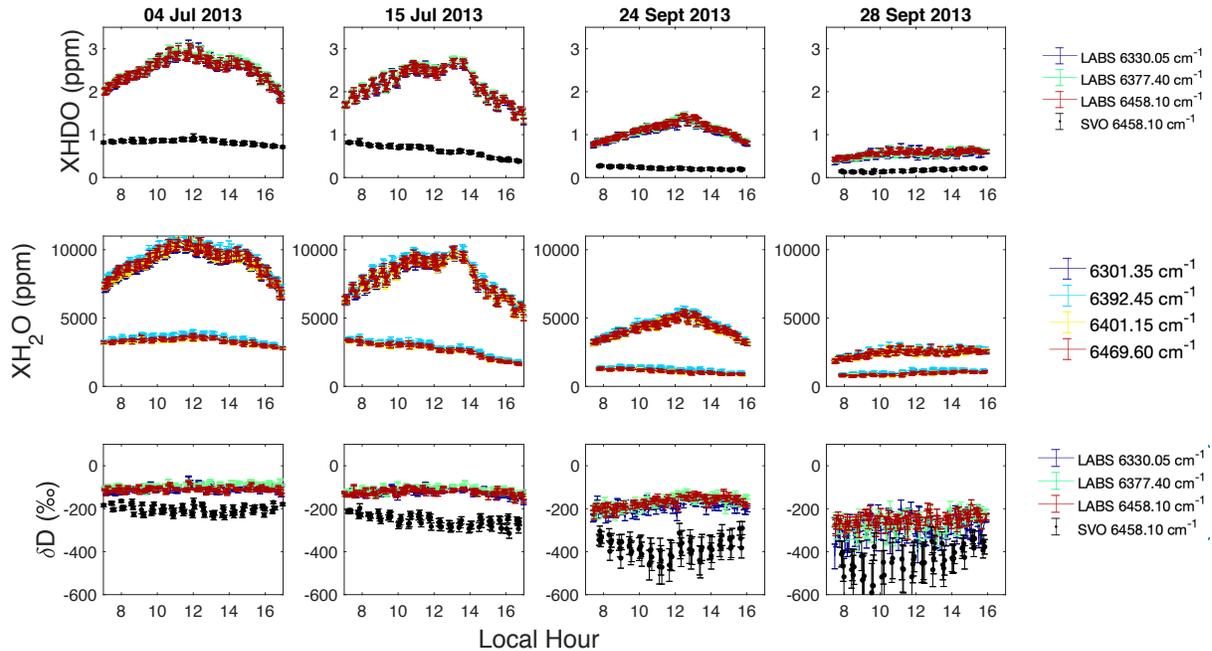


343 **Figure 4.** Histograms of RMS error from spectral fitting of HDO windows for the entire data set from 2011  
 344 to 2019. The data are separated into histograms according to observation mode (LABS and SVO) and three  
 345 spectral windows ( $6330.05\text{ cm}^{-1}$ ,  $6377.40\text{ cm}^{-1}$ , and  $6458.10\text{ cm}^{-1}$ ). Similar figures for  $\text{H}_2\text{O}$  windows are  
 346 shown in the Appendix **Figure B3**.  
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### 348 **3.2 HDO, $\text{H}_2\text{O}$ , and $\delta\text{D}$ retrieval from CLARS-FTS**

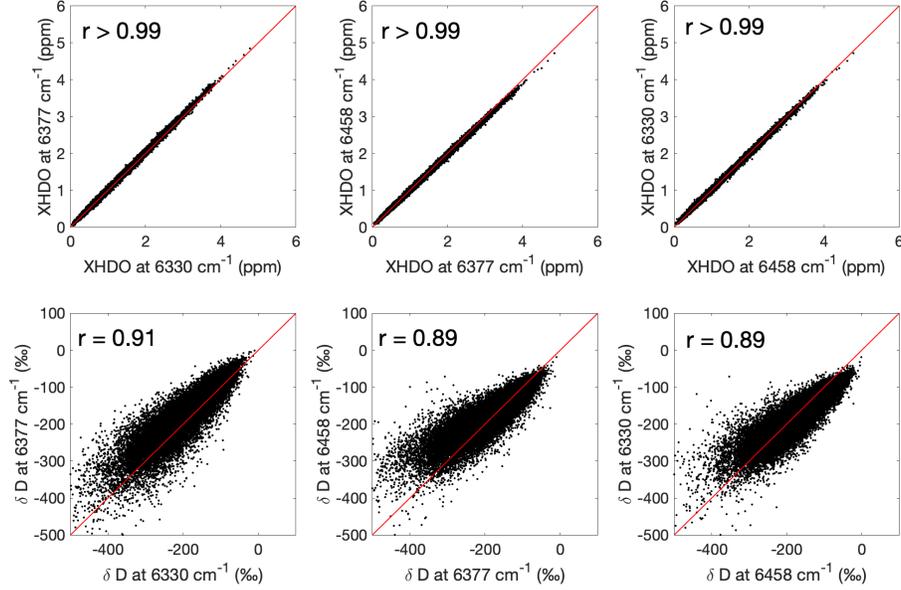
349 **Figure 5** shows the diurnal observations of XHDO, XH<sub>2</sub>O, and  $\delta\text{D}$  values on the sampled  
 350 set of retrievals from four days in 2013. Over these four days, CLARS LABS was targeting two  
 351 closer surface reflection targets: West Pasadena and Santa Anita. The LABS XHDO results from  
 352 three spectral windows are differentiated by color and the SVO XHDO results are plotted in  
 353 black. XH<sub>2</sub>O results from four spectral windows are shown. To derive  $\delta\text{D}$  shown in **Figure 5**,  
 354 weighted averaged XH<sub>2</sub>O are calculated first over all the spectral windows. Equation (1) is then  
 355 applied to obtain the  $\delta\text{D}$  values for each of the three XHDO retrievals. From **Figure 5**, we  
 356 observe relatively good agreement in the retrievals from the three spectral windows, which also  
 357 suggests these retrieval results are robust. Furthermore, uncertainties in retrieval results, shown  
 358 with error bars, are relatively small for all spectral windows on each of the four days. The time  
 359 series indicates that XHDO and XH<sub>2</sub>O values exhibit both diurnal and seasonal variability for the  
 360 LABS observations. For all four days, XHDO and XH<sub>2</sub>O increase consistently until mid-day and  
 361 then decrease in the late afternoon into the evening, as is reasonable based on diurnal  
 362 temperature patterns. In July, we note that the XHDO ranges between 1–4 ppm and XH<sub>2</sub>O  
 363 ranges between 5,000 and 10,000 ppm whereas in September, the concentrations of both species  
 364 are smaller. For SVO, both XHDO and XH<sub>2</sub>O are lower than for LABS. The results of  $\delta\text{D}$  are  
 365 also shown for LABS and SVO. The  $\delta\text{D}$  shows a smaller diurnal variability compared to XHDO  
 366 and XH<sub>2</sub>O. The SVO  $\delta\text{D}$  is smaller than the LABS data. This is consistent with the vertical  
 367 distribution of  $\delta\text{D}$ , in which  $\delta\text{D}$  generally decreases (more depleted) with elevation (**Galewsky**

368 et al., 2014). The seasonal change of  $\delta D$  is associated with the changes in specific humidity,  
 369 which is illustrated in Section 3.5.



370 **Figure 5.** Diurnal observations of XHDO, XH<sub>2</sub>O, and  $\delta D$  values on the sampled set of retrievals from four  
 371 days in 2013, where LABS results from the different spectral windows are differentiated by color. These  
 372 data are combined observations from West Pasadena and Santa Anita surface targets, which were two closer  
 373 ones targeted over these four days by CLARS-FTS. For XHDO and  $\delta D$ , the SVO data are from band  
 374 6458.10 cm<sup>-1</sup> only, while for XH<sub>2</sub>O, the SVO data (lower time series) are from all available bands as LABS.  
 375 Some data are not visible because of overlapping. Error bars show uncertainty values in individual  
 376 observations. To derive  $\delta D$ , weighted averaged XH<sub>2</sub>O are calculated first over all the spectral windows.  
 377 Equation (1) is then applied to obtain the  $\delta D$  values for each of the three XHDO retrievals.

378  
 379 HDO and H<sub>2</sub>O SCDs, XHDO and XH<sub>2</sub>O, and  $\delta D$  values are computed separately from  
 380 each spectral window. **Figure 6** shows the correlation results of LABS XHDO and  $\delta D$  from one  
 381 spectral window versus another. The three columns show the three correlations that can be done  
 382 using the set of three spectral windows. The three rows are HDO and  $\delta D$  correlation results. As  
 383 one can see from **Figure 6**, there is very good agreement between mixing ratio results among the  
 384 three spectral windows. Visually, each of the correlations appears as a straight line with a slope  
 385 of approximately 1. This is verified quantitatively by the high correlation coefficients indicated  
 386 for each correlation plot. Furthermore, the  $\delta D$  values from different spectral windows are also  
 387 significantly correlated. The reason the correlation for  $\delta D$  becomes weaker is because relative  
 388 difference in  $\delta D$  between spectral windows is amplified when ratioing against XH<sub>2</sub>O following  
 389 equation (1), especially when XH<sub>2</sub>O is small. The point-by-point absolute difference is about  
 390 (25.7±24.1) % on average. The window-to-window XH<sub>2</sub>O retrievals (correlation plots not shown  
 391 here) are also highly consistent, as can be observed from examples in **Figure 5**. These results  
 392 suggest that averaging the LABS retrieval results from the three spectral windows will not bias  
 393 the overall determination of mixing ratio, and in turn will reduce overall uncertainty in the  $\delta D$   
 394 values.



395 **Figure 6.** Correlations of HDO and  $\delta D$  retrievals from all the LABS observations for one spectral window  
 396 versus another. In order of column, the correlations are:  $6377\text{ cm}^{-1}$  vs  $6330\text{ cm}^{-1}$ ,  $6458\text{ cm}^{-1}$  vs  $6330\text{ cm}^{-1}$ ,  
 397  $6458\text{ cm}^{-1}$  vs  $6377\text{ cm}^{-1}$ , respectively. The two rows show XHDO correlations and  $\delta D$  correlations,  
 398 respectively. The 1:1 line in red and the correlation coefficients ( $r$ ) are also indicated for  $\delta D$ . For XHDO,  
 399 the correlation coefficients all larger than 0.99.

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### 3.3 Seasonal cycles of HDO, $\text{H}_2\text{O}$ , and $\delta D$

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Given the LABS observations that retrieval uncertainties are not necessarily identical for all spectral windows (**Figure 5**), a weighted average rather than a simple average is used to compute a final mixing ratio value for each observation, given by (using XHDO as an example):

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$$\overline{\text{XHDO}} = \frac{\sum_{i=1}^3 \text{XHDO}_i \times w_i}{\sum_{i=1}^3 w_i} \text{ where } w_i = \frac{1}{\sigma_i^2} \quad (4)$$

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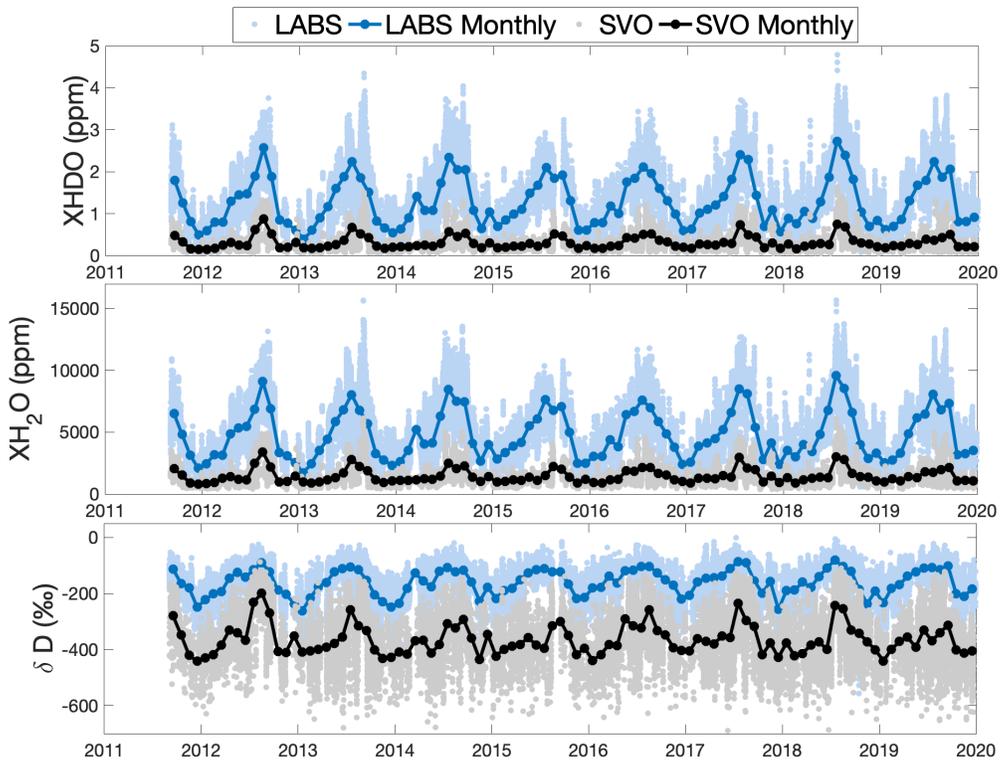
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The weight for the mixing ratio value associated with each spectral window is defined as the reciprocal of the retrieval uncertainty, which in turn is defined using conventional error propagation of SCD uncertainties for the water vapor isotopologue and  $\text{O}_2$  according to **Equation 3**. The same calculations (**Equation 4**) are made for  $\text{XH}_2\text{O}$  and  $\delta D$  for both LABS and SVO observations. **Figure 7** shows the time series of weighted means of XHDO,  $\text{XH}_2\text{O}$ , and  $\delta D$ . For SVO XHDO, only the  $6458.10\text{ cm}^{-1}$  spectral window results are shown, as explained in **Section 3.1**. The averaged  $\delta D$  retrievals are  $(-156.1 \pm 60.0)\text{‰}$  for LABS and  $(-344.7 \pm 95.0)\text{‰}$  for SVO. The  $\delta D$  retrieval uncertainties are  $(6.1 \pm 10.2)\text{‰}$  for LABS and  $(42.4 \pm 31.6)\text{‰}$  for SVO. From **Figure 7**, we can observe first that XHDO and  $\text{XH}_2\text{O}$  tend to be smaller for SVO observations compared with LABS observations. This is reasonable given the fact that  $\text{XH}_2\text{O}$  decreases with altitude due to decreasing temperatures and pressures. More evaporation also implies decreasing concentrations of the heavy isotopologue, which is again consistent with the plot in that XHDO and  $\delta D$  are smaller for SVO observations. On average, the difference between monthly averaged LABS and SVO measurements is  $(204.9 \pm 35.2)\text{‰}$ . We also find that the SVO  $\delta D$  values are more variable compared to the LABS measurements. This may be because of the much thinner atmosphere above PBL and therefore they more sensitive to any perturbations to the water vapor abundance.

423 From the first two panels, we see that XHDO and XH<sub>2</sub>O values reach minimum values in  
 424 the winter and maximum values in the summer months. This is again reasonable based on  
 425 interannual patterns of temperature. From the last panel, we see that δD also reaches a minimum  
 426 in the winter, which is interpreted as the time of greatest depletion in HDO. The δD value  
 427 increases to around zero in the summer, implying when observed fractionation levels in  
 428 atmospheric water vapor are approximately equivalent to that of the standard VSMOW. On  
 429 average, the δD for the peak months from July to September is -112.4‰ and bottom months  
 430 from December to February is -208.9‰. According to Rayleigh distillation (Rayleigh et al.  
 431 1902), there is preferential condensation of the heavier isotopologues (HDO here) compared with  
 432 H<sub>2</sub>O when the water vapor mixing ratio is lower. As a result, seasonal variations in δD are driven  
 433 by the variations in humidity to the first order. This pattern fits with the seasonal temperature  
 434 change for Los Angeles, higher temperature and absolute humidity (water content) from spring  
 435 to summer and lower temperature and absolute humidity in the wintertime. The seasonal  
 436 differences may also be due to varied contributions from different sources. For instance, there is  
 437 likely more surface evaporation with higher δD when it is hotter during summer seasons.  
 438 Therefore, lower absolute humidity in winter is associated with greater depletion in the heavy  
 439 isotopologue, as is expected. The correlation between absolute humidity and δD is further  
 440 explored in Section 3.5. Overall, CLARS-FTS provides continuous and robust estimations of  
 441 XH<sub>2</sub>O, XHDO, and δD for the LA basin.



442 **Figure 7.** Time series of XHDO, XH<sub>2</sub>O, and δD observed by CLARS-FTS from Sept. 2011 to Dec. 2019.  
 443 LABS measurements are the weighted means from retrievals of the three spectral windows (6330.05 cm<sup>-1</sup>,  
 444 6377.40 cm<sup>-1</sup>, and 6458.10 cm<sup>-1</sup>). The weights are calculated by their retrieval uncertainty. The SVO  
 445 measurements are from retrievals using 6458.10 cm<sup>-1</sup> spectral window only. The monthly means of these  
 446 measurements are also shown.  
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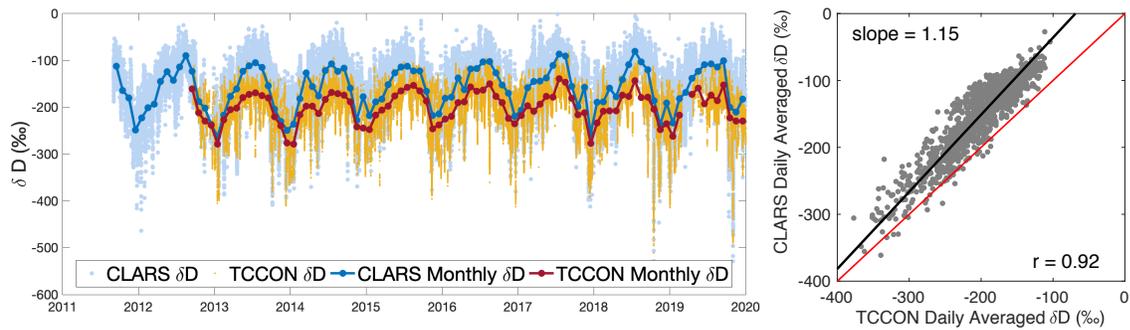
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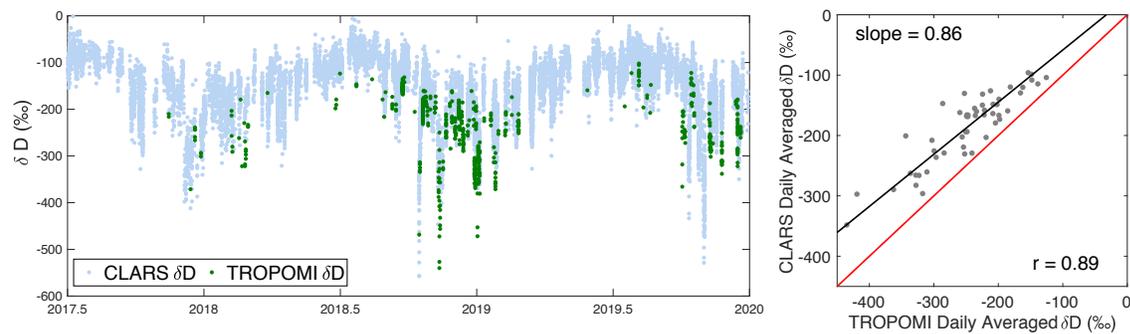
### 3.4 Comparison of $\delta D$ with TCCON and TROPOMI

Measurements of  $\delta D$  from CLARS, TCCON, and TROPOMI recorded at the same time and location are expected to differ due to many differences which include observation geometry, vertical profile sensitivity, radiative transfer modeling, and data filtering. However, these measurements should show similar temporal patterns due to the seasonal water cycles that drive the concentration of water isotopologues. As shown in **Figure 8(a)**, CLARS-FTS and TCCON  $\delta D$  measurements show high consistency in seasonal cycles with peaks in summers and troughs in winters. The seasonal amplitudes from monthly variations are 138.2‰ and 106.9‰ for CLARS-FTS and TCCON, respectively. On a daily basis, both data sets show high correlation (correlation coefficient ( $r$ ) = 0.92) for daily averaged values. However, there is a systematic offset between the two measurements, with monthly averaged CLARS-FTS data higher than TCCON data by  $(40.4 \pm 18.0)\%$  on average. Also, data for summer months (59.2‰ on average) have a higher average difference than winter months (30.5‰ on average). This systematic difference is expected because CLARS-FTS, compared to TCCON, measures an extra reflected light path in the PBL, where the heavy water isotopologue is more abundant than high altitudes above the PBL. As a result, CLARS-FTS measures higher  $\delta D$  values compared to TCCON. Further discussion to reconcile this difference is described in **Section 4.2**. Since TROPOMI and TCCON have similar measurement geometry, we see a similar comparison result between CLARS-FTS and TROPOMI (**Figure 8(b)**). Unfortunately, very few TROPOMI observations (in total 635) are available for comparison during the two-year period of overlapping data. However, their correlation is still high ( $r=0.89$ ) for the daily averaged values. The offset between these two daily averaged datasets is about  $-61.7\%$ , with TROPOMI data being more negative than CLARS-FTS data.

(a) CLARS-FTS and TCCON



(b) CLARS-FTS and TROPOMI



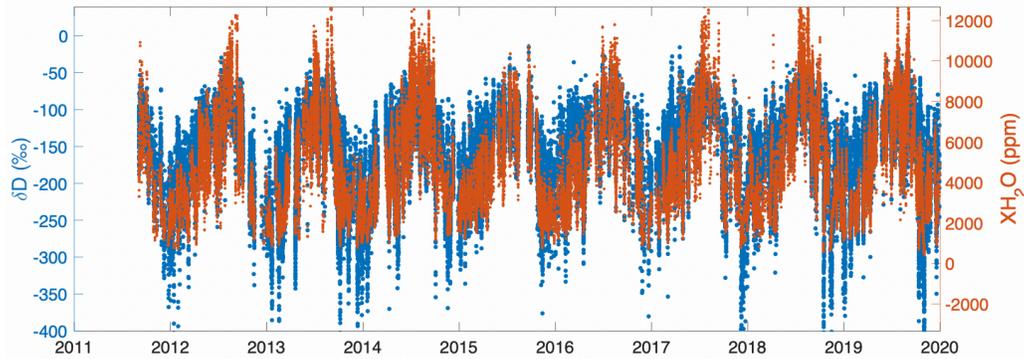
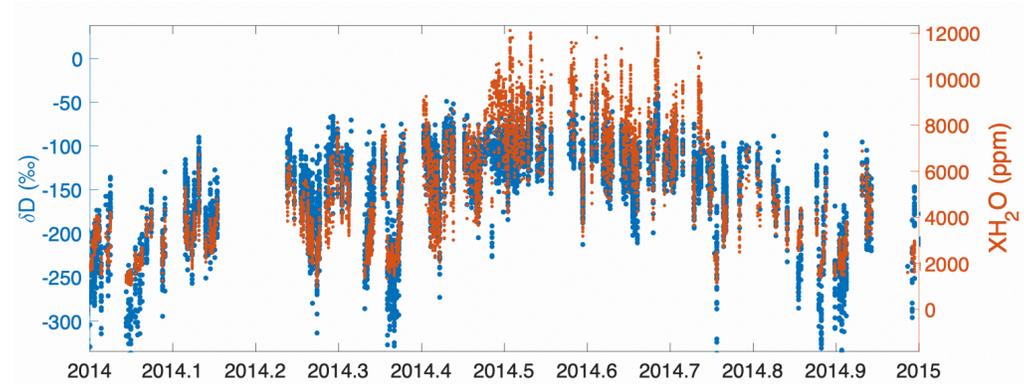
474 **Figure 8.** (a) left: comparisons of time series of  $\delta D$  measurements between CLARS-FTS and TCCON;  
 475 right: scatter plot of their  $\delta D$  measurements; (b) left: comparisons of time series of  $\delta D$  measurements  
 476 between CLARS-FTS and TROPOMI; right: scatter plot of their  $\delta D$  measurements.

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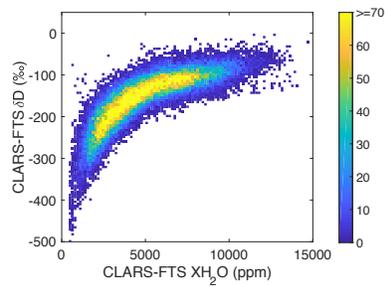
478 **3.5 Correlation between  $\delta D$  and humidity**

479 According to Rayleigh distillation theory (**Rayleigh, 1902**), when water vapor abundance  
480 is lower, heavier isotopologues, such as HDO, are more likely to condense compared with H<sub>2</sub>O.  
481 As a result, heavier isotopologues are more depleted (lower  $\delta D$ ) in the air when humidity is  
482 lower. Therefore, the seasonal variations in  $\delta D$  are found to be primarily driven by variations in  
483 humidity. However, small departures from this humidity– $\delta D$  correlation as predicted by  
484 Rayleigh distillation could provide new insights into secondary processes of the hydrological  
485 cycle related to evaporation and condensation. In this section, we will focus on the  
486 characterization of the humidity– $\delta D$  correlations from measurements and leave further  
487 comparison with theoretical calculations to future studies.

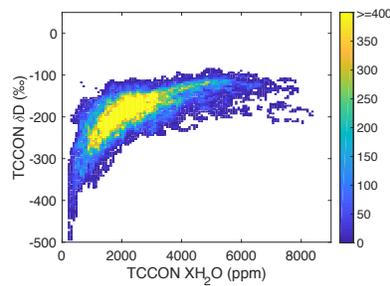
488 In **Figure 9(a)**, we show the comparison of time series of XH<sub>2</sub>O and  $\delta D$  from CLARS-  
489 FTS. The seasonal cycles show highly consistent seasonal patterns. Even on a daily time scale,  
490 both datasets closely track one another, as shown in **Figure 9(b)**. For example, the many low  
491 value anomalies during winter months shown in XH<sub>2</sub>O are also shown in  $\delta D$ . Such a strong  
492 correlation can be seen in the scatter plots in **Figure 9(c-e)** from all CLARS-FTS, TCCON-  
493 Caltech, and TROPOMI data. This nonlinear correlation from remote sensing data has been  
494 reported by many other studies (e.g., **Worden et al., 2007; Noone, 2012; Schneider et al.,**  
495 **2020**). The progressive decrease of  $\delta D$  with a decreasing water vapor mixing ratio clearly  
496 demonstrates the preferential condensation of HDO compared with H<sub>2</sub>O. Another possible  
497 process that may explain the  $\delta D$  and humidity correlation is the air mass mixing model, as  
498 described in **Noone (2012)**, that represents an exchange between two reservoirs with different  
499 H<sub>2</sub>O mixing ratio and  $\delta D$ . For example, the mixing of the atmosphere with evaporation from the  
500 surface. A thorough exploration of these theoretical models to explain the measurements is  
501 beyond the scope of this paper and will be our future works.  
502

(a) CLARS XH<sub>2</sub>O and CLARS  $\delta$ D(b) CLARS XH<sub>2</sub>O and CLARS  $\delta$ D in 2014

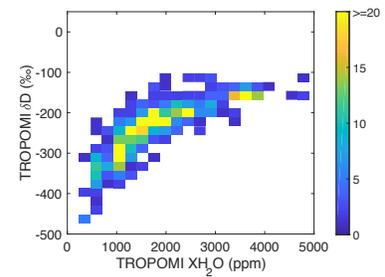
(c) CLARS-FTS



(d) TCCON-Caltech



(e) TROPOMI



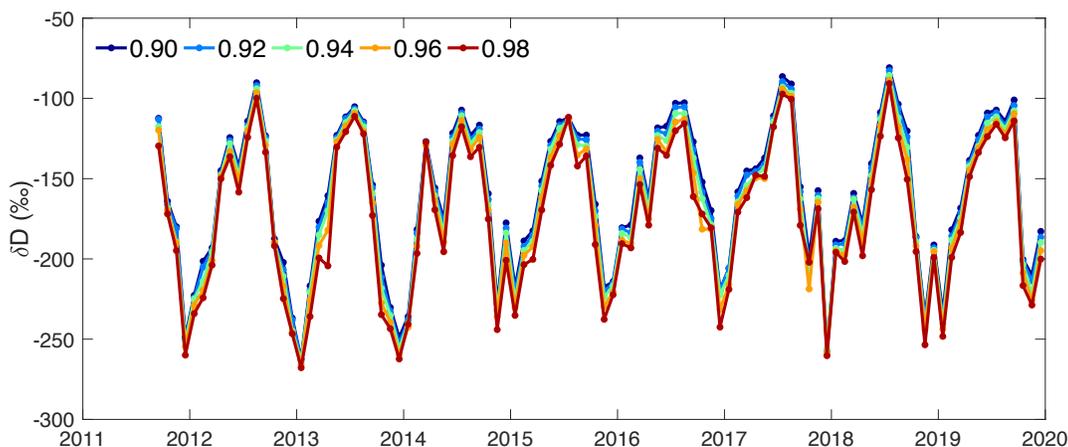
504 **Figure 9.** (a) The time series of  $\delta$ D and XH<sub>2</sub>O from CLARS-FTS. For illustration purposes, some of the  
 505 high XH<sub>2</sub>O values above 12000 ppm are not shown; (b) The same as (a) but zooming into the year 2014;  
 506 (c-e) Density plots between XH<sub>2</sub>O and  $\delta$ D from CLARS-FTS, TCCON-Caltech, and TROPOMI data,  
 507 respectively.

## 508 4. Discussion

### 509 4.1 Uncertainty in CLARS-FTS $\delta D$ retrievals due to the aerosol scattering effect

510 Since aerosol scattering is not incorporated into the GFIT algorithm, we have used a set  
511 of tight filters (**Table 2**) to screen the data that may be affected by clouds and aerosols. For  $\delta D$   
512 retrievals, the effects of aerosol scattering and surface pressure variations largely cancel out since  
513  $\delta D$  is derived from the ratio of HDO and H<sub>2</sub>O columns using spectral windows in the same  
514 wavelength range. This was demonstrated by **Boesch et al. (2013)** from a series of retrieval  
515 sensitivity tests. Here we conducted a sensitivity study for CLARS-FTS retrievals by tightening  
516 the aerosol filters and evaluating the impact on the seasonal cycles of the observed monthly  
517 averaged  $\delta D$ . As shown in **Figure 10**, we can observe a high consistency between retrievals with  
518 different O<sub>2</sub> ratio filters, as defined in **Section 3.1**. Since the O<sub>2</sub> ratio is an effective indicator of  
519 the aerosol scattering effect, this result indicates that the temporal variabilities in  $\delta D$  retrievals  
520 are only slightly affected by the impacts of aerosols. This also confirms the conclusions from  
521 **Boesch et al. (2013)**.

522



523 **Figure 10.** Time series of  $\delta D$  measurements from CLARS-FTS under different threshold for the O<sub>2</sub> ratio  
524 filters (0.90, 0.94, 0.94, 0.96, and 0.98) for screening retrievals affected by aerosol scattering effect.  
525

### 526 4.2 Reconciling the difference between CLARS and TCCON

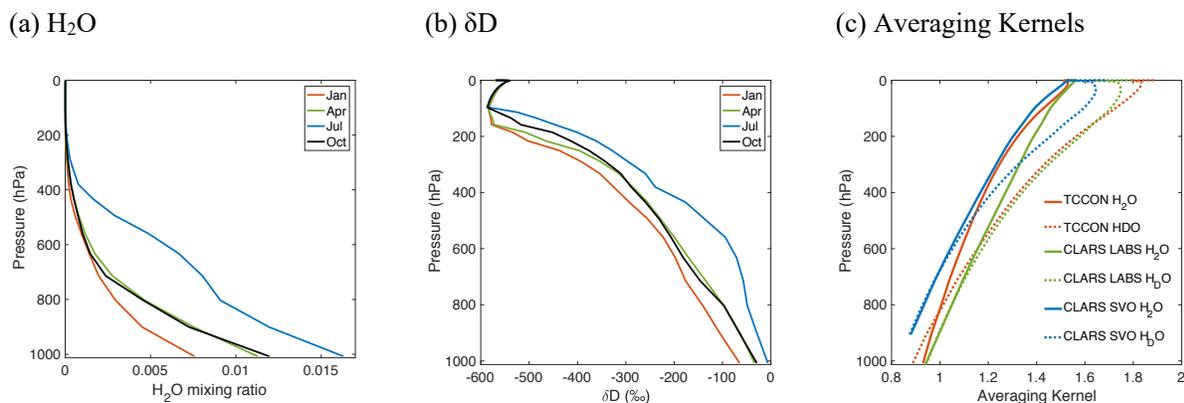
527 As shown in **Figure 8(a)**, a systematic offset is observed in  $\delta D$  retrievals between  
528 CLARS and TCCON. Here we specifically investigate the effects of differences in observation  
529 geometries and averaging kernels on the measured offset. The averaging kernel defines the  
530 sensitivity of the retrievals to the true column. Since different observing systems have different  
531 averaging kernels, their impacts on the retrievals should be evaluated for interpreting the  
532 retrieval discrepancies.

533 Our strategy to investigate the causes of the difference between CLARS and TCCON is  
534 to first construct a set of “true” H<sub>2</sub>O, HDO and  $\delta D$  profiles. These “true” profiles are created  
535 based on a priori GFIT profiles, shown in **Figure 11(a)** and (b), which are then scaled in a way  
536 that they will generate similar column  $\delta D$  retrievals with TCCON and partial column  $\delta D$   
537 retrievals with free tropospheric measurements from CLARS-SVO. We then apply the  
538 observation operators and averaging kernels (**Figure 11(c)**) from CLARS-FTS and TCCON,

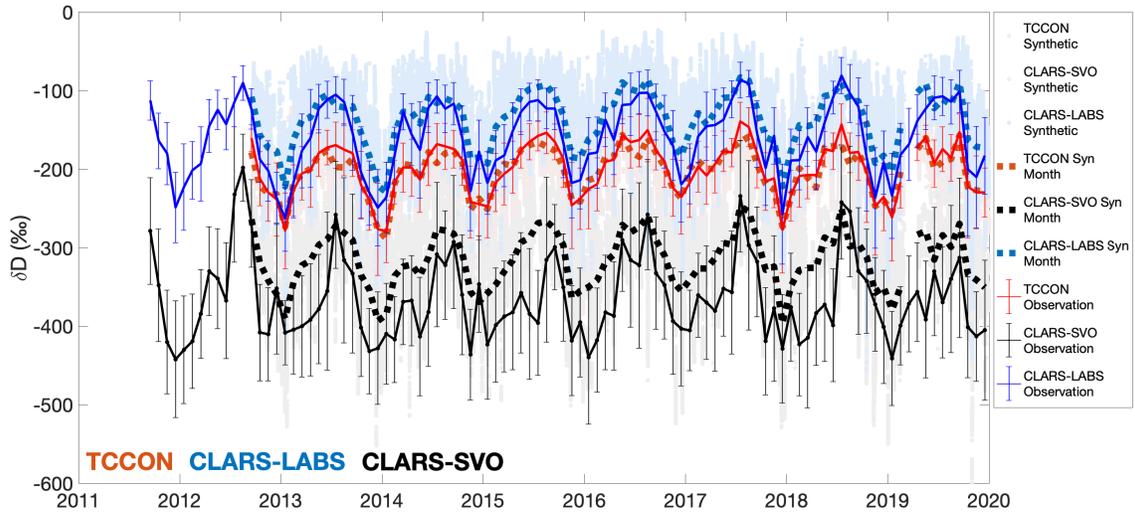
539 separately, on the “truth” profiles to simulate what (synthetic)  $\delta D$  values can be measured from  
 540 both instruments. Finally, we assess the differences of the synthetic  $\delta D$  values from TCCON and  
 541 CLARS and compare with the differences from real measurements.

542 Examples of monthly averaged a priori profiles selected from different seasons for  $H_2O$   
 543 and  $\delta D$  are shown in **Figure 11(a)** and **(b)**. The a priori  $\delta D$  profiles are generated by GFIT with  
 544 assumed fractionation parameters for the troposphere and stratosphere. After the a priori profiles  
 545 are scaled to match the TCCON total  $\delta D$ , we found they significantly overestimate (not shown  
 546 here) the free troposphere partial  $\delta D$  column when compared with CLARS-SVO. This indicates  
 547 that the free troposphere above PBL has been relatively overestimated. We therefore adjusted the  
 548 HDO profiles by scaling all levels above CLARS (1.6 km) by 0.9 and the levels below CLARS  
 549 by 1.1 in order to match the CLARS and TCCON observations within their uncertainties. Using  
 550 these scaled profiles, we applied the CLARS-LABS observation operator and averaging kernel to  
 551 generate synthetic  $\delta D$  values. All synthetic  $\delta D$  values for TCCON, CLARS-SVO, and CLARS-  
 552 LABS are shown in **Figure 12**. We can see that the observed time series of  $\delta D$  values can be  
 553 very well reproduced from the synthetic profiles. This consistency suggests that the  
 554 discrepancies between CLARS and TCCON are primarily driven by the difference in the  
 555 observation geometries and averaging kernels. We conducted a further experiment using the  
 556 observation operator only and assuming TCCON and CLARS have the same averaging kernel.  
 557 The results (not shown here) show a very small difference, which suggests that the averaging  
 558 kernel has a smaller contribution than the observation operator. Moreover, the fact that in **Figure**  
 559 **8** and **Figure 12** summer months have higher difference (between CLARS-LABS and TCCON)  
 560 than winter months is because of the geometries. Winter months have larger solar zenith angle  
 561 and therefore larger air mass in the incident light path compared to reflected light path. As a  
 562 result, the contribution from the reflected path becomes smaller and so CLARS-FTS and  
 563 TCCON are getting closer estimates.

564



565 **Figure 11.** Examples of monthly averaged a priori profiles from GFIT selected from different seasons for  
 566 (a)  $H_2O$  and (b)  $\delta D$ . The original a priori profiles are available on a daily basis. (c) Examples of column  
 567 averaging kernels from TCCON, CLARS-SVO, and CLARS-LABS observations. These column averaging  
 568 kernels are averaged profiles from all available retrieval windows. For CLARS, they are from the three  
 569 HDO windows ( $6330.05\text{ cm}^{-1}$ ,  $6377.40\text{ cm}^{-1}$ , and  $6458.10\text{ cm}^{-1}$ ) and four  $H_2O$  windows ( $6301.35\text{ cm}^{-1}$ ,  
 570  $6392.45\text{ cm}^{-1}$ ,  $6401.15\text{ cm}^{-1}$ , and  $6469.60\text{ cm}^{-1}$ ). For TCCON, they are from the spectral windows as listed  
 571 in Table 3 of **Wunch et al. (2015)**, which includes 15 spectral windows for  $H_2O$  and 6 spectral windows  
 572 for HDO.



573 **Figure 12.** The reconstructed  $\delta D$  time series for TCCON and CLARS-FTS (including SVO and LABS  
 574 modes) from applying the observation operators, which is associated with the observation geometries, and  
 575 the averaging kernels to the reconstructed “truth”  $H_2O$  and  $HDO$  profiles. For comparison, the observed  $\delta D$   
 576 time series are also shown. The “truth” profiles are constructed from the a priori  $H_2O$  and  $HDO$  profiles  
 577 generated from the GFIT program. The profiles are first scaled to match the TCCON  $\delta D$  retrievals. Then a  
 578 scale factor (0.9) is further applied to all levels above CLARS (1.6 km) and another scale factor (1.1) to  
 579 levels below CLARS in order to match the CLARS and TCCON  $\delta D$  observations within their uncertainties.  
 580  
 581  
 582

583 **5. Conclusions**

584 In this study, we retrieved XHDO, XH<sub>2</sub>O, and  $\delta$ D using the high resolution NIR  
585 observations from CLARS-FTS that demonstrates high sensitivity to the PBL atmosphere in the  
586 LA basin. Diurnal observations of XHDO, XH<sub>2</sub>O, and  $\delta$ D with high temporal resolution (as high  
587 as 3 minutes) are generated over 33 surface reflection targets covering the LA basin from 2011 to  
588 2019. The temporal variabilities in  $\delta$ D data between CLARS-FTS and a collocated TCCON  
589 observatory and TROPOMI observations are highly correlated. CLARS-FTS observes higher  
590 values due to its longer path along the PBL where HDO is more abundant. The difference  
591 between CLARS and TCCON or TROPOMI  $\delta$ D retrievals can be attributed to the difference  
592 primarily in the observation geometries and secondarily in the averaging kernels. From CLARS  
593 measurements, the XH<sub>2</sub>O and XHDO time series show strong seasonal cycles that are associated  
594 with the seasonal variation of temperature that drives the evaporation. The  $\delta$ D shows low values  
595 in winter (more depletion of HDO) and high values in summer (less depletion of HDO), mainly  
596 driven by the change of atmospheric humidity. HDO and  $\delta$ D from CLARS-FTS provide high  
597 spatial and temporal resolution datasets for further study of hydrological processes in southern  
598 California.

599 The data set resulting from this work possesses a large amount of potential for future  
600 study. Immediate next steps include examining the spatial variability in the  $\delta$ D data. Analysis of  
601 spatial variability can be performed by mapping  $\delta$ D values according to surface reflection targets  
602 location associated with the observation. In doing so, one may be able to identify regions of the  
603 LA basin which show similar  $\delta$ D patterns throughout the year and therefore are potentially  
604 influenced by similar water vapor sources. If such regions can be identified, additional analysis  
605 of temporal patterns can be performed in order to attempt to understand the relative importance  
606 of various water vapor sources in that region. The time series can also be examined for  
607 correlation with other meteorological patterns, such as precipitation and wind patterns, in order  
608 to further examine climatological trends or weather anomalies in LA. Finally, there is a large  
609 amount of potential to use this data set in modeling studies regarding meteorological processes,  
610 such as cloud formation, and atmospheric water vapor sources for LA.  
611

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621 <https://megacities.jpl.nasa.gov>. The TROPOMI HDO data set from this study is available for  
622 download at [ftp://ftp.sron.nl/open-access-data-2/TROPOMI/tropomi/hdo/9\\_1/](ftp://ftp.sron.nl/open-access-data-2/TROPOMI/tropomi/hdo/9_1/) (last access: 12  
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719

720 **Appendix A: Comparison of Spectral Windows**

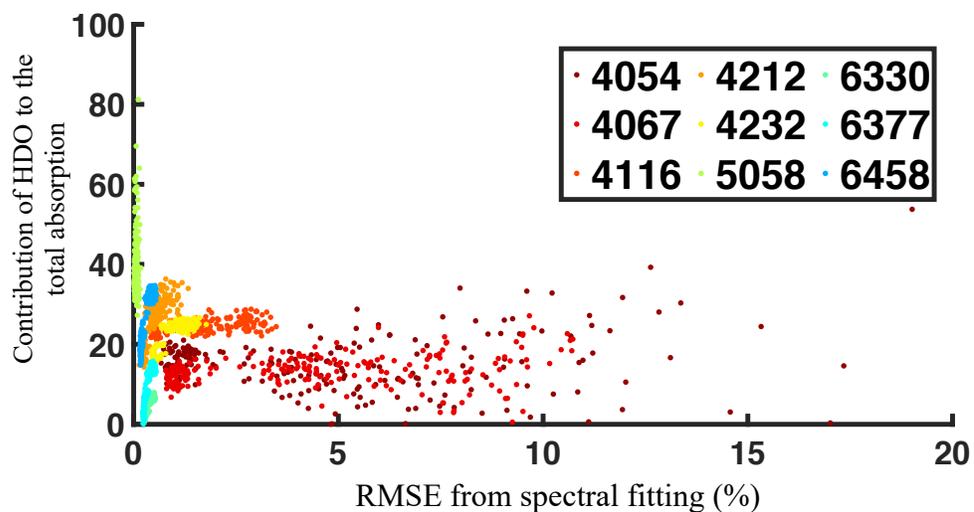
721 We tested 9 spectral windows (**Table A1**) using CLARS-FTS spectral observations on a  
 722 sample of 4-6 days in which the instrument took measurements for two surface reflection point  
 723 locations throughout the course of the day. The results from each spectral window were compared  
 724 according to overall quality of spectral fittings and resulting retrieval uncertainties. In addition to  
 725 recommended TCCON windows, a micro-window around 5058  $\text{cm}^{-1}$ , was tested based on a survey  
 726 of successful spectral windows from prior TCCON studies (**Rokotyan et al., 2014**). **Figure A1**  
 727 shows a plot of the RMSE from the spectral fitting for each spectral window versus the contribution  
 728 from HDO to the overall absorption using a set of  $\sim 4000$  observations from 4 distinct days in 2013.  
 729 From this plot, one can see that the lower frequency spectral windows, which are plotted  
 730 red/orange/yellow colors, have higher spectral fitting error, implying a worse quality fit for these  
 731 spectral windows. Comparing VSF errors for this set observations, the lower spectral windows  
 732 also tended to yield observations with higher overall uncertainties. Therefore, the spectral windows  
 733 in the 4000-5000  $\text{cm}^{-1}$  were eliminated as possible candidates. The green points, referring to the  
 734 5058  $\text{cm}^{-1}$  micro-window, indicate lower spectral fitting error and also relatively high contribution  
 735 due to HDO. However, examining the spectra for this micro-window indicated that it contains only  
 736 one main absorption line, compared to the 6000-7000  $\text{cm}^{-1}$  spectral windows, which contain many  
 737 more. Therefore, the results using 6000-7000  $\text{cm}^{-1}$  were considered more robust. From this  
 738 examination of spectral fit quality from the 9 tested spectral windows, the three spectral windows  
 739 with central wavenumbers in the 6000-7000  $\text{cm}^{-1}$  range were selected as primary candidates for  
 740 further investigation.

741

**Table A1. Tested Spectral Windows with Associated Parameters**

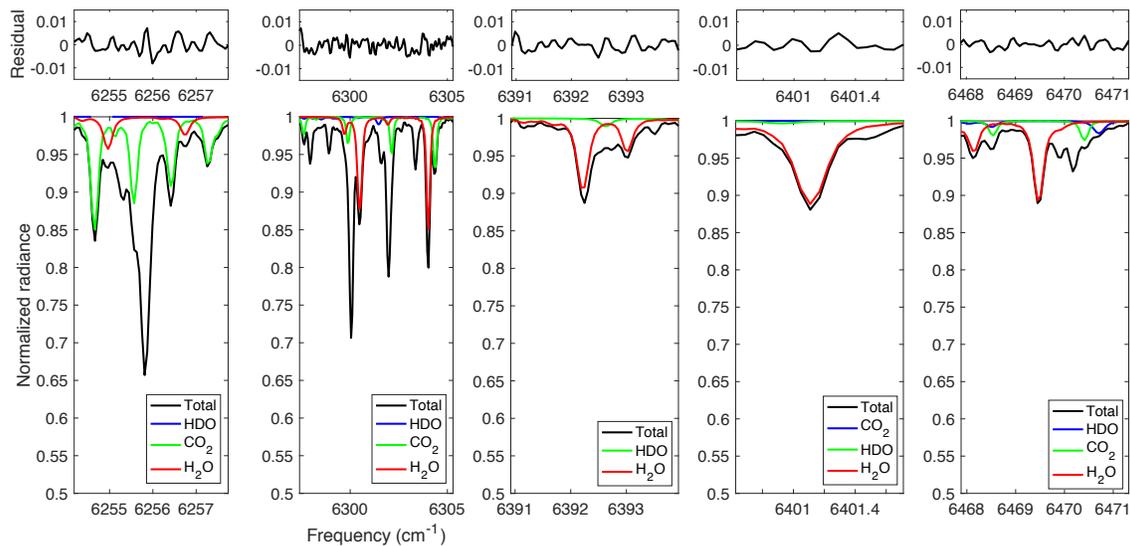
| Center ( $\text{cm}^{-1}$ ) | Width ( $\text{cm}^{-1}$ ) | Gases to fit                                | continuum basis functions   |
|-----------------------------|----------------------------|---|-----------------------------|
| 4054.60                     | 3.30                       | HDO, H <sub>2</sub> O, CH <sub>4</sub>      |                             |
| 4067.60                     | 8.80                       | HDO, H <sub>2</sub> O, CH <sub>4</sub>      |                             |
| 4116.10                     | 8.00                       | HDO, H <sub>2</sub> O, CH <sub>4</sub>      |                             |
| 4212.45                     | 1.90                       | HDO, H <sub>2</sub> O, CH <sub>4</sub>      | 2-order polynomial for      |
| 4232.50                     | 11.00                      | HDO, H <sub>2</sub> O, CH <sub>4</sub> , CO | continuum; frequency shift, |
| 5058.95                     | 1.60                       | HDO, H <sub>2</sub> O, CO <sub>2</sub>      | solar lines                 |
| 6330.05                     | 45.50                      | HDO, H <sub>2</sub> O, CO <sub>2</sub>      |                             |
| 6377.40                     | 50.20                      | HDO, H <sub>2</sub> O, CO <sub>2</sub>      |                             |
| 6458.10                     | 41.40                      | HDO, H <sub>2</sub> O, CO <sub>2</sub>      |                             |

742



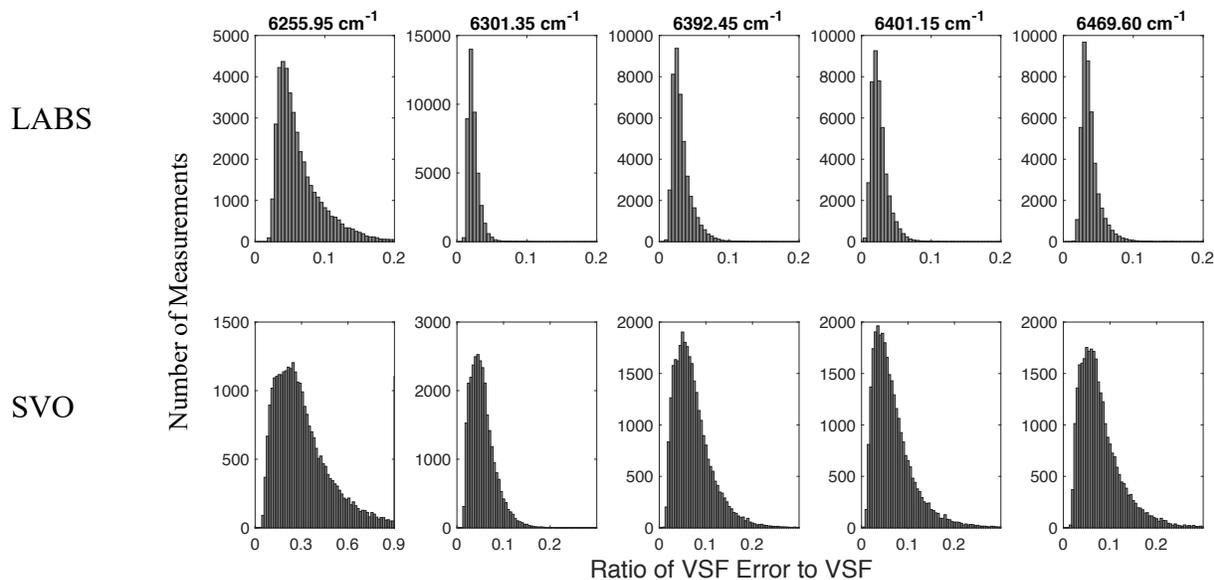
743 **Figure A1:** A plot of RMSE from the spectral fitting versus the contribution from HDO to the  
 744 overall absorption for a sample for ~4000 observations from 4 days in 2013. Results from  
 745 various spectral windows are indicated by color moving from red to blue as the central  
 746 wavenumber of the spectral window increases.  
 747  
 748  
 749

750 **Appendix B. H<sub>2</sub>O spectral windows and their fitting residuals and retrieval errors**

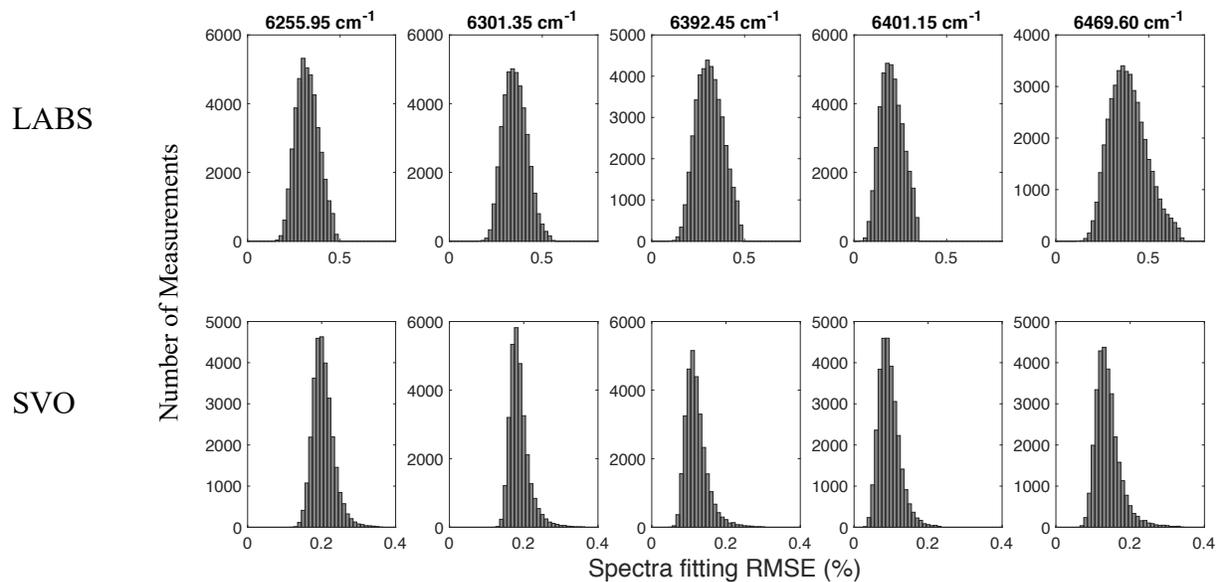


751 **Figure B1.** Examples of CLARS-FTS spectra windows of (from left to right) 6255.95 cm<sup>-1</sup>, 6301.35 cm<sup>-1</sup>,  
 752 6392.45 cm<sup>-1</sup>, 6401.15 cm<sup>-1</sup>, and 6469.60 cm<sup>-1</sup> for retrieving H<sub>2</sub>O in this study. These samples of normalized  
 753 spectra are taken from a mid-day observation on 7/14/2013 over the West Pasadena surface target. The  
 754 lower panel shows the full spectral fit with total contribution, contribution from H<sub>2</sub>O, and from other  
 755 interfering gases. The upper label shows the residuals of the spectral fits, defined as the difference in total  
 756 measured and total calculated radiance.  
 757

758



759 **Figure B2.** Retrieval error, in ratio of VSF error to VSF, for the entire filtered XH<sub>2</sub>O datasets from 2011 to  
 760 2019. Both VSF and VSF error values are calculated from CLARS-GFIT. The data are separated into  
 761 histograms according to observation modes (LABS and SVO) and five spectral windows (6255.95 cm<sup>-1</sup>,  
 762 6301.35 cm<sup>-1</sup>, 6392.45 cm<sup>-1</sup>, 6401.15 cm<sup>-1</sup>, and 6469.60 cm<sup>-1</sup>). Note that the SVO at 6255.95 cm<sup>-1</sup> has a  
 763 different x-axis range.  
 764



765 **Figure B3.** Histograms of RMS error from spectral fitting for the entire data set from 2011 to 2019. The  
 766 data are separated into histograms according to observation mode (LABS and SVO) and five spectral  
 767 windows ( $6255.95\text{ cm}^{-1}$ ,  $6301.35\text{ cm}^{-1}$ ,  $6392.45\text{ cm}^{-1}$ ,  $6401.15\text{ cm}^{-1}$ , and  $6469.60\text{ cm}^{-1}$ ).