# Machine learning bridges microslips and slip avalanches of sheared granular gouge

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#### Abstract

Understanding the origin of stress avalanche of fault gouges may offer deeper insights into many geophysical processes such as earthquakes. Microslips of sheared granular gouges were found to be precursors of large slip events, but the documented relation between local and global avalanches remains largely qualitative. We examine the stick-slip behavior of a slowly sheared granular system using discete element method simulations. The microslips, i.e., local avalanche events, are found to demonstrate significantly different statistical and spatial characteristics between the stick and slip states. We further investigate the correlation between the global stress fluctuations and the features extracted from microslips based on the machine learning (ML) approach. The data-driven model that incorporates the information of the spatial distribution of microslips can robustly predict the magnitude of stress fluctuation. A further feature importance analysis confirms that the spatial patterns of microslips manifest key information governing the global stress fluctuations.

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16	Key Points:
17 18	• The microslips of sheared granular gouge demonstrate different statistical and spatial characteristics in stick and slip states
19	• Clustering of large microslips is correlated closely to large stress drop
20 21	• Machine learning establish a quantitative relationship between microslips and global stress fluctuations
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# 24 Abstract

Understanding the origin of stress avalanche of fault gouges may offer deeper insights into many 25 geophysical processes such as earthquakes. Microslips of sheared granular gouges were found to 26 be precursors of large slip events, but the documented relation between local and global 27 avalanches remains largely qualitative. We examine the stick-slip behavior of a slowly sheared 28 29 granular system using discete element method simulations. The microslips, i.e., local avalanche events, are found to demonstrate significantly different statistical and spatial characteristics 30 between the stick and slip states. We further investigate the correlation between the global stress 31 fluctuations and the features extracted from microslips based on the machine learning (ML) 32 approach. The data-driven model that incorporates the information of the spatial distribution of 33 microslips can robustly predict the magnitude of stress fluctuation. A further feature importance 34 analysis confirms that the spatial patterns of microslips manifest key information governing the 35 36 global stress fluctuations.

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# 38 Plain Language Summary

Frictional instability of natural fault gouges may play a key role in numerous geophysical 39 processes, such as earthquakes and debris flows. Direct investigation of natural faults is difficult 40 owing to their burial depth and broad distribution beneath the earth. Structural and statistical 41 similarities between granular materials and fault gouges render the former an ideal model system 42 for understanding the mechanism of natural fault gouges. When slowly deformed, granular 43 44 materials generate cycles of friction increase and reduction, i.e., stick-slip cycles, analogous to the earthquake cycles. Microslips are found to be precursors of large slip events, but their 45 correlations are mostly qualitative. This study uses numerical simulations to generate a series of 46 stick-slip cycles of slowly sheared granular gouge. Distinctive differences are observed in the 47 statistical and spatial characteristics of microslips between stick and slip stages. A trained 48 machine learning model is further used to predict the global slip avalanche from the features 49 extracted from microslips and its prediction accuracy can be significantly improved when 50 considering the spatial information of microslips. This work suggests that the microslips detected 51 inside natural gouge faults (e.g., local acoustic emission signal or local seismic wave) and their 52 locations can be used to assess their frictional stability. 53

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# 55 **1 Introduction**

The frictional stability of fault gouge layers underpins key understandings to many geophysical 56 processes, including but not limited to earthquakes, debris flows, and landslides (Song et al., 57 2017; Ren et al., 2019; Nanjo, 2020). A granular gouge subjected to slow shearing demonstrates 58 a typical stick-slip behavior, which plays a crucial role in triggering the frictional stability of the 59 fault (Byerlee & Brace, 1966; Marone et al., 1991; Aharonov & Sparks, 2004; Denisov et al., 60 61 2016; Dorostkar & Carmeliet, 2019). Therefore, the stick-slip behavior of sheared granular gouges has been studied extensively in both laboratory experiments (Marone, 1998; Niemeijer et 62 al., 2010; Scuderi et al., 2015, 2016; Leeman et al., 2016; Tinti et al., 2016; Rivière et al., 2018) 63 and numerical simulations (Aharonov & Sparks, 2004; Mair & Hazzard, 2007; Ferdowsi et al., 64 2014b; Dorostkar et al., 2017a; Gao et al., 2018; Ma et al., 2020). Particular attention has been 65 placed on the influences of controlling factors on the stick-slip dynamics of granular gouge, such 66

as the wall geometry and friction (Rathbun et al., 2013), presence of liquids (Dorostkar et al.,
2017b, 2018), particle characteristics (Mair et al., 2002; Dorostkar & Carmeliet, 2019), boundary
vibration (Ferdowsi et al., 2014a), normal pressure (Gao et al., 2018), particle size polydispersity
(Ma et al., 2020), and particle breakage (Wang et al., 2020). These studies offer novel insights
into the complex dynamic behaviors of natural fault gouges and earthquake physics.

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However, the microscopic origin of slip avalanche of slowly deformed granular gouge remains 73 74 poorly understood. To address this issue, Johnson et al. (2013) employed a biaxial shear apparatus to investigate the physics of laboratory earthquake and found that the acoustic 75 emission and microslip exhibit an exponential increase in the rate of occurrence, reaching a peak 76 at the onset of slip avalanche. The corresponding DEM simulations comfirmed that the microslip 77 event rate correlates well with large slip event onset (Ferdowsi et al., 2013). Microslip or local 78 avalanche is essentially a result of the localized particle rearrangements (Ma et al., 2021). Due to 79 80 the disordered structure of granular materials, a microslip may trigger nearby microslips, and the accumulation of these microslips may give rise to a global stress avalanche (Castellanos & 81 Zaiser, 2018; Cao et al., 2019). Thus, microslips are widely regarded as precursors of large slip 82 events and can be used to predict frictional weakening (Bolton et al., 2019, 2020; Trugman et al., 83 2020). 84

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Furthermore, the statistics of local and global avalanches reveal a simple relation between the 86 87 number of local avalanches and the global avalanches (Barés et al., 2017). The spatial characteristics of microslips are also closely correlated with the stress avalanche, where large 88 stress drop is accompanied by a series of connected localized zones spanning the entire system, 89 whereas during the elastic regime, the microslip events occur with low concentration and are 90 spatially dispersed (Cao et al., 2018). Other particle scale metrics, such as coordination number, 91 sliding contact ratio, potential energy, kinetic energy, evolves correspondingly during the stick 92 phase and slip instability (Ferdowsi et al., 2015; Barés et al., 2017; Dorostkar & Carmeliet, 2018; 93 94 Ma et al., 2020). Thus, studying the microscopic structure and dynamics of a granular gouge may 95 help to unveil its stick-slip behaviors (Cipelletti et al., 2019).

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Unfortunately, existing findings on the relation between microslips and global stress avalanche 97 remains largely qualitative, whereas further advance on the subject matter demands quantitative 98 correlations to be established. In this letter, we employ the machine learning (ML) approach to 99 bridge the microslips and global stress fluctuations, including both stress recharge (stick regime) 100 and stress drop (slip regime). ML offers data-driven approaches to automatically investigate the 101 underlying relations between variables and facilitate the process of revealing complex and 102 inexplicit patterns of large datasets (Marone, 2018; Bergen et al., 2019; Ren et al., 2020). 103 Particularly, ML has gained increasing popularity in recent years and has been widely used in 104 many areas of geoscience, such as predicting the timing and size of laboratory earthquakes 105 (Rouet-Leduc et al., 2017; Corbi et al., 2019), revealing the frictional state of granular fault 106 (Rouet-Leduc et al., 2018; Ren et al., 2019), estimating earthquake magnitude and GPS 107 displacement rate (Rouet-Leduc et al., 2019; Mousavi & Beroza, 2020), and performing 108 earthquake early warning and earthquake detection (Rouet-Leduc et al., 2018; Hulbert et al., 109 2019; Mousavi et al., 2020; Trugman et al., 2020). 110

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To do so, we perform the discrete element method (DEM) simulations of quasi-static shear of 112 granular gouge to achieve stick-slip dynamics. The microslip is manifested as the particle 113 rearrangements and quantified by the nonaffine particle motion. Then we compare between the 114 statistical and spatial characteristics of microslips in the stick and slip regimes. We use a two-115 step scheme for feature selection to consider both the statistical and spatial characteristics of 116 microslips in the ML model training. The trained XGBoost model can well predict the global 117 stress fluctuation from the features extracted from the microslips. Finally, we analyze the feature 118 importance of the trained ML model and conclude that the spatial patterns of microslips contain 119 key information about the stick-slip dynamics of granular gouge. 120

#### 121 2 Materials and Methods

DEM simulations of simple shear tests were performed to obtained data of microslips and global stress fluctuations during the stick-slip cycles of granular gouge. Figure 1a shows the simple shear model setup of the granular gouge, which consists of 9,134 particles with diameters

uniformly distributed from  $0.8 d_{50}$  to  $1.2 d_{50}$ , where the average particle diameter  $d_{50}$  =1.25 mm.

126 The size of the granular gouge sample is  $32 d_{50}$  (length) ×  $16 d_{50}$  (depth) ×  $16 d_{50}$  (height). The 127 granular gouge is confined by two rough particle walls used to apply the shear loading and 128 normal pressure. The top wall is fixed in the shear direction, while the normal pressure is 129 maintained constant by a servo-control at 10 MPa. The granular gouge is sheared by moving the 130 bottom wall in the *x* direction with a constant velocity while the vertical movement is 131 constrained. The shear rate  $\dot{\gamma}$ , defined as the ratio of shear velocity to the undeformed sample

height, is set to 0.05 to achieve stick-slip dynamics.

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The numerical simulation is performed by the DEM code LIGGGHTS (Kloss et al., 2012). The 134 Hertz-Mindlin contact model with Coulomb sliding friction is employed to simulate the contacts 135 and deformation between particles. The particles have a density of 2900 kg/m<sup>3</sup>, a Poisson's ratio 136 of 0.25, Young's modulus of 65 GPa, a friction coefficient of 0.1, and a restitution coefficient of 137 0.87 (Ma et al., 2020). The wall particles adopt the same material properties as those in the shear 138 body. The friction coefficient between the particle walls and the shear body is set to 0.9 to 139 enhance surface friction. To collect enough data for the subsequent machine learning, we shear 140 the granular gouge up to a shear strain of 50%. The evolution of normalized shear stress, defined 141 as the ratio of shear stress  $\sigma$  to the applied normal pressure p, is shown in Figure 1b. When it is 142 sheared into the steady-state regime, the gouge is found to undergo typical intermittent dynamics 143 and serrated plastic flow. This phenomenon is seen to be universal in many amorphous solids 144

like metal glasses (Sun et al., 2012; Cao et al., 2018), and porous materials (Baró et al., 2013).

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The enlarged view of the dotted box shown in Figure 1b demonstrates that each stick-slip cycle starts with a nonlinear recharge of shear stress and is followed by a rapid drop. The recharge and drop of shear stress correspond to the stick and slip stages, respectively. We define stress fluctuation as the change of shear stress at the start and end of the recharge/drop events. Thus,

the stress fluctuation of a drop event is positive, and the recharge event negative. Only the

magnitude of stress fluctuation greater than a threshold of  $10^{-5}$  is considered. During the slow

- shearing of granular gouge, we recorded 3,191 stress drop and recharge cycles.
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155 Shear strain,  $\gamma$ 156 **Figure 1**. (a) Setup of the DEM experiment. Normal pressure and shear displacement are respectively applied on the 157 top and bottom particle walls. Periodic boundary conditions are applied in the shear and depth directions. (b) Stress-158 strain curve resulted from the DEM simulation. The *y* axis denotes the shear stress  $\sigma$  normalized by the normal 159 pressure *p*. The inset shows the enlarged stick-slip cycle which consists of stress recharge and drop stage 160 represented by the blue and red shaded region, respectively.

161

#### 162 **3 Results**

#### 163 **3.1 Statistical and spatial characteristics of microslips**

The microslips that occurred during the recharge and drop events are manifested as irreversible 164 particle rearrangements which are hereby quantified by the nonaffine particle displacements 165  $D_{\min}^2$  (see Text S1) (Ma et al., 2021). It should be noted that many other quantities, such as local 166 displacement, local energy fluctuation (Barés et al., 2017; Zheng et al., 2018), granular 167 temperature (Ma et al., 2018, 2019), local acoustic emission (Trugman et al., 2020), force chain 168 bulking (Gao et al., 2019; Liu et al., 2020) can also be used for characterizing microslips. Figure 169 2a and 2b show the spatial distributions of  $D_{min}^2$  during the recharge stage and the drop stage of a 170 typical stick-slip cycle. Particles with higher  $D_{\min}^2$  are colored in red. Due to the discrete nature 171 and corporative particle motion of granular materials, the deformation of granular gouge occurs 172 as a succession of localized micro-slips distributed within the system. Intuitively, the microslips 173 are scattered throughout the granular gouge during the recharge stage. During the drop stage, the 174 microslips are more spatially concentrated and tend to establish large stress avalanches inside the 175 granular gouge. 176

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178 179 Figure 2. Statistical and spatial characteristics of microslips occurred during recharge and drop stages. Spatial maps 180 of  $D_{\min}^2$  occurred during the (a) recharge state and (b) drop stage of a typical stick-slip cycle. (c) Comparison of the statistical quantities of  $D_{\min}^2$  during the recharge and drop stage. (d) Normalized spatial correlation function of  $D_{\min}^2$ 181 between two particles separated by distance r where r is in unit of the mean particle diameter. The data points are 182 183 averaged over the recharge or drop stages falling into each bin. Solid lines are fits to the Ornstein-Zernike function. The data points and fitting lines of the different bin are shifted vertically for better visualization. (e) Evolutions of 184 correlation length  $\xi_r$  and Moran' I with the magnitude of stress fluctuation. The error bar represents the standard 185 deviation. Note that (d) and (e) are calculated over all stick-slip cycles. 186

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Figure 2c compares the statistical features of microslips that occurred during the recharge stage 188 (blue) and drop stage (red) of a stick-slip cycle. The microslips demonstrate significantly 189 different statistical characteristics at the recharge stage and drop stage. For example, the 99.5<sup>th</sup> 190 percentile, max, variance, skewness, and kurtosis are larger for drop stage. The difference in 191 statistics of microslips may suggest different underlying mechanism for stress recharge and stress 192 drop. The stick-slip dynamics of granular materials can be seen as the jamming-unjamming 193 process accompanied by the formation and buckling of force chains, which are triggered by 194 localized particle rearrangements known as microslips or local avalanches (Barés et al., 2017; 195 Gao et al., 2019). 196

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The spatial distributions of microslips that occurred during the recharge and drop stage can be further quantified using the normalized spatial correlation function (see Text S2) (Ma et al.,

2019, 2021). We group the recharge and drop events according to the magnitude and sign of 200 stress fluctuation. Logarithmic binning is used. Figure 2d shows the normalized correlation 201 functions  $C_{D_{aux}^2}(r)$  for recharge and drop events of different magnitudes. The spatial 202 autocorrelation decays rapidly within a short distance of several  $d_{50}$ , showing a short-range 203 ordering. Solid lines indicate that the decay of correlations with r are reasonably well fitted by 204 the Ornstein-Zernike function as  $C_{D^2_{rm}}(r) \propto r^{-0.5} \exp(-r/\xi_r)$ . We can see that the correlation 205 length of microslips  $\xi_r$  remains nearly unchanged for recharge events and increases rapidly for 206 large stress drop (see red line and left axis of Figure 2e). This trend indicates that a more 207 cooperative and concentrated distribution of microslips constitutes the microscopic origin of 208 global slip avalanche. 209 210

The spatial autocorrelation of microslips can also be quantified by global Moran's I (see Text 211 S2) (Ma et al., 2019, 2021). The Moran's I of particle  $D_{\min}^2$  for recharge and drop events of

212 different magnitudes show a very similar trend as the correlation length  $\xi_r$  (see blue line and 213 right axis of Figure 2e). The spatial correlation analysis of microslips indicates that the spatially 214 correlated microslips forming large shear transition zones are responsible for the stress drop and 215 frictional weakening. The stress drop increases with the increasing degree of aggregation of 216 microslips. The spatial distribution of microslips during recharge stages shows on average a 217

- plateau over different bins. 218
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#### **3.2 Machine learning predicts the stress fluctuations** 220

221 In order to establish the quantitative relation between microslips and the magnitude of global stress fluctuation, we resort to use the Extreme Gradient Boosting (XGBoost) technique to 222 interrogate the data (Chen & Guestrin, 2016). Different from the Deep Learning, we need to 223 224 extract physically reasonable features from the raw data for input for the Machine Learning (ML). The above analysis demonstrates a clear difference of microslips between recharge and drop 225 events. Therefore, it is necessary to consider both the statistical and spatial characteristics of 226 microslips in the feature extraction. We first calculate the maximum, mean, variance, skewness, 227 and kurtosis of particle  $D_{\min}^2$  within each particle's second-neighbor shell (see Figure 3a), 228 corresponding to the second minimum of pair correlation function shown in Figure S1a. These 229 statistics contain information on how particle  $D_{\min}^2$  distributes in space. We then calculate the 230 statistical features of each particle's medium range statistics (see Figure 3b). The statistical 231 operator includes mean, max, variance, percentiles, and various higher-order moments. These 232 233 statistical features are connected as the medium-range feature vector (MRF).

234

To highlight the importance of the spatial pattern of microslips in the prediction of global stress 235

- fluctuation, we also calculate the statistical features of particles  $D_{\min}^2$  as the input vector for 236
- XGBoost model training. This feature vector does not contain any information about the spatial 237
- distribution of microslips, and is referred to as particle-scale feature vector (PSF). MRF and PSF 238

are extracted for each recharge/drop event, and the corresponding output of XGBoost is the

global stress fluctuation of the recharge/drop event. A typical structure of a XGBoost is depicted
 in Figure 3c. The XGBoost modeling process is briefly introduced in Text S3.

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The shuffled dataset is divided into training set, test set, and validation set, with a proportion of 243 60%, 20%, and 20%, respectively. The three sets do not overlap each other to avoid "information 244 leakage". The loss function of XGBoost for regression problems is the mean square error (MSE). 245 The hyperparameters of XGBoost are tuned using Bayesian Optimization (Snoek et al., 2012). 246 The performance of XGBoost models using PSF and MRF as inputs are shown in Figure 3d and 247 Figure 3e, respectively. As can be seen, the trained XGBoost models not only classify the 248 recharge and drop event from the microslips, but also predict the magnitude of stress fluctuation 249 with good accuracy. By taking into account both statistical and spatial characteristics of the 250 microslips, the trained XGBoost model exhibits better performance with a coefficient of 251 determination  $R^2 = 0.78$ . 252

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Figure 3. Machine learning builds the bridge between microslips and global stress fluctuation. (a) The statistical characteristics of particle  $D_{\min}^2$  within each particle's second-neighbor shell. (b) Feature extraction process: particlescale feature vector (red column) and medium-range feature vector (blue columns). These two feature vectors are fed as input to the downstream XGBoost model to predict global stress fluctuation. (c) Schematic of XGBoost (a supervised ML approach) based on the gradient boosting decision. Performance of XGBoost model trained by (d) PSF and (e) MRF, respectively.

Figure 3e depicts that certain success can be achieved in learning the complex relations between local and global avalanches for prediction. We further analyze the feature importance of the 264 XGBoost model trained by MRF. The feature importance is quantified by Shapley Additive 265 Explanation (SHAP) value (Lundberg & Lee, 2017). The SHAP value for each feature is the 266 average marginal contribution of a feature value across all possible coalitions, representing their 267 contribution towards a higher or lower final prediction. Figure 4a shows the mean absolute 268 SHAP values of the top 10 important features. The mean value of  $\phi_{kurt}$  is the most important 269 feature, changing the predicted absolute stress fluctuation on average by 0.6 percentage points.

270

 $\phi_{kurt}$  measures the tail-heaviness of  $D_{\min}^2$  of a particle's second nearest neighbors (Westfall, 271 2014). The smaller  $\phi_{kurt}$  indicates the considered particle and its neighbors move in a corporative 272 manner, i.e., particles with either high  $D_{\min}^2$  or low  $D_{\min}^2$  are spatially clustered. To investigate 273 how the mean  $\phi_{kurt}$  affects the model prediction, we present the SHAP dependence plot in Figure 274 4b. Each dot denotes a recharge/drop event in the ML dataset, and the scatters are colored 275 according to the global Moran's I of particle  $D_{\min}^2$ . The higher mean  $\phi_{kurt}$  results in smaller and 276 negative SHAP value, pushing the XGBoost prediction towards a recharge event. In contrast, 277 microslips of a drop event demonstrate stronger spatial correlation and thus have smaller  $\phi_{kurt}$ . 278 This feature helps XGBoost to distinguish between the recharge and drop events and predict the 279 magnitude of global stress fluctuation. This study reveals that the spatial distribution of 280 microslips contains key information on the stress state of a granular gouge such that microslips 281 (e.g., local acoustic emission signal and local seismic wave) detected inside the natural gouge 282

faults may also serve useful to predict its frictional stability.



284 Mean(ISHAP value)(average impact on model output) Mean $\varphi_{kurt}$ 285 **Figure 4**. Feature importance analysis. (a) SHAP values for the top 10 important features. (b) Dependence plot for 286 the mean value of  $\phi_{kurt}$ , colored by the global moran's *I*.

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#### 288 4 Conclusions

We numerically investigated the relations between microslips and global stress fluctuation of a slowly sheared granular gouge. The microslip is manifested as irresible particle rearrangement and is quantified by nonaffine particle motion. The statistical features and spatial distributions of microslips that occurred during the rechage and drop stages of a stick-slip cycle demonstrate apprantely different characteristics. Both the Moran's *I* and the correlation length of particle  $D_{min}^2$  indicate that microslips in the drop stage are spatially correlated to form large local avalanches, leading to large stress drop and frictional weakening. The difference in the
 microscopic dynamics of recharge and drop events suggest that we may quantitatively connect
 the microslips and global stress fluctuation.

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The use of XGBoost boosts to build the bridge between microslips and macro stress fluctuation. 299 Two sets of input fractures are extracted from the raw data to train the ML models. By using the 300 input feature vector containing both statistical and spatial information of microslips, the trained 301 XGBoost model can not only distinguish btween recharge and drop events but also predict the 302 magnitude of stress fluctuation with good accurancy. The feature importance analysis by SHAP 303 values reveals that the kurtosis of  $D_{\min}^2$  within each particle's first and second nearest neighbors 304 is the most important feature, which characterize the local spatial autocorrelation of microslips. 305 We conclude that the spatial distributions of microslips contain key information about the stress 306 307 state of granular gouge fault and its frictional stability. It should be noted that there are many other ways to extract the spatial patterns of microslips, such as Convolutional Neural Network, 308 Graph Embedding for feature extraction, and complex network analysis. This study may shed 309 lights into the mechanisms governing earthquake nucleation, microslips, friction fluctuations, 310 and their connection during the stick-slip dynamics of earthquake cycles. 311

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