# Energetic charged particles in the terrestrial magnetosphere: Cluster/RAPID results

Elena A. Kronberg<sup>1</sup>, Patrick W. Daly<sup>2</sup>, Elena E. Grigorenko<sup>3</sup>, Artem Smirnov<sup>4</sup>, Berndt Klecker<sup>5</sup>, and Andrey Yu. Malykhin<sup>3</sup>

<sup>1</sup>Ludwig Maximilian University of Munich <sup>2</sup>Max-Planck Institut fuer Sonnensystemforschung <sup>3</sup>Space Research Institute of RAS <sup>4</sup>Helmholtz Centre Potsdam GFZ German Research Centre for Geosciences <sup>5</sup>MPI

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# Abstract

Energetic charged particles  $\sim>40$  keV can affect the plasma temperature and pressure and as a consequence the whole dynamics of the magnetosphere. How do particles get to such energies is a fundamental science question. Energetic particles are also potentially hazardous for the space observations. It is, therefore, necessary to study the origin of energetic plasma, its acceleration, distribution and consequences on the magnetospheric dynamics. In this work we review the related results based on observations from the Cluster/RAPID energetic charged particle detector. These results represent new insights in plasma acceleration, unexpected features in its distributions, effects on substorm and geomagnetic storm dynamics and remediated observations in the radiation belts during approximately 1.5 solar cycles.

# Energetic charged particles in the terrestrial magnetosphere: Cluster/RAPID results

# E. A. Kronberg<sup>1</sup>, P. W. Daly<sup>2</sup>, E. E. Grigorenko<sup>3,4</sup>, A. G. Smirnov<sup>5,6,7</sup>, B. Klecker<sup>8</sup> and A. Yu. Malykhin<sup>3</sup>

5	<sup>1</sup> Department of Earth and Environmental Sciences, Ludwig Maximilian University of Munich, Munich,
6	Germany.
7	<sup>2</sup> Max Planck Institute for Solar System Research, Göttingen, Germany.
8	<sup>3</sup> Space Research Institute, Russian Academy of Sciences, Russia.
9	<sup>4</sup> Moscow Institute of Physics and Technology, Moscow, Russia.
10	<sup>5</sup> Helmholtz-Centre Potsdam - GFZ German Research Centre for Geosciences, Potsdam, Germany.
11	<sup>6</sup> Institute of Physics and Astronomy, University of Potsdam, Potsdam, Germany.
12	<sup>7</sup> Geophysical Center of the Russian Academy of Sciences, Moscow, Russia.
13	<sup>8</sup> Max Planck Institute for Extraterrestrial physics, Garching, Germany.

# Key Points:

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15	•	Energetic particles at energies ${>}40~{\rm keV}$ have to be considered for the plasma tem-
16		perature and pressure calculations
17	•	Effective acceleration is due to interaction of charged particles with multi-scale
18		magnetic structures and/or electromagnetic fluctuations
19	•	Direction of Interplanetary Magnetic Field leads to ion distribution asymmetries

between Northern and Southern hemispheres

 $Corresponding \ author: \ Elena \ Kronberg \verb|@geophysik.uni-muenchen.de|$ 

### 21 Abstract

Energetic charged particles  $\sim >40$  keV can affect the plasma temperature and pressure 22 and as a consequence the whole dynamics of the magnetosphere. How do particles get 23 to such energies is a fundamental science question. Energetic particles are also poten-24 tially hazardous for the space observations. It is, therefore, necessary to study the ori-25 gin of energetic plasma, its acceleration, distribution and consequences on the magne-26 tospheric dynamics. In this work we review the related results based on observations from 27 the Cluster/RAPID energetic charged particle detector. These results represent new in-28 sights in plasma acceleration, unexpected features in its distributions, effects on substorm 29 and geomagnetic storm dynamics and remediated observations in the radiation belts dur-30 ing approximately 1.5 solar cycles. 31

# 32 Plain Language Summary

Acceleration of plasma in the universe is a fundamental science question. Under-33 standing of acceleration at supernovae shocks, of cosmic jets or laboratory plasmas still 34 has many open questions. Near-Earth space environment is an excellent laboratory to 35 investigate plasma dynamics and to reveal fundamental laws it obeys. Energetic plas-36 mas are also hazardous for space satellites and play a key role in space weather. It is, 37 therefore, necessary to study the energization of space plasmas, their distribution and 38 consequences on the magnetospheric dynamics. In situ observations in the near-Earth 39 space by Cluster satellites and energetic particle detector, RAPID, reveal new insights 40 in plasma acceleration, unexpected features in its distributions and effects on substorm 41 and geomagnetic storm dynamics. These observations also help space weather applica-42 tions to determine the level of energetic particle intensities. 43

## 44 1 Introduction

The terrestrial magnetosphere is an effective plasma accelerator. There are two main sources for the magnetospheric plasma: the solar wind and the ionosphere.

The solar wind ion population consists mainly of protons. The abundance of  $\mathrm{He}^{++}$ 47 is in the range of 3% to 6% with variations observed at times of up to 30% (Neugebauer, 48 1981; Aellig et al., 2001; Kasper et al., 2007; Borrini et al., 1982). The abundance of other 49 heavy ions is <0.1% (including  $O^{6+}$ ) (Gloeckler & Geiss, 1989). The number of solar wind 50 electrons maintains the charge neutrality in the plasma media. The energy of the pris-51 tine solar wind at  $\sim 1$  Astronomical Unit is about 1 keV with a Maxwellian core and steeply 52 falling power law part of the spectral distribution for protons (see Figure 1) and about 53 10s of eV for the electrons which have harder distribution at higher energies than pro-54 tons. The solar wind interacts with the magnetic field of the Earth, forms the bow shock 55 and then enters in to the magnetosphere via magnetic reconnection and Kelvin-Helmholtz 56 instability. However, in the upstream region of the Earth's bow shock and in the plasma 57 sheet region we observe strong hardening of the power law part of the proton spectrum 58 compared to those in the solar wind, see Figure 1. This implies strong proton energiza-59 tion. Also electrons are strongly energized in the plasma sheet. 60

The ionosphere delivers to the magnetosphere protons,  $He^+$ ,  $N^+$ ,  $O^+$ , and the molec-61 ular ions NO<sup>+</sup> and  $O_2^+$  (Yau et al., 1991). The dominant heavy ion is O<sup>+</sup> and it is con-62 sidered as a main tracer of the ions of the ionospheric origin. Electrons flow out from 63 the ionosphere to keep overall the charge neutrality. However, their plasma character-64 istics are not well studied yet (Borovsky et al., 2020). The ionospheric charged particles 65 outflow on the open and closed magnetic field lines into the magnetosphere. The tem-66 perature of the outflowing ionospheric ions in the lobes is about 7 eV (Kronberg et al., 67 2014). The observed temperature in the plasma sheet is about 20 keV (Malykhin et al., 68 2020). This means that the oxygen ions are strongly energized in the plasma sheet. 69



Figure 1. The proton spectra in the solar wind measured by the Parker Solar Probe is indicated by the blue line. The slope of the black bar is equivalent to the spectral slope in the plasma sheet during quiet time taken from Kronberg et al. (2010). Adapted from McComas et al. (2016).

Energized particles play an important role in the dynamics of the terrestrial magnetosphere. Here under energetic we mean electrons in the range from 40 to 400 keV and
ions from ~30 keV to 4 MeV. These energy ranges are determined by the capabilities
of the energetic particle detector RAPID on the Cluster mission. Such energetic particles can affect the plasma temperature and pressure and as a consequence the whole dynamics of the magnetosphere starting from small scales such as wave properties and plasma
instabilities towards the large scale dynamics of the magnetic storms and substorms.

Another crucial aspect is the effects of such energetic particles on space observa tions and space weather. Electrons and protons at these energies can lead, for example,
 to spacecraft anomalies related to surface charging (Matéo-Vélez et al., 2018). The pro tons at energies of ~100 keV can enter inside X-Ray telescopes and produce erroneous
 signal leading to loss of astrophysical and cosmological observations (Fioretti et al., 2016).

It is, therefore, necessary to study the origin of energetic plasma which is mainly related to investigation of acceleration mechanisms leading to the observed spectra hardening, distribution of energetic plasma and consequences of its presence on the magnetospheric dynamics.

In this work we present the related results from the observations by the Cluster/RAPID 86 instrument which is described in Section 2. We discuss the origin of the energetic par-87 ticles in the region upstream of the bow shock during solar maximum and declining phase 88 of a solar cycle in Section 3. The acceleration of charged particles in the plasma sheet 89 is considered in Section 4. We discuss the distribution of the energized ions on closed 90 magnetic field lines beyond the radiation belts in Section 5. Features of substorm dy-91 namics revealed using energetic particle observations are shown in Section 6. Influence 92 of energetic ions on the dynamics of the magnetic storms is demonstrated in Section 7. 93 Caveats in measurements of the radiation belts and their dynamics are discussed in Sec-94 tion 8. A summary of this review can be found in Section 9. 95

# <sup>96</sup> 2 Instrumentation and methods

The Cluster mission consists of 4 identical satellites launched July-August 2000 into a highly elliptical polar orbit (initially  $4 \times 19.5 R_E$ ) to explore all regions of the Earth's magnetosphere. The mission and its payload are described by Escoubet et al. (1997).

RAPID (Research with Adaptive Particle Imaging Detectors) (Wilken et al., 1997) is an energetic particle spectrometer, measuring the 3-D distribution of electrons (39 to 400 keV) and ions (protons 28 keV to 4 MeV, helium 137 keV to 4 MeV, oxygen 270 keV to 4 MeV).

The electrons are detected by 3 sets of 3 microstrip solid state detectors, each set being a "pin-hole camera" with a single entrance aperture, covering a total angular range of 60°. The sets, or heads, thus cover the full 180° in the half-plane containing the spin axis, with an angular resolution of 20° in the polar direction. Through the rotation of the spacecraft (period  $\sim 4$  s) the measurements can be sorted into up to 16 azimuthal sectors of 22.5°. The energy of the electrons is determined by the signal strength in the absorbing solid state detector, and is sorted into up to 8 energy channels.

The ions are measured in an independent detection system consisting of again 3 111 heads each covering 60°. Each head has an entrance aperture, a time-of-flight system, 112 and a solid state detector to measure the incoming particle energy. Together, flight time 113 and particle energy allow the ion mass and thus species to be identified. In addition, the 114 incoming direction can be determined electronically to within 15°. Thus each ion head 115 can be subdivided into 4 subdirections, for a total of 12 directions in the half-plane con-116 taining the spin axis. As for the electrons, the ion data can be sorted into 16 azimuthal 117 sectors through the spinning of the spacecraft. The ion energies are sorted into up to 8 118 energy channels. 119

Because of the relatively wide energy channels, it is not immediately obvious to which energy the measured differential flux is to be assigned. The answer is in fact dependent on the steepness of the spectrum. This issue of spectral analysis with wide energy channels is treated in depth by Kronberg and Daly (2013).

The energies measured by RAPID represent the high energy tail of the full spectrum; the thermal energies are covered by the CIS experiment (Rème et al., 2001). How these datasets are combined is treated by Kronberg et al. (2010).

All the Cluster data are readily available online at the Cluster Science Archive (https://
 csa.esac.esa.int/csa-web/). An overview of the RAPID data products is given by
 Daly and Kronberg (2010).

# <sup>130</sup> 3 Origin of energetic particles in the region upstream of the Earth's bow shock

# 3.1 Solar maximum

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Interaction of the supersonic solar wind with the terrestrial magnetic field forms 133 the bow shock. Upstream of the Earth's bow shock one can observe the generation of 134 low frequency hydrodynamic waves, see Figure 2 on the left. Furthermore, the so-called 135 upstream events are observed in this region. These are events during which the enhanced 136 intensities of energetic charged particles are observed, see Figure 3 on the right. It is be-137 lieved that these particles are scattered by the upstream waves leading to diffusive trans-138 port and Fermi-acceleration of the first order (Lee, 1982). It has been shown that, to 139 be observed, stable Interplanetary Magnetic Field (IMF) conditions with direct connec-140 tion to the bow shock are required, e.g., (Lin et al., 1974; Scholer et al., 1979). In Kronberg 141 et al. (2009) we have examined whether this theory can be confirmed by observations 142 from the Cluster mission. An analysis of the ion density gradients has shown that they 143

fall-off exponentially with the distance from the bow shock independent on the species 144 and charge. This means that the dynamics of ions of different species is determined by 145 the same mechanisms. The diffusion coefficient of the ion transport is linear dependent 146 on energy per charge in accordance with the theory from Blandford and Ostriker (1978). 147 The mechanism leads to ion acceleration in the region upstream of the bow shock. This 148 region crosses, for example, XMM Newton X-Ray telescope (Jansen et al., 2001). The 149 upstream ions can affect the observations of the telescope leading to their losses (Kronberg 150 et al., 2020). This mechanism is particularly efficient for long-living stable IMF direc-151 tion and a strong shock. Such conditions are typical for Coronal Mass Ejections (CMEs). 152 CMEs are tremendous eruptions of plasma from the Sun's upper atmosphere which mainly 153 occur during high solar activity. The associated with CME emerging plasma bubbles have 154 strong, relatively stable magnetic field and high propagation speed forming strong shock. 155

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# 3.2 Declining phase of the solar cycle

During the declining phase of a solar cycle the terrestrial magnetosphere often en-157 counters fast solar winds, called High Speed Streams (HSSs). The HSSs overtake reg-158 ular slow solar wind and form shocks named Corotating Interaction Regions (CIRs). HSSs 159 are coupled with Alfvén waves. The interaction between CIRs (associated with shocks 160 and therefore enhanced solar wind dynamic pressure), the Alfvén waves in HSSs (asso-161 ciated with change of IMF  $B_z$  (Desai et al., 2008)) and the terrestrial magnetic field may 162 trigger magnetospheric substorms. Substorms are associated with the magnetic field re-163 configuration, charged particle acceleration and occurrence of polar aurora. Moreover, 164 they lead to a lost of charged particles from the magnetosphere. In 2007 upstream events 165 were discovered at a distance of up to  $1750 R_{\rm E}$  upstream of the Earth's bow shock by 166 the mission STEREO-A (Desai et al., 2008). The origin of such events was at first un-167 clear. We have studied 117 such events with the help of STEREO-A, STEREO-B, ACE, 168 Geotail and Cluster satellites (Kronberg et al., 2011), see Figure 3. Thereby three the-169 ories were tested: (1) leakage of particles from the magnetosphere during substorms, e.g. 170 (Sarris et al., 1978); (2) first order Fermi-acceleration of particles upstream of the bow 171 shock (Lee, 1982); (3) Fermi-acceleration between Alfvén waves in CIR and the bow shock 172 (Desai et al., 2008). It was found that more than 50% of the upstream events have mag-173 netospheric origin. The presence of oxygen ions, energetic electrons, particle flow in the 174 sunward direction and simultaneous occurrence of geomagnetic activity were the main 175 characteristics to support this conclusion. This and also by Kronberg et al. (2015) stud-176 ies showed that the magnetosphere loses a significant amount of plasma also at the day-177 side. In Savin et al. (2014) it was demonstrated that solar wind and foreshock distur-178 bances triggered by resonant oscillations in the outer magnetospheric regions (e.g., in 179 the bow shock surface and the magnetosheath (cavity modes)) modulate oxygen outflow 180 from the magnetosphere into the solar wind. 181

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# 4 Charged particle acceleration in the plasma sheet

Particles trapped on closed magnetospheric magnetic field lines form the plasma 183 sheet at the nightside. This is the main reservoir of the magnetospheric plasma. In the 184 plasma sheet the energy of ions and electrons can increase many orders of magnitude. 185 Adiabatic mechanisms such as betatron and Fermi acceleration lead to electron energiza-186 tion in the plasma sheet during charge particle's drift towards the Earth. However, these 187 acceleration mechanisms have limits. For example, Malykhin, Grigorenko, Kronberg, and 188 Daly (2018) showed that the betatron acceleration works for electrons at energies up to 189 90 keV. The ions are effectively accelerated by non-adiabatic processes. Ion drift along 190 quasi-steady dawn-dusk electric field leads to acceleration of ions up to energies of  $\sim 100$ 191 keV (value of the cross-tail potential drop in the magnetotail during highly disturbed 192 geomagnetic activity). Therefore, much effort has been put into investigations of addi-193



**Figure 2.** Left: Overview of the upstream region in a simulation with out-of-plane magnetic field fluctuations, indicating the extent of the ultra low frequency (ULF) waves in the upstream region. Black curves show magnetic field lines. Taken from Battarbee et al. (2020). Middle: Trajectories of Cluster satellites along the magnetic fields which are shown by arrows. Right: Diffusion coefficient as a function of energy per charge for protons and helium in the region upstream of the Earth's bow shock. Taken from Kronberg et al. (2009).



Figure 3. Left: Cartoon is showing Alfvén waves embedded within the rarefaction regions of the HSSs that follow CIRs. The rarefaction regions are associated with upstream events. Taken from Desai et al. (2008). Right: 1 h averaged proton intensity from the STEREO-A/SEPT antisunward sensor (Müller-Mellin et al., 2007) shows CIRs profiles on the scale of weeks and upstream events as small spikes on scales of hours. Taken from Kronberg et al. (2011).

tional acceleration mechanisms that can boost particle energies from  $\sim 30$  to 100s of keV range.

Many studies indicate that the energy increase is associated with the reconnection 196 of magnetic fields and its consequences. Related mechanisms of electron acceleration were 197 discussed, e.g., in Hoshino (2005); Birn et al. (2012). Cluster/RAPID observations demon-198 strated that electrons can be effectively accelerated while they are flowing into the X-199 line along the separatrix (Wang et al., 2013). These results imply that the electron ac-200 celeration region is significantly larger than the size of the predicted electron diffusion 201 region in the classical Hall magnetic reconnection model. Retinò et al. (2008) showed 202 that the presence of thin current sheets and associated electromagnetic field fluctuations 203 are essential for the acceleration of energetic electrons during reconnection. Electromag-204 netic waves can transfer their energy to particles. 205

A consequence of magnetic field reconnection or other plasma instabilities is the 206 dipolarization which leads to significant increase of the meridional magnetic field com-207 ponent on the scale of 10 minutes, see Figure 4, upper right plot. During development 208 of the dipolarization energetic electron fluxes increase. Malykhin, Grigorenko, Kronberg, 209 and Daly (2018) showed that this is mainly the effect of the betatron acceleration for en-210 ergies up to 90 keV. Asano et al. (2010) and Malykhin, Grigorenko, Kronberg, and Daly 211 (2018) found that there is a need for non-adiabatic acceleration mechanisms to explain 212 electron intensities at higher energies. When the increase of the meridional magnetic field 213 component is sudden and on scales of seconds, such an increase is called dipolarization 214 front (Nakamura et al., 2002; Fu, Khotyaintsev, Vaivads, André, & Huang, 2012). It can 215 be also superimposed on the dipolarization, see Figure 4, upper right plot. Fu, Khotyaint-216 sev, Vaivads, André, Sergeev, et al. (2012); Liu et al. (2017) demonstrated that betatron 217 and Fermi acceleration act behind the dipolarisation fronts. 218

Electromagnetic fluctuations in a wide frequency range play a role in particle heating and acceleration during dipolarizations. For example, high-frequency whistler waves with frequencies between proton plasma frequency and electron gyrofrequency may lead to electron heating and to a decrease in the ratio of hydrogen to electron temperature  $(T_p/T_e)$  – an important parameter for the current sheet dynamics (Grigorenko et al., 2016). In this study it was important to include energetic protons in the calculations of the temperature to avoid its underestimation.

Luo et al. (2014) using Cluster observations showed that the ion acceleration is related to the presence of X-lines near the Earth. The ion diffusion region is not known as an effective ion acceleration region (Birn et al., 2012). Therefore, consequences of reconnection such as plasmoids and magnetic field dipolarizations may lead to ion acceleration. Grigorenko et al. (2015) showed that heavy ions will be more effectively accelerated than protons when they are trapped in a plasmoid associated with certain level of electromagnetic fluctuations.

Ions interacting with dipolarizations associated with low frequency waves in the 233 range between proton and oxygen cyclotron frequencies can be effectively accelerated (Grigorenko 234 et al., 2017). Different particle species such as protons, helium and oxygen are acceler-235 ated by resonant interaction with electromagnetic fluctuations at the corresponding ion 236 component gyrofrequencies. Protons are effectively accelerated by the dipolarization fronts 237 superimposed on the dipolarization. The thickness of the fronts is less than a gyrora-238 dius of thermal protons (Runov et al., 2009; Balikhin et al., 2014) which ensures the nona-239 diabatic acceleration of protons, see Figure 4 (Malykhin, Grigorenko, Kronberg, Koleva, 240 et al., 2018; Parkhomenko et al., 2018). 241

Malykhin et al. (2019) demonstrated that some fraction of light ions can be accelerated up to energies  $\leq 600$  keV and some fraction of oxygen ions can be accelerated up to  $\sim 1.2$  MeV during dipolarization. Such strong energy gains cannot be explained by



Figure 4. Left: Observation of two dipolarizations (shaded in gray) of the magnetic field in the near-Earth tail. Top to bottom: time profiles of the thermal pressure of  $H^+$  and  $O^+$  ions calculated based on CIS/CODIF observations (<40 keV) (shown by the gray lines) and time profiles of the pressure calculated considering the energetic component of the spectrum based on RAPID observations (shown by the black lines); time profile of the  $B_z$  magnetic field component. Adapted from Malykhin et al. (2020). Right: Trajectories and energy variation with time of electrons, protons and oxygen (bottom) during the magnetic field dipolarization associated with multiple dipolarization fronts (top) for times  $t_d$  and  $t_{elf}$ . Taken from Parkhomenko et al. (2018).

acceleration at a single propagating dipolarization front associated with dipolarization.
A possibility of multistage ion acceleration in the course of interaction with multiple dipolarization fronts during dipolarizations (see illustration of such magnetic structures in the upper right Figure 4) is suggested to explain acceleration of ions to such high energies.

In summary, interaction of charged particles with multi-scale magnetic field structures and/or electromagnetic fluctuations which are the consequences of reconnection or other plasma instabilities, leads to effective acceleration. Multiple interaction with such structures boosts the energies even higher.

#### 4.1 Contribution to the plasma pressure

The accelerated ions strongly affect the plasma pressure. The plasma pressure dur-255 ing magnetic field dipolarizations can be underestimated by a factor of >2 if the ions at 256 energies >40 keV are neglected (Kronberg et al., 2017; Malykhin et al., 2020), see Fig-257 ure 4 on the left. This is especially related to contribution of heavy energetic ionospheric 258 ions. Nakamura et al. (2013) have shown that the pressure gradient force exceeding the 259 Earthward directed  $j \times B$  led to the flow bouncing during a dipolarization affecting the 260 dynamics of the associated substorm. In this work it was essential to consider contribu-261 tion to the plasma pressure of the energetic particles measured by the RAPID instru-262 ment. Influence of the energetic ions on the magnetic storm dynamics is described in Sec-263 tion 7. 264

## <sup>265</sup> 5 Distribution on closed magnetic field lines

A considerable part of the charged particles from the solar wind and the ionosphere are trapped on closed magnetospheric magnetic field lines. We studied the dependence of the distributions of the hydrogen and oxygen ions on such closed magnetic field lines



Figure 5. Left: Comparison of the oxygen pressure before the substorm and during in simulations by (Fok et al., 2006) with observations of the oxygen intensities by Cluster from Kronberg et al. (2015). Taken from Kronberg et al. (2014). Right: Map of the cold ion sources in the polar region of the Northern hemisphere and associated sinks in the plasma sheet during different geomagnetic activities. Taken from Li et al. (2013).

on the geomagnetic activity, the solar wind dynamic pressure and the direction of the 269 IMF using 7 years of Cluster observations (Kronberg et al., 2012, 2015). The distribu-270 tion of energetic ions (>274 keV) shows duskward asymmetry which is different from the 271 rather symmetric distribution of the ions at bulk energies ( $\sim 10 \text{ keV}$ ) during geomagnet-272 ically active times (Kronberg et al., 2017). Especially strong asymmetry is seen in re-273 lation with AE-index and for oxygen ions, see Figure 5. The AE-index measures the strength 274 of the currents in the auroral zone and indicates the strength of substorms (Nose et al., 275 2017). Substorms are associated with magnetic field dipolarizations and plasmoid for-276 mations which effectively accelerate charged particles as described in Section 4. An ac-277 celeration of the ions by non-adiabatic acceleration mechanisms (see Section 4 and Kronberg 278 et al. (2015)) and gradient and curvature drifts in the duskward direction can explain 279 duskward asymmetry. The observations agree well with the models, see, e.g., Figure 4. 280 Another explanation of the duskward asymmetry is the distribution of cold ions with to-281 tal energy <70 eV. During geomagnetically active times large amount of cold ions pop-282 ulate the plasma sheet closer to the Earth and more duskward than during geomagnet-283 ically quiet times due to stronger convection, see Figure 5 (Li et al., 2013). These cold 284 ions can get effectively accelerated to many 100s of keV just within 10s of seconds, see 285 e.g. Grigorenko et al. (2015), leading to duskward distribution of ions. This can also ex-286 plain why duskward asymmetry is not observed at bulk energies for oxygen ions at dis-287 tance >10  $R_E$  during geomagnetic active times: cold ions are accelerated so fast that 288 they do not spend much time at  $\sim 10$  keV energies. 289

It is remarkable that the distribution of the ions also depends on the direction of 290 the IMF within the ecliptic plane. One can observe higher ion intensities in the regions 291 in which reconnection between the IMF and the terrestrial magnetic field is expected. 292 This leads to asymmetric ion distributions not only between dawn and dusk but also be-293 tween the Northern and Southern hemispheres, see Figure 6. This was discovered by Luo 294 et al. (2017) using 11 years of Cluster energetic ion (>274 keV) observations. On the day-295 side such asymmetry can be caused by the local acceleration in diamagnetic cavities (Nykyri 296 et al., 2012) or entering of the particles accelerated in the upstream region of the bow 297 shock via reconnection into the magnetosphere (Trattner et al., 2011), see Figure 6. There-298



**Figure 6.** Left: Oxygen intensities with energy from 274 to 962 keV under southward IMF with different IMF By directions in both the Northern and Southern Hemispheres. The gray circles with antiparallel arrows indicate quadrants where the reconnection and location of diamagnetic cavity is expected. Right: Arrows indicate the Earth's magnetic field and the draped IMFs for each orientation. Shaded gray regions show under each IMF orientation where the most antiparallel components are located relative to the cusp region and where the diamagnetic cavity can form and trap the particles. Taken from Luo et al. (2017).

fore, not only substorm related acceleration processes lead to asymmetries on the closed magnetic field lines.

#### 301 6 Substorm dynamics

Energetic particles help to reveal different aspects related to substorm dynamics, 302 for example causes of associated dispersionless injections of electrons at geostationary 303 orbits. A classic substorm event with several dispersionless injections was investigated 304 using data from 15 satellites by Kronberg et al. (2017). These were data from 4 Clus-305 ter satellites, 3 THEMIS satellites, 6 LANL satellites and 2 Van Allen Probes. These data 306 were expanded by the magnetic field observations measured by many stations at high 307 latitudes from INTERMAGNET (St-Louis, 2008) and SuperMAG (Gjerloev, 2012) net-308 works. Such injections feed with charged particles the ring current and the radiation belts, 309 the most important regions for the space weather effects. The origin of dispersionless in-310 jections is controversial. In this study it was shown that during different phases of the 311 substorm event the reasons for the injections can be different: dynamic motion of plasma 312 sheet structures during the growth phase, current wedge formation during the substorm 313 onset and magnetic dipolarization fronts during the expansion phase. Only the injection 314 associated with the substorm onset and formation of the current wedge filled the radi-315 ation belts and ring current with a substantial number of charged particles. 316

Another important result of this study is that the currents along the magnetic fields formed as a result of the substorm and related to the magnetosphere-ionosphere coupling, can be responsible for the electron acceleration at geostationary orbits. The acceleration leads to the formation of a so-called bump on the tail of the electron intensity spectra at the geosynchronous orbit but not in the tail region observed by Cluster. It is suggested that this effect occurs because of drift-resonant interaction with ULF waves generated by the field-aligned currents.

# <sup>324</sup> 7 Magnetic storms

The accelerated particles move in the direction of the Earth and then drift around, 325 forming the ring current. The ring current generates a magnetic field in the southward 326 direction at the Earth's surface, namely a field in the opposite direction to the terres-327 trial magnetic field, produced by the Earth's dynamo. The terrestrial magnetic field will 328 be reduced and this disturbance is called magnetic storm, one of the most important char-329 acteristics of the space weather. The influence of the ionospheric and energetic particles 330 on the strength of the magnetic storm was investigated in Kronberg et al. (2017). Both 331 332 of these components are often neglected in the calculation of the plasma pressure as in the plasma sheet as well in the ring current. Haaland et al. (2010) showed that the power 333 law part of the proton spectra (energetic part) becomes significantly harder in the plasma 334 sheet during magnetic storms. Our calculations demonstrated that the plasma pressure 335 and as a consequence the Dst-index (represents the strength of the magnetic storm near 336 the equator (Nose et al., 2017)) are strongly underestimated when the ionospheric ions 337 and ions at energies > 40 keV are neglected. The Space Weather Modeling Frame (Tóth 338 et al., 2012) simulations with two different ionospheric ion outflow models as an input 339 have shown similar to observations results. The space weather modeling, however, still 340 can be improved by considering energetic part of the particle spectra. 341

## 342 8 Radiation belts

The Van Allen radiation belts are a region with highly-energetic (>100 keV) charged 343 particles, mainly electrons and protons, trapped by the terrestrial magnetic field. The 344 radiation belts are hazardous for satellite operation. The electrons with energies of tens 345 of keV, injected into the inner magnetosphere as a result of substorm activity, produce 346 waves, and can be accelerated to higher energies. The particles with energies of hundreds 347 of keV can accumulate at the surface of the spacecraft, leading to surface-charging ef-348 fects, and can penetrate deep into instrument electronics causing satellite anomalies and 349 instrument malfunctions by deep dielectric charging, e.g. (Baker et al., 1994, 2004). 350

Due to their importance for satellite operation, a large number of instruments have 351 provided observations of the radiation belt populations. One of the longest continuous 352 data sets of electron flux at energies 40 to 400 keV has been supplied by the RAPID/IES 353 detector. The IES instrument was not specifically designed for radiation belt studies, and 354 therefore is subject to background contamination by higher energy electrons and pro-355 tons. The effect of the charged particle radiation belt environment on the RAPID/IES 356 detector was investigated in Kronberg et al. (2016). For this the IES detector was sim-357 ulated in the Geant4 environment (Agostinelli et al., 2003). The percentages of the to-358 tal electron intensity attributed to contamination were calculated for L-shells up to 9. 359 The background-corrected measurements were compared to the observations by the MagEIS 360 detector aboard the Van Allen Probes mission (Blake et al., 2013) during October-December 361 2012, and the ratio between the two data sets was close to 1. 362

Smirnov et al. (2019) compared the background-corrected electron flux measurements by the IES detector to those by RBSP/MagEIS in 2012-2016, both in terms of omnidirectional flux and flux at 90° pitch angle. It was found that the corrected Cluster measurements were close to the MagEIS observations during the 2012-2016 period. These background-corrected observations can be used in various scientific applications. Breuillard et al. (2015) derived energetic electron anisotropy to model the chorus wave power distribution in the outer radiation belt. Smirnov et al. (2019) presented a long-term analysis of electron flux dynamics during  $\sim 1.5$  solar cycles in 2001-2016. It was found that the long-term trend of electron intensities at energies 40 to 400 keV in the outer belt exhibited a high positive correlation with the solar wind dynamic pressure and AE index (see Figure 7). A simple generalized empirical model which connects electron intensi-



Figure 7. (a) The solar wind dynamic pressure, and (b) yearly averaged electron flux intensities by RAPID/IES. Adapted from Smirnov et al. (2019).

ties at L-shells from 4 to 6 to the solar wind dynamic pressure was derived:

$$\log_{10} y = 0.35(\pm 0.1) + 0.5(\pm 0.18)x,\tag{1}$$

where x stands for the normalized solar wind dynamic pressure and y represents the logarithm of electron flux, normalized to their minimum and maximum values.

Trapped energetic particles in planetary magnetospheres undergo three types of 365 periodic motion, namely the gyro-, bounce, and drift motion. Each type is associated 366 with the corresponding adiabatic invariant ( $\mu$ , K, and  $L^*$ , respectively), which are con-367 served when the magnetic field changes are slow compared to the motion periods. Elec-368 tron flux, measured as a function of energy, position, and time, can be converted to phase 369 space density (PSD), as a function of three adiabatic invariants. Analyzing the PSD be-370 havior under the fixed values of  $\mu$ , K, and  $L^{\star}$ , allows for a separation between adiabatic 371 and non-adiabatic processes, and helps to utilize observations from different magnetic 372 latitudes. This is of particular interest for the Cluster constellation, due to its polar or-373 bit. The method to calculate the adiabatic invariants and convert the RAPID/IES elec-374 tron flux to PSD was developed in Smirnov et al. (2020). Figure 8 shows the long-term 375 behavior of low- $\mu$  electrons along approximately 1.5 solar cycles and the 27-day averaged 376 F10.7 and AE indices. F10.7 index (Tapping, 2013) follows the sunspot cycle evolution, 377 while the AE index intensifies with higher occurrence of coronal holes and HSSs during 378 the declining phase. At L-values from 4 to 6, the general PSD behavior closely resem-379 bles the variation of the AE index. The PSD values are higher during the declining phase 380 of the solar cycle 23 (2004–2007), and exhibit minimum values during the solar minimum 381 in 2009, also referred to as the "radiation belts desert". RAPID/IES mainly samples low-382  $\mu$  electrons, for which the substorm activity is of paramount importance. These results 383 agree well with previous findings by Zhao et al. (2017) that observed a high correlation 384 of PSD enhancements with auroral geomagnetic indices, suggesting a direct transport 385 of electrons from substorm injections. 386

#### 387 9 Summary

The observations of energetic charged particles by the Cluster/RAPID instrument substantially enhanced our understanding of dynamics of the terrestrial magnetosphere. The following results were obtained:



Figure 8. (a) PSD for Cluster-1 to Cluster-4 under the fixed values of adiabatic invariants in 2004-2019; (b) AE and F10.7 indices. Adapted from Smirnov et al. (2020)

Charged particle energization mechanisms were thoroughly investigated. The results demonstrate that interaction of charged particles with multi-scale magnetic field structures and/or electromagnetic fluctuations which are the consequences of reconnection or other plasma instabilities leads to an effective acceleration. Multiple interaction with such magnetic structures increases energy even higher.

The energetic particles affect the plasma temperature and pressure. The Cluster results demonstrated that the energetic part >40 keV has to be considered during geomagnetically disturbed times in the calculations of the ion plasma temperature and pressure. Otherwise, the related physics, e.g., in flow breaking region or in the ring current is incorrectly interpreted. Therefore, it is recommended to include the energetic part and ionospheric heavy ions when estimating these parameters.

The distribution of the energetic ions on the closed magnetic field lines is strongly affected by the solar wind and geomagnetic activity. A substantial amount of ions escape the magnetosphere populating the region upstream of the bow shock. Surprisingly the distribution is also affected by the IMF  $B_y$  direction leading to dawn-dusk asymmetries as well as to asymmetries between Northern and Southern hemispheres. This also implies the importance of reconnection processes effects on the energetic ion distribution not only at the night side but also at the dayside of the magnetosphere.

The radiation belts are effectively filled with electrons during substorms onsets. Clear relation of electron intensities in the radiation belts at L-shells 4 to 6 and substorm activity was demonstrated for approximately 1.5 solar cycles.

A large data set of electron intensities in the radiation belts and ring current region is available for scientific investigations, thanks to the recently developed backgroundcorrection technique. The Cluster constellation is one of the longest ongoing missions in the Earth's magnetosphere, providing more than 20 years of observations. The IES detector, similar to that on Cluster, was carried by the Polar mission and is also included in the payload of the recently launched BeiDou GNSS constellation. The combined data set thus covers more than 2 solar cycles. It can be useful for the long-term statistical anal-

419 ysis of electron flux behavior in the outer radiation belt, and can be incorporated into

 $_{420}$  the empirical models of the medium energy electron flux, e.g, by Smirnov et al. (2020).

421 Moreover, the electron fluxes converted to PSD can be used to set up boundary condi-

tions for radiation belts modeling, as well as for data assimilation purposes.

The vast amount of RAPID data collected within the last 20 years, available at the Cluster Science Archive, is welcoming future explorations.

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