# Strong super-rotation during the 2018 Martian Global Dust Storm

Kylash Rajendran<sup>1,1</sup>, Stephen Lewis<sup>2,2</sup>, James Andrew Holmes<sup>2,2</sup>, Paul Michael Streeter<sup>2,2</sup>, Anna A. Fedorova<sup>3,3</sup>, and Manish Patel<sup>2,2</sup>

<sup>1</sup>The Open University <sup>2</sup>Open University <sup>3</sup>Space Research Institute

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#### Abstract

Super-rotation affects - and is affected by - the distribution of dust in the martian atmosphere. We modelled this interaction during the 2018 global dust storm (GDS) of Mars Year 34 using data assimilation. Super-rotation increased by a factor of two at the peak of the GDS, as compared to the same period in the previous year which did not feature a GDS. A strong westerly jet formed in the tropical lower atmosphere, with strong easterlies above 60 km, as a result of momentum transport by thermal tides. Enhanced super-rotation is shown to have commenced 40 sols before the onset of the GDS, due to equatorward advection of dust from southern mid-latitudes. The uniform distribution of dust in the tropics resulted in a symmetric Hadley cell with a tropical upwelling branch that could efficiently transport dust vertically; this may have significantly contributed to the rapid expansion of the storm.

# Enhanced Super-rotation Before and During the 2018 Martian Global Dust Storm

# Kylash Rajendran<sup>1</sup>, Stephen R. Lewis<sup>1</sup>, James A. Holmes<sup>1</sup>, Paul M. Streeter<sup>1</sup>, Anna A. Fedorova<sup>2</sup>, and Manish R. Patel<sup>1,3</sup>

<sup>1</sup>School of Physical Sciences, The Open University, Milton Keynes, UK <sup>2</sup>Space Research Institute of the Russian Academy of Sciences (IKI RAS), Moscow, Russia <sup>3</sup>Space Science and Technology Department, Science and Technology Facilities Council, Rutherford Appleton Laboratory, Oxfordshire, United Kingdom

#### Key Points:

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| 10 | • | The martian atmosphere was already in a state of enhanced super-rotation prior   |
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| 11 |   | to the onset of the Mars Year 34 global dust storm                               |
| 12 | • | Super-rotation doubled during the peak of the storm and tropical easterlies were |
| 13 |   | strongly enhanced above 60 km  |
| 14 | • | Equatorward transport of dust from southern mid-latitudes led to enhanced ver-   |
| 15 |   | tical transport in the tropics in the lead-up to the storm                       |

Corresponding author: Kylash Rajendran, kylash.rajendran@open.ac.uk

#### 16 Abstract

Super-rotation affects - and is affected by - the distribution of dust in the martian at-17 mosphere. We modelled this interaction during the 2018 global dust storm (GDS) of Mars 18 Year 34 using data assimilation. Super-rotation increased by a factor of two at the peak 19 of the GDS, as compared to the same period in the previous year which did not feature 20 a GDS. A strong westerly jet formed in the tropical lower atmosphere, with strong east-21 erlies above 60 km, as a result of momentum transport by thermal tides. Enhanced super-22 rotation is shown to have commenced 40 sols before the onset of the GDS, due to equa-23 torward advection of dust from southern mid-latitudes. The uniform distribution of dust 24 in the tropics resulted in a symmetric Hadley cell with a tropical upwelling branch that 25 could efficiently transport dust vertically; this may have significantly contributed to the 26 rapid expansion of the storm. 27

## <sup>28</sup> Plain Language Summary

Dust plays a major role in driving the behaviour of the atmosphere of Mars. Dur-29 ing a global dust storm, winds lift and transport dust throughout the atmosphere; in turn, 30 dust affects wind direction and strength by heating and cooling the surrounding air. Us-31 ing a technique that combines satellite observations with simulations of the martian at-32 mosphere, we demonstrate that winds at tropical latitudes were greatly strengthened dur-33 ing the 2018 global dust storm as a result of heating by dust. We show that tropical winds 34 were strengthened even before the onset of the storm, as a result of a dust-driven mod-35 ification to the tropical circulation pattern. This change in the tropical circulation may 36 have played a role in the formation of the global dust storm. 37

### 38 1 Introduction

A planetary atmosphere is said to be in a state of *super-rotation* if the total ax-30 ial angular momentum of the atmosphere is greater than the angular momentum of an 40 atmosphere in a state of pure solid body rotation (Read, 1986). Super-rotation usually 41 takes the form of an equatorial jet directed in the same sense as the planet's rotation, 42 and it has been shown that such a state can only be maintained by the equatorward trans-43 port of angular momentum by non-axisymmetric eddy motions (Hide, 1969). Identify-44 ing the exact mechanisms that initiate and maintain super-rotating jets remains a sig-45 nificant challenge in planetary atmospheric research, and is crucial for understanding the 46 atmospheric dynamics of slow-rotating terrestrial planets such as Venus and Titan (Read 47 & Lebonnois, 2018). The Martian atmosphere is also a candidate for super-rotation, as 48 first identified by Lewis and Read (2003). 49

Global dust storms (GDS) are the most dramatic of all dust-related phenomena 50 in the martian atmosphere, and occur every few martian years. During a GDS, the planet 51 is encircled by a shroud of dust that can persist for several months at a time, drastically 52 affecting the atmospheric state. Such storms have all been observed to occur during the 53 high dust loading season ( $L_{\rm S}=180^{\circ}-360^{\circ}$ ), when the planet's orbit brings it closer to the 54 Sun (Kahre et al., 2017). The formation mechanisms of a GDS are still not well under-55 stood; in particular, it has not been established why a GDS forms out of regional storms 56 in some years, but not others. Posited hypotheses include the re-distribution of dust be-57 tween finite reservoirs between years (Mulholland et al., 2013; Newman & Richardson, 58 2015), competition between the hemispheric circulations during northern winter (Haberle, 59 1986), weak coupling between orbital and rotational angular momentum on Mars (Shirley, 60 2017; Newman et al., 2019) and enhanced surface drag due to constructive interference 61 of tides (Montabone et al., 2008; Martinez-Alvarado et al., 2009). 62

Lewis and Read (2003) have shown that dust-driven heating can excite a superrotating jet in the Martian atmosphere. In turn, the strengthening of the equatorial jet by the presence of dust affects dust transport; in this way, there is an intimate link between dust-driven heating and the strength of the equatorial jet which may have implications for GDS evolution.

The most recent GDS occurred in Mars Year (MY) 34 (using the naming conven-68 tion of Clancy et al. (2000)). The storm initiated shortly after the northern hemisphere 69 autumn equinox at  $L_{\rm S}=187^{\circ}$ , and its dynamical evolution was monitored by multiple or-70 bital, surface and Earth-based instruments, e.g. (Guzewich et al., 2019; Hernández-Bernal 71 et al., 2019; Kass et al., 2019; Smith, 2019; Sánchez-Lavega et al., 2019; Shirley et al., 72 73 2020). Bertrand et al. (2020) simulated the MY34 GDS using the NASA Ames Mars Global Climate Model (MGCM) and noted that eastward tropical expansion was a dominant 74 feature of both the MY25 and MY34 GDS, highlighting the link between super-rotation 75 and dust transport in the early stages of GDS formation. Gillespie et al. (2020) used 76 the EMARS reanalysis to study dust encirclement of the northern hemisphere in the early 77 stages of the GDS, and showed that a significant portion (16%) of the encirclement could 78 be attributed to wind-driven dust advection. Super-rotation during the MY34 GDS was 79 further examined by Montabone et al. (2020), who simulated the GDS period using the 80 MGCM developed at the Laboratoire de Météorologie Dynamique (LMD). They reported 81 a global super-rotation value of S = 16% during the peak of the MY34 GDS, as com-82 pared to a pre-dust storm value of 5% (super-rotation metrics are defined in section 2.3). 83

In this paper we perform an analysis of super-rotation during MY33–34, with a spe-84 cial focus on the period leading up to the onset of the MY34 GDS ( $L_{\rm S}=160^{\circ}-187^{\circ}$ ) which 85 has not been studied by previous authors. MY33 was used for comparisons as it is an 86 adjacent year that did not have a GDS. For this work we analyse the outputs from a data 87 assimilation scheme (Lewis et al., 2007) that was used to assimilate temperature and col-88 umn dust data from the Mars Climate Sounder (MCS) and Atmospheric Chemistry Suite 89 (ACS) instruments into a MGCM. As data assimilation uses observational data to con-90 strain the evolution of model dynamics, it provides an excellent tool to perform such anal-91 ysis. In Section 2, we describe the satellite observations and give a description of the model 92 and assimilation scheme used. We also introduce metrics to quantify super-rotation. We 93 document our results in Section 3, and discuss our findings in Section 4. 94

### 95 **2** Data and Methods

#### 2.1 Observation data

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MCS is a passive 9-channel limb-scanning radiometer aboard the Mars Reconnais-97 sance Orbiter (MRO) (McCleese et al., 2007). Measurements are taken at approximately 98 3 am and 3 pm local time along a Sun-synchronous orbit, across a wavelength range of 99  $0.3-45\,\mu\mathrm{m}$ . Radiance profiles are generated from the surface to around 80 km with an 100 intrinsic vertical resolution of  $5 \,\mathrm{km}$ , enabling the retrieval of profiles of temperature, dust 101 and water ice (Kleinböhl et al., 2009, 2017). We used version v5.2 for all MCS data out-102 side of the MY34 GDS period. During the GDS period  $(L_{\rm S}=180^{\circ}-240^{\circ})$ , v5.3.2 was used 103 instead, as it extracts additional dust information using a water ice channel. This is pos-104 sible in GDS conditions due to the absence of water ice clouds in the lower atmosphere 105 (Montabone et al., 2020). 106

ACS (Korablev et al., 2018) is an array of three infrared spectrometers aboard the 107 ExoMars Trace Gas Orbiter (TGO), that together provide spectral coverage over a wave-108 length range of  $0.7-17 \,\mu$ m. The near-infrared (NIR) channel uses an echelle grating with 109 an Acousto-Optical Tunable filter with a resolving power of 25,000, to retrieve atmospheric 110 density and temperature profiles based on the 1.57  $\mu m \text{ CO}_2$  band (Fedorova et al., 2020). 111 The orbit of TGO allows for solar occultations that cross the terminator over a range 112 of local times, a key feature of these datasets. ACS-NIR temperature profiles were in-113 cluded in our assimilation of MY34 over the period  $L_{\rm S}=163^{\circ}-360^{\circ}$ . 114

### 2.2 Global Climate Model and data assimilation scheme

The Martian atmosphere is modelled using the UK version of the LMD MGCM 116 (Forget et al., 1999). The model utilises a spectral dynamical core (Hoskins & Simmons, 117 1975), together with a vertical finite difference scheme that conserves energy and angu-118 lar momentum (Simmons & Burridge, 1981) and a semi-Lagrangian tracer advection scheme 119 (Newman et al., 2002). Physical processes on Mars are represented using the physics pack-120 age from the LMD model, which includes schemes for  $CO_2$  condensation and sublima-121 tion, radiative transfer, dust transport and boundary layer processes (Forget et al., 1999; 122 123 Madeleine et al., 2011; Colaïtis et al., 2013) amongst many others. Dust transport is via a radiatively active two-moment scheme that advects the dust mass mixing ratio and num-124 ber density, which are then used to infer the particle size distribution at each point for 125 radiative transfer calculations (Madeleine et al., 2011). 126

The data assimilation scheme is an implementation of the Analysis Correction (AC) 127 scheme, originally developed for terrestrial applications (Lorenc et al., 1991) and later 128 modified and re-tuned for the Martian context (Lewis et al., 1996, 2007). The scheme 129 employs a form of successive corrections, with analysis increments interlaced between dy-130 namical timesteps. It has been successfully employed in assimilating profiles of temper-131 ature and column dust opacity from the Thermal Emission Spectrometer (TES) (Lewis 132 & Barker, 2005; Montabone et al., 2005; Lewis et al., 2007) MCS (Steele et al., 2014; Holmes 133 et al., 2019, 2020; Streeter et al., 2020) and ACS (Streeter et al., 2021) instruments. 134

We performed the assimilation over MY33–34, and ran the model at T42 spatial 135 resolution (triangular truncation at horizontal wavenumber 42), which corresponds to 136 a physical grid resolution of 3.75°. Fifty-five unevenly-spaced, terrain-following vertical 137 levels were used, with the model top at approximately 110 km. The water cycle was not 138 simulated, in order to isolate the dynamical impact of the dust. Available temperature 139 profiles from the MCS and ACS instruments were assimilated at each timestep, and dust 140 opacity at each model level was rescaled so as to match the total column dust optical 141 depth (CDOD) in the model to MCS observations. Details of the temperature and dust 142 assimilation schemes are given in Lewis et al. (2007) and in the Supplementary Mate-143 rial of Streeter et al. (2020). The inclusion of ACS temperature profiles in the assimi-144 lation did not produce any significant changes in the overall structure of the circulation 145 or the strength of the jet during MY34, as compared to an assimilation that only included 146 MCS temperature profiles; this confirms the mutual coherence of the different data sources, 147 and the robustness of the assimilation procedure. 148

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#### 2.3 Super-rotation metrics

Super-rotation can be quantified by defining a global super-rotation index S (Lewis & Read, 2003) as

$$S = \frac{\iiint \rho u a \cos \phi \, dV}{\iiint \rho \Omega a^2 \cos^2 \phi \, dV} \times 100\%,\tag{1}$$

where  $\rho$  is density, u is the zonal wind velocity, and a,  $\phi$  and  $\Omega$  are the planetary radius, latitude and rotation rate respectively. The volume element is  $dV = a^2 \cos \phi \, d\lambda \, d\phi \, dz$ , where  $\lambda$  and z are east longitude and geometric height respectively. Each integral is performed over the whole atmospheric volume. S is a measure of the ratio between the atmospheric and solid-body components of angular momentum; S > 0 only if the atmosphere as a whole has more angular momentum than a pure solid body rotation.

Super-rotation can also be diagnosed by defining a local super-rotation index s as

$$s = \left[\bar{m}/\Omega a^2 - 1\right] \times 100\%. \tag{2}$$

Here  $\bar{m}$  is the zonally averaged value of axial angular momentum m, which is defined as

$$m = a\cos\phi \left[\Omega a\cos\phi + u\right]. \tag{3}$$

<sup>156</sup> Whereas S measures global super-rotation, s is a measure of its spatial distribution. s<sup>157</sup> compares the angular momentum of a ring of fluid at a given latitude and height against <sup>158</sup> the angular momentum of an equal-mass fluid ring at rest at the equatorial surface (Read <sup>159</sup> & Lebonnois, 2018). Positive values of s represent local super-rotation.

#### 160 3 Results

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#### 3.1 Atmospheric super-rotation

Fig. 1 shows the variation of the global super-rotation metric S through MY33–34 162 (lower panel), along with corresponding equatorial (10°N–10°S) CDOD values (upper 163 panel). The large peak in S centered at  $L_{\rm S}=200^{\circ}$  in MY34 clearly corresponds to the 164 mature phase of the GDS (Kass et al., 2019). There are also several smaller local max-165 ima that occur on both curves. The peaks at  $L_{\rm S}=105^{\circ}$  in MY33 and  $L_{\rm S}=135^{\circ}$  in MY34 166 have been identified as spurious and are due to short absences in MCS data availabil-167 ity over those periods. The MY34 curve has additional peaks at  $L_{\rm S}=175^{\circ}$  and  $L_{\rm S}=325^{\circ}$ . 168 The peak at  $L_{\rm S}=175^{\circ}$  occurs just prior to the initiation of the GDS, and will be discussed 169 further in section 3.3. The peak at  $L_{\rm S}=325^{\circ}$  corresponds to the timing of the regional 170 'C' storm that occurred late in the year, as can be seen from the concurrent increase in 171 CDOD values. 172

Outside of the global dust storm period, the values of S are broadly similar in both years. Global super-rotation variations have a semi-annual structure, with broad peaks occurring during equinoxes and troughs occurring during the solstices. These changes reflect the changes in the Hadley cell structure over the course of the year (Lewis & Read, 2003). S remains largely positive throughout both years of study, except during northern summer solstice. The average values of S are 3.4% and 3.9% for MY33 and MY34 respectively.

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## 3.2 Super-rotation during the MY34 GDS

The major feature in Fig. 1 is a large peak in S corresponding to the mature phase 181 of the MY34 GDS. Global super-rotation in MY34 first diverges significantly from MY33 182 values at  $L_{\rm S}=165^{\circ}$ , when the dust opacity in the tropics first begins to increase. After 183 an initial peak to S = 7.2% at  $L_{\rm S} = 173^{\circ}$ , global super-rotation increases sharply upon 184 storm onset in the Acidalia corridor at  $L_{\rm S}=187^{\circ}$ . After reaching a plateau at S=11.3%, 185 the value of S was boosted again by the onset of intense dust lifting in the Tharsis re-186 gion at  $L_{\rm S}=197^{\circ}$  (Bertrand et al., 2020; Montabone et al., 2020), reaching a maximum 187 of S = 12.6% at  $L_{\rm S} = 201^{\circ}$ . The peak value is twice as large as the corresponding MY33 188 value of S = 6.2%, indicating a significant increase in atmospheric angular momentum. 189 After the peak, S decays at a uniform rate, and returns to background values at  $L_{\rm S}=240^{\circ}$ . 190

Fig. 2 depicts values of zonal-mean zonal winds and the local super-rotation index across latitude and height, during the MY34 GDS period in both years. The data have been averaged in time over 20° periods of solar longitude. These broadly cover the initiation, peak and decay phases of the storm. In all periods, super-rotation is stronger in MY34 than in MY33.

In the time period of storm initiation (Figs. 2(a)–(b)), the atmosphere is in an equinoctial state. There are extratropical westerly jets in both hemispheres, and easterlies over much of the tropics in MY33. There is a substantial difference in tropical wind morphology between the years; this is reflected by the local super-rotation values in the tropics, which peak at  $s\approx 12\%$  in MY34 compared to  $s\approx 1\%$  in MY33. Tropical westerlies extend up to 50 km in MY34.

Super-rotation is strongest in the period  $L_{\rm S} = 200^{\circ} - 220^{\circ}$  of MY34 (Fig. 2(d)), which corresponds to the peak of the GDS. The super-rotating jet continues to domi-



Figure 1. Values of the global super-rotation metric S during MY33 (blue) and MY34 (red). Positive values indicate global super-rotation. The upper panel shows the corresponding zonal mean CDOD values averaged over  $10^{\circ}$ N- $10^{\circ}$ S.



Figure 2. Zonal-mean zonal wind (black contours) and local super-rotation index (filled contours) during the MY34 GDS period for MY33 (left) and MY34 (right). Wind contours are drawn at  $10 \text{ m s}^{-1}$  intervals, with westerlies (easterlies) denoted by solid (dashed) contours. The bold solid contour is the zero wind line. Maximal wind values are marked in red.



Figure 3. Profiles of tropical zonal-mean zonal winds (averaged diurnally and across  $30^{\circ}$  S –  $30^{\circ}$  N) for different times during the GDS period for MY33 (left) and MY34 (right).

| 204 | nate the tropical band, and extends up to 60 km. There is a strong peak in local super-            |
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| 205 | rotation of $s = 22\%$ at 20 km, reflecting a significant acceleration of tropical winds by        |
| 206 | the increased dust presence. Easterly winds above 60 km are also enhanced in MY34, and             |
| 207 | increase with height to $-100 \text{ m s}^{-1}$ near the top of the model. In comparison, tropical |
| 208 | easterly winds in MY33 peak at $-40 \text{ m s}^{-1}$ at 50 km, and remain relatively constant at  |
| 209 | higher altitudes (Fig. $2(c)$ ).   |

In the decay phase of the storm (Figs. 2(e)-(f)), easterly winds dominate the tropical regions above 20 km and increase monotonically in MY34 to a peak of  $-140 \text{ m s}^{-1}$ . Super-rotation is confined to the lower atmosphere below 30 km, and is greatly reduced in intensity.

Fig. 3 shows the evolution of zonal-mean zonal wind profiles in the tropics over the course of the GDS period for MY33 (left) and MY34 (right). Wind profiles at each solar longitude were averaged diurnally and between 30° S-30° N.

The tropical zonal wind profile does not change substantially between  $L_{\rm S}=165^{\circ}-$ 217  $210^{\circ}$  of MY33. Easterly winds in the lowermost atmosphere are overlain by westerly winds 218 with velocities of around  $5 \,\mathrm{m\,s^{-1}}$  (Fig. 3(a)). The wind profile for  $L_{\mathrm{S}}=165^{\circ}$  in MY34 is 219 similar to the winds in MY33, but by  $L_{\rm S}=170^{\circ}$  zonal wind speeds have increased at al-220 most all heights (Fig. 3(b)). The wind profile weakens slightly at  $L_{\rm S}=185^{\circ}$ , before strength-221 ening substantially to reach peak values of  $35 \,\mathrm{m\,s^{-1}}$  at  $L_{\mathrm{S}}=210^{\circ}$ , with westerly winds pen-222 etrating all the way down to the lowest model level (5 m above the surface). In the de-223 cay phase of the GDS period, wind profiles from MY34 weaken and start to converge with 224 profiles from MY33. 225

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# 3.3 Atmospheric circulation prior to GDS onset

Fig. 4 shows the evolution of the meridional circulation and the dust distribution in the assimilation for both years over the period  $L_{\rm S}=158^{\circ}-185^{\circ}$ , prior to GDS initiation at  $L_{\rm S}=187^{\circ}$ . There are no substantive differences between the years at  $L_{\rm S}=158^{\circ}$  (Fig. 4(a)– (b)). At  $L_{\rm S}=161^{\circ}$  in MY34, turbulent dynamics along the southern polar front causes a mass of dust from the southern mid-latitudes to be transported northward and to en-



Figure 4. Zonal mean dust distribution and dynamical structure at four times (rows) in the lead-up to the GDS for MY33 (left) and MY34 (right). Light (dark) orange shading shows regions of greater than  $1 \times 10^8$  ( $3 \times 10^8$ ) dust particles per kg. Blue-green filled contours indicate the local super-rotation index. Solid (dashed) lines show positive (negative) contours of the mass streamfunction, drawn at  $\pm (0.5, 1, 2, 3, 4, 5, 10, 20, 30, 40, 50) \times 10^8 \text{ kg m}^{-1} \text{ s}^{-1}$ .

croach into the tropical region. Some of this dust is entrained into the Hadley cell circulation and is transported to higher levels in the atmosphere as well as into the northern hemisphere (Fig. 4(d)). The resulting dust distribution is more uniformly spread across
tropical latitudes compared to MY33, where the bulk of the dust remains in the southern hemisphere (Fig. 4(c)).

The impact of the different dust distributions on the circulation can be seen in Figs. 4(e)-237 (h). In MY33, dust-driven heating is enhanced in the southern mid-latitudes compared 238 to the tropics, causing an inversion of the diabatic heating gradient in the southern hemi-239 sphere between 20–40 km. This weakens the thermally direct circulation in the south-240 ern hemisphere, resulting in the formation of a lopsided Hadley cell with a wavy upwelling 241 branch that is less efficient at vertical tracer transport (Figs. 4(e),(g)). In contrast, the 242 relatively uniform dust distribution in MY34 engenders a symmetrical Hadley cell with 243 an upwelling branch closely aligned to the vertical, as well as enhanced local super-rotation 244 (Figs. 4(f),(h)). The circulation patterns in both years do not change substantially be-245 tween  $L_{\rm S} = 170^{\circ} - 185^{\circ}$ , suggesting that they are relatively stable. 246

### <sup>247</sup> 4 Discussion

Our MGCM with data assimilation predicted a peak value of S = 12.6% during the peak of the storm in MY34, as compared to a value of S = 6.3% in MY33 (Fig. 1). While the findings are qualitatively similar to those of Montabone et al. (2020), our results show that super-rotation increased by a factor of two during the GDS, rather than by a factor of three as found by Montabone et al. (2020). This quantitative difference between the works can be partly attributed to the impact of temperature assimilation in our model, which provides a robust observational constraint on model dynamics.

The morphology of tropical winds in our model during the GDS (Fig. 2) is consis-255 tent with the dust-driven enhancement of thermal tides, as explained by Lewis and Read 256 (2003). As westward-propagating tides are excited by dust in the lower atmosphere they 257 propagate vertically and induce westerly super-rotating winds at their source regions (Fels 258 & Lindzen, 1974). Furthermore, the tides eventually break in the upper atmosphere and 259 deposit easterly momentum into the background winds, resulting in strengthened east-260 erly flow at higher levels. Such tidally-driven changes in upper atmospheric winds will 261 have a significant impact on the distribution of trace species such as water that are trans-262 ported into the upper atmosphere during global dust storms (Fedorova et al., 2018, 2020; 263 Aoki et al., 2019). By triggering the breaking levels of gravity waves, these winds also 264 play a role in controlling the amount and spectral distribution of gravity waves enter-265 ing the thermosphere, which has implications for water loss as enhanced gravity wave 266 activity increases the hydrogen escape flux (Yiğit et al., 2021). 267

Enhanced tropical winds played a significant role in the horizontal transport and 268 distribution of dust during the initial phase of the GDS (Bertrand et al., 2020; Gillespie 269 et al., 2020). We have shown that the martian atmosphere in MY34 was already in a state 270 of enhanced super-rotation prior to the onset of the GDS at  $L_{\rm S}=187^{\circ}$  (Figs. 1, 2(d), 3, 271 4(f). One consequence of strengthened pre-GDS tropical winds is that eastward dust 272 advection during the initial stages of the GDS would have been more rapid in MY34 than 273 in years with weaker pre-storm super-rotation. This rapid transport would have enabled 274 the timely activation of the secondary Tharsis lifting center, which was crucial in mak-275 ing the dust storm global in extent (Bertrand et al., 2020; Montabone et al., 2020). Fur-276 thermore, we observe that super-rotating winds create regions of enhanced vertical wind 277 shear in the lower atmosphere (Fig. 3(b)). As strongly sheared winds are subject to Kelvin-278

Helmholtz instability, we speculate that jet-induced wind shear could have enhanced the
generation of turbulence and surface wind stresses prior to the storm. This would lead
to enhanced dust lifting, especially in the daytime when atmospheric static stability is
low. However, further work will be required to quantify the impact of such shear-induced
turbulence on vertical dust transport.

Our analysis of the pre-storm circulation indicates that advection of dust from the 284 southern mid-latitudes into tropical regions played a key role in inducing the circulation 285 pattern of a symmetric Hadley cell with a vertically-aligned tropical upwelling branch 286 (Fig. 4). As a result of this pattern, vertical transport was much more efficient in MY34 287 in the lead-up to the storm, as compared to MY33. Such a pre-storm circulation pat-288 tern could be an important pre-requisite condition that enabled the rapid lifting of dust 289 during GDS initiation (Haberle et al., 1993; Wilson et al., 2008; Shirley et al., 2020). There-290 fore, even though the MY34 GDS is considered to have initiated in the northern hemi-291 sphere (Sánchez-Lavega et al., 2019; Bertrand et al., 2020), our results indicate that dust 292 transport from the southern hemisphere may still have played a crucial role in GDS ini-293 tiation.

In conclusion, we have conducted a data assimilation study of super-rotation in MY33– 295 34. We found that super-rotation increased by a factor of two during the MY34 GDS, 296 with substantial changes to tropical wind profiles at all levels. The atmosphere was found 297 to have entered a state of super-rotation more than forty sols prior to GDS onset, as a 298 result of tropical heating induced by dust transported into the tropics from southern mid-299 latitudes. The resulting uniform redistribution of dust across tropical latitudes resulted 300 in a symmetrical Hadley cell circulation with a tropical upwelling branch that was closely 301 aligned to the vertical, thereby enhancing the efficiency of vertical transport of dust in 302 MY34 and providing conditions conducive to the rapid global-scale expansion of the dust 303 storm. 304

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- Assimilation data from this study may be accessed at https://figshare.com/s/aa4ee8eab7510d0081dc.
- ACS raw data products are available from the ESA Planetary Science Archive (https://
- archives.esac.esa.int/psa/). MCS data is available from the NASA Planetary Data
- 315 System (https://pds-atmospheres.nmsu.edu/).

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