

Critically-stressed reservoir stimulation direction via stress preconditioning in horizontal EGS doublets

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Abstract

In this work, it is suggested numerically that it is possible to direct shear stimulation treatments in critically-stressed reservoirs. This would aid in the creation of Enhanced Geothermal Systems by promoting hydraulic connectivity in doublet-well systems. In this case, the stimulation treatment is directed using only the poroelastic stress changes associated with a previous stimulation treatment to precondition the stress field. This methodology is shown for reverse, strike-slip, and normal faulting stress regimes.

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12 stimulation treatment to precondition the stress field. This methodology is
13 shown for reverse, strike-slip, and normal faulting stress regimes.

14 **Keywords** EGS · Reservoir Stimulation · Reservoir Engineering · Poroelastic
15 stress · Hydraulic Shearing

16 **Highlights**

- 17 1. It is suggested that shear stimulation treatments in EGS reservoirs can be
18 directed
19 2. Injection-induced poroelastic stress changes are significant in a critically-
20 stressed crust
21 3. The methodology is shown for reverse, strike-slip, and normal faulting
22 stress regimes

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1 Introduction

Low permeability and inter-well connectivity are common problems preventing Enhanced Geothermal Systems (EGS) from reaching their potential (Tester et al., 2006; Ziagos et al., 2013). Indeed, it has previously been pointed out that the optimal distribution of permeable pathways is critical for the successful development of sufficient productivity for commercial EGS power generation (Robinson et al., 1971; Ziagos et al., 2013). For this reason, the ability to guide reservoir stimulation treatments such that specific areas of the reservoir can be targeted for stimulation would represent a significant development. A further advantage of this kind of stimulation targeting would be the ability to avoid the reactivation of large faults; either directly in the case that fault locations are known, or indirectly, in that only the most crucial parts of a reservoir are stimulated, thus moderating the risk that the stimulation treatment encounters large faults (e.g., Kim et al. (2018)). This should aid in the mitigation of induced seismicity. In fact, it is thought that the development of alternate stimulation concepts is integral to the mitigation of seismic risk from hydraulic stimulation (Håring et al., 2008) and that the engineering of reservoir connectivity would represent a key development for EGS (Rybach, 2010).

Directed reservoir stimulation techniques have been investigated before. For example, in Soultz-sous-Forêts, Baria et al. (2004) showed the positive effect of the contemporaneous stimulation of two wells in the context of an EGS project in crystalline rock. Their focus was primarily on the effect that an elevated pore pressure would have on the stimulation of a second well. However, the idea of altering the stress field in order to benefit a stimulation treatment has also been suggested as long ago as 1977 when Shuck (1977) filed a patent which involved injecting fluid to alter the plane of the maximum principal stress for use in hydraulic fracturing. Boutéca et al. (1983) investigated, both numerically and experimentally, the possibility of using fluid injection to alter the stress state such that a hydraulic fracturing treatment would connect two wells. This idea has been expanded upon by, for example, Warpinski and Branagan (1989), who were able to show stress changes of over 2 MPa due to the opening of a hydraulic fracture in lenticular reservoirs with the intent of reorienting potential hydraulic fractures such that they would intersect natural ones. Warpinski and Branagan (1989) estimated that larger pre-stimulation treatments would be able to induce stress changes of over 4 MPa, which, in this case, was a stress change large enough to swap the directions of the principal horizontal stresses. Warpinski and Branagan (1989) primarily considered their results relevant for single-well systems. Certainly, the effect of stress shadowing due to fracture opening has been widely discussed (e.g., Fisher et al. (2004); Vermylen and Zoback (2011)). Other relevant works include the effects of fluid-production-induced poroelastic stress changes on refracturing (Elbel and Mack, 1993), the work by Minner et al. (2002), which showed that injection and production can result in poroelastic stress changes that can dramatically alter fracture geometry on infill wells,

68 and Berchenko and Detournay (1997); Gao et al. (2019) who used models to
69 analyze the deviation of hydraulic fractures associated with poroelastic stress
70 changes resulting from production and injection.

71 Although there have been a number examples of EGS in sedimentary rocks
72 (e.g., Evans et al. (2012)), the focus here will be on EGS in crystalline rocks,
73 which tend to be deeper and therefore typically offer higher temperatures.
74 Various configurations exist for EGS wells (e.g., Chen and Jiang (2015)), but
75 a typical EGS setup might employ a doublet well configuration (e.g., Jupe
76 et al. (1992); Dorbath et al. (2009); Kim et al. (2018)) whereby fluid is circu-
77 lated between an injection and a production well, where these wells can either
78 be vertical or directional in nature. Crystalline rock and high temperatures
79 do pose new challenges for directional drilling, but improvements are being
80 made. Certainly, a number of EGS wells have been drilled directionally (e.g.,
81 Tester et al. (2006); Kwiatek et al. (2008); Dorbath et al. (2009); Kwiatek et al.
82 (2014); Kim et al. (2018); Norbeck et al. (2018); Kwiatek et al. (2019)) and hor-
83 izontally drilling in hard, high temperature rock is possible (albeit potentially
84 cost inhibitive) (Shiozawa and McClure, 2014). In fact, recent publications
85 are beginning to consider the multi-stage stimulation of horizontal wells for
86 EGS (e.g., Meier et al. (2015); Kumar and Ghassemi (2019)). It has even been
87 suggested that the multi-stage horizontal well stimulation employed in the oil
88 and gas industry should act as a model for the EGS industry (Ziagos et al.,
89 2013; U.S. Department of Energy, 2019).

90 Typically, for EGS in crystalline rock, the reservoirs are primarily thought
91 to be stimulated in shear (Evans et al., 2005b; Zang et al., 2014). Coulomb
92 faulting theory is a typical way to assess shear failure potential. From Coulomb
93 faulting theory, it is clear that an increase in pore pressure reduces the effective
94 stress on a shear plane and brings the shear plane closer to failure. Indeed,
95 in many instances of shear stimulation in crystalline rock, it is thought that
96 the increase in pore pressure was the dominant contributor to the induced
97 shear displacement (Pearson, 1981; Pine and Batchelor, 1984; Jupe et al.,
98 1992; Deichmann and Giardini, 2009). From Coulomb faulting theory it is
99 clear that it is possible to stimulate EGS reservoirs with injection pressures
100 below the minimum principal stress. This is a fundamental difference between
101 EGS stimulation and hydraulic fracturing operations, as hydraulic fracturing
102 operations occur at injection pressures above the minimum principal stress
103 in order open fractures in a tensile manner. However, changes in total stress
104 can also cause shear failure. For example, poroelastic stress changes, or the
105 stress changes resulting from pore pressure-induced deformation of reservoir
106 rock, have been shown to be significant in induced seismicity, where they have
107 at times been largely responsible for fluid production (e.g., Segall (1989)),
108 injection (e.g., Chen et al. (2017)), and hydraulic fracturing (e.g., Deng et al.
109 (2016)) operation-induced seismicity. Poroelastic stressing differs from changes
110 in pore pressure in that it does not necessarily lead to isotropic changes in
111 effective stress. A simple increase or decrease in pore pressure will not directly
112 lead to a change in the differential stress; however, importantly, the resulting
113 poroleastic changes to total stress can be, and frequently are, anisotropic.

114 This induced anisotropy allows poroelastic stress changes to have a significant
115 influence on a shear plane’s potential for failure, even when small in magnitude,
116 as these changes are capable of either increasing or decreasing differential
117 stress.

118 In this work, the stimulation of an EGS doublet well system will be in-
119 vestigated. Specifically, an investigation will be made into the possibility of
120 guiding the stimulation from one well to another, as previously discussed by
121 Baria et al. (2004). Unlike in Baria et al. (2004), however, this work will
122 consider poroelastic stress changes, which have been shown to be relevant in
123 EGS stimulations (Jacquey et al., 2018), as well as address the three main
124 stress regimes in generic scenarios. This investigation will be carried out by
125 first stimulating one of the doublet wells according to normal stimulation pro-
126 cedure. The stress changes associated with this first stimulation treatment
127 will then encourage stimulation in a certain direction, allowing the stimula-
128 tion treatment of the second well to be guided toward the stimulated region
129 surrounding the first well. In this way the stress field is ”preconditioned” be-
130 fore the stimulation of the second well. The advantage of this methodology is
131 that it (1) helps ensure connectivity between the two doublet wells and (2)
132 reduces the stimulation of less useful rock mass, which decreases the chance
133 of accidentally inducing a large magnitude event on a nearby fault. This in-
134 vestigation will be performed with a poroelastic reservoir simulator where
135 the permeability enhancement is based on the results of field studies. Even
136 if, as mentioned above, further technological advancement may be necessary
137 to allow horizontal EGS wells to be readily and cost-effectively drilled, here
138 the investigation will concern the stimulation of horizontal EGS doublet wells
139 drilled in critically-stressed crystalline rock. This investigation will also have
140 implications for directionally-drilled wells; however, in these cases the results
141 would depend on the inclination of the wells. Although significant temperature
142 differences may typically be present between the injected fluid and reservoir
143 during EGS stimulations, the analysis here will be isothermal to isolate the
144 effects of poroelasticity.

145 **2 Methodology**

146 In order to model the pressure and stress changes resulting from either fluid
147 production or injection, a sequentially coupled 2-D plane strain poroelastic
148 reservoir simulator is employed. Although the model is 2-D plane strain, it will
149 be appropriate for modelling 3-D stress changes due to fluid production and
150 injection activities from horizontal wells which are parallel to a principal stress
151 direction (Cheng, 2016). An equivalent continuum approach will be employed,
152 meaning that fractures will not be explicitly modelled, a previously explored
153 approach for modelling fractured media (e.g., Oda (1986); Miller (2015); Gan
154 and Elsworth (2016)). This approach was taken because the fractured rock
155 mass bulk behaviour is the focus and scale of the paper.

2.1 Flow Model

The combination of the conservation of mass of a single phase and Darcy's Law,

$$\frac{\partial(\phi\rho)}{\partial t} - \nabla \cdot \left(\frac{k}{\mu} \rho (\nabla P - \nabla(\rho g z)) \right) = q, \quad (1)$$

is used as the foundation of the flow model. Here ϕ is the porosity, ρ the fluid density, k the permeability, μ the fluid's dynamic viscosity, P the pore pressure, g the acceleration due to gravity, z the depth, and q the mass source terms. A fully implicit finite difference in time, finite volume in space framework (Aziz and Settari, 2002) is used to discretize the equation, which is then solved for the primary variable of pressure.

2.2 Mechanical Model

The mechanical model is based on the conservation of momentum,

$$\nabla \cdot \sigma' + \nabla(\alpha P) = -f, \quad (2)$$

where σ' is the effective stress, α the Biot coefficient, and f represents the body forces. The sign convention is such that tension and extension are negative. This equation is then combined with the linear theory of poroelasticity (Biot, 1941; Rice and Cleary, 1976; Wang, 2000),

$$S_{ij} - \alpha P \delta_{ij} = \frac{E}{(1+\nu)} \epsilon_{ij} + \frac{E\nu}{(1+\nu)(1-2\nu)} \epsilon_{kk} \delta_{ij}, \quad (3)$$

in a finite element framework such that the stresses and strains associated with fluid production and injection can be solved for. Here, the total stress is represented by S , the Kronecker delta by δ_{ij} , the drained Young's Modulus by E , the drained Poisson's ratio by ν , and the strain by ϵ .

3 Problem Setup

The horizontal wells in this investigation will penetrate granitic basement rock, all at 4500 m depth, a similar depth to the EGS program of Soultz, France (Dorbath et al., 2009). The granitic basement rock is assumed to extend up to 2500 m depth, not unlike Basel EGS, Switzerland (Ladner and Häring, 2009). The overburden, however, will not be modelled and will be instead replaced with a constant applied stress based on a reasonable lithostatic pressure gradient. The model boundaries are chosen such that the wells are far enough away to limit their effect on the simulations. As shown in Figure 1, the entire set-up will be modelled in 2-D plane strain, an appropriate approach to model horizontal wells (Cheng, 2016). The investigation of the effects of preconditioning

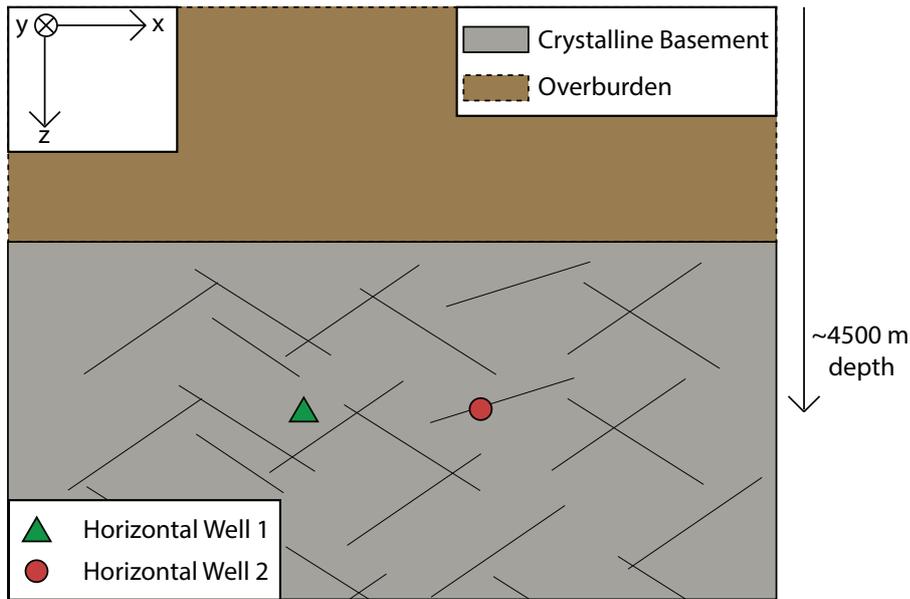


Fig. 1 Schematic of the problem setup for reverse and strike-slip faulting stress regimes. This represents a side view of two horizontal wells. The plane is normal to the orientation of the wells. The overburden is not modelled. Not to scale.

186 will be investigated for reverse, strike-slip, and normal faulting stress regimes.
 187 In each case, all three wells will be drilled parallel to S_{hmin} . Note that there
 188 are conflicting results regarding the orientation of reservoir creation during the
 189 stimulation of crystalline rock, with some operations indicating nearly parallel
 190 to S_{Hmax} (Häring et al., 2008; Kwiatek et al., 2019) and others indicating an
 191 offset such that reservoir creation occurs in the direction of strike of optimally
 192 oriented shear planes (Evans et al., 2005b; Kim et al., 2017, 2018) with it at
 193 times being difficult to determine exactly what happened in each case. It is
 194 further likely that the created reservoir geometry depends on the pre-existing
 195 discontinuities (Häring et al., 2008). For this reason and given observational
 196 inconsistencies, it is difficult to determine the optimum orientation of the wells
 197 with respect to the stress field.

198 The initial pore pressure and vertical stress are calculated using typical hy-
 199 drostatic and lithostatic gradients, respectively. The assumption that the crust
 200 is critically stressed (Brudy et al., 1997; Townend and Zoback, 2000; Zoback
 201 and Townend, 2001; Zoback et al., 2002) is then used alongside the notion
 202 that the frictional strength of pre-existing faults is what limits the differential
 203 stress in the crust (Zoback and Healy, 1992; Brudy et al., 1997; Zoback and
 204 Harjes, 1997). This allows for the direct calculation of the minimum principal
 205 stress in the normal faulting stress regime and the maximum principal stress
 206 in the reverse faulting stress regime. For example, in a normal faulting stress
 207 regime, the minimum possible horizontal stress that could be present on a sup-
 208 posed optimally-oriented fault can be calculated using the vertical stress, the

209 pore pressure, and the assumed coefficient of friction. This process is repeated
 210 at all depths to calculate the initial minimum principal stress everywhere in
 211 the model. In the strike-slip faulting stress regime, the maximum principal
 212 stress is calculated using this methodology after the minimum principal stress
 213 is assumed to be 0.8 times the vertical stress. The coefficient of friction will
 214 be assumed to 0.6, although there have been indications that the coefficient of
 215 friction in granitic rock may be higher (e.g., Blanpied et al. (1995)). It will be
 216 assumed that the frictional coefficient remains constant during stimulation, in
 217 agreement with laboratory studies (e.g., Ishibashi et al. (2018)).

218 A reasonable value of the intact Young's Modulus for granite is 36 GPa
 219 (Villeneuve et al., 2018). However, the rock is assumed to be fractured, mean-
 220 ing that, depending on the density of fractures, it may not be possible to use
 221 the intact Young's Modulus to describe the bulk behaviour (Villeneuve et al.,
 222 2018). Using a moderate fracture density and geological strength index, the
 223 rock mass Young's Modulus was taken as 50 % of the intact Young's Modulus
 224 based on findings by Villeneuve et al. (2018). This results in an equivalent
 225 Young's Modulus of 18 GPa. The Biot coefficient of this fractured granite is
 226 taken as 0.76, similar to that found by (Evans et al., 2003) for a fractured gran-
 227 ite. The Poisson's ratio of the granite rock will be taken as 0.15, a relatively
 228 low value for granite due to its fractured nature (Walsh, 1965). A summary of
 229 the parameters used can be found in Table 1.

Table 1 *Model parameters*

Variable	Value	Unit
Fluid reference density (STP), ρ_f	1000	$\frac{kg}{m^3}$
Fluid compressibility, c_f	$5e - 10$	$\frac{1}{Pa}$
Fluid dynamic viscosity, μ	0.001	$Pa \cdot sec$
Granite drained Young Modulus, E_g	18e9	Pa
Granite drained Poisson's Ratio, ν_g	0.15	—
Granite initial bulk porosity, ϕ_g	0.02	—
Granite Biot coefficient, $\alpha_{s,g}$	0.76	—
Coefficient of friction, μ_f	0.6	—

230 In order to avoid the compounding effects of thermal strains and to more
 231 clearly illustrate the effects of the stress preconditioning, the production and
 232 stimulation phases will be assumed to be isothermal (i.e., the reservoir will
 233 be assumed to be stimulated with water at reservoir temperature). This is
 234 obviously not a realistic scenario for a typical geothermal stimulation, and the
 235 probable effects of the thermal strains will be discussed in a later section.

236 3.1 Initial Bulk Permeability

237 Granite fractures can be assumed to have a high permeability (on the order
 238 of $10^{-12} m^2$ (Ishibashi et al., 2018)) compared to granitic matrix, which gen-

erally has a permeability on the order of 10^{-21} to $10^{-20} m^2$ (Morrow et al., 1986). For this reason, the matrix permeability will be assumed to be negligible compared to the fracture permeability, meaning that flow will be principally in the fractures, equivalent to the level B distinction suggested by Cornet (2016), where flow is dominated by flow through reactivated fractures. In highly fractured and faulted crystalline rocks, the permeability of critically-stressed faults is much higher than that of faults which are poorly oriented for failure in the modern-day stress field (Barton et al., 1995). Evans et al. (2012) found in a study of European case histories, that all crystalline rock masses investigated were critically stressed. Therefore, the optimally-oriented faults and fractures in the granite investigated here will be assumed to be initially at least somewhat permeable, even if they need further shear stimulation to produce or inject fluid at rates sufficient for their given operational goal. This is supported by, for example, the pre-stimulation tests performed in granite in the Soultz HDR site and the Basel 1 enhanced geothermal system which yielded effective permeabilities of $3 \cdot 10^{-16} m^2$ and $10^{-17} m^2$ respectively (Evans et al., 2005b; Häring et al., 2008; Ladner and Häring, 2009). These tests also agree with the findings of Zoback and Townend (2001), who found that bulk permeability in the upper crust is high ($10^{-17} m^2$ to $10^{-16} m^2$) due to critically-stressed faults. For this reason a starting value of $10^{-17} m^2$ is used for bulk permeability, a value on the low end of bulk permeabilities seen in the field as mentioned above. The actual initial value of the permeability seen in the simulation will be lower than this value due its dependence on pressure and stress addressed in Section 3.2.

3.2 Shear Stimulation

Although stimulation in Enhanced Geothermal Systems may well be mixed mode between the creation of new fractures and the shearing of old fractures and faults (McClure and Horne, 2014; Norbeck et al., 2018) (especially with injection pressures above the minimum principal stress), it is thought that shear failure is the dominant and most promising mechanism of reservoir creation in hard rock formations in EGS stimulation (Evans et al., 2005b; Ziagos et al., 2013; Zang et al., 2014). Indeed, it has been previously shown in laboratory (e.g., Chen et al. (2000); Ishibashi et al. (2018)) and field (e.g., Jupe et al. (1992); Evans et al. (2005b); Ladner and Häring (2009); Guglielmi et al. (2015)) studies of granitic rock that fracture permeability increases with shear displacement. In this study specifically, it will be assumed that the fractures and faults optimally oriented for shear in the prevailing stress field will be the planes upon which shear failure occurs, as seen, for example at Soultz (Evans et al., 2005b).

Shear stimulation of granitic reservoirs results in a permeability increase that can vary depending on the site, even varying within the same well (Evans et al., 2005a). For example, permeability was increased by three orders of magnitude at the Fjällbacka Hot Dry Rocks Project, Sweden following stimulation

(Jupe et al., 1992), but Soultz, France only saw an increase in transmissivity of a factor of fifteen when the effect of the stimulation is evaluated over the entire wellbore (Evans et al., 2005a). Here, stimulation will be assumed to ultimately result in a permeability increase of a factor of 200, similar to the results of stimulation at Basel (Ladner and Häring, 2009) and the 1993 stimulation of a 550 m section of hole at Soultz (Evans et al., 2005a).

The permeability used in the numerical model will be based on the notion of a changing aperture width with effective normal stress and a stepwise change in permeability occurring after a failure condition is reached (Miller and Nur, 2000; Miller, 2015). As in Miller (2015), permeability is assumed to take the form

$$k = k_0 e^{\frac{-\sigma_n}{\sigma^*}}, \quad (4)$$

where k_0 is the initial permeability defined in Section 3.1, σ_n is the effective normal stress acting on the assumed shear plane, and σ^* is a normalizing constant taken as 100 MPa. The normalizing constant is picked as a large value such that the initial individual values of permeability in the reservoir are not significantly lower than the overall values of bulk permeability seen in the EGS reservoirs in Section 3.1. Again following the model used by Miller (2015), the failure planes (with one assumed orientation for each reservoir block) will follow an unbounded normal distribution with a mean orientation corresponding to the optimal orientation in the given stress regime and a standard deviation of 0.02 radians. As the standard deviation is small, this is, in essence, equivalent to using a von Mises distribution with a large concentration coefficient. Note that this model implies, based on the assumed critically-stressed nature of the reservoir, that minuscule changes in stress or pore pressure could result in shear failure if a given cell is optimally oriented. In fact, however, it will be assumed that all cells require a Coulomb stress increase of 0.1 MPa before failure in addition to any stress increase required due to a non-optimal orientation. A 0.1 MPa Coulomb stress increase is a reasonable value for the initiation of slip (Stein, 1999). Coulomb stress, τ , is defined as

$$\tau = \tau_s - \mu_f (S_n - P), \quad (5)$$

where τ_s and S_n are the shear stress and normal stress on a potential shear plane (for calculations of Coulomb stress this plane is assumed to be optimally oriented in the prevailing stress regime) and μ_f is the static coefficient of friction. Generally, the Coulomb stress will increase when the maximum principal total stress increases, the minimum principal total stress decreases, or the pore pressure increases.

If the Coulomb failure criteria for a given cell is reached, a stepwise change in permeability will occur (Miller and Nur, 2000; Miller, 2015) such that k_0 in Equation 4 will be replaced by k'_0 ; where k'_0 is defined as

$$k'_0 = xk_0. \quad (6)$$

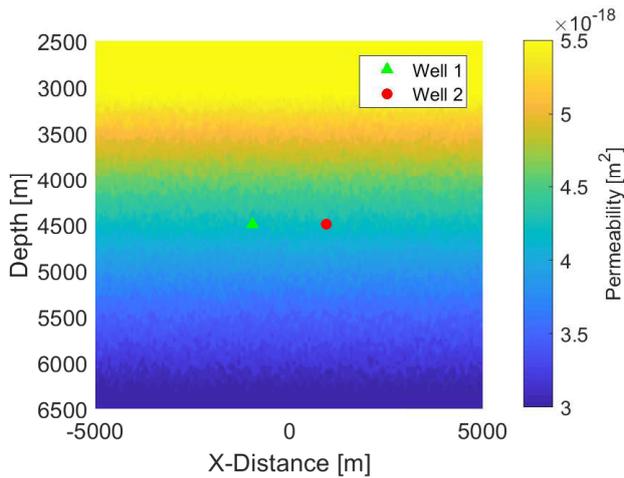


Fig. 2 The initial permeability field used in the reverse faulting case. The heterogeneity is due to the randomness associated with the permeability model.

320 Here, x is a multiplication factor taken to be equal to 200 based on the
 321 reviewed stimulation treatments mentioned above. This methodology for mod-
 322 elling permeability enhancement due to shear stimulation implies that perme-
 323 ability enhancement largely remains after high pressure is stopped. This is
 324 representative, for example, of the shear stimulation at the Soultz HDR site
 325 (Evans et al., 2005b). Note that porosity is kept constant throughout the sim-
 326 ulation, reflecting, for example, the methodologies of Miller and Nur (2000)
 327 and Baisch et al. (2010). This means that the coupling between the mechanical
 328 model and the flow model is entirely contained in the change of permeability.
 329 An example of the permeability field, Figure 2, is shown for the reverse fault-
 330 ing case. Although the permeability fields of each run will vary slightly due
 331 to the randomness associated with the permeability model, this variation is
 332 limited and the general trends predicted by the results are repeatable.

333 4 Results

334 This section will be subdivided into three subsections, one subsection for each
 335 stress regime, Table 2. Beginning with a reverse faulting stress regime, two
 336 wells will be stimulated with the goal of connecting the stimulated regions of
 337 each well to create a doublet system. In the reverse faulting case, the compar-
 338 ison will be made between the case where the first well is flowback after its
 339 stimulation and the case where this first well is not flowback after stimula-
 340 tion and instead the second well is stimulated immediately. For the remaining
 341 stress regimes, however, the flowback case will not be presented and instead
 342 the effect of the first stimulation treatment on the second will be shown by
 343 comparing the average propagation lengths of the stimulated region outside of
 344 the two wells and inside the two wells.

Table 2 *Principal stress orientations.* The wells are drilled in the y-direction; however, the orientations of the principal stresses change depending on the stress regime. Note that S_x is S_{Hmax} and S_y is S_{hmin} in each case.

Regime	S_1	S_2	S_3
Reverse Faulting	S_x	S_y	S_z
Strike-Slip Faulting	S_x	S_z	S_y
Normal Faulting	S_z	S_x	S_y

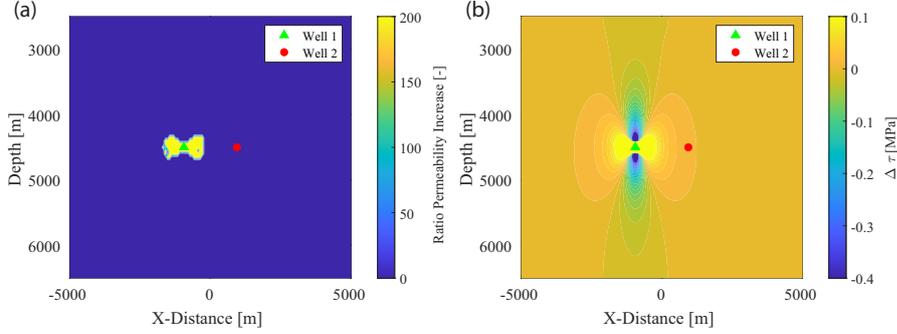


Fig. 3 The result of the stimulation of well 1 in the reverse faulting stress regime case. (a) The permeability enhancement associated with the stimulation treatment (t=3 days). (b) The Coulomb stress changes resulting from the stimulation treatment (t=3 days).

345 4.1 Reverse Faulting

346 In a reverse faulting stress regime, the maximum principal stress is horizontal and the minimum principal stress is vertical. Therefore, increases in the
 347 total horizontal stress (specifically the maximum horizontal stress, S_{Hmax})
 348 and decreases in the total vertical stress will generally result in an increase in
 349 Coulomb stress.
 350

351 In this case, the two wells will be located at a depth of 4500m and separated
 352 by 1884m, with the midpoint between the two wells having an X-Distance
 353 coordinate of 0m, Figure 2. The stimulation treatment procedure is begun by
 354 first stimulating the left-most of the two wells with an injection rate of 0.014
 355 $\frac{kg}{msec}$, which corresponds to $7.0 \frac{kg}{sec}$ for a 500 m long well length section, over
 356 a period of three days. This stimulation treatment would be similar to, but
 357 slightly smaller than, the 2000 stimulation of GPK2 at Soultz-sous-Forêts, for
 358 example (Dorbath et al., 2009). The permeability increases and Coulomb stress
 359 changes associated with this stimulation treatment are shown in Figure 3.

360 Next, the first well is flown back with a rate of $2.33 \frac{kg}{sec}$ over a period of 9
 361 days, resulting in the entire fluid mass that was injected with the stimulation
 362 treatment being reproduced. Note that it is probably unlikely that the entire
 363 injected mass would be reproduced in reality; however, the purpose here is
 364 simply to illustrate the effect of flowback on the far-field poroelastic stresses.
 365 The permeability above and below the previously stimulated region increases
 366 slightly during this flow back period, Figure 4a. This is due to production in-

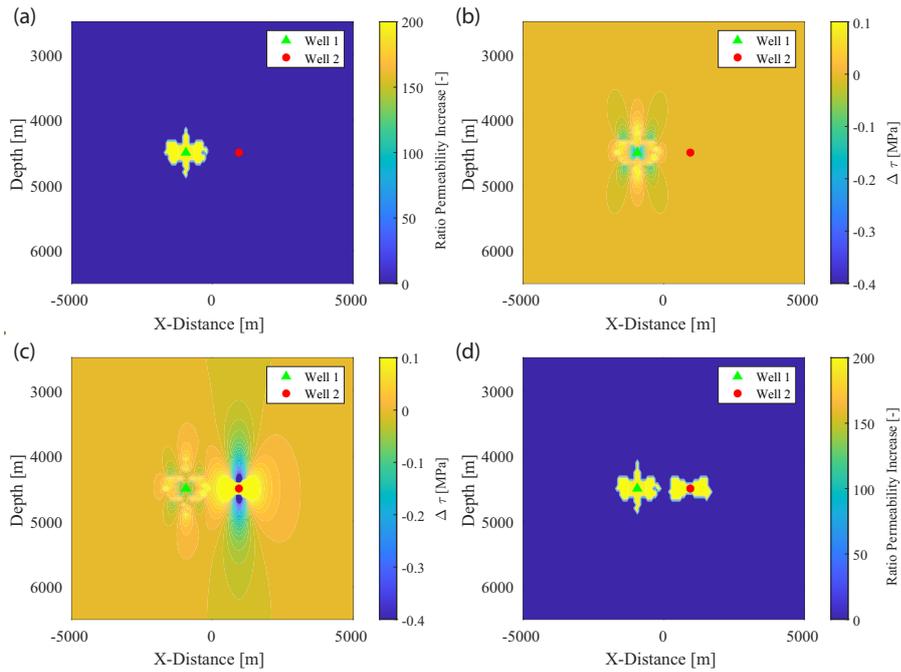


Fig. 4 For the reverse faulting stress regime case, the result of flowing back the first well before stimulating the second well. (a) The permeability enhancement at a time immediately after the flowback period ($t=12$ days), note the enhancement that has occurred above the initially stimulated region. (b) The Coulomb stress at a time immediately after the flowback period of the first well ($t=12$ days). The Coulomb stresses in-between the two wells has been reduced since the initial stimulation treatment, when compared to Figure 3b. (c) The Coulomb stresses after the stimulation of the second well ($t=15$ days). (d) The permeability enhancement at the end of the entire procedure ($t=15$ days). The two wells are not connected with a separation of the two stimulated zones of 362m.

367 ducing increased total horizontal stresses and decreased total vertical stresses
 368 in this region. The Coulomb stress changes associated with this flowback peri-
 369 od, Figure 4b, show the result of these stress changes with increases above
 370 and below the previously stimulated region. This type of increased Coulomb
 371 stress and shear failure occurring above production zones is analogous to the
 372 reverse faulting sometimes seen during hydrocarbon production (e.g., Segall
 373 (1989)). Figure 4b also indicates that the Coulomb stress in-between the two
 374 wells has decreased since flowback began, when compared to Figure 3b. The
 375 changes to pore pressure and the maximum and minimum principal stresses
 376 are shown in Figure 5.

377 At this stage, the second well is stimulated using the same stimulation
 378 treatment that was used in the first well. The Coulomb stress changes, Fig-
 379 ure 4c, and permeability field enhancements, Figure 4d, indicate that the two
 380 stimulated zones were not connected in this case, being still separated by 362m
 381 of unstimulated rock mass.

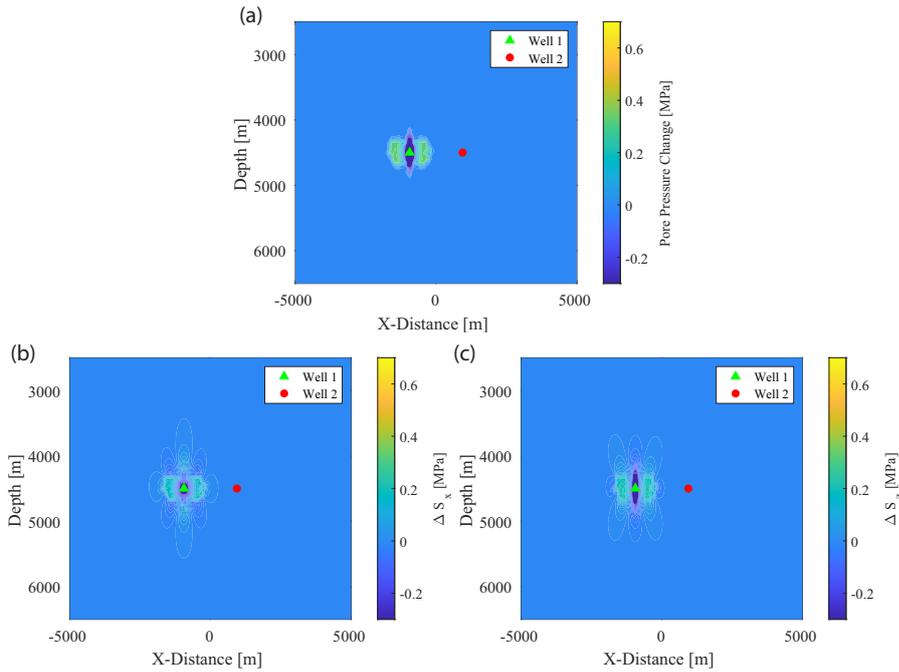


Fig. 5 For a reverse faulting stress regime, the (a) pore pressure changes, (b) maximum principal total stress changes (ΔS_x for reverse faulting), and (c) minimum principal total stress changes (ΔS_z for reverse faulting) after stimulation and flowback of the first well.

382 If, instead of flowing back the well, the well is simply shut-in and the second
 383 stimulation begun immediately after the termination of the first, the Coulomb
 384 stress changes associated with the first stimulation will largely remain during
 385 the second stimulation. As these Coulomb stress changes are encouraging
 386 failure and are larger closer to the first well, they may potentially cause the
 387 stimulation treatment of the second well to be directed towards the stimulated
 388 zone of the first well.

389 To test this procedure, the first well is stimulated as before with an injection
 390 rate of $7.0 \frac{kg}{sec}$ over three days. Following this, the second well is stimulated
 391 immediately after the first stimulation treatment with no flowback period.
 392 The stimulation treatment again consists of an injection rate of $7.0 \frac{kg}{sec}$
 393 over three days. In this way, the Coulomb stress changes associated with the first
 394 stimulation treatment remain and help to ensure connection between the two
 395 wells' stimulated regions.

396 At the midpoint between the two wells, the Coulomb stress just before the
 397 second stimulation had increased by 0.056 MPa, Figure 3b. However, at the
 398 location of equivalent distance from well 2 but in the opposite direction (a
 399 depth of 4500m and an X-distance of 1884m), the Coulomb stress just before
 400 the second stimulation has only increased by 0.0056 MPa. These differences in
 401 Coulomb stress change are what ultimately cause the stimulation of well 2 to

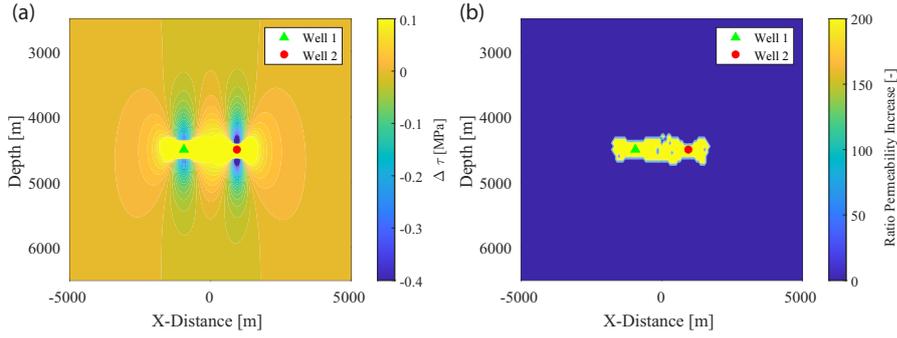


Fig. 6 For the reverse faulting stress regime case, the result of not flowing back the first well before beginning the stimulation treatment of the second well. (a) The Coulomb stresses after the stimulation of the second well ($t=6$ days). (b) The permeability enhancement at the end of the entire procedure ($t=6$ days). The stimulated zone of each well extends and average 761m away from the other doublet well and 942m towards it.

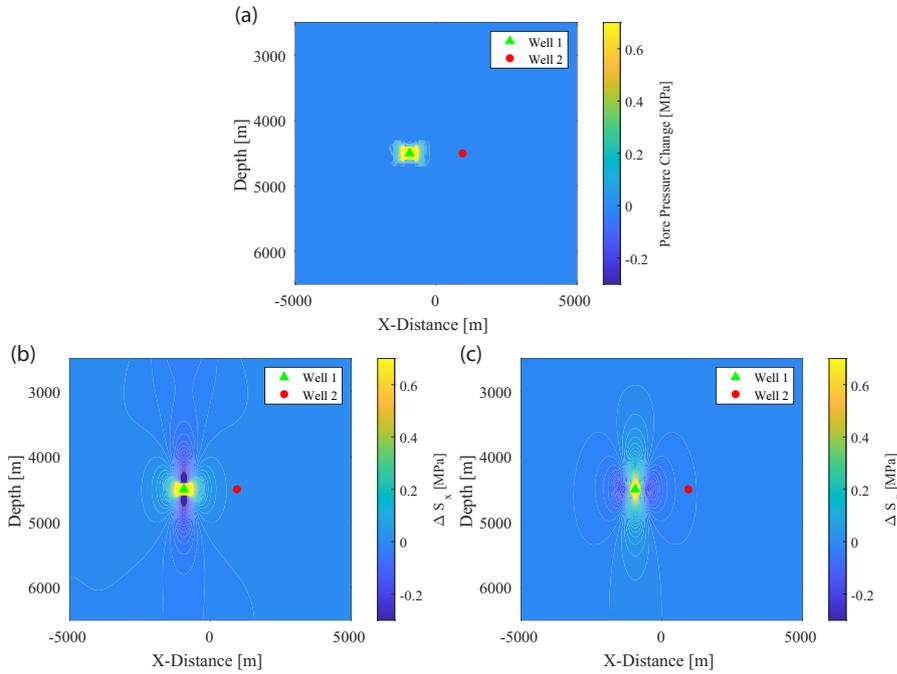


Fig. 7 For a reverse faulting stress regime, the (a) pore pressure changes, (b) maximum principal total stress changes (ΔS_x for reverse faulting), and (c) minimum principal total stress changes (ΔS_z for reverse faulting) associated with the stimulation of the first well without flowback. Note how, in the region between the two wells, the maximum principal total stress increases, the minimum principal total stress decreases, and the pore pressure remains unchanged. These changes explain the Coulomb stress changes seen in Figure 3b and indicate that the stress is being preconditioned due to total stress changes, not pore pressure changes. It is useful to compare this figure to Figure 5.

402 be directed towards the stimulated region of well 1 as opposed to propagating
403 equal distances in both directions. In fact, the stimulation treatments of both
404 wells, on average, propagate 761 m away from the other doublet well and 942
405 m towards it, Figure 6b, meaning that the stimulated zones extend over 20%
406 farther in-between the two wells than they do on the outside of the two wells.
407 Note that the average stimulation length of each well here is similar to, for
408 example, the seismicity cloud resulting from stimulation at Soultz-sous-Forêts,
409 which extended over 1000 meters horizontally and 500 meters vertically (Evans
410 et al., 2005b).

411 Unlike the results shown by Baria et al. (2004), the direction of the stim-
412 ulation treatment here is accomplished entirely by changes in stress, not pore
413 pressure. At an X-Distance of 0 (the center point between the two wells - 942m
414 from each well), the pore pressure change after the stimulation of the first well
415 is zero. The change in the S_{Hmax} , however, is 0.045 MPa, and results in over
416 half of the Coulomb stress change required for failure, Figure 7.

417 For the remaining two stress regimes, a flowback case will not be shown.
418 Instead, the average distances of propagation will be used to demonstrate the
419 degree to which the stimulation treatment was effectively directed.

420 4.2 Strike-Slip Faulting

421 In a strike-slip faulting stress regime, the maximum and minimum principal
422 stresses are both horizontal. The stress changes induced by injection through a
423 horizontal well will be anisotropic. For example, during the stimulation of the
424 first well, Figure 8a and b, the horizontal stress perpendicular to the first well
425 will experience greater compressive changes than the horizontal stress parallel
426 to it at large distances. Assuming the well is drilled parallel to the minimum
427 principal stress, this means that the maximum principal stress will increase
428 (becoming more compressive) more than the minimum principal stress, result-
429 ing in an increase in differential stress and Coulomb stress. These Coulomb
430 stress changes will be more pronounced near the stimulated region of the first
431 well, meaning that the stimulation treatment of the second well will be more
432 likely to propagate towards the first well than in the other direction. In this
433 case, the two wells will be located at a depth of 4500m and separated by 1450m,
434 with the midpoint between the two wells having an X-Distance coordinate of
435 0m.

436 The stimulation treatment procedure is begun by first stimulating the left-
437 most of the two wells with a stimulation rate of $0.0247 \frac{kg}{msec}$, which corresponds
438 to $12.37 \frac{kg}{sec}$ for a 500 m long well length section, over a period of three days,
439 Figure 8a. This stimulation treatment would be similar to, but slightly smaller
440 than, the 2000 stimulation of GPK2 at Soultz-sous-Forêts, for example (Dor-
441 bath et al., 2009). Next, injection into the first well is stopped and the second
442 well is stimulated with exactly the same stimulation treatment. The first well
443 does not undergo a flowback period before the stimulation of the second well.

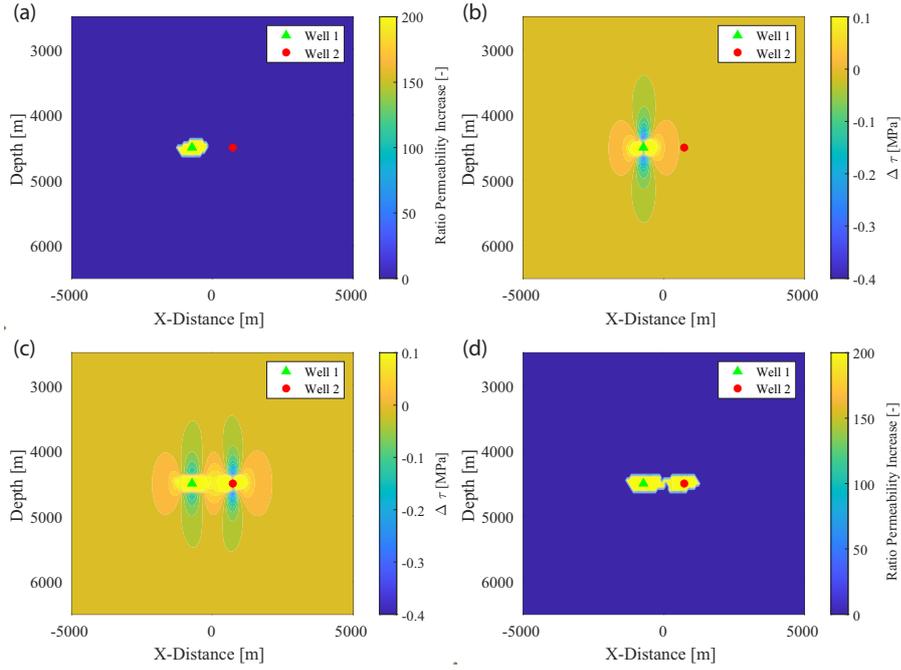


Fig. 8 The result of the stimulation treatment in a strike-slip faulting stress regime. (a) The permeability enhancement associated with the stimulation treatment of the first well ($t=3$ days). (b) The Coulomb stress changes resulting from the stimulation treatment of the first well ($t=3$ days). (c) The Coulomb stresses after the stimulation of the second well ($t=6$ days). (d) The permeability enhancement at the end of the entire procedure ($t=6$ days). The stimulated zone of each well extends an average 543m away from the other doublet and 725m towards it.

444 At the midpoint between the two wells, the Coulomb stress just before the
 445 second stimulation has increased by 0.042 MPa, Figure 8b. However, at the
 446 location of equivalent distance from well 2 but in the opposite direction (a
 447 depth of 4500m and an X-distance of 1450m), the Coulomb stress just before
 448 the second stimulation has only increased by 0.004 MPa. These differences in
 449 Coulomb stress change are what ultimately cause the stimulation of well 2 to
 450 be directed towards the stimulated region of well 1 as opposed to propagating
 451 equal distances in both directions. In fact, the stimulation treatments of both
 452 wells, on average, propagate 543 m away from the other doublet well and 725
 453 m towards it, Figure 8d, meaning that the stimulated zones extend over 30%
 454 farther in-between the two wells than they do on the outside of the two wells.

455 This change in Coulomb stress that guides the stimulation treatment of
 456 the second well towards the first well is caused by changes in total stress, not
 457 changes in pore pressure. At the midpoint of the two wells, the change in the
 458 maximum horizontal stress just before the second stimulation is 0.145 MPa
 459 whereas the change in the pore pressure is $4e-7$ MPa, Figure 9.

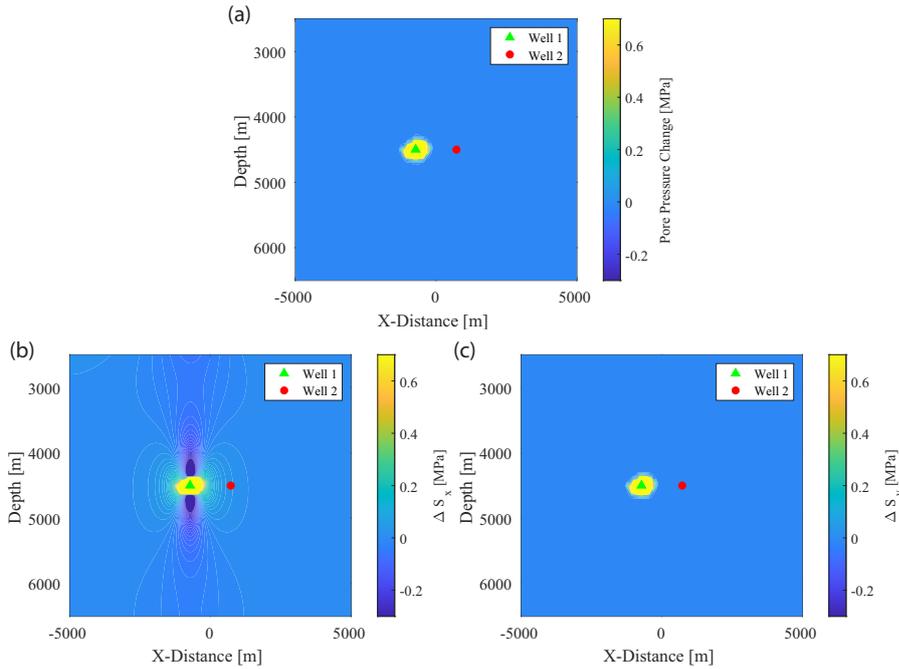


Fig. 9 For a strike-slip faulting stress regime, the (a) pore pressure changes, (b) maximum principal total stress changes (ΔS_x for strike-slip faulting), and (c) minimum principal total stress changes (ΔS_y for strike-slip faulting) associated with the stimulation of the first well without flowback. Note how, in the region between the two wells, the maximum principal total stress increases, whereas the minimum principal total stress and pore pressure remain unchanged. These changes explain the Coulomb stress changes seen in Figure 8b and indicate that the stress is being preconditioned due to total stress changes, not pore pressure changes.

460 4.3 Normal Faulting

461 In a normal faulting scenario, the vertical stress is the maximum principal
 462 stress. In the case that two doublet wells are drilled horizontally in a direc-
 463 tion parallel to the minimum principal stress, the injection-induced poroelastic
 464 stress changes caused by the stimulation of the first well will be expected to
 465 increase the total vertical stress primarily in locations above and below the
 466 stimulated well. This implies that the poroelastic stress changes will primar-
 467 ily encourage shear failure in locations which are vertically in-line with the
 468 well and not those horizontally in-line. For this reason, the wells are aligned
 469 vertically with a separation of 1000m at depths of 3500m and 4500m.

470 The stimulation treatment procedure is begun by first stimulating the shal-
 471 lower of the two wells with a stimulation rate of $0.0125 \frac{kg}{msec}$, which corresponds
 472 to $6.25 \frac{kg}{sec}$ for a 500 m long well length section, over a period of three days, Fig-
 473 ure 10a. This stimulation treatment would be similar to, but slightly smaller
 474 than, the 2000 stimulation of GPK2 at Soultz-sous-Forêts, for example (Dor-
 475 bath et al., 2009). Next, injection into the first well is stopped and the second

476 well is stimulated with exactly the same stimulation treatment. The first well
 477 does not undergo any flowback period before the stimulation of the second
 478 well.

479 At the midpoint between the two wells (a depth of 4000m and an X-distance
 480 of 0m), the Coulomb stress just before the second stimulation has increased by
 481 0.047 MPa, Figure 10b. However, at the location of equivalent distance from
 482 well 2 but in the opposite direction (a depth of 5000m and an X-distance of
 483 0m), the Coulomb stress just before the second stimulation has only increased
 484 by 0.002 MPa. These differences in Coulomb stress change are what ultimately
 485 cause the stimulation of well 2 to be directed towards the stimulated region of
 486 well 1 as opposed to propagating equal distances in both directions. In fact,
 487 the stimulation treatments of both wells, on average, propagate 400 m away
 488 from the other doublet well and 500 m towards it, Figure 10d, meaning that
 489 the stimulated zones extend 25% farther in-between the two wells than they
 490 do on the outside of the two wells.

491 This change in Coulomb stress that guides the stimulation treatment of
 492 the second well towards the first well is caused by changes in total stress, not
 493 changes in pore pressure. At the midpoint of the two wells, the change in the
 494 vertical stress just before the second stimulation is 0.159 MPa whereas the
 495 change in the pore pressure is 1.65e-5 MPa, Figure 11.

496 5 Discussion

497 5.1 Assumptions

498 5.1.1 Isothermal Simulations

499 The influence of temperature has not been considered in the analyses although
 500 temperature-induced stresses may play a significant role during EGS stimu-
 501 lation (e.g., Ghassemi and Tao (2016)). This was primarily done to simplify
 502 the analyses and more clearly illustrate the effects of stress preconditioning. In
 503 case-specific applications of this methodology, temperature effects should be
 504 considered. In fact, it may even be possible to design a stimulation procedure
 505 such that temperature-change induced stresses further precondition the stress
 506 field in a beneficial manner.

507 To evaluate the influence of the temperature-change induced stresses such
 508 that their neglect can be justified, the flow model was extended to include
 509 the conservation of energy,

$$\frac{\partial H_m}{\partial t} + \nabla \cdot \Gamma + \nabla \cdot f_T = q_T, \quad (7)$$

510 where H_m is the enthalpy of the entire medium, Γ is the heat conduction, f_T is
 511 the convection, and q_T represents the source terms. The equation is discretized
 512 and solved fully implicitly with the mass conservation equation, yielding both
 513 pressure and temperature. Equilibrium is assumed between the fluid and rock

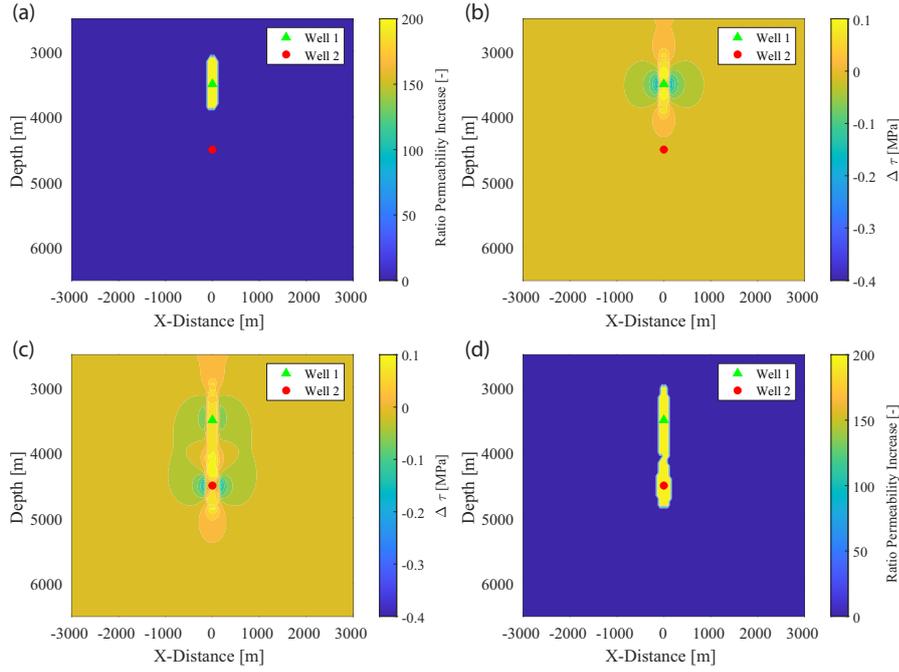


Fig. 10 The result of the stimulation treatment in a normal faulting stress regime. (a) The permeability enhancement associated with the stimulation treatment of the first well ($t=3$ days). (b) The Coulomb stress changes resulting from the stimulation treatment of the first well ($t=3$ days). (c) The Coulomb stresses after the stimulation of the second well ($t=6$ days). (d) The permeability enhancement at the end of the entire procedure ($t=6$ days). The stimulated zone of each well extends and average 400m away from the other doublet and 500m towards it.

514 temperature in each cell. To simplify the analysis, the fluid density and vis-
 515 cosity are assumed to remain constant with change in temperature. In the
 516 mechanical model, the thermal strain, ϵ_T ,

$$\epsilon_T = \alpha_T \Delta T, \quad (8)$$

517 is added to the mechanical strains before the computation of stress changes.
 518 Here, α_T is the coefficient of linear thermal expansion and T is the tempera-
 519 ture. A surface temperature of 30°C and a thermal gradient of $0.035 \frac{^\circ\text{C}}{\text{m}}$
 520 are assumed. The thermal conductivity of the water and granite are assumed
 521 to be 0.67 and $2.5 \frac{\text{W}}{\text{mK}}$ respectively. The heat capacity of the water and granite
 522 are assumed to be 4183 and $950 \frac{\text{J}}{\text{kgK}}$ respectively. The coefficient of linear
 523 expansion of granite is taken as $40 \cdot 10^{-6} \frac{1}{^\circ\text{C}}$, and the fluid enters the reservoir
 524 at a temperature of 47°C .

525 Using this updated model, the reverse faulting case was rerun up to the
 526 point just before the second stimulation. At the midpoint between the two
 527 wells, the difference in change in Coulomb stress in this case and in the case
 528 presented previously where temperature effects were not considered is 0.0002

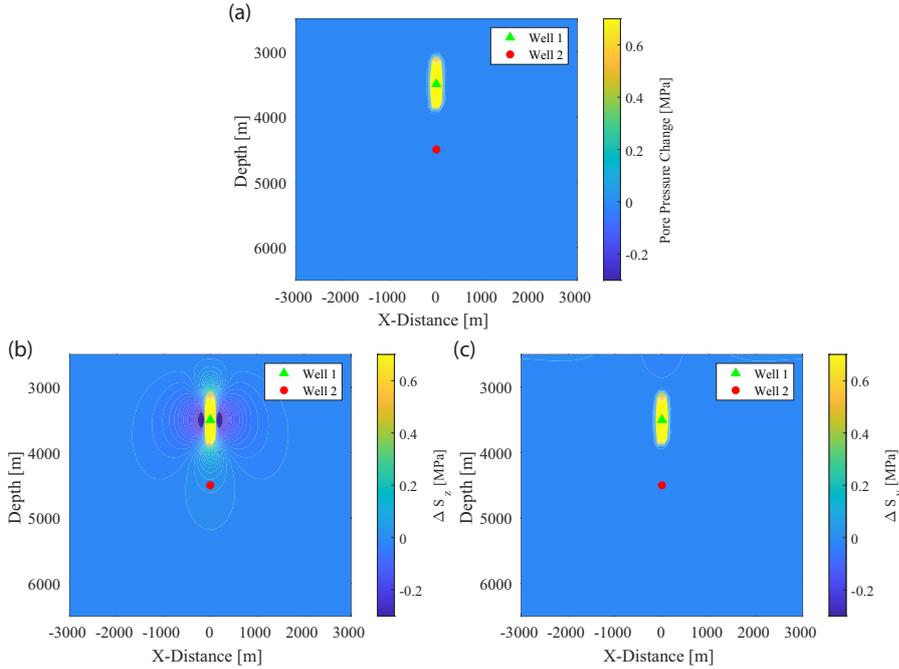


Fig. 11 For a normal faulting stress regime, the (a) pore pressure changes, (b) maximum principal total stress changes (ΔS_z for normal faulting), and (c) minimum principal total stress changes (ΔS_y for normal faulting) associated with the stimulation of the first well without flowback. Note how, in the region between the two wells, the maximum principal total stress increases, whereas the minimum principal total stress and pore pressure remain unchanged. These changes explain the Coulomb stress changes seen in Figure 10b and indicate that the stress is being preconditioned due to total stress changes, not pore pressure changes.

529 MPa. Considering that the Coulomb stress change due only to poroelastic
 530 effects was 0.056 MPa, this justifies not including stress changes due to tem-
 531 perature in the model. These Coulomb stress changes are small due to the
 532 small temperature changes of the system. The small temperature changes are
 533 due to the relatively small injection volumes (approximately $1620 m^3$ over 500
 534 m of wellbore over 3 days). In order to keep the model as simple as possi-
 535 ble and better illustrate the effects of poroelastic stress changes, the effect of
 536 temperature is therefore not included in the model.

537 5.1.2 Stress Criticality

538 In the simulations presented here, the reservoir was assumed to be critically
 539 stressed based on findings by Evans et al. (2012). In reality, however, knowledge
 540 of the in situ stress state is very important for reservoir stimulation activities,
 541 and the state of stress should ideally be investigated in each case before the
 542 start of operations. If the crust is less critically stressed than in the presented

543 cases, larger pore pressure changes will be needed to stimulate the reservoir
544 as shear failure would become more pore-pressure dominated. The poroelastic
545 stresses that guide the second stimulation treatment will make up less of the
546 required changes for failure. For instance, in this case it was assumed that
547 a Coulomb stress change of 0.1 MPa is required to induce shear failure. The
548 Coulomb stress change induced by the preconditioning at the mid-point be-
549 tween the wells was approximately 0.05 MPa, meaning that it made up half
550 of the required stress change. However, if the required Coulomb stress change
551 was instead 0.5 MPa, this preconditioning stress would only make up ten per-
552 cent of this value and it would presumably play a smaller role in directing the
553 second stimulation treatment. Conversely, this would mean larger pore pres-
554 sure changes would have been needed to stimulate the first well. This would
555 result in larger induced poroelastic stress changes.

556 5.1.3 Stress Redistribution

557 The model used here does not include stress redistribution associated with
558 shear failure occurring during stimulation. Previous studies (e.g., Catalli et al.
559 (2013)) have shown that stress redistribution associated with shear failure
560 during hydraulic stimulation can have a significant impact on future events.
561 Indeed, stimulation treatments of granitic rock have been shown to be ca-
562 pable of altering the stress field through aseismic slip occurring within the
563 stimulated zone (e.g., Cornet and Julien (1989); Schoenball et al. (2014)). Al-
564 though these stress changes have been shown to be large (on the order of
565 ten MegaPascals), they are thought to be largely confined to the stimulated
566 region (Schoenball et al., 2014). It is possible to come up with a far-field esti-
567 mate of this effect, if, for example, the Coulomb stress changes occurring near
568 the location of the second well can be calculated assuming the energy release
569 equivalent to a dynamic earthquake of M_w 3.0 occurring at the wellbore of the
570 first well. This amount of energy release due to aseismic slip is approximately
571 equal to that which occurred at the Le Mayet de Montagne granitic test site
572 (Cornet, 2016). An M_w 3.0 earthquake corresponds to a fault length of ap-
573 proximately 330 m according to typical earthquake scaling laws (Stein and
574 Wysession, 2003). King et al. (1994) found unclear correlations between after-
575 shocks and Coulomb stress changes after distances of about 3 fault lengths,
576 which is less than the separation between the two wells in each of the three
577 cases presented. Given that correlation between stress changes and aftershocks
578 was seen for positive Coulomb stress changes of the order of 0.01 MPa (King
579 et al., 1994), the stress changes associated with aseismic slip in the reservoir
580 at the location of the second well are most likely not significantly larger than
581 the poroelastic stress changes induced by the treatment itself (0.05 MPa at
582 the midpoint of the two doublet wells in the reverse faulting case where the
583 well separation is the largest). Therefore, although it would be unreasonable
584 to claim that stress changes associated with slip in the stimulated zone of the
585 first well are negligible for the stimulation of the second well, it can be con-
586 cluded that the poroelastic stress changes are significant in their own right.

587 For this reason, the poroelastic stress changes shown here may still be sig-
588 nificant enough to direct a given stimulation treatment. However, in order to
589 better evaluate the possibility of directing a stimulation treatment, the effect
590 of the stress redistribution associated with the events occurring during the
591 first stimulation treatment on the far-field stresses should be investigated, for
592 example with a Mohr-Coulomb plasticity model. Regardless of whether this
593 stress preconditioning methodology is employed or not, stress redistribution
594 associated with shear failure in the stimulated region of the first well is likely
595 to occur.

596 *5.1.4 Use of an Equivalent Continuum Plane Strain Elastic Model*

597 These investigations could have been performed with a discontinuum model
598 instead of the equivalent continuum approach implemented here. Discontin-
599 uum models, where fractures are explicitly modelled, represent a large body of
600 literature with many recent technical developments and applications to EGS
601 modelling (e.g., McClure and Horne (2014); Tene et al. (2017)). These models
602 are better equipped than equivalent continuum models to predict small-scale
603 behaviour and are generally able to more realistically replicate a specific site's
604 response to fluid injection. However, these models are generally more computa-
605 tionally expensive and require longer simulation times than equivalent contin-
606 uum models. Indeed, equivalent continuum models are capable of investigating
607 the effects of fluid injection in a fractured media, and can be especially useful
608 for larger-scale simulations, such as those performed here. These use of equiv-
609 alent continuum models for fractured-media simulations has been addressed
610 previously (Oda, 1986; Miller, 2015; Gan and Elsworth, 2016).

611 The use of a 2-D plane strain model over a 3-D or generalized plane strain
612 model (e.g., Cheng (1998)) is valid when the wellbore is long compared to
613 its diameter and in-line with one of the principal stress directions (Cheng,
614 2016). It is possible, however, that out-of-plane displacements, especially near
615 the heel and toe of the wells, might alter the results slightly. In these regions,
616 during the fluid injection, it is likely that changes to the principal total stress
617 parallel to the wellbore will be slightly reduced if this effect is included.

618 The mechanical model used here is also entirely elastic. It is probable that a
619 more rigorous approach would alter the magnitude of the stress changes found.
620 For example, Pijnenburg et al. (2018) recently showed that the use of an elastic
621 simulator during the modelling of fluid production in a sandstone likely under-
622 predicts strains and over-predicts total stress changes in the case that the
623 deformation is inelastic. Essentially, the use of a linear elastic simulator here
624 corresponds to the assumption that the non-linear responses of the system
625 remain localized such that the mechanical behaviour of the rock mass as a
626 whole can be well represented by such a linear elastic model (Cornet, 2016).

627 Further, certain parameters are likely to change throughout the stimulation
628 procedure. For example, the relatively low Poisson's ratio chosen due to the
629 fractured nature of the rock is likely to increase as shear failure occurs (Walsh,
630 1965). This would have implications for the magnitude of the changes to each

631 component of the stress tensor. Deformation-induced porosity changes were
632 also not accounted for here; an effect which may quantitatively influence the
633 results. It is also probable in reality that many of the poroelastic parameters
634 used here vary with effective stress (e.g., Walsh (1965); Bernabé (1986)). This
635 variation was not accounted for in the analyses performed here, unlike in other
636 equivalent continuum models applied to EGS (e.g., Gan and Elsworth (2016)).

637 5.2 Implications

638 Variations on this approach could be imagined. For example, stimulating both
639 wells at the same time would allow for both wells to benefit from advanta-
640 geous stress changes. However, each well would experience less preconditioning
641 Coulomb stress changes than the second well experienced during these simula-
642 tions. This is due to the fact that the pore pressures will not yet have reached
643 their post-stimulation values. Additionally, this approach would require suffi-
644 cient pumping power to stimulate two wells at once. Another possibility would
645 be to use the poroelastic and thermoelastic stress changes associated with an
646 existing doublet-well system to direct the stimulation treatment of a third
647 well. This would presumably incur larger stress changes than those used here
648 and would allow for the more efficient direction of the stimulation of the third
649 well.

650 It should be noted that one possible drawback to not flowing back the
651 wells is the possibility of inducing a large seismic event. Frequently these large
652 magnitude events occur after stimulation activities have been stopped (e.g.,
653 (Häring et al., 2008; Kim et al., 2018)), and it has even been suggested that
654 flowing the wells back could help prevent seismicity (McClure, 2015). Despite
655 this, the methodology proposed here is designed to use the built up poroe-
656 lastic stresses due to the increased pore pressure associated with injection to
657 facilitate the stimulation of another well. As shown in Section 4.1, flowing the
658 well back makes this process significantly less effective.

659 The successful implementation of this methodology would yield a num-
660 ber of advantages. Engineers would have higher confidence in connecting two
661 wells separated by a given distance when using this methodology as opposed
662 to the case where the wells are flown back before the next stimulation. Alter-
663 natively, wells could be separated by a larger distance, reducing the risk of
664 short-circuiting and increasing the contact time of the circulating fluid with
665 the reservoir. Additionally, because this methodology encourages the second
666 stimulation treatment to advance towards the first well, it seems less likely
667 that this stimulation treatment will stimulate a large fault as the total stimu-
668 lated reservoir volume is reduced for a given well separation distance. Further,
669 it can be imagined that this type of technique could be implemented in com-
670 bination with other similar techniques, such as fluid production, to provide
671 reservoir engineering solutions for large-scale reservoir creation. Of course, the
672 ability to influence the direction of a stimulation treatment does not mean that
673 operators have total control over how the stimulation treatment propagates,

674 simply that the stimulation treatment is guided such that it is more likely to
675 advance in a certain direction.

676 5.3 Future Outlook

677 These results potentially have implications for hydraulic fracturing. Although
678 not directly applicable, it has been shown that poroelastic stress changes dur-
679 ing injection are able to alter the stress field and affect a shear stimulation.
680 Mode I fracturing depends on the pore pressure overcoming the minimum prin-
681 cipal stress. It can therefore be imagined that both injection and production
682 are capable of altering the minimum stress such that mode I fracture propaga-
683 tion is either attracted to or repelled from a particular region of a reservoir. In
684 fact, it has already been shown that hydraulic fracture propagation is affected
685 by pre-existing injection and production wells (e.g., Berchenko and Detournay
686 (1997); Gao et al. (2019)). Further investigations should be performed on how
687 to purposefully use these stress changes to help direct hydraulic fracturing
688 treatments.

689 The numerical results found here indicate that a shear stimulation treat-
690 ment can be directed in a critically-stressed crust. Following this, experimental
691 work should be carried out to try to achieve these results in a real experimental
692 rock laboratory. Should those experiments be successful, other methodologies
693 for directing stimulation treatments should be investigated, especially ones ca-
694 pable of directing stimulation treatments in less critically-stressed reservoirs.

695 6 Conclusion

696 In this work, shear stimulation treatments in critically-stressed fractured granitic
697 rock from horizontal wells have been directed via the stress changes associ-
698 ated with a previous stimulation to preconditioning the stress field for the
699 next stimulation. These stress changes increase the Coulomb stress primarily
700 in the region between the two wells which results in the stimulation treat-
701 ment of the second well preferentially propagating towards the first. These
702 results have implications for reservoir engineering applications in EGS reser-
703 voirs. Further research should be performed to both confirm the results in
704 meso-scale field demonstrations and develop methodologies for directing stim-
705 ulation treatments in less critically-stressed reservoirs.

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