# Forcing of the MJO-related Indian Ocean Heating on the Intraseasonal Lagged NAO

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#### Abstract

Based on the experiments with the Coupled Forecast System version 2 (CFSv2), the mechanism by which the Madden-Julian Oscillation (MJO) modulates the North Atlantic Oscillation (NAO) is investigated. To isolate the cyclic MJO heating with an eastward propagation over the Indian Ocean and western Pacific, three sets of experiments are conducted with spatio-temporal varying heating added to the model's internally generated heating. The results suggest that the anomalous MJO heating over the Indian Ocean, rather than the western Pacific, dominates the formation of the NAO anomaly in the following 10-20 days. The MJO heating triggers a westward propagation of the storm track that influences Europe, complementing the eastward pathway of influence via the North Pacific. Both pathways contribute to an enhanced storm track over the North Atlantic and the positive NAO anomaly. The Eurasian pathway is less important for the formation of the negative NAO anomaly.

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16	Key Points:					
17	• The lagged connection of the NAO with the MJO is reproduced well in the CFSv2 model					
18	with prescribed cyclic MJO heating forcings					
19	• The MJO-related heating (cooling) over the Indian Ocean alone plays a dominant role in					
20	the formation of the positive (negative) NAO anomaly					
21	• The MJO affects the NAO via a secondary Eurasian (westward oriented) pathway, and					
22	the primary North Pacific (eastward oriented) pathway					
23						

# 24 Abstract

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## 36 Plain Language Summary

As the dominant source of predictability for subseasonal forecasts, the Madden-Julian 37 Oscillation (MJO) has a great impact on the leading winter extratropical mode, the North 38 Atlantic Oscillation (NAO). Previous studies suggested that the impact is through the 39 downstream propagation of the MJO-excited Rossby wave across the North Pacific and its 40 interaction with synoptic eddies over the North Atlantic. In addition to the traditional Pacific 41 pathway, here we find another pathway by which the MJO convection over the Indian Ocean 42 triggers a westward propagation of the Asian storm track that penetrates Eurasia and influence 43 the wave-eddy interaction related to the NAO. 44

# 45 **1 Introduction**

The North Atlantic Oscillation (NAO), one of the dominant mode of the extratropical atmospheric variability, has a significant influence on weather and climate over the wide northern hemisphere (Hurrel et al., 2003). The NAO varies temporally with periods spanning a broad spectrum from synoptic to decadal, and its activity at intraseasonal timescale is one substantial source of extended-range weather forcecasting. To predict NAO activity at intraseasonal timescales is thus of substantial importance.

The Madden–Julian Oscillation (MJO), the prominent mode of tropical intraseasonal variability, is characterized with large-scale convection propagating eastward through the Indian Ocean and the western Pacific along with baroclinic atmospheric circulation anomalies (Madden & Julian, 1994; Zhang, 2005). Since the associated heating drives atmospheric teleconnections, the MJO has a great impact on the extratropical atmospheric circulations (Garfinkel et al., 2014; Matthews et al., 2004; Seo & Son, 2012; Stan et al., 2017; Zheng & Chang, 2019).

Previous studies have revealed a lagged connection of the NAO with preceding MJO 58 59 events (Cassou, 2008; Jiang et al., 2017; Lin et al., 2009). The positive (negative) NAO tends to 60 follow the enhanced (reduced) convection over the Indian Ocean at a lag of 10–15 days, which corresponds to the MJO phase 3 (phase 6). Such a lagged connection is reproduced well both in 61 simplified general circulation models (GCMs) forced with the fixed MJO heating (Lin & Brunet, 62 2018; Shao et al., 2019) and in fully coupled atmosphere-ocean models forced with the eastward-63 moving MJO heating (Straus et al., 2015; Yadav et al., 2019). This suggests that the MJO's 64 leading impact on the NAO is a robust intrinsic feature of the atmosphere. Thus it provides an 65 important source for subseasonal predictability of the NAO (Lin et al., 2010; Tseng et al., 2018). 66

Physically, the MJO's impact on the NAO may originate from the MJO-excited Rossby 67 wave propagation toward the North Altantic and its subsequent interaction with synoptic eddies 68 locally (Cassou, 2008; Lin et al., 2009). An initiation of Rossby wave trains in the Pacific and 69 the interactions among the transient eddies from the eastern Pacific during the MJO phase 3 70 71 (phase 6) are considered to be responsible for setting the positive (negative) NAO phase (Rivière 72 & Drouard, 2015; Fromang & Rivière, 2020). Another pathway of influence was suggested by Lin et al. (2015), who indicated that the MJO convection anomaly over the Indian Ocean may 73 induce a disturbance in the South Asian jet that propagates along the midlatitude waveguide 74 (e.g., Branstator, 2002) and affects the NAO. In addition to the wave train across the North 75 Pacific, Shao et al. (2019) found that the MJO heating initiates an upstream and westward 76 propagation of disturbances to western Asia-Europe along the Asian subtropical jet. Lin and 77 Brunet (2018) also revealed an early influence of a westward propagating Rossby wave excited 78 by the MJO. This suggests the possibility of multiple pathways for the mechanism of the MJO's 79 influence on the NAO. 80

To further explore the possible pathways and understand the mechanism for the MJO–NAO 81 connection, it is helpful to isolate the individual contribution of the MJO-related regional and 82 moving heating over the Indian Ocean and the western Pacific. On one hand, the positive heating 83 anomaly over the western Pacific tends to induce a negative-phased NAO (Li et al., 2006b; Lin 84 et al., 2005; Yadav and Straus, 2017), and this influence is sensitively dependent on the 85 longitudinal position of the heating. On the other hand, the tropical Indian Ocean heating induces 86 an NAO tendency toward its positive phase (Hoerling et al., 2004; Li et al., 2006a; Lin et al., 87 2015; Yu & Lin, 2016). Note that there is potential interference between the effects of the MJO's 88 Indian Ocean and western Pacific heating, because one excited Rossby wave can alter the 89 background flow and affect the route of a Rossby wave excited by its counterpart. 90

In order to isolate the individual and combined roles of the MJO-linked regional heating, we use a coupled ocean–atmosphere–land model, CFSv2, to conduct sensitivity experiments. Since the initial state is still important for the extended-range weather forecasting, we use observed initial conditions, but add idealized, temporally evolving heating, designed to mimic realistic propagating MJO heating, to the full coupled model. Such a scheme preserves the information embedded in the initial state containing the preceding MJO's signal, and thus has advantages over that with fixed heating forcing (Fromang & Rivière, 2020).

# 98 2 Model and Methods

99 2.1 Model Experiments

The coupled model, Climate Forecast System model version 2 (CFSv2) of the National Centers for Environmental Prediction (NCEP), is used. Its atmospheric component model (the Global Forecast System) has a T126 horizontal resolution with 64 vertical levels. The outputs are interpolated to 1°×1° resolution and 11 pressure levels. The oceanic component is the Modular Ocean Model version 4 (MOM4) from the Geophysical Fluid Dynamics Laboratory (GFDL). The details of the CFSv2 are described in Saha et al. (2014).

The Control runs consist of four-month reforecasts (1 December to 31 March) beginning from three sets of perturbed initial conditions (ICs) on 1 December of each year over the 31-yr period (1980–2010), thus yielding a total of 93 simulations of 121 days. The ICs for atmosphere, ocean, sea-ice, and land surface are obtained from the CFS Reanalysis (CFSR) (Saha et al., 2010), of which the atmosphere ICs are perturbed by blending different atmospheric states
through a weighted average (e.g., Yadav et al., 2019).

An idealized MJO evolving heating is added to model generated temperature tendency to 112 mimic realistic MJO cycles. Here the idealized MJO heating is prescribed with the spatial pattern 113 and propagating features by combining the technique developed by Jang and Straus (2012) and 114 Straus et al. (2015), and with the vertical profile based on estimates of the total diabatic heating 115 from the tropical rainfall measuring mission (TRMM) as Straus et al. (2015). The technique was 116 also recently applied for model heating bias correction by Swenson et al. (2019). The magnitude 117 and temporal structure of the idealized heating are overall based on observations of Yadav and 118 Straus (2017). Such an added heating, even with a relatively small magnitude, can effectively 119 organize the evolution of tropical diabatic heating in the model to give a robust MJO (Straus et 120 121 al., 2015).

To understand the model response to isolated MJO cycles, three sets of MJO-heating 122 experiments with heating forcing over different regions are conducted, including the Indo-123 western Pacific Ocean (MJO IP) runs, Indian Ocean (MJO IO) runs, and western Pacific 124 125 (MJO\_WP) runs. All these experiments are listed in Table S1 in the supporting information. Note that the MJO\_IP runs represent the typical MJO cycles, similar to the Slow Repeated 126 Cycles runs in Yadav et al. (2019), except for a smaller ensemble size. A maximum added 127 heating rate of 0.8 K  $d^{-1}$  was chosen with the overall magnitude scaled, so that the *total* (model 128 generated plus added) heating was roughly close to the observations. A fixed phase speed of 3° 129 longitudes per day was adapted to match the slow cases, yielding a period of 60 days. 130

Figures 1a–1f display the spatial structure and temporal evolution of the idealized heating 131 for these experiments. The vertical structure of the heating (Figures 1a-1c) indicates an evolution 132 133 from shallow to deep gradually (Yadav et al., 2019), accompanied by an eastward propagation with two complete MJO cycles (Figures 1d–1f). The added heating decays off the equator as a 134 Gaussian form with an e-folding scale of 15° latitudes (Figure S1). The vertically averaged total 135 diabatic heating anomalies are shown in Figures 1g-1i. The diabatic heating in the MJO\_IP runs 136 matches the observational slow MJO cases of Yadav and Straus (2017), except for an 137 138 amplification by a factor of two. The heating in the MJO\_WP runs is additionaly induced over the eastern Indian Ocean centering around 90°E, lagging that in MJO\_IO runs by about 10 days. 139 140 This may be caused by the feedback of the mid-latitude waves to tropical convective heating

141 (Frederiksen & Lin, 2013; Moore et al., 2010). Overall, the prescribed heating mirrors the
142 horizontal structure of the MJO diabatic heating reasonably well (Figure S2).

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Figure 1. The pressure-longitude cross-section of added heating anomalies on day 12 (a–c), vertically averaged (over 1000–50 hPa) added heating anomalies (d–f) and total diabatic heating anomalies (g–i) at each reforecast day. Dispalyed is the value averaged over latitudes 15°S–15°N. From left to right, it corresponds to MJO\_IP, MJO\_IO, and MJO\_WP runs, respectively. Units are K day<sup>-1</sup>.

149 2.2 Analysis Methods

The NAO pattern is defined as the first empirical orthogonal function (EOF) mode of the daily 500-hPa geopotential height fields (Z500) over the North Atlantic region (20°N–85°N,

90°W-30°E). The daily NAO indices are calculated as the projection of the daily Z500 152 anomalies onto the NAO pattern; this time series of projections is subsequently normalized. The 153 occurrence probabilities of the positive (negative) phase of the NAO at each reforecast day are 154 obtained by counting the number of days with the daily NAO index greater (less) than 0.5 (-0.5) 155 and then dividing by the total ensemble size of 93. When comparing the contribution of the MJO 156 heating imposed over different regions, the identical NAO pattern in Control runs is used for 157 calculating the NAO index. For each experiment, anomalies are defined as departures from the 158 annual cycle in the Control runs. This model annual cycle is estimated by projecting the daily 159 time series of each winter separately onto the first three Legendre polynomials in time as 160 described in Straus et al. (2003), and then averaging over all winters. We use Fourier 161 decomposition as a spatial filter method to isolate the synoptic-scale (zonal waves 5–12) eddies. 162

# 163 **3 Results**

The NAO pattern in Control runs (Figure 2a) agrees well with that in the ERA-Interim 164 reanalysis (Dee et al., 2011) as shown in Yadav et al. (2019). For the MJO\_IP runs, the NAO 165 pattern is similar to that in Control runs, with a slightly stronger amplitude over the positive 166 sector (Figure 2b). As the model simulates both the evolving MJO cycles and the NAO pattern 167 well, we first investigate whether the observed MJO-NAO linkage can be reproduced in the 168 MJO IP runs. Figure S3 illustrates the occurrence probabilities of the NAO at each reforecast 169 day in the MJO\_IP runs. For the first (second) cycle, the occurrence probabilities of the positive 170 NAO are increased about 10–15 days after the MJO heating reaches roughly the 90°E longitude 171 at around day 15 (80). Conversely, the increase of the negative NAO occurs 15-20 days 172 following the MJO heating (cooling) over the western Pacific (Indian Ocean) at around day 40. 173 This suggests that the MJO-NAO connection can be captured well by forcing this coupled model 174 with the anomalous heating associated with the MJO cycles. 175

In order to compare the contribution of the MJO heating over different regions, the evolution of NAO index from day 1 to 121 in each experiment is displayed (Figure 2c). Similar to the occurrence probabilities, the averaged NAO indices show fluctuations of roughly two cycles. During the first 10 days, the NAO indices in each MJO-heating runs are close to that in Control runs. After day 15, the NAO index in MJO\_IO runs begins to increase until around day 30. Generally, the increase of the NAO index in MJO\_IO runs leads that in MJO\_WP runs for about 10 days during the first cycle. This indicates that the positive NAO anomaly may be firstly initiated by the Indian Ocean heating of the MJO. Note that the NAO indices in the second cycles display relatively weaker amptitudes to those in the first cycles, in coincidence with the somewhat weaker MJO heating anomalies over the Indian Ocean. As complex tropicalextratropical interactions are involved in the second cycle, we mainly focus on the NAO-like response in the first cycle.

It is difficult to directly interpret the NAO-like response to the MJO-related regional 188 189 forcing because the additional heating over the Indian Ocean (western Pacific) are induced in the MJO\_WP (MJO\_IO) runs. To isolate the contribution of the MJO forcing over the western 190 Pacific, we examine the differences in the heating (see Figure S4a) and response between the 191 MJO\_IP and MJO\_IO runs. The differences of Z500 over the North Atlantic between the 192 MJO\_IP and MJO\_IO runs, and between the MJO\_IP (MJO\_IO) and Control runs, are projected 193 onto the NAO pattern in Control runs. Figure 2d displays the projection indices of the Z500 194 differences for the MJO\_IP minus MJO\_IO runs (green curve), MJO\_IP minus Control runs (red 195 curve), and MJO IO minus Control runs (blue curve). The green curve shows that the MJO 196 197 heating over the western Pacific only contributes at a later stage. In contrast, the blue curve indicates a dominant role of the anomalous Indian Ocean heating on the NAO anomaly. The 200-198 199 hpa streamfunction together with wave activity vectors for the raw MJO IP runs demonstrate that the MJO-initiated Rossby waves originate mostly from the Indian Ocean heating/cooling 200 201 region (Figure S5). Seo and Son (2012) showed that the Indian Ocean heating can generate a much broader and stronger Rossby wave source because of its position relative to the subtropical 202 jet, where the relative vorticity is large. This further verifies the crucial role the anomalous MJO 203 heating over the Indian Ocean. Thus, we will concentrate on the response to the heating forcing 204 205 in the MJO\_IO runs.





Figure 2. The 500-hpa geopotential height (Z500) anomalies (unit: gpm) regressed onto the NAO index in (a) Control runs and (b) MJO\_IP runs, together with (c) ensemble-averaged NAO index at each reforecast day for the MJO\_IP (red curve), MJO\_IO (blue curve), MJO\_WP (green curve) and Control runs (black curve), and (d) projection indices of the Z500 differences for the MJO\_IP minus MJO\_IO runs (green curve), the MJO\_IP minus Control runs (red curve), and the MJO\_IO minus Control runs (blue curve), respectively.

The differences of the Z500 between the MJO\_IO and Control runs from day 7 to day 21 are shown in Figure 3, which illustrates the evolution of the extratropical response to the MJO heating. Since the diabatic heating anomaly over the Indian Ocean does not appear until around day 7 (see Figure S4b), there is no evident response before day 11. From day 13 to day 15, negative anomalies of the geopotential height are built up over the subtropical Asia and extend northwestward. There is also formation of an anomalous anticyclone over the western North

Pacific, along with a north-south dipole over the eastern North Pacific. Over the high latitudes of 219 the North Atlantic, a negative height anomaly is formed, which deepens over the Greenland 220 region until day 17. Meanwhile, a positive height anomaly to the northwest of the Indian Ocean 221 extends westward, merging with the one over the subtropical North Atlantic, together forming 222 the southern part of the NAO-like pattern. Thus, a positive NAO-like pattern is developed. 223 Delayed by 15-20 days from the cooling over the Indian Ocean (see Figure S4b), a strong 224 negative NAO-like response is developed at around days 55-61 (Figure S6). Prior to the 225 formation of negative NAO anomaly, the positive height anomalies over the Eurasia and negative 226 height anomalies over the North Pacific are established respectively, which are generally 227 reversed to the situation following the heating over the Indian Ocean. 228



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Figure 3. Composite differences of the 500-hPa geopotential height (unit: gpm) between the MJO\_IO and Control runs from day 7 to day 21. Dotted is significant at the 90% confidence level according to a Student's t test.

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Because the subtropical jet can act as a waveguide to modulate the wave/eddy propagation (Lee et al., 2019; Zheng & Chang, 2020), we further investigate the differences of the zonal wind at 300 hpa between the MJO IO and Control runs (Figure S7). At the initial stage (days 7–12), a

positive zonal wind anomaly is induced over the subtropical Asian region and enhances with both westward and eastward extension. At days 13–15, a positive zonal wind anomaly is formed over the mid-latitudes of North Atlantic. Then it strengthens and exhibits a poleward shift at days 16–18, along with a negative anomaly built up to its south. This suggests that the North Atlantic mid-latitude jet is pushing northward, in favor of the formation of the positive NAO anomaly.

To understand how the anomalous subtropical Asian jet related to the MJO modulates the 242 North Atlantic jet, we examine the shift in the storm tracks. Figure S8 shows the evolution of the 243 300-hPa synoptic-scale v'v' anomaly (the difference between the MJO IO and Control runs). At 244 days 10–12, the positive v'v' anomalies are established along the South Asia jet, which extend to 245 the western Europe at days 13-15, indicating an expanding of Asian storm track influencing the 246 western Europe directly. Subsequently, the storm tracks are enhanced over the eastern North 247 248 Atlantic. In addition, the strengthening storm tracks over the North Pacific exihibit an eastward extension, which seem to influence the North Atlantic at days 22-24. Both eastward and 249 westward extension of the enhanced storm tracks are likely to contribute to the positive NAO 250 anomaly. On the contrary, the weakened subtropical Asian jet (at days 43-45) and the 251 252 subsequent weakening storm tracks (at around days 49-57) over the North Atlantic are conducive to the negative NAO anomaly. 253

254 In order to detect the pathways via which the MJO-induced disturbances propagate into the North Atlantic, we further calculate the one-point lead-lag correlation of the 300-hpa synoptic-255 256 scale meridional wind anomaly (v'), where a base point is chosen nearby the subtropical Asian jet at 60°E, 30°N. The lead-lag correlation coefficients are presented as a function of time and 257 longitude, averaged from 30°N to 40°N (Figure 4). Consistent with the westerly jet stream, the 258 synoptic eddies generally exhibit eastward movement. The westward propagation indicated by 259 the dashed line in Figure 4a suggests that the synoptic eddies over the North Atlantic at around 260 days 12–15 may be affected by those from the South Asia jet. From day 16 to 18, the synoptic 261 eddies over the North Atlantic in MJO\_IO runs exhibit a faster eastward phase speed with 262 respect to those in Control runs, along with a northeast-southwest tilt (Figure S9), favoring the 263 development of the positive NAO phase (e.g., Ren et al., 2012; Song, 2016). As such, the 264 anomalies along the Asian subtropical jet may firstly influence the positive NAO through the 265 westward Asian-Europe pathway. 266

On the other hand, the disturbances induced along the subtropical Asian jet propagate 267 across the North Pacific and downstream into the North Atlantic from day 12 to day 28. The 268 increase of the projection index later at around days 22-30 (see Figure 2d) could be largely 269 attributed to the Pacific pathway. This also suggests that it takes a longer time for the MJO 270 heating over the Indian Ocean to affect the NAO via the Pacific pathway than the westward 271 272 pathway. The eastward tilt of the lines in Figure 4b shows a more traditional Pacific pathway (at days 45-60) for the negative NAO-like responses at around days 55-60, while the westward 273 274 pathway is much weaker. This indicates that the Eurasian pathway may play a secondary role.



Figure 4. One-point lead-lag correlation coefficients of 300-hpa synoptic-scale v' for a base point at 30°N, 60°E on day 13 (a) and day 45 (b), averaged over the latitudes 30°N–40°N. Here, v' is the difference of the meridional wind (waves 5–12) between the MJO\_IO and Control runs. The ordinate is reforecast time. The interval of the contours is 0.1.

# 280 **4 Summary and Discussions**

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The mechanism for the observed lagged-connection of the NAO with the MJO is investigated by forcing a coupled model (CFSv2) with idealized cyclic MJO heatings. Three sets of experiments, each having an integration of two eastward propagating MJO cycles with a period of about 60 days, are conducted to isolate the contribution of the MJO-related heating anomalies over the Indian Ocean and western Pacific, respectively. The results demonstrate an increase of the positive (negative) NAO occurrence following the MJO heating (cooling) over the Indian Ocean at a lag of 10–15 (15–20) days. Furthermore, the anomalous MJO heating over

the Indian Ocean, rather than that over the western Pacific, dominates the NAO-like response. 288 The MJO heating over the Indian Ocean can exert its influence via the traditional pathway across 289 the North Pacific, however part of the impact may be transmitted via the Eurasia, which takes a 290 shorter time (around 7-10 days). The shift of the strengthened subtropical Asian jet causes an 291 expanding Asian storm track to propagate westward, influencing the Europe and enhancing the 292 storm tracks over the eastern North Atlantic, thus favoring the development of the positive NAO. 293 On the other hand, the induced disturbances in the subtropical Asian jet propagate across the 294 North Pacific and downstream into the North Atlantic, enhancing the NAO after around 10–15 295 days. The situation is essentially reversed following the MJO's cooling over the Indian Ocean for 296 the negative NAO-like response, although the Eurasian pathway does not seem to play a role 297 here. 298

Since the MJO-related heating anomalies over the Indian Ocean and the western Pacific are not independent, and presumably feed back on each other, the dominant effect of the anomalous MJO heating over the Indian Ocean does not rule out the potential role of the heating over the western Pacific. The western Pacific heating anomaly has minor contribution, probably because the affected initial state of the portion of the subtropical jet is less effective in transporting the disturbances along the jet waveguide to the North Atlantic (e.g., Lin & Brunet, 2018).

While we have shed light on the direct MJO-related tropical forcing of the North Atlantic 306 307 region, other mechanisms are likely to play a role, such as the possible stratosphere pathway (e.g., Barnes et al., 2019; Lee et al., 2019). There are also tropical-extratropical interactions of 308 intraseasonal oscillations (Frederiksen & Lin, 2013), which may play a role in the model 309 experiments. For example, the synoptic eddies from the North Atlantic jet may affect those along 310 311 the subtropical jet, which in turn have an influence on the North Atlantic jet through downstream propagation. The mechanism for how the extratropics influence the tropical heating linked to the 312 MJO deserves further investigation. 313

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# **@AGU**PUBLICATIONS

2	[Geophysical Research Letters]
3	Supporting Information for
4	[Forcing of the MJO-related Indian Ocean heating on the Intraseasonal Lagged NAO]
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11	Contents of this file
12	Text S1
13	Figures S1 to S9
14	Table S1
15	Introduction
16	This supporting information provides the method for calculating the horizontal components of
17	the wave activity flux (Text S1), and description of the model experiments (Table S1). Figures
18	S1–S9 are provided to support the analysis presented in the main text.
19	Text S1.
20	We calculate the horizontal components of the phase-independent wave activity
21	flux derived by the Takaya and Nakamura (2001), and assume that the group velocity
22	is zero. Using their equation (38), a two-dimensional (horizontal) wave activity flux is

23 expressed as follows:

$$W = \frac{p\cos\phi}{2|U|} (W_{\lambda}, W_{\phi}) + C_U M$$
(1)

24 with

25 
$$W_{\lambda} = \frac{U}{a^2 \cos^2 \phi} \left[ \left( \frac{\partial \psi'}{\partial \lambda} \right)^2 - \psi' \frac{\partial^2 \psi'}{\partial \lambda^2} \right] + \frac{V}{a^2 \cos \phi} \left[ \frac{\partial \psi'}{\partial \lambda} \frac{\partial \psi'}{\partial \phi} - \psi' \frac{\partial^2 \psi'}{\partial \lambda \partial \phi} \right]$$
(2)

26 
$$W_{\phi} = \frac{U}{a^2 \cos \phi} \left[ \frac{\partial \psi'}{\partial \lambda} \frac{\partial \psi'}{\partial \phi} - \psi' \frac{\partial^2 \psi'}{\partial \lambda \partial \phi} \right] + \frac{V}{a^2} \left[ \left( \frac{\partial \psi'}{\partial \lambda} \right)^2 - \psi' \frac{\partial^2 \psi'}{\partial \phi^2} \right]$$
(3)

where  $\psi'$  is the perturbation streamfunction, (U, V) are the basic state zonal and meridional winds,  $(\varphi, \lambda)$  are latitude and longitude, *a* is the earth's radius, and *p* is pressure.

30

31



32

Figure S1. The vertically averaged (over 1000–50 hPa) added heating anomalies (unit: K day<sup>-1</sup>)
in the MJO\_IP (top), MJO\_IO (middle), and MJO\_WP runs (bottom).



Figure S2. The vertically averaged (over 1000–50 hPa) diabatic heating anomalies (unit: K day<sup>-1</sup>)
in the MJO\_IP (top), MJO\_IO (middle), and MJO\_WP runs (bottom).



Figure S3. Occurrence probabilities (%) of the positive (red curve) and negative (black curve)
NAO from day 1 to day 121 in the MJO\_IP runs. The red (black) reference line represents the
climatological mean of the positive (negtive) NAO with value 33.7 (29.9).



Figure S4. The differences of the vertically averaged (over 1000–50 hPa) diabatic heating (unit: K day<sup>-1</sup>) averaged over latitudes 15 S–15 N between the MJO\_IP and MJO\_IO runs (a), and between the MJO\_IO and Control runs (b).



Figure S5. Composites of the 200-hPa streamfunction (contours), horizontal components of phase-independent wave activity flux (arrows) and diabatic heating anomalies (shaded) in the raw MJO\_IP runs. Contour interval is  $1.0 \times 10^{-6} \text{ m}^2 \text{ s}^{-1}$  for streamfunction, plotted north of 15 N. Units of wave activity flux are m<sup>2</sup> s<sup>-2</sup>, and only vectors larger than 0.5 are shown. Units of heating anomalies are K day<sup>-1</sup>.





Figure S6. Composite differences of the 500-hPa geopotential height (unit: gpm) between the
MJO\_IO and Control runs from day 40 to day 61. The composite results at the 95% confidence
level according to a Student's t test are dotted.



Figure S7. Composite differences of the 300-hPa zonal wind (unit: m s<sup>-1</sup>) between the MJO\_IO
and Control runs averaged every 3 days.



**Figure S8.** Composite differences of the 300-hPa synoptic v'v' (unit: m<sup>2</sup> s<sup>-2</sup>) between the MJO\_IO and Control runs averaged every 3 days. The composite results at the 90% confidence level according to a Student's t test are dotted.



68 Figure S9. One-point lagged regression coefficients of the 300-hPa anomalous synoptic-scale

69 streamfuction for the base point at 45 N, 30 E on day 15 in Control runs (shading) and MJO\_IO

- runs (contours) respectively, from day 15 to day 18. The interval of the contours is 0.2.
- 71

72	Table S1.	Description	of Model Experiment	ts
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Experiments	Added heating region	Ensemble size	Duration	Simulations
Control	none	3	December to March in 1980–2010	93
MJO_IP	45 °-180 E, 20 S-10 N	3	December to March in 1980–2010	93
MJO_IO	45 °-130 °E, 20 °S-10 °N	3	December to March in 1980–2010	93
MJO_WP	115 °-180 °E, 20 °S-10 °N	3	December to March in 1980–2010	93