

Statistical Relationship between Interplanetary Magnetic Field Conditions and the Helicity Sign of Flux Transfer Event Flux Ropes

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Abstract

Flux Transfer Events (FTEs) are transient phenomena produced by magnetic reconnection at the dayside magnetopause typically under southward interplanetary magnetic field (IMF) conditions. They are usually thought of as magnetic flux ropes with helical structures forming through patchy, unsteady, or multiple X-line reconnection. While the IMF often has a non-zero B_Y component, its impacts on the FTE flux rope helicity remain unknown. We survey Magnetospheric Multiscale (MMS) observations of FTE flux ropes during the years 2015 – 2017 and investigate the solar wind conditions prior to the events. By fitting a force-free flux rope model, we select 84 events with good fits and obtain the helicity sign (i.e., handedness) of the flux ropes. We find that positive (negative) helicity flux ropes are mainly preceded by a positive (negative) B_Y component. This finding is compatible with flux ropes formed through a multiple X-line mechanism.

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Key Points:

- The helicity sign of 84 Flux Transfer Events is studied using force-free flux rope model fitting
- Right-handed (left-handed) FTE flux ropes are mostly preceded by positive (negative) IMF B_Y
- This IMF B_Y control of the helicity sign is compatible with a multiple X-line formation mechanism

Abstract

Flux Transfer Events (FTEs) are transient phenomena produced by magnetic reconnection at the dayside magnetopause typically under southward interplanetary magnetic field (IMF) conditions. They are usually thought of as magnetic flux ropes with helical structures forming through patchy, unsteady or multiple X-line reconnection. While the IMF often has a non-zero B_Y component, its impacts on the FTE flux rope helicity remain unknown. We survey Magnetospheric Multiscale (MMS) observations of FTE flux ropes during years 2015 – 2017 and investigate the solar wind conditions prior to the events. By fitting a force-free flux rope model, we select 84 events with good fits and obtain the helicity sign (i.e., handedness) of the flux ropes. We find that positive (negative) helicity flux ropes are mainly preceded by positive (negative) B_Y component. This finding is compatible with flux ropes formed through a multiple X-lines mechanism.

Plain Language Summary

The Earth’s near-space environment is very dynamic, with transient phenomena triggered by interaction between the solar wind, a wind of accelerated ions and electrons flowing outward from the Sun, and the Earth’s magnetopause, the magnetospheric boundary that shields us from the solar wind by deflecting it around. The solar wind carries along an Interplanetary Magnetic Field (IMF) whose orientation determines the dynamics of the interaction. When the IMF is southward, magnetic reconnection, a phenomenon that allows sheared magnetic fields to rearrange and release magnetic field energy into particle energy, can be triggered at the Earth’s magnetopause on the dayside. A Flux Transfer Event (FTE) is a transient portal that allows the bursty transfer of solar wind into the Earth’s magnetosphere. FTEs are believed to form due to patchy, transient reconnection or in between multiple reconnection sites. An FTE is envisaged as a twisted magnetic field structure with helical field that looks like a rope. For the first time in space, we study the relationship between the IMF orientations and the twist direction of an ensemble of FTEs observed by NASA’s Magnetospheric Multiscale mission, by modelling FTEs as magnetic flux ropes. We found that the flux rope twist direction is controlled by the IMF orientation, such that the rope is twisted in the left-handed or right-handed sense depending on the east-west component of the IMF. This result supports the formation of FTEs by a multiple reconnection mechanism.

1 Introduction

A Flux Transfer Event (FTE) is a transient phenomenon generated at magnetospheric magnetopauses, and has been most studied at the Earth. It is recognised in spacecraft data as a bipolar magnetic field variation in the direction normal to the magnetopause (B_N), with enhanced core field (C. Russell & Elphic, 1978). This magnetic field profile suggests a magnetic flux rope structure with helicoidal field. Various formation mechanisms have been proposed for FTEs, such as transient and patchy dayside reconnection (C. Russell & Elphic, 1978), single X-line with unsteady reconnection rate (Scholer, 1988; D. J. Southwood et al., 1988), and multiple X-line reconnection (Lee & Fu, 1985; Raeder, 2006). In recent years, there are growing evidence supporting multiple X-line reconnection mechanisms (e.g., Hasegawa et al., 2010; Øieroset et al., 2011). In the Lee & Fu’s model, three reconnection X-lines are assumed to simultaneously exist in the presence of non-zero B_Y , leading to production of two helical flux tubes (with the same helical sense). In the Raeder’s model, FTEs only develop when the dipole tilt is large; they are formed as a result of non-stationary, sequential generation of new X-lines.

Since solar wind conditions control magnetic reconnection at the Earth’s magnetopause, they should control the nature and properties of FTEs. Early spacecraft surveys revealed that FTEs are strongly associated with southward IMF conditions (Berchem & Russell, 1984; C. Russell et al., 1996) consistent with generation from reconnection

at low latitudes (e.g., Paschmann et al., 1982). There is no strong control from other solar wind parameters such as plasma beta, dynamic pressure, and Mach number on the FTE occurrence (Kuo et al., 1995; Wang et al., 2006). The occurrence of FTEs is found dependent on the IMF orientation but not on its magnitude (Wang et al., 2006). The effect of the IMF B_Y component was studied in relation to the spatial distribution and motion of FTEs (e.g., Fear et al., 2012; Karlson et al., 1996). However, direct studies on relationships between the IMF B_Y and FTE topologies themselves are still limited.

FTEs are known to have twisted interior field (e.g., Cowley, 1982; Saunders et al., 1984) with a field-aligned core field and an azimuthal field increasing away from the core. Twisting features of FTEs have been theoretically evaluated in terms of magnetic helicity (Song & Lysak, 1989; Wright & Berger, 1989, 1990). Magnetic helicity is a measure that can quantify magnetic field topology into twist, shear, linking, and kinking of magnetic fields. It is defined as $H = \int_V \mathbf{A} \cdot \mathbf{B} d^3r$, where H is the total helicity of the entire magnetic field in a volume V , $\mathbf{B}$ is the magnetic field, $\mathbf{A}$ is the vector potential of $\mathbf{B}$ (i.e., $\mathbf{B} = \nabla \times \mathbf{A}$), and d^3r is the differential volume element. Here we use the definition of magnetic helicity to describe the twist of an FTE flux rope and we will only consider its sign. The twist direction around the core field can be characterized as the “handedness” or “sense/sign of the helicity” of the flux rope. The sign of flux rope helicity was studied mostly in magnetic clouds (e.g., Bothmer & Schwenn, 1998) and in flux ropes at Venus, Mars, and Titan (e.g. Martin et al., 2020; C. Russell, 1990; Wei et al., 2010) to understand their formation mechanisms. At Earth, magnetic helicity was studied in magnetotail flux ropes (Zhang et al., 2010). A few FTE flux ropes were observed in the magnetotail flank with the positive sign of helicity under southward and duskward IMF conditions, indicating that they originated on the dayside and survived far downstream (Eastwood et al., 2012). Here we present a first dedicated study of the sign of helicity of FTEs at the Earth’s dayside magnetopause.

Based on topological consideration, the helicity sign of FTEs should be controlled by the IMF. Fig 1 shows a schematic illustration of FTE formation by the multiple X-line reconnection mechanism under southward IMF with a non-zero B_Y component. In 2-D (Figs 1a,1c), as viewed from the dusk side, multiple reconnection between the magnetospheric and magnetosheath fields would produce a magnetic island (shown in purple) with an anti-clockwise field rotation. In 3-D, depending on the B_Y (out-of-plane) component, the magnetic island becomes a magnetic flux rope with an axial component pointing outward (Fig 1a) or inward (Fig 1c) from the plane. The magnetic field rotation (tangential component) with respect to the axial direction of the flux rope determines its handedness or helicity sign. In this picture, the southward IMF with positive B_Y would produce right-handed (RH) flux ropes (Figs 1a,1b) while the southward IMF with negative B_Y would produce left-handed (LH) flux ropes (Figs 1c,1d). Such topological consideration has yet to be statistically tested.

We present a statistical study of FTEs observed by NASA’s Magnetospheric Multiscale mission (MMS, Burch et al. (2015)) and characterize the twist profiles of FTEs by fitting into a flux-rope model with systematic effort. We first introduce selections of FTEs, instrumentations, and illustration of events. We then present statistical analyses of the solar wind conditions. Finally, discussion and conclusions are presented.

2 Data and Methods

2.1 Event selections and instrumentations

We first obtain a list of FTEs as observed by MMS. Fargette et al. (2020) published a list of MMS1 observations of 229 FTEs consisting of 186 flux-rope-type structures and 43 flux-rope-type with reconnection at a central current sheet (e.g., interlinked flux tubes). An FTE is selected based on visual inspection of data plots in the Geocentric Solar Eclip-

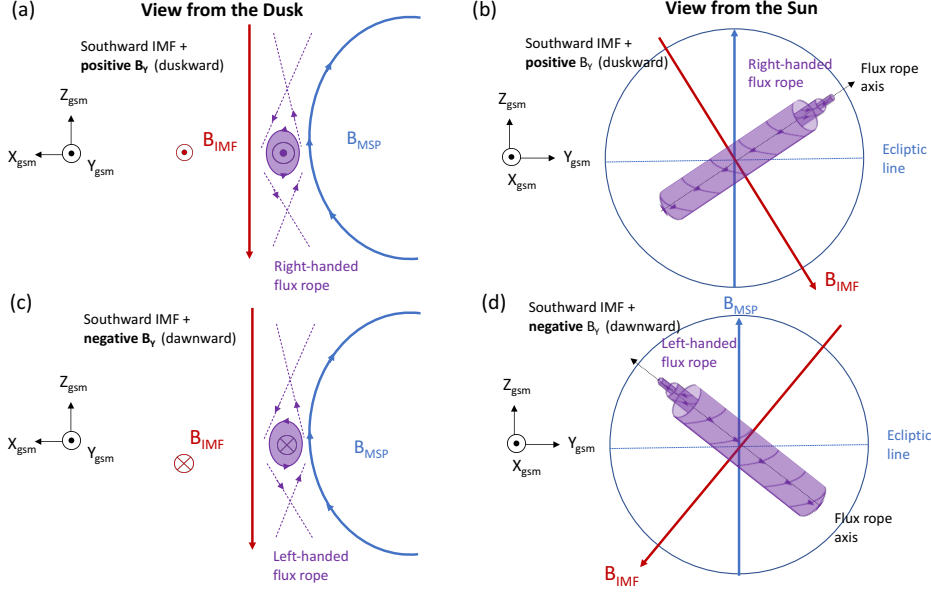


Figure 1. Schematic illustration of the generation of an FTE flux rope through the multiple X-line reconnection mechanism at the dayside magnetopause when the southward IMF has a significant (a, b) positive B_Y component and (c, d) negative B_Y component. The flux rope is shown in purple with arrows indicating the sense of twist when viewed from the dusk (a, c) and the Sun (b, d). When the flux rope is generated under southward IMF with positive B_Y , it has a right-handed sense of twist corresponding to a positive helicity. In contrast, when the flux rope is generated under southward IMF with negative B_Y , it has a left-handed sense of twist corresponding to a negative helicity.

tic (GSE) coordinates characterised by (1) a bipolar signature in one of the magnetic field components, and (2) an increase in the total (plasma and magnetic) pressure. We use Flux Gate Magnetometer (C. T. Russell et al., 2016) and Fast Plasma Instrument (Pollock et al., 2016) data in burst mode for the FTE intervals, and we only focus on events without reconnection (at a central current sheet) for our analyses (events with such sharp central current sheets typically do not fit a coherent flux rope structure). To analyse solar wind conditions preceding the FTEs, we obtain magnetic and velocity fields in GSE coordinates, IMF clock and cone angles in Geocentric Solar Magnetospheric (GSM) coordinates, plasma number density, dynamic pressure, Mach number, and plasma beta from the High-Resolution OMNI database (King & Papitashvili, 2005). The IMF clock and cone angles are defined as $\arctan(B_y/B_z) \in [0^\circ, \pm 180^\circ]$ and $\arccos(B_x/|\mathbf{B}|) \in [0^\circ, 180^\circ]$, respectively.

2.2 Flux rope fitting

We perform a model fitting onto the data using a model first introduced by Burlaga (1988) to describe the magnetic flux rope structure of magnetic clouds in the solar wind (see also Lepping et al. (1990)). The model assumes a cylindrically symmetric and force-free ($\nabla \times \mathbf{B} = \alpha \mathbf{B}$) configuration with a constant α in which the solution satisfying $\nabla^2 B = -\alpha^2 B$ was found by Lundquist (1950). The solution is in terms of the zeroth- and first-order Bessel functions; the axial component is modelled as $B_A = B_0 J_0(\alpha R)$, the tangential (azimuthal) component as $B_T = B_0 H J_1(\alpha R)$, and the radial component as $B_R = 0$, where $H = \pm 1$ is the helicity sign, B_0 is the maximum field strength within

the flux rope interval, and R is the radial distance from the axis. From this model, we obtain a set of fit parameters $(\theta_0, \phi_0, Y_0, H)$ for each flux rope in the local observation frame $(\mathbf{x}_v, \mathbf{y}_v, \mathbf{z}_v)$ coordinates (i.e., the flux rope's frame). To aid understanding, we reproduced an illustration from Burlaga (1988) in Fig S1. The x_v is defined to be opposite to the flux rope motion direction such that $\mathbf{x}_v = -\mathbf{V}_{AV}/|\mathbf{V}_{AV}|$, where $\mathbf{V}_{AV}$ is the average velocity vector across the flux rope. The $\mathbf{z}_v$ is calculated from $\pm\mathbf{x}_v \times \mathbf{n}$, where $\mathbf{n}$ is the normal to the model magnetopause obtained from Shue et al. (1997), the positive (negative) sign is applied when the Y -component of $\mathbf{V}_{AV}$ is positive (negative) to keep the $\mathbf{z}_v$ pointing northward (see Fig S1). Finally, $\mathbf{y}_v = \mathbf{z}_v \times \mathbf{x}_v$ completes the orthonormal system. The angle $\theta_0 \in [-90^\circ, 90^\circ]$ is the angle of the flux rope axis from the ecliptic plane where $\theta_0 = -90^\circ$ is southward and $\theta_0 = 90^\circ$ is northward. The angle $\phi_0 \in [0, \pm 180^\circ]$ is the angle of the flux rope axis from the Sun-Earth line where the positive angle is duskward and negative angle is dawnward, and Y_0 is the impact parameter, which is set to range from $-2R_E$ to $2R_E$. The sense of helicity H is $+1$ for right-handed (RH) or -1 for left-handed (LH) flux ropes. This handedness corresponds to the sense of rotation of the azimuthal (tangential) field around the flux rope axis: the rotation is anti-clockwise for RH flux ropes and it is clockwise for LH flux ropes when viewed from above (i.e. the axial field is pointing towards you). The four parameters are fitted onto the data by trial and error. An optimised set of parameters yield the minimum value of χ^2 defined as $\chi^2 = \sum_i (|\mathbf{B}_{data,i} - \mathbf{B}_{model,i}|^2)/N$ where N is the number of vectors of magnetic field measurements. Examples of the model fitting results are shown in Fig 2.

Each flux rope is fitted for both helicity signs. The sign of helicity is then manually chosen based on visual inspection and comparison of the χ^2 values of the two cases. Among the 186 flux ropes, we found that not all of them can be fitted well to the model, plausibly due to the fact that those flux ropes are not totally force-free. Also, since we will investigate the solar wind conditions preceding these events, we exclude events for which OMNI data are missing. We select 84 flux ropes that are well fitted to the model based on visual inspection (i.e., low χ^2 value). Note that all events are in the northern winter hemisphere (September - February) due to the MMS orbit that samples data near the subsolar region during this time of year. Table S1 lists the time intervals of these flux ropes along with their fit parameters $(\theta_0, \phi_0, Y_0, H)$ and χ^2 .

3 Event illustrations and Statistical Analyses

3.1 Event illustrations

Fig 2 shows examples of LH (left) and RH (right) flux ropes, observed by MMS1 on 23 January 2016 at 23:45 UT and 29 January 2017 at 1:57 UT, respectively. MMS1 was located at $[7.3, -9.4, -1.1]_{GSE} R_E$ for the first event and $[9.9, -5.5, 1.2]_{GSE} R_E$ for the second event. The average 15 minutes of IMF clock angles preceding the first and second events are $-162^\circ \pm 3$ and $114^\circ \pm 4$, respectively (i.e. southward). The bipolar magnetic variations are seen in the B_X component (as expected for a magnetopause normal orientation) while the enhanced core field is seen in B_Y component for both events (Figs 1a, 1a') in the GSE coordinates. However, the senses of rotation of B_X are opposite in each case. To move to the flux rope's frame, we obtain the $(\mathbf{x}_v, \mathbf{y}_v, \mathbf{z}_v)$ coordinates as described in Section 2.2 (see also Fig S1) from the average ion bulk velocity during the flux rope intervals in Figs 2b, 2b', bounded by the vertical dotted lines. The magnetic fields are then transformed to this local observation frame and normalised with the maximum magnetic field strength, called (B_{xv}, B_{yv}, B_{zv}) , in Figs 1c, 1c' for the purpose of fitting into the model. The fitting results to the Burlaga model are plotted as dashed lines in Figs 1c, 1c', along with the fit parameters in text in the same panels. The model fitting in Fig 2c shows that it has a negative helicity (LH) while in Fig 1c' it has a positive helicity (RH); the opposite sense of twist is seen in B_{yv} component. Figs 2d (2d'), 2e (2e'), and 2f (2f') show variations in ion number density, ion temperature, and plasma, magnetic, and total pressure across the two flux rope intervals, respectively.

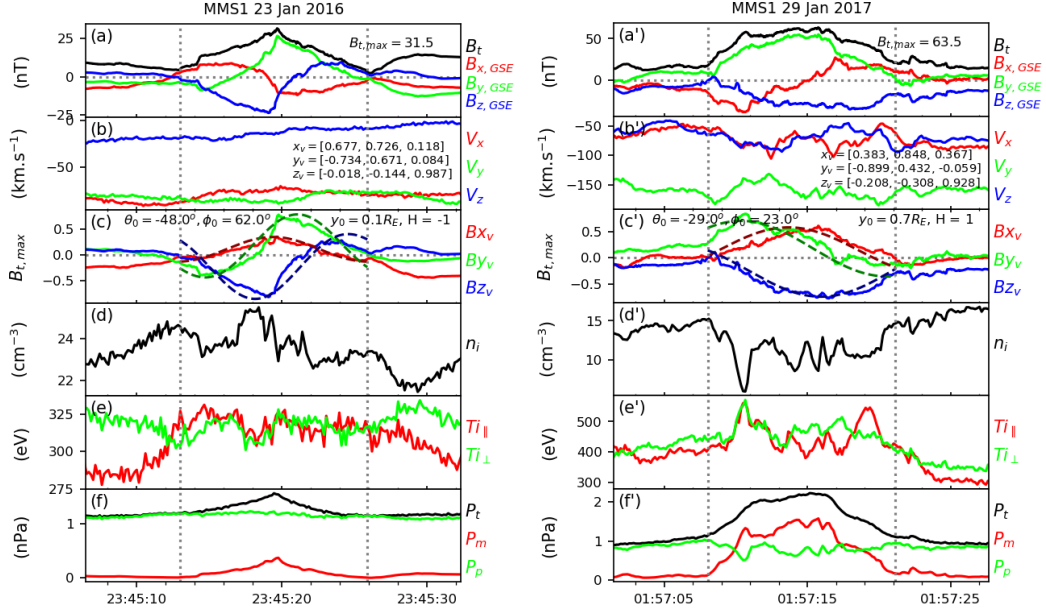


Figure 2. Overview of FTEs with (a-f) LH and (a'-f') RH flux rope structures. (a,a') Magnetic fields in GSE coordinates; (b, b') ion velocity in GSE coordinates with the transformations in text for the local frame coordinates; (c, c') magnetic fields in the local frame coordinates with the model parameters (θ_0, ϕ_0, Y_0, H) and flux rope model; (d, d') ion number density; (e, e') ion temperature; and (f, f') magnetic pressure, plasma pressure, and the total pressure.

3.2 Spatial distribution of FTEs

Among the 84 FTE flux ropes, we found that there are 59 (70%) RH flux ropes and 25 (30%) LH flux ropes. Fig 3 shows the spatial distribution of FTE locations in the X-Z and X-Y planes in the GSE coordinates for the RH (blue cross) and LH (orange cross) flux ropes. As seen in Fig 3, the positive (RH) and negative (LH) helicity flux ropes uniformly distribute on the dayside magnetopause with their positions being at low-latitudes. In other words, there is no spatial preference for FTE flux ropes' handedness. This suggests that the sense of twist is not related to these local properties but should be associated with remote or upstream parameters.

It is important to note that the handedness is different from the sequence of polarity of the bipolar variation of FTEs that is observed dependent on the hemisphere (e.g., Rijnbeek et al., 1984; D. Southwood et al., 1986). The bipolar variation is observed to be outward followed by inward to the magnetopause for spacecraft located in the northern hemisphere; this order is reversed for spacecraft located in the southern hemisphere. Both sequences can have the same helicity sign as the order of sequence depends on the spacecraft trajectory; the direction of the core field with respect to the bipolar variation is what determines the helicity sign of the flux rope.

3.3 Solar wind IMF conditions

We studied solar wind conditions using averages over interval 15 minutes prior to the events, for all FTE flux ropes. Figs 4a and 4b show histograms of the averaged IMF clock angle 15 minutes before the two types of events (RH and LH). It is found that the RH flux ropes are mainly preceded by IMF clock angles from 90° to 180° , which is the

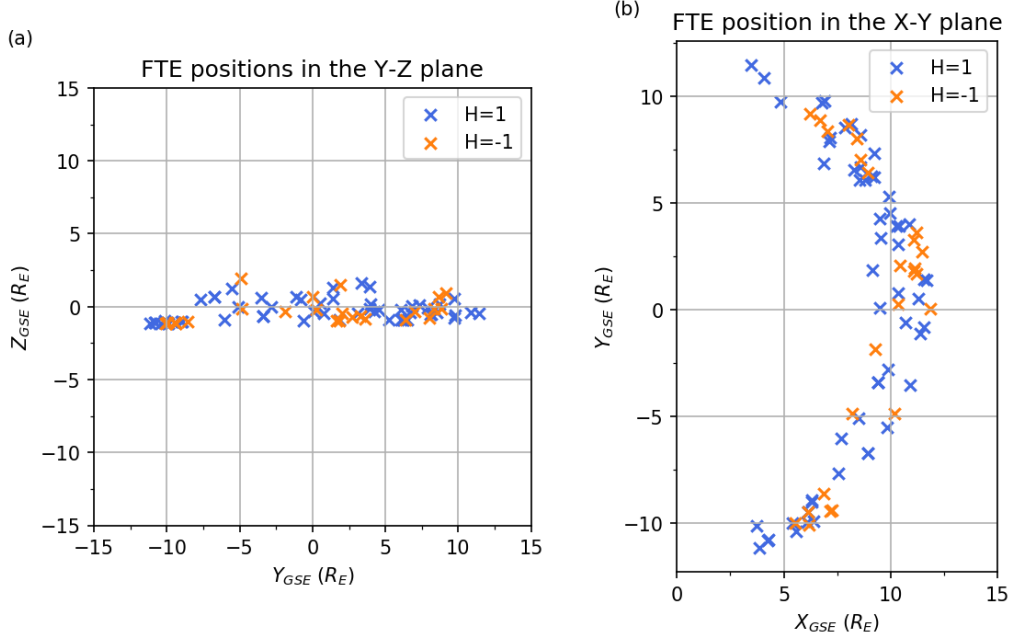


Figure 3. Spatial distribution of the FTE flux ropes in the (a) Y-Z and (b) X-Y planes with data points color-coded by the helicity signs $H = 1$ (RH) in blue and $H = -1$ (LH) in orange. Both types of flux rope uniformly distribute on the dayside magnetopause.

duskward and southward direction. In contrast, the LH flux ropes are mainly preceded by IMF clock angles of -90° to -180° , which is dawnward and southward. Note that the averages over 10, 20, and 25 minutes give similar results. We do not find any clear correlation between other solar wind parameters and the helicity sign of flux ropes (see Fig S2). This suggests that the helicity sign of FTE flux ropes is mainly controlled by the IMF clock angle (e.g., the IMF B_Y component). Nevertheless, there are clearly some LH flux ropes (9 out of 25 cases) that are preceded by duskward IMF (IMF $B_Y > 0$) and some RH flux ropes (5 out of 59 cases) that are preceded by dawnward IMF (IMF $B_Y < 0$). The smaller outlier population for the RH flux ropes may relate to the seasonal (i.e., the dipole tilt) effects because all events are observed during September - February which are near the winter solstice (December). Note that not all of the outlier events are preceded by southward IMF; a significant population of the outlier RH (4 out of 5) and LH (3 out of 9) flux ropes are preceded by northward IMF. Excluding the northward IMF events, it is still unclear whether this outlier group is due to statistical uncertainties, such as IMF propagation errors, or to some unknown physical mechanism controlling the helicity in addition to the IMF B_Y .

For the purpose of discussion and further analyses, we define two groups of LH flux ropes to be (1) “regular” for LH flux ropes that are preceded by dawnward IMF and (2) “outlier” for LH flux ropes that are preceded by duskward IMF. Note that we study only the LH flux group due to the significant outlier population. Comparing the solar wind conditions between the two groups, we found that they have different IMF cone angle as shown in Figs 4c, 4d. The regular group is mostly preceded by IMF cone angle $< 90^\circ$ (sunward) while the outlier group is mainly preceded by IMF cone angle $> 90^\circ$ (anti-sunward). The magnitudes of IMF $B_x/|\mathbf{B}|$ are also different (see panel (c) of Fig S3). The outlier group has mostly negative IMF $B_x/|\mathbf{B}|$ values and large magnitude. How-

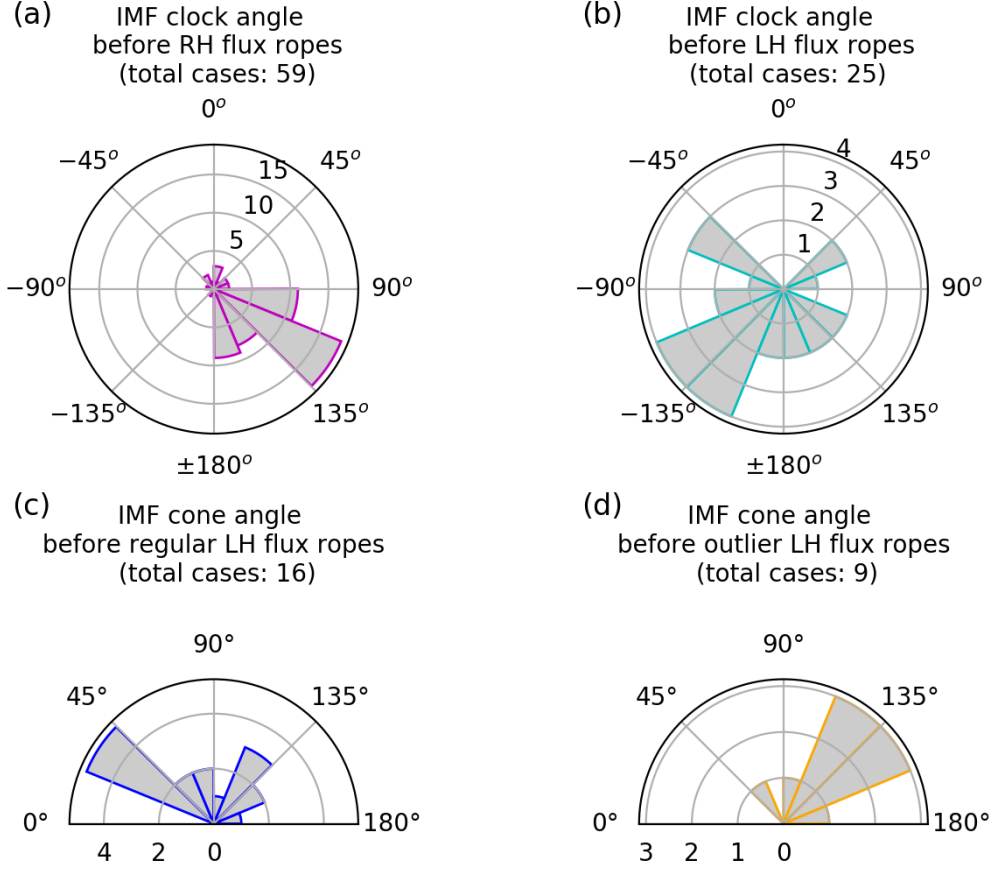


Figure 4. (top) Distribution of the averaged IMF clock angle 15 minutes before the FTE observations obtained from the OMNI database for (a) RH and (b) LH flux ropes. The IMF clock angle is mainly in the 90° to 180° clock angle range (duskward-southward) before the RH flux ropes and in the -90° to -180° range (dawnward-southward) before the LH flux ropes. (bottom) Distribution of the averaged IMF cone angle 15 minutes before LH flux ropes for (c) regular group and (d) outlier group. The regular LH flux ropes are mostly preceded by sunward-tilted IMF B_x while the outlier LH flux ropes are mainly preceded by antisunward-tilted IMF B_x .

ever, the regular group has a weak, positive IMF $B_x/|\mathbf{B}|$. There is no significant difference in other solar wind parameters between these groups.

4 Discussion

We have analysed the helicity sign of 84 FTE flux ropes observed by MMS1 (from a list by Fargette et al. (2020)) at the dayside magnetopause through model fitting. We found that there are 59 RH (70%) and 25 (30%) LH flux ropes. We also analysed the solar wind conditions preceding the events. We found a correlation between the IMF B_Y sign and the helicity sign: RH flux ropes ($H = 1$) are mainly preceded by IMF $B_Y > 0$ while LH flux ropes ($H = -1$) are mostly preceded by IMF $B_Y < 0$. This shows that the twist direction of the FTE flux ropes is controlled by the IMF B_Y component.

Our main results place constraints on the FTE generation mechanism. Indeed, as we illustrate in Fig 1, the helicity sign of an FTE can be predicted as a function of the IMF B_Y in the context of a multiple, sequential X-line formation mechanism. The sense of rotation of the azimuthal field of the flux ropes from our statistical analyses can be explained by this picture where the FTEs are generated by multiple component reconnection X-lines on the dayside magnetopause as predicted by Lee and Fu (1985). Additionally, most events are observed near the winter solstice, i.e., when the dipole tilt is large, consistent with the FTE production due to sequential, multiple X-line mechanism proposed by Raeder (2006). Indeed, the Maximum Magnetic Shear Model (Trattner et al., 2007) predicted that the component reconnection should be dominant on the dayside when the southward IMF has a significant B_Y component because the draped IMF in the magnetosheath region makes a first contact with the subsolar region. The FTEs should be generated from this region through multiple X-line reconnection where the sign of IMF B_Y across the neutral line controls the sense of twist and core field as depicted in Fig 1. Even though our finding does not rule out other FTE formation mechanisms, it is consistent and compatible with the multiple X-line mechanism.

Statistical analyses on the solar wind conditions prior to the LH flux ropes preceded by IMF $B_Y > 0$ (outlier) and those preceded by IMF $B_Y < 0$ (regular) show that the outlier group has a strong, negative IMF B_X while the regular group has a small, positive IMF B_X . It is unclear whether why the magnitude and polarity of IMF B_X should control the helicity sign of FTE flux ropes. When the IMF is due south and B_X is negative, the magnetic merging line is found to shift southward at the dayside (Peng et al., 2010). In addition, with all our events being in the winter hemisphere, there plausibly be a combined effect between the IMF B_X and the dipole tilt (e.g., Palmroth et al., 2012; Hoilijoki et al., 2014) that can complicate reconnection at the dayside and thus the FTE formation. The IMF B_X component was found to impact the north-south hemispheric asymmetry of FTE occurrence, properties, sizes, and motions (Hoilijoki et al., 2019) as a result of a reduction of the reconnection rate at the dayside due to the smaller tangential magnetic field to the magnetopause. The FTE generation may also be complicated by processes downstream of bow shock when the IMF cone angle is small. We leave this as an open question that should be addressed in future work.

5 Conclusions

We have surveyed the helicity sign of 84 FTE flux ropes observed by MMS near the winter solstice during years 2015 – 2017 that can be fitted well to a cylindrically force-free flux rope model with a constant α (Burlaga, 1988). We found that 59 (70%) flux ropes are RH and 25 (30%) of them are LH. Investigations of the IMF conditions show that the RH flux ropes are mainly preceded by southward IMF with positive B_Y while the LH flux ropes are mostly preceded by southward IMF with negative B_Y . This control of FTE flux rope helicity sign by the IMF B_Y component is consistent with its formation through sequential, multiple X-line reconnection. We also found an outlier group

of flux ropes whose helicity sign is inconsistent with the IMF B_Y sign. There are 14 out of 84 flux ropes that are preceded by unexpected IMF B_Y polarity. Investigation of the solar wind conditions preceding LH flux ropes show that the outlier group is associated with strong and negative IMF B_X . This shows that the presence of IMF B_X further complicates the formation of FTE flux ropes at the dayside magnetopause. Future work would be desirable for a fuller understanding of FTE helicity generation of this outlier group.

Acknowledgments

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Statistical Relationship between Interplanetary Magnetic Field conditions and the Helicity of Flux Transfer Event Flux Ropes

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Table S1

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Introduction

This supplementary information includes details of the FTE flux rope fit parameters and detailed analyses of the solar wind conditions preceding the events. The FTEs are listed in Table S1 with their beginning and end time. Table S1 includes the flux rope model fit parameters as introduced in Section 2.2 in the main text. The fit parameters are obtained from the cylindrically symmetric force-free flux rope model with a constant α as in Burlaga (1988). The model fit parameters consist of two angles (θ_0, ϕ_0) for characterizing the orientation of the flux rope axis, the impact parameter y_0 , and the helicity sign $H = \pm 1$, where $H=1$ is right-handed (RH) and $H=-1$ is a left-handed (LH) flux ropes. We illustrate the geometry of magnetic flux rope for the purpose of fitting into the model in Fig S1, adapted from Burlaga (1988). The quality of the model fit to the data is indicated by the deviation of the observed data from the model and

defined as $\chi^2 = \sum_i (|\mathbf{B}_{data,i} - \mathbf{B}_{model,i}|)^2 / N$ where N is the number of vectors of magnetic field measurements.

The analyses of the solar wind conditions preceding all of the FTEs are shown in Fig S2. Using the model fitting, the FTEs are categorized into RH and LH flux ropes. There are 59 RH and 25 LH flux ropes. Fig S2 shows the normalized distributions of the averaged solar wind conditions 15 minutes before the FTEs. It is clearly seen that the IMF clock angle is negative (e.g., IMF $B_Y < 0$) for LH flux ropes while the IMF clock angle is positive (e.g., IMF $B_Y > 0$) for RH flux ropes. The distributions of other solar wind parameters do not show significant difference between the two.

For the LH flux ropes, we found that there are 9 out of 25 events that are preceded by duskward IMF. As per the scenario, described in the paper, wherein the sign of helicity is controlled by the IMF B_Y sign which fixes the topology of the magnetic field in-between the sequential X-lines forming FTEs, such cases constitute outliers. In search for a controlling factor, we analyze the solar wind conditions preceding the events specifically for LH flux ropes only. Fig S3 shows normalized distributions of the average solar wind conditions 15 minutes before the LH flux ropes. The LH flux ropes are divided into the regular group (16 events) that are preceded by IMF $B_Y < 0$ and the outlier group (9 events) that are preceded by IMF $B_Y > 0$. It can be seen that the IMF cone angle (panel (c)) between two groups are different. The ratio of IMF $B_x/|\mathbf{B}|$ for the outlier group is mainly negative while the IMF B_x of the normal group is mostly positive. The magnitude of the IMF $B_x/|\mathbf{B}|$ for the outlier group is also stronger. The distributions of other solar wind parameters show slight differences between the two groups, which may be enhanced for a larger data set.

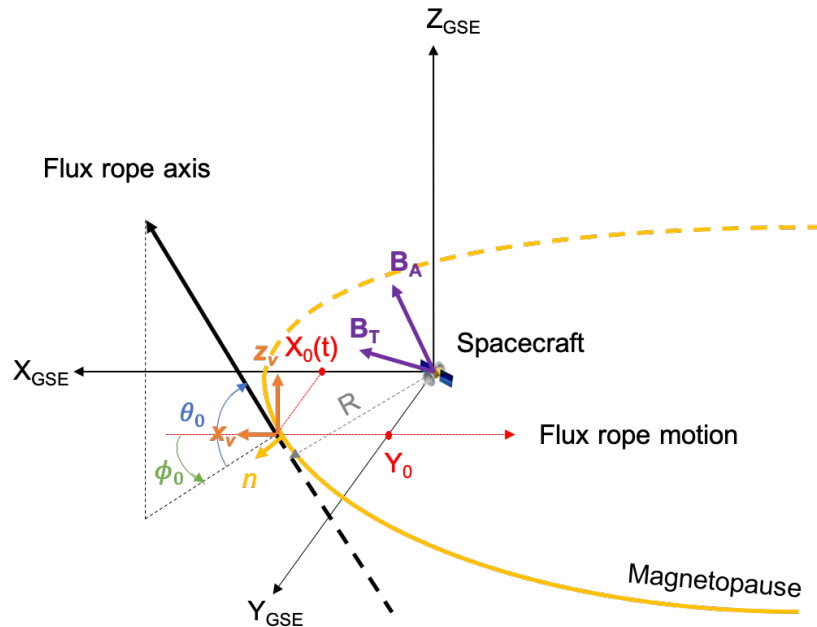


Fig S1. The geometry of magnetic flux rope passing a spacecraft in which we model following Burlaga (1988), adapted from Figure 2 of Burlaga (1988).

Normalized Distributions of Average Solar Wind Conditions (15 min) before RH and LH Flux Ropes

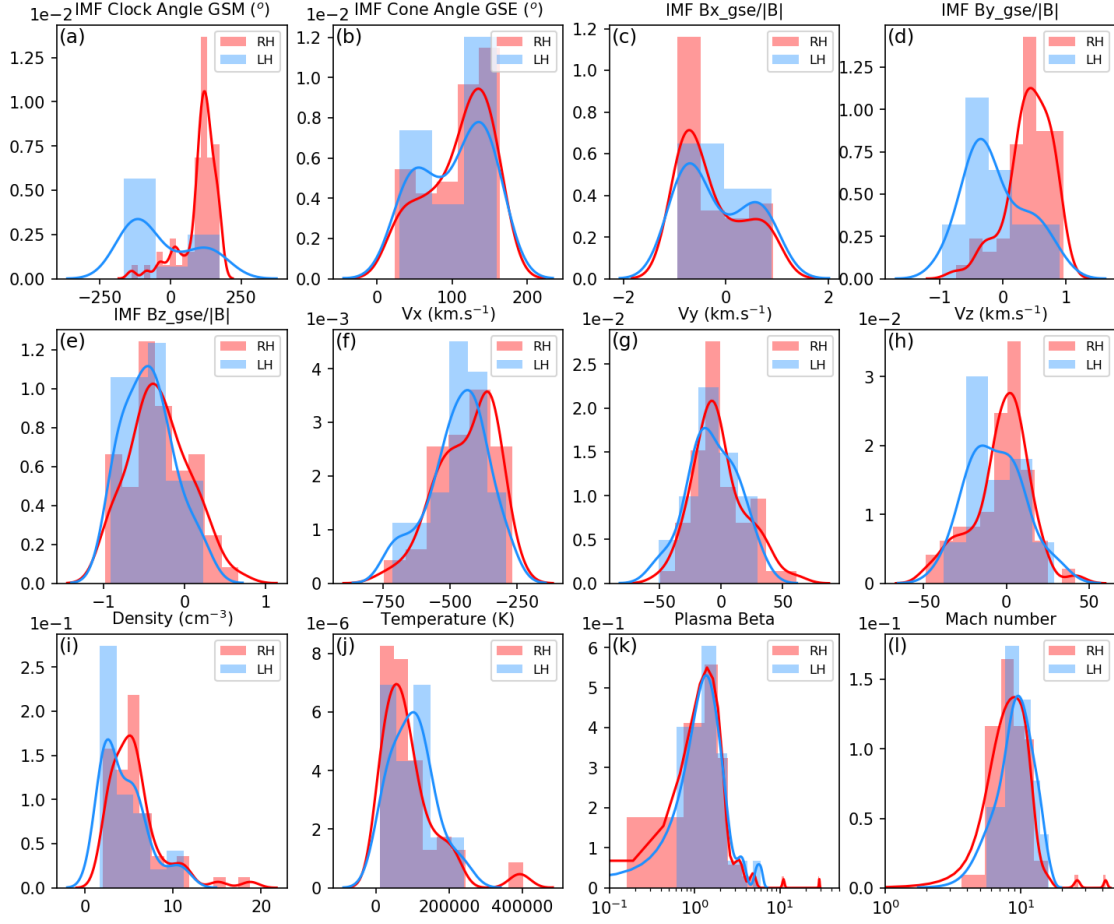


Fig S2. Normalized distributions of the average solar wind conditions 15 minutes before the right-handed (RH) and left-handed (LH) flux ropes. The solar wind parameters are shown as the following. (a) IMF clock angle in the GSM coordinate system. (b) IMF cone angle in the GSE coordinate system. (c,d,e) IMF B_x , B_y , and B_z in the GSE coordinates normalized by the IMF magnitude. (f,g,h) Ion bulk flow velocity V_x , V_y , and V_z components in the GSE system. (i) Ion number density. (j) Ion temperature. (k) Plasma beta. (l) Alfvén Mach Number.

Normalized Distributions of Average Solar Wind Conditions (15 min) before LH Flux Ropes

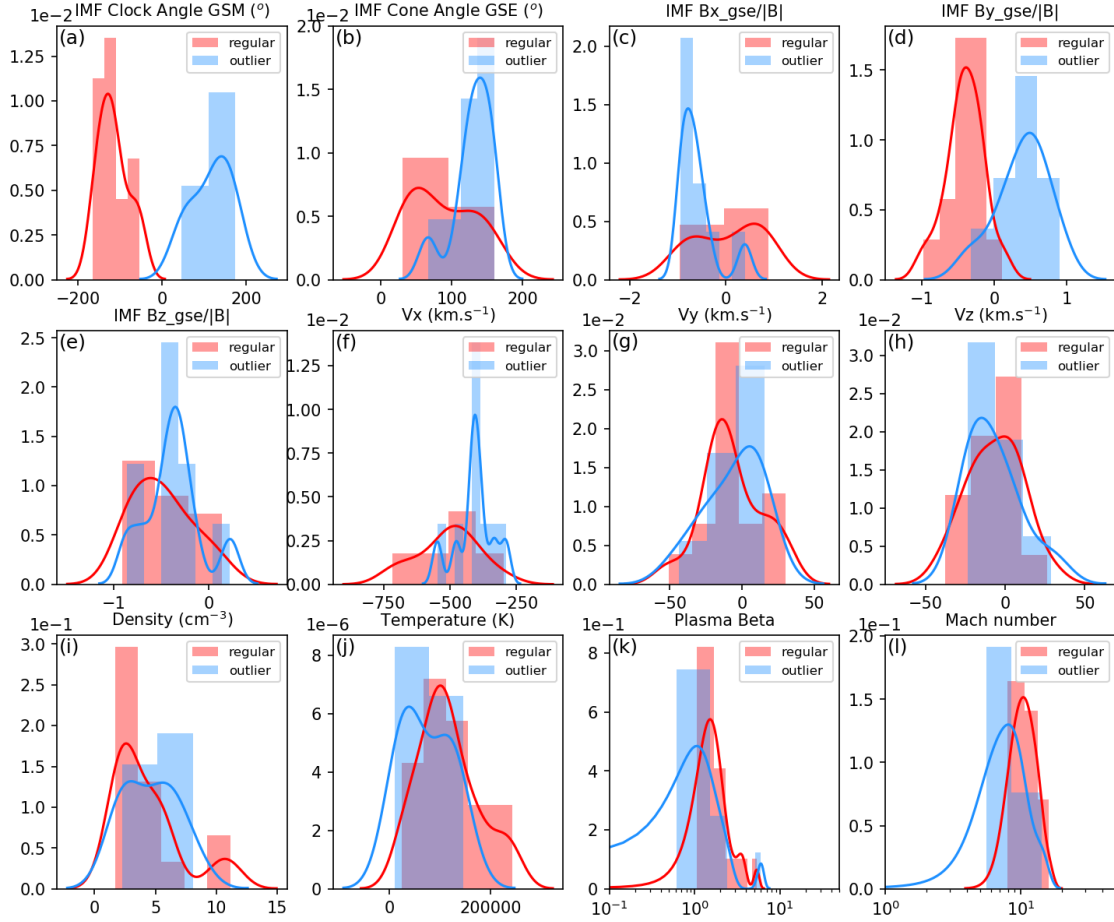


Fig S3. Normalised distributions of the average solar wind conditions 15 minutes for the LH flux ropes that are preceded by IMF $B_y < 0$ (regular group) and IMF $B_y > 0$ (outlier group). The panel format is the same as in Figure S1 except for (b,c,d) that are the IMF B_x , B_y , and B_z components in the GSE coordinates normalized by the total IMF magnitude. (f,g,h) Ion bulk flow velocity V_x , V_y , and V_z components in the GSE system. (i) Ion number density. (j) Ion temperature. (k) Plasma beta. (l) Alfvén Mach Number.

Table S1. List of FTE flux ropes observed by MMS1 with their model fit parameters. The start and end times delineate the FTE intervals based on the bipolar component variation and enhanced core field signatures. The fit parameters (θ_0 , ϕ_0 , Y_0 , H) are obtained from optimization of Burlaga (1988)'s cylindrically symmetric force-free flux rope model with a constant α . The χ^2 is the deviation of the observational data from the most optimized modeled flux rope data, normalized by the number of data points in each event.

Event	Start time	End time	θ_0 (deg)	ϕ_0 (deg)	Y_0 (R_E)	H	χ^2
1	2015-09-14 16:06:51.000	2015-09-14 16:07:33.000	38.0	-175.0	1.0	1	0.0033
2	2015-09-23 08:15:03.000	2015-09-23 08:15:36.000	18.0	171.0	-0.6	1	0.0038
3	2015-09-23 10:56:28.000	2015-09-23 10:56:58.000	-26.0	14.0	-0.1	-1	0.0033
4	2015-09-25 09:24:22.793	2015-09-25 09:24:52.977	41.0	147.0	-0.5	1	0.0053
5	2015-09-25 09:57:43.000	2015-09-25 09:58:28.000	-20.0	-38.0	1.0	-1	0.0037
6	2015-10-03 10:45:51.000	2015-10-03 10:46:34.000	17.0	179.0	1.0	1	0.0037
7	2015-10-03 13:27:12.396	2015-10-03 13:27:42.044	14.0	177.0	0.7	1	0.0034
8	2015-10-06 17:28:07.567	2015-10-06 17:28:17.015	50.0	8.0	-1.0	1	0.0081
9	2015-10-08 07:41:32.693	2015-10-08 07:41:59.549	20.0	174.0	1.0	1	0.0046
10	2015-10-11 11:05:13.000	2015-10-11 11:05:57.000	29.0	173.0	1.0	1	0.0046
11	2015-10-20 06:16:03.000	2015-10-20 06:17:26.000	41.0	180.0	-0.5	1	0.0029
12	2015-10-22 13:40:19.000	2015-10-22 13:40:48.000	-17.0	-6.0	-0.8	-1	0.0034
13	2015-10-31 05:45:52.000	2015-10-31 05:46:27.000	19.0	173.0	0.9	1	0.0037
14	2015-11-05 04:58:41.290	2015-11-05 04:58:51.487	-37.0	-19.0	-1.0	-1	0.0079
15	2015-11-05 14:07:07.131	2015-11-05 14:07:44.639	-45.0	16.0	-1.0	-1	0.0048
16	2015-11-05 14:36:39.263	2015-11-05 14:36:44.695	35.0	-180.0	1.0	1	0.0093
17	2015-11-06 06:57:42.000	2015-11-06 06:58:26.000	-21.0	-4.0	-1.0	-1	0.0033
18	2015-11-06 13:23:48.000	2015-11-06 13:24:29.000	33.0	171.0	-0.3	1	0.0055
19	2015-11-06 13:26:17.915	2015-11-06 13:26:32.647	31.0	6.0	-1.0	1	0.0066
20	2015-11-08 14:02:51.000	2015-11-08 14:03:23.000	20.0	-170.0	1.0	1	0.0034
21	2015-11-09 10:06:54.162	2015-11-09 10:07:02.160	36.0	112.0	0.1	1	0.0128

22	2015-11-10 02:43:43.908	2015-11-10 02:44:10.768	30.0	152.0	-0.8	1	0.0061
23	2015-11-11 03:56:21.000	2015-11-11 03:57:18.000	20.0	-176.0	0.4	1	0.0028
24	2015-11-12 07:06:01.695	2015-11-12 07:06:06.617	-60.0	-180.0	-0.9	-1	0.013
25	2015-11-12 07:20:20.201	2015-11-12 07:20:35.685	-46.0	-21.0	-1.0	-1	0.0055
26	2015-12-02 10:00:42.427	2015-12-02 10:01:26.951	-15.0	-10.0	0.9	-1	0.0036
27	2015-12-02 10:21:35.000	2015-12-02 10:21:59.000	-12.0	-7.0	1.0	-1	0.003
28	2015-12-05 00:40:35.058	2015-12-05 00:40:43.465	54.0	38.0	0.7	1	0.0113
29	2015-12-05 00:40:50.216	2015-12-05 00:41:04.843	24.0	7.0	-1.0	1	0.0078
30	2015-12-08 10:30:06.246	2015-12-08 10:30:17.494	40.0	-154.0	0.1	1	0.0096
31	2015-12-11 12:23:27.054	2015-12-11 12:23:35.626	32.0	-175.0	1.0	1	0.0098
32	2015-12-19 09:27:02.581	2015-12-19 09:27:18.115	45.0	-180.0	0.1	1	0.006
33	2016-01-23 23:26:20.490	2016-01-23 23:26:37.970	-13.0	-75.0	-1.0	-1	0.0068
34	2016-01-23 23:45:12.961	2016-01-23 23:45:25.871	-48.0	62.0	0.1	-1	0.0055
35	2016-01-27 22:17:19.000	2016-01-27 22:18:02.000	-24.0	3.0	0.8	1	0.0023
36	2016-01-27 22:49:42.164	2016-01-27 22:49:48.419	-30.0	-38.0	0.4	1	0.01
37	2016-01-29 22:38:50.825	2016-01-29 22:38:59.880	-37.0	77.0	-1.0	-1	0.0142
38	2016-01-31 05:54:46.000	2016-01-31 05:55:47.000	-30.0	4.0	1.0	1	0.0046
39	2016-02-01 22:26:45.004	2016-02-01 22:26:51.898	54.0	-175.0	0.1	1	0.008
40	2016-02-07 03:06:51.154	2016-02-07 03:07:33.233	-11.0	3.0	1.0	-1	0.0027
41	2016-02-11 01:56:07.266	2016-02-11 01:56:26.135	-34.0	-7.0	1.0	-1	0.0088
42	2016-02-11 01:56:30.049	2016-02-11 01:56:39.605	14.0	-45.0	-0.4	1	0.0076
43	2016-02-11 01:57:06.171	2016-02-11 01:57:16.309	25.0	-96.0	-1.0	1	0.0108
44	2016-02-11 02:00:20.668	2016-02-11 02:00:44.815	-20.0	134.0	-0.7	1	0.0063
45	2016-02-11 02:39:04.000	2016-02-11 02:40:45.000	-30.0	33.0	1.0	1	0.0046
46	2016-02-11 02:46:21.943	2016-02-11 02:46:53.569	-34.0	25.0	-0.7	1	0.004

47	2016-02-14 01:25:36.846	2016-02-14 01:25:50.260	-60.0	70.0	-1.0	-1	0.0092
48	2016-02-15 01:28:58.000	2016-02-15 01:29:49.000	-31.0	-5.0	-0.1	1	0.003
49	2016-02-19 22:55:42.433	2016-02-19 22:55:53.393	27.0	-125.0	1.0	1	0.0115
50	2016-02-19 23:53:43.419	2016-02-19 23:54:06.796	-26.0	-51.0	0.8	1	0.0047
51	2016-02-19 23:54:35.258	2016-02-19 23:54:58.486	-35.0	-48.0	-1.0	1	0.0082
52	2016-02-28 00:58:54.553	2016-02-28 00:59:16.930	24.0	-10.0	-1.0	1	0.0058
53	2016-10-10 15:43:20.000	2016-10-10 15:43:54.000	18.0	172.0	0.0	1	0.003
54	2016-10-27 12:00:43.000	2016-10-27 12:01:09.000	61.0	-172.0	-0.1	1	0.0056
55	2016-11-06 16:52:52.000	2016-11-06 16:53:30.000	32.0	5.0	-1.0	1	0.006
56	2016-11-08 10:49:25.000	2016-11-08 10:50:24.000	70.0	3.0	-1.0	1	0.0045
57	2016-11-08 11:19:54.000	2016-11-08 11:20:42.000	-34.0	-1.0	-1.0	-1	0.004
58	2016-11-08 13:55:34.000	2016-11-08 14:01:03.000	-18.0	-10.0	-0.4	1	0.0012
59	2016-11-12 18:19:44.000	2016-11-12 18:20:11.000	16.0	178.0	-1.0	-1	0.0044
60	2016-11-15 12:20:40.201	2016-11-15 12:20:49.706	31.0	174.0	1.0	1	0.0104
61	2016-11-15 15:49:46.000	2016-11-15 15:50:07.000	19.0	174.0	-1.0	-1	0.0049
62	2016-11-23 09:03:15.000	2016-11-23 09:03:45.000	1.0	164.0	0.6	1	0.0053
63	2016-11-27 08:39:08.000	2016-11-27 08:40:05.000	28.0	-174.0	1.0	1	0.004
64	2016-12-02 09:30:09.603	2016-12-02 09:30:19.524	33.0	-177.0	-0.8	1	0.0082
65	2016-12-14 05:30:40.000	2016-12-14 05:31:07.000	9.0	1.0	0.9	-1	0.0025
66	2016-12-19 07:42:04.000	2016-12-19 07:43:34.000	13.0	-180.0	-0.1	1	0.0026
67	2016-12-19 09:15:40.000	2016-12-19 09:17:46.000	6.0	-180.0	-1.0	1	0.0013
68	2016-12-19 13:54:51.000	2016-12-19 13:56:05.000	36.0	175.0	1.0	1	0.004
69	2016-12-23 03:16:48.000	2016-12-23 03:17:08.000	51.0	-128.0	0.9	-1	0.0093
70	2016-12-26 14:50:36.000	2016-12-26 14:52:36.000	16.0	-171.0	-0.8	1	0.002
71	2016-12-27 08:02:30.190	2016-12-27 08:02:35.208	29.0	-180.0	0.1	-1	0.0135

72	2016-12-28 06:32:50.089	2016-12-28 06:32:57.921	32.0	-180.0	1.0	1	0.0135
73	2016-12-29 11:12:00.000	2016-12-29 11:12:26.000	26.0	1.0	-0.8	1	0.0058
74	2016-12-29 12:18:11.000	2016-12-29 12:19:03.000	-20.0	-1.0	-0.9	-1	0.0036
75	2017-01-01 03:01:02.993	2017-01-01 03:01:13.648	-33.0	-32.0	-0.6	1	0.0061
76	2017-01-01 06:27:39.512	2017-01-01 06:27:41.575	30.0	-162.0	-0.1	1	0.0186
77	2017-01-11 04:22:47.969	2017-01-11 04:22:59.126	20.0	155.0	0.5	1	0.0094
78	2017-01-13 00:58:21.296	2017-01-13 00:58:47.966	22.0	-175.0	-1.0	-1	0.0033
79	2017-01-15 01:11:04.000	2017-01-15 01:11:57.000	-25.0	1.0	1.0	1	0.0032
80	2017-01-29 01:52:19.844	2017-01-29 01:52:32.496	-32.0	3.0	0.3	1	0.0055
81	2017-01-29 01:52:32.578	2017-01-29 01:52:46.879	-31.0	-32.0	-0.1	1	0.0084
82	2017-01-29 01:57:08.049	2017-01-29 01:57:21.139	-29.0	23.0	0.7	1	0.0055
83	2017-02-04 00:14:40.907	2017-02-04 00:14:56.063	-15.0	14.0	0.6	1	0.0071
84	2017-02-04 07:49:16.574	2017-02-04 07:49:21.439	38.0	154.0	1.0	-1	0.0148