Electrostatic Waves and Electron Heating Observed over Lunar Crustal Magnetic Anomalies

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Abstract

Above lunar crustal magnetic anomalies, large fractions of solar wind electrons and ions can be reflected and stream back towards the solar wind flow, leading to a number of interesting effects such as electrostatic instabilities and waves. These electrostatic structures can also interact with the background plasma, resulting in electron heating and scattering. We study the electrostatic waves and electron heating observed over the lunar magnetic anomalies by analyzing data from the Acceleration, Reconnection, Turbulence, and Electrodynamics of Moon's Interaction with the Sun (ARTEMIS) spacecraft. Based on the analysis of two lunar flyby events in 2011 and 2013, we find that the electron two-stream instability (ETSI) and electron cyclotron drift instability (ECDI) may play an important role in driving the electrostatic waves. We also find that ECDI, along with the modified two-stream instability (MTSI), may provide the mechanisms responsible for substantial isotropic electron heating over the lunar magnetic anomalies.

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Key Points:

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10	• Two types of electrostatic instabilities are observed over the lunar crustal mag-
11	netic anomalies during ARTEMIS flyby
12	• Electron two-stream instability and electron cyclotron drift instability may play
13	an important role in driving the electrostatic waves
14	• Electron cyclotron drift instability, along with modified two-stream instability, may
15	cause isotropic electron heating

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16 Abstract

Above lunar crustal magnetic anomalies, large fractions of solar wind electrons and ions 17 can be reflected and stream back towards the solar wind flow, leading to a number of 18 interesting effects such as electrostatic instabilities and waves. These electrostatic struc-19 tures can also interact with the background plasma, resulting in electron heating and scat-20 tering. We study the electrostatic waves and electron heating observed over the lunar 21 magnetic anomalies by analyzing data from the Acceleration, Reconnection, Turbulence, 22 and Electrodynamics of Moon's Interaction with the Sun (ARTEMIS) spacecraft. Based 23 on the analysis of two lunar flyby events in 2011 and 2013, we find that the electron two-24 stream instability (ETSI) and electron cyclotron drift instability (ECDI) may play an 25 important role in driving the electrostatic waves. We also find that ECDI, along with 26 the modified two-stream instability (MTSI), may provide the mechanisms responsible 27 for substantial isotropic electron heating over the lunar magnetic anomalies. 28

²⁹ Plain Language Summary

Without a global magnetic field or a thick atmosphere, the solar wind directly impacts 30 the surface of the Moon. However, over regions where the lunar crust is strongly mag-31 netized, the charged particles in the solar wind can be reflected and travel back towards 32 the incoming solar wind, generating interesting features like electrostatic waves. These 33 waves can also in turn affect the solar wind by increasing the temperature of its charged 34 particles. To understand the mechanisms causing the waves and heating, we analyze data 35 from the Acceleration, Reconnection, Turbulence, and Electrodynamics of Moon's In-36 teraction with the Sun spacecraft. Our results indicate that the lunar environment be-37 comes unstable because of the reflected charged particles, thereby creating free energies 38 that lead to the waves and heating. 39

40 1 Introduction

In the absence of a global magnetic field and a thick atmosphere, unlike the case of the Earth, the surface of the Moon directly interacts with the incident solar wind plasma. Traditionally, the Moon has been thought to act as a simple barrier to the solar wind flow, causing the absorption of plasma at the upstream surface and formation of a plasma wake in the downstream. However, recent observations from Chandrayaan-1, Kaguya, and Chang'E-1 reveal that Moon-solar wind interaction is in fact much more complicated

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and dynamic, capable of creating a variety of interesting effects around the Moon. For 47 example, the surface of the Moon, immersed in the solar wind plasma, charges to an elec-48 trostatic potential in order to balance the total incident currents (Whipple, 1981; J. S. Halekas 49 et al., 2002, 2011). Moreover, solar wind sputtering from the lunar surface and ioniza-50 tion of the tenuous neutral exosphere can produce heavier lunar pickup ions, which can 51 then be accelerated downstream from the Moon by the motional electric field (Yokota 52 et al., 2009; J. S. Halekas, Poppe, Delory, et al., 2012; Cao et al., 2020). Some other ex-53 amples of lunar interaction include backscattering of solar wind ions and photoelectron 54 emission from the lunar surface (Reasoner & Burke, 1972; Goldstein, 1974; Lue et al., 55 2014; Bhardwaj et al., 2015; Harada et al., 2017). 56

One of the most interesting and unique Moon-solar wind interactions happens over 57 the lunar crustal magnetic anomalies. Previous studies have shown that the lunar crustal 58 magnetic fields can perturb the solar wind flow, causing plasma compressions at the lu-59 nar limb (Russell & Lichtenstein, 1975). More recent measurements from Kaguya sug-60 gest that "mini-magnetospheres" can form over strong magnetic anomaly regions, par-61 tially shielding the lunar surface from the solar wind (Saito et al., 2010). In addition, 62 local crustal magnetic fields are found capable of reflecting solar wind ions and electrons 63 from the lunar surface (Lue et al., 2011; Saito et al., 2012; J. Halekas et al., 2012; J. S. Halekas, 64 Poppe, Farrell, et al., 2012). Using Chandrayaan-1 data, Lue et al. (2011) reported that 65 on average 10% of the incident solar wind ions reflect over large-scale magnetic anoma-66 lies. The reflection efficiency can reach up to 50% for ions and as much as 100% for elec-67 trons above regions of strongest crustal fields (J. S. Halekas et al., 2001). The magnet-68 ically reflected ions and electrons, along with backscattered particles and photoelectrons, 69 can then form counter-streaming beams towards the incoming solar wind flow, result-70 ing in a number of fundamental plasma processes such as electrostatic instabilities and 71 waves. These electrostatic structures can also in turn have an impact on the lunar plasma 72 environment, leading to substantial electron heating and scattering. 73

A variety of plasma instabilities and waves of different origin have been previously observed above the lunar crustal magnetic anomalies (Nakagawa, 2016; Harada & Halekas, 2016). Tsugawa et al. (2011) reported that monochromatic, left-hand polarized (in the spacecraft frame) whistler waves with frequencies close to 1 Hz were detected by Kaguya near the Moon. A further statistical analysis suggested that the waves were generated by the solar wind interaction with lunar magnetic anomalies. In addition, broadband elec-

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trostatic mode, resulting from counter-streaming electron beams, is another type of waves
commonly observed in the lunar upstream plasma (Harada et al., 2014).

In this paper, we investigate two types of electrostatic instabilities observed over 82 the lunar crustal magnetic anomalies by Acceleration, Reconnection, Turbulence, and 83 Electrodynamics of Moon's Interaction with the Sun (ARTEMIS) spacecraft. We report 84 for the first time on a class of electrostatic waves propagating perpendicular to the am-85 bient magnetic field, possibly driven by electron cyclotron drift instabilities. This type 86 of electrostatic waves is analogous to those observed in the foot region of perpendicu-87 lar shocks. In the end, we also discuss the mechanisms of electron heating observed along 88 with the electrostatic waves over the magnetic anomalies. 89

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2 Instrumentation and Observations

NASA's ARTEMIS spacecraft, consisting of two satellites (P1 and P2) originally 91 from the THEMIS (Time History of Events and Macroscale Interactions During Sub-92 storms) mission, occupies stable 26-h period elliptical near-equatorial orbits around the 93 Moon (Angelopoulos, 2011). To investigate the plasma environment above the dayside 94 lunar surface, we utilize measurements from four of the instruments: the Electrostatic 95 Analyzer (ESA; McFadden et al., 2008), Electric Field Instrument (EFI; Bonnell et al., 96 2008), Search Coil Magnetometer (SCM; Roux et al., 2008), and Fluxgate Magnetome-97 ter (FGM; Auster et al., 2008). The ESA measures electron energies in the range of 2 98 eV to 32 keV and ion energies from 1.6 eV to 25 keV (McFadden et al., 2008). The EFI 99 is capable of measuring the three components of the ambient electric fields from ~ 10 100 Hz up to 8 kHz (Bonnell et al., 2008). 101

We select two events best suited for studying the electrostatic waves and electron 102 heating over the magnetic anomalies, where the observations can be made at altitudes 103 below 50 km above the lunar surface. On 26 November 2011, ARTEMIS P2 flew by the 104 Moon at average GSE coordinates of [53, 16, 0] earth radii (R_E) , located in the solar wind 105 well upstream of the Earth's bow shock. The data of the flyby obtained from the above 106 four instruments are shown in Figure 1. The probe is found to briefly fly over the mag-107 netic anomaly region between 10:10 UT and 10:12 UT, indicated by an enhancement of 108 the fluctuations in the ambient magnetic field in Figure 1b. When the altitude of P2 de-109 scends below 50 km, two counter-streaming electron beams along the ambient magnetic 110

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Figure 1. Data from an ARTEMIS P2 lunar flyby over magnetic anomalies on 26 November 2011. (a) Altitude of the probe as a function of time. P2 reaches a periselene at an altitude of 23.8 km at 10:12 UT. (b) Ambient magnetic field in SSE coordinates. (c) Electron pitch angle spectrum for energies ranging from 2 eV to 32 keV. The differential energy flux has unites of $eV/(eV \text{ cm}^2 \text{ sr s})$. (d) Electron temperatures parallel (Z axis) and perpendicular (X and Y axis) to the magnetic field. (e) Wave burst data in magnetic field-aligned coordinates, Z axis being parallel to the magnetic field. (f)–(g) FFT wave spectra of electric and magnetic field, respectively. Text labels indicate time of day in UT, solar zenith angle (SZA), and spacecraft (X, Y, Z) SSE coordinates in units of lunar radii (R_L).

field can be seen intermittently from the electron pitch angle spectrum in Figure 1c. Since the magnetic field is $+B_x$ dominated in SSE coordinates, the X axis being the direction pointing from the Moon toward the Sun, the beam with pitch angles around 180° can therefore be identified as incoming solar wind electrons. The other beam with ~ 0° pitch angles results from the primary electrons reflected from the magnetic anomalies, as well as photoelectrons emitted from the dayside lunar surface (Whipple, 1981; J. S. Halekas, Poppe, Farrell, et al., 2012).

During the time period of the flyby, we observe high frequency electrostatic fluc-118 tuations ranging from ~ 0.1 to 8 kHz, as shown in the electric field FFT spectrum in 119 Figure 1f. The fluctuations become more intense between 10:10 UT and 10:12 UT. Fig-120 ure 1e shows the high-resolution wave burst data (which has been rotated to magnetic 121 field-aligned coordinates, with E_z parallel to the field line), revealing that the electro-122 static fluctuations are mostly perpendicular to the magnetic field between 10:10 UT and 123 10:11 UT. We also find broadband magnetic fluctuations extending from tens of Hz down 124 to near-DC levels in magnetic field FFT spectrum (Figure 1g). These waves are most 125 likely to be whistler mode, as there are really no other electromagnetic modes that can 126 propagate in this frequency range. Figure 1d shows the electron temperatures parallel 127 $(T_{e,z})$ and perpendicular $(T_{e,x} \text{ and } T_{e,y})$ to the magnetic field, where perpendicular elec-128 tron heating is observed between 10:10 UT and 10:11 UT. In addition, strong isotropic 129 heating is seen between 10:11 UT and 10:12 UT, accompanied by the intense electrostatic 130 fluctuations. 131

Figure 2 shows an overview of another flyby event (ARTEMIS P1) that occurred 132 on 10 February 2013, when P1 was at average GSE coordinates of $[57, 5, 5]R_E$. The sig-133 natures we see are very similar to the previous event. Electrostatic fluctuations are ob-134 served in the electric field FFT spectrum (Figure 2f) between 16:28 UT and 16:36 UT, 135 although the field aligned wave burst data (Figure 2e) indicate that the electrostatic fluc-136 tuations are mainly parallel to the magnetic field this time. Strong isotropic heating is 137 also observed in the electron temperature profile (Figure 2d) between 16:31 UT and 16:37 138 UT, coinciding with the intense electrostatic fluctuations. In addition, a recent analy-139 sis by J. S. Halekas et al. (2014) pointed out that this flyby event has many of the as-140 pects of a collisionless shock, despite the small scale size. 141

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Figure 2. Data from an ARTEMIS P1 lunar flyby over magnetic anomalies on 10 February 2013. P1 reaches the lowest altitude of 19.5 km above the lunar surface at 16:34 UT. All panels same as Figure 1.

¹⁴² 3 Origin of the Electrostatic Waves

We now consider the generation mechanisms for the electrostatic fluctuations, which 143 we suspect may result from a combination of different plasma processes in the complex 144 lunar environment over the magnetic anomalies. The Moon acts as a barrier to the in-145 coming solar wind flow. Due to the influence of crustal magnetic fields, large fractions 146 of the solar wind electron and ion populations reflect above the lunar surface and stream 147 towards the solar wind flow, resulting in varieties of plasma instabilities that could pro-148 duce the electrostatic fluctuations shown in Section 2. Two possible drivers for the waves 149 in Figures 1 and 2 are proposed: electron two-stream instability (ETSI) that could cause 150 electrostatic fluctuations parallel to the ambient magnetic field, and electron cyclotron 151 drift instability (ECDI), which can generate the electrostatic waves in the perpendicu-152 lar direction. 153

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3.1 Electron Two-Stream Instability

Electron two-stream instability driven by counter-streaming electron beams is one 155 of the most commonly found electrostatic instabilities in space plasmas. For example, 156 ETSI has been reported in the solar wind (Malaspina et al., 2013), Earth's magnetotail 157 (Matsumoto et al., 1994), and at the bow shock (Bale et al., 1998). The nonlinear evo-158 lution of ETSI often leads to the formation of time domain structures (Mozer et al., 2015), 159 such as electrostatic solitary waves (Jao & Hau, 2014; Graham et al., 2016), electron phase-160 space holes (Franz et al., 2005; Hutchinson, 2017; Holmes et al., 2018; Steinvall et al., 161 2019), and double layers (Andersson et al., 2002; Ergun et al., 2003). 162

We show an example of the time domain structures observed during the ARTEMIS 163 P1 lunar flyby in Figure 3a and the corresponding electric field FFT spectrum in Fig-164 ure 3b. Since the incident and counter-streaming electron beams are mostly adiabatic, 165 the intense electric field spikes that result from the ETSI have significant field-aligned 166 components E_z . If we zoom in on the time scale, a series of isolated bipolar electric field 167 structures, known as electron phase-space holes, can be seen in Figure 3c. Similar bipo-168 lar structures have also been previously observed near the Moon by Kaguya (Hashimoto 169 et al., 2010). Electron phase-space holes are often responsible for heating and scatter-170 ing background electrons through wave-particle interaction in many space plasmas (Mozer 171 et al., 2016; Vasko et al., 2017). 172



Figure 3. (a) An example of the time domain structures observed during the ARTEMIS P1 lunar flyby on 10 February 2013. (b) Corresponding electric field FFT spectrum showing the broadband electrostatic fluctuations. (c) Zoom-in on the time scale over the blue-colored region in (a) to demonstrate electron phase-space holes.



Figure 4. (a) An example of a reflected ion beam traversing the solar wind plasma perpendicular to the background magnetic field. This sample ion velocity distribution cut, in plasma frame, was obtained at 10:11:30 UT during the P2 lunar flyby event as shown in Figure 1. (b) Schematic illustration of the magnetic field (-Y direction) and incoming solar wind (-X direction) geometry during the flyby in SSE coordinates.

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3.2 Electron Cyclotron Drift Instability

Most of the electrostatic instabilities driven by either field-aligned beams or currents can only generate electrostatic fluctuations along the magnetic field. However, electron cyclotron drift instability, which arises from a relative drift between ions and electrons across the magnetic field, can lead to electrostatic waves in the electron cyclotron frequency in the perpendicular direction (Forslund et al., 1970).

ECDI results from reactive coupling between electron cyclotron Bernstein modes 179 and ion beam modes (Muschietti & Lembège, 2013). The linear dispersion relation of 180 the ECDI can be found in Janhunen et al. (2018). This type of instability has been ob-181 served in many laboratory plasmas (Ripin & Stenzel, 1973; Stenzel & Ripin, 1973) and 182 occasionally in space (Wu & Fredricks, 1972; Wilson et al., 2010). ECDI plays an im-183 portant role in particle acceleration and heating in the foot of supercritical quasi-perpendicular 184 shocks, where a fraction of incoming ions are reflected at the steep front (Matsukiyo & 185 Scholer, 2006). Similar conditions favoring the occurrence of ECDI can be found in the 186

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lunar plasma environment tens of kilometers above the magnetic anomalies. The electrons in these regions are strongly magnetized; however, the ions are considered to be
unmagnetized due to their very large gyroradii in comparison to the scale of the lunar
crustal magnetic fields. The ion beam reflected from the magnetic anomalies therefore
can stream in any direction, in particular, across the magnetic field, triggering the ECDI.

ECDI provides a mechanism capable of driving the perpendicular electrostatic fluc-192 tuations during the P2 flyby as shown in Figure 1e. Reflected ion beams traversing the 193 solar wind plasma perpendicular to the background magnetic field are observed from ion 194 velocity distributions, a good example of which is presented in Figure 4a. Two ion beams 195 can be correspondingly identified: a dense solar wind ion core close to the origin and the 196 reflected ion beam streaming at ~ 200 km/s across the magnetic field. Figure 4b illus-197 trates the geometry of the magnetic field (-Y direction) and incoming solar wind (-X198 direction) during the flyby in SSE coordinates. Once the solar wind ions are reflected 199 from the magnetic anomaly, they are accelerated by the motional electric field and stream 200 towards the -Z direction. The perpendicular configuration of the magnetic field and re-201 flected ion beam therefore favors the occurrence of ECDI, resulting in the time-domain 202 structures with perpendicular electric field in Figure 1e. 203

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3.3 Discussion

As discussed earlier, ECDI can result in electrostatic waves propagating perpen-205 dicular to the magnetic field. However, so far we have not considered the effect of the 206 electron parallel motion in the dispersion relation of the ECDI. If we allow a finite wave 207 vector along the magnetic field in the dispersion relation, then a new type of instabil-208 ity naturally arises in the solution: modified two-stream instability (MTSI) (Janhunen 209 et al., 2018). Due to the parallel component of the wave fields, previous studies have shown 210 that the MTSI can cause strong parallel heating of electrons (McBride et al., 1972; Mat-211 sukiyo & Scholer, 2006). In section 4, we will show that ECDI, along with MTSI, may 212 provide a mechanism responsible for substantial isotropic electron heating as observed 213 over lunar magnetic anomalies. 214

4 Wave-Particle Interaction and Electron Heating

When waves are traveling through a plasma, the fluctuating wave fields can inter-216 act with the charged particles of the plasma, resulting in many interesting nonlinear ef-217 fects. As shown in section 2, one of the important features we notice in the two flyby events 218 is the significant electron heating observed over lunar crustal magnetic anomalies (Fig-219 ures 1d and 2d). Furthermore, the electric and magnetic field FFT wave spectra (Fig-220 ures 1f-1g and 2f-2g) reveal that the electron heating is always accompanied by signif-221 icant electrostatic and/or electromagnetic wave power, suggesting that the wave-particle 222 interaction may play an important role in heating the electrons. 223

To demonstrate the mechanisms causing the electron heating above the magnetic 224 anomaly regions, here we only focus on the ARTEMIS P2 lunar flyby as shown in Fig-225 ure 1. We plot the perpendicular electron temperature and the electromagnetic wave en-226 ergy density together as a function of time in Figure 5a. We present a similar figure show-227 ing the parallel electron temperature and the electrostatic wave energy density as a func-228 tion of time in Figure 5b, where the perpendicular temperature is also shown for com-229 parison. We notice that there are two peaks (A and B) in the perpendicular electron tem-230 perature and one peak (C) in the parallel temperature. Further investigation of the cor-231 relation between the wave power and the electron temperature shows that these peaks 232 are caused by different heating mechanisms. 233

We note that peak A in the perpendicular electron temperature is very well aligned 234 with the electromagnetic wave power in Figure 5a, suggesting that the heating likely re-235 sults from the cyclotron resonance with the whistler modes. When the electrons encounter 236 the whistler waves Doppler-shifted to their cyclotron frequency or higher harmonics, they 237 can strongly interact with the wave fields, gaining perpendicular energy and causing the 238 waves to damp (Tsurutani & Lakhina, 1997; Stenzel, 2016). Similar perpendicular heat-239 ing by electromagnetic waves near the Moon has been reported in J. S. Halekas, Poppe, 240 Farrell, et al. (2012) – though in this case it happens in the magnetotail. In addition to 241 cyclotron damping, we also note that the peak A coincides with the peaks of the elec-242 trostatic wave power in Figure 5b, suggesting that the perpendicular electrostatic waves 243 driven by ECDI are likely to be another source contributing to the perpendicular heat-244 ing in peak A. This electron heating mechanism resulting from ECDI is also observed 245 in the foot region of perpendicular shocks (Matsukiyo & Scholer, 2006). 246



Figure 5. (a) Perpendicular electron temperature and electromagnetic wave energy density (frequencies ranging from near-DC levels to tens of Hz) as a function of time for the P2 lunar flyby event shown in Figure 1. (b) Parallel electron temperature and electrostatic wave energy density (frequencies ranging from 10 Hz to 8 kHz) as a function of time for the same flyby. The perpendicular temperature is also shown in the background for comparison. Inset shows the same electric field as Figure 1e. The largest two peaks in the perpendicular temperature is denoted by A and B, respectively. The largest peak in the parallel temperature is denoted by C.

As to peak B, since it is only accompanied by the peaks of the electrostatic wave 247 power in Figure 5b, this heating therefore is likely to be caused by ECDI as well. Last 248 but not the least, peak C may seem to be quite puzzling at first. Even though it is aligned 249 with peak B and accompanied by intense electrostatic wave power, the perpendicular elec-250 tric fields resulting from ECDI cannot heat the electrons in the parallel direction. How-251 ever, as discussed in section 3.3, MTSI can be driven unstable in the similar conditions 252 as ECDI. In fact, ECDI and MTSI can often be excited simultaneously, allowing for sub-253 stantial electron heating both perpendicular and parallel to the magnetic field (Wu et 254 al., 1984; Muschietti & Lembège, 2013; Janhunen et al., 2018). Since ETSI can also lead 255 to parallel heating, therefore, the isotropic heating seen in peaks B and C (Figure 5b) 256 may be caused by a combination of contributions from ECDI, MTSI, and ETSI. 257

5 Conclusions

In conclusion, we have investigated two types of electrostatic instabilities observed over the lunar crustal magnetic anomalies during ARTEMIS flyby. The electrostatic waves

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propagating parallel to the ambient magnetic field are attributed to upward electron beams 261 reflected by the crustal magnetic fields. We have also reported for the first time on ob-262 servations of another class of electrostatic waves propagating perpendicular to the mag-263 netic field. A proposed free-energy source is associated with reflected ion beams stream-264 ing across the background magnetic field. Finally, our analysis suggests that the perpen-265 dicular electron heating observed above the magnetic anomalies is mainly caused by cy-266 clotron damping and ECDI. The isotropic heating, on the other hand, may result from 267 joint effects due to ECDI, MTSI, and ETSI. 268

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