

High temperature and pressure rheological experiments on felsic granulite

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Abstract

Rheological properties of continental lower crust (CLC) are essential for understanding many geodynamical processes in the Earth. Here we performed a series of deformation experiments on synthetic felsic granulite (65% plagioclase + 15% quartz + 9% clinopyroxene + 6% orthopyroxene + 5% amphibole) using a 5 GPa modified Griggs-type deformation apparatus at 827-927 {degree sign}C and 1 GPa. All experiments reached a steady-state creep at strain rates ranging from 3.0×10^{-6} s⁻¹ to 1.0×10^{-4} s⁻¹, yielding a stress exponent of 4.2 ± 0.1 , a pre-exponential factor of $10^{-4.4 \pm 0.2}$ MPa^{-4.2} s⁻¹, and an activation energy of 260 ± 30 kJ/mol. Microstructural observations show that plagioclase and pyroxene in deformed samples develop noticeable intracrystalline plasticity and shape preferred orientation (SPO). Electron backscatter diffraction (EBSD) measurements demonstrate more significant crystallographic preferred orientations (CPOs) of pyroxene and plagioclase in deformed granulites compared to those in hot-pressed ones. By contrast, quartz develops nearly random fabrics probably due to the simultaneous activation of multiple slip systems during deformation. These features indicate that dislocation creep dominates the deformation of these minerals and felsic granulite, consistent with the obtained stress exponent. Our data imply that a CLC consisting of felsic granulite is weaker rheologically compared to the quartz-dominated upper crust and the olivine-dominated uppermost mantle, which supports the ‘jelly sandwich’ model for the strength of continental lithosphere. Additionally, extrapolations of flow laws of our felsic granulite and formerly-reported mafic granulite strongly favor the hypothesis of delamination-induced decratonization of the North China Craton.

High temperature and pressure rheological experiments on felsic granulite

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Key Points:

- The rheology of felsic granulite can be described by $\dot{\epsilon} = 10^{-4.4 \pm 0.2} \sigma^{4.2 \pm 0.1} \exp[-(260 \pm 30)/(RT)]$.
- Dislocation creep dominates the deformation of granulite and its component minerals.
- Our data favor delamination-induced decratonization of the North China Craton.

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28 understanding many geodynamical processes in the Earth. Here we performed a series
29 of deformation experiments on synthetic felsic granulite (65% plagioclase + 15%
30 quartz + 9% clinopyroxene + 6% orthopyroxene + 5% amphibole) using a 5 GPa
31 modified Griggs-type deformation apparatus at 827-927 °C and 1 GPa. All
32 experiments reached a steady-state creep at strain rates ranging from $3.0 \times 10^{-6} \text{ s}^{-1}$ to
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34 $10^{-4.4 \pm 0.2} \text{ MPa}^{-4.2} \text{ s}^{-1}$, and an activation energy of $260 \pm 30 \text{ kJ/mol}$. Microstructural
35 observations show that plagioclase and pyroxene in deformed samples develop
36 noticeable intracrystalline plasticity and shape preferred orientation (SPO). Electron
37 backscatter diffraction (EBSD) measurements demonstrate more significant
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41 multiple slip systems during deformation. These features indicate that dislocation
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1. Introduction

The continental lower crust (CLC) represents where extensive mass and energy exchanges between the upper crust and the uppermost mantle occur. Rheological property of the CLC is an important parameter that concerns the interpretation of many geological processes and phenomena in the deep Earth, such as delamination and underplating (e.g., [Meissner & Mooney, 1998](#)), surface topography (e.g., [Bendick et al., 2008](#)), lower-crustal earthquakes (e.g., [Reyners et al., 2007](#); [Shi et al., 2018](#)), seismic lamination and anisotropy of the CLC (e.g., [Meissner et al., 2006](#); [Ross et al., 2004](#); [Shapiro et al., 2004](#)), and coupling and decoupling of crust and mantle (e.g., [Royden, 1996](#)). Understanding the rheological properties of CLC is thus fundamental to fully decipher the geodynamics of the Earth.

There are, however, still hot debates regarding the rheology of CLC. Some researchers suggest that the CLC is mechanically weaker than the upper mantle and the upper crust (i.e., the “jelly sandwich” model) (e.g., [Chen & Molnar, 1983](#)), while some argue that the strength of continental lithosphere resides entirely in the crust and that the upper mantle is significantly weaker compared to the CLC [i.e., the “crème brûlée” model ([Burov & Watts, 2006](#))] (e.g., [Jackson, 2002](#); see [Bürgman & Dresen, 2008](#) for a review). Both hypotheses find supports from many geophysical and geodetic observations and numerical modellings (e.g., [Beaumont et al., 2001](#); [Clark et al., 2005](#); [Copley et al., 2011](#); [Ross et al., 2004](#); [Royden et al., 1997](#)). These disputes arise, on the one hand, due to the variety of rheology in response to heterogeneous lithologies and thermodynamic conditions, including temperature and fluid, in the CLC. On the other hand, geophysical and geodetic techniques usually give non-unique, qualitative but not quantitative constraints on the mechanical properties of the CLC. In contrast to these inferences generally deduced from natural observations, laboratory experiments on representative CLC rocks under controlled thermodynamic conditions allow us to resolve these disputes in a more

79 straightforward way.

80
81 Based on studies on exposed lower crustal sections and basalt-borne lower crustal
82 xenoliths and seismic velocities of the CLC, it is generally suggested that the CLC is
83 intermediate to mafic in bulk composition ([Gao et al., 1998](#); [Rudnick & Fountain,](#)
84 [1995](#); [Rudnick & Gao, 2003](#)). A recent review ([Hacker et al., 2015](#)), however, argues
85 that much of the CLC is relatively felsic instead. Whatever the bulk composition is,
86 mafic/felsic granulites are commonly observed as an important rock type in lower
87 crustal sections and xenoliths (e.g., [Chen et al., 2006](#); [Franěk et al., 2011](#); [Griffin &](#)
88 [O'Reilly, 1986](#); [Martelat et al., 1999](#)). Therefore, a thorough knowledge of rheology
89 of the CLC requires an understanding of the mechanical behavior of mafic/felsic
90 granulites. Only a few studies have touched the rheology of granulites up to now, in
91 stark contrast to numerous studies on the proxy of the upper mantle, i.e. olivine (see a
92 review by [Hirth & Kohlstedt, 2003](#)). [Wilks and Carter \(1990\)](#) first investigated the
93 rheology of CLC through experiments on selected natural mafic and felsic granulites.
94 Unfortunately, due probably to the coarse-grained ($\gg 100 \mu\text{m}$) nature of their samples,
95 most of their experiments show semibrittle to cataclastic behavior and do not reach
96 steady-state creep. Therefore, these results cannot be reliably extrapolated to nature.
97 Recently, [Wang et al. \(2012\)](#) performed experiments on synthetic mafic granulites
98 using a 3 GPa modified Griggs-type apparatus, which greatly improves the precision
99 of measured stress level with molten salt as the confining pressure medium. They
100 showed that the steady-state creep strength of the studied mafic granulite is weaker
101 than olivine aggregates, and thus supports the “jelly sandwich” model of the
102 continental lithosphere strength envelope.

103
104 Since pyroxene and plagioclase are the most common minerals that make up the CLC,
105 many studies have alternatively made efforts to determine the flow laws of pyroxene
106 and/or plagioclase aggregates under various thermodynamic conditions (e.g.,

Bystricky & Mackwell, 2001; Dimanov & Dresen, 2005; Dimanov et al., 1999, 2003; Rybacki & Dresen, 2000; Rybacki et al., 2006) to understand the rheology of CLC. Theoretically, these single-phase flow laws can be used to reasonably deduce the flow law of a multiphase aggregate that is made up of these phases (e.g., Ji et al., 2003; Tullis et al., 1991). Nevertheless, this application is generally limited by the facts that 1) the deformation mechanism for each component mineral may not be the same or may change with progressive deformation, especially when dynamic recrystallization occurs; 2) possible reactions may happen between the component minerals; and 3) mechanical anisotropy is not taken into account in such theoretical models. Therefore, a direct experimental study on granulite is necessary.

In that the constitutive equation for a synthetic mafic granulite (Wang et al., 2012) has been established recently, we therefore aim to determine in this study the rheological parameters of felsic granulite under conditions appropriate to the CLC. To overcome the limitations of previous experiments on natural felsic granulites by Wilks and Carter (1990), we used synthetic fine-grained (32-53 μm) felsic granulites, which have been well characterized in terms of water concentrations and chemical compositions of the major component minerals, as experimental samples. In addition, the deformation apparatus used in this study improves greatly the accuracy of stress measurement during deformation compared to the one used in Wilks and Carter (1990) by using molten salt as the confining pressure medium (Green & Borch, 1989).

2. Experimental Samples and Methods

2.1. Starting Material

As starting material, we used a natural, fresh felsic granulite (09TPZ-1) collected from Taipingzhai in the Archean lower crust terrane of the North China Craton (NCC). This felsic granulite is primarily composed of 65% plagioclase, 15% quartz, 9% clinopyroxene, 6% orthopyroxene, and 5% amphibole, with accessory K-feldspar and

ilmenite. P-T estimates using major element- and REE-based thermobarometers yielded a temperature of 1025-1060 °C and a pressure of 0.9 GPa for the clinopyroxene-orthopyroxene pair in the Taipingzhai granulites (Yang & Wei, 2017). Mineral separates were prepared using heavy liquid and magnetic separation and then ground to fine powders using a tungsten carbide mortar. These powders were then sorted to a desired grain size of 32-53 µm using a sedimentation method based on the Stokes' law of settling. The sorted powders were subsequently well mixed in terms of the modal composition of the starting material. Prior to being loaded into sample capsules, the powders were dried in a drying oven at 100 °C for more than 48 hours to remove possible absorbed water.

2.2. Deformation Apparatus and Experiments

All experiments were performed in a 5 GPa modified Griggs-type apparatus at the State Key Laboratory of Geological Processes and Minerals Resources, China University of Geosciences (Wuhan) (GPMR-CUG). Sample assembly and preparation procedure used in this study generally followed those detailed in Zhang and Green (2007) and Wang et al. (2012), except that a slight modification to the sample dimension (3.2 mm ID instead of 3.0 mm ID) was made in this study. The sample assemblies were kept in a drying oven at 100 °C for > 12 h before they were put into the high-pressure apparatus. Constant piston-velocity deformation experiments were carried out at a pressure of 1 GPa and temperatures of 827-927 °C immediately following a hot-pressing of the powder sample under 1 GPa and 927 °C conditions for 12 h. For those experiments not performed at 927 °C, load was not applied until temperature was adjusted to the desired value and kept for 1 h while maintaining the confining pressure constant. To determine the initial sample length and microstructures for each deformation experiment, three hot-pressing experiments were performed separately under 1 GPa and 927 °C conditions for 12 h. At the end of each experiment, temperature was quenched to 300 °C within a few seconds to retain

the microstructures. Experimental temperatures are measured to a precision of 1 °C using the thermocouple located near the top of the sample. This thermocouple generally gives higher (up to 25 °C) temperature readings than that located near the bottom of the sample. Data of pressure, temperature, load, strain rate, and displacement were recorded automatically every 10 seconds by a computer.

2.3. Mineral Composition Analyses

To examine whether any modifications to chemical compositions occur during experiments, we measured the major-element compositions of component minerals in the starting material and the hot-pressed and deformed samples by a JEOL JXA-8230 electron probe microanalyser (EPMA) at the Center for Global Tectonics, School of Earth Sciences, China University of Geosciences (Wuhan). The analyses were conducted at an accelerating voltage of 15 kV, a probe current of 20 nA, and a beam size of 3 µm. Dwell times were 10 s on element peaks and 5 s on background locations adjacent to peaks. Raw X-ray intensities were corrected using a ZAF correction procedure. Natural and synthetic SPI standards were used as monitor standards: sanidine for K, pyrope for Fe and Al, diopside for Ca and Mg, jadeite for Na, rhodonite for Mn, olivine for Si, and rutile for Ti. Furthermore, we also collected semi-quantitative X-ray mineral chemical mapping of deformed samples using an energy-dispersive spectrometer mounted on a Quanta 450 field emission gun scanning electron microscopy (SEM) at GPMR-CUG.

2.4. Microstructural Characterization

Thin sections and/or polished sections were prepared by cutting the experimental specimens longitudinally and polishing to within 0.05 µm using colloidal alumina. Microstructures of hot-pressed and deformed samples were examined by a petrographic microscope and the SEM at GPMR-CUG. Crystallographic preferred orientation (CPO) of minerals were determined by an Oxford Instruments Nordlys-II

electron backscatter diffraction (EBSD) system at GPMR-CUG. To reveal the intracrystalline deformation features, we also performed EBSD mapping on selected grains from deformed samples. The measurements were carried out with an accelerating voltage of 20 kV, a spot size of 6, and a working distance of 23-25 mm. Diffraction patterns were manually collected and indexed using the Channel 5 + software. Only those measurements with mean angular deviation (MAD) values below 1.0 (between detected and simulated Kikuchi bands) were recorded. The strength of CPO is quantified by the misorientation index (M-index) developed by Skemer et al. (2005) and the J-index using the program of *SuperJctfPC.exe* courtesy of David Mainprice. The M-index is 0 for a random fabric and 1 for a single crystal, while the J-index is 1 for a random distribution and infinity for a single crystal orientation.

2.5. Fourier Transform Infrared (FTIR) Analyses

To determine the initial water concentrations of minerals in the experimental samples, we prepared a thick section of 225 μm thickness by doubly polishing a rock slab randomly cut from the natural sample (09TPZ-1). Unpolarized FTIR spectroscopic measurements were taken at GPMR-CUG on a Nicolet 6700 spectrometer coupled with a Continuum IR microscope using a KBr beam splitter and a liquid-nitrogen cooled MCT-A detector. 128 scans were accumulated for each spectrum with a resolution of 4 cm^{-1} . Depending on grain size, the aperture size was set to be 50 $\times$ 50 μm to 100 $\times$ 100 μm . Measurements were made on optically clean, crack- and inclusion-free regions. More than 10 different grains for each mineral were analyzed. Water concentrations were calculated using a modified form of the Beer-Lambert law

$$c = \Delta / (I \times t)$$

where c is the water concentration (ppm H_2O by weight), Δ is the integrated area (cm^{-1}) of absorption bands in the region of interest, t is the sample thickness (cm), and I is the integral specific absorption coefficient ($\text{ppm}^{-1} \text{ cm}^{-2}$). For each spectrum,

baseline correction was manually performed and the OH absorption bands were integrated between 3000 and 3800 cm⁻¹. The obtained area beneath the OH peaks was multiplied by 3 to give the Δ values. The integral specific absorption coefficients of 7.09 ppm⁻¹ cm⁻² for clinopyroxene (Bell et al., 1995), 14.84 ppm⁻¹ cm⁻² for orthopyroxene (Bell et al., 1995), 13.50 ppm⁻¹ cm⁻² for quartz (Thomas et al., 2009), and 15.30 ppm⁻¹ cm⁻² for plagioclase (Johnson & Rossman, 2003) were used. The thickness of the thick section was measured using a digital micrometer and reported as an average of 10 measurements covering the whole section. The uncertainties arising from use of an unpolarized IR beam on unoriented anisotropic minerals can be minimized by measurements on more than ten different grains of each mineral (e.g., Kovács et al., 2008; Withers, 2013).

3. Results

3.1. Experimental Results

The three hot-pressing experiments yield quite similar axial compression ratio, ranging between 35% and 38% (Table 1). The retrieved hot-pressed and deformed samples remain nearly cylindrical in shape (Figure 1). Measurements on one of the hot-pressed samples show that it has a density of 3.14 g/cm³ and a porosity of ~ 1%.

At high temperatures and pressures, a power-law creep often dominates the plastic deformation of rocks. The constitutive relation is well described by an empirical equation of the form:

$$\dot{\epsilon} = A\sigma^n \exp[-Q/RT]$$

where $\dot{\epsilon}$ is the strain rate, A is the pre-exponential factor related with material, σ is the steady-state flow stress, n is the stress exponent related to deformation mechanism ($n = 3-5$ for dislocation creep and $n = \sim 1$ for diffusion creep), Q is the activation energy, R is the gas constant, and T is the absolute temperature. 10 deformation experiments (Table 1) in total have been carried out at a confining pressure of 1 GPa

to determine the rheological parameters for the studied felsic granulite. Representative stress versus strain curves are illustrated in Figure 2a, which shows that a steady-state deformation generally reaches after a strain of about 20%. The deformation experiments can be classified into two groups: one group is performed at constant temperature but varying strain rates of $3.6 \times 10^{-6} \text{ s}^{-1}$, $1.1 \times 10^{-5} \text{ s}^{-1}$, $9.4 \times 10^{-5} \text{ s}^{-1}$, respectively, while the other group is performed at constant strain rate but varying temperatures of 827 °C, 877 °C, and 927 °C, respectively. As shown in Figure 2b, a log-log plot of strain rate versus stress for the first group of experiments yields a stress exponent of $n = 4.2 \pm 0.1$ and a pre-exponential factor of $A = 10^{-4.4 \pm 0.2} \text{ MPa}^{-4.2} \text{ s}^{-1}$. The creep activation energy was determined to be $260 \pm 30 \text{ kJ/mol}$ using least square fits based on the second group of experiments (Figure 2c). The rheology of felsic granulite over the range of conditions explored here can thus be described by the following constitutive equation:

$$\dot{\varepsilon} = 10^{-4.4 \pm 0.2} \text{ MPa}^{-4.2} \text{ s}^{-1} \sigma^{4.2 \pm 0.1} \exp\left(-\frac{260 \pm 30 \text{ kJ/mol}}{RT}\right)$$

where $\dot{\varepsilon}$ is in s^{-1} , σ in MPa, and T in Kelvin.

3.2. Microstructures

Under petrographic microscope and SEM, the hot-pressed samples show nearly homogeneous distribution of each component mineral; no compositional layering or single-phase clusters have been observed (Figures 3a and 3b). A weak shape preferred orientation (SPO) as shown by the partial alignment of anisometric pyroxene and amphibole grains (Figures 3a and 3b) can be discerned. EBSD results show that all minerals from the hot-pressed samples have very weak CPOs (Figure 4a). Pyroxene CPOs are similar, featured by the [100]-axis approximately parallel to while the [001]-axis nearly normal to the compression direction. Although being more diffused, the [100]-axis of plagioclase grains appear to be nearly normal to the compression direction. The *c*-axes of quartz form two weak, broad point maxima with one nearly normal to and the other oblique at an angle of $\sim 45^\circ$ to the compression direction.

274

275 In contrast to the hot-pressed samples, the deformed samples developed more
276 significant SPO manifested by the strong alignment of elongate pyroxene, amphibole,
277 plagioclase and sometimes quartz (Figures 3c and 3d), which also defines the foliation
278 plane in deformed samples. Undulose extinction is frequently observed in plagioclase
279 and quartz. EBSD mappings show that subgrain structures are common in plagioclase,
280 pyroxene and quartz (Figure 5). Due probably to low finite strain ($\epsilon < 50\%$; Table 1),
281 mineral CPOs in the deformed samples are still weak. However, pyroxene and
282 plagioclase CPOs in the deformed samples are stronger than their counterparts in the
283 hot-pressed samples (Figure 4). Pyroxene CPOs are characterized by a clear point
284 maximum of the [100]-axis nearly normal to the foliation plane and a great girdle of
285 the [001]-axis nearly parallel to the foliation plane. The [100] axes of plagioclase
286 form a girdle in the foliation plane, while poles to the (001)-plane concentrate
287 approximately normal to the foliation plane. The quartz CPOs are nearly random with
288 J-index ranging from 2.82 to 3.44 (Figures 4b-c).

289

290 Finally, it is noteworthy that the phase boundaries and triple junctions among the
291 component minerals in deformed samples are clear-cut and devoid of any amorphous
292 phases (i.e., melt), even at high magnifications (Figure 6).

293

294 3.3. Chemical Composition and Water Concentration

295 The major-element compositions of minerals in the starting material and selected
296 hot-pressed and deformed samples are shown in Table 2. Compared to the starting
297 material, no significant compositional change occurs during hot-pressing and
298 deformation. In addition, X-ray maps of Si, Al, Fe, Mg and Ca show no discernable
299 compositional zoning in each mineral and sharp grain/phase boundaries (Figure 7).

300

301 Figure 8 illustrates the typical infrared spectra for plagioclase, quartz, clinopyroxene,

and orthopyroxene from the starting material. All spectra are normalized to 1 cm thickness and offset vertically for clarity. A summary of FTIR results is listed in [Table 3](#). Most FTIR spectra of plagioclase are characterized by two absorption bands ([Figure 8a](#)): a strong sharp band at 3620 cm^{-1} and a broad band centered around 3440 cm^{-1} , which may be related to structural OH or water molecule ([Johnson & Rossman, 2004](#)). The calculated water concentrations in plagioclase vary considerably, ranging from 17-400 wt ppm with an average of 209 ± 126 wt ppm ([Table 3](#)). Infrared spectra of quartz are consistently similar and characterized by the presence of two small sharp peaks at 3380 cm^{-1} and 3440 cm^{-1} superimposing on a broad band at 3400 cm^{-1} ([Figure 8b](#)). The sharp peaks are generally attributed to structurally bound hydrogen associated with Al^{3+} or monovalent cations ([Aines & Rossman, 1984](#); [Johnson, 2006](#)), while the broad band is considered to be due to the O-H stretching vibrations of “liquid-like” molecular water (e.g., [Aines & Rossman, 1984](#); [Ito & Nakashima, 2002](#); [Kronenberg & Wolf, 1990](#); [Nakashima et al., 1995](#)). Quartz contains an average water concentration of 355 ± 127 wt ppm ([Table 3](#)). Clinopyroxene displays significant absorption bands at 3620 cm^{-1} , 3530 cm^{-1} , and 3450 cm^{-1} , with intensity of 3620 cm^{-1} generally being the strongest ([Figure 8c](#)). Such bands are analogous to those reported by previous studies for clinopyroxene and commonly considered as structural OH (e.g., [Peslier et al., 2002](#); [Rossman, 1996](#); [Skogby et al., 1990](#); [Xia et al., 2006](#); [Yang et al., 2008](#)). The calculated average water concentration in clinopyroxene is 937 ± 521 wt ppm ([Table 3](#)). The most prominent feature of the FTIR spectra in orthopyroxene is a stronger sharp absorption band at around $3620\text{--}3605\text{ cm}^{-1}$ accompanied by two weaker, broad bands at $3430\text{--}3415\text{ cm}^{-1}$ and $3250\text{--}3235\text{ cm}^{-1}$, respectively ([Figure 8d](#)). These bands are related to structural OH in orthopyroxene and have been frequently reported by previous studies (e.g., [Skogby et al., 1990](#); [Xia et al., 2006](#); [Yang et al., 2008](#)). The average water concentration is 1132 ± 555 wt ppm in orthopyroxene ([Table 3](#)). According to mineral mode and measured water concentration in each component mineral, the studied felsic granulite thus has a bulk

water concentration of 344 ± 182 wt ppm, excluding the contribution of amphibole. This value falls within the range reported previously for granulites from the eastern China (Xia et al., 2006; Yang et al., 2008).

4. Discussion

4.1. Deformation Mechanism(s)

Owing to being volumetrically predominant, plagioclase should govern the rheological behavior of the studied felsic granulite. Microstructural observations showed that plagioclase in deformed samples developed distinct SPO (Figure 2d), which is nearly normal to the compressive stress direction. EBSD analyses on deformed samples illustrated that plagioclase has significant intracrystalline plasticity (Figure 5a) and stronger CPOs (Figures 4b-d). These facts suggest that dislocation creep is the dominant deformation mechanism in plagioclase, consistent with the dislocation-creep deformation of the felsic granulite as indicated by the stress exponent of 4.2 ± 0.1 . In general, the dominant slip system during dislocation-creep deformation of a mineral can be deduced from its CPO. Strictly speaking, this practice is inapplicable to mineral deformation by axial compression (like in this study) as only the slip plane can be unambiguously determined. Nevertheless, for those low-symmetry crystal system minerals, like plagioclase and pyroxene, the most probable slip direction can generally be inferred if the slip plane is determined, because the common slip systems in these minerals are well known and limited in number. In this study, the CPO patterns of deformed plagioclase are characterized by both the [100]-axis and the (001)-plane subparallel to the foliation plane, suggesting that (001)-plane is the slip plane with [100]-axis being the most probable slip direction, i.e., the (001)[100] slip system. Such a slip system has been frequently inferred to be activated during dislocation creep of plagioclase in naturally and experimentally deformed rocks (e.g., Heidelberg et al., 2000; Ji et al., 2004; Mehl & Hirth, 2008; Wang et al., 2012; Xie et al., 2003).

358

359 Microstructural observations showed that pyroxene has distinct SPO being consistent
360 with that of plagioclase. EBSD analyses revealed significant intracrystalline plasticity
361 and enhanced CPOs in pyroxene after deformation (Figures 4b-d, 5b, and 5c). The
362 CPOs of pyroxene, featured by a point maximum of [100]-axis nearly normal to the
363 foliation and a girdle distribution of [001]-axis nearly parallel to the foliation, point to
364 a dominant slip system of (100)[001]. This slip system is commonly reported to be the
365 one in experimentally and naturally deformed pyroxene by a dislocation-creep
366 deformation (e.g., Farla et al., 2017; Lund et al., 2006; Mauler et al., 2000; Wang et
367 al., 2012; Zhang & Green, 2007). These observations suggest that dislocation creep
368 controls the deformation of pyroxene, which is in agreement with the
369 dislocation-creep deformation of the felsic granulite as well.

370

371 Quartz in deformed samples shows weak SPO and nearly random fabrics. Such
372 features may be interpreted as deformation by grain boundary sliding or diffusion
373 creep. However, these deformation mechanisms do not involve significant
374 intracrystalline plasticity, which is inconsistent with the observed undulose extinction
375 and intracrystalline misorientation in quartz (Figure 5d). In addition, since Si is the
376 slowest diffusing species in quartz, the diffusion of Si will thus be the rate-limiting
377 factor in diffusion creep of quartz. Based on the Arrhenius equation for Si diffusion
378 compiled by Cherniak (2010), we estimated the time required to diffuse Si for 15 μm
379 (i.e., the radius of the smallest quartz grains) at a temperature of 927 °C. Our
380 estimation suggests that the required time is far beyond that of the experimental
381 timescale, indicating that intracrystalline diffusion creep is impossible in our
382 experiments. The absence of fluid or melt along grain boundaries does not favor the
383 operation of grain boundary diffusion creep either, which is identical with the lack of
384 chemical zoning in each mineral (Figure 7).

385

Alternatively, the nearly random CPO of quartz can be ascribed to its small volume fraction (15%) in the samples, which is not enough to form a stress-supporting framework, and the dispersion of quartz grains between those of strong minerals such as plagioclase and clinopyroxene. However, pyroxene in our samples has also a low volume fraction (15%) and is well mixed with other minerals. If the above hypothesis is valid, then it is expected to see nearly random fabrics of pyroxene too. This is obviously not the case—pyroxene in our deformed samples shows well-defined CPO patterns (Figure 4b-d). In addition, as had been demonstrated in natural observations (Martelat et al., 1999), even if quartz does not form a stress-supporting framework, it can still develop non-random CPOs. Zhang and Green (2007) conducted deformation experiments at 927-1227 °C on quartz eclogites. Their data show that quartz develops distinct SPO but no significant CPOs. The authors suggested that multiple slip systems were activated for quartz during deformation, leading to weak to random CPOs. We consider that this hypothesis may well account for the random quartz CPOs in our study. Therefore, dislocation creep dominates the deformation of quartz as well, which is consistent with the observed intracrystalline plasticity in quartz.

In a word, our mechanical data and microstructural observations point to dislocation creep as the dominant deformation mechanism in the deformed felsic granulite and its major component minerals.

4.2. Comparison with Previous Studies

Since plagioclase is the primary component mineral in the studied felsic granulite and commonly considered to be representative of the CLC in terms of rheology, it is thus interesting to compare the rheological strength of our felsic granulite with that of plagioclase in the dislocation-creep regime. Prior to this discussion, however, it is necessary to address whether water in our samples has been lost during deformation experiments. Considering that H has an elevated mobility at high temperatures and

can go through the platinum capsule easily (e.g., [Wilks et al., 1990](#)), it must have been partially lost during experiments. However, an accurate determination of water concentrations of minerals after deformation is hampered due to the following two reasons. First, the grain sizes (32-53 μm) of minerals in our experimental samples are very small. In particular, compression- and decompression-induced microcracks are common in many grains ([Figure 6](#)). These facts make measurements of water concentrations within each mineral after experiments using a conventional FTIR, which is usually applied with an aperture size of $\geq 50 \mu\text{m}$, impractical. Second, although the secondary ion mass spectrometry (SIMS) technique has a better spatial resolution (5-10 μm ; e.g., [Huari et al., 2002](#)) than a conventional FTIR, its application is currently restricted by lacking suitable standards, especially for plagioclase and quartz. However, as suggested below, water loss from our experimental samples must be minimal.

In addition, it is important to notice that the dominant slip system {i.e., the (001)[100] slip system} in plagioclase in this study differs from the one that has been commonly identified based on TEM observations and CPOs of plagioclase from naturally deformed middle-lower crustal rocks (e.g., [Ji & Mainprice, 1988, 1990](#); [Kruse et al., 2001](#); [Miranda & Klepeis, 2016](#); [Montardi & Mainprice, 1987](#); [Olsen & Kohlstedt, 1984](#); [Svahnberg & Piazzolo, 2010](#)). [Ji et al. \(2000, 2004\)](#) proposed that the relative importance of the [001] or [100] Burgers vectors may be a function of temperature, strain rate, water activity or confining pressure. In particular, water is potentially the most important parameter that may influence the slip system transitions in plagioclase, considering that the range of temperature and pressure conditions in the middle-lower crust is limited. Laboratory experiments ([Jung & Karato, 2001](#); [Manthilake et al., 2013](#); [Ohuchi et al., 2012](#)) have clearly shown that, at least for olivine and orthopyroxene, water concentration in these minerals has a pronounced influence on the dominant slip system and their CPOs. The water concentration in plagioclase from

the lower crustal rocks can vary significantly from several tens to nearly one thousand wt ppm (e.g., [Xia et al., 2006](#); [Yang et al., 2008](#)). However, how water concentration may influence the activity of these slip systems in plagioclase remains largely unknown due to lack of experimental constraints. Moreover, information about the water concentration in plagioclase from previous studies on naturally deformed middle-lower crustal rocks is generally lacking, thus making the discussion of the relative activity and strength of the (010)[001] and (001)[100] slip systems in plagioclase inappropriate.

When extrapolated to a natural strain rate of 10^{-14} s^{-1} , our felsic granulite is more than 3 times weaker than dry anorthite ([Rybacki et al., 2006](#)) under the lower-crustal temperature (800-900 °C) conditions ([Figure 9a](#)). Because the presence of water can significantly reduce the mechanical strength of minerals and rocks (e.g., [Chen et al., 2006](#); [Griggs & Blacic, 1965](#); [Hirth & Kohlstedt, 2003](#); [Mei & Kohlstedt, 2000](#); [Wang et al., 2012](#)) and plagioclase in our starting material contains a water concentration of $\sim 0.021 \text{ wt\% H}_2\text{O}$ ([Table 3](#)), it is thus no wonder that our felsic granulite is significantly weaker than the dry anorthite aggregates ([Rybacki et al., 2006](#)). In addition, FTIR results show that quartz (15%), clinopyroxene (9%) and orthopyroxene (6%) in the studied felsic granulite all contain a large amount of water ([Table 3](#)). Extrapolations of published dislocation-creep flow laws for these minerals under wet conditions ([Bystricky et al., 2016](#); [Chen et al., 2006](#); [Rutter et al., 2004](#)) ([Figure 7a](#)) reveal that quartz and clinopyroxene are both much weaker while orthopyroxene is slightly stronger than dry anorthite. Hence, the net effects of these minor minerals also tend to drive the felsic granulite weaker than the dry anorthite. Using the flow law for dislocation creep of anorthite under hydrous conditions ([Rybacki et al., 2006](#)), we also calculated the flow strength of wet anorthite (0.02 wt% H_2O) as a function of temperature at a strain rate of 10^{-14} s^{-1} ([Figure 9a](#)). It is clear from [Figure 9a](#) that our felsic granulite is only slightly (up to nearly 2 times) stronger

than wet anorthite at 800-900 °C. Such a small strength contrast implies that hydrogen diffusional loss from the component minerals during deformation, if any, must be insignificant. Alternatively, this strength contrast could be attributed to the combined effects of quartz and pyroxene, which are respectively weaker and stronger than the wet anorthite (Figure 9a), if we assume that the influence of anorthite content in plagioclase on rheology is negligible.

In Figure 9b, we also compare the rheological strength of our felsic granulite with those of granulites reported by previous studies by extrapolation to a natural strain rate of 10^{-14} s^{-1} . It is clear that our felsic granulite is 6-7 times stronger than the Adirondack felsic granulite (Wilks & Carter, 1990), but comparable to that of the amphibole-rich Pikwitonei mafic granulite (Wilks & Carter, 1990) at 800-900 °C. However, the reason behind the similarities and differences of the strength between our sample and those of Wilks and Carter (1990) is difficult to address currently due to significantly different mineralogy [65% Pl + 15% Py + 15% Qz + 5% Amp in this study versus 45% Kfs + 21% Qz + 20% Cpx + 1% Pl + 1% Amp and 58 Hbl + 40% Pl + 1% Cpx + 1% Opx for felsic and mafic granulite, respectively, in Wilks and Carter (1990)]. Furthermore, water concentrations in minerals from granulites studied by Wilks and Carter (1990) are unstated, which renders such a discussion impossible as well.

Compared with the mafic granulite (57% Pl + 14% Opx + 24% Cpx + 5% Opx) studied by Wang et al. (2012), the felsic granulite studied here is much (6-7 times) stronger (Figure 9b). At first glance, this seems unreasonable since the mafic granulite has more pyroxene which would make it stronger than the pyroxene-poor felsic granulite (Figure 9a). Two reasons could probably explain this discrepancy. First, the mafic granulites have water concentrations (0.16-0.28 wt% H₂O) far more than those (0.034 wt% H₂O) of the felsic granulites. More significant water-weakening effect

caused by more water concentrations in the mafic granulites could explain why the mafic granulites are weaker than the felsic granulites. Second, the anisometric shape of pyroxene and plagioclase grains as a result of fracturing along cleavage planes during powder preparation leads to a weak SPO (i.e., foliation) and fabric (Figures 3c, 3d and 4a) during hot-pressing, which in turn gives rise to mechanical anisotropy (e.g., Hansen et al., 2012; Ji et al., 2000) and increases the strength of the felsic granulite.

4.3. Geological Implications

Our microstructural observations, EBSD results and stress exponent ($n = 4.2 \pm 0.1$) all point to dislocation creep as the dominant deformation mechanism in the studied felsic granulite at lower-crustal P - T conditions. Dislocation creep also plays an important role in the deformation of many naturally deformed lower-crustal rocks (e.g., Lund et al., 2006; Martelat et al., 1999). Therefore, the flow law of felsic granulite obtained here can be applied reliably to natural conditions.

When extrapolated to a natural strain rate of 10^{-14} s^{-1} , our felsic granulite is much stronger than wet quartz aggregates (Figure 9a), which are commonly considered to be rheologically representative of the upper crust, under the same temperature and pressure conditions. However, due to low temperature and dominantly brittle deformation, the upper crust is generally stronger than the lower crust (Bürgmann & Dresen, 2008; Kohlstedt et al., 1995; Ranalli & Murphy, 1987). In addition, dry olivine aggregates (Hirth & Kohlstedt, 2003) and eclogite (Zhang & Green, 2007), which are representative of the upper mantle and the thickened lower crust, respectively, are nearly one order of magnitude stronger than the studied felsic granulite at the Moho depth (Figure 10), even if the felsic granulite is stressed along a ‘hard’ direction. We also calculated the flow strength of water-saturated olivine aggregates at Moho depth (1 GPa, 900 °C) as a function of temperature (Figure 10). The water concentration in water-saturated olivine is estimated to be $\sim 57 \text{ wt ppm}$

H₂O using the solubility equation of [Zhao et al. \(2004\)](#). As a matter of fact, water concentration in mantle olivine is generally far below its saturation level. It is obvious from [Figure 10](#) that the felsic granulite is nearly 3 times weaker than water-saturated olivine aggregates at Moho. These results suggest that the felsic granulite is weaker than the upper crust and the uppermost mantle, supporting the “jelly sandwich” model of the continental lithosphere strength envelope (e.g., [Chen & Molnar, 1983](#)). This viewpoint is consistent with that of [Wang et al. \(2012\)](#), and is supported by many geodetic and geological observations worldwide (e.g., [Homburg et al., 2010](#); for a review, see also [Bürgmann & Dresen, 2008](#)). It is worth mentioning that this viewpoint holds true only when the CLC is composed mainly of felsic granulite similar in mode composition to the one studied here. However, numerous previous studies have shown that the CLC is lithologically complex, with mineral species and modes varying from locality to locality (e.g., [Chen et al., 2006](#); [O'Brien & Rötzler, 2003](#)). This compositional variation explains in part why the “jelly sandwich” model fails somewhere (e.g., [Jackson, 2002](#)) in the Earth.

As summarized by [Wang et al. \(2012\)](#), a weak CLC is helpful to understanding many issues related to deep geodynamics, such as laminated seismic reflectors and seismic anisotropy observed in the CLC (e.g., [Meissner et al., 2006](#); [Ross et al., 2004](#)) and surface topography and crustal shortening (e.g., [Bendick et al., 2008](#); [Chen et al., 2017](#); [Gerbault et al., 2005](#)). For instance, [Gerbault et al. \(2005\)](#) suggested through numerical modeling that horizontal flow of the lower crust from Puna towards Altiplano of the central Andes driven by excess density-induced stress from the Altiplano mantle could explain the different amounts of crustal shortening in Altiplano and Puna. More importantly, rheology of the studied felsic granulite, which samples the Archean lower crust of the NCC, is immediately applicable to and has particular implications for the evolution of the NCC. Several lines of evidence have proved that the NCC, one of the oldest continental nuclei in the world, had lost its

keel or the NCC had been destroyed since the Mesozoic (e.g., Fan et al., 2000; Menzies et al., 1993; Xu et al., 2001). Delamination of the thickened, eclogitic lower crust, whose existence had been inferred from many geochemical and geological evidence (e.g., Gao et al., 2004; Xu et al., 2002, 2008), and probably the lithospheric mantle has been put forward as a viable mechanism to explain the destruction of the NCC (e.g., Gao et al., 2009). In addition to gravity-triggered instability, the onset of delamination requires mechanical decoupling between the lower CLC (eclogite)/lithospheric mantle and the upper CLC (consisting mainly of granulite in the NCC) as well (Meissner & Mooney, 1998). The mafic granulite (Wang et al., 2012) and the felsic granulite studied here are nearly one to two orders of magnitude softer than eclogite (Zhang & Green, 2007; see Figure 10). This strength contrast will surely contribute to mechanical decoupling between the upper CLC and the lower CLC, which in turn favors the hypothesis of delamination-induced destruction of the NCC.

5. Conclusions

We present here high-temperature and -pressure rheological experiments on synthetic felsic granulites. We varied strain rate and temperature to determine their effects on the measured stress. Our experiments and data reduction result in a stress exponent of $n = 4.2 \pm 0.1$ and an activation energy of $Q = 260 \pm 30$ kJ/mol. Microstructural analyses revealed enhanced CPO strength of pyroxene and plagioclase in the deformed samples compared to their counterparts in the hot-pressed ones, whereas quartz CPOs become random after deformation. In addition, all these minerals display distinct intracrystalline plasticity like undulose extinction and subgrain structure. These microstructural observations combined with the obtained stress exponent indicate that deformation was dominated by dislocation creep in both the rock and the component minerals. Extrapolations of our flow law imply a weaker continental lower crust compared to the upper crust and the uppermost mantle, supporting the “jelly sandwich” continental lithosphere strength model. In particular, this conclusion favors

the delamination model for the destruction of the NCC.

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848

FIGURE CAPTIONS

Figure 1. Microphotographs showing the typical shape of hot-pressed and deformed samples cutting along their longitudinal sections. (a) Hot-pressed (sample GA289); (b) Deformed (sample GA370). It can be inferred from these cross-sectional shapes that both the hot-pressed and the deformed samples should be nearly cylindrical.

Figure 2. Mechanical results of high-temperature creep of felsic granulite. (a) Typical stress-strain curves. Note that a steady-state deformation is generally reached after a strain of about 20%. (b) Log-log plot of strain rate versus stress for the steady-state flow of felsic granulite at 827 °C, 877 °C and 927 °C, yielding a stress exponents of $n = 4.2 \pm 0.1$ and a pre-exponential factor of $A = 10^{-4.4 \pm 0.2} \text{ MPa}^{-4.2} \text{ s}^{-1}$. (c) Semi-log plot of stress versus inverse temperature for constant strain rates of $9.4 \times 10^{-5} \text{ s}^{-1}$, $1.1 \times 10^{-5} \text{ s}^{-1}$ and $3.6 \times 10^{-6} \text{ s}^{-1}$, yielding an activation energy of $Q = 260 \pm 30 \text{ kJ/mol}$.

Figure 3. Representative microstructures of hot-pressed (a, b) and deformed (c, d) felsic granulite samples. (a and b) The hot-pressed sample (GA279) was annealed under a pressure of 1 GPa and 927 °C for 12 h. A weak SPO of aligned elongate pyroxene and amphibole grains can be discerned. (c and d) A more pronounced SPO of elongated pyroxene and amphibole, and sometimes quartz can be observed in a deformed sample (GA370). The green opposite arrows represent the longitudinal shortening direction during hot-pressing and deformation. (a and c) Backscattered electron images; (b and d) Photomicrographs in cross-polarized light. All scale bars are 100 μm . Mineral abbreviations here and hereinafter are after [Whitney and Evans \(2010\)](#).

Figure 4. Crystallographic preferred orientations of clinopyroxene, orthopyroxene, plagioclase, and quartz from a hot-pressed (a), and three deformed (b, c, and d) felsic granulite samples. Pole figures are lower hemisphere equal area projections and

contoured with a half width of 20° and a cluster size of 5°. For the hot-pressed sample, the longitudinal shortening direction (as indicated by the green opposite arrows) is vertical in each pole figure. For the deformed samples, the foliation plane (S) is E-W and normal to the paper in each pole figure. The contour interval is one multiple of uniform distribution. *n*, number of measured grains; *J*, J-index; *M*, M-index; U, upper-hemisphere projection; L, lower-hemisphere projection.

Figure 5. EBSD mappings showing intracrystalline deformation features (subgrain structure) in (a) plagioclase, (b) clinopyroxene, (c) orthopyroxene, and (d) quartz.

Figure 6. SEM micrographs showing typical grain boundaries of quartz with other minerals and triple junctions among plagioclase, quartz and pyroxene in the deformed samples at high magnifications. No melt is clearly visible along the grain boundaries and triple junctions with these magnifications.

Figure 7. Compositional mapping of Si, Al, Fe, Mg, and Ca in a deformed felsic granulite. Warmer or brighter colors indicate higher concentrations. Color scales in different X-ray maps do not correspond to the same element concentration. Note the clear-cut grain/phase boundaries and the absence of compositional zoning in minerals.

Figure 8. Representative FTIR spectra of plagioclase (a), quartz (b), clinopyroxene (c), and orthopyroxene (d) from the starting material. All spectra were normalized to 1 cm thickness and offset vertically for clarity.

Figure 9. Flow strength comparison of felsic granulite in this study (black dashed lines) with common crustal mineral aggregates (a) and other granulites (b) in the dislocation creep regime by extrapolating flow laws to a natural strain rate of 10^{-14} s^{-1} . Flow laws: wet and dry anorthite aggregates (Rybacki & Dresen, 2000); wet quartz

aggregates ([Rutter & Brodie, 2004](#)); wet clinopyroxene aggregates ([Chen et al., 2006](#));
wet enstatite aggregates ([Bystricky et al., 2016](#)); felsic and mafic granulite (WC1990)
([Wilks & Carter, 1990](#)).

Figure 10. Comparison of flow strength of our felsic granulite (black solid line) with
those of eclogite and olivine aggregates in the dislocation creep regime by
extrapolating flow laws to a natural strain rate of 10^{-14} s^{-1} . Mafic granulite from [Wang
et al. \(2012\)](#) is also shown to facilitate comparison. Flow laws for eclogite and
dry/wet olivine aggregates are taken from [Zhang and Green \(2007\)](#) and [Hirth and
Kohlstedt \(2003\)](#), respectively. The water concentration ($\sim 57 \text{ wt ppm H}_2\text{O}$) for wet
olivine aggregates was estimated using the water solubility equation of olivine from
[Zhao et al. \(2004\)](#) by assuming a fayalite number (X_{Fa}) of 0.1 under a pressure of 1
GPa and a temperature of 900 °C. This water concentration represents the maximum
water concentration that can be incorporated into olivine close to the Moho.

Figure 1.

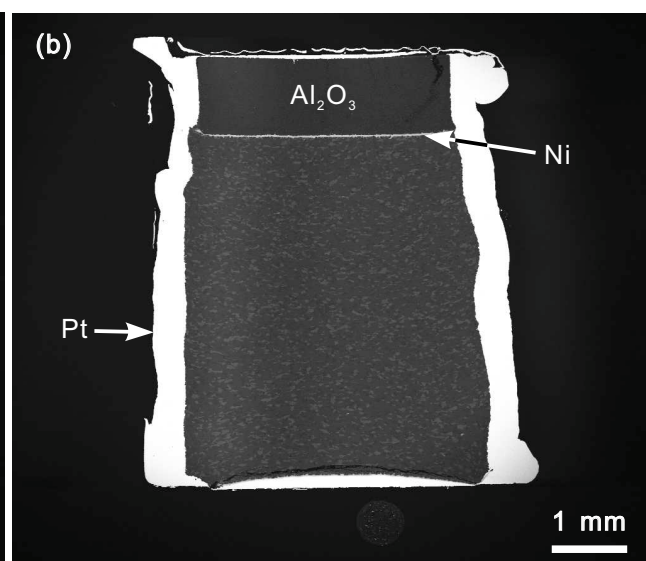
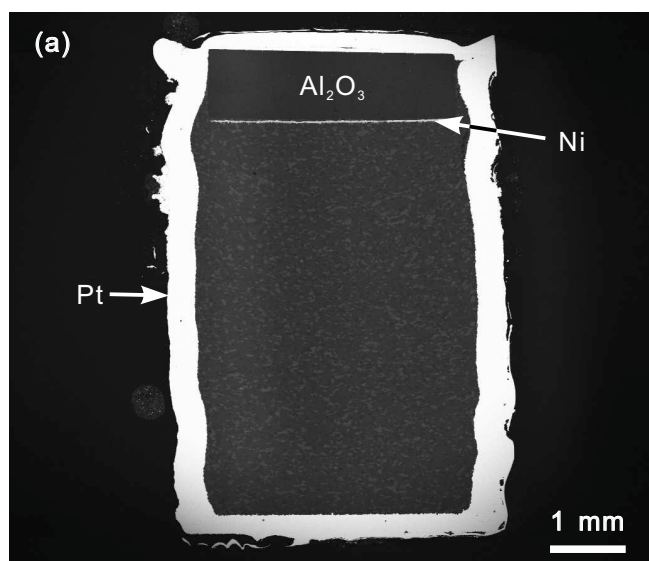


Figure 2.

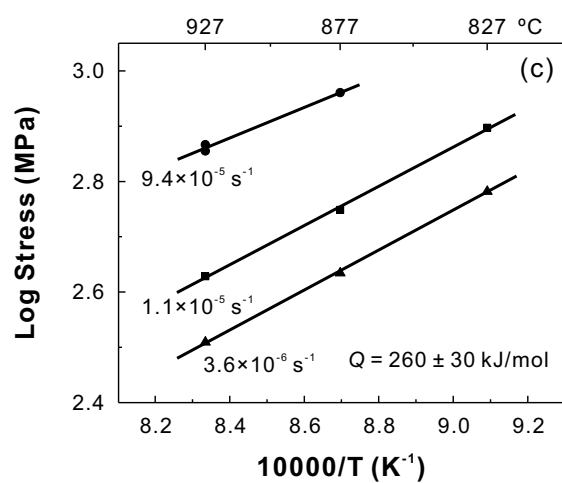
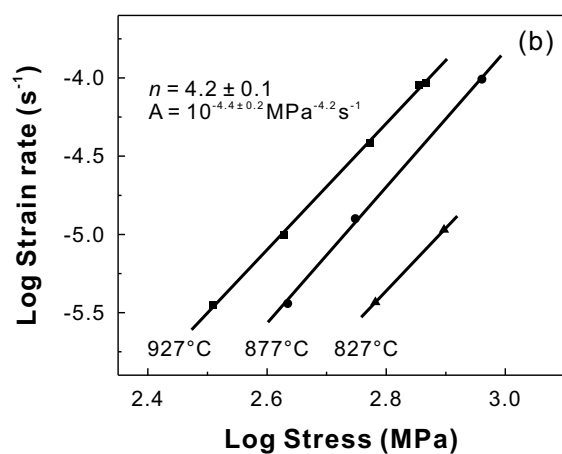
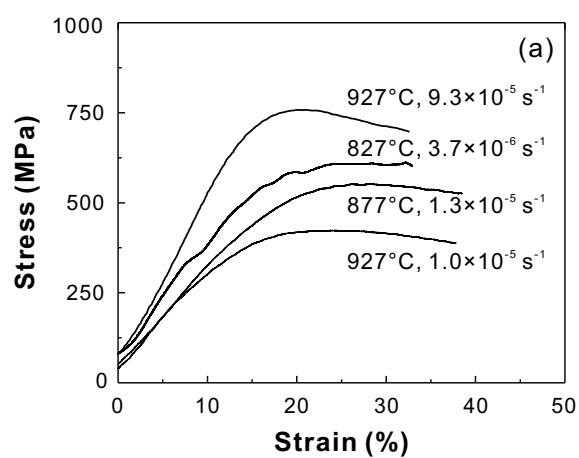


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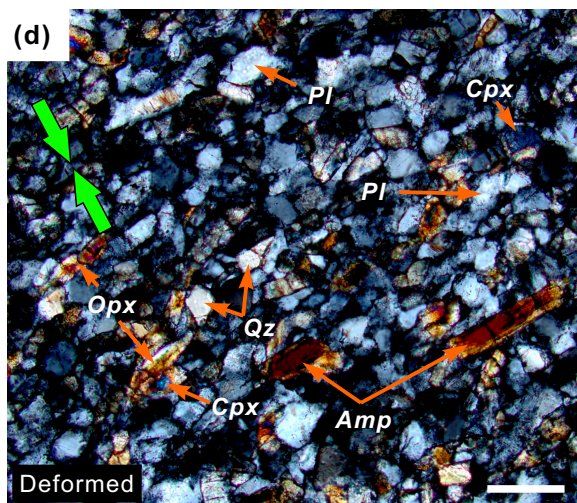
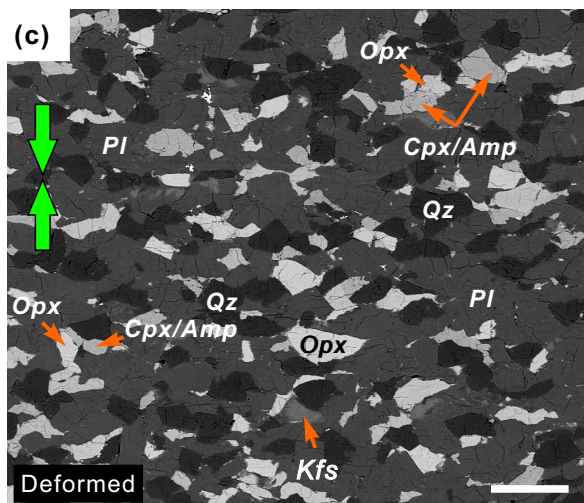
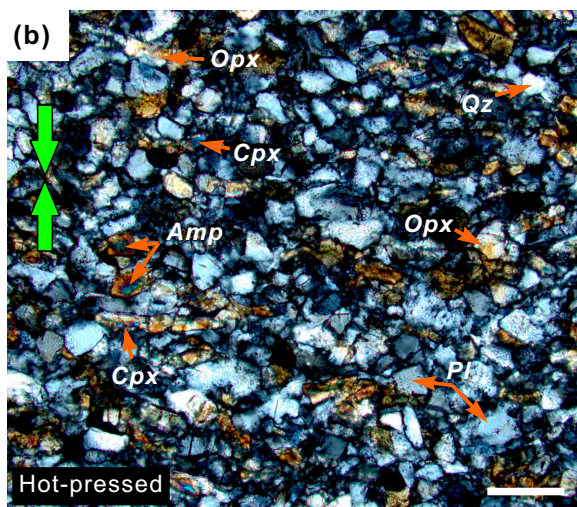
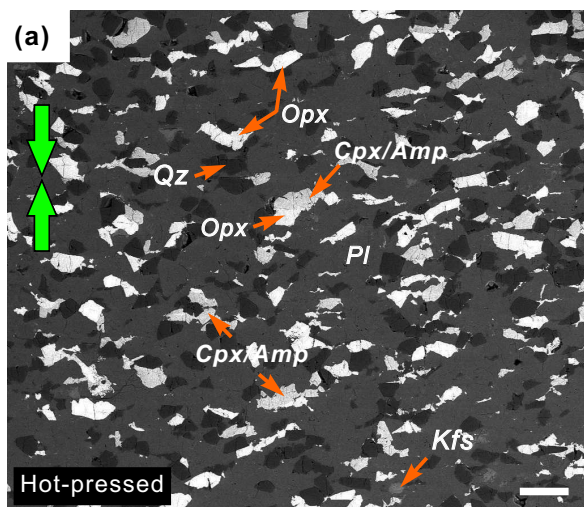


Figure 4.

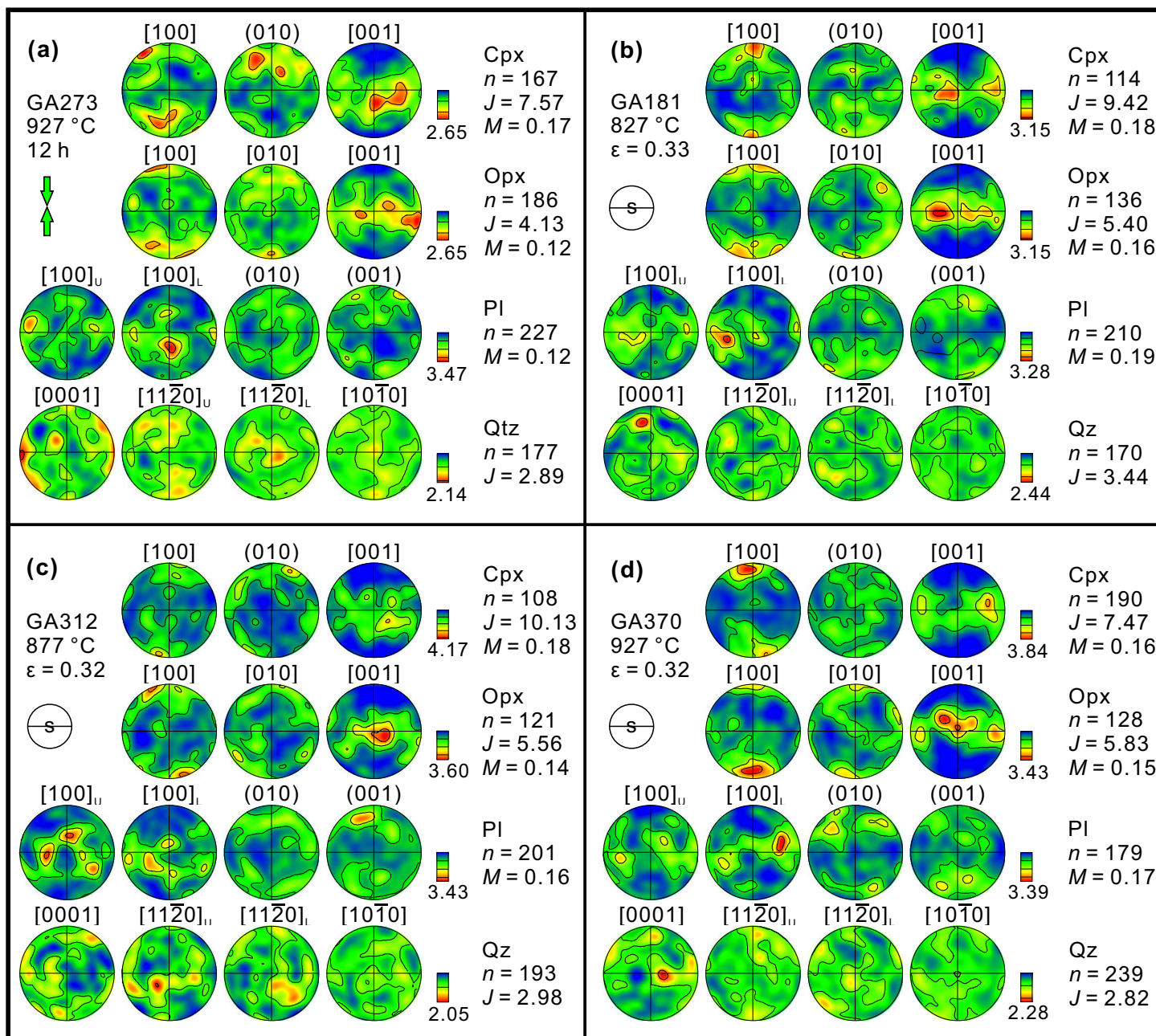


Figure 5.

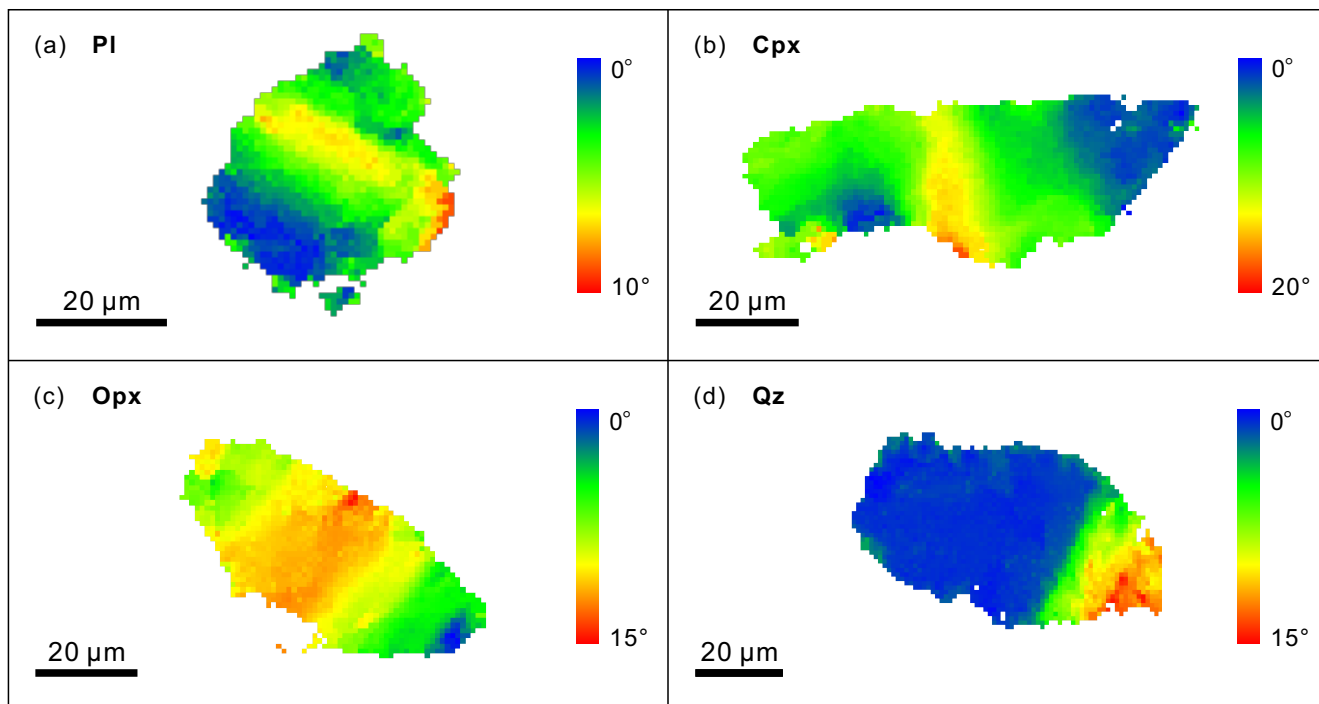


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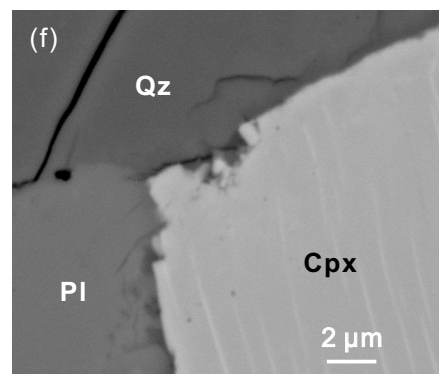
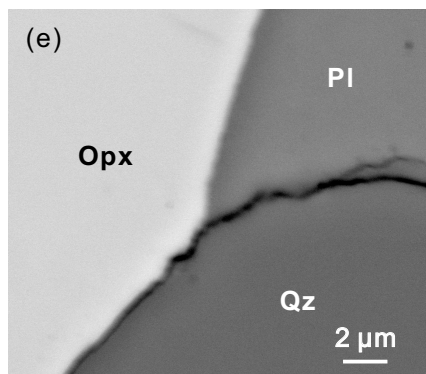
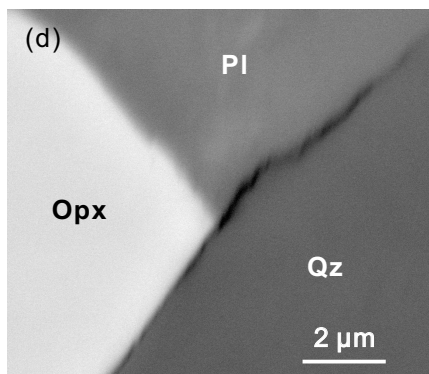
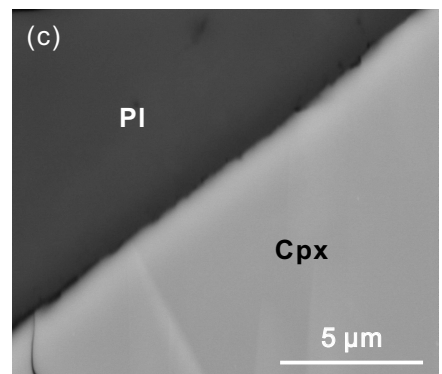
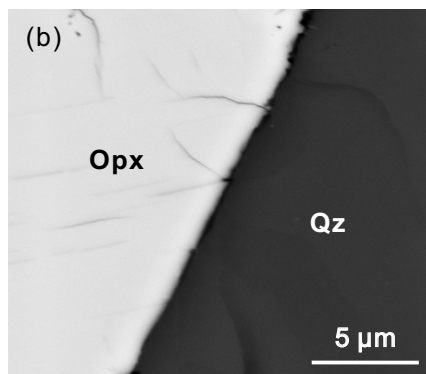
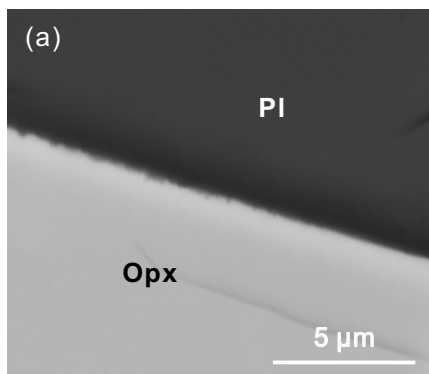


Figure 7.

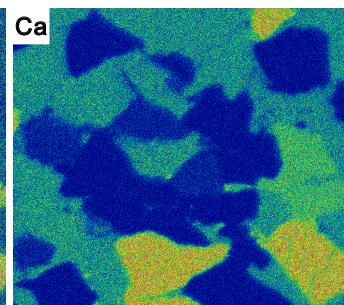
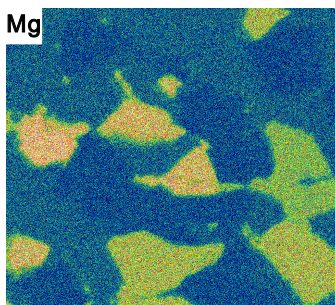
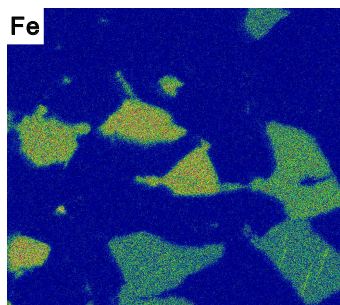
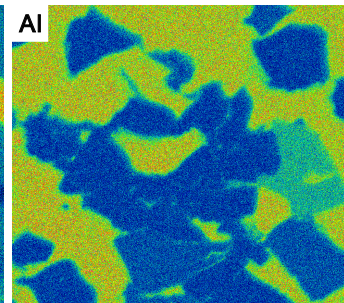
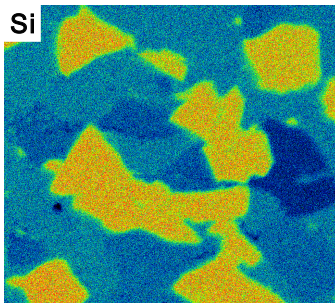
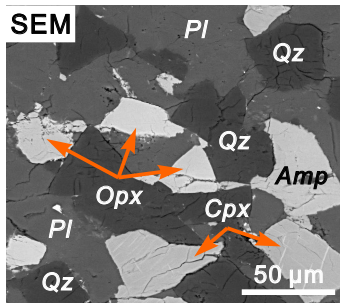


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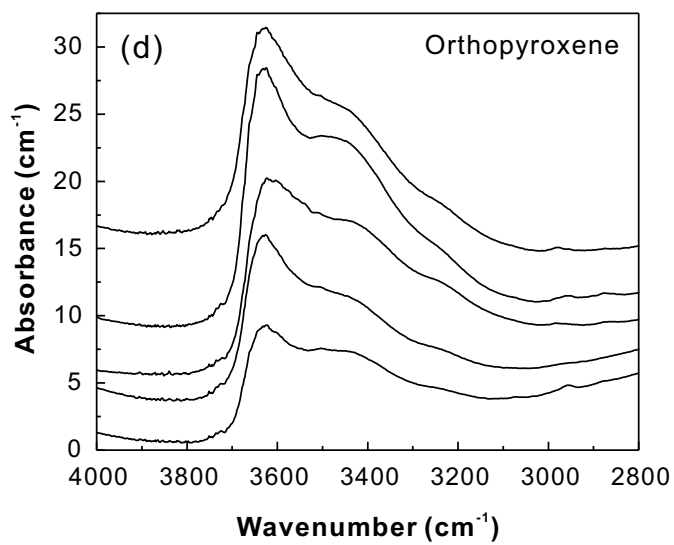
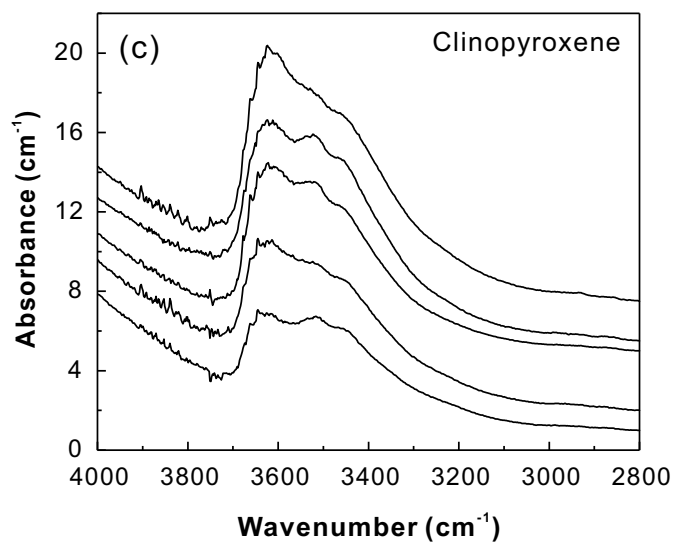
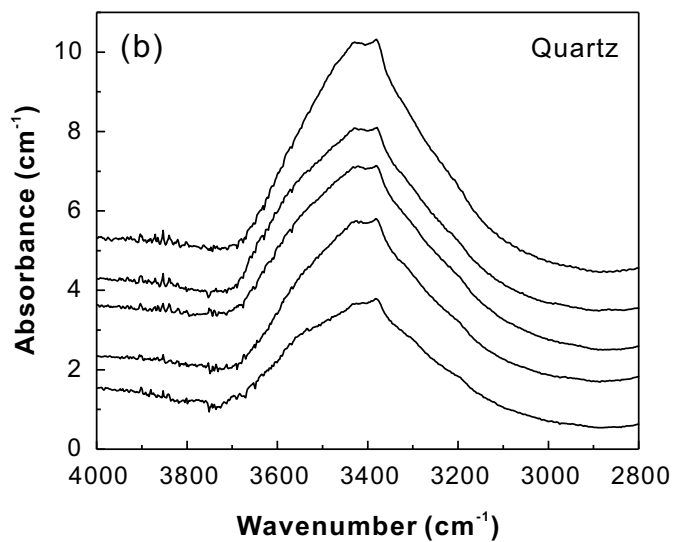
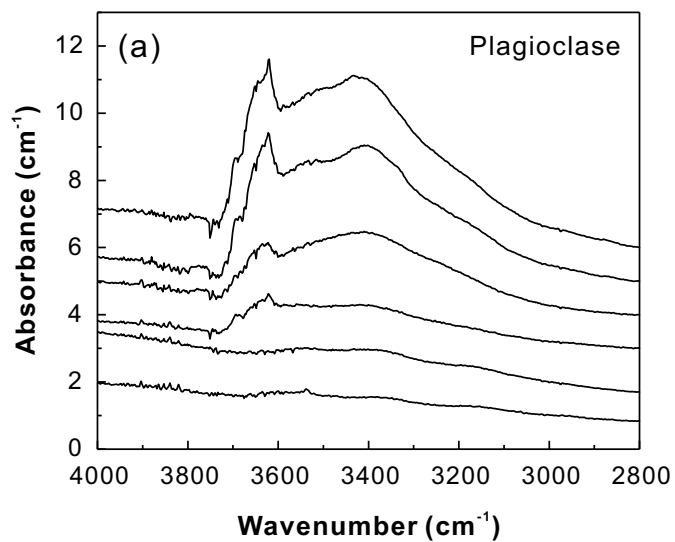


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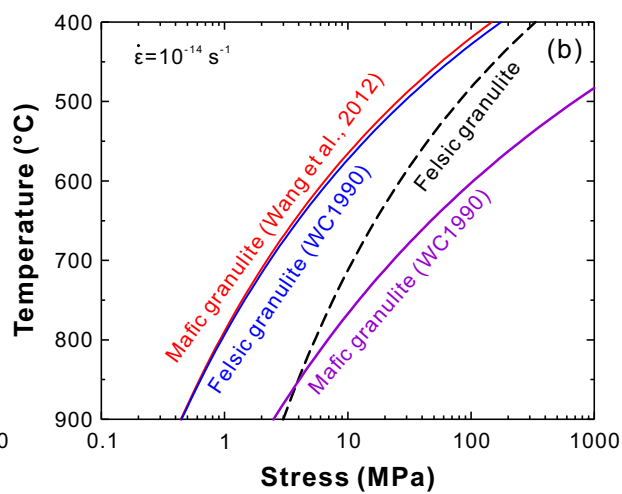
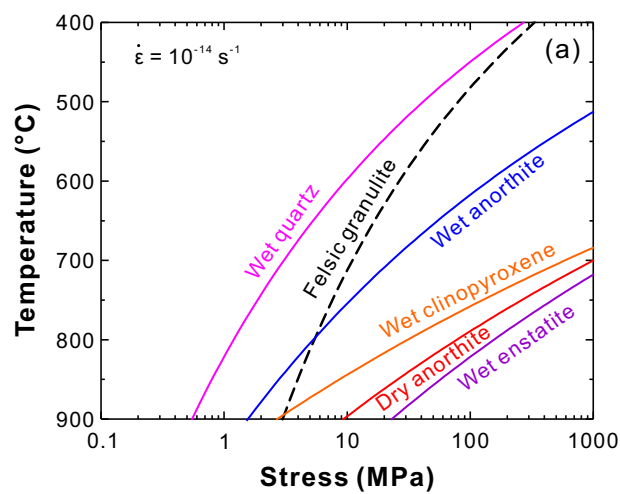


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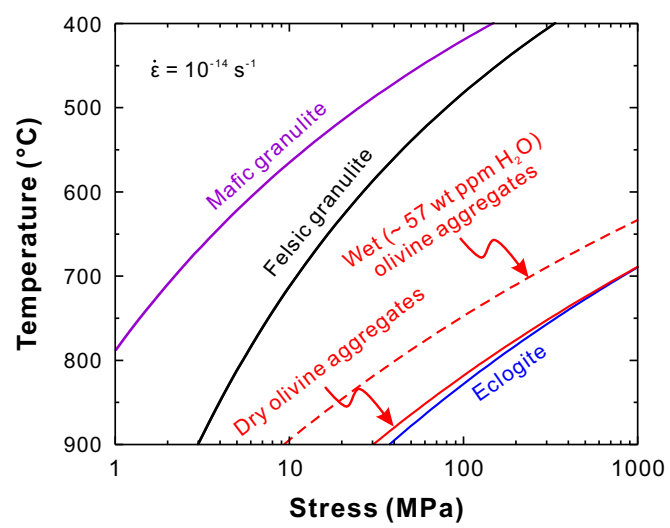


Table 1.
Experimental Conditions and Mechanical Data for Felsic Granulite

Expt. #	<i>P</i> (GPa)	<i>T</i> (°C)	<i>R</i> or ε (%) ^a	$\dot{\varepsilon}$ (s ⁻¹)	σ (MPa)
<i>Hot-pressing Experiments</i>					
GA273	1	927	35	-	-
GA279	1	927	37	-	-
GA283	1	927	38	-	-
<i>Deformation Experiments</i>					
GA331	1	927	34	3.6×10 ⁻⁶	323
GA321	1	927	38	1.0×10 ⁻⁵	425
GA139	1	927	49	3.8×10 ⁻⁵	592
GA327	1	927	42	9.0×10 ⁻⁵	716
GA370	1	927	33	9.3×10 ⁻⁵	735
GA317	1	877	37	9.8×10 ⁻⁵	913
GA312	1	877	38	1.3×10 ⁻⁵	560
GA206	1	877	36	3.6×10 ⁻⁶	431
GA181	1	827	33	3.7×10 ⁻⁶	605
GA161	1	827	33	1.1×10 ⁻⁵	789

^aAxial compression ratio and strain for deformed and hot-pressed samples, respectively.

Table 2.

Chemical Compositions (wt%) of Major Minerals in the Starting Material and Selected Hot-pressed and Deformed Granulite

Sample	Mineral	SiO ₂	TiO ₂	Al ₂ O ₃	Cr ₂ O ₃	FeO	MnO	MgO	CaO	Na ₂ O	K ₂ O	Total	End Member
09TPZ-1 Starting material	Pl	61.09	0.02	23.56	-	0.09	0.01	-	6.83	7.20	0.47	99.28	An ₃₃ Ab ₆₄ Or ₃
	Cpx	52.42	0.14	2.09	0.01	11.38	0.37	11.92	20.41	0.64	0.01	99.34	Wo ₄₅ En ₃₆ Fs ₁₉
	Opx	51.74	0.06	1.24	0.02	27.79	0.81	17.05	0.59	0.03	0.01	99.42	Wo ₁ En ₅₂ Fs ₄₇
	Amp	42.29	2.24	11.18	0.03	16.28	0.17	9.87	10.93	1.39	1.85	96.25	
GA283 Hot-pressed	Pl	60.18	0.04	24.29	-	0.13	0.02	0.01	7.65	6.30	0.65	99.28	An ₃₉ Ab ₅₇ Or ₄
	Cpx	52.50	0.22	1.94	0.02	11.63	0.38	11.66	20.50	0.47	0.01	99.35	Wo ₄₅ En ₃₆ Fs ₁₉
	Opx	52.27	0.07	0.92	0.03	27.32	0.68	17.62	0.59	0.03	0.01	99.53	Wo ₁ En ₅₃ Fs ₄₆
	Amp	42.35	2.55	10.90	0.07	17.78	0.19	8.95	11.12	1.45	1.94	97.31	
GA161 Deformed	Pl	60.30	0.04	24.39	-	0.19	0.01	0.01	7.51	6.36	0.53	99.35	An ₃₈ Ab ₅₉ Or ₃
	Cpx	52.63	0.22	1.99	0.03	11.74	0.36	11.50	20.38	0.51	0.00	99.38	Wo ₄₅ En ₃₅ Fs ₂₀
	Opx	52.66	0.10	0.99	0.04	26.18	0.65	17.98	0.66	0.01	0.01	99.29	Wo ₁ En ₅₄ Fs ₄₅
	Amp	43.03	2.19	11.19	0.01	18.02	0.18	8.76	11.12	1.23	1.86	97.59	
GA370 Deformed	Pl	59.78	-	24.70	0.01	0.20	-	-	7.83	6.23	0.48	99.24	An ₄₀ Ab ₅₇ Or ₃
	Cpx	52.17	0.22	1.98	0.02	12.04	0.34	11.63	20.43	0.48	0.01	99.31	Wo ₄₅ En ₃₅ Fs ₂₀
	Opx	51.48	0.11	0.94	0.05	27.75	0.58	17.50	0.67	0.02	0.01	99.12	Wo ₁ En ₅₃ Fs ₄₆
	Amp	43.40	1.80	12.02	0.01	14.82	0.16	10.18	11.40	1.24	1.95	96.98	

“-”, under detection limit.

Table 3.

Summary of FTIR Results for Major Component Minerals in the Starting Material

Mineral	Measured Grains	A_i (cm ⁻¹) ^a		C_{H2O} (wt ppm)	
		Range	Average ^b	Range	Average ^b
Pl	15	2.0-46.1	24.0 ± 14.4	17-401	209 ± 126
Qz	11	21.0-64.5	35.9 ± 12.9	208-636	355 ± 127
Cpx	11	10.3-97.8	49.8 ± 27.7	193-1839	937 ± 521
Opx	12	51.1-248.2	126.1 ± 61.8	459-2229	1132 ± 555

^a A_i represents the absolute integrated absorbance.^bErrors represent one standard deviation.