# A Comparison of Six Transport Models of the MADE-1 Experiment implemented with Different Types of Hydraulic Data

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#### Abstract

Six conceptually different models of steady groundwater flow and conservative transport are applied to the heterogeneous MADE aquifer. Their predictive capability is assessed by comparing the modelled and observed longitudinal mass distributions at different times of the plume in the MADE-1 experiment, as well as at a later time. The models differ in their conceptualization of the heterogeneous aquifer structure, computational complexity, and use of permeability data obtained from various observation methods (DPIL, Grain Size Analysis, Pumping Tests and Flowmeter). Models depend solely on aquifer structural and flow data, without calibration by transport observations. Comparison of model results by various measures, i.e. peak location, bulk mass and leading tail, reveals that the predictions of the solute plume agree reasonably well with observations if the models are underlined by a few parameters of close values: mean velocity, a parameter reflecting log-conductivity variability and a horizontal length scale related to conductivity spatial correlation. From practitioners perspective the robustness of the models is an important and useful property. The model comparison provides insight into relevant features of transport in heterogeneous aquifers. After further validation by additional field experiments or by numerical simulations, the results can be used to provide guidelines for users in selecting conceptual aquifer models, characterization strategies, quantitative models and implementation for particular goals.

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17	•	Predictions of transport in highly heterogeneous aquifer are achieved with six mod-
18		els of different conceptualization and field data input
19	•	Models' predictions are reliable if they share similar values of mean velocity and
20		permeability degree of variability and correlation scale
21	•	The MADE site plume longitudinal mass distribution is a robust transport mea-
22		sure

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Key Points:

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Six conceptually different models of steady groundwater flow and conservative transport 25 are applied to the heterogeneous MADE aquifer. Their predictive capability is assessed 26 by comparing the modelled and observed longitudinal mass distributions at different times 27 of the plume in the MADE-1 experiment, as well as at a later time. The models differ 28 in their conceptualization of the heterogeneous aquifer structure, computational com-29 plexity, and use of permeability data obtained from various observation methods (DPIL, 30 Grain Size Analysis, Pumping Tests and Flowmeter). Models depend solely on aquifer 31 structural and flow data, without calibration by transport observations. Comparison of 32 model results by various measures, i.e. peak location, bulk mass and leading tail, reveals 33 that the predictions of the solute plume agree reasonably well with observations if the 34 models are underlined by a few parameters of close values: mean velocity, a parameter 35 reflecting log-conductivity variability and a horizontal length scale related to conductiv-36 ity spatial correlation. From practitioners perspective the robustness of the models is 37 an important and useful property. The model comparison provides insight into relevant 38 features of transport in heterogeneous aquifers. After further validation by additional 39 field experiments or by numerical simulations, the results can be used to provide guide-40 lines for users in selecting conceptual aquifer models, characterization strategies, quan-41 titative models and implementation for particular goals. 42

#### 43 **1** Introduction

Modelling contaminant transport by groundwater is a topic of great interest that stimulated intensive research in the last four decades due to its relevance to aquifer pollution (Dagan, 1989; Gelhar, 1993; Fetter et al., 2018). The task of predicting transport faces a few difficulties: the processes are of long duration, measurements are scarce, the subsurface medium is of complex heterogeneous structure subjected to uncertainty and many times the geometry and the mass content of the contaminant source is also not known with certainty.

<sup>51</sup> Under these circumstances models play an important role: they help understand-<sup>52</sup> ing the involved processes, analysing field data and making long range prediction. Mod-<sup>53</sup> els developed in the past differ in conceptualization of the aquifer structure, in the re-<sup>54</sup> quired data, in quantification of transport, in the formulation of the governing equations <sup>55</sup> and mechanisms they represent, in computational complexity and in models goals.

We focus on transport of plumes of conservative solutes in steady natural gradient flow, driven by a constant mean head gradient. Quantification of the spatial distribution is by m(x,t), mass per unit length, where x is the mean flow direction and t the time. It encapsulates the process of longitudinal spreading in space and time. In practice it allows, for instance, to estimate the mass of solute pumped by wells or flowing into rivers or reservoirs. It also serves as a first step toward achieving other goals like determining the local concentration  $C(\mathbf{x}, t)$ .

There is general agreement that spatial variability of the hydraulic conductivity  $K(\mathbf{x})$  is the main growth mechanism responsible for plume spreading in aquifers, termed macrodispersion (Zech et al., 2015). The effect increases with higher K heterogeneity as quantified for instance by the log-conductivity variance  $\sigma_Y^2$ .

A few elaborate field experiments were conducted in the past in order to gain understanding and validate models. The most challenging one is at the *MADE* site (Boggs et al., 1992), situated in a highly heterogeneous aquifer, making it of relevance to many actual aquifers (Gomez-Hernandez et al., 2017). An important feature of *MADE* was the application of different observation methods in order to characterize the K spatial distribution (e.g. Boggs et al. (1990); Rehfeldt et al. (1992); Bohling et al. (2016)). Longterm tracer tests provide graphs of observed mass distribution m(x, t) in space at a few fixed t values (Boggs et al., 1990, 1995). It has motivated a flurry of works on structure
characterization by field data and different modelling strategies (see e.g. Zheng et al. (2011)).
We will elaborate on MADE in section 2.

We examine the ability of six conceptually different models to predict the observed 77 mass m. Subsequently, we examine the predictive power of the models by extending the 78 time (t = 1000 days) beyond MADE observations ( $t_{\text{max}} = 503$  days). The selected 79 models, outlined in Section 3, differ in conceptualizations of formation structure and trans-80 port, in the use of field data, in initial conditions, in computational methodology and 81 82 in effort. A few of the models were developed in the past while the others were formulated for use in the present paper. The models cover a wide spectrum of configurations. 83 We concentrate here only on models that can predict transport based on field data of 84 aquifer properties and flow; we do not consider models calibrated on prolonged trans-85 port tests which are generally of large duration and cost. 86

<sup>87</sup> We believe that the comparison of the models which differ in type and underlying <sup>88</sup> field data, conceptualization of K spatial structure and complexity of computations by <sup>89</sup> using *MADE* as a platform is important in helping the research community to grasp trans-<sup>90</sup> port issues and the users in selecting characterization strategies, goals of models and method <sup>91</sup> implementation.

The plan of the paper is as follows: Sect. 2 recapitulates the MADE aquifer characterization and transport experiment; Sect. 3 describes the methodology of the models, including their application to MADE; Sect. 4 is devoted to the model prediction for solute mass m in comparison to MADE observations as well as time beyond. Sect. 5 contains the general discussion on data comparison while Sect. 6 concludes the paper.

#### 97 2 MADE Transport Field Experiment

The *MADE* experiment was the object of a large body of publications dealing with the aquifer properties data collection and analysis as well as transport observations and interpretation (see for instance reviews Zheng et al. (2011); Gomez-Hernandez et al. (2017)). We recall in the following only those aspects of *direct* relevance to the examined models.

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#### 2.1 Hydraulic Conductivity Spatial Distribution

The MADE site aquifer is composed of highly heterogeneous alluvial terrace de-104 posits. Measurements of hydraulic conductivity at multiple locations (see Fig. 1) were 105 performed by granulometry of soil samples, flow meter, slug test, and Direct Push In-106 jection Logger (DPIL) (Boggs et al., 1990; Rehfeldt et al., 1992; Bohling et al., 2016). 107 Besides, two pumping tests provide equivalent conductivities  $K_{eq}$  of the volume surround-108 ing the wells (Boggs et al., 1992). The use of different techniques at same site and sub-109 sequent application by different models offers an unique opportunity to examine their 110 impact on transport prediction. 111

The spatially distributed observations (Fig. 1), carried out at different depths resulted in a large volume of data which served for geostatistical analysis. Fig. 2 summarizes their outcomes, which is of interest for the application of the different models (Sect. 3). The most reliable and extensively used data are those based on flowmeter (Boggs et al., 1990; Rehfeldt et al., 1989) with a total number of N = 2611 observations and more recently DPIL with N = 31123 (Bohling et al., 2012, 2016).

The differences in the geostatistical parameters, especially in the geometric mean  $K_{\rm G}$ , are a result of the different properties of the observation methods as well as the density and locations of observation points. Particularly, the difference between flowmeter and DPIL can be explained by the inability of flowmeter to detect low K values. While



**Figure 1.** Left: Map of MADE site according to Boggs et al. (1990); Bohling et al. (2016): Locations of hydraulic conductivity measurements devices (coloured dots); tracer test source area (black outline); and sampling network boundary (dashed outline) where bromide samples were collected. Right: Potentiometric surface map of head measurements according to Boggs et al. (1990). Black dot marks tracer test injection location.

the maximal values  $K_{\text{max}}$  are close (around  $0.005 \,\text{m/s}$ ) the minimal values  $K_{\text{min}}$  differ 122 by two orders of magnitude. The effect manifested in Fig. 2 is that flowmeter  $K_{\rm G}$  is larger 123 and  $\sigma_V^2$  is smaller. The difference in the longitudinal integral scales  $I_h$  suggests that the 124 zones of low K values are less connected than those of higher magnitude. The analysis 125 based on soil samples (N = 214) is less reliable. Still, it provided input data for con-126 ductivity conceptualization based on the lithofacies approach (e.g. Carle & Fogg (1996)) 127 as presented by Bianchi & Zheng (2016) and herein for a simplified binary structure ap-128 proach. The impact of these difference upon flow and transport are discussed in Sect. 4. 129

#### 2.2 Flow

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In Fig. 1 we reproduce the head contour lines map (Boggs et al., 1992, Fig. 3). The head gradient is not constant, but slowly varying in space. The mean head gradient is between  $J \in [0.003, 0.0036]$  depending on the choice of boundary locations. The nonuniformity of the head contour density indicates the presence of large scale mean hydraulic conductivity trends.

#### 2.3 Transport Experiment

We focus on the first tracer transport experiment, which was conducted in years 1986–1988 (Boggs et al., 1990, 1992; Rehfeldt et al., 1992; Adams & Gelhar, 1992). The tracer plume displayed a non-Gaussian longitudinal solute mass distribution with the bulk of the mass staying near the source, but with lower amounts spreading downgradient extensively.



Figure 2. Geostatistical measures for MADE from DPIL (direct push injection logging) (Bohling et al., 2016), flowmeter, grain size analysis, slug tests (Rehfeldt et al., 1992): log-conductivity variance  $\sigma_{\ln K}^2$ , horizontal and vertical integral scale (correlation length) *I* and  $I_v$ , respectively. Visualization of geometric mean conductivity  $K_G$ , range of observed values from minimal to maximal ( $[K_{\min}, K_{\max}]$ ), and range of one variance around mean ( $[K_G \cdot e^{-\sigma^2}, K_G \cdot e^{\sigma^2}]$ ).  $K_{eq}$  denotes the equivalent conductivity for two large scale pumping tests (Boggs et al., 1990).

Initial Conditions: A quantity of  $M_0 \cong 25$  kg of Bromide dissolved in 10 m<sup>3</sup> 142 of water was injected over a period of 48 hours through 5 wells of half a meter screen length. 143 While the subsequent transport took place under practically mean uniform flow condi-144 tions, the tracer solution was forced radially into the aquifer during injection. As a re-145 sult, the initial tracer body was much larger than the body adjacent to the wells screens, 146 which can be seen in the early tracer plume snapshot at 9 days after injection (Adams 147 & Gelhar, 1992, Fig.5). This is important as far as ergodicity and setup of initial con-148 ditions are concerned. The apparent upstream tracer spread cannot be interpreted as 149 a result of upstream dispersion. The injection mode also implies that the initial condi-150 tion was flux proportional, with a preference of mass flowing in high conductivity chan-151 nels. 152

Plume Detection and Data Aggregation (Upscaling): We reproduce the longitu-153 dinal mass distribution  $\bar{m}(x,t)$  of Adams & Gelhar (1992, Fig.7) at six times T: 49, 126, 154 202, 279, 370, and 503 days after beginning of injection. The computation of  $\bar{m}$  is based 155 on concentration C(x, y, z, t) sampled in a dense MLS network, which thins out with dis-156 tance to the source (Fig. 1). Subsequently, C was numerically integrated over transverse 157 planes (y, z), accumulated and averaged over slices of 10 m length in the x direction to 158 obtain the upscaled longitudinal mass  $\overline{m}(x,t)$  (Adams & Gelhar, 1992). The reported 159 mass is displayed at the centers of the slices at  $x = -5 \,\mathrm{m}, 5 \,\mathrm{m}, 15 \,\mathrm{m}, \dots$  (see discus-160 sion of difference between the fine scale m and the upscaled one  $\bar{m}$  in the sequel). 161

The fact that the reported mass is an upscaled/aggregated quantity, was overlooked by many previous studies which compared  $\bar{m}$  with modeled mass at fine scale, as mentioned by Fiori et al. (2019). The significance of data aggregation is discussed in section 4 herein.

166 Mass Recovery The reported mass  $\bar{m}$  does not obey the mass conservation re-167 quirement  $\int \bar{m}(x,t) dx = 1$  except at 126 days after injection. Mass apparently decreases 168 over time after 126 days with recovery rates of 2.06, 0.99, 0.68, 0.62, 0.54, and 0.43, for 169 the T = 49, 126, 202, 279, 370, and 503 days, respectively.

As discussed in Adams & Gelhar (1992), the excessive mass recovery at t = 49days could be due to spurious hydraulic connections among the multilevel samplers. This is a result of the method of installation and is enhanced by the pressure of injection. Preferential sampling from high conductivity regions is also possible. In such case, the assumption of uniform tracer distribution employed in the spatial integration easily leads
to an overestimated mass recovery.

We attribute the apparent mass loss at later times to insufficient sampling in the downstream zone in line with Fiori (2014). The high heterogeneity implies possible channeling effects at MADE. In combination with the much lower sampling density downstream, it is generally difficult to sample the leading edge of the plume, which may contain a significant fraction of the plume mass (Fiori et al., 2019).

<sup>181</sup> Unlike some authors, we do not normalize observed mass  $\bar{m}$  by the total reported <sup>182</sup> mass i.e.  $\tilde{m} = \bar{m}(x,t) / \int \bar{m}(x,t) dx$ , but use  $\bar{m}$  for comparison with theoretical mod-<sup>183</sup> els, as it is appropriate for a conservative solute. Normalizing would imply that mass ap-<sup>184</sup> parent loss is proportional to the observed mass.

#### **3** Subsurface Transport Models

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#### 3.1 Common Features and Prerequisites

We compare various models which were developed independently by the authors in the past or devised recently in the frame of the present study. Before describing the models specific properties, we recapitulate a few common features:

190	1.	We examine only predictive models. They rely solely on structural and flow data
191		which can be measured independently of transport. Models with parameters cal-
192		ibrated by transport experiments are not considered.
193	2.	Flow is steady and uniform in the mean (natural gradient), driven by the steady
194		head gradient $J$ in the $x$ direction.
195	3.	Transport is advective and spreading is caused by the spatial variability of $K(\mathbf{x})$ .
196		The effective porosity $\theta$ is assumed to be constant.
197	4.	A plume of mass $M_0$ of a conservative solute is injected initially in the aquifer at
198		$t = 0$ . Spreading is quantified by the mass arrival at a control plane at x: $M_{total}(x, t) = 0$
199		$\theta \int_x^\infty \int \int C(x', y, z, t) dx' dy dz$ , where C is the concentration. For a fixed $x, M =$
200		$M_{total} / M_0$ is the BTC whereas for a fixed t it is the relative mass accumulated
201		beyond $x$ .
202	5.	We determine the fine scale relative mass spatial distribution $m(x,t) = -\partial M/\partial x$
203		at a few times $t$ . However, in line with $MADE$ observations, we calculate the (up-
204		scaled) relative mass averaged over a medium slice of length $\Delta = 10 \mathrm{m}$ centered
205		at x, which is given by $\overline{m}(x,t;\Delta) = (1/\Delta)[M(x-\Delta/2,t)-M(x+\Delta/2,t)]$ such
206		that $m = \lim_{\Delta \ll I} \bar{m}$ .
207	6.	The upscaled relative mass $\bar{m}$ is derived for the <i>MADE</i> conditions at $t = 49, 126$ ,
208		202, 279, 370, and 503 days after injection toward comparison with measured $\bar{m}$ .
209		Additionally, models are applied to prediction of $\bar{m}$ at $t = 1000$ days for inter-
210		comparison. As useful additional quantification we also consider the mass flux through
211		the control plane, $\mu(x,t) = \partial M / \partial t$ .
212	1 • •	Remark on additional MADE transport Models: Several other transport models
213	which	are not discussed here have been presented in the literature for <i>MADE</i> . We do not
214	consid	ler models calibrated on transport observations since they are not predictive, such
215	as the	work of Barlebo et al. (2004). This further includes dual-domain models (e.g. Har-

vey & Gorelick (2000) and Feehley et al. (2000)) and the continuous time random walk

(CTRW) model of Berkowitz & Scher (2001). The lithofacies approach of Bianchi & Zheng
(2016) would have been a candidate but results are only available for the *MADE-2* tracer

experiment setting rather than *MADE*-1 considered here. The same holds for the work

of Salamon et al. (2007). The conceptual framework is however underlying the facies model herein. Similarly, the fractional ADE model of Benson et al. (2001) is applied only to

MADE-2 data. Furthermore, not all parameters, such as the skewness parameter  $\beta$ , are 222 fully predictive based on structural data only. Dogan et al. (2014) presented a predic-223 tive fully numerical transport model for the MADE-1 experiment based on flowmeter 224 and DPIL measurements. They generated detailed representations of the K field con-225 ditioned to observations in a small sector of the MADE site aquifer and subsequently 226 solved the flow and transport equations. However, due to the computational effort, trans-227 port simulation results are limited to a fraction of the total plume transport distance. 228 Thus, their results are unsuitable for model comparison, particularly with prediction be-229 yond observation times. 230

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#### 3.2 Brief Description

#### 3.2.1 First Order Approximation (FOA)

Background The solution of flow and transport in heterogeneous formations of 233 a random  $Y = \ln K$  structure by a first order approximation in the log-conductivity 234 variance  $\sigma_Y^2$  was the topic of intensive research in the last four decades (see e.g. the mono-235 graphs Dagan (1989); Gelhar (1993); Rubin (2003)) leading to various analytical solu-236 tions. We briefly recall past results and recent advances relevant to this work. 237

K-Structure The random Y field is regarded as stationary and multi-Gaussian. 238 It is characterized completely by  $K_G$ ,  $\sigma_Y^2$ , and the two point covariance  $C_Y$  of horizon-239 tal integral scale I and vertical one  $I_v$ . The anisotropy ratio  $e = I_v/I$  is generally smaller 240 than unity. 241

Flow The mean velocity is given by  $U = K_{\text{eff}} J/\theta$ , with the effective conductiv-242 ity  $K_{\text{eff}}/K_G = \text{func}(\sigma_Y^2, e)$  determined from the solution of the flow equations for the 243 random velocity field. The latter is obtained by expanding the mass conservation equa-244 tion and Darcy's law in power series in  $Y' = Y - \langle Y \rangle$ . 245

Transport Traditionally, the mean relative mass distribution  $\langle m(x,t) \rangle$  was derived for conditions of given initial deterministic resident concentrations  $C_0(\mathbf{x}, 0)$  either in a volume  $V_0$ , with  $C_0 = M_0/(\theta V_0)$ , or with mass concentrated on the plane x = 0over an area  $A_0$  and quantified by  $m_0 = M_0/A_0$  (Kreft & Zuber, 1978). Similarly, detection was in the resident mode with the mean  $\langle M \rangle$  and  $\langle m \rangle$  satisfying the ADE

$$\frac{\partial \langle M \rangle}{\partial t} + U \frac{\partial \langle M \rangle}{\partial x} = D_L(t) \frac{\partial^2 \langle M \rangle}{\partial x^2} \tag{1}$$

The macrodispersion coefficient  $D_L = U\alpha_L$ , with longitudinal macrodispersivity  $\alpha_L$ , 246 was determined in the Lagrangean framework with the aid of the solute particles tra-247 jectories. At first order  $\alpha_L(t)$  grows from zero at t=0 to an asymptotic constant value 248  $\alpha_L = \sigma_Y^2 I$  (Dagan, 1989). The transient  $\alpha_L(t)$  was determined by a quadrature for an 249 exponential covariance and is approximated accurately by an analytical expression as 250 function of mean flow velocity, time and aquifer statistics:  $\alpha_L(t)/\sigma_V^2 I = \text{func}(Ut/I, e)$ 251 (Dagan & Cvetkovic, 1993, Eq. 20). The asymptotic value is attained after a travel dis-252 tance of a few integral scales I. The Gaussian solution approximates  $\langle m \rangle$  satisfactorily 253 for field experiments in weakly heterogeneous aquifers like Cape Cod (Hess et al., 1992) 254 and Borden Site (Sudicky, 1986) and/or far from the injection zone. However, it failed 255 to model the highly skewed mass distribution observed at MADE close to the injection 256 zone. This finding has motivated development of new nonlinear models. 257

Fiori et al. (2017) presented the solution of the ADE (1) for the more realistic conditions of flux proportional injection. The solution for the BTC is given by the inverse Gaussian distribution:

$$\langle M \rangle = \frac{1}{2} \left\{ \operatorname{erfc}\left(\frac{x - Ut}{2\sqrt{D_L t}}\right) + \exp\left(\frac{Ux}{D_L}\right) \operatorname{erfc}\left(\frac{x + Ut}{2\sqrt{D_L t}}\right) \right\}$$
(2)

Unlike the Gaussian solution of Eq. (1),  $\langle m(x,t) \rangle = -\partial \langle M \rangle / \partial x$  from Eq. (2) displays a skewed shape and lack of upstream dispersion. Jankovic et al. (2017) showed that the solution (2) with dispersion by first-order approximation  $D_L(x)$  represents accurately the results of numerical simulations even for the large value of  $\sigma_Y^2 = 8$  except an underestimation of a few percents of the mass in the tail of late arrival. We apply Eq. (2) to predict the *MADE* plume in the sequel as part of models comparison.

Application to MADE Application of Eq. (1) to MADE was based on the param-264 eters derived from Bohling et al. (2016) by DPIL as follows:  $K_{\rm G} = 0.58 m/d$ ,  $\theta = 0.31$ , 265  $I = 9.1m, I_v = 1.8m, \sigma_Y^2 = 5.9$ . Rather than the first order approximation we used 266 the more general formula by Zarlenga et al. (2018, Eq. 5) to arrive at  $K_{\text{eff}} = 2.28 \text{ m/d}$ . 267 After identifying the representative mean head gradient in the plume zone (Fig. 1) as 268 J = 0.0036, the mean velocity is given by  $U = K_{\text{eff}} J/\theta = 0.026 \,\text{m/d}$ . Subsequently, 269 the transient regime is taken into account by using a preasymptotic Dispersion  $D_{\rm L}(x)$ 270  $=U\alpha_{\rm L}(x)$ , calculated according to Fiori et al. (2019, Eqs. C1,C2) based on  $\alpha_{\rm L}$  of (Da-271 gan & Cvetkovic, 1993, Eq. 20), with the asymptotic value  $\alpha_{\rm L} = \sigma_V^2 I = 53.7m$ . This 272 was the information needed in order to derive  $\langle \bar{m} \rangle$  based on Eq. (1) at the t and x val-273 ues pertinent to MADE. 274

#### 3.2.2 Multi-indicator Model and Self Consistent Approximation (MIM-SCA)

Background MIMSCA was developed in the last 15 years as an approximate model of flow and transport for aquifers of arbitrary degree of heterogeneity (Dagan & Fiori, 2003; Fiori et al., 2006; Cvetkovic et al., 2014). Its outcome has been compared with accurate numerical solutions for  $\sigma_Y^2 \leq 8$  and applied to MADE. Fiori et al. (2019) recently applied the model to assess the uncertainty of prediction due to non-ergodic conditions or parametric uncertainty.

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K - Structure The aquifer is modeled as an ensemble of rectangular inclusions tessellating the space similarly to layers of bricks. The elements are of dimension  $2I \times 2I \times 2I_v$ . The block K values are assigned independently with random values from a univariate pdf f(K) (multi-indicator model). As usual, f(K) was chosen to be lognormal of parameters  $K_G, \sigma_Y^2$ . This way the random K field is completely defined in terms of 4 parameters, similarly to the FOA (section 3.2.1).

Flow The mean velocity is given by  $U = K_{\text{eff}}J/\theta$ . The mean effective conductivity  $K_{\text{eff}}$  is derived by the self consistent approximation (SCA), a well established method in the literature on heterogeneous aquifers (Dagan, 1989). Here, it consists in solving the flow equations for a generic inclusion of conductivity K, submerged in a homogeneous matrix of the unknown  $K_{\text{eff}}$  and determining the latter by the SCA argument.  $K_{\text{eff}}$  follows as solution of a simple integral equation. Suribhatla et al. (2011) compared it with accurate numerical simulations with satisfactory agreement.

Transport Transport for the MIM was solved also with the SCA. Fiori et al. (2003) 296 determined analytically the travel time required for a solute particle to move over an in-297 clusions of conductivity K, surrounded by a matrix of  $K_{eff}$ . The BTC  $\langle M \rangle$  follows as 298 sum over the travel time random residuals pertaining to the inclusions of different K ly-299 ing between the injection plane x = 0 and the control plane at x and the initial con-300 dition was of flux proportional injection at x = 0 while the mean BTC was derived by 301 fast Fourier transform. Jankovic et al. (2003) showed a satisfactory agreement between 302 the semi-analytical results and accurate 3D numerical simulations. Fiori et al. (2017) showed 303 that the bulk of the BTC is well approximated by the FOA, while MIMSCA captures 304 also the long tail of a few percents of the solute mass observed in numerical simulations. 305

Application to MADE Fiori et al. (2013) applied the method to the MADE site transport setting for  $\langle m \rangle$  and for  $\langle \bar{m} \rangle$ , with an update in Fiori et al. (2019) motivated by an update in geostatistical input parameters by Bohling et al. (2016). The needed parameters values are the same as those given above for FOA.

#### 3.2.3 TDRW

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Background Time-domain random walk and continuous time random walk ap-311 proaches have been used extensively over the past two decades for the modeling of trans-312 port in heterogeneous porous media (Noetinger et al., 2016). Within this framework and 313 based on the stochastic TDRW method of Comolli et al. (2019), Dentz et al. (2020) de-314 rived a predictive upscaled model that avoids calibration by transport observations. The 315 basic idea is to quantify particle motion in spatially variable flow fields through a Markov 316 processes for equidistant particle velocities, whose steady state distribution is given by 317 the flux-weighted distribution of flow velocities. While details can be found in Dentz et 318 al. (2020), we describe here the main features of the model. This modeling approach has 319 been used and verified for the prediction of the evolution of particle velocity statistics, 320 particle distributions, dispersion and breakthrough curves in pore and Darcy scale het-321 erogeneous porous media (Hakoun et al., 2019; Comolli et al., 2019). 322

<sup>323</sup> *K-Structure* Hydraulic conductivity is represented by a three-dimensional log-<sup>324</sup> normally distributed multi-Gaussian spatial random field. Thus, the random *K*-field is <sup>325</sup> characterized in terms of 4 parameters, similarly to the *FOA* (section 3.2.1): geometric <sup>326</sup> mean conductivity  $K_{\rm G}$ , ln *K* variance, and correlation lengths  $\ell_h$  and  $\ell_v$ . Random re-<sup>327</sup> alizations were filtered such that the spatial mean and variance of the log hydraulic con-<sup>328</sup> ductivity are within a 5% tolerance interval around the target values.

Flow Model Groundwater flow is the result of the steady state groundwater flow 329 equation, which is solved numerically on multi-Gaussian hydraulic conductivity fields char-330 acterized by a log-normal marginal distribution. A unit head drop between inlet and out-331 let is considered and no-flux boundaries are specified at the horizontal domain bound-332 aries. The target variable is the magnitude of the Eulerian velocity  $v_e(\mathbf{x})$  (absolute value 333 of the Eulerian velocity), which is characterized by a uni-variate distribution. It is ob-334 tained from the numerically obtained magnitude of the Darcy velocity  $q(\mathbf{x})$  by multipli-335 cation with the magnitude J of the head gradient, and geometric mean conductivity  $K_{\rm G}$ , 336 and division by (constant) porosity  $\theta$ , thus  $v_e(\mathbf{x}) = q(\mathbf{x})K_G J/\theta$ . 337

Transport Model Transport is modeled by a continuous time random walk. Thus, 338 particles move at constant space increment at transition times that are obtained from 339 the particle velocity. The plume mass distribution at a given time is equivalent to the 340 particle distribution. The particle velocity is modeled as a stationary Markov process, 341 whose steady state distribution  $p_v(v)$  is given by the flux-weighted Eulerian flow veloc-342 ity. The flux-weighting is due to the fact that in this framework particle velocities sam-343 ple the flow velocity equidistantly along path lines. This is in contrast to isochrone sam-344 pling in classical Lagrangian frameworks, for which the steady state distribution of par-345 ticle velocities is equal to the Eulerian velocity distribution (Dentz et al., 2016; Comolli 346 et al., 2019). 347

The Markov process of particle velocities is modeled through an Ornstein-Uhlenbeck process for the normal scores. The normal scores are obtained by mapping the velocities first to a uniform and then to a unit Gaussian random variable. The model requires the Eulerian velocity distribution and advective tortuosity as inputs. The latter is given by the ratio of the mean Eulerian velocity and the mean Eulerian velocity component along the mean hydraulic gradient.

Application to MADE For application to MADE, the model is parameterized based 354 on the description of experimental conditions and aquifer properties as head gradient of 355 J = 0.0036 and porosity of 0.31 according to Boggs et al. (1992); Adams & Gelhar (1992) 356 The retardation coefficient is set equal to one. The distribution of Eulerian velocity mag-357 nitude as the propagator of the upscaled transport model is derived using the geosta-358 tistical parameters of log-normal hydraulic conductivity reported by Bohling et al. (2016) 359 (Fig. 2). The average velocity component in mean flow direction is given by  $\overline{v}_1 = \overline{q}_1 K_G J/\theta =$ 360  $1.942 \times 10^{-7}$  m/s = 0.0167 m/d. The average Eulerian velocity magnitude is  $\overline{v}_e = \overline{q} K_G J/\theta =$ 361 2.234 m/s = 0.0193 m/d.362

A point source particle distribution is assumed. Following Boggs et al. (1992) and Fiori et al. (2013), the initial mass distribution is approximately flux-weighted. Thus, in this modeling framework, the initial distribution of particle velocities is set equal to the flux-weighted Eulerian velocity distribution (Dentz et al., 2020).

#### 3.2.4 Binary Facies

367

*Background* An alternative approach to adopting continuous univariate Y dis-368 tribution consists in representing the media as an assemblage of hydrofacies. Among the 369 geostatistical methodologies adopted for this purpose, the one based on the combined 370 use of transition probability and Markov chain has been mostly employed since its in-371 troduction by Carle & Fogg (1996). It enjoys flexibility in handling justapositonal ten-372 dencies among hydrofacies and the availability of the software T-PROGS (Carle, 1999) 373 In the context of the studies of MADE, Bianchi & Zheng (2016) presented an applica-374 tion to the MADE-2 experiment by adopting representation by 5 hydrofacies. Here we 375 apply the methodology to the MADE-1 experiment in a more parsimonious way. We main-376 tain the highly conductive hydrofacies and combine the remaining four, of low conduc-377 tivity, into a single hydrofacies. 378

K-Structure The porous media is modeled by using two or more facies, each one 379 of constant hydraulic conductivity. The spatial distribution of the facies is generated ran-380 domly based on transition probabilities, e.g. by using T-PROGS. Hydrofacies indenti-381 fication is usually based on granulometry analysis. Here we focus on two hydrofacies of 382 probabilities of occurrence  $p_1$  and  $p_2$ , respectively. Hydraulic conductivities  $K_{1,2}$  are the 383 weighted arithmetic means of the hydraulic conductivity of all samples belonging to the 384 same granulometry class, i.e. hydrofacies. The hydraulic conductivity of the samples are 385 calculated from the characteristic diameters  $d_{10}$  and  $d_{25}$  of the sediments by using a mod-386 ified version of the Kózeni-Carman expression proposed by Riva et al. (2010). Transi-387 tion probabilities between hydrofacies are obtained by fitting a Markov model to the experimental transition probabilities (Carle & Fogg, 1996). They are expressed with the 389 aid of the characteristic thickness and length for each facies, denoted as  $L_{z,1}, L_{z,2}, L_{x,1}$ 390 and  $L_{x,2}$  respectively. Commonly, samples density is relatively low in the horizontal di-391 rections, leading to uncertain experimental transition probabilities in these directions, 392 even after assuming isotropy in the horizontal plane. 393

Flow and Transport Flow and transport are solved repetitively in a Monte Carlo frame by making use of numerical solvers, Modflow 2005 in combination with particle tracking (Pollock, 2012). Mean velocity U and mean relative mass distribution  $\langle \bar{m} \rangle$  are obtained by ensemble averaging.

Application to MADE The binary hydrofacies model was based on the granulometry of 214 samples taken from the 38 boreholes at MADE (Boggs et al., 1990), which is a data set completely independent from those of DPIL used by the FOA and MIM-SCA models. The parameters values pertinent to MADE were identified as:  $p_1 = 0.145$ ,  $p_2 = 0.855$ ,  $K_1 = 190 \text{ m/d}$ ,  $K_2 = 1.49 \text{ m/day}$ ,  $L_{z,1} = 0.702 \text{ m}$ ,  $L_{z,2} = 4.15 \text{ m}$ . The low borehole density prevented reliable estimate of the transition probability in the horizon-

tal directions. Thus, isotropy is assumed in horizontal direction and the characteristic 404 lengths were based on the integral scales identified by Rehfeldt et al. (1992) and assum-405 ing that the relationship  $L_{1,z}/L_{1,h} = L_{2,z}/L_{1,h} = e$  apply, with e being the anisotropy 406 ratio. Thus, horizontal length scales are obtained by dividing the vertical ones, by the 407 estimate of e = 0.0437 (Rehfeldt et al., 1992), arriving at  $L_{h,1} = 16.0 m$ , and  $L_{h,2} =$ 408 94.8 m. In vertical direction the Markov model of transition probability fitted to the ex-409 perimental one was used. As already mentioned above, the horizontal length scales val-410 ues shall be regarded as estimates, affected by uncertainty. A check of the results ob-411 tained with different, smaller, values (not shown) led to similar plume mass distributions 412 except for the long distance tail of minute mass. An ensemble of over 200 independent 413 Monte Carlo realizations of the hydrofacies distribution were thus generated. 414

Flow was solved numerically with Modflow 2005 in a computational domain of  $300 \times$ 415  $100 \times 10 \,\mathrm{m^3}$  with grid spacing of  $1 \,m$  in horizontal and  $0.5 \,m$  in vertical directions for 416 each hydrofacies realization. Constant head boundary conditions were applied with a mean 417 head gradient of J = 0.003. The adopted porosity value was  $\theta = 0.31$ . Advective trans-418 port was simulated by tracking 1000 particles, distributed according to the mass distri-419 bution measured at day 9 since the beginning of the tracer test and advected by the ran-420 dom velocity (in the absence of local dispersion) by means of the ModPath 6 package. 421 The resulting ensemble mean velocity is  $U = 0.079 \, m/day$  with a standard deviation 422 of  $0.0046 \,\mathrm{m/d}$ . 423

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### 3.2.5 Binary Inclusions

Background Further simplification of the previous binary facies model is achieved by representing the high conductivity zones as rectangular inclusions submerged in the low conductivity matrix. Furthermore, for the high length to thickness ratio of the inclusions it was found that flow and transport can be modeled approximately as two-dimensional. We describe here briefly the application of the model to *MADE* presented in Zech et al. (2020), in which the stochastic conceptualization of the binary hydraulic conductivity took large scale deterministic information into account.

<sup>438</sup> Flow and Transport Flow and transport is solved for every random K realization numerically making use of Darcy's Law and ADE solvers. Mean velocity U and mean relative mass distribution  $\langle \bar{m} \rangle$  are obtained by ensemble averaging.

Application to MADE Zech et al. (2020) adopted the K-structure for MADE and 441 its characterizing parameter values based on the inspection of the piezometric surface 442 map of Boggs et al. (1992, Fig. 3), on two large scale pumping tests (Table 1) and on 443 a few flow meter logs (Boggs et al., 1992). According to observations, the area near the 444 source is dominated by low conductivity  $K_2$  with inclusions of high  $K_1$ . The relative area 445 of the inclusions is p = 15% determined from flow meter log analysis. Beyond a dis-446 tance of 20 m downstream of the injection location the distribution is inverted: the bulk 447 is dominated by the high conductivity  $K_1$  with p = 15% inclusions of low  $K_2$ . The in-448 clusions are of thickness  $\ell_v = 0.5 \,\mathrm{m}$  and length  $\ell_h \in [5, 10, 20] \,\mathrm{m}$ . The former was de-449 termined from flowmeter observations while the latter are subjected to parametric un-450 certainty. The range was determined from the ratio of thickness and observed anisotropy 451 value. The random component of the model comprises of the locations of the three ver-452 tical inclusions of 0.5 m thickness, while length and horizontal position of inclusions are 453

fixed for every realization. Given the inclusion occurrence of p = 15%, an aquifer thickness of 10 m and a thickness of 0.5 m, a total number of 3 inclusions per block are randomly placed in vertical position with equal probability. Altogether, an ensemble of 600 conductivity structures was created with random inclusion structures, with groups of 200 realizations of same length inclusion length  $\ell_h$  of 5 m, 10 m or 20 m.

Flow and transport was calculated for every random K-structure solving the ground-459 water flow equation and subsequently the ADE with the FEM software OpenGeoSys in 460 a 2D cross section of  $220 \, m \times 10 \, m$ . The boundary conditions were of constant head on 461 the vertical boundaries defined by the selected head gradient J = 0.003. The initial condition for transport was imposed by injecting a solute discharge  $Q_{in} = 1.15 \cdot 10^{-5} \,\mathrm{m}^3/\mathrm{s}$ 463 during a period of  $T_{in} = 48.5$  h in the injection well with a screen length of 0.6 m (Boggs 464 et al., 1992). Porosity is  $\theta = 0.32$ , local dispersivity was  $\alpha_L = 0.01$  m, not impacting 465 the overall mass distribution. The resulting ensemble mean velocity was U = 0.0254466 m/d with a standard deviation of  $0.02 \,\mathrm{m/d}$ . 467

#### 468 3.2.6 Reactors

*Background* The reactor model conceptualizes subsurface solute transport as series of reactors which are linked to aquifer structure. The concept allows to model transport in a setting where tracer is injected into a low permeability zone, capturing strongly constrained downstream movement and skewed plume shapes. The concept was applied e.g. by Molin & Cvetkovic (2010).

K-Structure Hydraulic conductivity is conceptualized as log-normal spatial random function Y with continuous two-point correlation structure. The multi-Gaussian
geostatistical parameters are those reported by Bohling et al. (2016) based on DPIL for
the model application to MADE of same values as those adopted in the First Order and
MIMSCA models (sections 3.2.1 and 3.2.2).

<sup>479</sup> Flow Mean uniform flow velocity is calculated analytically based on Darcy's Law: <sup>480</sup>  $U = K_{\text{eff}} \cdot J/\theta$ .  $K_{\text{eff}}$  for *MADE* is derived analogously to the *MIMSCA* model from <sup>481</sup> geostatistical parameters (Zarlenga et al., 2018), resulting in U = 0.026 m/d.

Transport Transport is modelled analytically as series of flow reactors, each described by an exponential function  $\exp(-t/\Delta\tau)$  of the mean turnover time  $\Delta\tau$ . While this model cannot be related directly to the permeability distribution, it is of interest to examine the outcome of a conceptually different model. The number of reactors is a function of the ratio  $x/\Delta x$  of the transport domain size x and the velocity fluctuations length scale  $\Delta x$ . Each reactor has a mean turnover time of  $\Delta\tau = \Delta x/U$  and thus a residence time pdf of  $\exp(-tU/\Delta x)$ .

The residence time pdf in the  $x/\Delta x$  series of reactors is then a Gamma function:

$$f(t,x) = \frac{e^{-\frac{tU}{\Delta x}} t^{\frac{x}{\Delta x} - 1} \left(\frac{U}{\Delta x}\right)^{x/\Delta x}}{\Gamma\left(\frac{x}{\Delta x}\right)}$$
(3)

The cumulative density function of residence time is obtained by integration as

$$F(t,x) = \int_0^t f(\theta,x) \, d\theta = 1 - \frac{\gamma \left( x/\Delta x, tU/\Delta x \right)}{\Gamma \left( x/\Delta x \right)} \tag{4}$$

489 where  $\gamma$  is the lower incomplete Gamma-function.

The spatial tracer distribution, as the tracer position pdf p(x;t) [1/L] at time t follows as:

$$p(x,t) = -\frac{\partial F(t,x)}{\partial x}$$
(5)  
=  $\frac{1}{\Delta x \Gamma\left(\frac{x}{\Delta x}\right)} \left[ G_{2,3}^{3,0}\left(\frac{tU}{\Delta x} | \begin{array}{c} 1,1\\ 0,0,\frac{x}{\Delta x} \end{array}\right) + \Gamma\left(\frac{x}{\Delta x},\frac{tU}{\Delta x}\right) \left( \log\left(\frac{tU}{\Delta x}\right) - \psi^{(0)}\left(\frac{x}{\Delta x}\right) \right) \right]$ (6)

where G is the Meijer g-function, and  $\psi^0$  is the Polygamma function. With unit tracer mass, we have m(x,t) = p(x,t).

<sup>492</sup> Note that in the limit  $\Delta x \to 0$ , the number of reactors tends to infinity  $x/\Delta x \to \infty$ , and we recover plug flow as  $f(t, x) \to \delta(t - x/U)$ .

Like in the *MIMSCA* model, the length parameter  $\Delta x$ , the integral scale of the velocity fluctuations, is assumed to be in the range of two to three log-conductivity horizontal correlation length *I*.

<sup>497</sup> Application to MADE The value of  $U = 0.026 \ m/d$  is identical with the one used <sup>498</sup> in the FOA or MIMSCA models and similarly I = 9.1m is based on Bohling et al. (2016).

#### 499 4 Prediction and Inter-comparison using MADE Data

#### 4.1 Results Presentation

500

The visual presentation of results concerning the mass spatial distribution signif-501 icantly influences the perception of model performance. Critical aspects are the display 502 scale, data upscaling (aggregation) and normalization. We present longitudinal mass dis-503 tributions  $\langle \bar{m} \rangle$  at linear and logarithmic scales as well as in a cumulative form  $\langle M \rangle$ , to 504 achieve a comprehensive display. The various display modes of the spatial mass distri-505 butions allow to interpret a few plume's specific features: (i) mass peak location and bulk 506 behavior; (ii) the tails (forefront and trailing zones) and (iii) mass recovery. They are 507 relevant to specific goals such as risk assessment and remediation. 508

In line with the *MADE* experimental results, we remind that  $\langle \bar{m}(x,t) \rangle$  is aggregated over intervals of 10 m. We illustrate the effect by displaying the model outcomes at T = 126 days in both forms, the fine scale  $\langle m(x,t) \rangle$  and the upscaled  $\langle \bar{m}(x,t) \rangle$  (Fig. 3) as well as  $\langle M \rangle$ . All longitudinal mass distributions are normalized with respect to the injected mass such that the area beneath m or  $\bar{m}$  is unity. For the *MADE* data this is only true for t = T = 126 days where mass recovery is around 99%.

Spatial moments are commonly used to quantify the comparison between different models and measurements. However, this is not appropriate because of the skewed shape of  $\langle \bar{m} \rangle$  and the large impact of the uncertainty of the tail. Instead, we compare in Fig. 4 recovery locations at T = 126 days for 5%, 50%, and 95% as predicted by models relative to *MADE* values of x = 0.7, 8.6, 42.8 meters. The definition of the recovery location is for instance for  $x_{95\%}$ , the position for which 95% of the total mass is upstream of  $x_{95\%}$ , i.e.  $\langle M \rangle = 0.95$ .

Spatial moments are commonly used to quantify the comparison between differ-522 ent models and measurements. However, this is not appropriate because of the skewed 523 shape of  $\langle \bar{m} \rangle$  and in particular the large impact of the minute and uncertain mass frac-524 tion in the forefront tail, upon the second spatial moment (Fiori et al., 2017). Instead, 525 we compare in Fig. 4 recovery locations at T = 126 days for 5%, 50%, and 95% as pre-526 dicted by models relative to MADE values of x = 0.7, 8.6, 42.8 meters. The definition 527 of the recovery location is for instance for  $x_{95\%}$ , the position for which 95% of the to-528 tal mass is upstream of  $x_{95\%}$ , i.e.  $\langle M \rangle = 0.95$ . 529



Figure 3. Longitudinal mass distribution for models at T = 126 d in fine model resolution (left column) and upscaled (aggregated) form ( $\Delta x = 10$  m, right column) against *MADE-1* experiment data at linear scale (1st row), log-scale (2nd row) and in cumulative mass (3rd row). Vertical lines in 3rd column indicate locations of 5% (dotted), 50% (dashed), and 95% (dashed dotted) recovered mass. Note that for the *Binary Facies* model the fine scale refers to a grid resolution of 2 meters.



Figure 4. Model-experiment-comparison of recovery locations x for 5%, 50%, and 95% mass recovery (columns) at 126 days after injection where mass recovery was 99% in the experiment. *MADE* values at the x-axis [in m] and model values as absolute [in m] and relative [in %] difference in coloured bars and in numbers.

#### 4.2 Comparison between Models Prediction and MADE Experiment

Fig. 3 displays the longitudinal mass distribution for all models and the *MADE* experimental data at T = 126 days after injection:  $\langle m \rangle$  at model's fine resolution and  $\langle \bar{m} \rangle$ aggregated over 10 m intervals, including *MADE*. Direct comparison is most revealing at that time since the experimental recovery rate is 99%.

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The various models display some differences in their mass distribution  $\langle m \rangle$  at fine scale. Particularly, the peak value is higher for the *Flow Reactors* by a factor of 2 than the other models. The *Binary Facies* model displays plume tailing downstream of the other models, with  $x_{95\%} \cong 60m$  while for all the others models  $30m \lesssim x_{95\%} \lesssim 40m$ (see Fig. 3).

The comparison between  $\langle m \rangle$  and the upscaled  $\langle \bar{m} \rangle$  in Fig. 3 reveals a few inter-540 esting features: upscaling reduces the peak values of  $\langle m \rangle$  by a factor of around 2, the spread-541 ing zone is expanded, and the differences between models prediction are greatly reduced. 542 In particular, the predicted  $\langle \bar{m} \rangle$  agrees quite well with *MADE*, much better than  $\langle m \rangle$ , 543 as far as visual inspection reveals. This is expectable though in the past models predic-544 tion at fine scale were compared with MADE (e.g. Harvey & Gorelick (2000); Dogan et 545 al. (2014)). The upstream spread of the aggregated  $\langle \bar{m} \rangle$  for MADE is partly an artifact 546 of upscaling: it smears the upstream forced injected mass of the initial solute body over 547 10 meters. It could be erroneously interpreted as upstream dispersion. 548

The quantitative results in Fig. 4 on recovery locations relative to MADE strengthen the conclusions from the visual inspection. All models are in good agreement with MADEfor the location  $x_{50\%}$ . We attribute that to the closeness of the mean velocity U of various models (see previous Section). The *Binary Facies* differs by 15% due to the differ-



Figure 5. Longitudinal mass distribution for models and *MADE-1* experiment at 49, 202, 279, 370, and 503 days after injection at linear scale (1st column), log-scale (2nd column) and in cumulative form (3rd column) with recovery locations (Fig. 3). Observe the different recovery rates of 2.06, 0.68, 0.62, 0.54, and 0.43, respectively.



**Figure 6.** Longitudinal mass prediction of models at 1000 days at linear scale (1st column), log-scale (2nd column) and in cumulative form (3rd column) with recovery locations.

ence in the value of U. The agreement is still good for  $x_{5\%}$  with some deviations for the Binary Inclusion model. Last, the models prediction relative to MADE of  $x_{95\%}$ , reflecting the "fast" moving solute, is more variable, but still within acceptable differences in practice.

Fig. 5 summarizes the model and MADE results of  $\langle \bar{m} \rangle$  for all times for which MADEexperimental data is available. The comparison at these times with MADE is more difficult than for T = 126 days because of the variable mass recovery, which is visible in Fig. 5 in the cumulative mass panel  $\langle M \rangle$ .

Our interpretation of the apparent "mass loss" for later times is the less dense sam-561 pling in the downstream zone as illustrated by Fig. 1 on one hand and the not-sampled 562 solute quickly moving in high conductivity channels on the other hand, as already men-563 tioned before. This is clearly visible in Fig. 5 displaying the larger predicted mass than 564 the measured one downstream of the peak. Despite that, it is remarkable that for T =565 202, 270, 370 days all models agree quite with both measured peak value and its distance 566 from the injection zone. At the largest time T = 503 days the low mass recovery (43) 567 %) makes the comparison between data and models prediction quite problematic. Still, 568 the peak location and even the value are within acceptable differences in practice. As 569 for inter-comparison of models prediction, inspection of the robust cumulative mass  $\langle M \rangle$ 570 at different times shows remarkable closeness except for *Binary Facies*. The latter over-571 estimates the location  $x_{\%}$  pertaining to fixed values of  $\langle M \rangle > 0.4$ . This enhanced tail-572 ing is attributed in part to the larger integral horizontal scale identified from granulom-573 etry analysis (Fig. 2) as compared to that resulting from DPIL, which was used by the 574 most of the models. In addition, it may be related to channels of high conductivity present 575 in some realizations of the binary K field. It is remarkable that for both  $\langle \bar{m} \rangle$  and  $\langle M \rangle$ , 576 the predictions by FOA, MIMSCA and TDRW are very close for all x (see discussion 577 in Sect 5). 578

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#### 4.3 Prediction Beyond MADE Experiment

The important role of models is to provide prediction of future solute plumes de-580 velopment. In order to compare the outcome of the 6 different models considered in the 581 present study, we have used them for the same MADE conditions, but at larger time than 582 T=503 days. Thus, the long term (at T=1000 days) predicted plume mass spatial dis-583 tribution is displayed in Fig. 6. All models agree in the peak travel distance as a con-584 sequence of similar flow velocities U. However, the differences between the peak value 585 are more pronounced but still within a factor of two and even less for all models, except 586 the *Reactors*, which predicts a higher peak and reduced tail. As for prediction of fore-587 front tailing associated with fast moving solute, both Binary Facies and Binary Inclu-588 sions display longer tails with  $80 \lesssim x_{95\%} \lesssim 140$  meters. Inspection of the cumulative 589

distribution reveals again that prediction of models except *Binary Facies* are within a relatively narrow distribution.

#### 592 5 Discussion

#### 5.1 Modelling

The main aim of the study is to compare the prediction of six different models, with MADE serving as a platform. However, based on conceptual similarity as well as prediction, the six models can be divided in three groups as follows: (i) FOA (First Order Approximation), MIMSCA (Multi Indicator and Self Consistent Approximation), and TDRW (Time Domain Random Walk); (ii) BF (Binary Facies) and BI (Binary Inclusions); and (iii) Reactors. Herein a discussion of the main results for each group.

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#### (i) FOA, MIMSCA, and TDRW

For the three models K is modeled as a multi-Gaussian stationary random field, completely characterized by  $K_{\rm G}$ ,  $\sigma_Y^2$ , I,  $I_v$ , the flow mean velocity being given by  $U = K_{\rm eff} J/\theta$ . They lead to approximate distributions of  $\langle m \rangle$ ,  $\langle \bar{m} \rangle$  and  $\langle M \rangle$  as functions of x, t.

The three models differ in conceptualization and computational complexity. FOAis analytical, leading to an Inverse Gaussian  $\langle m \rangle$  which satisfies an ADE with macrodispersivity derived analytically by FOA. MIMSCA is semi-analytical, based on summation of travel time through inclusions of random K. TDRW is semi-numerical, with the velocity field derived by *Monte Carlo* simulations while transport is based on an approximation of the *Lagrangian* velocity field.

<sup>610</sup> One of our main result is that the solutions for  $\langle \bar{m} \rangle$  and  $\langle M \rangle$  by the three models are very close and in good agreement with the bulk of *MADE* experimental data. Thus they are are very robust and prediction depends primarily on U Fiori et al. (2017), as well as  $\sigma_Y^2$  and I and much less on models methodology.

#### (ii) BF and BI

The hydraulic conductivity heterogeneity is modeled by two values  $K_1, K_2$  of volume fractions  $p_1$  and  $p_2 = 1 - p_1$ . Two length scales  $L_x$ ,  $L_z$  characterize the K-facies, whose geometry has random elements. Flow and transport are solved numerically and repeatedly, by *Monte Carlo* simulations; besides the mean values  $U, \langle \bar{m} \rangle$  and  $\langle M \rangle$ , the statistical moments of these parameters can be also obtained.

The two models differ in few respects: BF is three-dimensional; the random facies 620 geometry is generated by transitional probability using the *TPROGS* code, which re-621 quires the knowledge of two more length scales for the second K-facies being different 622 from the first. All the structural parameters are obtained from granulometry measure-623 ments. In contrast, the simpler BI model is two-dimensional and consists of identical 624 rectangular inclusions of conductivity  $K_1$  submerged in the  $K_2$  matrix for specified de-625 terministic regions. The inclusions lengths assume 3 different values of same probabil-626 ity, their elevation being random. The structural parameters  $K_1, K_2$  and  $p_1$ , as well as 627 the length scales, are derived from pumping tests and few flowmeter measurements. 628

The main results for these two models are as follows: the agreement with *MADE1* data is reasonable; while the characterization effort is less demanding than for the previous models, the numerical solutions and the *Monte Carlo* simulations of both flow and transport are quite involved, especially for *BF*.

The main conclusions are: the favorable comparison with *MADE1* is an additional proof of models robustness; the simplified structures and characterization are adapted to the particular features of the *MADE* site for which two dominant zones could be delineated.

#### 637 *(iii)* Reactors

This model has a different conceptualization from the previous ones and it was mo-638 tivated by its general use in engineering and convenience for reactive transport modelling. 639 The model parameters are U, which is derived from the solution of flow via  $K_{\text{eff}}$ , and  $\Delta x$ , 640 the velocity longitudinal correlation length which is related to I. The degree of hetero-641 geneity quantified by  $\sigma_Y^2$  is not included, instead the series of reactors aggregate the in-642 herent dispersion of a flow reactor; thus the model is not a general candidate for mod-643 eling advective transport in a heterogeneous aquifer. Still, it was found of interest to com-644 pare the analytical solution for  $\langle \bar{m} \rangle$  with MADE plume. The surprising result is that  $\langle m \rangle$ 645 and even more so  $\langle \bar{m} \rangle$  agree reasonably well with *MADE*, though the predicted peak of 646  $\langle \bar{m} \rangle$  is larger by a factor of 1.5 than prediction by other models for T = 1000 days. 647

The finding strengthens the conclusion about the robustness of  $\langle \bar{m} \rangle$  and  $\langle M \rangle$  in predicting the measured *MADE* plume and its future development, with the predominant role of two parameters, the mean velocity U and correlation scale I; whether the reactors model can be used for prediction requires its further development and comparison with more cases.

5.2 Data Selection

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The K distribution at the MADE site as inferred from different characterisation 654 methods is summarized in Fig. 2: (i) DPIL measurements are a novel, affordable tech-655 nique for shallow aquifers (Dietrich et al., 2008) by which the most comprehensive data 656 set was obtained suitable for a geostatistical interpretation, including two-point statis-657 tics; (ii) Granulometry (or grain size analysis) is a standard method in hydrogeology that 658 yields highly uncertain conductivity estimates (Vienken & Dietrich, 2011); (iii) flowme-659 ter measurements and pumping tests are standard methods for hydrogeological charac-660 terisation, the main limitation being accuracy of low K values (Fig. 2). 661

The (semi-)analytical models FOA, MIMSCA, TDRW, and Reactors used DPIL 662 observations to estimate the flow velocity and the plume spreading. The model *Binary* 663 Facies (BF) used granulometry data and the Binary Inclusions (BI) model used pump-664 ing test estimates, information from head maps and a few flowmeter data. In essence, 665 BF simplifies the facies approach applied by Bianchi and Zheng (2016) MADE2 by con-666 sidering 2 facies instead of 5, where the binary K values are inferred from granulome-667 try. The BI approach simplifies further the 3D random structure by considering regu-668 lar inclusions in two dimensions. Thus the semi-analytical models are relatively simple 669 for computations but use a more extensive DPIL data set whereas BI and BF modelling 670 approaches are heavier on computation but use more readily accessible data sets. 671

The derivation of the mean velocity U by the different models imply the use of es-672 timates of the measured mean head gradient J and the effective porosity  $\theta$ , which are 673 approximate. The models differ primarily in the use of the K data. Still, the resulting 674 estimates of the mean velocity U are relatively close as revealed by the values appear-675 ing in Sect. 3.2 (0.026m/d for FOA, MIMSCA, 0.019m/d for TDRW, 0.079m/d for BF 676 and 0.025 m/d for BI). Even the deviation for BF is within an acceptable range, vari-677 ous approximations notwithstanding. Thus, the estimates of U are quite robust, which explains the relative closeness of the predicted and measured locations of the peak of  $\langle \bar{m} \rangle$ 679 in Figures 3, 5, 6. 680

Similarly, the prediction of spreading as quantified by the  $\langle \bar{m} \rangle$  and  $\langle M \rangle$  distributions is quite robust, as already discussed above. An interesting finding which may somewhat explain the relative closeness of the distributions for different data characterization methods is the magnitude of the *FOA* asymptotic longitudinal macrodispersivity  $\alpha_L = \sigma_Y^2 I$  based on the different values of  $\sigma_Y^2$  and *I* in Fig. 2. The resulting values of  $\alpha_L$  are 53.7m, 54.1m and 49.6m for *DPIL*, *Flowmeter* and *Grain Size*, respectively.

Although above observations strictly apply to the *MADE* site only, they are nevertheless encouraging and motivate similar comparative analysis e.g., for less heterogeneous aquifers for which experimental data are available.

#### 690 6 Summary and Conclusions

With a variety of hydraulic data available, MADE provides a unique opportunity 691 for a comparative analysis of predictive modelling, from a wide range of (semi)-analytical 692 models (FOA, MIMSCA and TDRW) that utilise extensive DPIL data for inferring geo-693 statistical parameters, to numerical models (BF and BI) with relatively simple (binary) 694 structures that utilise much less extensive data sets (granulometry and pumpoing tests+ 695 flowmeter). The present paper takes advantage of these possibilities offered by MADE, 696 and focuses on comparing predictions of the plume spreading by six different models. The 697 models differ in theoretical formulations, in the conceptualization of aquifer structure, 698 in the field data input, and in the computational effort. Common features of the mod-699 els are: flow is steady, uniform in the mean and driven by a head gradient J; solute spread-700 ing is caused by aquifer conductivity heterogeneity; models rely on structural data and 701 flow data with no calibration on transport observations i.e. the models are predictive; 702 plume mass behavior is assumed ergodic, i.e. the mean relative mass distribution  $\langle \bar{m} \rangle$ 703 derived by the model is compared with the measured mass  $\bar{m}$  at a few times T. More-704 over, the apparent loss of mass of measured mass at MADE1 for T > 126 days (attributed 705 to limited sampling) is not taken into account by the models. Model comparison at T =706 1000 days, i.e., beyond the period of measurements, is also included. 707

The main and encouraging result for practitioners is that all model prediction agree 708 reasonably well with MADE1 mass distributions and the same for the comparison at T=1000709 days. Thus, the measures of the solute plumes are robust and models are reliable as long 710 as they are underlined by a few basic parameters: mean velocity U, a parameter reflect-711 ing log-conductivity variability and one taking horizontal correlations in conductivity into 712 account. However, the reasonable agreement in model results is also related to the par-713 ticular quantity under examination: the longitudinal mass distribution which is aggre-714 gated over spatial intervals is quite robust itself. If other measures are employed, such 715 as local concentrations, results might differ. 716

To render the above conclusions of general validity, the study shall be extended by application to other cases than *MADE*. As a first step synthetic examples can be considered like formations of log-normal conductivity, with different connectivities as analyzed for instance by Fiori et al. (2017), or typical facies structures (e.g. Carle & Fogg (1996)). Similarly, one needs to test the hypothesis that a consistent comparison would have been obtained even for a less heterogeneous aquifer.

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els and MADE observations can be found in the references. Scripts used in the paper
 are available upon request from the corresponding author.

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