First studies of mesosphere and lower thermosphere dynamics using a multistatic specular meteor radar network over southern Patagonia

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Abstract

This paper presents for the first time results on winds, tides, gradients of horizontal winds, and momentum fluxes at mesosphere and lower thermosphere (MLT) altitudes over southern Patagonia, one of the most dynamically active regions in the world. For this purpose, measurements provided by SIMONe Argentina are investigated. SIMONe Argentina is a novel multistatic specular meteor radar system that implements a SIMONe (Spread Spectrum Interferometric Multistatic meteor radar Observing Network) approach, and that has been operating since the end of September 2019. Average counts of more than 30000 meteor detections per day result in tidal estimates with statistical uncertainties of less than 1 m/s. Thanks to the multistatic configuration, horizontal and vertical gradients of the horizontal winds are obtained, as well as vertical winds free from horizontal divergence contamination. The vertical gradients of both zonal and meridional winds exhibit strong tidal signatures. Mean momentum fluxes are estimated after removing the effects of mean winds using a four-hour, eight-kilometer window in time and altitude, respectively. Reasonable statistical uncertainties of the momentum fluxes are obtained after applying a 28-day averaging. Therefore, the momentum flux estimates presented in this paper represent monthly mean values of waves with periods of four hours or less, vertical wavelengths shorter than eight kilometers, and horizontal scales less than 400 km.

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Key Points:

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13	•	First observations of MLT dynamics over one of the most dynamically active re-
14		gions in the world
15	•	Estimates of mean horizontal winds and their gradients are possible, thanks to the
16		multistatic configuration.
17	•	Mean momentum fluxes are estimated with vertical velocity estimates free of hor-
18		izontal divergence contamination.

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19 Abstract

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³⁸ 1 Introduction

The mesosphere and lower thermosphere (MLT) is the atmospheric region that cou-39 ples the lower and upper parts of the terrestrial atmosphere. For this reason, knowledge 40 of its dynamics is of great importance in order to understand the behavior of the atmo-41 sphere as a whole. The coupling is accomplished mainly via propagation of three dom-42 inant types of waves: planetary waves (PWs), tides, and gravity waves (GWs). PWs are 43 waves with scales of thousands of kilometers and periods of up to ~ 30 days. They are 44 mainly generated in the troposphere by land-sea discontinuities, or triggered in-situ by, 45 e.g., baroclinic instabilities and filtered gravity waves (e.g., Rossby, 1939; McCormack 46 et al., 2014; H.-L. Liu & Roble, 2002). Tides are also waves with horizontal scales of thou-47 sands of kilometers, but periods that are sub-harmonics of the solar and lunar days. Ther-48 mal tides are mainly a consequence of solar radiation absorption by water vapour in the 49 troposphere and ozone in the stratosphere, while the lunar tide results from the grav-50 itational pull of the Moon (e.g., Lindzen & Chapman, 1969; Forbes, 1984). GWs are small 51 to medium scale waves with periods ranging from about 5 minutes to many hours. They 52 can be triggered by a myriad of different sources, e.g., the orography, thunderstorms, shear 53 instabilities, convection, etc. (e.g., Hines, 1988; Piani et al., 2000; Fritts & Alexander, 54 2003).55

During the last decades, specular meteor radars (SMRs) have been extensively used 56 to study winds and atmospheric waves in the MLT (e.g., Hocking, 2005; Clemesha et al., 57 2009; Hoffmann et al., 2010; A. Z. Liu et al., 2013; Laskar et al., 2016; Jia et al., 2018, 58 and references therein). They have also been used to study GWs, which are known to 59 play an important role in determining the wind and thermal structure of the MLT (e.g., 60 Fritts, 1984). Particularly, some studies have focused on extracting information about 61 GW-driven momentum fluxes from SMR measurements (e.g., Fritts et al., 2010; Placke 62 et al., 2011a; Andrioli et al., 2015). However, understanding the results on momentum 63 flux estimates based on SMR winds is not trivial, mainly because of the uncertainties 64 associated with the estimation procedure (e.g., Fritts et al., 2012b). In fact, Vincent et 65 al. (2010) showed that the accuracy in the momentum flux estimation is highly depen-66 dent on the number of meteor detections. Consequently, the usage of multistatic meteor 67 radar systems represents one way to reduce the uncertainties of the momentum flux es-68 timates (e.g., Spargo et al., 2019). Furthermore, by detecting more meteors and being 69

able to observe them from different viewing points, multistatic SMR systems also allow
 for more reliable estimations of horizontal wind gradients.

The MLT over the southern part of Argentina and Chile is considered to be one 72 of the most dynamically active regions in the globe. Satellite-based studies have revealed 73 that GW-driven momentum fluxes increase considerably at both stratosphere and MLT 74 altitudes over Patagonia (e.g., Trinh et al., 2018; Vadas et al., 2019a). Numerical model 75 simulations have reported generation of secondary GWs with horizontal scales of up to 76 2000 km at mesospheric altitudes over the southern Andes (Vadas & Becker, 2019). Nev-77 78 ertheless, wave coupling processes in the MLT region over the Patagonian sector are still not well understood, partly because the installation of ground-based instruments has not 79 been possible, either due to logistics challenges or instrument requirements. In this work, 80 we present preliminary results of a multistatic SMR network that allows, for the first time, 81 measurements of MLT dynamics in the Patagonian region. Besides the local support, 82 our success has been possible thanks to a novel approach that we call SIMONe (Spread 83 Spectrum Interferometric Multistatic meteor radar Observing Network) (Chau et al., 2019). 84 SIMONe makes use of modern radar practices like spread-spectrum, MIMO (Multiple-85 Input, Multiple-Ouput), and compressed sensing applied to atmospheric radars (Vierinen 86 et al., 2016; Urco et al., 2018, 2019). This allows for much easier installation, operation 87 and expansion of the network than previous equivalent systems. 88

The paper is organized as follows. Section 2 introduces the SIMONe Argentina system. Section 3 provides a detailed description of the different analyses performed to the data. The main results and findings are presented and discussed in Section 4. Finally, the concluding remarks are provided in section 5.

⁹³ 2 SIMONe Argentina

SIMONe Argentina is a state-of-the-art network of multistatic specular meteor radars 94 that was installed in September of 2019 in the southern province of Santa Cruz, Argentina. 95 It is comprised of one single transmitting site with five linearly polarized Yagi antennas 96 in a pentagon configuration, and five receiving sites with one dual-polarization Yagi an-97 tenna each. The receivers are placed between 30 and 270 km of distance from the trans-98 mitting site, which is located at 49.6° S, 71.4° W (see Figure 1 for details on the geo-99 graphical distribution of the sites). This type of network configuration is known as MISO 100 (Multiple-Input, Single-Output), since only one antenna is used on reception (e.g., Chau 101 et al., 2019). 102

SIMONe Argentina is the result of an effort led by the Leibniz Institute of Atmo-103 spheric Physics (Germany) in collaboration with the Universidad Nacional de la Patag-104 onia Austral (Argentina), and the Arctic University of Norway. A similar system has been 105 installed in Peru (SIMONe Peru). SIMONe systems use coded spread spectrum on trans-106 mission (Vierinen et al., 2016). A phase coded signal based on pseudo-random sequences 107 is generated and transmitted on each antenna independently. Transmission is done at 108 a frequency of 32.55 MHz and with an average power of 400 W per antenna. All five trans-109 mitted codes are simultaneously decoded at each receiving site by means of compressed 110 sensing (e.g., Urco et al., 2019). Hardware and software details of both systems, i.e., SI-111 MONe Peru and SIMONe Argentina, can be found in Chau et al. (2020). 112

SIMONe Argentina started operations by the end of September 2019 and has been running since then with almost no interruptions. Figure 2 shows a summary of the detection statistics for the first seven months of operations. The upper panel indicates the normalized percentage of meteor counts for each individual link. The bottom panel is used to present the average daily total counts for each month. Problems with the local power supply at the transmitting site resulted in fewer meteor detections during April of 2020. Besides, the receiving site at Gobernador Gregores was out of operation dur-



Figure 1. Map of SIMONe Argentina. The transmitter is indicated with a filled red pentagon, and the receivers are indicated with blue crosses.

ing December of 2019 due to damage in the antenna cables. Nevertheless, for most of
the time, the number of counts was much higher than in standard specular meteor radars.
The links concentrating most of the meteor detections are Tres Lagos-El Calafate and
Tres Lagos-La Estela. Starting in March 2020, the link Tres Lagos-Rio Gallegos exhibits
a considerable increase in the counts, as a result of having rotated the transmitting antennas by 90 degrees. By month, January presents the largest counts, with an average
of more than 50000 meteor detections per day.

¹²⁷ **3 Data analysis**

Specular meteor radars (SMRs) are used to measure the Doppler shift of meteor
trails due to their drifting with the mesospheric neutral winds (e.g., Jones et al., 1998).
In order to extract the wind information from the measurements, one may implement
an all-sky fit of the Doppler velocities measured during a certain period of time and within
a given altitude interval (e.g., Hocking et al., 2001; Holdsworth et al., 2004). In other
words, one must solve the following equation:

$$\mathbf{u} \cdot \mathbf{k} = 2\pi f + \zeta,\tag{1}$$

where $\mathbf{u} = (u, v, w)$ is the neutral wind vector, with u, v and w being its zonal (east-134 west), meridional (north-south) and vertical (up-down) components, respectively. $\mathbf{k} =$ 135 (k_u, k_v, k_w) is the Bragg wave vector (scattered minus incident) in the meteor-centered 136 east-north-up coordinate system (perpendicular to the meteor trail); f is the Doppler 137 shift; and ζ is the Doppler shift uncertainty. For this equation to be valid, one must make 138 the assumption that the winds at each given height interval are uniform during the se-139 lected period of time (homogeneous method). The results using the homogeneous method 140 have been obtained assuming w = 0. 141

Monostatic SMRs commonly allow for hourly horizontal wind estimations every 2-3 km in altitude (e.g., Jacobi et al., 1999; Hoffmann et al., 2010). The wind estimates are representative of mean values over an area of ~ 200 km in radius. These horizon-

SIMONe Argentina detection statistics



Figure 2. Upper panel: normalized percentage of meteor counts, color coded for each individual link. Bottom panel: monthly daily mean counts for all the links combined (see text for details).

tal winds can be further processed in order to obtain information about large scale waves
such as planetary waves and tides (e.g., Murphy et al., 2006; Chau et al., 2015; Conte
et al., 2018).

One of the advantages of using multistatic SMR networks is that the amount of 148 meteor detections is significantly increased (e.g., Stober & Chau, 2015). During most of 149 the time since it started operations, SIMONe Argentina has been capable of detecting 150 more than 30000 meteors per day (see Figure 2). With such amounts of meteor detec-151 tions, one can not only reliably estimate horizontal winds with higher time and verti-152 cal resolutions (~ 15 min and 1 km, respectively), but also investigate second order pa-153 rameters such as the squares of the perturbation components of the radial velocities, i.e., 154 the momentum flux components (e.g., Spargo et al., 2019). Another advantage multi-155 static SMR networks present is that they are capable of sampling the observed volume 156 from different viewing angles. In this scenario, the increased number of meteor detec-157 tions can be further exploited in order to estimate first-order derivatives of the horizon-158 tal winds, and in this way include in the estimation a mean vertical wind free of hori-159 zontal divergence contamination (e.g., Chau et al., 2017). 160

Below, we describe the main procedures used in this work, i.e., tidal estimation, 161 wind estimation using a gradient method, and mean momentum flux estimation. For each 162 of the three procedures, a first wind estimation is carried out in order to remove outliers. 163 That is, Equation 1 is solved in bins of the same size as that used later on in each given 164 procedure (e.g., 4 hour, 4 km in the gradient method); and then the radial velocities the 165 values of which have a corresponding residual of more than 3 standard deviations are 166 removed. This first wind estimation is carried out only in those bins containing a min-167 imum of 10 meteor detections. 168

3.1 Tidal estimation

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Horizontal winds obtained from meteor radar measurements have been used for several decades to investigate tides and planetary waves in the MLT (e.g., Hocking & Thayaparan, 1997; Fritts et al., 2012a). Different mathematical techniques such as least squares
or wavelet analysis can be applied to the wind data in order to extract the tidal infor-

mation (e.g., Stening et al., 1997; Sandford et al., 2006; He et al., 2017). To avoid zero padding or interpolating when encountering data gaps, a least squares approach was se lected for this study.

After removing the outliers, Equation 1 was again fitted to the Doppler shift mea-177 surements. For this purpose, a weighted least squares (WLS) technique was implemented 178 using bins of 1 hour and 2 km (in altitude), shifted by half an hour and 1 km, respec-179 tively. The inverse of the squared Doppler shift uncertainties (ζ in Eq. 1) were used as 180 weights. The WLS was carried out only in those bins containing a minimum of 10 me-181 182 teor detections. Then, under the assumption that the obtained hourly horizontal winds are the result of the superposition of a mean wind and different period oscillations, the 183 following equation was fitted to the zonal (u) and meridional (v) wind components 184

$$[u + \psi_u, v + \psi_v] = [U_0, V_0] + \sum_{i=1}^4 [A_u, A_v]_i \cos\left(2\pi \frac{(t - [\phi_u, \phi_v]_i)}{T_i}\right).$$
(2)

Here, ψ_u and ψ_v are the Doppler shift uncertainties (error) propagated into the estimated 185 winds; U_0 and V_0 are the mean zonal and meridional winds; A_{u_i} (A_{v_i}) and ϕ_{u_i} (ϕ_{v_i}) are 186 the amplitude and phase, respectively, of the zonal (meridional) component of each con-187 sidered wave; $T_1 = 2$ day; T_i for i > 1 is the period of each considered tide ($T_2 = 24$ 188 h, $T_3 = 12$ h and $T_4 = 8$ h, for the diurnal, semidiurnal and terdiurnal solar tides, re-189 spectively); and t is the Universal Time (UT) in hours. The cosine of a sum was used 190 to linearize Equation 2, which then was solved by applying the WLS method using a run-191 ning window of 4 days shifted by 1 day. 192

3.2 Wind field gradient method

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If one relaxes the assumption of homogeneity, the wind field inside the observed area may be estimated using the gradient method. This method consists in approximating the horizontal winds with their first-order Taylor expansion terms (e.g., Burnside et al., 1981; Browning & Wexler, 1968; Chau et al., 2017). This means introducing the following expression into Equation 1

$$\mathbf{u}(x, y, z, t) = \mathbf{u}_0(z, t) + \mathbf{u}_x(z, t)(x - x_0) + \mathbf{u}_y(z, t)(y - y_0) + \mathbf{u}_z(z, t)(z - z_0),$$
(3)

where $\mathbf{u}_0 = (u_0, v_0, w_0)$ represents the mean wind; $\mathbf{x}_0 = (x_0, y_0, z_0)$ is a reference point; and

$$\mathbf{u}_{x} = \left(\frac{du}{dx}, \frac{dv}{dx}, \frac{dw}{dx}\right)$$
$$\mathbf{u}_{y} = \left(\frac{du}{dy}, \frac{dv}{dy}, \frac{dw}{dy}\right)$$
$$\mathbf{u}_{z} = \left(\frac{du}{dz}, \frac{dv}{dz}, \frac{dw}{dz}\right).$$

The mean wind and the first order derivatives depend only on time and height. The co-201 ordinates (x, y, z) are in km, and calculated taking into consideration the latitude, lon-202 gitude and altitude of each meteor detection and the radius of the Earth at the refer-203 ence point. \mathbf{x}_0 is determined using the latitude and longitude of the transmitting site, 204 and the altitude of each height level considered in the WLS fit. The latter was imple-205 mented using bins of 4 hour and 4 km (in altitude), shifted by 15 minutes and 1 km, re-206 spectively. For this study, we have assumed that (dw/dx, dw/dy, dw/dz) = 0, which means 207 solving for nine unknowns. In this method, a condition of having a minimum of forty 208 meteor detections per bin was selected. 209

Mathematically speaking, the zero order approximation (\mathbf{u}_0) as well as the first order derivatives $(\mathbf{u}_x, \mathbf{u}_y \text{ and } \mathbf{u}_z)$ are defined at the reference point, \mathbf{x}_0 . However, the dependence of our nine estimates on this reference point is not explicitly written in Equation 3 for the following reason. The mean winds and the gradients are obtained after fitting measurements made at points randomly distributed within an area of roughly ~ 200 km of radius around \mathbf{x}_0 . Thus, it is assumed that they represent a "mean" zero order approximation, and "mean" first order derivatives, not only at the reference point, but over the whole area aforementioned.

3.3 Momentum flux estimates

The procedure followed in this study to estimate the GW momentum flux is based on the works by Thorsen et al. (1997) and Hocking (2005). It consists in applying a least square method to solve the following equation

$$(\mathbf{u}' \cdot \mathbf{k})^2 = (2\pi (f - \hat{f}))^2. \tag{4}$$

In this expression, **k** and f are, respectively, the same Bragg wave vector and same Doppler shift as in Equation 1, but **u'** represents the perturbed wind vector instead. \hat{f} is a socalled mean Doppler shift,

$$\hat{f} = \mathbf{u} \cdot \mathbf{k} / 2\pi,\tag{5}$$

where **u** is the mean wind that results from solving Equation 1 or 3. To guarantee more reliable values of the six unknowns ($\langle u'u' \rangle$, $\langle v'v' \rangle$, $\langle u'w' \rangle$, $\langle u'v' \rangle$, $\langle u'v' \rangle$ and $\langle w'w' \rangle$), the fit was performed only in those bins containing forty meteor detections or more.

From Equations 4 and 5, it follows that different momentum flux estimates may be obtained depending on the mean wind that is used to determine \hat{f} . In this study, the wind estimates that result from applying the gradient method were used to calculate \hat{f} (see previous section). That is, Equation 3 (with dw/dx = dw/dy = dw/dz = 0) was introduced into Equation 1, and the latter then solved using bins of 4 hour and 8 km, shifted by 15 min in time and 1 km in altitude. Finally, Equation 4 was solved in bins of 4 hour and 4 km (in altitude), shifted by 30 min and 1 km, respectively.

²³⁵ 4 **Results and discussion**

One of the main goals of this work is to provide, for the first time, information on
 the MLT dynamics over southern Patagonia obtained using SIMONe Argentina. For this
 reason, our results are discussed as they are presented.

4.1 Mean winds and tides

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In Figure 3, we present the mean zonal (U_0) and meridional (V_0) winds, and the 240 total amplitudes of the quasi two-day planetary wave (Q2DW) and the diurnal (D1), semid-241 iurnal (S2) and terdiurnal (T3) solar tides. The vertical black dashed line indicates Jan-242 uary 1st 2020. The term total amplitude refers to the magnitude of the vector sum of 243 the corresponding zonal and meridional components of each fitted wave. The statisti-244 cal uncertainties of the estimated parameters are shown in the right column panels. All 245 quantities were obtained after applying the procedure detailed in Section 3.1. Data gaps 246 are shown in white. 247

From inspection of Figure 3, two features stand out: the S2 tide is the dominant 248 wave, with amplitudes in the order of 40-65 m/s, and the Q2DW exhibits strong enhance-249 ments after 4 January 2020. It is well known that the semidiurnal solar tide at middle 250 latitudes dominates over all other tidal components (e.g., Andrews et al., 1987; Pancheva 251 & Mukhtarov, 2011). Furthermore, many studies of tides in the northern hemisphere have 252 reported that S2 decreases significantly around the onset of a sudden stratospheric warm-253 ing (SSW) event, to later on recover and reach even larger amplitudes than those exhib-254 ited prior to the SSW (e.g., Chau et al., 2015; Siddiqui et al., 2018; Conte et al., 2019). 255 In September of 2019, approximately 12 days before the 27th (first day of available data 256 from SIMONe Argentina), a SSW event was registered in the southern hemisphere (e.g., 257

Yamazaki et al., 2020). Interestingly enough, the largest amplitudes of S2 are seen be-258 tween 27 September and 12 October, which may be an indicative of the recovery phase 259 of S2 after the weakening associated with a SSW event. On the other hand, it is also pos-260 sible that the reported enhancement of S2 is simply the result of changes in the prop-261 agation conditions of this tidal component. For example, the S2 amplitudes may change 262 due to different responses of the symmetric and antisymmetric Hough modes to the weaker 263 and more hemispherically symmetric mean zonal wind typical of the equinox times. Pre-264 vious studies of the S2 tide at middle latitudes of the southern hemisphere have shown 265 a decrease of S2 in the beginning of September, and a later recovery by the beginning 266 of October (e.g., Conte et al., 2017). However, this decrease/increase behavior of S2 dur-267 ing the September-October months is not obvious every year. Lastly, note that for the 268 entire dataset analyzed in this study, S2 presents significant intraseasonal variability, which 269 becomes evident in the many, although weaker, enhancements observed after ~ 31 Oc-270 tober. 271

The Q2DW at middle latitudes has been reported to reach maximum amplitudes 272 during the summer (e.g., Kumar et al., 2018). In our results, the Q2DW is active mostly 273 in summer, in agreement with previous studies. Even more, it becomes the dominant wave 274 by the end of January 2020, with amplitudes larger than those corresponding to S2. Offermann 275 et al. (2011) showed that the Q2DW exhibits a triple peak structure in the NH during 276 summer. Although it may not be obvious at first, after a more through inspection of Fig-277 ure 3, it can be noticed that the largest amplitudes of the Q2DW are distributed in three 278 subsequent enhancements, around 7, 13 and 19 January. A fourth enhancement can be 279 seen around 9 February, but the latter is significantly weaker than the previous three. 280

Both the diurnal and terdiurnal solar tides exhibit considerable intraseasonal variability. In the case of D1, its activity becomes more evident mostly below ~ 90 km and during summer. Above ~ 92 km, and mainly during equinox times, T3 becomes more noticeable, with amplitudes similar to those corresponding to D1.

Compared to five-year average values at 54° S presented by Conte et al. (2017), the 285 summer reversal of the mean zonal wind shown in Figure 3 is observed at altitudes ap-286 proximately two km lower. This is consistent with previous studies reporting a decrease 287 with latitude of the height of the zonal wind summer reversal (e.g., Hoffmann et al., 2010; 288 Wilhelm et al., 2019). Besides, it might seem that U_0 starts the transition into summer 289 conditions relatively early, around 3 October. However, above $\sim 92 \text{ km}, U_0$ experiences 290 a late reversal to westward conditions around 24 October, to finally go into summer con-291 ditions (i.e., eastward above the mesopause) about five days later. V_0 blows mainly to-292 wards the equator, and only after 10 March poleward values start to dominate. 293

The statistical uncertainties of all the fitted parameters presented in Figure 3 are very small. Only above 103 km and below 77 km, values of $\sim 2-3$ m/s are obtained (not shown). The low statistical uncertainties are a consequence of the large amount of meteor detections provided by SIMONe Argentina. By solving Equation 1 in bins of one hour and two km, one guarantees wind estimates with very low uncertainties. The latter, combined with the fact the WLS method used to solve Equation 2 is applied to a very well conditioned matrix, results in small statistical uncertainties.

4.2 Gradients and vertical wind

In Figure 4, we present one week of 4-hour, 4-km mean winds and gradients obtained after applying the procedure described in Section 3.2. The first row shows, from left to right, the mean zonal (u_0) , meridional (v_0) and vertical (w_0) winds. The panels in the middle row are used to show the zonal eastward, northward and upward (first order) derivatives, i.e., u_x , u_y and u_z , respectively. The bottom row shows the same as the middle one, but for the meridional component, i.e., v_x , v_y and v_z . The mean winds are given in m/s, and the gradients in m/s/km. The mean vertical wind and the gradients



SIMONe Argentina - Mean winds & tides

Figure 3. Mean zonal (U_0) and meridional (V_0) winds, the quasi two-day planetary wave (Q2DW), and the diurnal (D1), semidiurnal (S2) and terdiurnal (T3) solar tides during September 2019 - April 2020 over SIMONe Argentina. The corresponding statistical uncertainties (in m/s) are shown in the right column. All parameters were estimated using a four-day running window shifted by one day. The vertical black dashed line indicates January 1st 2020.

are shown only between 80 and 97 km of altitude, where the data coverage is good enough
 to estimate them more reliably.

Before starting the description of the main features observed in Figure 4, it is im-311 portant to stress here that the variability seen in the mean winds is representative of large 312 scale structures, with periods greater than four hours and vertical wavelengths larger than 313 four km. Both u_0 and v_0 exhibit a clear diurnal and semidiurnal periodicity, the latter 314 being more evident, something expected given the latitude of SIMONe Argentina. The 315 semidiurnal periodicity is also very clear in the vertical gradients, u_z and v_z , which in-316 dicates a strong influence of the S2 tide in the vertical changes of the horizontal winds. 317 Although they may not be evident in Figure 4, other periodicities are present in the ver-318 tical gradients. For example, GW effects can be noticed below 90 km. At these altitudes, 319 u_z is mainly positive (with mean values in the order of 20 m/s/km). In fact, during most 320 of the summertime, below ~ 90 km of altitude, and when averaged over longer periods 321 of time, e.g., 1 day, u_z is strongly dominated by positive amplitudes (not shown here). 322 During the summer, westward zonal winds are dominant in the stratosphere and meso-323 sphere. This condition allows eastward propagating GWs to easily reach mesospheric al-324 titudes, where they break and deposit momentum and energy. This deposition of mo-325 mentum creates an eastward drag that decelerates the zonal (westward) wind. Due to 326 the Coriolis effect, the deceleration of the zonal wind introduces an equatorward merid-327 ional wind component, which in turn leads to an upward motion and a subsequent (adi-328 abatic) cooling of the mesopause region (e.g., Smith, 2012). Finally, due to the thermal 329 wind equation, $u_z \approx -\alpha(\phi) \partial T / \partial y$ (where y is positive towards north, T is the temper-330 ature, and α is a factor that depends on the sine of the latitude, ϕ), the adiabatic cool-331 ing results in the positive values of u_z . In the case of v_z , the amplitudes are in the or-332 der of 3-5 m/s/km weaker than those of u_z , but the semidiurnal periodicity is seen at 333

all altitudes. No evident pattern or structure in time can be noticed when v_z is averaged over longer periods of time.

In the case of u_x and v_x , amplitude values are in the order of 0.1-0.3 m/s/km. The 336 larger amplitudes of u_y (compared to u_x) may be an indicative of zonal wind latitudi-337 nal changes related to the Antarctic polar vortex, part of which usually locates below 338 the region seen by SIMONe Argentina (e.g., Figure 8 in Orte et al. (2019)). The upward 339 component of the relative vorticity can be coarsely approximated by $v_x - u_y$ (for a pre-340 cise calculation, one needs to include the latitude information; see Equation A16 in Chau 341 et al. (2017)). Below ~ 84 km, $v_x - u_y$ values are mainly positive, which suggests coun-342 terclockwise vortical effects. Besides, differences between the amplitudes of u_y , v_x and 343 u_x, v_y , suggest that changes in the horizontal gradients due to GWs depend on the prop-344 agation direction of the these waves. 345

The errors in both the horizontal and vertical gradients are in the order of 5-10 %. For example, for the week of data shown in Figure 4, the root mean square errors of u_x , u_y , u_z , v_x , v_y , and v_z are 0.011, 0.023, 1.1, 0.022, 0.013, and 0.99 m/s/km, respectively.

Previous observational studies have investigated the horizontal gradients in the zonal 349 and meridional winds (e.g., Conde & Smith, 1998; Meriwether et al., 2008; Chau et al., 350 2017). However, the present study is the first one to show results on both the horizon-351 tal and vertical gradients of MLT horizontal winds over southern Patagonia. Chau et 352 al. (2017) showed that horizontal divergence values of $\sim 0.1 \text{ m/s/km}$ are large enough 353 to introduce a mean apparent vertical wind of 1-2 m/s. After calculating $u_x + v_y$, one 354 obtains values in the order of 0.4-0.5 m/s/km. Consequently, one can now understand 355 the importance of estimating the vertical wind together with the gradients. By doing so, 356 one eliminates the effects of biases introduced in w_0 by the mean horizontal variability 357 of u and v. However, the horizontal variability within scales smaller than the observed 358 volume is not removed, and will affect the vertical wind estimates. 359

The four-hour mean vertical winds (w_0) exhibit large amplitudes and considerable 360 variability both in time and height. The errors in these estimates are in the order of 1-361 2 m/s. After computing daily averages, vertical wind amplitudes are reduced to a few 362 m/s, which are still considerably large values. Previous studies have investigated MLT 363 mean vertical winds using specular meteor radars. For example, Babu et al. (2012) and 364 Egito et al. (2016) reported that the vertical winds at low latitudes may reach magni-365 tudes of 6-10 m/s. Using VHF radar measurements, Fritts et al. (1990) observed that 366 at summer mesopause altitudes vertical velocities may reach values of up to 10 m/s, but 367 within time scales of 5-30 minutes. That is, vertical winds should exhibit values in the 368 order of a few to tens cm/s when averaged over periods of 1 day or more (e.g., Lu et al., 369 2017).370

There are several factors that adversely affect our vertical wind estimates. (1) The 371 daily variability in the meteor counts (meteor detections maximize early in the morn-372 ing and minimize in the evening) can result in irregular time sampling at altitudes where 373 the detections are lower, i.e., above 97 km and below 80 km, in the case of SIMONe Ar-374 gentina. For this reason, we have limited our analysis of w_0 to altitudes between 80 and 375 97 km. (2) Despite the multistatic configuration, sometimes the meteor detections are 376 not well distributed among all five links. This affects in a similar way every altitude ob-377 served by SIMONe Argentina. (3) Small-scale horizontal variability is not removed af-378 ter applying the gradient method. Hence, horizontal changes of the zonal and meridional 379 winds occurring within the observed volume can easily leak into the vertical winds, and 380 thus result in the large w_0 values presented in Figure 4. To better understand and quan-381 tify the adverse impact of points (2) and (3), we are currently working in applying the 382 gradient method to forward modelled data using different geometric configurations. This 383 will shed light on the reliability of the vertical winds presented in this study as well as 384 other studies using multistatic SMR configurations. 385



(4h, 4km) mean winds and gradients over SIMONe Argentina

Figure 4. SIMONe Argentina estimates of: (first row) mean zonal (u_0) , meridional (v_0) and vertical (w_0) winds; (second row) zonal eastward (u_x) , meridional (u_y) and upward (u_z) derivatives; (third row) meridional eastward (v_x) , northward (v_y) and upward (v_z) derivatives. All parameters were estimated using a 4 hour, 4 km (in vertical) bin, shifted by 15 min and 1 km, respectively.

Despite the large amplitudes, coherent phase structures can be seen in w_0 . For example, around 17 February one can observe a clear diurnal periodicity. Diurnal wave signatures in w_0 have been found throughout the entire SIMONe Argentina dataset (not shown). Moreover, similar features have been observed in the mean vertical winds over Peru and northern Germany (Chau et al., 2020; Charuvil Asokan et al., 2020). Hence, it is possible that our results on w_0 are also revealing real geophysical features, e.g., diurnal tidal effects. This will be carefully investigated in future studies.

393

4.3 Gravity-wave-driven momentum flux

We now present and discuss the momentum flux estimates obtained after subtract-394 ing the mean winds calculated following the gradient method, i.e., the u_0, v_0 and w_0 that 395 were estimated together with the gradients. In Figure 4, it can be seen that the hori-396 zontal wind gradients and the vertical mean wind exhibit considerable variability, both 397 in time and altitude. Part of this variability is leaked into the mean horizontal winds when 398 one solves Equation 1 without including the gradients and w_0 in the wind vector **u**. This 300 implies that the u_0 and v_0 obtained using the gradient method constitute a better rep-400 resentation of the real mean wind, provided that enough meteors are detected in order 401 to accomplish a robust wind estimation. Besides, it is important to have in mind that 402 the subtracted winds were estimated in bins of 4 hour and 8 km (in altitude). This means 403 that the corresponding momentum flux estimates are representative mostly of waves with 404 temporal and vertical scales of less than 4 hour and 8 km, respectively. 405

In Figure 5, we present 28-day averages of the momentum flux estimates that result from subtracting the \hat{f} that was calculated using the u_0 , v_0 , w_0 mentioned in the previous paragraph (see Equation 4). The averages were calculated over 28 days in order to obtain estimates that are statistically more significant. The upper panels present U and V, which correspond to the 28-day averages of the mean zonal and meridional winds, respectively. The middle panels are used to show the 28-day mean horizontal momentum vertical fluxes, $\langle u'w' \rangle$, $\langle v'w' \rangle$. The statistical uncertainties of $\langle u'w' \rangle$ and $\langle v'w' \rangle$, i.e., $\sigma(u'w')$ and $\sigma(v'w')$, respectively, are shown in the bottom panels.

Momentum flux estimations based on SMR observations were first presented by Hocking 414 (2005). Since then, several studies have investigated momentum fluxes using meteor radar 415 winds (e.g., Fritts et al., 2010; Placke et al., 2011b; Fritts et al., 2012b; Andrioli et al., 416 417 2013; Placke et al., 2015, and references therein). More recently, de Wit et al. (2016) observed a modulation by the quasi-biennial oscillation (QBO) of the momentum fluxes 418 over Tierra del Fuego, an island south of Santa Cruz province. Typical amplitudes of <u'w'> 419 and $\langle v'w' \rangle$ reported in these previous studies rarely reach values larger than 60-70 m²/s². 420 Using a multistatic meteor radar network over southern Australia, Spargo et al. (2019) 421 observed values for $\langle u'w' \rangle$ and $\langle v'w' \rangle$ in the order of 40-60 m²/s², which are larger 422 than momentum flux absolute values obtained from satellite measurements (e.g., Ern et 423 al., 2011; Trinh et al., 2018). The latter is most likely due to observational filtering in-424 herent to satellites. 425

From Figure 5, it can be seen that our momentum flux estimates have amplitudes 426 in the order of 25-45 m²/s². Some large amplitudes of about 90 m²/s² can be observed 427 above 96 km of altitude, where the corresponding statistical uncertainties are in the or-428 der of 10 % or more (see bottom panels). Both momentum flux estimates exhibit vari-429 ability in time and height, although it is more evident in the case of $\langle u'w' \rangle$. Notice that 430 despite the 28-day averaging, the magnitudes of $\langle u'w' \rangle$ and $\langle v'w' \rangle$ are considerable 431 throughout the entire period of time analyzed in this study. Momentum flux estimates 432 averaged over 20 days presented by Placke et al. (2015) exhibit maximum amplitudes 433 of 6-8 m^2/s^2 during the summer of the northern hemisphere. The largest amplitudes of 434 ten-day momentum flux estimates shown by Spargo et al. (2019) reach values of 30-40 435 m^2/s^2 in the beginning of the spring of the southern hemisphere. On the other hand, de 436 Wit et al. (2017) observed monthly mean momentum fluxes over Tierra del Fuego with 437 amplitudes larger than 70 m^2/s^2 during the summer. These amplitude values are in the 438 same order as those obtained in our study. Hence, it is clear that the wave activity in 439 the MLT region over the Patagonian sector is very strong. This is also consistent with 440 previous studies based on satellite measurements (e.g., Ern et al., 2011), and numerical 441 simulations (e.g., Lund et al., 2020). 442

As GWs propagate upwards, they transfer momentum and energy into the mean 443 flow. Consequently, a decrease in the vertical flux of zonal momentum should correspond 444 with an increase in the zonal wind speed (e.g., Fan et al., 1991). In other words, a de-445 crease of the momentum flux indicates that a wave drag has been introduced into the 446 mean flow, thus triggering an acceleration of the zonal wind. Our results are consistent 447 with this reasoning, since for altitudes below 90-92 km, <u'w'> amplitudes, which are 448 mostly positive, exhibit a clear vertical gradient: as the altitude increases, they progres-449 sively decrease from values of around 70-75 m^2/s^2 below 82 km, to values of ~ 10 m^2/s^2 450 above 90 km. The latter is also consistent with the results on u_z presented in the pre-451 vious section. Besides, note that after 22 February, negative values of <u'w'> start to 452 develop below 90-91 km. At that time of the year, U has become eastward at all alti-453 tudes observed by SIMONe, a condition that allows westward propagating GWs to reach 454 higher altitudes and, most likely, induce the aforementioned <u'w'> negative values. 455

In the case of $\langle v'w' \rangle$, an upward movement of southward momentum dominates mostly below ~ 91 km, from the beginning of October until the end of March. During the last week of December 2019, positive values of $\langle v'w' \rangle$ start to dominate above 92-94 km, an altitude range that had been dominated by negative values of $\langle v'w' \rangle$ since the beginning of November 2019. During the latter, and approximately above 96 km of altitude, an upward movement of eastward momentum can be noticed again. These positive values of <u'w'> develop very abruptly around 26 November, and remain dominant above 96 km for more than 25 days. We wonder if this might be an indicative of
eastward momentum deposition by GWs that were in-situ generated at altitudes above
90 km.

Momentum flux estimates obtained after subtracting the mean winds determined 466 with the homogeneous method were also analyzed (not shown). In this case, only the 467 4-hour, 8-km mean horizontal winds are used to determine f (see Equation 5). These 468 momentum flux estimates have similar amplitudes to those presented in Figure 5. How-469 470 ever, they exhibit some differences in their variability in time and altitude, and their statistical uncertainties are larger than those shown in Figure 5. It is for the latter and for 471 consistency reasons, that the present study has been focused on analyzing the momen-472 tum fluxes that result from subtracting the mean horizontal and vertical winds obtained 473 after applying the gradient method. 474

To finalize, we discuss the procedure followed to estimate $\sigma(\mathbf{u}'\mathbf{w}')$ and $\sigma(\mathbf{v}'\mathbf{w}')$, and 475 the reason for calculating 28-day averages. Unless one knows with some degree of cer-476 tainty that a given wave event has occurred, the effects of GWs should be treated as stochas-477 tic processes. In other words, the mean momentum flux estimates are highly dependent 478 on the effects of geophysical variability. Kudeki and Franke (1998) showed that in or-479 der to obtain statistically significant momentum flux estimates at mesospheric heights, 480 one must consider averaging intervals of more than 25 days. Specifically, they found that 481 the statistical uncertainty of $\langle u'w' \rangle$ can be approximated with: 482

$$\sigma(u'w') = \sqrt{\frac{(\langle u'u' \rangle)(\langle w'w' \rangle)}{T/\tau}},$$
(6)

where, $\langle u'u' \rangle$ and $\langle w'w' \rangle$ are averaged over the interval of time T; and τ is equal to 483 half of the mesosphere Brunt-Väisälä period (~ 7 min). $\sigma(v'w')$ is obtained using same 484 Equation 6, but replacing $\langle u'u' \rangle$ by $\langle v'v' \rangle$. Selection of an averaging window T =485 28 days resulted in the values presented in Figure 5. In this way, the $\langle u'w' \rangle$ and $\langle v'w' \rangle$ 486 obtained from our study must be understood as representatives of a monthly mean mo-487 mentum flux due to waves with periods of 4 h or less, and horizontal scales less than 400 488 km. Besides, for those estimates corresponding to altitudes lower than \sim 98 km, statis-489 tical uncertainties between 2 and 15-16 m^2/s^2 should be taken into consideration. 490

The momentum flux estimates are also affected by the correlated Doppler shift er-491 rors. In other words, because the Doppler shift uncertainties are squared when introduced 492 into Equation 4, the resulting momentum flux estimates are in fact an overestimation 493 of the real $\langle u'w' \rangle$ and $\langle v'w' \rangle$ (Vierinen et al., 2019). However, given that the error 494 propagated Doppler shift uncertainties we obtained are small (\sim one order of magnitude 495 smaller than the statistical uncertainties) and that the amplitudes of $\langle u'w' \rangle$ and $\langle v'w' \rangle$ 496 are large, the effects of the correlated errors were neglected in this study. We plan to fur-497 ther investigate this issue by extending our analysis to include non-zero lag second-order 498 statistics of the wind velocities. 499

500 5 Concluding remarks

The first multistatic meteor radar based studies of mean winds, tides, gradients and 501 momentum flux over the southern part of Patagonia have been presented in this paper. 502 By doing this, we have demonstrated the ability of SIMONe Argentina to obtain not only 503 information on typical MLT parameters such as mean winds and tides, but also to suc-504 cessfully estimate previously little investigated parameters, such as horizontal and ver-505 tical gradients of the horizontal winds. Using the latter, one can estimate, e.g., the hor-506 izontal divergence and the relative vorticity, parameters from which global circulation 507 models can benefit and in this way help to further the understanding of MLT dynam-508 ics. 509



After removing (4h, 8-km) mean winds (w/ gradients) / 28-day average

Figure 5. 28-day averages of 4-h, 8-km mean zonal and meridional winds (U and V, respectively), horizontal momentum vertical fluxes ($\langle u'w' \rangle$, $\langle v'w' \rangle$) and their corresponding statistical uncertainties ($\sigma(u'w')$, $\sigma(v'w')$). 4 hour, 8 km (in altitude) mean horizontal and vertical winds estimated in combination with the gradients (i.e., the u_0 , v_0 , w_0 described in Section 3.2) were subtracted before estimating the momentum flux. The vertical black dashed line indicates 1 January 2020.

Our results show a strong positive vertical gradient in the zonal wind during the summer, in agreement with the residual mean meridional circulation. Besides, the northward zonal gradient (u_y) is stronger than the eastward zonal gradient (u_x) , which indicates latitudinal changes of the zonal wind most likely due to the influence of the southern hemisphere polar vortex. Mean vertical winds (w_0) have also been estimated, but only when the horizontal and vertical gradients were taken into account.

Momentum fluxes, $\langle u'w' \rangle$ and $\langle v'w' \rangle$, have been estimated after removal of hor-516 izontal and vertical mean winds that were fitted together with the horizontal wind gra-517 dients. Compared to some previous studies, our momentum flux estimates exhibit larger 518 amplitudes, which indicates that the GW activity in the MLT over southern Patagonia 519 is very strong. The statistical uncertainties of $\langle u'w' \rangle$ and $\langle v'w' \rangle$ were also included 520 in the analysis. The latter revealed that in order to have statistically significant momen-521 tum flux estimates, one should consider averages of at least 28 days. In this way, our re-522 sults must be considered as representative of monthly mean momentum fluxes, driven 523 by waves with periods shorter than 4 hours, vertical wavelengths shorter than 8 km, and 524 horizontal scales less than 400 km. 525

We are confident that SIMONe Argentina has also the potential to investigate nonzero lag second-order statistics of MLT wind velocities, e.g., by using correlation function techniques such as those presented in Vierinen et al. (2019). This will be explored in future studies. Besides, we also plan to investigate momentum flux estimates without averaging over long periods of time, provided there is evidence of specific (deterministic) wave events occurring in the troposphere/stratosphere.

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