# Solar and magnetic control of minor ion peaks in the dayside Martian ionosphere

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#### Abstract

The Neutral Gas and Ion Mass Spectrometer of the Mars Atmosphere and Volatile Evolution provides a large data set to explore the ion composition and structure of the Martian ionosphere. Here the dayside measurements are used to investigate the minor ion density profiles with distinctive peaks above 150 km, revealing a systematic trend of decreasing peak altitude with increasing ion mass. We specifically focus on a subset of species including  $O\$^+\$$ ,  $N\$_2^+\$$ ,  $C\$^+\$$ ,  $C\$^+\$$ ,  $N\$_2^+\$$ ,  $N\$_2^+\$$ ,  $C\$^+\$$ ,  $N\$_2^+\$$ ,  $N\$_2^+\$$ ,  $C\$^+\$$ ,  $N\$_2^+\$$ ,  $N\$_2^+$ ,  $N\$_2^+\$$ ,  $N\$_2^+$ ,  $N\ast_2^+$ ,

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#### **Key Points:** 14

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15	٠	The peak density and altitude for most minor ion species produced via direct pho-
16		toionization show weak or no SZA variation.
17	•	The minor ion peak density and altitude tend to increase significantly with increas-
18		ing solar activity.
19	•	The minor ion peak density and altitude show clear difference between the strongly
20		and weakly magnetized regions.

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#### 21 Abstract

The Neutral Gas and Ion Mass Spectrometer of the Mars Atmosphere and Volatile 22 Evolution provides a large data set to explore the ion composition and structure of the 23 Martian ionosphere. Here the dayside measurements are used to investigate the minor 24 ion density profiles with distinctive peaks above 150 km, revealing a systematic trend 25 of decreasing peak altitude with increasing ion mass. We specifically focus on a subset 26 of species including O<sup>+</sup>, N<sup>+</sup><sub>2</sub>/CO<sup>+</sup>, C<sup>+</sup>, N<sup>+</sup>, He<sup>+</sup>, and O<sup>++</sup>, all of which are mainly pro-27 duced via direct photoionization of parent neutrals. Our analysis reveals weak or no vari-28 29 ation with solar zenith angle (SZA) in both peak density and altitude, which is an expected result because these ion peaks are located within the optically thin regions sub-30 ject to the same level of solar irradiance independent of SZA. In contrast, the solar cy-31 cle variations of peak density and altitude increase considerably with increasing solar ac-32 tivity, as a result of enhanced photoionization frequency and atmospheric expansion at 33 high solar activities. He<sup>+</sup> serves as an exception in that its peak density increases to-34 wards large SZA and meanwhile shows no systematic variation with solar activity. The 35 thermospheric He distribution on Mars should play an important role in determining these 36 observed variations. Finally, the peak altitudes for all species are elevated by at least sev-37 eral km within the weakly magnetized regions, possibly attributable to the suppression 38 of vertical diffusion by preferentially horizontal magnetic fields in these regions. 39

#### 40 **1** Introduction

Mars possesses a well-defined ionosphere on the sunlit side that is mainly produced
by solar Extreme Ultraviolet (EUV) and X-ray ionization (Withers, 2009). The electron
density distribution in the dayside Martian ionosphere has been extensively studied over
the past few decades, thanks to a rich data set accumulated by both radio occultation
and radar sounding experiments made on board several spacecrafts such as the Mars Global
Surveyor (e.g. Tyler et al., 2001) and the Mars Express (e.g. Gurnett et al., 2005; Pätzold
et al., 2005).

Existing analyses reveal that the electron density,  $N_e$ , near the main ionospheric peak varies with altitude, z, according to the Chapman theory formulated as

$$N_e = N_m \exp\left(\frac{1}{2}\left(1 - \frac{z - z_m}{H} - \exp\left(-\frac{z - z_m}{H}\right)\right)\right),\tag{1}$$

where  $N_m$  is the peak electron density,  $z_m$  is the peak electron altitude, and H is the 50 51 scale height of the background neutral atmosphere (Chapman, 1931a, 1931b). This idealized theory predicts systematic variations of the peak electron density and altitude with 52 both solar zenith angle (SZA) and solar EUV and X-ray flux, as fully supported by nu-53 merous studies available in the literature (e.g. Hantsch & Bauer, 1990; Morgan et al., 54 2008; Fox & Yeager, 2009; Fox & Weber, 2012; Yao et al., 2019). It is well-known that 55 the peak electron altitude corresponds to where unit optical depth is reached due to at-56 mospheric photoabsorption, implying that its location also responds to features in the 57 background atmosphere such as the non-migrating tides (e.g. Bougher et al., 2001; Mendillo 58 et al., 2017) and planet-encircling dust storms (e.g. Wang & Nielsen, 2003; Fang et al., 59 2020). Meanwhile, many studies have indicated clearly that the presence of strong crustal 60 magnetic anomalies has an appreciable impact on the electron density distribution (e.g. 61 Ness et al., 2000; Diéval et al., 2015; Venkateswara Rao et al., 2017; Diéval et al., 2018; 62 Mohanamanasa et al., 2018; Fallows et al., 2019). 63

In contrast, the ion composition of the dayside Martian ionosphere was historically very limited, with only two individual measurements made by the Retarding Potential Analyzers (RPA) on board the Vikings 1 and 2, suggesting  $O_2^+$  as the dominant ion species followed by O<sup>+</sup> and CO<sub>2</sub><sup>+</sup> (Hanson et al., 1977). Such a situation has been greatly improved with the arrival of the Mars Atmosphere and Volatile Evolution (MAVEN) space-

craft at the red planet in September 2014 (Jakosky et al., 2015), with its Neutral Gas 69 and Ion Mass Spectrometer (NGIMS) capable of measuring a rich variety of species in 70 the Martian ionosphere covering a broad mass range of 2-150 Da (Mahaffy et al., 2015). 71 The preliminary NGIMS results revealed the presence of more than a dozen species per-72 sistently seen in the dayside Martian ionosphere (Benna, Mahaffy, Grebowsky, Fox, et 73 al., 2015) and the presence of transient metallic ion species when Mars was perturbed 74 by the near collision with Comet C/2013 A1 (Siding Spring) (Benna, Mahaffy, Grebowsky, 75 Plane, et al., 2015). 76

With the aid of the large NGIMS data set, several recent studies have focused on 77 the structural variability of various ion species in the dayside Martian ionosphere. Dur-78 ing the nominal mission phase, the MAVEN periapsis was typically at 150-160 km, i.e., 79 above the peak altitude of most ion species (e.g. Fox & Weber, 2012) and characteriz-80 ing the topside ionosphere only. For instance, the NGIMS analysis of Wu et al. (2019) 81 revealed a near constant density scale height of 100 km for all ion species on the day-82 side and meanwhile a clear impact of the ambient magnetic field configuration. A sim-83 ilar magnetic control of the ion distribution was obtained by Withers et al. (2019). Girazian, 84 Halekas, et al. (2019) further reported that high Solar Wind (SW) dynamical pressures 85 led to the depletion of all species in the topside Martian ionosphere, a trend persistently 86 seen at all SZAs and in both strong and weak magnetic field regions. 87

Several ion species do show clear peak structures at high altitudes. For instance, 88 the dayside  $O^+$  peak was observed to be at 220-300 km and its variations with season, 89 SZA, and solar ionizing flux were characterized by Girazian, Mahaffy, et al. (2019). Oc-90 casionally, the MAVEN spacecraft made Deep Dip (DD) campaigns down to a periap-91 92 sis altitude as low as 120-130 km, allowing the properties near the main ionospheric peak, including the ion composition, to be investigated (Vogt et al., 2017). Despite the exist-93 ing efforts, a variety of minor ion species in the dayside Martian ionosphere, with clear 94 density peaks at sufficiently high altitudes to be sampled by the NGIMS during the nom-95 inal MAVEN mission phase, have not been explored in detail. This serves as the main 96 motivation of the present study. 97

The paper is organized as follows. In Section 2, the peak densities and altitudes of minor ion species in the dayside median sense are derived, and the variability among different species is discussed. For a selected subset of minor ion species, the variations of their derived peak parameters, along with possible interpretations, are then presented. Specifically, we focus on the variations with SZA and solar activity in Section 3, as well as the variations with magnetic field configuration in Section 4. Finally, we discuss and draw conclusions in Section 5.

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### 2 Parameterization of minor ion peaks

The analysis presented in this study relies mainly on the NGIMS ion density mea-106 surements made in the Open Source Ion mode, which are ideally suited for character-107 izing the structure and composition of the Martian ionosphere (Mahaffy et al., 2015). 108 Here we include a total number of  $\sim 1100$  dayside MAVEN orbits from October 2014 109 to June 2018, with periapsis SZA below  $75^{\circ}$  and irrespective of latitude, longitude, as 110 well as solar activity. The dayside median density profiles of various ion species follow-111 ing the initial identification of Benna, Mahaffy, Grebowsky, Plane, et al. (2015) are dis-112 played in Figure 1, from the typical MAVEN periapsis of 150 km during the nominal mis-113 sion phase up to 450 km. 114

The figure reveals clearly the structural diversity of the Martian ionosphere with the detected ion species falling into two broad categories. On the one hand, the peak altitudes of a variety of ion species are located below the displayed altitude range, including the most dominant ion species,  $O_2^+$ , as well as some other ion species including  $CO_2^+$ ,



Figure 1. The structure of the dayside Martian ionosphere in terms of the median density profiles of all ion species measured by the NGIMS with SZA below  $75^{\circ}$  and reported in Benna, Mahaffy, Grebowsky, Plane, et al. (2015). While the peak altitude of the most abundant species,  $O_2^+$ , is clearly located below the lower boundary of 150 km, a variety of minor ion species present distinctive peaks at sufficiently high altitudes to be sampled by the NGIMS during the MAVEN nominal mission phase.

NO<sup>+</sup>, and OCOH<sup>+</sup>. On the other hand, many other minor ion species do present dis-119 tinctive peaks above 150 km, including  $O^+$ ,  $N_2^+/CO^+$ ,  $C^+$ ,  $N^+$ ,  $Ar^+$ ,  $O^{++}$ ,  $He^+$  in the 120 top panel and  $HCO^+/HOC^+/N_2H^+$ ,  $OH^+$ ,  $H_2O^+$ ,  $ArH^+$ ,  $H_2^+$ ,  $H_3O^+$ ,  $NH^+$ ,  $CH^+$  in the 121 bottom panel, both in the order of declining peak density. In addition, two species in 122 the figure,  $HNO^+$  and  $HO_2^+$ , marginally show the appearance of a layer structure peaked 123 near 150 km. Note that  $\overline{CO^+}$  and  $N_2^+$  cannot be distinguished by the NGIMS, with a 124 mass resolution of 1 Da, due to their near equality in mass per charge (Mahaffy et al., 125 2015).  $HCO^+$ ,  $HOC^+$ , and  $N_2H^+$  cannot be distinguished for the same reason. 126

For the purpose of this study, we derive the peak parameters, including the peak 127 density and peak altitude, from the observed distribution of each of the 15 minor ion species 128 quoted above with clearly observed layer structures. This is implemented by an empir-129 ical fitting of the density profile within an altitude width of 60 km centered at the ob-130 served maximum using the idealized Chapman function given by Equation 1. The peak 131 parameters are then straightforwardly given by the best-fit values of  $N_m$  and  $z_m$  in the 132 equation for each species involved. We do not use directly the density and altitude of 133 the observed maximum as the peak parameters in order to eliminate possible fluctua-134 tions in the ion density distribution forced either by gravity waves from below (e.g. Eng-135 land et al., 2017; Siddle et al., 2019) or by SW interactions from above (e.g. Kopf et al., 136 2008). While this may not be necessary for the situation depicted in Figure 1 since the 137 ion density fluctuations have been effectively removed by combining all dayside measure-138 ments, we persist in applying the Chapman fitting throughout this study because we will 139 encounter, in the following section, many cases for which the ion density profiles are not 140 sufficiently smooth with localized irregularities superimposed on their large scale trends. 141 A similar procedure has also been utilized to obtain the peak electron density and al-142 titude from radio occultation measurements (e.g. Yao et al., 2019). 143



Figure 2. The dayside median density profiles of various ion species (grey dots), superimposed by the best-fit Chapman profiles (dashed orange lines). Each Chapman profile is constrained by the NGIMS measurements centered around the observed maximum with a common width of 60 km (black solid lines). The identified peaks are marked for clarification (red circles), with the best-fit peak parameters provided in the figure legend. Note that the dayside median  $H_2^+$  distribution presents two peaks which are fitted separately.



Figure 3. The derived minor ion peak altitude in the dayside median sense as a function of the ion mass, for all the 15 minor ion species displayed in Figures 1 and 2. The dayside median  $H_2^+$  distribution presents two distinctive layer structures and only the peak altitude for the upper layer is indicated.

In Figure 2, we show the best-fit Chapman profiles superimposed on the dayside 144 median NGIMS observations for all minor ion species with distinctive peak structures. 145 In each panel, the portion of the NGIMS density profile used for constraining the peak 146 parameters is indicated by the thick solid line, whereas the extension of the Chapman 147 profile to the full displayed altitude range is indicated by the dashed line to demonstrate 148 the restricted validity of the Chapman formulism in describing the minor ion distribu-149 tion. The best-fit peak density and altitude appropriate for the dayside median situa-150 tion, as well as their uncertainties, are provided in the figure legend of each panel for ref-151 erence, along with the respective  $R^2$  goodness-of-fit. Note that the  $H_2^+$  distribution in 152 the dayside Martian ionosphere is characterized by two separate peaks, including a lower 153 one at 167 km and an upper one at 300 km, both with a peak density of around  $2 \text{ cm}^{-3}$ . 154 The best-fit Chapman profiles for both peaks are indicated in Figure 2. 155

For the 15 species examined here, the peak density covers a wide range from  $0.3 \text{ cm}^{-3}$ 156 for  $CH^+$  to 720 cm<sup>-3</sup> for  $O^+$ , which should rely on the abundances of their parent species 157 in the ambient atmosphere and ionosphere as well as the efficiencies of their dominant 158 chemical production and loss pathways involved (e.g. Krasnopolsky, 2002; Fox & Yea-159 ger, 2006; Fox, 2009; Matta et al., 2013; Fox, 2015). Of more interest is the observation 160 of minor ion peak altitude that obviously decreases with increasing ion mass, as depicted 161 in Figure 3. For  $H_2^+$ , only the peak altitude of the upper layer is indicated. According 162 to the figure, the derived peak altitude ranges from around 300 km for relatively light 163 ion species such as  $H_2^+$  and  $He^+$  to around 160 km for heavy ones such as ArH<sup>+</sup>. Mean-164 while, the observation that even heavier ion species such as  $\rm CO_2^+$  and  $\rm OCOH^+$  present 165 layer structures below the altitude range displayed in Figures 1 and 2, along with the 166 well established fact that the peak altitude of the dominant ion species,  $O_2^+$ , is at 130-167 140 km in the dayside median sense (e.g. Fox & Weber, 2012), is fully compatible with 168 the aforementioned trend. 169

The observed mass dependence of ion peak altitude could be interpreted as follows. Unlike O<sub>2</sub><sup>+</sup> with peak altitude located within regions under near photochemical equilibrium (PCE), most of the species displayed in Figure 3 are peaked at high altitudes where the condition of PCE is violated (e.g. Mendillo et al., 2011). For such a situation, we may assume for simplicity that the ion peak altitude corresponds to where the ion diffusion and chemical loss timescales are identical. The ion diffusion timescale is proportional to the square of the ion scale height divided by the ion diffusion coefficient, of which



Figure 4. The distribution of the MAVEN orbits with respect to the periapsis SZA and solar ionizing flux, both as a function of the date of observation. The portion of the entire available data set included in this study is marked by red. The solar ionizing flux is obtained by integrating the solar EUV and X-ray model spectrum of Thiemann et al. (2017) over the wavelength range from 0.5 nm to 90 nm.

the former scales as  $M_i^{-2}$  and the latter scales as  $M_i^{-1}$  where  $M_i$  is the ion mass. With-177 out loss of generality, we may further assume that for each ion species, the dominant chem-178 ical loss pathway is its reaction with  $CO_2$ , indicating that the  $CO_2$  density right at the 179 ion peak should be proportional to  $M_i$ . This naturally implies a higher ambient CO<sub>2</sub> den-180 sity at the peak and consequently a lower peak altitude for heavy ion species as com-181 pared to light ones. The above line of reasoning is subject to several over-simplifications 182 such as the neglect of the diversity of pathways for ion chemical loss and the neglect of 183 the mass dependence of binary ion collision frequency. In addition, the peak altitudes 184 of some species such as  $CO_2^+$  and  $OCOH^+$  are located within regions under PCE, im-185 plying a different mechanism responsible for the formation of their peaks from the mech-186 anism addressed above. While a robust interpretation of the NGIMS observations shown 187 in Figure 3 clearly relies on detailed photochemical model calculations, the simplified ar-188 gument presented here is able to provide useful insights into the underlying physics and 189 highlight the role of mass dependent ion diffusion in controlling the location of the ion 190 density peak in the dayside Martian ionosphere. 191

#### <sup>192</sup> 3 Solar control of minor ion peak parameters

To investigate the variations of the minor ion peak structure with both SZA and 193 solar activity, the Chapman fitting procedure outlined above is applied to the median 194 ion density profiles obtained from a selected group of orbits in our sample, with the de-195 tailed scheme of grouping dependent on the variation that we intend to seek. The dis-196 tribution of the MAVEN orbits used in this study is depicted in Figure 4 with respect 197 to the periapsis SZA and solar ionizing flux, both as a function of the date of observa-198 tion. The latter is obtained by integrating the solar EUV and X-ray model spectrum con-199 structed with the aid of the MAVEN Extreme Ultraviolet Monitor band irradiance data 200 (Eparvier et al., 2015; Thiemann et al., 2017). The integration is performed from 0.5 nm 201 up to a maximum wavelength of 90 nm corresponding to the  $CO_2$  ionization potential 202



**Figure 5.** The SZA variations of the derived peak densities for various minor ion species in the dayside Martian ionosphere.

(Masuoka, 1994). The portion of the available data set included in this study is marked
 by red in the figure. Figure 4 demonstrates that the available NGIMS data set is not evenly
 sampled with low SZA measurements preferentially made at relatively high solar activ ities.

We consider in this study the SZA and solar cycle variations of a subset of minor 207 ion species including  $O^+$ ,  $N_2^+/CO^+$ ,  $C^+$ ,  $N^+$ , He<sup>+</sup>, and  $O^{++}$ . A common feature of these 208 species is that the dominant production channel is direct solar EUV and X-ray ioniza-209 tion. For instance,  $O^+$  is mainly produced from both single photoionization of O and dis-210 sociative photoionization of CO or  $CO_2$ . Here  $Ar^+$  is excluded as an exception because 211 its peak is close to the lower boundary which, along with the broad appearance of the 212 peak (see Figures 1 and 2), does not allow the peak parameters to be accurately deter-213 mined in some cases. Several other species such as  $OH^+$  and  $H_2O^+$ , though with clear 214 peaks well characterized by the data, are not included in our investigation because they 215 are mainly produced via ion-neutral reactions instead of direct photoionization. As an 216 example, the reaction between  $O^+$  and  $H_2$ , instead of the direct photoionization of  $H_2O$ , 217 is the dominant channel producing ionospheric  $OH^+$  on Mars due to the low  $H_2O$  abun-218 dance in the Martian upper atmosphere (e.g. Fox et al., 2015). As presented in Cui et 219 al. (2020), the density variation of each of these species exhibits a very complicated pat-220 tern and possibly a strong dawn-dusk asymmetry. 221

We start with the SZA variations of the derived minor ion peak parameters. For 222 this purpose, we divide the NGIMS data set into several subsamples, each covering a lim-223 ited SZA range with a width of 5°. To avoid contamination by possible solar cycle vari-224 ation (see below), we restrict our analysis to those measurements made with the solar 225 ionizing flux in the range of  $0.7-0.9 \text{ mW m}^{-2}$ , appropriate for the low solar activity con-226 dition. The SZA variations of the derived peak parameters for all the 6 minor ion species are displayed in Figure 5 for peak density and Figure 6 for peak altitude, respectively. 228 The uncertainties in these peak parameters, which are not displayed in the figures, are 229 typically of comparably small amount as those quoted in Figure 2 legend due to the large 230 number of orbits available for each subsample. 231

Despite the considerable scattering, both figures suggest weak or no SZA variation in either peak parameter over the SZA range from 25° to 75°. This is an expected result because the atmosphere should be optically thin near the peaks of the minor ion species involved here, implying that regions at different SZA feel roughly the same level of so-



**Figure 6.** The SZA variations of the derived peak altitudes for various minor ion species in the dayside Martian ionosphere.



Figure 7. Similar to Figure 5 but for the minor ion peak density as a function of the solar ionizing flux (scaled by  $0.8 \text{ mW m}^{-2}$ ). The dashed lines correspond to the best-fit power law relations given by Equation 2 in the text.

lar EUV and X-ray irradiance. The conclusion of either weak or no SZA variation for
 the O<sup>+</sup> peak parameters was also reported by Girazian, Mahaffy, et al. (2019).

In contrast, He<sup>+</sup> appears to be an exception with its peak density exhibiting a strong 238 increasing trend towards the near terminator regions, despite that the peak altitude still 239 remains roughly independent of SZA. Quantitatively, Figure 5 suggests that the peak 240 density variation for the remaining 5 species is no more than 25% whereas the enhance-241 ment in the He<sup>+</sup> peak density at large SZA could reach more than a factor of 3 as com-242 pared to the low SZA value. This feature is indicative of an enhanced He abundance at 243 large SZA in the ambient atmosphere caused by the subsidence in regions of horizontal 244 wind convergence and the subsequent buildup of minor atmospheric species with large 245 vertical scale heights such as  $H_2$  and  $H_2$  (e.g. Elrod et al., 2017). The above discussions 246 indicate that the SZA variation of a minor ion species in the dayside Martian ionosphere 247 is strongly modulated by the variation of the background neutral atmosphere (e.g. Mendillo 248 et al., 2017). Similar observations have also been reported for the diurnal variations of 249 several protonated ion species (Cui et al., 2020). 250



Figure 8. Similar to Figure 6 but for the minor ion peak altitude as a function of the solar ionizing flux (scaled by  $0.8 \text{ mW m}^{-2}$ ). The dashed lines correspond to the best-fit linear relations given by Equation 3 in the text.

We next move on to the solar cycle variations of minor ion peak parameters by di-251 viding the NGIMS data set into consecutive subsamples in increasing order of solar ion-252 izing flux from 0.7 mW m<sup>-2</sup> to 2.1 mW m<sup>-2</sup> with a common interval of 0.2 mW m<sup>-2</sup>. 253 Due to the absence of any strong SZA variation in either peak parameter, we do not dis-254 tinguish between different SZAs, except for He<sup>+</sup> which is restricted to the SZA range up 255 to  $50^{\circ}$ . The variations of the derived peak parameters are shown in Figure 7 for peak 256 density and Figure 8 for peak altitude, respectively. Our analysis reveals a systematic 257 trend of increasing peak density and peak altitude with increasing solar ionizing flux for 258 each species, again with He<sup>+</sup> being the only exception that does not reveal any unam-259 biguous variation in its peak density. To be more quantitative, the corresponding Pear-260 son correlation coefficients, denoted as  $R_i$ , are computed for all the 6 minor ion species 261 and provided in Table 1 for reference. 262

The presence of strong solar cycle variation of the minor ion peak density is clearly 263 linked to a higher photoionization frequency when subject to a more intense solar ion-264 265 izing flux. A higher abundance of the parent neutrals likely makes a further contribution to the observed solar cycle variation, which is driven by enhanced photolysis of neu-266 trals in the dayside Martian upper atmosphere. Meanwhile, the Martian atmosphere ex-267 pands in response to increasing solar EUV and X-ray irradiance as indicated by a ris-268 ing exobase altitude (Fu et al., 2020). This naturally moves the ionosphere to higher al-269 titudes and causes the elevation of all minor ion peak altitudes, a mechanism that also 270 accounts for the known effect of planet-encircling dust storms on the Martian ionospheric 271 structure (e.g. Wang & Nielsen, 2003). The abnormal variation for  $He^+$  peak density 272 as displayed in Figure 7 is likely indicative of a lower He abundance in the dayside Mar-273 tian upper atmosphere at higher solar activities, which counterbalances a higher He pho-274 toionization frequency. Without showing the details, we mention that our conjecture is 275 verified by the NGIMS observation of a reduced He density from  $7 \times 10^4$  cm<sup>-3</sup> to  $4 \times$ 276  $10^4$  cm<sup>-3</sup> when the solar ionizing flux increases from 0.8 mW m<sup>-2</sup> to 2.0 mW m<sup>-2</sup>, both 277 referring to the respective He<sup>+</sup> peak altitude. 278

For a more quantitative parameterization of the observed solar cycle variations, they are described by a power law relation for peak density and a linear relation for peak altitude, in the forms of

$$N_{m,i} = \tilde{N}_{m,i} \left(\frac{I}{\tilde{I}}\right)^{\alpha_i},\tag{2}$$

**Table 1.** The best-fit parameters for various minor ion species, including the linear Pearson correlation coefficients,  $R_i$ , the peak densities,  $N_{m,i}$ , the peak altitudes,  $z_{m,i}$ , as well as the power law indexes,  $\alpha_i$ , and the linear slopes,  $\beta_i$ , that characterize the extent to which  $N_{m,i}$  and  $z_{m,i}$  vary with the solar ionizing flux (see Equations 2 and 3 in the text). Uncertainties for all parameters except for  $R_i$  are also provided. The parameters,  $N_{m,i}$  and  $\alpha_i$ , for He<sup>+</sup> peak density are not provided due to the absence of visible solar cycle variation.

Ion Species	Peak Altitude Parameters			Peak Density Parameters		
	$R_i$	$z_{m,i} \ (\mathrm{km})$	$\beta_i \ (\mathrm{km})$	$R_i$	$N_{m,i} \; ({\rm cm}^{-3})$	$\alpha_i$
0+	0.982	$209 \pm 9$	$29\pm5$	0.948	$562 \pm 103$	$0.93 \pm 0.28$
$N_2^+/CO^+$	0.990	$185\pm7$	$29 \pm 4$	0.998	$93 \pm 4$	$1.09\pm0.07$
$C^{\tilde{+}}$	0.903	$210\pm22$	$28\pm12$	0.985	$13 \pm 2$	$1.45\pm0.23$
$N^+$	0.938	$219\pm17$	$28\pm9$	0.871	$11 \pm 3$	$0.73\pm0.37$
$\mathrm{He^{+}}$	0.891	$246\pm34$	$28\pm19$	0.008	N/A	N/A
$O^{++}$	0.963	$226\pm14$	$31\pm8$	0.952	$0.50\pm0.12$	$1.25\pm0.36$

282 and

$$z_{m,i} = \tilde{z}_{m,i} + \beta_i \left(\frac{I}{\tilde{I}}\right),\tag{3}$$

where I is the solar ionizing flux defined in Section 2,  $N_{m,i}$  and  $z_{m,i}$  are the peak density and altitude for species i,  $\tilde{N}_{m,i}$  and  $\tilde{z}_{m,i}$  are the respective values for a reference ionizing flux of  $\tilde{I} = 0.8 \text{ mW m}^{-2}$ ,  $\alpha_i$  and  $\beta_i$  are species dependent free parameters to be constrained by data-model comparison. The power law relation is implemented here to reflect the desired limiting behavior of zero peak density when the solar EUV and X-ray irradiance is switched off. The best-fit power law and linear models are superimposed in Figures 7 and 8 for comparison, with the respective best-fit parameters provided in Table 1.

Despite the considerable variability in minor ion peak altitude suggested by the NGIMS 291 measurements (see Section 2), different species are characterized by a comparable increase 292 in peak altitude, which is on average 44 km over the solar ionizing flux range displayed 293 in Figure 8. In contrast, the extent to which the peak density varies with solar activity 294 differs substantially from species to species. Neglecting He<sup>+</sup>, the variation is maximized 295 for  $C^+$  with a power index of ~ 1.5 and minimized for  $N^+$  with a power index of ~ 0.7. 296 The derived power indexes suggest that over the available range of solar ionizing flux, 297 the  $C^+$  peak density increases by a factor of nearly 4 and the  $N^+$  peak density increases 298 by a factor of 2 only. It is interesting to note from Table 1 that both the  $O^+$  and  $N_2^+/CO^+$ 299 peak densities are almost linearly correlated with the solar ionizing flux. 300

#### <sup>301</sup> 4 Magnetic control of minor ion peak parameters

In this section, we further investigate for  $O^+$ ,  $N_2^+/CO^+$ ,  $C^+$ ,  $N^+$ ,  $He^+$ , and  $O^{++}$ 302 the variations of their peak parameters with the ambient magnetic field configuration, 303 characterized by both draped fields formed via SW interactions (e.g. Brain et al., 2006) 304 and crustal fields that tend to distribute over the southern hemisphere of Mars (e.g. Langlais 305 et al., 2019). Again to avoid contamination by the strong solar cycle variation, our anal-306 vsis is restricted to NGIMS measurements made with the solar ionizing flux in the range 307 of 0.7-0.9 mW m<sup>-2</sup> appropriate for the low solar activity condition. For our purpose, two 308 subsamples are defined, one for weak magnetic field regions with crustal field intensity 309 below 10 nT and the other one for strong magnetic field regions with crustal field inten-310 sity above 30 nT, where the crustal magnetic field model at a fixed altitude of 400 km 311



**Figure 9.** A comparison of the minor ion density profiles between strongly and weakly magnetized regions, defined with the crustal field intensity below 10 nT and above 30 nT, respectively, at a reference altitude of 400 km according to the model of Langlais et al. (2019).

<sup>312</sup> based on Langlais et al. (2019) is used. The minor ion density profiles for the two sub<sup>313</sup> samples are compared in Figure 9, where both ion peaks are indicated by the solid cir<sup>314</sup> cles. The best-fit Chapman profiles are not shown in the figure to avoid over-crowdedness,
<sup>315</sup> but the best-fit peak parameters are indicated in the figure legend for reference.

The figure shows clearly that the density profiles for both regions are similar be-316 low 200 km, whereas at higher altitudes, the minor ion densities, including the peak den-317 sities, within the strongly magnetized regions tend to be considerably reduced up to at 318 least 350 km as compared to the weakly magnetized regions. Meanwhile, the magnetic 319 control of minor ion peak altitude is also visible in Figure 9 in that the peak altitude tends 320 to be higher when the ambient magnetic fields are weaker. To be more quantitative, the 321 variation of peak density with magnetic field intensity is maximized for  $C^+$  and  $N^+$  with 322 an enhancement of more than 15% near weak magnetic fields, whereas the variation of 323 peak altitude is maximized for  $He^+$  with an elevation by 27 km in the same regions. A 324 similar magnetic control of  $O^+$  density above 200 km has recently been reported by Withers 325 et al. (2019). 326

The effect of ambient magnetic fields on the Martian ionosphere occurs mainly via modification of plasma diffusion (e.g. Shinagawa & Cravens, 1989; Matta et al., 2015). Below 200 km, the effect of diffusion is usually negligible (e.g. Mendillo et al., 2011; Mukundan et al., 2020), which is responsible for the absence of magnetic control at these al-

titudes in Figure 9. At higher altitudes, the effect of diffusion is critically dependent on 331 the ambient magnetic field configuration. Recent investigations (e.g. Xu et al., 2017) in-332 dicate that the magnetic field lines tend to be more horizontal in the weakly magnetized 333 regions, implying a reduced effect of vertical diffusion relative to the strongly magnetized 334 regions with preferentially vertical field lines (see also Wu et al., 2019). This naturally 335 leads to an elevated minor ion peak altitude as observed. Meanwhile, when subject to 336 the same level of solar EUV and X-ray irradiance, the total content of ions is unchanged 337 but they are redistributed by vertical diffusion which is responsible for a reduced den-338 sity profile from 200 km to at least 350 km incorporating the peak region. The obser-339 vation reported here has to be distinguished from that of Wu et al. (2019) focusing on 340 the dayside regions well above the minor ion peaks and suggesting instead an enhanced 341 distribution near strong magnetic fields. The combination of the two works provides a 342 more thorough picture of the structural variability of minor ion distribution over the full 343 altitude range sampled by the NGIMS during the MAVEN nominal mission phase. 344

#### **5** Discussions and Conclusions

Historically, the ion composition of the Martian ionosphere was only available from 346 the RPA measurements made on board the Vikings 1 and 2 (Hanson et al., 1977). The 347 extensive data set accumulated by the MAVEN NGIMS has provided a unique oppor-348 tunity to explore the structural variability of the Martian ionosphere in terms of the mi-349 nor ion distribution (Mahaffy et al., 2015). The dayside median density profiles of many 350 ion species detected by the NGIMS show distinctive layers peaked above 150 km, the typ-351 ical periapsis altitude during the nominal mission phase (Benna, Mahaffy, Grebowsky, 352 Fox, et al., 2015). Whenever possible, the peak parameters for each ion species, includ-353 ing the peak density and peak altitude, are derived from the idealized Chapman fitting 354 to the measured densities over a restricted altitude range centered at the observed max-355 imum. 356

The derived peak altitude shows a clear anti-correlation with the ion mass, from around 160 km for heavy species such as  $ArH^+$  to more than 300 km for light species such as  $He^+$  and  $H_2^+$ . Assuming that the minor ion peak altitude is controlled by where the ion chemical loss and diffusion timescales become equal, we propose a simplified argument to interpret the above observation, predicting that the ambient atmospheric density at the ion peak should be roughly proportional to the ion mass. This naturally accounts for an elevated peak altitude for relatively light ion species as revealed by the NGIMS measurements.

We further investigate the solar control of a selected subset of minor ion species 365 with (1) distinctive peak structures above 150 km and (2) chemical production pathways 366 dominated by the direct photoionization of their parent neutrals. These ions include  $O^+$ , 367  $N_2^+/CO^+$ ,  $C^+$ ,  $N^+$ ,  $He^+$ , and  $O^{++}$  in the order of decreasing dayside median peak den-368 sity. For most species, both their peak densities and altitudes show either weak or no 369 variation with SZA, which is an expected result because the atmosphere is optically thin 370 near the peaks, implying that regions at different SZA feel roughly the same level of so-371 lar EUV and X-ray irradiance. The conclusion of weak or no SZA variation for the ion 372 peak parameters was also reported by Girazian, Mahaffy, et al. (2019), but for  $O^+$  only. 373 For He<sup>+</sup> as an exception, a substantial increase in peak density towards large SZA is sug-374 gested by the data, despite that the peak altitude still remains roughly constant. The 375 strong SZA variation in He<sup>+</sup> peak density, which is compatible with the diurnal varia-376 tions of several protonated species as recently reported by Cui et al. (2020), could be in-377 terpreted as driven by the variation of He in the ambient atmosphere, as a result of the 378 subsidence in regions of horizontal wind convergence and the subsequent buildup of mi-379 nor atmospheric species with large vertical scale heights including He (Elrod et al., 2017). 380

The solar cycle variations of minor ion peak density and altitude are characterized 381 by a clear increase in either parameter with increasing solar ionizing flux, which is per-382 sistently seen for all species investigated here except for He<sup>+</sup>. The variation in peak den-383 sity is clearly the result of an enhanced photoionization frequency at high solar activ-384 ities, whereas the variation in peak altitude is caused by the expansion of the Martian 385 upper atmosphere and the consequent elevation of the ionospheric layer structure also 386 at high solar activities. It is interesting to note that the increase in peak altitude with increasing solar ionizing flux is of comparable amount for all species, by 44 km over the 388 solar ionizing flux range from  $0.8 \text{ mW m}^{-2}$  to  $2.0 \text{ mW m}^{-2}$ , but the increase in peak den-389 sity differs substantially among various minor ion species. For He<sup>+</sup> as an exception, no 390 clear solar cycle variation in peak density is observed despite that the peak altitude still 391 increases with increasing solar activity as normal. This is likely indicative of a reduc-392 tion in thermospheric He abundance at high solar activities which counterbalances a si-393 multaneous enhancement in He photoionization frequency. 394

Finally, the difference in minor ion peak parameters between strongly and weakly 395 magnetized regions is examined. Despite that no magnetic control of minor ion distri-396 bution is found below 200 km, the difference in ion distribution between different mag-397 netic field intensities is distinctive above 200 km, manifest as an enhanced ion distribu-398 tion up to at least 350 km in weakly magnetized regions. A similar magnetic control was 399 recently reported by Withers et al. (2019) for the three most abundant species,  $O_2^+$ ,  $O^+$ , and  $CO_2^+$ , in the Martian ionosphere. In addition, the minor ion peak altitude tends to 401 be elevated by at least several km in weakly magnetized regions. We propose that such 402 an observation is related to the preference of more horizontal field lines in weakly mag-403 netized regions (e.g. Xu et al., 2017; Wu et al., 2019), which suppresses the effect of vertical ion diffusion relative to local chemical loss and thus moves the ion peak to a higher 405 altitude. 406

The results presented here are useful towards establishing an overall picture of the 407 structural variability of minor ion distribution in the dayside Martian ionosphere, high-408 lighting the roles of solar illumination and magnetic field configuration as controlling fac-409 tors. These results are well suited for follow-up comparisons with realistic photochem-410 ical model calculations (e.g. Shinagawa & Cravens, 1989; Krasnopolsky, 2002; Fox & Yea-411 ger, 2006; Fox, 2009; Matta et al., 2013; Fox, 2015; Matta et al., 2015). However, we cau-412 tion that the present study relies on the analysis of a small portion of minor ion species 413 detected by the NGIMS. A thorough investigation of additional species, of which the chem-414 ical production channels are diverse and not restricted to direct photoionization of par-415 ent neutrals, may demonstrate a more complicated pattern of spatial and temporal vari-416 ations. 417

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Figure 1.



Figure 2.



Figure 3.



Figure 4.



Figure 5.



Figure 6.



Figure 7.



Figure 8.



Figure 9.

