

# Meteotsunami generation by tropical cyclone rainbands: nearshore effects of rainband dynamics and storm surge

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## Abstract

During tropical cyclones, meteotsunami waves can be triggered by atmospheric disturbances accompanying tropical cyclone rainbands (TCRs). Due to a paucity of high resolution field data along open coasts during these extreme events, relatively little is known about meteotsunami generation by TCRs and the coastal impact of these wave phenomena. Here we link high-resolution field measurements of sea-level and air pressure from Hurricane Harvey (2017) with a numerical model to show that large drops in air pressure accompanying trains of very narrow TCRs can initiate meteotsunami  $O(40\text{ cm})$  in height along open coasts distant from the storm center ( $>200\text{ km}$ ). The resonant-amplification and propagation of meteotsunami generated by pressure forcing is highly dependent on oceanographic (storm surge, bathymetry, and coastal morphology) and atmospheric factors (variable TCR forward speed, TCR path of translation). We discover that meteotsunami hazard can extend several days before and after hurricane landfall, and that meteotsunami are more ubiquitous along the open coast than tidal gauge records suggest, likely due to the highly-localized propagation and inherent structure of TCRs. This combined field and numerical study identifies the potential, but sometimes highly localized conditions necessary, for meteotsunami to modify storm processes (e.g., overwash, beach erosion) and serve as a coastal flood hazard during hurricane impact.

1                   **Meteotsunami generation by tropical cyclone**  
2                   **rainbands: nearshore effects of rainband dynamics and**  
3                   **storm surge**

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9                   **Key Points:**

- 10                   • Meteotsunami are more ubiquitous in the nearshore during tropical cyclones (TCs)  
11                    than tidal gauges suggest  
12                   • Air pressure anomalies accompanying outer TC rainbands can trigger meteotsunami  
13                     $O(0.4\text{ m})$  in height and with periods  $O(20\text{ min})$   
14                   • Meteotsunami can be highly localized and dependent on atmospheric and oceano-  
15                    graphic factors, including storm surge

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16 **Abstract**

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 25 amplification and propagation of meteotsunami generated by pressure forcing is highly  
 26 dependent on oceanographic (storm surge, bathymetry, and coastal morphology) and at-  
 27 mospheric factors (variable TCR forward speed, TCR path of translation). We discover  
 28 that meteotsunami hazard can extend several days before and after hurricane landfall,  
 29 and that meteotsunami are more ubiquitous along the open coast than tidal gauge records  
 30 suggest, likely due to the highly-localized propagation and inherent structure of TCRs.  
 31 This combined field and numerical study identifies the potential, but sometimes highly  
 32 localized conditions necessary, for meteotsunami to modify storm processes (e.g., over-  
 33 wash, beach erosion) and serve as a coastal flood hazard during hurricane impact.

34 **Plain Language Summary**

35 During tropical cyclones, spiral rainbands distant from the storm center can trig-  
 36 ger small variations in sea level known as meteorological tsunami (“meteotsunami”). Rel-  
 37 atively little is known about the forcing that initiates these waves and the beach haz-  
 38 ards they pose as it is challenging to collect field data during storm impact. This paper  
 39 combines new field measurements collected during Hurricane Harvey (2017) with model  
 40 simulations to show that large drops in air pressure accompanying passage of narrow spi-  
 41 ral rainbands can initiate meteotsunami approximately 40 cm in height along open coasts  
 42 more than 200 km away from storm landfall. Simulations show that the size of these waves  
 43 depends on oceanographic and storm-specific factors, including the amount of storm surge.  
 44 We discover that meteotsunami occur frequently along open coasts during tropical cy-  
 45 clones. The beach hazard associated with this phenomenon can be highly localized and  
 46 extend several days before and after hurricane landfall.

47 **1 Introduction**

48 Observational evidence has shown that tropical cyclones (TCs) can initiate tsunami-  
 49 like variations in sea level with periods of several minutes to hours (Mercer et al., 2002;  
 50 Mecking et al., 2009; Olabarrieta et al., 2017; Dusek et al., 2019; Shi et al., 2020). Tidal  
 51 gauge records suggest that meteorological tsunami (“meteotsunami”) are common in the  
 52 Gulf of Mexico (GOM) and along Florida’s Atlantic Coast during TCs, triggered by at-  
 53 mospheric disturbances accompanying tropical cyclone rainbands (TCRs) (Shi et al., 2020).  
 54 Meteotsunami contributions to total water levels at tidal gauge locations during TCs can  
 55 be significant (e.g., maximum crest elevation of 0.78 in a mean water depth of  $\sim 6$  m, Shi  
 56 et al. (2020)), and may also represent a hazard along beaches through initiation of ex-  
 57 treme wave runoff (e.g., the Daytona Beach 1992 meteotsunami, Churchill et al. (1995)),  
 58 rip currents (Linares et al., 2019), and dune and beach erosion. However, given the shel-  
 59 tered locations of tidal gauges along U.S. coasts (i.e., within estuaries, harbors, and bays),  
 60 and scarcity of field data close to shore during extreme events, relatively little is known  
 61 about meteotsunami hazard along open coasts during TCs.

62 Meteotsunami generation on open coasts is typically a multi-resonant process (Mon-  
 63 serrat et al., 2006). First, sudden changes in air pressure and/or wind speed associated  
 64 with a moving-atmospheric disturbance (e.g., storm, squall, frontal passage, or atmospheric

65 gravity wave) produce a small water level perturbation on the continental shelf due to  
 66 the inverse barometer effect. This perturbation can become strongly amplified as it prop-  
 67 agates with the disturbance, first due to external resonance processes and later as a free  
 68 wave due to nearshore wave transformations such as shoaling, refraction, and superpo-  
 69 sition of incident and reflected waves. Maximum energy transfer from the atmosphere  
 70 to the ocean occurs when the speed of the disturbance approaches the shallow water wave  
 71 celerity (Froude number near 1), a phenomenon called Proudman resonance (Proudman,  
 72 1929). Upon reaching the coast, additional resonance effects, such as the matching of har-  
 73 bor seiche periods, are required for the free wave to reach destructive heights (up to sev-  
 74 eral meters) - a phenomenon that has been observed in coastal basins around the world  
 75 (Rabinovich, 2019).

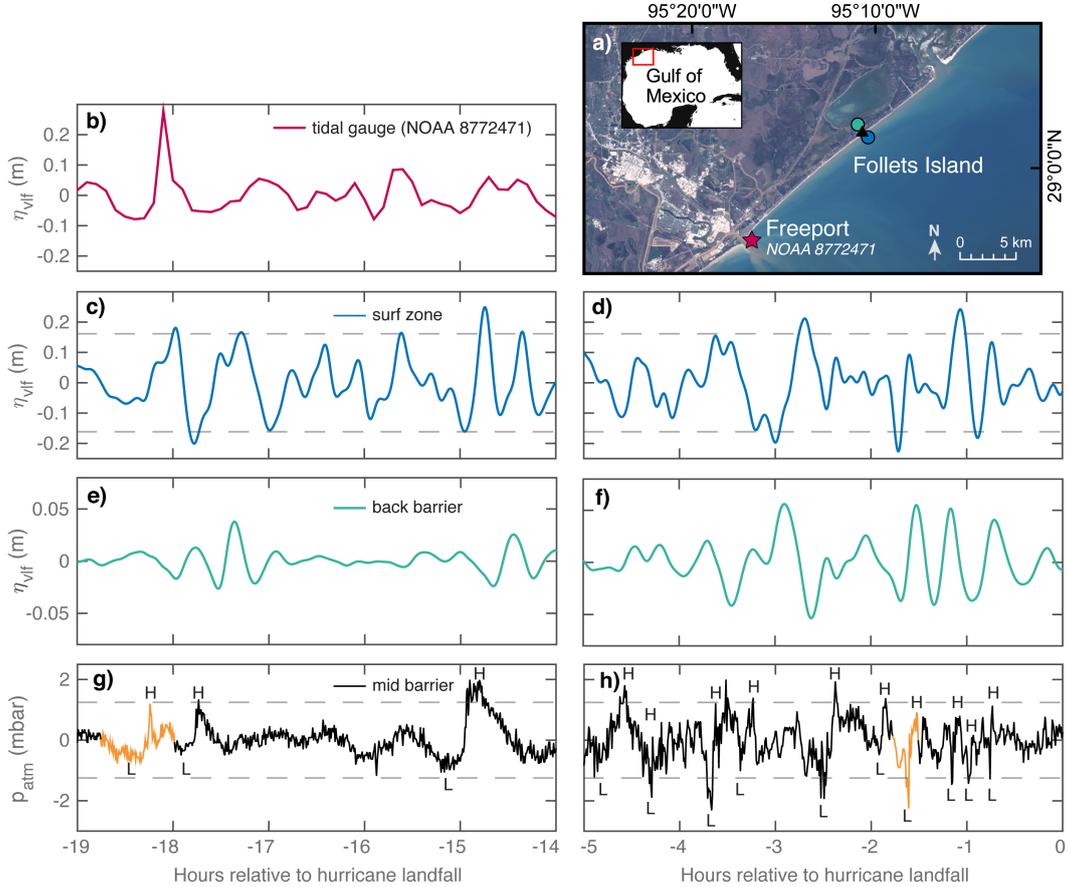
76 Using a coupled ocean and atmospheric modeling framework, Shi et al. (2020) showed  
 77 that for an idealized TC (and simplified bathymetry), wind stress dominates over pres-  
 78 sure forcing for meteotsunami generation by TCRs. The numerically simulated meteot-  
 79 sunami were similar to those observed in tidal gauge records, both in frequency (1-2 hr  
 80 periods), maximum crest elevation ( $\sim 0.2$ - $0.4$  m), and sequence (either single peak or se-  
 81 quential meteotsunami). Single peak meteotsunami were found to be driven by “outer”  
 82 TCRs – that is, rainbands that are distant from the inner core of the hurricane – and  
 83 sequential meteotsunami by trains of principal and secondary rainbands (“inner” TCRs)  
 84 within the inner core region. Notably, when using only pressure forcing to simulate TCR  
 85 propagation, the modeled sea-level anomalies generally followed the inverse barometer  
 86 effect and were not amplified through resonance processes.

87 During Hurricane Harvey (2017), in-situ measurements of sea level from nearshore  
 88 environments along the Texas coast showed variability in the meteotsunami frequency  
 89 band during passage of outer TCRs (here,  $\sim 200$  km from the storm center), albeit with  
 90 periods shorter than that modeled by Shi et al. (2020) (here,  $\sim 8$  to 45 min periods). For  
 91 the time period proximate to hurricane landfall, trains of outer TCRs with very small  
 92 horizontal scales ( $< 50$  km in arc-length) passed frequently over one field site (typically  
 93 every  $\sim 30$  minutes) and co-located observations of air pressure showed that they were  
 94 accompanied by large drops in air pressure. Here we test the hypothesis that many of  
 95 these very-low frequency sea-level anomalies are meteotsunami initiated by atmospheric  
 96 disturbances accompanying TCRs. In the absence of co-located measurements of wind  
 97 forcing, we numerically model the generation potential of meteotsunami by a very-narrow  
 98 outer TCR using pressure forcing alone. By linking high resolution field measurements  
 99 with numerical simulation, we show that large drops in air pressure  $O(2$  mbar) associ-  
 100 ated with trains of outer TCRs can trigger meteotsunami similar to observations of sea-  
 101 level anomalies in the nearshore, both in period  $O(20$  min) and height  $O(0.4$  m), via Proud-  
 102 man resonant wave growth.

## 103 2 Observations

104 Continuous measurements of sea level from two nearshore environments along the  
 105 Texas coast showed variability at frequencies  $f$  below infragravity waves ( $f < 3$  mHz,  
 106  $> 5.6$  min periods), but above known tidal constituents and storm surge ( $f > 0.1$  mHz,  
 107  $< 2.8$  hr periods) episodically during Hurricane Harvey (2017). For a complete descrip-  
 108 tion of both field sites, instrumentation, as well as a synopsis of Hurricane Harvey’s im-  
 109 pact, the reader is directed to Anarde et al., (*in review*) and Blake & Zelinsky (2018).  
 110 Herein, we examine observations of very-low frequency (“VLF”, 5.6 min to 2.8 hr pe-  
 111 riods) wave phenomena over a period of  $\sim 4$  days in the surf zone and back barrier at Fol-  
 112 lets Island (Figure 1a), a barrier island located  $\sim 200$  km northeast of hurricane land-  
 113 fall along the upper Texas Gulf coast. Sea-level fluctuations with periods upwards of 10  
 114 min have been observed in surf zones on beaches elsewhere, with generation mechanisms  
 115 typically attributed to shear instability of the alongshore current (Oltman-Shay et al.,  
 116 1989) or forcing from wave groups (Haller et al., 1999) (i.e., shear waves). The VLF sea-

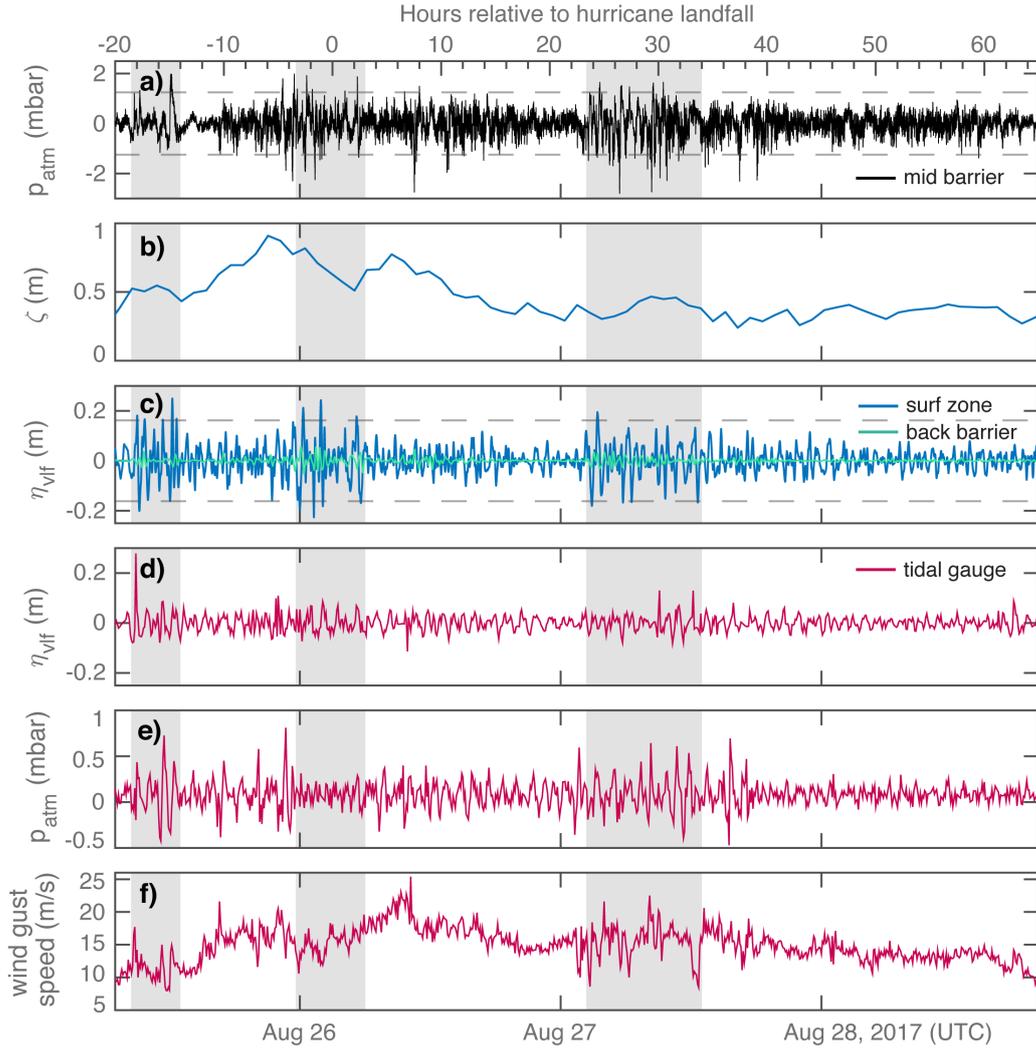
117 level anomalies observed in the surf zone during Hurricane Harvey were relatively large  
 118 (maximum peak-to-trough wave height = 42 cm, Figure 1d) and much lower in frequency  
 119 and span a larger frequency range ( $\sim 8$  to 45 min periods) than is typical of shear waves,  
 120 suggesting an alternative or additional generation mechanism(s).



**Figure 1.** a) Field site location; b-f) the VLF sea-level anomaly  $\eta_{vlf}$  measured at NOAA tidal gauge 8772471 (located within the Freeport harbor 14 km southwest of the site, NOAA (2017)) and in the surf zone and back-barrier environments at Follets Island for two time intervals of elevated  $\eta_{vlf}$ ; and high-frequency air pressure  $p_{atm}$  ( $f > 0.1$  mHz) recorded at Follets Island. The dashed lines in (c-d) and (g-h) denote 3-standard deviations ( $\sigma$ ) for  $\eta_{vlf}$  and  $p_{atm}$  (respectively). Air pressure disturbances are identified in (g-h) as pressure couplets with a peak (“H”) or trough (“L”) amplitude that exceeds  $3\sigma$  ( $\pm 1.25$  mbar). The disturbances highlighted orange correspond to the land-falling TCRs (bands of high reflectivity) at Follets Island in Figure 3a and b (respectively).

121 As shown in Figure 2c, relatively large sea-level anomalies ( $\eta$ ) were observed at VLF  
 122 (subscript  $vlf$ ) during three time intervals over the study period at the Follets Island  
 123 field site: -18.5 to -14 hrs, -3 to 3 hrs, and 23.5 to 34 hrs relative to hurricane landfall.  
 124 These time periods were identified by sequential instances of  $\eta_{vlf}$  in the surf zone that  
 125 exceeded (in absolute value) a threshold of 3 times the standard deviation ( $\sigma$ ) of  $\eta_{vlf}$   
 126 (dashed lines, Figure 2c), where time series were bandpass filtered with a low-frequency  
 127 cutoff of 0.1 mHz (2.8 hrs) and high-frequency cutoff of 3 mHz (5.6 min) to isolate  $\eta_{vlf}$ .  
 128 While arbitrary, this threshold (0.16 m) allows for a more targeted examination of forc-

129 ing mechanisms responsible for generation of the largest VLF oscillations. Relatively large  
 130 instances of  $\eta_{vlf}$  were also observed in the back barrier during each of these time peri-  
 131 ods, albeit these sea-level anomalies were small in magnitude ( $<5$  cm). Notably, the back  
 132 barrier environment was only hydraulically connected to the nearshore through a tidal  
 133 inlet located 8 km northeast of the field site (San Luis Pass) and via the Freeport har-  
 134 bor located 15 km southwest of the site (Figure 1a) for the duration of the storm (i.e.,  
 135 no storm overwash or island breaching occurred in the vicinity of the site). Tidal gauges  
 136 operated by the National Oceanic and Atmospheric Administration (NOAA) within Freeport  
 137 harbor (Figure 2d) and at San Luis Pass (not shown) also showed water level variabil-  
 138 ity at VLF, however, only the Freeport harbor gauge measured a VLF anomaly of similar  
 139 magnitude to the surf zone oscillations (0.35 m in height at -18 hrs).

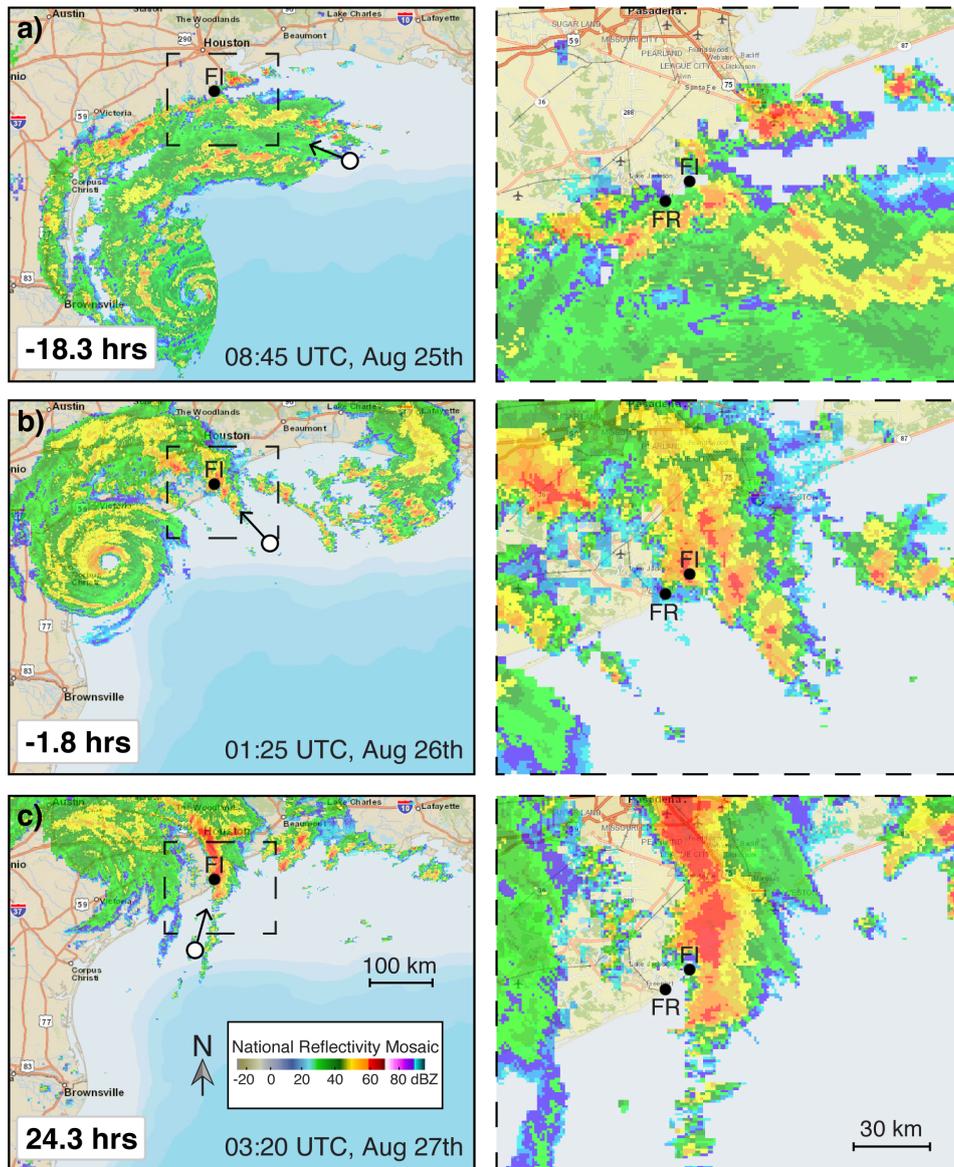


**Figure 2.** a) Time series of high-frequency air pressure  $p_{atm}$ , b) storm surge  $\zeta$ , and c) VLF sea-level anomalies  $\eta_{vlf}$  in the surf zone and back-barrier at the Follets Island field site; d)  $\eta_{vlf}$ , e)  $p_{atm}$ , and f) wind gust speed measured at the NOAA Freeport tidal gauge. Note the different plot scales in (a) and (e). The dashed lines in (a) and (c) denote 3-standard deviations ( $\sigma$ ) for  $p_{atm}$  (Follets Island) and  $\eta_{vlf}$  (surf zone only). Shaded intervals reflect time periods with sequential instances of  $\eta_{vlf}$  (surf zone) above the  $3\sigma$  threshold.

140 The three time periods distinguished by surf zone exceedances of  $|\eta_{vlf}|$  above the  
 141  $3\sigma$  threshold were characterized by a wide range of local water levels (0.2-0.9 m storm  
 142 surge  $\zeta$  – calculated as the surf zone mean water level minus the predicted astronom-  
 143 ical tide – Figure 2b). High resolution measurements of atmospheric (barometric) pres-  
 144 sure recorded by a subaerial mounted pressure transducer in the mid-barrier environ-  
 145 ment at the Follets Island field site (30 sec sampling frequency) revealed that 87% of the  
 146  $3\sigma$  exceedances of  $|\eta_{vlf}|$  (surf zone) were preceded within 30 minutes by air pressure dis-  
 147 turbances. An air pressure disturbance is here defined as an alternating low-to-high pres-  
 148 sure couplet with a peak or trough amplitude that exceeds  $|\pm 1.25|$  mbar (dashed lines  
 149 in Figure 2a). This threshold defines 3 standard deviations of the high-frequency air pres-  
 150 sure  $p_{atm}$  – that is, the barometric pressure high-pass filtered to remove signals repre-  
 151 sentative of the inverted barometer effect associated with the storm-scale tropical depres-  
 152 sion (i.e., storm surge  $f < 0.1$  mHz). As elaborated upon below, this definition allows  
 153 for detailed analysis of temporal changes in meteorological forcing throughout the study  
 154 period. Notably, time series of high-frequency air pressure measured at the Freeport tidal  
 155 gauge (6-min sampling frequency) did not show pressure anomalies in excess of  $\pm 1.25$   
 156 mbar (as elaborated upon in Section 4).

157 On a more regional scale, mosaics of atmospheric radar reflectivity (a measure of  
 158 precipitation intensity) show that air pressure disturbances measured during all three  
 159 time periods were concomitant with bands of high reflectivity associated with passage  
 160 of TCRs. The radar mosaics in Figure 3 depict land-falling TCRs characteristic of each  
 161 of the three time periods of elevated  $\eta_{vlf}$ . For all three time periods, the radially-propagating  
 162 TCRs were oriented approximately perpendicular to the coast offshore Follets Island at  
 163 landfall, however the direction of propagation (from first identification of the TCR in deep  
 164 water, circles in Figure 3) and structure of the TCRs notably differed for each time pe-  
 165 riod. Early in the storm (-18.5 to -14 hrs), Follets Island is located in the upper right  
 166 quadrant of the TC, approximately 300 km NE of the TC eyewall (Figure 3a). The land-  
 167 falling TCRs at the field site during this period were infrequent and radially propagated  
 168 to the WNW as sub-components of larger outer rainbands that spanned most of the cen-  
 169 tral Texas coast ( $>200$  km in arc-length). As the storm center moved onshore and closer  
 170 to the field site (eyewall  $\sim 200$  km SW), trains of rainbands with very small horizontal  
 171 scales ( $<50$  km in arc-length, -3 to 3 hrs, Figure 3b) passed frequently over the field site  
 172 (typically every  $\sim 30$  minutes), propagating towards the NW. We classify these very-narrow  
 173 TCRs as outer rainbands as they are located far outside the inner core region – that is,  
 174 beyond 3 times the radius of maximum wind as defined by Wang (2009) – which dur-  
 175 ing this time period was very compact (Alford et al., 2019) ( $\sim 60$  km as estimated from  
 176 RMS HWind by Brown-Giammanco et al. (2018)). Narrow arc-length TCRs ( $<75$  km)  
 177 were also observed late in the study period (23.5 to 34 hrs, Figure 3c) when the storm  
 178 was stalled inland  $\sim 180$  km west of the field site. During this time, the land-falling TCRs  
 179 at Follets Island were oriented slightly more oblique to the coastline and again sub-components  
 180 of larger outer rainbands ( $>100$  km in arc-length) propagating toward the NE. The for-  
 181 ward translation speed of the TCRs across the continental shelf during each of the three  
 182 time periods was highly variable, ranging from 6 to 29 m/s along the path of a single TCR  
 183 (as estimated from the reflectivity mosaics). However, the mean forward speeds of the  
 184 three land-falling TCRs shown in Figure 3 were similar (a,b: 17 m/s and c: 14 m/s).

185 Surface pressure fluctuations of land-falling TCRs are often characterized by lead-  
 186 ing pressure troughs followed by pressure ridges, or alternatively low-to-high pressure  
 187 couplets (e.g. Ligda, 1955; Ushijima, 1958; Hamuro et al., 1969; Yu & Tsai, 2010). The  
 188 troughs (ridges) of air pressure disturbances measured at Follets Island are delineated  
 189 with an “L” (“H”) in Figure 1a and b. Whilst the identification of pressure couplets is  
 190 subjective, there is clearly a significant increase in the number of air pressure disturbances  
 191 for the period proximate to hurricane landfall over early storm conditions, a result of more  
 192 frequent passage of trains of outer TCRs over the field site. The air pressure disturbances  
 193 proximate to landfall are also characterized by larger pressure trough amplitudes and



**Figure 3.** Radar reflectivity mosaics generated by the NOAA National Center of Environmental Information (NCEI-NOAA, 2019) showing bands of high reflectivity associated with land-falling tropical cyclone rainbands (TCRs) preceding instances of elevated  $\eta_{vlf}$  (above the  $3\sigma$  threshold) during each of the three time periods identified in Figure 2c. Note that the plots on the right depict the TCRs shown in the left at a finer resolution proximate to the Follets Island (FI) field site and the Freeport harbor tidal gauge (FR). The circle and arrow symbols show the direction of propagation of the land-falling TCRs from first identification in deep water.

194 shorter trough periods than the more infrequent outer TCRs typical of early-storm con-  
 195 ditions, despite similar estimated forward speeds at landfall.

196 The air pressure disturbances highlighted orange in Figure 1g-h were concomitant  
 197 with passage of the TCRs shown in Figure 3a and b (respectively). Due to the highly-  
 198 localized nature of the land-falling TCRs during both the period proximate to hurricane  
 199 landfall (-5 to 0 hrs) and late in the storm (23.5 to 34 hrs), the measured peak and trough  
 200 amplitudes of the air pressure disturbances may be underestimated. However, for the  
 201 disturbance highlighted in Figure 1h, radar reflectivity mosaics show that the TCR tra-  
 202 versed directly over the Follets Island field site (Figure 3b), which gives high credence  
 203 to the structure of this pressure couplet. In the following section, we employ a numer-  
 204 ical model to explore the generation potential of meteotsunami to a simplified represen-  
 205 tation of the measured air pressure disturbance highlighted in Figure 1h. This time pe-  
 206 riod is of particular interest since large VLF variability in sea level was only observed  
 207 in the surf zone and not at nearby tidal gauges. The model is then used in an exploratory  
 208 framework to better understand potential factors that influence the sea-level response  
 209 to direct pressure forcing by a very narrow outer TCR.

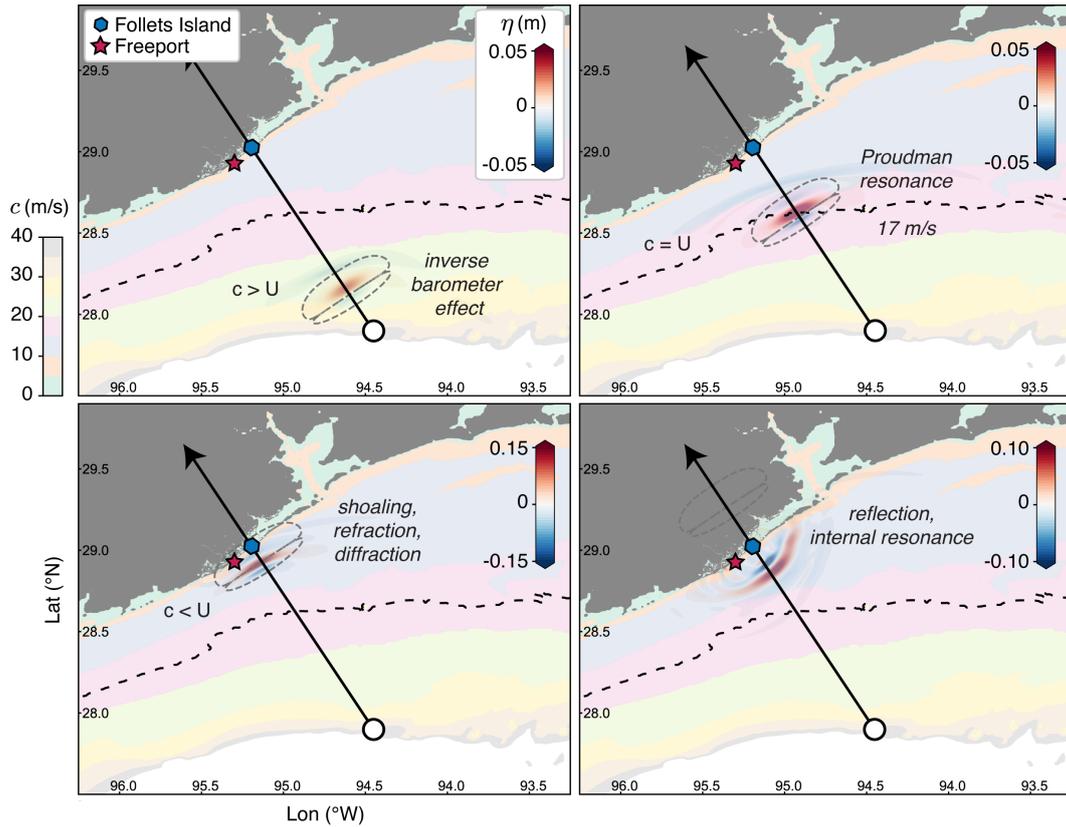
### 210 3 Numerical Modeling

211 The numerical model utilized in this study is a 2D (depth-integrated) hydrostatic  
 212 model in spherical coordinates built on the nonlinear shallow water equations, modified  
 213 to include spatially-dependent air pressure. The governing equations, staggered grid setup,  
 214 and numerical solution scheme are outlined in Kowalik et al. (2005). The modeling do-  
 215 main encompasses the continental shelf along the eastern GOM and extends landward  
 216 to include harbors and bays within the greater Freeport and Houston/Galveston region  
 217 (Figure 4). Bathymetry data was created with a base layer of the NOAA Etopo1 dataset  
 218 (Amante & Eakins, 2009) (27.5 to 29.9 N, -96.25 to -93.25 W) interpolated to a grid res-  
 219 olution of 6 arcseconds ( $\sim 185$  m) and referenced to the mean high water. This grid spac-  
 220 ing was selected to achieve reasonable simulation time, while satisfying the requirement  
 221 of having at least 20 grid points to numerically represent the wavelength of the air pres-  
 222 sure disturbance. For the Freeport area (28.75 to 29.25 N, -95.5 to -95 W), the baseline  
 223 grid is superimposed by a 1/9 arcsecond grid (NCEI, 2015) averaged to 6 arcseconds,  
 224 which allows for high resolution of bathymetric variability within the Freeport Harbor,  
 225 albeit subaerial structures such as the Freeport jetties are not resolved. The still water  
 226 level was modified to incorporate the effect of storm surge on water depth for select sim-  
 227 ulations herein. Bottom friction is based on the Manning model with a Manning coef-  
 228 ficient of  $0.025 \text{ sm}^{(1/3)}$ . A coastal wall is set at a water depth of 0.3 m to avoid runup  
 229 on the  $\sim 185$ -m wide land cells as well as numerical instabilities at zero water depth. Out-  
 230 flow boundary conditions are applied to all the boundaries of the model domain. Model  
 231 outputs are recorded at the Follets Island surf zone measurement location and the Freeport  
 232 harbor tidal gauge every 20 seconds. Astronomical tides are not included in the com-  
 233 putations.

Due to a lack of spatial information on the characteristics of the air pressure dis-  
 turbance, we assume a surface pressure function where the amplitude of the crest  $A_c$  (0.83  
 mbar) and trough  $A_t$  (2.25 mbar) decay exponentially along the length  $L$  of the TCR  
 (23 km)

$$P(x, y) = \begin{cases} A_c * x * \exp(-(y)^2 - \left(\frac{x}{L_c}\right)^2), & x < 0 \\ A_t * x * \exp(-(y)^2 - \left(\frac{x}{L_t}\right)^2), & x > 0 \end{cases} \quad (1)$$

234 where  $(x, y)$  are the longitude and latitude excursions along length  $L$  (estimated from  
 235 radar reflectivity) and  $L_c$  and  $L_t$  are the wavelength of the pressure crest and trough,  
 236 respectively. As discussed in more detail below, for the range of forward speeds simu-  
 237 lated here,  $L_c$  ( $L_t$ ) spanned 1.8-8.7 km (4.7-22.6 km). Although the path of a TCR is



**Figure 4.** Overview of the model domain and simulated generation mechanism of nearshore meteotsunami by continuous pressure forcing over the continental shelf. The colored contours represent the shallow water wave celerity  $c = \sqrt{gh}$  in terms of the water depth  $h$  and the blue-red colormap the instantaneous sea-level anomaly  $\eta$ . The bold dashed contour line identifies the location of Proudman resonance – that is, when the speed of the air pressure disturbance  $U$  matches  $c$  – and maximum energy transfer (here, 17 m/s); the solid black line shows the path along which the disturbance travels (offshore to onshore); and the dashed gray lines indicate the spatial extent of the pressure disturbance.

238 generally radial, the translation of the pressure disturbance is here simplified to a lin-  
 239 ear path (single direction) beginning at the edge of the continental shelf and extending  
 240 inland past the Follets Island field site.

241 Figure 4 provides an overview of the processes responsible for meteotsunami genera-  
 242 tion, amplification, and propagation in the numerical simulations, illustrated for the  
 243 scenario of a 17 m/s TCR forward translation speed  $U$  (i.e., the average forward speed  
 244 of the TCR from radar reflectivity) and 0.71 m of storm surge (as measured in the surf  
 245 zone during this time period). In deep water, the air pressure disturbance acts on the  
 246 water surface following the inverse barometer effect. In contrast to storm surge genera-  
 247 tion (inverse barometer effect acting over large oceanic regions), the subsequent sea-  
 248 level anomaly is small, on the order of several cm. The celerity of the forced wave (i.e.,  
 249 the Proudman resonance contours) is faster than the speed of the disturbance, which is  
 250 faintly visible by the blue sea-level trough located ahead of the leading edge of the pres-  
 251 sure disturbance. As the air pressure disturbance travels across the shelf, a sea-level per-  
 252 turbation initiated by the inverse barometer effect can grow in height due to Proudman

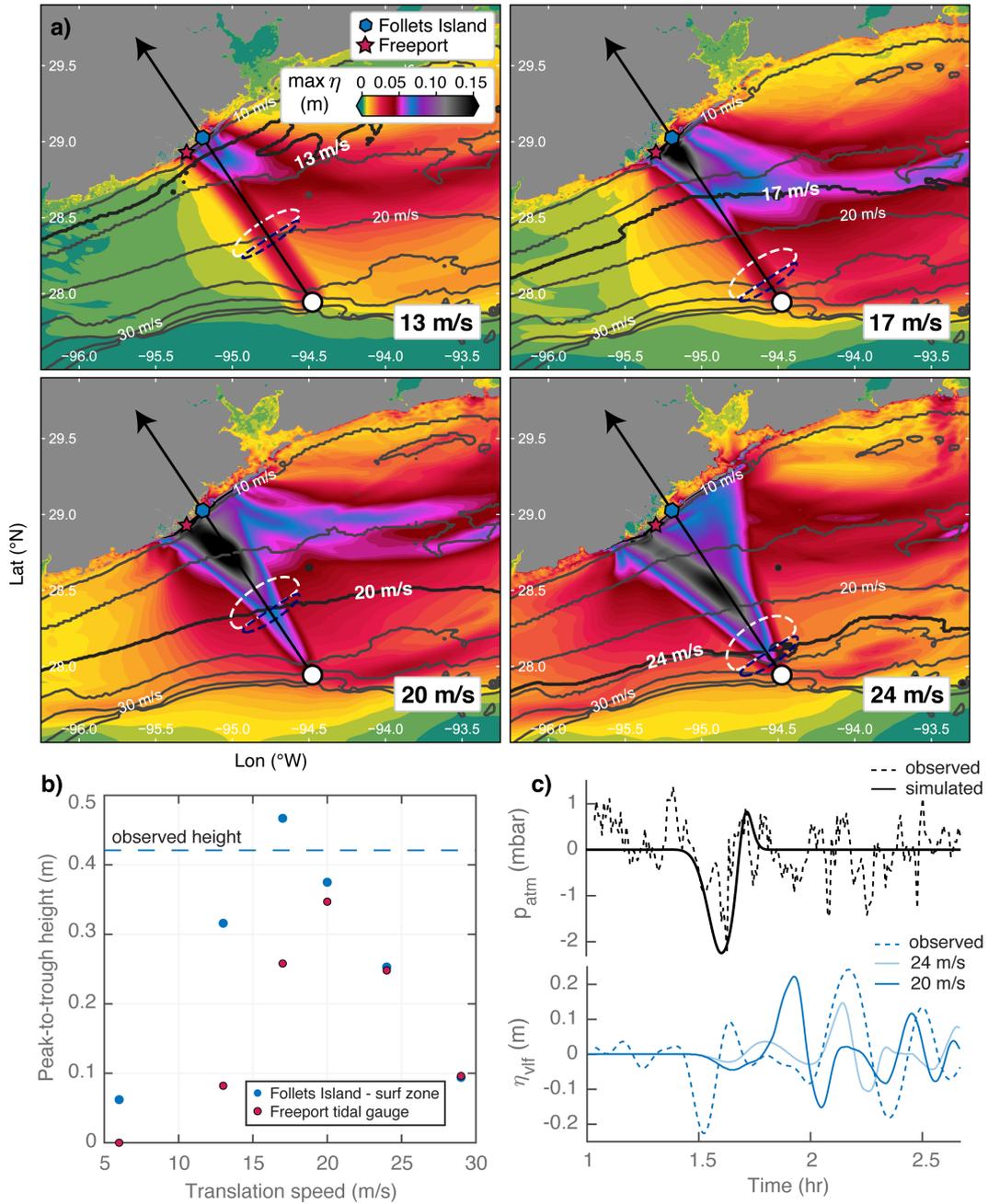
253 resonance. The forward speed of the TCR determines the region (water depth) where  
 254 resonant amplification occurs (Proudman resonance) – that is, when the speed of the air  
 255 disturbance matches the shallow water wave celerity of the region – which in Figure 4  
 256 is delineated by a dashed bathymetry contour corresponding to a shallow water wave celer-  
 257 ity of 17 m/s (30-m depth). Importantly, the contour of Proudman resonance also dem-  
 258 marcates the detachment location at which the resonantly-amplified wave becomes slower  
 259 than the air pressure disturbance and can thereafter propagate as a free wave. The trans-  
 260 formation of the resonantly-amplified wave in the coastal zone through processes such  
 261 as shoaling, refraction, diffraction, and reflection results in a maximum crest elevation  
 262 (herein referred to as the “maximum sea-level anomaly”) of 0.31 m (peak-to-trough height  
 263 of 0.47 m) in the surf zone for the 17 m/s scenario. Lastly, while the model can simu-  
 264 late wave amplification by internal resonance processes in coastal bays and harbors, this  
 265 phenomenon was not clearly evident in any of the numerical simulations described herein.

### 266 3.1 Sensitivity to TCR forward speed

267 Radar reflectivity measurements indicate that the forward translation speed of the  
 268 TCR varied across the continental shelf between 6 and 29 m/s, which with regard to me-  
 269 teotsunami generation, changes the depth and location of Proudman resonance (Figure  
 270 5a). For the bathymetry and coastal morphology offshore Follets Island, TCR forward  
 271 translation speeds above 17 m/s allow the resonantly-amplified wave to detach from the  
 272 air pressure disturbance far offshore and thereafter refract towards Freeport. Consequently,  
 273 Figure 5b shows that for the imposed cross-shelf perpendicular trajectory of the TCR  
 274 in this study, relatively large meteotsunami (peak-to-trough wave heights  $>0.25$  m) are  
 275 only observed in the surf zone at Follets Island (Freeport harbor) for a subset of TCR  
 276 forward speeds, namely 13-24 m/s (17-24 m/s).

277 The time evolution of simulated and observed pressure and sea level at Follets Is-  
 278 land are shown in Figure 5c for the 20- and 24-m/s forward speed scenarios. These sce-  
 279 narios were selected for comparison as they most closely matched 1) the observed lag be-  
 280 tween passage of the air pressure disturbance (pressure trough) and arrival of the peak  
 281 sea-level anomaly (the 24 m/s scenario, not shown), and 2) the peak meteotsunami wave  
 282 height (the 20 m/s scenario, Figure 5b). Note that in order to match the period of the  
 283 air pressure disturbance at landfall with field observations, the wavelength of the dis-  
 284 turbance was varied for each forward speed simulation (6.5-31.3 km, see the changing  
 285 spatial extent of the dashed lines in Figure 5a). Although the 20 m/s forward speed sce-  
 286 nario produces a sea-level anomaly of similar magnitude (observed = 0.24 m, simulated  
 287 = 0.22 m) and period (observed = 23 min, simulated = 20 min) to the observed VLF  
 288 anomaly in the surf zone, the lag is shorter than observed (observed = 32 min, simulated  
 289 = 20 min). Conversely, the simulated peak-to-trough height for the 24 m/s scenario is  
 290 smaller than observed, despite similar lag times. Lastly, VLF wave heights at the Freeport  
 291 tidal gauge were  $<10$  cm (Figure 2d) whereas numerically simulated meteotsunami reached  
 292 24 and 35 cm in height for the 24- and 20-m/s scenarios, respectively. The source of this  
 293 mismatch is likely due to model bathymetry, in that the Freeport jetties that shelter the  
 294 tidal gauge from direct wave impact are not resolved.

295 While the measured air pressure disturbances can reproduce meteotsunami of sim-  
 296 ilar height and period to the surf zone observations using this simplified model frame-  
 297 work, as elaborated upon below, discrepancies between simulated and measured wave  
 298 characteristics may stem from model uncertainties (bathymetry), model simplification  
 299 (idealized and temporally-constant representation of the pressure waveform, Williams  
 300 et al. (2020)), and missing physical processes (wind forcing, Shi et al. (2020)).



**Figure 5.** Sensitivity of the simulated a) maximum sea-level anomaly  $\eta$  and b) the maximum peak-to-trough wave height to select TCR forward translation speeds within the range of observations (6-29 m/s). The still water was modified to incorporate the effect of storm surge (0.71 m) on water depth and the dashed line in (b) demarcates the observed peak-to-trough wave height in the surf zone at Follets Island. c) Comparison of the time evolution of the simulated and observed air pressure disturbance ( $p_{atm}$ ) and sea-level response (20 and 24 m/s only) at Follets Island.

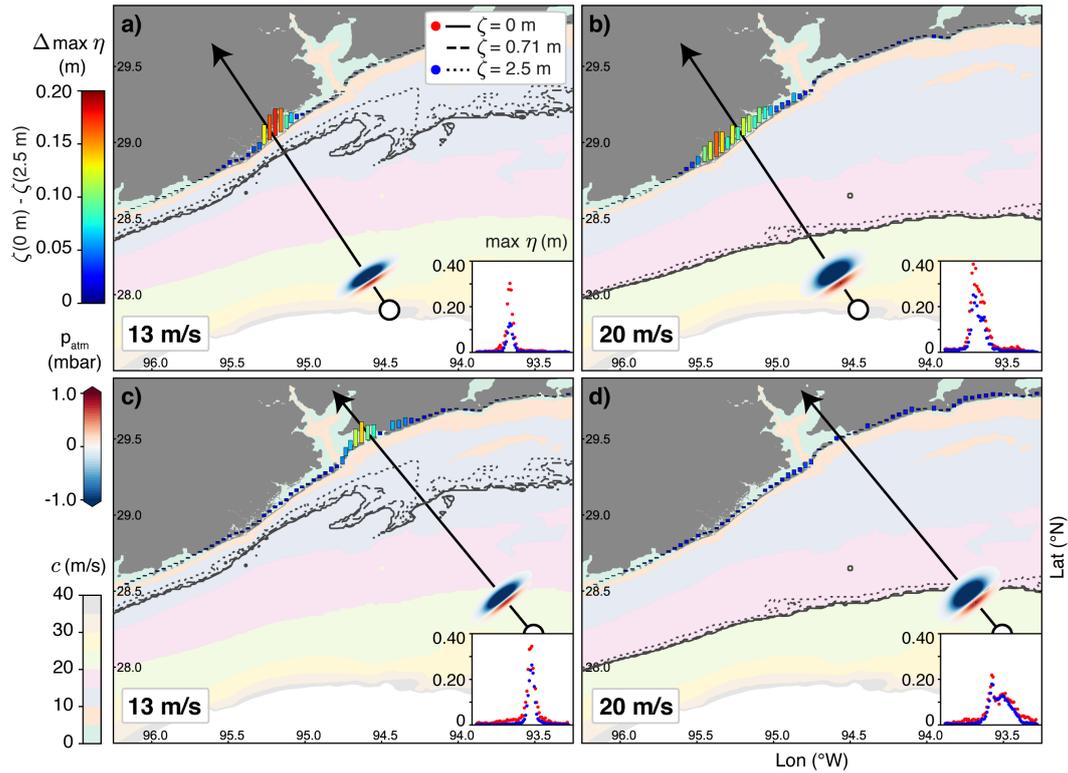
301

### 3.2 Sensitivity to storm surge

302

Using the model setup as an exploratory framework, the sensitivity of meteotsunami generation and propagation to changes in storm surge was examined for variable TCR

303



**Figure 6.** Sensitivity of meteotsunami generation, amplification, and propagation to storm surge  $\zeta$  for variable TCR forward speeds and landfall locations. Meteotsunami generation varies with surge via the location of Proudman resonance (bold contour lines). Changes in meteotsunami amplification and propagation are depicted by the vertical colored bars which show the difference in the maximum sea-level anomaly  $\eta$  between the case of no surge and the largest simulated surge (2.5 m) along the entire coastline, with the maximum values for each scenario shown in the subplot for direct comparison.

304 forward translation speeds at both the original Follets Island landfall location (Figure  
 305 6a-b) and 80 km to the northeast where the continental shelf is wider (Figure 6c-d). Here  
 306 we impose a larger surge of  $\zeta=2.5$  m (via an increase in still water), which is compar-  
 307 able to the surge at Follets Island during Hurricane Ike in 2008 (2.6 m, Harter & Figlus  
 308 (2017)).

309 The vertical colored bars in each pane of Figure 6 show the difference in the max-  
 310 imum sea-level anomaly between the case of zero surge and the 2.5 m surge scenario along  
 311 the entire coastline for each simulation. In all cases, the effect of storm surge on meteot-  
 312 sunami generation is to increase water depth and thereby move the location of Proud-  
 313 man resonance landward. Hence, the effect of surge on meteotsunami amplification and  
 314 propagation is complex, and varies both with TCR forward translation speed as well as  
 315 the offshore bathymetry and coastal morphology. For example, comparing Figure 6a and  
 316 b, the effect of an increase in surge on meteotsunami hazard (via an increase in the max-  
 317 imum sea-level anomaly at the coast) is larger for a relatively slow moving TCR (13 m/s)  
 318 than for a relatively fast moving TCR (20 m/s) for TCR landfall at Follets Island. Al-  
 319 though the meteotsunami surge response is also sensitive to changes in forward speed  
 320 farther up the coast (Figures 6c-d), comparison of the two landfall locations shows that  
 321 the decrease in meteotsunami hazard with an increase in surge is larger at Follets Island

322 due to a slightly steeper sloping continental shelf (i.e., narrower meteotsunami enhance-  
 323 ment region). Given the sensitivity of meteotsunami surge response to offshore bathymetry  
 324 and coastal morphology for our simplified linear TCR trajectory and constant forward  
 325 speed, it is likely that the effect of surge on meteotsunami hazard is made further com-  
 326 plex by more radial TCR translation paths and/or varying speed.

## 327 4 Discussion

328 Many of the VLF sea-level anomalies observed in the nearshore environment along  
 329 the upper Texas Gulf coast during Hurricane Harvey appear to be initiated by moving-  
 330 atmospheric disturbances associated with radially-propagating outer TCRs. This hypoth-  
 331 esis is supported by numerical modeling of the sea-level response to measured air pres-  
 332 sure forcing proximate to hurricane landfall, a time period characterized by frequent pas-  
 333 sage of very narrow outer TCRs (<50 km in arc-length, Figure 3b) and large air pres-  
 334 sure disturbances (1-2 mbar pressure troughs, Figure 1h). The numerical model repro-  
 335 duces meteotsunami similar to the observed VLF sea-level anomalies in the surf zone,  
 336 albeit the simulated magnitude, period, and spatial extent of meteotsunami hazard (i.e.,  
 337 the maximum sea-level anomaly at the coast) along the open coast is highly dependent  
 338 on the shelf bathymetry as well as the forward speed and path of translation of the air  
 339 pressure disturbance (Figure 5). These findings are consistent with other numerical in-  
 340 vestigations of Proudman resonant wave growth (e.g. Ličer et al., 2017; Shi et al., 2020;  
 341 Williams et al., 2020). Here we find that the meteotsunami hazard additionally depends  
 342 on the magnitude of storm surge, which acts to move the location of Proudman resonance  
 343 landward (Figure 6). For the atmospheric (TCR structure, forward speed, path of trans-  
 344 lation) and oceanographic (bathymetric configuration) conditions explored in this study,  
 345 an increase in storm surge results in a decrease in meteotsunami hazard. These numer-  
 346 ical results suggest that along this open coastline, meteotsunami hazard is largest for rel-  
 347 atively low surge events, like Hurricane Harvey.

348 It is unknown how increases in model complexity to incorporate additional oceanog-  
 349 raphic (astronomical tides and currents) and atmospheric factors (radial TCR propa-  
 350 gation, variable TCR forward speed, temporal-modifications to the air pressure wave-  
 351 form, wind forcing) will effect meteotsunami hazard, and therefore the relative effect of  
 352 storm surge for the cases simulated here. Coupled ocean-atmosphere simulations are needed  
 353 for this purpose, however, to numerically reproduce the high frequency (18 min period)  
 354 and short wavelength (6.5-31 km, for the range of potential forward speeds) air pressure  
 355 disturbances measured during Hurricane Harvey, hydrodynamic models need to be forced  
 356 with higher temporal (<1 min) and spatial (<500 m) meteorological data than used in  
 357 previous studies (e.g., the 3 km grid spacing and 5 min temporal resolution of the Weather  
 358 Research and Forecast (WRF) model used by Shi et al. (2020)).

359 In this study, we focused on the generation potential of meteotsunami by narrow  
 360 trains of outer TCRs. Shi et al. (2020) showed that meteotsunami can also be triggered  
 361 by wind and pressure forcing accompanying outer TCRs that are squall-line like in struc-  
 362 ture. Hence, it is likely that many of the VLF sea-level anomalies identified in the surf  
 363 zone that were coincident with passage of TCRs that are squall-line like in structure far  
 364 before (-18.5 to -14 hrs, Figure 3a) and after storm landfall (23.5 to 34 hrs, Figure 3c)  
 365 are meteotsunami. These observations suggest that, for Hurricane Harvey, meteotsunami  
 366 hazard spanned several days before and after the peak storm surge (Figure 2) along open  
 367 coast located to the right of the TC track.

368 Figure 1 demonstrates that not all atmospheric disturbances were followed by large  
 369 sea-level anomalies in the surf zone at Follets Island, and notably, only a single sea-level  
 370 anomaly of comparable magnitude to surf zone observations was measured at regional  
 371 tidal gauges (0.35 m in peak-to-trough height, Figure 1b). Evaluation of radar reflectiv-  
 372 ity and meteorological data at the Freeport tidal gauge during this time shows that a

373 squall-line like TCR traversed directly (and nearly perpendicular) over the harbor and  
 374 it was accompanied by a sharp increase in wind gusts (17 m/s, Figure 2f). Given the ori-  
 375 entation of the harbor jetties, it is likely that meteotsunami propagation into this shel-  
 376 tered harbor would require normal incidence, a hypothesis that should be explored fur-  
 377 ther through numerical modeling. Regardless, it is clear from the surf zone observations  
 378 presented here that during TCs, meteotsunami are more ubiquitous along the open coast  
 379 than tidal gauge records suggest. This finding likely also stems from the inherent nar-  
 380 row structure of some outer TCR rainbands, which as demonstrated numerically (Fig-  
 381 ure 5a), trigger meteotsunami with highly-localized propagation and spatial extent (10s  
 382 of kms). Therefore, efforts to produce flood-risk forecasts that incorporate meteotsunami  
 383 generation potential by TCRs will require accurate forecasts of individual TCRs during  
 384 storms, which remains challenging. As the pressure couplets accompanying TCRs are  
 385 significantly influenced by convective precipitation (Yu & Tsai, 2010; Yu et al., 2018),  
 386 climate induced changes to TCR convection due to a warming climate may magnify the  
 387 associated pressure anomalies, and thereby meteotsunami generation via Proudman res-  
 388 onance.

389 Although meteotsunami were typically small in height in the surf zone during Hur-  
 390 ricane Harvey (<42 cm), meteotsunami contributions to the total water level variance  
 391 in the nearshore were large, reaching a maximum of 23% in the surf zone and 78% in the  
 392 back-barrier bay at Follets Island. In some cases, meteotsunami wave heights were com-  
 393 parable to the increase in total water level by storm surge (e.g., ~30-40 cm from -18.5  
 394 to -15 hrs and +25 to +35 hrs, Figure 2b). VLF sea-level anomalies were also observed  
 395 during storm-driven overwash at Matagorda Peninsula, a barrier peninsula located ~85  
 396 km southwest of Follets Island and closer to storm landfall. Anarde et al., (*in review*)  
 397 hypothesize that these VLF anomalies are likewise meteotsunami triggered by TCRs,  
 398 and using field data show that the slow variation of total water depth associated with  
 399 this phenomenon slightly modulate infragravity wave heights during overwash. Although  
 400 meteotsunami clearly modify storm processes in very shallow water, it is unknown whether  
 401 meteotsunami are important contributors to sediment suspension and flux. A higher den-  
 402 sity of field measurements is needed to characterize meteotsunami transformation in the  
 403 nearshore and the relative contribution of this phenomena to morphological change (i.e.,  
 404 beach and dune erosion) during hurricane impact.

## 405 5 Conclusions

406 Measurements of hydrodynamic and meteorological forcing in the nearshore envi-  
 407 ronment during Hurricane Harvey provide new insights into processes that contribute  
 408 to meteotsunami hazard along open coasts during TCs. Co-located measurements of sea  
 409 level and air pressure along the upper Texas coast show that sea-level anomalies in the  
 410 meteotsunami frequency band (~8 to 45 min periods) occur with large changes in air  
 411 pressure accompanying passage of TCRs. It is demonstrated using numerical modeling  
 412 that drops in air pressure  $O(2$  mbar) concomitant with trains of very narrow outer TCRs  
 413 (<50 km in arc-length) can trigger meteotsunami similar in period  $O(20$  min) and height  
 414  $O(0.4$  m) to surf zone sea-level anomalies. Hence, we find that pressure forcing (air dis-  
 415 turbances with periods typically <30 min) accompanying trains of outer TCRs can re-  
 416 sult in large resonant amplification of sea-level anomalies. This finding is in direct con-  
 417 trast with idealized numerical modeling studies which have shown that the sea-level re-  
 418 sponse to pressure forcing (air disturbances with 1-2 hr periods) accompanying passage  
 419 of both outer and inner TCRs follows the inverse barometer effect. Our numerical sim-  
 420 ulations show that the region of Proudman resonance, and thereafter meteotsunami prop-  
 421 agation, is highly sensitive to both oceanographic and atmospheric conditions, includ-  
 422 ing TCR forward speed, TCR path of translation, bathymetric configuration, coastal mor-  
 423 phology, and storm surge. For this open coast and the very narrow trains of outer TCRs  
 424 explored here, meteotsunami hazard via Proudman resonance is largest for fast moving

TCRs (13-24 m/s) and low-levels of surge (<1 m). Lastly, comparison of field observations in the surf zone and at neighboring tidal gauges show that meteotsunami are more ubiquitous along the open coast during TCs than tidal gauge records suggest due to the inherent structure and highly-localized propagation of TCRs which likely limits meteotsunami propagation into tidal inlets, harbors, and coastal bays.

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